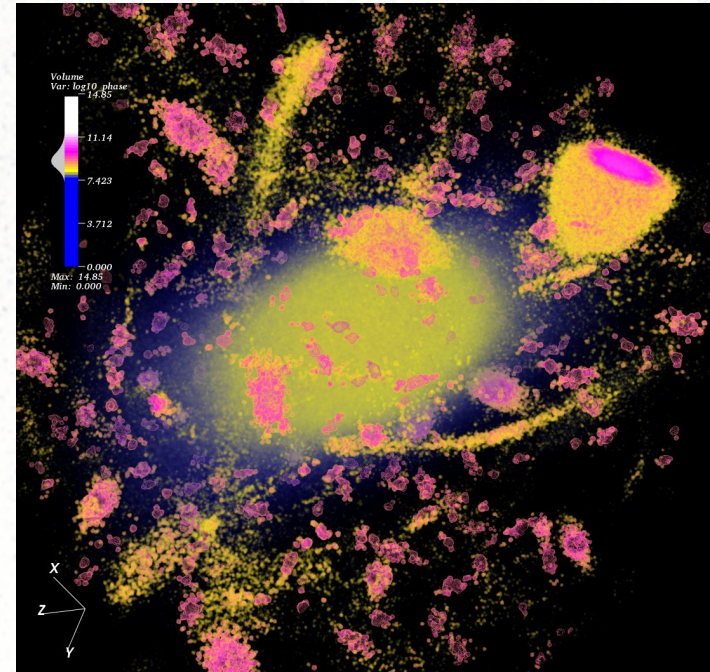
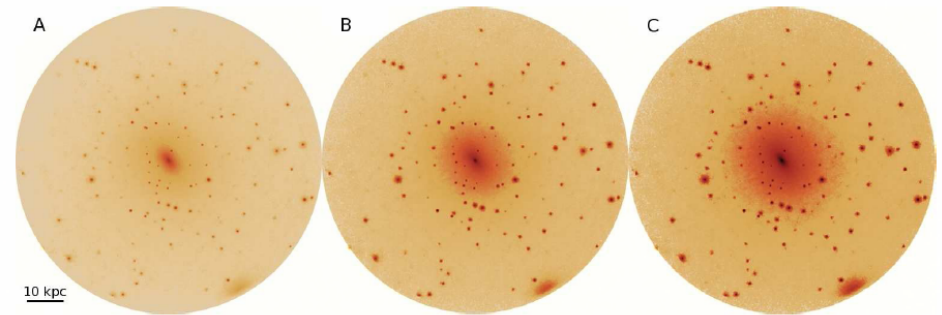


Numerical Simulations of Galactic DM Substructure And Implications for Direct and Indirect Detection

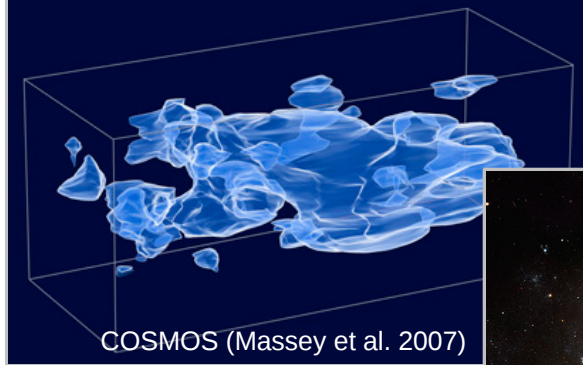
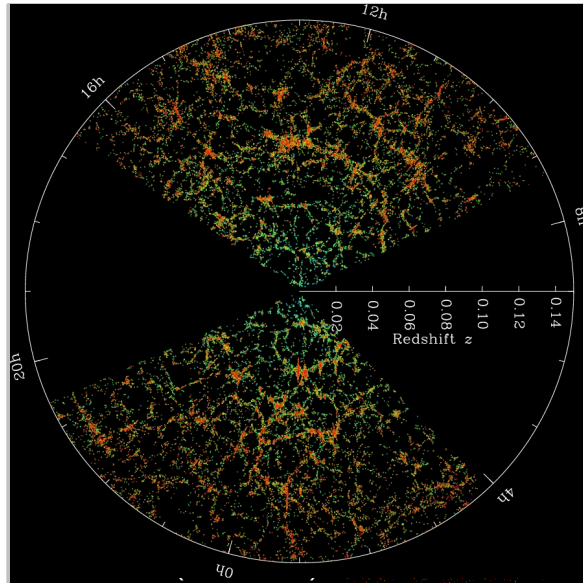
Michael Kuhlen, UC Berkeley



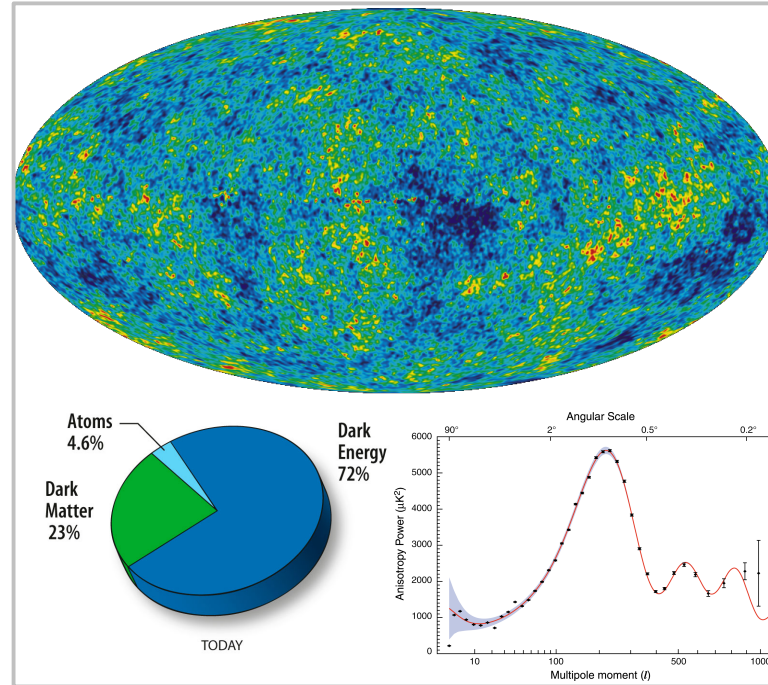
Collaborators: P. Madau (UCSC), J. Diemand (Zurich), M. Zemp (Michigan), B. Moore (Zurich), J. Stadel (Zurich), D. Potter (Zurich), N. Weiner (NYU), B. Anderson (SCIPP), R. Johnson (SCIPP), M. Kamionkowski (Caltech), S. Koushiappas (Brown)

There's evidence for dark matter on many scales...

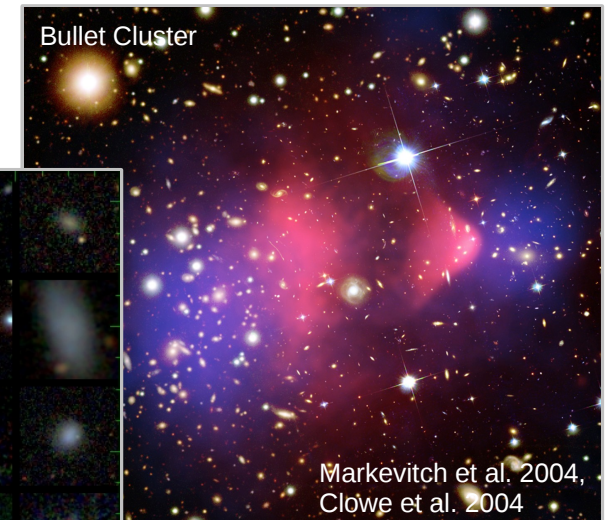
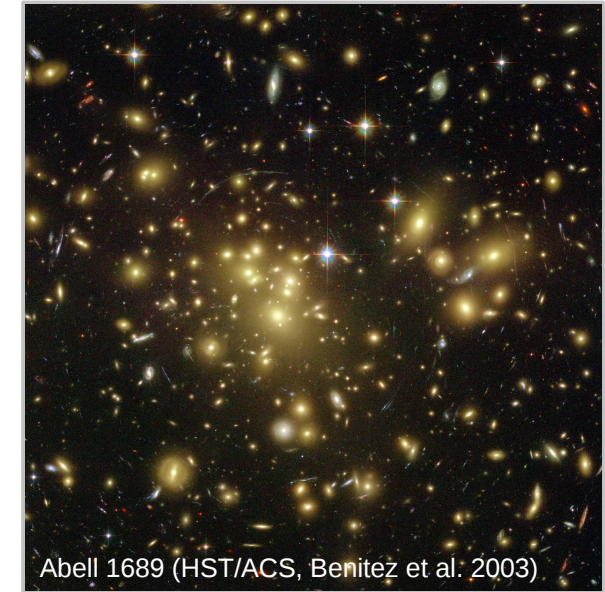
Large Scale Structure



Cosmic Microwave Background



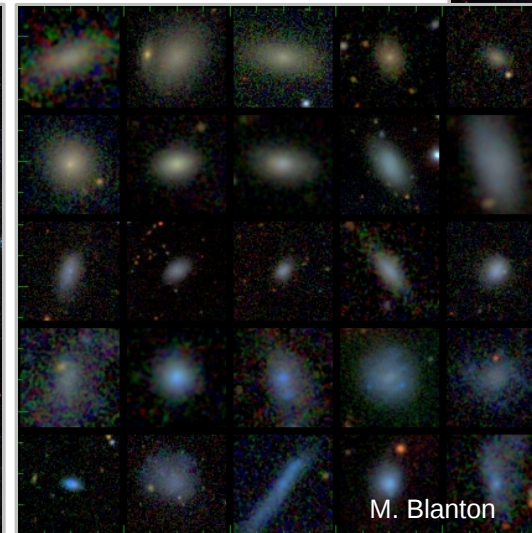
Galaxy Clusters



Galaxies



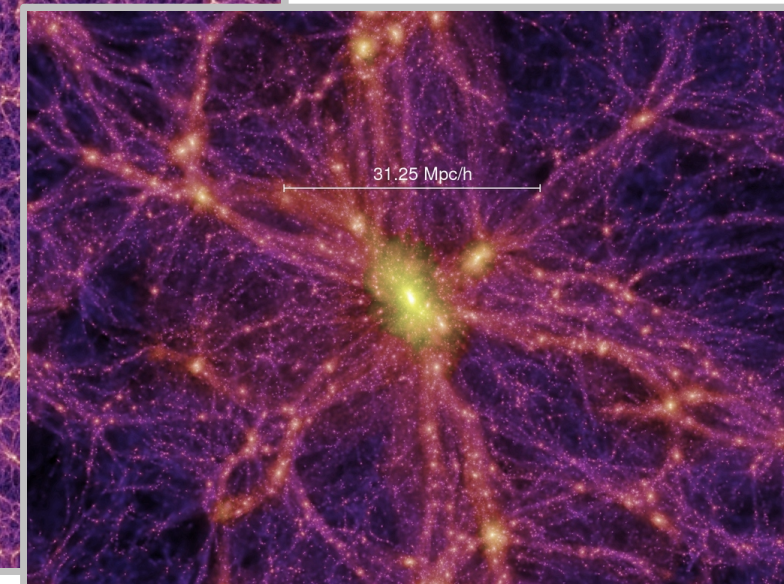
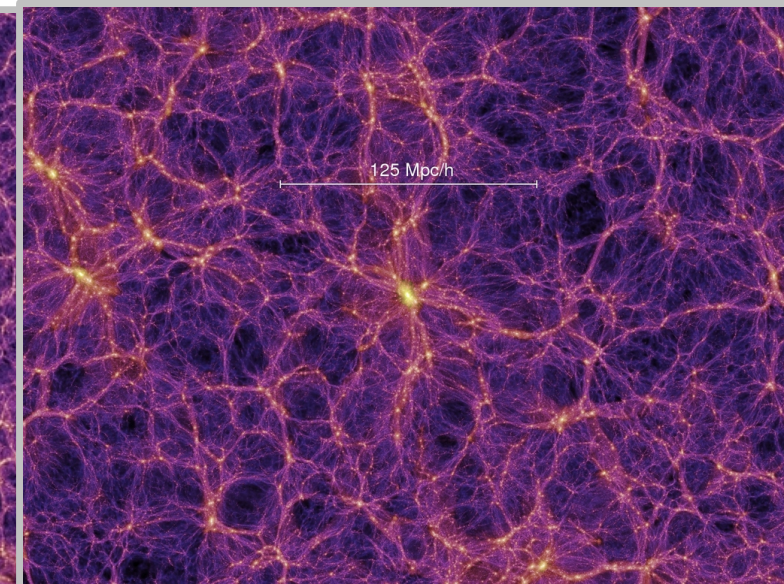
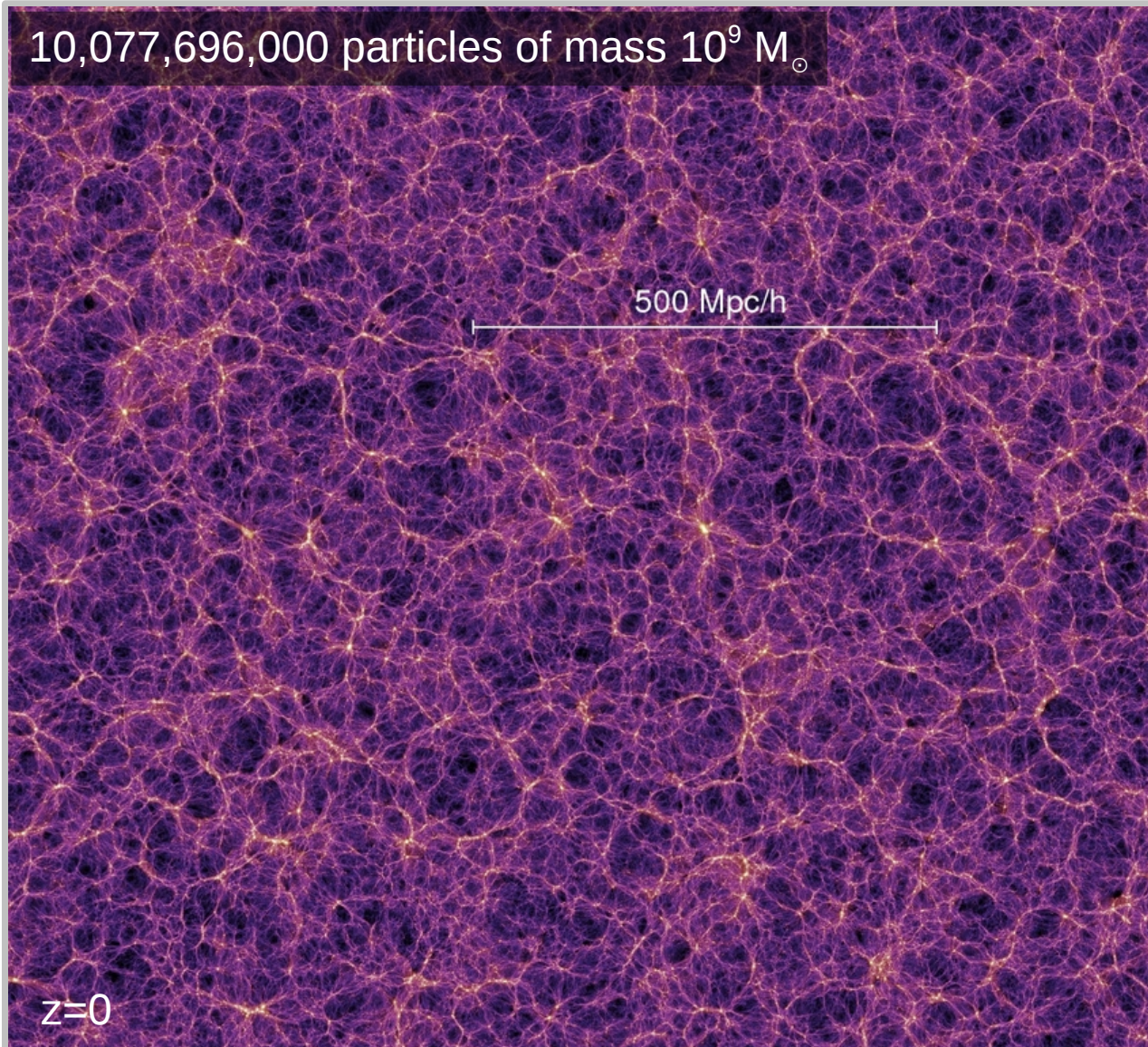
Dwarf Galaxies



Numerical Simulations: Universe Scale

The Millenium Run (Springel et al. 2005)

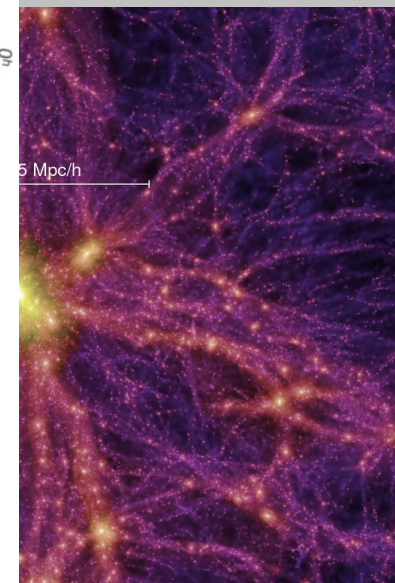
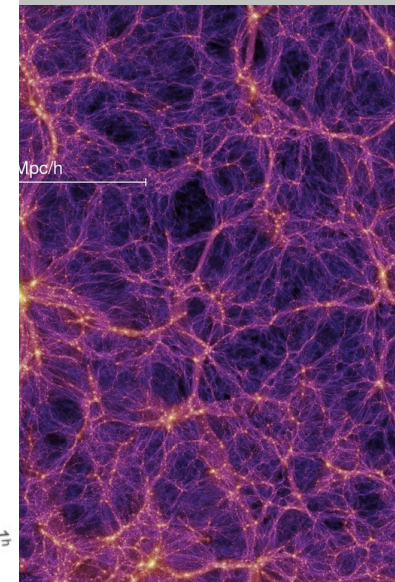
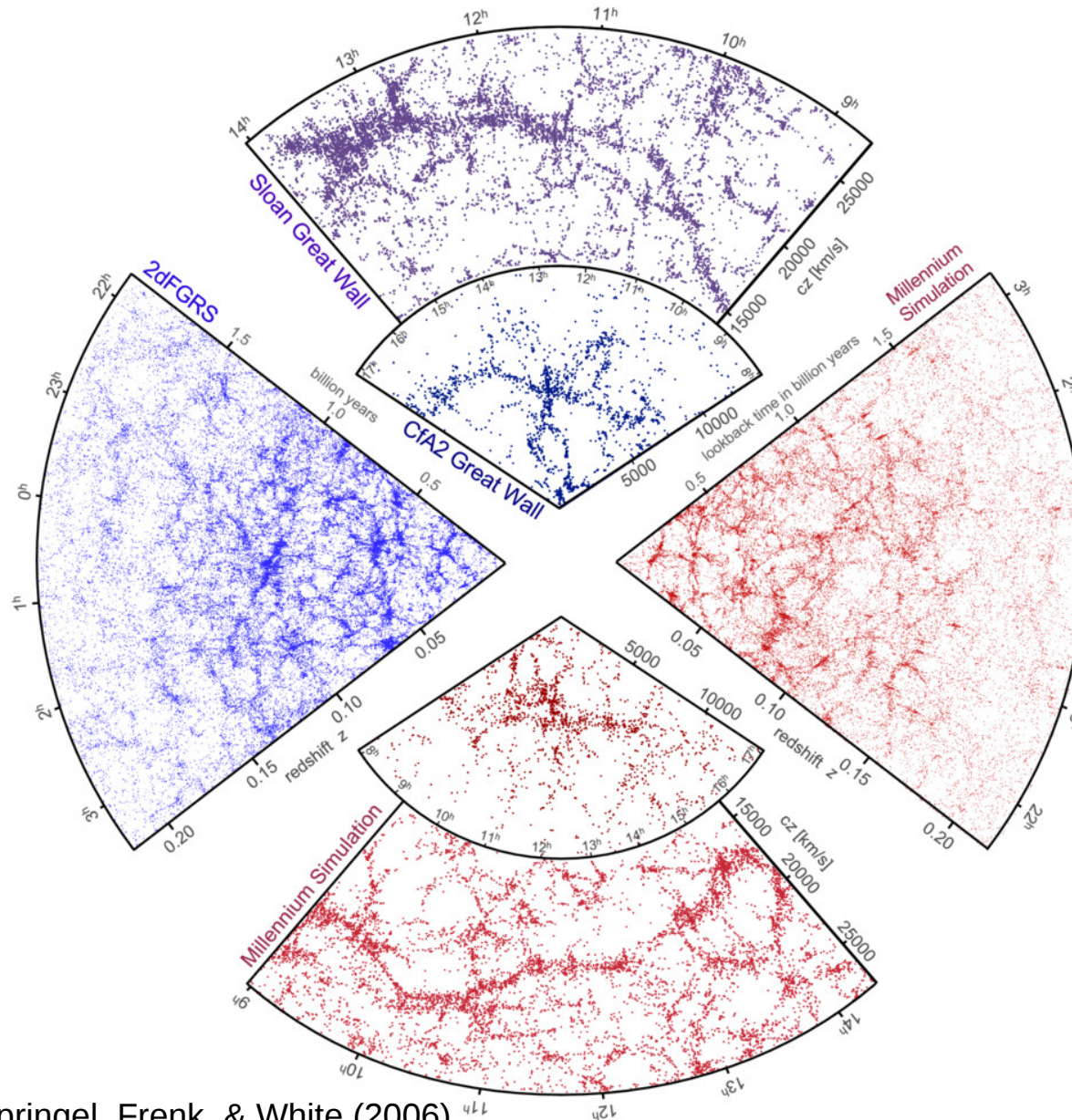
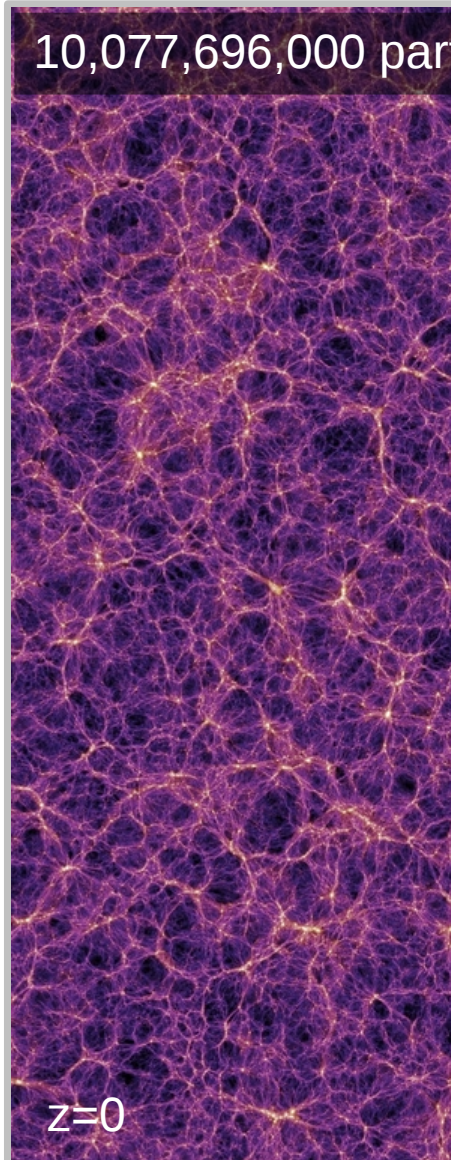
Max-Planck-Institut
für Astrophysik



Numerical Simulations: Universe Scale

The Millenium Run (Springel et al. 2005)

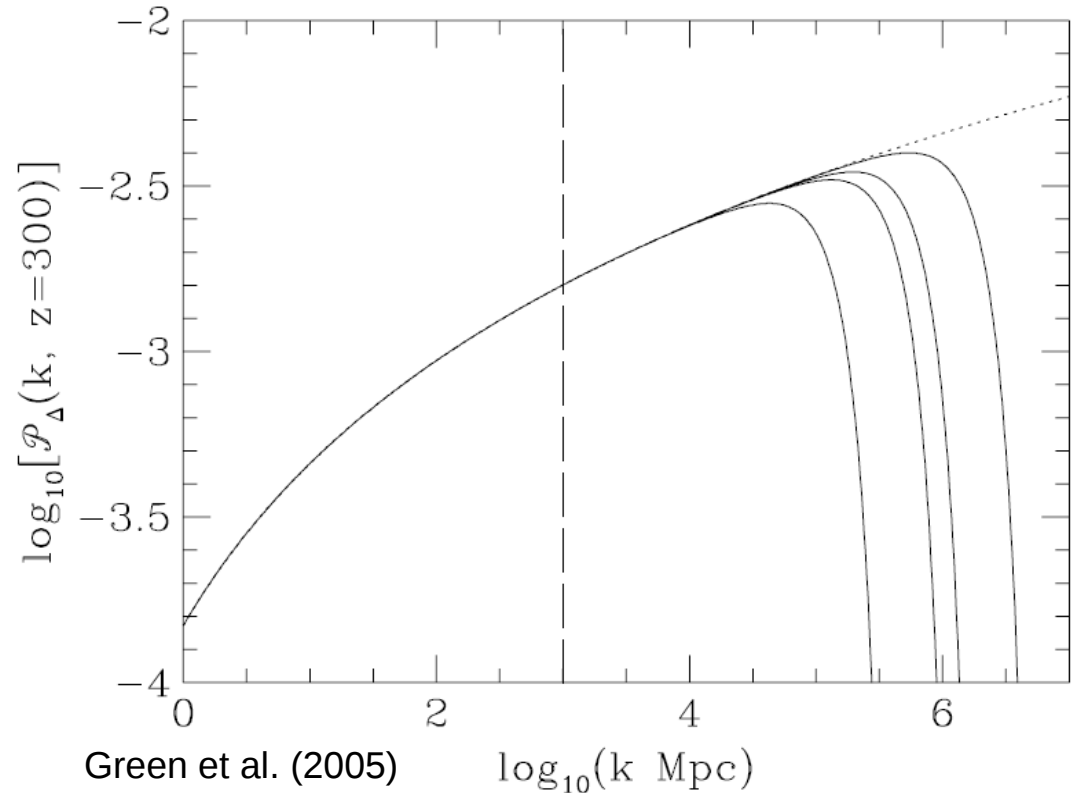
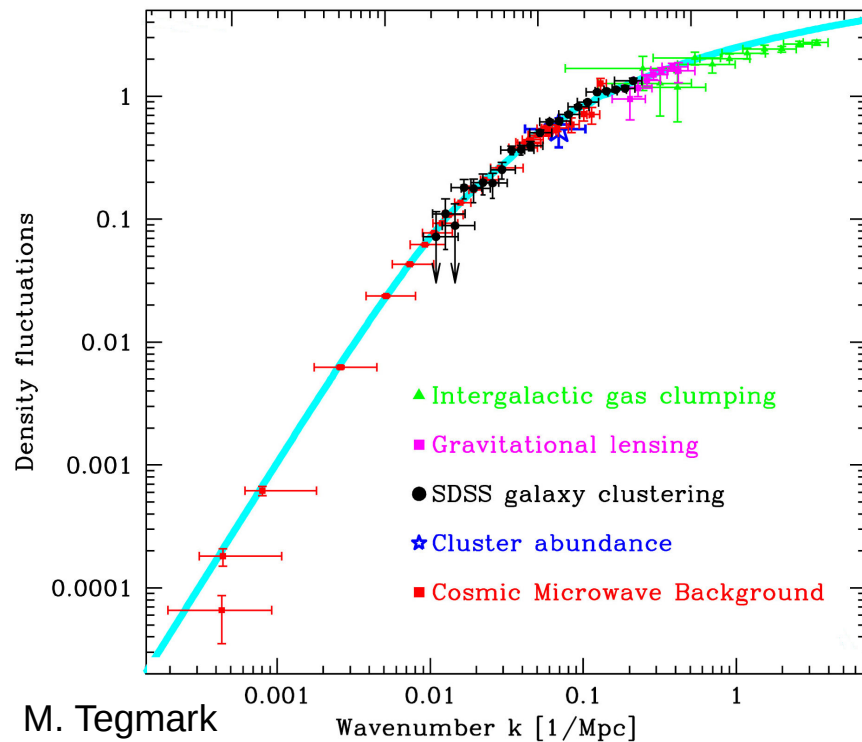
Max-Planck-Institut
für Astrophysik



Springel, Frenk, & White (2006)

What mass are the smallest DM structures?

A cutoff in the dark matter density fluctuation power spectrum is set by collisional damping between chemical decoupling (abundance freeze-out) and kinetic decoupling (end of elastic scattering), and by free-streaming afterwards.



The density fluctuations on (sub-)galactic scales grow to be **strongly non-linear**. Following their evolution and dynamics accurately requires **extremely high resolution numerical simulations**.

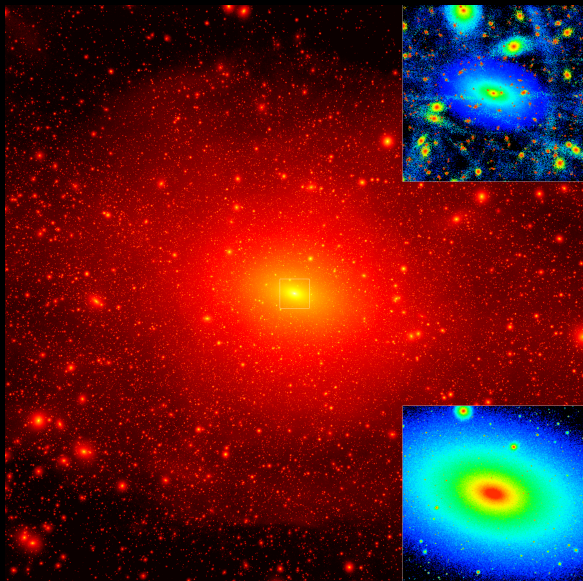
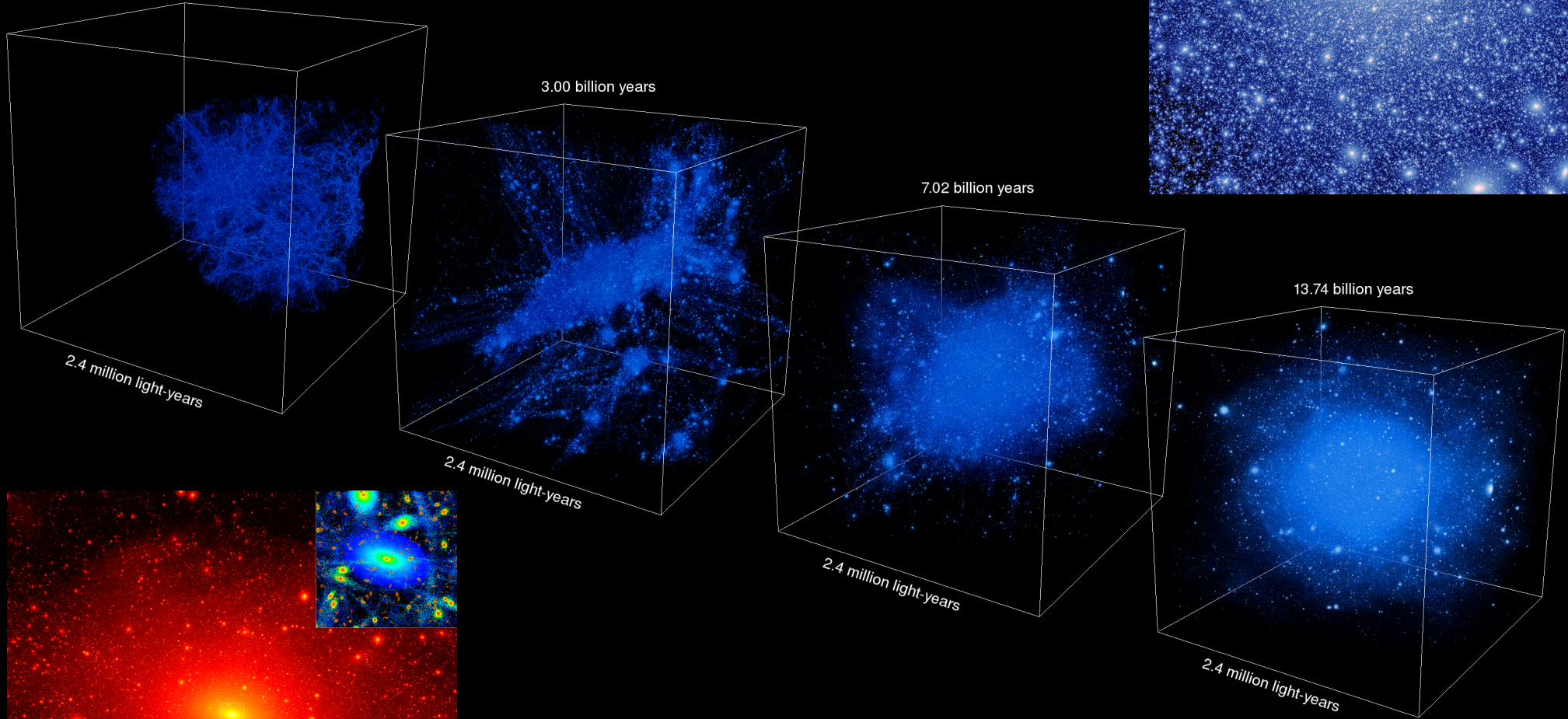
The Via Lactea Project

J. Diemand – M. Kuhlen – P. Madau
(& B. Moore, D. Potter, J. Stadel, M. Zemp)

GHALO
Stadel et al. (2009)
2.1 billion particles, 1,000 M_{\odot}



Time since Big Bang: 0.50 billion years



VIA LACTEA II
Diemand, Kuhlen et al. 2008
1.1 billion particles, 4,000 M_{\odot}



The Via Lactea Project

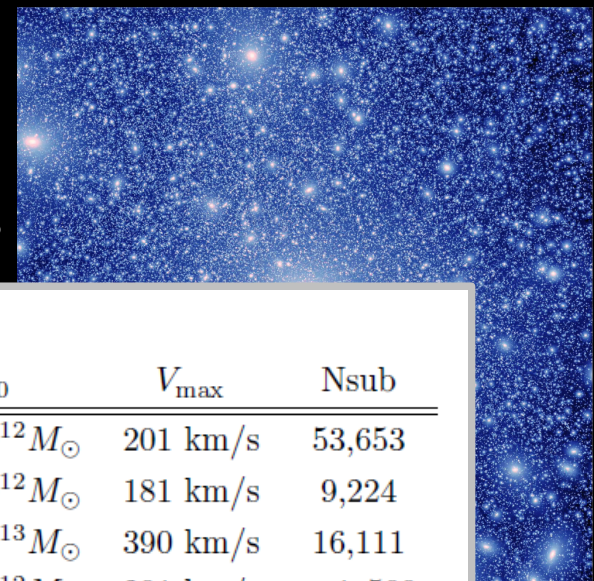
J. Diemand – M. Kuhlen – P. Madau

(& B. Moore, D. Potter, J. Stadel, M. Zemp)

GHALO

Stadel et al. (2009)

2.1 billion particles, 1,000 M_{\odot}



Via Lactea Series

Name	N_{hires}	M_{hires}	z_i	z_f	cosmology	r_{200}	M_{200}	V_{max}	Nsub
VL-II	1.09×10^9	$4.1 \times 10^3 M_{\odot}$	104.3	0	WMAP-3	402 kpc	$1.93 \times 10^{12} M_{\odot}$	201 km/s	53,653
VL-I	2.34×10^8	$2.1 \times 10^4 M_{\odot}$	48.8	0	WMAP-3	389 kpc	$1.77 \times 10^{12} M_{\odot}$	181 km/s	9,224
1e8Ell	1.36×10^8	$1.8 \times 10^5 M_{\odot}$	59.5	0.49	WMAP-1	512 kpc	$1.34 \times 10^{13} M_{\odot}$	390 km/s	16,111
VL-IIIm	2×10^7	$2.6 \times 10^5 M_{\odot}$	104.3	0	WMAP-3	402 kpc	$1.93 \times 10^{12} M_{\odot}$	201 km/s	$\sim 1,500$

GHALO Series

Name	N_{hires}	M_{hires}	z_i	z_f	cosmology	r_{200}	M_{200}	V_{max}	Nsub
GHALO ₂	2.1×10^9	$1,000 M_{\odot}$	48.8	0	WMAP-3	347 kpc	$1.3 \times 10^{12} M_{\odot}$	153 km/s	$\sim 10^5$
GHALO ₃	$\sim 10^8$	$2.7 \times 10^4 M_{\odot}$	48.8	0	WMAP-3	347 kpc	$1.3 \times 10^{12} M_{\odot}$	150 km/s	$\sim 2,500$
GHALO ₄	$\sim 5 \times 10^6$	$7.3 \times 10^5 M_{\odot}$	48.8	0	WMAP-3	347 kpc	$1.3 \times 10^{12} M_{\odot}$	150 km/s	–

Silver River

Name	N_{hires}	M_{hires}	z_i	z_f	cosmology	r_{200}	M_{200}	V_{max}	Nsub
SR1	50×10^9	$100 M_{\odot}$	180	currently at $z = 6.3$	WMAP-7	-	-	-	-

Time since Big

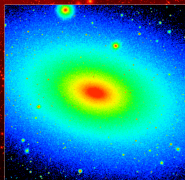
2.4 mil

2.4 million light-years

VIA LACTEA II

Diemand, Kuhlen et al. 2008

1.1 billion particles, 4,000 M_{\odot}



How to transfer 140 TB of data across the country

Sector – High Performance Distributed File System and Parallel Data Processing Engine

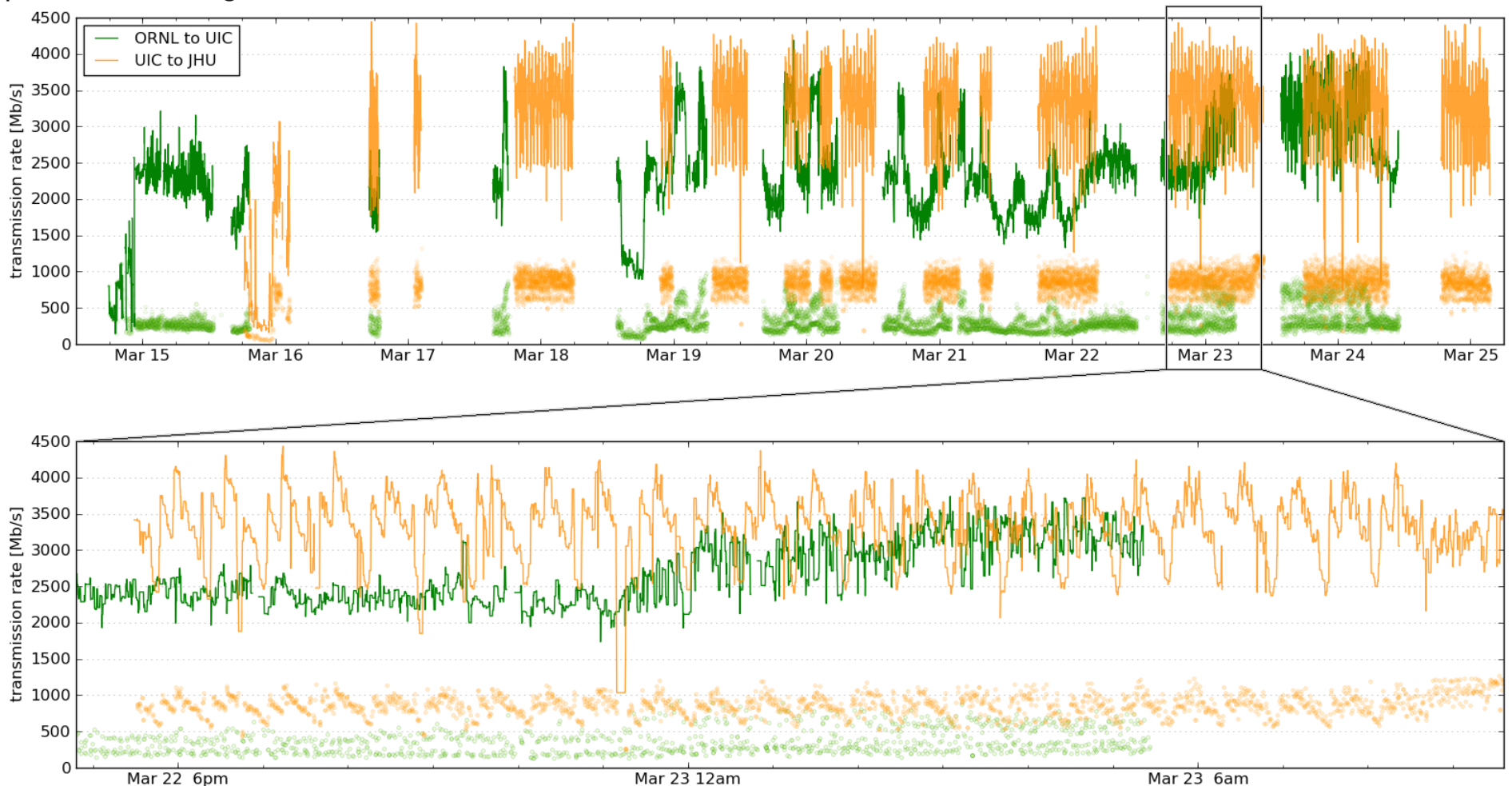
<http://sector.sourceforge.net/>

Yunhong Gu @ UIC

UDT – UDP-based Data Transfer

<http://udt.sourceforge.net/>

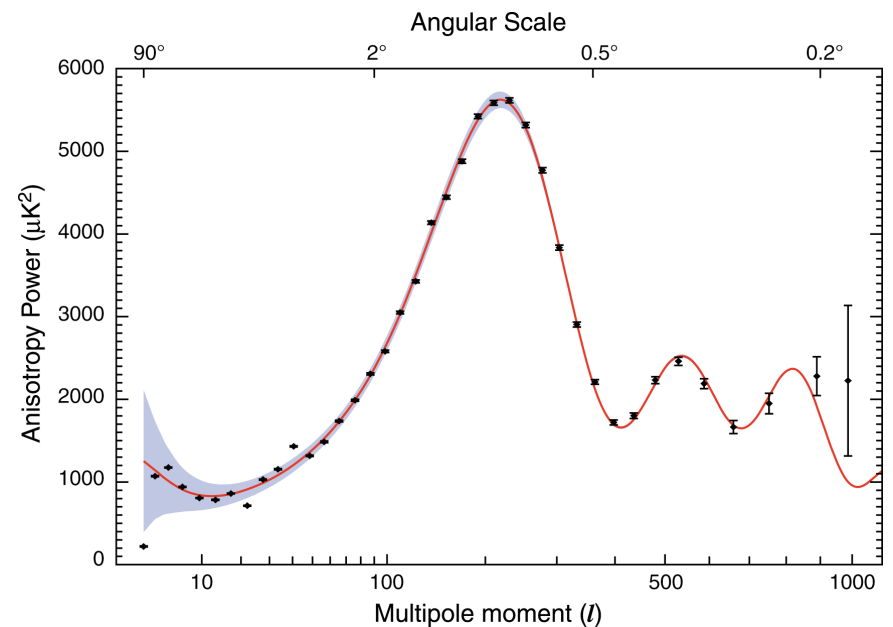
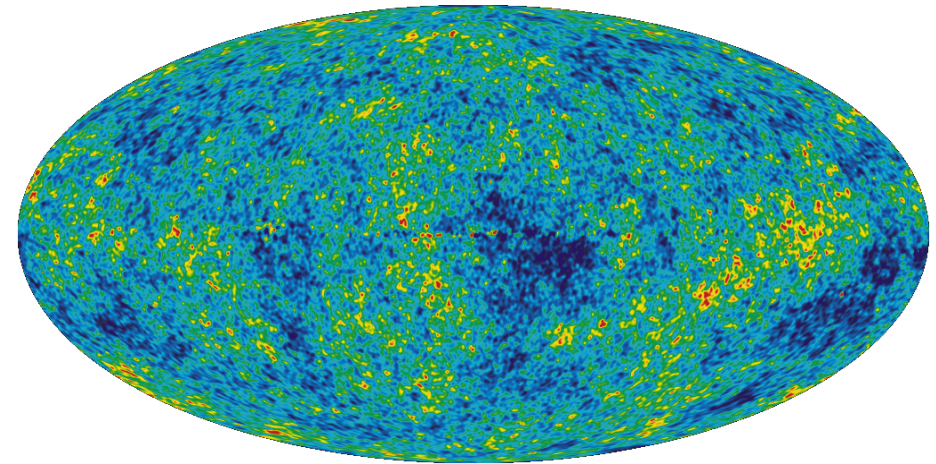
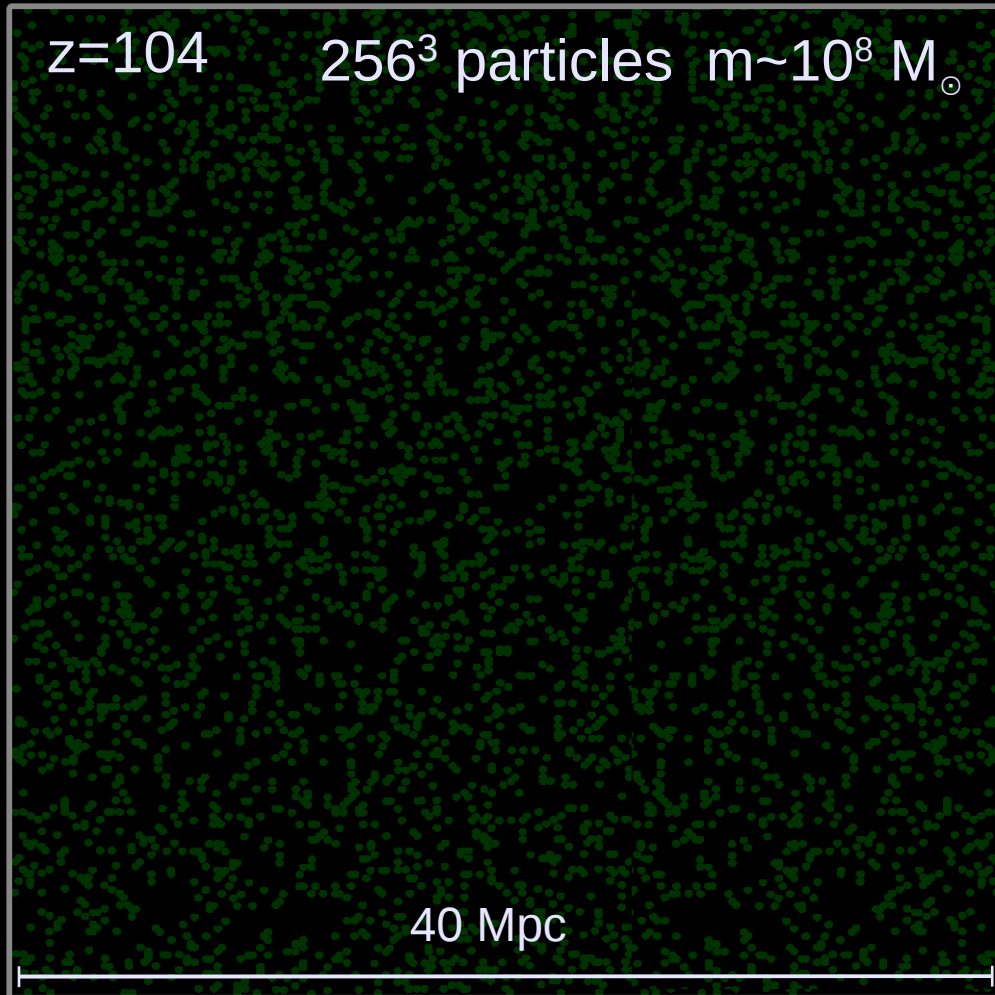
Used 10 Gbps links from ORNL to UIC and from UIC to JHU (storage server): 2500 – 3500 Mbps



Multi-mass initial conditions

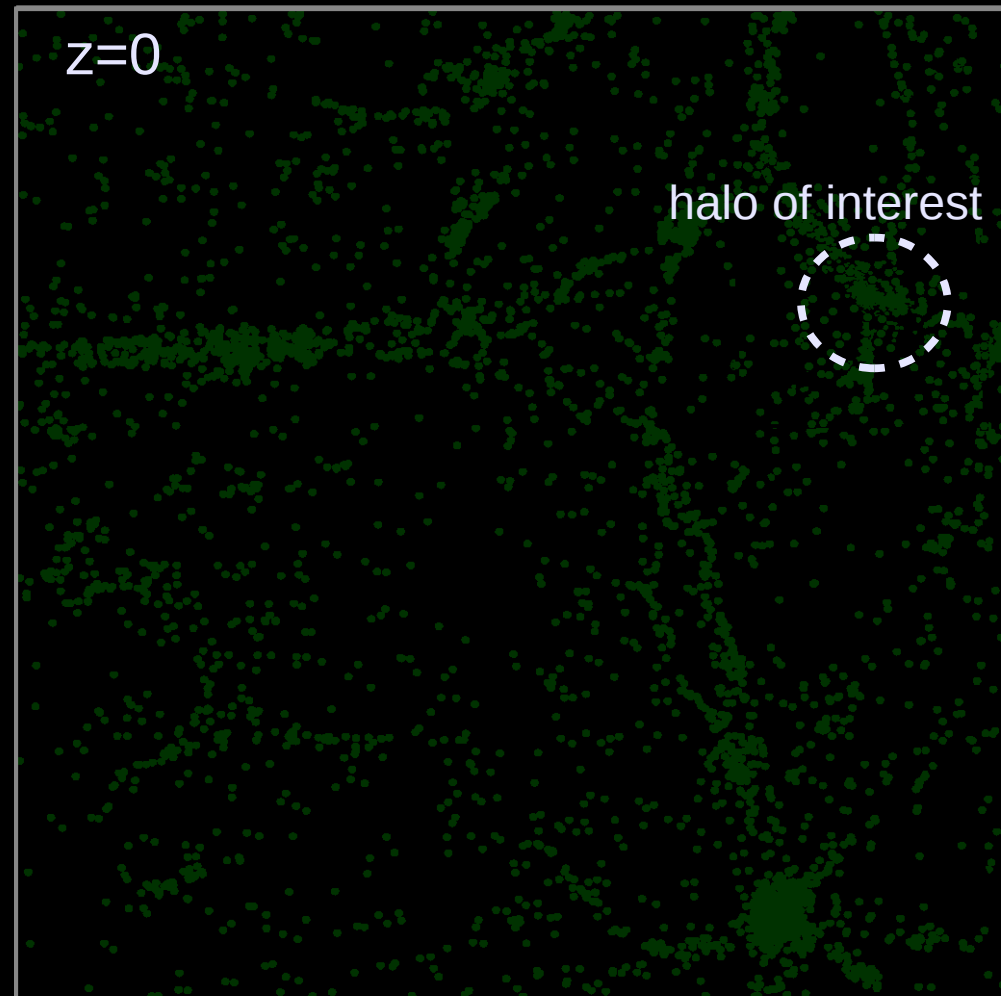
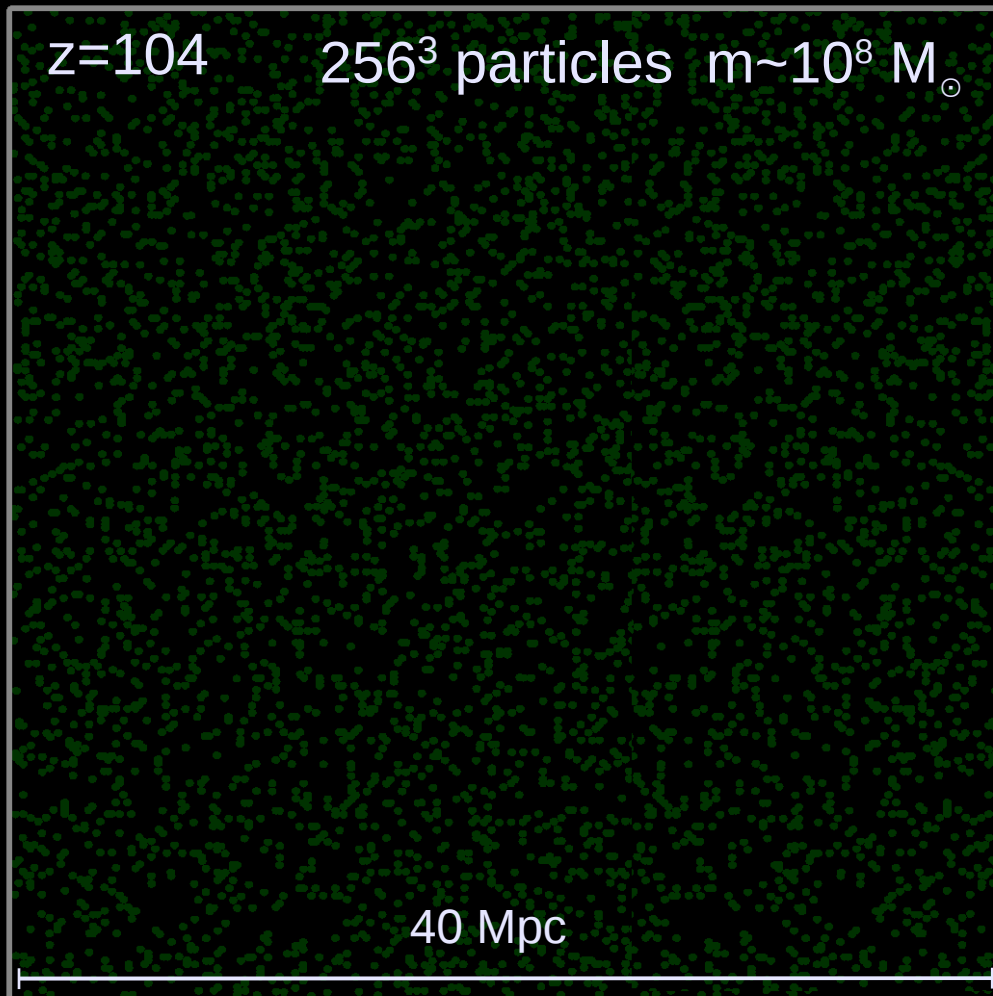
Initial conditions at redshift $z=104$, ~15 million years after the Big Bang.

A low resolution simulation is initialized with a Gaussian random field and a Cold Dark Matter power spectrum, normalized by CMB.



Multi-mass initial conditions

This low resolution simulation is quickly run to $z=0$ (13.7 billion years after the Big Bang), and a halo of interest (e.g. mass, V_{\max} equal to Milky Way, no recent major mergers) is selected.



Multi-mass initial conditions

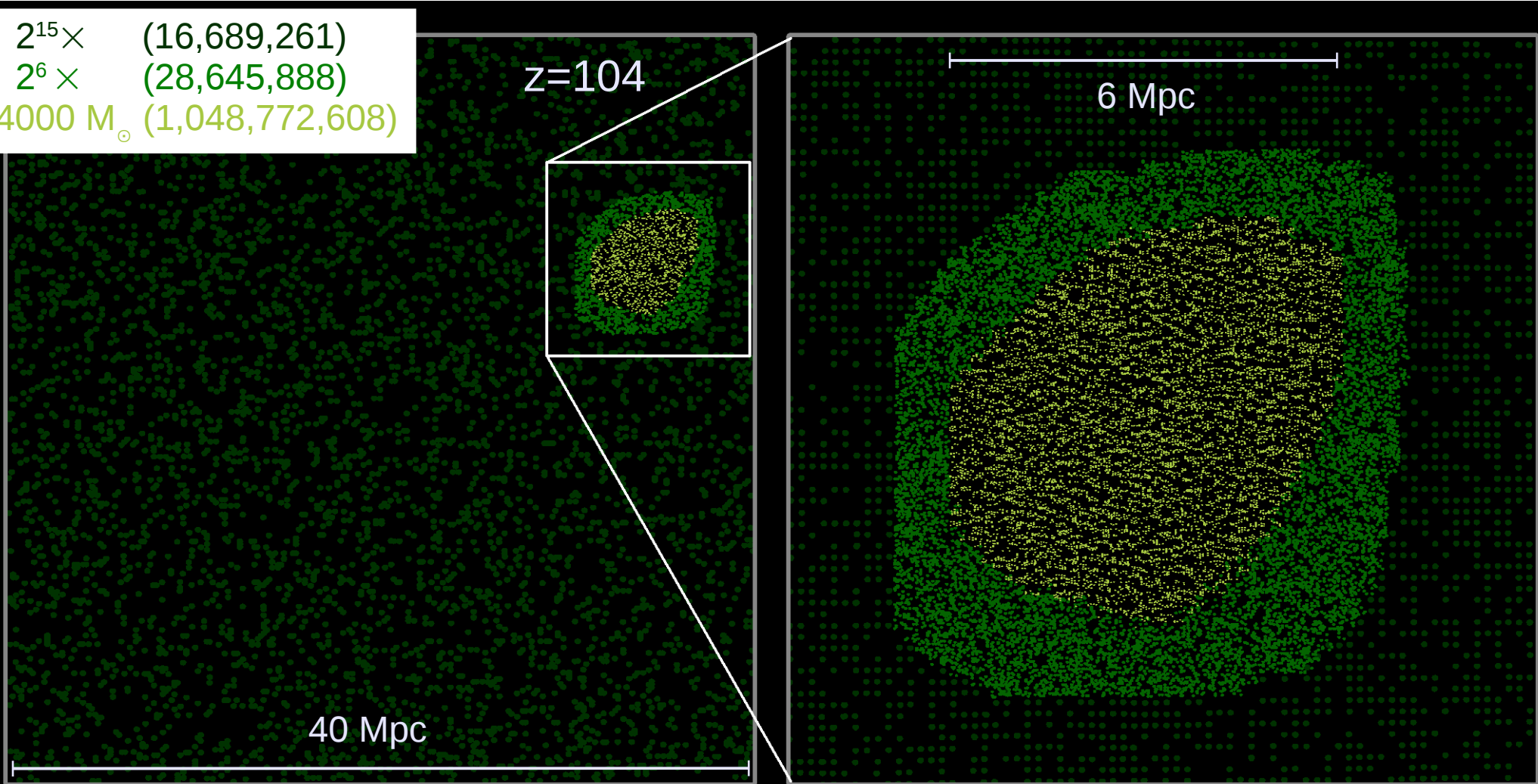
Back in the initial conditions, cover the Lagrangian volume of the halo of interest with a billion (or more) high resolution particles (yellow), which have a mass of $4,000 M_{\odot}$ in VL2.

$2^{15} \times$	(16,689,261)
$2^6 \times$	(28,645,888)
$4000 M_{\odot}$	(1,048,772,608)

$z=104$

40 Mpc

6 Mpc



Multi-mass initial conditions

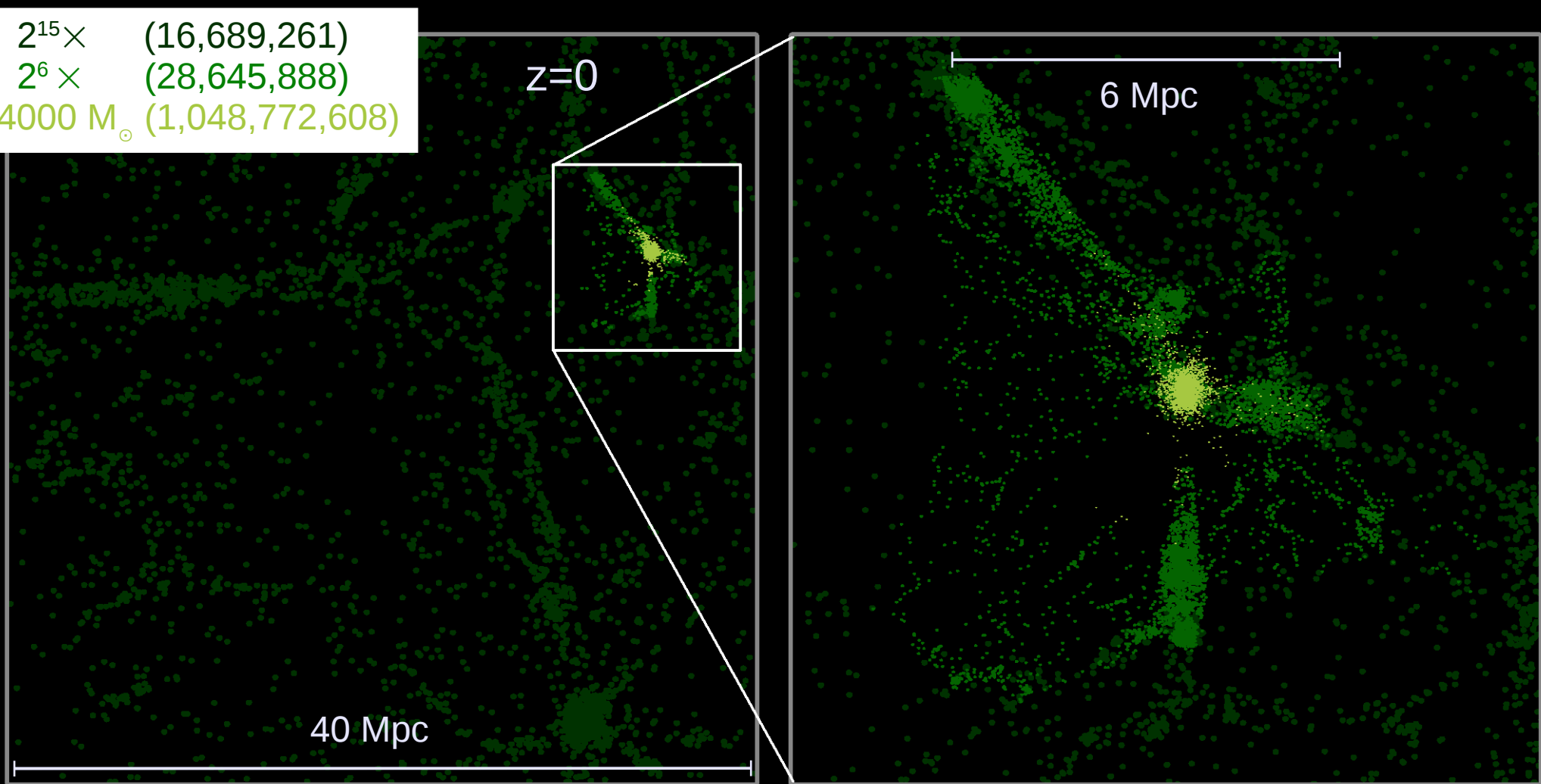
Gravity has caused the dark matter to clump. Only the region covered by high resolution particles is of interest. The low resolution particles provide the larger scale environment, e.g. the tidal field.

$2^{15} \times$	(16,689,261)
$2^6 \times$	(28,645,888)
$4000 M_{\odot}$	(1,048,772,608)

$z=0$

40 Mpc

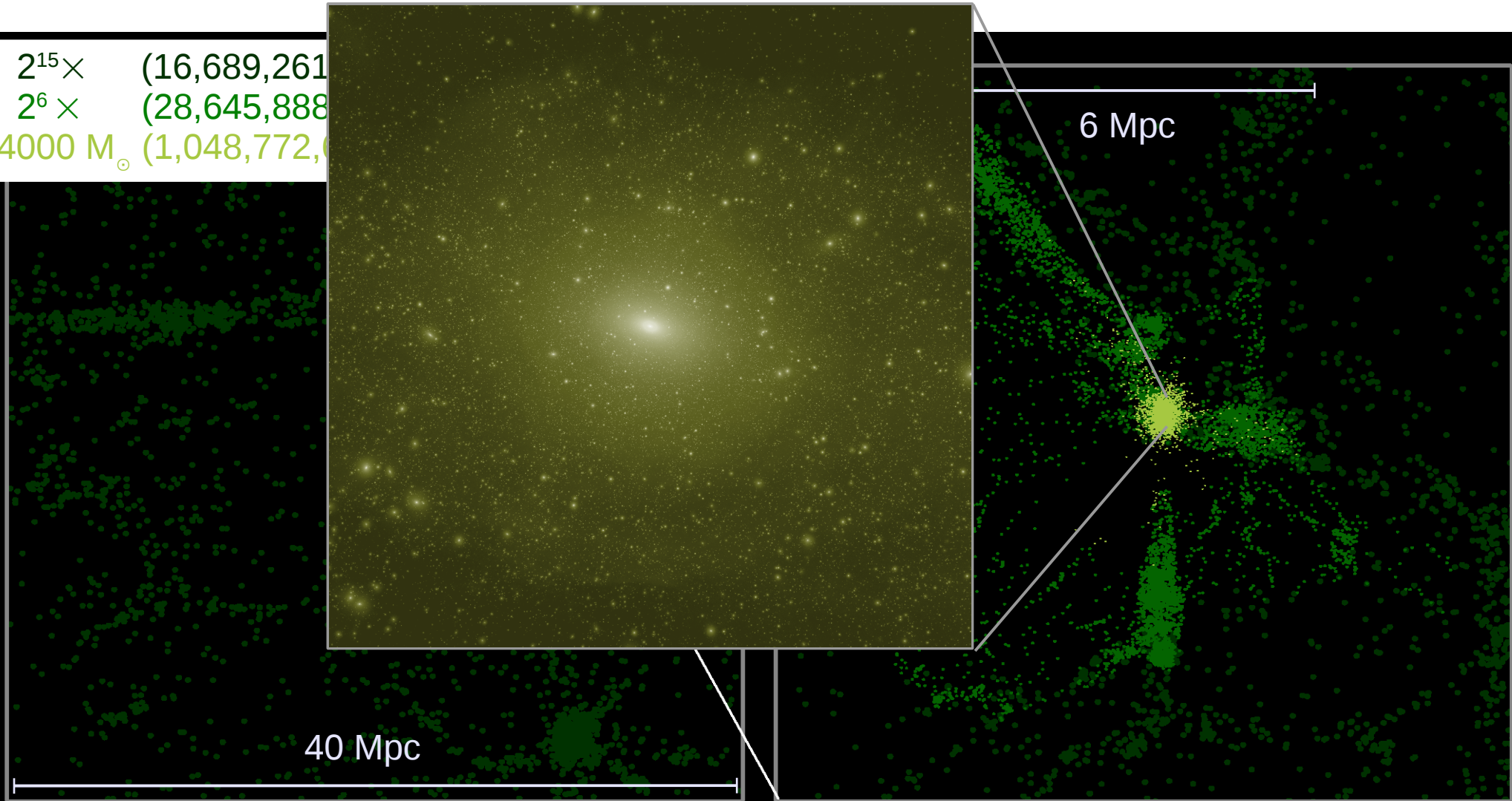
6 Mpc



Multi-mass initial conditions

Gravity has caused the dark matter to clump. Only the region covered by high resolution particles is of interest. The low resolution particles provide the larger scale environment, e.g. the tidal field.

$2^{15} \times$	(16,689,261
$2^6 \times$	(28,645,888
$4000 M_{\odot}$	(1,048,772,



PKDGRAV – a hierarchical n-body tree code

Written by Joachim Stadel (U. Zurich) and Tom Quinn (U. Washington)

Newton's equations of motion in co-moving coordinates

$$\frac{d\vec{x}}{dt} = \vec{v}$$

$$\frac{d\vec{v}}{dt} + 2H(a)\vec{v} = -\frac{1}{a^2}\vec{\nabla}\phi.$$

Cosmology, the expansion of the universe

$$H(a) = \dot{a}/a \text{ (Hubble constant)}$$

$$\frac{\ddot{a}}{a} = -\frac{4}{3}\pi G\rho_b(t) + \frac{\Lambda}{3} \text{ (2nd Friedman equation)}$$

Gravitational potential

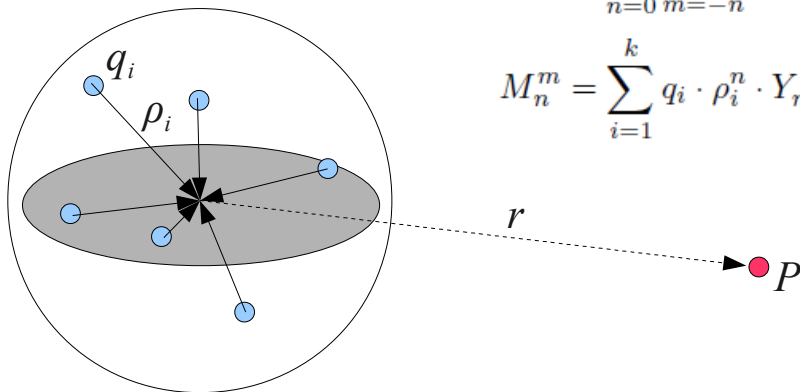
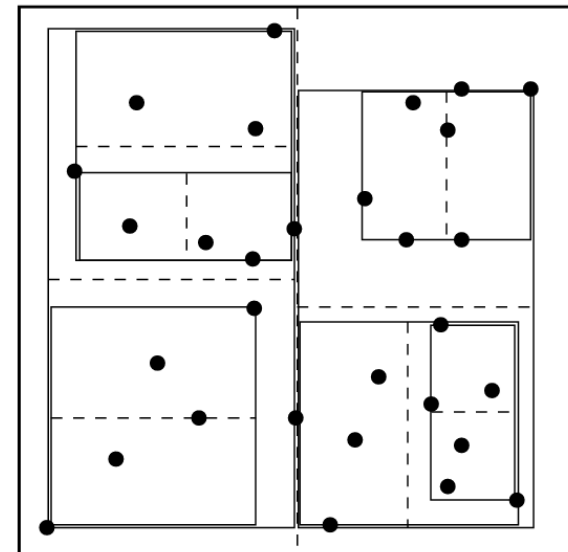
$$\nabla^2\phi = 4\pi G\rho a^2 - \Lambda a^2 + 3a\ddot{a}$$

$$= 4\pi G(\rho - \rho_b)a^2$$



The exact spatial configuration of distant sources is unimportant. Can use a **multipole expansion of the potential** instead.

Spatially bisected binary tree:
gravity calculation $O(N^2) \rightarrow O(N \log N)$



$$\phi(P) = \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{M_n^m}{r^{n+1}} \cdot Y_n^m(\theta, \phi)$$

$$M_n^m = \sum_{i=1}^k q_i \cdot \rho_i^n \cdot Y_n^{-m}(\alpha_i, \beta_i)$$

Fast Multipole Method: $O(N \log N) \rightarrow O(N)$

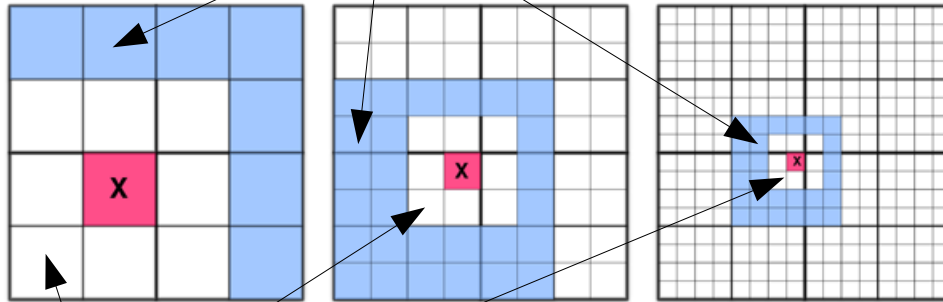
Algorithm: Greengard & Rohklin (1985)

"A short course in fast multipole methods" by Beatson & Greengard (1997)

"An overview of fast multipole methods" by Ihler (2004)

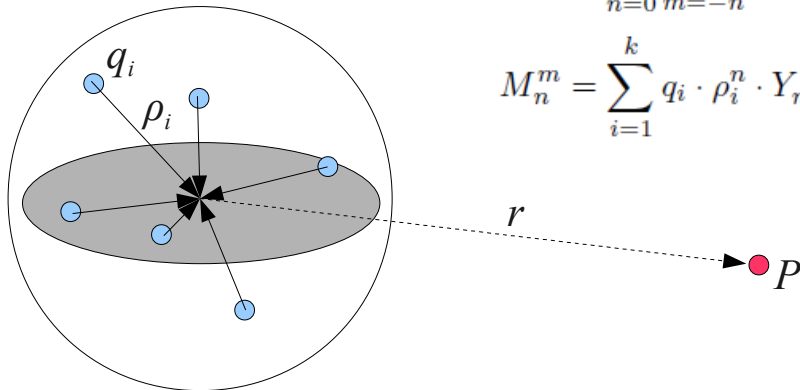
Interaction list

use multipole expansion



Near neighbors

evaluate at lower level,
 $O(N^2)$ direct summation at finest level.



$$\phi(P) = \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{M_n^m}{r^{n+1}} \cdot Y_n^m(\theta, \phi)$$

$$M_n^m = \sum_{i=1}^k q_i \cdot \rho_i^n \cdot Y_n^{-m}(\alpha_i, \beta_i)$$



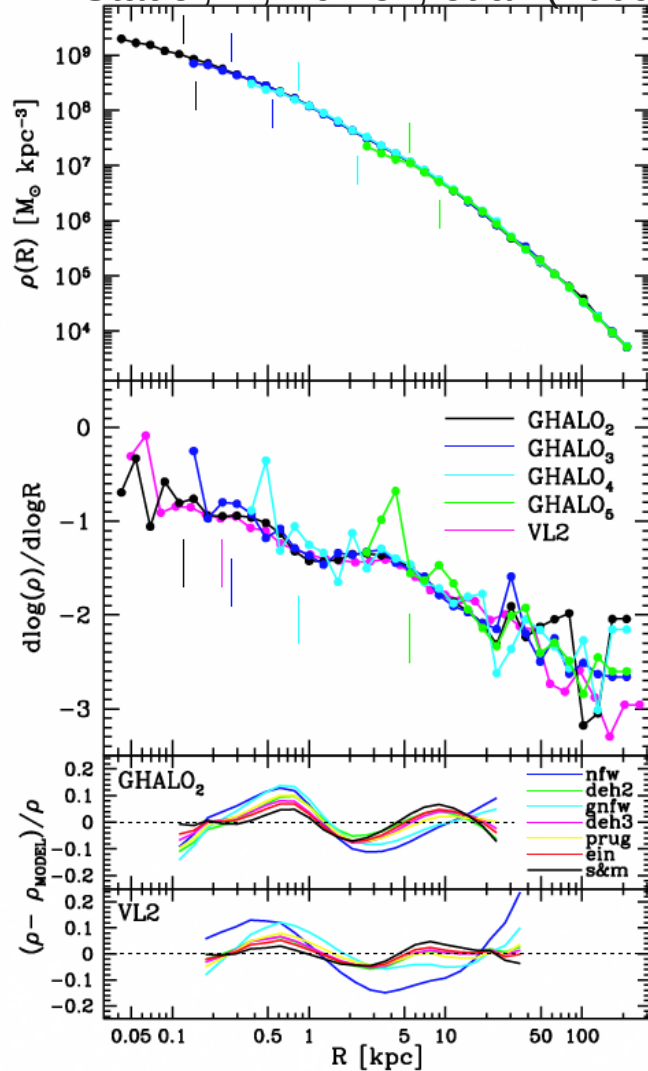
The exact spatial configuration of distant sources is unimportant. Can use a **multipole expansion of the potential** instead.

A regular tree code (e.g. Barnes & Hut) is **$O(N \log N)$** , because **every particle (N) is visited at every level (logN)** in order to construct and evaluate the multipole expansion.

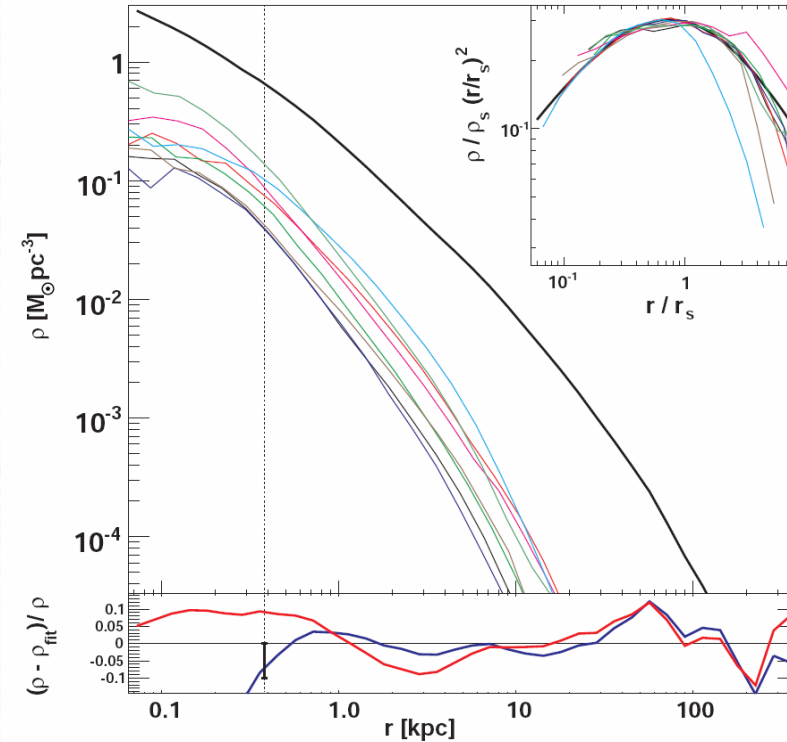
In the **Fast Multipole Method**, on every level **only every cell is visited once**, in order to construct a single expansion for the potential from all particles outside a cell's near neighbors. **$O(N)$**

Dark Matter Halos Are “Cuspy”

Stadel, ..., Kuhlen, et al. (2009)



Diemand, Kuhlen, et al. (2008)



The host halo and the subhalos exhibit a “universal” density profile.

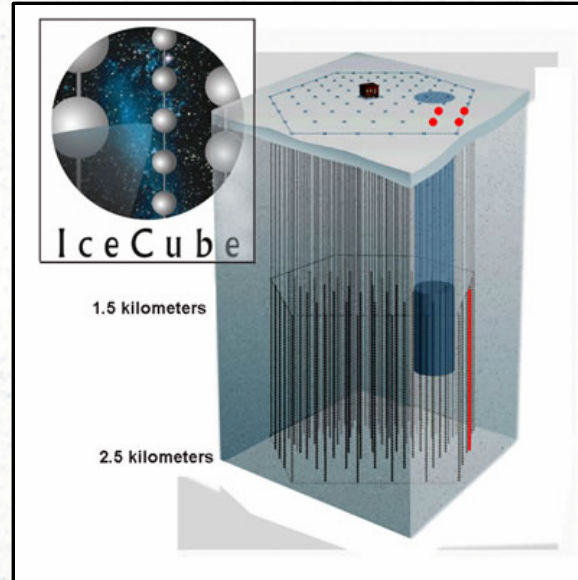
- Steep outer slope: $\rho \sim r^{-\gamma}$ $\gamma \simeq 3 - 4$
- Rolls over to isothermal slope: $\gamma = 2$
- Continues to become shallower, but remains cuspy $\gamma \simeq 0.8$ as far in as simulations can resolve.

Substructure Relevance for Indirect Detection

Gamma-rays
(Fermi, A.C.T.'s)



Neutrinos
(e.g. IceCube)



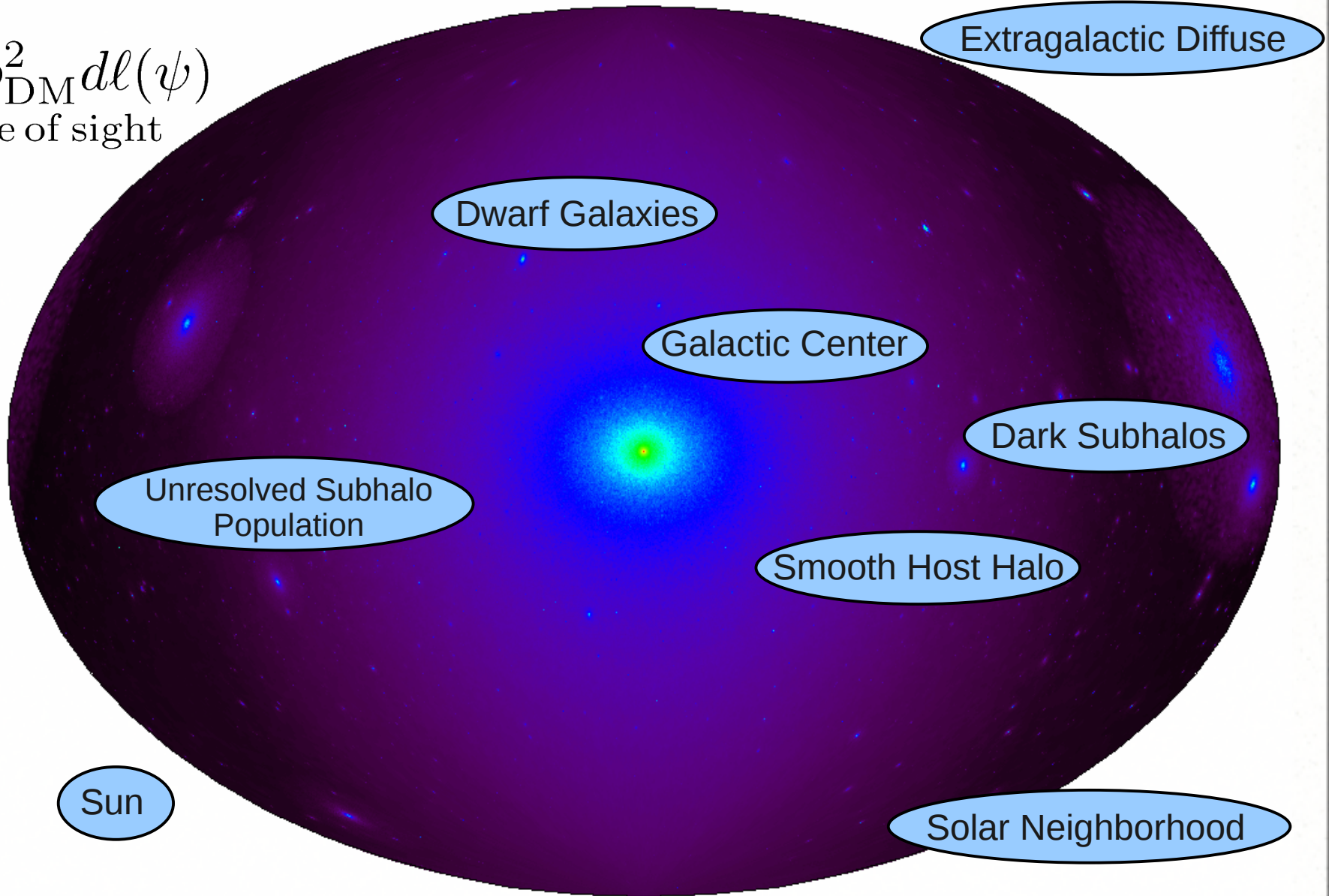
Positrons/Anti-protons
(e.g. Pamela)



$$N_{\gamma} = \left[\int_{\text{line of sight}} \rho_{\text{DM}}^2 dl(\psi) \right] \frac{\langle \sigma v \rangle}{2M_{\chi}^2} \left[\int_{E_{th}}^{M_{\chi}} \left(\frac{dN_{\gamma}}{dE} \right) A_{\text{eff}}(E) dE \right] \frac{\Delta\Omega}{4\pi} \tau_{\text{exp}}$$

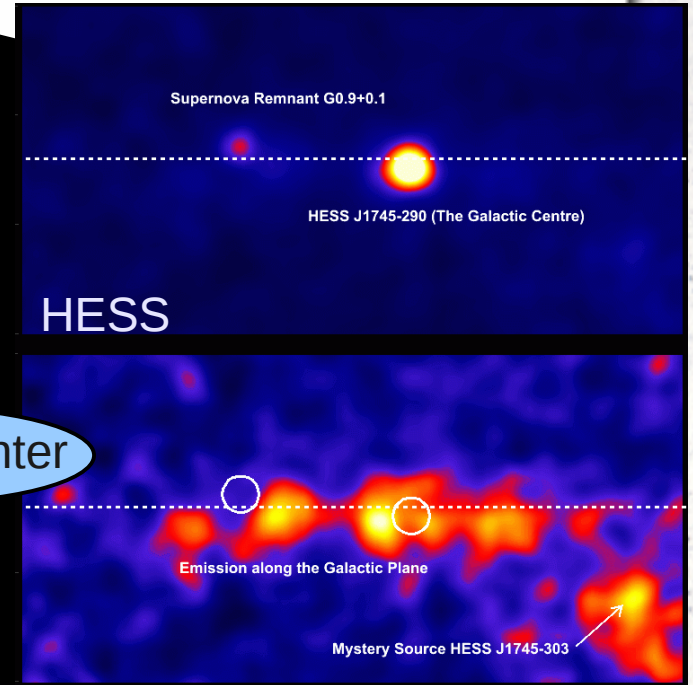
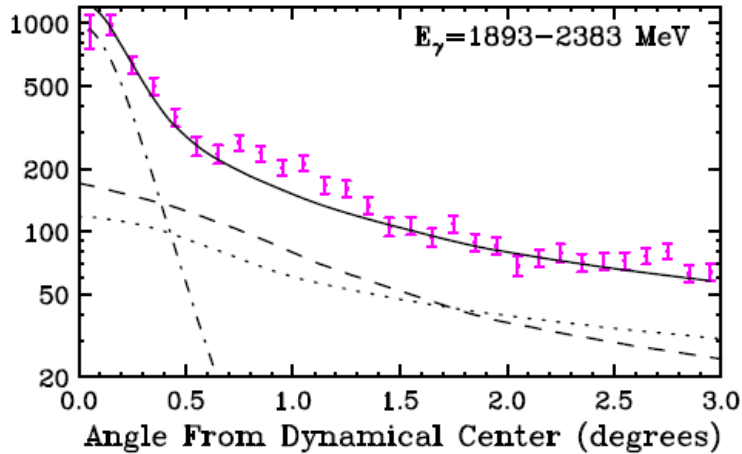
Substructure Relevance for Indirect Detection

$$\int_{\text{line of sight}} \rho_{\text{DM}}^2 dl(\psi)$$

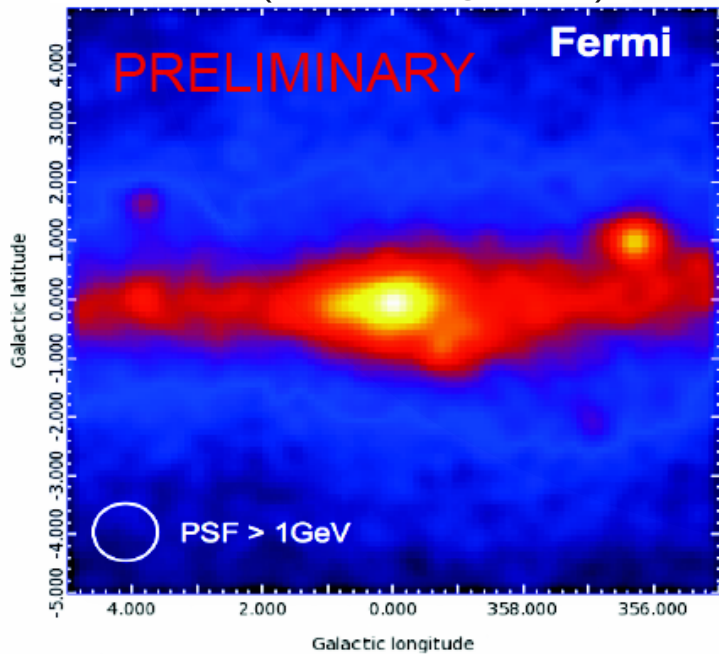


Substructure Relevance for Indirect Detection

Hooper & Goodenough (2010)

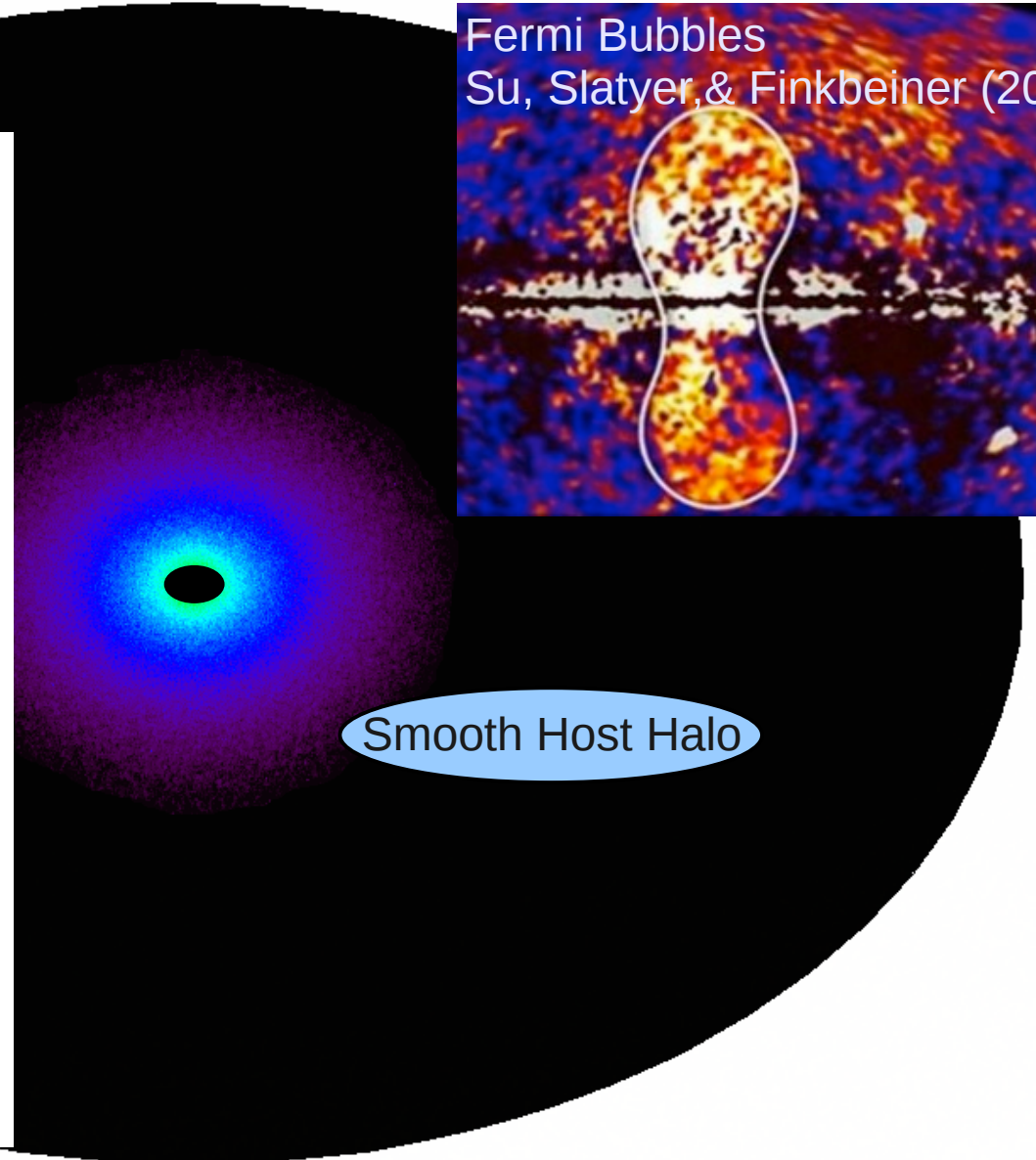
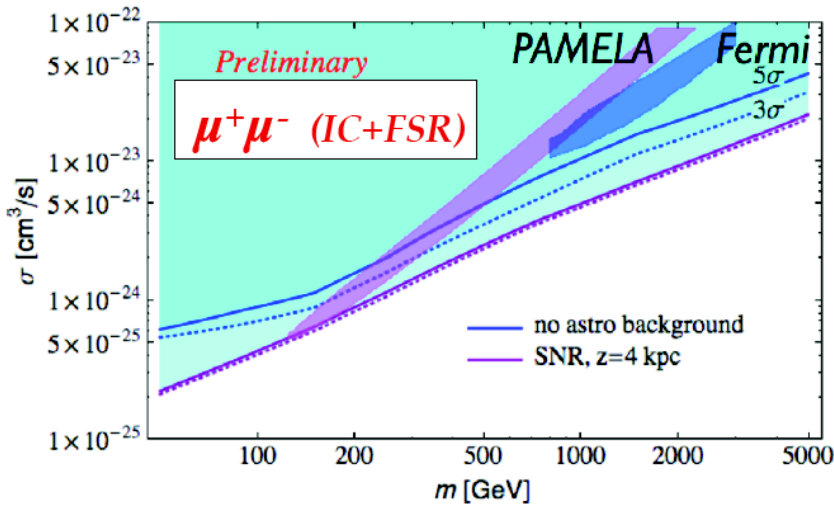
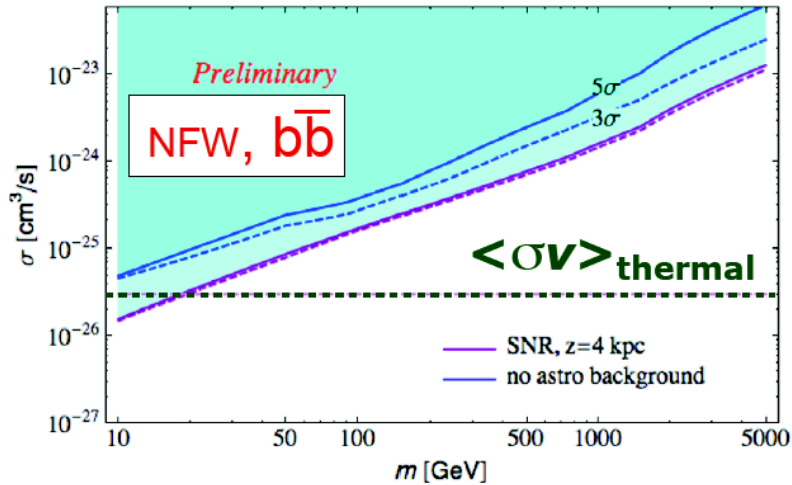


Gustaffson (PPC 2011 @CERN)



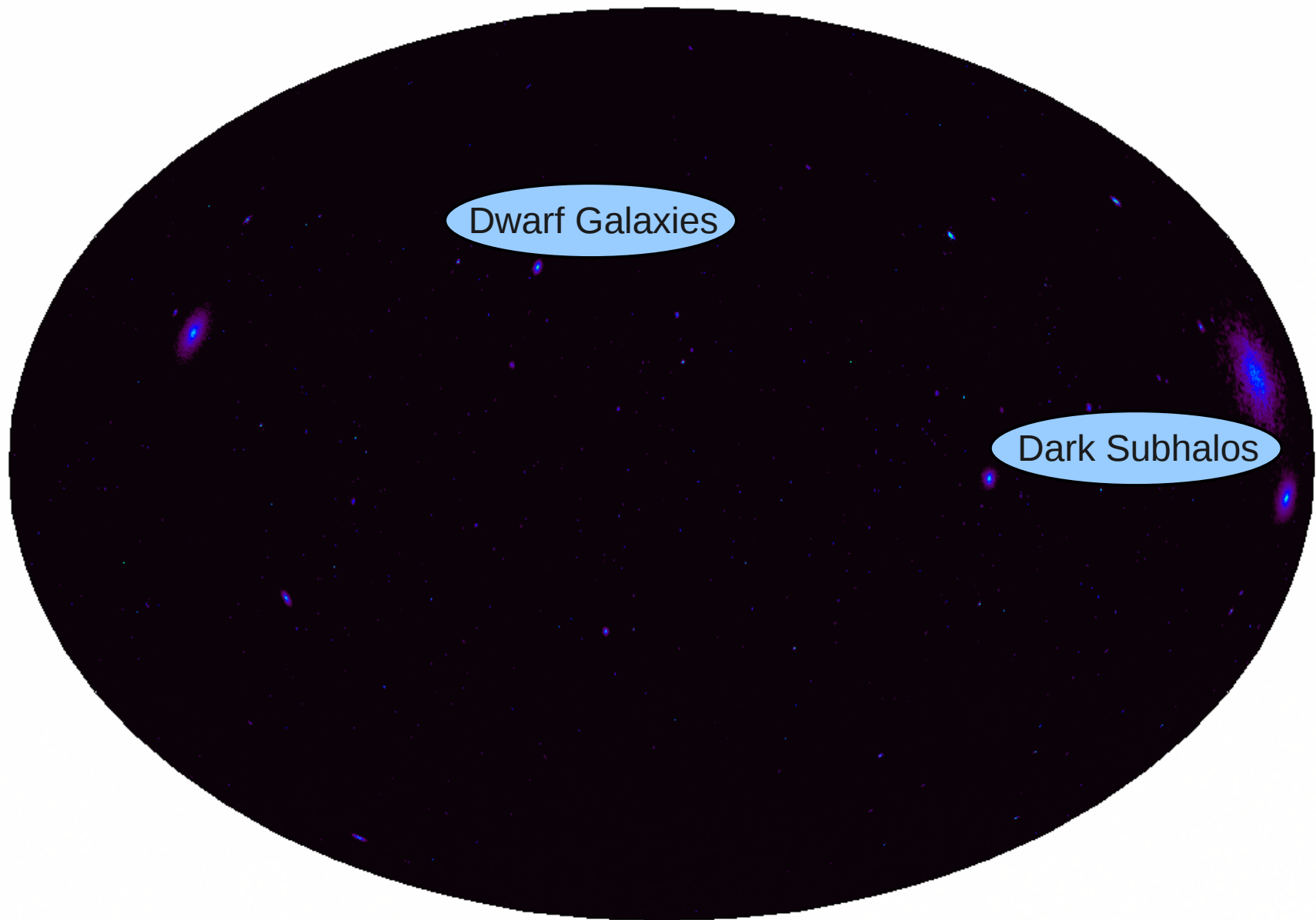
Substructure Relevance for Indirect Detection

Fermi Bubbles
Su, Slatyer, & Finkbeiner (2010)

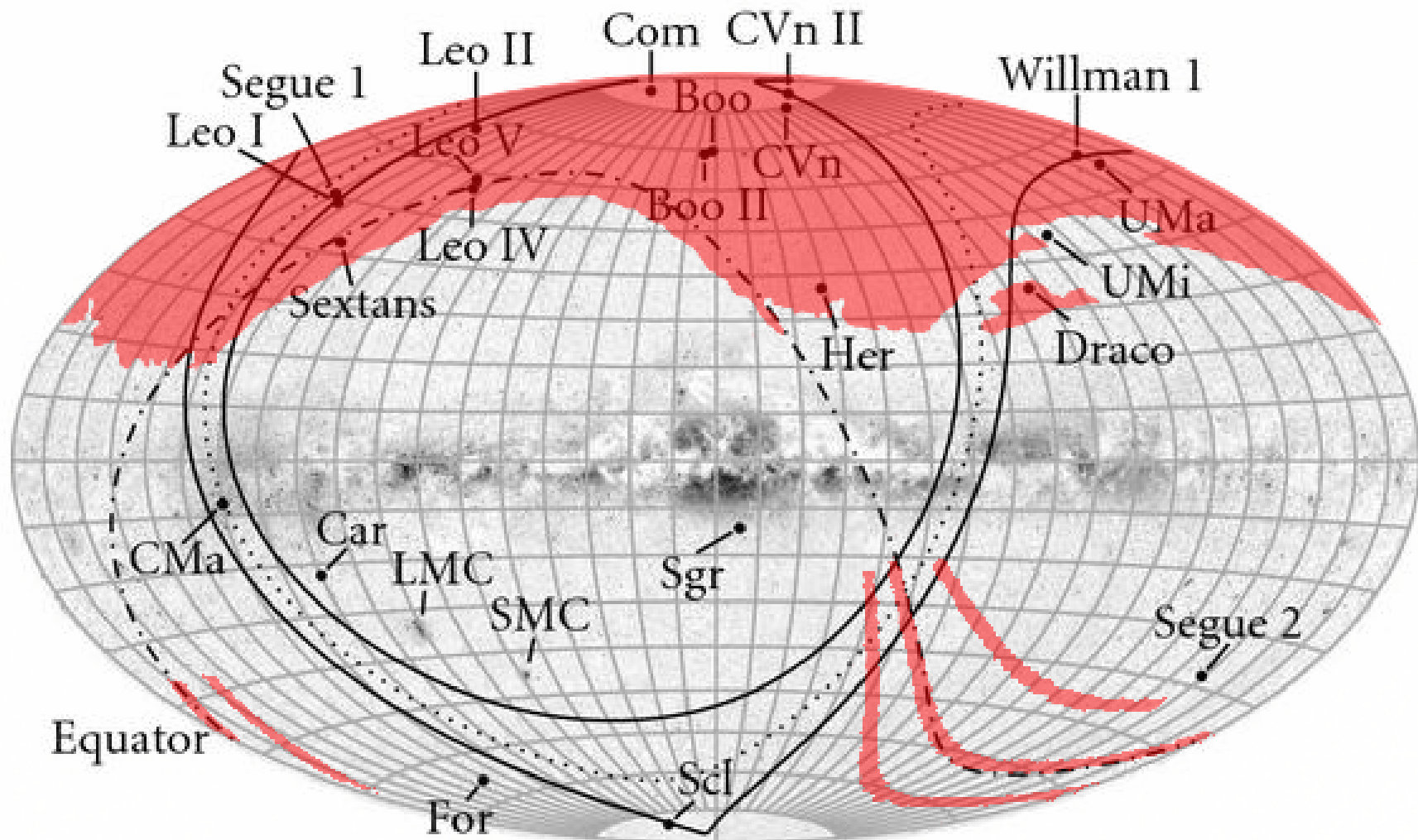


Gustaffson (PPC 2011 @CERN)

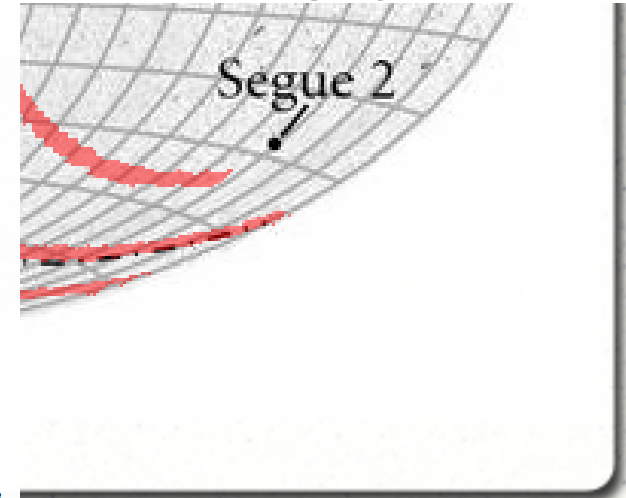
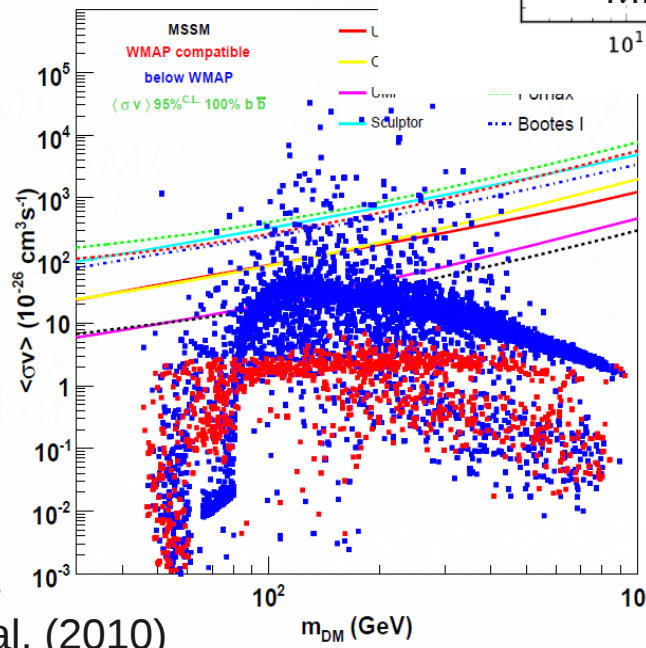
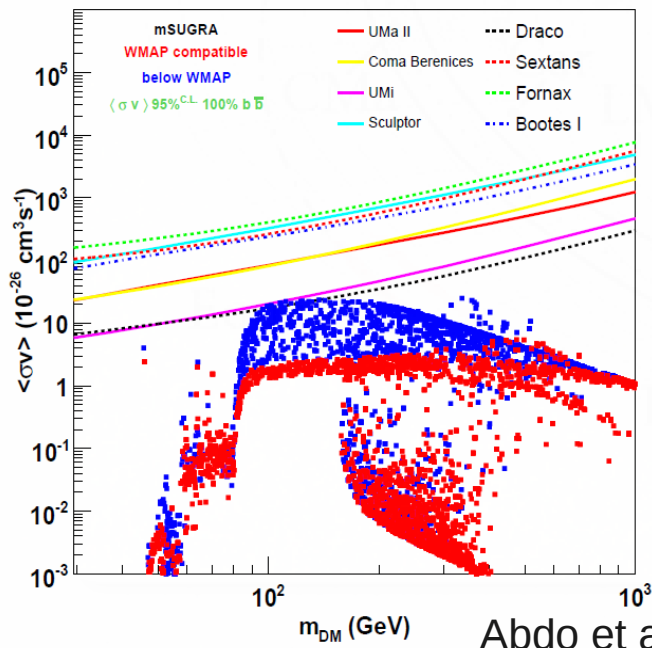
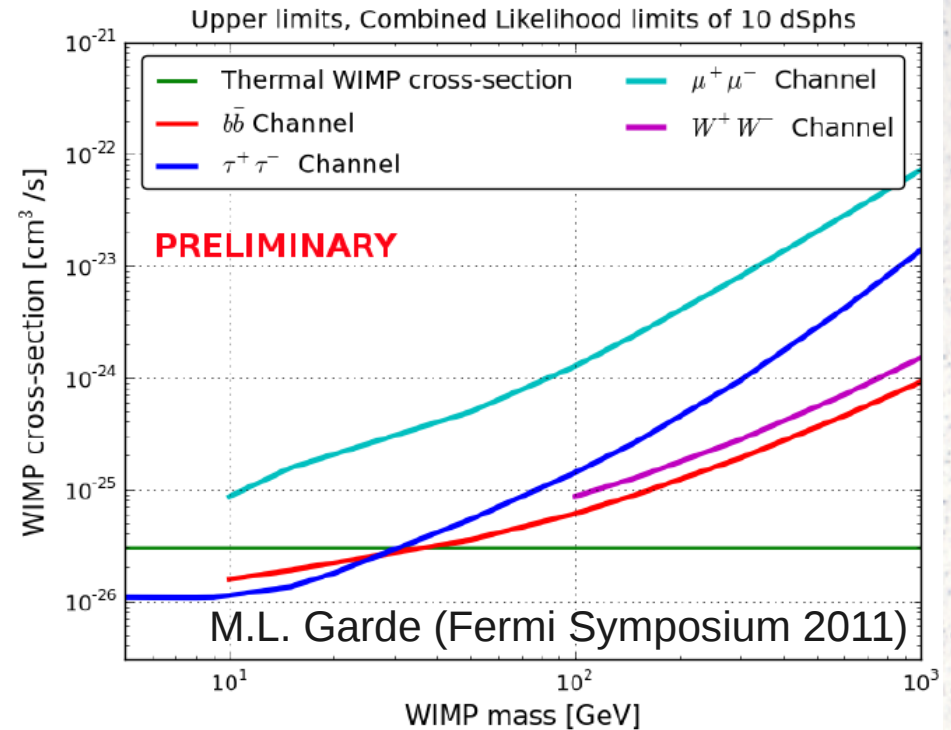
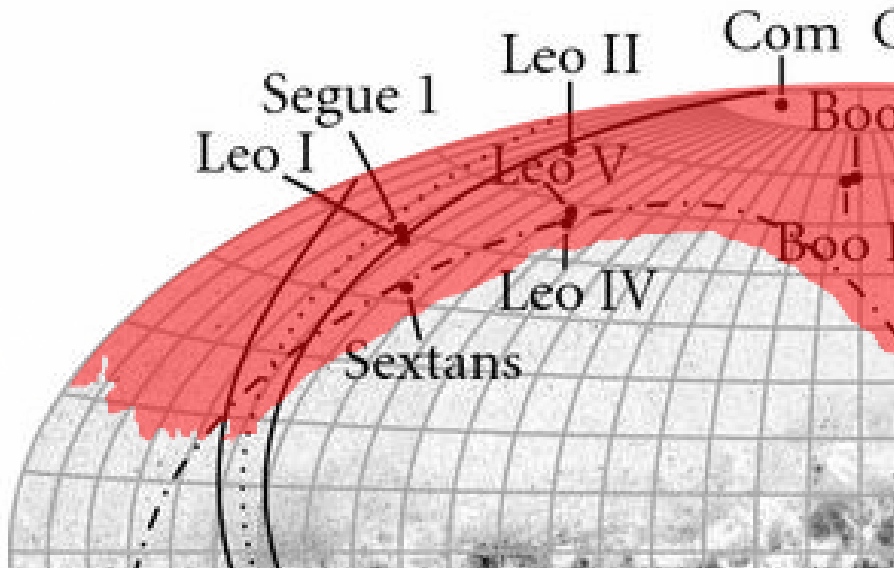
Substructure Relevance for Indirect Detection



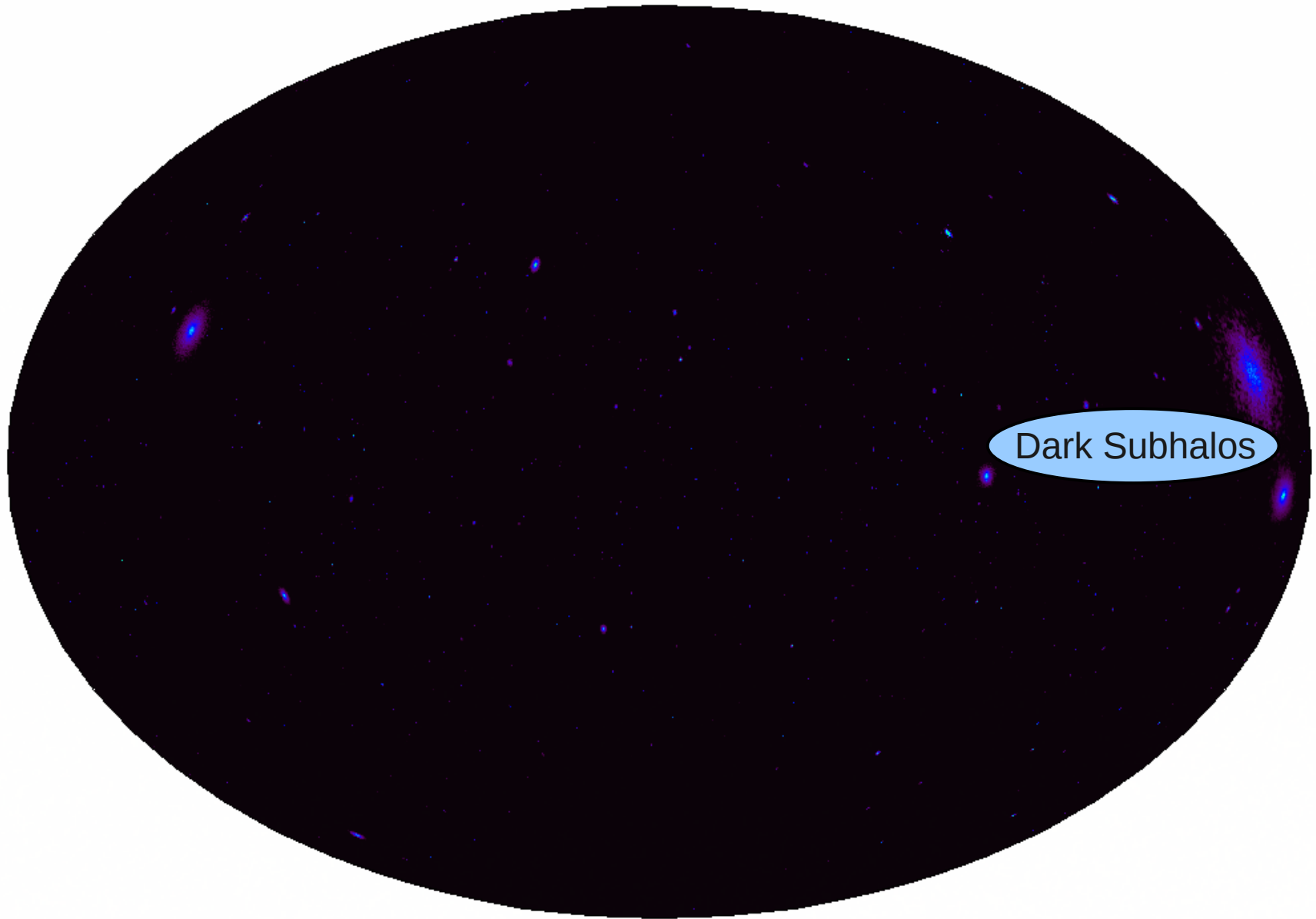
Substructure Relevance for Indirect Detection



Substructure Relevance for Indirect Detection

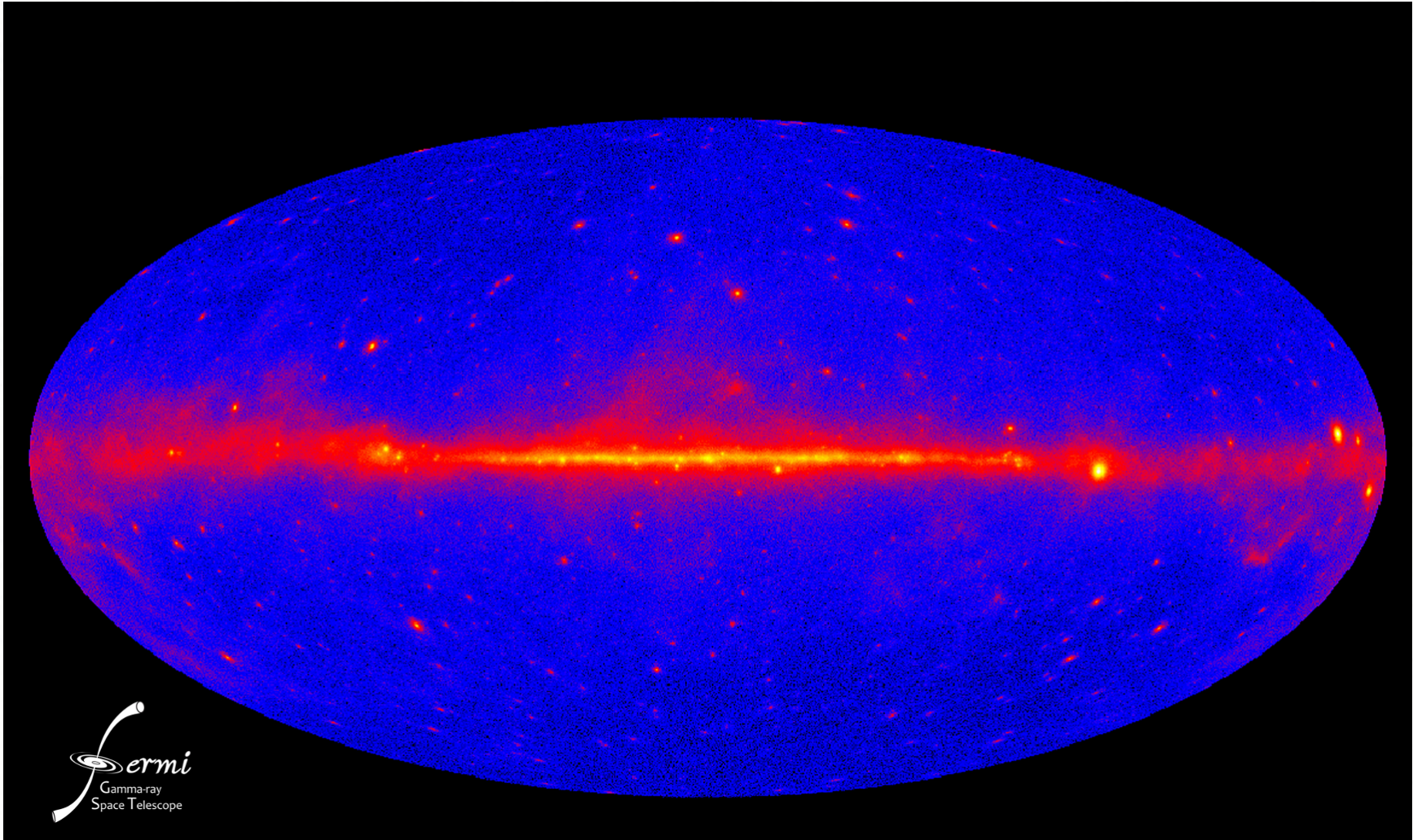


Substructure Relevance for Indirect Detection



Indirect Detection of Subhalos

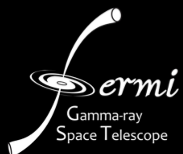
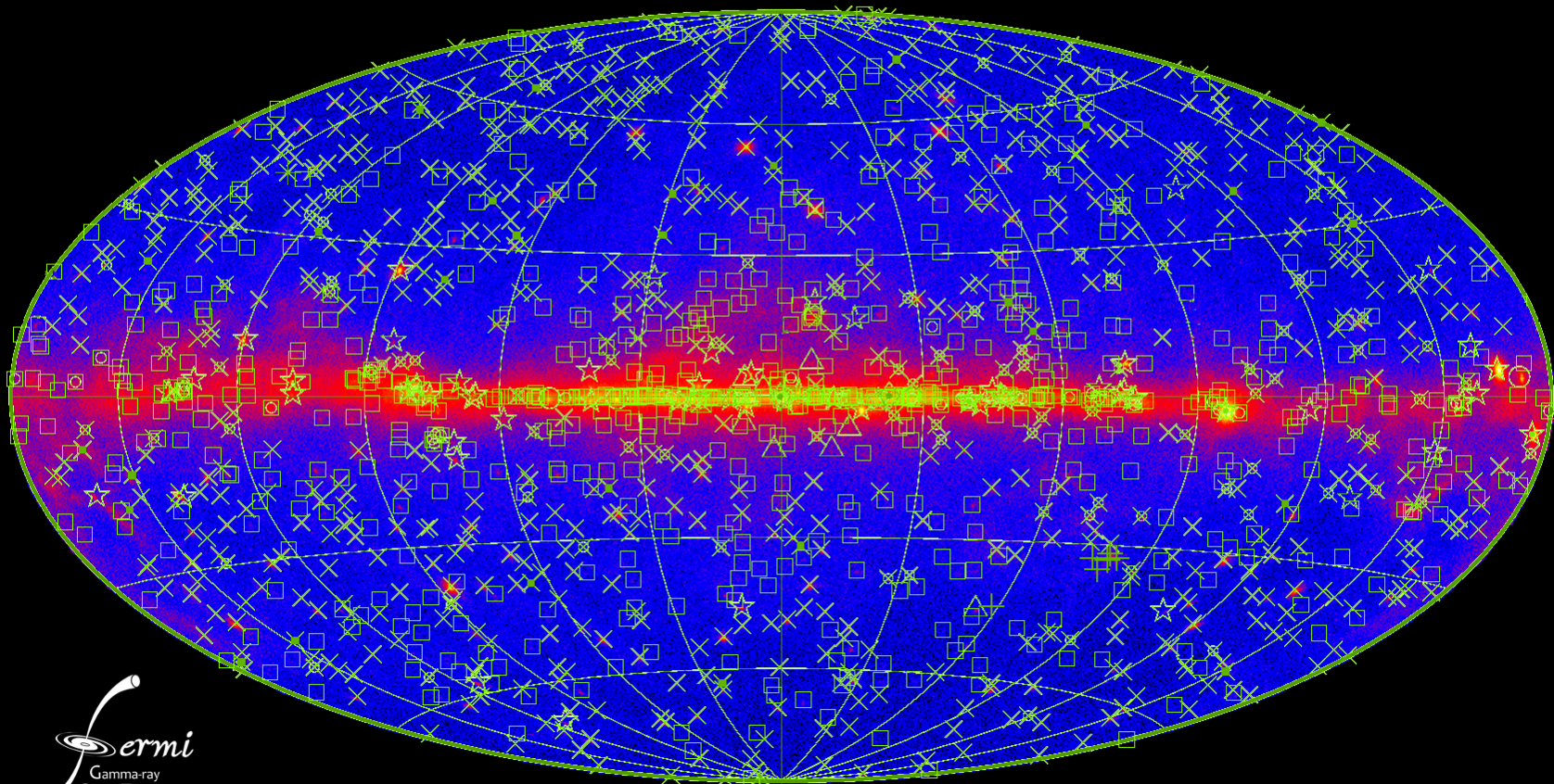
The Fermi Gamma-ray Space Telescope was launched on June 11th 2008 and has been observing the sky for more than 2 years.



Indirect Detection of Subhalos

So far, now dark matter signal has been detected. ☹ Stay tuned...

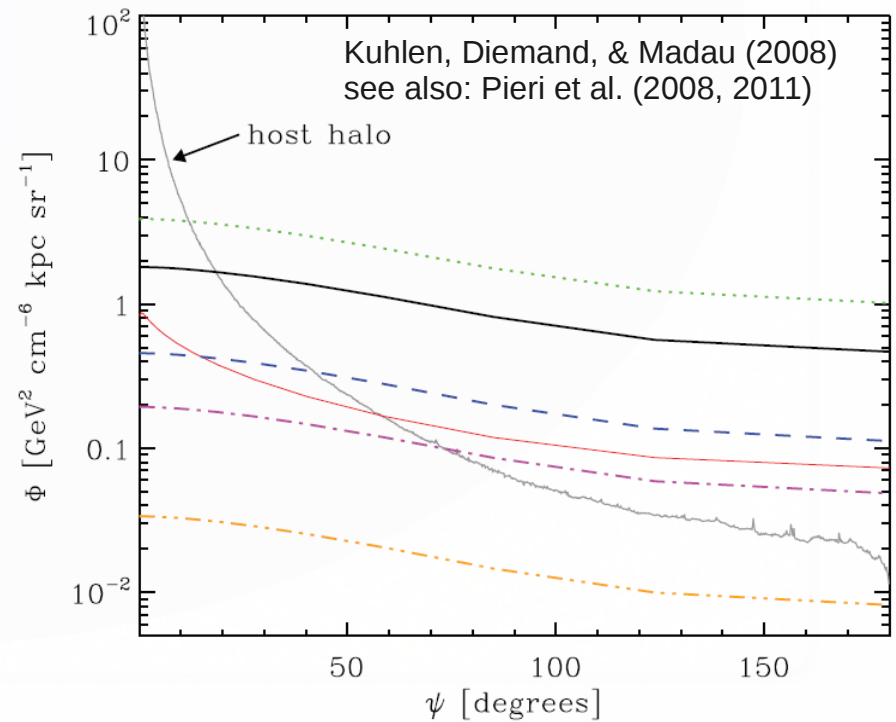
1FGL Source Catalog
(Abdo et al. 2010)



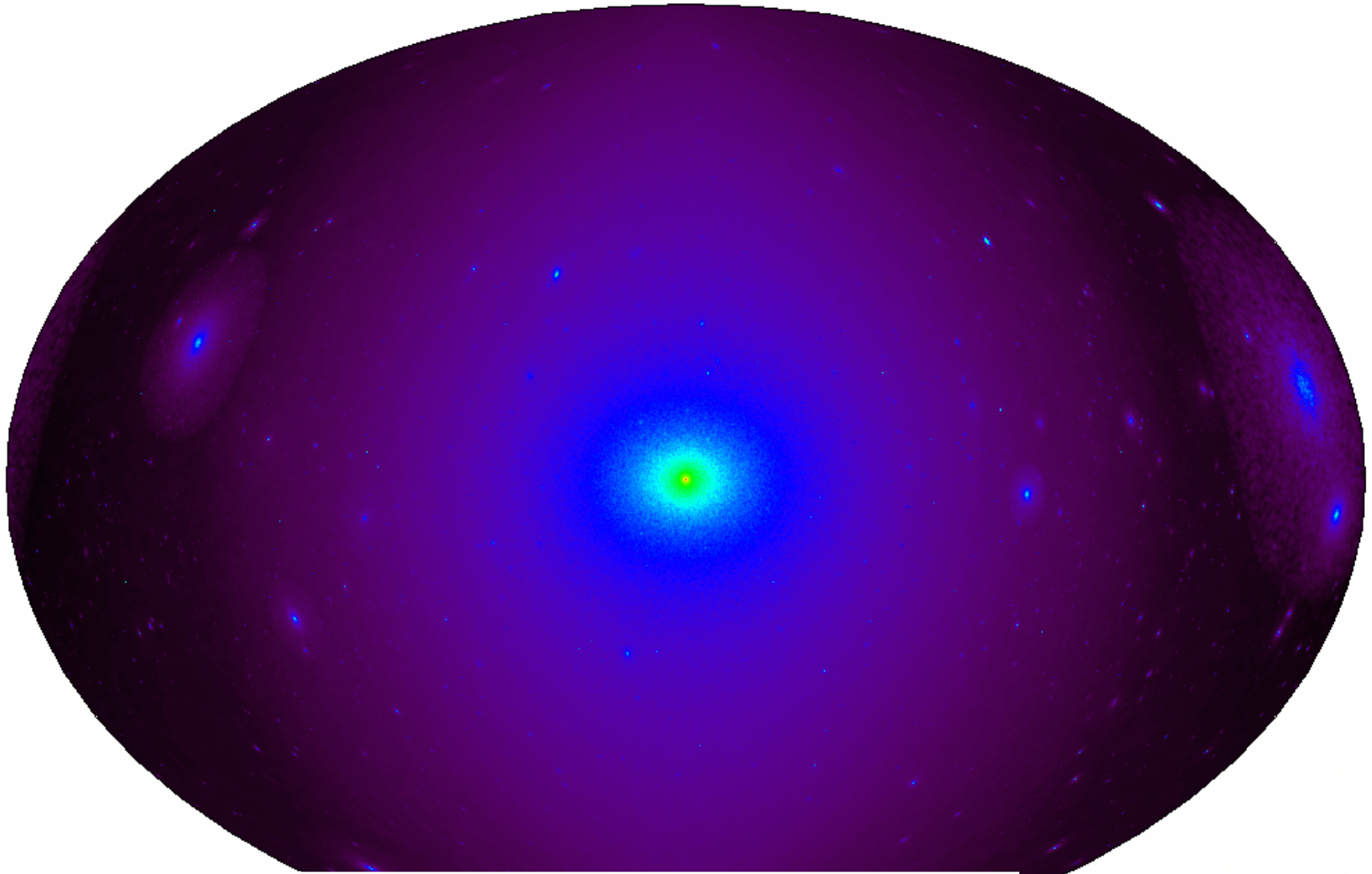
- | | | | |
|--------------------|---|--------------|--------------------|
| □ No association | ▣ Possible association with nearby SNR or PWN | | |
| × AGN - blazar | * Starburst Gal | ☆ Pulsar | ★ Pulsar w/PWN |
| ⊠ AGN - unknown | + Galaxy | ◇ PWN | △ Globular cluster |
| ⊗ AGN - non blazar | ○ SNR | ⊠ XRB or MQO | |

Substructure Relevance for Indirect Detection

Unresolved Subhalo Population

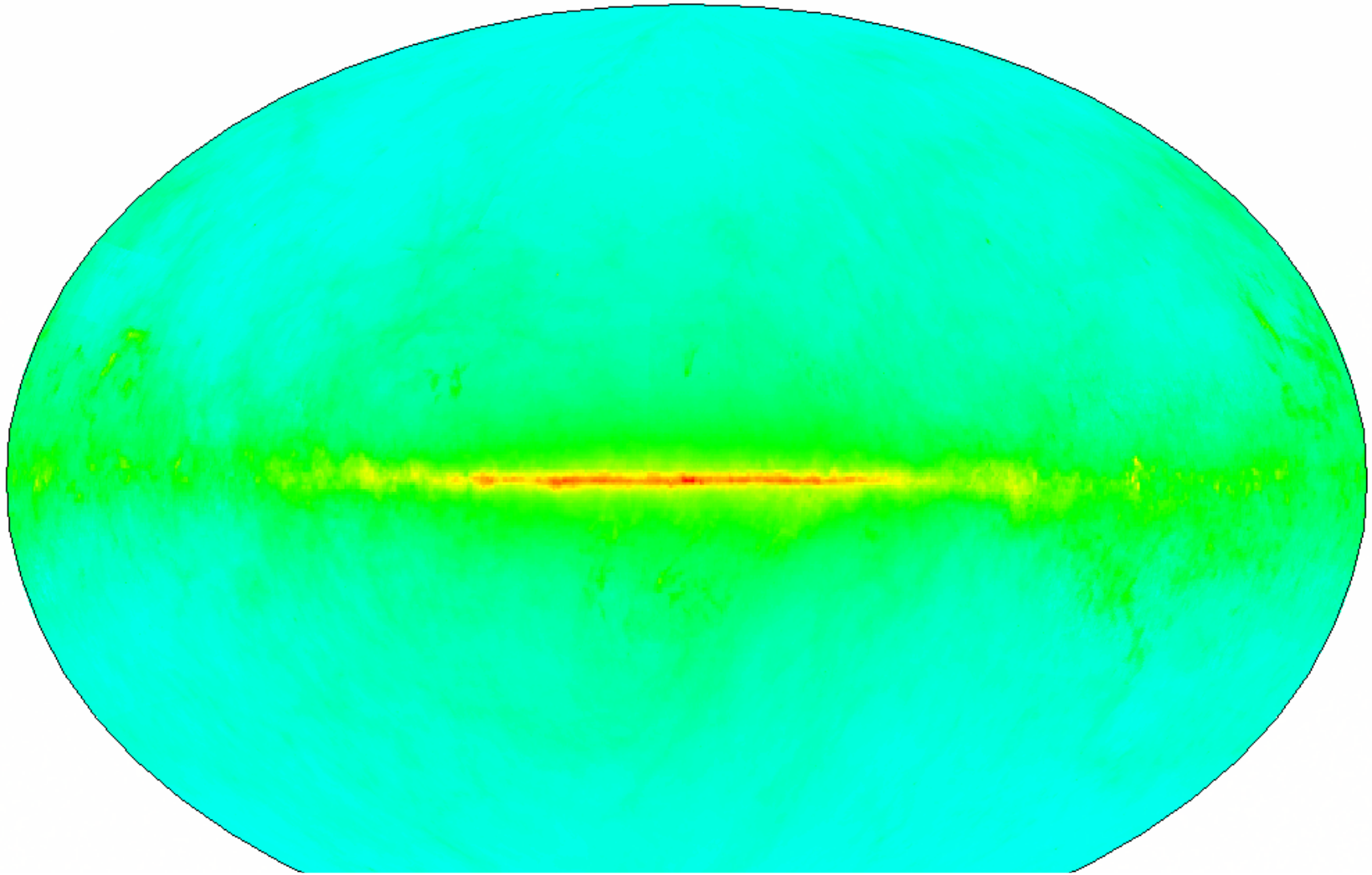


Substructure Relevance for Indirect Detection



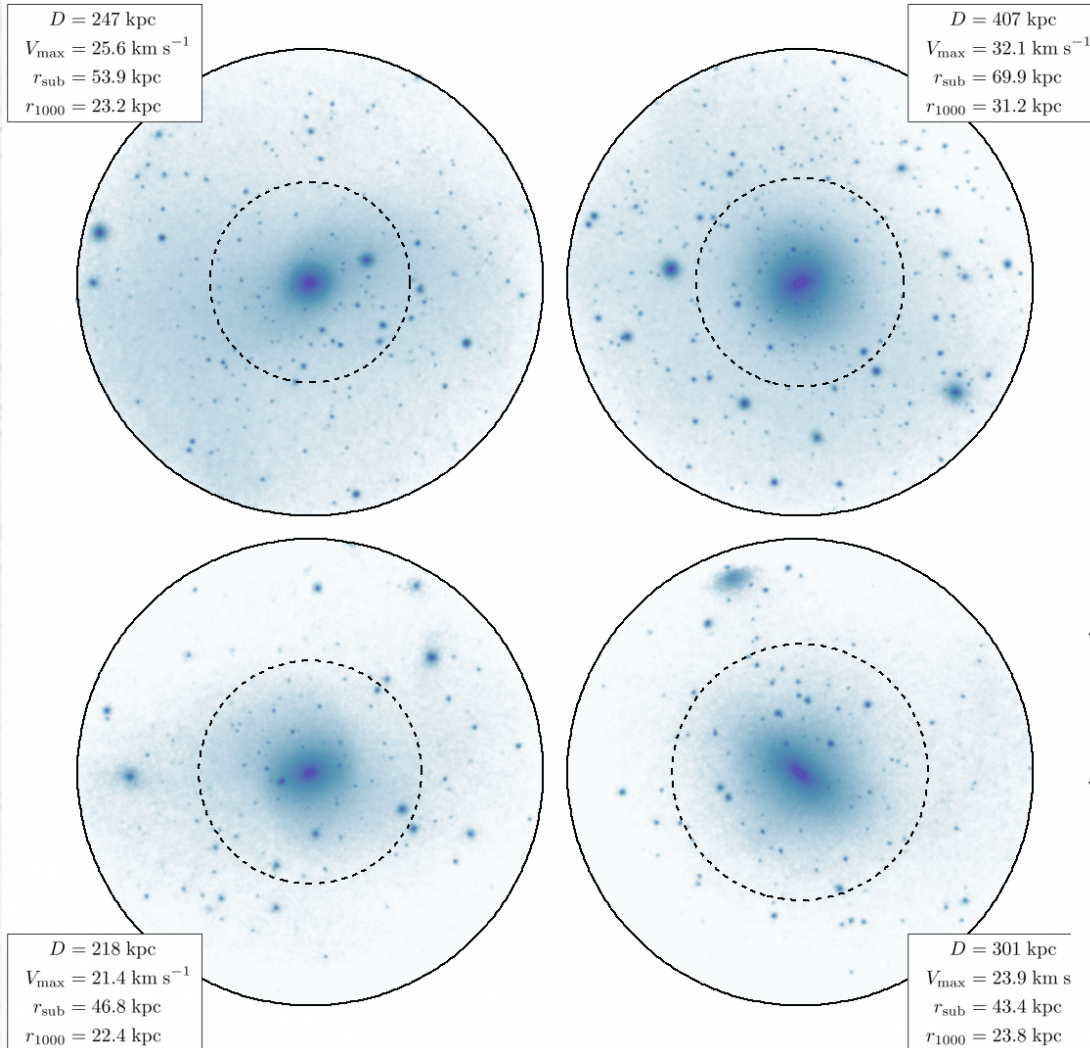
host halo + resolved subhalos + diffuse subhalo component

Substructure Relevance for Indirect Detection



host halo + resolved subhalos + diffuse subhalo component + GALPROP

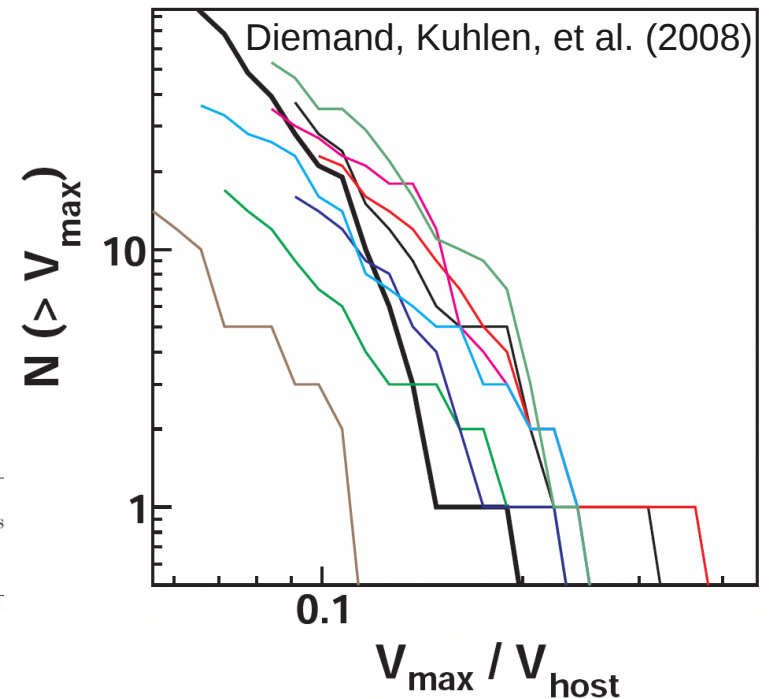
Substructure, Sub-substructure, Sub-sub-sub...



Kuhlen, Diemand, & Madau (2008)

$$\langle \rho \rangle^2 \neq \langle \rho^2 \rangle$$

A clumpy DM distribution leads to a luminosity enhancement – the **substructure boost factor**

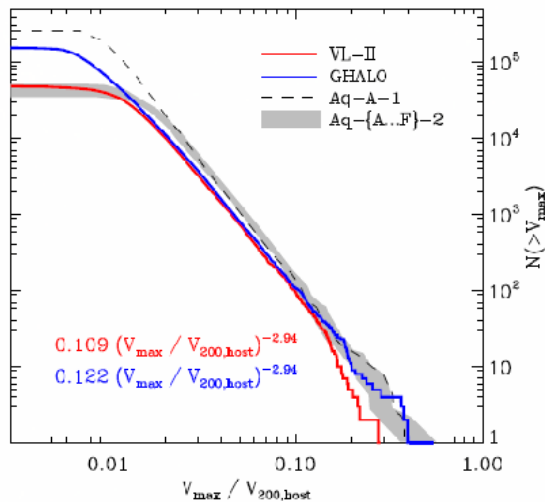


Substructure Boost: Subhalo Modeling Approach

Calibrate to numerical simulations.

Must extrapolate below resolution limit (“only” 12 orders of magnitude...)

Mass/ V_{\max} function:

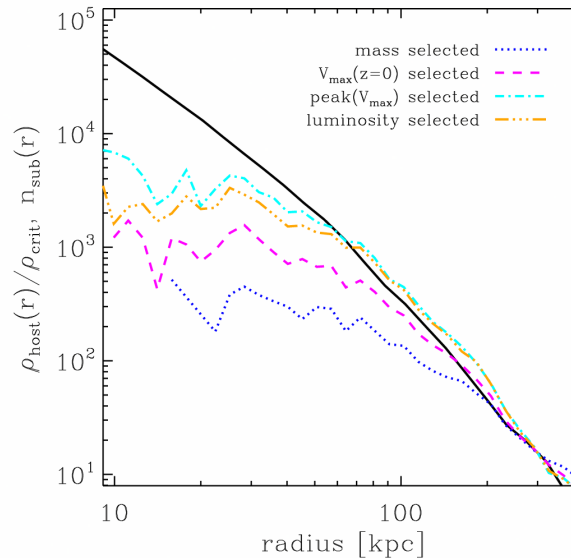


The subhalo mass function is steeply rising towards low masses.

$$\frac{dN}{dM} \sim M^{-\gamma} \quad \gamma \approx 1.9$$

$$N(>V_{\max}) \sim V_{\max}^{-\delta} \quad \delta \approx 3$$

Radial Distribution:

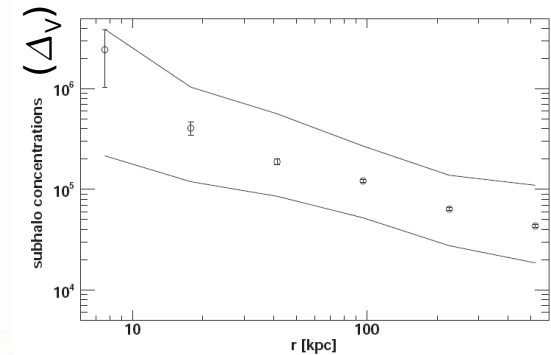


The mass-selected distribution is “anti-biased” but a luminosity-selected sample less so.

Concentration:

$$L_{\text{ann}} \sim \frac{M^2}{R^3} \frac{c^3}{f(c)^2}$$

$$\sim V_{\max}^3 \sqrt{\Delta_V}$$



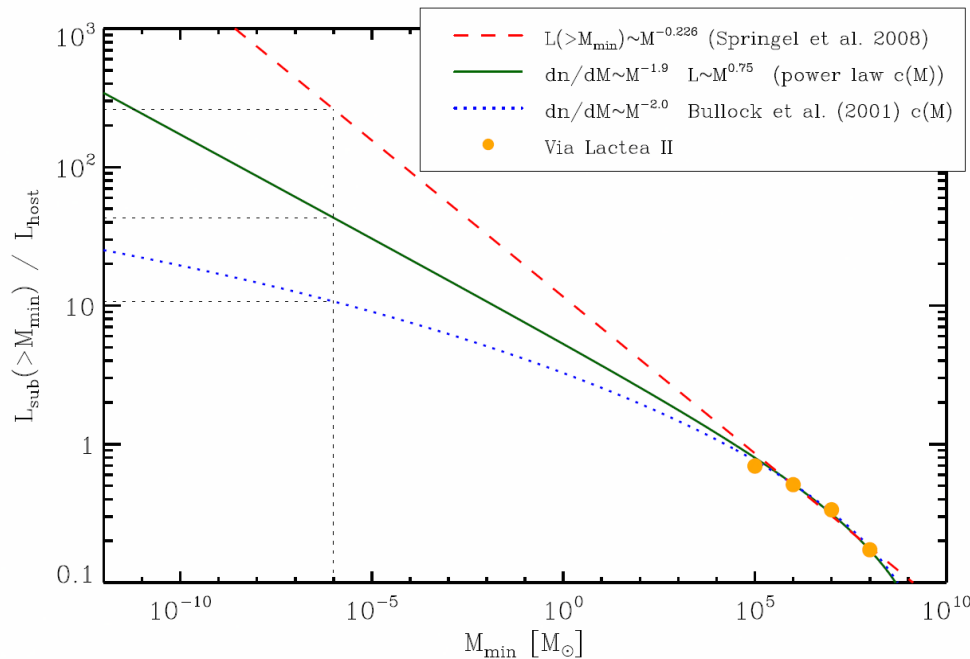
$$\Delta_V = \langle \rho(<R_{V_{\max}}) \rangle / \rho_{\text{crit}}$$

is a measure of subhalo concentration. It rises towards the center.

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Must extrapolate below resolution limit (“only” 12 orders of magnitude...)



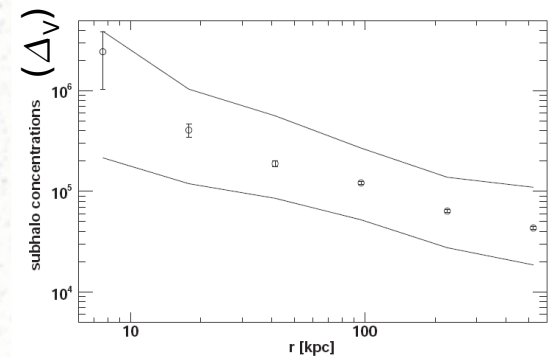
Depends **critically** on what one assumes for the concentration-mass relation for subhalos below the simulations' resolution limit!

See also: Martinez, Bullock, Kaplinghat, Strigari, & Trota (2010)

Concentration:

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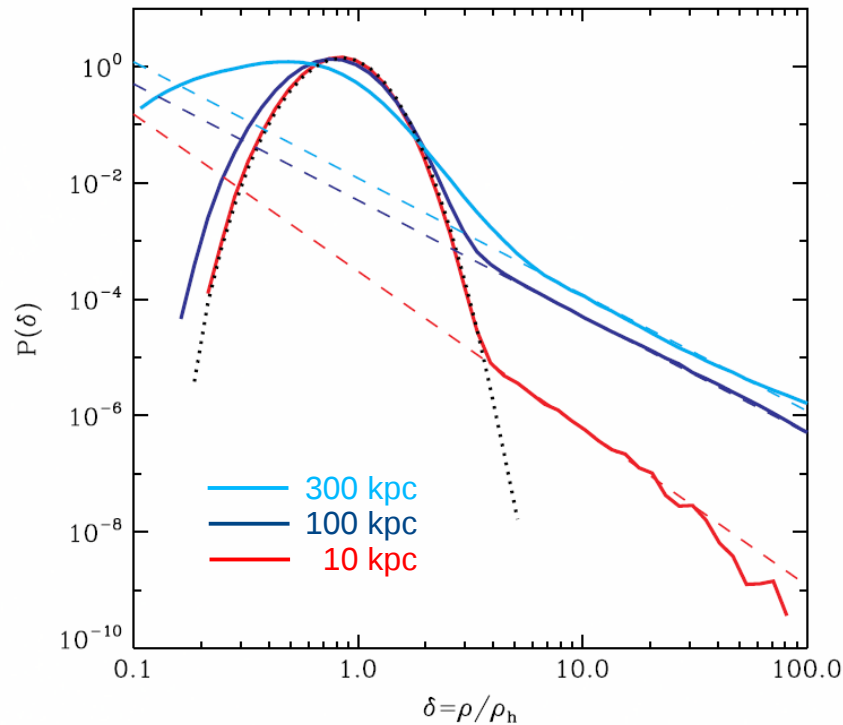
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Substructure Boost: Density Distribution Approach

Measure the PDF of ρ/ρ_{host} in the simulation.

It's fit well by a **log-normal** plus a **powerlaw** tail due to substructure.



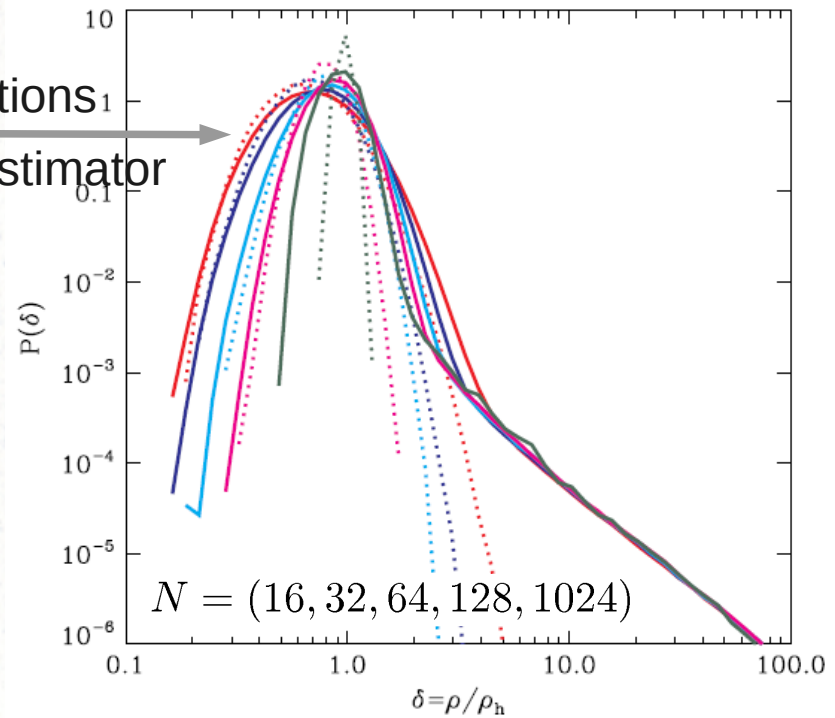
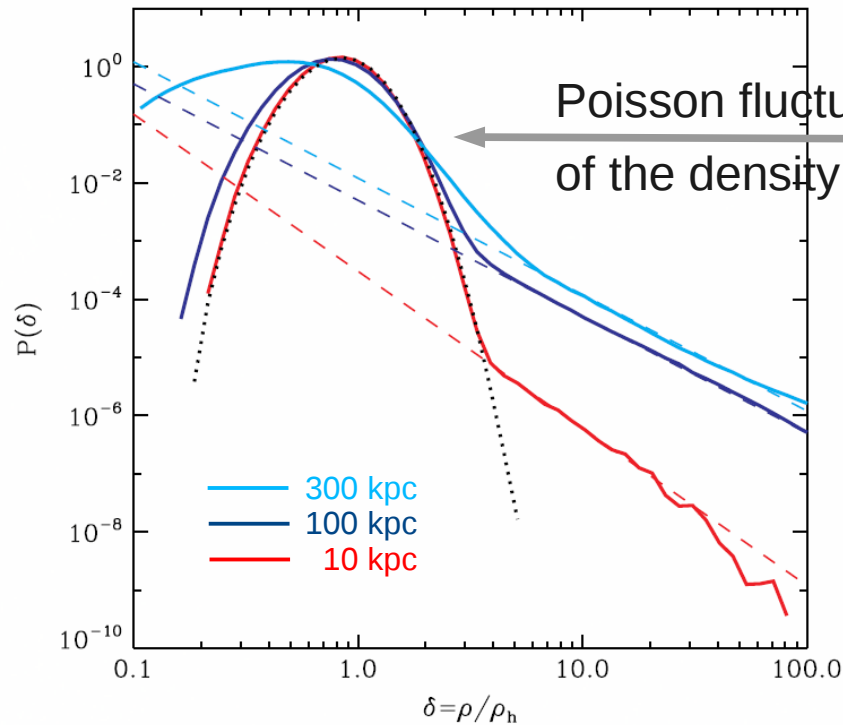
$$P(\rho; r) = \frac{f_s}{\sqrt{2\pi\Delta^2}} \frac{1}{\rho} \exp\left\{-\frac{1}{2\Delta^2} \left[\ln\left(\frac{\rho}{\rho_h} e^{\Delta^2/2}\right)\right]^2\right\} + (1 - f_s) \frac{1 + \alpha(r)}{\rho_h} \Theta(\rho - \rho_h) \left(\frac{\rho}{\rho_h}\right)^{-(2+\alpha)}$$

Kamionkowski, Koushiappas & Kuhlen (2010)

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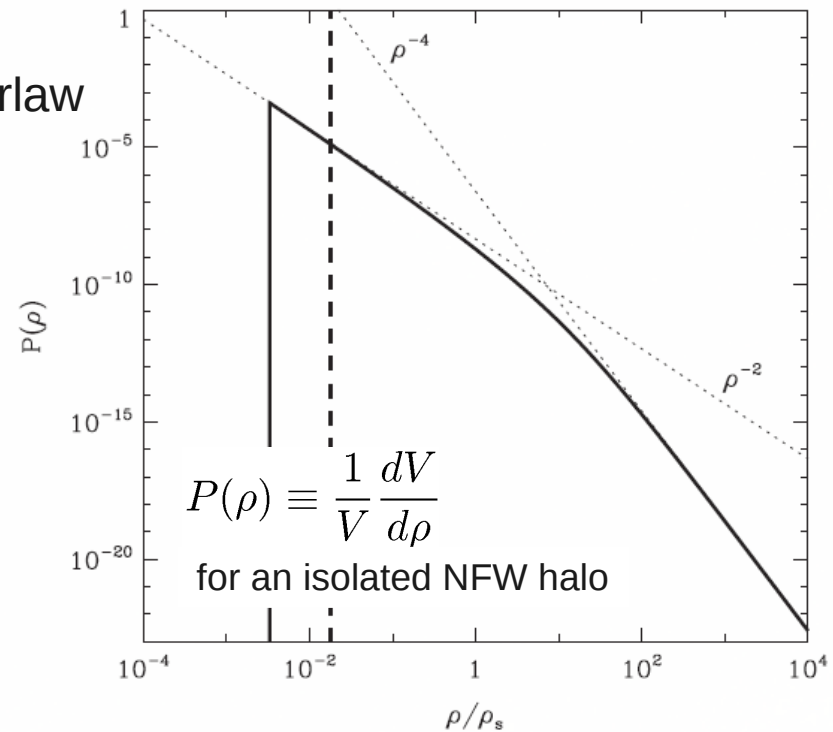
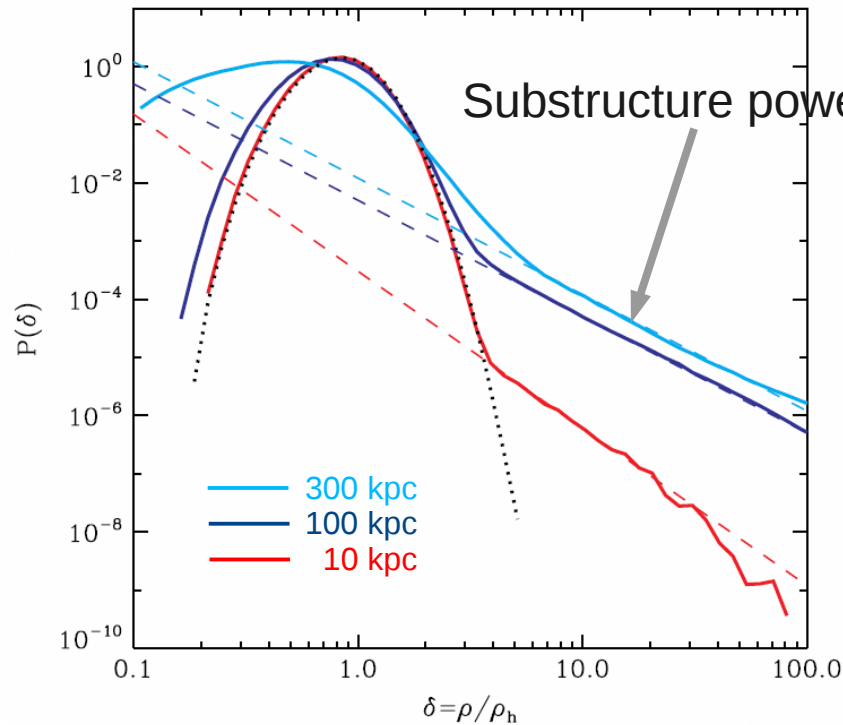
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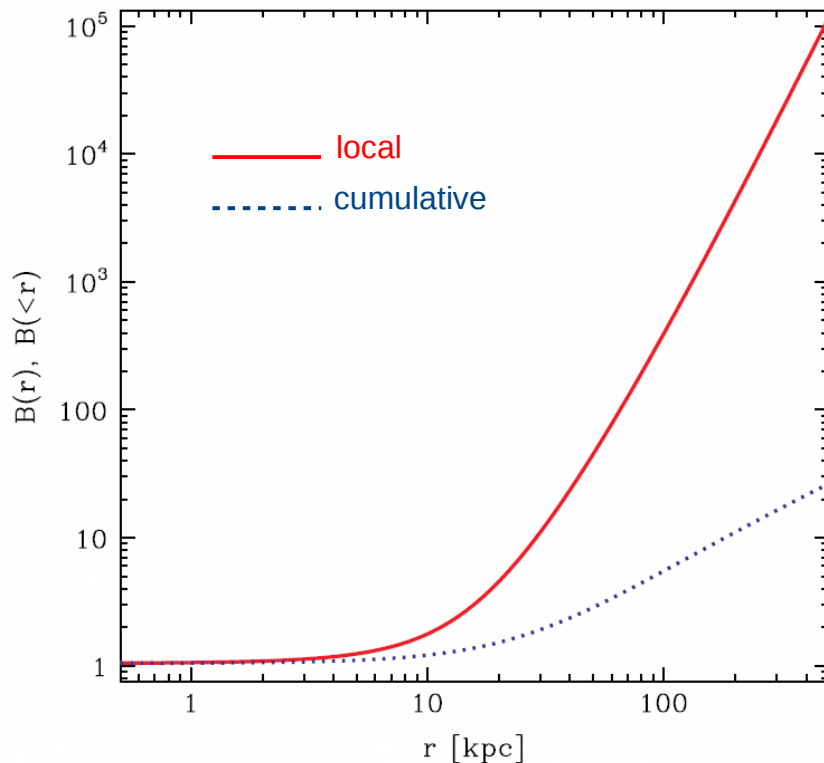
Kamionkowski, Koushiappas & Kuhlen (2010)

Substructure Boost: Density Distribution Approach

Use this distribution to calculate a boost factor as a function of radius.

$$B(r) = \frac{\int \rho^2 dV}{\int [\bar{\rho}(r)]^2 dV} = \int_0^{\rho_{\max}} P(\rho, r) \frac{\rho^2}{[\bar{\rho}(r)]^2} d\rho,$$

$$= f_s e^{\Delta^2} + (1 - f_s) \frac{1 + \alpha}{1 - \alpha} \left[\left(\frac{\rho_{\max}}{\rho_h} \right)^{1-\alpha} - 1 \right].$$



Depends on ρ_{\max} , which is set by the halo collapse epoch:

$$\rho_{\max} = 80 \text{ GeV cm}^{-3} \left(\frac{z_c}{40} \right)^3 \left(\frac{c}{3.5} \right) \frac{f(3.5)}{f(c)}$$

[Here c is the **natal** concentration!]

The biggest uncertainty is in f_s and its GC radius dependence.

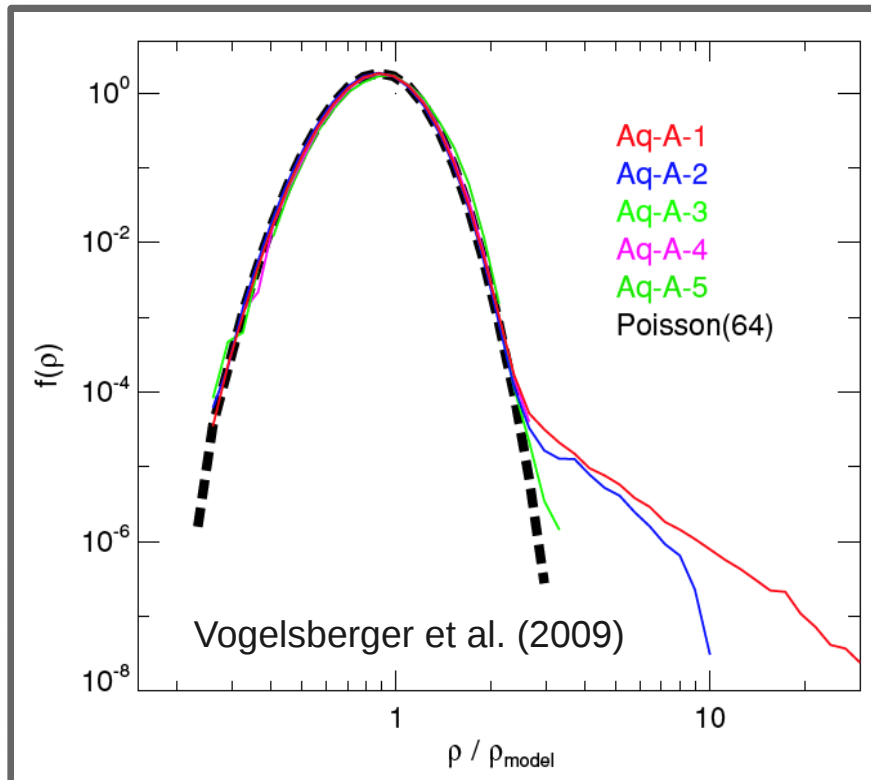
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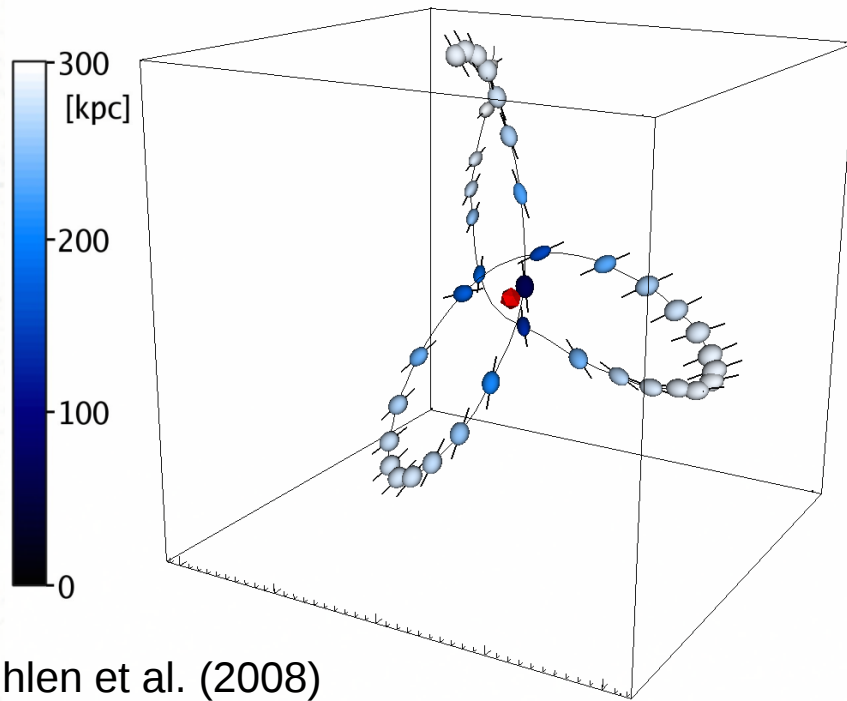
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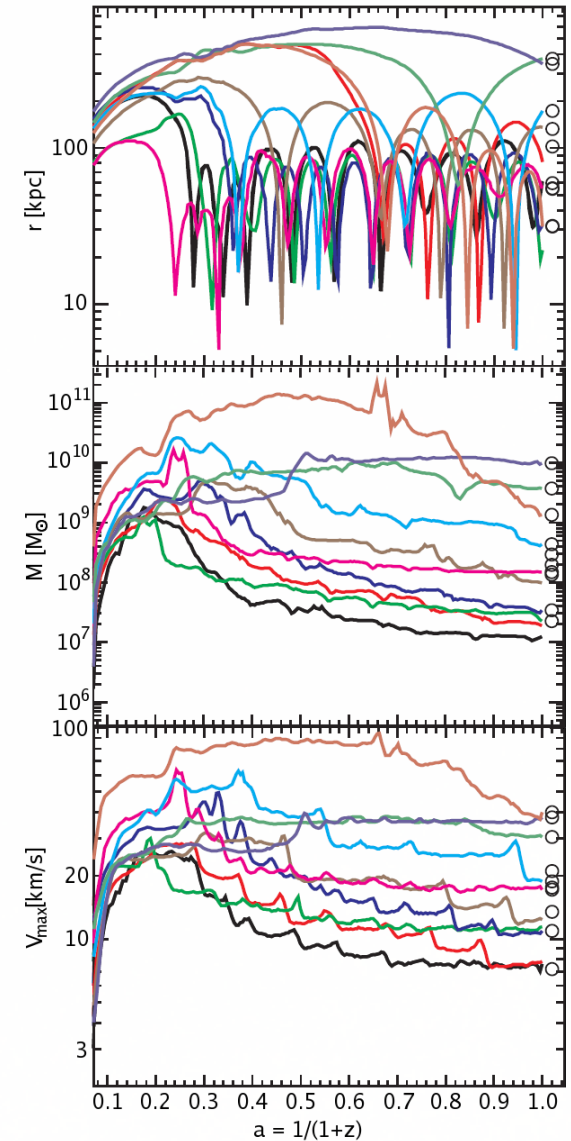
Kamionkowski, Koushiappas & Kuhlen (2010)

Tidal Interactions with the Host Halo

- Subhalos orbit through host halo and are subject to tidal interactions.
- Strongest during peri-center passage.
- Tidal mass loss from outside in.
- Diverse amount of tidal mass loss.

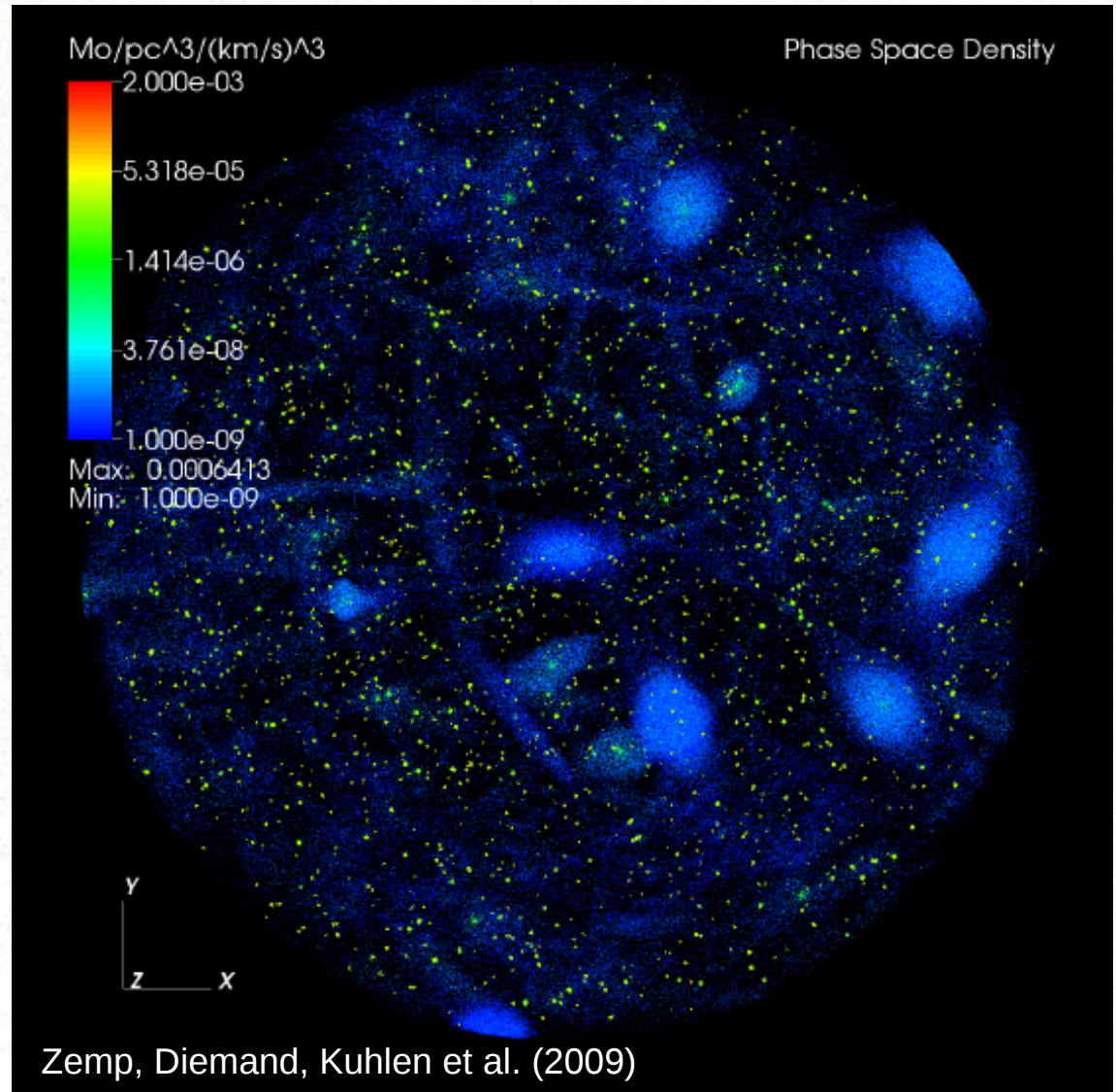


Diemand, Kuhlen, et al. (2008)



Velocity Space Substructure

When viewed in **phase-space-density**, many additional unbound substructures become apparent: dark matter tidal streams from disrupted subhalos.

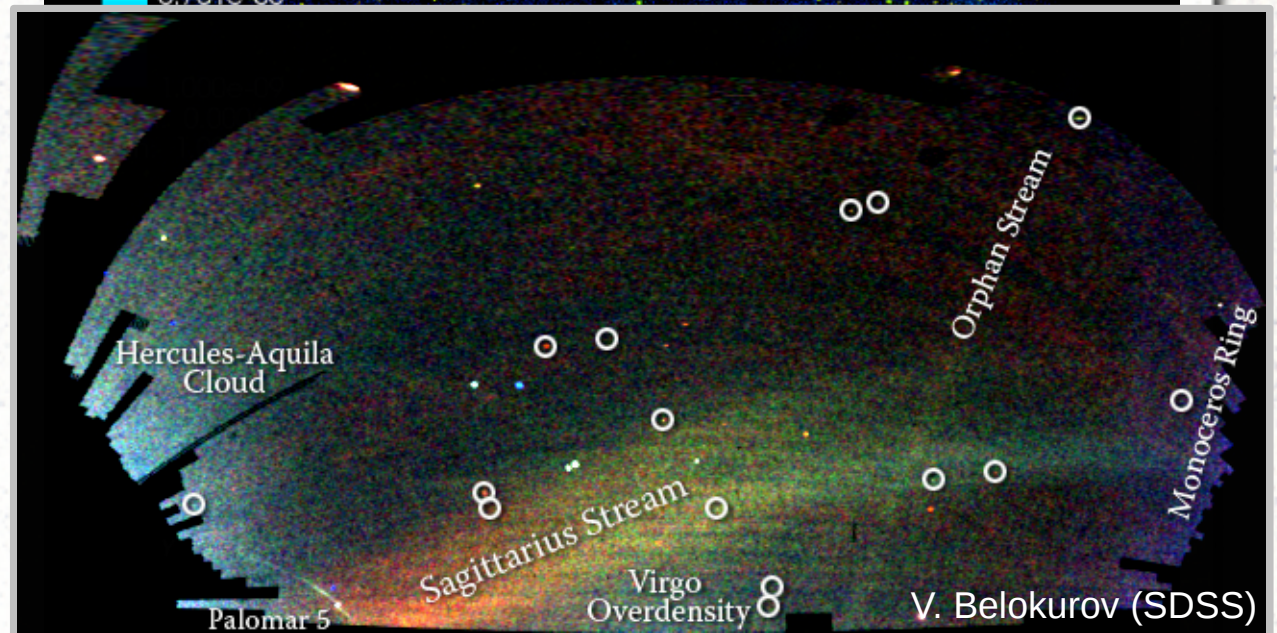
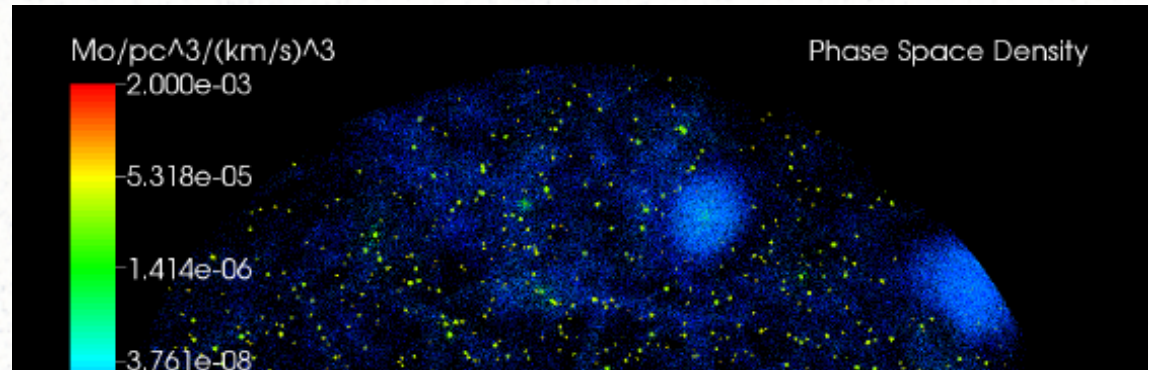


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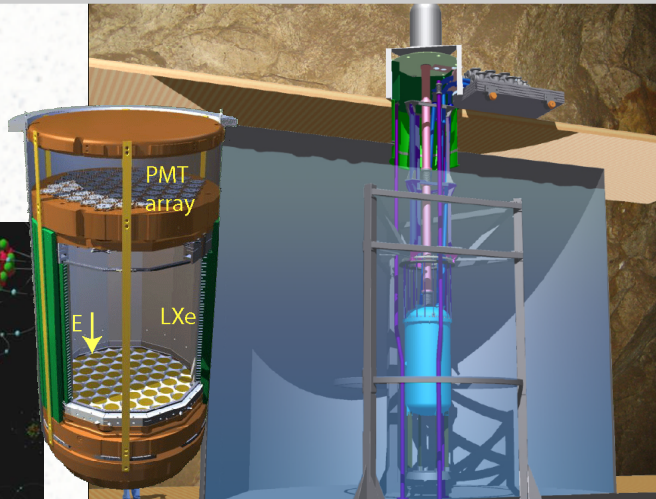
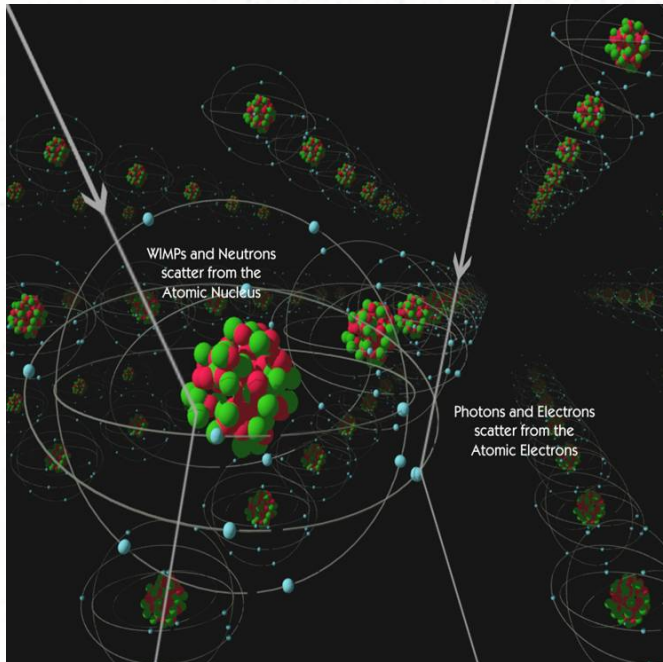
Direct counterparts to the stellar streams from disrupted satellites (e.g. SDSS Field of Streams).

In the future will there be a Missing Streams Problem?

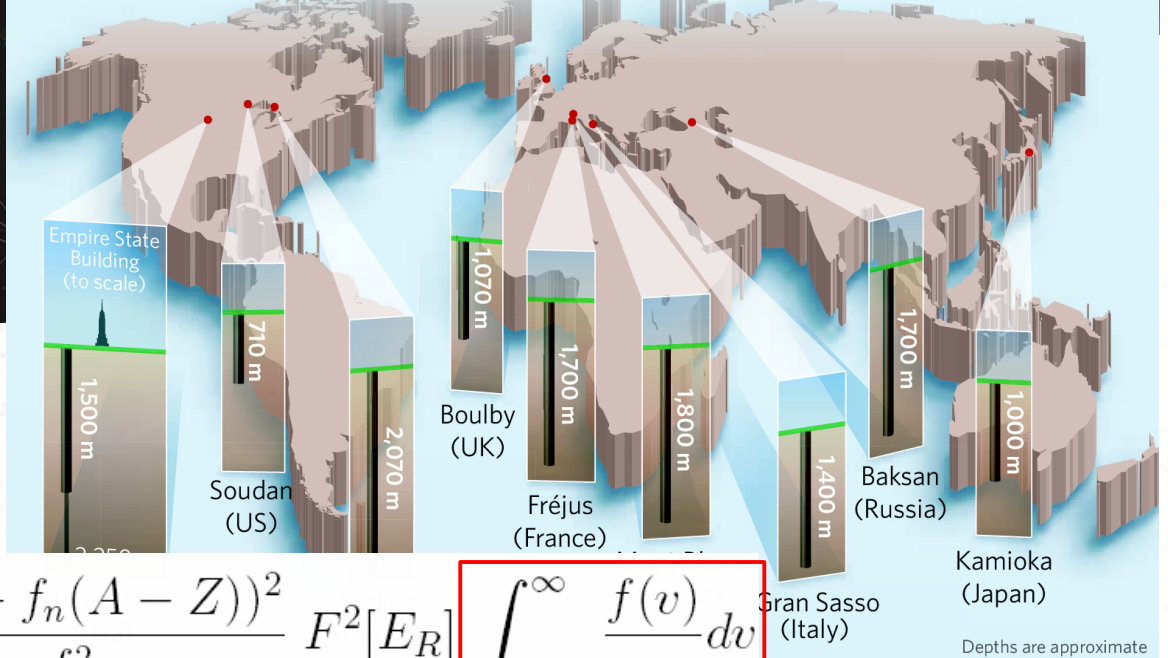


Zemp, Diemand, Kuhlen et al. (2009)

Velocity Space Substructure and Direct Detection

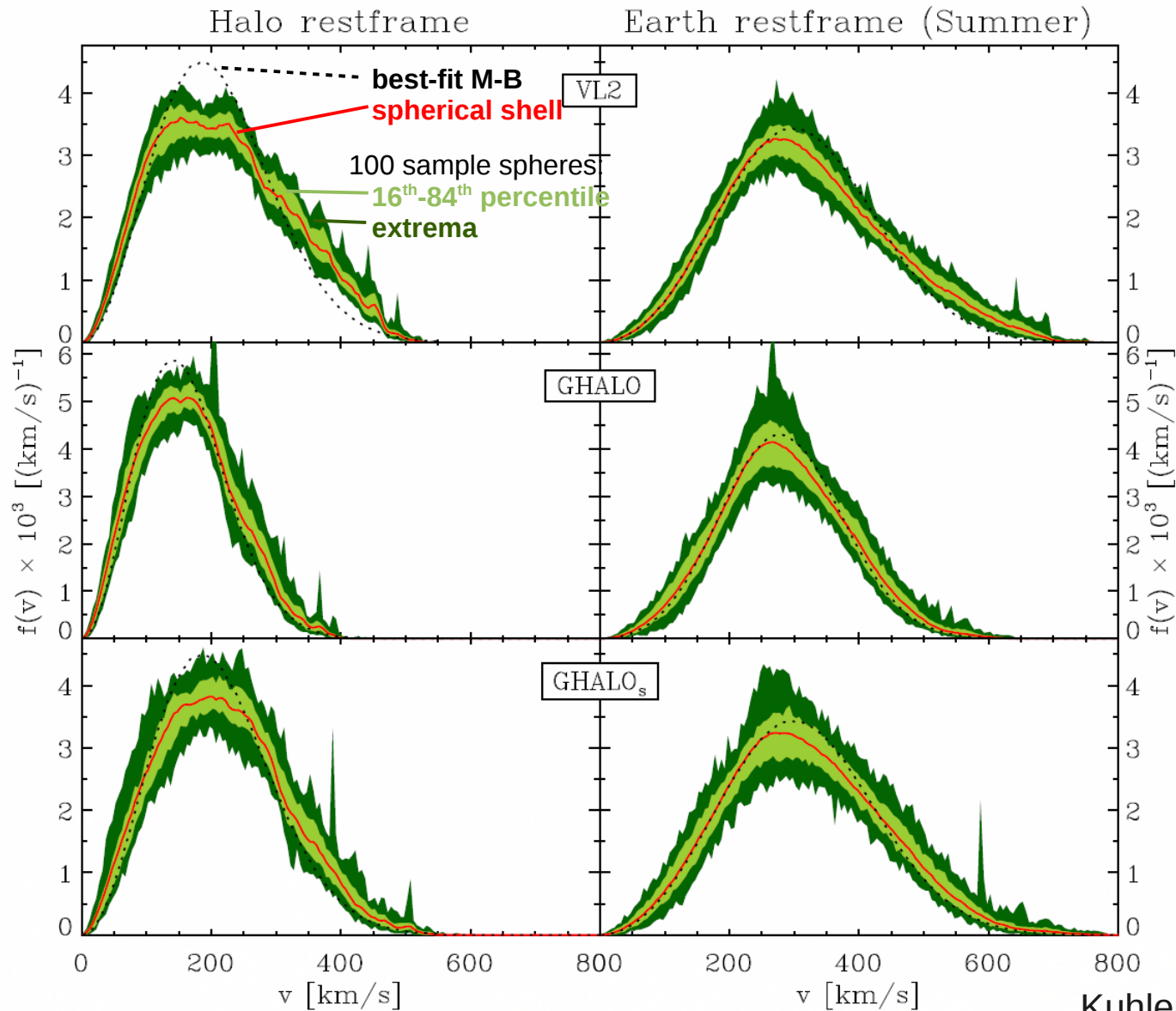


UNDERGROUND LABS AROUND THE WORLD



$$\frac{dR}{dE_R} = N_T M_N \frac{\rho_\chi \sigma_n}{2m_\chi \mu_{ne}^2} \frac{(f_p Z + f_n (A - Z))^2}{f_n^2} F^2[E_R] \int_{\beta_{min}}^{\infty} \frac{f(v)}{v} dv$$

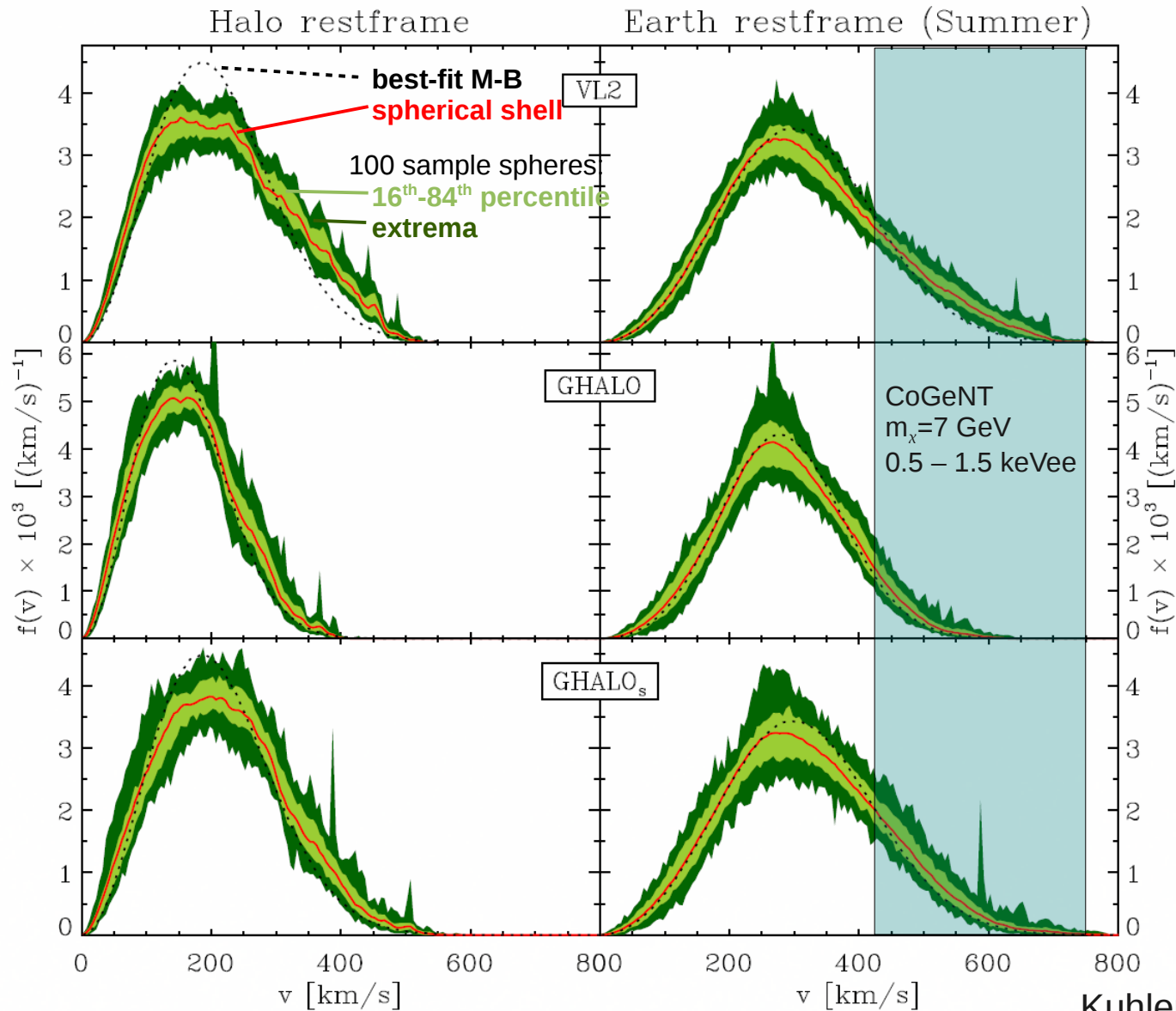
Velocity Space Substructure and Direct Detection



Kuhlen et al. (2010)

See also: Hansen et al. (2005), Vogelsberger et al. (2009)

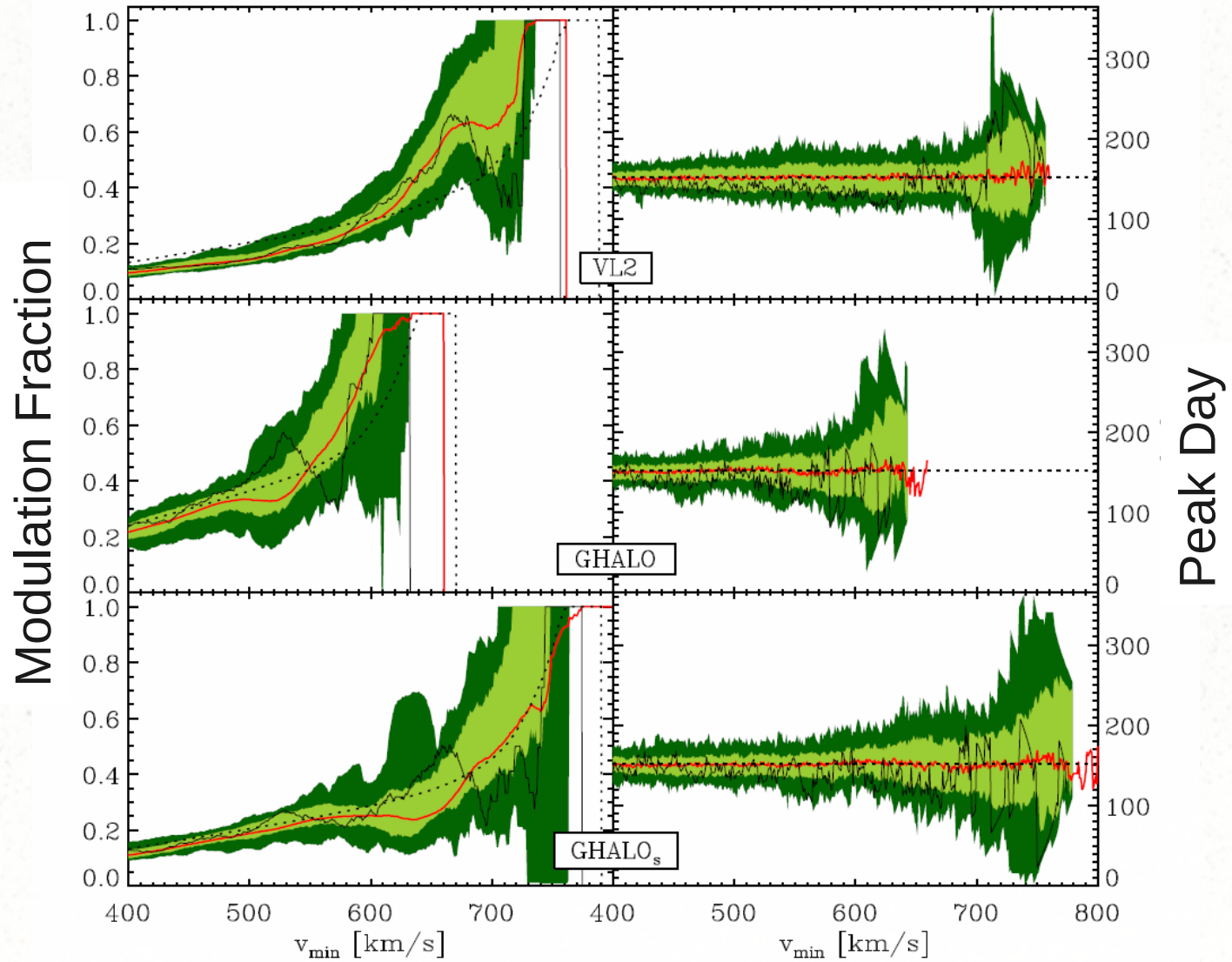
Velocity Space Substructure and Direct Detection



Kuhlen et al. (2010)

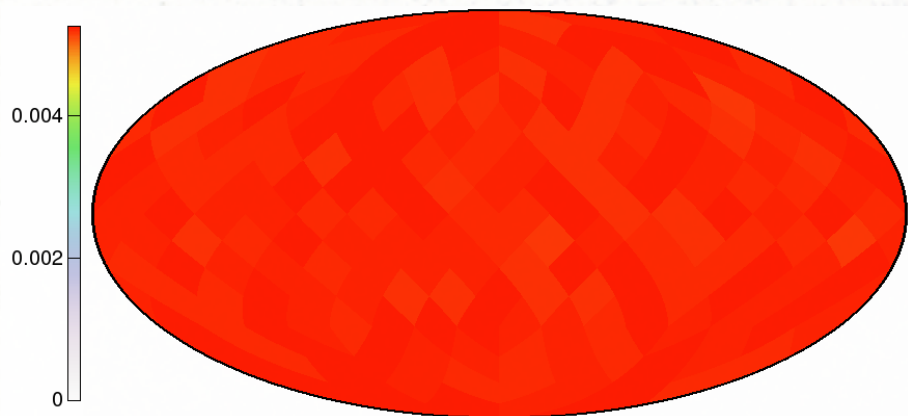
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Velocity Space Substructure and Direct Detection

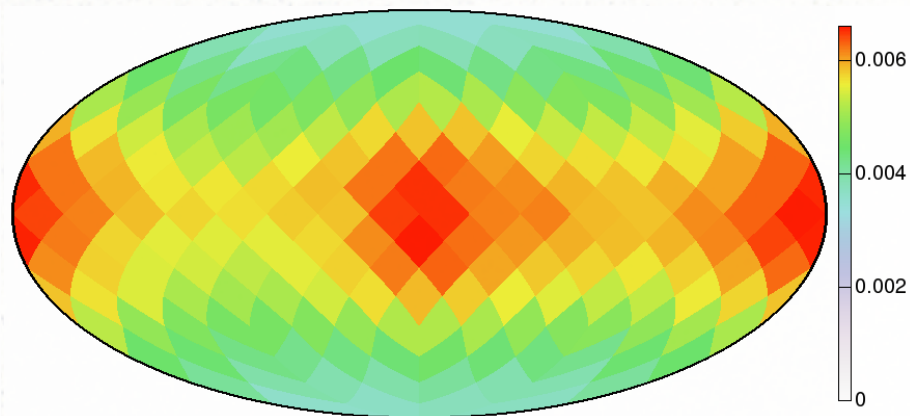


In Halo Restframe ($v_{\min} = 0$ km/s)

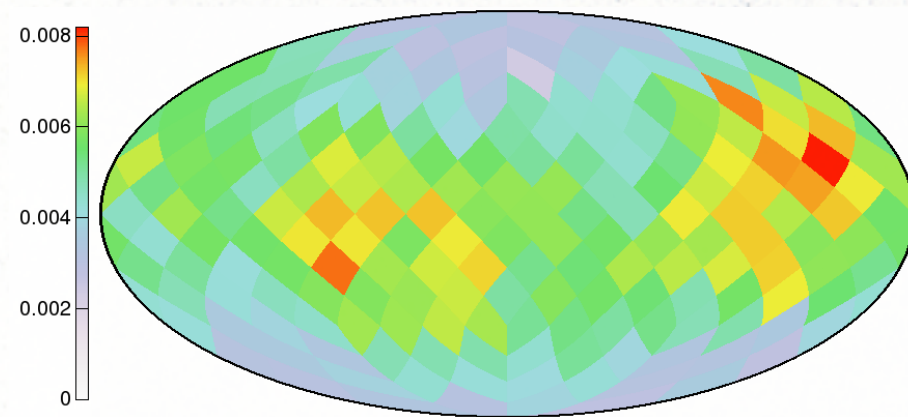
Maxwell-Boltzmann (isotropic)



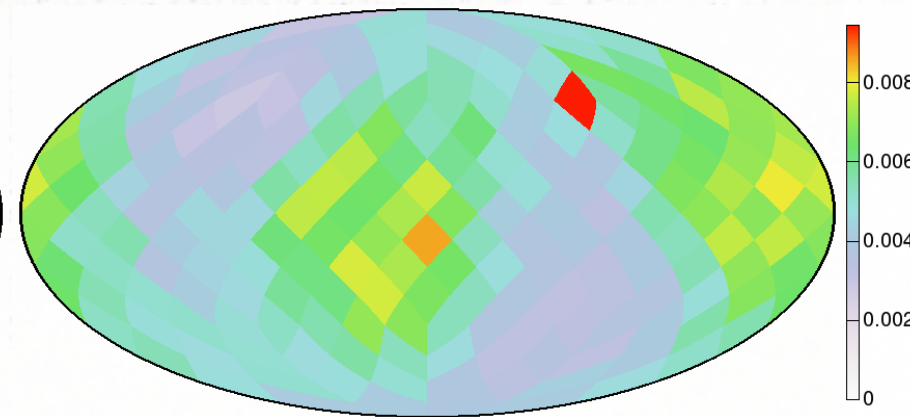
Spherical Shell ($8 \text{ kpc} < R < 9 \text{ kpc}$)



Sample Sphere #001

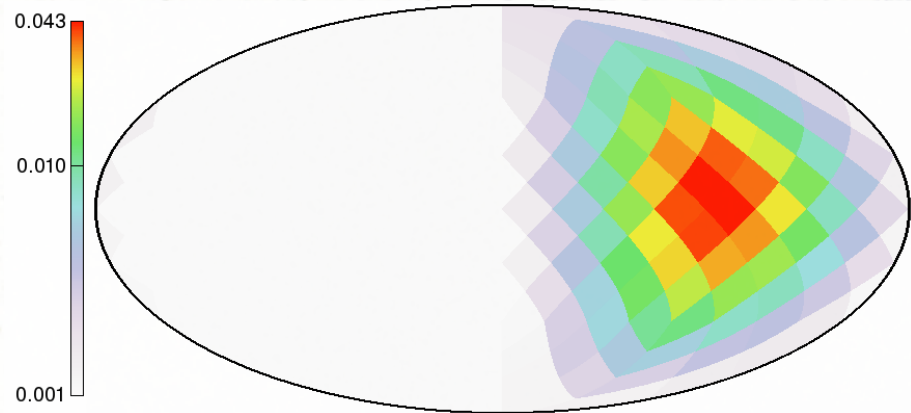


Sample Sphere #004 (containing a subhalo)

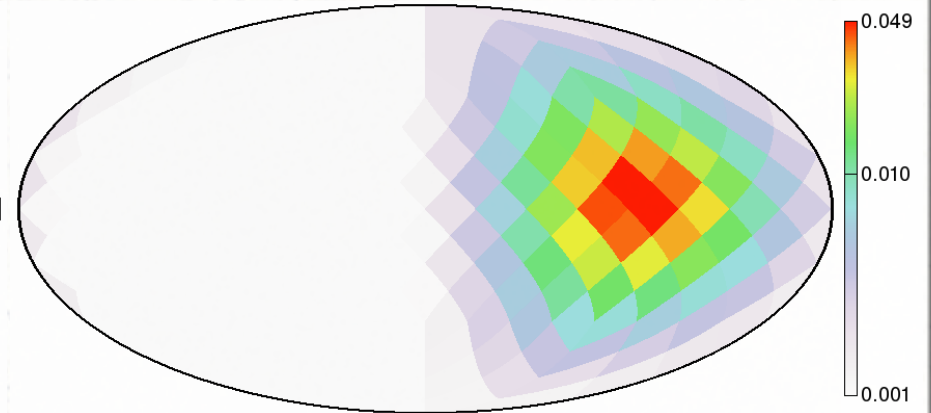


In Earth Restframe ($v_{\min} = 0$ km/s)

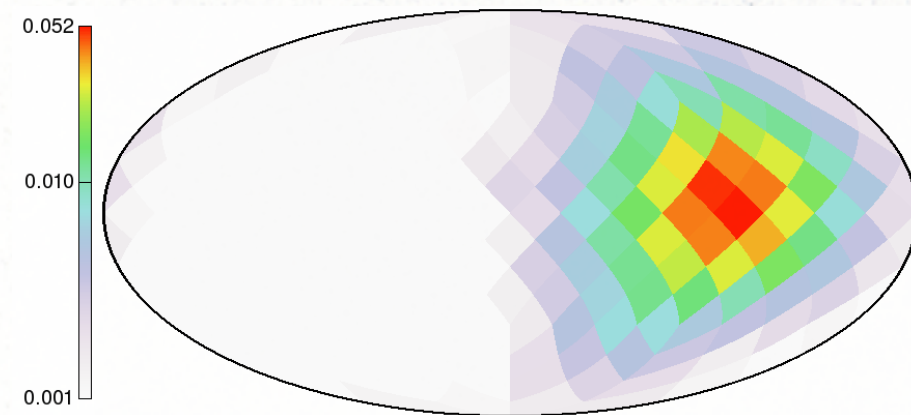
Maxwell-Boltzmann (isotropic)



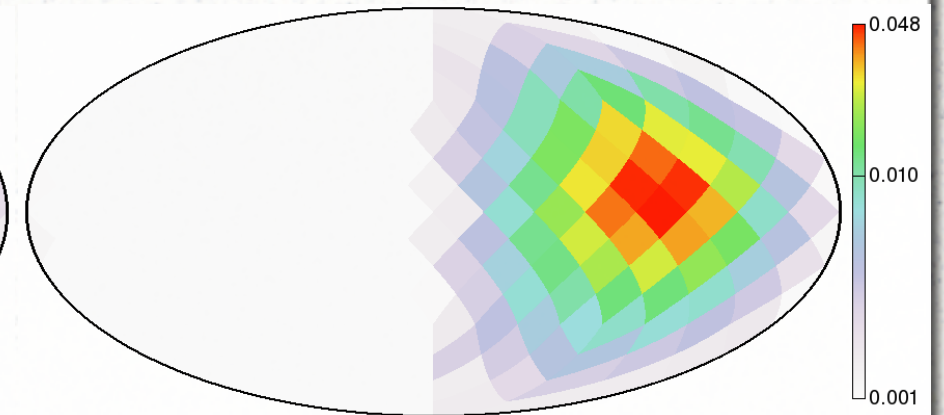
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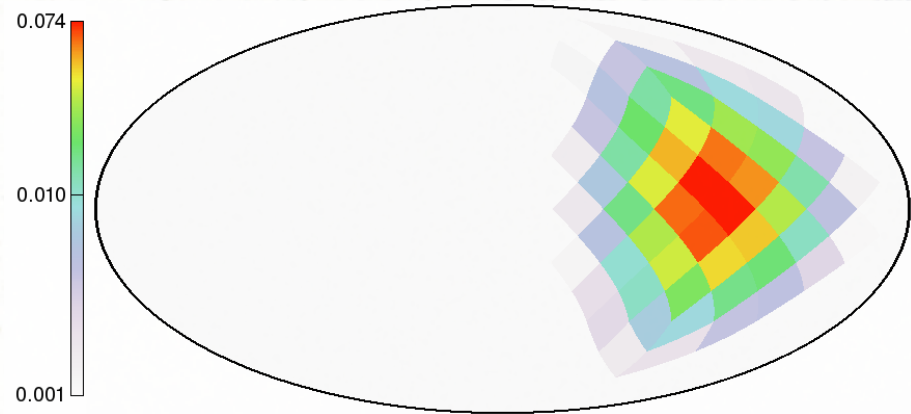


Sample Sphere #004 (containing a subhalo)

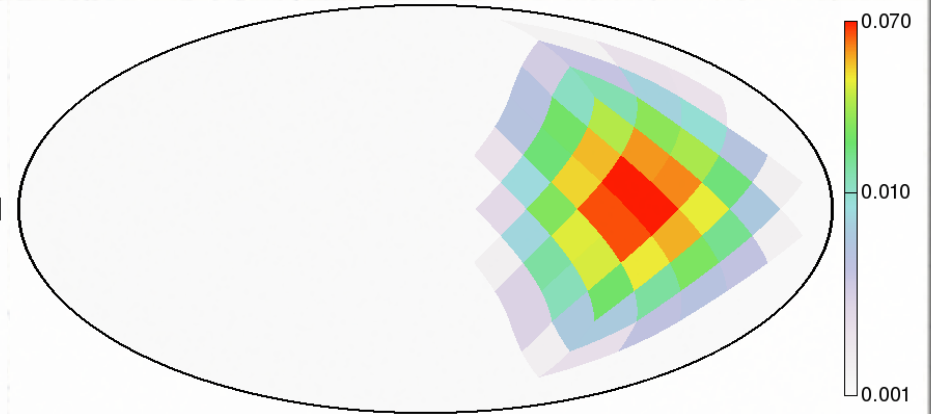


In Earth Restframe ($v_{\min} = 500$ km/s)

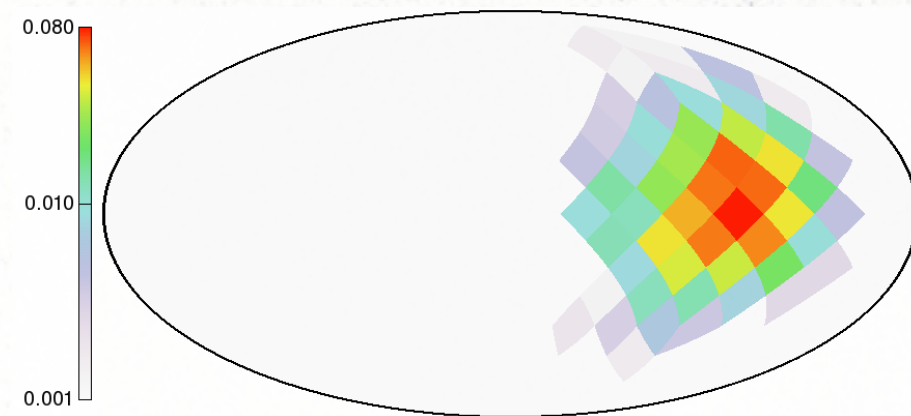
Maxwell-Boltzmann (isotropic)



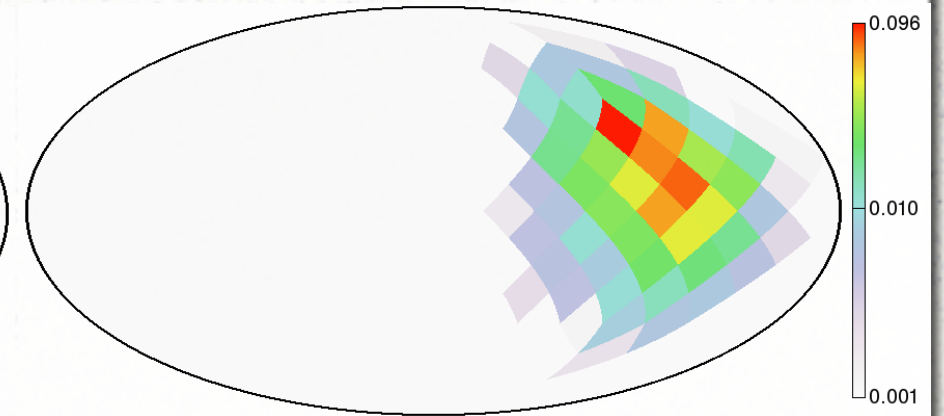
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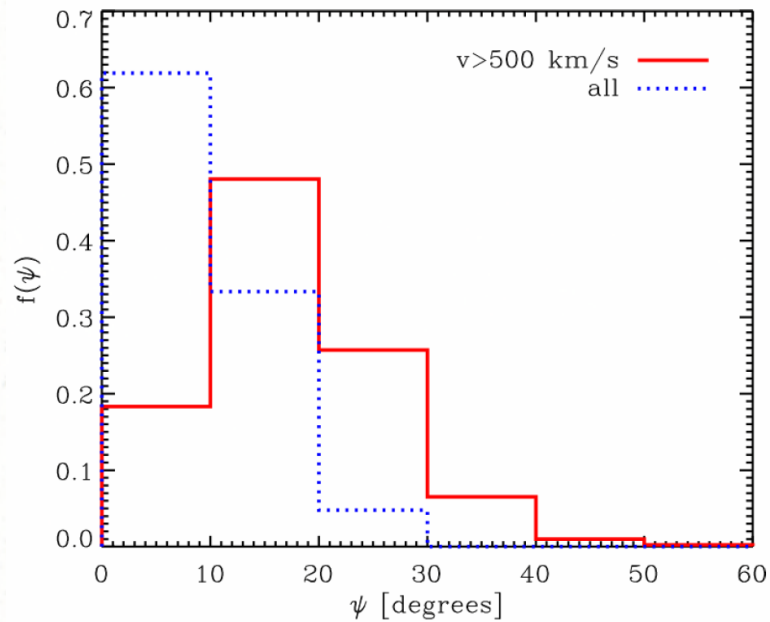
Sample Sphere #001



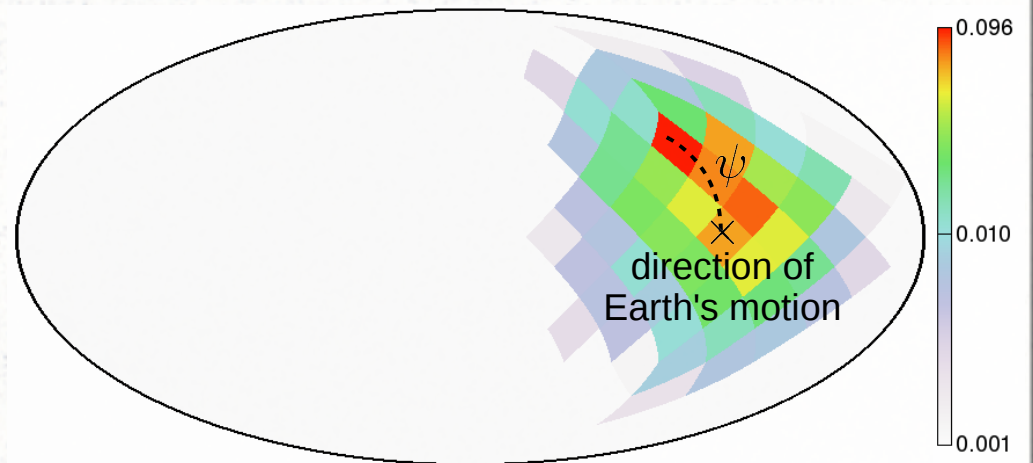
Sample Sphere #004 (containing a subhalo)



Hotspot Direction



Sample Sphere #004 (containing a subhalo)



At $v_{\min} = 500$ km/s the hotspot is more than 10° away from the direction of Earth's motion in $\sim 80\%$ of all cases!

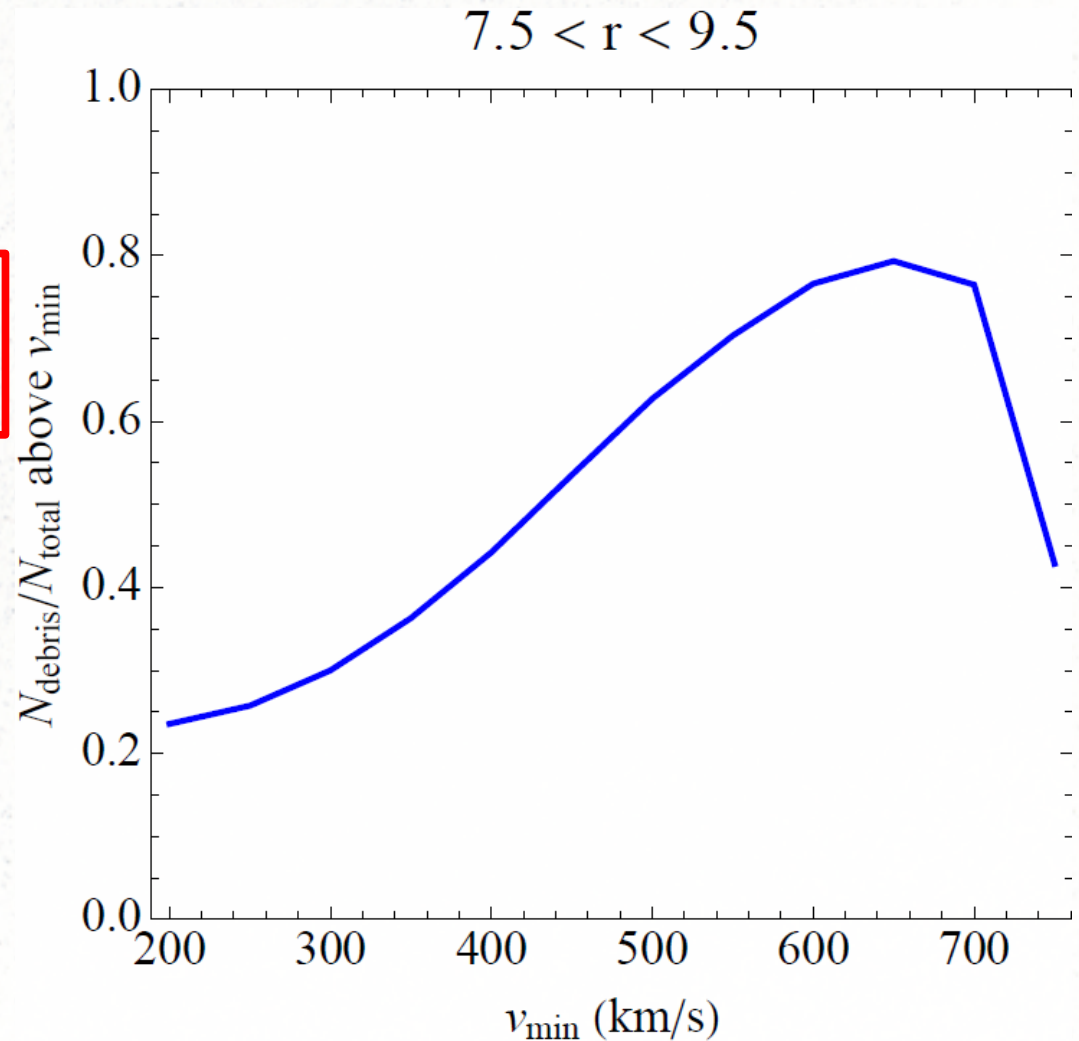
Velocity Space Substructure and Direct Detection

Recent work with M. Lisanti – **PRELIMINARY!!**

Study of the origin of the velocity space structure by tracking (sub)halo particles through time in the simulation.

Debris \equiv particles that were at some earlier time bound to a subhalo but are unbound at $z=0$.

1) The fraction of debris particles is quite large at high velocities!



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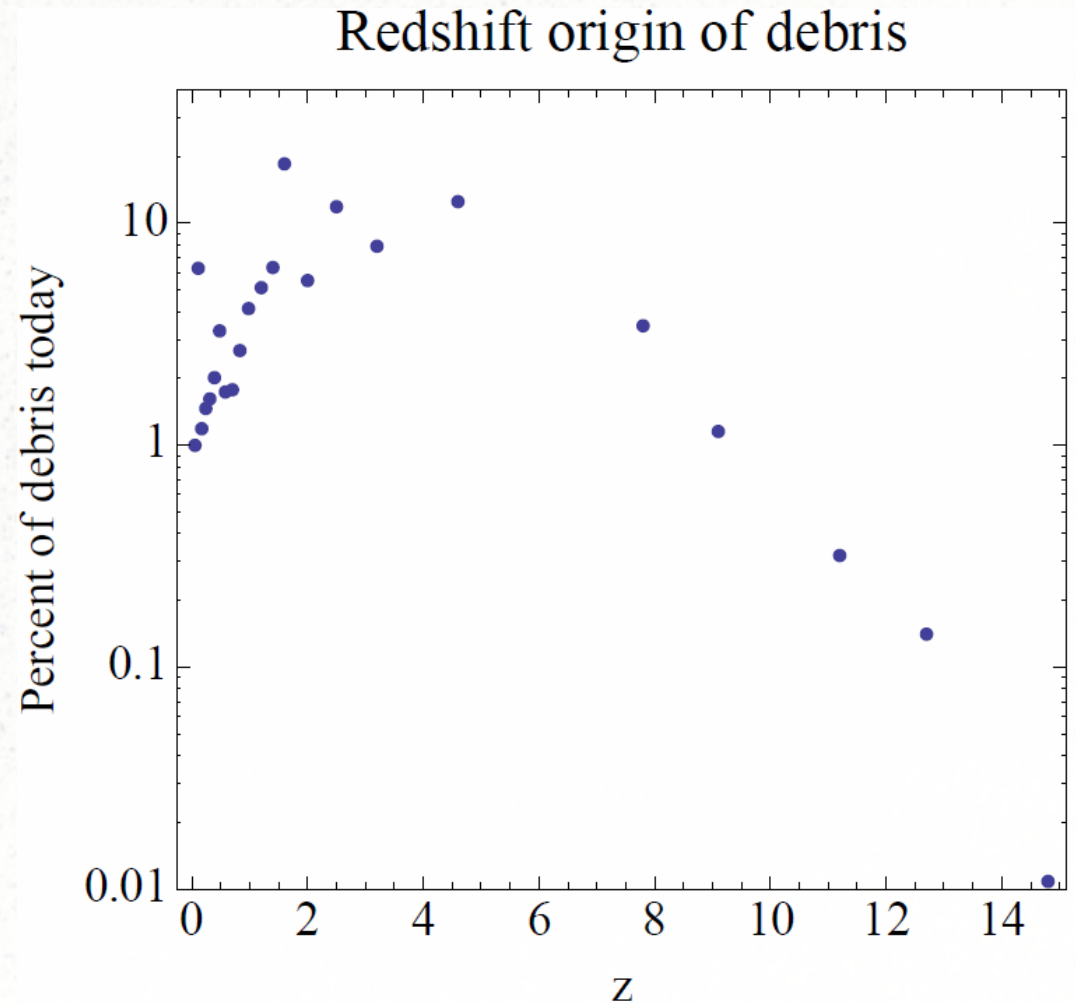
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2) Majority of debris is accreted after $z=3$ (in last 10 Gyr).



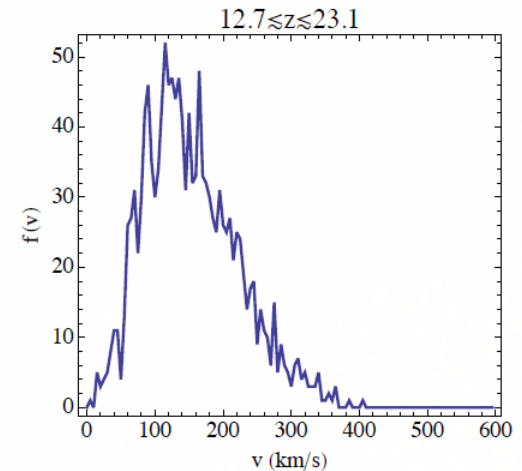
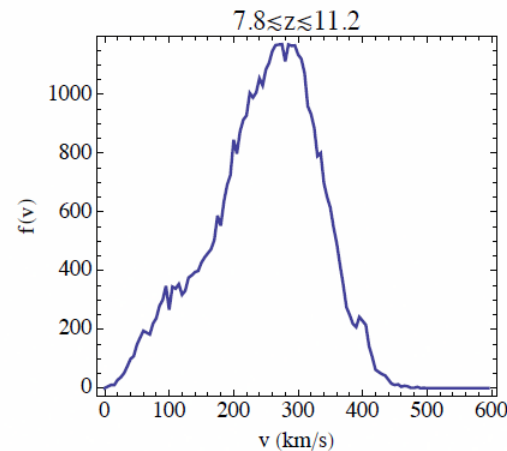
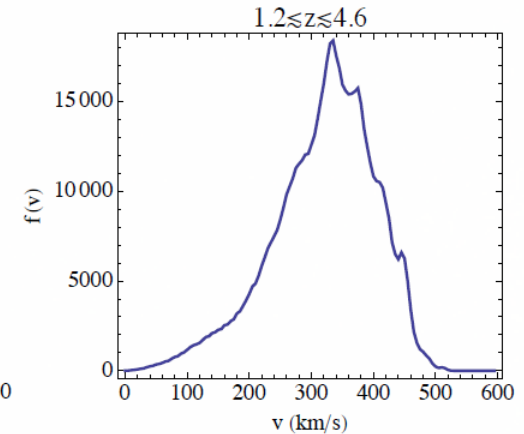
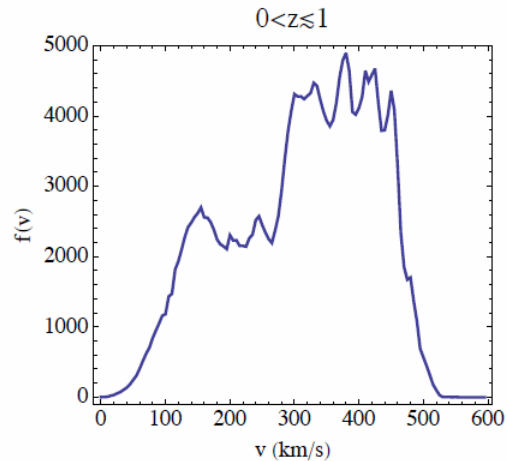
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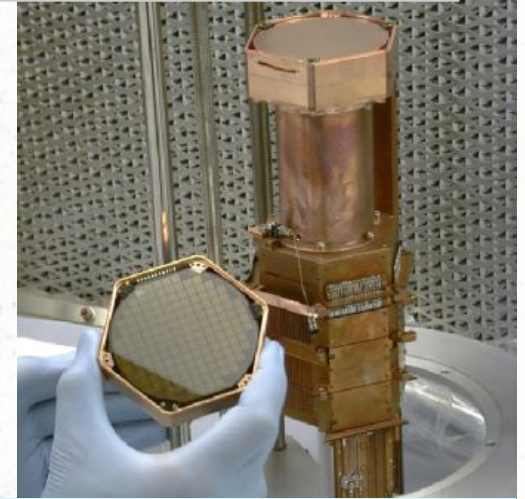
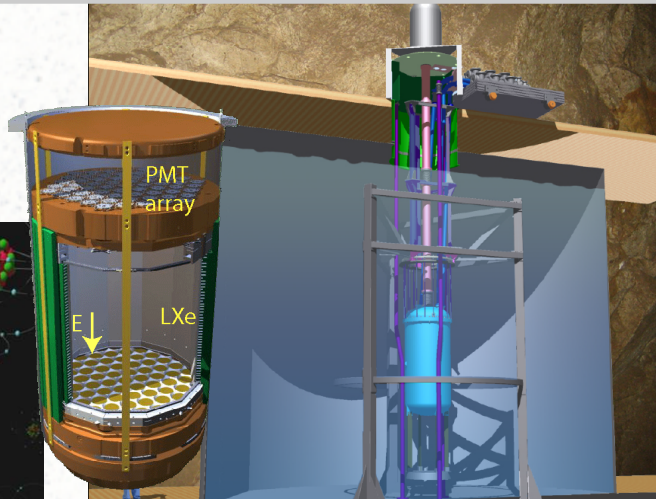
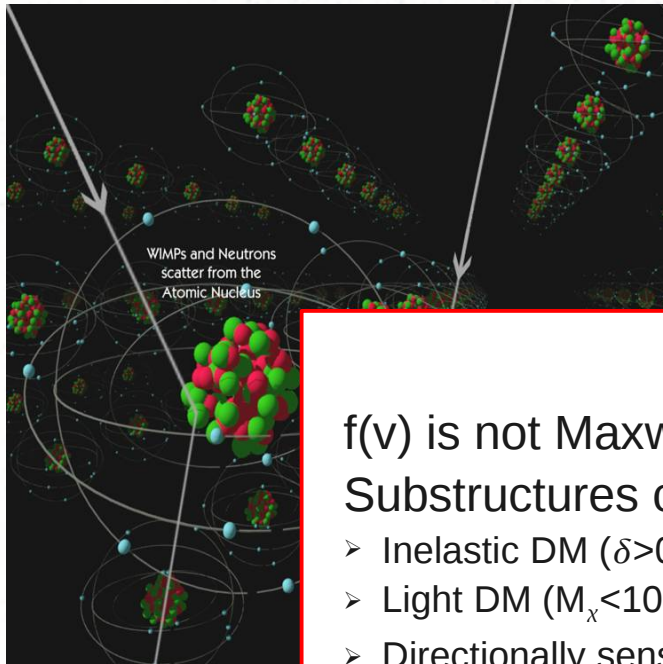
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Debris \equiv particles that were at some earlier time bound to a subhalo but are unbound at $z=0$.

- 1) The fraction of debris particles is quite large at high velocities!
- 2) Majority of debris is accreted after $z=3$ (in last 10 Gyr).
- 3) Early accreted material is virialized. Later material shows large departures from Maxwellian.



Velocity Space Substructure and Direct Detection

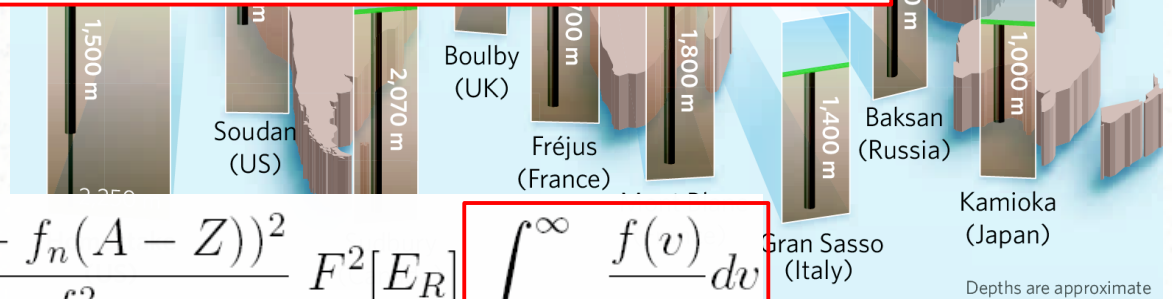


$f(v)$ is not Maxwellian.

Substructures can be important at high velocities!

- Inelastic DM ($\delta > 0$)
- Light DM ($M_\chi < 10 \text{ GeV}$)
- Directionally sensitive experiments often require high E_{recoil} , large β_{min} .

$$\beta_{\text{min}} = \sqrt{\frac{1}{2m_N E_R} \left(\frac{m_N E_R}{\mu} + \delta \right)}$$

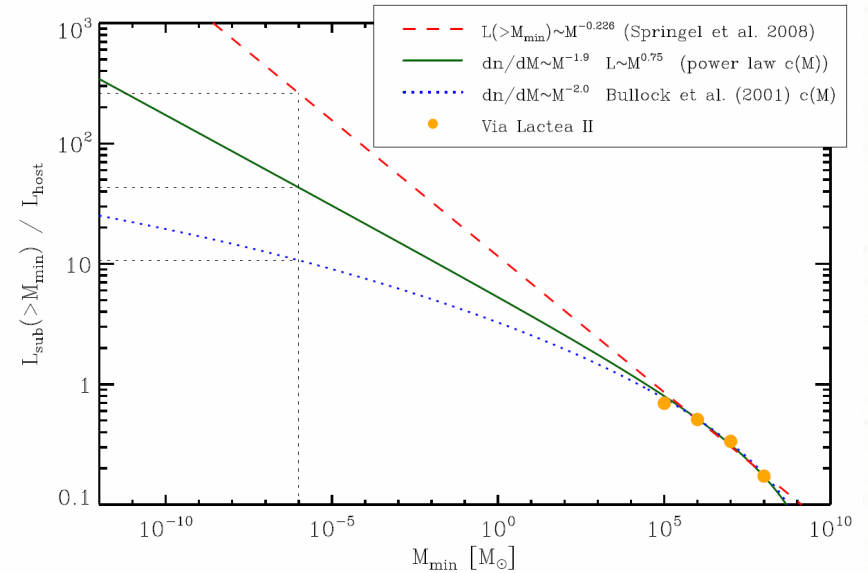


$$\frac{dR}{dE_R} = N_T M_N \frac{\rho_\chi \sigma_n}{2m_\chi \mu_{ne}^2} \frac{(f_p Z + f_n (A - Z))^2}{f_n^2} F^2[E_R] \int_{\beta_{\text{min}}}^{\infty} \frac{f(v)}{v} dv$$

Uncertainties and Questions

★ for CDM simulations

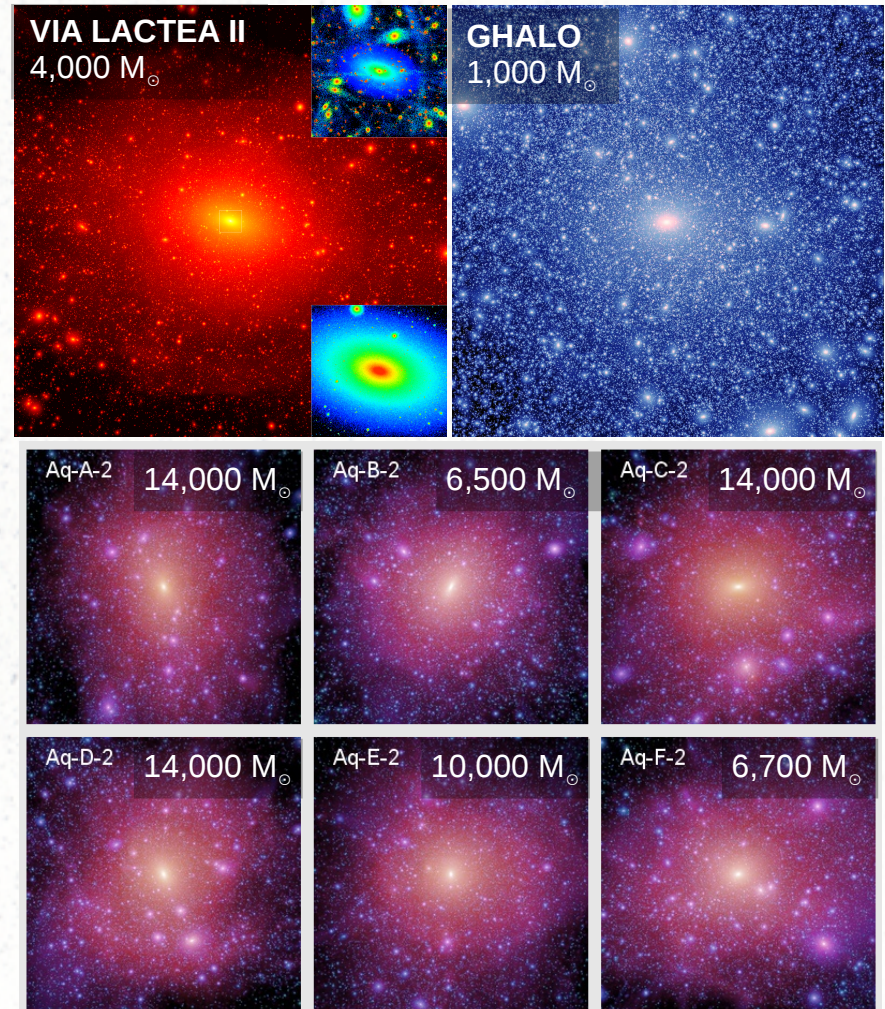
- extrapolation to lower masses
 - concentration-mass relation
 - tidal stripping & disruption
 - local phase-space structure



Uncertainties and Questions

★ for CDM simulations

- extrapolation to lower masses
 - concentration-mass relation
 - tidal stripping & disruption
 - local phase-space structure
- cosmic variance & cosmology
 - The 6 Aquarius host halos show considerable variation in merging history
 - σ_8 , n_s affects the typical subhalo collapse time and hence properties



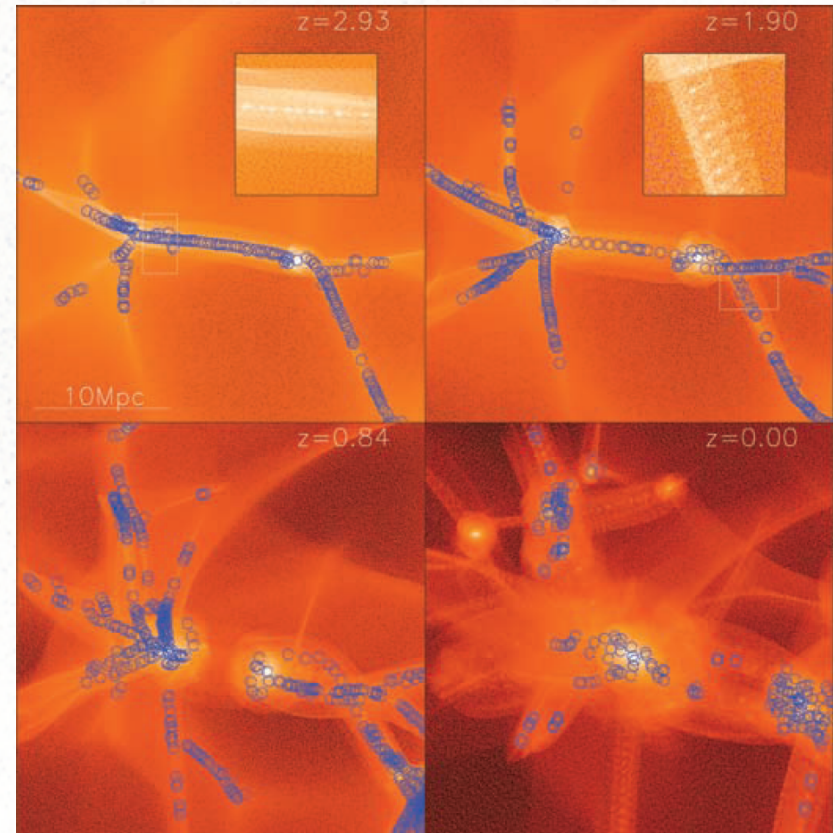
Uncertainties and Questions

★ for CDM simulations

- extrapolation to lower masses
- cosmic variance & cosmology

★ beyond “Cold Dark Matter”

- Warm (Tepid?), Self-interacting, ...
 - $P(k)$ cutoff in the scales sampled by particles. How to avoid spurious fragmentation?
 - Non-negligible thermal velocities in the IC's.
 - No longer collisionless: Monte-Carlo method for δv (Davé et al. 2001)
- Model dependent



Uncertainties and Questions

★ for CDM simulations

- extrapolation to lower masses
- cosmic variance & cosmology

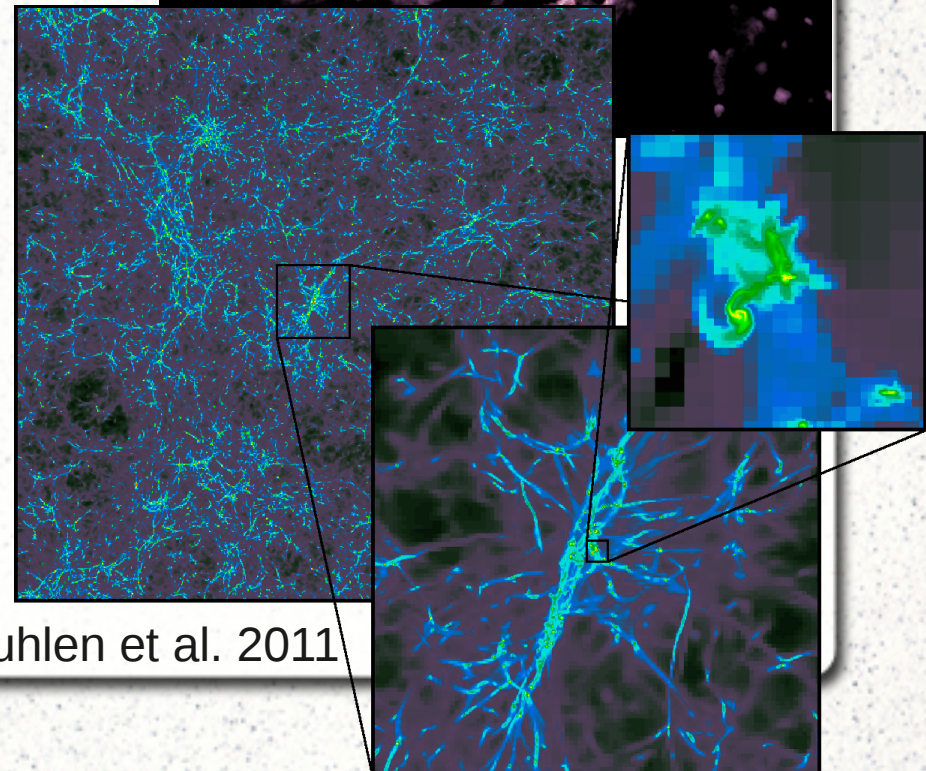
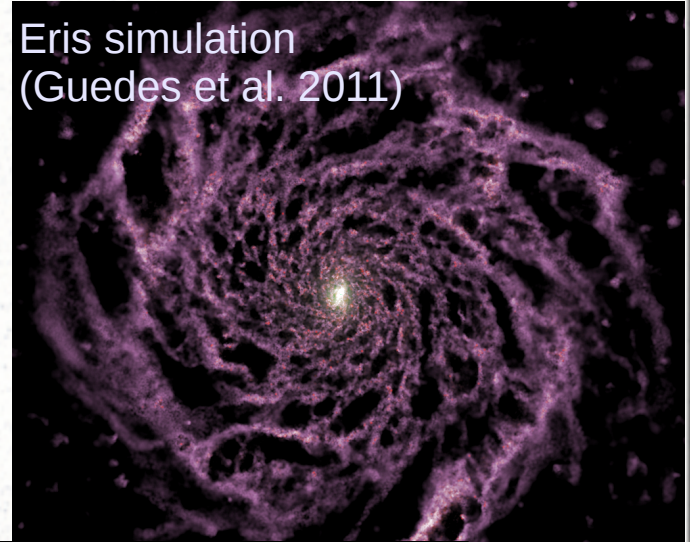
★ beyond “Cold Dark Matter”

- Warm (Tepid?), Self-interacting, ...
- Model dependent

★ baryonic physics!

- “So, when are you going to include baryons in your simulations?”
- Challenges:
 - expensive!
 - many ways to implement known important physics (star formation, supernova feedback, metal enrichment, interstellar chemistry, etc., etc.)
 - unknown important physics?

Eris simulation
(Guedes et al. 2011)



Kuhlen et al. 2011