

Humboldt Kolleg, Kitzbühel 2022

Gravitational Quantum Physics:

*how to avoid the appearance of the
classical world in gravity experiments?*

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arxiv:2203.05587

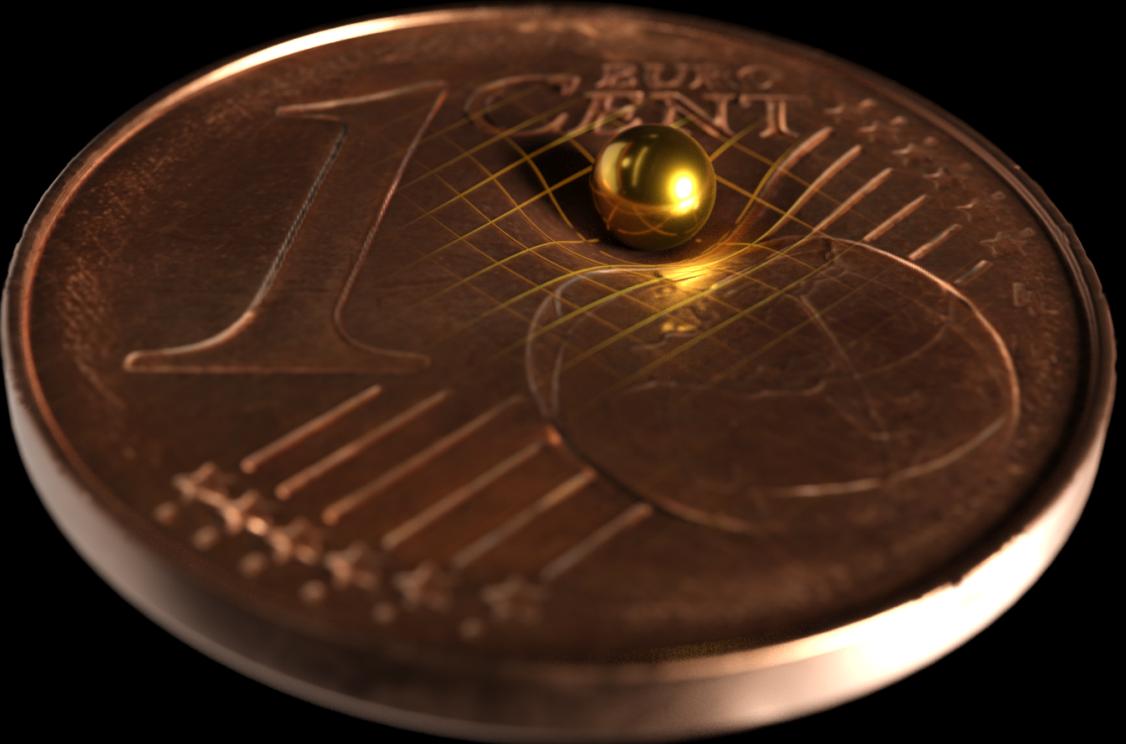


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Gravitational Quantum Physics: *how to avoid the appearance of the classical world in gravity experiments?*

Peter Asenbaum

Talk on Friday:
Gravity in large quantum states and
the Aharonov-Bohm effect



IQOQI, Austrian Academy of Sciences

arxiv:2203.05587

Prelude: The 1957 Chapel Hill Conference

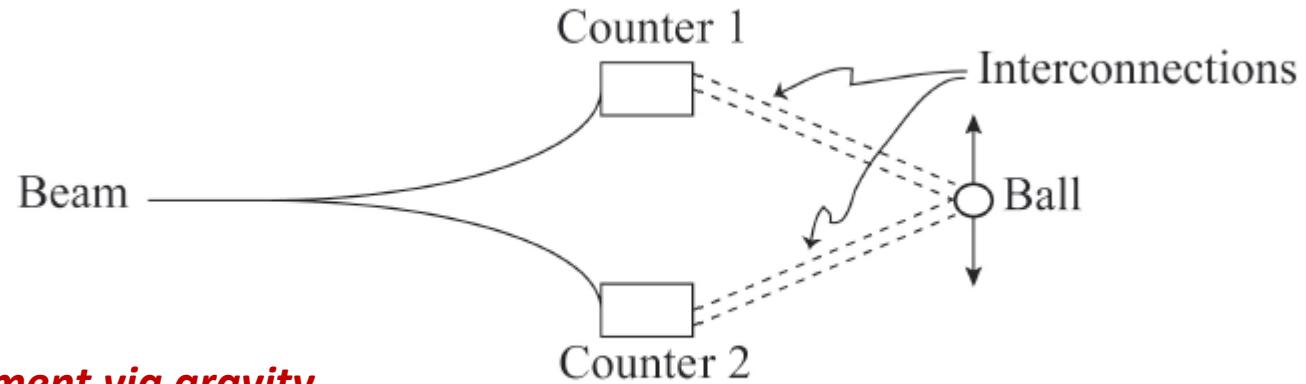


“One should think about designing an experiment which uses a gravitational link and at the same time shows quantum interference”

Chapel Hill Conference 1957

Quantum entanglement via gravity

R P Feynman

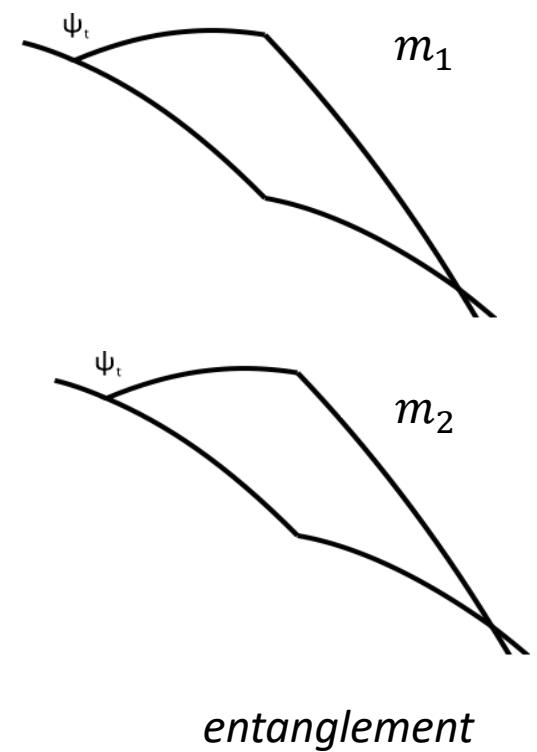
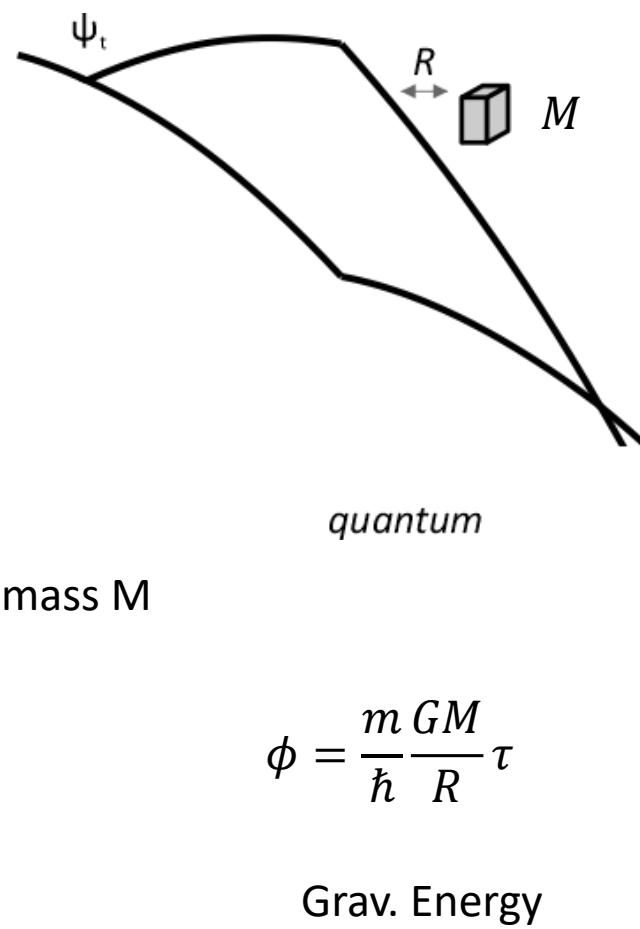
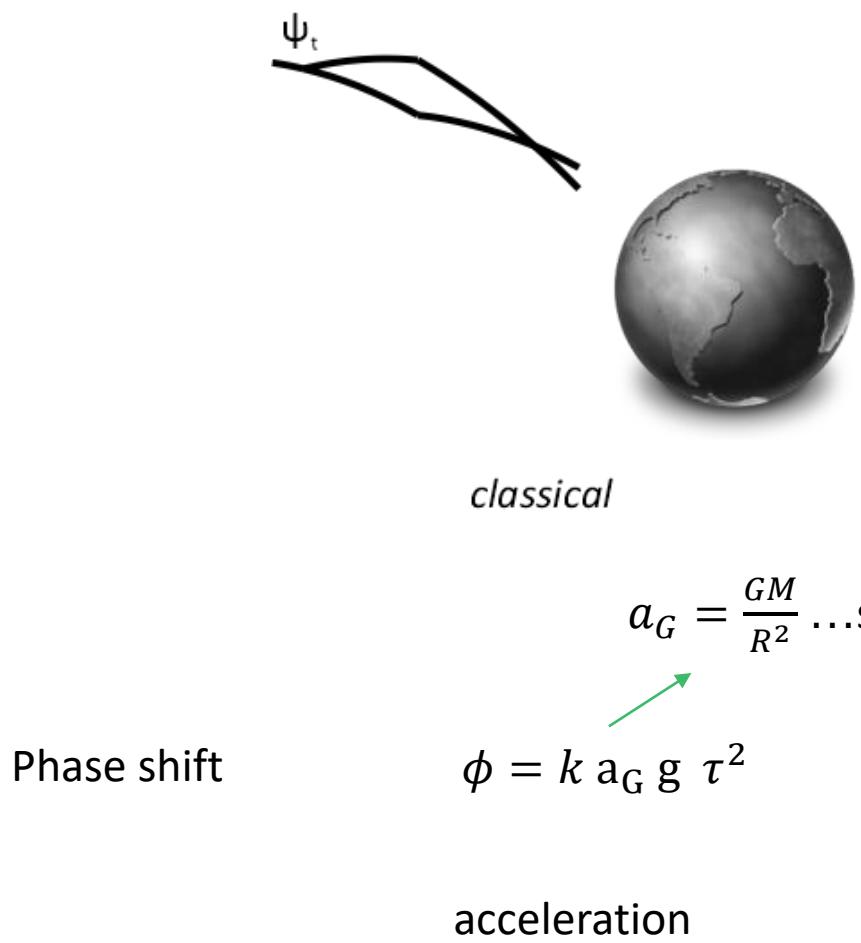


“...that we’re in trouble if we believe in quantum mechanics but don’t quantize gravitational theory”

„But aside from that possibility [that quantum mechanics fails], if you believe in quantum mechanics up to any level then you have to believe in gravitational quantization in order to describe this experiment.”

Here he explicitly refers to the ability to apply the superposition principle, i.e. probability amplitudes, to different gravitational field configurations, and not the necessity to invoke quanta of the field (gravitons).

Gravity in Quantum Systems



For two large quantum states:

$$\phi = \frac{m_1 G m_2}{\hbar R} \tau$$

Observation of Gravitationally Induced Quantum Interference*

R. Colella and A. W. Overhauser

Department of Physics, Purdue University, West Lafayette, Indiana 47907

and

S. A. Werner

Scientific Research Staff, Ford Motor Company, Dearborn, Michigan 48121

(Received 14 April 1975)

We have used a neutron interferometer to observe the quantum-mechanical phase shift of neutrons caused by their interaction with Earth's gravitational field.

Atom-Interferometric Test of the Equivalence Principle at the 10^{-12} Level

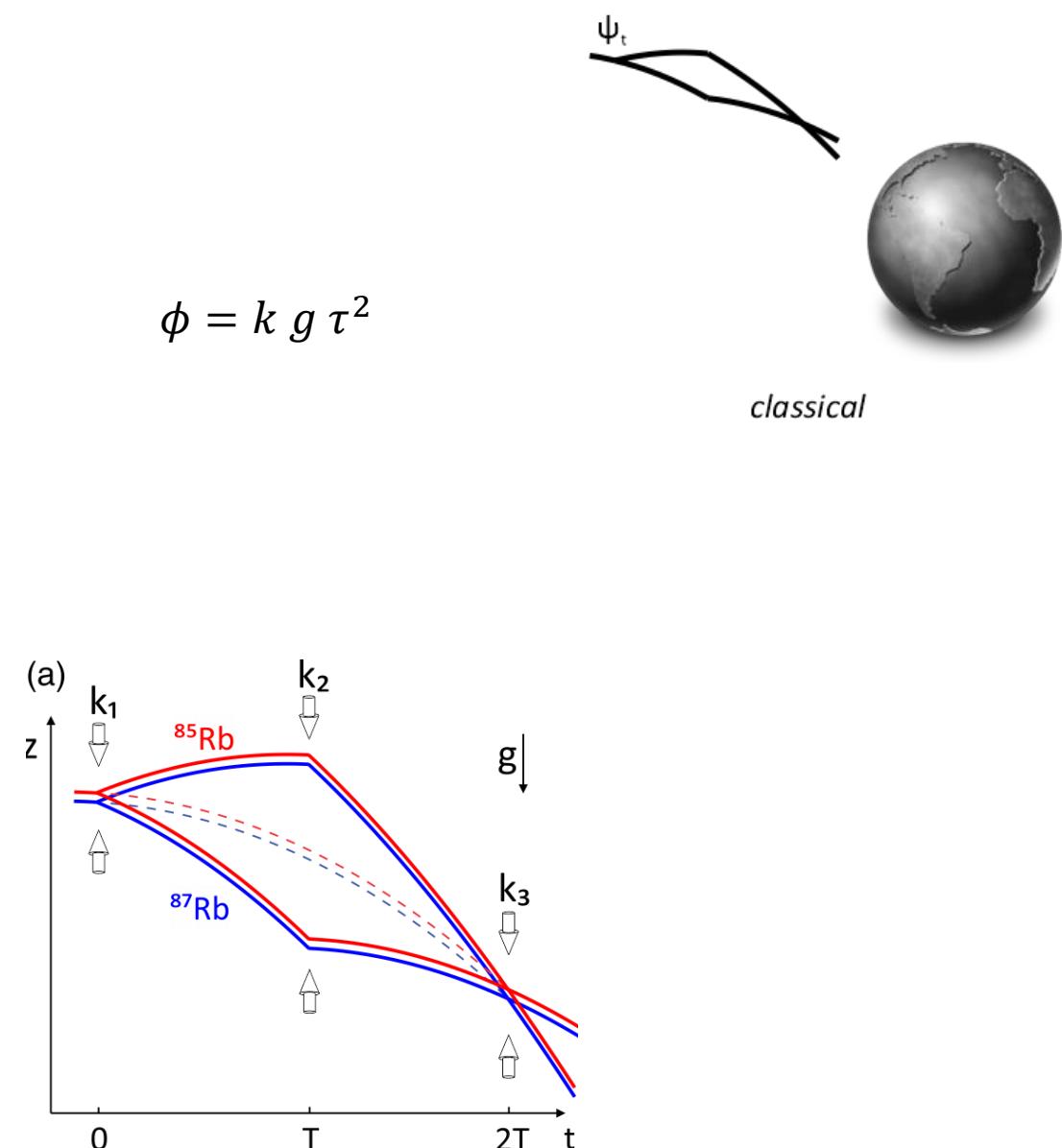
Peter Asenbaum[✉], Chris Overstreet[✉], Minjeong Kim[✉], Joseph Curti, and Mark A. Kasevich[†]

Department of Physics, Stanford University, Stanford, California 94305, USA

(Received 26 June 2020; accepted 5 October 2020; published 2 November 2020)

We use a dual-species atom interferometer with 2 s of free-fall time to measure the relative acceleration between ^{85}Rb and ^{87}Rb wave packets in the Earth's gravitational field. Systematic errors arising from kinematic differences between the isotopes are suppressed by calibrating the angles and frequencies of the interferometry beams. We find an Eötvös parameter of $\eta = [1.6 \pm 1.8(\text{stat}) \pm 3.4(\text{syst})] \times 10^{-12}$, consistent with zero violation of the equivalence principle. With a resolution of up to $1.4 \times 10^{-11} \text{ g}$ per shot, we demonstrate a sensitivity to η of $5.4 \times 10^{-11}/\sqrt{\text{Hz}}$.

No mass dependence! QM satisfies Equivalence Principle!

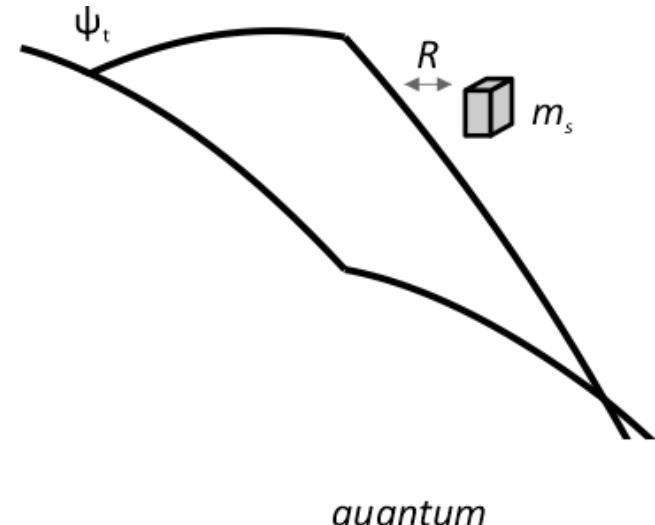


PHYSICS

Observation of a gravitational Aharonov-Bohm effect

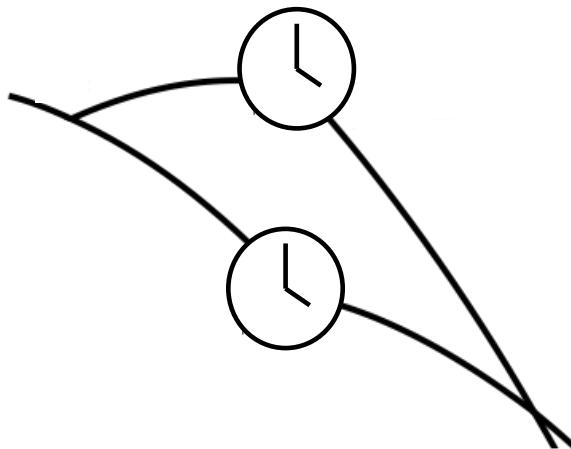
Chris Overstreet^{1†}, Peter Asenbaum^{1,2†}, Joseph Curti¹, Minjeong Kim¹, Mark A. Kasevich^{1*}

Gravity curves space and time. This can lead to proper time differences between freely falling, nonlocal trajectories. A spatial superposition of a massive particle is predicted to be sensitive to this effect. We measure the gravitational phase shift induced in a matter-wave interferometer by a kilogram-scale source mass close to one of the wave packets. Deflections of each interferometer arm due to the source mass are independently measured. The phase shift deviates from the deflection-induced phase contribution, as predicted by quantum mechanics. In addition, the observed scaling of the phase shift is consistent with Heisenberg's error-disturbance relation. These results show that gravity creates Aharonov-Bohm phase shifts analogous to those produced by electromagnetic interactions.



$$\phi = \frac{m}{\hbar} \frac{GM}{R} \tau$$

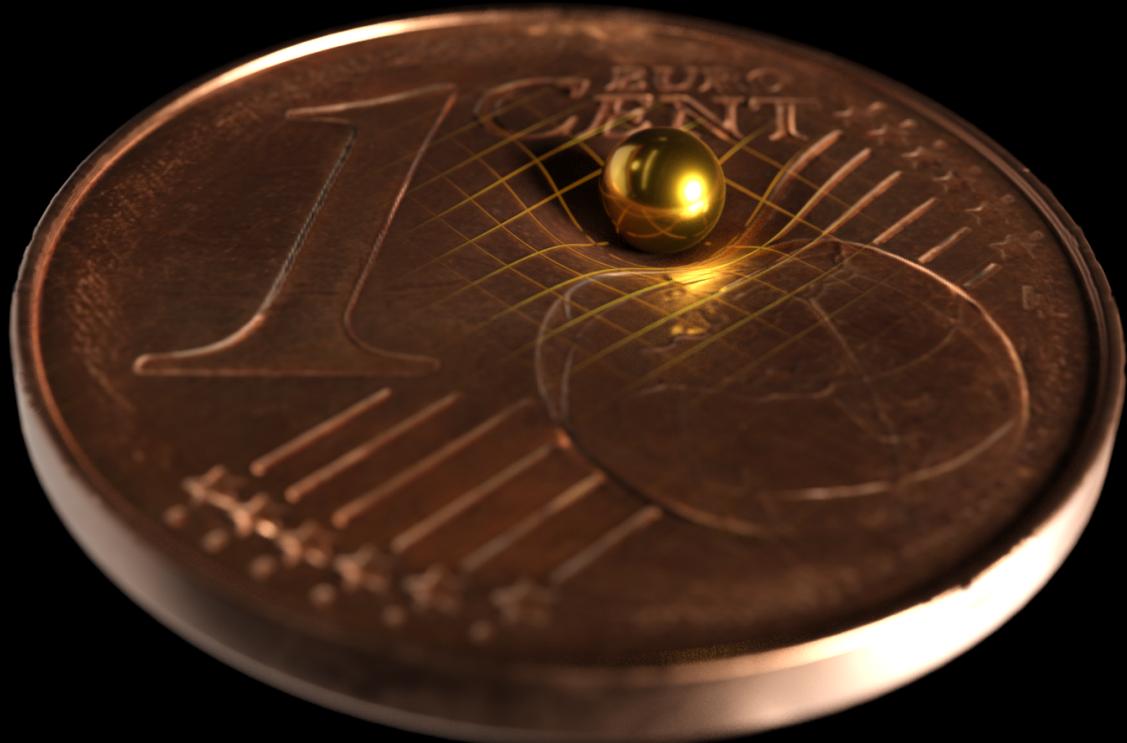
Two clocks measure hybrid
between acceleration and
energy



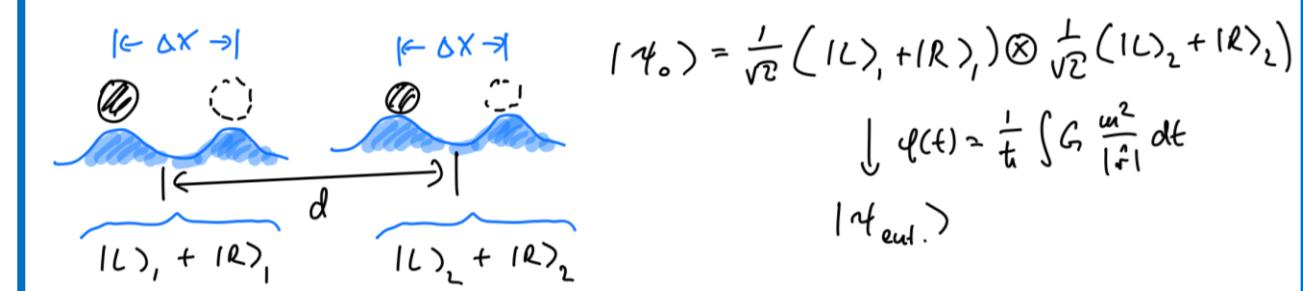
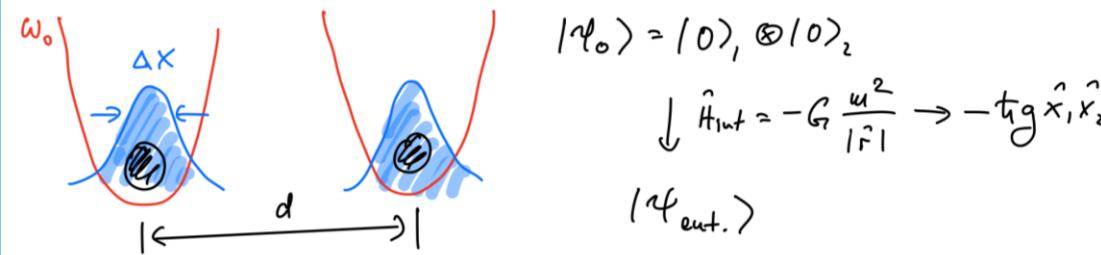
not high-sensitivity $\frac{\Delta m}{m} = 10^{-11}$

What about quantum systems as gravitational SOURCE masses?

- Source vs test mass
- In gravitational entanglement regime:
each particle is sensor and source



This rather be solid...



Al Balushi et al., PRA 98, 043811(2018),
 Krisnanda et al., npj Quantum
 Information 6, 12 (2020),
 Cosco et al., PRA 103, L061501 (2021)
 Weiss et al., PRL 127, 023601 (2021)

Bose et al., PRL 119,
 240401 (2017),
 Marletto et al., PRL
 119, 240402 (2017)

ENTANGLEMENT RATE $\dot{g} = \frac{G}{t} \frac{\mu^2}{d} \left(\frac{\Delta x}{d}\right)^2 > \Gamma_{\text{decoherence}}$

$$\left. \begin{aligned} d &\sim 10^{-6} \text{ m} \\ \frac{\Delta x}{d} &\sim 10^{-1} \end{aligned} \right\} \dot{g} \approx 2\pi \cdot 10^{23} \cdot 10^6 \cdot 10^{-2} \frac{\mu^2}{\text{m}^2}$$

$$\approx 6 \cdot 10^{27} \cdot \mu^2 \geq \Theta(1) \leftrightarrow \mu^2 \geq (1.2 \cdot 10^{-14})^2$$

$$\rightarrow g \approx \frac{\mu}{d^3} \geq \frac{10^{-14}}{(10^{-6})^3} \sim 10^4 \frac{\text{kg}}{\text{m}^3} \equiv \text{SOLID-STATE}$$

Louis Witten: "What prevents this from becoming a practical experiment?"

The decoherence challenge...

see also

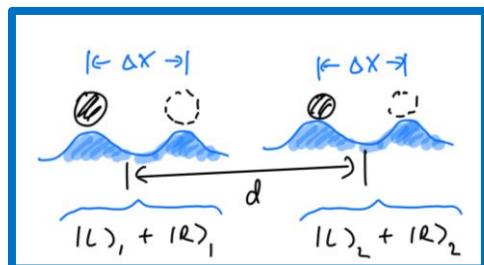
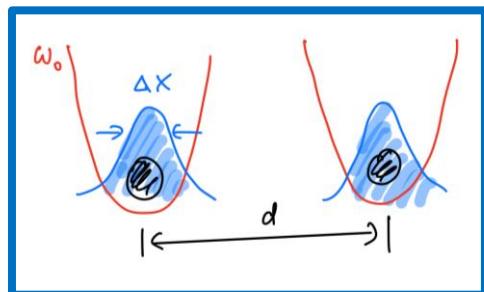
O. Romero-Isart et al., PRL 107, 020405 (2011)

O. Romero-Isart, PRA 84, 052121 (2011)

S. Rijavec et al., New J. Phys. 23, 043040 (2021)

T. Weiss et al., PRL 127, 023601 (2021)

Examples:



	ΔX	squeezing	T_i	T_e
75 nm glass sphere	$> 2 \mu\text{m}$	$> 60 \text{ dB}$	4K	4K
Planch mass Pb sphere	$> 12 \mu\text{m}$	$> 60 \text{ dB}$	8K	6K
LIGO ($\theta(10\text{cm})$)	$> 1 \mu\text{m}$	$> 60 \text{ dB}$	4K	100 mK

for $\bar{n} \sim 0.1$, $\rho < 10^{-17}$ under

assuming GAS SCATTERING & BLACK BODY RADIATION

Decoherence & the appearance of a classical

Joos & Zeh, Caldeira & Leggett, Unruh & Zurek
Paz & Zurek, Hu & Paz & Zhang, Milburn, ...

$$\dot{\rho} = -\frac{i}{\hbar}[H, \rho]$$

DECOHERENCE

$$(i) \text{ GAS DAMPING } \Gamma_g \approx 2 \cdot 10^{26} \rho R^2 T_e^{-1/2}$$

$$(ii) \text{ BLACK BODY SCATTERING } \Gamma_{g,s} \approx 3 \cdot 10^{36} R^6 T_e^9 \Delta x^2$$

$$\text{ABSORPTION } \Gamma_{b,s} \approx 7 \cdot 10^{25} R^3 T_e^6 \Delta x^2$$

$$\text{EMISSION } \Gamma_{b,e} \approx 7 \cdot 10^{25} R^3 T_i^6 \Delta x^2$$

$$(iii) \text{ TRAP FLUCTUATIONS } \Gamma_x = \frac{\pi \omega_0^2}{4 \sigma_*^2} \zeta_x(\omega_0)$$

$$(iv) \text{ FREQUENCY FLUCTUATIONS } \Gamma_\nu = \frac{\pi \omega_0^2}{16} \zeta_\nu(2\omega_0)$$

$$(v) \text{ INTERNAL DECOHERENCE } \Gamma_k = \frac{Q_k T_e}{\hbar \omega_0} \gamma \approx 10^{11} \cdot \frac{T_e}{\omega_0} \cdot \gamma$$

BB emission
atmosphere

$$\Gamma_{\text{gas}} = \frac{\lambda_k}{t_c} \frac{16\pi}{3} \rho R^2$$

$$\frac{8! \cdot 8 \zeta(9) c R^6}{9\pi} \left(\frac{Q_k T_e}{t_c} \right)^9 \text{Re} \left\{ \frac{\epsilon - 1}{\epsilon + 2} \right\}^2$$

$$\lambda_{e(a)} = \frac{16\pi^5 c R^3}{189} \left(\frac{Q_k T_e(c)}{t_c} \right)^6 \text{Im} \left\{ \frac{\epsilon - 1}{\epsilon + 2} \right\}^2$$

Decoherence from

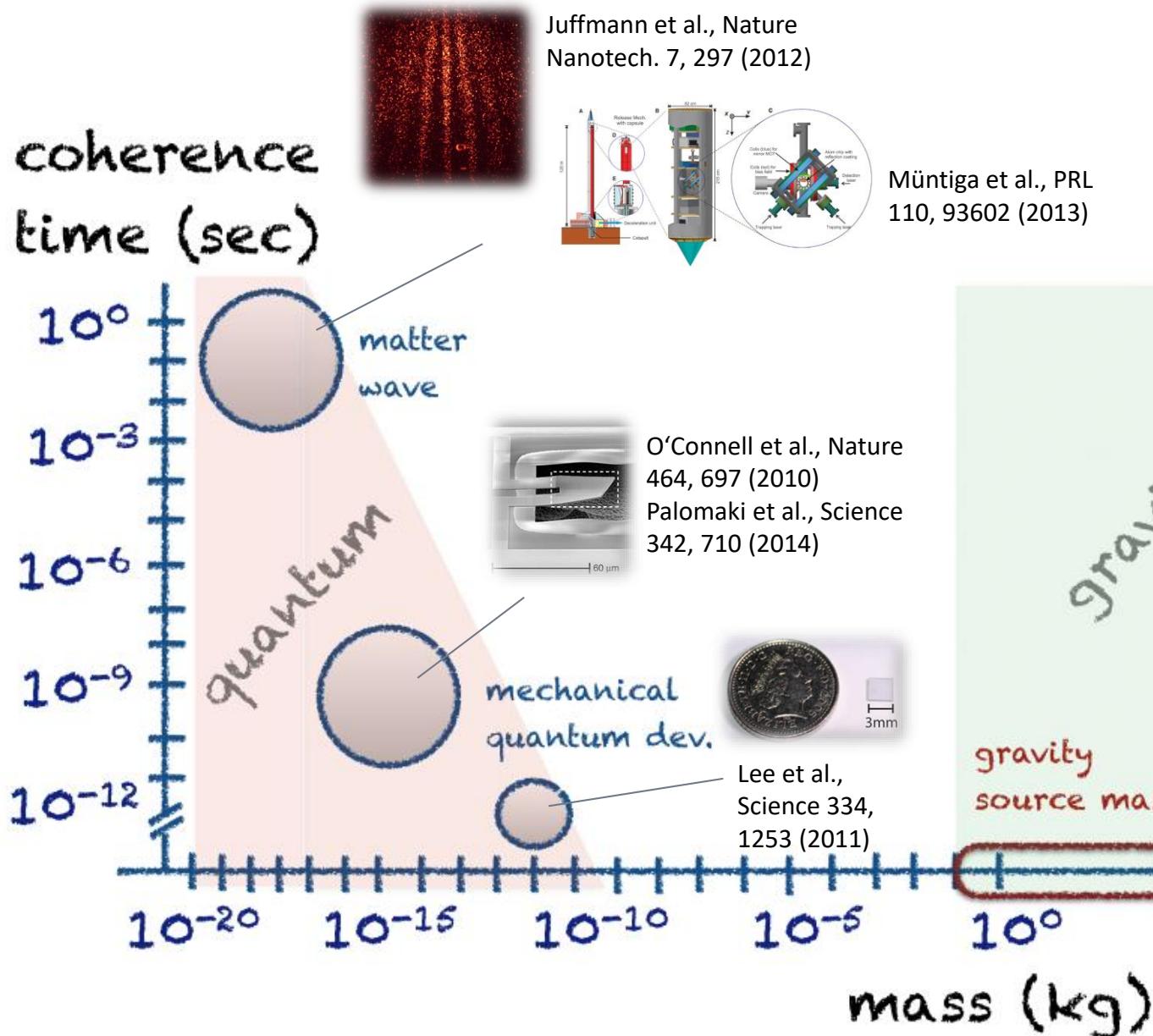
$$\lambda_k = 2\pi t_c (\sqrt{2r})$$

Decoherence from

$$\lambda_k = \frac{\pi^2 \beta t_c}{Q_k T_e} \sim \delta C$$

Thermal decoherence

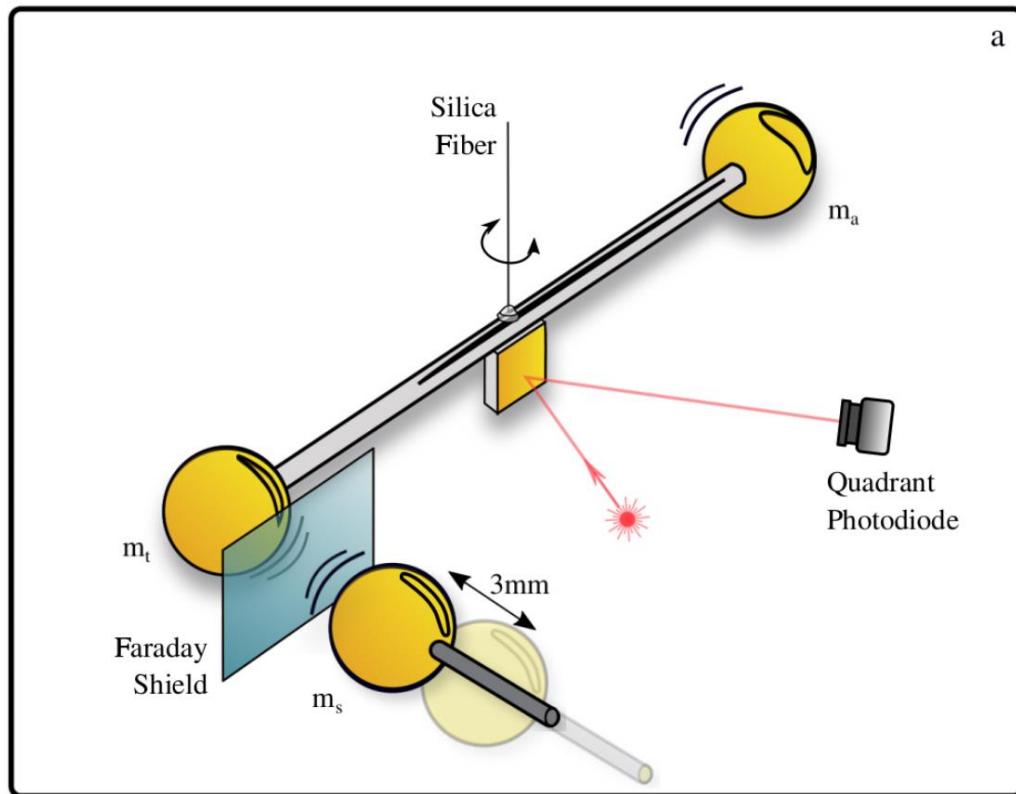
Reality check: quantum systems as gravitational source masses?



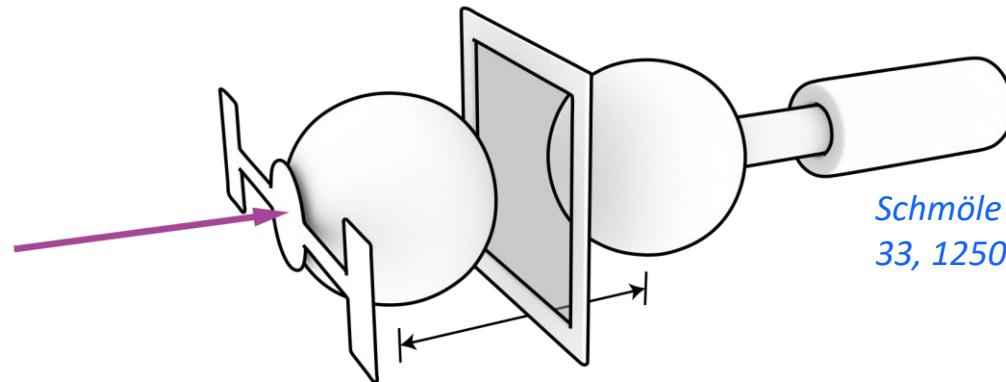
- How small can we make a source mass?
- How massive can we make a quantum system?

1. Go small!
2. Go coherent!

How small can we go? The idea:



Westphal et al., Nature 591, 225 (2021)



Schmöle et al., Class. Quant. Grav. 33, 125031 (2016)



- periodic modulation of a source mass @ f_{mod}
- generates a test mass acceleration @ $n \times f_{\text{mod}}$ ($n=1, 2, \dots$)
- fundamental limit: thermal noise of test mass oscillator

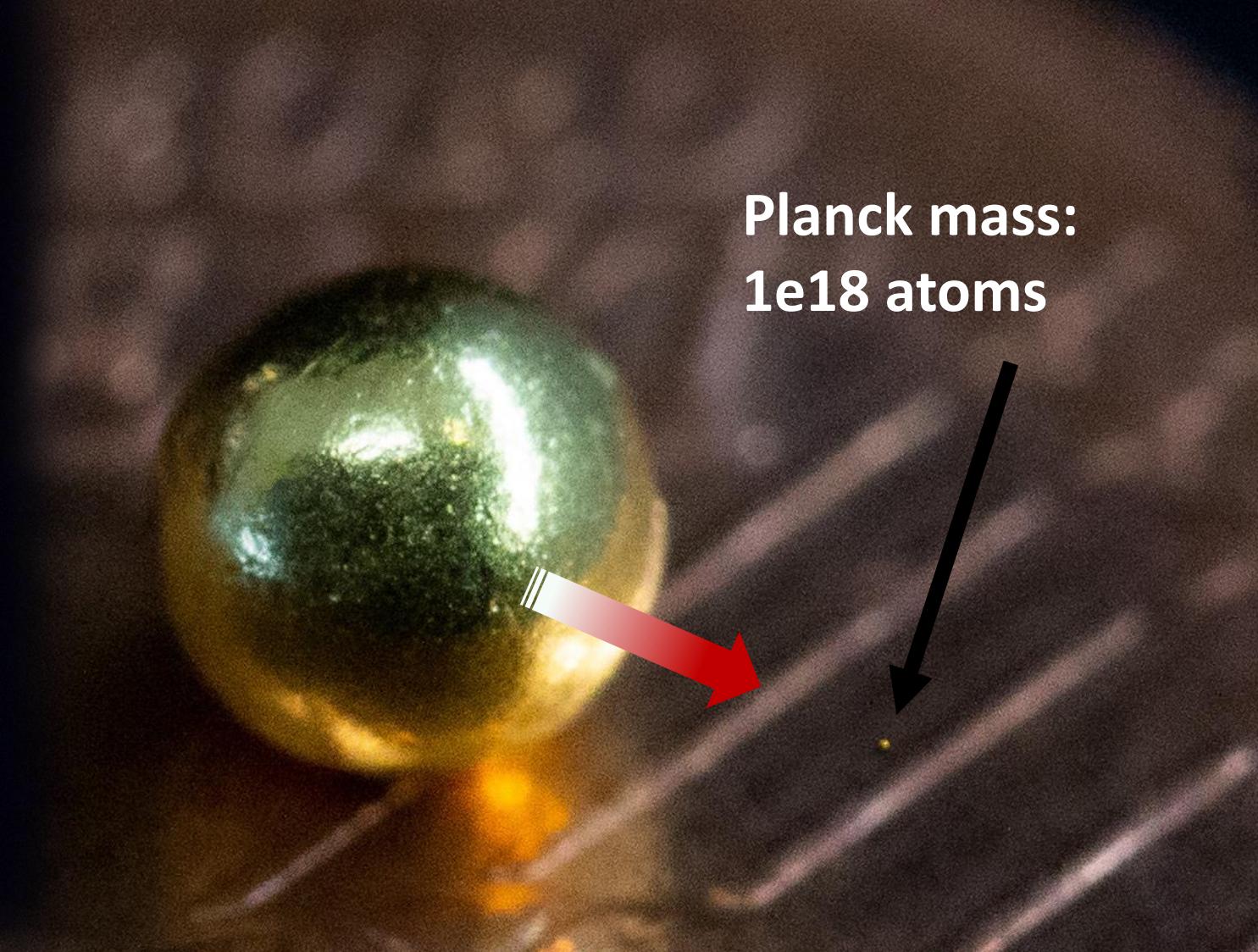
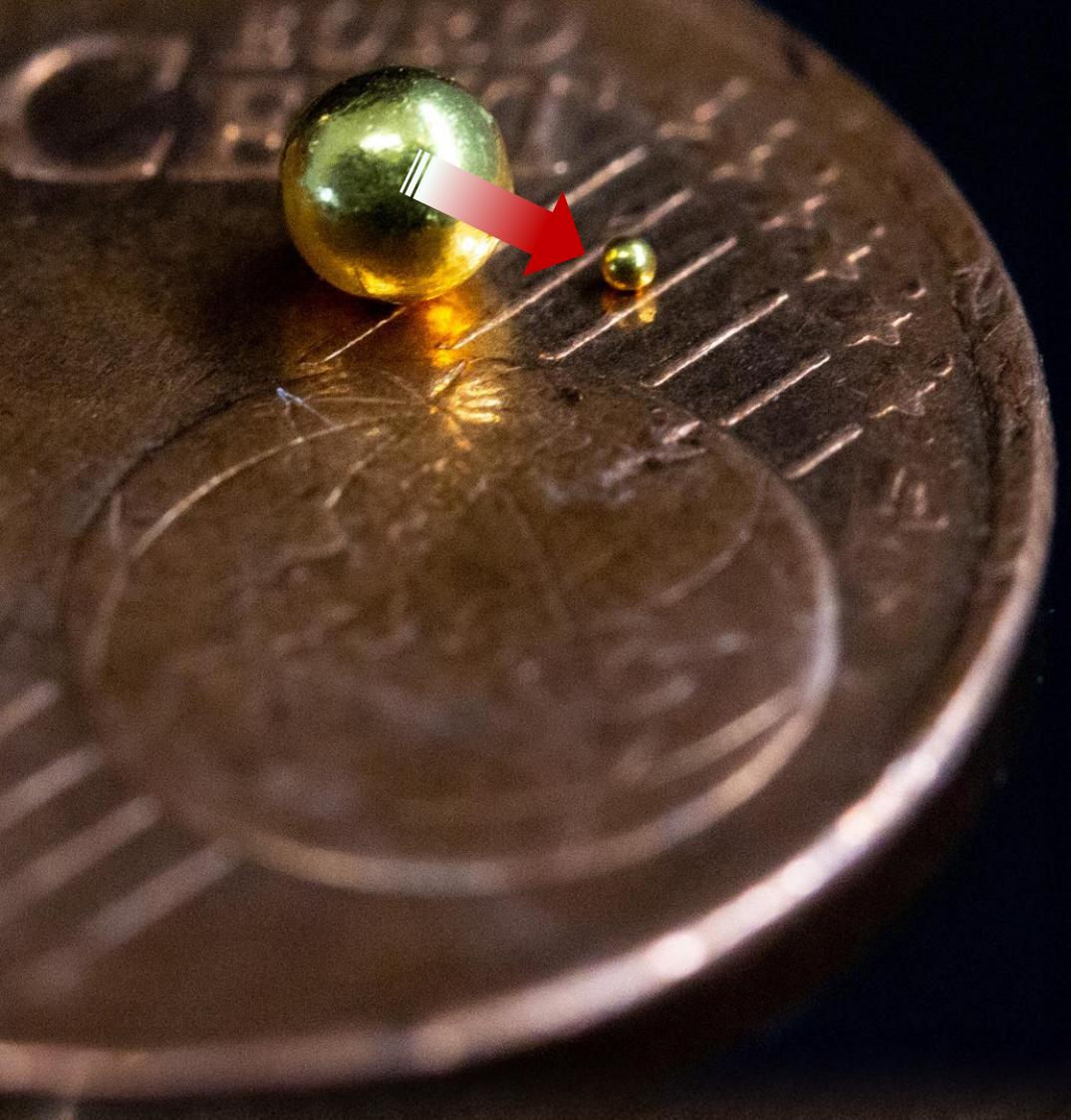
$$d(t) = d_0 + d_m \omega_s \sqrt{2} t$$

$$a(t) = G \frac{m}{d(t)^2} = G \frac{m}{d_0^2} \left(1 - 2 \frac{d_m}{d_0} \omega_s \sqrt{2} t + 3 \left(\frac{d_m}{d_0} \right)^2 \omega_s^2 t^2 + \dots \right)$$

$$\left. \begin{aligned} \omega_s &\sim 9 \cdot 10^{-5} \text{ rad/s} \\ d_0 &\sim 4 \mu\text{m} \\ d_m &\sim 1.6 \mu\text{m} \end{aligned} \right\} a \sim 3 \cdot 10^{-10} \frac{\text{m}}{\text{s}^2}$$

M=100 mg

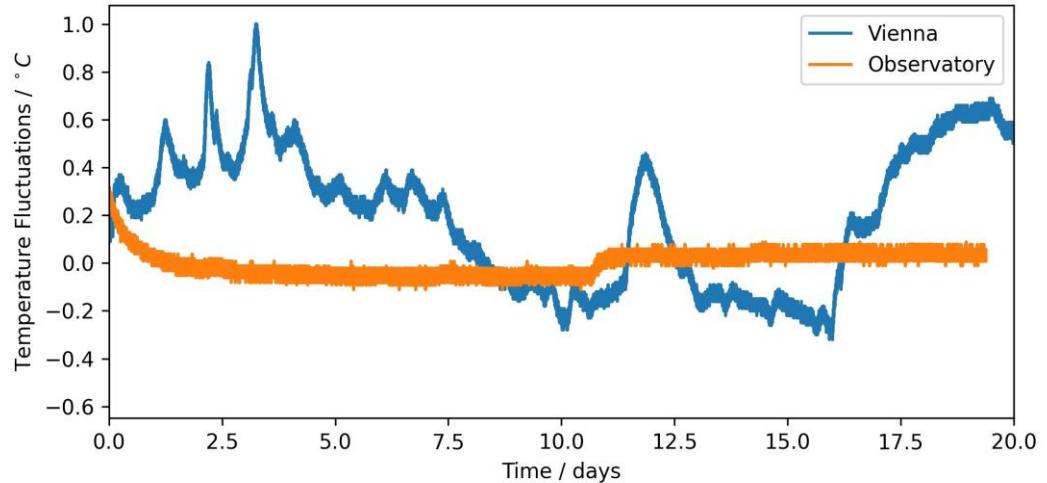
Next steps: going smaller in mass...



Planck mass:
1e18 atoms

Improved noise performance...

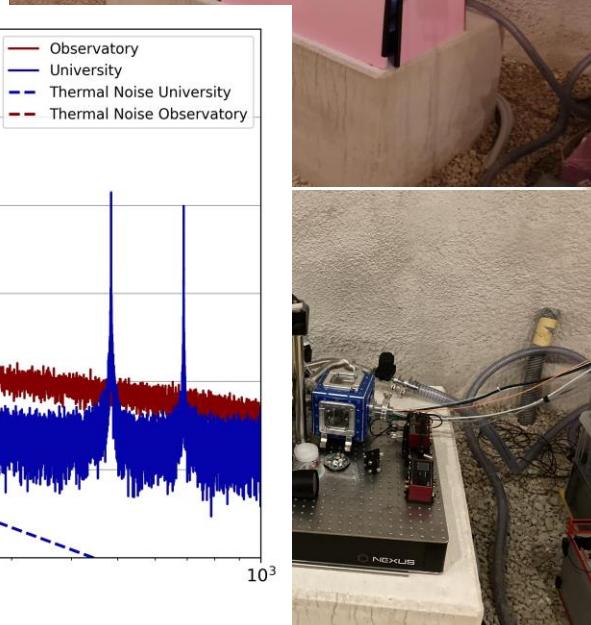
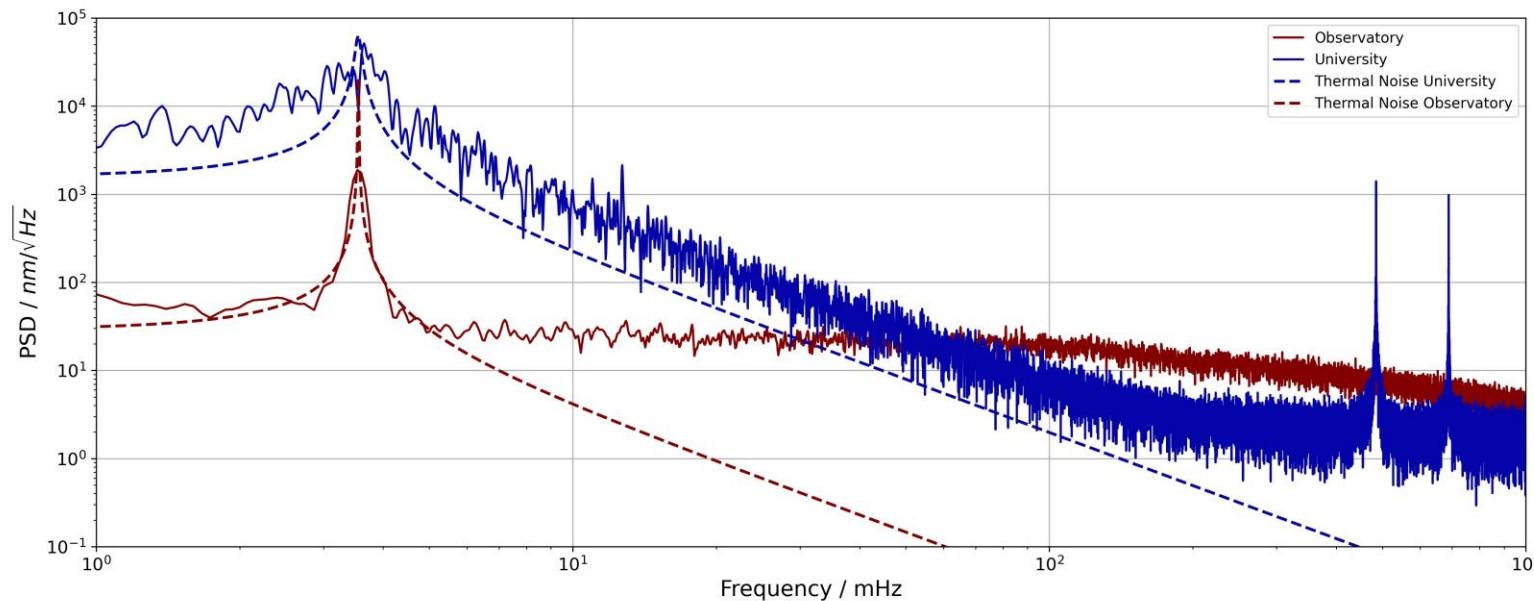
Temperature



City

Mine

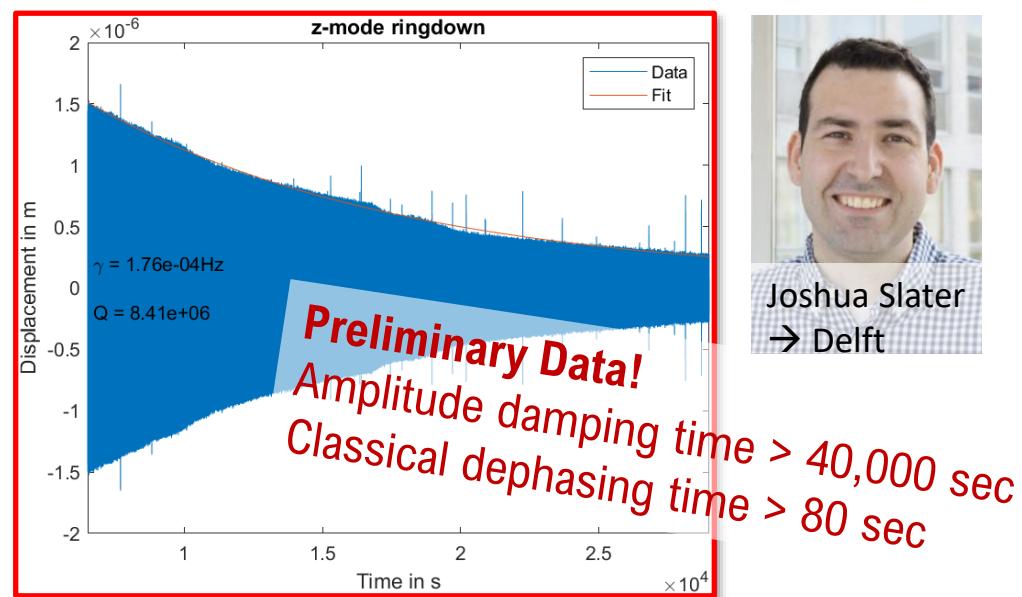
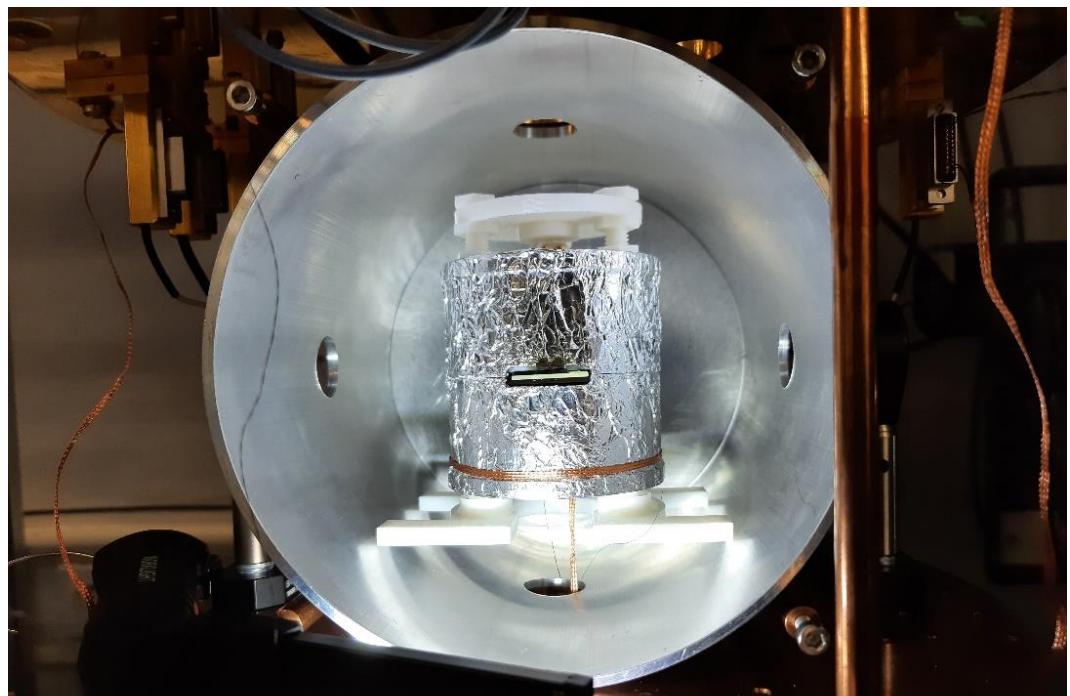
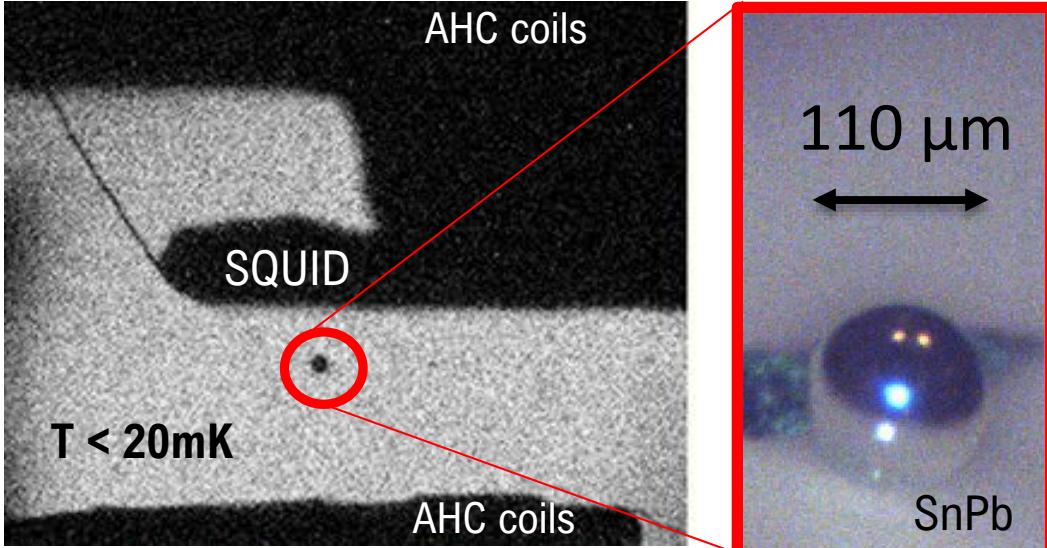
Displacement sensing



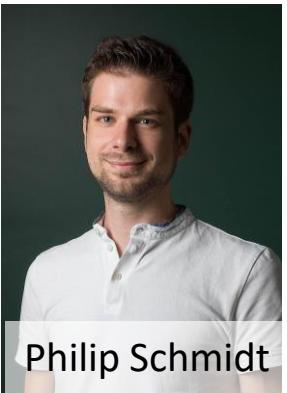
Conrad Observatory
Trafulberg, Austria



Go coherent: superconducting levitation



Gerard Higgins



Philip Schmidt



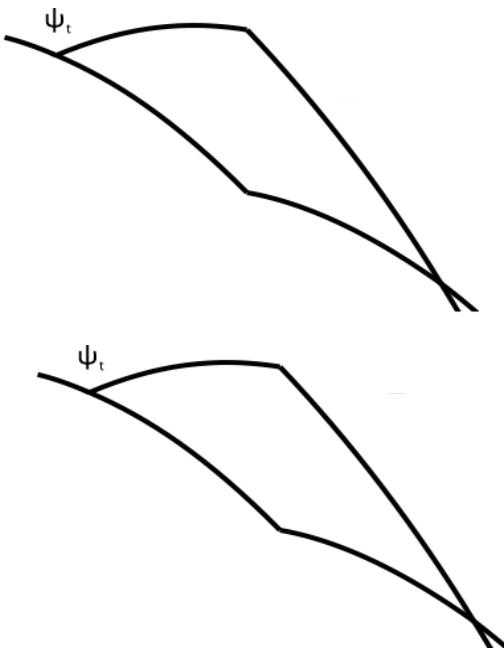
Witlef Wieczorek
→ Chalmers



Michael Trupke

Summary

Goal:



Grav. entanglement

Exp. parameters:

- Mass: 10^{-10} kg
- Pressure: 10^{-17} mbar
- Temperature: 1 K
- Sensitivity: $10^{-17}g$
- Systematics:....

Implications:

- Gravity plays by quantum rules
 - Grav. field clearly non-classical
- BUT:
- Field can be ignored to predict experiment
 - Superposition of field cannot be observed independently



Quantum-“Mechanics” in Vienna: The Levitation Team 2022

+ our collaboration partners:

The ERC Synergy team: Lukas Novotny, Romain Quidant (ETH) / Oriol Romero-Isart (Innsbruck)
Eric Adelberger (UWash) / Caslav Brukner (Vienna) / Rudolf Gross (WMI) / Andreas Kugi (TU Wien) / Nikolai Kiesel (Vienna) /
Monika Ritsch-Marte (Innsbruck) / Vladan Vuletic (MIT) / Robert Wald (UChicago) / Witold Wieczorek (Chalmers)



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**OPTOMECHANICAL
TECHNOLOGIES**