



Probing the Higgs sector with gravitational waves

GLÁUBER CARVALHO DORSCH

ICEX 
*Departamento
de Física*

UF  MG

LHCP 2022
Taipei (online), 16th May 2022

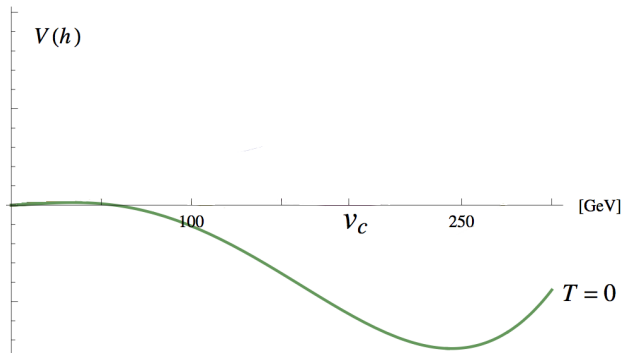
Probing the Higgs sector with gravitational waves

- Electroweak phase transition
The bridge between the Higgs and gravitational wave physics
- LISA: the Laser Interferometer Space Antenna
...and the possibility of detecting these GWs
- Probing BSM physics with GWs: a 2HDM example
- Conclusions and outlook

The electroweak phase transition

**Cold, nearly empty
Universe ($T \approx 0$)**

Higgs mechanism proceeds
as in textbooks



The electroweak phase transition

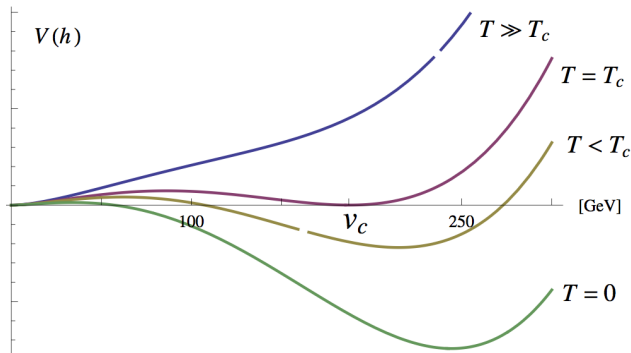
Cold, nearly empty
Universe ($T \approx 0$)

Higgs mechanism proceeds
as in textbooks

But the early Universe is hot!

Higgs immersed in plasma
Thermal corrections to *effective potential*

$$V_{\text{eff}} \stackrel{\text{high } T}{\approx} V_0 + \sum_i g_i \frac{T^2}{24} m_i^2 - \underbrace{\frac{T}{12\pi} m_i^3}_{\text{bosons only}} + \dots$$



The electroweak phase transition

Cold, nearly empty
Universe ($T \approx 0$)

Higgs mechanism proceeds
as in textbooks

But the early Universe is hot!

Higgs immersed in plasma
Thermal corrections to *effective potential*

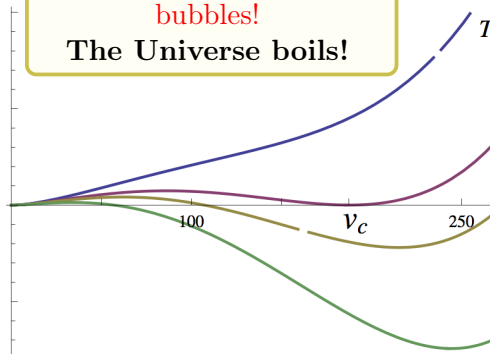
Barrier between vacua



bubbles!

The Universe boils!

$$V_{\text{eff}}^{\text{high } T} \approx V_0 + \sum_i g_i \frac{T^2}{24} m_i^2 - \underbrace{\frac{T}{12\pi} m_i^3}_{\text{bosons only}} + \dots$$



$T \gg T_c$
 $T = T_c$

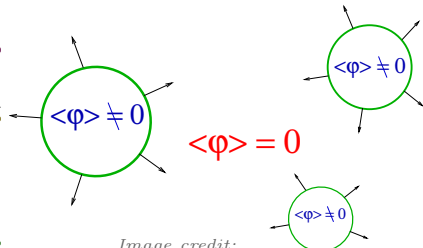
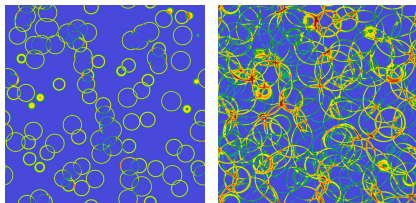


Image credit:
Morrisey and Ramsey-Musolf,
New J. Phys. **14** (2012) 125003

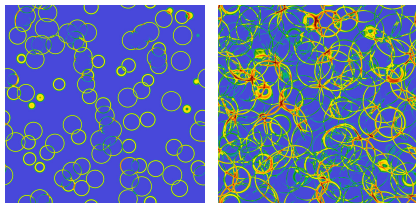
Gravitational waves



*Hindmarsh, Huber, Rummukainen and Weir
Phys. Rev. D* **92**, 123009

Collisions break spherical symmetry
↓
quadrupole moment (sources GWs)!
(kinetic energy, sound waves, turbulence)

Gravitational waves



*Hindmarsh, Huber, Rummukainen and Weir
Phys. Rev. D 92, 123009*

Collisions break spherical symmetry
 \Downarrow
 quadrupole moment (sources GWs)!
 (kinetic energy, sound waves, turbulence)

Long distances

sources causally disconnected
 white noise $\Rightarrow \Omega_{\text{GW}} \sim k^3$

Total GW energy is finite:
 spectrum decreases for $k \gtrsim k_*$

$$f_{\text{peak}} = 16.5 \times 10^{-3} \text{ mHz} \left(\frac{g_{s*}}{100}\right)^{1/6} \frac{T_*}{100 \text{ GeV}} \left(\frac{f_*}{H_*}\right)$$

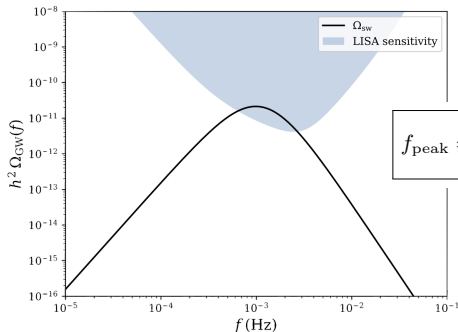
Typically

$$f_*/H_* \simeq 100 - 1000 \Rightarrow f_{\text{peak}} \simeq \mathcal{O}(\text{mHz})$$

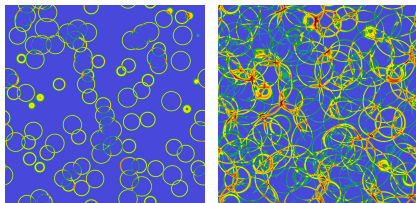
$$T_* \simeq 100 \text{ GeV}$$

Density parameter

$$\frac{d\Omega_{\text{GW}}}{d \ln k} = \frac{1}{12H^2} \frac{k^3}{2\pi^2} P_h(\mathbf{k}, t)$$



Gravitational waves



*Hindmarsh, Huber, Rummukainen and Weir
Phys. Rev. D 92, 123009*

Collisions break spherical symmetry
 \Downarrow
 quadrupole moment (sources GWs)!
 (kinetic energy, sound waves, turbulence)

Long distances

sources causally disconnected
 white noise $\Rightarrow \Omega_{\text{GW}} \sim k^3$

Total GW energy is finite:
 spectrum decreases for $k \gtrsim k_*$

$$f_{\text{peak}} = 16.5 \times 10^{-3} \text{ mHz} \left(\frac{g_{s*}}{100}\right)^{1/6} \frac{T_*}{100 \text{ GeV}} \left(\frac{f_*}{H_*}\right)$$

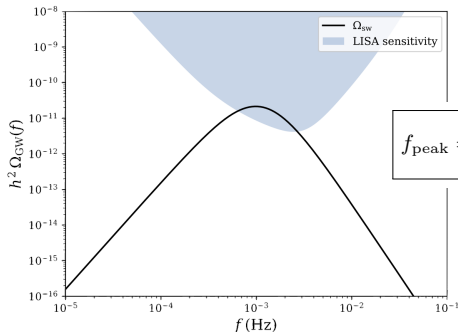
Typically

$$f_*/H_* \simeq 100 - 1000 \Rightarrow f_{\text{peak}} \simeq \mathcal{O}(\text{mHz})$$

$$T_* \simeq 100 \text{ GeV}$$

Density parameter

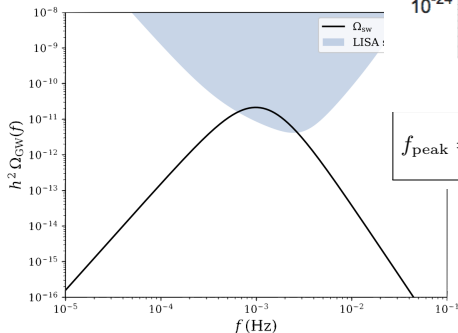
$$\frac{d\Omega_{\text{GW}}}{d \ln k} = \frac{1}{12H^2} \frac{k^3}{2\pi^2} P_h(\mathbf{k}, t)$$



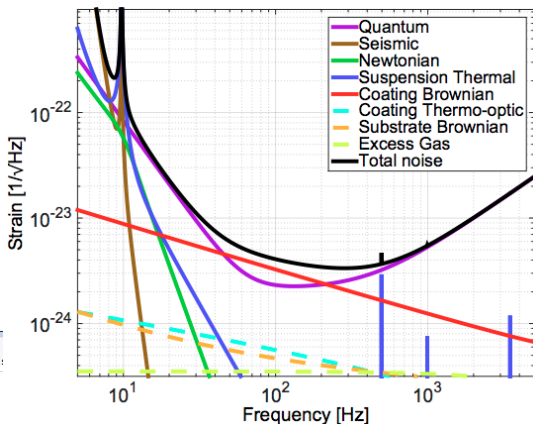
Gravitational waves

Ground-based interferometers

Low f sensitivity limited by seismic, thermal, Earth's gravity gradients



LIGO noise budget (LIGO Document T1800044-v5)



$$f_{\text{peak}} = 16.5 \times 10^{-3} \text{ mHz} \left(\frac{g_{s*}}{100} \right)^{1/6} \frac{T_*}{100 \text{ GeV}} \left(\frac{f_*}{H_*} \right)$$

Typically

$$f_*/H_* \simeq 100 - 1000 \Rightarrow f_{\text{peak}} \simeq \mathcal{O}(\text{mHz})$$

$$T_* \simeq 100 \text{ GeV}$$

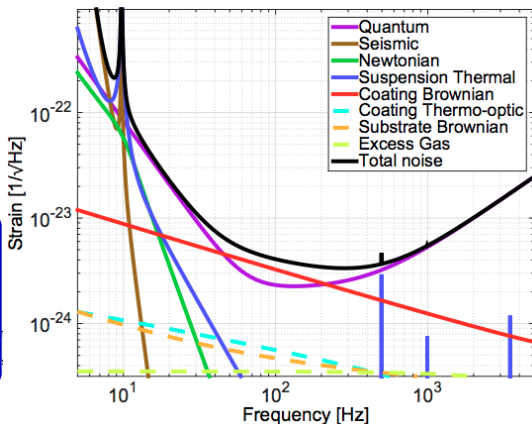
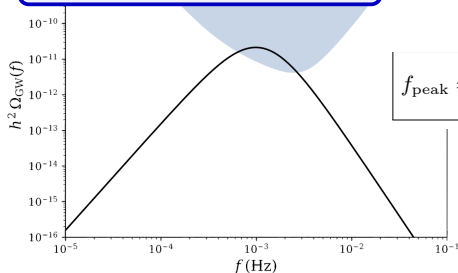
Gravitational waves

LIGO noise budget (LIGO Document T1800044-v5)

Ground-based interferometers

Low f sensitivity limited by seismic, thermal, Earth's gravity gradients

Must go to space!



$$f_{\text{peak}} = 16.5 \times 10^{-3} \text{ mHz} \left(\frac{g_{s*}}{100} \right)^{1/6} \frac{T_*}{100 \text{ GeV}} \left(\frac{f_*}{H_*} \right)$$

Typically

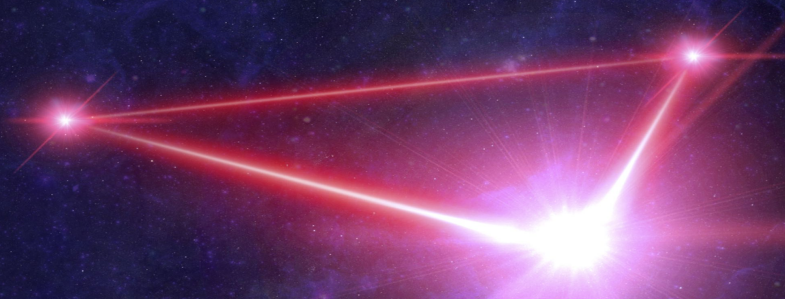
$$f_*/H_* \simeq 100 - 1000 \Rightarrow f_{\text{peak}} \simeq \mathcal{O}(\text{mHz})$$

$$T_* \simeq 100 \text{ GeV}$$

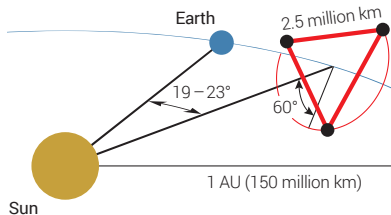


LISA

Laser Interferometer Space Antenna



LISA: Laser Interferometer Space Antenna

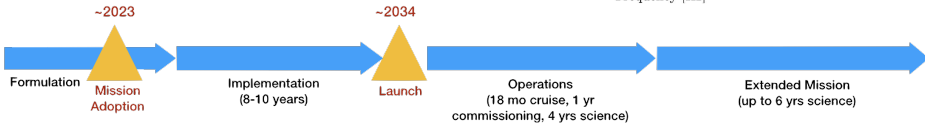
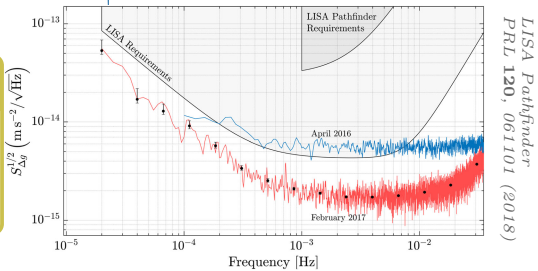


3 free-falling spacecraft preserving equilateral shape while orbiting the sun(!)

Sensitive to variations of 1 part in 10^{22} arm lengths

Gravitons decouple much earlier than photons from plasma!

Detection of GWs from EWPT give us a “picture” of the Universe at 10^{-12} s after “Big Bang” (CMB: 380,000 years!)



Astro2020 Whitepaper, arXiv:1907.06482 [astro-ph-IM]

GWs as complementary probes of BSM physics

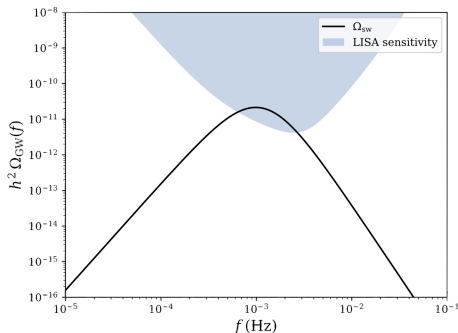
Accurate prediction of spectrum from BSM



test BSM physics with GWs!

Key parameters:

- $\alpha \equiv \frac{\text{energy released}}{\text{total energy in radiation}}$
- $\beta^{-1} \equiv \begin{array}{l} \text{characteristic timescale} \\ \text{of phase transition;} \\ \Gamma = \Gamma_* e^{-\beta(t-t_*)} \end{array}$
- $T_* \equiv \text{temperature}$
- $v_w \equiv \text{velocity of bubble expansion}$



A CONCRETE EXAMPLE: THE 2HDM CASE

Two-Higgs-doublet models

GCD, Huber, Mimasu, No, PRL 113 (2014) 21, 211802

GCD, Huber, Mimasu, No, PRD 93 (2016) 11, 115033

Caprini, Chala, GCD et al., JCAP 03 (2020) 024

- Minimal SM extensions:

- ▶ Two $SU(2)_L$ scalar doublets: Φ_1 and Φ_2 .
- ▶ Motivated by many SM extensions (e.g. SUSY, Composite Higgs).
- ▶ Various heavy scalars (h_0, H_0, A_0, H^\pm) increase EWPT strength.
- ▶ 7 parameters: 4 scalar masses, overall mass scale of Φ_2 , 2 mixing angles

Two-Higgs-doublet models

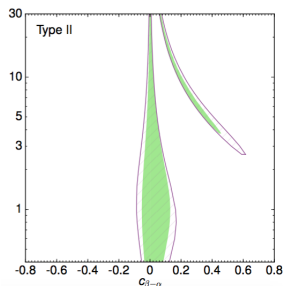
GCD, Huber, Mimasu, No, PRL 113 (2014) 21, 211802

GCD, Huber, Mimasu, No, PRD 93 (2016) 11, 115033

Caprini, Chala, GCD et al., JCAP 03 (2020) 024

● Minimal SM extensions:

- ▶ Two $SU(2)_L$ scalar doublets: Φ_1 and Φ_2 .
- ▶ Motivated by many SM extensions (e.g. SUSY, Composite Higgs).
- ▶ Various heavy scalars (h_0, H_0, A_0, H^\pm) increase EWPT strength.
- ▶ 7 parameters: 4 scalar masses, overall mass scale of Φ_2 , 2 mixing angles



Fix $m_{h_0} = 125$ GeV

Higgs measurements: $h_0 \approx h_{\text{SM}}$

EWPO \implies approx. degeneracy

$$m_{H^\pm} \approx m_{A^0} \text{ or } m_{H^0}$$

Two-Higgs-doublet models

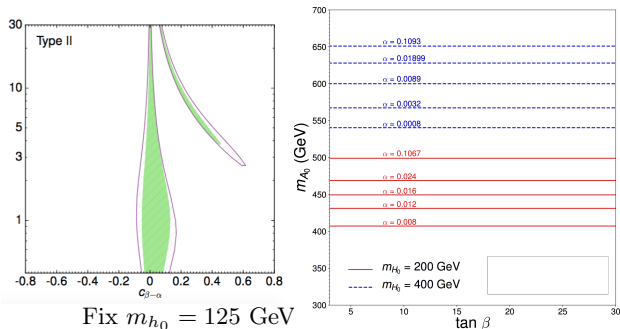
GCD, Huber, Mimasu, No, PRL **113** (2014) 21, 211802

GCD, Huber, Mimasu, No, PRD **93** (2016) 11, 115033

Caprini, Chala, GCD et al., JCAP **03** (2020) 024

Minimal SM extensions:

- ▶ Two $SU(2)_L$ scalar doublets: Φ_1 and Φ_2 .
- ▶ Motivated by many SM extensions (e.g. SUSY, Composite Higgs).
- ▶ Various heavy scalars (h_0, H_0, A_0, H^\pm) increase EWPT strength.
- ▶ 7 parameters: 4 scalar masses, overall mass scale of Φ_2 , 2 mixing angles



Fix $m_{h_0} = 125$ GeV

Higgs measurements: $h_0 \approx h_{\text{SM}}$

EWPO \implies approx. degeneracy

$$m_{H^\pm} \approx m_{A^0} \text{ or } m_{H^0}$$

Two-Higgs-doublet models

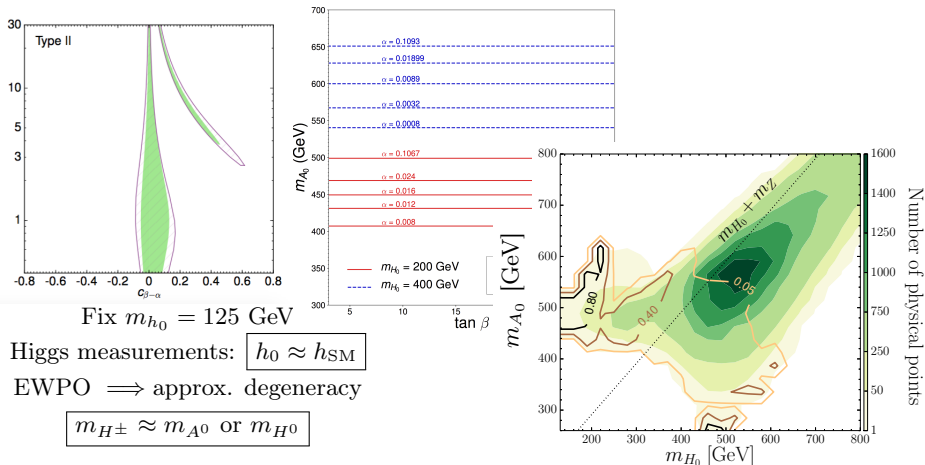
GCD, Huber, Mimasu, No, PRL **113** (2014) 21, 211802

GCD, Huber, Mimasu, No, PRD **93** (2016) 11, 115033

Caprini, Chala, GCD et al., JCAP **03** (2020) 024

Minimal SM extensions:

- ▶ Two $SU(2)_L$ scalar doublets: Φ_1 and Φ_2 .
- ▶ Motivated by many SM extensions (e.g. SUSY, Composite Higgs).
- ▶ Various heavy scalars (h_0, H_0, A_0, H^\pm) increase EWPT strength.
- ▶ 7 parameters: 4 scalar masses, overall mass scale of Φ_2 , 2 mixing angles



Two-Higgs-doublet models

GCD, Huber, Mimasu, No, PRL **113** (2014) 21, 211802

GCD, Huber, Mimasu, No, PRD **93** (2016) 11, 115033

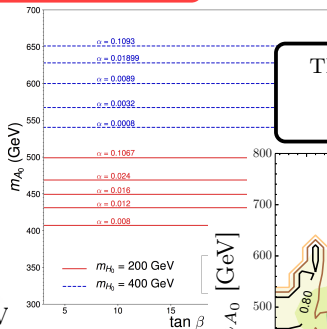
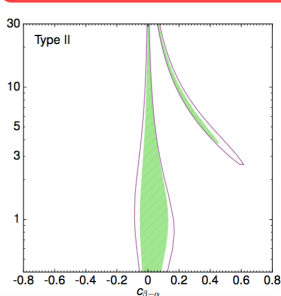
Caprini, Chala, GCD et al., JCAP **03** (2020) 024

Dialogical aspect

Cosmological phase transition



implications for collider searches



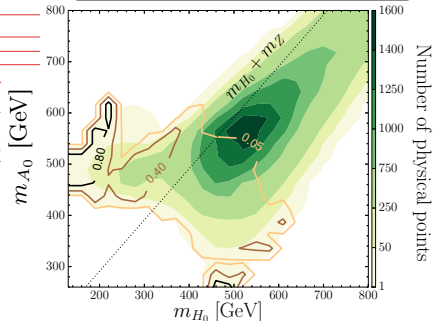
This motivated a search for $A/H \rightarrow ZH/A$ by CMS and ATLAS

Fix $m_{h_0} = 125$ GeV

Higgs measurements: $h_0 \approx h_{SM}$

EWPO \implies approx. degeneracy

$$m_{H^\pm} \approx m_{A^0} \text{ or } m_{H^0}$$



Two-Higgs-doublet models

GCD, Huber, Mimasu, No, *PRL* **113** (2014) 21, 211802

GCD, Huber, Mimasu, No, *PRD* **93** (2016) 11, 115033

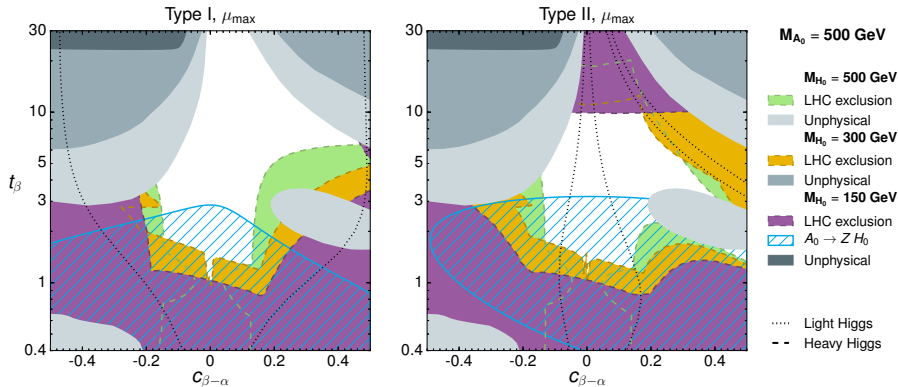
Caprini, Chala, GCD et al., *JCAP* **03** (2020) 024

Dialogical aspect

Cosmological phase transition



implications for collider searches

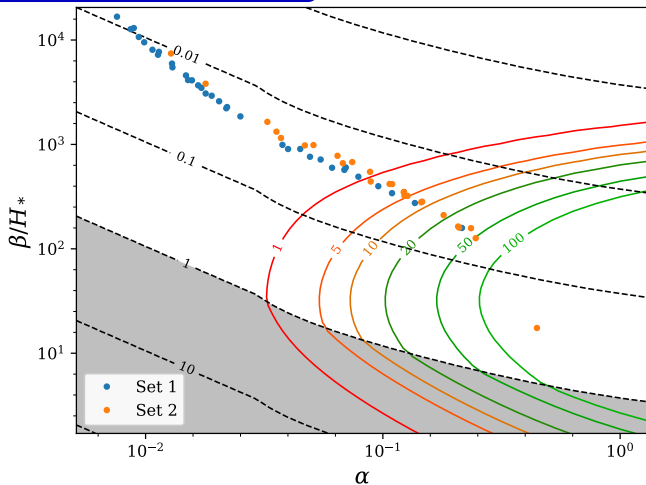


See also Kling, Su and Su, *JHEP* **06** (2020) 163

LISA as a probe for BSM physics

Complementary aspect

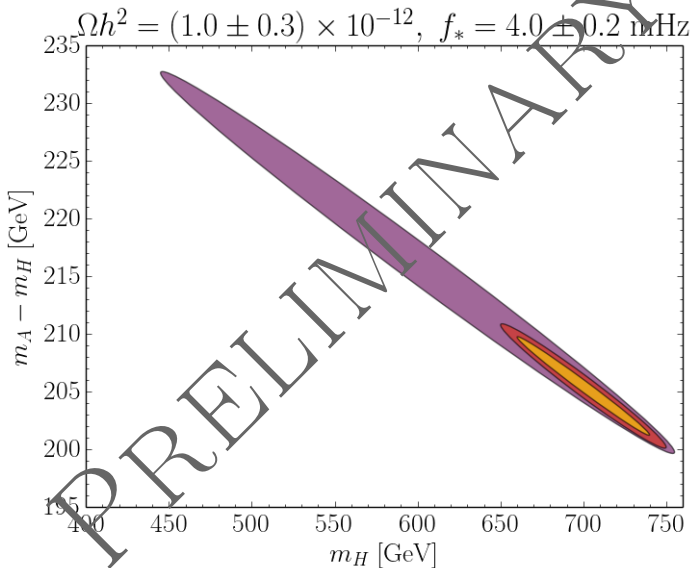
Caprini, Chala, GCD et al., JCAP **03** (2020) 024



Orange: 2HDM Type-II excluded at LHC (but not Type-I)

Blue: Not excluded by LHC

LISA as a probe for BSM physics



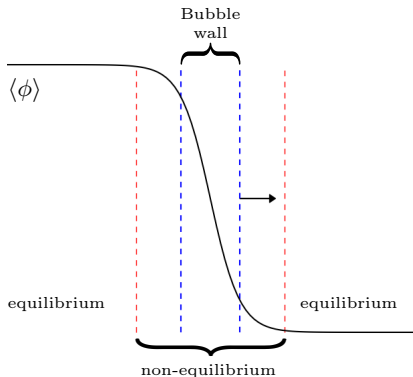
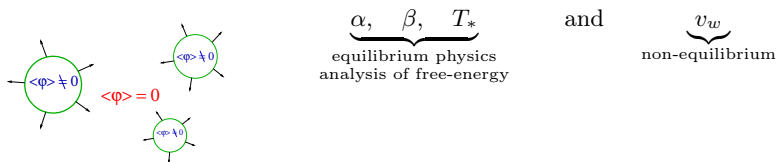
BUBBLE WALL VELOCITY:

A COMMENT ON RECENT DEVELOPMENTS

Bubble wall velocity and Boltzmann equation

GCD, Huber, Konstandin, JCAP 04 (2022) 04, 010

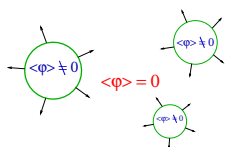
GW spectrum determined by 4 thermodynamical parameters:



Bubble wall velocity and Boltzmann equation

GCD, Huber, Konstandin, JCAP 04 (2022) 04, 010

GW spectrum determined by 4 thermodynamical parameters:



α, β, T_*
equilibrium physics
analysis of free-energy

v_w
non-equilibrium

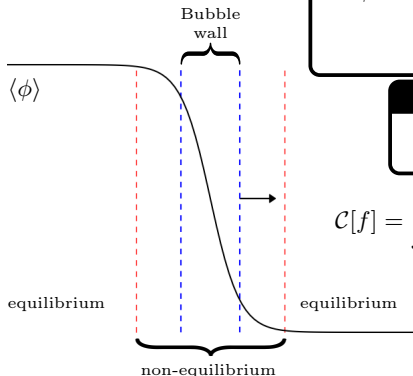
v_w computed from Klein-Gordon equation

$$-\phi'' + \frac{dV_T}{d\phi} + \underbrace{\sum_i \frac{dm_i^2}{d\phi} \int \frac{d^3p}{(2\pi)^3 2E_i} \delta f_i(x^\mu, p^\mu)}_{\text{friction}} = 0$$

δf from Boltzmann equation

$$p^\mu \partial_\mu \delta f_i(x^\mu, p^\mu) = \mathcal{S}[f_j] + \mathcal{C}[f_j]$$

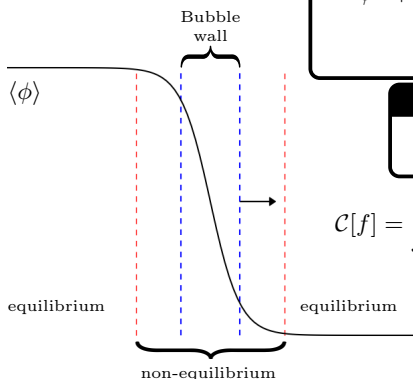
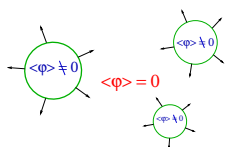
$$\mathcal{C}[f] = \int_k \int_{p'} \int_{k'} |\mathcal{M}|^2 (2\pi)^4 \delta^4(p + k - p' - k') \mathcal{P}[f]$$



Bubble wall velocity and Boltzmann equation

GCD, Huber, Konstandin, JCAP 04 (2022) 04, 010

GW spectrum determined by 4 thermodynamical parameters:



α, β, T_*
equilibrium physics
analysis of free-energy

v_w
non-equilibrium

v_w computed from Klein-Gordon equation

$$-\phi'' + \frac{dV_T}{d\phi} + \underbrace{\sum_i \frac{dm_i^2}{d\phi} \int \frac{d^3p}{(2\pi)^3 2E_i} \delta f_i(x^\mu, p^\mu)}_{\text{friction}} = 0$$

δf from Boltzmann equation

$$p^\mu \partial_\mu \delta f_i(x^\mu, p^\mu) = \mathcal{S}[f_j] + \mathcal{C}[f_j]$$

$$\mathcal{C}[f] = \int_k \int_{p'} \int_{k'} |\mathcal{M}|^2 (2\pi)^4 \delta^4(p + k - p' - k') \mathcal{P}[f]$$

How can we solve this
integro-differential
equation?

Current debate: continuous *vs.* discontinuous friction

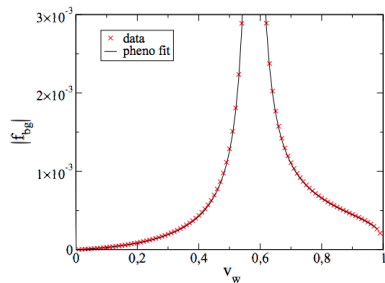
Konstandin, Nardini, Rues, *JCAP* **09** (2014) 028

Fluid Ansatz (90's)

$$f(x, p) = \frac{1}{e^{\beta p^\mu (u_\mu + \delta u_\mu) - \mu} \pm 1}$$

Moore, Prokopec, *PRD* **52** (1995) 7182

3 fluctuations: μ , $\delta u_0 \equiv \delta T/T$, $\delta u_z \equiv \delta v$
Discontinuous friction



Current debate: continuous *vs.* discontinuous friction

Fluid Ansatz (90's)

$$f(x, p) = \frac{1}{e^{\beta p^\mu (u_\mu + \delta u_\mu) - \mu} \pm 1}$$

Moore, Prokopec, PRD 52 (1995) 7182

3 fluctuations: μ , $\delta u_0 \equiv \delta T/T$, $\delta u_z \equiv \delta v$
Discontinuous friction

Recently proposed

$$f(x, p) = \frac{1}{e^{\beta p^\mu u_\mu - \mu} \pm 1} + \delta f$$

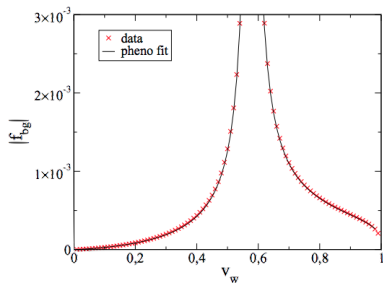
Cline, Kainulainen
PRD 101 (2020) 063525

ad hoc factorization property:

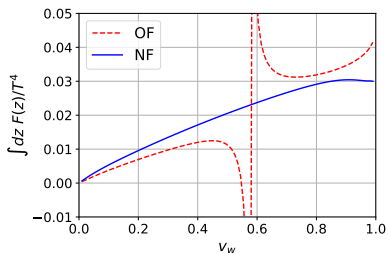
$$\langle X \delta f \rangle \sim \left\langle \frac{p_z}{E} \delta f \right\rangle \times \int d^3 p \frac{E}{p_z} X f_{\text{eq}}$$

Leaves 2 fluctuations: μ and $u \equiv \langle (p_z/E) \delta f \rangle$

Konstandin, Nardini, Rues, JCAP 09 (2014) 028



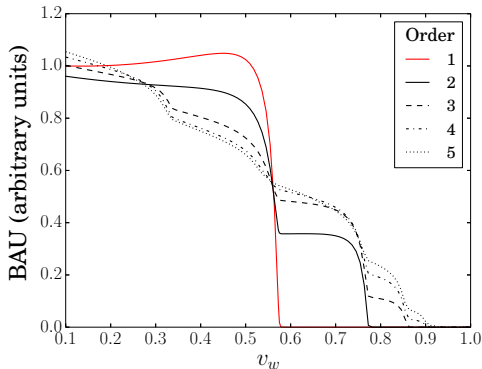
Laurent, Cline, PRD 102 (2020) 063516



Could the discontinuity be an artifact
of truncating the fluid Ansatz
(i.e. including only 3 perturbations)?

This happens for the
baryon asymmetry

GCD, Huber, Konstandin
JCAP 08 (2021) 020



Extended fluid Ansatz

$$f(x, p) = \frac{1}{e^{\beta p^\mu u_\mu + \delta} \pm 1}$$

with

$$\delta = w^{(0)} + p^\mu w_\mu^{(1)} + p^\mu p^\nu w_{\mu\nu}^{(2)} + \dots$$

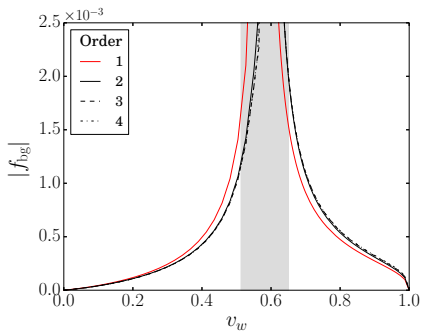
Linearized Boltzmann equation

$$p^\mu \partial_\mu \delta f_i(x^\mu, p^\mu) = \mathcal{S}[f_j] + \mathcal{C}[f_j]$$

\Downarrow

$$A \cdot q' + \Gamma \cdot q = S$$

$$\text{with } q = (w^{(0)}, w_0^{(1)}, w_z^{(1)}, \dots)^T$$



Extended fluid Ansatz

$$f(x, p) = \frac{1}{e^{\beta p^\mu u_\mu + \delta} \pm 1}$$

with

$$\delta = w^{(0)} + p^\mu w_\mu^{(1)} + p^\mu p^\nu w_{\mu\nu}^{(2)} + \dots$$

Linearized Boltzmann equation

$$p^\mu \partial_\mu \delta f_i(x^\mu, p^\mu) = \mathcal{S}[f_j] + \mathcal{C}[f_j]$$

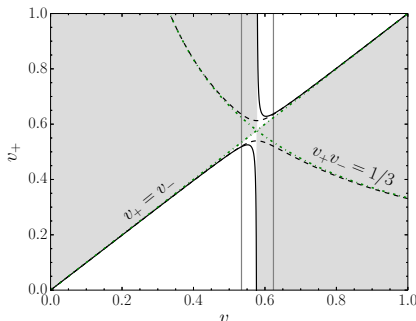
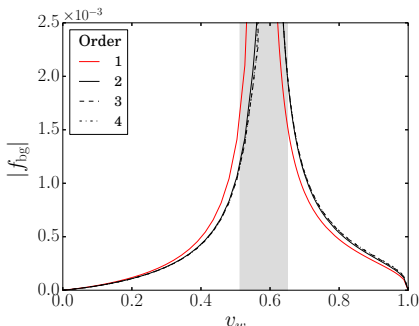
\Downarrow

$$A \cdot q' + \Gamma \cdot q = S$$

$$\text{with } q = (w^{(0)}, w_0^{(1)}, w_z^{(1)}, \dots)^T$$

Important!

A discontinuity is expected from energy-momentum conservation!



$$f(x, p) = \frac{1}{e^{\beta p^\mu u_\mu + \delta} \pm 1}$$

with

$$\delta = w^{(0)} + p^\mu w_\mu^{(1)} + p^\mu p^\nu w_{\mu\nu}^{(2)} + \dots$$

Linearized Boltzmann equation

$$p^\mu \partial_\mu \delta f_i(x^\mu, p^\mu) = \mathcal{S}[f_j] + \mathcal{C}[f_j]$$

\Downarrow

$$A \cdot q' + \Gamma \cdot q = S$$

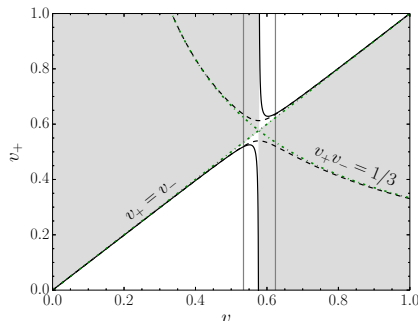
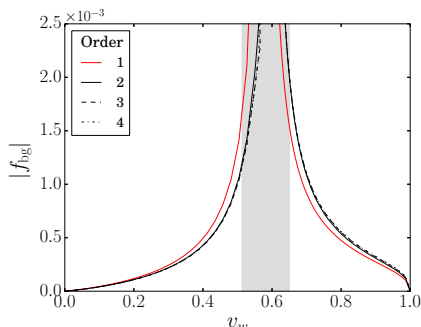
$$\text{with } q = (w^{(0)}, w_0^{(1)}, w_z^{(1)}, \dots)^T$$

Important!

A discontinuity is expected from energy-momentum conservation!

Discontinuity recovered in non-linearized approaches

Laurent, Cline, arXiv:2204.13120



Conclusions

- We live in a golden age for Cosmology
(precision measurements, GW detection, new experiments in sight...)
 - Cosmological observations can be used to constrain BSM Particle Physics.
The Early Universe reached higher energies than any accelerator we could dream of building in the foreseeable future.
 - The study of the EWPT provides a rich Particle-Cosmology interface.
EW epoch \longleftrightarrow physics @ EW scale (relevant at current colliders!)
-
- Room for improvement in estimate of thermodynamical parameters
 - Application of extended fluid Ansatz to specific models
(computation of v_w from first principles)
 - Improvement in implementation (collision terms) and solution of Boltzmann equation

Conclusions

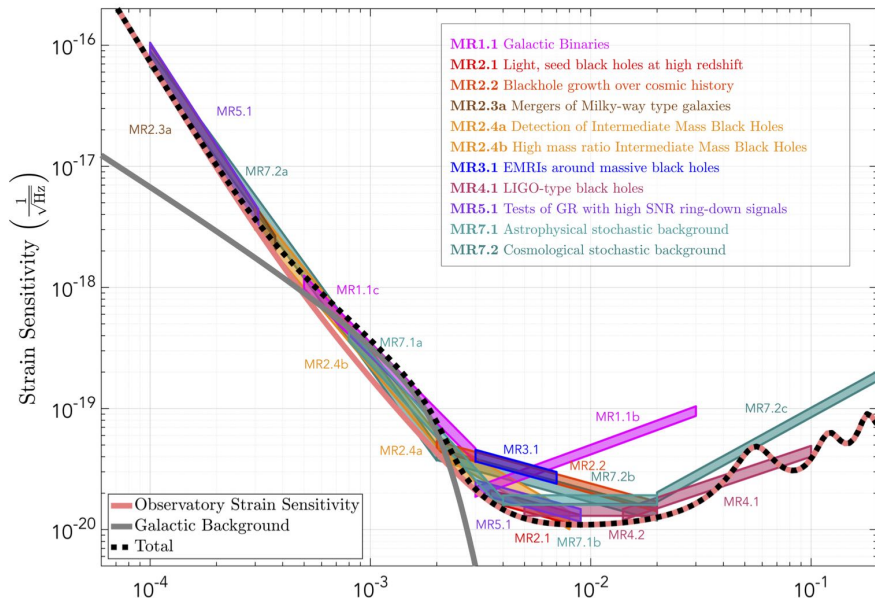
- We live in a golden age for Cosmology (precision measurements, GW detection, new experiments in sight...)
 - Cosmological observations can be used to constrain BSM Particle Physics. The Early Universe reached higher energies than any accelerator we could dream of building in the foreseeable future.
 - The study of the EWPT provides a rich Particle-Cosmology interface. EW epoch \longleftrightarrow physics @ EW scale (relevant at current colliders!)
-
- Room for improvement in estimate of thermodynamical parameters
 - Application of extended fluid Ansatz to specific models (computation of v_w from first principles)
 - Improvement in implementation (collision terms) and solution of Boltzmann equation

THANK YOU!

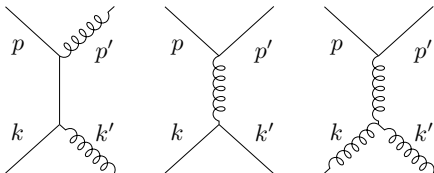
APPENDICES

LISA Science Objectives

LISA Mission Proposal L3



Collision terms



Singular behaviour ← kinetic term
but
collisions important for convergence

$$\text{coll.} \sim \delta_p + \delta_k - \delta_{p'} - \delta_{k'}$$

Annihilations

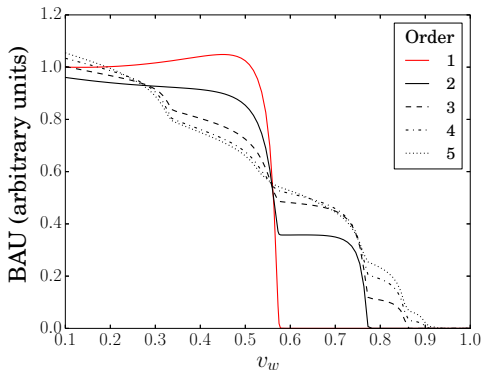
$$|\mathcal{M}|^2 \sim -g_s^4 \frac{st}{(t - m_q^2)^2}$$

$$t = -2p \cdot p' = -2|\mathbf{p}||\mathbf{p}'| \cos \theta_{pp'}$$

$$\int_p p^\mu \dots p^\nu C[f] \simeq \int_p \int_k \int_{p'} \int_{k'} \frac{st}{(t - m_q^2)^2} \delta^4(\dots) p^\mu \dots p^\nu f_p f_k (1 \pm f_{p'}) (1 \pm f_{k'}) \times$$

$$\times \left[\dots + w_{\rho\sigma}^{(2)} (p^\rho p^\sigma + k^\rho k^\sigma) + \dots \right]$$

Results: Supersonic baryogenesis



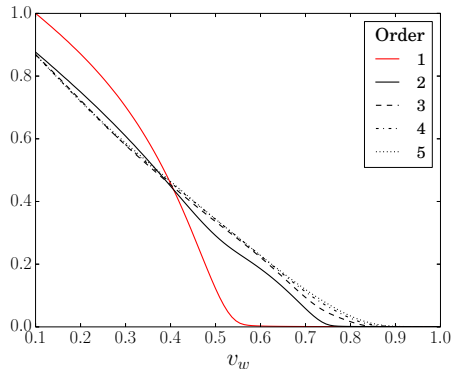
$\alpha_s = 0.01$

Continuous along v_w
similar to found recently in
Cline & Kainulainen
PRD 101 (2020) no. 6, 063525

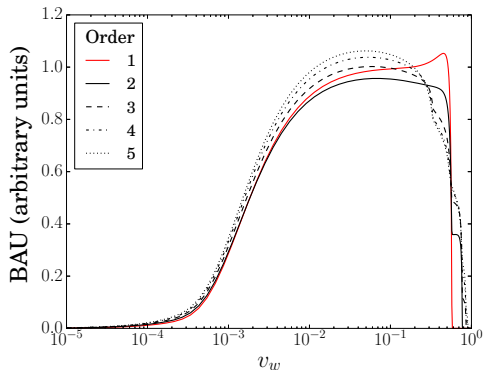
BAU suppressed for $v_w > c_s$, but
not prohibitively small!
(except for $v_w \rightarrow 1$)

convergence
parameter $\sim \frac{T}{\Gamma} \sim DT$

$\alpha_s = 0.06$



Results: small v_w



$$\alpha_s = 0.01$$

Either way the discrepancy is
 $\sim \mathcal{O}(20\%)$ for subsonic walls



3-fluid reasonably reliable
in this regime

BAU can be either
enhanced or suppressed
relative to 1st order

$$\alpha_s = 0.06$$

