

**Clues to a mysterious Universe
— exploring the interface of particle,
gravity and quantum physics**

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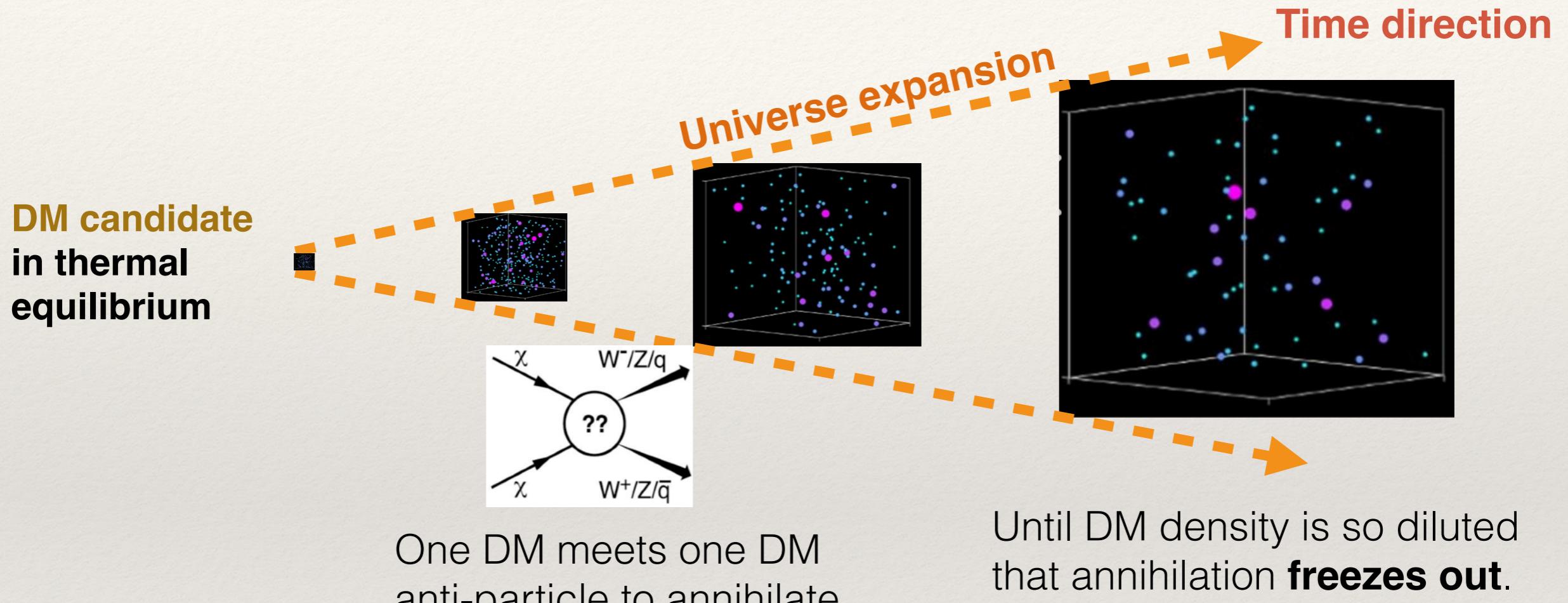
Towards a full description of MeV dark matter decoupling

Xiaoyong Chu

In collaboration with Jui-Lin Kuo, Josef Pradler
— arXiv:2205.05714 and follow-up

Standard Cosmology is established, where for Dark Matter (DM):

the conventional scenario: **Thermal freeze-out** (weakly-interacting massive particle)



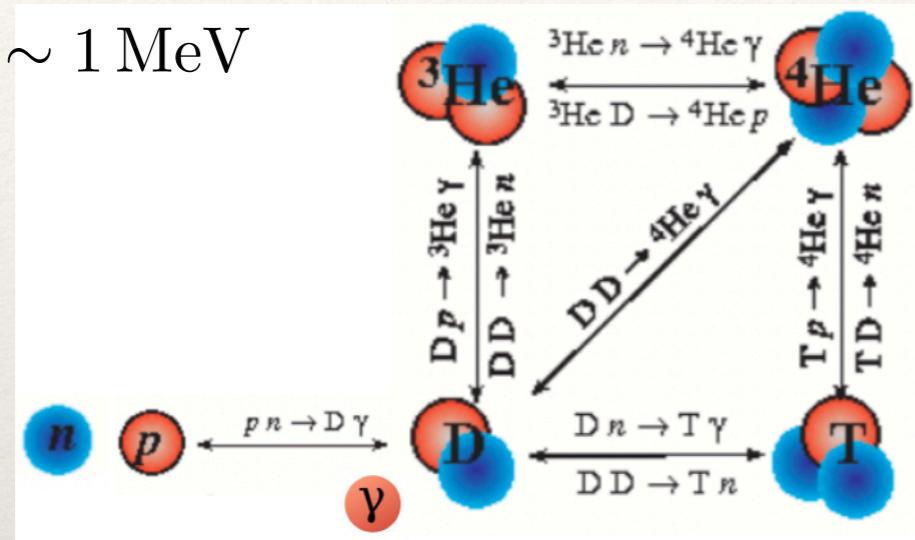
That is:

- DM contributes to the **radiative energy of Universe** while being relativistic;
- DM ejects **energy (visibly or invisibly)** when it becomes non-relativistic.

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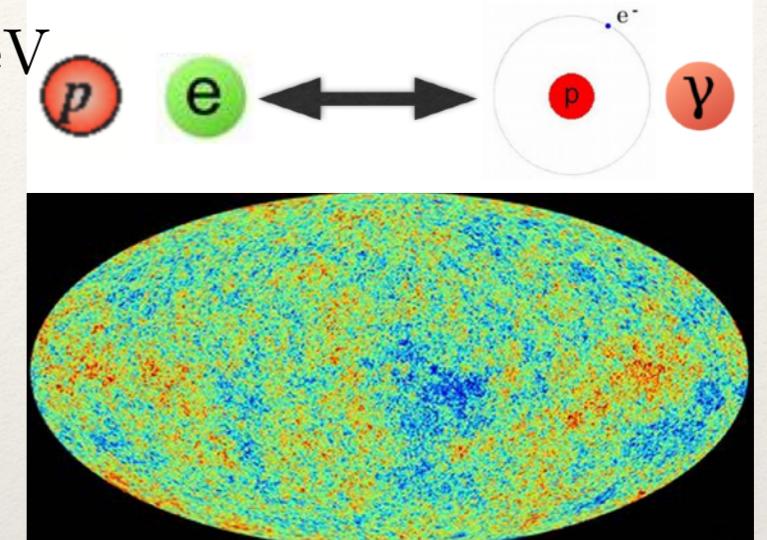
Big Bang Nucleosynthesis (BBN)

with $T_\gamma \sim 1 \text{ MeV}$



Cosmic Microwave Background (CMB)

with $T_\gamma \sim 0.3 \text{ eV}$



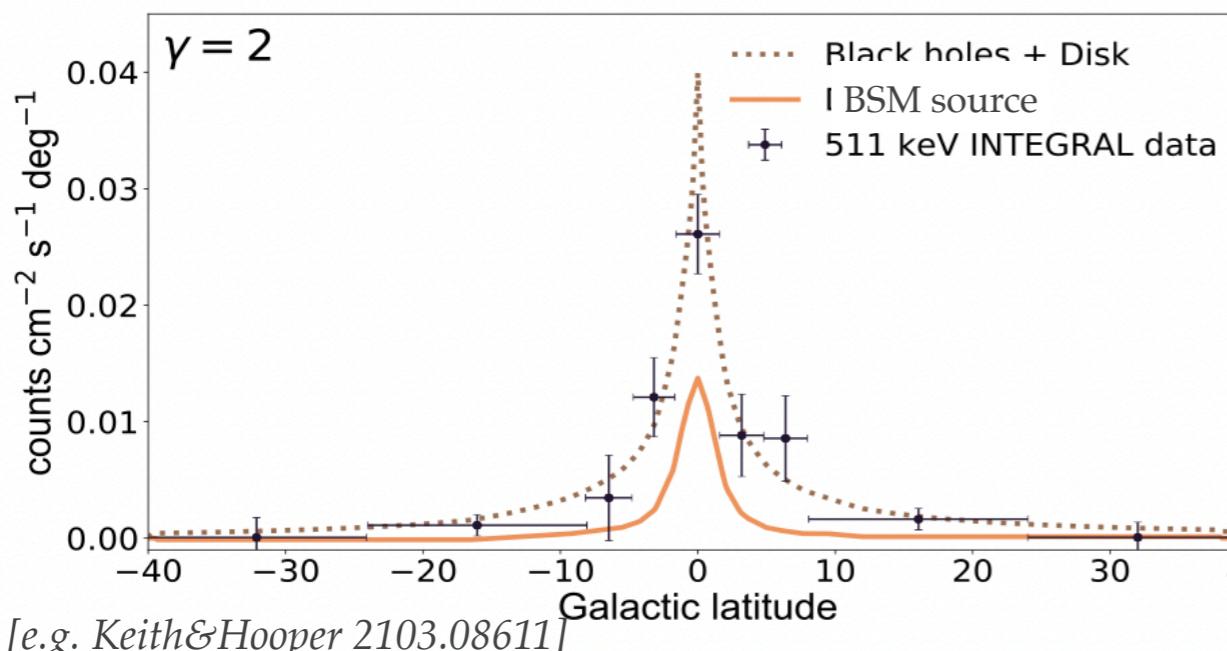
Both constrain the additional EM and **background energy densities**
(around the two photon temperatures above).

At a result, observations suggest: **thermal DM mass** should be above MeV.

But, **what value** exactly,
if DM couples to both electromagnetic (EM) bath and neutrino?

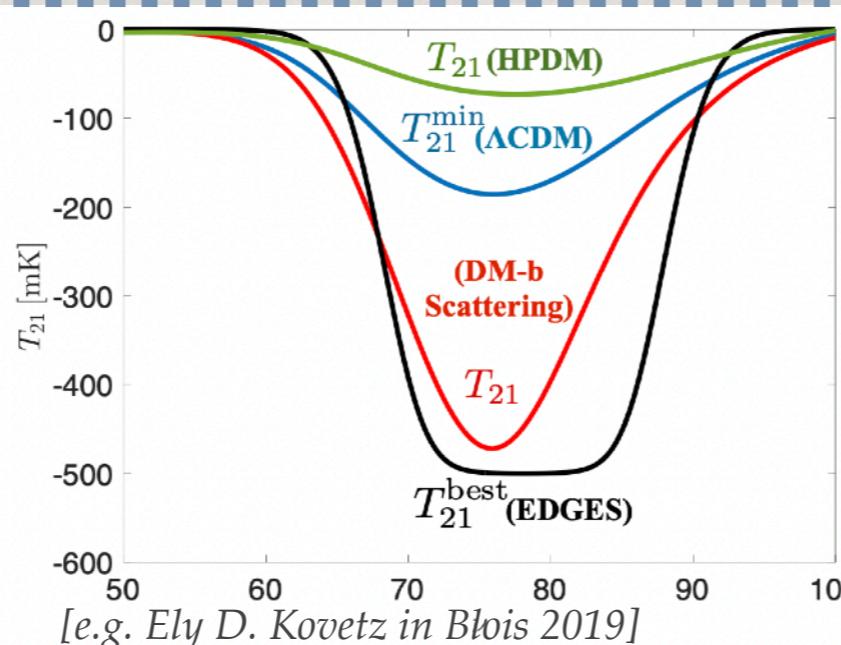
Solving MeV dark particle freeze-out is also relevant for the phenomenology of light dark physics.

Galactic 511 keV-line excess

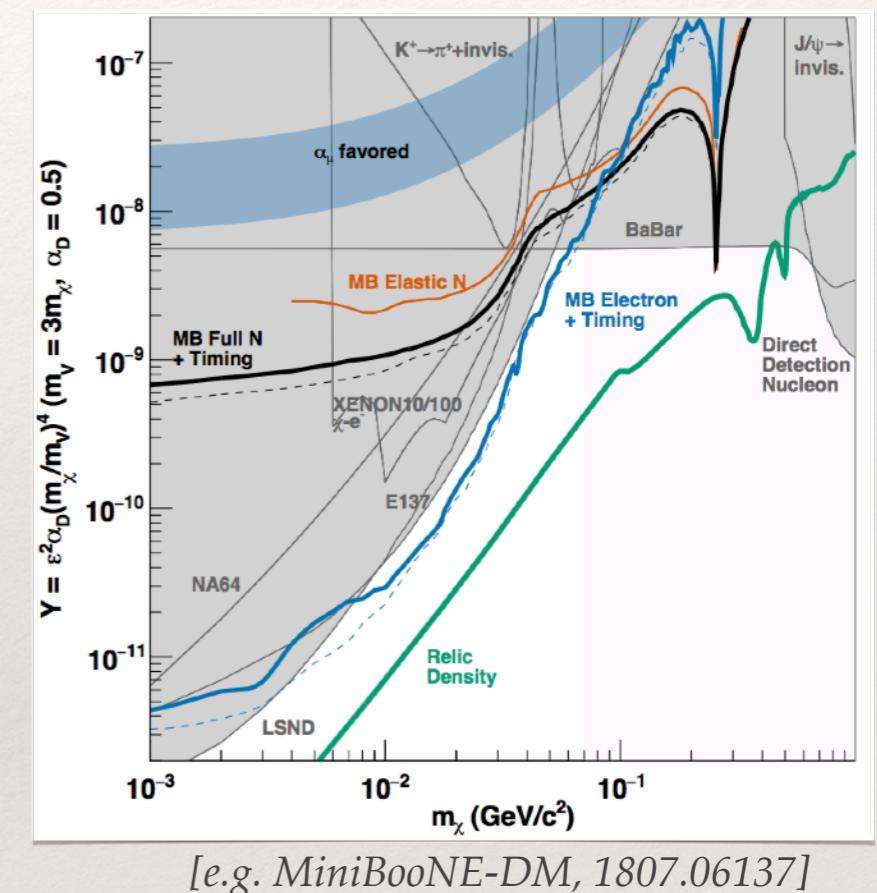


EDGES explanation

Baryon scatters
with sub-leading
MeV DM to further
cool down?



DM
annihilation
produces
low-energy
positrons?



A lower bound from cosmology?

Other anomalies, etc...

To describe DM freeze-out

A three-sector system with a benchmark model

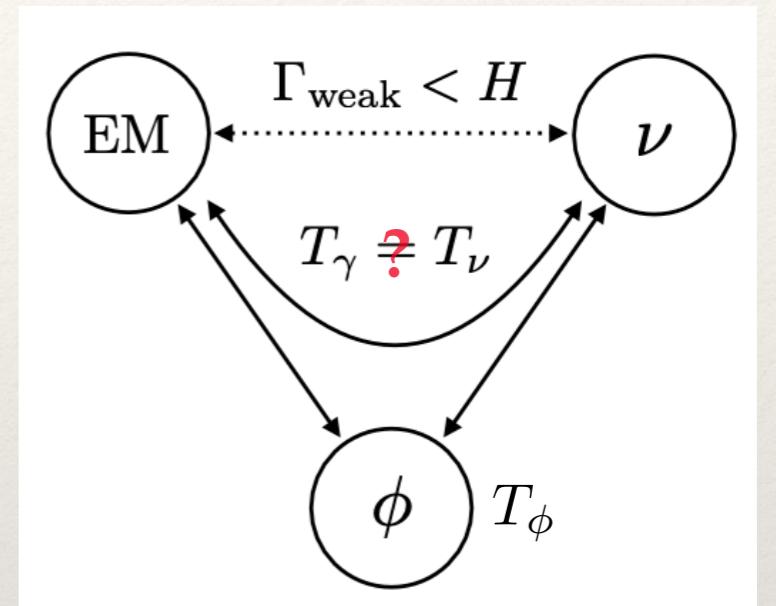
$$\frac{1}{\Lambda_{Z'}^2} (\bar{l} \gamma^\mu l) (\phi^* \overset{\leftrightarrow}{\partial}_\mu \phi)$$

In earlier works

[M.Escudero 1812.05605 & N.Sabti, M.Escudero, et al, 1910.01649]:

- Actual DM annihilation cross section was never obtained for such a three-sector system;
- Neglecting scattering processes, only Maxwell-Boltzmann statistics,

[complex scalar DM, with p-wave annihilation]



A three-sector system with a benchmark model

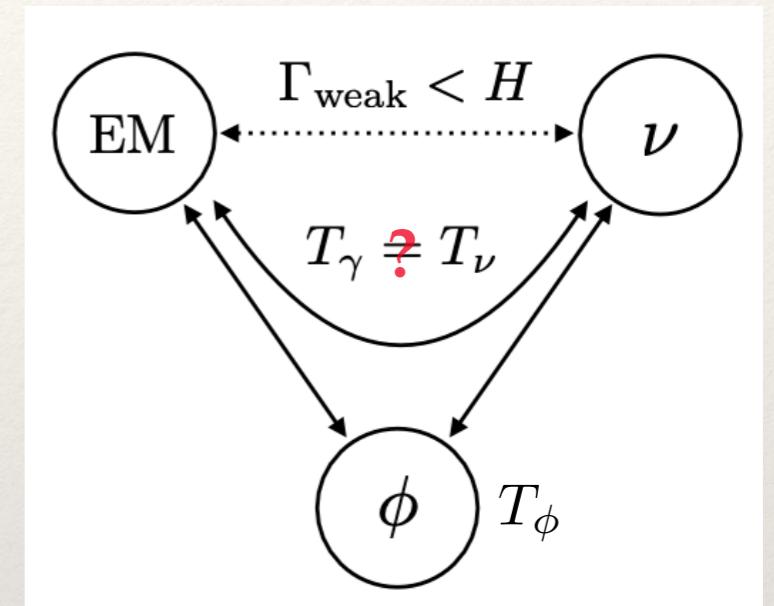
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[complex scalar DM, with p-wave annihilation]



We work out the actual DM annihilation cross section, and, at the same time, obtain the ratio of neutrino-to-photon temperature consistently.

Solving the momentum distribution of each particle species is extremely time-consuming, we have to make some simplifications too.

Two simplifications introduced:

$$\frac{1}{\Lambda_{Z'}^2} (\bar{l} \gamma^\mu l) (\phi^* \overset{\leftrightarrow}{\partial}_\mu \phi)$$

1. **Kinetic equilibrium** for each sector: $f_i(E_i, \mu_i) = \frac{1}{e^{(E_i - \mu_i)/T_i} \mp 1} \equiv \frac{1}{e^{\tilde{E}_i - \tilde{\mu}_i} \mp 1}$

A) the EM sector

Obviously satisfied, with
null chemical potentials
(up to asymmetries).

B) the neutrino sector

Non-equilibrium neutrinos
contribute negligibly even in
SM-only case.

C) the dark sector

depending on **DM self-interaction.**

Evidences for strongly self-interacting DM?

2. **Further separate T and μ for:**

B) the relativistic neutrino

$$f_\nu(\tilde{E}_\nu, \tilde{\mu}_\nu) \simeq \frac{1}{e^{\tilde{E}_\nu} + 1} + \tilde{\mu}_\nu \frac{1}{e^{\tilde{E}_\nu} + e^{-\tilde{E}_\nu} + 2} \\ = f^{(0)}(E_\nu) + \tilde{\mu}_\nu f^{(1)}(E_\nu).$$

C) the non-relativistic DM

$$f_\phi(\tilde{E}_\phi, \tilde{\mu}_\phi) \simeq \begin{cases} \frac{1}{e^{\tilde{E}_\phi} \mp 1} & \text{before freeze-out} \\ \frac{e^{\tilde{\mu}_\phi}}{e^{\tilde{E}_\phi}} & \text{after freeze-out} \end{cases} \\ \rightarrow e^{\tilde{\mu}_\phi} f_\phi^{(0)}(\tilde{E}_\phi)$$

To calculate exactly MeV DM freeze-out:

Solving the number/energy densities by including all two-body processes:

$$\frac{dn_i}{dt} + 3Hn_i = g_i \int \frac{d^3 p_i}{(2\pi)^3 E_i} C[f_i] \equiv \frac{\delta n_i}{\delta t},$$

$$\frac{d\rho_i}{dt} + 3H(\rho_i + p_i) = g_i \int \frac{d^3 p_i}{(2\pi)^3 E_i} \delta E C[f_i] \equiv \frac{\delta \rho_i}{\delta t},$$

For each two-body process $1 + 2 \leftrightarrow 3 + 4$, **the phase space factor:**

$$J = f_1 f_2 (1 \pm f_3) (1 \pm f_4) (1 - e^{-\tilde{\mu}_1 - \tilde{\mu}_2 + \tilde{\mu}_3 + \tilde{\mu}_4} e^{\tilde{E}_1 + \tilde{E}_2 - \tilde{E}_3 - \tilde{E}_4})$$

3f and 4f terms not included yet.

$$\frac{\delta n_i}{\delta t} = \sum_{i \neq j} a_{ij} \beta_{ij}(\tilde{\mu}_i, \tilde{\mu}_j) \gamma_{ij}(T_i, T_j),$$

$$\frac{\delta \rho_i}{\delta t} = \sum_{i \neq j} b_{ij} \beta_{ij}(\tilde{\mu}_i, \tilde{\mu}_j) \zeta_{ij}(T_i, T_j),$$

Scan the two temperatures to obtain all numerical values of multi-integrals.

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All described by **5 variables**:

$$T_\gamma, (T_\nu, \tilde{\mu}_\nu), (T_\phi, \tilde{\mu}_\phi)$$

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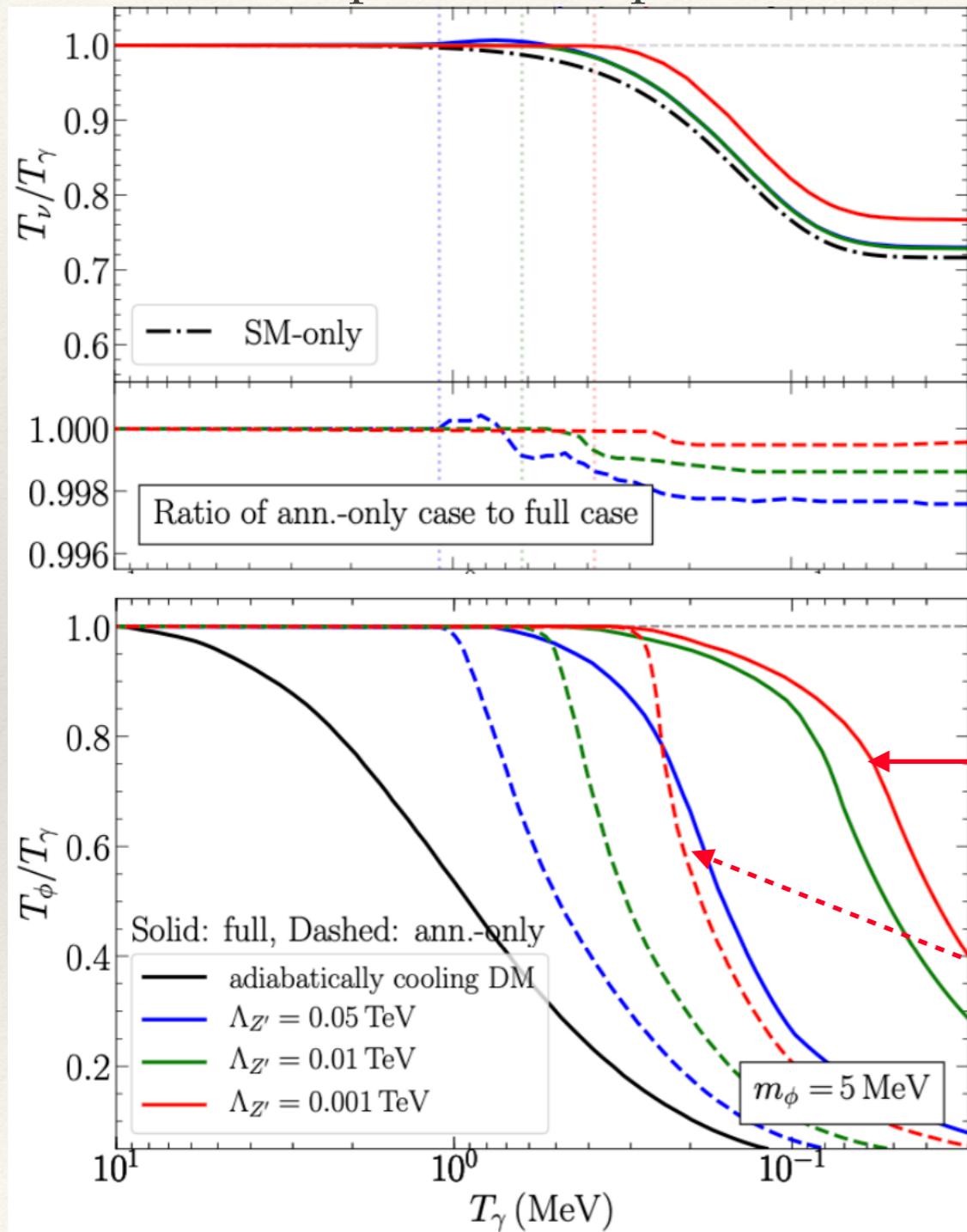
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~~Scan the two temperatures to obtain all numerical values of multi-integrals.~~

Our numerical results

How the system evolves:

Temperature evolution
w.r.t. to photon temperature



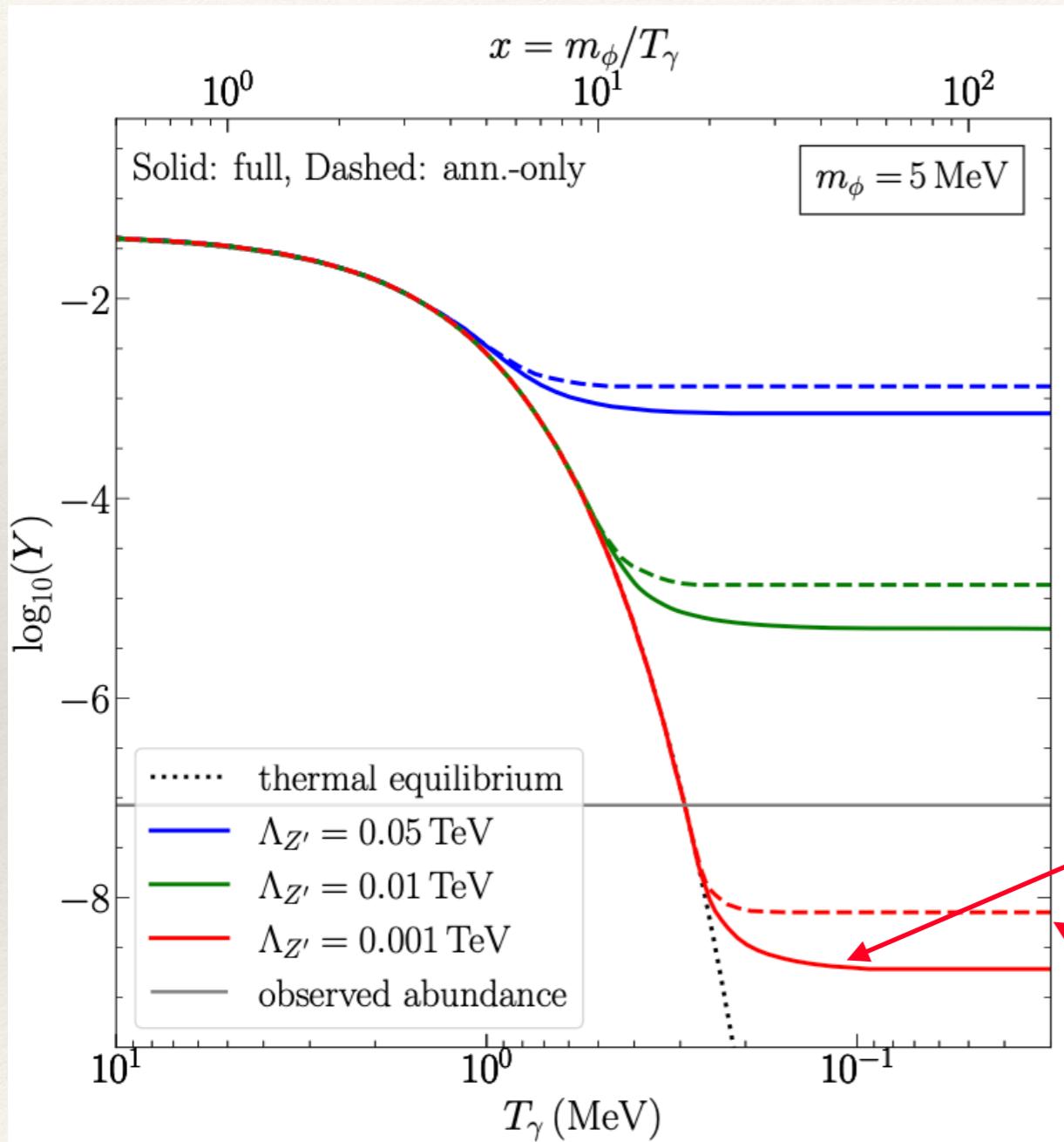
$$\frac{1}{\Lambda_{Z'}^2} (\bar{l} \gamma^\mu l) (\phi^* \overset{\leftrightarrow}{\partial}_\mu \phi)$$

- The standard value with **only weak interaction** for neutrino:
- **Larger coupling** allows photon bath to heat up neutrino via DM.
- Energy transfer is mostly induced by **DM pair annihilation/creation**.
- **DM temperature** is in between photon and neutrino temperatures, before it gradually decouples.
(DM-SM scattering is important here)

How the system evolves:

DM abundance evolution
w.r.t. to photon temperature

$$\frac{1}{\Lambda_{Z'}^2} (\bar{l} \gamma^\mu l) (\phi^* \overset{\leftrightarrow}{\partial}_\mu \phi)$$



- **Tracking DM temperature** is important for p-wave annihilation.

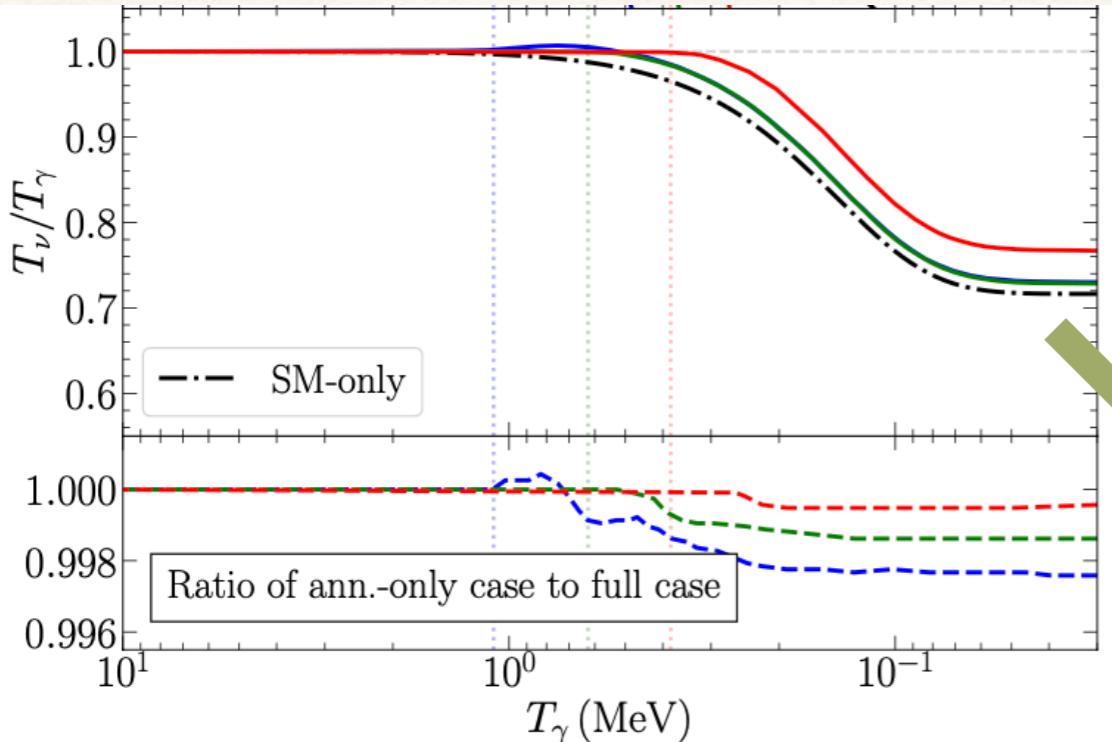
$$\langle \sigma_{\text{ann}} v_M \rangle_{\text{p-wave}} = b \langle v_{\text{rel}}^2 \rangle = b (6T_\phi/m_\phi)$$

(DM-SM scattering is important here)

Taking current Planck bounds:

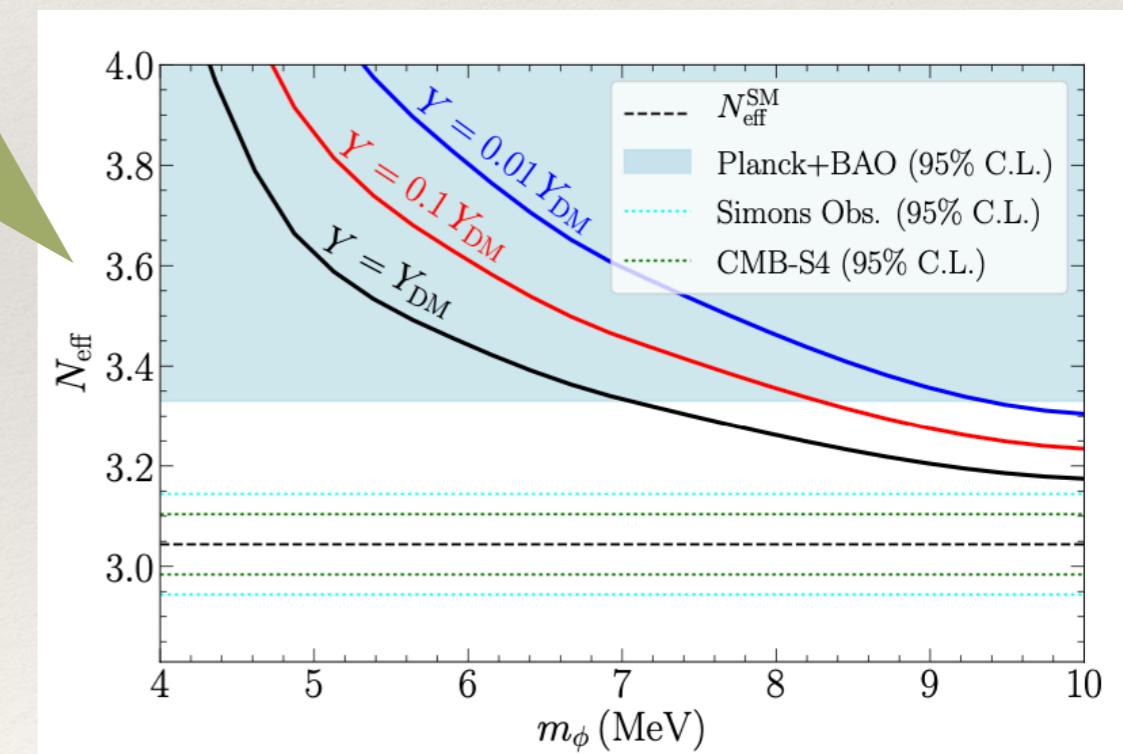
$$2.66 \leq N_{\text{eff}} \leq 3.33$$

around CMB, $0.686 \leq T_\gamma/T_\nu \leq 0.739$



Our bound: $m_\phi > 7$ MeV (p-wave)

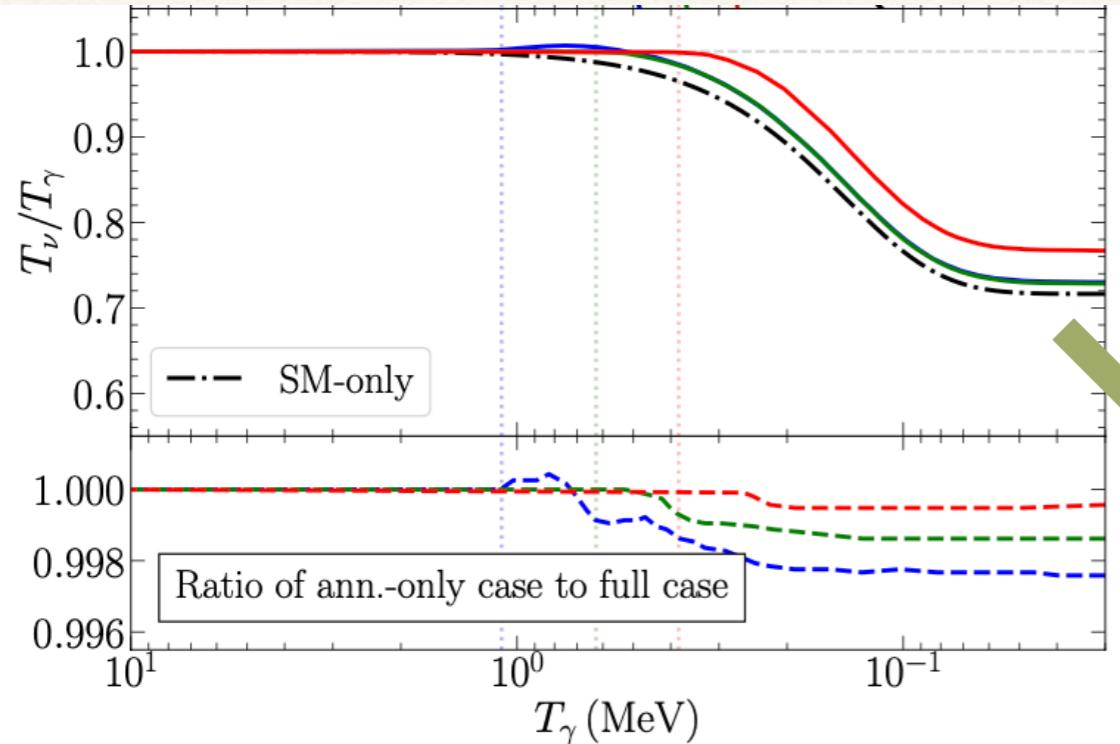
- 8.5/4.6 MeV for e/ ν -only case [Neff, s-wave], 1910.01649;
- 70/200 GeV for e/ γ -line case [all, s-wave], e.g. 1511.08787.



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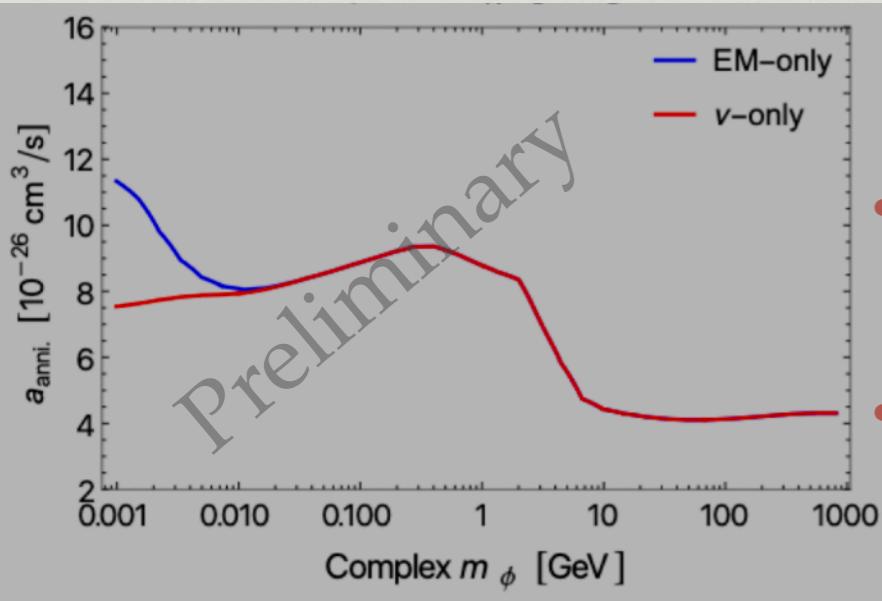
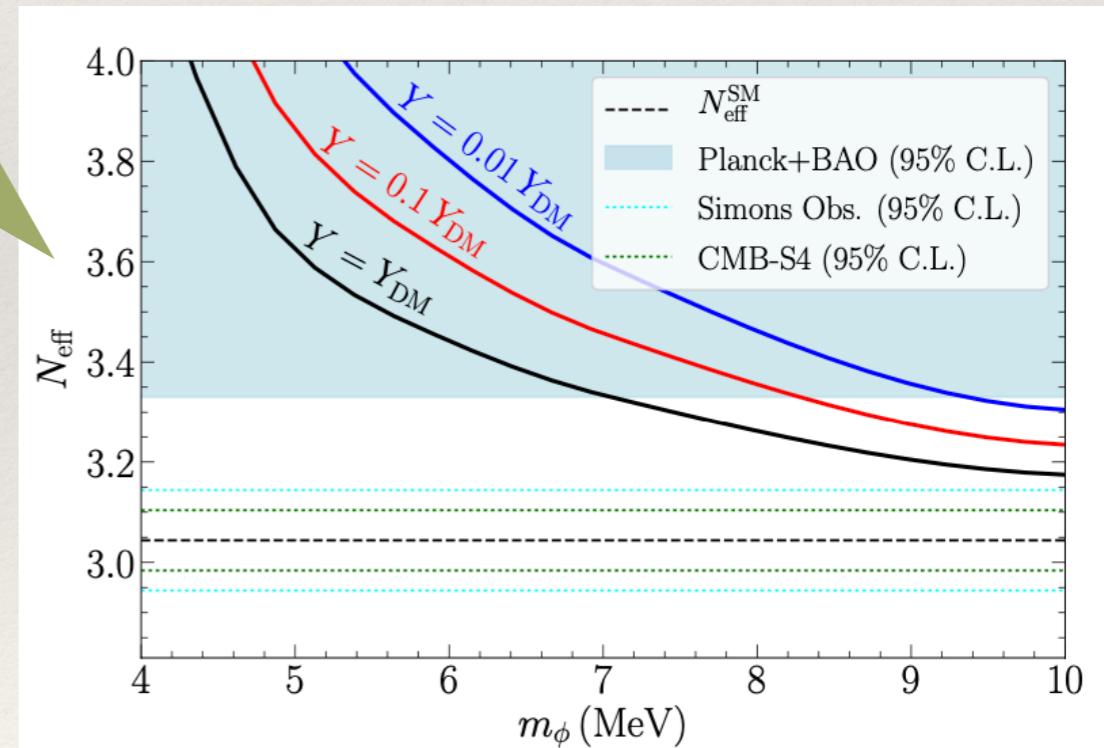
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stronger than previous results (which use pre-fixed **constant cross sections, 4.5 MeV** [N. Sabti, et al, 1910.01649]).

- This is due to:
- **Larger cross section** needed for MeV DM;
 - Enhancement in **p-wave** before freeze-out.



Conclusions

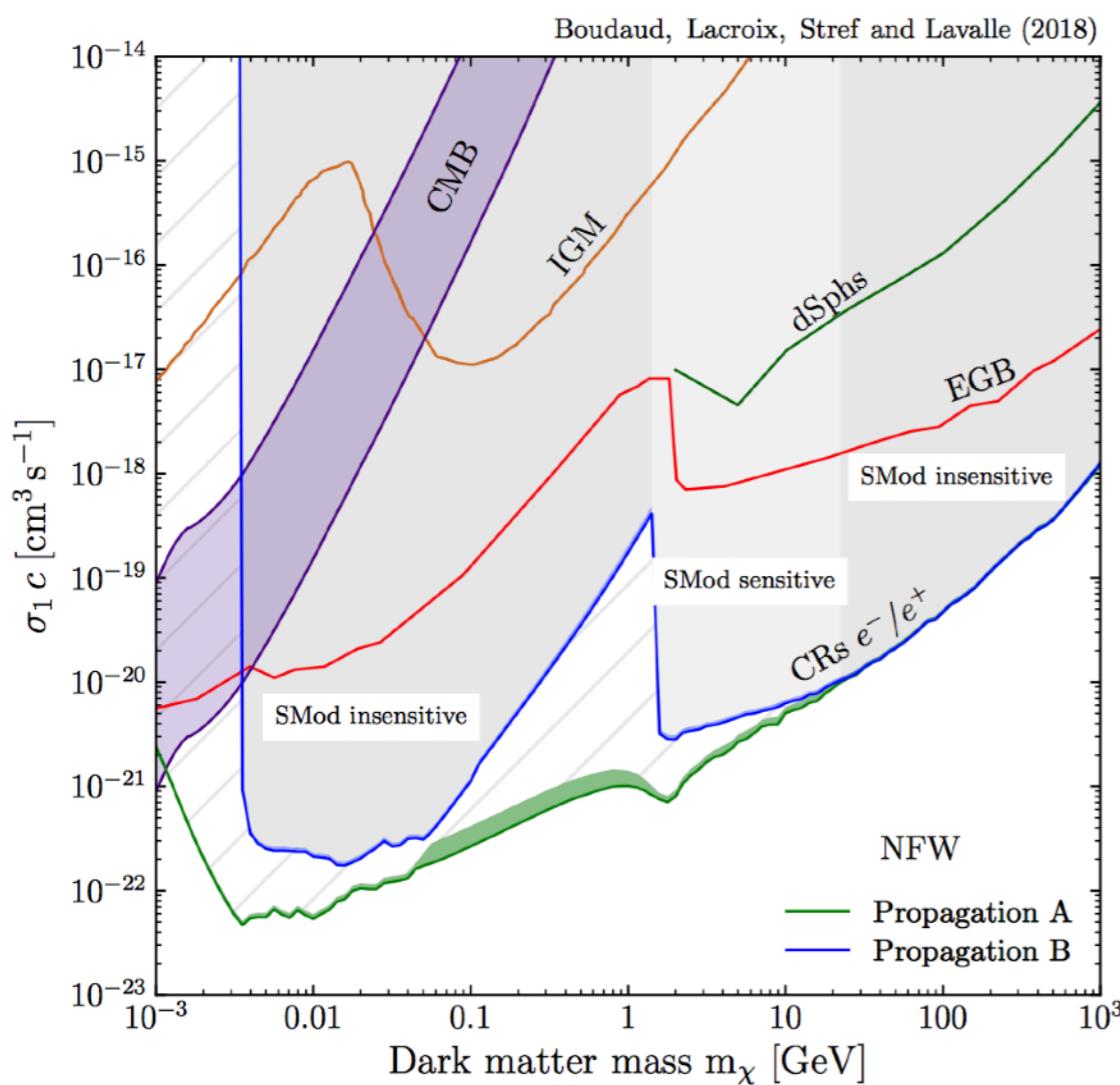
Conclusion and Prospects

- ❖ With better BBN/CMB measurements, we should improve the precision of **theoretical calculations in MeV physics** too.
- ❖ A consistent numerical treatment of **MeV DM freeze-out (into both EM/neutrino)** is provided here, without solving the exact momentum distribution functions.
- ❖ We can now obtain the **full history** of MeV DM decoupling, and will apply it to more cases (DM spins, **varying branching ratio**).
 - **BBN constraints** are crucial for fine-tuned branching ratios.

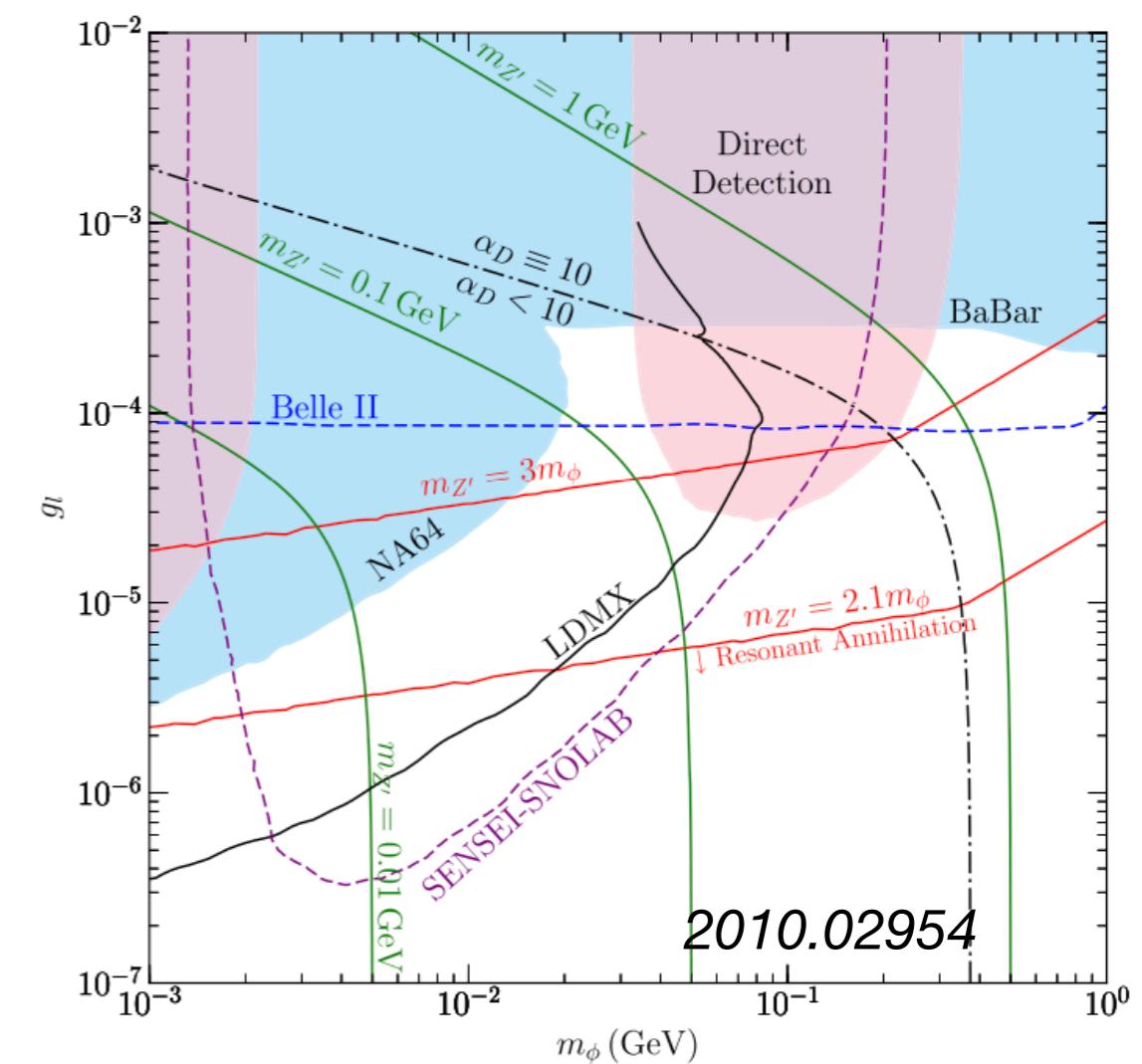
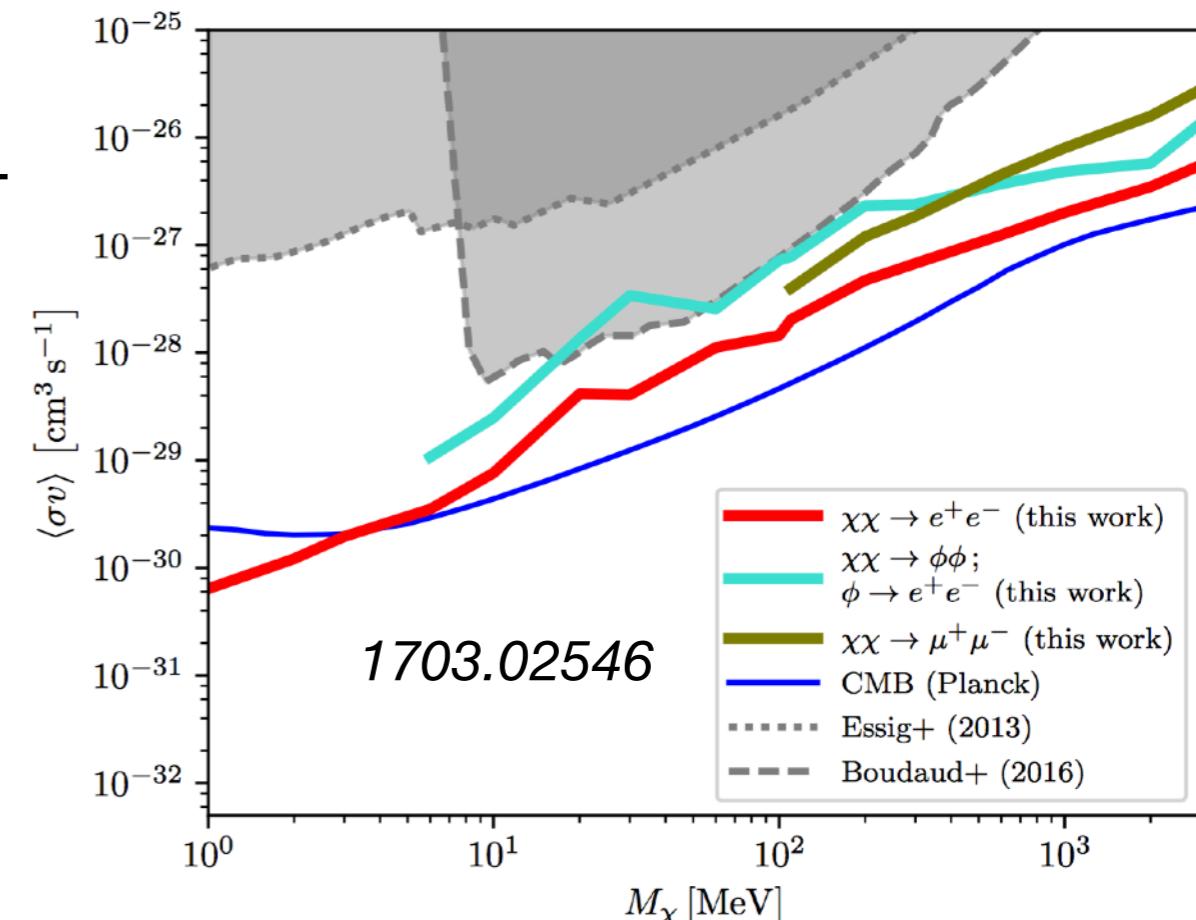
Thanks!

Non-Neff Bounds on DM

$$\begin{aligned}\langle \sigma v \rangle &= \langle \sigma v \rangle_{s\text{-wave}} + \langle \sigma v \rangle_{p\text{-wave}} + \text{higher orders} \\ &= \sigma_0 c + \sigma_1 c \left\langle \frac{v_r^2}{c^2} \right\rangle + \mathcal{O} \left(\frac{v_r^4}{c^4} \right),\end{aligned}$$

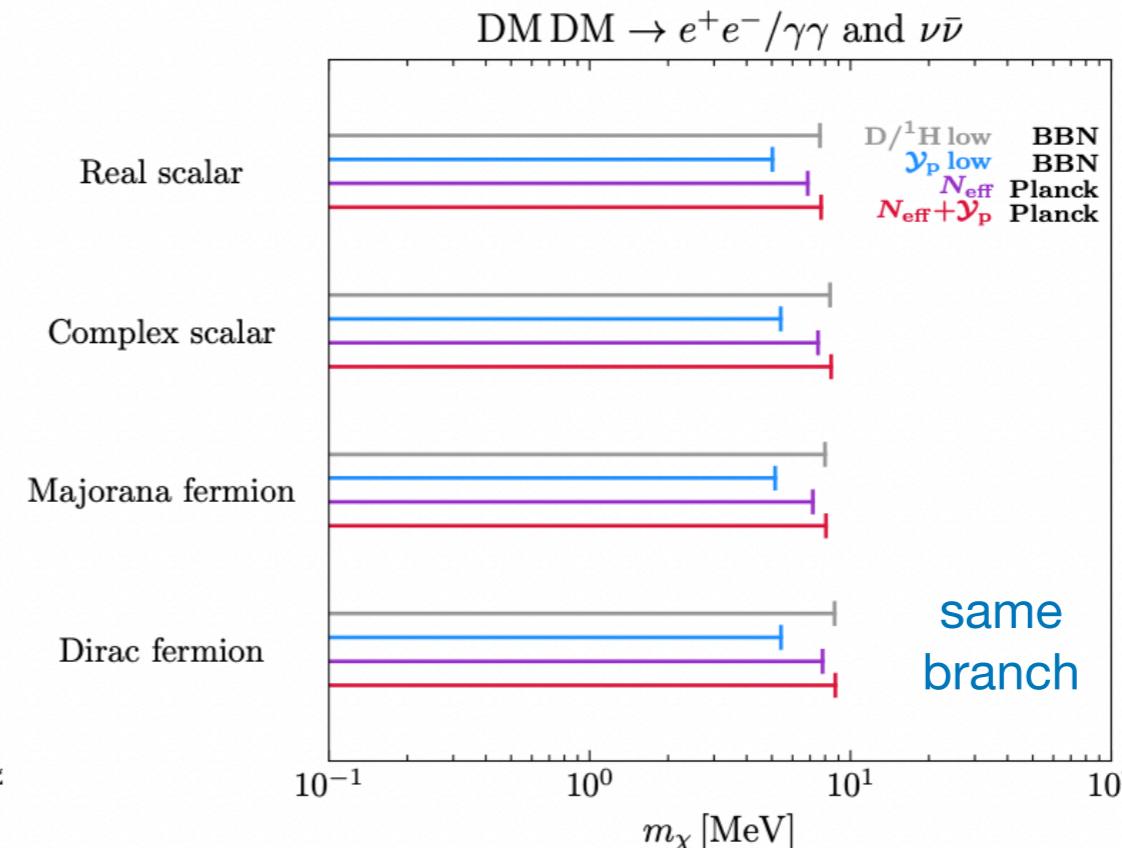
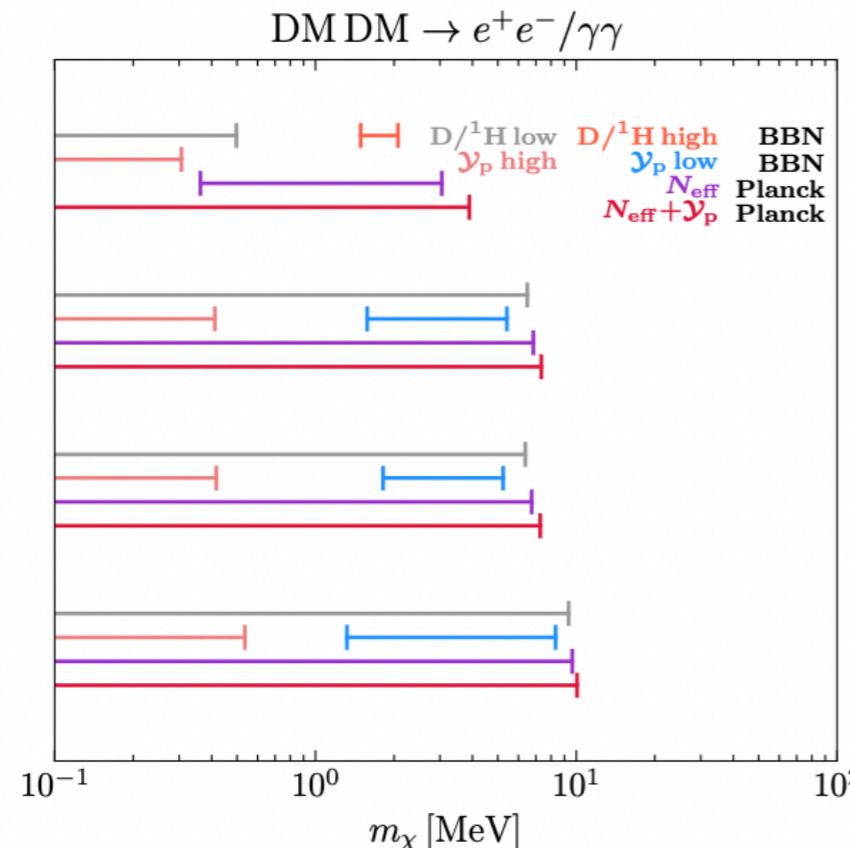


$$\begin{aligned}C[f_1] &= -\frac{Sg_2}{2} \int d\Pi_{i=2,3,4} (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - p_4) \\ &\times J \frac{1}{g_1 g_2} \sum_{\text{spins}} |\mathcal{M}_{12 \leftrightarrow 34}|^2.\end{aligned}$$



Previous CMB/BBN on s-wave DM

1901.06944: sudden decoupling induced by DM annihilation into e/v



Thermalised and non-mu neutrinos + MB
statistics in collision rates + zero-mass electron
[with DM: 1812.05605, 1910.01649]

