

Continuum Dark Matter



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With Hong, Kurup, Lee, work in progress With Perelstein, Ferrante, work in progress With Csaki, Ismail, Nakli, Suzuki, work in progress

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- Introduction
- Gapped Continuum
- Gapped Continuum QFT
- Equilibrium and Non-equilibrium Thermodynamics
- Z-portal Model for Gapped Continuum DM
- Summary

Outline

DM-what we don't know

- **Composition of Dark Matter**
- Interaction of Dark Matter

Mass of Dark Matter (range:10⁻²² to 10⁶⁷ eV)

continuum states

is a localized excitation of quantum field (i.e. particle)

• Our Proposal: Dark Matter is made of an ensemble of gapped

- It's not even clear whether the DM that provide successful explanations to the rotation curve of disk galaxies, CMB, and large structure formation

continuum states

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Mills theory in large-N limit, etc

• Our Proposal: Dark Matter is made of an ensemble of gapped

- It's not even clear whether the DM that provide successful explanations to the rotation curve of disk galaxies, CMB, and large structure formation
- continuum with a mass gap is not so uncommon in condensed matter physics: e.g. edge state in fractional quantum hall effect, topological superfluid, 2D Ising model, 2d SU(2) Thirring model, 2d SU(N) Yang-

scales (no mass gap).

Unparticles (Georgi): another example of gapless continuum

is dual to N = 4 SUSY broken to N = 2 via masses for two chiral adjoints)

(before Higgs discovery) by Terning et al, Falkowski et al et al (SL): for off-shell form factor (by gapped continuum) for Higgs EFT Quiros et al)

- The appearance of a continuum is very common in QFT's: e.g. spectrum of CFTs necessarily forms a continuum since the theory does not admit any mass
- String Theory (e.g. Gubser et al, Kraus, Trivedi et al, etc): gapped continuum
- shows up when one has a large number of D3 branes distributed on a disc (which
- Gapped Continuum in particle physics: -Softwall model (Higgs with a small mass gap)
- -Quantum Critical Higgs (Higgs pole + gappend continuum: after Higgs discovery) by Csaki
- -Continuum Naturalness (for solving little hierarchy by Csaki et al (SL), and also by

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The appearance of a continuum is very common in QFT's: e.g. spectrum of CFTs necessarily forms a continuum since the theory does not admit any mass

- So, how about DM as a gapped continuum? (never explored before!)
- shows up when one has a large reason of D3 branes distributed on a disc (which

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Gapped Continuum, instead of ordinary particles

Continuum DM: singly-excited states are characterized by a continuous parameter μ^2 , in addition to the usual 3-momentum p

 $\rho(\mu^2) d\mu^2$, where ρ is the spectral density of the theory

$$\langle 0|\Phi(p)\Phi(-p)|0\rangle = \int \frac{d\mu^2}{2\pi} \frac{i\rho(\mu^2)}{p^2 - \mu^2 + i\epsilon}$$

- The parameter μ^2 plays the role of mass in the kinematic relation $p^2 = \mu^2$ for each state. The number of states is properly





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CFT Continuum vs Gapped Continuum (IR deformation)



Physics of Gapped Continuum DM

CFT continuum case:

- It's often stated that CFT's and theories with continuum spectra do not have a particle interpretation and no S-matrix can be defined: interactions leading to a non-trivial fixed point are also essential for producing the continuum spectrum of the theory
- by turning off the interactions, the spectrum changes from continuum into that of an ordinary free particle, hence the asymptotic states defined in the usual manner would not capture the physics of the system properly
- this means that one needs to find an alternative approach for defining scattering processes

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- Our theoretic description of gapped continuum: Generalized Free
- Continuum (continuum analog of Generalized Free Fields: Greenberg
 - Also: Polyakov, early '70s- skeleton expansions'
 - CFT completely specified by 2-point function-rest vanish

CFT Continuum

Generalized Free Fields

- CFT completely specified by 2-point function rest vanish
- Scaling 2-point function:

Can be generated from:

$$G(p) \sim \int_{\mu^2}^{\infty} dM^2 \frac{\rho}{p^2}$$

Polyakov, early '70s- skeleton expansions

$$\begin{split} G(p^2) &= -\frac{i}{\left(-p^2 + i\epsilon\right)^{2-\Delta}} \\ \mathcal{L}_{\rm GFF} &= -\hbar^{\dagger} \left(\partial^2\right)^{2-\Delta} \hbar \quad \stackrel{\rm Georgi}{\underset{\rm hep-ph/0703260}{\text{Georgi}}} \end{split}$$

Branch cut starting at origin - spectral density purely a continuum:



there is a non-trivial continuum (with a mass gap), and described by:

$$S = \int \frac{d^4p}{(2\pi)^4} \, \Phi$$

Generalized Free Continuum

- -consider the case that the effects of the strong interactions can be captured by the fact that
 - $\Phi^{\dagger}(p)\Sigma(p^2)\Phi(p)$

there is a non-trivial continuum (with a mass gap), and described by:

$$S = \int \frac{d^4p}{(2\pi)^4} \Phi$$

which is designed to properly reproduce the two-point function of theory $\int d^4x \; e^{ip(x-y)} \langle 0|T\Phi(x)\Phi^{\dagger}(y)|0\rangle = \langle 0|\Phi(x)\Phi^{\dagger}(y)|0\rangle = \langle 0|\Phi(x)\Phi^{\dagger}(y)|0\rangle$

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Generalized Free Continuum

- -consider the case that the effects of the strong interactions can be captured by the fact that
 - $\Phi^{\dagger}(p)\Sigma(p^2)\Phi(p)$

$$\Phi(p)\Phi^{\dagger}(-p)|0\rangle = \frac{i}{\Sigma(r^{2})} \int du^{2} i \rho(\mu^{2})$$

This picture is supported by the concrete
eakly coupled (reextra dimensional construction!
of free field"

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Dark Matter Continuum Spectral Density from 5D Model

♦ modeling generalized free continuum by Warped 5D model $ds^2 = e^{-2A(y)}dx^2 + dy^2$

- warped 5D setup we will have a 3-brane placed at the position z = R, which from the point of view of the gapped continuum field will be a UV brane cutting off the space

The 5D action of the coupled scalar-gravity system

$$S = \int d^5x \sqrt{g} \left(-M^3R + \frac{1}{2}(\partial\phi) - V(\phi) \right) - \int d^4x \sqrt{g^{ind}} V_4(\phi)$$



 y_s = finite distance location of the curvature singularity where the spacetime ends in the y coordinates

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- The superpotential (w/ relation $V = 3W'^2 - 12W^2$) leading to the desired 5D (fully includes the backreaction of the metric to the presence of the scalar field) background : $W = k(1 + e^{\phi})$



In the modeling generalized free continuum by Warped 5D model

Solution
$$A(y) = -\log\left(1 - \frac{y}{y_s}\right) + ky,$$

 y_s = finite distance location of the curvature singularity where the spacetime ends in the y $\phi(y) = -\log\left(k(y_s - y)\right),\,$ coordinates

- Scalar gapped continuum: In conformally flat coordinate, Schröding EOM: $\left(-\partial_y^2 + \hat{V}(y)\right)\Psi(p,y) = e^{2A(y)}$ "Schrödinger Eqn" $-\ddot{\psi}+V(z)$ $V(z) = rac{e^{-2ky}}{4y_s^2}$ if $V \rightarrow \mu_0^2 = \text{finite} => C$ $z \rightarrow \infty$

=> continuum begins at:

$$\mu_0^2 =$$

soft wall & continuum

$$\mathcal{L} = \sqrt{g} \left[\frac{1}{2} g^{MN} D_M \Phi^{\dagger} D_N \Phi - V(\Phi) \right]$$
nger form of eom:

$$\Psi(p,y) = e^{-2A(y)}\Phi(p,y)$$

$$z)\psi = p^{2}\psi$$

$$\frac{ky}{2} \left[4m^{2}(y_{s} - y)^{2} + 15(1 + k(y_{s} - y))^{2} - 6\right]$$
Continuum!
$$ID QM$$
problem
$$\frac{9}{4y_{s}^{2}}e^{-2ky_{s}}$$

$$V^{(z)}$$

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 $v^{0} p^{2} \Psi(p, y)$
 $\psi^{0} \psi^{0} \psi^{0} \psi^{0} \psi^{0} \psi^{0} \psi^{0} \psi^{0} \psi^{0} \psi^{0$

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What's so new for gapped continuum as a DM?

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One may not simply plug gapped continuum into formalism developed for particle DM: need a new theoretical framework for dealing with gapped continuum in order to calculate the relic density of DM, and to deal with the finite temperature physics necessary for describing general features of cosmological history of DM

-requires a systematic development of theories of gapped continuum DM

What's so new for gapped continuum as a DM?

- signatures in colliders and cosmic microwave background measurements

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Gapped Continuum as a DM can give striking new experimental

What's so new for gapped continuum as a DM?

- signatures in colliders and cosmic microwave background measurements
- version of WIMP models).

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Gapped Continuum as a DM can give striking new experimental

• The strong suppression of direct detection signals (will show later) reopens the possibility of a Z-mediated dark sector again (and also other continuum

Theories of DM ?



Theories of DM ?



Theories of DM ?



THE WIMP MIRACLE Insensitive to the initial conditions of the Universe:

due to the thermal equilibrium between the DM and SM gases in the early Universe

Relic abundance \propto

Correct relic abundance for dark matter mass around the TeV scale and weak-force interactions



WIMP Dark Matter

 Original idea of WIMP Miracle



WIMP Dark Matter

- Original idea of WIMP Miracle
- => now pushed to a conner by the null results from DM direct detection experiments

Moore's Law works in DM!





WIMP Dark Matter

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Moore's Law works in DM!

 Z boson exchange excluded except for finetuned corners of parameter space, and requiring tuning for Higgs
 Total 10-43 10-44 10-45 10-45 10-45 10-45 10-45 10-46 10-45 10-46 10-48 10-46 10-48 10-48 10-48 10-48 10-48 10-49 10-48 10-49 10-48 10-49 10-48 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49 10-49





Is there another DM paradigm that gives qualitatively different signatures, but still provide the same level of simple, elegance and compelling explanation as WIMP?

Been searching for WIMPs...

The dominant paradigm is being challenged.



Freeze-Out of Gapped Continuum DM



annihilation: $DM+DM \leftrightarrow SM+SM$

quasi-elastic scattering (QES): DM+SM ↔ DM+SM

Freeze-Out of Gapped Continuum DM



Freeze-Out of Gapped Continuum DM


Freeze-Out of Gapped Continuum DM



Freeze-Out of Gapped Continuum DM



Gapped Continuum Z-portal DM

- ◆ Z-portal Model (with Z₂ symmetry) Consider a complex scalar field Φ with no SM gauge quantum numbers (this plays the role of DM field, and is lifted to 5D)
 - Add another complex scalar field χ which is a doublet under SU(2)_L and carries $U(1)_{\rm Y}$ charge -1/2

components of χ . The mass eigenstates are

- $\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_{\Phi} + \mathcal{L}_{\chi} + \mathcal{L}_{int}$ includes couplings to the SM Z and $U(1)_{Y}$ $\mathcal{L}_{\Phi} = \Phi^{\dagger}(p)\Sigma(p^2)\Phi(p)$ $\mathcal{L}_{\chi} = (D_{\mu}\chi)^{\dagger} (D^{\mu}\chi) - m_{\gamma}^{2}\chi^{\dagger}\chi$ $\mathcal{L}_{int} = -\lambda \Phi \chi H + c.c.$ $\rho(p^2) = \frac{1}{\pi} \mathrm{Im} \Sigma^{-1}(p^2)$ spectral density:
- When the Higgs gets a vev, L_{int} -term induces mass mixing between Φ and the neutral
 - $\tilde{\Phi} = \cos \alpha \, \Phi + \sin \alpha \, \chi^0, \qquad \tilde{\chi}^0 = -\sin \alpha \, \Phi + \cos \alpha \, \chi^0.$

Gapped Continuum Z-portal DM



The mixing angle is given by



$$\overline{\Phi} g'^2 \sin^2 lpha \left(ilde{\Phi}_2 \partial_\mu ilde{\Phi}_1 - ilde{\Phi}_1 \partial_\mu ilde{\Phi}_2
ight) Z^\mu$$

$$an 2lpha \ = \ rac{2\lambda v}{m_{\phi}^2 - m_{\chi}^2}$$



Gapped Continuum Z-portal DM



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♦ Late decay ● decay within the continuum state: $DM(\mu_1) \rightarrow DM(\mu_2) + SM$

- Since all continuum states carry the same quantum number, such decays will necessarily occur continuously throughout the history of the universe.

Late decay decay within the continuum state: $DM(\mu_1) \rightarrow DM(\mu_2) + SM$

continuously throughout the history of the universe.

- In the early universe: DM in thermal and chemical equilibrium with the SM
- As temperature drops below the gap scale μ_0 , DM decouples from the SM and the total number of DM states is frozen out, just like for the usual thermal-relic particle DM

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- In the early universe: DM in thermal and chemical equilibrium with the SM
- As temperature drops below the gap scale μ_0 , DM decouples from the SM and the total number of DM states is frozen out, just like for the usual thermal-relic particle DM
- However, the mass distribution of the DM states continues to evolve, thanks to the above decays
- The decays shift the distribution towards lower masses, closer to the gap scale.

Since all continuum states carry the same quantum number, such decays will necessarily occur

♦ Late decay decay within the continuum state: $DM(\mu_1) \rightarrow DM(\mu_2) + SM$



$$n = \int \frac{d\mu^2}{2\pi} \int \frac{d^3p}{(2\pi)^3} \rho(\mu^2) e^{-\beta E_{\mu}}$$
$$\rightarrow \mu^2$$

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Late decay decay within the continuum state: $DM(\mu_1) \rightarrow DM(\mu_2) + SM$



- the fact that there are fewer states for it to decay into.

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$$\rightarrow \mu^2$$

The decays shift the distribution towards lower masses, closer to the gap scale.

Lifetime of a DM state increases with decreasing mass, due to both phase-space suppression and

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e.g. in our model, DM states are currently clustered within a few hundred keV above the gap scale (on average, each DM state undergoes roughly one decay per Hubble time)

_ S









 $\xi = \log(t/t_d)$, where $\tau = \Gamma_0 t$ and t_d is the time at decoupling Δ=μ- μ₀



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To avoid this, late decay (to e^+e^-) needs to end by T>T_{CMB} => gives a lower bound on the effective coupling

Late decay (T > T_{CMB}) 10^{0} $e^{+}(v)$ 107 $\mathbf{t}_{\mathrm{toc}}$ 2 e- (v) 10⁻¹ μι = µ2 105 2 10⁻² \sin^2 potentially observable effect cosmology. If the SM particl produced in the decay intera electromagnetically (i.e. all § 10⁻³ particles except neutrinos), decays that occur after CM decoupling can reionize hydrc drastically changing the opti 10⁻⁴ depth for CMB photons

> To avoid this, late decay (to e⁺e⁻) needs to end by T>T_{CMB} => gives a **lower bound** on the effective coupling



Direct detection

- even after freeze out, distribution of DM state keeps evolving: distribution is peaked at the mass gap (μ_0) at very late time (these decays are important for CMB physics), and DM states pass through the earth with non-relativistic speed ($v \sim 10^{-3}$)
- quasi elastic scattering (QES): $DM(\mu_1) + SM_1 \rightarrow DM(\mu_2) + SM_2$

Direct detection

- even after freeze out, distribution of DM state keeps evolving: distribution is peaked at the mass gap (μ_0) at very late time (these decays are important for CMB physics), and DM states pass through the earth with non-relativistic speed ($v \sim 10^{-3}$)
 - - => If incoming DM state has $\mu_1 = \mu_0 + \Delta$, accessible final continuum modes are in very narrow window μ_2 $\in [\mu_0, \mu_0 + \Delta + Q]$. For weak scale μ_0 , this basically means that the integral appearing in the QES cross section is constrained to a tiny interval in μ , and leads to a significant suppression of the rate

quasi elastic scattering (QES): $DM(\mu_1) + SM_1 \rightarrow DM(\mu_2) + SM_2$

$$\sigma \sim \int \frac{d\mu_2^2}{2\pi} \rho(\mu_2^2) \,\hat{\sigma}(\mu_1,\mu_2)$$

Direct detection Q is	s the kinetic ene
• quasi elastic scatter $\Delta \ll$	< μ_0 in today's untivistic.
 even after freez 	
distribution is p	
decays are impo	e.g. Δ ~ 10
the earth with n	μ_0 a positive number that
=> If incoming	
final continuu	
\in [μ_0 , μ_0 + Δ +Q]	
means that the	
section is con	
leads	s to a sign

nergy of the collision in the center-of-mass frame universe, while $Q \ll \mu_0$ as long as ambient DM is non-

$$\sigma_{\rm cont} \sim \left(\frac{\Delta+Q}{\mu_0}\right)^{1+r} \sigma_{\rm particles}$$

100 keV at the present time, while $Q \sim 1 \text{ keV}$ at the weak scale $\rightarrow \sim 10^9 \text{ suppression}$

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at depends on the behavior of the spectral density near the gap (r=1/2 for XD)

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Ge, Si, NaI, LXe, ...



num Nature of DM



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Indirect Detection

these processes are unsuppressed :

 $DM(\mu_1) + DM(\mu_2) \rightarrow SM_1 + SM_2$

Since there is no continuum state in the final state, the rates of

Indirect Detection

these processes are unsuppressed :

> $\mu_1 \approx \mu_2 \approx \mu_0$ in the current universe \rightarrow both rates and kinematics of annihilation in the galactic halos are basically identical to those of particle DM

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Since there is no continuum state in the final state, the rates of

Colliders Phenomenology

bound for Z-portal WIC:



- for low energy experiments (low compared to gap scale): e.g. LEP
- Same suppression mechanism (by continuum kinematics) as in Direct Detection appl

Colliders Phenomenology

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Colliders Phenomenology



 $SM1+SM2 \rightarrow DM(\mu_1) + DM(\mu_2)$

 $\sigma \sim \int \frac{d\mu_1^2}{2\pi} \rho(\mu_1^2) \int \frac{d\mu_2^2}{2\pi} \rho(\mu_2^2) \hat{\sigma}(\mu_1, \mu_2) \, d\mu_2 \,$

 SM_1





Colliders Phenomenology





for high enough energy: (no suppression, an rich pheno)

Mass Distributions of Cascade Decay



120

130

Dark Matter Continuum Spectral Density from 5D Model

Warped 5D model

- Sca

In cont

alar gapped continuum:
$$\mathcal{L} = \sqrt{g} \left[\frac{1}{2} g^{MN} D_M \Phi^{\dagger} D_N \Phi - V(\Phi) \right]$$

formally flat coordinate, Schrödinger form of eom: $\psi = e^{-\frac{3}{2}A} \Phi$
 $\left(-\partial_z^2 + \hat{V}(z) \right) \Psi(z) = p^2 \Psi(z)$
 $V(z) = \frac{e^{-2ky}}{4y_s^2} \left[4m^2(y_s - y)^2 + 15\left(1 + k(y_s - y)\right)^2 - 6 \right]$

after integrating out bulk:

Solk:

$$S = \int \frac{d^4p}{(2\pi)^4} \, \Phi^{\dagger}(p) \Sigma(p^2) \Phi(p)$$

$$\rho(p^2) = -2 \operatorname{Im} \frac{1}{\Sigma(p^2)}$$

$$\Sigma(p) = \frac{1}{R} \left. \frac{f'(z,p)}{f(R,p)} \right|_{z=R}$$



Dark Matter Continuum Spectral Density from 5D Model

Warped 5D model

- Scalar gapped continuum near the gap:

In conformally flat coordinate, Schrödinger form of eom:

$$\left(-\partial_z^2 + \hat{V}(z) \right) \Psi(z) = p^2 \Psi(z)$$
$$\lim_{z \to \infty} \hat{V}(z) = \mu_0^2 \left(1 + e^{-2z(2\mu_0/3)} + \frac{8}{3} e^{-z(2\mu_0/3)} \right)$$

$$\Psi(z,\mu) = C L_m^n (3e^{-2z\mu_0/3}) \exp\left(\frac{3}{2}\sqrt{1-\frac{\mu^2}{\mu_0^2}}\log\left(e^{-\frac{2\mu_0 z}{3}}\right) - \frac{3}{2}e^{-\frac{2\mu_0 z}{3}}\right)$$

can expand the arguments of the Laguerre polynomial around the mass gap

$$\rho(p) = \frac{1}{\pi} \operatorname{Im} G(R, R; p).$$

$$ho(\mu^2) \propto \left(rac{\mu^2}{\mu_0^2}-1
ight)^{1/2}$$

Dark Matter Continuum Spectral Density from 5D Model



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$$ho(\mu^2) \propto \left(rac{\mu^2}{\mu_0^2}-1
ight)^{1/2}$$

Summary



- 1. Gapped Continuum DM = theoretically and phenomenologically motivating!
- 2. Continuum Kinematics : late decay, relaxation of direct detection bound
- 3. Revival of Weakly Interacting Massive Continuum (WIC) !
- 4. Many possible models + many detailed pheno study to be done.
- 5. Continuum Collider Physics = totally new \rightarrow needs a systematic investigations a
- 6. Many more (including continuum freeze-in DM, etc)

- Weakly Interacting Massive Particle
 - Weakly Interacting Massive Continuum


Back-up

Electromagnetic Energy Injection After Recombination

-important constraints on the model arises from re-ionization of Hydrogen by DM decays after recombinition

-at most ~0.1% H ionization fraction for points on the boundary obtained by the simple criterion $\Gamma = H(t_{CMB})$





Electromagnetic Energy Injection After Recombination

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DM mass distribution as a function of time for a sample parameter point. Only about 0.2% of DM particles lie on the tail above this threshold.



$$S_{\text{eff}} = \int_{\text{UV}} d^4 x \ e^{-4A} \Phi^{\dagger}(x, y) \partial_y \Phi(x, y)|_{y=0}$$
$$= \int_{\text{UV}} \frac{d^4 p}{(2\pi)^4} \ \hat{\Phi}^{\dagger}(p) \left(e^{-3A(z)} \frac{f'(z, p)}{f(R, p)} \right)_{z=R} \hat{\Phi}(p)$$

$$\int d^4x \sqrt{g} (D_{\mu}HD^{\mu}He^{2A} + D_{\mu}\chi D^{\mu}\chi e^{2A} - \hat{\lambda}k^{\frac{1}{2}}\Phi\chi H + h.c.)_{z=R}$$

In order to get the proper 4D effective action with $\lambda = \hat{\lambda} k e^{-A}$ of the order of the electroweak scale we need the field redefinitions $H \to He^{-A}, \chi \to \chi e^{-A}, \Phi \to \Phi e^{-\frac{3}{2}A}\sqrt{R}.$

$$\Sigma(p) = \frac{1}{R} \left. \frac{f'(z,p)}{f(R,p)} \right|_{z=R},$$

Normalization of Spectral Density

