

Dark matter freeze-in from semi-production

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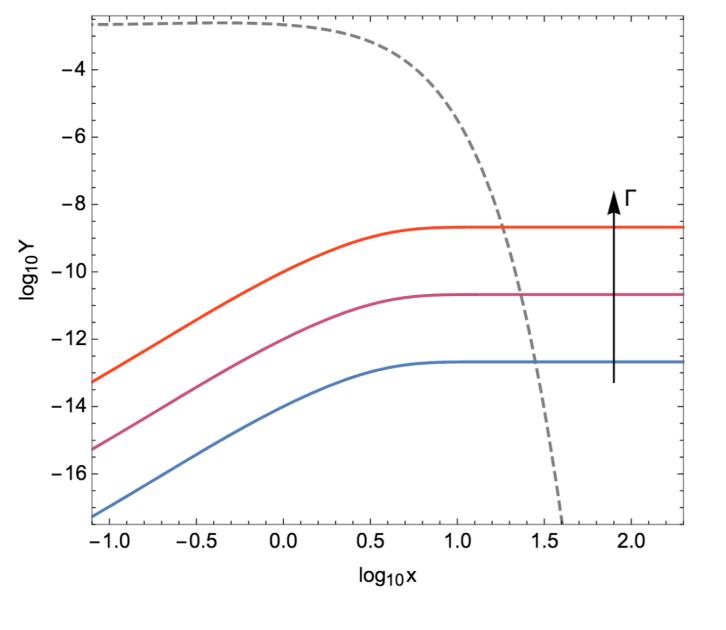
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Freeze-in production of DM

From 1706.07442



 $x = \frac{1}{T}$

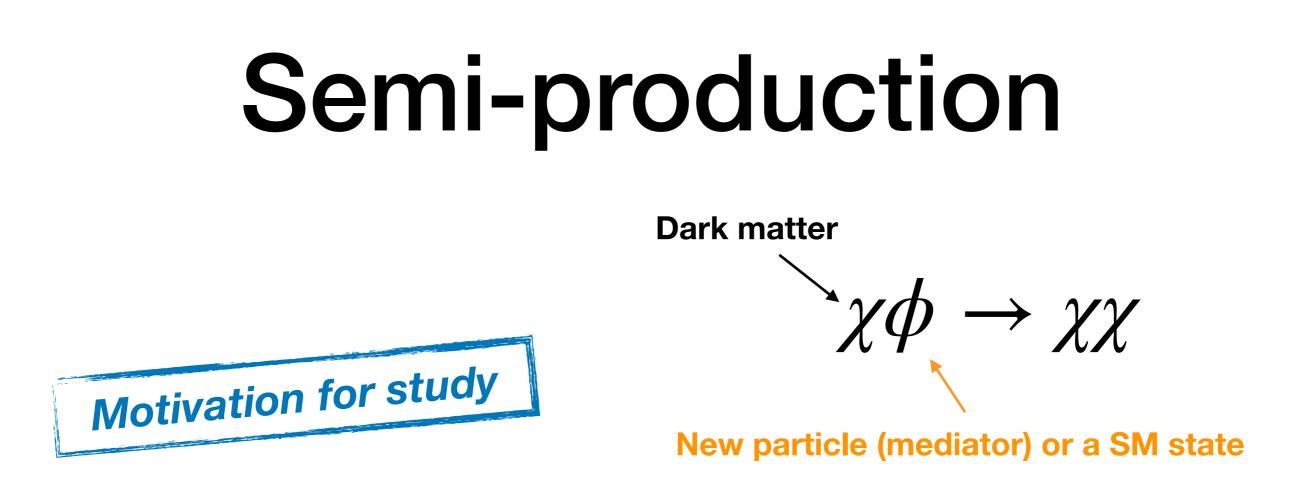
- Dark matter interacts feebly with SM (not in equilibirum) and it's initial concentration is tiny
- DM density freezes in when the density of SM is suppressed and the production is inefficient
- The production process can be a decay or annihilation

Typical annihilation cross section for freeze-in

$\langle \sigma v \rangle \lesssim 10^{-40} \, \mathrm{cm}^3 / s$

This is too weak for a near-future direct or indirect detection

IS IT POSSIBLE TO HAVE AN OBSERVABLE SIGNAL FROM FREEZE-IN?



What is different from the pair-annihilation freeze-in?

- The production rate is proportional to the DM density.
 Smaller initial abundance → larger cross section
- Semi-production modifies the energy of DM particles in a non-trivial way, so the temperature evolution can affect the relic density

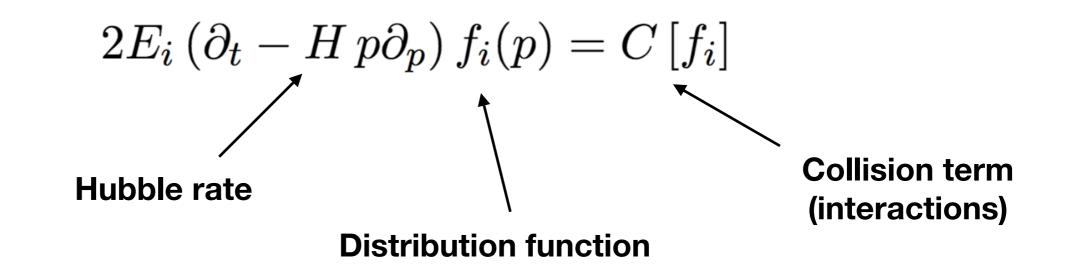
Toy model

We start the investigation with a simple toy model

$$\mathcal{L}_{int} = \mathcal{L}_{SM} + \mathcal{L}_{\phi-SM} + \frac{\lambda}{2}\phi\left(\chi^3 + (\chi^*)^3\right)$$
Z3 symmetry

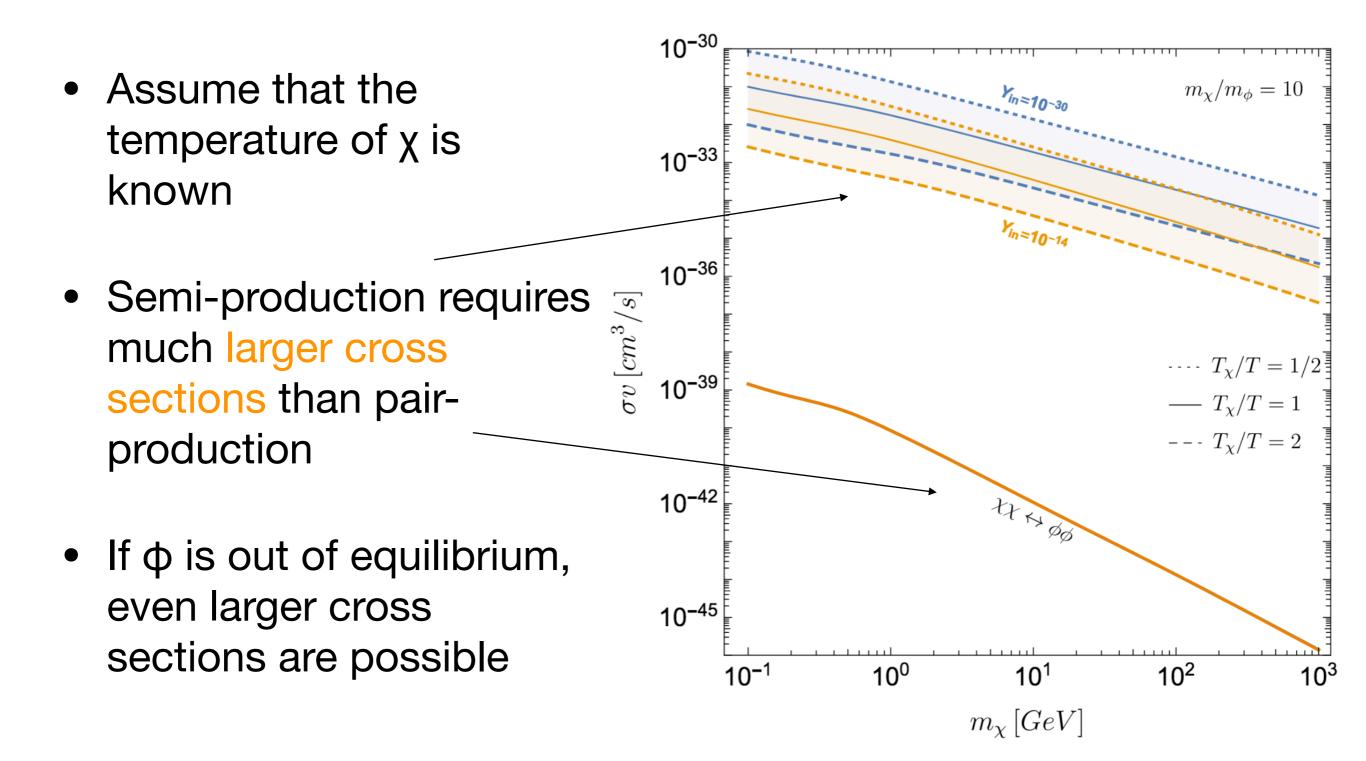
- Assume that φ is in equilibrium with SM and χ has some tiny initial abundance (e.g. from the gravitational production)
- The abundance of χ is determined by the Boltzmann equation (we search for the model parameters that can reproduce the observed relic density of DM)

Boltzmann equation (BE)



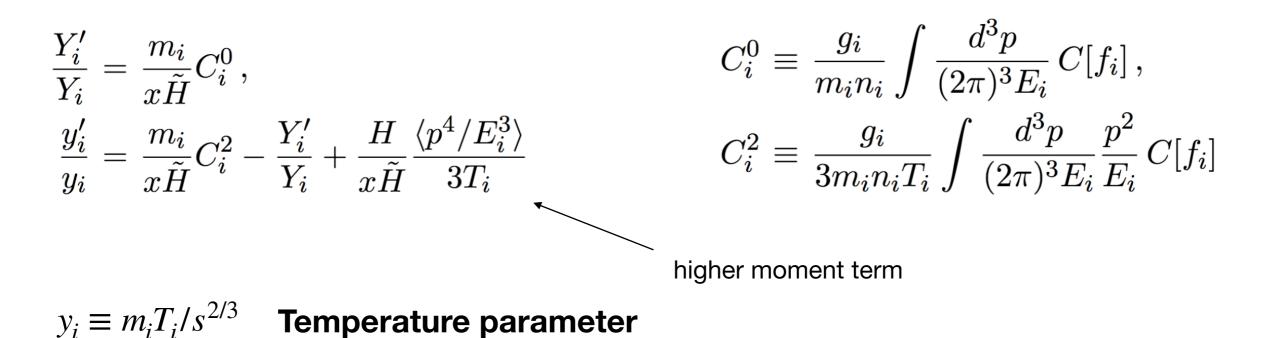
Integrating over the phase-space of the DM particle (and introducing comoving variables) one gets an equation for the abundance

Toy model



Coupled system of BEs

- In reality we don't know the temperatures of χ and φ apriori, but they can be important!
- We assume that φ and χ are both self-thermalized (have an equilibrium shape) and solve for their density and temperature



Setting

 We now consider a more detailed example model, where φ is a scalar singlet coupled to the Higgs doublet

Higgs portal interactions

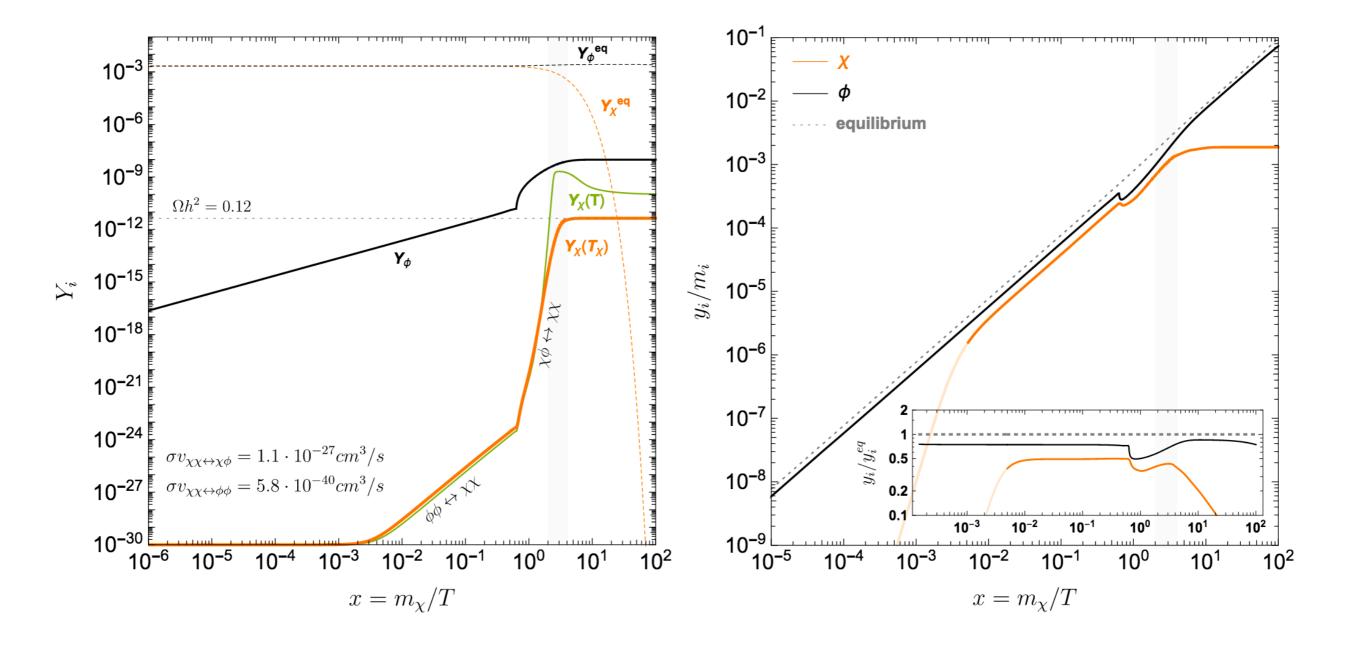
$$\mathcal{L}_{\phi-SM} = \overline{A\phi H^{\dagger}H + \frac{\lambda_{h\phi}}{2}\phi^{2}H^{\dagger}H - \mu_{h}^{2}H^{\dagger}H + \frac{\lambda_{h}}{2}(H^{\dagger}H)^{2}}$$

$$\mathcal{L}_{DS} = \frac{\mu_{\phi}^{2}}{2}\phi^{2} + \frac{\mu_{3}^{2}}{3!}\phi^{3} + \frac{\lambda_{\phi}}{4!}\phi^{4} + \mu_{\chi}^{2}\chi^{*}\chi + \frac{\lambda_{\chi}}{4}(\chi^{*}\chi)^{2}$$

$$+ \frac{\lambda_{1}}{3!}\phi\left(\chi^{3} + (\chi^{*})^{3}\right) + \frac{\lambda_{2}}{2}\phi^{2}(\chi^{*}\chi),$$
semi-production pair-production

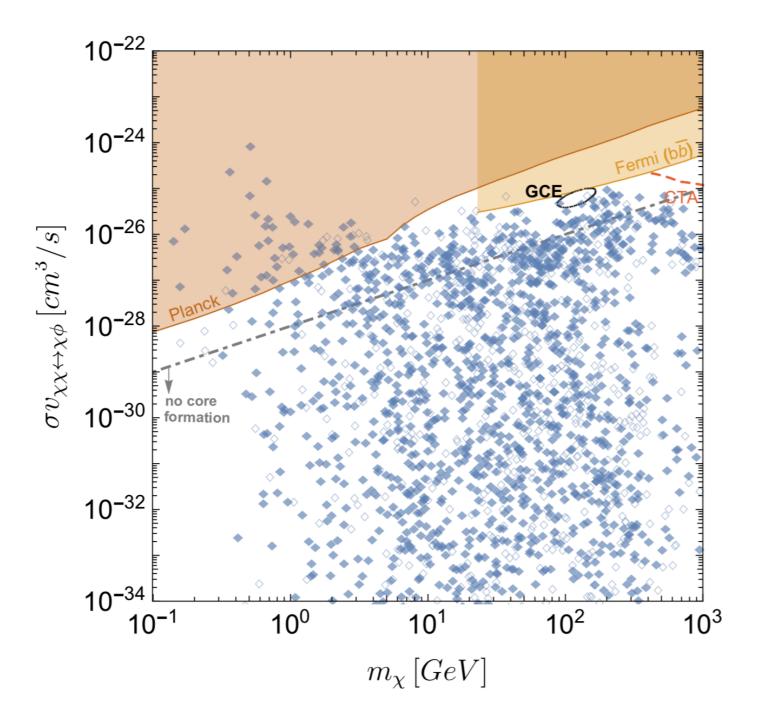
- ϕ gets a VEV, but χ doesn't
- $m_{\phi} < 3m_{\chi} \rightarrow \text{no decays}$

Evolution of density and temperature



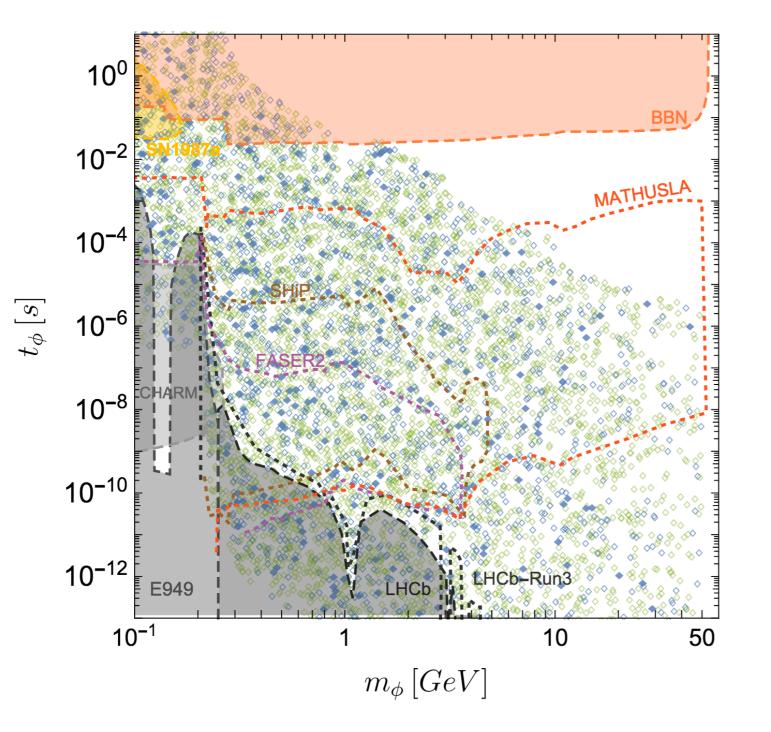
 $m_{\chi} = 100 \text{ GeV}, \ \mu_{\phi} = 1 \text{ GeV}, \ \lambda_1 = 1.1 \times 10^{-2}, \ \lambda_2 = 10^{-8}, \ \lambda_{h\phi} = 6 \times 10^{-11}$

Indirect detection constraints and predictions



- The results of the scan in the parameter space for the DM production dominated by the semi-annihilation processes.
- The coloured squares indicate the points, which are within the reach of the future searches for the mediator φ and the empty ones are beyond these prospects.
- The points above the grey dotdashed line can potentially explain the core formation in dSph [1803.09762]

Long-lived particle searches



- The constraints on the properties of the mediator φ and the prospects for its detection.
- The blue points correspond to the DM production dominated by the semi-annihilation, while the green ones – by the pair-annihilation.

Conclusion

- We have studied the freeze-in mechanism based on the semi-annihilation process
- Semi-production freeze-in requires larger cross sections than the pair-production freeze-in
- This mechanism can be incorporated in various models and promises an interesting phenomenology that is within the reach of near-future experiments:
 - Indirect detection
 - Collider searches (FPF)

