

Asymmetric dark matter

Part I - Models

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Part II by Hai-Bo Yu



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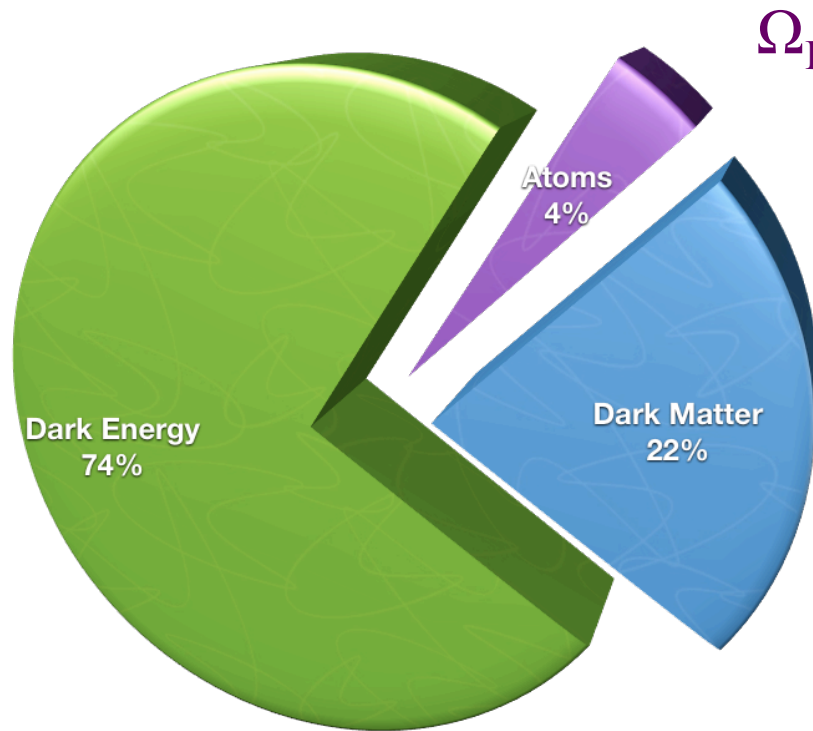
Mads Toudal Frandsen
Rudolf Peierls Centre for Theoretical Physics

Dark Matter Underground and in the Heavens, CERN, July 2011

Outline

- Motivate asymmetric dark matter genesis
- Overview of baryogenesis and ADM-genesis frameworks
- ADM from Technicolor - New Strong Dynamics

What is the world made of?



$$\Omega_B \sim 0.05$$

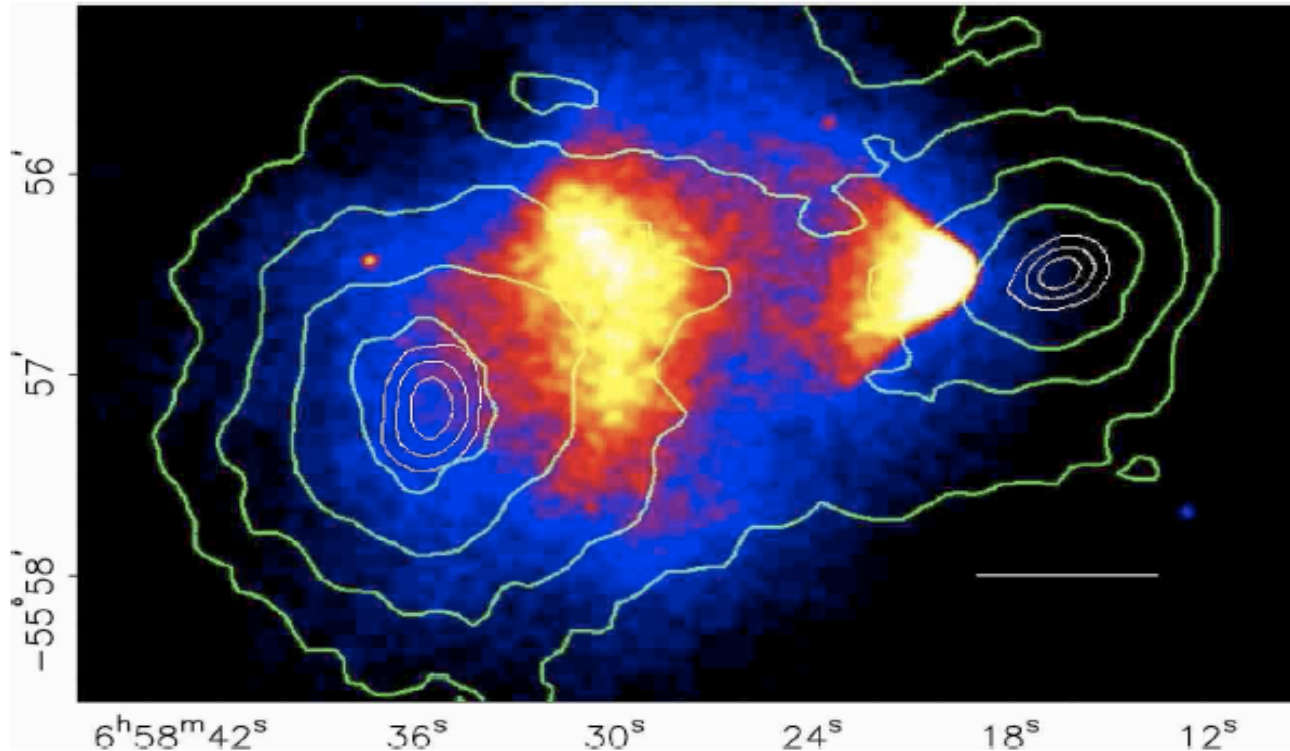
Baryons (but *no* antibaryons)

$$\Omega_{DM}/\Omega_B \approx 5$$

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
Λ_{QCD}	Baryons	U(1) Baryon number	$\tau > 10^{33}$ yr (dim-6 OK)	'freeze-out' from thermal equilibrium Asymmetry	$\Omega_B \sim 10^{-10}$ <i>cf. observed</i> $\Omega_B \sim 0.05$

What is the strength or scale of DM (self-interactions) Λ_χ ?

$\sigma_{\chi\chi} \leq 10^{-24} \text{ cm}^2/\text{GeV}$ from the ‘Bullet cluster’

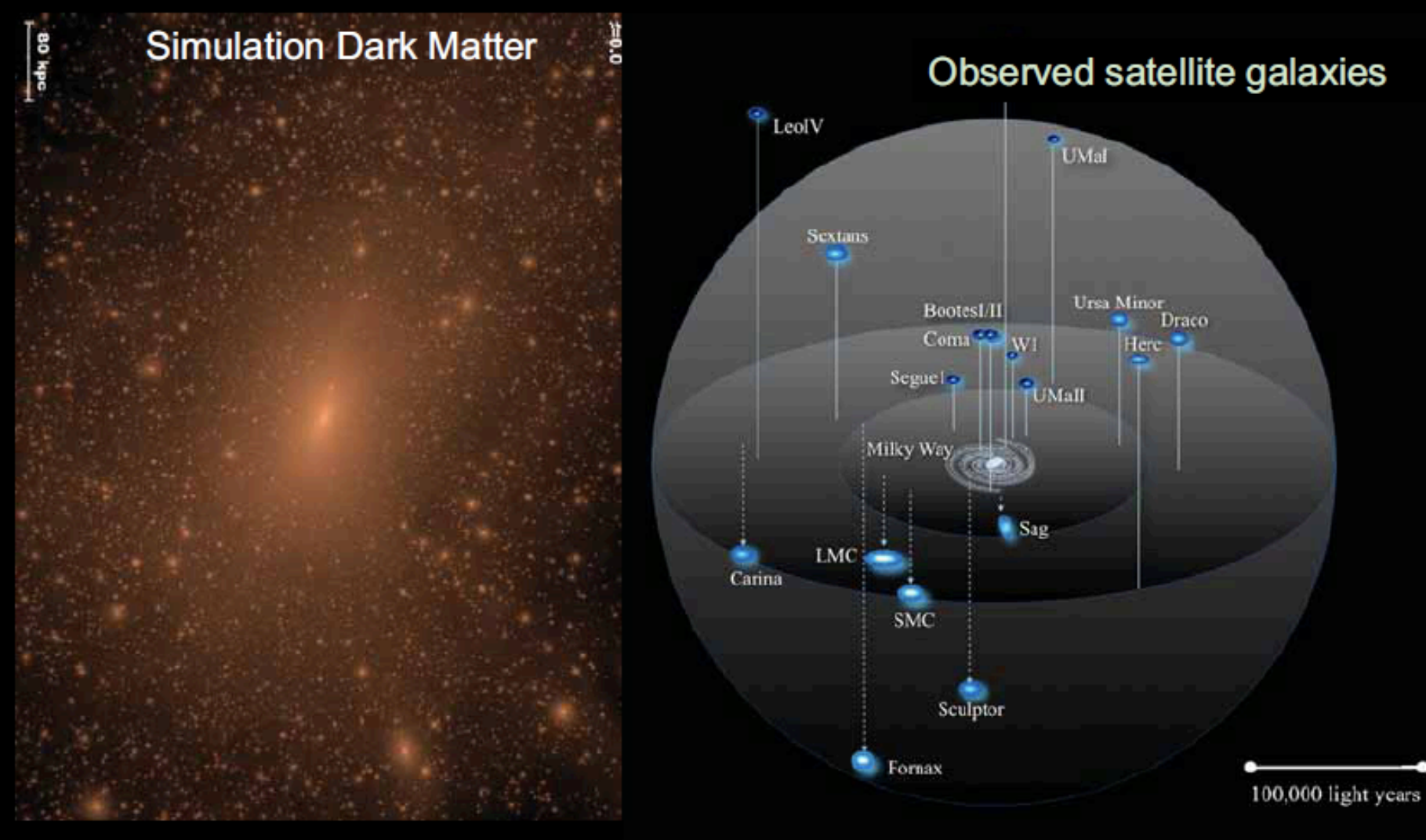


The **self-interaction** cross-section scaled up from QCD is $\sigma_{\chi\chi} \sim \sigma_{nn} (\Lambda_n/\Lambda_\chi)^2$, where $\sigma_{nn} \sim 10^{-23} \text{ cm}^2$, so in this case $\Lambda_\chi \geq 5 \text{ GeV}$

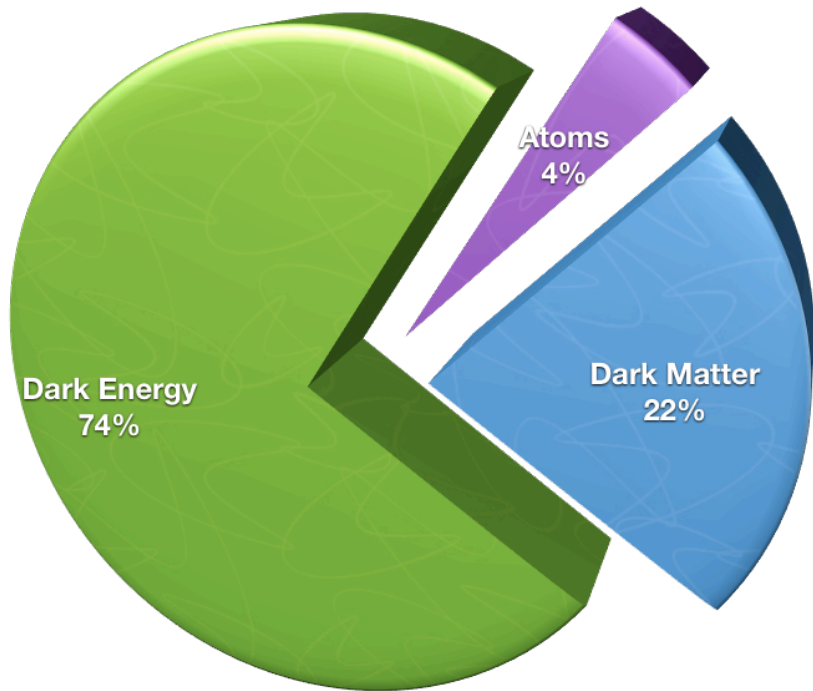
(However, see part-II of this talk by Hai-Bo Yu)

Amusingly $\Lambda_\chi \sim \sqrt{f_\pi v_{\text{EW}}} \approx 5 \text{ GeV}$

Self-interacting dark matter has been invoked to reduce excessive substructure in simulations of *collisionless* dark matter (Spergel & Steinhardt '99)



(See also talk by M. Kuhlen this workshop)



$$\Omega_B \sim 0.05$$

Baryons (but *no* antibaryons)

$$\Omega_{DM}/\Omega_B \approx 5$$

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
Λ_{QCD}	Baryons	Global U(1) Baryon number	$\tau > 10^{33}$ yr (dim-6 OK)	'freeze-out' from thermal equilibrium <i>Asymmetry</i>	$\Omega_B \sim 10^{-10}$ <i>cf. observed</i> $\Omega_B \sim 0.05$

What should the world be made of?

What do we expect for the *symmetric* thermal relic abundance of baryons?

$$\dot{n} + 3Hn = -\langle\sigma v\rangle(n^2 - n_T^2)$$

Chemical equilibrium is maintained as long as annihilation rate exceeds the Hubble expansion rate

‘Freeze-out’ occurs when annihilation rate:

$$\Gamma = n\sigma v \sim m_N^{3/2} T^{3/2} e^{-m_N/T} \frac{1}{m_\pi^2}$$

becomes comparable to the expansion rate

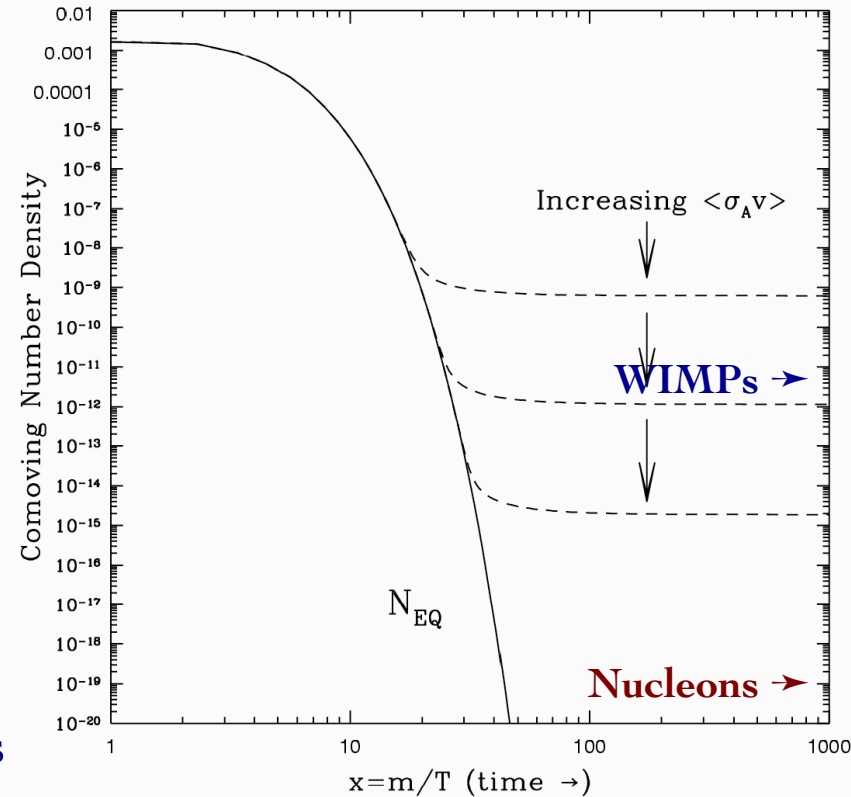
$$H \sim \frac{\sqrt{g}T^2}{M_P} \text{ where } g \Rightarrow \# \text{ relativistic species}$$

i.e. freeze-out occurs at $T \sim m_B/45$, with:

$$\frac{n_B}{n_\gamma} = \frac{n_{\bar{B}}}{n_\gamma} \sim 10^{-19}$$

However the observed ratio is 10^9 times *bigger* for baryons, and there are *no* antibaryons, so we must invoke an **initial asymmetry**:

$$\mathcal{N}_B \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-9}$$



What *should* the world be made of ?

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
Λ_{QCD}	Nucleons	Baryon number	$\tau > 10^{33}$ yr (dim-6 OK)	'freeze-out' from thermal equilibrium Asymmetric baryogenesis	$\Omega_{\text{B}} \sim 10^{-10}$ <i>cf. observed</i> $\Omega_{\text{B}} \sim 0.05$
$\Lambda_{\text{Fermi}} \sim G_{\text{F}}^{-1/2}$	Neutralino?	<i>R</i> -parity?	Violated? (matter parity adequate for p stability)	'freeze-out' from thermal equilibrium	$\Omega_{\text{LSP}} \sim 0.25$

The 'WIMP miracle' naturalised e.g. in (softly broken) **supersymmetry** :

$$\Omega_{\chi} h^2 \simeq \frac{3 \times 10^{-27} \text{cm}^{-3} \text{s}^{-1}}{\langle \sigma_{\text{ann}} v \rangle_{T=T_{\text{f}}}} \simeq 0.1 \quad \text{since } \langle \sigma_{\text{ann}} v \rangle \sim \frac{g_{\chi}^4}{16\pi^2 m_{\chi}^2} \approx 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$$

In this framework it is puzzling why $\Omega_{\text{DM}}/\Omega_{\text{B}} \approx 5$?

Natural to speculate that DMgenesis χ may be linked to baryogenesis with an asymmetric relic density related to that of baryons:

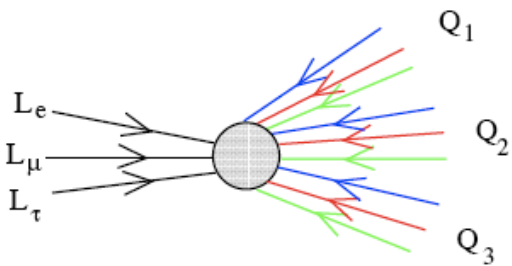
$$\Omega_{\chi} = (m_{\chi} \mathcal{N}_{\chi} / m_{\text{B}} \mathcal{N}_{\text{B}}) \Omega_{\text{B}}$$

Sakharov conditions for baryogenesis:

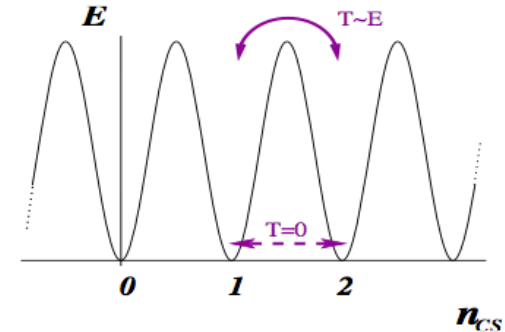
1. Baryon number violation
2. C and CP violation
3. Departure from thermal equilibrium

Classically baryon number can be violated by dim-6 operators in SM

When $T \gtrsim M_W$ Baryon number is also violated in the SM through sphaleron-mediated processes that preserve $B - L$, but violate $B + L \dots$



$$\partial_\mu j_i^\mu = \partial_\mu (\bar{\psi}^i \gamma^\mu \psi^i) = \frac{g^2}{8\pi} W^{a\mu\nu} \tilde{W}_{\mu\nu}^a$$



... CP -violation is *too weak* (and the electroweak phase transition is a 'cross-over')

The observed matter-antimatter asymmetry *requires* BSM physics
The matter we *know* originated *non-thermally* in the early universe

The same or a similar mechanism could generate the DM density

(See list of references in part-II of this talk by Hai-Bo Yu)

General features of ADM

- ADM is a complex field, either Dirac fermion or complex scalar
- Requires asymmetry generating or transferring operator and
- large annihilation cross-section to suppress symmetric relic density

(Griest & Seckel '87; Hooper, March-Russell & West '05; Belyaev, M.T.F, Sannino & Sarkar '10; Graesser, Shoemaker & Vecchi 11; Buckley 11)

Example: Symmetric relic density of composite scalar pNGB DM state ϕ

Overproduction of DM, below W threshold

$$\mathcal{L} = \partial_\mu \phi^* \partial_\mu \phi - m_\phi^2 \phi^* \phi + \frac{d_1}{\Lambda} H \partial_\mu \phi^* \partial_\mu \phi \quad (2)$$

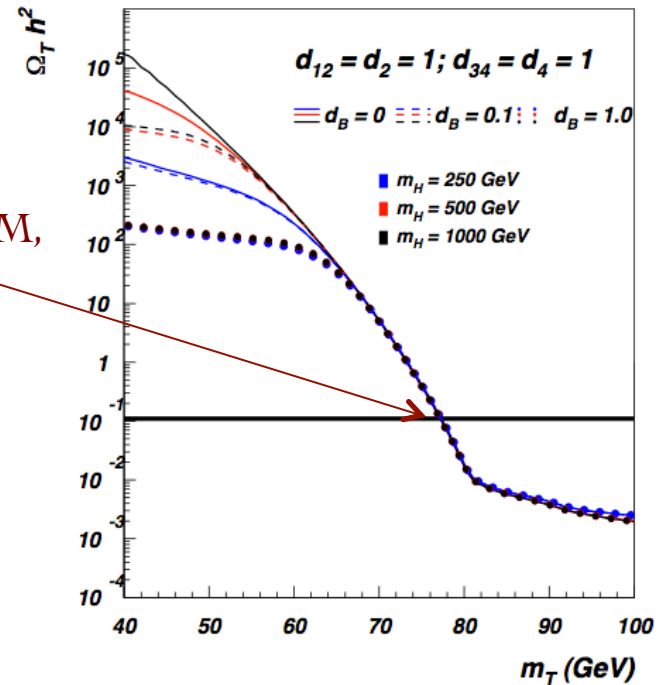
$$+ \frac{d_2}{\Lambda} m_\phi^2 H \phi^* \phi + \frac{d_3}{2\Lambda^2} H^2 \partial_\mu \phi^* \partial_\mu \phi + \frac{d_4}{2\Lambda^2} m_\phi^2 H^2 \phi^* \phi.$$

If constituents of ϕ has weak charges it has additional interactions, lets call it T !

$$\mathcal{L}_B = ie \frac{d_B}{\Lambda^2} T^* \overleftrightarrow{\partial}_\mu T \partial_\nu F^{\mu\nu}$$

$$L_{WW,ZZ} = -\frac{T^* T}{2} \text{Tr} [d_W W_\mu W^\mu + d_Z Z_\mu Z^\mu]$$

(M.T.F & Sannino '10; Belyaev, M.T.F, Sannino & Sarkar '10)



If ADM is light require new light *hidden sector* states

(e.g. Strassler & Zurek '06; Cohen, Phalen Pierce & Zurek '10; M.T.F, Sarkar & Schmidt-Hoberg '11)

Some frameworks for baryogenesis:

1. GUT scale baryogenesis (Decay of GUT scale coloured states) (Yoshimura '78; Dimopoulos & Susskind '78)
2. Electroweak baryogenesis (EWB) (sphalerons & 1st order phase transition) (Kuzmin, Rubakov & Shaposhnikov '85)
3. Leptogenesis (sphalerons, 1st or 2nd order phase transition) (Fukugita & Yanagida '86)
4. Affleck-Dine baryogenesis (Decays of coherent scalar fields in SUSY) (Affleck & Dine '84)
5. Spontaneous baryogenesis (Cohen & Kaplan '87)
6. Cold baryogenesis (Krauss & Trodden '99)
7. ... (Reviews in Dine & Kusenko '03; Cline '06; Schmidt '11)

DMgenesis can be linked to baryogenesis in *any of these frameworks*

Whatever the baryo/lepto-genesis mechanism, SM sphalerons may transfer an asymmetry if DM (constituents) carry weak charges

frameworks for DM-genesis linked to baryogenesis:

1. Via (e.g.) GUT scale baryogenesis

(Original technibaryon ADM from technicolor proposal) (Nussinov 85)

2. Via Electroweak baryogenesis (EWB)

(Kaplan 92)

(requires 1st order EW phase transition)

3. Via sphalerons/leptogenesis

(sphalerons, 1st or 2nd order phase transition)

(Barr, Chivukula & Farhi 90; Gudnason, Kouvaris & Sannino 05)

4. Via Affleck-Dine baryogenesis

(e.g. Enqvist & McDonald '98)

5. Via spontaneous baryogenesis

(March-Russell & McCullough '11)

6. Cold baryogenesis

(Konstandin & Servant '11)

7. ...

Common for 2., 3. is that DM carries $U(1)_X$ charge with $U(1)_X \times SU(2)_L^2$ anomaly

In 2. sphalerons produce the baryon/DM asymmetry – co-generation.

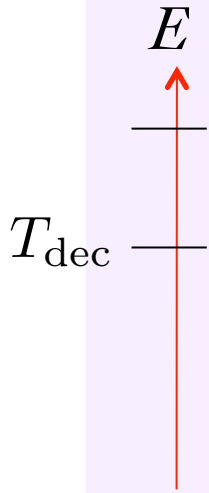
In 3. sphalerons transfer/process an existing asymmetry into DM – sharing.

Note that the DM or associated new physics sector may also change the EW phase-transition relevant for 2.

(Grojean, Servant & Wells '04; Cline, Jarvinen & Sannino 08,

Jarvinen, Rytto & Sannino 09,...)

Some classification of ADM models – models with asymmetry transfer

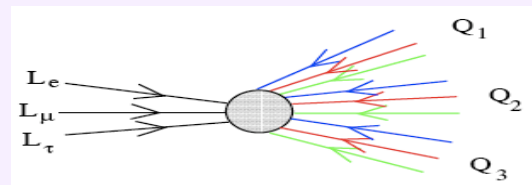


Asymmetry transfer/sharing:

- B or L asymmetry generated at a high scale, e.g. Leptogenesis.
- Below B-L is preserved, but transferred to DM
- Transfer operator decouples and asymmetry is 'frozen in'
- Symmetric component is annihilated away

Transfer via X and weak charges of DM (constituents)

$U(1)_X \times SU(2)_L^2$ anomaly, dim-4 (nonpert.) SM sphalerons, with $T_{\text{sph}} \sim v_{\text{EW}}$



(e.g. Barr, Chivukula & Farhi 90)

Light ADM via SM sphalerons challenging but possible

(M.T.F, Sarkar & Schmidt-Hoberg '11)

Transfer via induced L or B charges of DM

some higher dim. decay operator with model dependent T_{dec}

$$\mathcal{O}_{B-L} \mathcal{O}_X$$

$$\Delta W_{\text{eff}} = \frac{1}{M_i} \bar{X}^2 L_i H_u$$

(e.g. Farrar, Zaharijas '05; D.E Kaplan, Luty & Zurek '09)

Asymmetry co-generation:

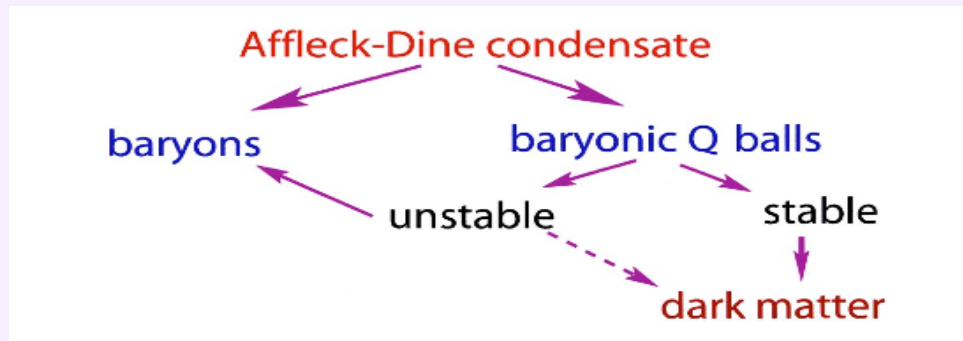
- B and DM asymmetries from same microphysics.
- Transfer operator decouples and either **B-L & B-X** or **B-L+X** is preserved

EWB co-generation via weak charges of DM (constituents)

(e.g Kaplan 92)

$U(1)_X \times SU(2)_L^2$ anomaly, sphalerons, 1st order phase transition, $T_{\text{sph}} \sim v_{\text{EW}}$

Or via new decaying states (e.g. Affleck-Dine co-generation)



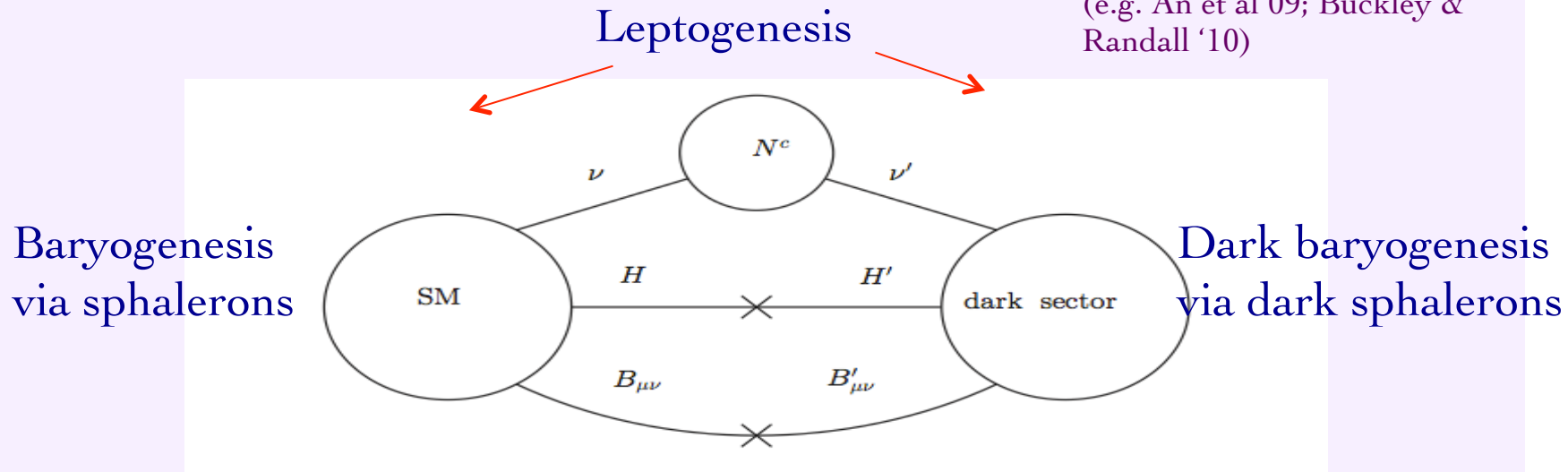
(e.g. Enqvist & McDonald '98;
Hall, March-Russell & West 10;
Hook 11; Cheung & Zurek 11;
March-Russell & McCullough
'11; Graesser, Shoemaker &
Vecchi '11)

ADM models with 'Hybrid' asymmetry generation

Asymmetry co-generation + transfer:

- L and L' (dark lepton number) co-generated at a high scale
- L transferred to B, and L' transferred to B' (dark baryon number B') by sphalerons and dark sphalerons

(e.g. An et al 09; Buckley & Randall '10)



- L, L' =DM charge X generated at a high scale
- L to B (X) by sphalerons

(e.g. Falkowski, Ruderman & Volansky '11)

ADM from Technicolor or New Strong Dynamics

Technicolour is a 'natural' framework for ADM:

(Nussinov 85, Barr, Chivukula and Farhi 90)

- Natural model of EWSB
- Composite neutral technibaryons charged under global $U(1)_{TB}$ part of unbroken chiral symmetries
- TB has weak charged constituents there is $U(1)_{TB} \times SU(2)_L^2$ anomaly
- Large annihilation cross-section so no symmetric cross-section (in general)

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
Λ_{QCD}	Nucleons	Baryon number $U(1)_B$	$\tau > 10^{33}$ yr (dim-6 OK)	'Freeze-out' from thermal equilibrium Asymmetric baryogenesis	$\Omega_B \sim 10^{-10}$ <i>cf.</i> observed $\Omega_B \sim 0.05$
$\Lambda_{\text{Fermi}} \sim G_F^{-1/2}$	Neutralino? Technibaryon?	R-parity? Technibaryon number $U(1)_{TB}$	violated? $\tau \sim 10^{18}$ yr e^+ excess?!	'Freeze-out' from thermal equilibrium Asymmetric (like the <i>observed</i> baryons)	$\Omega_{\text{LSP}} \sim 0.25$ $\Omega_{TB} \sim 0.25$

Technicolor (New Strong Dynamics)

- ① The SM gauge group is augmented:

$$G_{SM} \rightarrow SU(3)_c \times SU(2)_W \times U(1)_Y \times G_{TC} .$$

- ② The Higgs sector of the SM is replaced: $Q_L^a = \begin{pmatrix} U^a \\ D^a \end{pmatrix}_L$, $Q_R^a = (U_R^a, D_R^a)$,
 $a = 1, \dots, d(\mathcal{R}_{TC})$

$$\mathcal{L}_{Higgs} \rightarrow -\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} + i\bar{Q}_L \gamma_\mu D^\mu Q_L + i\bar{Q}_R \gamma_\mu D^\mu Q_R + \dots$$

Lightest Technibaryon stable due to unbroken global $U(1)_{TB}$ symmetry

Minimal chiral symmetries: 3 GB's + Custodial + DM.

$$SU_L(2) \times SU_R(2) \times U_{TB}(1) \rightarrow SU_V(2) \times U_{TB}(1) .$$

However, $U(1)_{TB}$ is anomalous due to the EW anomaly, when the technifermions are charged under $SU(2)_{EW}$

$$\partial_\mu J_{TB}^\mu = \frac{1}{2\sqrt{2}} \frac{g^2}{32\pi^2} \epsilon_{\mu\nu\rho\sigma} W^{\mu\nu} W^{\rho\sigma} , \quad \text{and} \quad J_{TB}^\mu = \frac{1}{2\sqrt{2}} (\bar{U} \gamma^\mu U + \bar{D} \gamma^\mu D)$$

Technibaryon ADM

(Nussinov 85)

The ADM relic density is simply related to the baryon relic density

$$\Omega_\chi = (m_\chi \mathcal{N}_\chi / m_B \mathcal{N}_B) \Omega_B$$

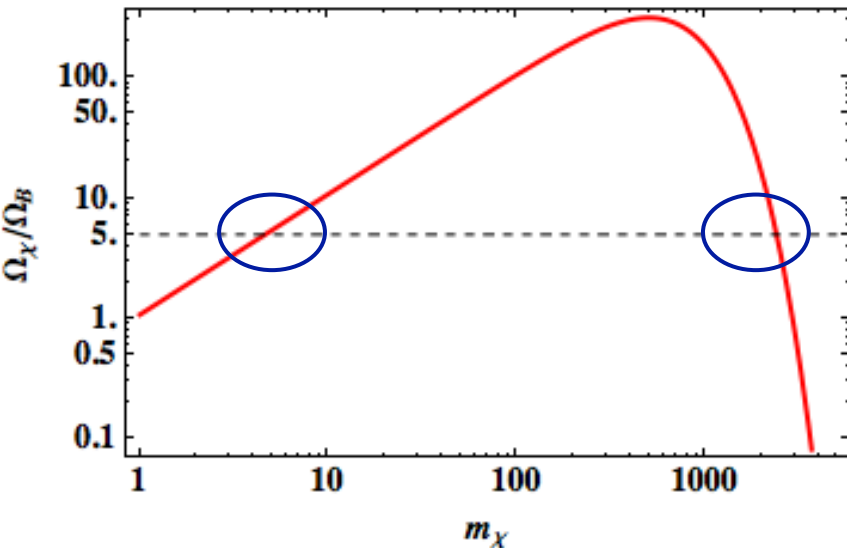
Nussinov thought of TeV scale technibaryon when relic density was less well known so $\mathcal{N}_\chi \sim \mathcal{N}_B$ $m_\chi \sim 1$ TeV was close enough...

Now similar asymmetries (such as from GUTs or from sphalerons)

$$\mathcal{N}_\chi \sim \mathcal{N}_B \Rightarrow m_\chi \sim 5 \text{ GeV}$$

In fact, there are 'two' solutions via Sphalerons (e.g. starting from Leptogenesis):

(Barr, Chivukula and Farhi 90)



$$\frac{\Omega_\chi}{\Omega_B} \sim \frac{m_\chi}{m_B} \times \begin{cases} 1 & \text{for } m_\chi \sim \text{GeV} \\ \left(\frac{m_\chi}{T_{\text{Sph}}}\right)^{3/2} e^{-m_\chi/T_{\text{Sph}}} & \text{for } m_\chi \sim \text{TeV} \end{cases}$$

Really Ω_χ/Ω_B depends on L/B (and possible new L') so more general masses are possible.

(Barr, Chivukula and Farhi 90; Gudnasone, Kouvaris & Sannino '05)

Note BCF did not include charge neutrality of universe but this does not change qualitative picture

(Harvey & Turner 90)

pNGB Technibaryons in Minimal Walking Technicolor

(Sannino and collaborators '05;...)

Example of more general solutions termed:
Techni-Interacting Massive Particles (TIMPs)

'iTIMP'

- \mathcal{R} real
- $T^0 \sim U_L D_L$
- Iso-singlet GB
- $M_{T^0} \sim g F_\Pi$

(Gudnason, Kouvaris & Sannino 05; M.T.F & Sannino 09)

'TIMP'

- 4 of $SU(4)$
- $U_L D_L U_L D_L$
- SM singlet
- $M_T \sim N_{TC}^{3/2} F_\Pi$

(Barr, chivukula & Farhi 90)

'TIMP'

- \mathcal{R} pseudo-real
- $T^0 \sim U_L D_L$
- SM singlet GB
- $M_{T^0}^2 \sim -g^2 F_\Pi^2$

(Ryttov and Sannino 09)

General analysis of signals and the annihilation cross-section

(Foadi, M.T.F & Sannino 08; Belyaev, M.T.F, Sarkar & Sannino 10, Del Nobile, Kouvars & Sannino '11)

'Minimal Dark Matter' in Technicolor theory space

(Cirelli, Fornengo Strumia '02)

Techni-Interacting Massive Particles

Minimal Technicolor Theory Space

Minimal Technicolor: 2 Dirac Flavors. No QCD charges.

$$Q_L = \left(U_L^{+1/2}, D_L^{-1/2} \right)^T, \quad U_R^{+1/2}, D_R^{-1/2}.$$

'Orthogonal TC'

- \mathcal{R} real
- F of $SO(N)$
- $SU(4)/SO(4)$
- $3_\pi \oplus 3 \oplus \bar{3}$

$$\begin{pmatrix} \pi & T_i \\ T_i^* & \pi^T \end{pmatrix}$$

'QCD TC'

- \mathcal{R} complex
- F of $SU(N)$
- G_{GB} : $SU(2)$
- 3_π

$$\pi = \begin{pmatrix} \pi^0 & \pi^+ \\ \pi^- & \pi^0 \end{pmatrix}$$

'Symplectic TC'

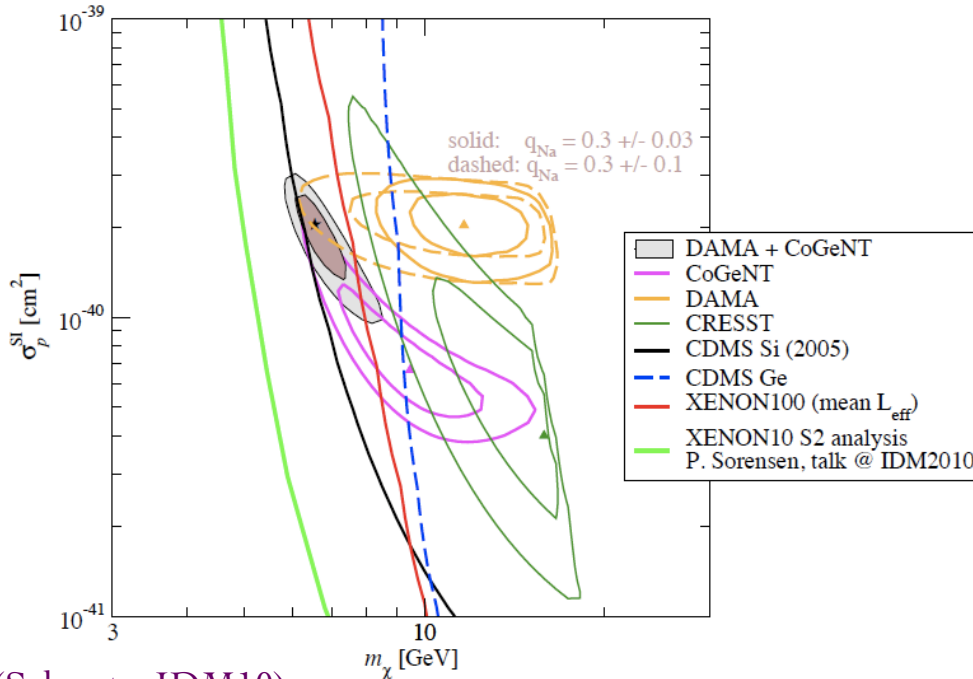
- \mathcal{R} pseudo-real
- F of $Sp(2N)$
- $SU(4)/Sp(4)$
- $3_\pi \oplus 1 \oplus \bar{1}$

$$\begin{pmatrix} \pi & T_s \\ T_s^* & \pi^T \end{pmatrix}$$

(M.T.F & Sannino 09)

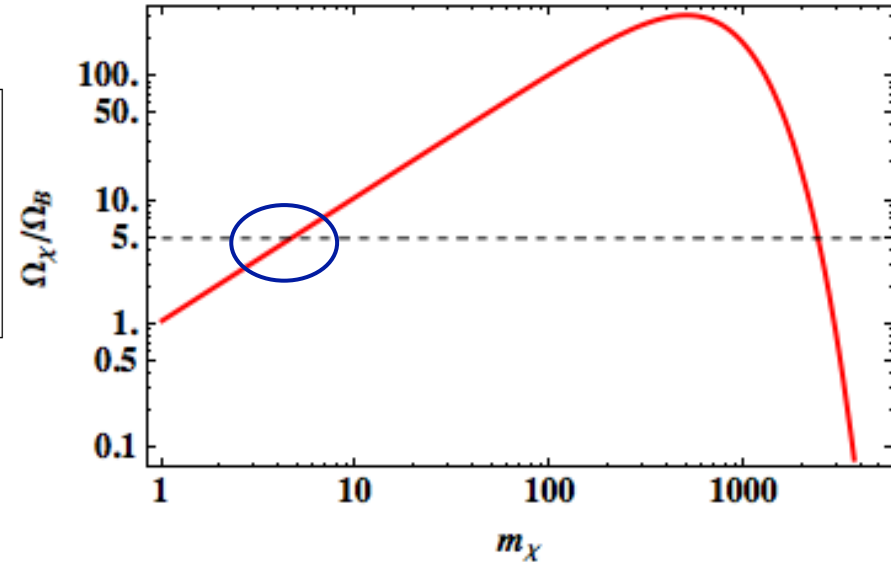
(Ryttov and Sannino 09)

GeV scale ADM



(Schwetz, IDM10)

$$\Omega_\chi = (m_\chi \mathcal{N}_\chi / m_B \mathcal{N}_B) \Omega_B$$



(Talks this workshop by Collar, Schmidt-Hoberg, Zupan)

Recent surge of interest in GeV scale ADM after hints of light DM in direct detection experiments

Can you get $\mathcal{N}_\chi \sim \mathcal{N}_B$ and light ADM via sphalerons in Technicolor (in agreement with constraints from e.g. W/Z measurements?)

GeV scale ADM and EWSB from two sectors

Assume 2 sectors $\mathcal{S}_1 \sim \text{TeV}$, $\mathcal{S}_2 \sim \text{GeV}$ (M.T.F, Sarkar and Schmidt-Hoberg 11;
M.T.F, Kahlhoefer, Sarkar & Schmidt-Hoberg 11 and work in progress)

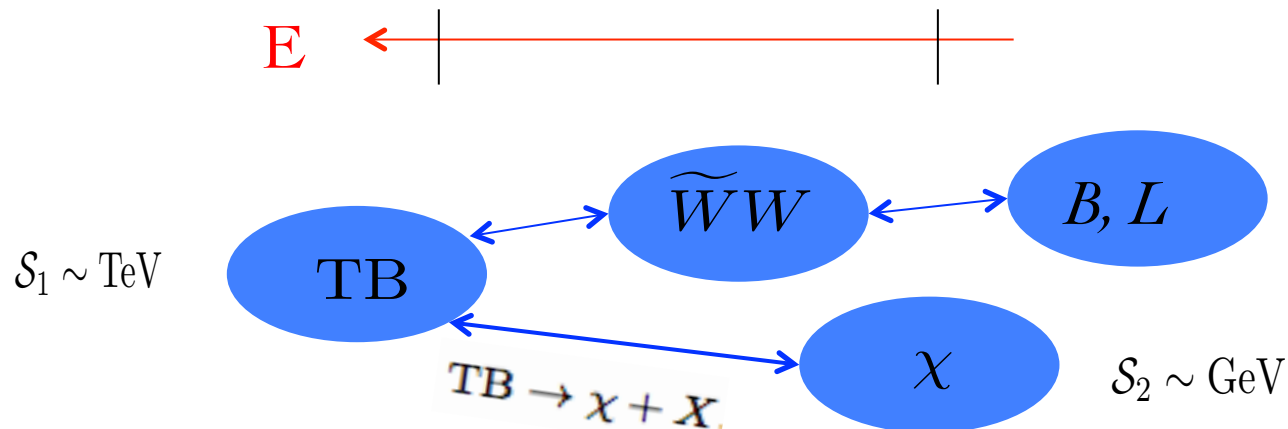
i) \mathcal{S}_1 breaks EW symmetry dynamically at a scale Λ_1

composite spectrum includes $\mathcal{O}(\text{TeV})$ mass particles TB carrying $U(1)_{\text{TB}}$ charge

\mathcal{S}_2 is SM singlet sector becoming strongly interacting at Λ_2 of $\mathcal{O}(\text{GeV})$

composite spectrum includes $\mathcal{O}(\text{GeV})$ mass particles χ carrying $U(1)_{\text{TB}}$ charge

$\mathcal{S}_{1,2}$ coupled via $U(1)_{\text{TB}}$ preserving decay operator $\text{TB} \rightarrow \chi + X$.



\mathcal{S}_1 States (constituents) carry weak charges and are connected to sphalerons

$\text{TB} \rightarrow \chi + X$ Is in equilibrium until $T \lesssim T_{\text{sph}}$

GeV scale ADM and EWSB from two scale strong dynamics

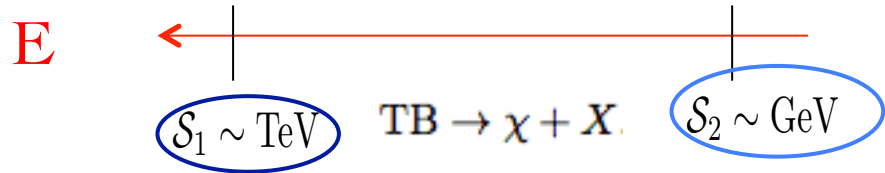
Assume new strong dynamics with weak charged fermions Q and SM Singlet fermions λ ... (Dietrich, Sannino & Tuominen 05; Sannino & Rytov 08; Galloway, Evans, Luty & Tacchi '10)

motivated also by constraints from EW precision measurements

We identify composite states $TB \sim Q \cdots Q$ and $\chi \sim \lambda \cdots \lambda$

$$Q_L^a = \begin{pmatrix} U^a \\ D^a \end{pmatrix}_L, U_R^a, D_R^a, \lambda^b$$

$a = 1, \dots, d(\mathcal{R}_1), b = 1, \dots, d(\mathcal{R}_2).$



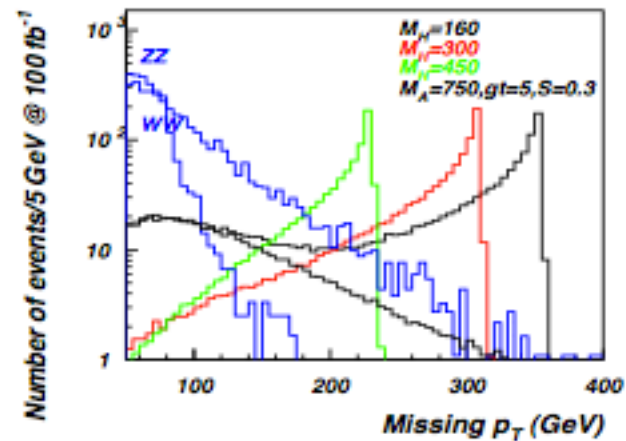
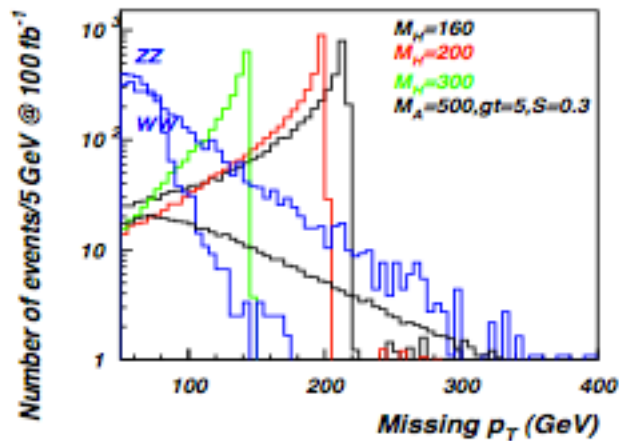
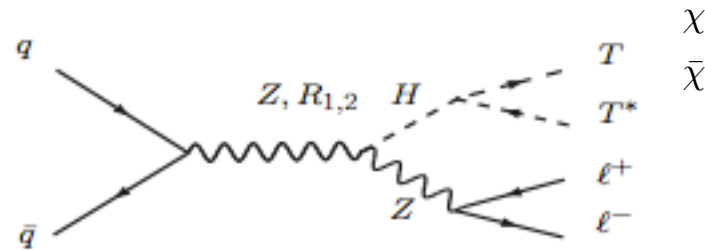
Scale separation can arise due do different Casimirs $\alpha_c = \frac{\pi}{3C_2(\mathcal{R})}$ (Raby, Dimopoulos & Susskind 80)

Can also arise/be enhanced in presence of 4-fermion interactions (M.T.F, Sarkar and Schmidt-Hoberg 11)

Recent 1-loop exact susy comp of a 2-scale model (Antipin, Mojaza & Sannino '11)

Model provides natural EWSB, DM stability and (putative) light ADM via weak sphalerons – decay operator can but need not arise in TC sector itself
 Achieving scale separation is obviously a challenge

(i) TIMP missing energy signals



(Foadi, M.T.F and Sannino 08; Godbole, Guchait, Mazumdar, Moretti and Roy 03).

Summary

Asymmetric dark matter is motivated by the observed asymmetry of baryonic matter and the desire to explain why

$$\Omega_{\text{DM}}/\Omega_{\text{B}} \approx 5$$

- ADM natural in the framework of New Strong Dynamics such as Technicolour

Recently vigorous ADM model building activity in a large variety of frameworks

- Interesting alternative to WIMP paradigm...Experiment will tell

Appendices

Two scales from strong dynamics

(Marciano 80; Lane & Eichten 89)

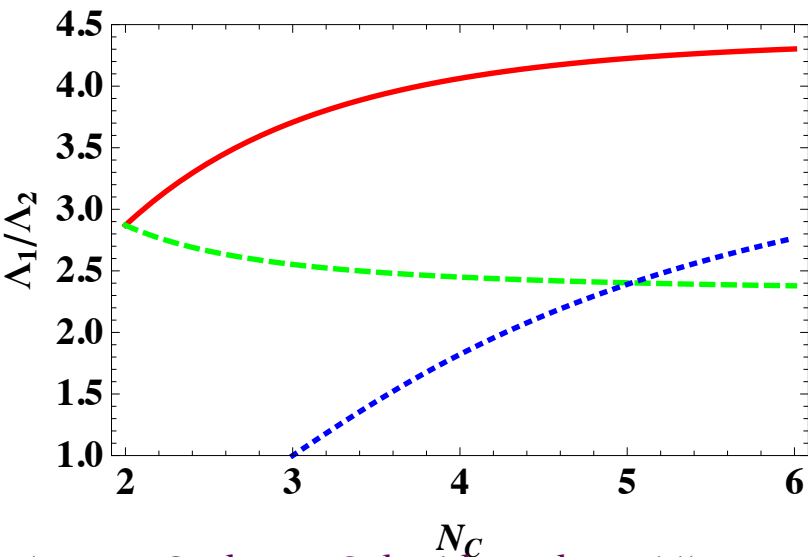
Critical coupling for ChiSB depends on the fermion rep...

$$\alpha_c = \frac{\pi}{3C_2(\mathcal{R})}, \quad \alpha_* = -4\pi \frac{\beta_0}{\beta_1}$$

(Raby, Dimopoulos & Susskind 80)

One-loop estimate of scale separation with 2 representations

$$\frac{\Lambda_1}{\Lambda_2} \simeq \exp \left[\frac{2\pi}{\beta_0(\mathcal{R}_2)} \left(\alpha_c(\mathcal{R}_1)^{-1} - \alpha_c(\mathcal{R}_2)^{-1} \right) \right]$$



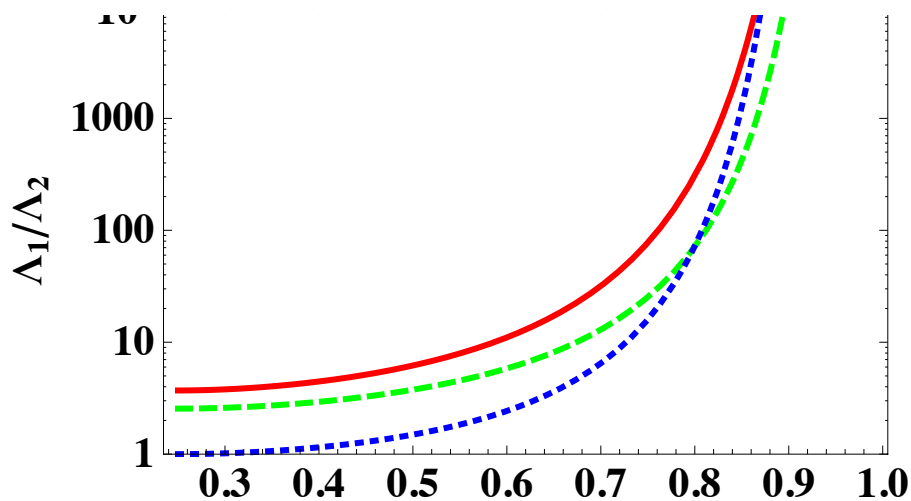
(M.T.F, Sarkar & Schmidt-Hoberg 11)

and in the NJL model on the 4-fermion coupling

(Appelquist et al 88, Kondo et al 89, Kurachi et al 07, Fukano & Sannino 09)

$$\mathcal{L} = \bar{Q}i\not{D}Q - \frac{1}{4}\text{Tr}F_{\mu\nu}^a F^{a\mu\nu} + \frac{4\pi^2 g_1}{\Lambda^2 N_1 d(\mathcal{R}_1)} [(\bar{Q}Q)^2 + (\bar{Q}i\gamma_5 T^a Q)^2]$$

$$\alpha_c(\mathcal{R}_1, g_1) = \begin{cases} 4(\sqrt{g_1} - g_1) \times \alpha_c(\mathcal{R}_1) & \text{for } \frac{1}{4} < g_1 \leq 1, \\ \alpha_c(\mathcal{R}_1) & \text{for } 0 \leq g_1 < \frac{1}{4}. \end{cases}$$



Recent 1-loop exact susy comp of a 2-scale model

(Antipin, Mojaza & Sannino 11)