

Searching for Ultralight Scalar Dark Matter with Muonium and Muonic Atoms

Yevgeny Stadnik

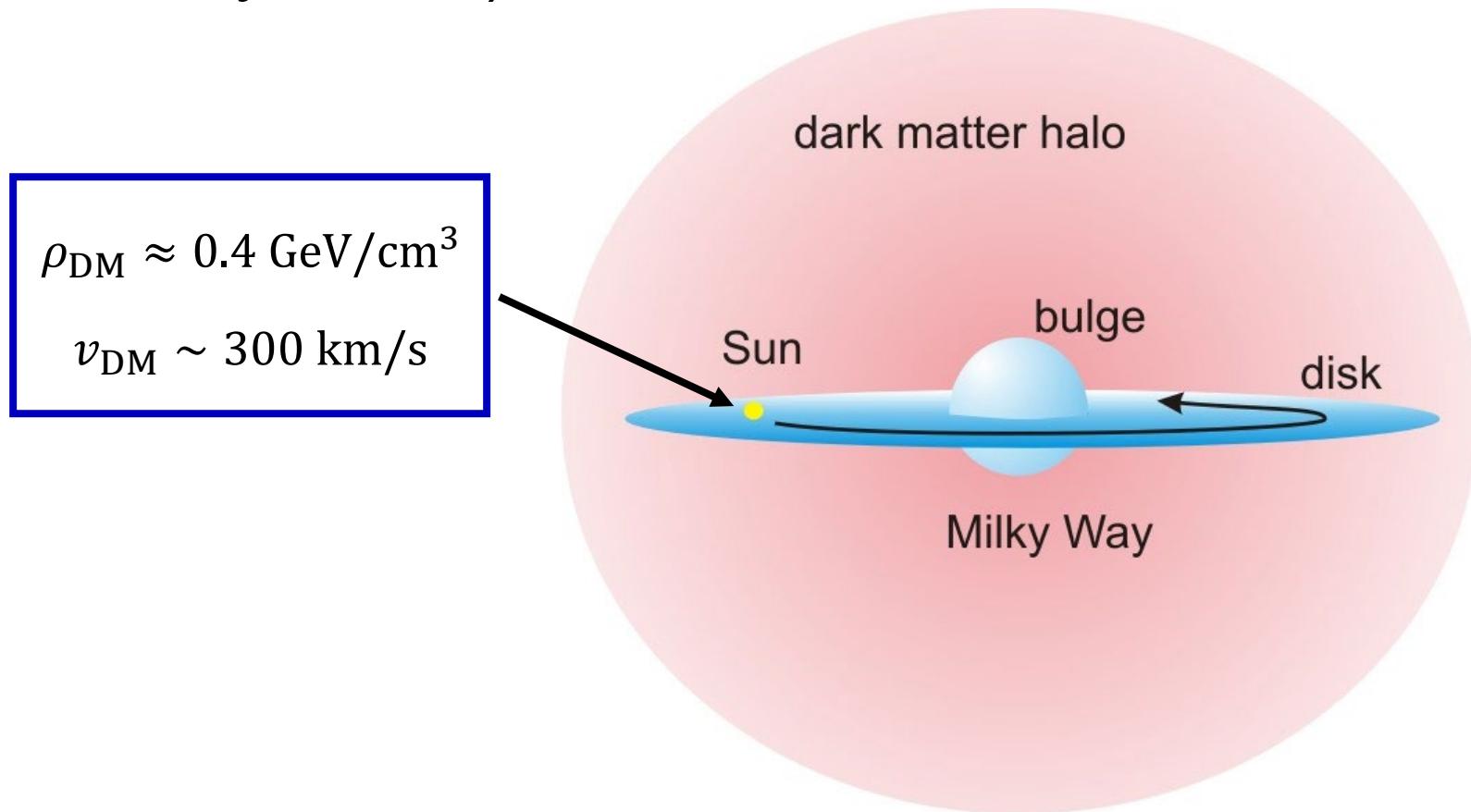
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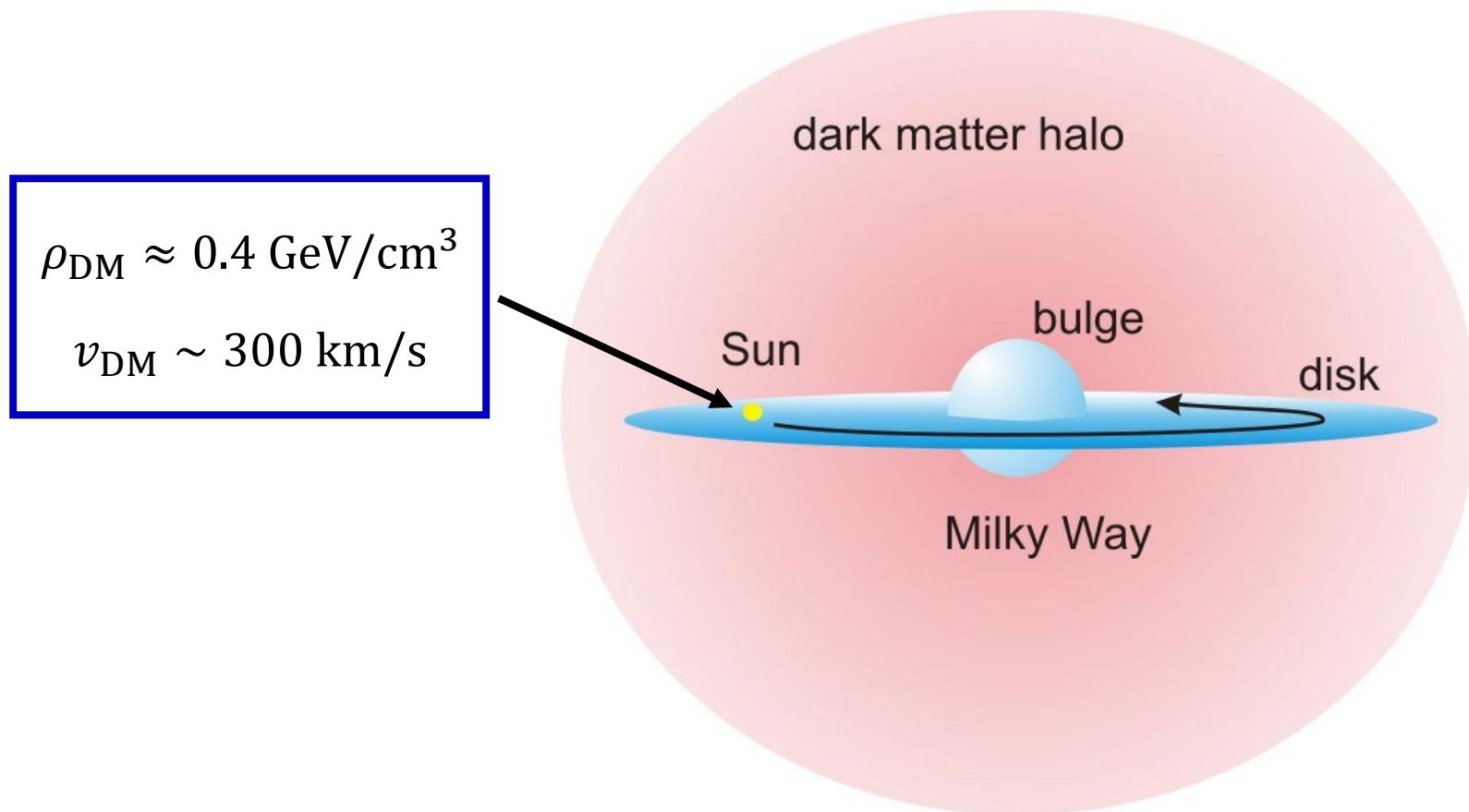
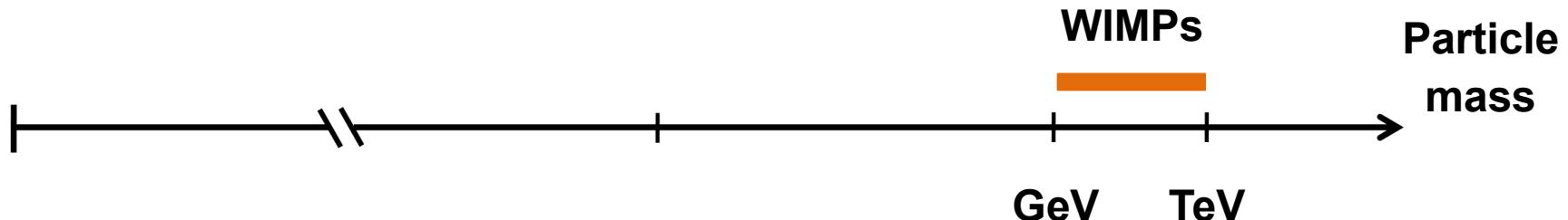
**The Dark Side of the Universe (DSU2022), University of New South Wales,
Sydney, Australia, 5th – 9th December 2022**

Dark Matter

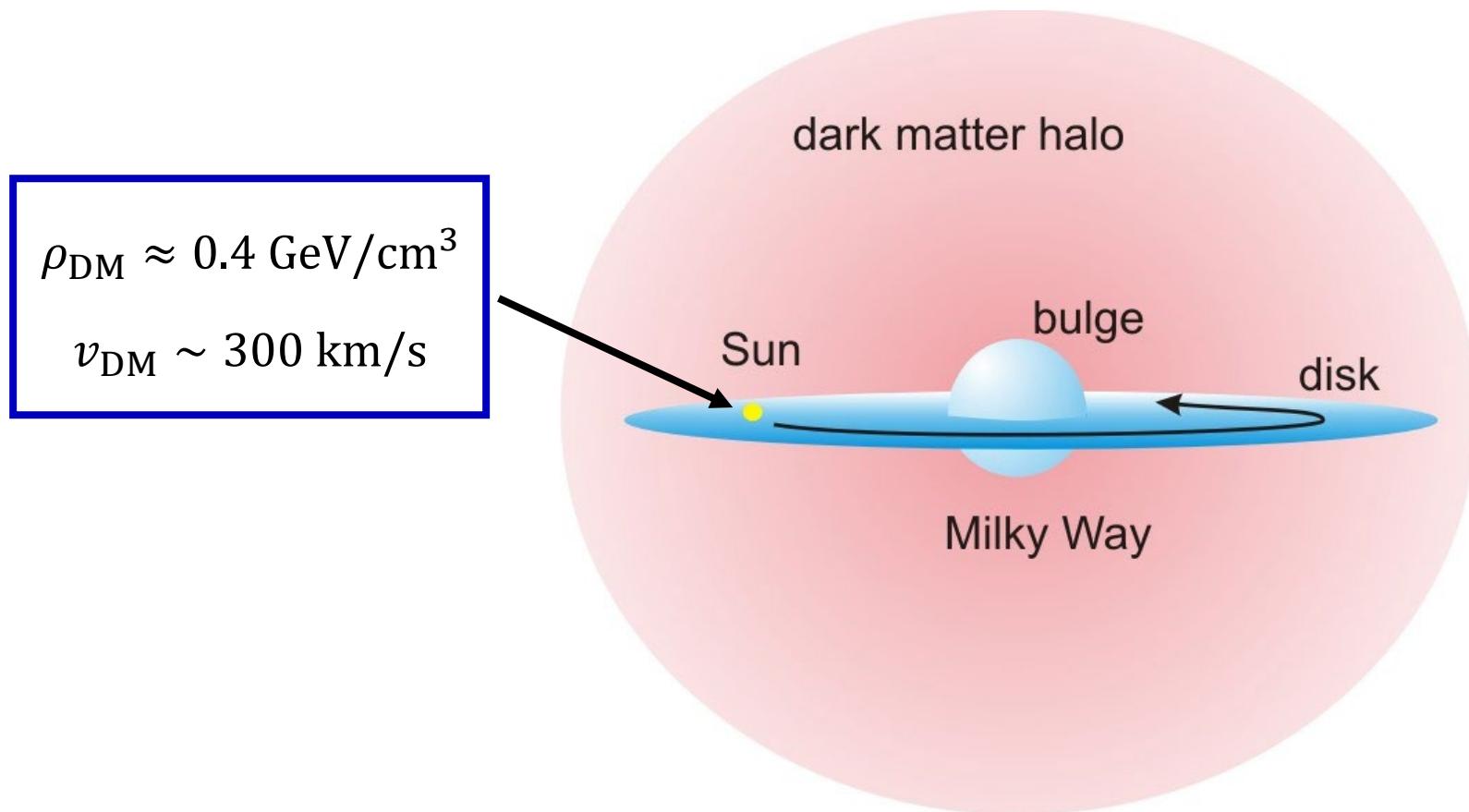
Strong astrophysical evidence for existence of **dark matter** (~5 times more dark matter than ordinary matter)



Dark Matter

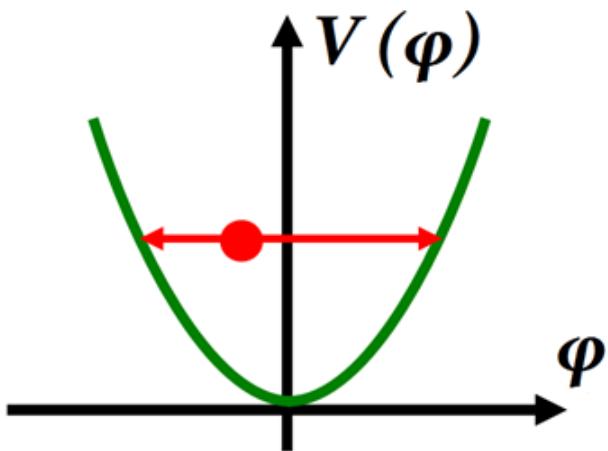


Dark Matter



Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)



$$V(\varphi) = \frac{m_\varphi^2 \varphi^2}{2}$$

$$\ddot{\varphi} + m_\varphi^2 \varphi \approx 0$$

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- $10^{-21} \text{ eV} \lesssim m_\varphi \lesssim 1 \text{ eV} \Leftrightarrow 10^{-7} \text{ Hz} \lesssim f_{\text{DM}} \lesssim 10^{14} \text{ Hz}$
 $T_{\text{osc}} \sim 1 \text{ month}$ **IR frequencies**

Lyman- α forest measurements [suppression of structures for $L \lesssim \mathcal{O}(\lambda_{\text{dB},\varphi})$]

[Related figure-of-merit: $\lambda_{\text{dB},\varphi} / 2\pi \leq L_{\text{dwarf galaxy}} \sim 100 \text{ pc} \Rightarrow m_\varphi \gtrsim 10^{-21} \text{ eV}$]

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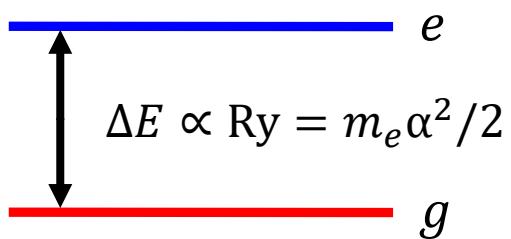
- *Wave-like signatures* [cf. *particle-like* signatures of WIMP DM]

Lyman- α forest measurements [suppression of structures for $L \lesssim \mathcal{O}(\lambda_{\text{dB},\varphi})$]

Probes of Ultralight Scalar DM

- Recent searches for ultralight scalar DM have focused on the electromagnetic (photon) and electron couplings

Atomic clocks



Optical cavities



$$L \propto a_B = 1/(m_e \alpha)$$

see, e.g., [[Stadnik, Flambaum, PRL 114, 161301 \(2015\)](#); [PRL 115, 201301 \(2015\)](#)] for details

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 - Proton radius puzzle
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- Extra motivation from persistence of various anomalies in muon physics, such as:
 - Proton radius puzzle
 - $(g - 2)_\mu$ puzzle
- No stable terrestrial sources of muons (unlike electrons)

Probing Oscillations of m_μ with Muonium Spectroscopy

[Stadnik, arXiv:2206.10808]

$$\mathcal{L}_{\text{lin}} = -\frac{\varphi}{\Lambda_\mu} m_\mu \bar{\mu} \mu \approx -\frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\mu} m_\mu \bar{\mu} \mu \Rightarrow \frac{\delta m_\mu}{m_\mu} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\mu}$$

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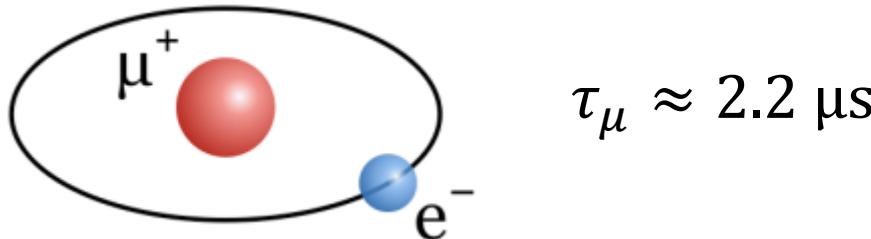
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Muonium = $e^- \mu^+$ bound state, $m_r = \frac{m_e m_\mu}{m_e + m_\mu} \approx m_e (1 - m_e/m_\mu)$



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Muonium = $e^- \mu^+$ bound state, $\textcolor{blue}{m_r} = \frac{m_e m_\mu}{m_e + m_\mu} \approx m_e (1 - \textcolor{red}{m_e/m_\mu})$

$$E_n^{\text{Rydberg}} = -\frac{\textcolor{blue}{m_r} \alpha^2}{2n^2} \Rightarrow \frac{\Delta v_{1S-2S}}{v_{1S-2S}} \approx 2 \frac{\Delta \alpha}{\alpha} + \frac{\Delta m_e}{m_e} + \frac{\textcolor{red}{m_e}}{m_\mu} \frac{\Delta m_\mu}{m_\mu}$$

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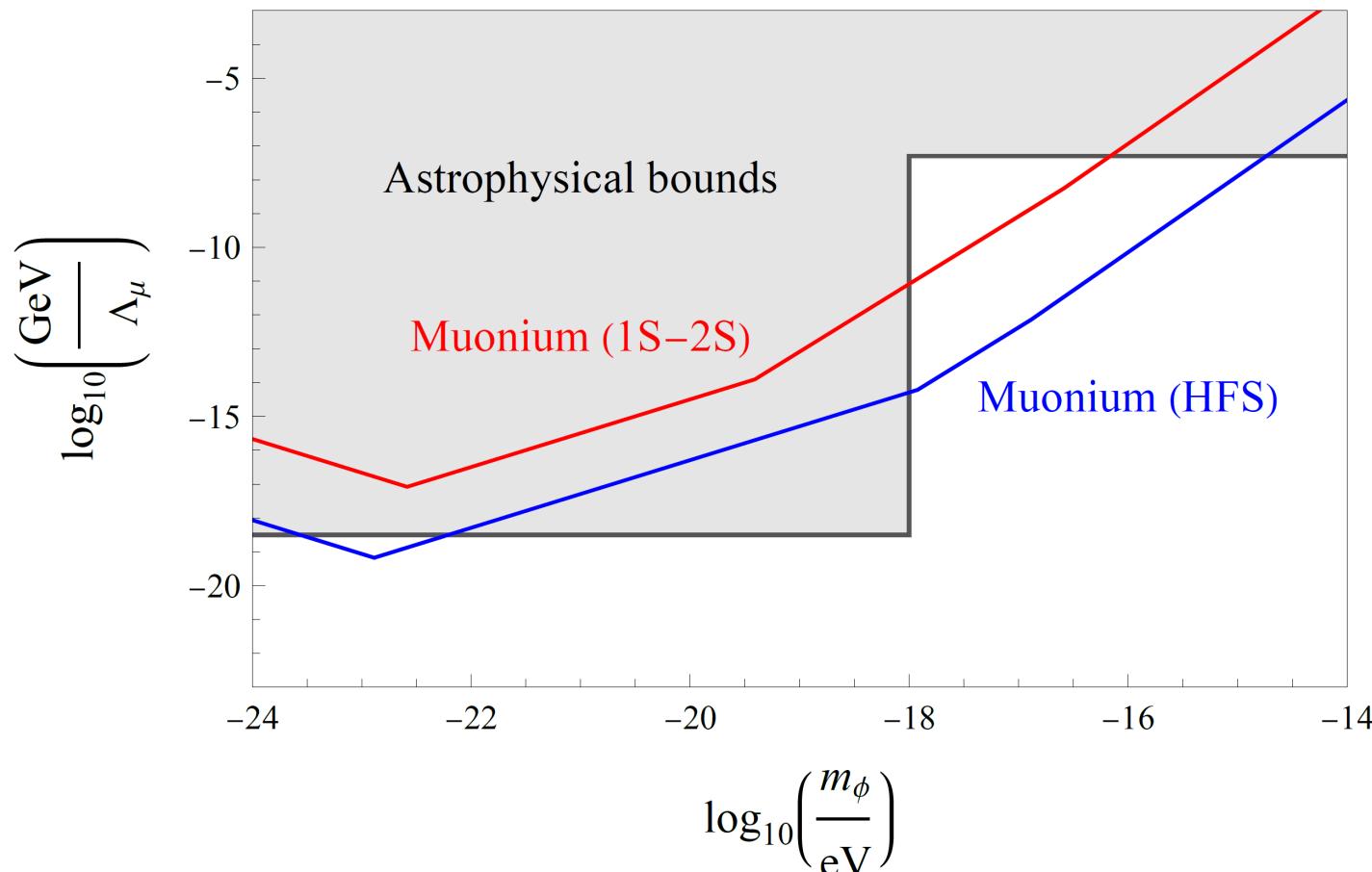
$$E_n^{\text{Rydberg}} = -\frac{m_r \alpha^2}{2n^2} \Rightarrow \frac{\Delta v_{1S-2S}}{v_{1S-2S}} \approx 2 \frac{\Delta \alpha}{\alpha} + \frac{\Delta m_e}{m_e} + \frac{m_e}{m_\mu} \frac{\Delta m_\mu}{m_\mu}$$

$$\Delta E_{\text{Fermi}} = \frac{8m_r^3 \alpha^4}{3m_e m_\mu} \Rightarrow \frac{\Delta v_{\text{HFS}}}{v_{\text{HFS}}} \approx 4 \frac{\Delta \alpha}{\alpha} + 2 \frac{\Delta m_e}{m_e} - \frac{\Delta m_\mu}{m_\mu}$$

Estimated Sensitivities to Scalar Dark Matter with $\varphi \bar{\mu} \mu / \Lambda_\mu$ Coupling

[Stadnik, arXiv:2206.10808]

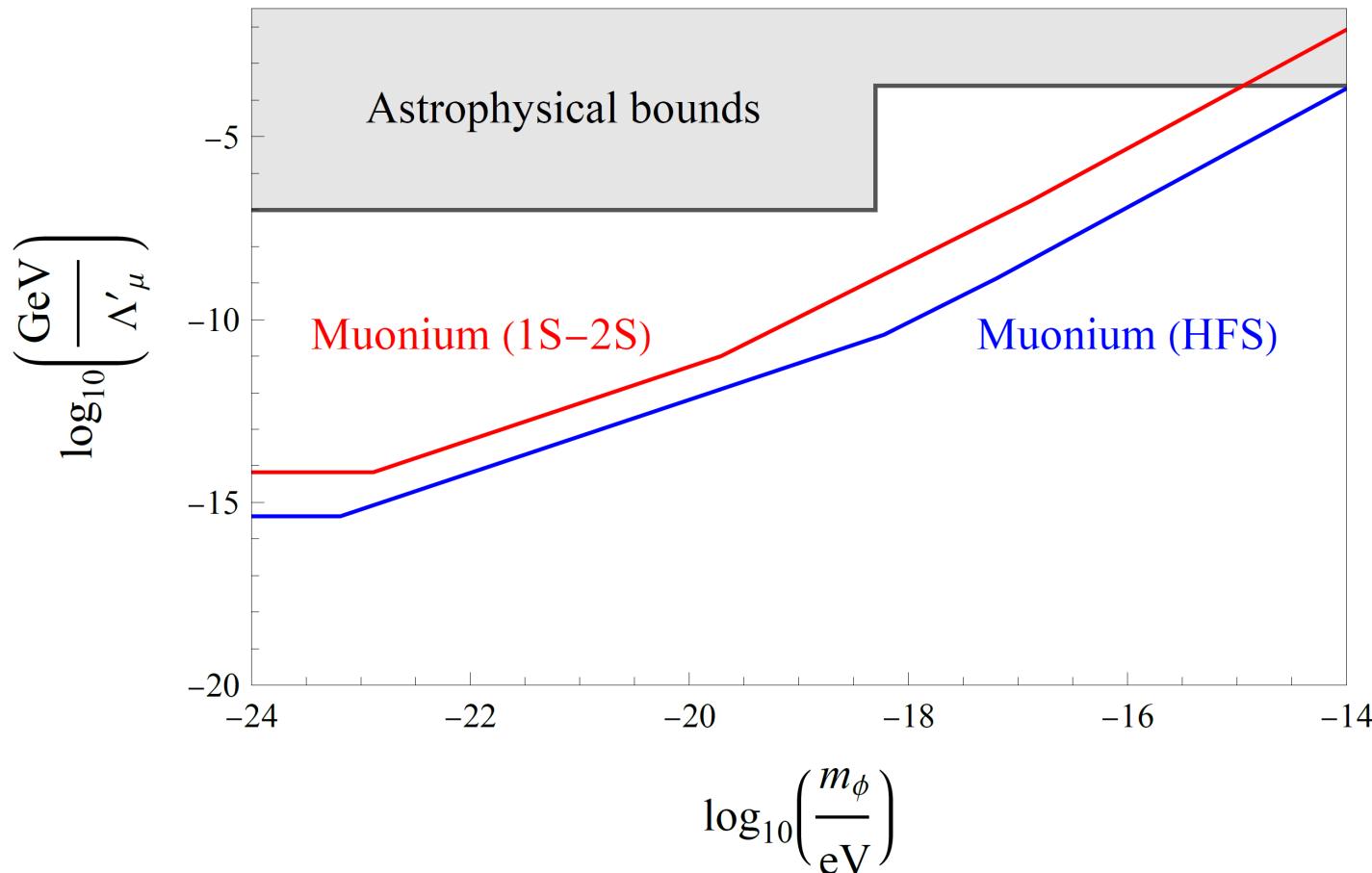
Up to 7 orders of magnitude improvement possible with existing datasets!



Estimated Sensitivities to Scalar Dark Matter with $\varphi^2 \bar{\mu} \mu / (\Lambda'_\mu)^2$ Coupling

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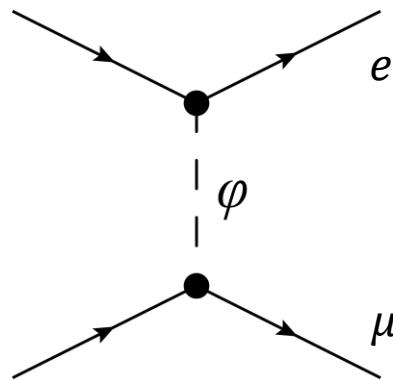
Up to 8 orders of magnitude improvement possible with existing datasets!



Probing Scalar-Muon Coupling with Muonium Free-fall

[Stadnik, arXiv:2206.10808]

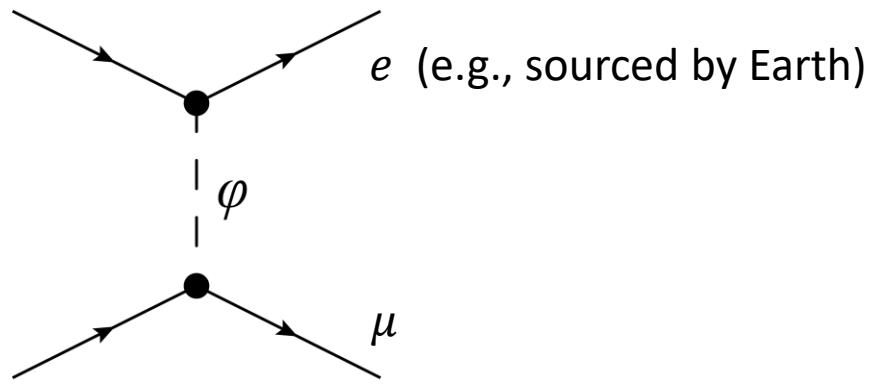
$$\mathcal{L}_{\text{lin}} = -\frac{\varphi}{\Lambda_e} m_e \bar{e} e - \frac{\varphi}{\Lambda_\mu} m_\mu \bar{\mu} \mu \Rightarrow V_{e\mu}(r) = -\frac{m_e m_\mu e^{-m_\varphi r}}{4\pi \Lambda_e \Lambda_\mu r}$$



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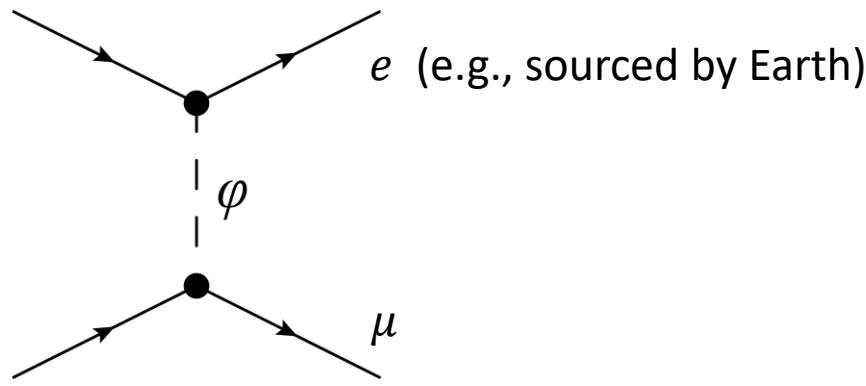


Local value of g measured in free-fall experiments using muonium would differ from experiments using non-muon-based test masses

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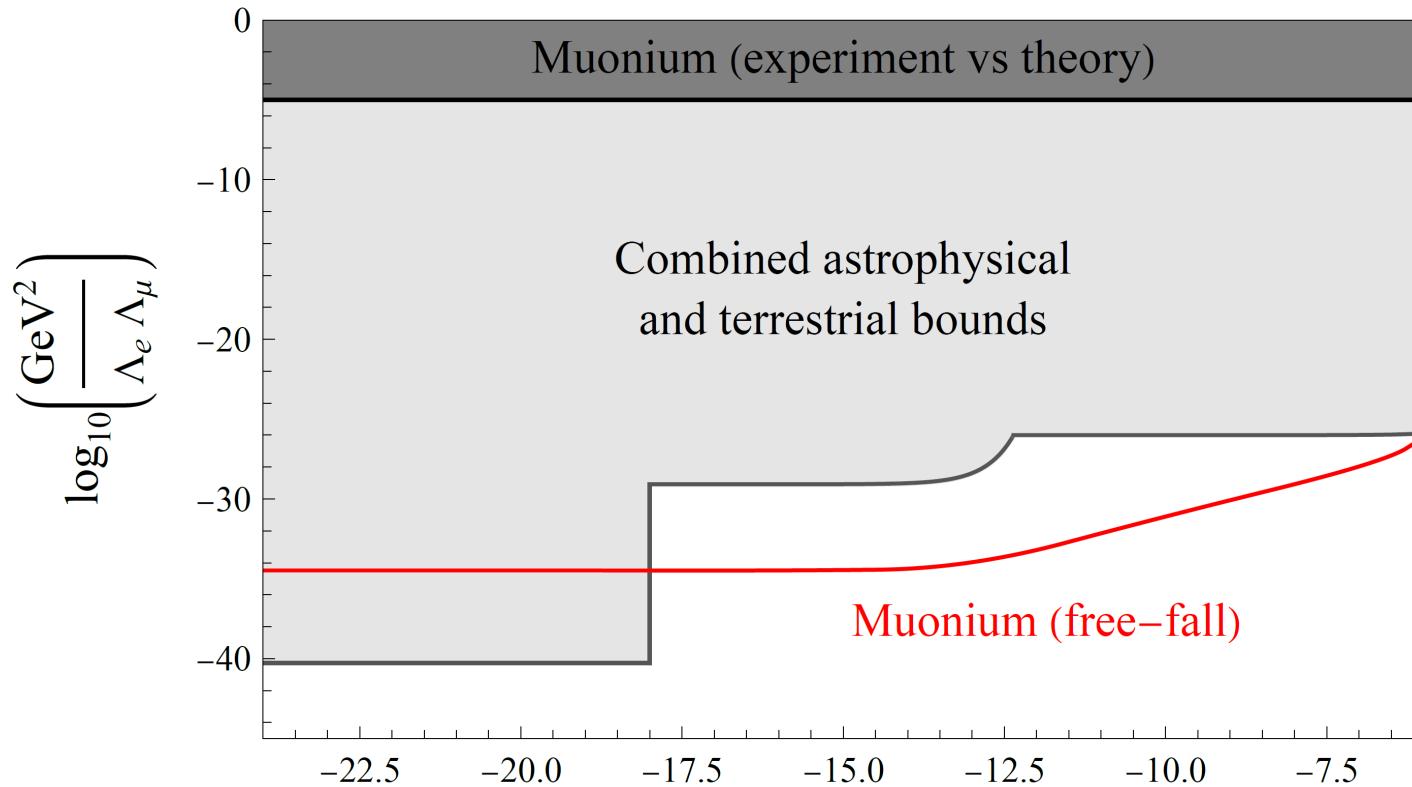
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Recently started LEMING experiment at the Paul Scherrer Institute aims to measure g with a precision of $\Delta g/g \sim 0.1$ using muonium

Probing Scalar-Muon Coupling with Muonium Free-fall

[Stadnik, arXiv:2206.10808]

Up to 5 orders of magnitude improvement possible with ongoing measurements!



Assume $\Lambda_\mu \ll \Lambda_e \ll \Lambda_{\text{other SM fields}}$

$$\log_{10} \left(\frac{m_\phi}{\text{eV}} \right)$$

Summary

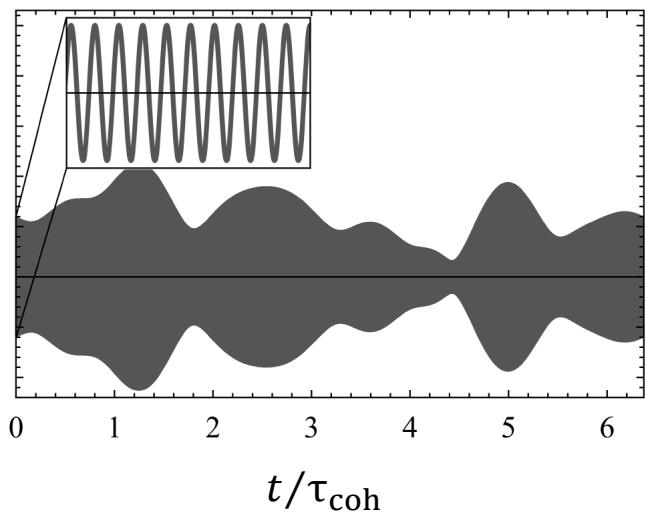
- Recent searches for ultralight scalar DM have focused on the electromagnetic (photon) and electron couplings
- Muonium spectroscopy offers a powerful probe of ultralight scalar dark matter via interactions with muons leading to apparent oscillations of muon mass
 - With existing datasets, up to $\sim 10^7$ improvement possible for $\varphi \bar{\mu} \mu$ coupling (up to $\sim 10^8$ for the $\varphi^2 \bar{\mu} \mu$ coupling over an even broader range of scalar DM masses)
- Ongoing muonium free-fall experiments to measure g offer up to $\sim 10^5$ improvement in sensitivity for the combination of $\varphi \bar{\mu} \mu$ and $\varphi \bar{e} e$ couplings by searching for φ -mediated forces

Back-Up Slides

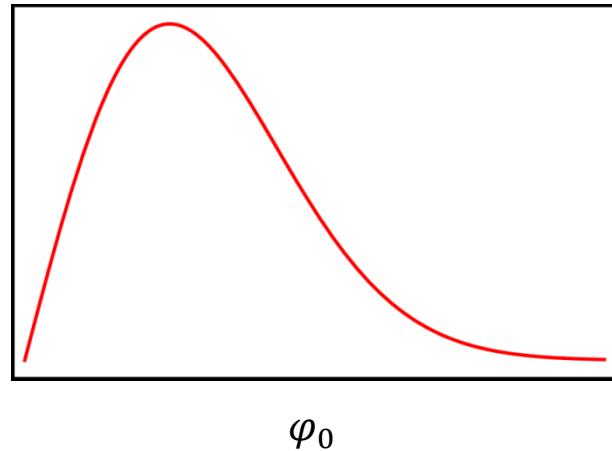
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Evolution of φ_0 with time



Probability distribution function of φ_0
(e.g., Rayleigh distribution)



Dark-Matter-Induced Variations of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],
 [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\begin{aligned} \mathcal{L}_\gamma &= \frac{\varphi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \Rightarrow \frac{\delta\alpha}{\alpha} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \\ \mathcal{L}_f &= -\frac{\varphi}{\Lambda_f} m_f \bar{f} f \approx -\frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} m_f \bar{f} f \Rightarrow \frac{\delta m_f}{m_f} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} \end{aligned}$$

$$\varphi = \varphi_0 \cos(m_\varphi t - \mathbf{p}_\varphi \cdot \mathbf{x}) \Rightarrow \mathbf{F} \propto \mathbf{p}_\varphi \sin(m_\varphi t)$$

$$\left. \begin{aligned} \mathcal{L}'_\gamma &= \frac{\varphi^2}{(\Lambda'_\gamma)^2} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \\ \mathcal{L}'_f &= -\frac{\varphi^2}{(\Lambda'_f)^2} m_f \bar{f} f \end{aligned} \right\} \Rightarrow \left\{ \begin{array}{l} \frac{\delta\alpha}{\alpha} \propto \frac{\delta m_f}{m_f} \propto \Delta\rho_\varphi \\ \mathbf{F} \propto \nabla\rho_\varphi \end{array} \right.$$

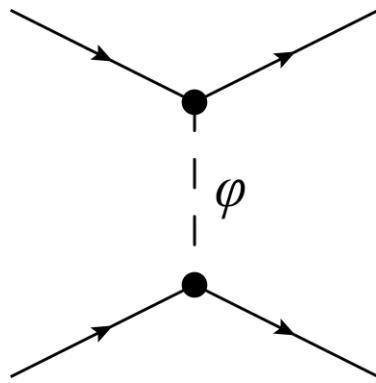
Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

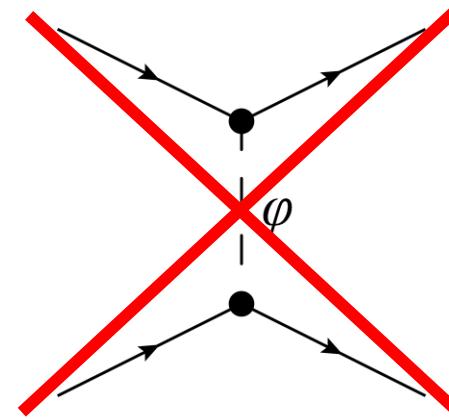
Linear couplings ($\varphi \bar{X}X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa \rho \quad \text{Source term}$$



Quadratic couplings ($\varphi^2 \bar{X}X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa' \rho \varphi \quad \text{Potential term}$$



$$\varphi = \varphi_0 \cos(m_\varphi t) \pm A \frac{e^{-m_\varphi r}}{r}$$

Profile outside of a spherical body

$$m_{\text{eff}}^2(\rho) = m_\varphi^2 \mp \kappa' \rho$$

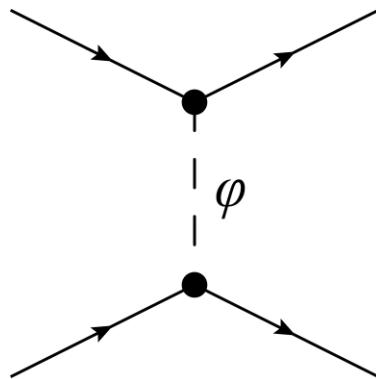
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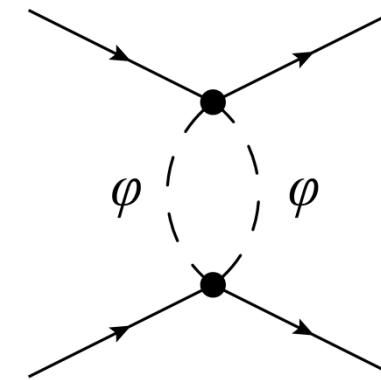
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Motional gradients: $\varphi_0 \cos(m_\varphi t - \mathbf{p}_\varphi \cdot \mathbf{x})$

“Fifth-force” experiments: torsion pendula, atom interferometry

$$\varphi = \frac{\varphi_0 \cos(m_\varphi t)}{\left(1 \pm \frac{B}{r}\right)} - \hbar C \frac{e^{-2m_\varphi r}}{r^3}$$

Gradients + amplification/screening