

Investigation of the Right-handed Vector Current via Unbinned Angular Analysis of $B \rightarrow D^*(D\pi)\ell\nu_\ell$

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Introduction

V_{cb} puzzle:

Inclusive decay: $B \rightarrow X_c \ell \nu$ ($X_c = D, D^*, D_0^* \dots$)

Exclusive decay: $B \rightarrow D^{(*)} \ell \nu$

in. 42.16(50) vs ex. 39.70(60) $\sim 3\sigma$ tension

Right-handed vector current:

$$\mathcal{O}_{V_L} = (\bar{c}_L \gamma^\mu b_L)(\bar{\ell}_L \gamma_\mu \nu_L), \quad \mathcal{O}_{V_R} = (\bar{c}_R \gamma^\mu b_R)(\bar{\ell}_L \gamma_\mu \nu_L).$$

$$\mathcal{H}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} V_{cb} [C_{V_L} \mathcal{O}_{V_L} + C_{V_R} \mathcal{O}_{V_R}] + \text{h.c.}$$

Un-binned angular analysis:

Existing binned analysis (projected χ^2 fit): Belle '17 '18 [1,2];

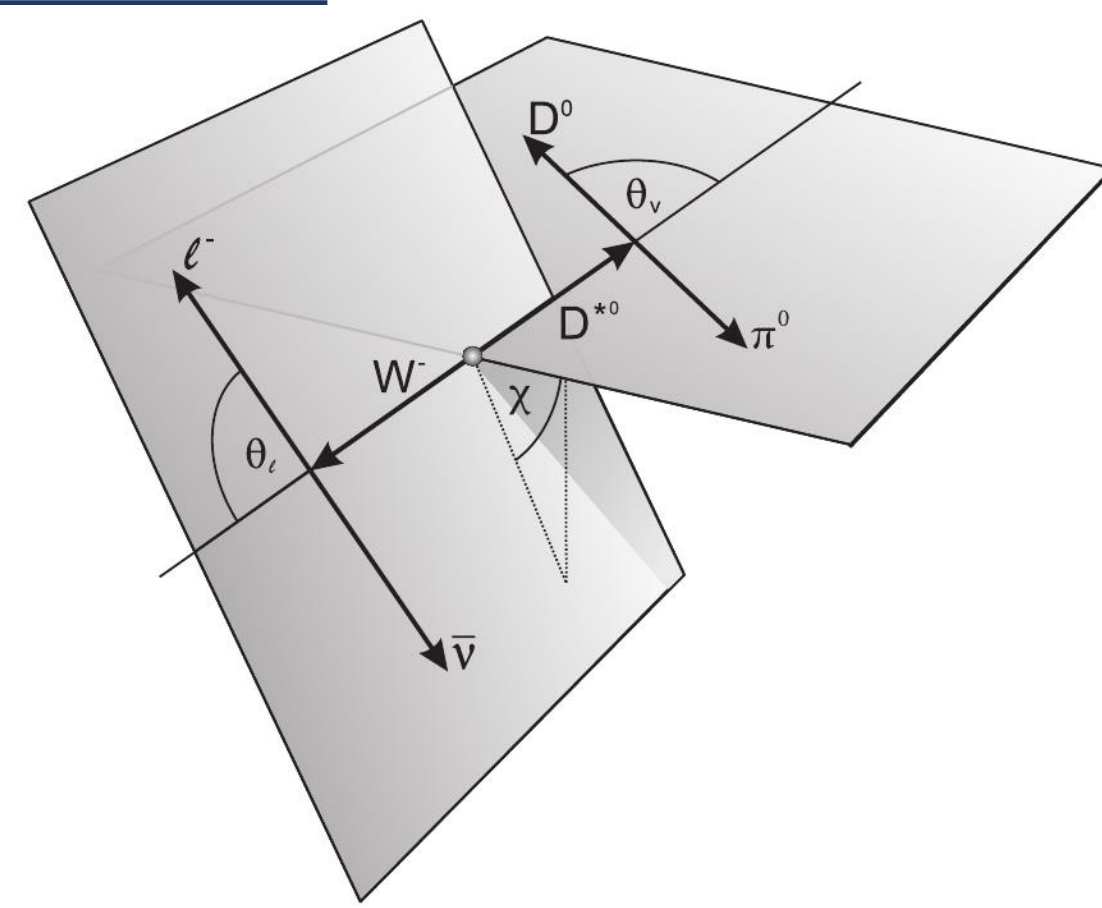
The experimental determination of $\langle g_i \rangle$ can be pursued by the maximum likelihood method:

$$\mathcal{L}(\langle g_i \rangle) = \sum_{i=1}^N \ln \hat{f}_{\langle g_i \rangle}(e_i)$$

Generation of Pseudo Data

Normalized probability density function (PDF):

$$\begin{aligned} \hat{f}_{\langle g_i \rangle}(\cos \theta_V, \cos \theta_\ell, \chi) = & \frac{9}{8\pi} \\ & \times \left\{ \frac{1}{6}(1 - 3\langle g_{1c} \rangle + 2\langle g_{2s} \rangle + \langle g_{2c} \rangle) \sin^2 \theta_V + \langle g_{1c} \rangle \cos^2 \theta_V \right. \\ & + (\langle g_{2s} \rangle \sin^2 \theta_V + \langle g_{2c} \rangle \cos^2 \theta_V) \cos 2\theta_\ell \\ & + \langle g_3 \rangle \sin^2 \theta_V \sin^2 \theta_\ell \cos 2\chi \\ & + \langle g_4 \rangle \sin 2\theta_V \sin 2\theta_\ell \cos \chi + \langle g_5 \rangle \sin 2\theta_V \sin \theta_\ell \cos \chi \\ & + (\langle g_{6s} \rangle \sin^2 \theta_V + \langle g_{6c} \rangle \cos^2 \theta_V) \cos \theta_\ell \\ & + \langle g_7 \rangle \sin 2\theta_V \sin \theta_\ell \sin \chi + \langle g_8 \rangle \sin 2\theta_V \sin 2\theta_\ell \sin \chi \\ & \left. + \langle g_9 \rangle \sin^2 \theta_V \sin^2 \theta_\ell \sin 2\chi \right\}, \end{aligned}$$



Kinematic angles for $B \rightarrow D^*(D\pi)\ell\nu_\ell$

CLN parametrization:

$$h_{A_1}(w) = h_{A_1}(1)(1 - 8\rho_{D^*}^2 z + (53\rho_{D^*}^2 - 15)z^2 - (231\rho_{D^*}^2 - 91)z^3),$$

$$R_1(w) = R_1(1) - 0.12(w-1) + 0.05(w-1)^2, \quad \mathcal{F}_1(z) = \frac{1}{P_{\mathcal{F}_1}(z)\phi_{\mathcal{F}_1}(z)} \sum_{n=0}^N a_n^{\mathcal{F}_1} z^n.$$

BGL parametrization:

$$g(z) = \frac{1}{P_g(z)\phi_g(z)} \sum_{n=0}^N a_n^g z^n, \quad f(z) = \frac{1}{P_f(z)\phi_f(z)} \sum_{n=0}^N a_n^f z^n,$$

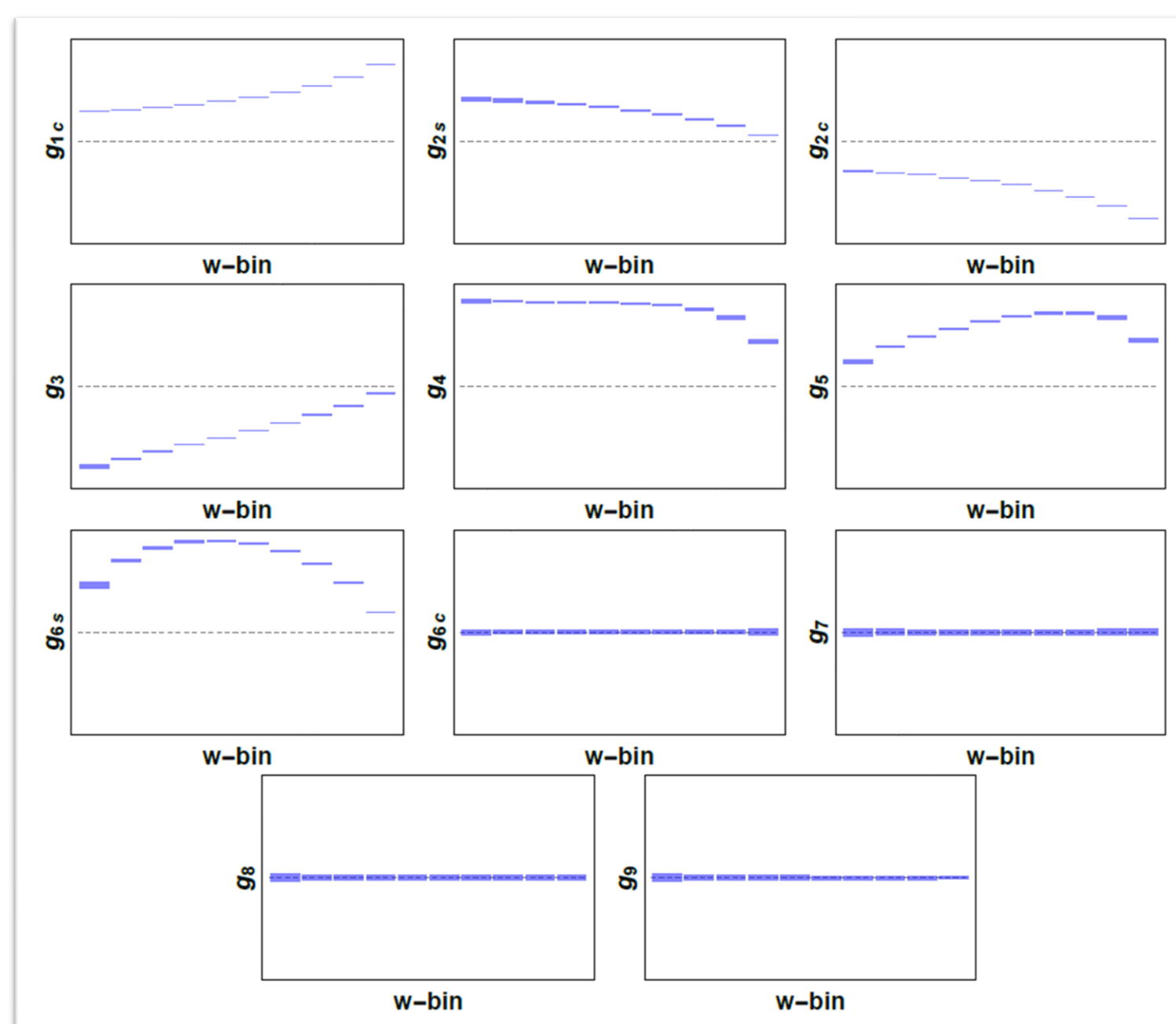
Pseudo data for $\langle g_i \rangle$ in the CLN parametrization:

$N_{\text{event}} = (5306, 8934, 10525, 11241, 11392, 11132, 10555, 9726, 8693, 7497)$

Pseudo data for $\langle g_i \rangle$ in the BGL parametrization:

$N_{\text{event}} = (5239, 8868, 10500, 11264, 11455, 11217, 10638, 9776, 8676, 7368)$

$\langle g_i \rangle$ generated in 10 w bins with covariance matrices by toy Monte-Carlo method



$\langle g_i \rangle$ generated in 10 w bins in the CLN parametrization

Results of Sensitivity Study

CLN fit:

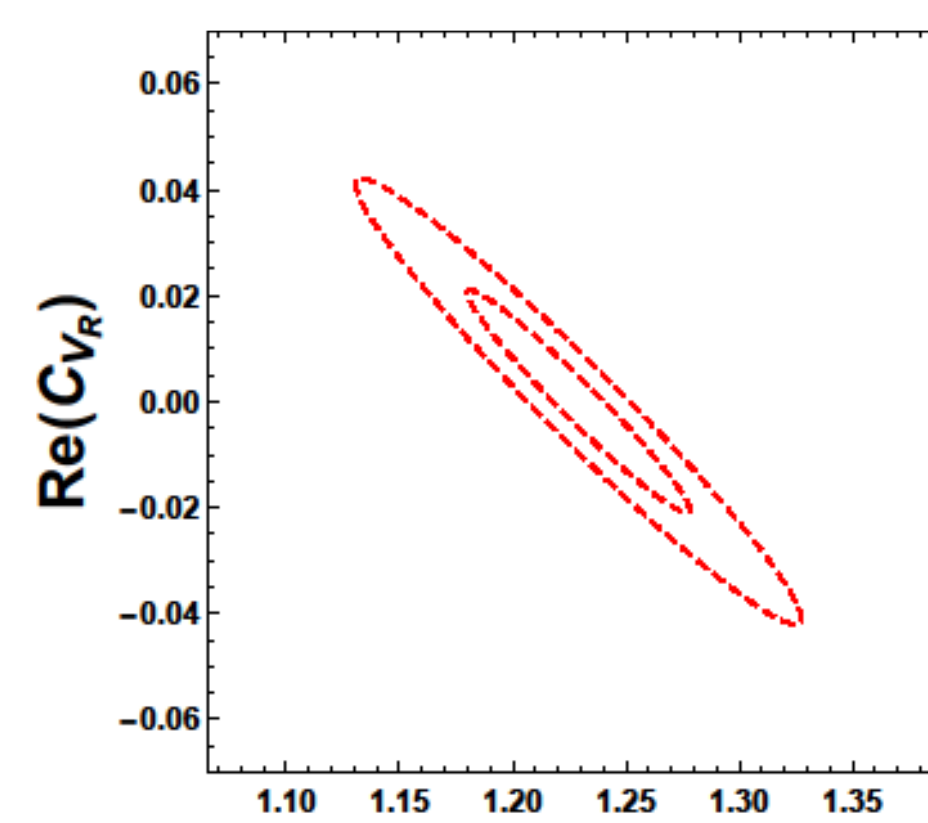
$$R_1(1) = \frac{h_V(1)}{h_{A_1}(1)}$$

$$\vec{v} = (\rho_{D^*}^2, R_1(1), R_2(1), C_{V_R}) = (1.106, 1.229, 0.852, 0)$$

$$\sigma_{\vec{v}} = (3.177, 0.049, 0.018, 0.021)$$

$$\rho_{\vec{v}} = \begin{pmatrix} 1. & -0.016 & -0.763 & 0.095 \\ -0.016 & 1. & 0.006 & -0.973 \\ -0.763 & 0.006 & 1. & -0.117 \\ 0.095 & -0.973 & -0.117 & 1. \end{pmatrix}$$

C_{V_R} can be determined to precision of ~ 2 (4)% in CLN (BGL) parametrization.



$R_1(1) - C_{V_R}$ and $h_{V_1}(1) - C_{V_R}$ contours

BGL fit:

$$h_V(1) = \frac{m_B \sqrt{r}}{P_g(0)\phi_g(0)} a_0^g$$

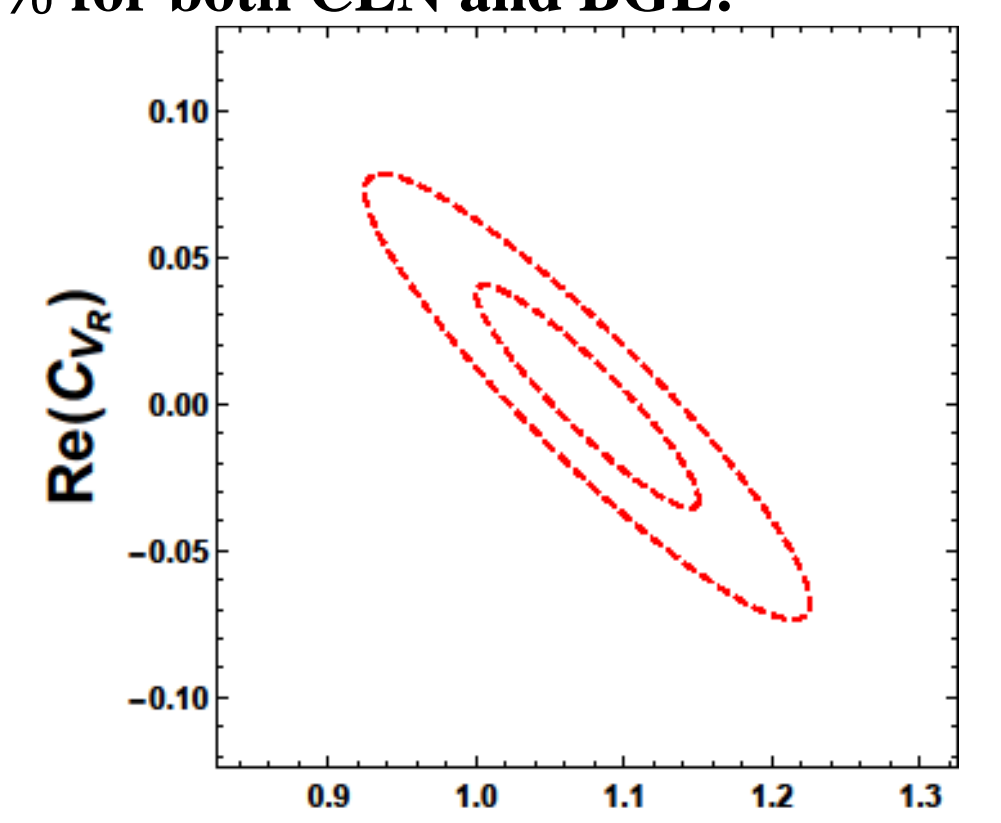
$$\vec{v} = (a_0^g, a_1^g, a_2^g, a_3^g, a_4^g, C_{V_R}) = (0.0132, 0.0169, 0.0070, -0.0852, 0.0241, 0.0024)$$

$$\sigma_{\vec{v}} = (0.0002, 0.0109, 0.0026, 0.0352, 0.0017, 0.0379)$$

$$\rho_{\vec{v}} = \begin{pmatrix} 1. & 0.022 & 0.039 & -0.035 & 0.000 & 0.189 \\ 0.022 & 1. & 0.860 & -0.351 & 0.000 & 0.316 \\ 0.039 & 0.860 & 1. & -0.762 & 0.000 & 0.283 \\ -0.035 & -0.351 & -0.762 & 1. & 0.000 & -0.119 \\ 0.000 & 0.000 & 0.000 & 0.000 & 1. & -0.923 \\ 0.189 & 0.316 & 0.283 & -0.119 & -0.923 & 1. \end{pmatrix}$$

C_{V_R} and the vector form factor are highly correlated!

$\text{Im}(C_{V_R})$ can also be determined at precision of 0.7% for both CLN and BGL!



Fit to the forward-backward asymmetry (A_{FB}) of the charged lepton:

Advantage: one angle measurement

$$\langle A_{FB} \rangle \equiv \frac{\int_0^1 \frac{d\Gamma}{d\cos\theta_\ell} d\cos\theta_\ell - \int_{-1}^0 \frac{d\Gamma}{d\cos\theta_\ell} d\cos\theta_\ell}{\int_0^1 \frac{d\Gamma}{d\cos\theta_\ell} d\cos\theta_\ell + \int_{-1}^0 \frac{d\Gamma}{d\cos\theta_\ell} d\cos\theta_\ell} = 3\langle g_{6s} \rangle$$

$$\vec{v} = (\rho_{D^*}^2, R_1(1), R_2(1), C_{V_R}) = (1.106, 1.229, 0.852, 0.000)$$

$$\sigma_{\vec{v}} = (2.200, 0.049, 0.031, 0.022)$$

C_{V_R} can be determined at a precision of 2.2% using A_{FB} alone! Almost as good as the full set of $\langle g_i \rangle$!

Summary

- The **un-binned angular analysis** is useful for the precision measurement of C_{V_R} by circumventing the V_{cb} puzzle.
- C_{V_R} is strongly correlated to the vector form factor.
- The real (imaginary) part of C_{V_R} can be determined at precision of 2-4 (1) % using the full set of $\langle g_i \rangle$.
- A_{FB} ($\langle g_{6s} \rangle$) alone can determine C_{V_R} at almost equally good precision as the full set of $\langle g_i \rangle$.

References

- [1] E. Waheed et al. [Belle], Phys. Rev. D 100, no.5, 052007 (2019) [erratum: Phys. Rev. D 103, no.7, 079901 (2021)]
- [2] A. Abdesselam et al. [Belle], [arXiv:1702.01521 [hep-ex]]

Questions and comments are welcome!

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