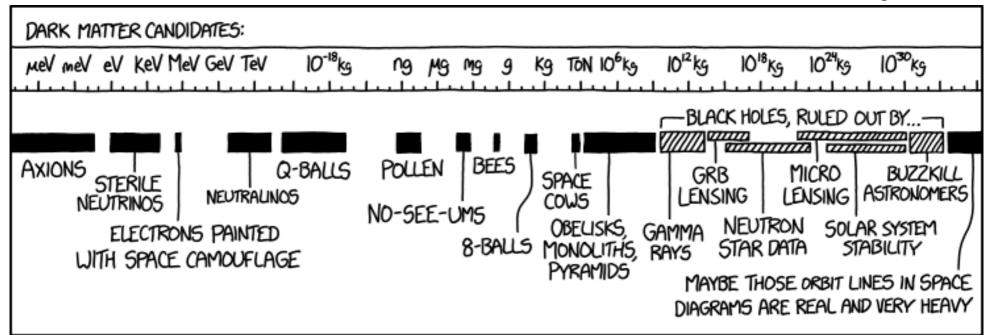
Observational signals of compact dark stars

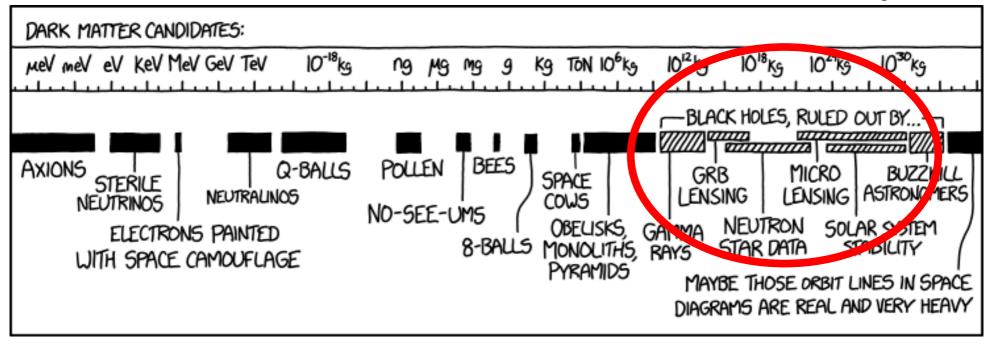
Alejandro Ibarra

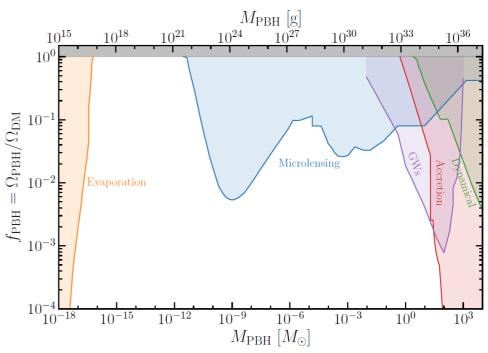


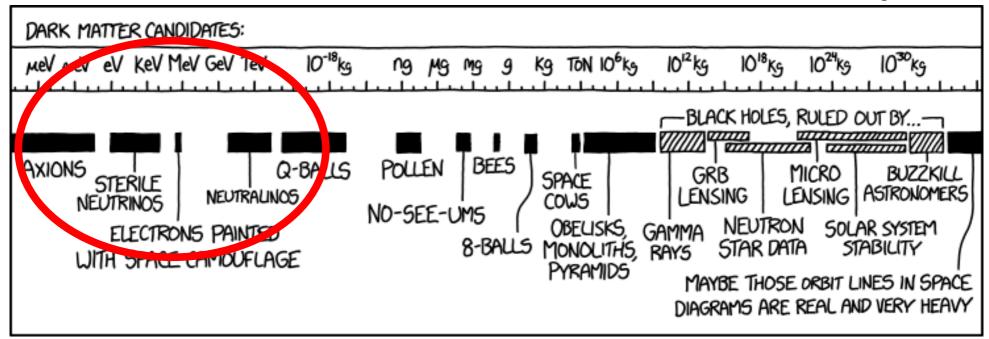


In collaboration with Boris Betancourt, Anja Brenner and Chris Kouvaris. arXiv: 2211.05845

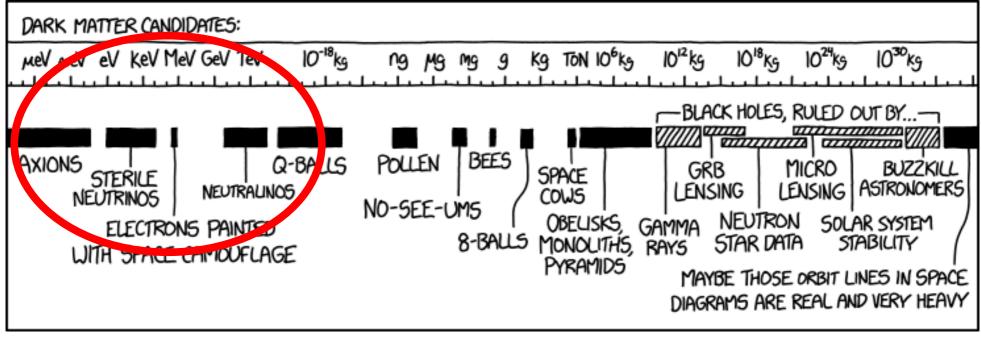




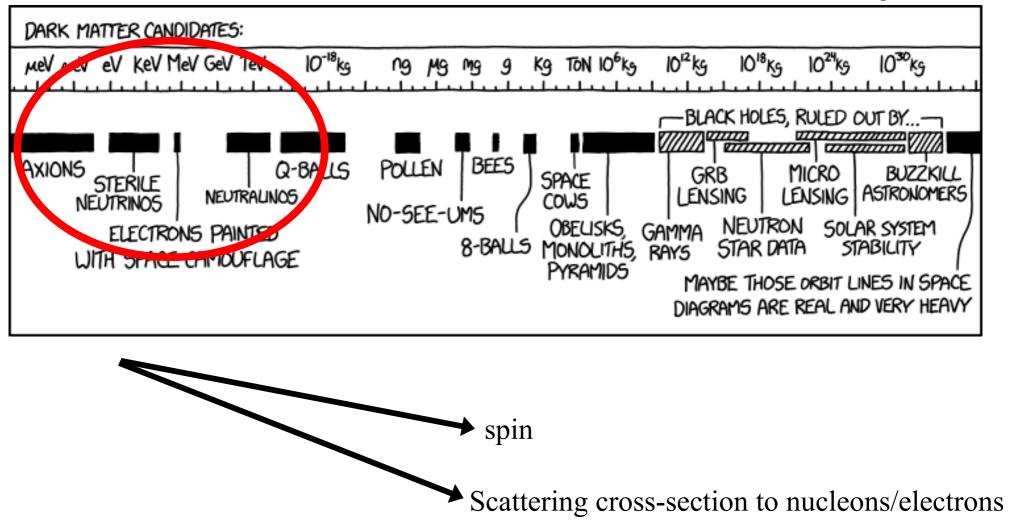


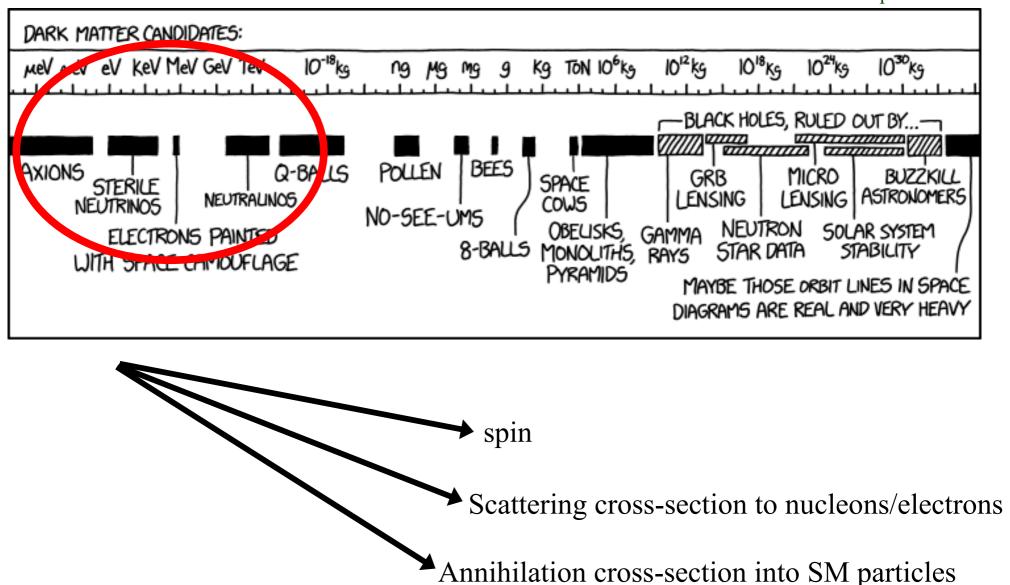


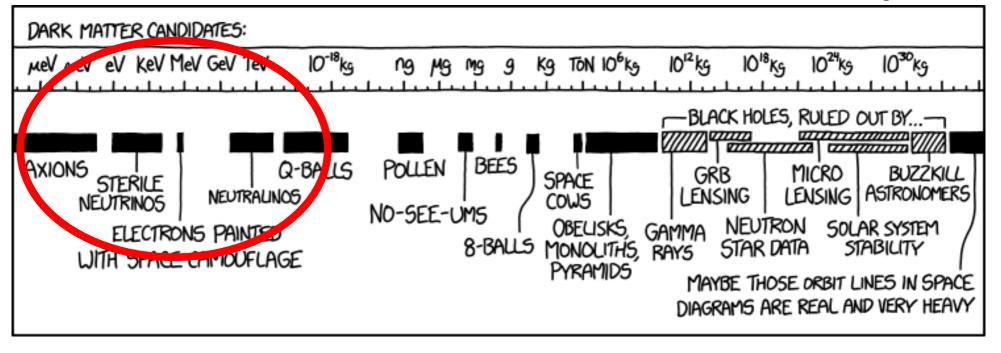
Explain xkcd

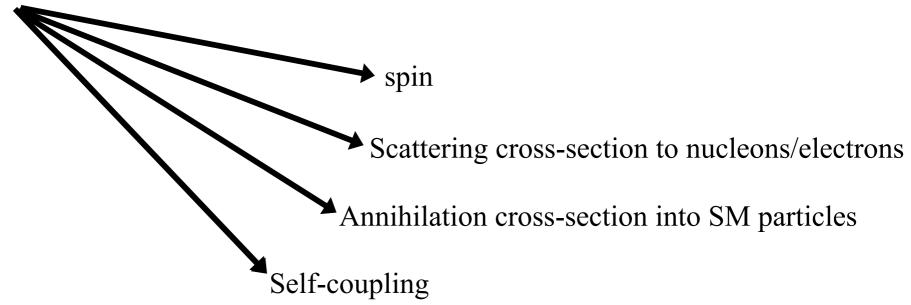


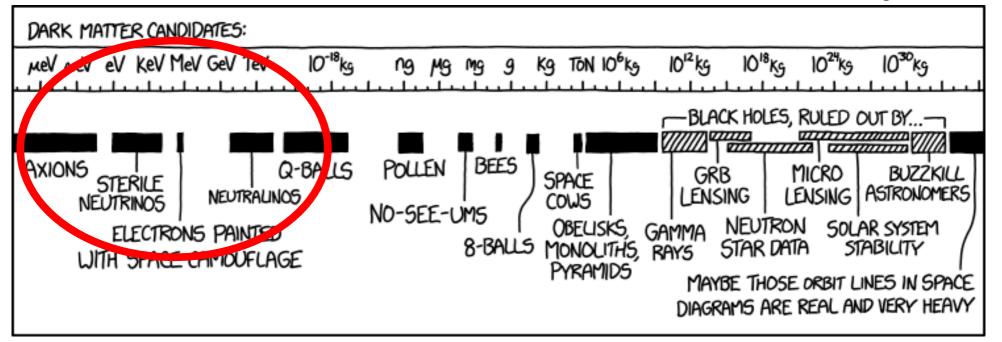
→ spin

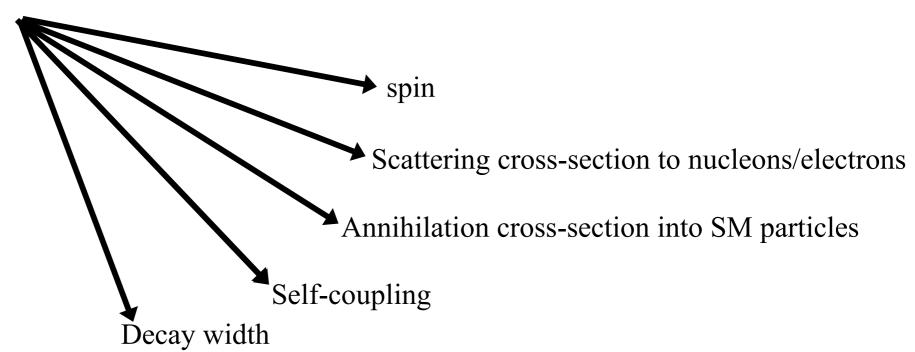


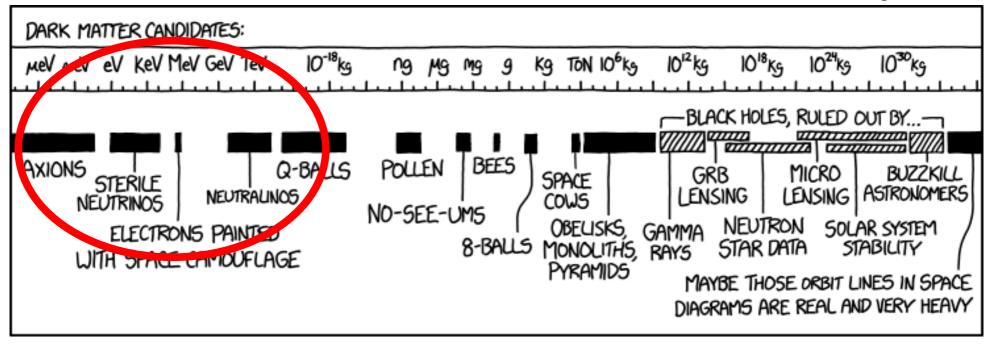


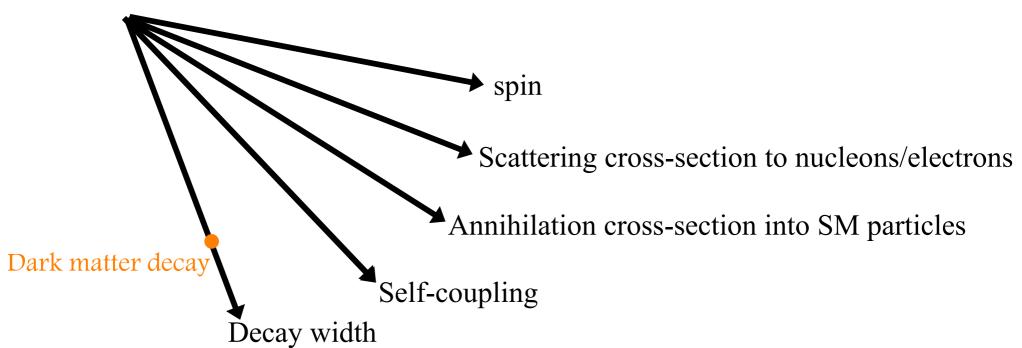


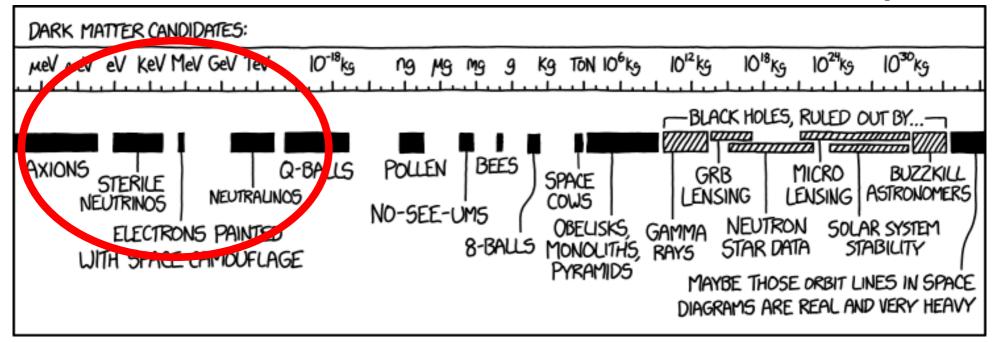


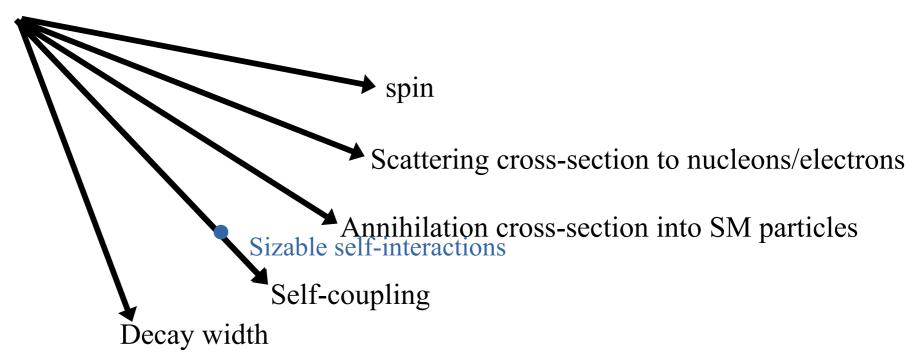


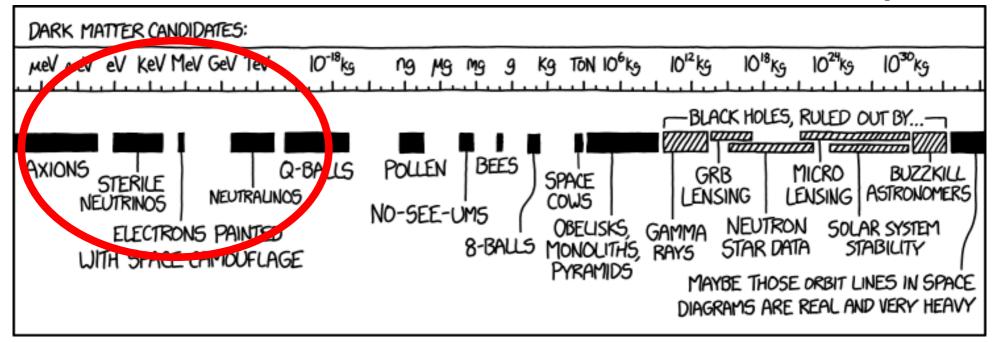


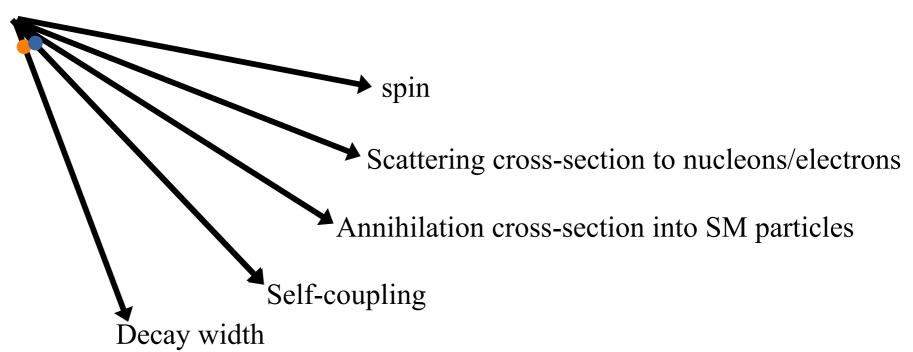


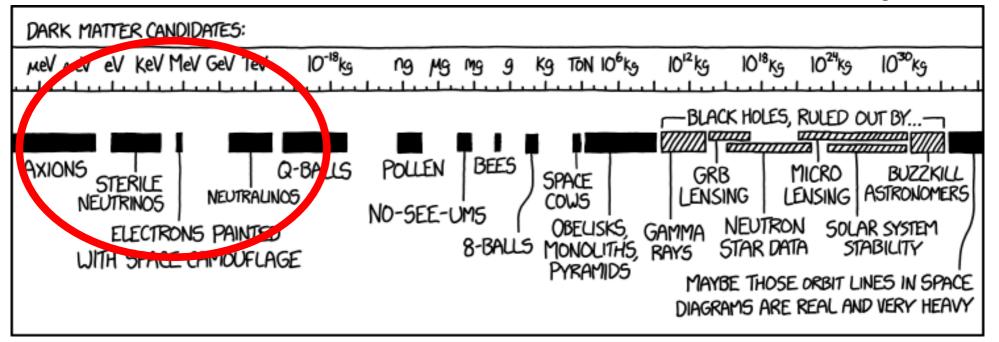


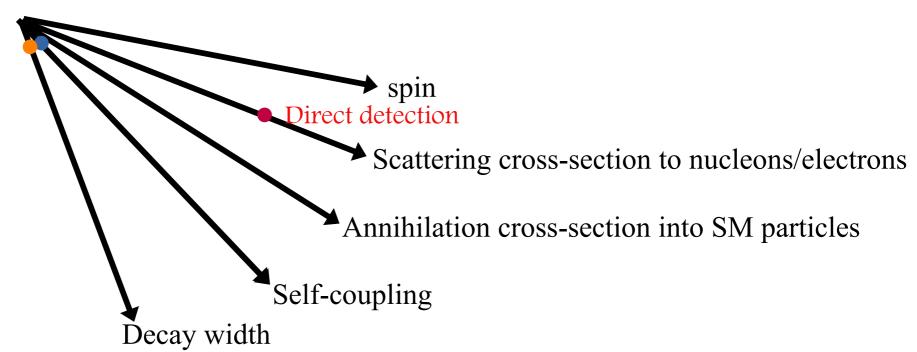


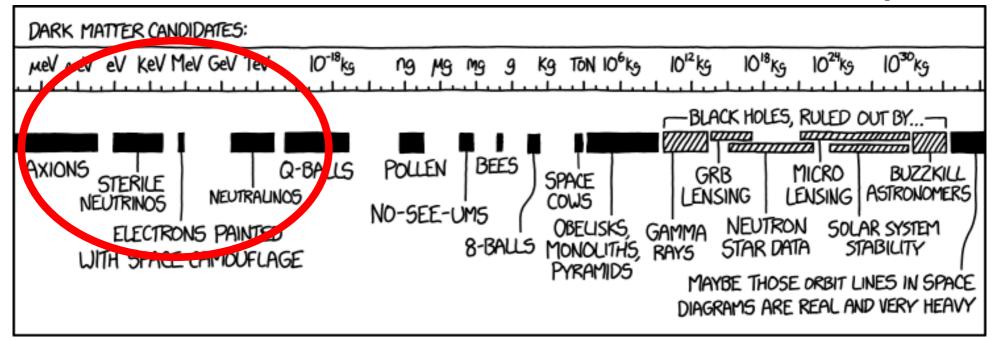


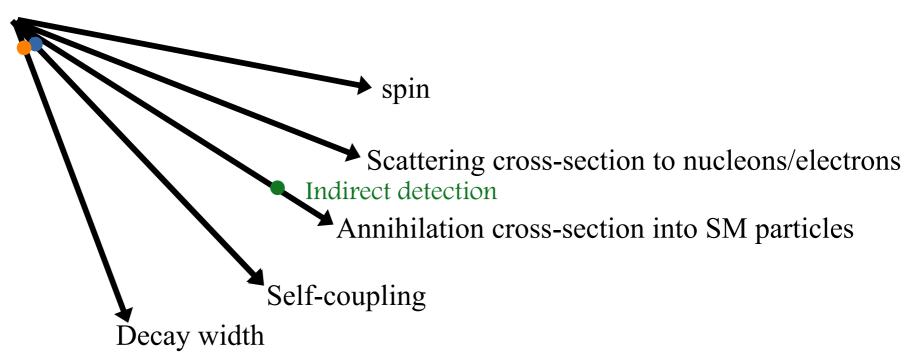


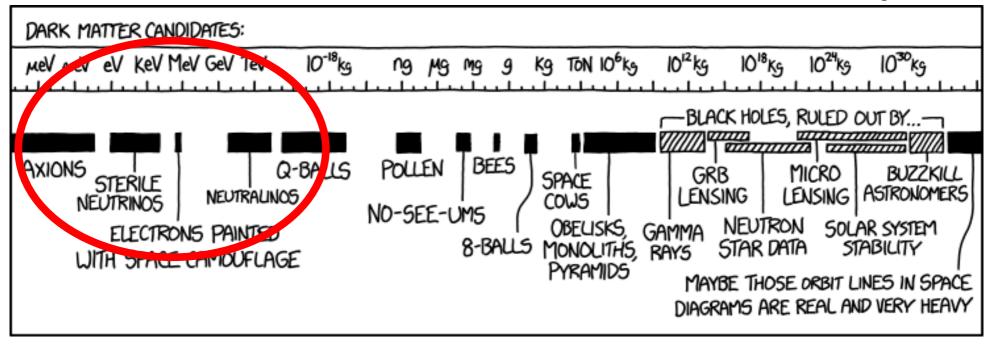


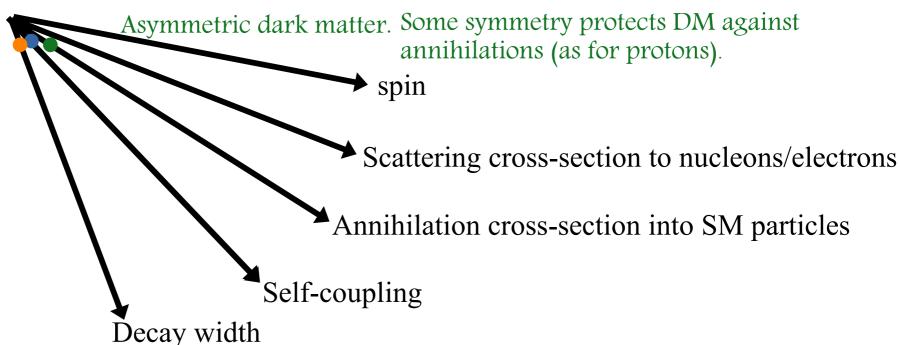


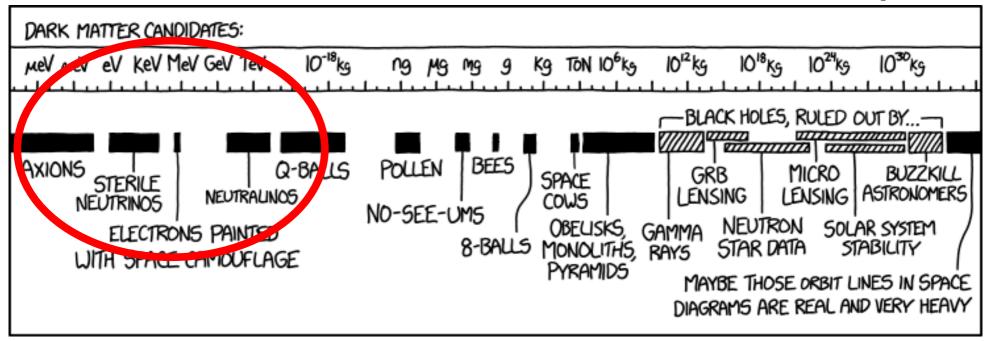


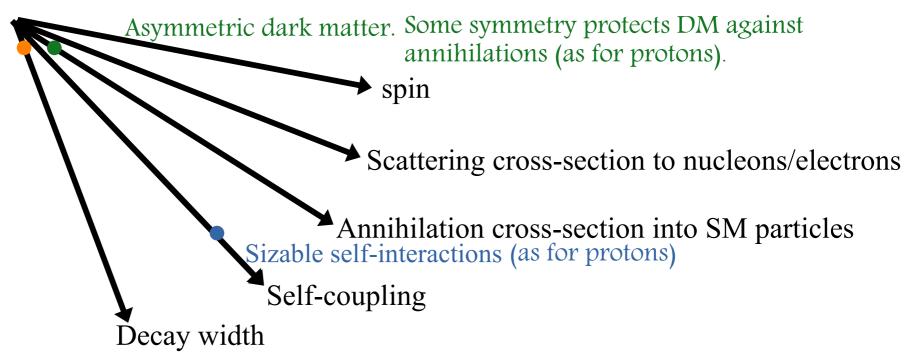








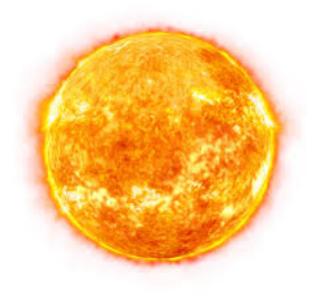




Protons do not annihilate.
Protons have strong self-interactions
Protons form stars



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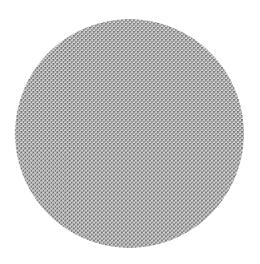




DM does not annihilate.

DM has strong self-interactions

DM form dark stars



Density profile of dark stars calculable from the Klein-Gordon equation in curved spacetime (for bosonic DM) and the Einstein equations:

Colpi et al'86

$$g^{\mu\nu}\nabla_{\mu}\nabla_{\nu}\phi - m^2\phi - \lambda|\phi|^2\phi = 0$$

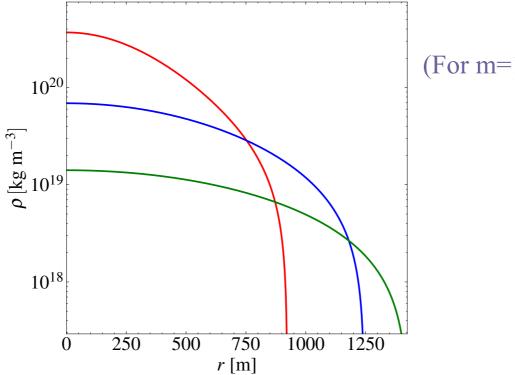
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi G T_{\mu\nu}$$

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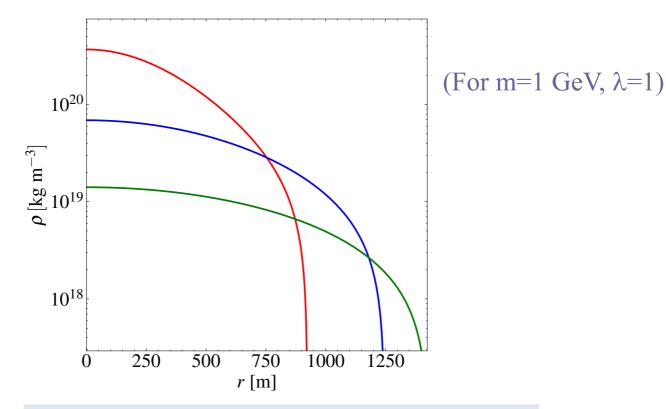
(For m=1 GeV, λ =1)

Density profile of dark stars calculable from the Klein-Gordon equation in curved spacetime (for bosonic DM) and the Einstein equations:

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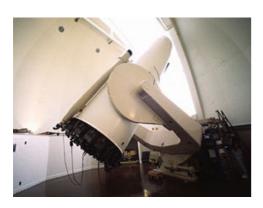
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi G T_{\mu\nu}$$



Dark stars are very compact objects

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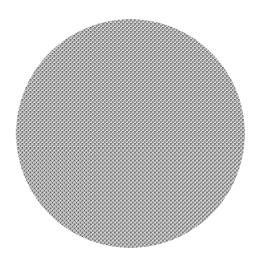


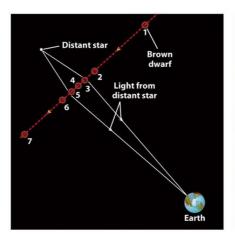


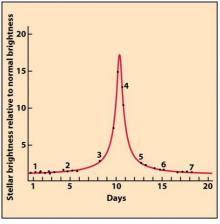
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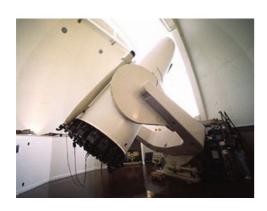






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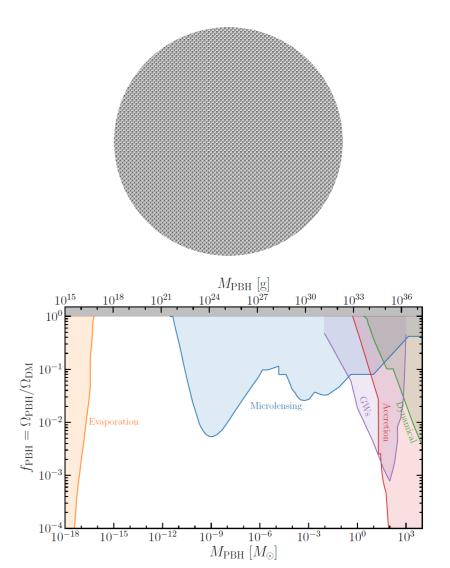




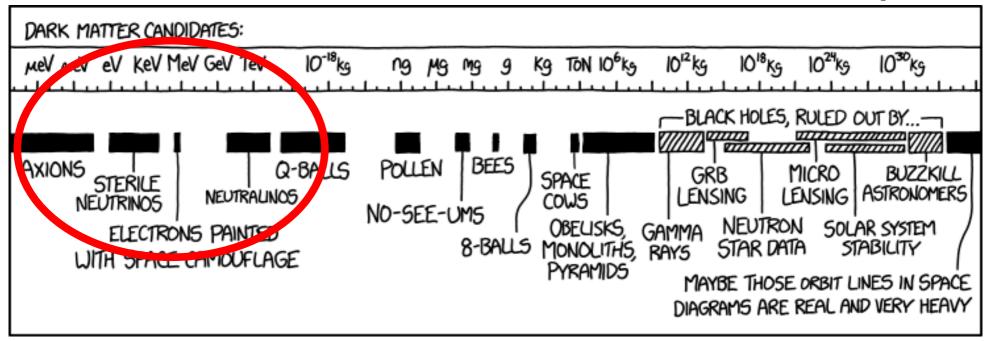
DM does not annihilate.

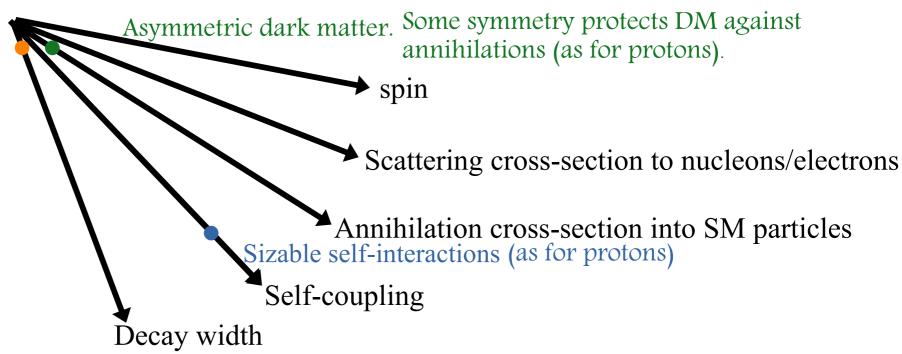
DM has strong self-interactions

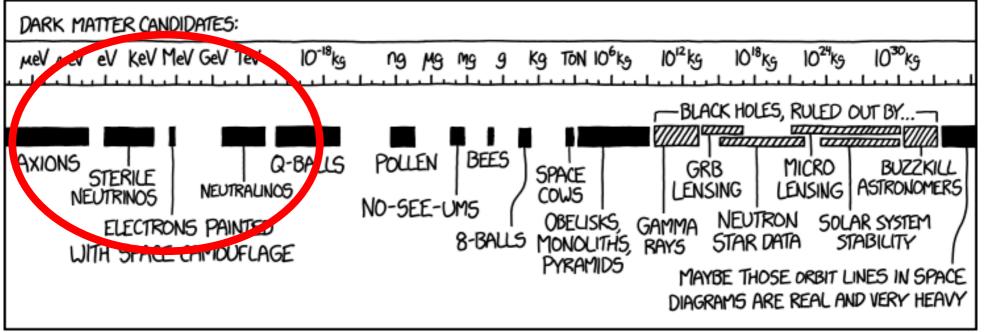
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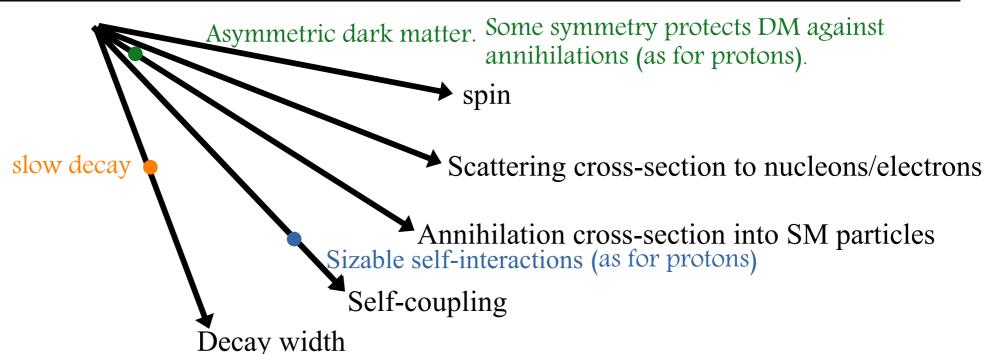


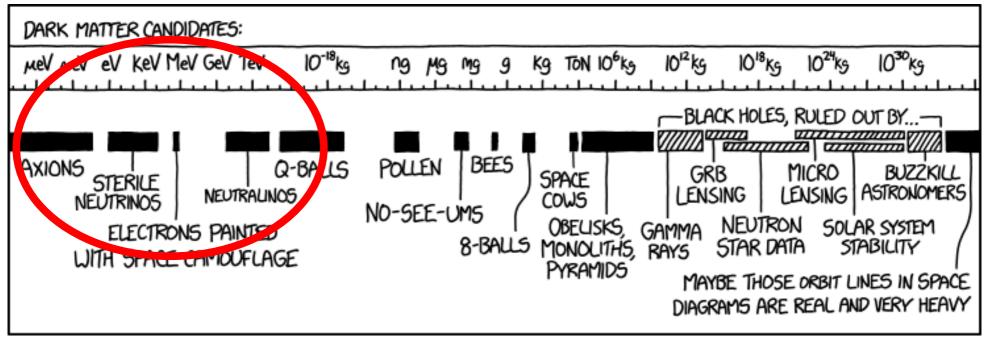
Are there other signals from dark stars?

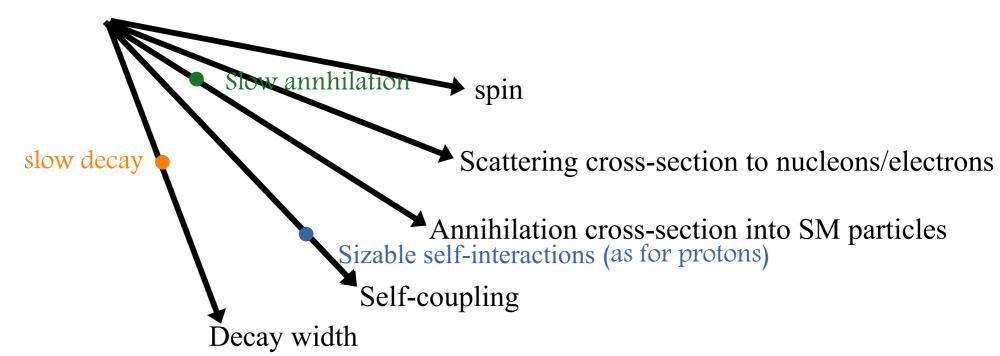


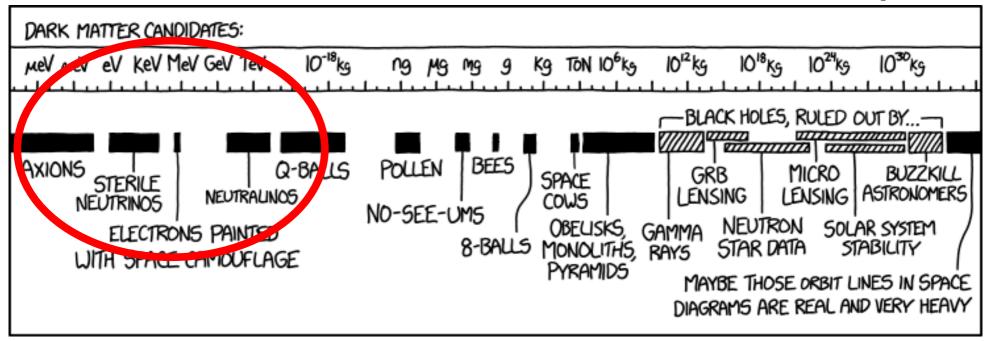


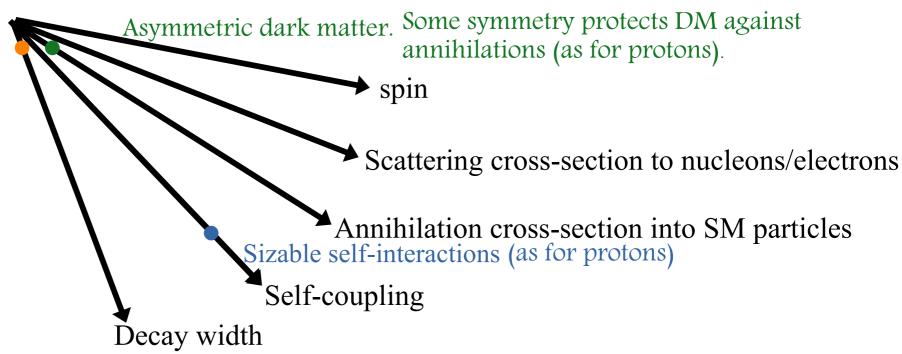


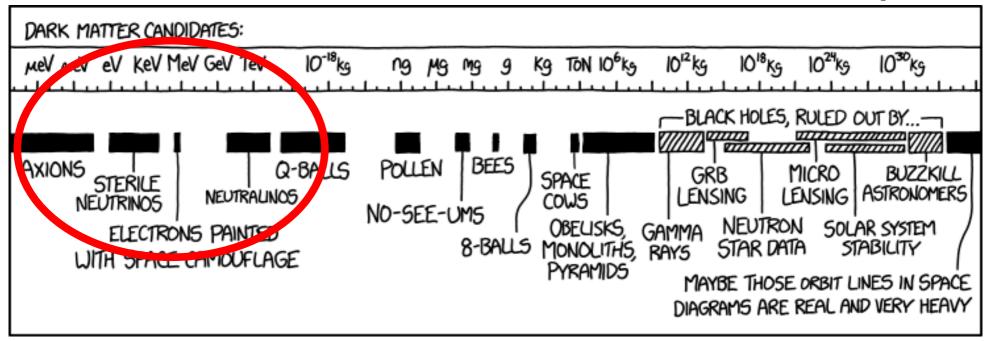


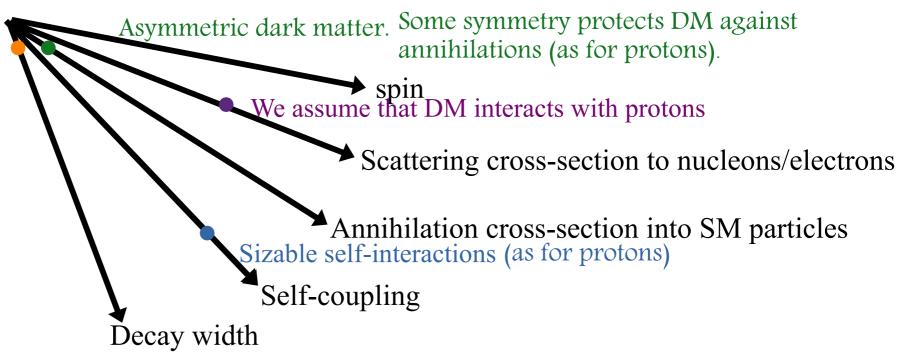


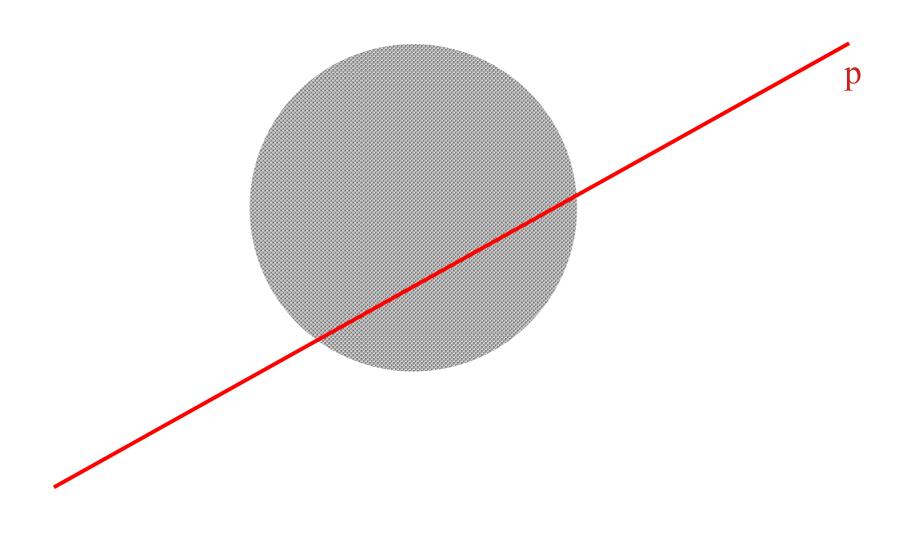


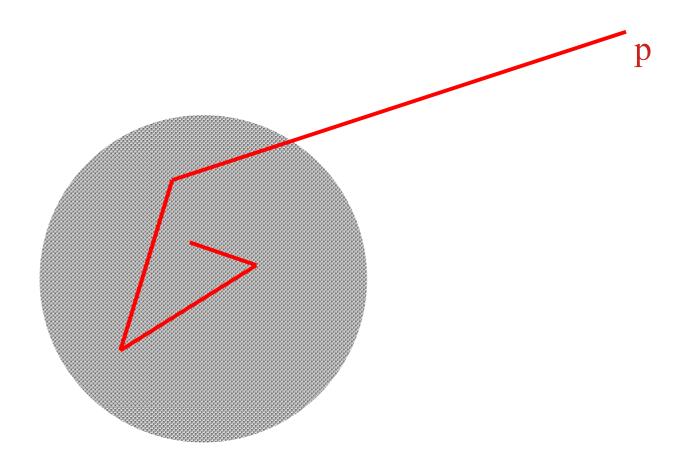


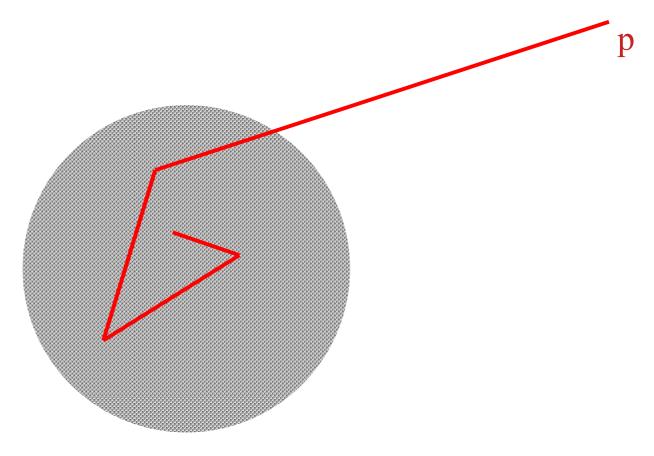




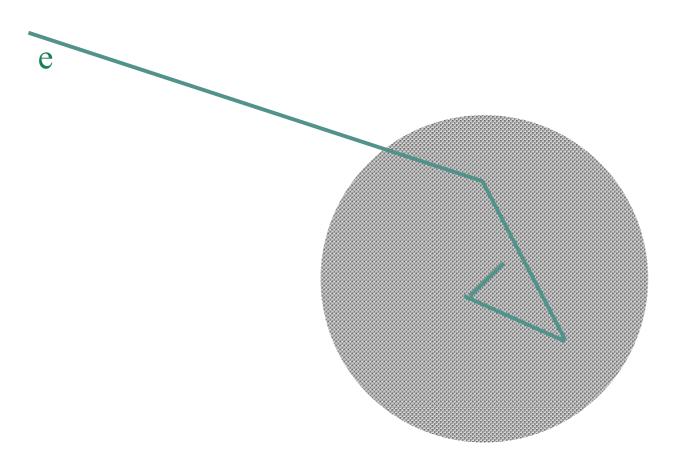


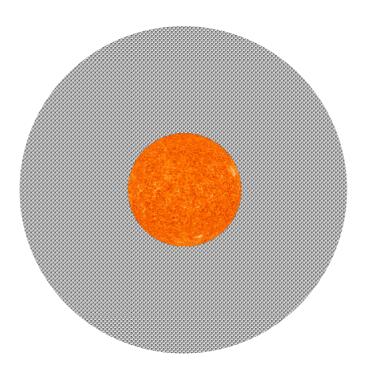


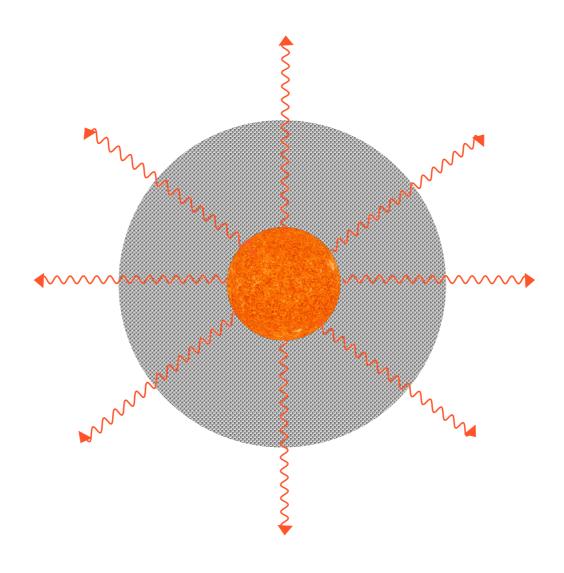




$$C_{v_s} = (1.3 \times 10^{15} \text{ s}^{-1}) \left(\frac{m}{1 \text{ GeV}}\right)^{-1} \left(\frac{R_{DS}}{1 \text{ km}}\right)^3 \left(\frac{\rho_{\text{core}}}{10^{18} \text{ gr/cm}^3}\right)$$
$$\times \left(\frac{\sigma}{10^{-50} \text{ cm}^2}\right) \left(\frac{n_p}{10^{-5} \text{ cm}^{-3}}\right) \left(\frac{v_s}{10^{-3}}\right)^{-1} \text{Erf} \left[\sqrt{\frac{3}{2}} \left(\frac{v_s}{\bar{v}}\right)\right]$$

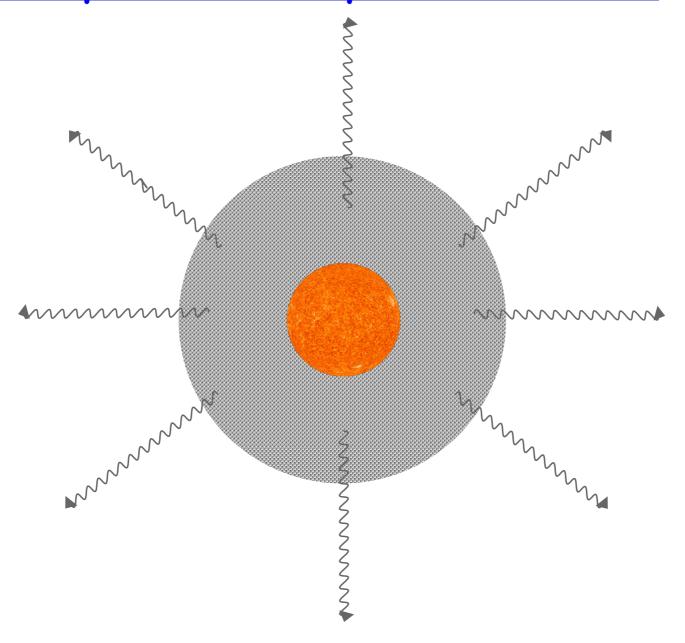




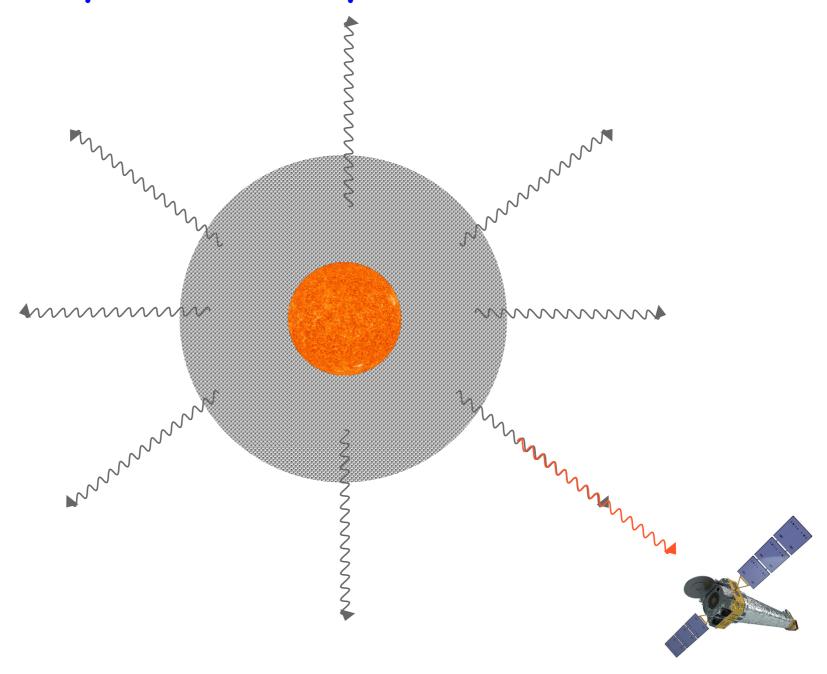




Proton capture in compact dark stars



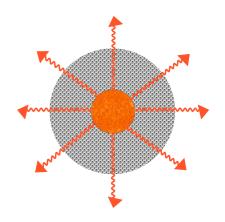
Proton capture in compact dark stars



$$\frac{dT}{dt} = -\frac{L_{\gamma} + L_{\gamma'}}{C_V}$$

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Photon luminosity:
$$L_{\gamma} = \left(4\pi r_{\rm th}^2\right) \int_0^{\infty} I(\nu) d\nu$$



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$$r_{\rm th}\sim T^{1/2} \qquad I(\nu)=\frac{2h}{c^2}\frac{\nu^3}{e^{\frac{h\nu}{k_BT}}-1}\left(1-{\rm e}^{-\tau(\nu)}\right)$$

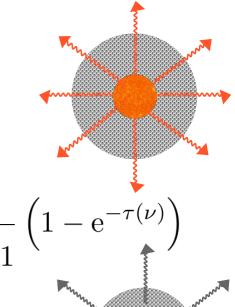
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Dark photon luminosity: $L_{\gamma'} = (4\pi R_{\rm DS}^2) \, \sigma_{\rm SB} T^4 {\rm e}^{-\frac{m_{\gamma'}}{T}}$



$$\frac{dT}{dt} = -\frac{L_{\gamma} + L_{\gamma'}}{C_V}$$

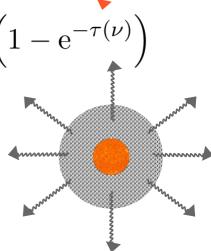
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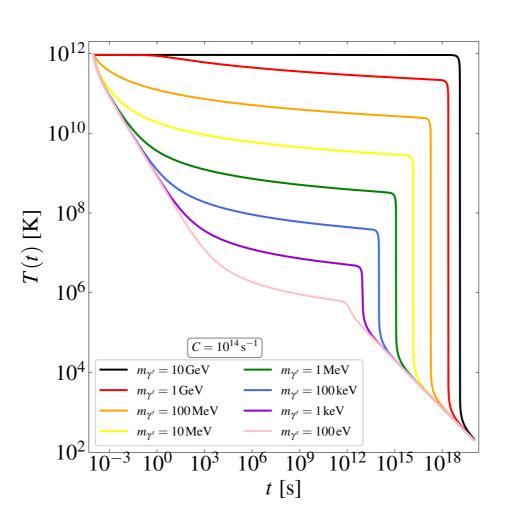
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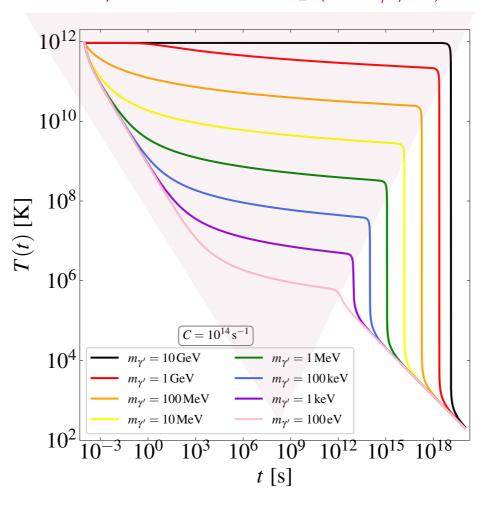


Heat capacity: The DM plausibly forms a Bose-Einstein condensate

$$C_{\rm V} \sim T^{3/2}$$

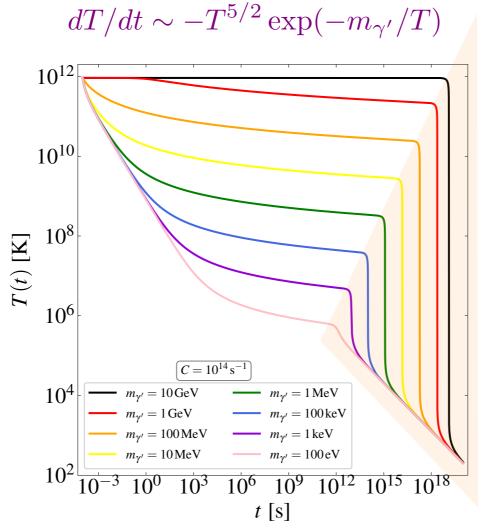


Proton gas optically thin. Cooling by dark photon emission $dT/dt \sim -T^{5/2} \exp(-m_{\gamma'}/T)$



Proton gas optically thin.

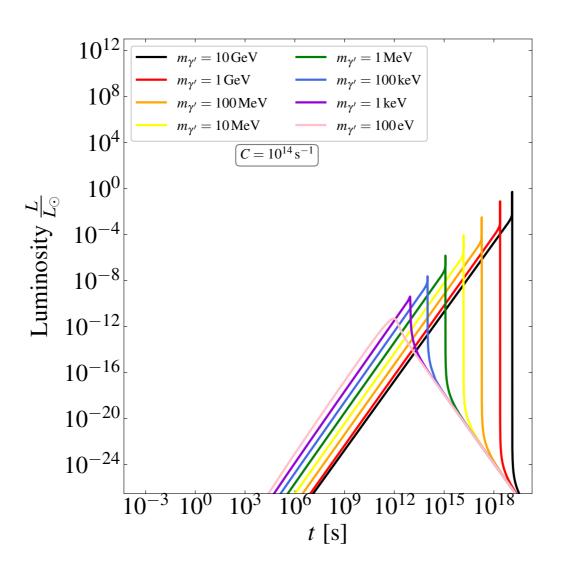
Cooling by dark photon emission



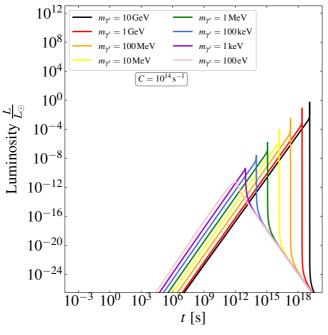
Proton gas optically thick.
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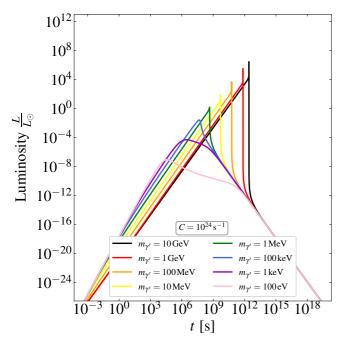
$$dT/dt \sim -T^{7/2}$$

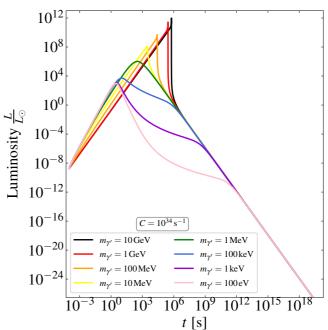
DS luminosity



DS luminosity

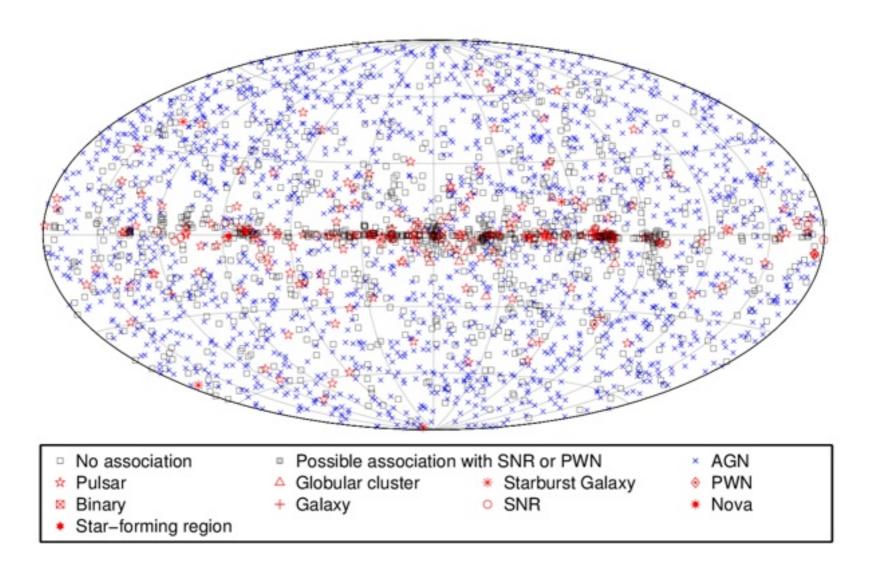




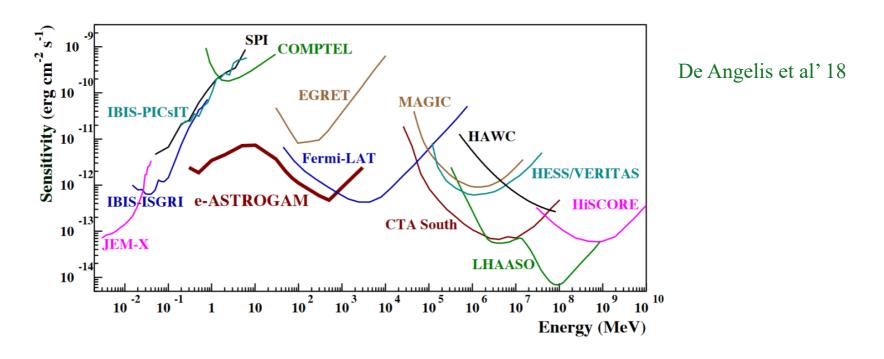


Dark stars could still be shining today. They could be detected as a point source in X-rays or γ -rays, with a black body spectrum (or bremsstrahlung), and with no optical counterpart.

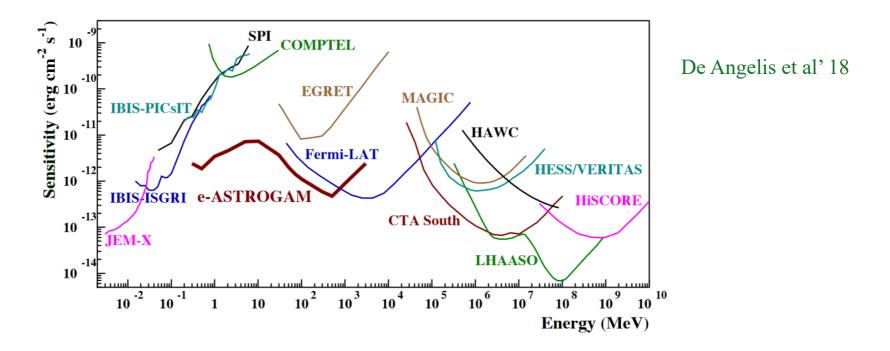
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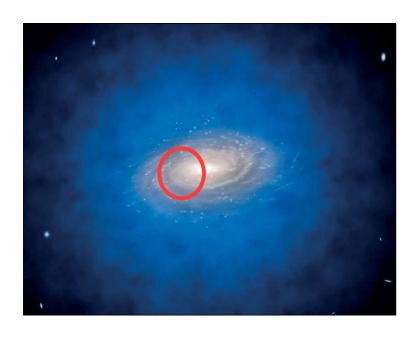
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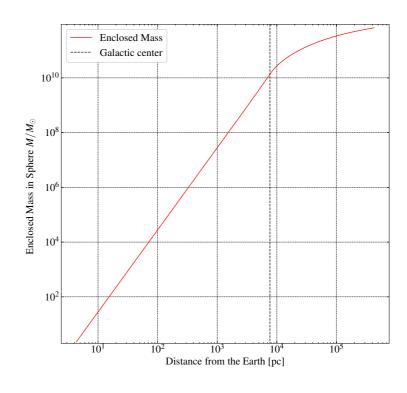


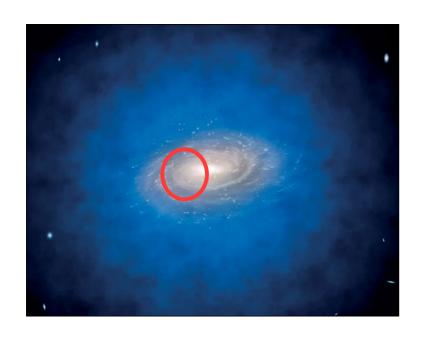
For a luminosity L, a dark star within a distance $d < (L/(4\pi S))^{1/2}$ is at the reach of experiments.

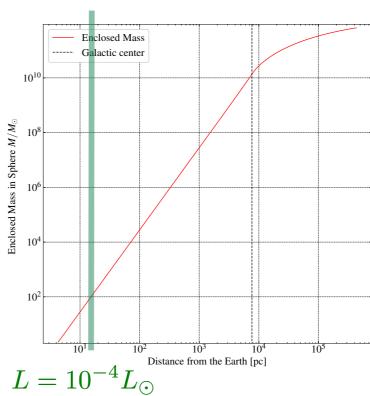
$$d < 1.8 \,\mathrm{kpc} \left(\frac{L}{L_{\odot}}\right)^{1/2} \left(\frac{S}{10^{-11} \,\mathrm{erg} \,\mathrm{cm}^{-2} \,\mathrm{s}^{-1}}\right)^{-1/2}$$





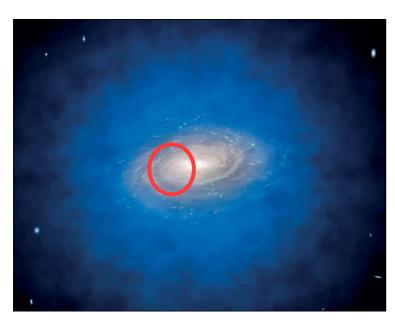






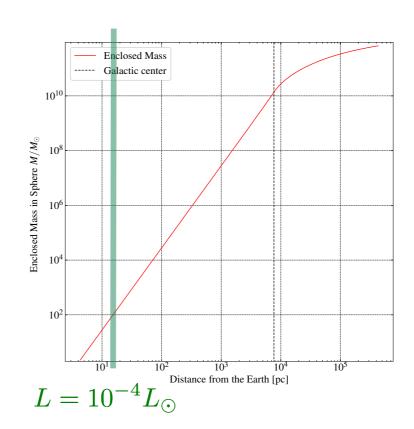
$$L=10^{-4}L_{\odot}$$
 Distance from the East

How many dark stars within a distance *d*?

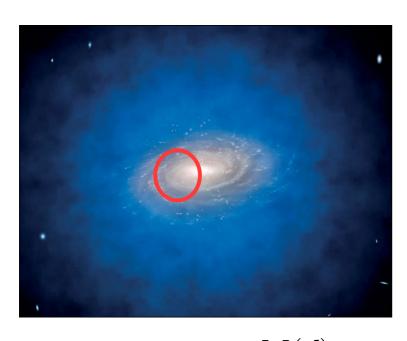


$$N_{\mathrm{DS}} = \mathcal{F}_{\mathrm{DS}} \frac{M(d)}{M_{\mathrm{DS}}}$$

Consider $M_{\rm DS}=M_{\odot},\,\mathcal{F}_{\rm DS}=10^{-2}$

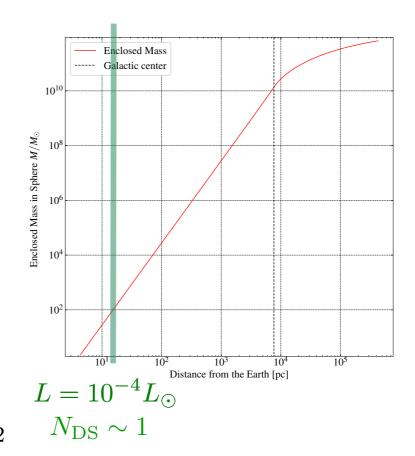


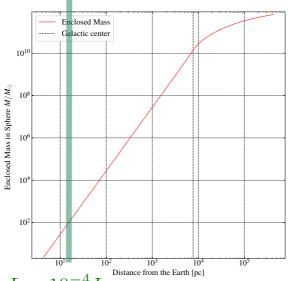
How many dark stars within a distance *d*?



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m DS}}$$

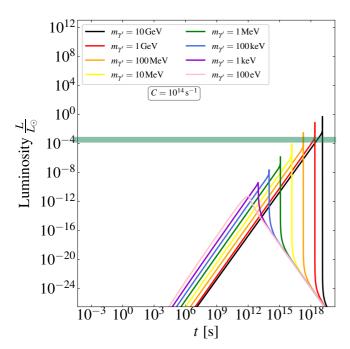
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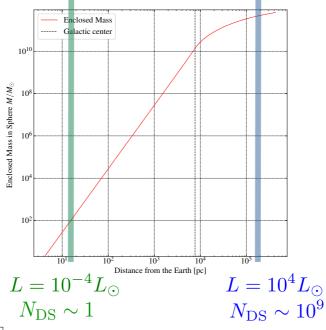


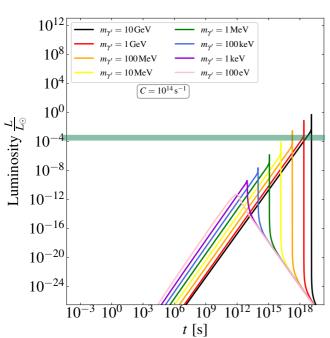
 $L = 10^{-4} L_{\odot}^{^{\text{D}}}$

 $N_{\rm DS} \sim 1$



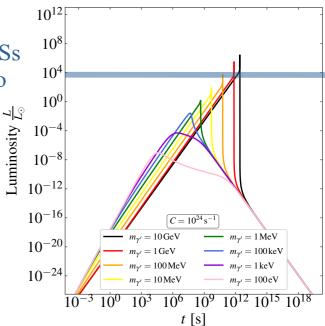
 \sim 1 event if *all* DSs formed \sim 10¹⁷ s ago

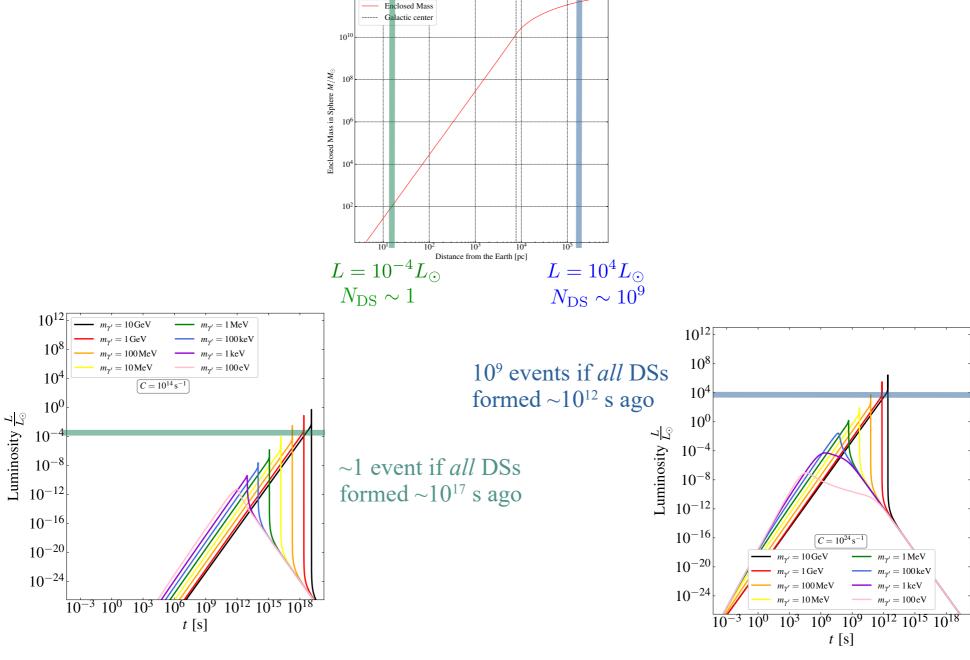




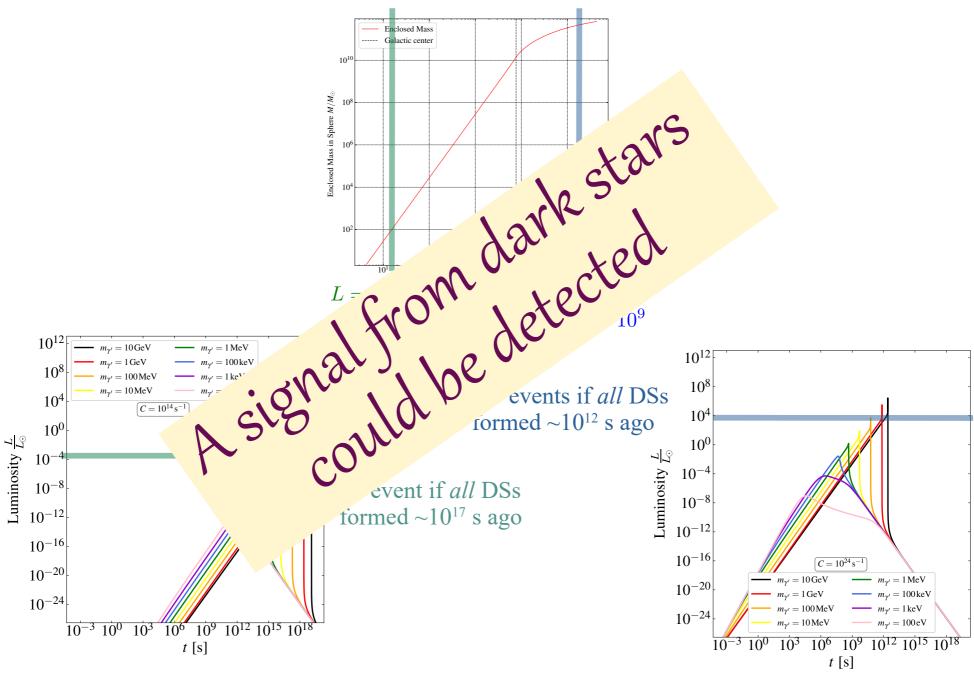
 10^9 events if *all* DSs formed $\sim 10^{12}$ s ago

 \sim 1 event if *all* DSs formed \sim 10¹⁷ s ago





Dark stars could form continuously over time.



Dark stars could form continuously over time.

Conclusions

- If the dark matter particle has strong self-interactions, it could form dark stars, that could be detected in MACHO searches.
- If the dark matter particle interacts with the proton, dark stars could capture protons from the interstellar medium. Electrons are also captured to keep the dark star electrically neutral. (Similar rationale if the dark matter interacts with the electron.)
- The captured electrons and protons form a hot gas that emits radiation with a characteristic spectrum.
- New target for indirect detection of asymmetric dark matter: point sources in X- or γ -rays (from scatterings with protons/electrons, from slow decays or from slow annihilations). These sources would be also detected as MACHOs.