

Superradiant Instability of Extremal Black Hole in STU Model

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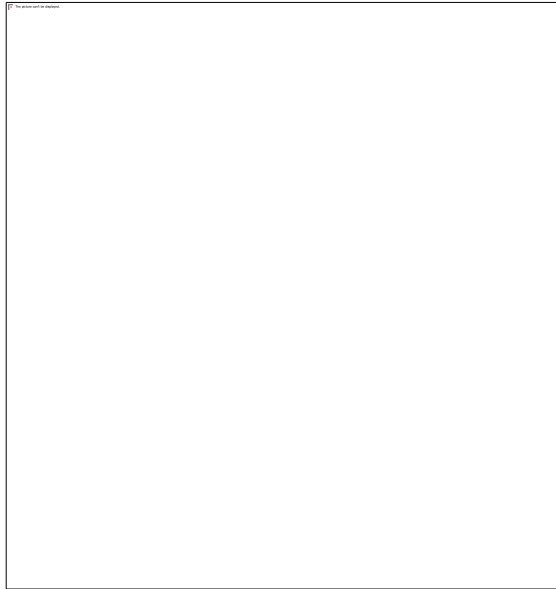
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Arxiv: : [2110.14942](https://arxiv.org/abs/2110.14942)

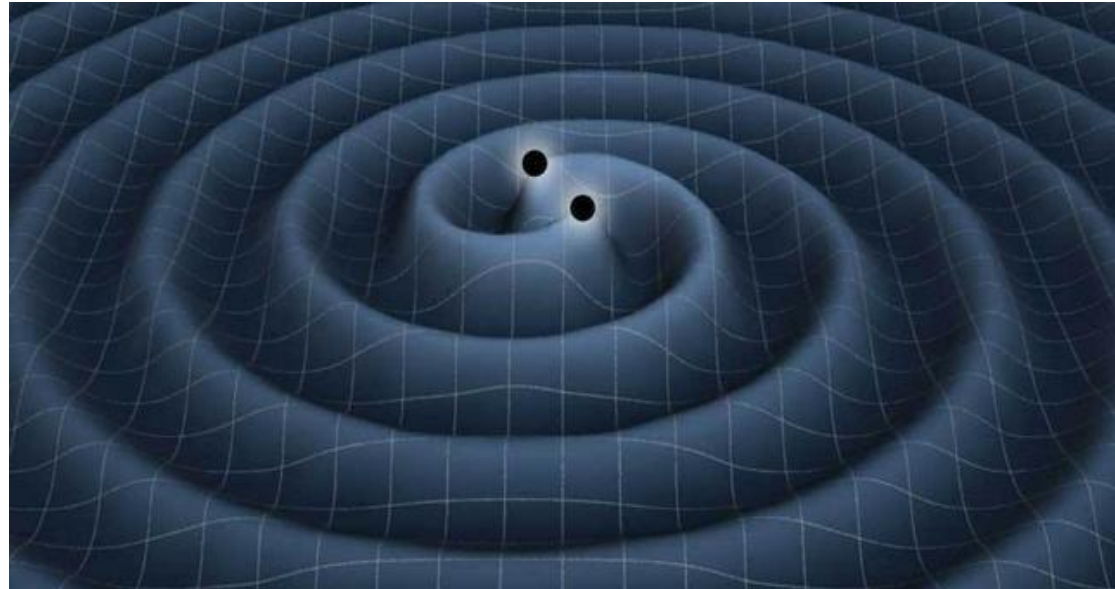
Collaborating with Run-Qiu Yang and Hong Lu

Black Hole Stability

- Black Hole is important objects in gravitational physics.



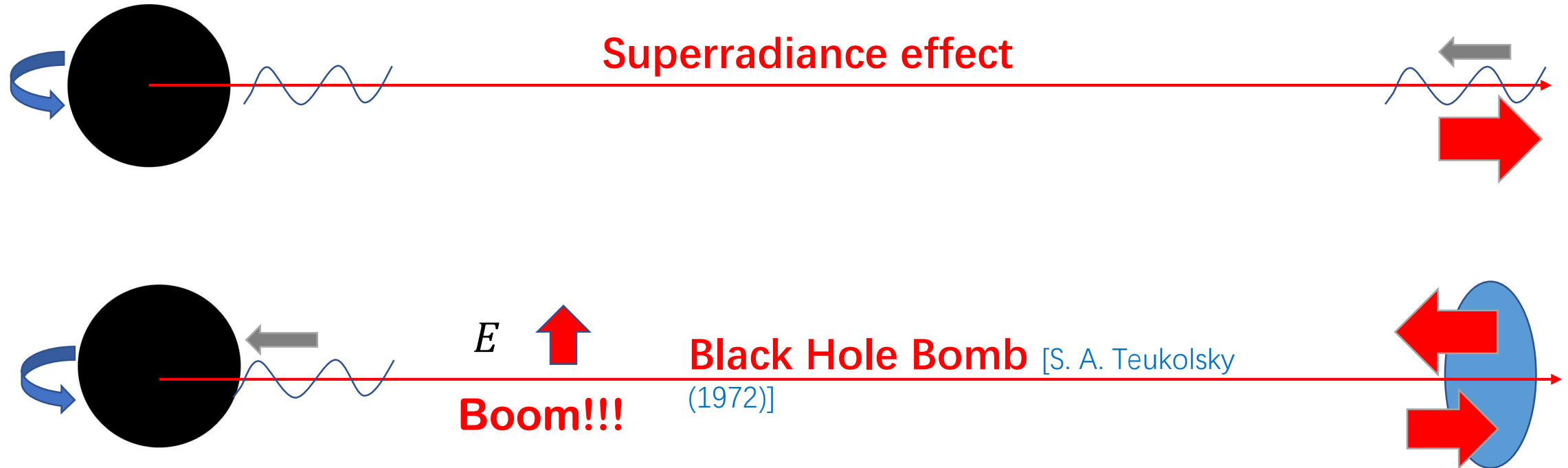
Black Hole image



Gravitational Wave

Motivation

What is superradiant instability?
Can it arise in static black hole ?



○ Kerr black hole [A.A. Starobinskil and S.M. Churilov (1974)] [S.M.Dolan (2007)]

○ Kerr-Newmann Black Hole [H. Furuhashi and Y. Nambu (2004)] [Y.Huang (2018)]

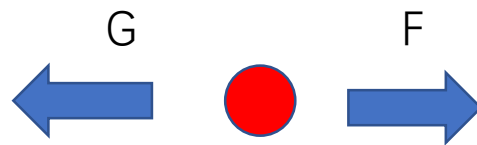
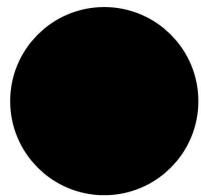
✗ RN Black Hole [S. Hod (2012)][S. Hod(2015)][J. H. Huang and Z. F. Mai (2016)]

Motivation: Why Extremal black hole in STU

$$ds^2 = -(\tilde{H}_1\tilde{H}_2\tilde{H}_3\tilde{H}_4)^{-\frac{1}{2}}dt^2 + (\tilde{H}_1\tilde{H}_2\tilde{H}_3\tilde{H}_4)^{\frac{1}{2}}(dr^2 + r^2d\Omega_2^2),$$
$$e^{\frac{1}{2}\vec{a}_i\cdot\vec{\varphi}} = \tilde{H}_i^{-1}(\tilde{H}_1\tilde{H}_2\tilde{H}_3\tilde{H}_4)^{\frac{1}{4}}, \quad \tilde{A}_i = (\tilde{H}_i^{-1} - 1)dt, \quad \tilde{H}_i = 1 + \frac{4\tilde{Q}_i}{r^2}.$$

Remark:

1. **No** Hawking Radiation.
2. It is a **static** black hole.
3. When four charges are **equal**, it **reduce** to Extremal RN black hole.



STU: rich structure

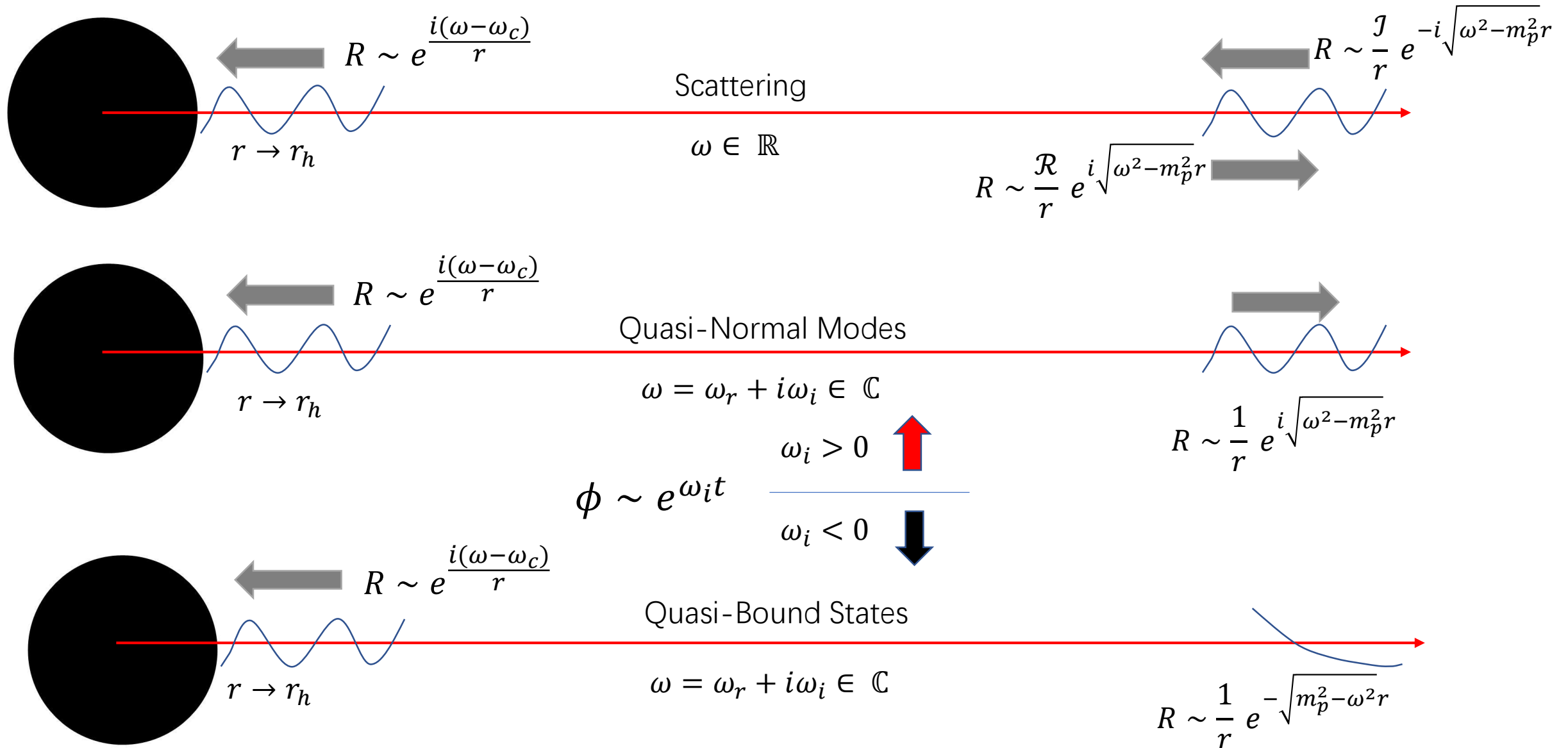
Linear Perturbation

KG
equation: $(g^{\mu\nu} D_\mu D_\nu - m_p^2) \Phi = 0, \quad D_\mu := \nabla_\mu - i\tilde{q}_1 \tilde{A}_{1\mu} - i\tilde{q}_2 \tilde{A}_{2\mu} - i\tilde{q}_3 \tilde{A}_{3\mu} - i\tilde{q}_4 \tilde{A}_{4\mu},$

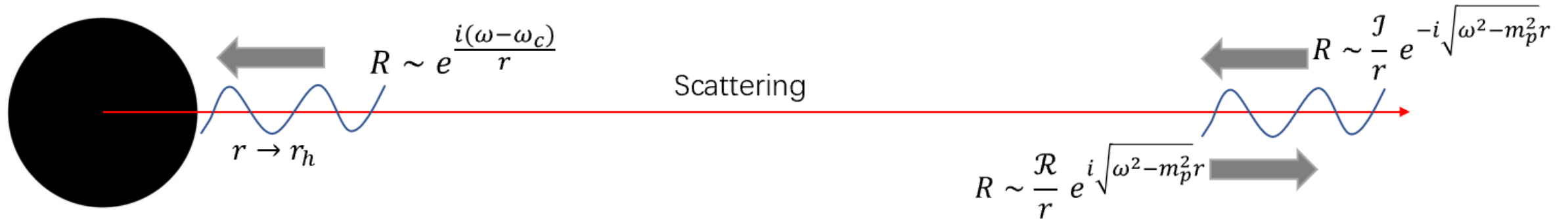
Anstaz: $\Phi = e^{-i\omega t} R(r) Y_{\ell m}(\theta, \varphi)$

Radial Equation $-r^2 \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + UR = 0,$

Boundary Condition



Superradiant Scattering



Wroskain Determinant

$$|\mathcal{R}|^2 = |\mathcal{J}|^2 - \frac{\omega - \omega_c}{\sqrt{\omega^2 - m_p^2}} |\mathcal{T}|^2$$

Conserved Current

Energy Momentum Tensor

$$\tilde{T}_{\mu\nu} = \sum_{j=r,i} 2\nabla_{\mu}\Phi_j\nabla_{\nu}\Phi_j - g_{\mu\nu} \left(\nabla^{\rho}\Phi_j\nabla_{\rho}\Phi_j + \left(\sum_i^4 \tilde{q}_i \tilde{A}_i^{\rho} \right)^2 \Phi_j^2 \right) - m_p^2 \Phi_j^2$$

Energy Current: $J_E^{\mu} = \tilde{T}_{\mu\nu} \left(\frac{\partial}{\partial t} \right)^{\nu}$

Total Energy: $E = \int_V \sqrt{-g} J_E^0$

Energy Flux: J_E^r

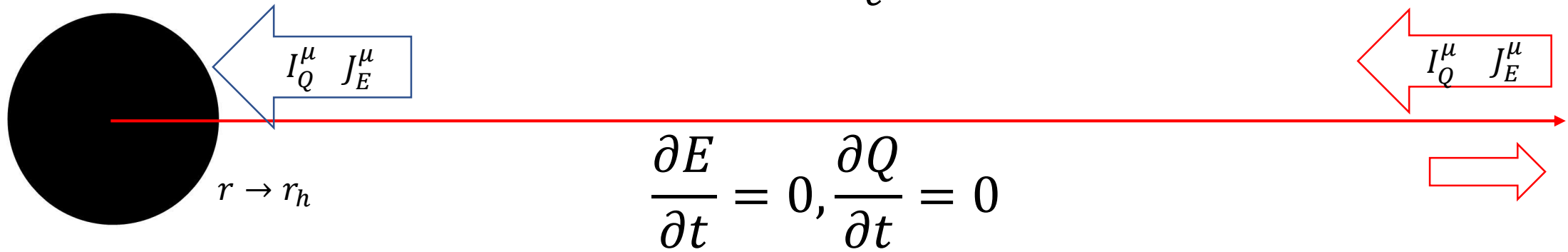
Charge Current $iI_Q^{\mu} = \Phi^{\dagger} D^{\mu}\Phi - \Phi(D^{\mu}\Phi)^{\dagger}$

Total Charge $Q = \int_V \sqrt{-g} I_Q^0$

Charge Current: I_Q^r

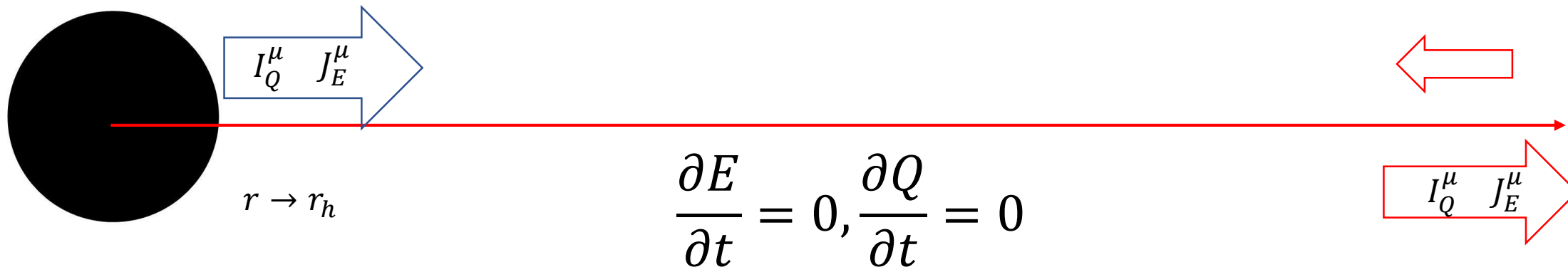
Current (Superradiant Scattering)

$$\omega > \omega_c$$



$$\omega \in \mathbb{R}$$

$$\omega < \omega_c$$



Superradiant Condition for Complex Frequency

$$\frac{\partial E}{\partial t} = 2\omega_i E, \quad \frac{\partial Q}{\partial t} = 2\omega_i Q$$



$$\omega_i > 0$$

E, Q



Unstable

$$\omega_i < 0$$

E, Q

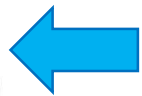
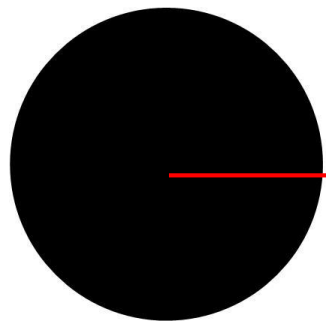


Stable

Conservation Law:

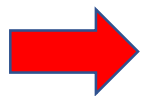
$$\frac{\partial E}{\partial t} = \int_S J_E^r \Big|_{r \rightarrow \infty}^{r \rightarrow 0}$$

$$\frac{\partial Q}{\partial t} = \int_S I_Q^r \Big|_{r \rightarrow \infty}^{r \rightarrow 0}$$



$$\omega_r > \omega_c$$

Quasi-Bound State



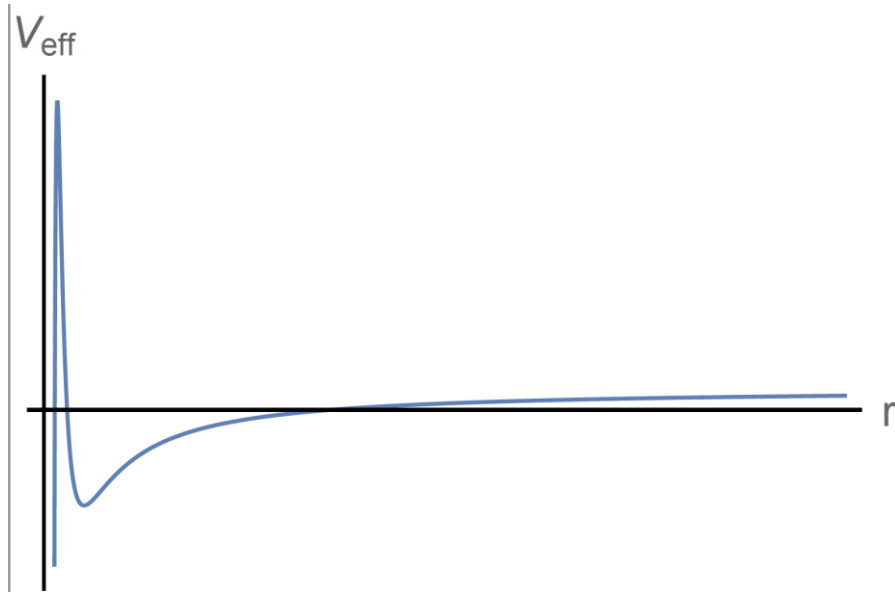
$$\omega_r < \omega_c$$

Quasinormal Modes



Superradiant Condition for instability: $\omega_i > 0 \rightarrow \omega_r < \omega_c$

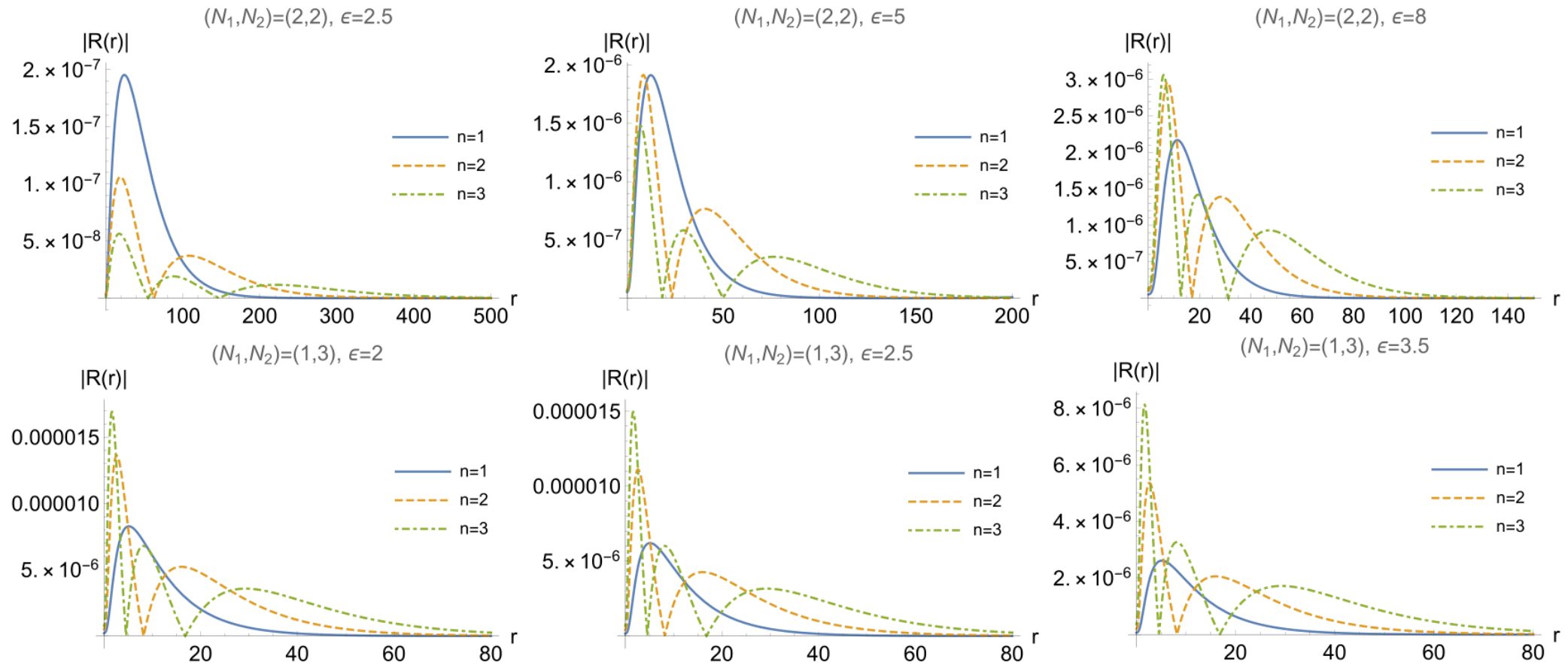
Numerical Result: Single Peak (Only $(N_1, N_2) = (2, 2)$)



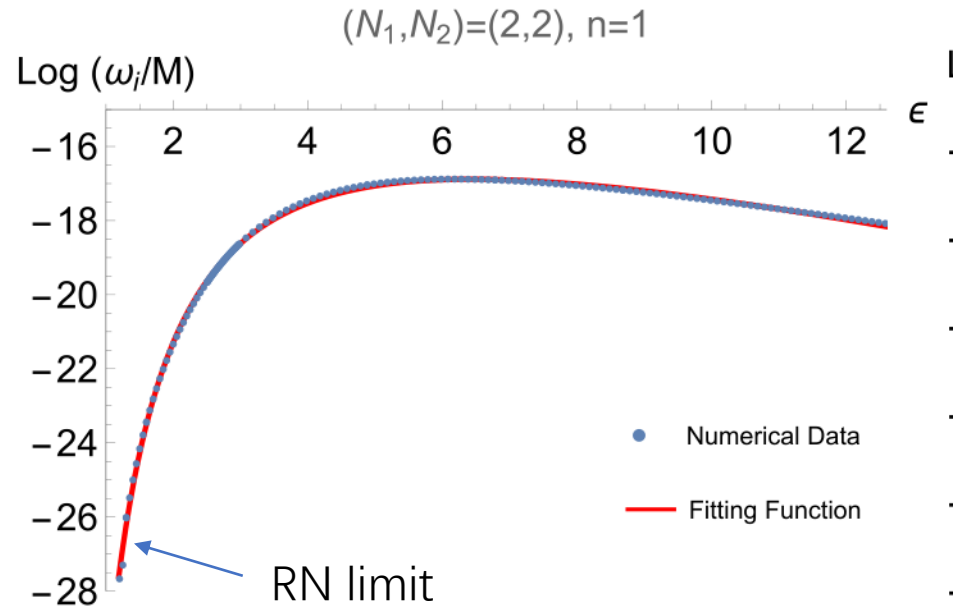
Parameters Choice:

$$\epsilon = \frac{\sqrt{N_1} Q_2}{\sqrt{N_2} Q_1} - 1, \ell = 1, Q_1 = 1, q_2 = 0, q_1 = \frac{2}{10}, m_p = \frac{2}{10}$$

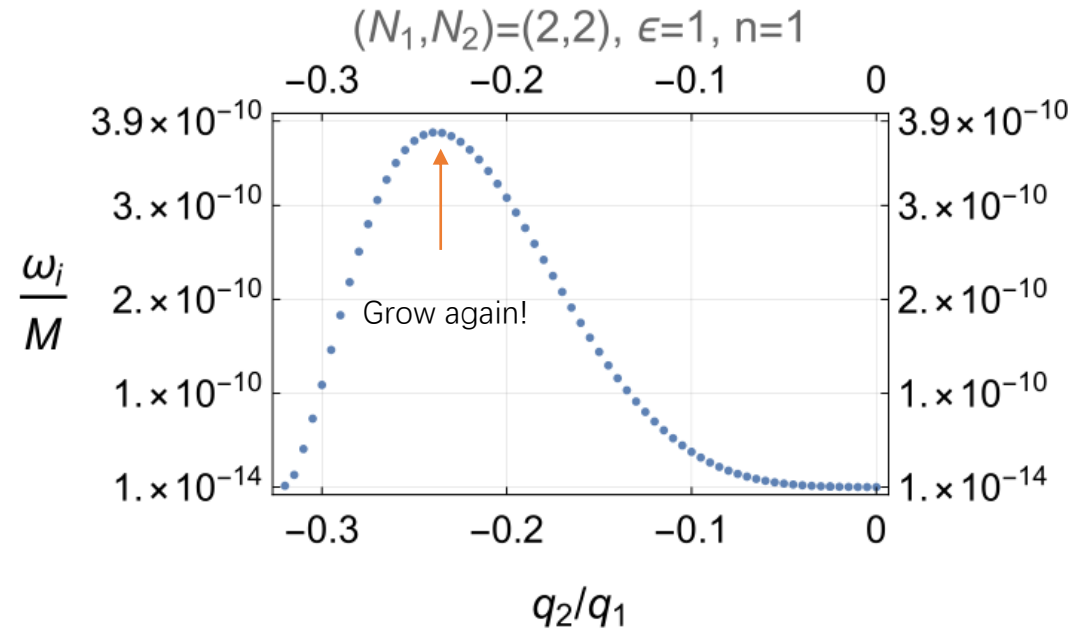
Numerical Result: Single Peak



Result



Superradiantly stable RN black hole is a fine-tuning case



$$\epsilon = 0.01, Q_1 = 1, m_p = \frac{2}{10}, q_1 = 20.2, q_2 = -20,$$

$$\omega_i = 0.1927 + 3.752 \times 10^{-8}i$$

Conclusion

- The relation between superradiant instability and unstable quasi-bound states or quasi-normal modes.
- RN black hole is a fine-tuning case in superradiant instability in the framework of STU supergravity model.

Thank you for your attention !