

SUSY $(g-2)_\mu$ With & Without Stable Neutralino

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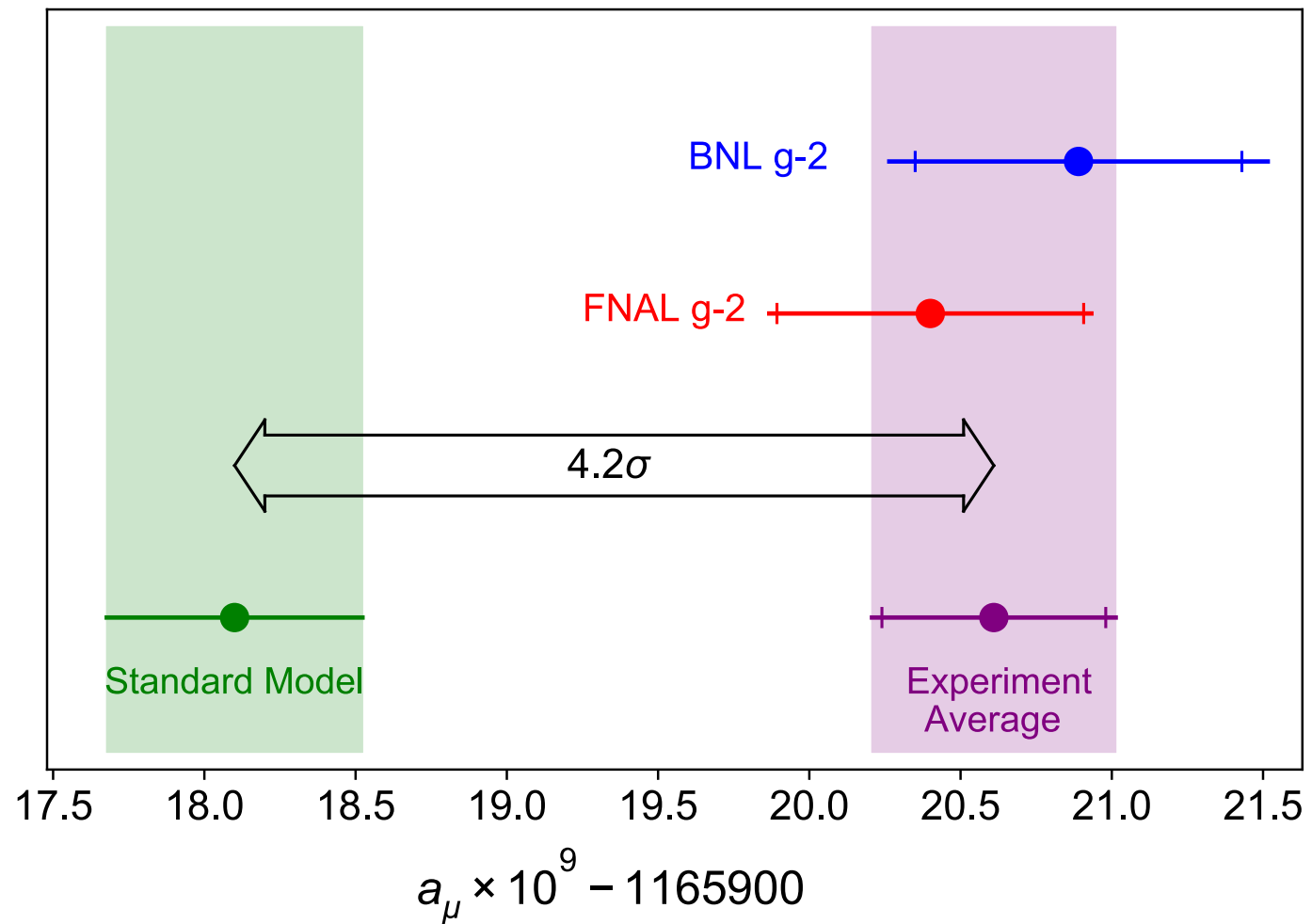
In collaboration with

Manimala Chakraborti, Sho Iwamoto, Jong Soo Kim, Rafał Masełek

Based on [2202.12928]

$(g - 2)_\mu$ anomaly

[Phys. Rev. Lett. 126 (2021) 14, 141801]



$$a_\mu^{\text{theo}} = 0.0011659\,1810(43)$$

$$a_\mu^{\text{exp}} = 0.0011659\,2061(41)$$

$$a_\mu^{\text{exp}} - a_\mu^{\text{theo}} \simeq (25 \pm 6) \times 10^{-10} \simeq \Delta a_\mu^{\text{BSM}} ?$$

Motivation

- There are many BSM scenarios that can explain the $(g-2)_\mu$ anomaly:

Leptoquarks, Z' , VLL, 2HDM, axion, ..

- Supersymmetry is particularly motivated since it offers:

Coupling Unification, Radiative EWSB, Baryogenesis, DM, ...

- There are many studies on SUSY $g-2$ already:

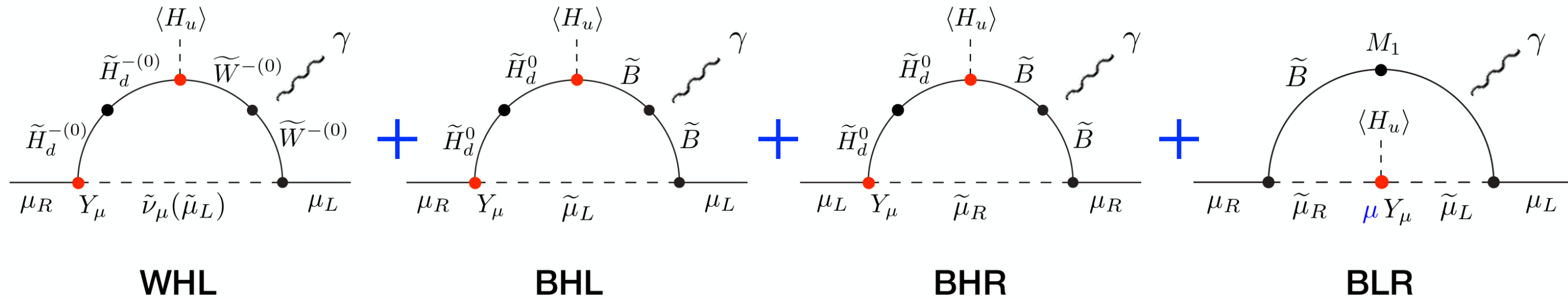
[Athrona, Balazsa, Jacoba, Kotlarskic, Stockinger, Stockinger-Kim]; [Chakraborti, Heinemeyer, Saha]; [Endo, Hamaguchi, Iwamoto, Kitahara]; [Cox, Han, Yanagida]; [Baum, Carena, Shah, Wagner]; [Badziak, KS]; [Hagiwara, Ma, Mukhopadhyay'18], ...

- Most studies assume the neutralino is the Lightest SUSY Particle (LSP) and stable.

Q: What happens if neutralino is unstable? (e.g. RPV, Gravitino LSP)

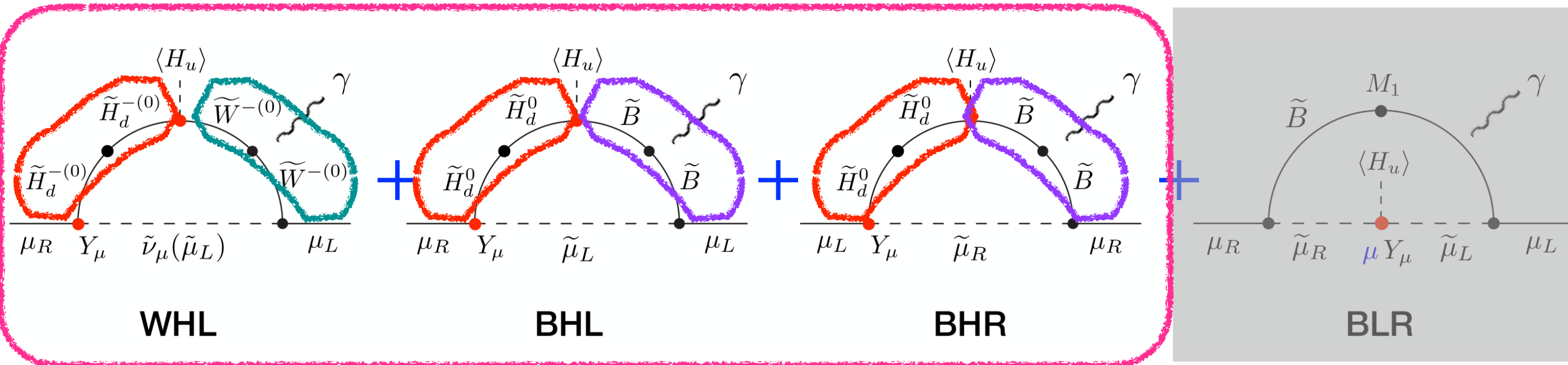
A: DM constraints go away, but LHC constraints change. **How?**

$$\Delta a_\mu^{\text{SUSY}} \simeq$$



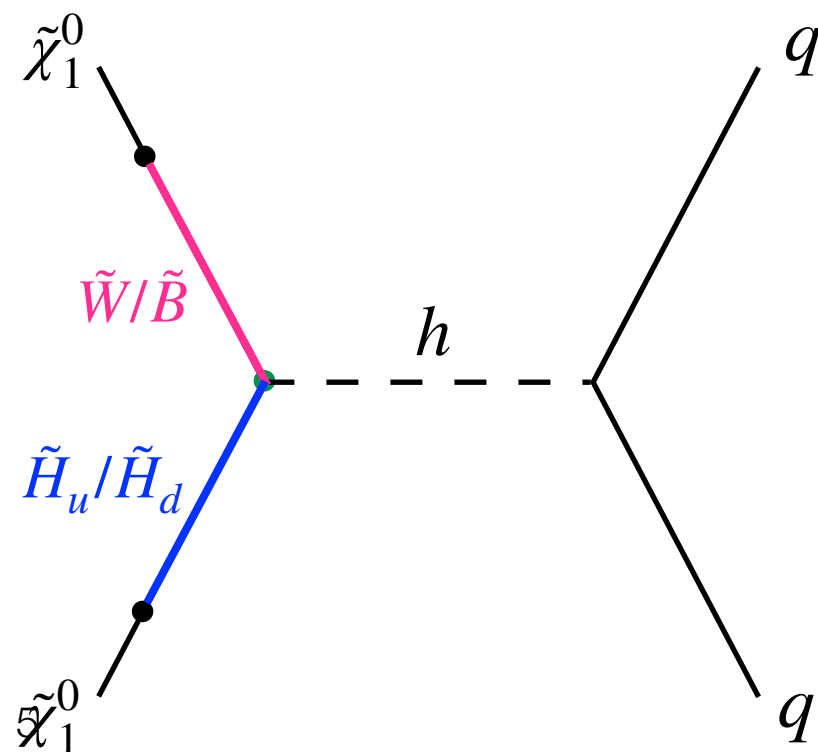
There are 4 important 1-loop (mass insertion) diagrams in large $\tan\beta$ regime

$$\Delta a_\mu^{\text{SUSY}} \approx$$

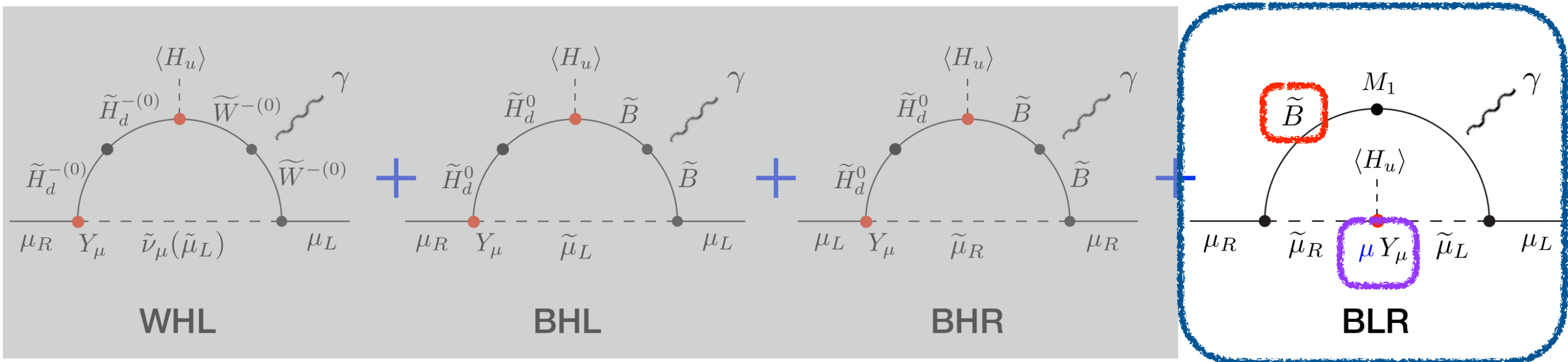


tension

Large gaugino-Higgsino mixing leads to a large cross-section for **DM Direct Detection:**



$$\Delta a_\mu^{\text{SUSY}} \approx$$



❖ **Bino** has very small annihilation cross-section

⇒ **Tend to produce too much DM** $\Omega_{\tilde{\chi}_1^0} > \Omega_{\text{DM}}$

← **tension**

❖ Large off-diagonal term in stau mass matrix:

- charge breaking vacuum: $m_{\text{stau}1}^2 > 0$

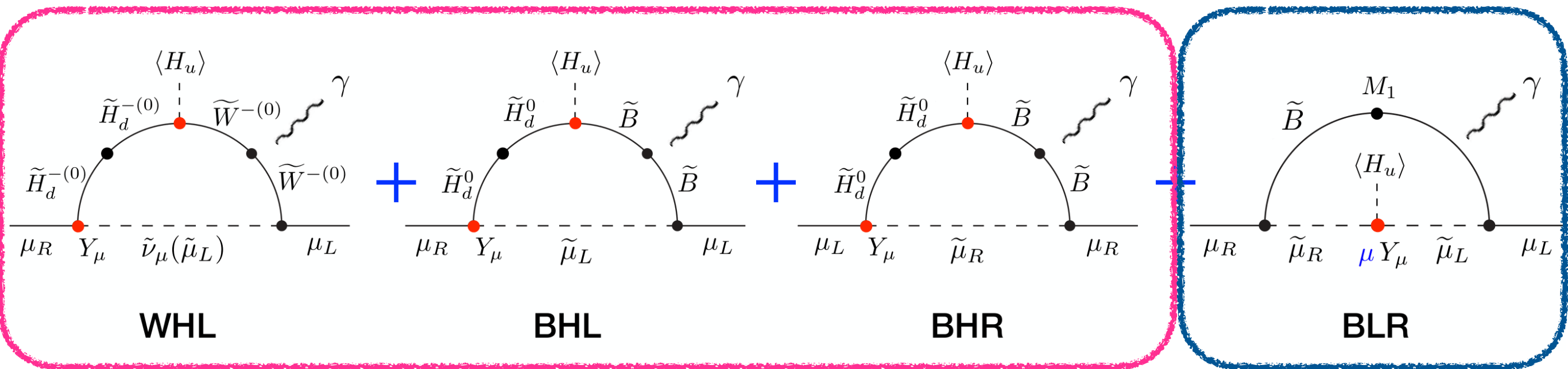
- LEP bound: $m_{\text{stau}1} > 81.9 \text{ GeV}$

- stau LSP: $m_{\text{stau}1} > m_{\text{neutralino}1}$

- **Vacuum (meta-)stability**

$$(\tilde{\tau} \text{ mass matrix}) \sim \begin{pmatrix} m_{\tilde{\tau}_R}^2 & Y_\tau \mu \langle H_u \rangle \\ Y_\tau \mu \langle H_u \rangle & m_{\tilde{\tau}_L}^2 \end{pmatrix}$$

$$\Delta a_{\mu}^{\text{SUSY}} \simeq$$



SUSY g-2 has a tension with:

- **DM** Direct Detection
- (Bino-like) **DM** overproduction
- lepton + **large E_T^{miss}** @ LHC
- Vacuum stability (for BLR)

consequence of **stable neutralino**

How the situation improves / deteriorates if **neutralino is unstable?**

Analysis

Model Parameters: (5 masses + $\tan\beta$)

M_1 : Bino mass

$$m_{\tilde{l}_R} \equiv \widetilde{m}_{\tilde{e}_R}^2 = \widetilde{m}_{\tilde{\mu}_R}^2 = \widetilde{m}_{\tilde{\tau}_R}^2$$

M_2 : Wino mass

$$m_{\tilde{l}_L} \equiv \widetilde{m}_{\tilde{\nu}_e} = \widetilde{m}_{\tilde{\nu}_\mu} = \widetilde{m}_{\tilde{\nu}_\tau} = \widetilde{m}_{\tilde{e}_L} = \widetilde{m}_{\tilde{\mu}_L} = \widetilde{m}_{\tilde{\tau}_L}$$

μ : Higgsino mass

$$\tan\beta \equiv \langle H_u \rangle / \langle H_d \rangle$$

MC
simulations

GM2Calc, CheckMATE2, MicrOmegas ...

constraints

a_μ^{SUSY}

DM abundance

$\sigma_{SI}^{\tilde{\chi}_1^0}$

LHC constraints

List of ATLAS & CMS searches included in our analysis

13 TeV

Name	E/TeV	$\mathcal{L}/\text{fb}^{-1}$	Description
atlas_1604_01306	13	3.2	Monophoton
atlas_1605_09318	13	3.3	3 b-jets + 0-1 lepton + MET
atlas_1609_01599	13	36	Monophoton
atlas_1704_03848	13	36	Monophoton
atlas_conf_2015_082	13	3.2	2 leptons (Z) + jets + MET
atlas_conf_2016_013	13	3.2	1 lepton + jets (4 tops, VVL quarks)
atlas_conf_2016_050	13	13.3	1 lepton + (b) jets + MET
atlas_conf_2016_054	13	13.3	1 lepton + (b) jets + MET
atlas_conf_2016_076	13	13.3	2 lepton + jets + MET
atlas_conf_2016_096	13	13.3	Multi-lepton + MET
atlas_conf_2017_060	13	36	Monojet
atlas_conf_2016_066	13	13.3	Photons, jets and MET
atlas_1712_08119	13	36	soft leptons (compressed EWKinos)
atlas_1712_02332	13	36	squarks and gluinos, 0 lepton, 2-6 jets
atlas_1709_04183	13	36	Jets + MET (stops)
atlas_1802_03158	13	36	search for GMSB with photons
atlas_1708_07875	13	36	EWKino search with taus and MET
atlas_1706_03731	13	36	Multilepton + Jets + MET (RPC and RPV)
atlas_1908_08215	13	36	2 leptons + MET (EWKinos)
atlas_1909_08457	13	139	SS lepton + MET (squark, gluino)
atlas_conf_2019_040	13	139	Jets + MET (squark, gluino)
atlas_conf_2019_020	13	139	3 leptons (EWKino)
atlas_1803_02762	13	36	2 or 3 leptons (EWKino)
atlas_conf_2018_041	13	80	Multi- <i>b</i> -jets (stops, sbottoms)
atlas_2101_01629	13	139	1 lepton + jets + MET
atlas_conf_2020_048	13	139	Monojet
atlas_2004_14060	13	139	$t\bar{t}$ + MET
atlas_1908_03122	13	139	Higgs bosons + <i>b</i> -jets + MET
atlas_2103_11684	13	139	4 or more leptons (RPV, GMSB)
atlas_2106_09609	13	139	Multijets + leptons (RPV)
atlas_1911_06660	13	139	Search for Direct Stau Production
cms_pas_sus_15_011	13	2.2	2 leptons + jets + MET
cms_sus_16_039	13	35.9	electrowekinos in multilepton final state
cms_sus_16_025	13	12.9	electroweakino and stop compressed spectra
cms_sus_16_048	13	35.9	two soft opposite sign leptons

8 TeV

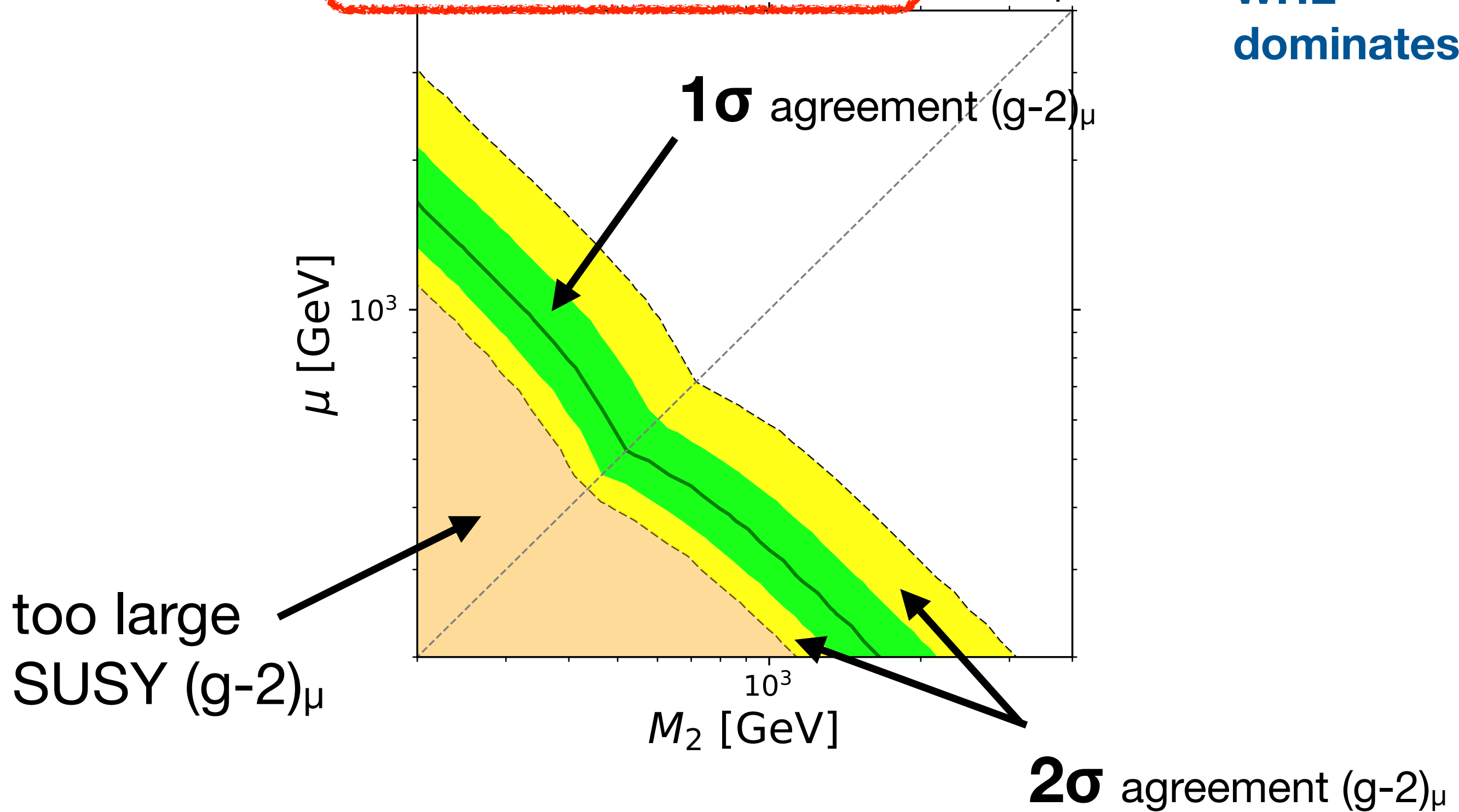
Name	E/TeV	$\mathcal{L}/\text{fb}^{-1}$	Description
atlas_1308_1841	8	20.3	0 lepton + ≥ 7 jets + MET
atlas_1308_2631	8	20.1	0 leptons + 2 b-jets + MET
atlas_1402_7029	8	20.3	3 leptons + MET (chargino+neutralino)
atlas_1403_4853	8	20.3	2 leptons + MET (direct stop)
atlas_1403_5222	8	20.3	stop production with Z boson and b-jets
atlas_1404_2500	8	20.3	Same sign dilepton or 3 lepton
atlas_1405_7875	8	20.3	0 lepton + 2-6 jets + MET
atlas_1407_0583	8	20.3	ATLAS, 1 lepton + (b-)jets + MET (stop)
atlas_1407_0608	8	20.3	Monojet or charm jet (stop)
atlas_1411_1559	8	20.3	monophoton plus MET
atlas_1501_07110	8	20.3	1 lepton + 125GeV Higgs + MET
atlas_1502_01518	8	20.3	Monojet + MET
atlas_1503_03290	8	20.3	2 leptons + jets + MET
atlas_1506_08616	8	20.3	di-lepton and 2b-jets + lepton
atlas_1507_05493	8	20.3	photonic signatures of gauge-mediated SUSY
atlas_conf_2012_104	8	20.3	1 lepton + ≥ 4 jets + MET
atlas_conf_2013_024	8	20.3	0 leptons + 6 (2 b-)jets + MET
atlas_conf_2013_049	8	20.3	2 leptons + MET
atlas_conf_2013_061	8	20.3	0-1 leptons + ≥ 3 b-jets + MET
atlas_conf_2013_089	8	20.3	2 leptons (razor)
atlas_conf_2015_004	8	20.3	invisible Higgs decay in VBF
atlas_1403_5294	8	20.3	2 leptons + MET, (SUSY electroweak)
atlas_higg_2013_03	8	20.3	2 leptons + MET, (invisible Higgs)
atlas_1502_05686	8	20.3	search for massive sparticles decaying to many jets
cms_1303_2985	8	11.7	α_T + b-jets
cms_1408_3583	8	19.7	monojet + MET
cms_1502_06031	8	19.4	2 leptons, jets, MET (only on-Z)
cms_1504_03198	8	19.7	1 lepton, ≥ 3 jets, ≥ 1 b-jet, MET (DM + 2 top)
cms_sus_13_016	8	19.5	OS lepton 3+ b-tags

compressed mass spectrum

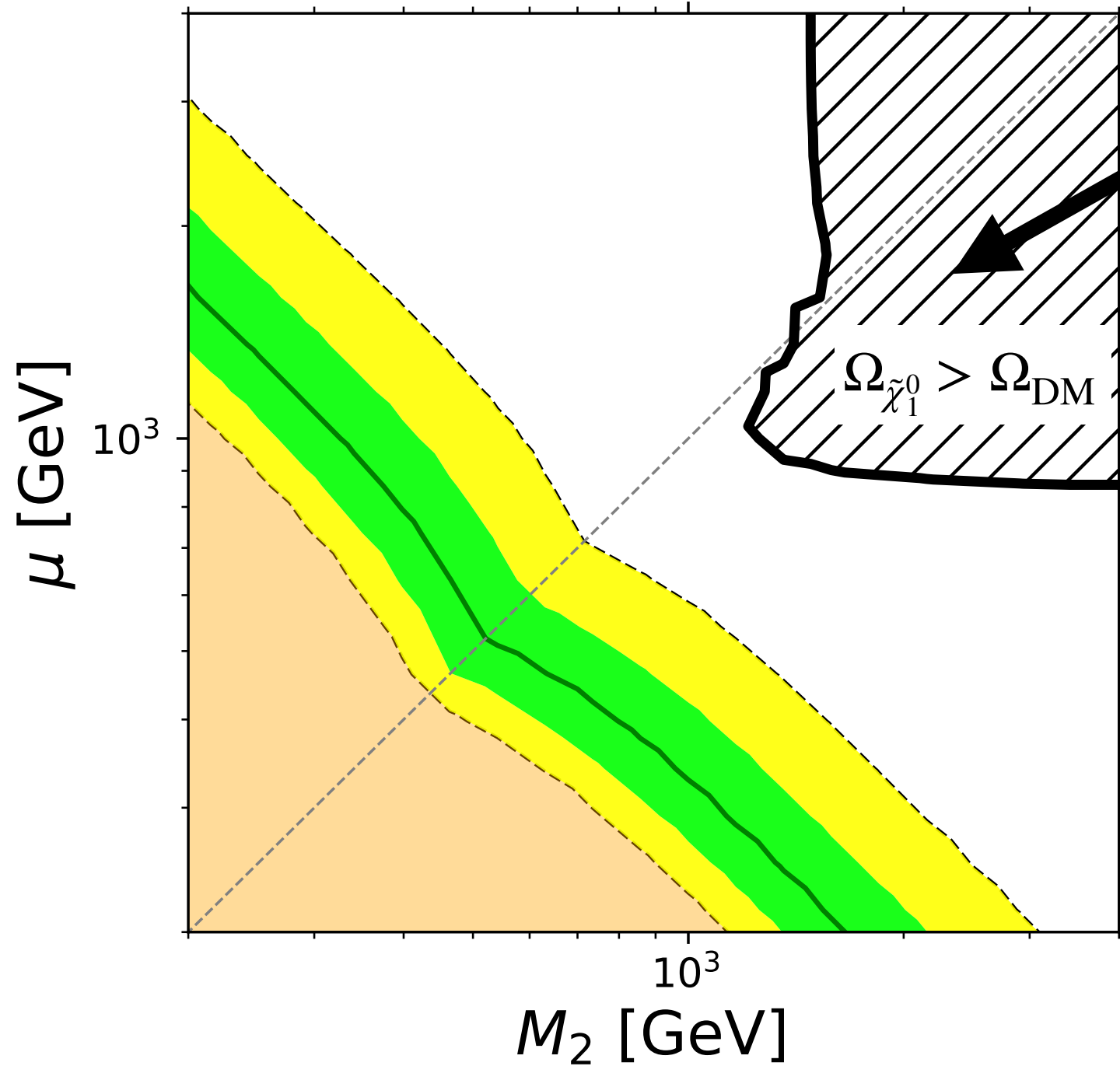
$$\tilde{m}_{l_R} = M_1 = 10 \text{ TeV}$$

$$\tilde{m}_{l_L} = \min(M_2, |\mu|) + 20 \text{ GeV}, \quad \tan \beta = 50$$

WHL
dominates



$$\tilde{m}_{l_L} = \min(M_2, |\mu|) + 20\text{GeV}, \quad \tan\beta = 50, \quad A = 0, \quad \tilde{m}_{l_R} = M_1 = 10\text{TeV}$$



**DM
overabundance**

compressed mass spectrum

$$\tilde{m}_{l_R} = M_1 = 10 \text{ TeV}$$

$$\tilde{m}_{l_L} = \min(M_2, |\mu|) + 20 \text{ GeV}$$

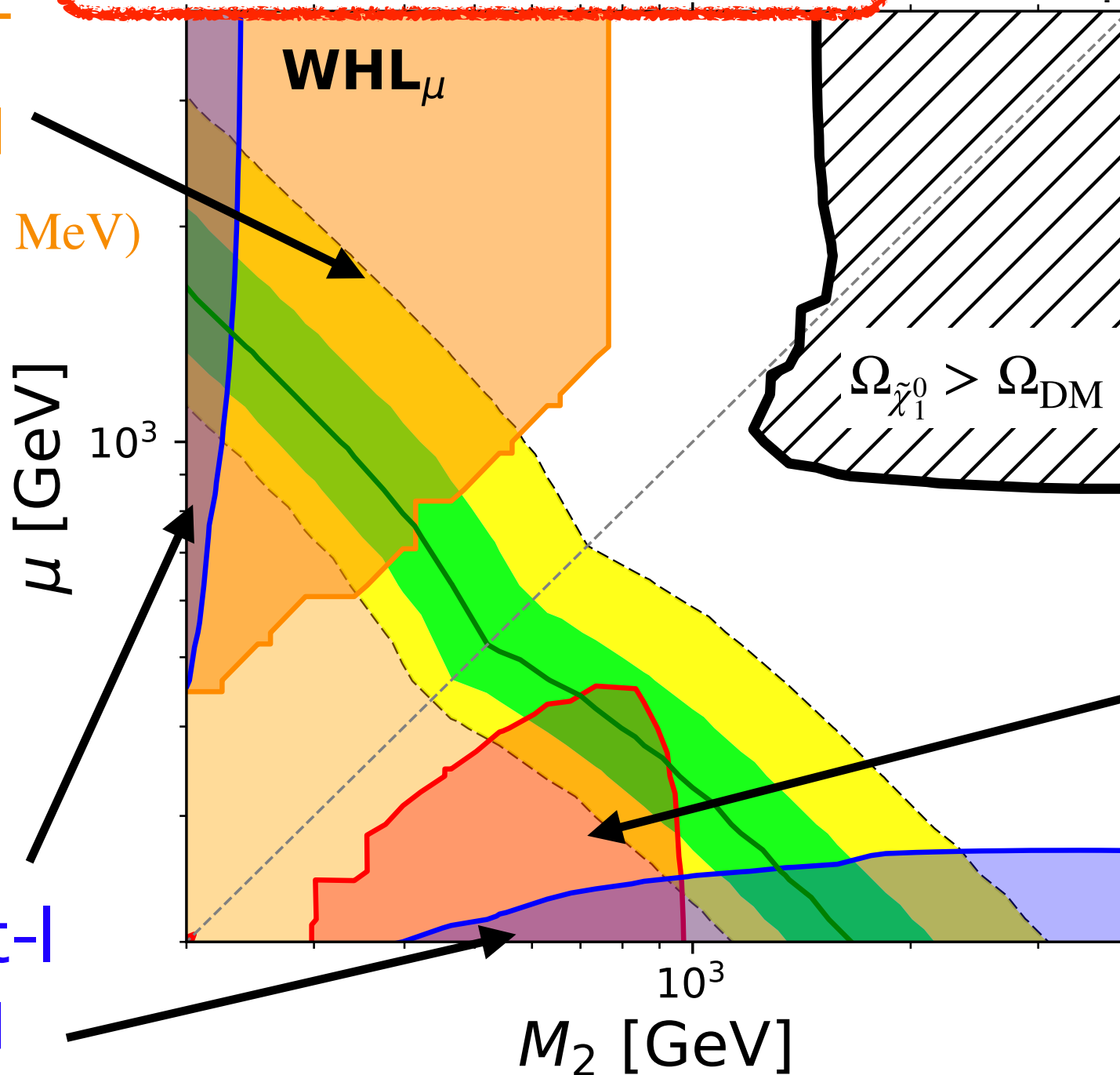
$$\tan \beta = 50$$

ATLAS DT
[2201.02472]

$$m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0} \sim \mathcal{O}(100 \text{ MeV})$$

$$c\tau_{\tilde{\chi}_1^\pm} \sim \mathcal{O}(1 \text{ cm})$$

$$\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 + X_{\text{soft}}^\pm$$



CMS I+I-
[2012.08600]

$$pp \rightarrow \tilde{W}^{+,0} \tilde{W}^{-,0}$$

$$W^\pm \rightarrow l^\pm \tilde{\nu}, W^0 \rightarrow l^\pm \tilde{l}^\mp$$

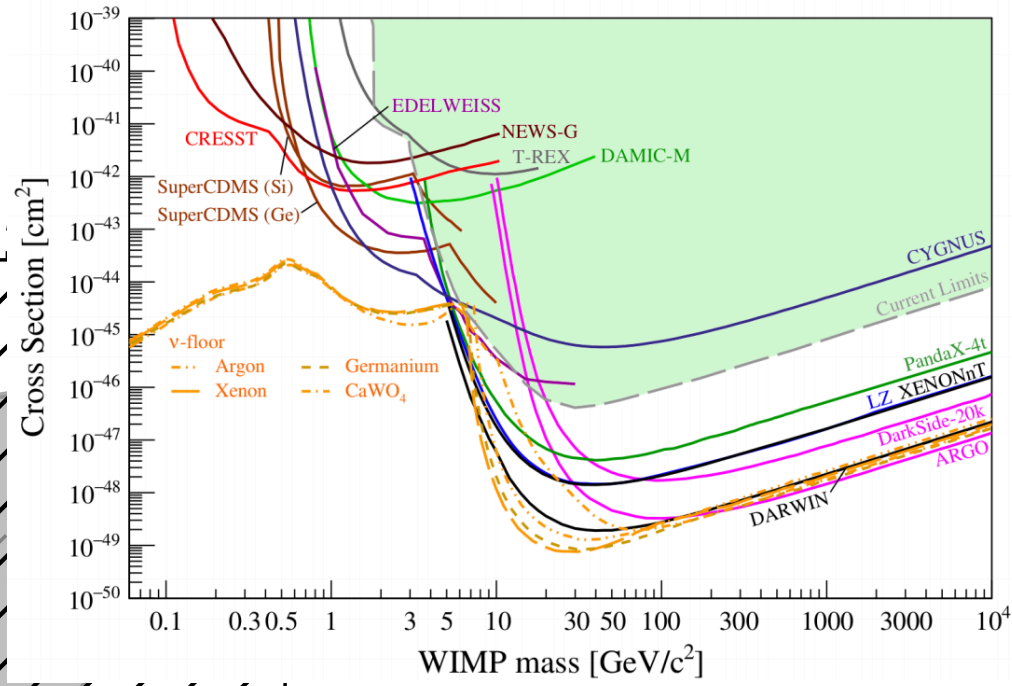
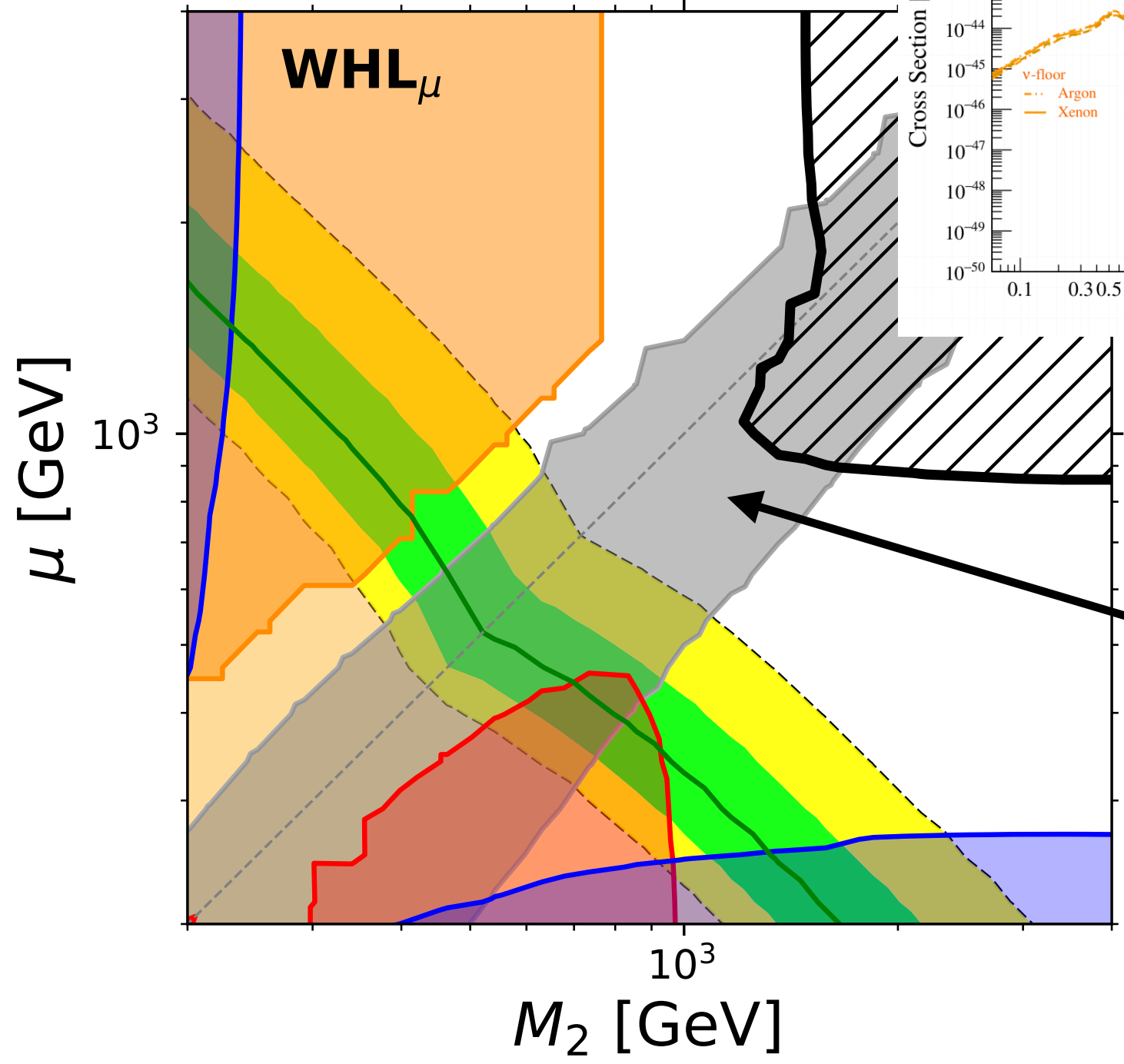
ATLAS soft-l
[1911.12606]

$$\tilde{\xi} \tilde{\xi}' \rightarrow (l^+ \tilde{\eta})(l^- \tilde{\eta}')$$

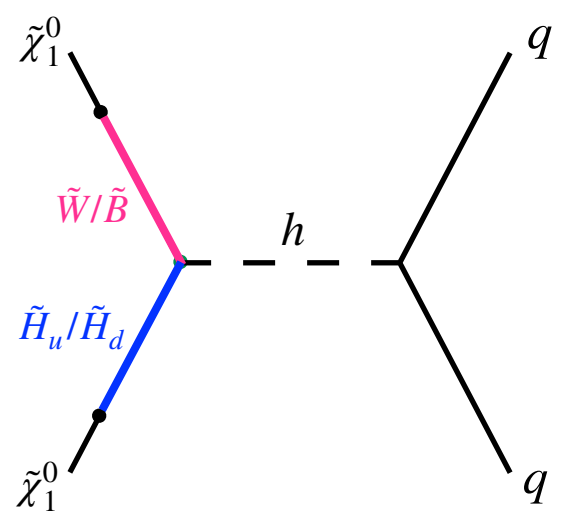
$$\tilde{\xi} \equiv \tilde{l} / \tilde{\nu}$$

$$\tilde{\eta} \equiv \tilde{\chi}_1^\pm / \tilde{\chi}_1^0$$

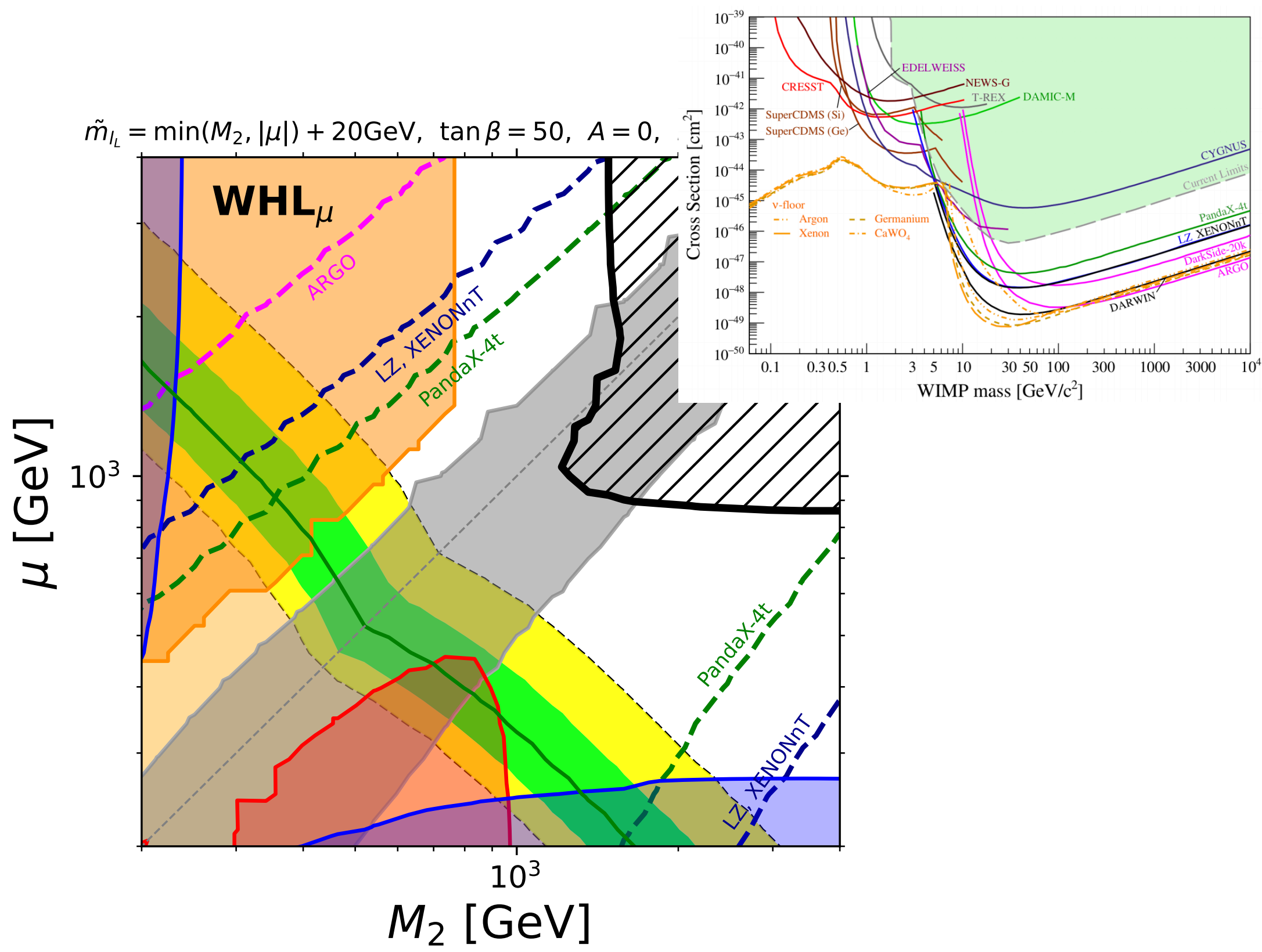
$$\tilde{m}_{l_L} = \min(M_2, |\mu|) + 20\text{GeV}, \quad \tan\beta = 50, \quad A = 0,$$



XENON1T
[1805.12562]



$$\sigma_{N-\tilde{\chi}_1^0}^{\text{SI}} \cdot \frac{\Omega_{\tilde{\chi}_1^0}}{\Omega_{\text{DM}}} > \sigma_{\text{XENON}}^{\text{SI}}$$

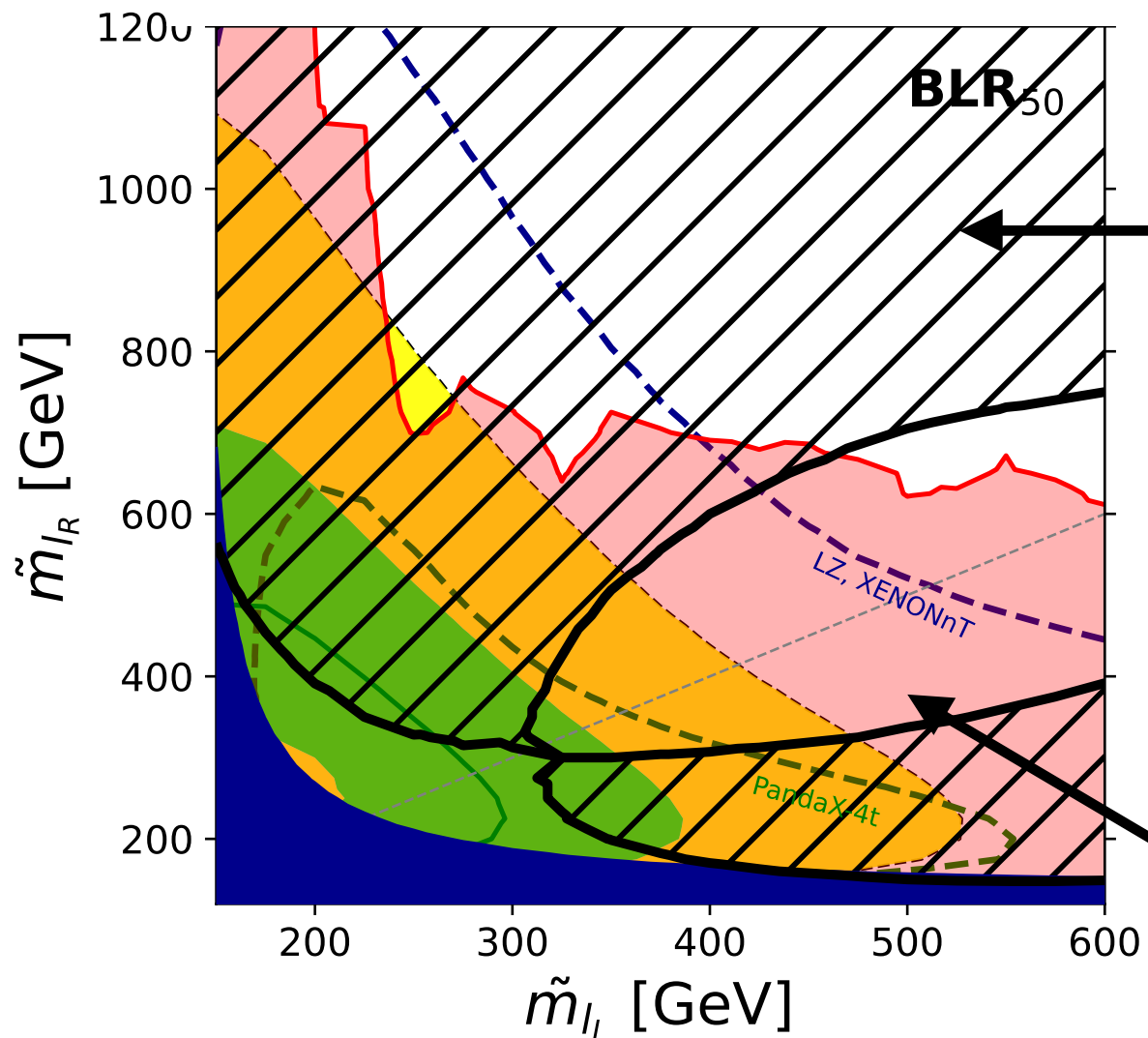
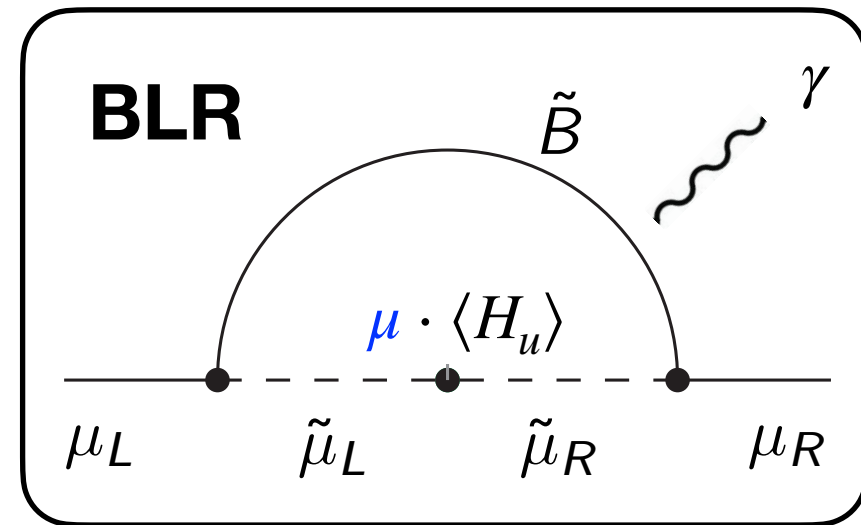


The entire SUSY g-2 region will be probed by the next generation DM-DD experiments

$M_1 = m_{\tilde{\tau}_1} - 20 \text{ GeV}$, \rightarrow compressed spectrum

$\mu = \mu_{\text{max}}$ \rightarrow vacuum stability upper limit

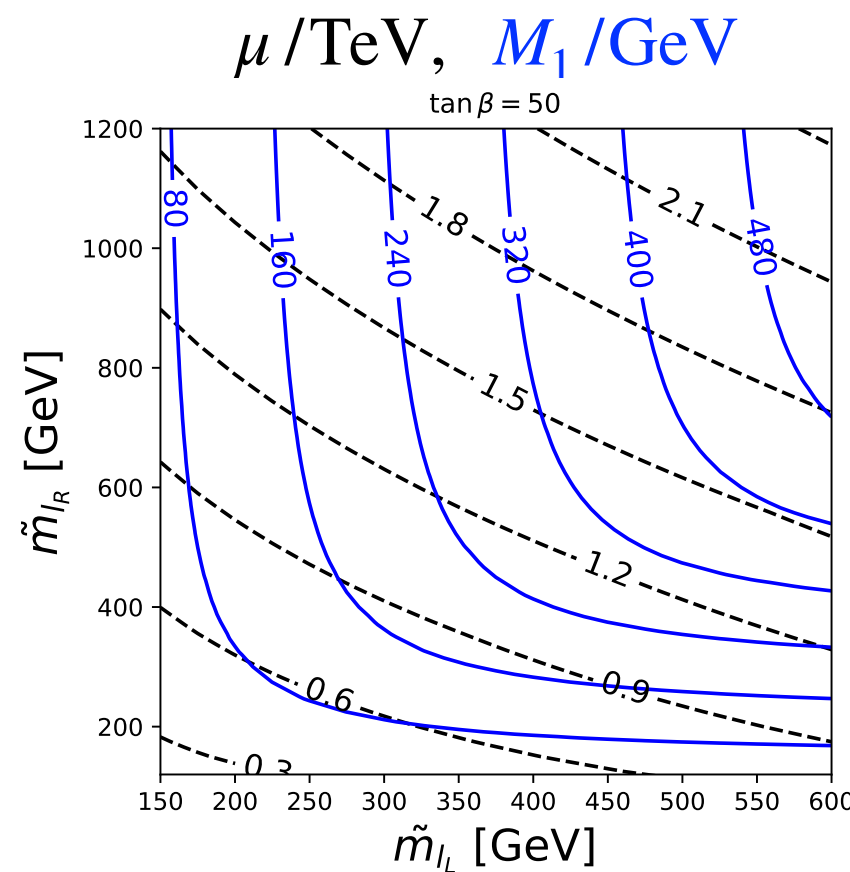
$\tan \beta = 50$ $M_2 = 10 \text{ TeV}$



DM overproduction due to Bino-like LSP

CMS I+I-
[2012.08600]

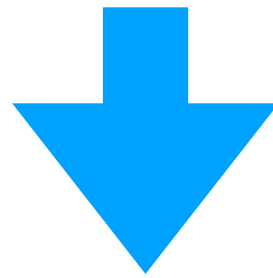
$pp \rightarrow \tilde{W}^{+,0} \tilde{W}^{-,0}$
 $W^\pm \rightarrow l^\pm \tilde{\nu}, W^0 \rightarrow l^\pm \tilde{l}^\mp$



Vacuum stability condition:

$$\left| m_{\tilde{\ell}_{LR}}^2 \right| \leq \left[1.01 \times 10^2 \text{ GeV} \sqrt{m_{\tilde{\ell}_L} m_{\tilde{\ell}_R}} + 1.01 \times 10^2 \text{ GeV} (m_{\tilde{\ell}_L} + 1.03 m_{\tilde{\ell}_R}) - 2.27 \times 10^4 \text{ GeV}^2 + \frac{2.97 \times 10^6 \text{ GeV}^3}{m_{\tilde{\ell}_L} + m_{\tilde{\ell}_R}} - 1.14 \times 10^8 \text{ GeV}^4 \left(\frac{1}{m_{\tilde{\ell}_L}^2} + \frac{0.983}{m_{\tilde{\ell}_R}^2} \right) \right]$$

SUSY $g-2$ solution is tightly constrained if the neutralino is stable, by DM and the missing energy signature at the LHC.



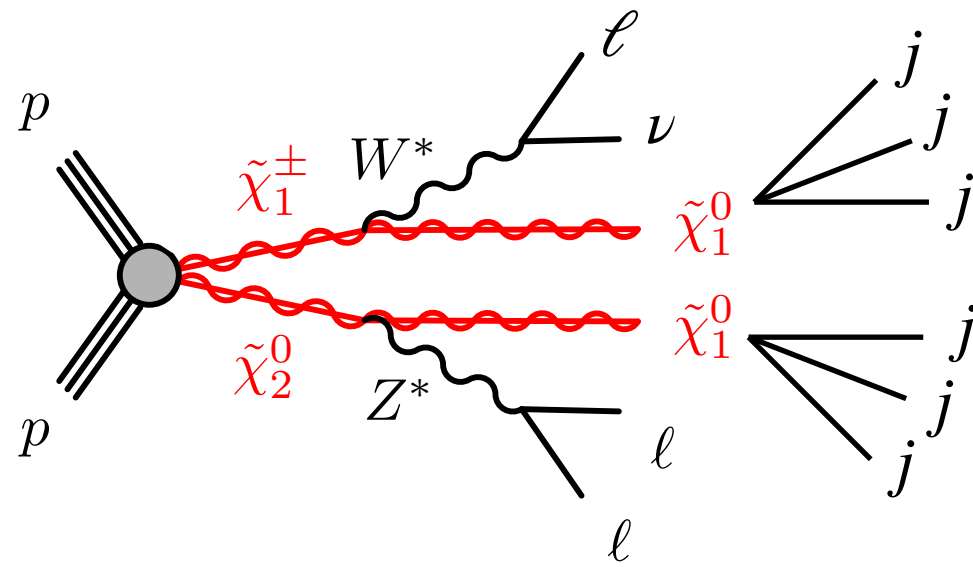
How the situation changes if the neutralino is **unstable**?

We study **two** such examples:

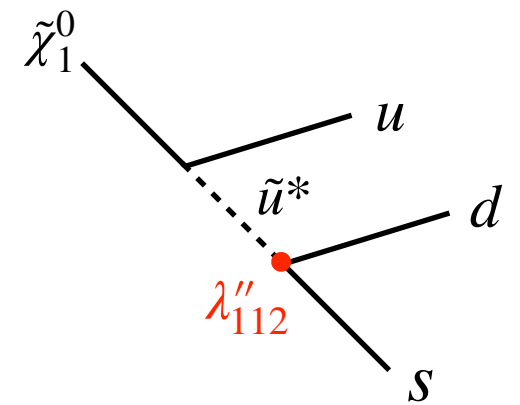
- **R-parity violation (RPV)**
- **Gravitino LSP**

RPV

(UDD-type)



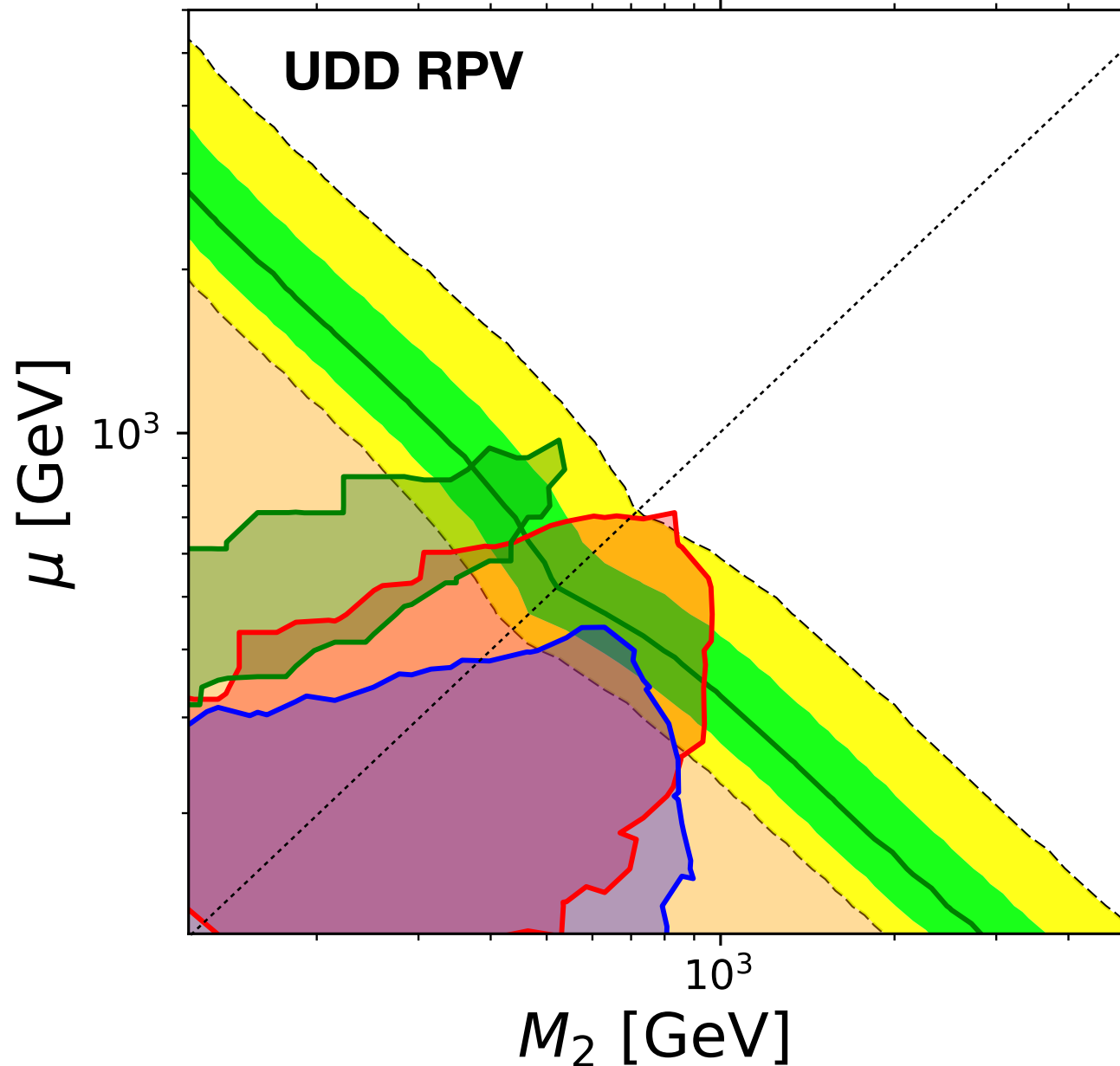
$$W_{\text{RPV}} = \lambda''_{ijk} U_i^c D_j^c D_k^c$$



UDD RPV

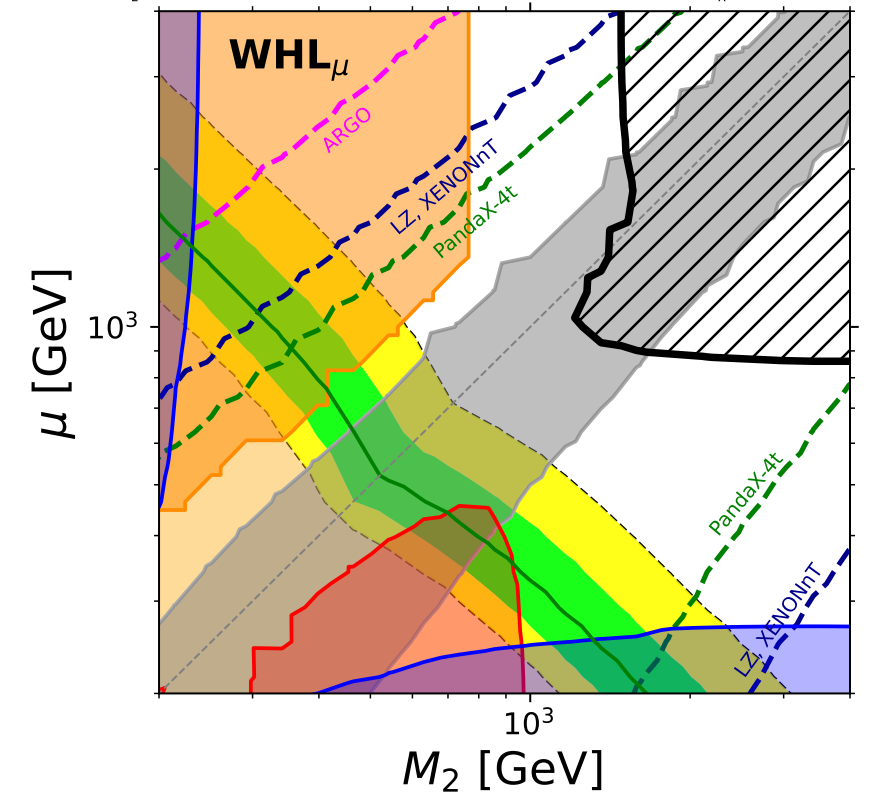
WHL dominated plane

$$m_{\tilde{l}_L} = \min(M_2, |\mu|) + 20\text{GeV}, \tan\beta = 50, A = 0, m_{\tilde{l}_R} = M_1 = 10\text{TeV}$$



stable neutralino

$$m_{\tilde{l}_L} = \min(M_2, |\mu|) + 20\text{GeV}, \tan\beta = 50, A = 0, m_{\tilde{l}_R} = M_1 = 10\text{TeV}$$



ATLAS multijet+l

[2106.09609]

CMS multilepton

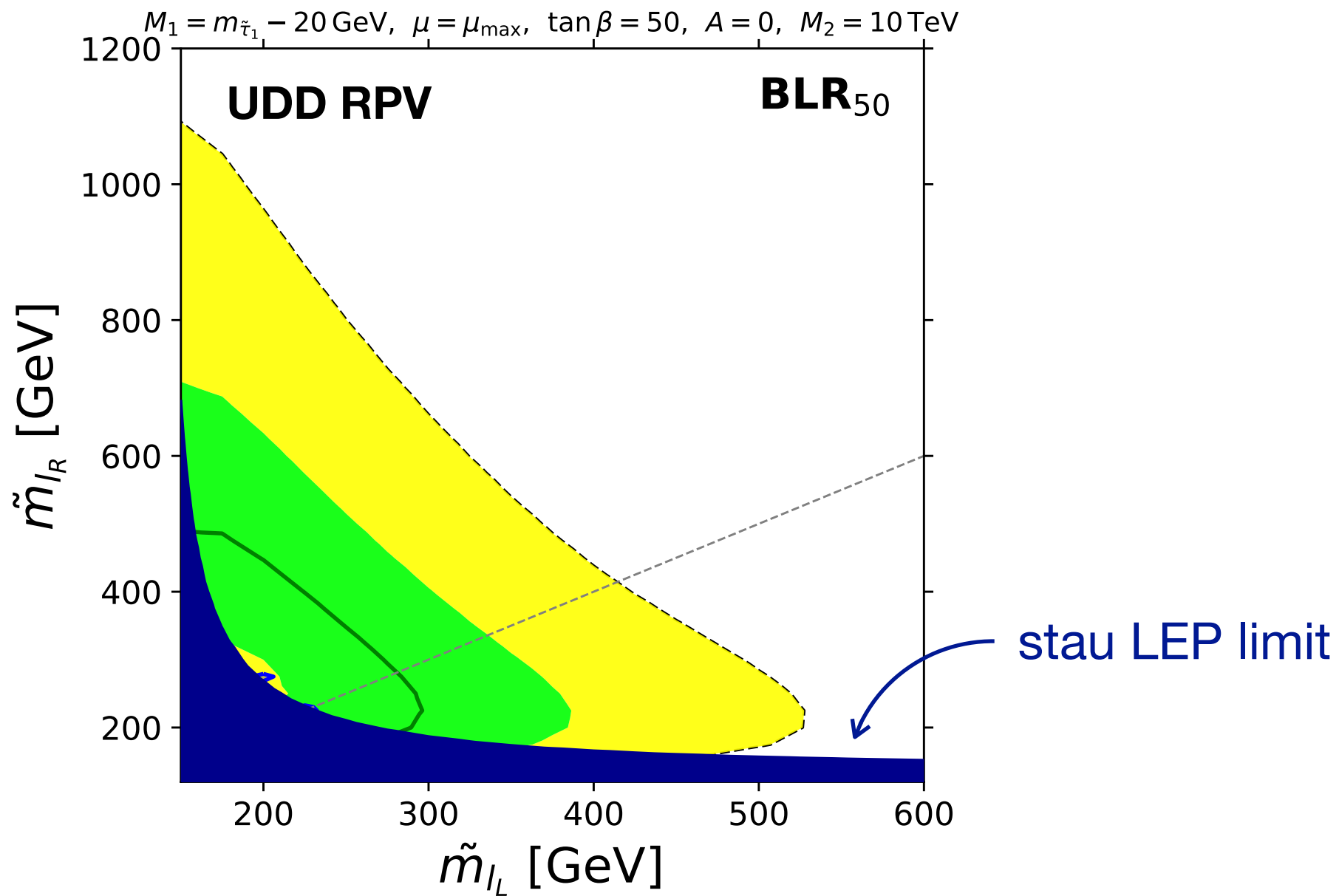
[1709.05406]

ATLAS jets + E_T^{miss}

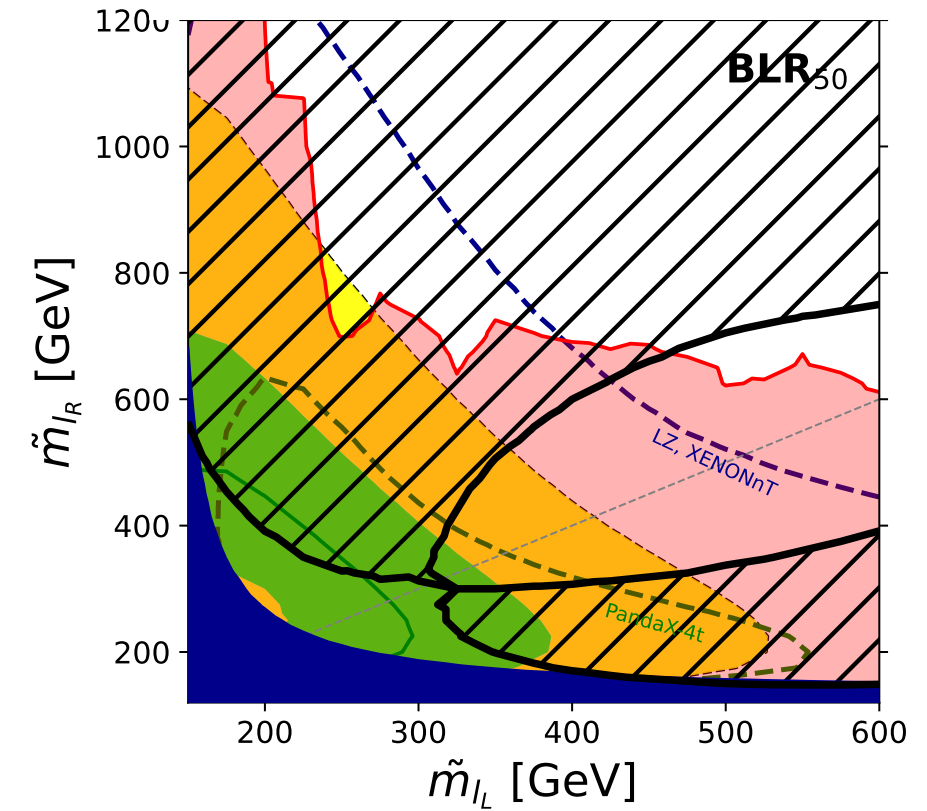
[ATLAS-CONF-2019-040]

UDD RPV

BLR dominated plane

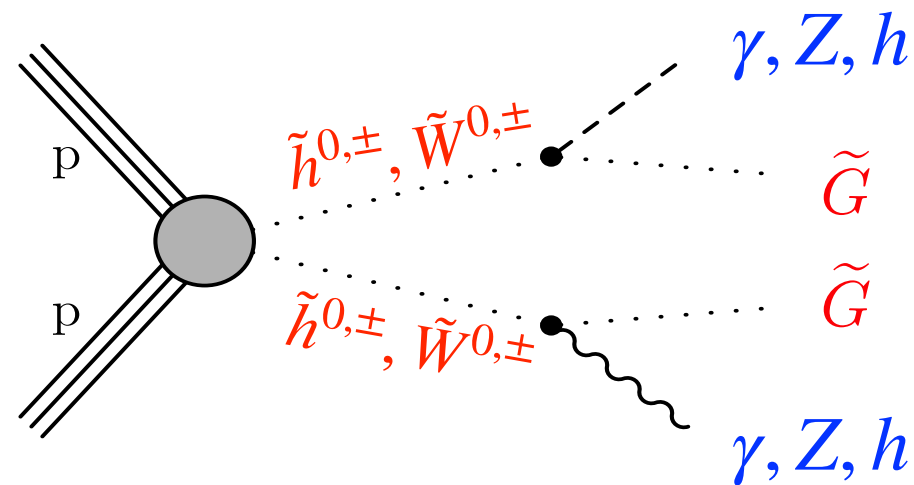


stable neutralino



Gravitino LSP

(general GMSB)



prompt decay into
“massless” gravitino

$$\Gamma(\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma) = |N_{11}c_W + N_{12}s_W|^2 \mathcal{A},$$

$$\Gamma(\tilde{\chi}_1^0 \rightarrow \tilde{G}Z) = \left(|N_{12}c_W - N_{11}s_W|^2 + \frac{1}{2} |N_{13}c_\beta - N_{14}s_\beta|^2 \right) \left(1 - \frac{m_Z^2}{m_{\tilde{\chi}_1^0}^2} \right)^4 \mathcal{A},$$

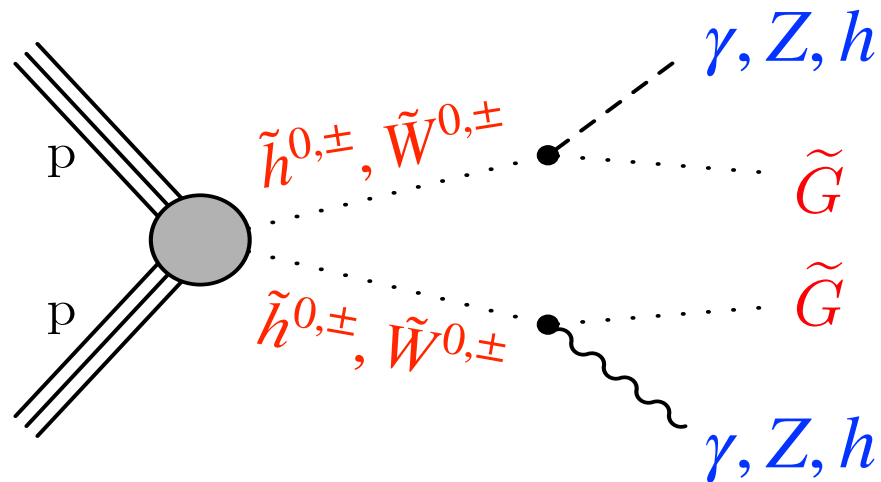
$$\Gamma(\tilde{\chi}_1^0 \rightarrow \tilde{G}h) = \frac{1}{2} |N_{13}c_\beta + N_{14}s_\beta|^2 \left(1 - \frac{m_h^2}{m_{\tilde{\chi}_1^0}^2} \right)^4 \mathcal{A},$$

$$\mathcal{A} = \frac{m_{\tilde{\chi}_1^0}^5}{16\pi m_{3/2}^2 M_{\text{pl}}^2} \sim \frac{1}{0.3 \text{ mm}} \left(\frac{m_{\tilde{\chi}_1^0}}{100 \text{ GeV}} \right)^5 \left(\frac{m_{3/2}}{10 \text{ eV}} \right)^{-2}$$

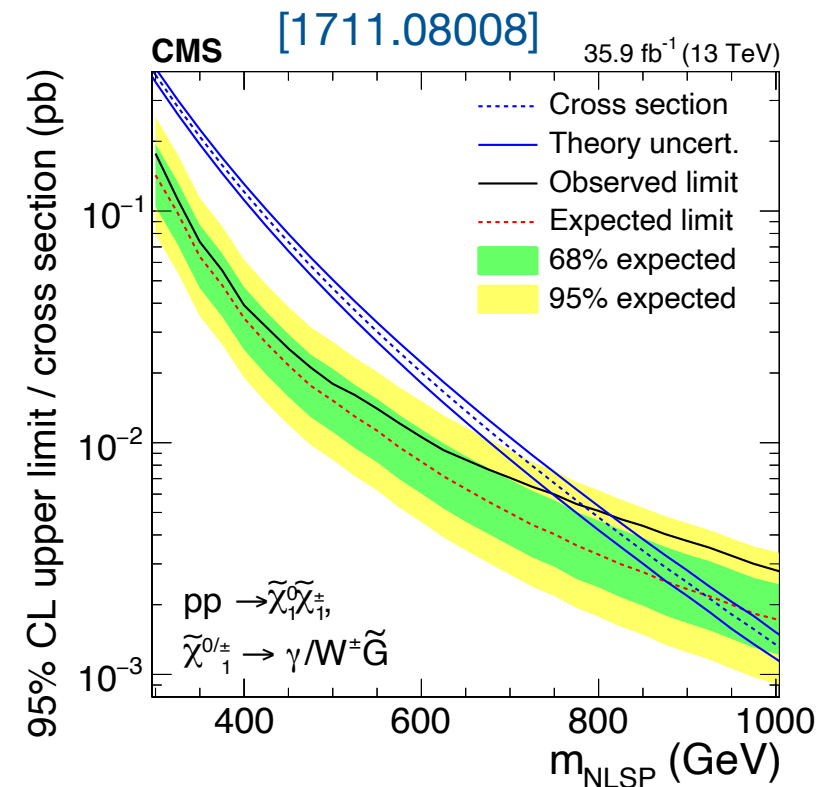
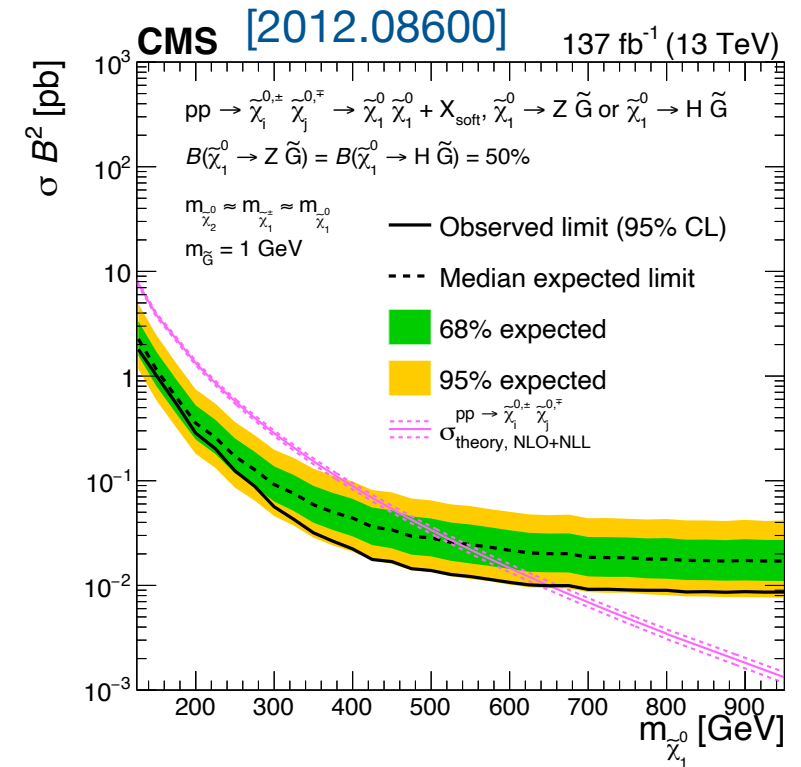
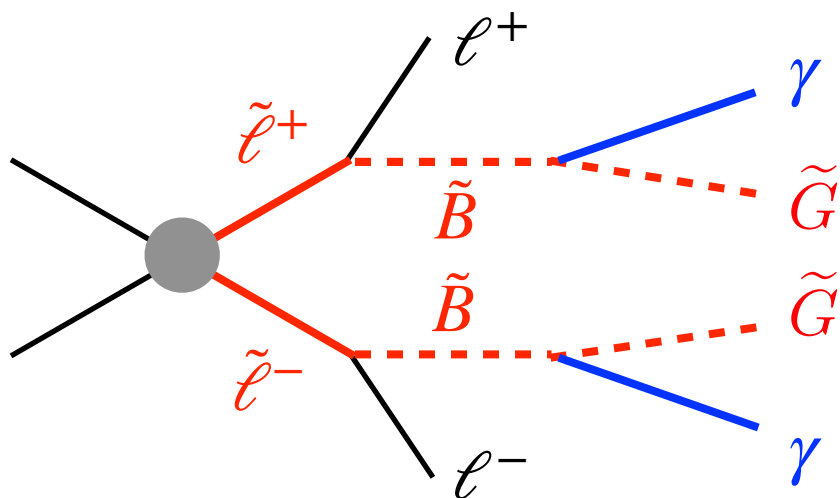
Gravitino LSP

We assume a **massless** gravitino ($m < 1\text{ GeV}$) and **prompt** neutralino decay.

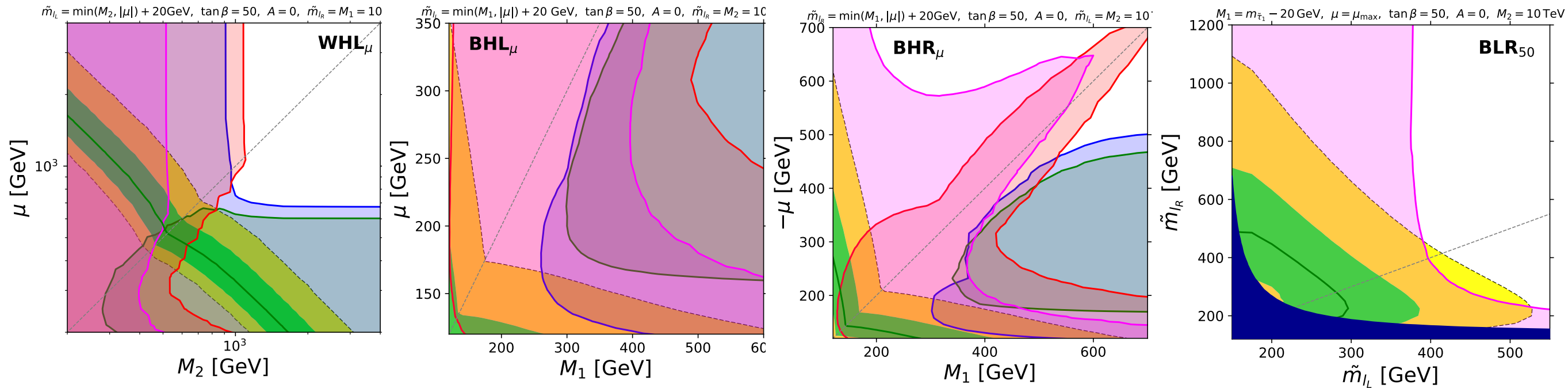
Wino, Higgsino LSP



Bino LSP



Gravitino LSP

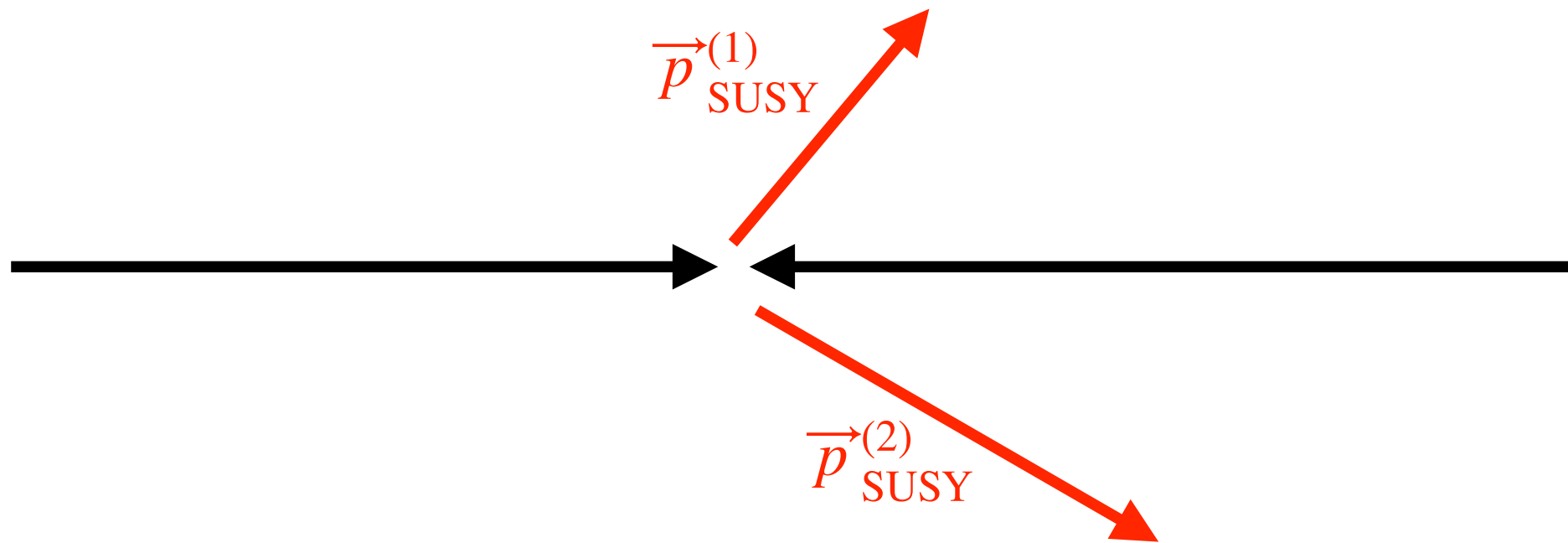


$[\gamma\tilde{G}][\gamma\tilde{G}]$ CMS $\gamma + E_T$ [1711.08008]
 $[\gamma\tilde{G}][Z(h)\tilde{G}]$ CMS $\gamma + E_T$ [1711.08008]

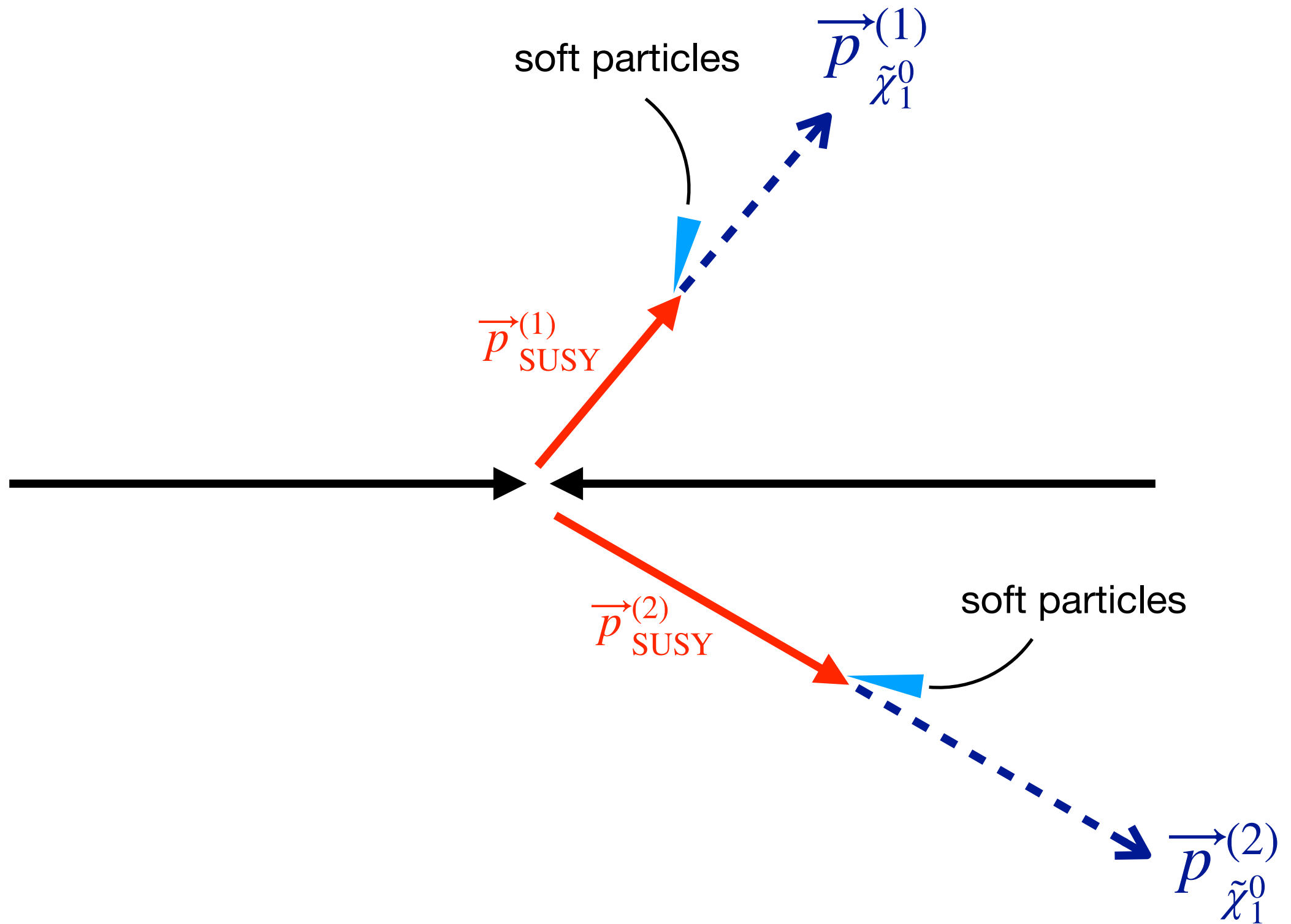
$[Z\tilde{G}][Z\tilde{G}]$ CMS $\ell^+\ell^-$ [2012.08600]
 $[h\tilde{G}][Z\tilde{G}]$ CMS $\ell^+\ell^-$ [2012.08600]

[(massless) gravitino LSP + neutralino NLSP] cannot explain muon g-2

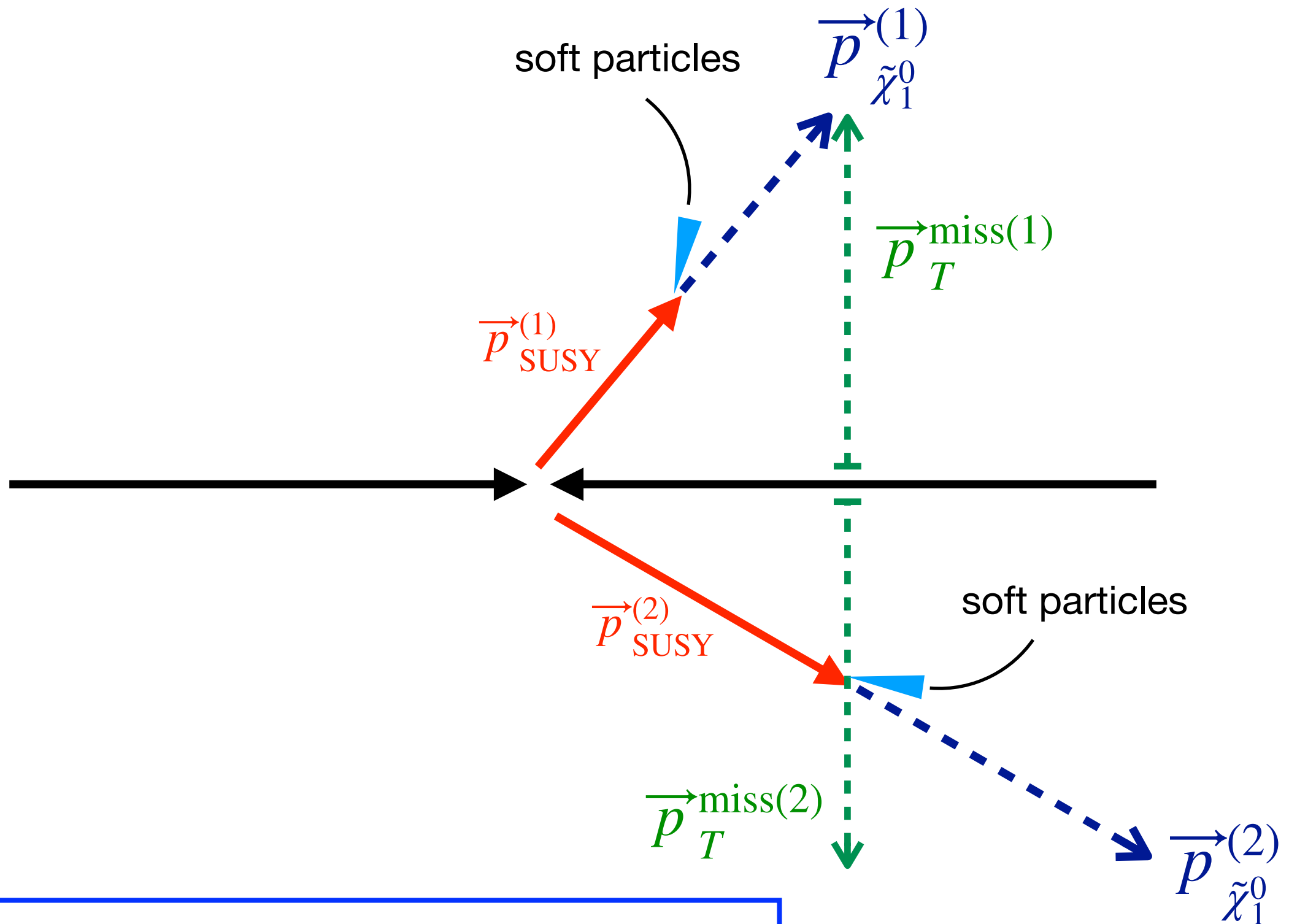
MSSM: compressed spectrum



MSSM: compressed spectrum



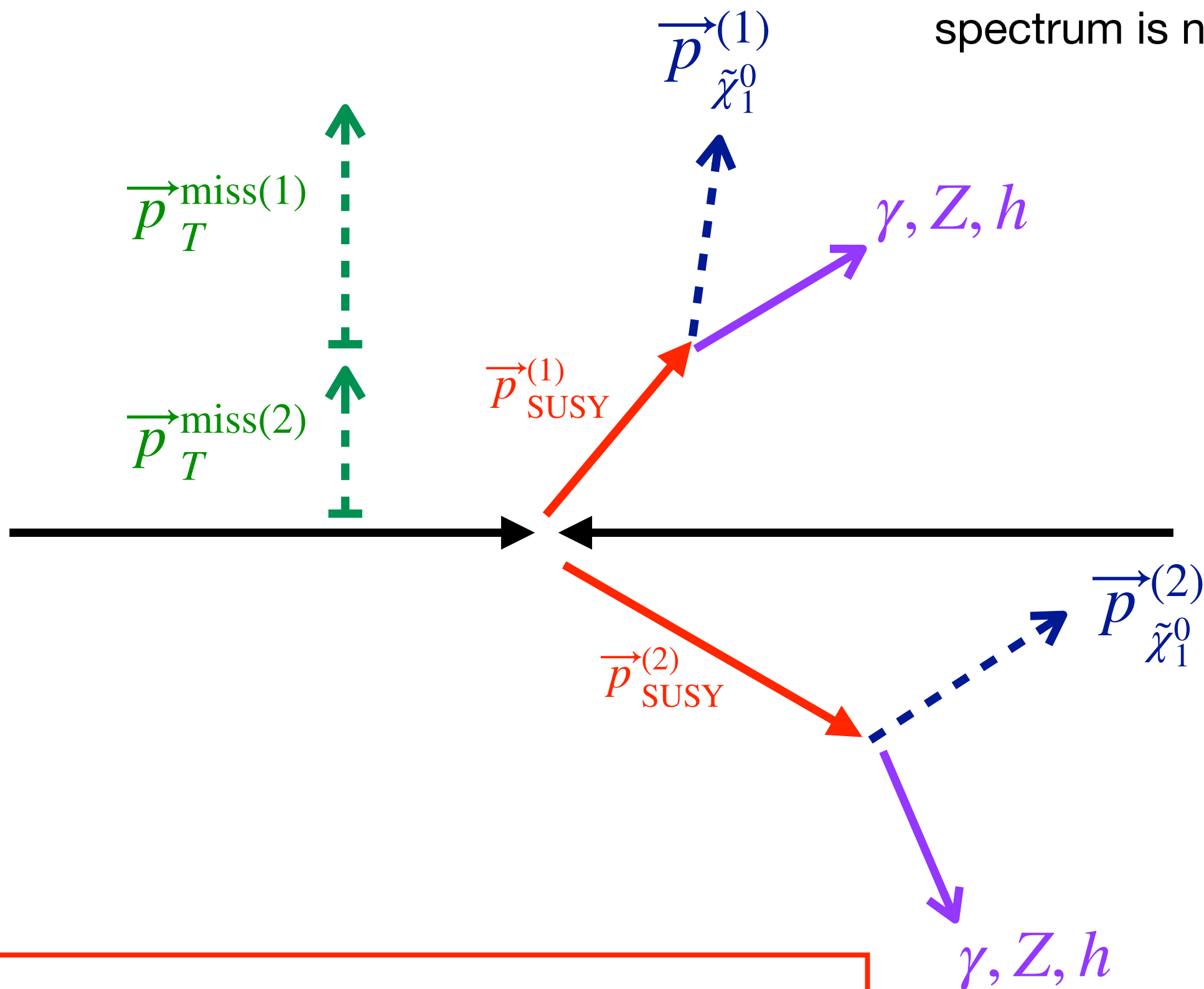
MSSM: compressed spectrum



$$\vec{p}_T^{\text{miss}} = \vec{p}_T^{\text{miss}(1)} + \vec{p}_T^{\text{miss}(2)} \simeq 0$$

Gravitino LSP

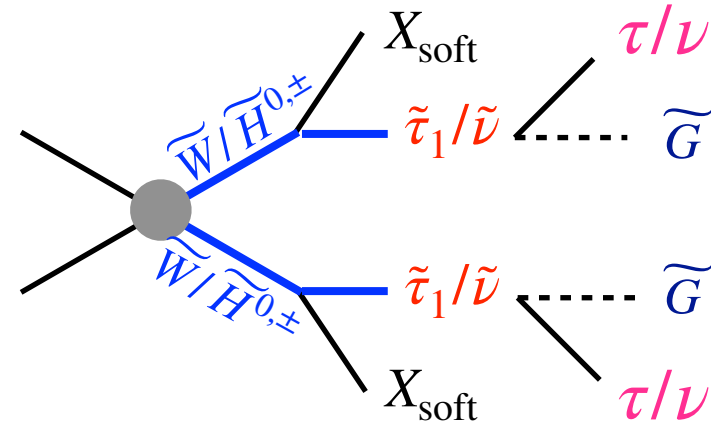
Gravitino is almost massless \rightarrow compressed spectrum is not possible



$$\vec{p}_T^{\text{miss}} = \vec{p}_T^{\text{miss}(1)} + \vec{p}_T^{\text{miss}(2)} \text{ is large}$$

Gravitino LSP with slepton NLSP

$\tilde{\tau}_1/\tilde{\nu}$ NLSP:

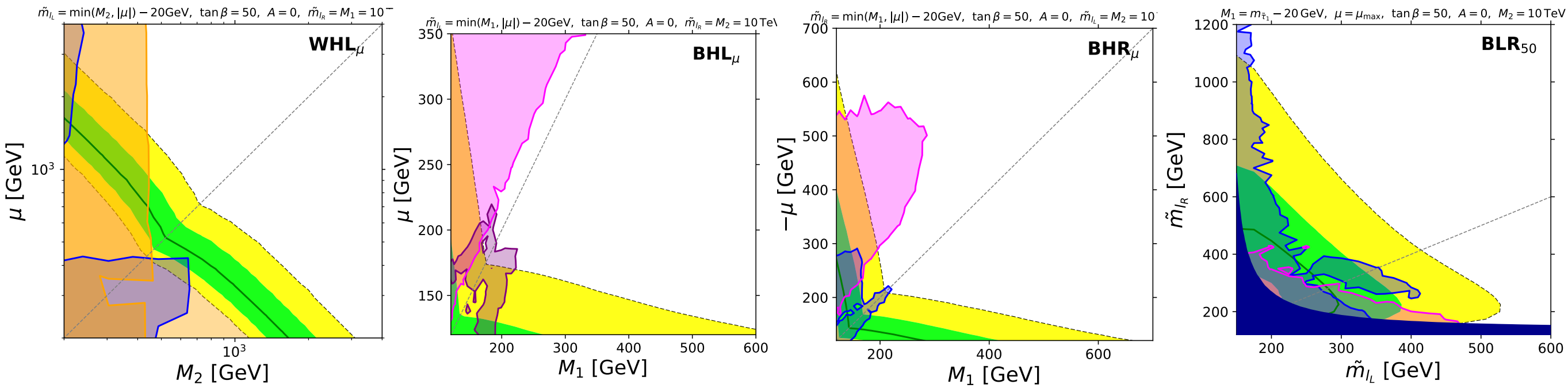


$$\left. \begin{array}{l}
 (M_2 \text{ vs } \mu) \text{ with } \tilde{m}_{l_L} = \min(M_2, \mu) + 20 \text{ GeV} \implies m_{l_L} = \min(M_2, \mu) - 20 \text{ GeV} \\
 \text{BHL plane:} \\
 (M_1 \text{ vs } \mu) \text{ with } \tilde{m}_{l_L} = \min(M_1, \mu) + 20 \text{ GeV} \implies m_{l_L} = \min(M_2, \mu) - 20 \text{ GeV}
 \end{array} \right\} \tilde{\nu}_L \text{ NLSP}$$

$$\left. \begin{array}{l}
 \text{BHR plane:} \\
 (M_1 \text{ vs } \mu) \text{ with } \tilde{m}_{l_R} = \min(M_1, |\mu|) + 20 \text{ GeV} \implies m_{l_R} = \min(M_1, \mu) - 20 \text{ GeV}
 \end{array} \right\} \tilde{e}_R, \tilde{\mu}_R, \tilde{\tau}_R \text{ NLSP}$$

$$\left. \begin{array}{l}
 \text{BLR plane:} \\
 (\tilde{m}_{l_L} \text{ vs } \tilde{m}_{l_R}) \text{ with } M_1 = m_{\tilde{\tau}_1} - 20 \text{ GeV} \implies M_1 = m_{\tilde{\tau}_1} + 20 \text{ GeV}
 \end{array} \right\} \tilde{\tau}_1 \text{ NLSP}$$

Gravitino LSP with slepton NLSP



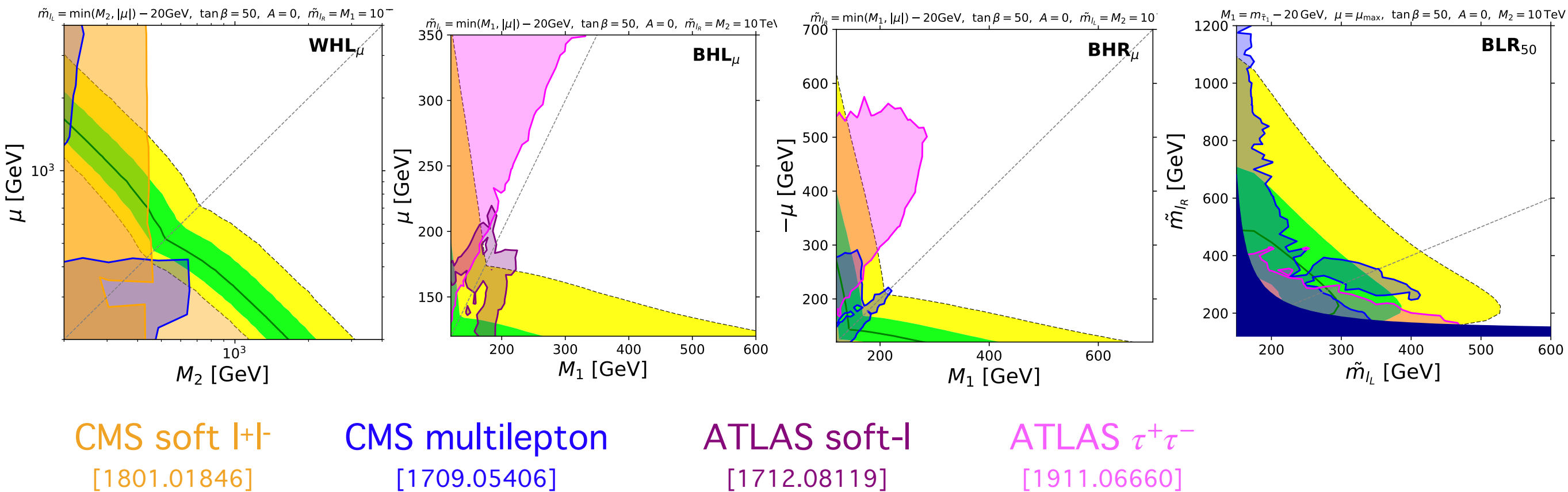
CMS soft $l+l^-$
[1801.01846]

CMS multilepton
[1709.05406]

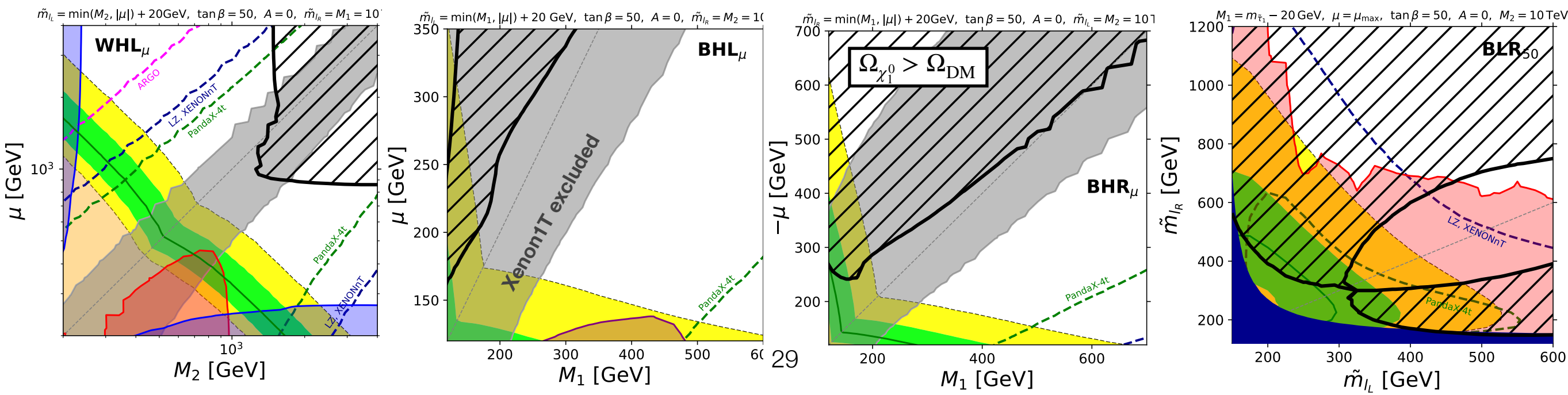
ATLAS soft-l
[1712.08119]

ATLAS $\tau^+\tau^-$
[1911.06660]

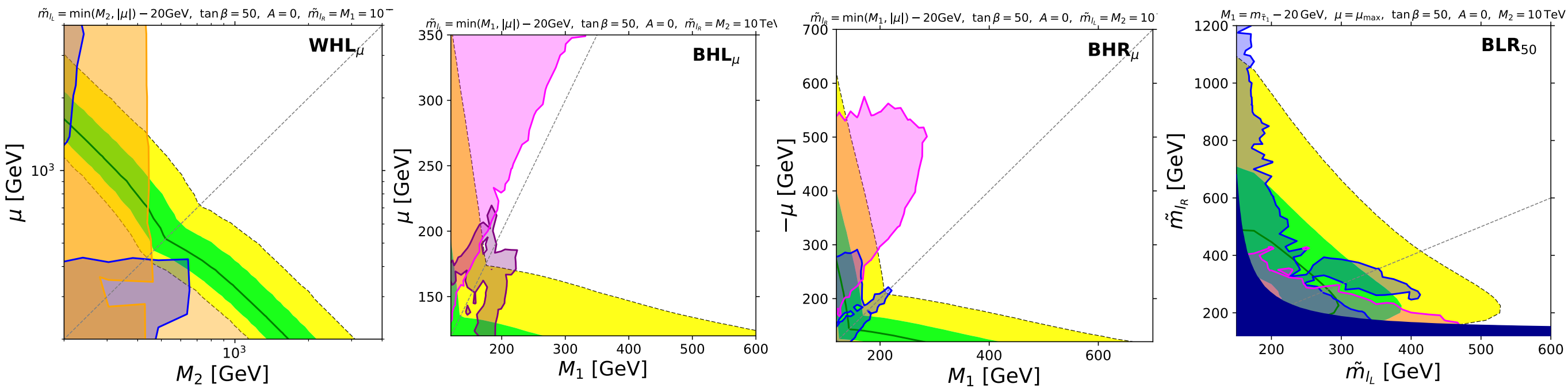
Gravitino LSP with slepton NLSP



MSSM with stable neutralino:



Gravitino LSP with slepton NLSP



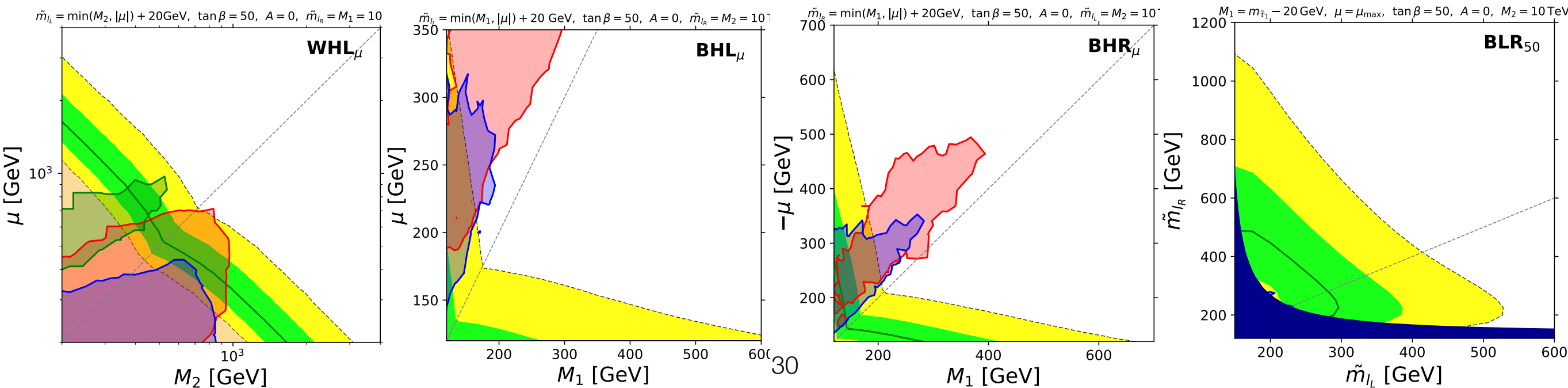
CMS soft $l+l-$
[1801.01846]

CMS multilepton
[1709.05406]

ATLAS soft- l
[1712.08119]

ATLAS $\tau^+\tau^-$
[1911.06660]

UDD RPV :



Summary

- SUSY $(g-2)_\mu$ solution requires either:

1. a large Wino(Bino)-Higgsino mixing or
2. a Bino-like LSP with large chirality flip in the slepton sector

These options are strongly constrained by the DM-DD and LHC due to large missing energy signature

- Those constraints are there because neutralino is stable.
- By decaying neutralino into 3-quarks via UDD RPV, the allowed parameter region opens up.
- Making neutralino decay into a massless gravitino **doesn't help**, because it becomes even more difficult to hide the missing energy signature at LHC.
- Gravitino LSP with slepton/stau NLSP is less constrained.



Norway grants

The research leading to the results presented in this talk has received funding from the Norwegian Financial Mechanism for years 2014-2021, grant nr 2019/34/H/ST2/00707



Understanding the Early Universe:
interplay of theory and collider experiments

Joint research project between the University of Warsaw & University of Bergen

Parameter planes definition

name	axes	range [TeV]	other parameters	$\tan \beta$
WHL $_{\mu}$	(M_2, μ)	$([0.2, 4], [0.2, 4])$	$\tilde{m}_{l_L} = \min(M_2, \mu) + 20 \text{ GeV}, M_1 = \tilde{m}_{l_R} = 10 \text{ TeV}$	50
WHL $_L$	(M_2, \tilde{m}_{l_L})	$([0.2, 4], [0.2, 2])$	$\mu = \min(M_2, \tilde{m}_{l_L}) - 20 \text{ GeV}, M_1 = \tilde{m}_{l_R} = 10 \text{ TeV}$	50
BHL $_{\mu}$	(M_1, μ)	$([0.12, 0.6], [0.12, 0.35])$	$\tilde{m}_{l_L} = \min(M_1, \mu) + 20 \text{ GeV}, M_2 = \tilde{m}_{l_R} = 10 \text{ TeV}$	50
BHL $_L$	(M_1, \tilde{m}_{l_L})	$([0.12, 0.8], [0.14, 0.22])$	$\mu = \min(M_1, \tilde{m}_{l_L}) - 20 \text{ GeV}, M_2 = \tilde{m}_{l_R} = 10 \text{ TeV}$	50
BHR $_{\mu}$	(M_1, μ)	$([0.12, 0.7], [0.12, 0.7])$	$\tilde{m}_{l_R} = \min(M_1, \mu) + 20 \text{ GeV}, M_2 = \tilde{m}_{l_L} = 10 \text{ TeV}$	50
BHR $_L$	(M_1, \tilde{m}_{l_R})	$([0.12, 0.8], [0.14, 0.25])$	$-\mu = \min(M_1, \tilde{m}_{l_R}) - 20 \text{ GeV}, M_2 = \tilde{m}_{l_L} = 10 \text{ TeV}$	50
BLR $_{50}$	$(\tilde{m}_{l_L}, \tilde{m}_{l_R})$	$([0.15, 0.6], [0.12, 1.2])$	$M_1 = m_{\tilde{\tau}_1} - 20 \text{ GeV}, \mu = \mu_{\max}, M_2 = 10 \text{ TeV}$	50
BLR $_{10}$	$(\tilde{m}_{l_L}, \tilde{m}_{l_R})$	$([0.15, 0.6], [0.12, 1.2])$	$M_1 = m_{\tilde{\tau}_1} - 20 \text{ GeV}, \mu = \mu_{\max}, M_2 = 10 \text{ TeV}$	10

Table 1: The parameter planes and choices of the other parameters. μ_{\max} is defined as the maximum value allowed by the vacuum stability constraint.

For GMSB we modify the planes to ensure that slepton/stau/sneutrino is the NLSP.

		QED	HVP	EW		
a_μ^{theo}	=	0.00	1165	91	810	(43)
a_μ^{exp}	=	0.00	1165	92	061	(41)

- The deviation is size of the EW correction in SM:

$$a_\mu^{\text{exp}} - a_\mu^{\text{theo}} \simeq (25 \pm 6) \times 10^{-10} \sim \mathcal{O} \left(\Delta a_\mu^{\text{SM,EW}} \right)$$

- We need **very light BSM particles** **OR** **enhancement from couplings**

$$\Delta a_\mu^{\text{BSM}} \sim \Delta a^{\text{SM,EW}} \cdot \underbrace{\left(\frac{m_W^2}{m_{\text{BSM}}^2} \right) \cdot \left(\frac{g_{\text{BSM}}}{g_{\text{SM}}} \right)}_{\mathcal{O}(1)}$$

Chiral ($\tan\beta$) enhancement in SUSY

- (g-2) operator requires chirality flip:

$$\mathcal{L}_{\text{eff}} \ni i\tilde{a}_\mu \cdot \bar{\psi}_L \sigma^{\mu\nu} \psi_R F_{\mu\nu}$$

$$\text{SM: } \tilde{a}_\mu^{\text{SM}} \propto Y_\mu \langle H \rangle = m_\mu$$

$$\vec{\mu} = g \left(\frac{e}{2m} \right) \vec{s}$$

$$a_\mu = \frac{(g-2)}{2} \equiv m_\mu \tilde{a}_\mu$$

Chiral ($\tan\beta$) enhancement in SUSY

- (g-2) operator requires chirality flip:

$$\mathcal{L}_{\text{eff}} \ni i\tilde{a}_\mu \cdot \bar{\psi}_L \sigma^{\mu\nu} \psi_R F_{\mu\nu}$$

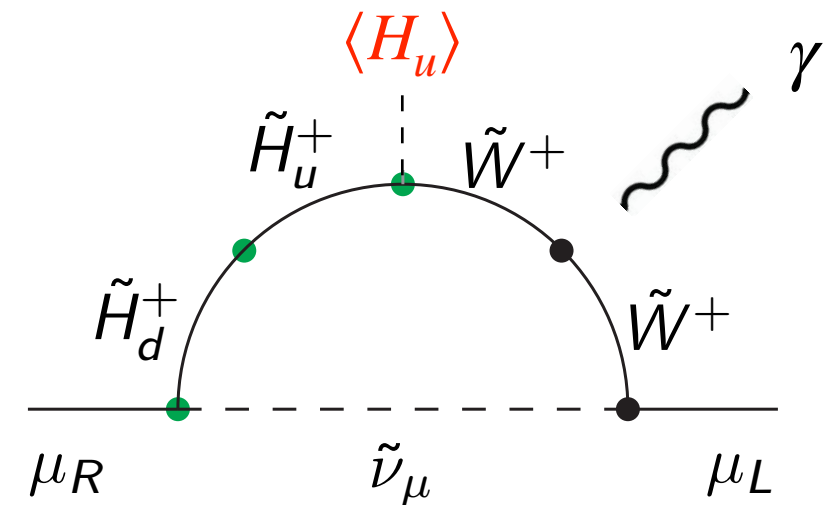
$$\vec{\mu} = g \left(\frac{e}{2m} \right) \vec{s}$$

$$a_\mu = \frac{(g-2)}{2} \equiv m_\mu \tilde{a}_\mu$$

SM: $\tilde{a}_\mu^{\text{SM}} \propto Y_\mu \langle H \rangle = m_\mu$

SUSY: $\Delta \tilde{a}_\mu^{\text{SUSY}} \propto Y_\mu \langle H_u \rangle = m_\mu \cdot \tan\beta$

$$m_\mu = Y_\mu \langle H_d \rangle \quad \tan\beta \equiv \frac{\langle H_u \rangle}{\langle H_d \rangle}$$



$$\langle H_u \rangle^2 + \langle H_d \rangle^2 = \langle H \rangle^2$$

↑
(246 GeV)²

Chiral ($\tan\beta$) enhancement in SUSY

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$$\mathcal{L}_{\text{eff}} \ni i\tilde{a}_\mu \cdot \bar{\psi}_L \sigma^{\mu\nu} \psi_R F_{\mu\nu}$$

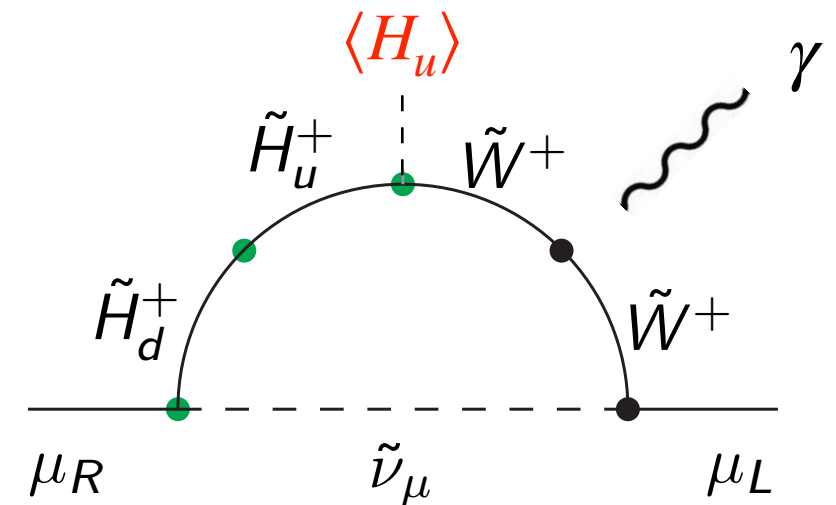
$$\vec{\mu} = g \left(\frac{e}{2m} \right) \vec{s}$$

$$a_\mu = \frac{(g-2)}{2} \equiv m_\mu \tilde{a}_\mu$$

SM: $\tilde{a}_\mu^{\text{SM}} \propto Y_\mu \langle H \rangle = m_\mu$

SUSY: $\Delta\tilde{a}_\mu^{\text{SUSY}} \propto Y_\mu \langle H_u \rangle = m_\mu \cdot \tan\beta$

$$m_\mu = Y_\mu \langle H_d \rangle \quad \tan\beta \equiv \frac{\langle H_u \rangle}{\langle H_d \rangle}$$



$$\langle H_u \rangle^2 + \langle H_d \rangle^2 = \langle H \rangle^2$$

↑
(246 GeV)²

$$\Delta a_\mu^{\text{BSM}} \sim \Delta a^{\text{SM,EW}} \cdot \left(\frac{m_W^2}{m_{\text{SUSY}}^2} \right) \cdot \tan\beta$$

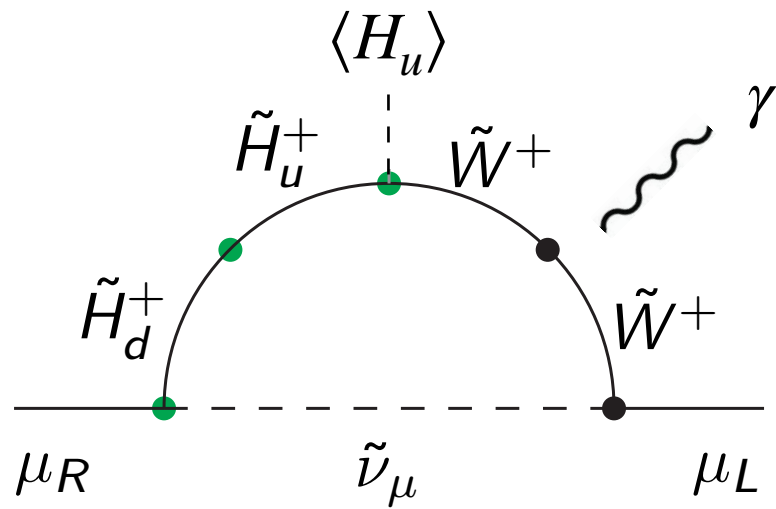
$\tan\beta \in [5 - 60] \rightarrow m_{\text{SUSY}} \in [200 - 600] \text{ GeV}$

- Due to strong LHC constraints, we *decouple coloured SUSY particles* (they do not contribute to $(g-2)_\mu$ anyway).
- a_μ^{SUSY} depends on **5 mass parameters** and **$\tan\beta$** :

$$\begin{array}{l}
 M_1 : \text{Bino mass} \\
 M_2 : \text{Wino mass} \\
 \mu : \text{Higgsino mass}
 \end{array}
 \left(\begin{array}{l}
 m_{\tilde{l}_R} \equiv \widetilde{m}_{\tilde{e}_R}^2 = \widetilde{m}_{\tilde{\mu}_R}^2 = \widetilde{m}_{\tilde{\tau}_R}^2 \\
 m_{\tilde{l}_L} \equiv \widetilde{m}_{\tilde{\nu}_e} = \widetilde{m}_{\tilde{\nu}_\mu} = \widetilde{m}_{\tilde{\nu}_\tau} = \widetilde{m}_{\tilde{e}_L} = \widetilde{m}_{\tilde{\mu}_L} = \widetilde{m}_{\tilde{\tau}_L} \\
 \tan\beta \equiv \langle H_u \rangle / \langle H_d \rangle
 \end{array} \right.$$

no LFV due to universal soft masses: avoid strong constraint from $\mu \rightarrow e \gamma$

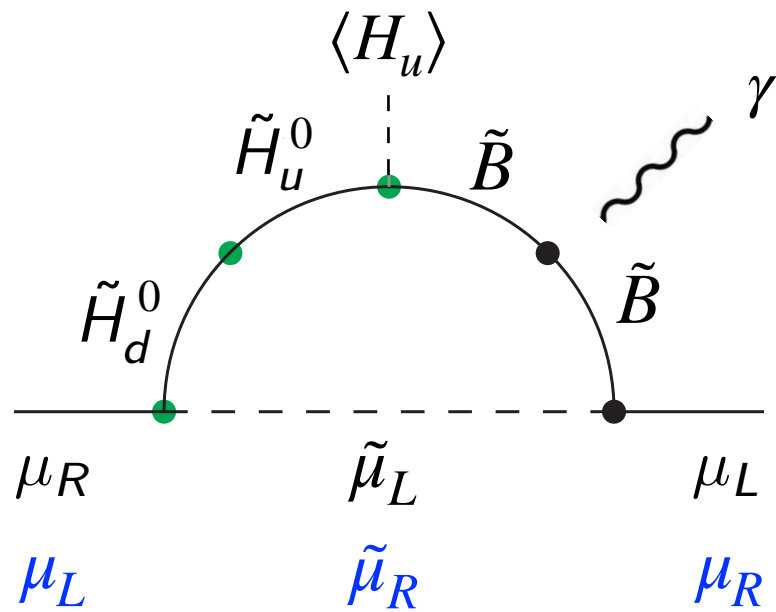
$$\Delta a_\mu^{\text{SUSY}} = \Delta a_\mu^{\text{WHL}} + \Delta a_\mu^{\text{BHL}} + \Delta a_\mu^{\text{BHR}} + \Delta a_\mu^{\text{BLR}}$$



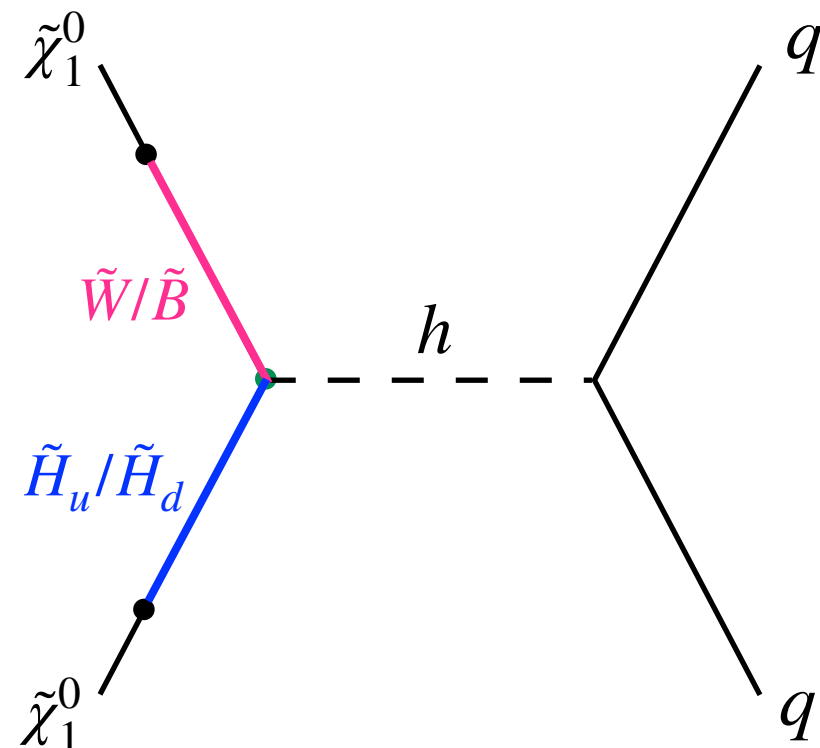
$$\Delta a_\mu^{\text{WHL}}(M_2, \mu, m_{\tilde{l}_L}) = \frac{\alpha_W}{8\pi} \frac{m_\mu^2}{M_2 \mu} \tan \beta \cdot f_W(\{\mathbf{m}\})$$

$$\Delta a_\mu^{\text{BHL}}(M_1, \mu, m_{\tilde{l}_L}) = \frac{\alpha_Y}{8\pi} \frac{m_\mu^2}{M_1 \mu} \tan \beta \cdot f_N(\{\mathbf{m}\})$$

$$\Delta a_\mu^{\text{BHR}}(M_1, \mu, m_{\tilde{l}_R}) = -\frac{\alpha_Y}{8\pi} \frac{m_\mu^2}{M_1 \mu} \tan \beta \cdot f_N(\{\mathbf{m}\})$$



Large gaugino-Higgsino mixing leads to a **large cross-section for DM Direct Detection:**

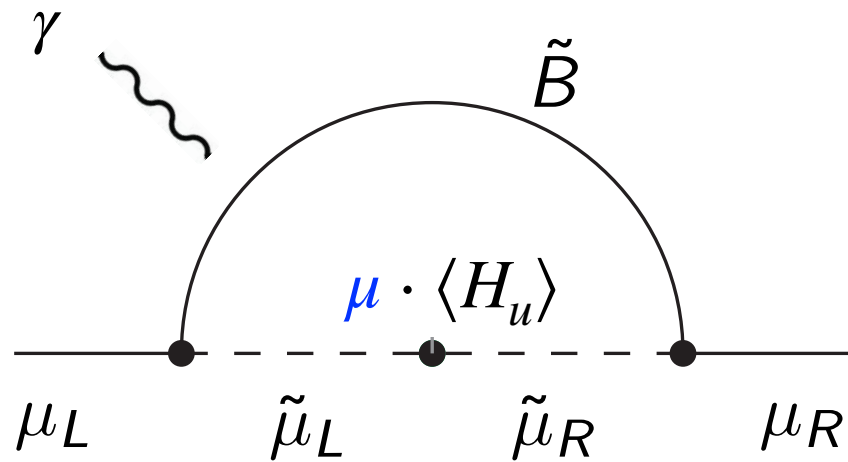


M_1 : Bino (\tilde{B}) mass

M_2 : Wino (\tilde{W}) mass

μ : Higgsino (\tilde{H}_u, \tilde{H}_d) mass

$$\Delta a_\mu^{\text{SUSY}} = \Delta a_\mu^{\text{WHL}} + \Delta a_\mu^{\text{BHL}} + \Delta a_\mu^{\text{BHR}} + \Delta a_\mu^{\text{BLR}}$$



$$\Delta a_\mu^{\text{BLR}}(M_1, m_{\tilde{\ell}_L}, m_{\tilde{\ell}_R}; \mu) = \frac{\alpha_Y m_\mu^2 M_1 \mu}{4\pi m_{\mu_L}^2 m_{\mu_R}^2} \tan \beta \cdot f_{\text{BLR}}(\{\mathbf{m}\})$$

↑
large μ needed

Constraints:

❖ Stau mass² becomes negative or too small!

- charge breaking vacuum: $m_{\text{stau}1}^2 > 0$

- LEP bound: $m_{\text{stau}1} > 90 \text{ GeV}$

- stau LSP: $m_{\text{stau}1} > m_{\text{neutralino}1}$

- Vacuum (meta-)stability:

$$(\tilde{\tau} \text{ mass matrix}) \sim \begin{pmatrix} m_{\tilde{\tau}_R}^2 & Y_\tau \mu \langle H_u \rangle \\ Y_\tau \mu \langle H_u \rangle & m_{\tilde{\tau}_L}^2 \end{pmatrix}$$

$$|m_{\tilde{\ell}_{LR}}^2| \leq \left[1.01 \times 10^2 \text{ GeV} \sqrt{m_{\tilde{\ell}_L} m_{\tilde{\ell}_R}} + 1.01 \times 10^2 \text{ GeV} (m_{\tilde{\ell}_L} + 1.03 m_{\tilde{\ell}_R}) - 2.27 \times 10^4 \text{ GeV}^2 + \frac{2.97 \times 10^6 \text{ GeV}^3}{m_{\tilde{\ell}_L} + m_{\tilde{\ell}_R}} - 1.14 \times 10^8 \text{ GeV}^4 \left(\frac{1}{m_{\tilde{\ell}_L}^2} + \frac{0.983}{m_{\tilde{\ell}_R}^2} \right) \right]$$

[Kitahara, Yoshinaga 13]; [Endo, Hamaguchi, Kitahara, Yoshinaga 13]

❖ **Overproduction of Bino-like neutralinos** in the early universe: $\Omega_{\tilde{\chi}_1^0} < \Omega_{\text{DM}}$

slepton-coannihilation needed $\Rightarrow m_{\text{slepton}} \sim m_{\text{Bino}}$

Unstable Neutralino (Gravitino, RPV)

$$\Delta a_\mu^{\text{SUSY}} = \Delta a_\mu^{\text{WHL}} + \Delta a_\mu^{\text{BHL}} + \Delta a_\mu^{\text{BHR}} + \Delta a_\mu^{\text{BLR}}$$

$$\Delta a_\mu^{\text{WHL}}(M_2, \mu, m_{\tilde{l}_L})$$

$$\Delta a_\mu^{\text{BHL}}(M_1, \mu, m_{\tilde{l}_L})$$

$$\Delta a_\mu^{\text{BHR}}(M_1, \mu, m_{\tilde{l}_R})$$

Higgsino, one gaugino, one slepton all must be light:

⇒ LHC constraint with large E_T ← Modified

gaugino-Higgsino mixing ⇒ ~~DM direct detection~~

$$\Delta a_\mu^{\text{BLR}}(M_1, m_{\tilde{l}_L}, m_{\tilde{l}_R}; \mu)$$

↑
large

Bino and both L and R sleptons must be light:

⇒ LHC constraint with large E_T ← Modified

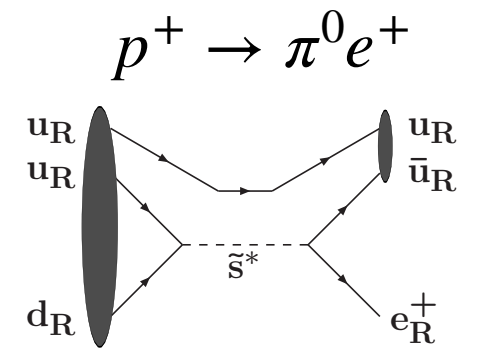
⇒ ~~Bino abundance~~ $\Omega_{\chi_1^0} < \Omega_{\text{DM}}$

⇒ Charged LSP, Vacuum stability

R-Parity Violation; UDD

$$W_{\text{RPV}} = \underbrace{\lambda''_{ijk} U_i^c D_j^c D_k^c}_{\cancel{B}} + \underbrace{\lambda_{ijk} L_i L_j E_k^c + \lambda'_{ijk} L_i Q_j D_k^c + \kappa_i L_i H_u}_{\cancel{L}}$$

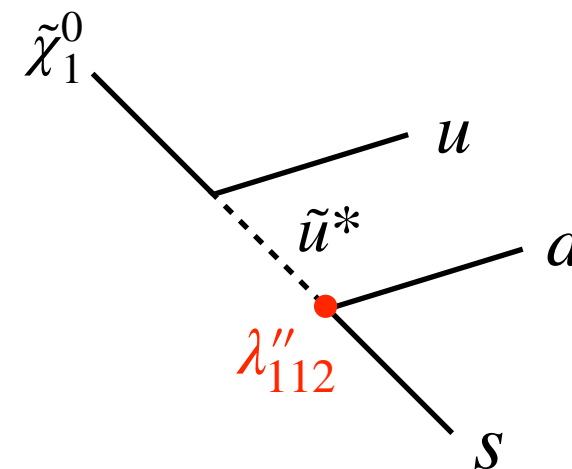
- Allowing both **B** and **L** violation leads to a rapid proton decay:
- We introduce only the **UDD** operator with: $\lambda''_{112} \neq 0$
- Constraint from K0-K0bar mixing can easily be satisfied:



$$|\lambda''_{112} \lambda''_{123}| \lesssim 2.8 \times 10^{-2} \left(\frac{m_{\tilde{s}_R, \tilde{u}_R}}{1 \text{ TeV}} \right) \quad [1810.08228]$$

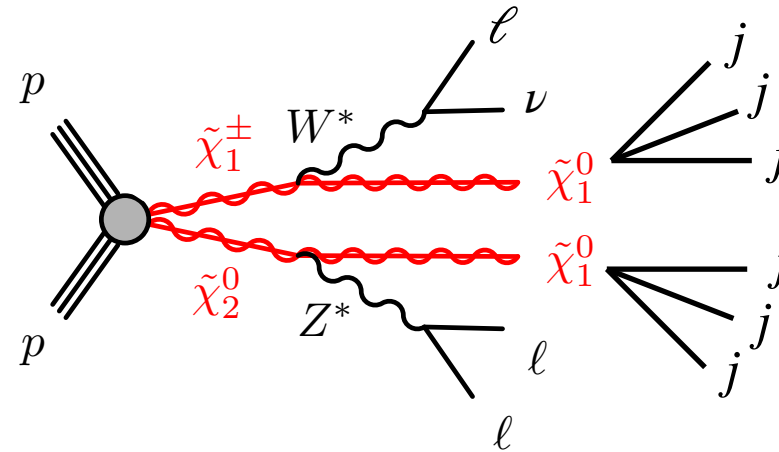
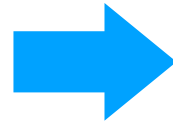
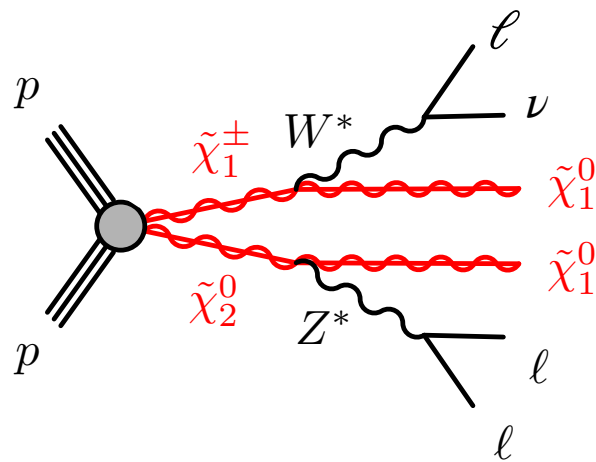
$$|\lambda''_{112} \lambda''_{113}| \lesssim 1.2 \times 10^{-1} \left(\frac{m_{\tilde{d}_R, \tilde{u}_R}}{1 \text{ TeV}} \right)$$

- **LHC signature is the most challenging:**
no leptons, no b-jets in the neutralino decay



R-Parity Violation; UDD

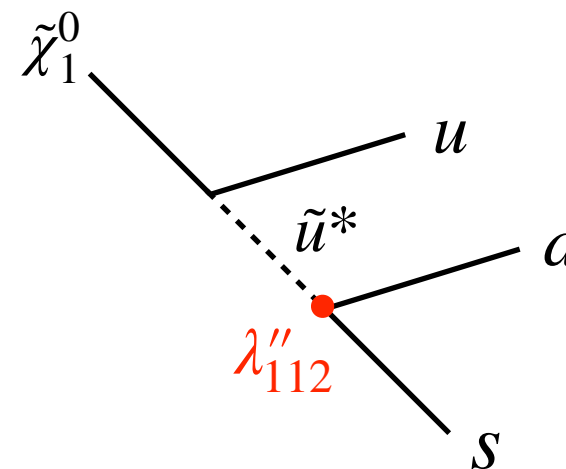
neutralino
LSP



RPV
(UDD-type)

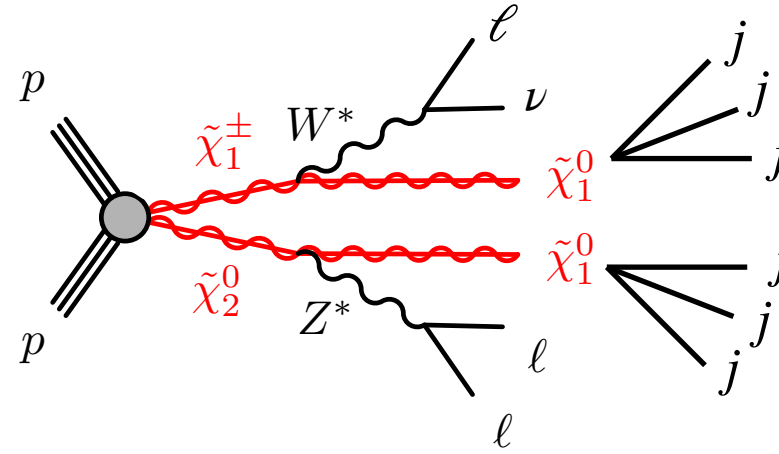
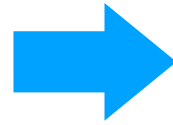
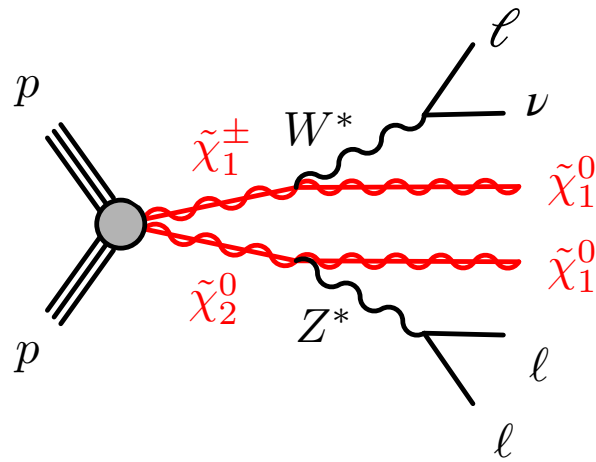
No missing energy, but multi-jet

- LHC signature is the most challenging:
no leptons, no b-jets in the neutralino decay



R-Parity Violation; UDD

neutralino
LSP

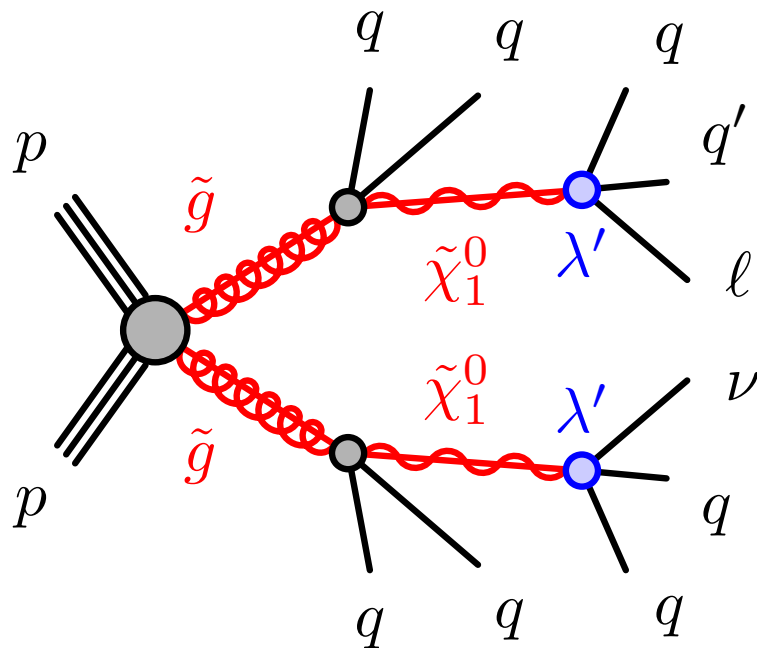


RPV
(UDD-type)

No missing energy, but multi-jet

- There exist ATLAS and CMS analyses sensitive to such final states:

ATLAS [2106.09609]



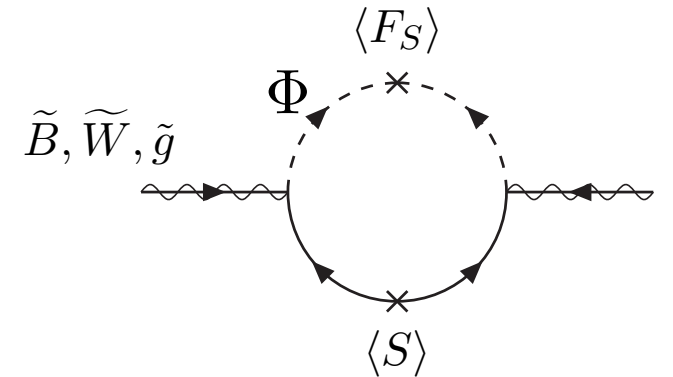
CMS [1709.05406]

Bin	Final state	Definition
1	2 SS leptons	0 jets, $M_T > 100$ GeV and $p_T^{\text{miss}} > 140$ GeV
2	2 SS leptons	1 jet, $M_T < 100$ GeV, $p_T^{\ell\ell} < 100$ GeV and $p_T^{\text{miss}} > 200$ GeV
3	3 light leptons	$M_T > 120$ GeV and $p_T^{\text{miss}} > 200$ GeV
4	3 light leptons	$p_T^{\text{miss}} > 250$ GeV
5	2 light leptons and 1 tau	$M_{T2}(\ell_1, \tau) > 50$ GeV and $p_T^{\text{miss}} > 200$ GeV
6	1 light lepton and 2 taus	$M_{T2}(\ell, \tau_1) > 50$ GeV and $p_T^{\text{miss}} > 200$ GeV
7	1 light lepton and 2 taus	$p_T^{\text{miss}} > 75$ GeV
8	more than 3 leptons	$p_T^{\text{miss}} > 200$ GeV

Gravitino LSP

- In the gauge-mediated SUSY breaking (GMSB) scenario, **light gravitino is motivated by naturalness:**

$$\delta m_h^2 \propto m_{SUSY}^2 \ln \left(\frac{\Lambda_{\text{mess}}}{M_{\text{PL}}} \right) \quad m_{3/2} = \frac{4\pi}{\sqrt{3}\alpha_W} M_2 \frac{\Lambda_{\text{mess}}}{M_{\text{PL}}}$$



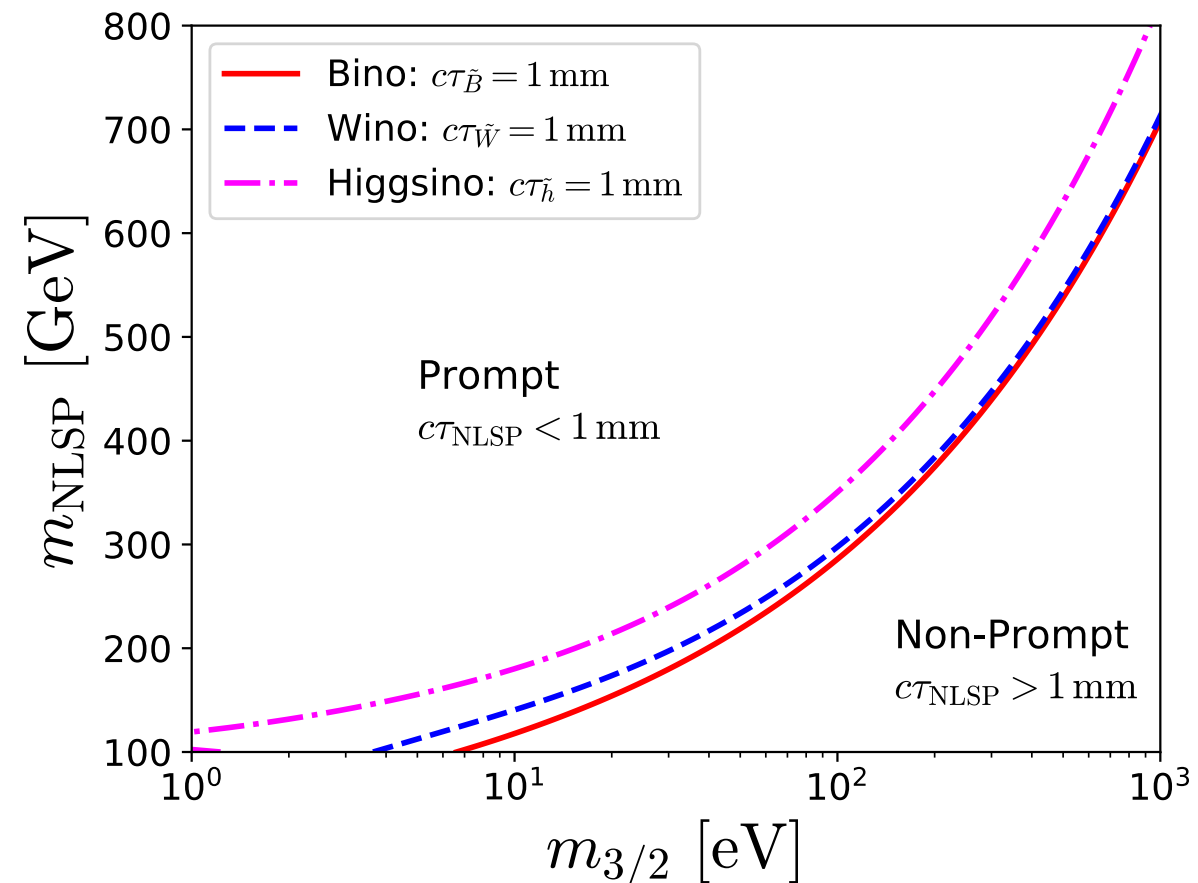
- The decay rate of the NLSP neutralino into the gravitino can be calculated. For light gravitinos ($< 10\text{-}100$ eV), the **neutralino decays are prompt.**

$$\Gamma(\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma) = |N_{11}c_W + N_{12}s_W|^2 \mathcal{A},$$

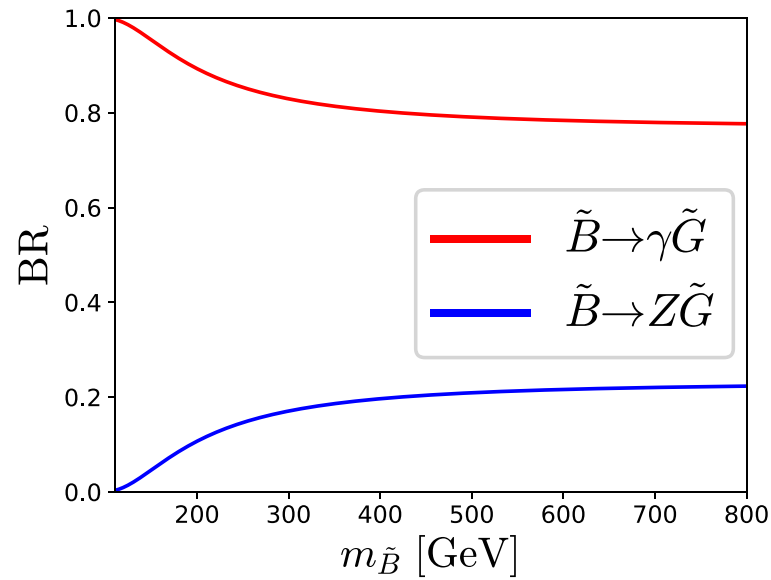
$$\Gamma(\tilde{\chi}_1^0 \rightarrow \tilde{G}Z) = \left(|N_{12}c_W - N_{11}s_W|^2 + \frac{1}{2}|N_{13}c_\beta - N_{14}s_\beta|^2 \right) \left(1 - \frac{m_Z^2}{m_{\tilde{\chi}_1^0}^2} \right)^4 \mathcal{A},$$

$$\Gamma(\tilde{\chi}_1^0 \rightarrow \tilde{G}h) = \frac{1}{2}|N_{13}c_\beta + N_{14}s_\beta|^2 \left(1 - \frac{m_h^2}{m_{\tilde{\chi}_1^0}^2} \right)^4 \mathcal{A},$$

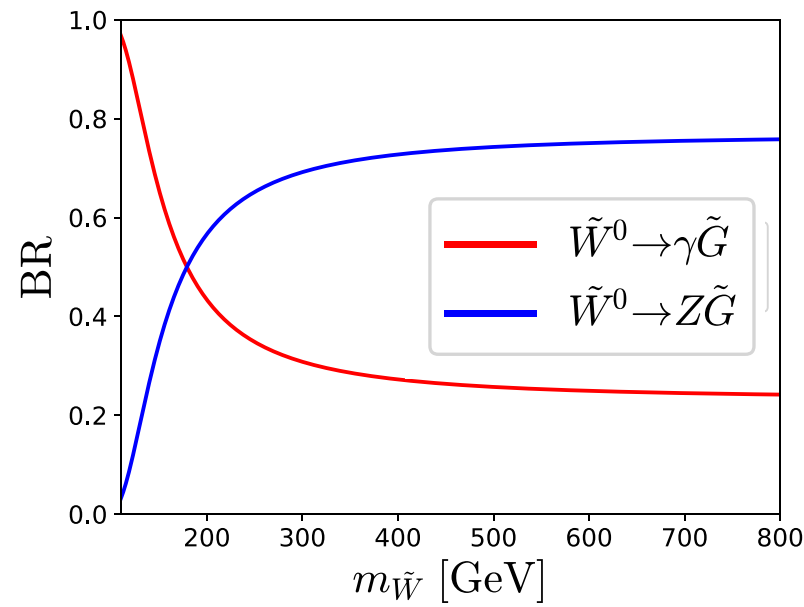
$$\mathcal{A} = \frac{m_{\tilde{\chi}_1^0}^5}{16\pi m_{3/2}^2 M_{\text{pl}}^2} \sim \frac{1}{0.3 \text{ mm}} \left(\frac{m_{\tilde{\chi}_1^0}}{100 \text{ GeV}} \right)^5 \left(\frac{m_{3/2}}{10 \text{ eV}} \right)^{-2}$$



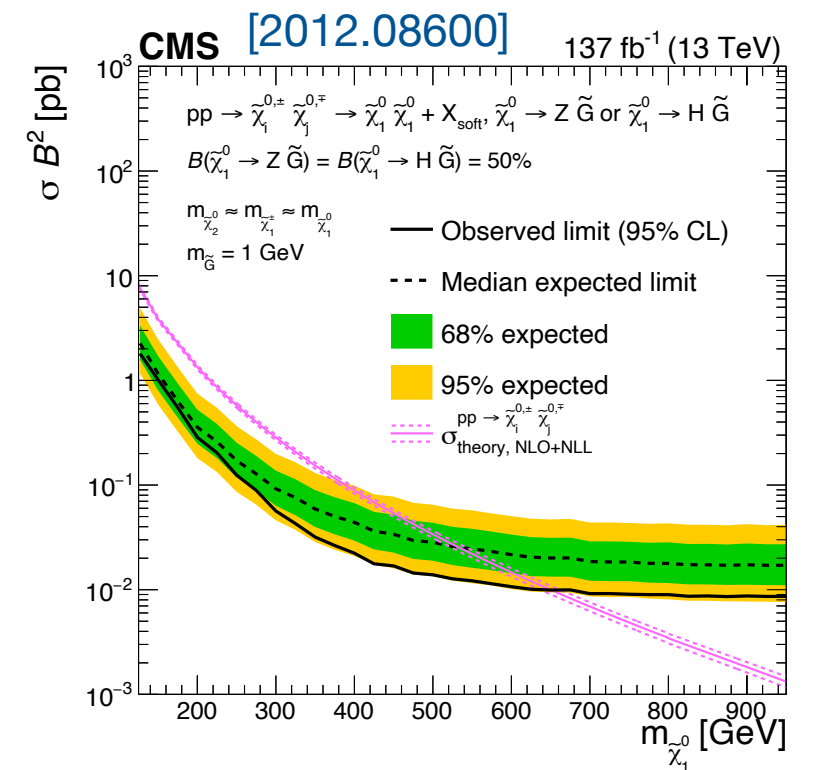
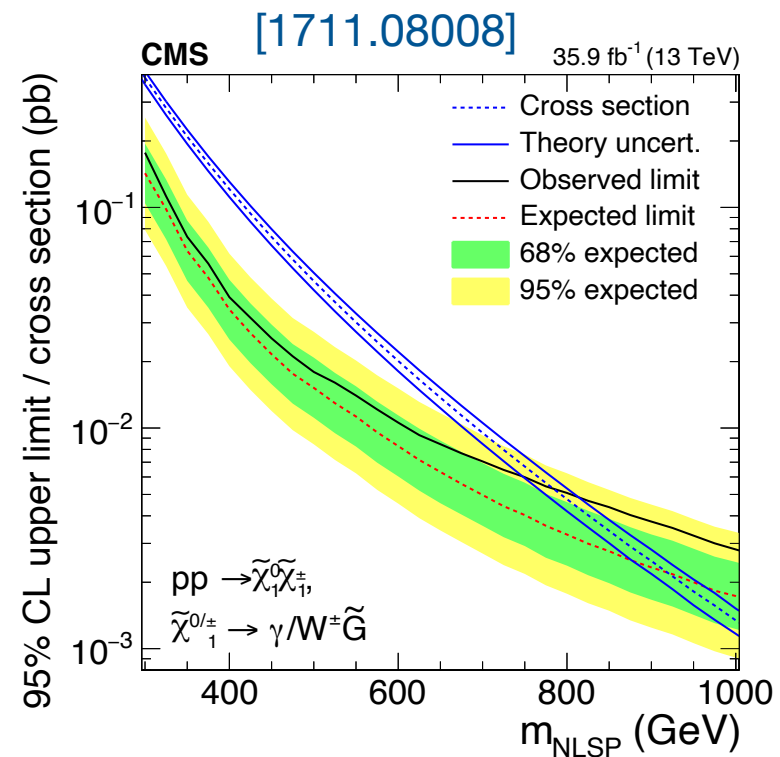
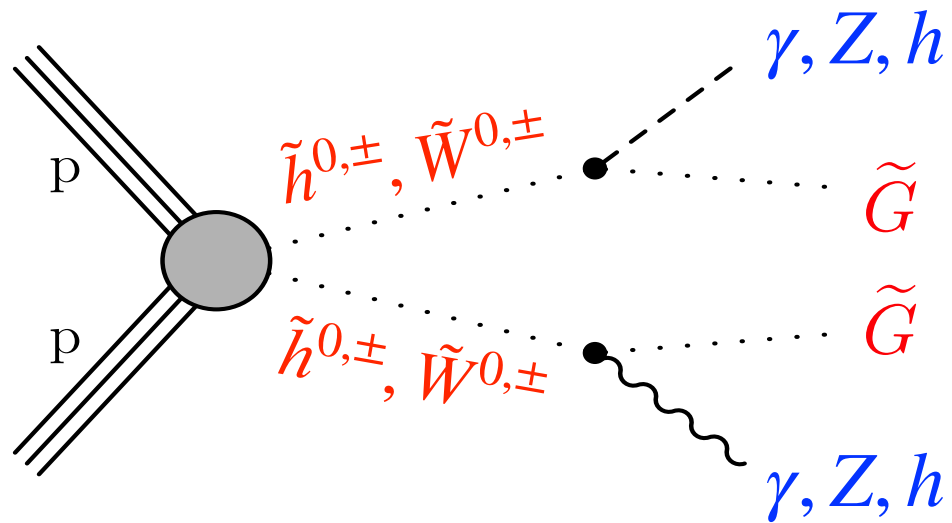
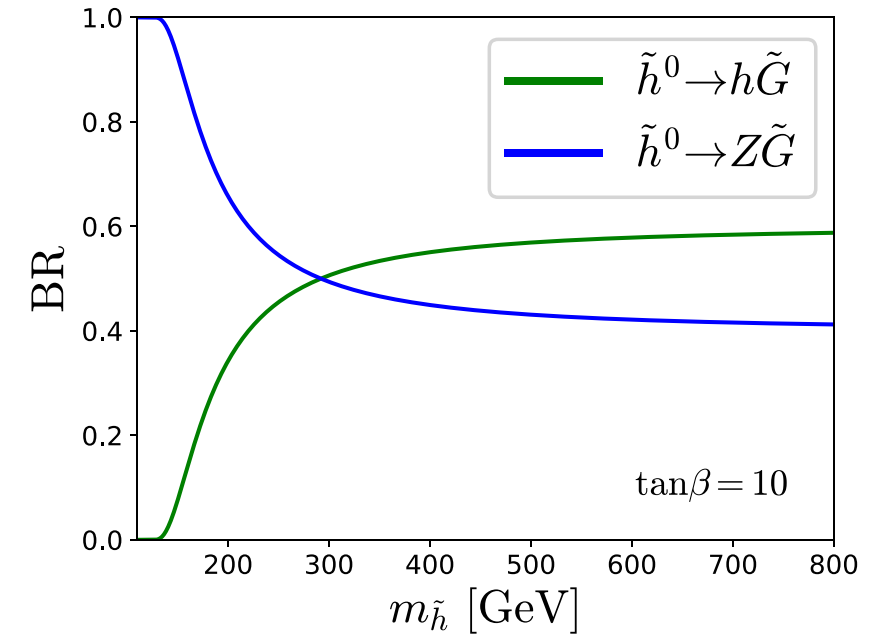
Bino-like



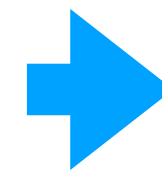
Wino-like



Higgsino-like



- Higgsino, Wino direct production excluded up to $\sim 700 \text{ GeV}$
- SUSY g-2 requires Higgsino or Wino with $m < 600 \text{ GeV}$



**SUSY (g-2)_μ
incompatible
with LHC**

Analysis Framework

SUSY g-2: 1-loop + leading 2-loop GM2Calc [Eur.Phys.J. C76 (2016) no.2, 62]

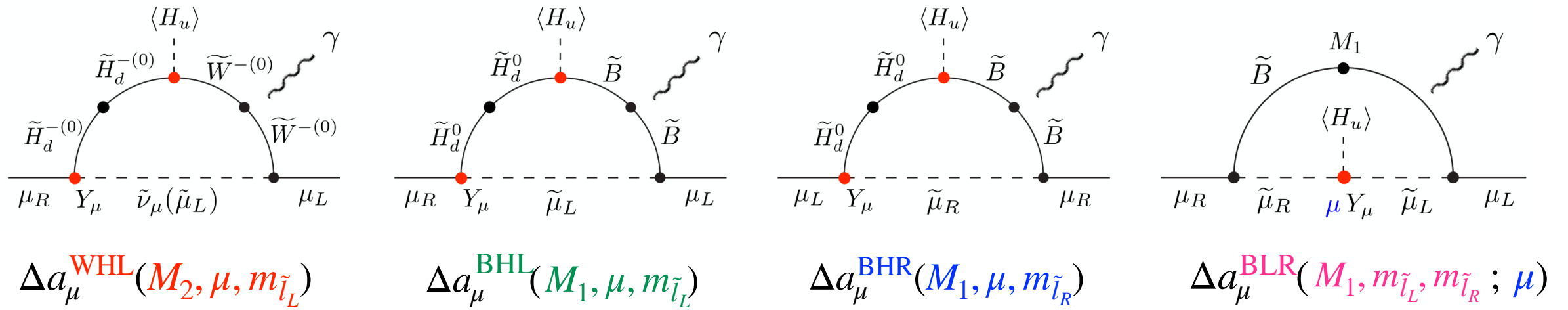
Neutralino abundance, Direct Detection: MicrOMEGAs [2003.08621]

Decay of SUSY particles: SUSY-HIT [hep-ph/0609292]

LHC constraints:

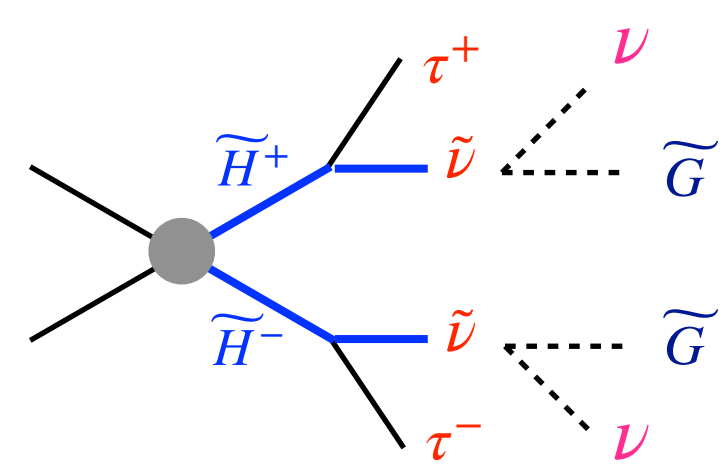
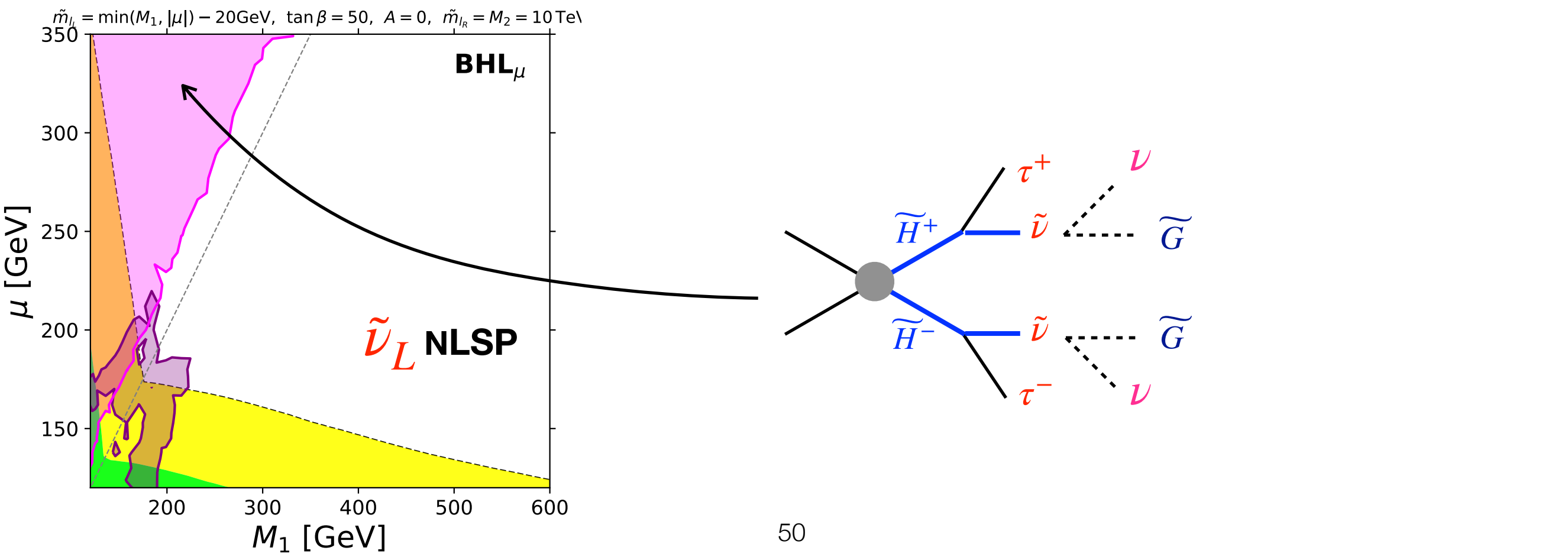
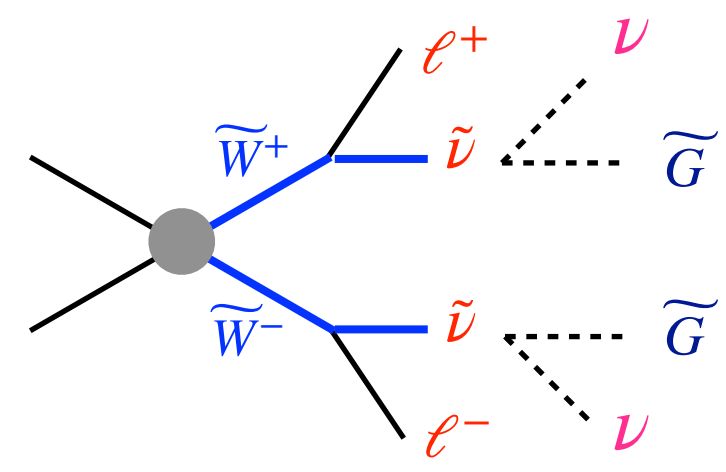
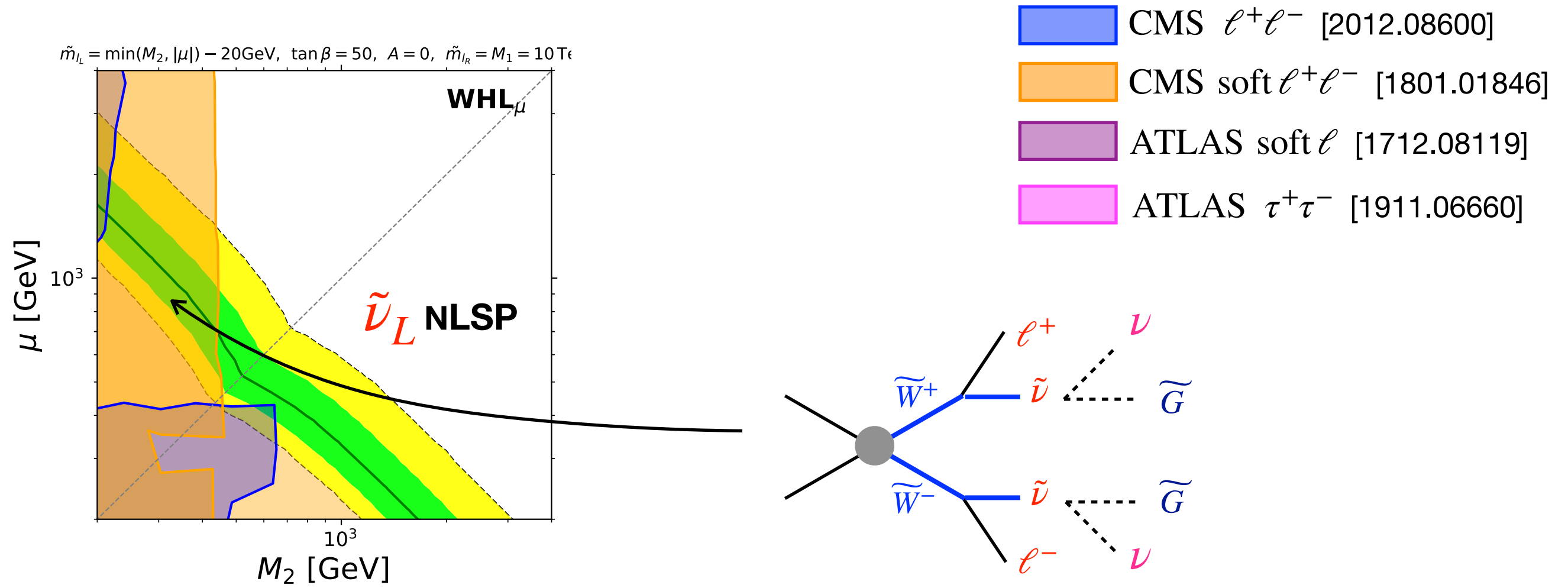
- **MSSM:** ① Mapping simplified model limits to the model point (σ BR)
- **RPV:** ② Pythia 8 + CheckMATE 2 [1907.09874], [1611.09856]
- **Gravitino LSP:** Both ① and ②

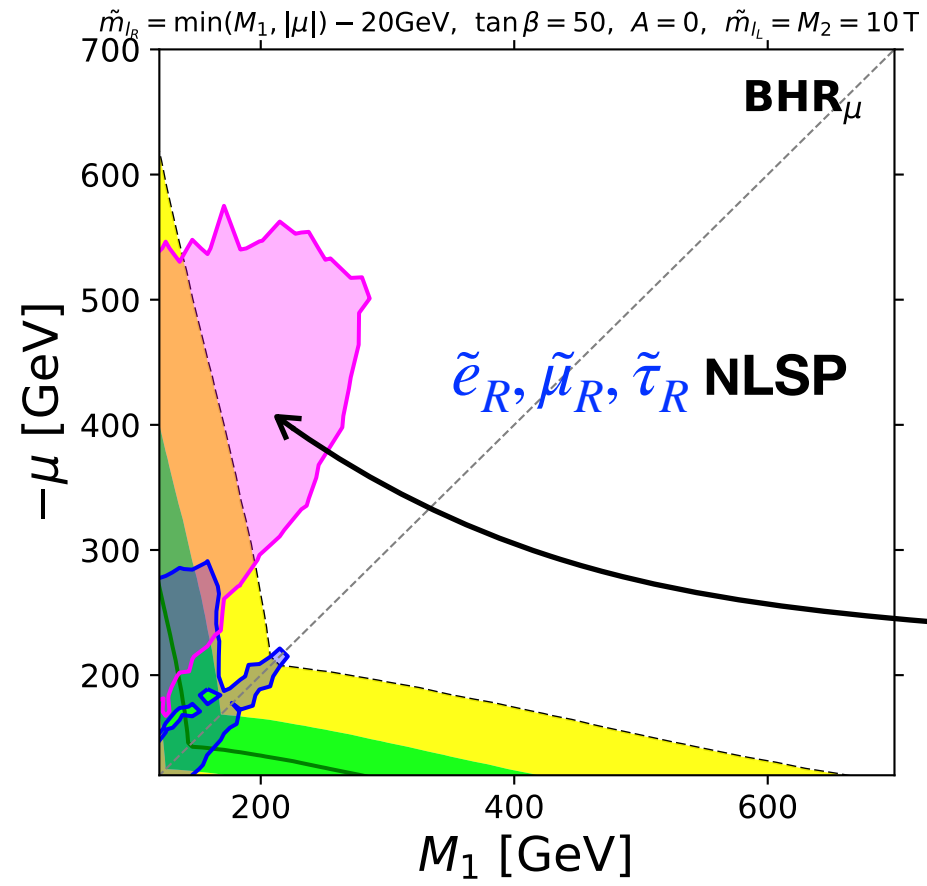
Parameter planes



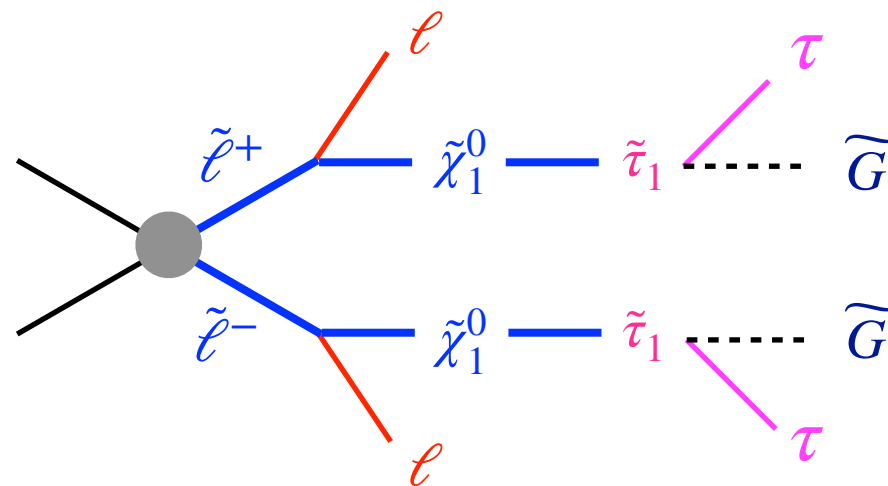
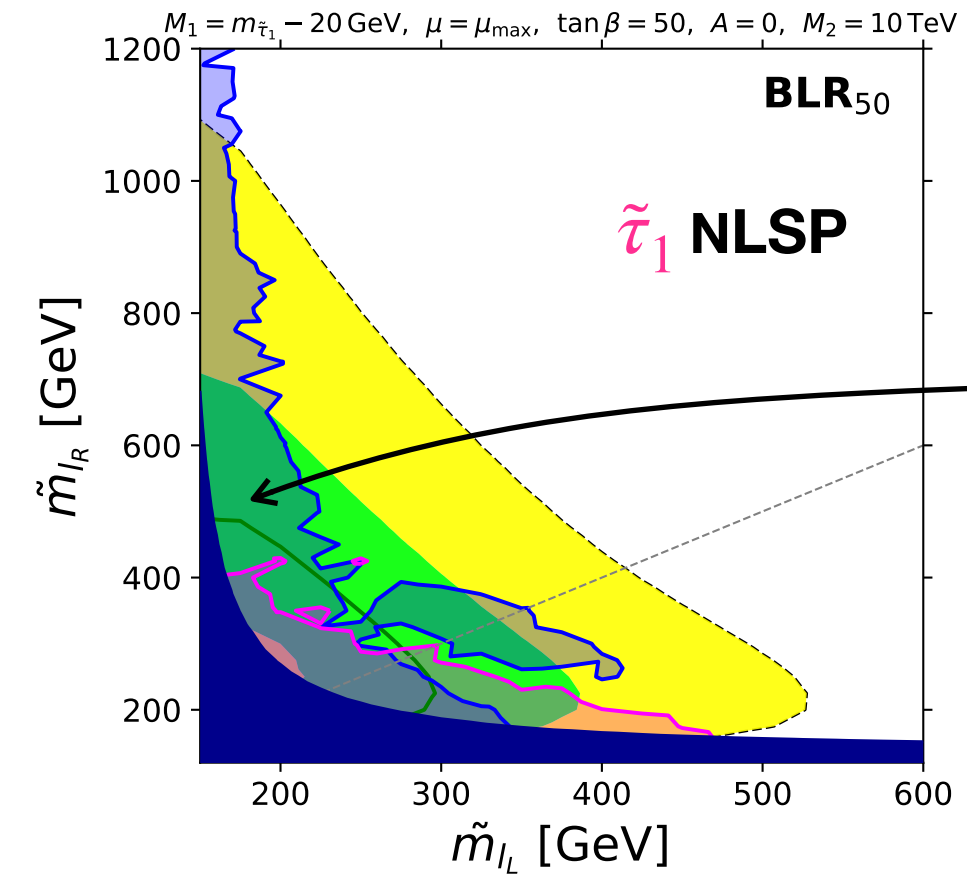
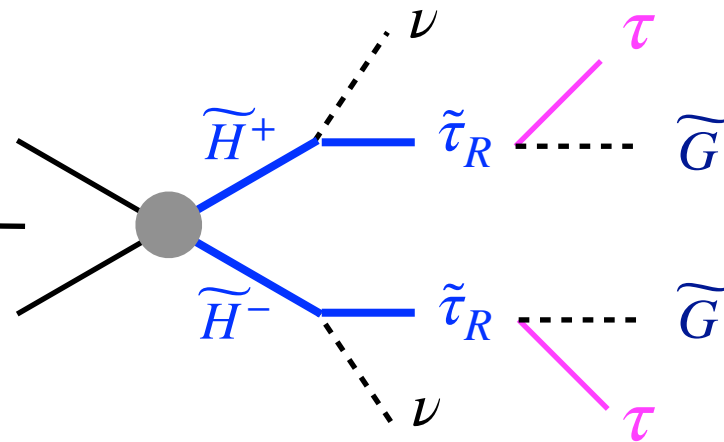
2D planes

name	axes	other parameters	$\tan \beta$
WHL	(M_2, μ)	$\tilde{m}_{l_L} = \min(M_2, \mu) + 20 \text{ GeV}$, $M_1 = \tilde{m}_{l_R} = 10 \text{ TeV}$	50
BHL	(M_1, μ)	$\tilde{m}_{l_L} = \min(M_1, \mu) + 20 \text{ GeV}$, $M_2 = \tilde{m}_{l_R} = 10 \text{ TeV}$	50
BHR	(M_1, μ)	$\tilde{m}_{l_R} = \min(M_1, \mu) + 20 \text{ GeV}$, $M_2 = \tilde{m}_{l_L} = 10 \text{ TeV}$	50
BLR	$(\tilde{m}_{l_L}, \tilde{m}_{l_R})$	$M_1 = m_{\tilde{\tau}_1} - 20 \text{ GeV}$, $\mu = \mu_{\text{max}}$, $M_2 = 10 \text{ TeV}$	50

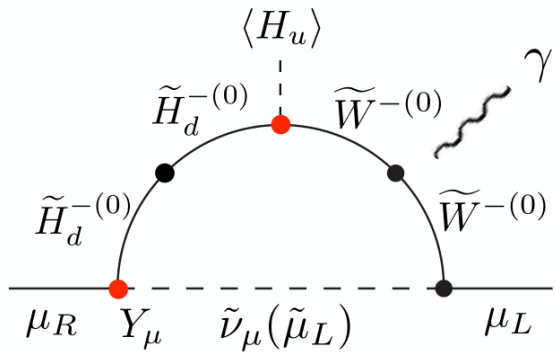




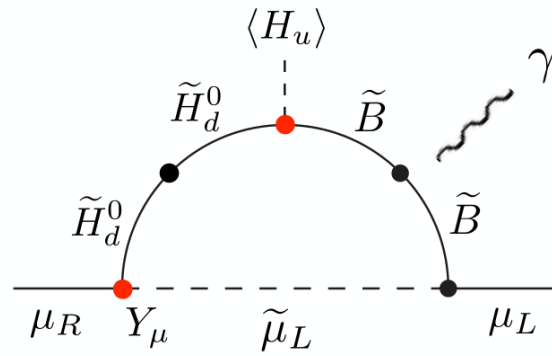
- CMS $\ell^+\ell^-$ [2012.08600]
- CMS soft $\ell^+\ell^-$ [1801.01846]
- ATLAS soft ℓ [1712.08119]
- ATLAS $\tau^+\tau^-$ [1911.06660]



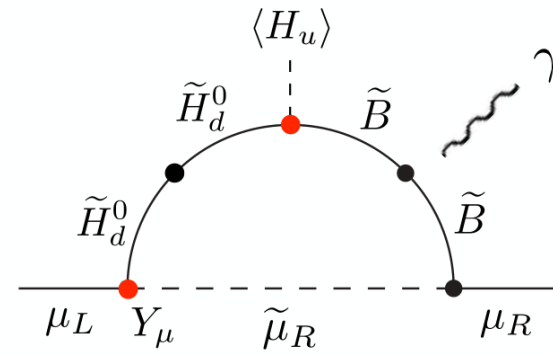
Non $\tilde{\chi}_1^0$ NLSP (Short Summary)



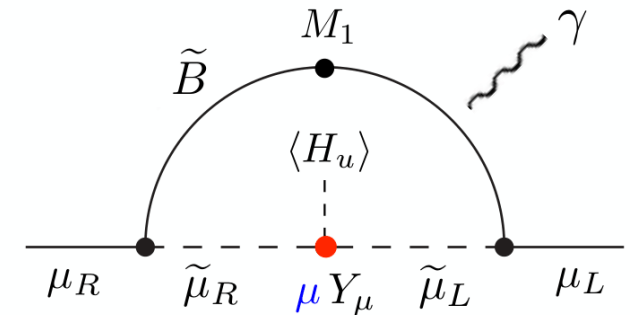
WHL



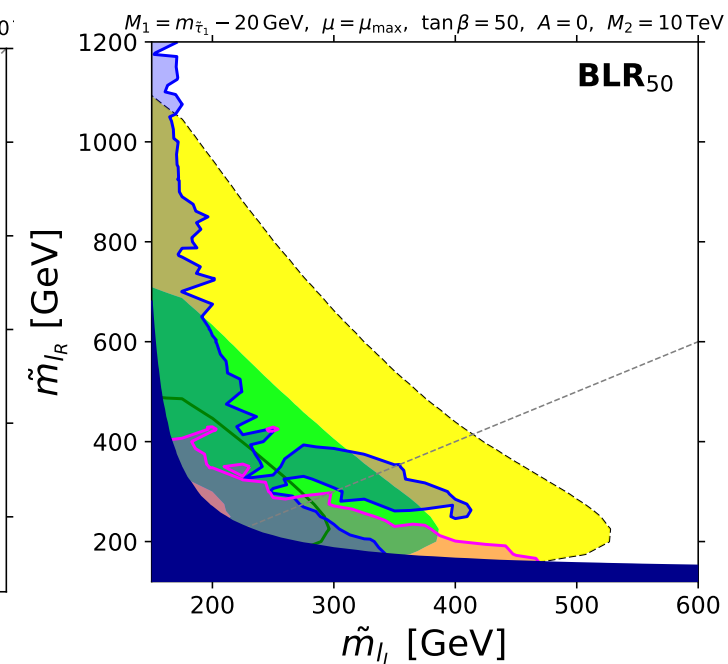
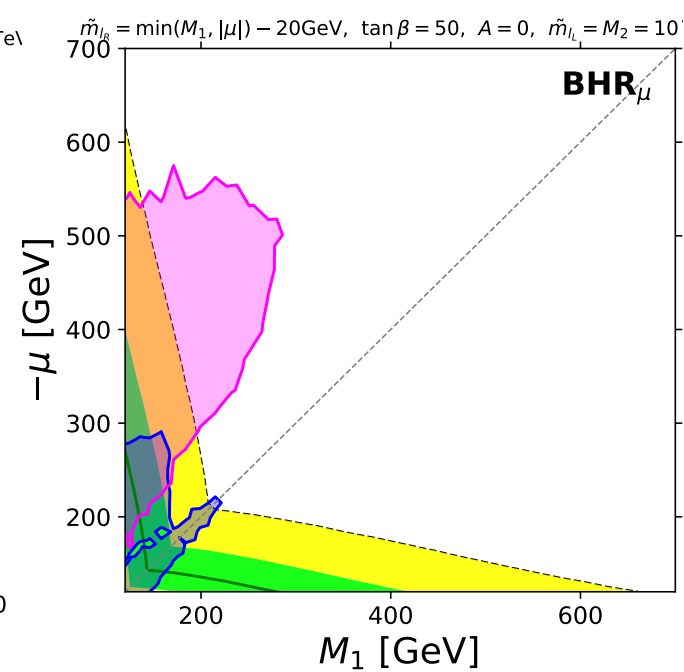
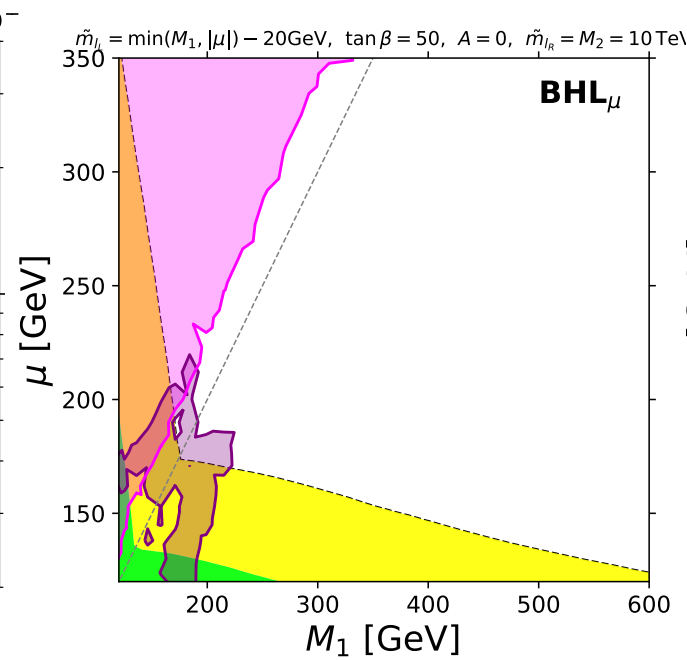
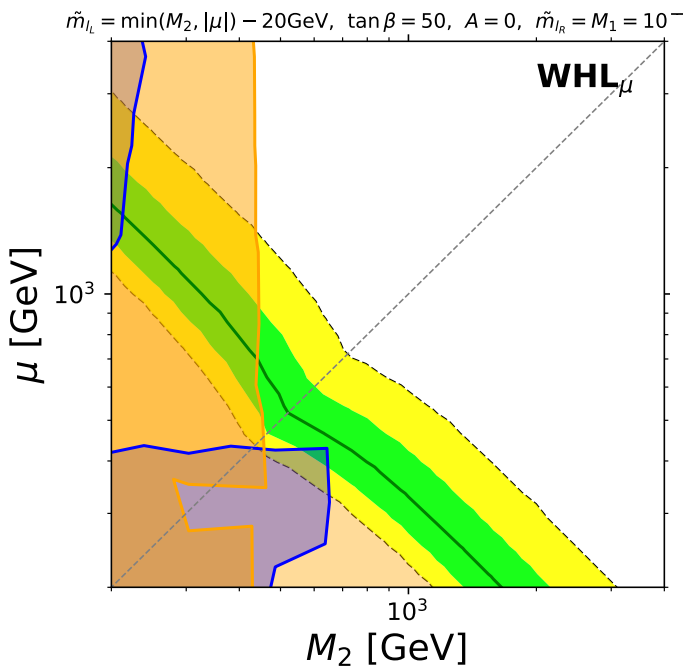
BHL



BHR



BLR



● small $|\mu|$ region is compatible with $(g-2)_{\mu}$

