

# Advancing globular cluster constraints on the axion-photon coupling

DSU2022 - UNSW

JCAP 10 (2022) 096, <u>arXiv:2207.03102</u> Matthew Dolan, **Frederick Hiskens**, Raymond Volkas





#### Part One





#### Part Two





#### Part Three











Can naturally be light and weakly interacting (invisible axions)

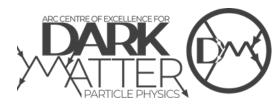






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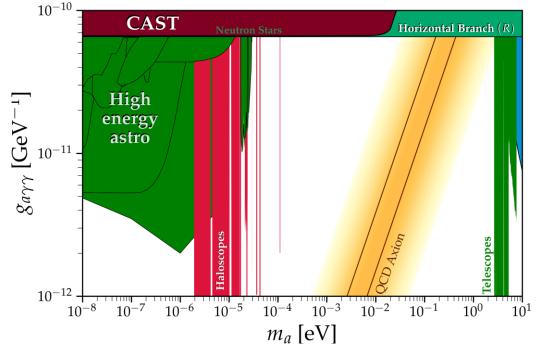


Advancing globular cluster constraints on the axion-photon coupling

cajohare/AxionLimits

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# Axions



Axions are pseudo Nambu-Goldstone bosons predicted in the Peccei-Quinn solution to the strong CP problem

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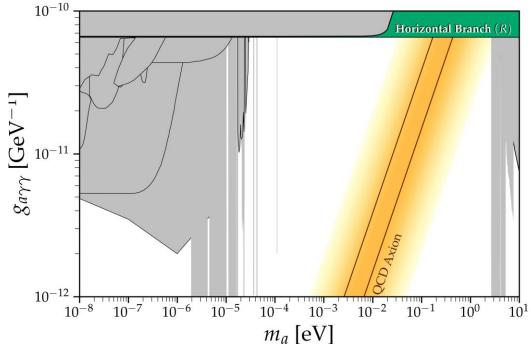
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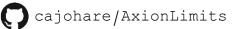
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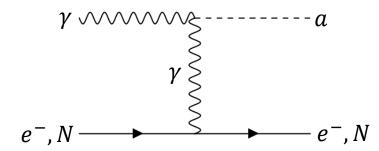








Axions could be produced in deep stellar interiors, e.g. Primakoff production ( $m_a \lesssim 1 \text{ keV}$ )

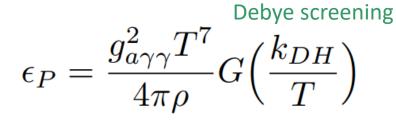




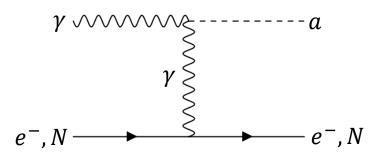


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If sufficiently light and weakly interacting, they can freely escape the local stellar region – new source of energy-loss



Raffelt G., *Phys. Rev. D* **33** (1986) 897

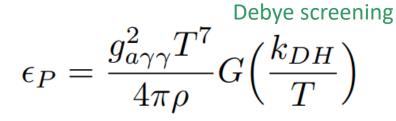






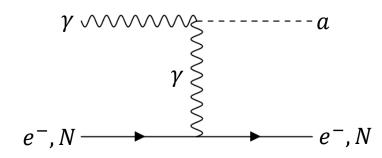
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Burning region loses energy

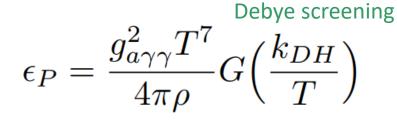




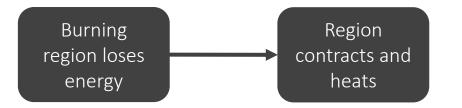


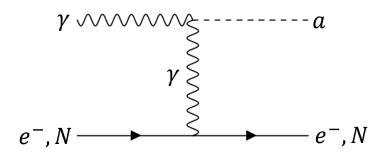
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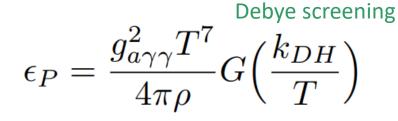


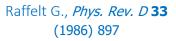


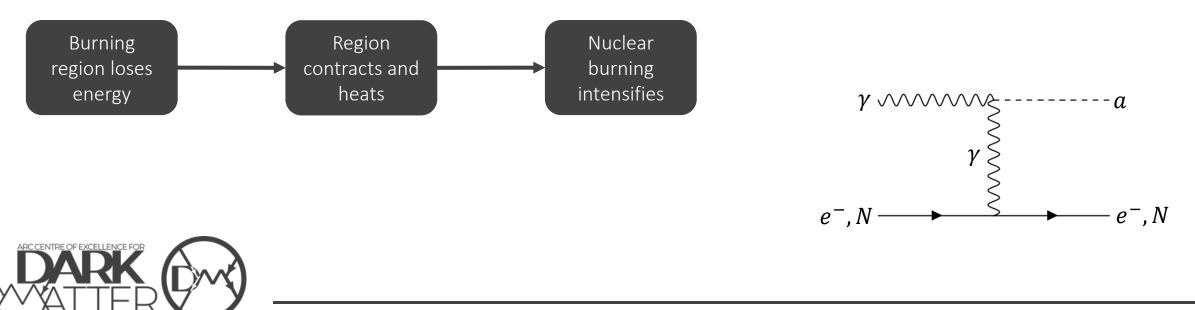


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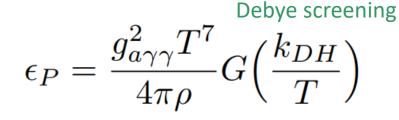


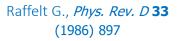


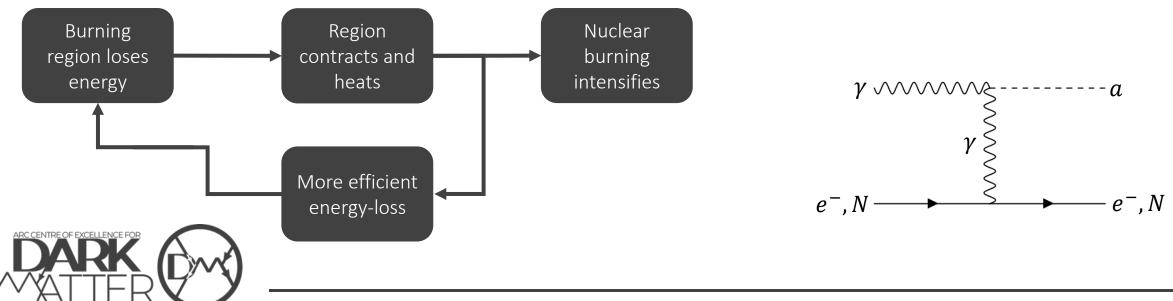


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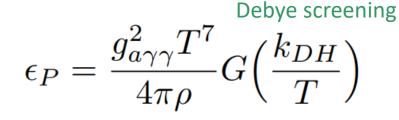


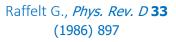


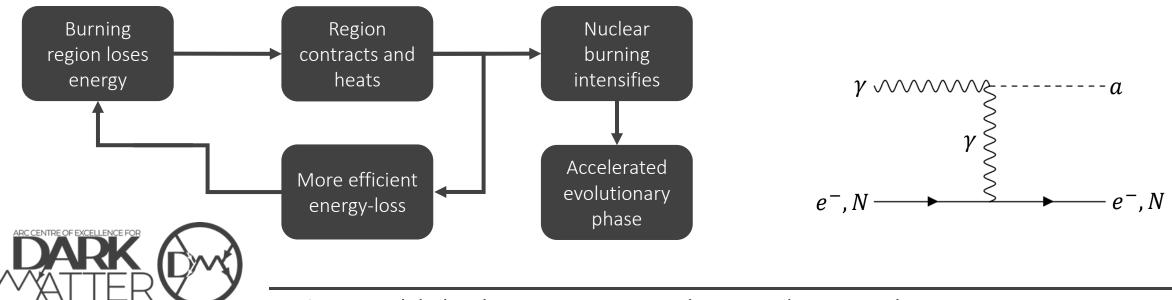


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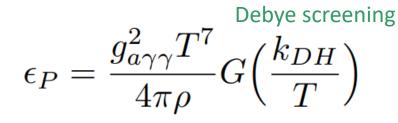




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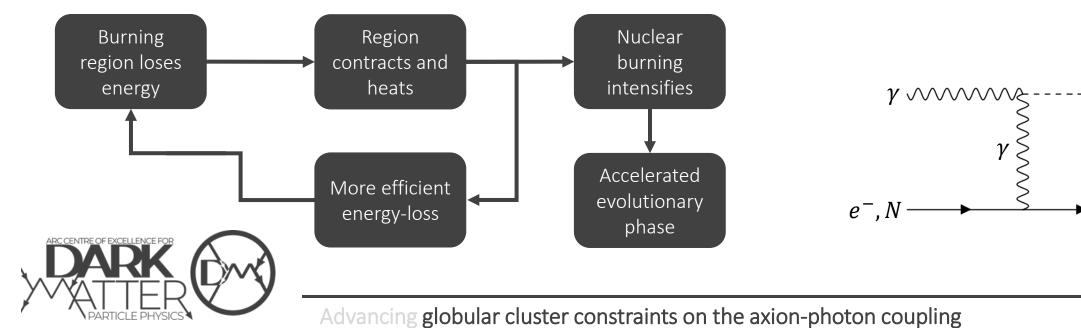
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Constraints derived in this manner are examples of **stellar cooling** bounds

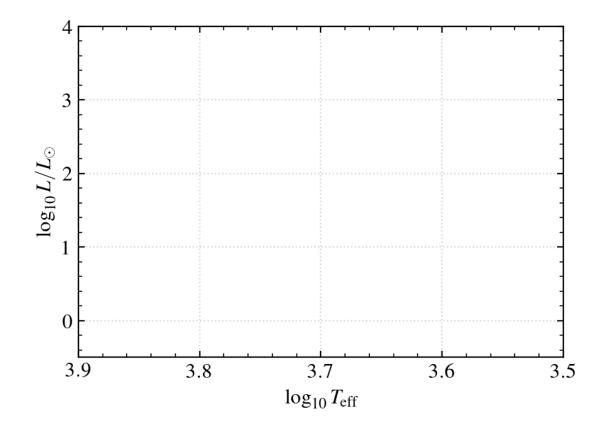


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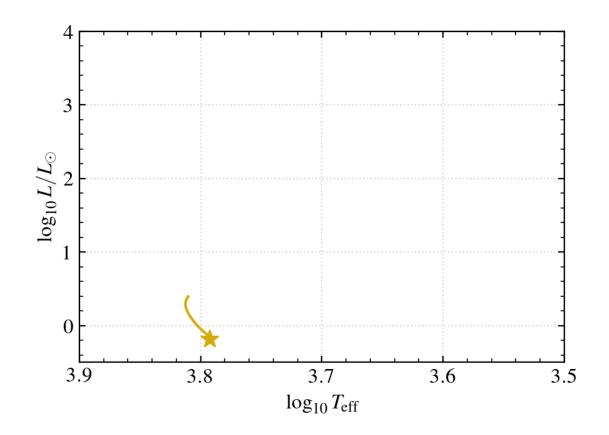






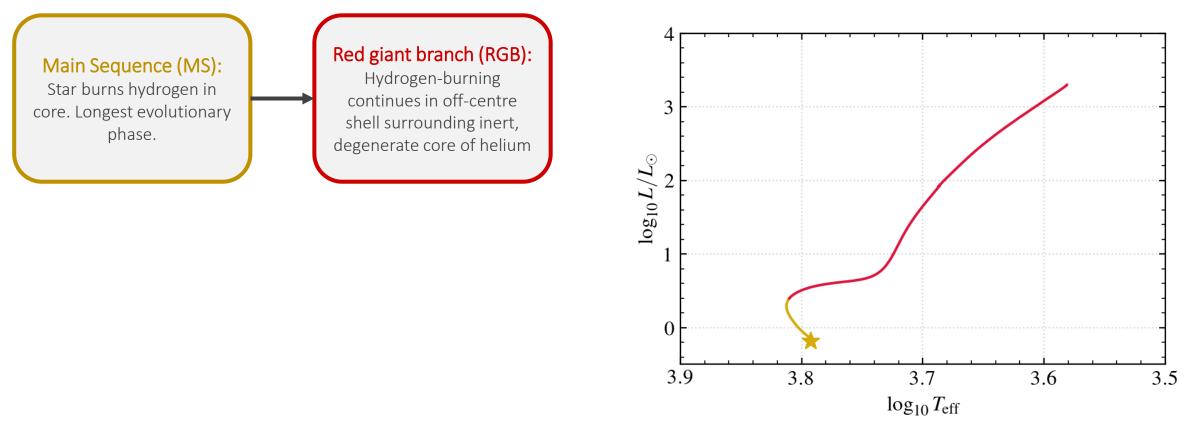


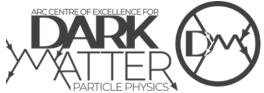
Main Sequence (MS): Star burns hydrogen in core. Longest evolutionary phase.



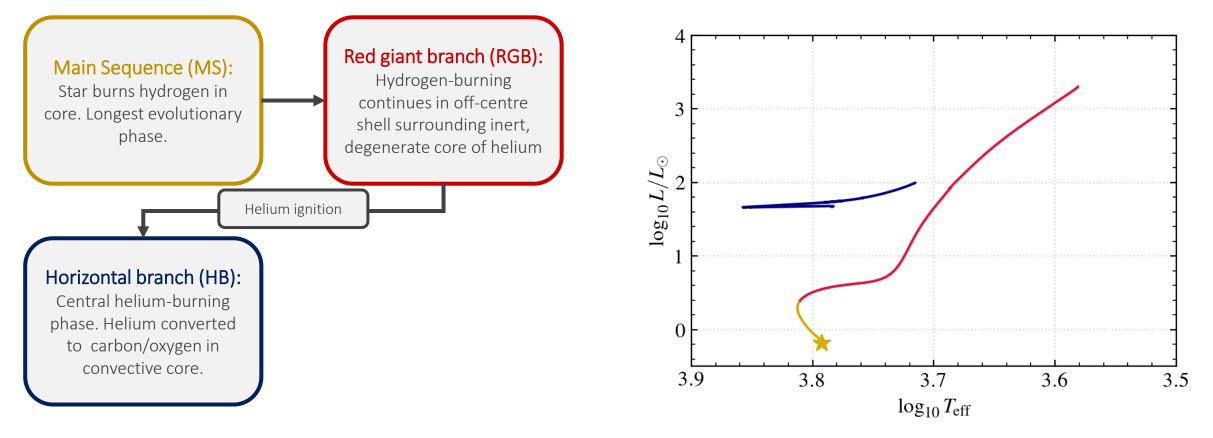






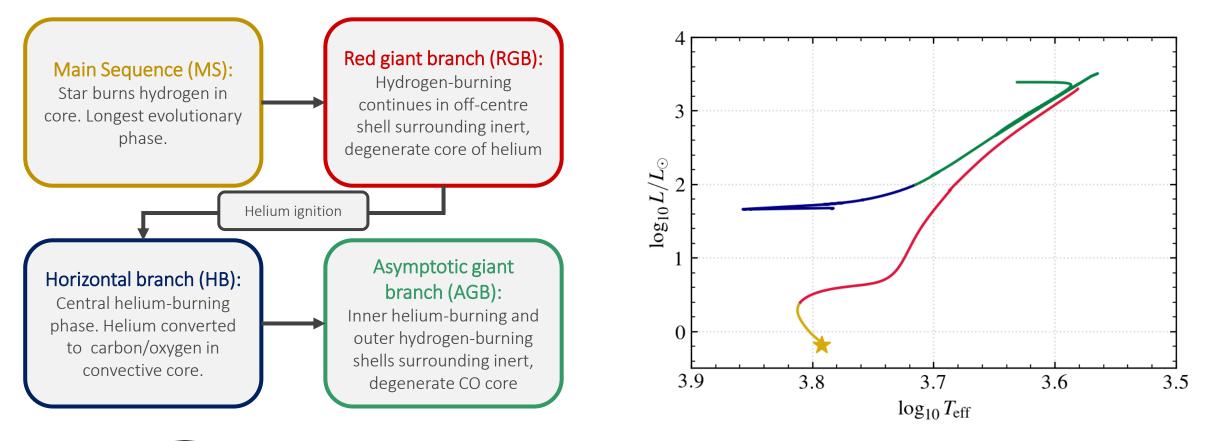


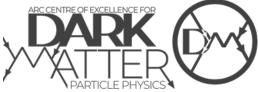




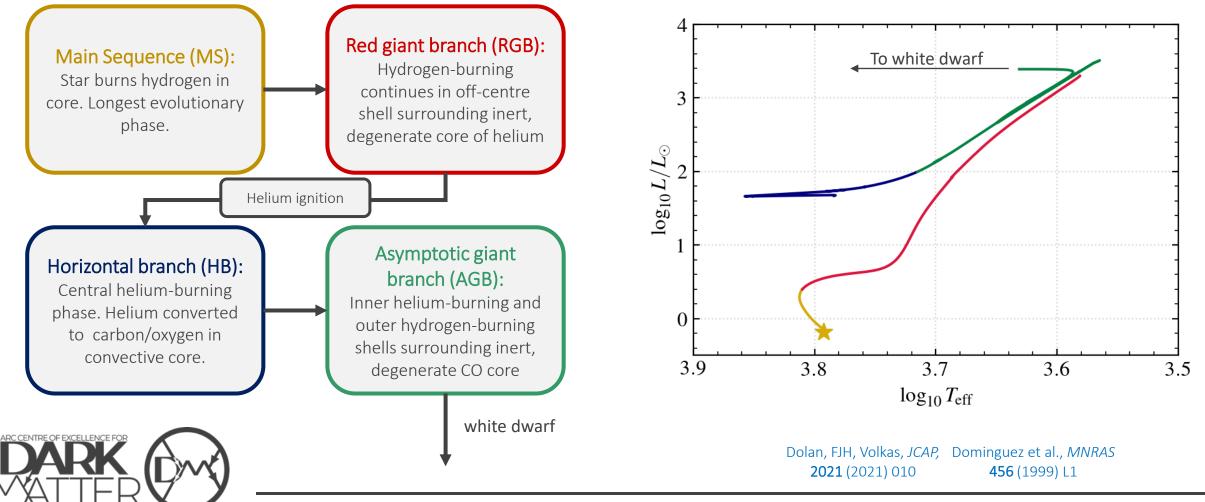




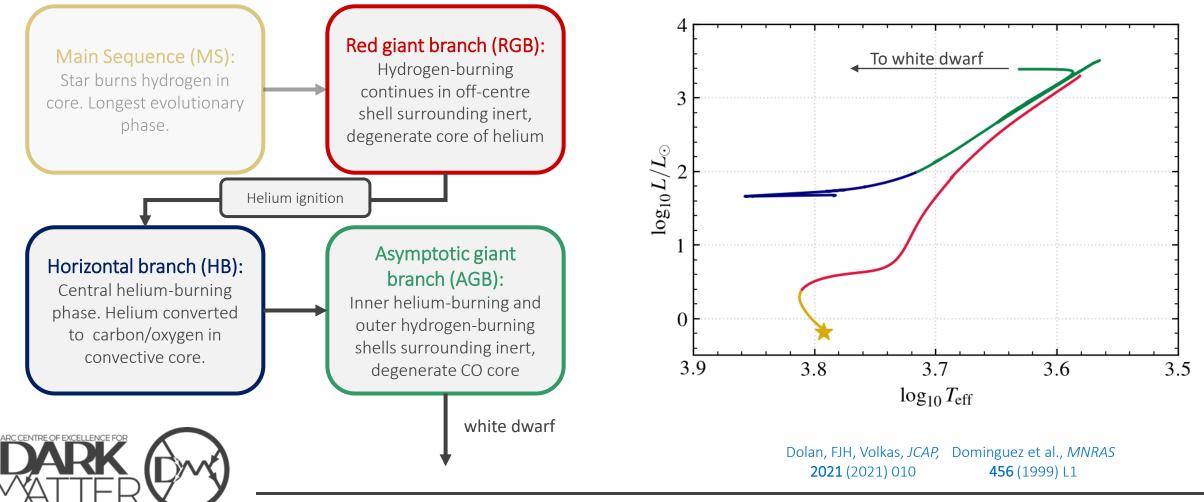




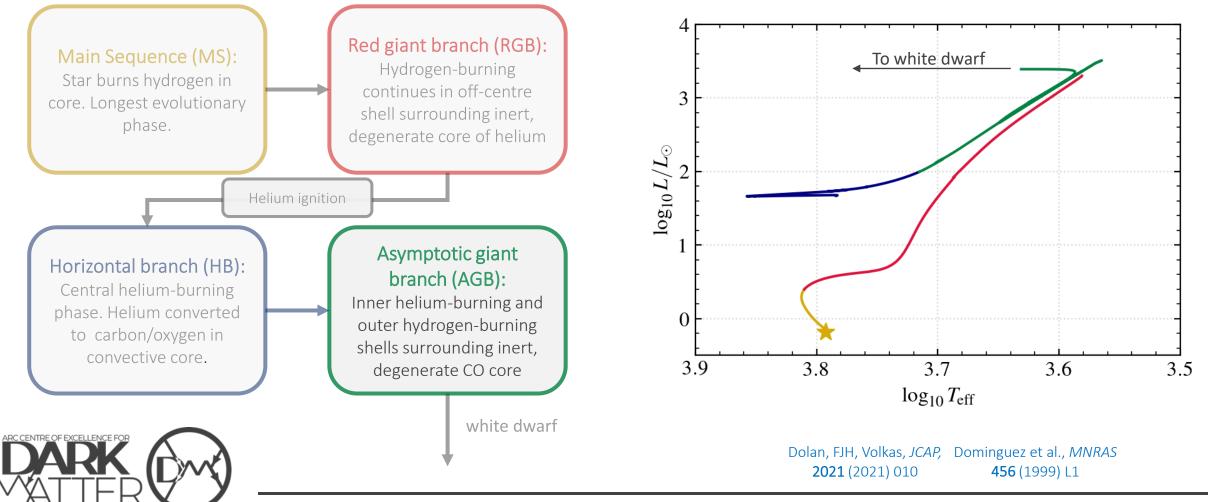




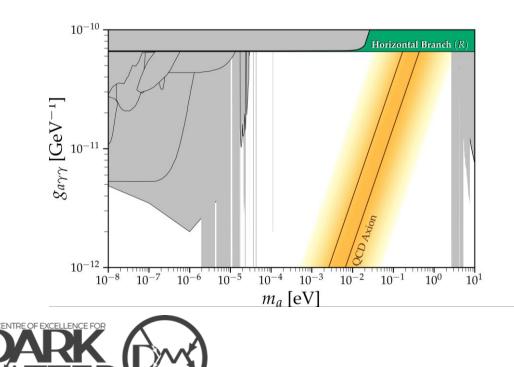








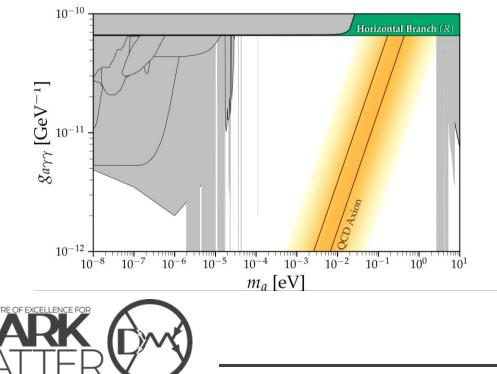
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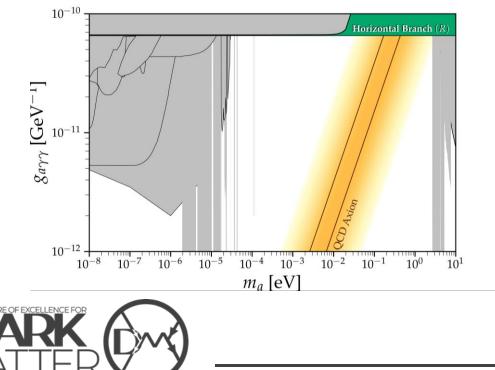




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*R***-parameter:** the ratio of **horizontal branch (HB)** to **red giant branch (RGB)** stars in globular clusters



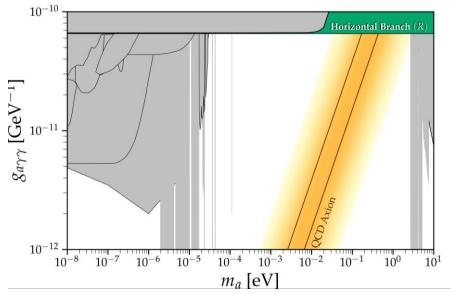




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#### **R**-parameter constraint

• In globular clusters, HB and RGB are populated by stars of approximately the same mass  $(0.8 M_{\odot})$ 

$$R = \frac{N_{\rm HB}}{N_{\rm RGB}} \approx \frac{\tau_{\rm HB}}{\tau_{\rm RGB}}$$

- For  $g_{10} = \frac{g_{a\gamma\gamma}}{10^{-10} \text{ GeV}^{-1}} \sim 1$  energy loss via the **Primakoff process** efficiently drains energy from HB cores, decreasing  $\tau_{HB}$
- Process is inefficient in RGB stars  $\Rightarrow \tau_{RGB}$  is unaffected
- For  $g_{10} \geq 0.66$  predictions contradict observed limit
- Include effects of electron degeneracy on  $\epsilon_P$

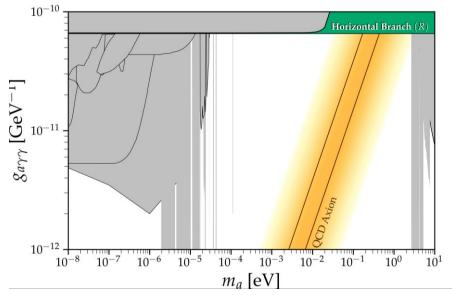
Raffelt & Dearborn., *Phys. Rev. D* **36** (1987) 2211 Ayala, et al., *Phys. Rev. Lett.* **113** (2014) 191302



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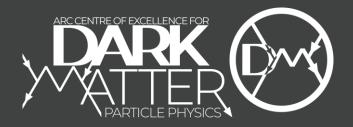
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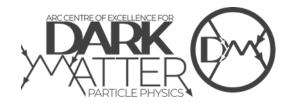


#### So what's the problem?





Simulations of the HB are subject to large systematic uncertainties from the modelling of **mixing across the boundary** of their convective core



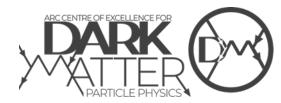
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Issue widely acknowledged in astrophysics literature, but seldom discussed in particle physics



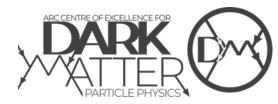
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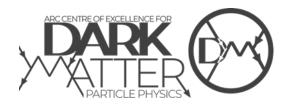
Convective core boundary defined as the point at which the *acceleration* of convective elements vanishes





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Elements arrive at the boundary with non-zero momentum ⇒ penetrate into stable region: **convective overshoot** 

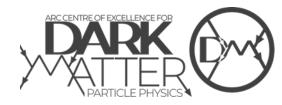




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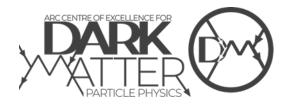


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Stellar evolution code *Modules for Experiments in Stellar Astrophysics* (MESA) is furnished with four different schemes



## Example





#### The issue with **R**





#### Example scheme: standard overshoot

- Time-dependent diffusive process
- Diffusion coefficients decrease exponentially with distance from the convective boundary
- Scale of exponential decrease set by free parameter  $f_{ov}$





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Simulate the HB

- Location of the boundary is unstable
- Can spark dramatic **core breathing pulses** large convective episodes which extend HB duration







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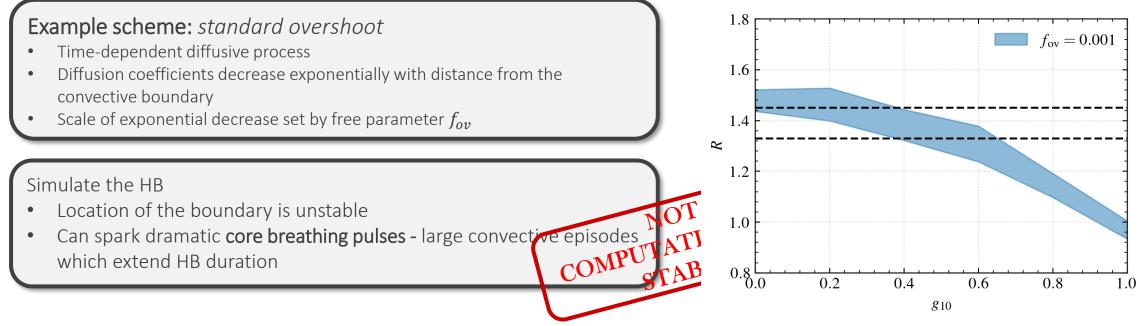
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- Can spark dramatic **core breathing pulses -** large convective episodes **NOT** which extend HB duration **COMPUTATIONALLY** ٠





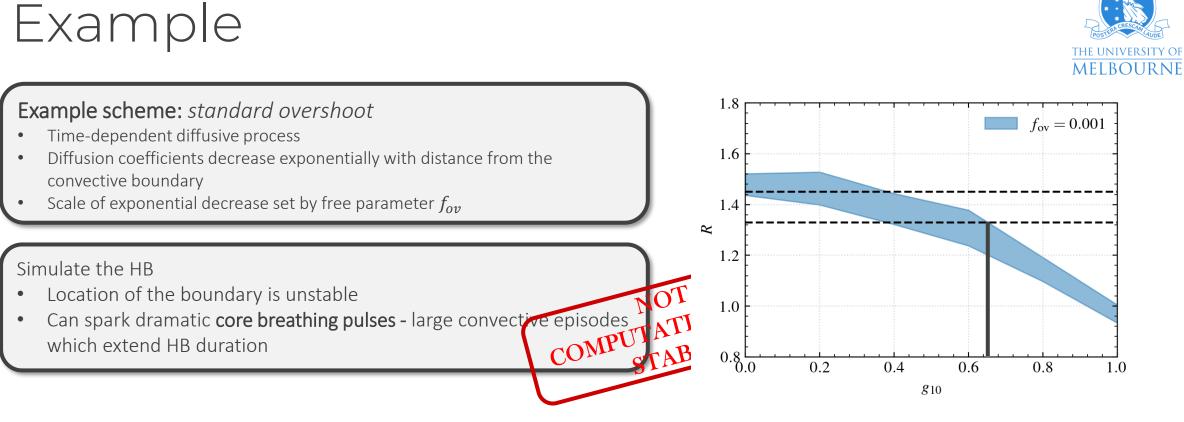






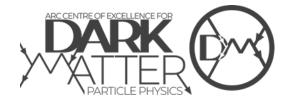
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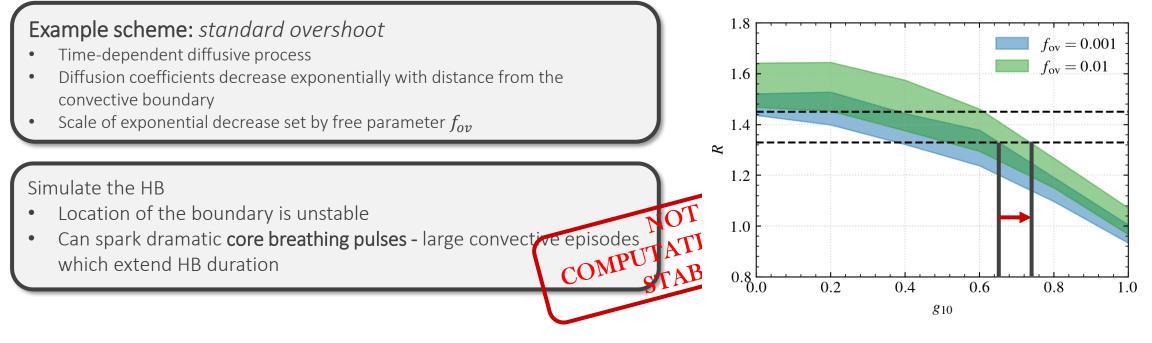




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Example

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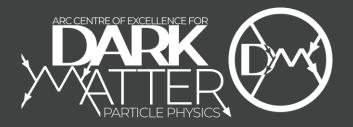
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Choice of  $f_{ov}$  systematically shifts R

Can only constrain when entire range falls below 95% CL

Which value of  $f_{ov}$  do we take?







Other globular cluster parameters exist which can provide complementary constraints on  $g_{a\gamma\gamma}$ 



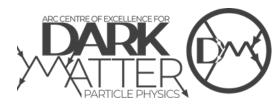




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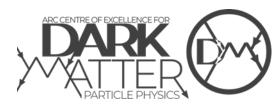
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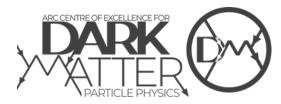
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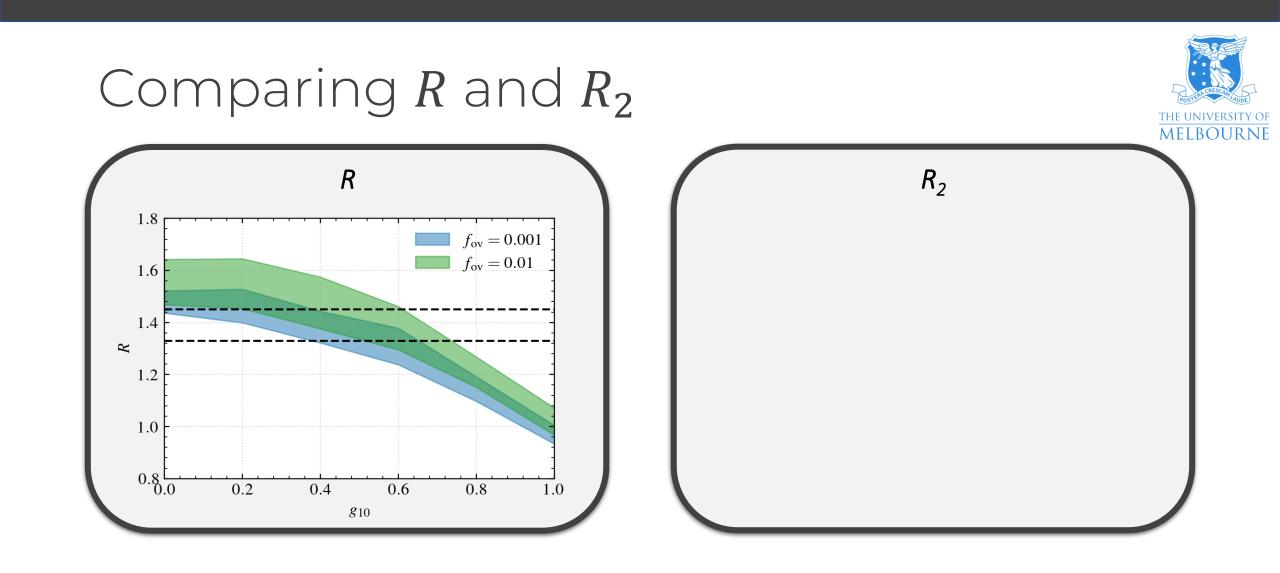


Historically used to constrain the effects of mixing across convective boundaries

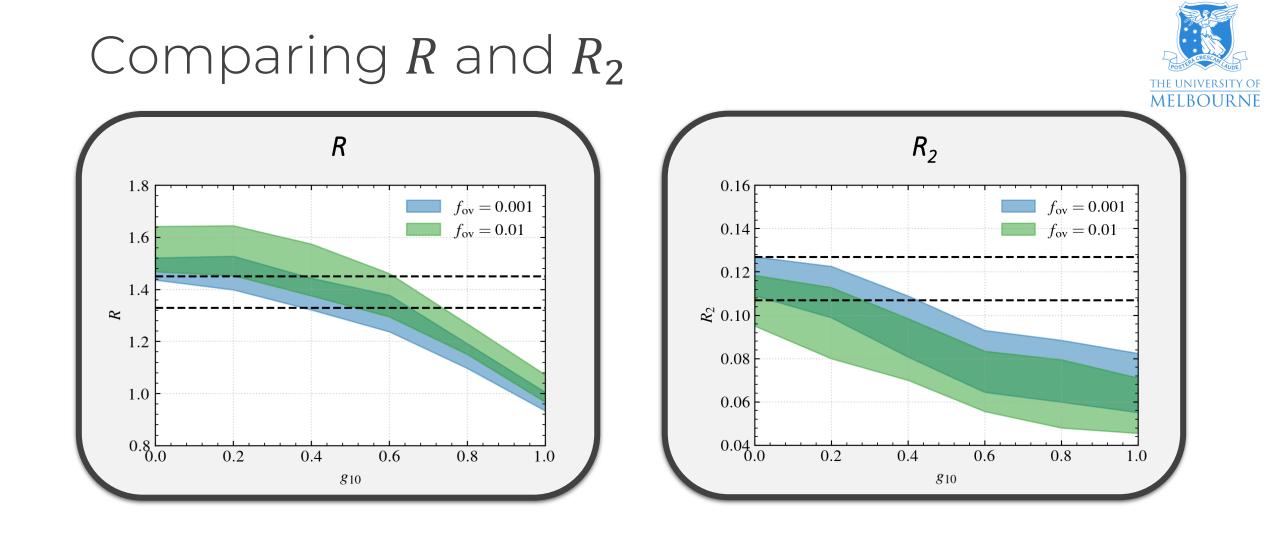
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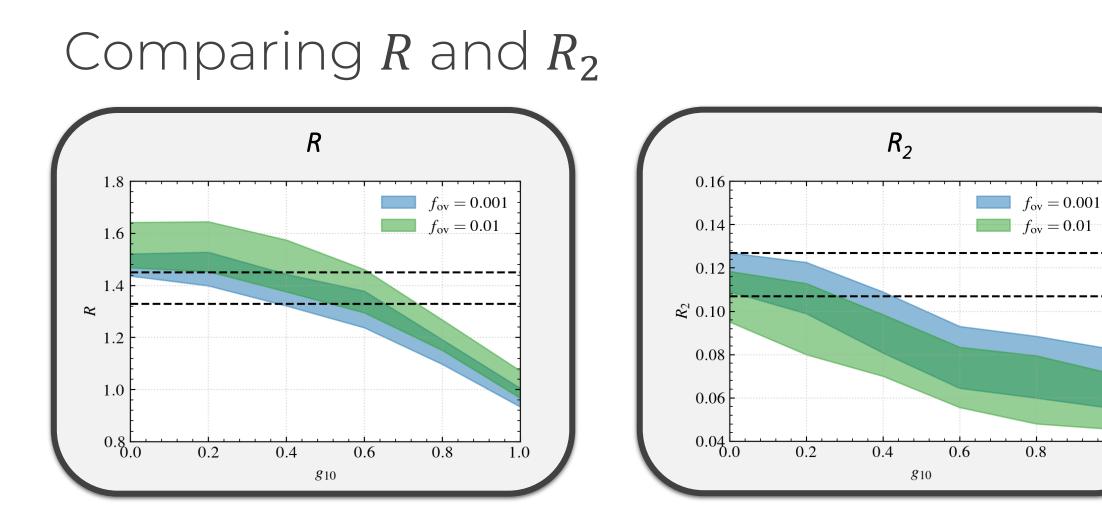










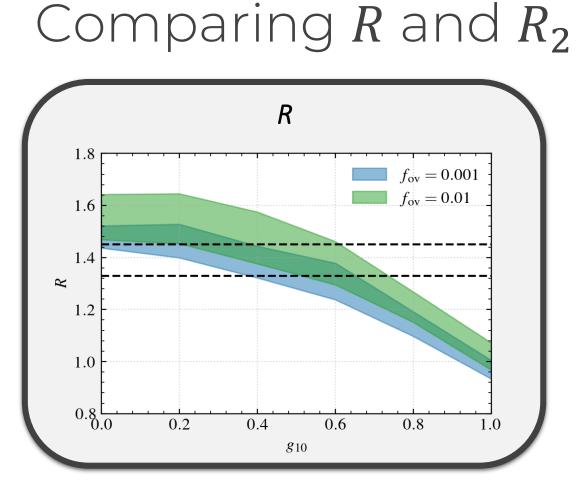


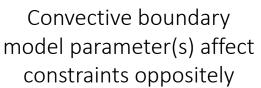
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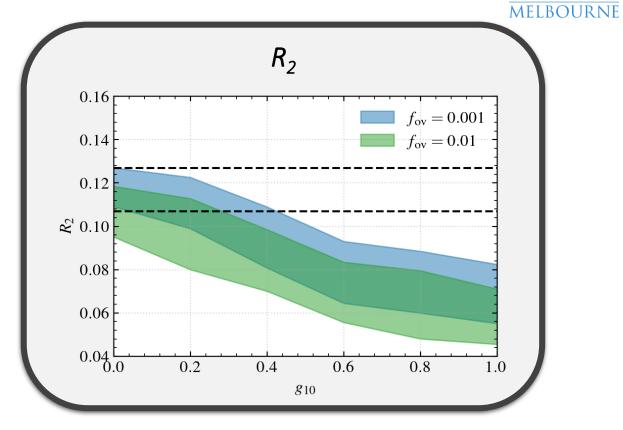
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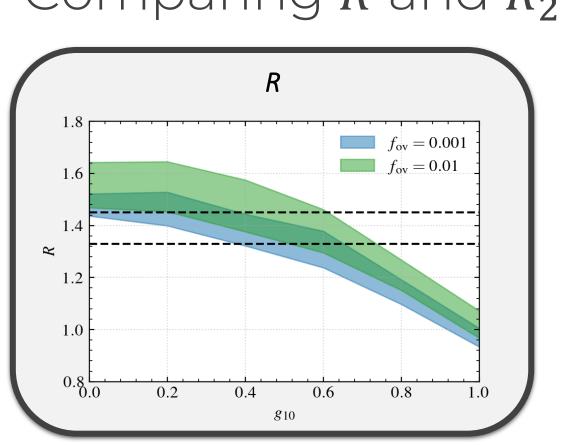


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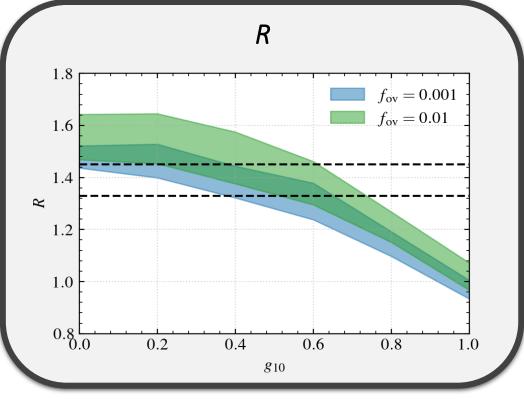
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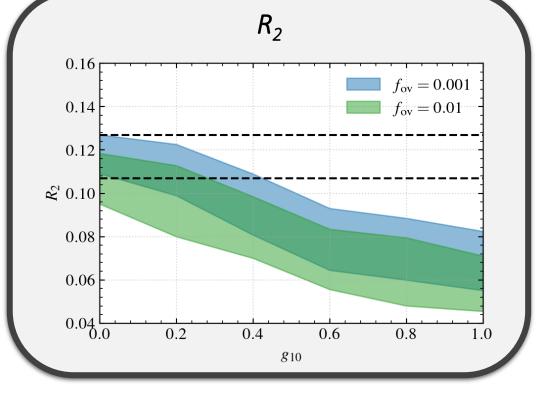
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## Comparing R and $R_2$



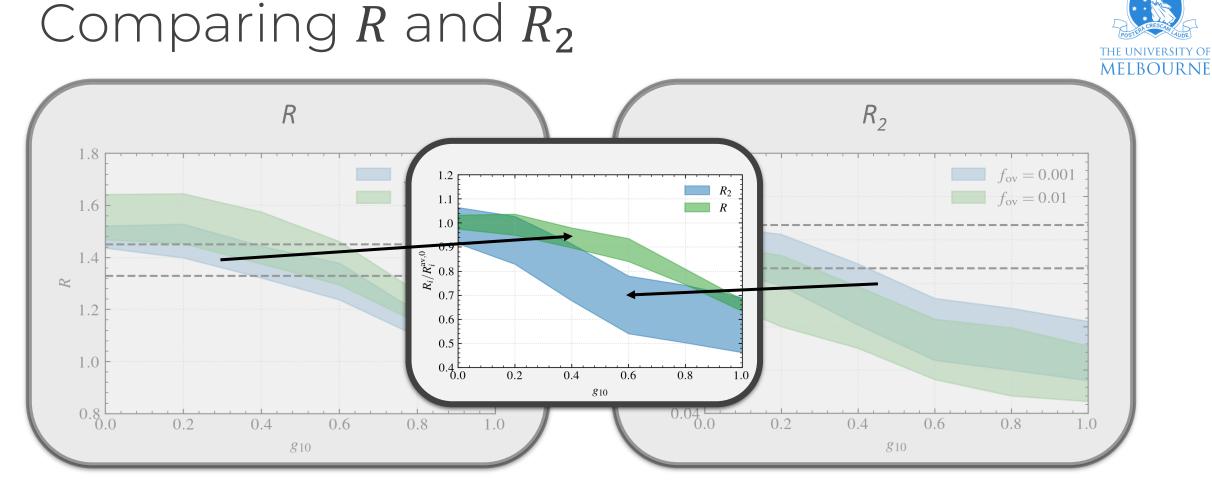


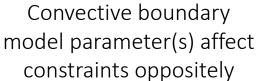
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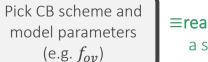
Pick CB scheme and model parameters (e.g.  $f_{ov}$ )





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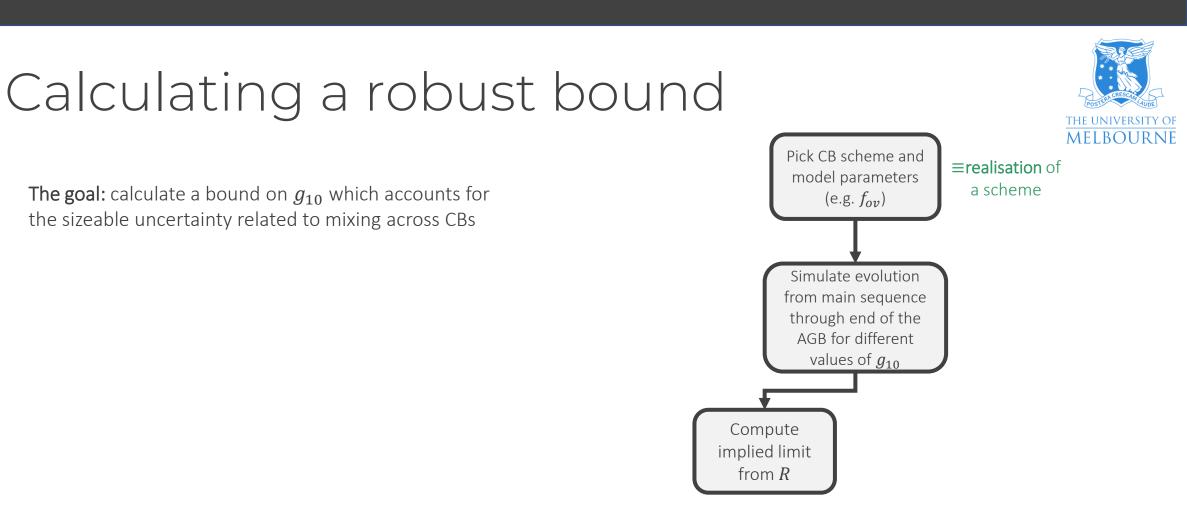


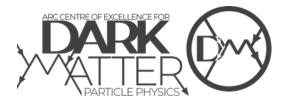
**≡realisation** cascheme



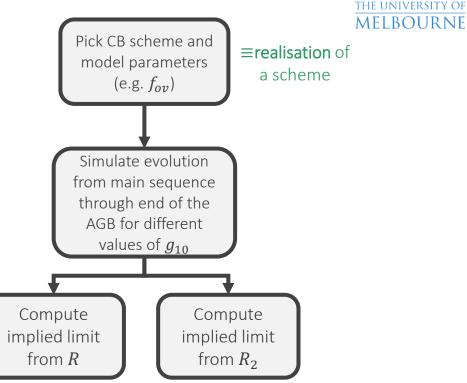
#### Calculating a robust bound The goal: calculate a bound on $g_{10}$ which accounts for the sizeable uncertainty related to mixing across CBs $\begin{bmatrix} Pick CB scheme and$ model parameters $(e.g. for) \\Fick CB scheme and$ $(e.g. for) \\Fick$





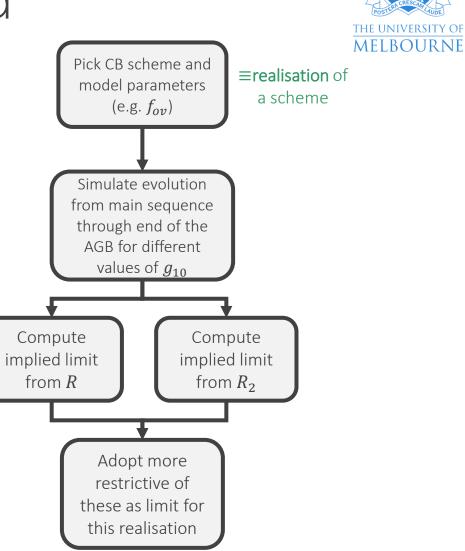


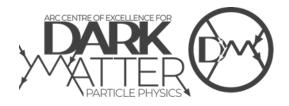
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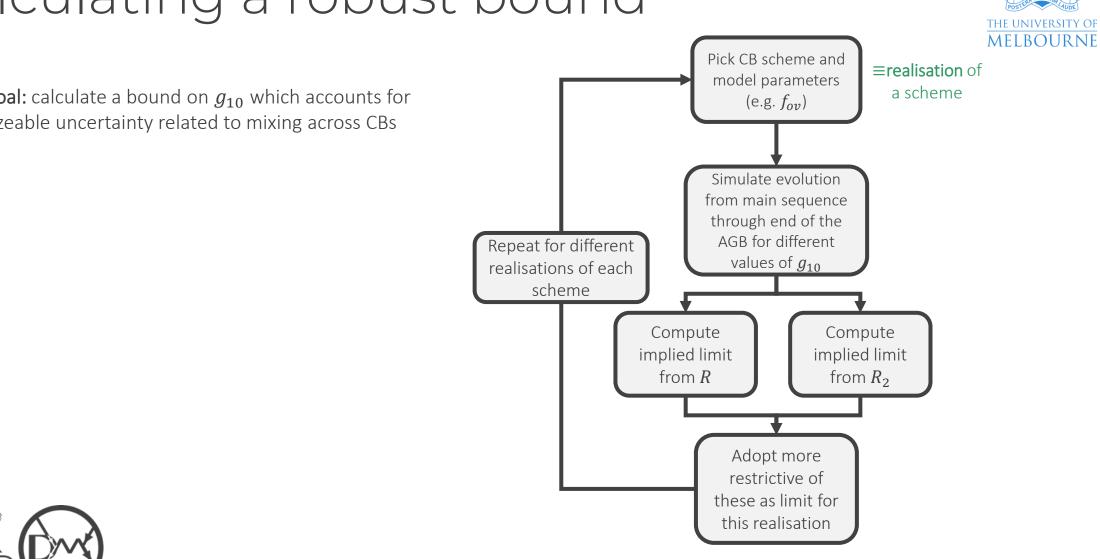




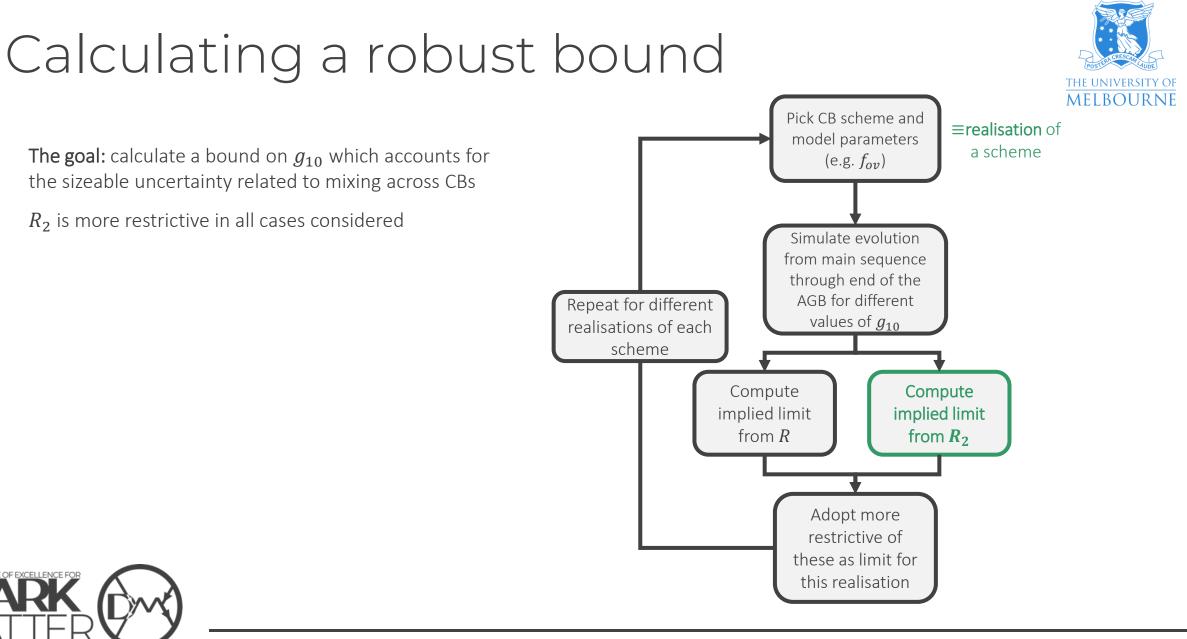
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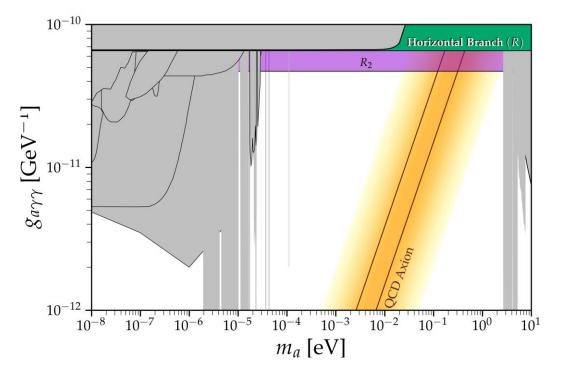




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Most conservative overall limit  $g_{10} \leq 0.47$  – an improvement of 30%





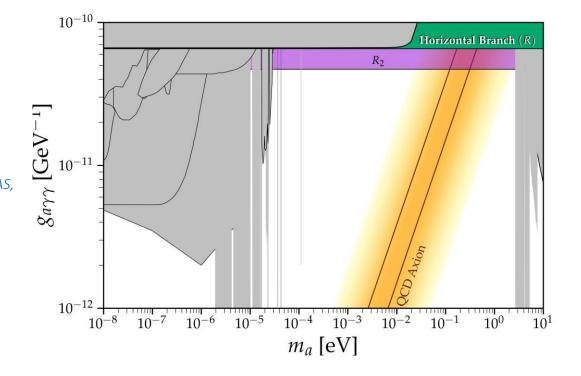


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Complementary constraints on HB convective structure from asteroseismology exist Constantino et al., MNRAS, 452 (2015) 123







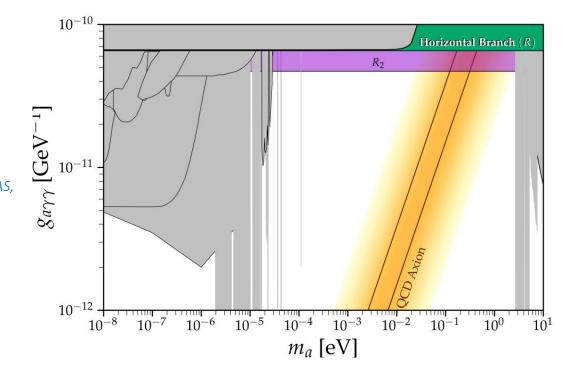
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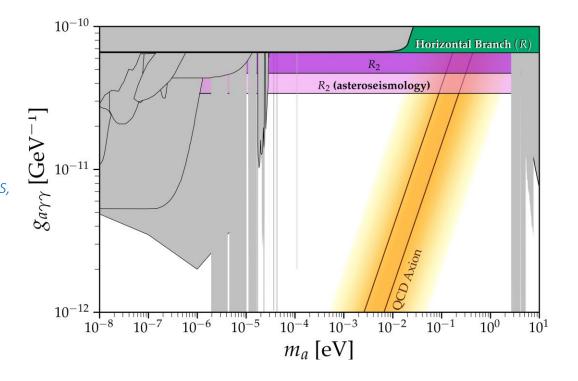
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Evidence not yet conclusive...











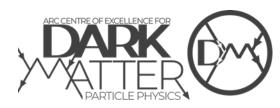
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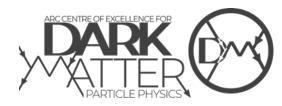




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This may improve to  $g_{10} = 0.34$  or better as uncertainty surrounding mixing across convective boundaries decreases (e.g. through asteroseismology)







#### Backup Slides



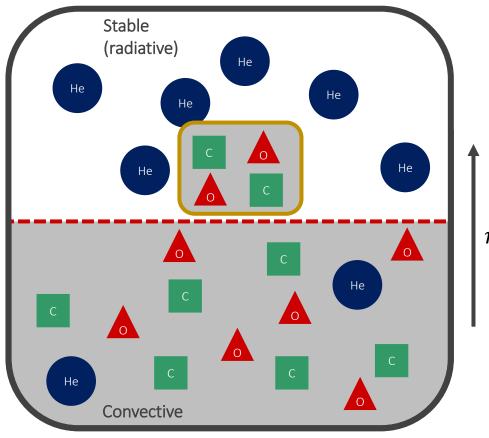
# Mixing across convective boundaries

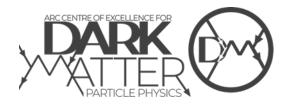
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**Convective elements** arrive at the boundary with non-zero momentum and penetrate the stable region: **convective overshoot** 

C/O are more opaque than  $He \Rightarrow$  convective region grows







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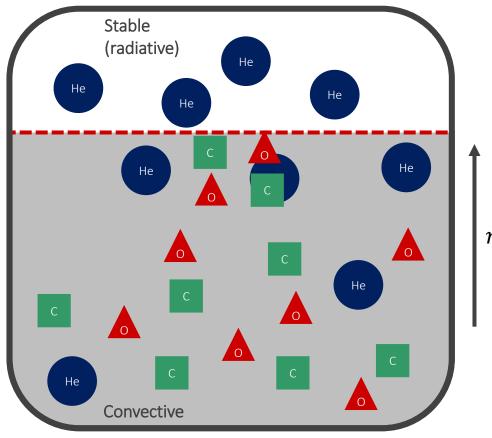
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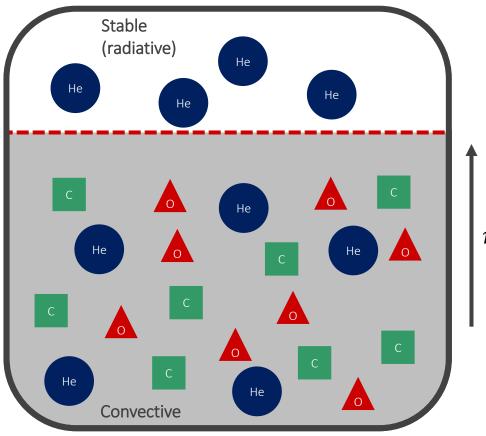
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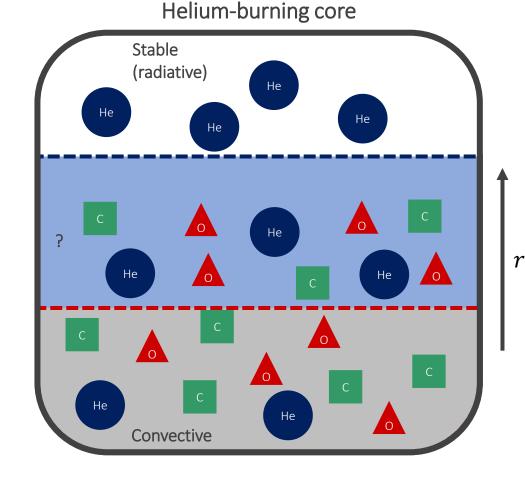
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**Result:** evolution of the core boundary is not stable

Effects are dire if they occur near the end of the HB  $\implies$  large convective episodes which significantly elongate the HB





# Calculating $R_2$



Simulate evolution of  ${\sim}0.8 M_{\odot}$  star through MS, RGB, HB and AGB

Convert results of simulation to probability density function of  $\Delta \log_{10} L_{HB} = \log_{10} L - \log_{10} L_{HB}$ 

$$P(\Delta \log L) = \frac{1}{\tau} \sum_{i=1}^{n} \frac{\Delta t_i}{\sigma \sqrt{2\pi}} \exp\left(-\frac{(\Delta \log L - \Delta \log L_i)^2}{2\sigma^2}\right)$$

A clear minimum exists between HB and AGB peaks

Calculate  $R_2$  as ratio of the areas either side of this minimum Repeat for non-zero values of  $g_{10}$ 

