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## Abstract

We present a reinterpretation study of existing results from the CMS Collaboration, specifically, searches for light BSM Higgs pairs produced in the chain decay  $pp \rightarrow H_{SM} \rightarrow hh(AA)$  into a variety of final states, in the context of the CP-conserving 2-Higgs Doublet Model (2HDM) Type-I. Through this, we test the LHC sensitivity to a possible new signature,  $pp \rightarrow H_{SM} \rightarrow ZA \rightarrow ZZh$ , with  $ZZ \rightarrow jj\mu^+\mu^-$  and  $h \rightarrow b\bar{b}$ . We perform a systematic scan over the 2HDM Type-I parameter space, by taking into account all available theoretical and experimental constraints, in order to find a region with a potentially visible signal. We show that such a signal is an alternative promising channel to standard four-body searches for light BSM Higgses at the LHC with an integrated luminosity of  $L = 300 fb^{-1}$ .

## General Two Higgs Doublets model (2HDM)

2HDM is one of the simplest and well motivated extension of the standard model (SM). The most general  $SU(2)_L \times U(1)_Y$  invariant scalar potential can be written as follows :

$$V = m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - [m_{12}^2 \Phi_1^\dagger \Phi_2 + h.c.] + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + \left[ \frac{\lambda_5}{2} (\Phi_1^\dagger \Phi_2)^2 + [\lambda_6 (\Phi_1^\dagger \Phi_1) + \lambda_7 (\Phi_2^\dagger \Phi_2)] \Phi_1^\dagger \Phi_2 + h.c. \right]. \quad (1)$$

- Assuming the CP-conservation in the 2HDM and following the hermiticity of the scalar potential,  $m_{11}^2, m_{22}^2, m_{12}^2, \lambda_{1,2,3,4}$  and  $\lambda_{5,6}$  are real. Invoking  $Z_2$  symmetry implies that  $\lambda_6 = \lambda_7 = 0$ .
  - 5 physical Higgses: two CP-even scalars ( $h, H$  with  $m_h < m_H$ ), one CP-odd ( $A^0$ ) and a charged Higgs pair ( $H^\pm$ )
  - The scalar potential has 10 independent parameters:  $m_{11}^2, m_{22}^2, m_{12}^2, v_1, v_2$  and  $\lambda_{1,...,5}$
  - 2 minimization conditions and the combination  $v_1^2 + v_2^2 \Rightarrow 7$  free parameters:
- $m_h < m_H, m_{H^\pm}, m_A, \alpha, \tan \beta = \frac{v_2}{v_1}$  and  $m_{12}^2$ .

## Alignment Limit

- $\cos(\beta - \alpha) \rightarrow 0, h \equiv H_{SM}$ ; standard hierarchy
- with decoupling  $m_{H^\pm} \sim m_A \sim m_H \gg v$
- without decoupling: current data do not at all push the model parameters to the decoupling
- $\sin(\beta - \alpha) \rightarrow 0, H \equiv H_{SM}$ ; inverted hierarchy

## Yukawa Couplings Type-I

One doublet couple to all fermions,

$$\begin{aligned} \kappa_h^{u, d, l} &= c_\alpha / s_\beta = s_{\beta-\alpha} + \cot \beta c_{\beta-\alpha} \\ \kappa_H^{u, d, l} &= s_\alpha / s_\beta = c_{\beta-\alpha} - \cot \beta s_{\beta-\alpha} \\ \kappa_A^{u, d, l} &= -\cot \beta, \kappa_A^u = \cot \beta \end{aligned}$$

with  $\alpha$  and  $\beta$  are the mixing angles.

## Theo & Exper constraints

There are several theoretical and experimental constraints that must be checked in order to obtain a viable model:

- Unitarity, Perturbativity, Vacuum Stability, EW Precision Observables (S, T and U); 2HDMC.
- Exclusion limits at 95% Confidence Level (CL) from Higgs searches at colliders (LEP, Tevatron and LHC); HiggsBounds.
- Constraints from the Higgs boson signal strength measurements; HiggsSignal.
- Constraints of flavour physics observables namely,  $B \rightarrow X_s \gamma, B_{s,d} \rightarrow \mu^+ \mu^-$  and  $\Delta m_{s,d}$ ; SuperIso.

## Numerical Analysis

We assume that the heaviest Higgs is the SM-like with mass of 125 GeV, whereas  $h$  and  $A$  would be lighter than  $H$ ,

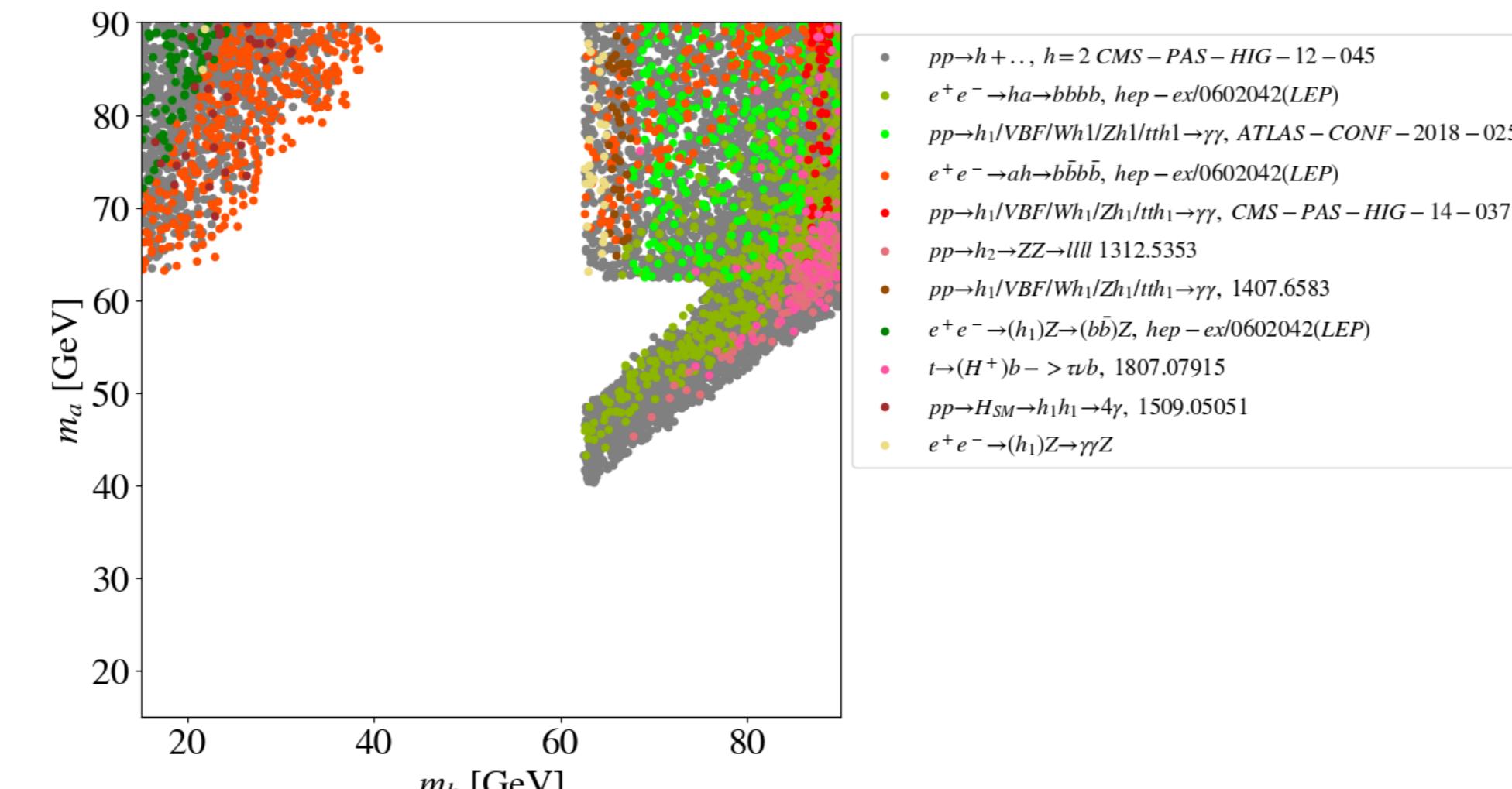


Figure 1: Allowed parameter space in Type-I at 95% CL vs. searches to which the region are sensitive to.

- Sensitivity in the region with low  $m_h$  and high  $m_A$  is mainly from LEP searches for processes such as  $e^+e^- \rightarrow ah \rightarrow b\bar{b}b\bar{b}$  and  $e^+e^- \rightarrow (h)Z \rightarrow (b\bar{b})Z$ .
- An update of LHC at Run3 is unlikely to rule out this mass combination  $(m_h, m_a)$

## Reinterpretation of exotic Higgs decays in Type-I

CMS search for  $pp \rightarrow H \rightarrow aa \rightarrow 2b2\tau$  at  $\sqrt{s} = 13$  TeV [1]:

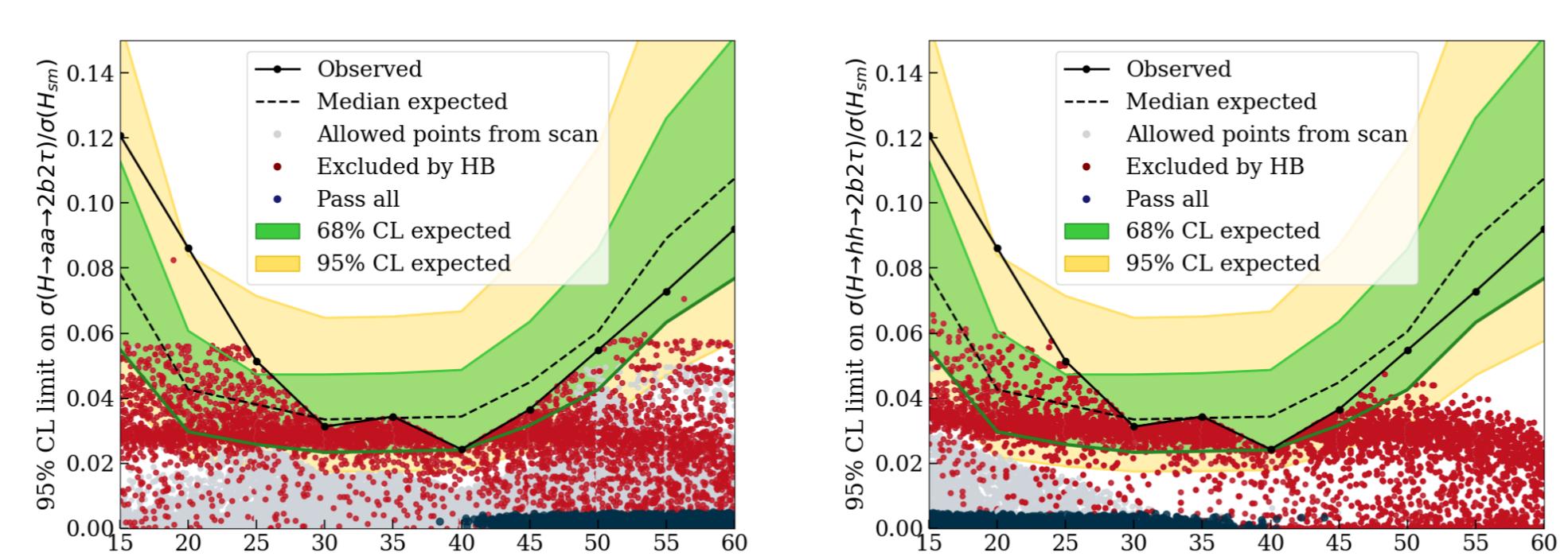


Figure 2: Observed and expected upper limits on  $\sigma(H \rightarrow aa(hh) \rightarrow 2b2\tau)/\sigma_{SM}(H)$  at 95% in 2HDM Type-I

- Recasting the  $2b2\tau$  search for  $H \rightarrow hh$ .
- Areas with sensitivity are excluded by previous searches.
- CMS searches at 13 TeV constrain  $B(H \rightarrow BSM) < 0.19$  at 95% CL.
- Type-I is more constrained by this search than  $H \rightarrow aa \rightarrow 2b2\mu, 2\tau2\mu$ .

CMS search for  $H^\pm \rightarrow W^\pm a \rightarrow \mu^+ \mu^-$  at  $\sqrt{s} = 13$  TeV [2]:

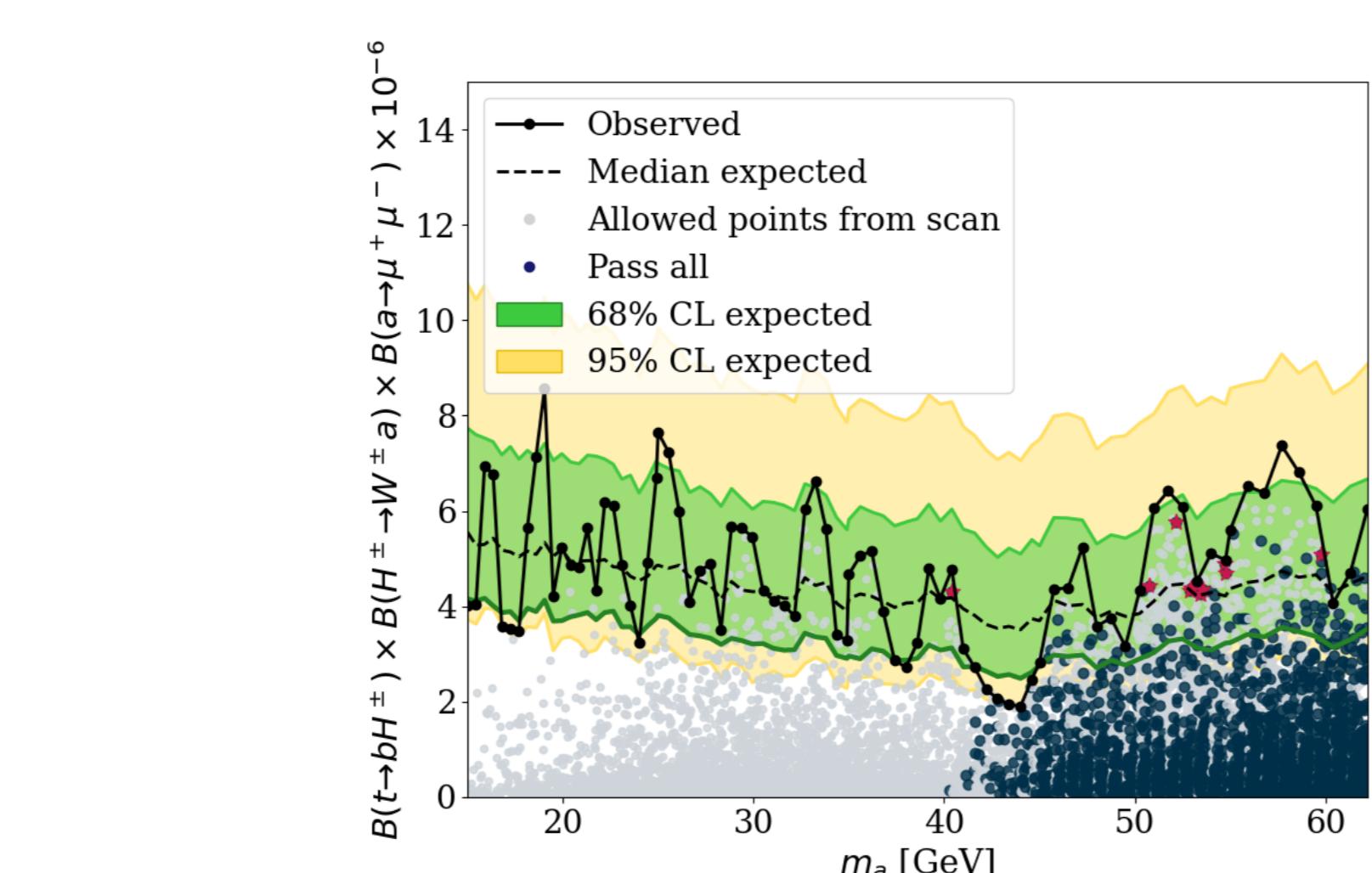


Figure 3: Observed and expected upper limits on  $B(t \rightarrow H+b) \rightarrow B(H^+ \rightarrow W^+ a) \times B(a \rightarrow \mu^+ \mu^-)$  at 95% in 2HDM type-I

- Bosonic decay of  $H^\pm$  can dominate fermionic channels in BSM when kinematically allowed.
- Type-I offers sufficient sensitivity, with an integrated luminosity of  $35.9 fb^{-1}$  (purple stars).

## New Signature: $pp \rightarrow H_{SM} \rightarrow Z^*Z^*h$

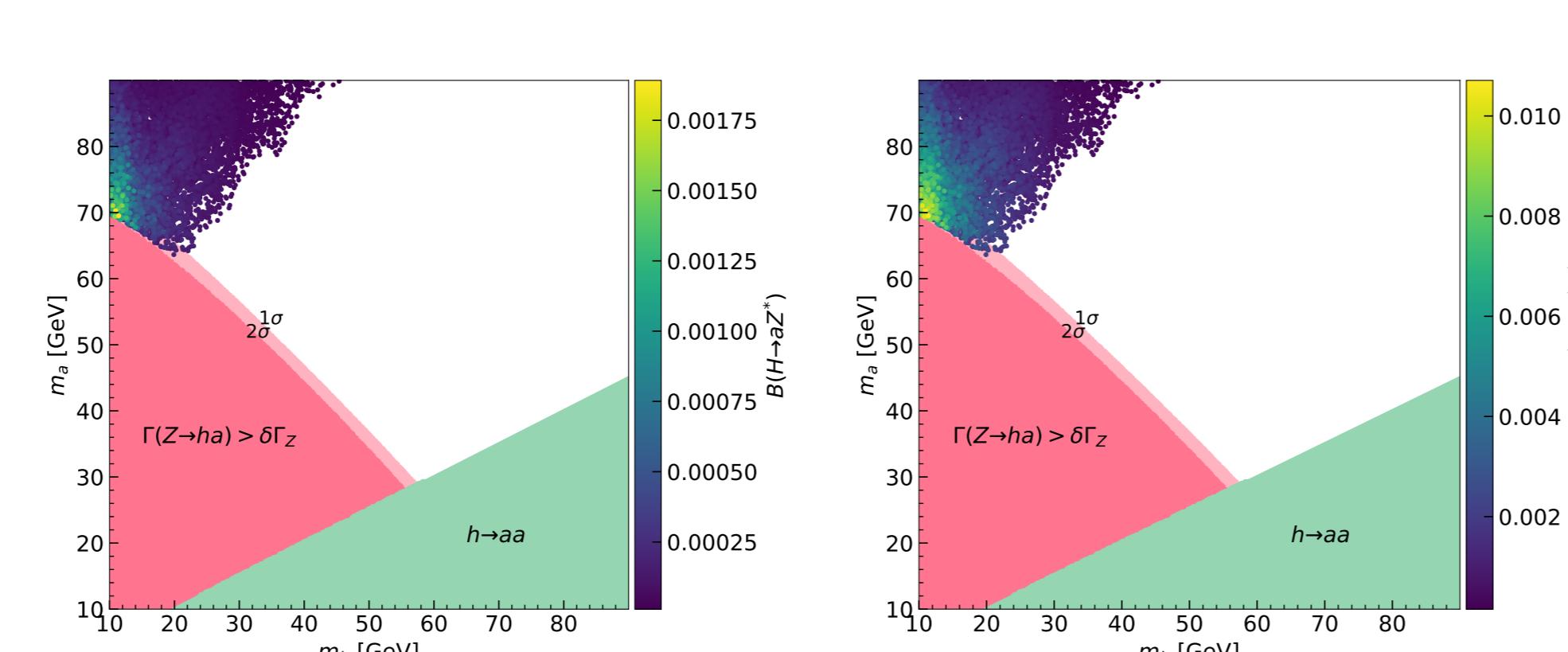


Figure 4:  $m_a$  and  $m_h$  vs.  $\sigma(gg \rightarrow H_{SM} \rightarrow Z^*a)$  (left panel) and  $\sigma(gg \rightarrow H_{SM} \rightarrow Z^*a \rightarrow Z^*Z^*h)$  (right panel) at 95% CL in 2HDM Type-I

- The subsequent decay of  $a$ , when the decay chain  $H \rightarrow aZ^*$  is open, could lead to  $a \rightarrow Z^*h$  with  $Z$  being off-shell and  $h$  decaying to fermions and/or  $\gamma\gamma$ .
- One could look for  $Z^*(\rightarrow 2\mu)Z^*(\rightarrow 2j)h(\rightarrow 2b)$  with di-muon trigger & standard  $|\eta(\mu)|$ .

## Signal vs. Background

- Few samples of BPs for the signal given by  $H \rightarrow aZ^* \rightarrow hZ^*Z^* \rightarrow \mu^+ \mu^- jj b\bar{b}$  are considered.
- Background processes with dominant contributions are top pair production in association with 2 initial state radiation (ISR) jets, and  $ZZ$  production with additional  $b\bar{b}$  quarks.

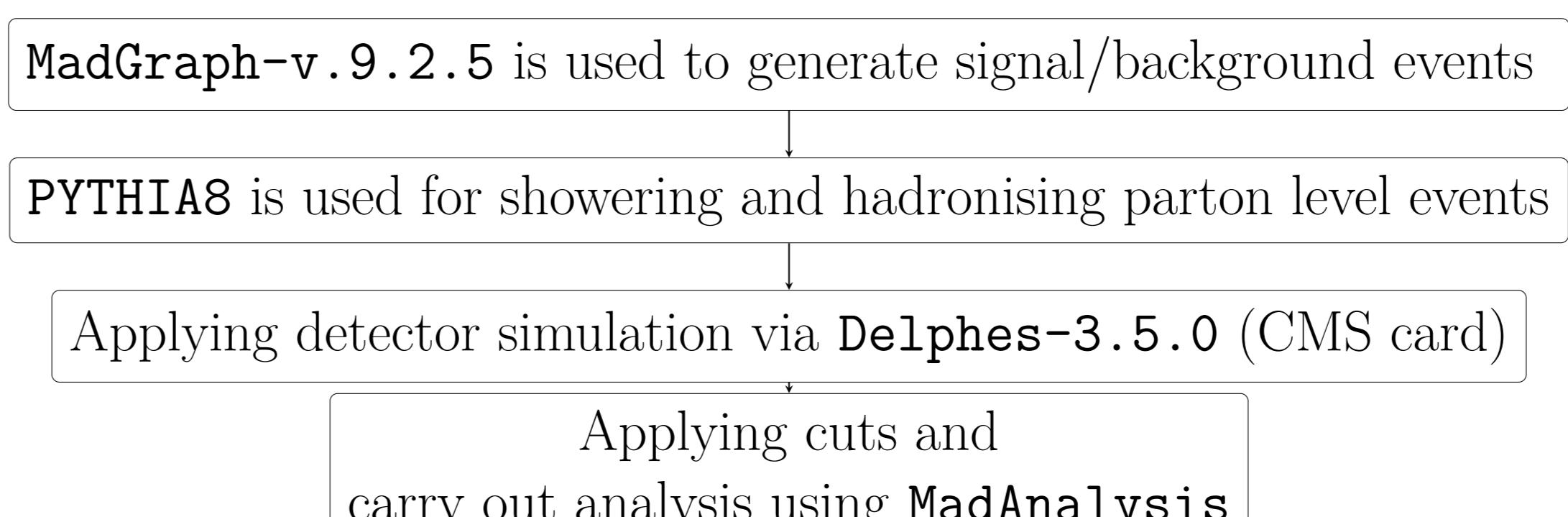
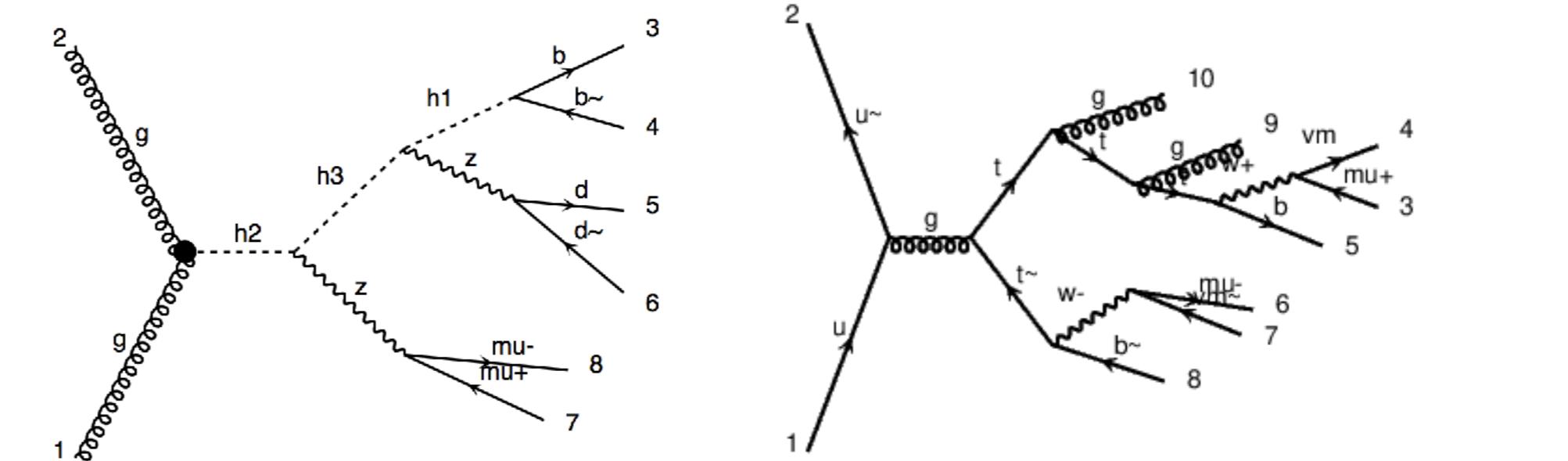


Figure 5: Description of the procedure used to generate and analyse MC events.

QCD corrections are considered through K-factor:

- NLO QCD correction to  $t\bar{t}$  production in association with 2-jets are of the order of -27%.
- NLO corrections to  $gg \rightarrow H$  are very large due to the contribution from  $gg$  radiation, whereas  $K_{NNLO}/NLO$  is much smaller than  $K_{NLO}/NLO$ .

Background	Cross section (pb)
$pp \rightarrow ZZb\bar{b}_{QCD} \rightarrow 6f$	$9.27 \times 10^{-3} \pm 2.4 \times 10^{-5}$
$pp \rightarrow ZZb\bar{b}_{QED} \rightarrow 6f$	$2.42 \times 10^{-4} \pm 5.5 \times 10^{-7}$
$pp \rightarrow gg\bar{t}\bar{t} \rightarrow gg\mu^+\mu^-jj$	$2.92 \pm 0.008$

Table 1: The background cross sections

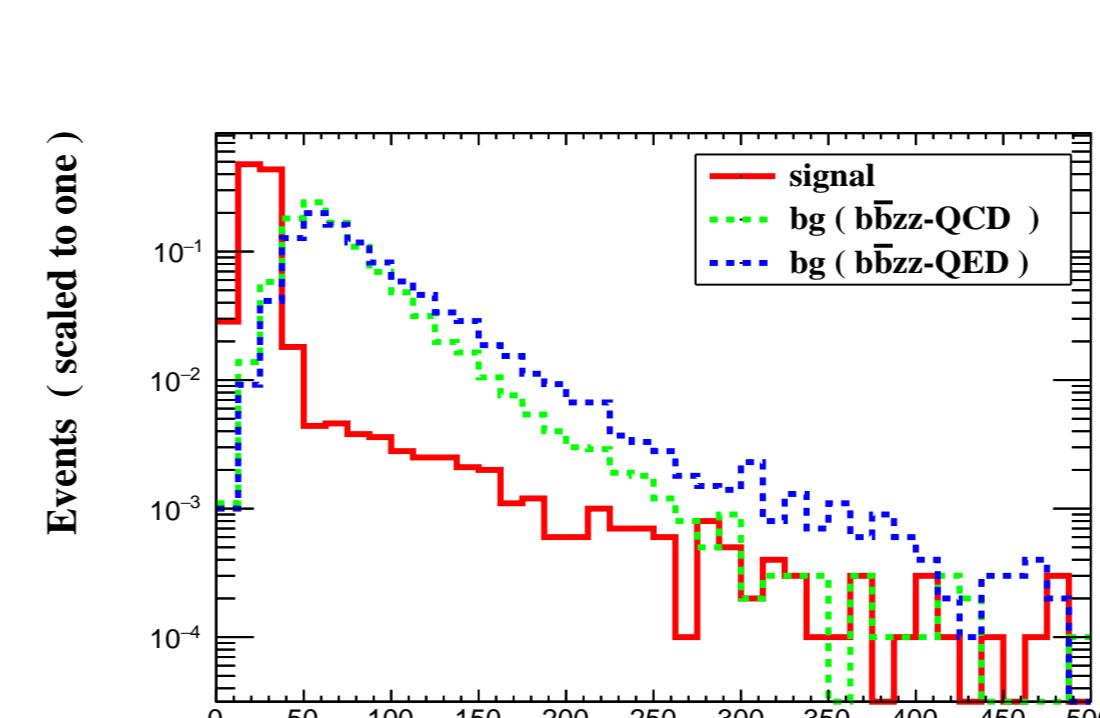


Figure 6:  $p_T(\mu_1)$  distribution at parton level

- Triggering on two muons reduces the required momenta to 17 GeV for the muon with higher  $p_T$ .
- Di-muon trigger increases the acceptance for numerous di-lepton processes and searches for new physics, compared with Isolated and Nonisolated single muon triggers.

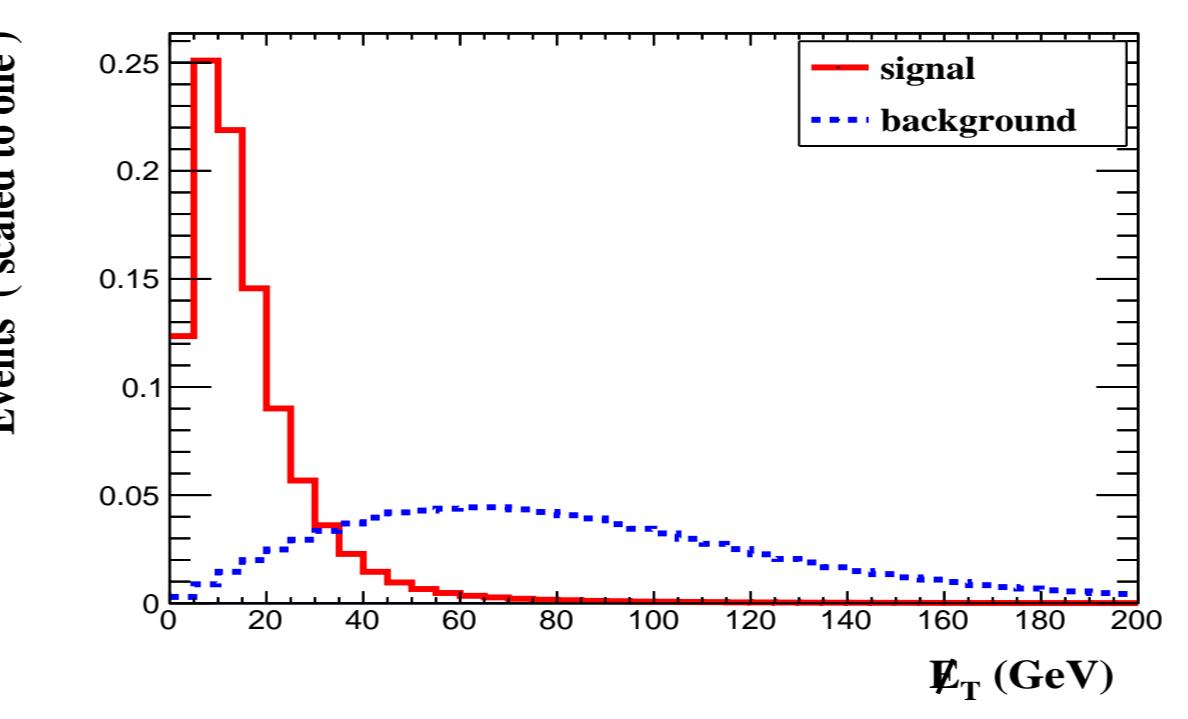


Figure 7:  $E_T$  distributions for the signal (red) and background (blue) at detector level

- $E_T$  distribution from simulated samples of background events is mainly from di-leptonic decay of  $t\bar{t}$  in association with 2 ISR jets, with  $t\bar{t} \rightarrow W^+bW^-b \rightarrow (\mu^+\nu_\mu b)(\mu^-\nu_\mu b)$ .
- The observed MET in signal at detector level is from semi-leptonic B-meson decay.

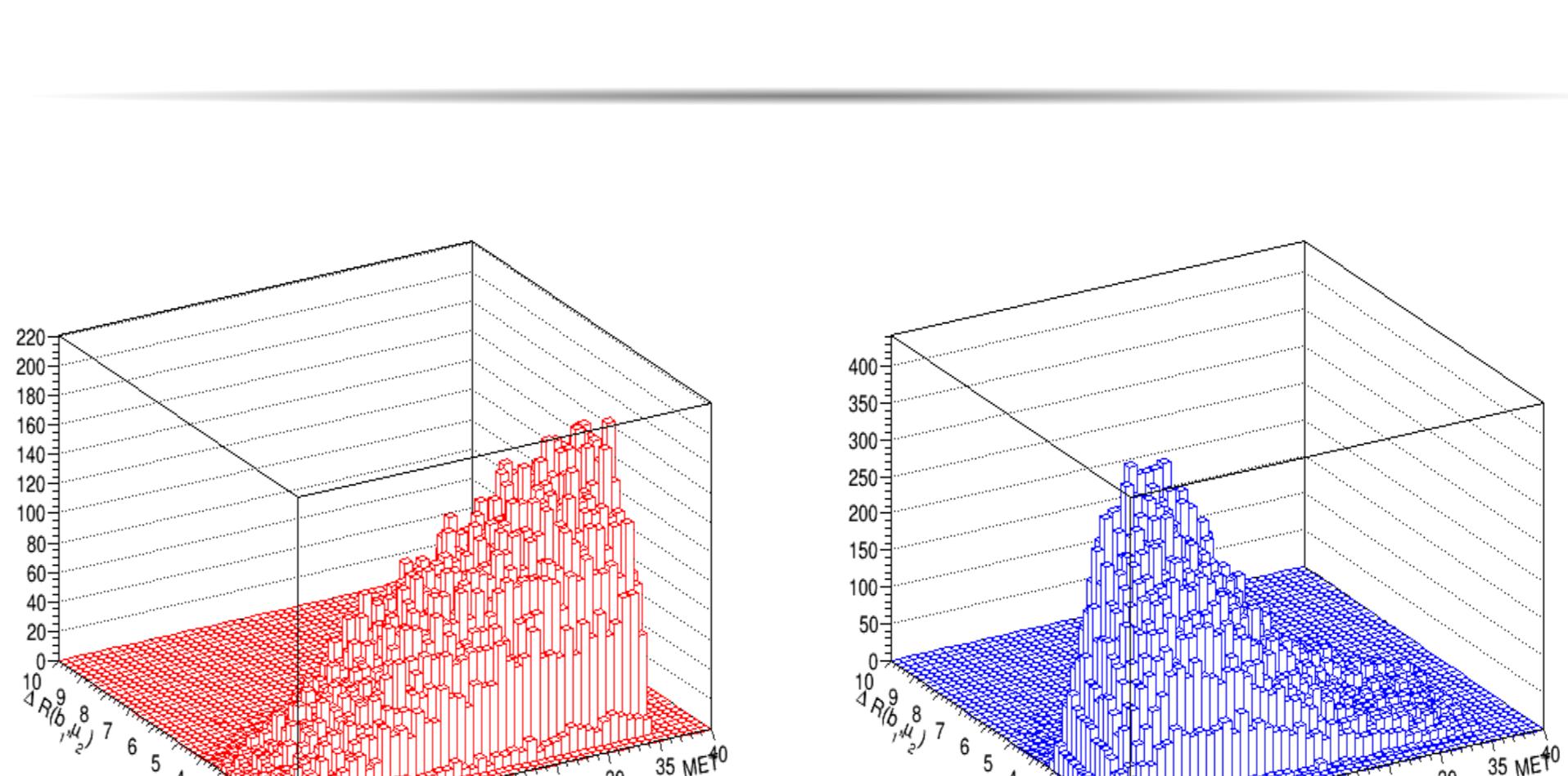


Figure 8:  $E_T$  vs.  $\Delta R(b_1, \mu_2)$ . Signal (blue) and background (red)

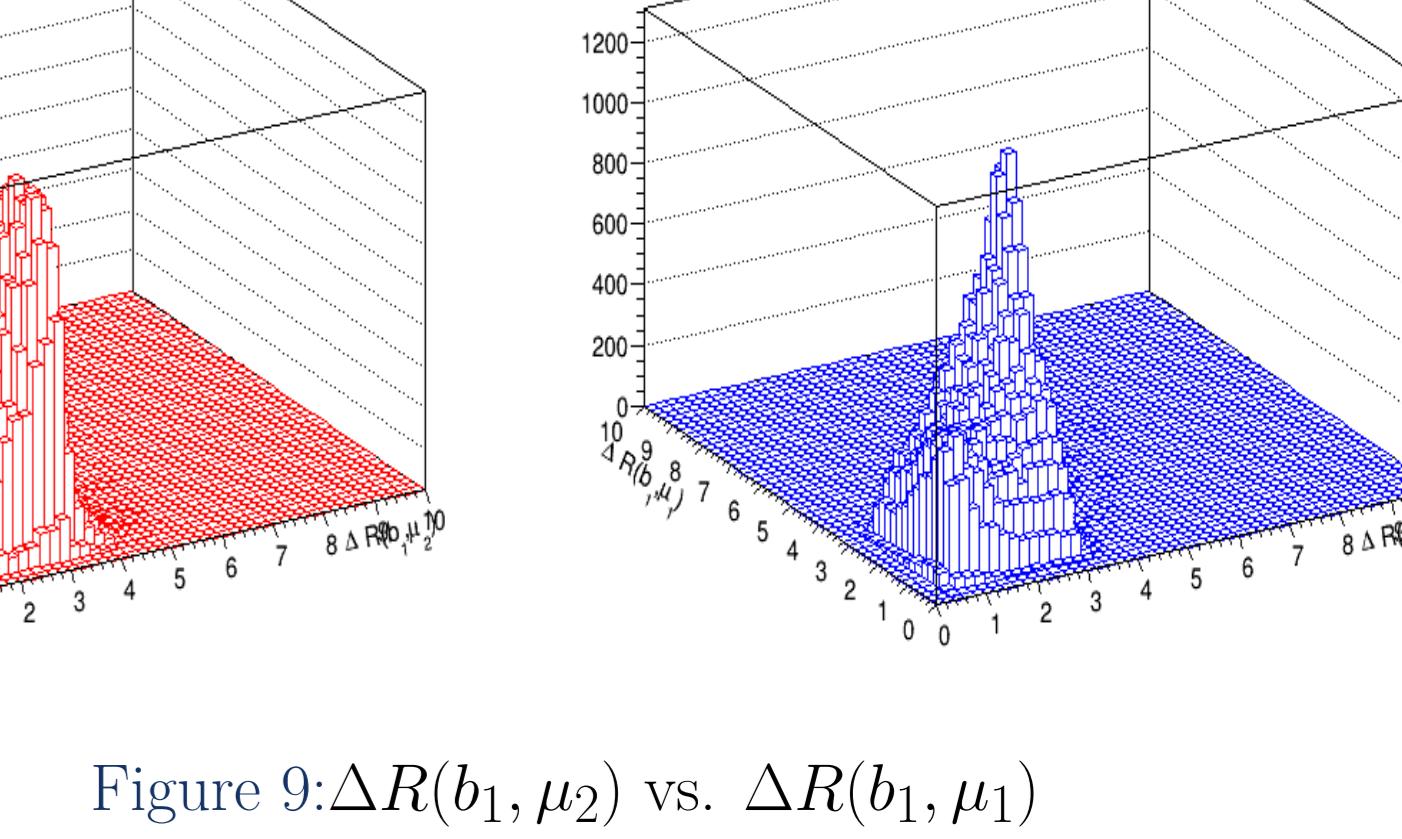


Figure 10: Significance and efficiency ( $\epsilon$ ) of each BP

- Some points on the grid have a significance larger than 3 on the  $(m_h, m_a)$  plane for  $\mathcal{L} = 300 fb^{-1}$ .
- A larger efficiency can be obtained in a parameter space with small  $m_a$ .

BP	$m_h$ (GeV)	$m_A$ (GeV)	$\sigma$ (pb)	k-factor	significance	$\epsilon$
BP1	11.85	72.75	$4.82 \times 10^{-4}$	2.69	4.88	0.0758
BP2	15.37	72.21	$3.28 \times 10^{-4}$	2.63	3.20	0.0757
BP3	17.15	76.24	$2.54 \times 10^{-4}$	2.63	2.29	0.0689
BP4	13.09	75.47	$3.538 \times 10^{-4}$	2.65	3.31	0.0709
BP5	14.15	74.35	$3.458 \times 10^{-4}$	2.62	3.29	0.0702
BP6	11.96	78.57	$3.557 \times 10^{-4}$	2.69	2.97	0.0662
BP7	12.60	77.17	$3.311 \times 10^{-4}$	2.66	2.87	0.065
BP9	14.30	76.77	$2.423 \times 10^{-4}$	2.63	2.31	0.0729
BP10	14.16	78.86	$2.572 \times 10^{-4}$	2.648	2.11	0.0602
BP20	11.83	74.06	$4.577 \times 10^{-4}$	2.69	4.51	0.073

Table 2: NNLO QCD correction to Higgs production, NLO QCD correction to  $t\bar{t}gg$ ,  $\epsilon = \sigma(cuts)/\sigma(no\ cuts)$  and  $\mathcal{L} = 300 fb^{-1}$ .

## Conclusion