The atmospheric neutrino flux and prompt neutrinos

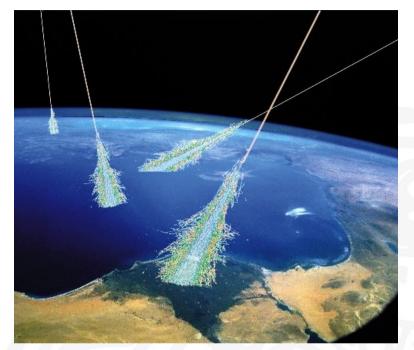
Rikard Enberg Uppsala University

4th Forward Physics Facility Meeting Feb 1, 2022

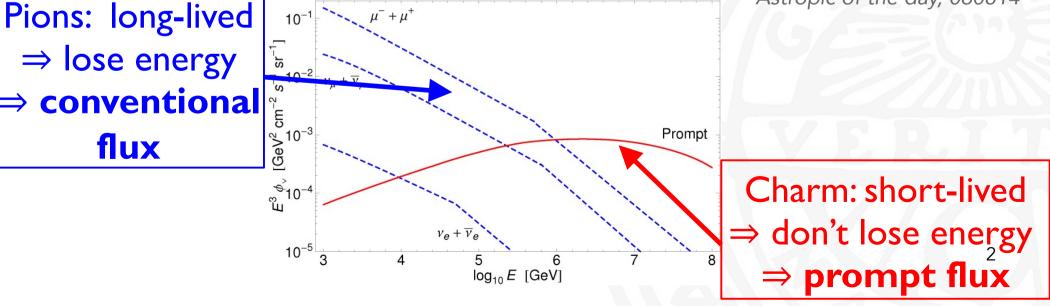


Atmospheric neutrinos

- Cosmic rays hit upper atmosphere, collide with air
- Very large energy → hadron production: pions, kaons, charm
- Semileptonic decays
 ⇒ neutrino flux



Astropic of the day, 060814



Conventional neutrino flux

- Pions (and kaons) are produced in more or less every inelastic collision
- π^+ always decay to neutrinos: $BR(\pi^+ \rightarrow \mu^+ v_\mu) = 99.98 \%$
- But π[±], K[±] are long-lived (cτ ~ 8 meters for π⁺)
 ⇒ lose energy through collisions before decay
 ⇒ neutrino energies are degraded
- This is called the *conventional neutrino flux*

Prompt neutrino flux

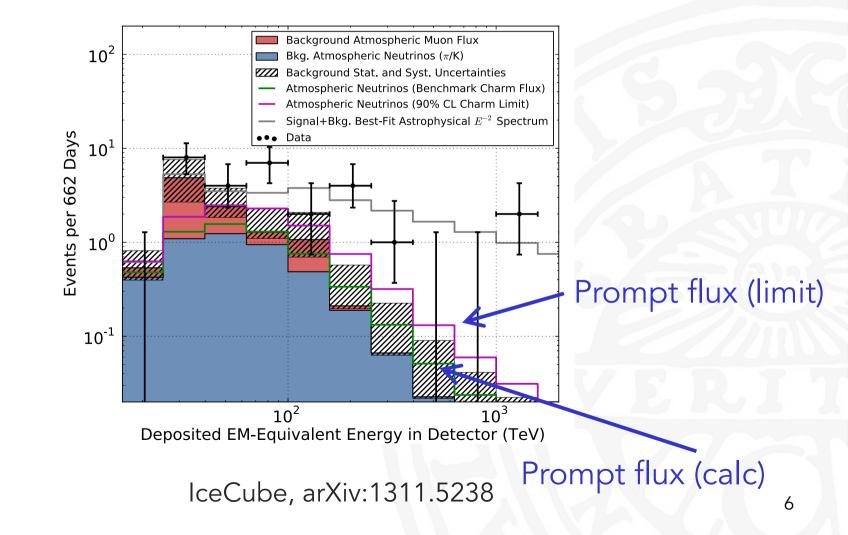
- Hadrons containing heavy quarks (charm or bottom) are extremely short-lived:
 - ⇒ decay before losing energy
 - ⇒ harder neutrino energy spectrum
- However, production cross-section is much smaller
- There is a cross-over energy above which prompt neutrinos dominate over the conventional flux
- This is called the **prompt neutrino flux**

Why are we interested?

- Atmospheric neutrinos are a large background to cosmic neutrinos at very high energies
- Thus need to understand atmospheric neutrinos in order to study astrophysical sources
- Learn about atmospheric cascades and the underlying production mechanism
- Higher energy pp collisions than in LHC: can maybe even learn something about QCD?

IceCube discovery of cosmic neutrinos from 2013

Significance was sensitive to the prompt flux prediction



Important message

QCD is crucial for prompt neutrinos:

- Small Bjorken-x (Need very small x)
- Forward region (Hard to measure at colliders)
- Fragmentation of quarks \rightarrow hadrons (Non-perturbative, hard to measure)
- Nuclear effects in pA hard interactions

FPF may help with some of these!

The calculation has many ingredients

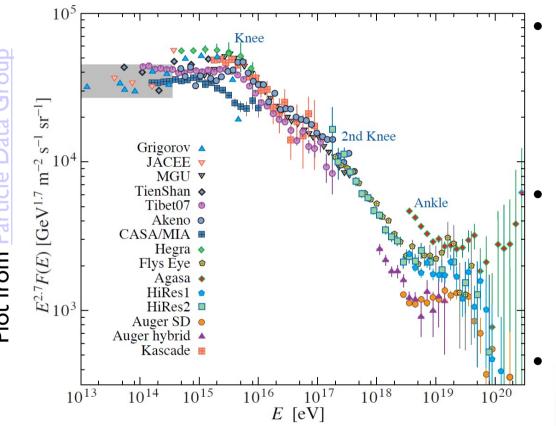
- Incident cosmic ray flux
- Forward cross section for heavy quarks in pp/pA collisions at extremely high energy (pQCD)
- Fragmentation of heavy quarks into hadrons
- Rescattering of nucleons, hadrons (hadronic xsecs) (scattering lengths)
- Decay spectra of charmed mesons & baryons (decay lengths)
- Cascade equations and their solution (Semi-analytic: spectrum-weighted Z-moments)

Calculations of the prompt flux

ERS: RE, Reno, Sarcevic, arXiv:0806.0418

Bhattacharya, RE, Reno, Sarcevic, Stasto, arXiv:1502.01076 Fedynitch, Engel, Gaisser, Riehn, Stanev, arXiv:1503.00544 Garzelli, Moch, Sigl, arXiv:1507.01570 Gauld, Rojo, Rottoli, Sarkar, Talbert, arXiv:1506.08025, 1511.06346 Halzen and Wille, arXiv:1605.01409 Bhattacharya, RE, Jeong, Kim, Reno, Sarcevic, Stasto, arXiv:1607.00193 PROSA Collaboration (Garzelli et al), arXiv:1611.03815, 1911.13164 Benzke, Garzelli, et al., arXiv:1705.10386 Goncalves, Maciula, Szczurek, arXiv:2103.05503 Jeong, Bai, Diwan, Garzelli, Kumar, Reno, arXiv:2107.01178 Arleo, Jackson, Peigné, arXiv:2112.10791

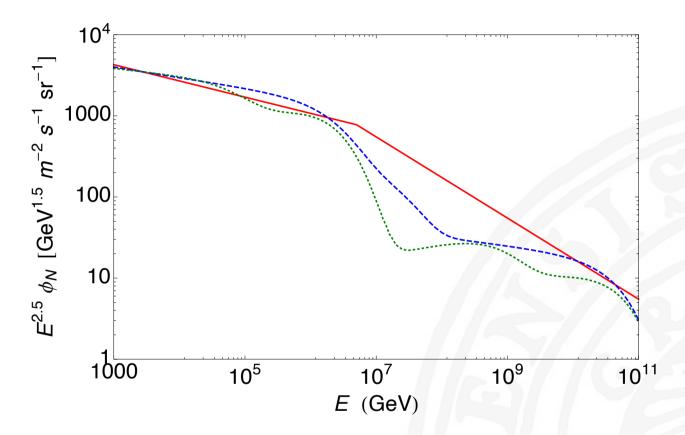
Cosmic rays (CR)



- Knees and ankles \rightarrow seems natural to associate different sources with different energy ranges of the CR flux
- Highest energies: Extragalactic origin? → GRBs, AGNs, or more exotic

Lower energies: Galactic origin? →SNRs etc

Incident cosmic ray flux: nucleons



Solid red = Broken power law (old standard) Dashed blue = Gaisser all proton (H3p) Dotted green = Gaisser, Stanev, Tilav (GST4)

Calculating the neutrino flux: Particle production

Particle physics inputs: energy distributions

$$\frac{dn(k \to j; E_k, E_j)}{dE_j} = \frac{1}{\sigma_{kA}(E_k)} \frac{d\sigma(kA \to jY, E_k, E_j)}{dE_j}$$
$$\frac{dn(k \to j; E_k, E_j)}{dE_j} = \frac{1}{\Gamma_k} \frac{d\Gamma(k \to jY; E_j)}{dE_j}$$

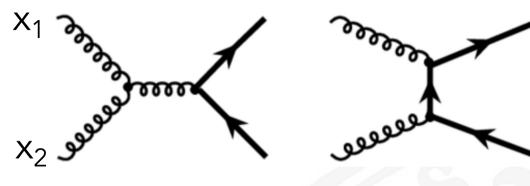
along with interaction lengths, or cooling lengths

$$\lambda_N(E) = \frac{\rho(h)}{\sigma_{NA}(E) n_A(h)}$$

 \rightarrow Need the charm production cross section d σ /dx_F

Problem with QCD

Charm production:



x_F = Feynman-x → ≃ momentum fraction of charm quark

where
$$x_{1,2} = \frac{1}{2} \left(\sqrt{x_F^2 + \frac{4M_{c\bar{c}}^2}{s}} \pm x_F \right)$$

CM energy is large: $s = 2E_pm_p$ so $x_1 \sim x_F$ and $x_2 \ll 1$

- → We need extremely small Bjorken-x, in the range 10⁻⁷ to 10⁻⁴
- \rightarrow Very asymmetric, x_F is very large

Problem with QCD at small x

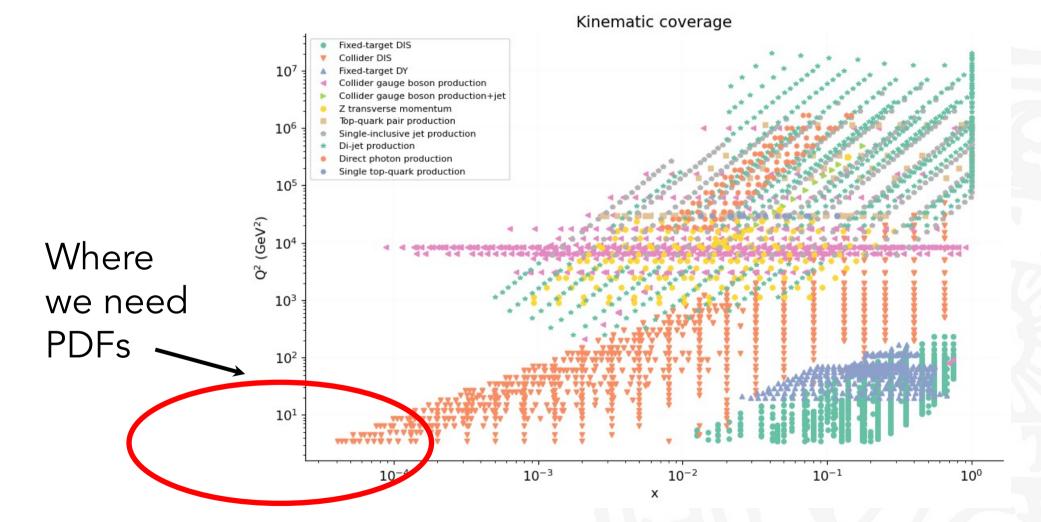
- Parton distribution functions poorly known at small x
- At small x, must resum large logs: $a_s \ln(1/x)$

- If logs are resummed (*BFKL*): power growth ~ $x^{-\lambda}$ of gluon distribution as $x \rightarrow 0$
- Unitarity might even be violated (T-matrix > 1)

How small x do we know?

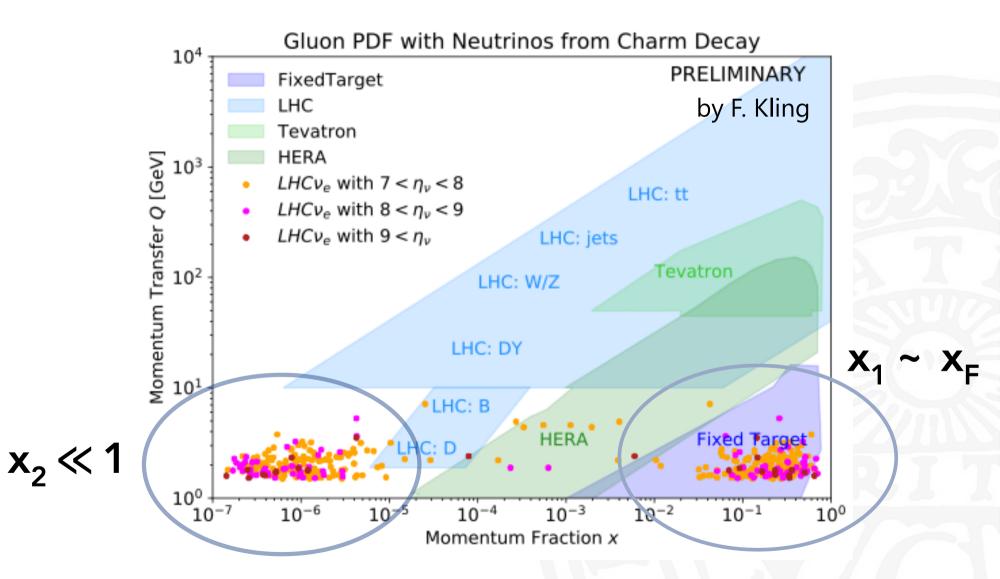
- We haven't measured anything at such small x
- E.g. the MSTW pdf has $x_{min} = 10^{-6}$
- But that is an extrapolation!
- HERA pdf fits: $Q^2 > 3.5 \text{ GeV}^2$ and $x > 10^{-4}$
- See Gao, Harland-Lang, Rojo, arXiv:1709.04922 for more on pdfs

Kinematic plane of NNPDF



NNPDF collaboration, https://nnpdf.mi.infn.it/research/data/ Also talk today by E. Nocera

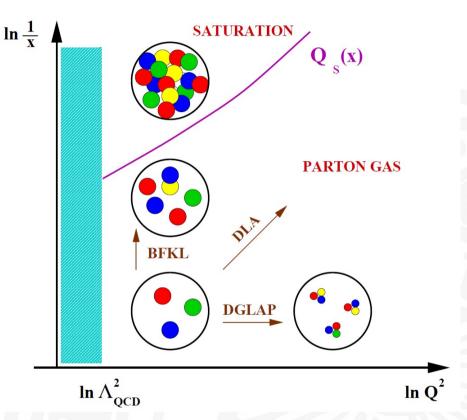
What can be done?



By F. Kling, from talks by T. Ariga and L. Harland-Lang

Parton saturation

- **Saturation** at small x:
 - Number of gluons in the nucleon becomes so large that gluons recombine
 - Reduction in the growth



- This is sometimes called the color glass condensate
- Non-linear QCD evolution: Balitsky-Kovchegov equation

Bhattacharya et al (2016): Redo QCD calculations in many ways

- Standard NLO QCD with newer PDFs
 - Earlier calc updated with RHIC/LHCb input, uses Nason, Dawson, Ellis and Mangano, Nason, Ridolfi
- Dipole picture with saturation
 - Approximate solution of Balitsky-Kovchegov equation
 - Update of ERS calc with new HERA fits + other dipoles

kT factorization with and without saturation

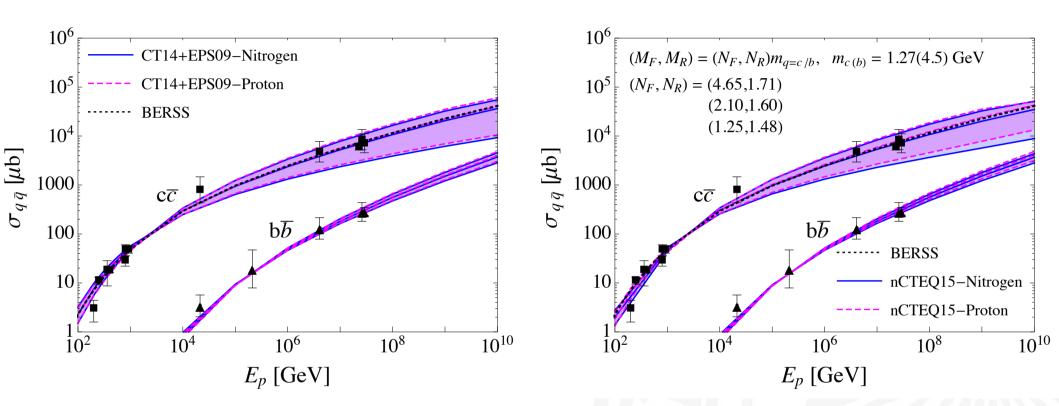
- Resums large logs, $a_s \log(1/x)$ with BFKL
- Off-shell gluons, unintegrated PDFs (+ subleading...)
- Kutak, Kwiecinski, Martin, Sapeta, Stasto (permutations)
 Include scale variations, PDF errors, charm mass, etc
 → Plausible upper and lower limits on xsec

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Also include nuclear shadowing

- Partons are not in a free nucleon, but in a nucleus!
- Estimate shadowing with nuclear PDFs (nCTEQ15 and EPS09)
- Reduces flux by 10–30% at the highest energies
- Larger effect on the flux than on the total $\sigma(cc)$ due to asymmetric $x_{1,2}$

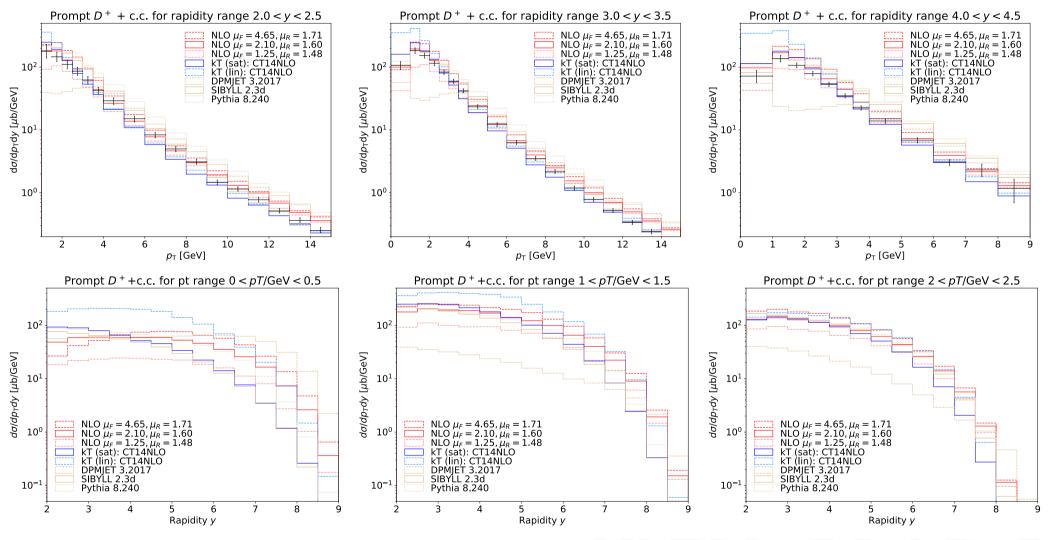
σ (cc) and σ (bb)



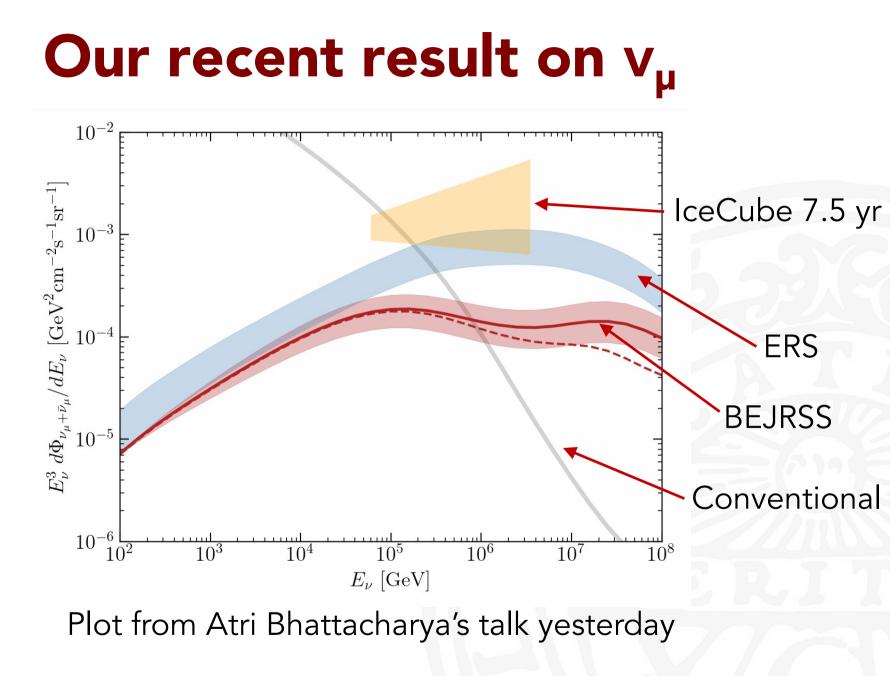
Data from RHIC, LHC and lower energies Total cross sections well described by all calculations (at high energies), nuclear shadowing small

(Error bands=scale variations and PDF uncertainties)

New work: D[±] meson production

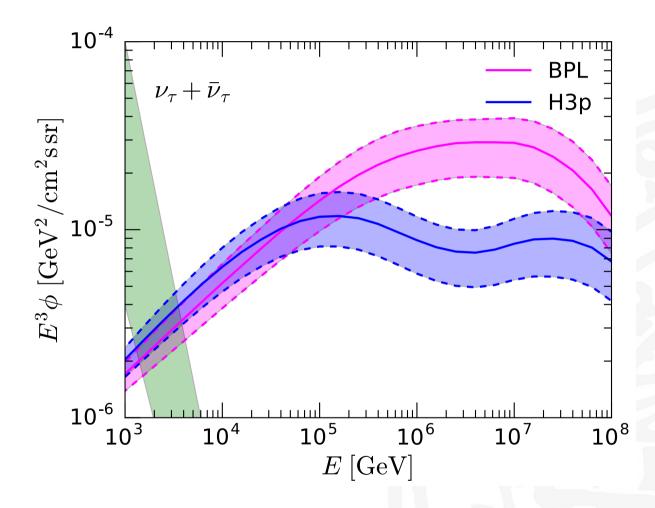


Some of our calculations compared with MC generators and LHCb data (plots by F. Kling)



Note that prompt $v_{\mu} \approx v_e$ but conventional $v_{\mu} \gg v_e$

Prompt flux important for v_{τ}



The conventional flux is much smaller for v_{τ}

Comparison of calculations

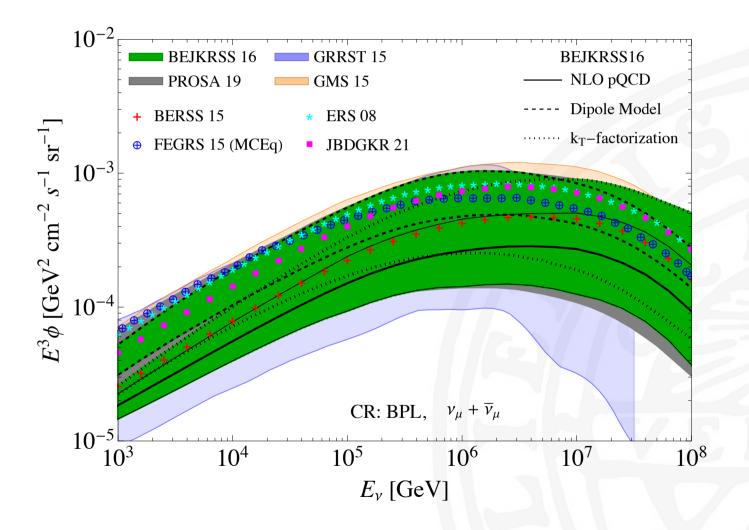
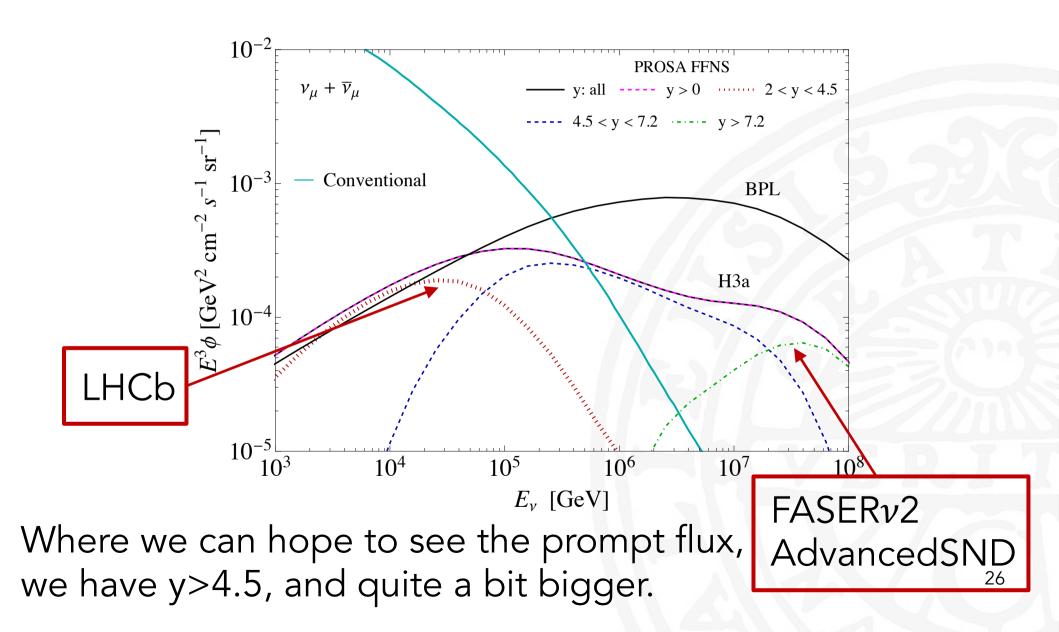
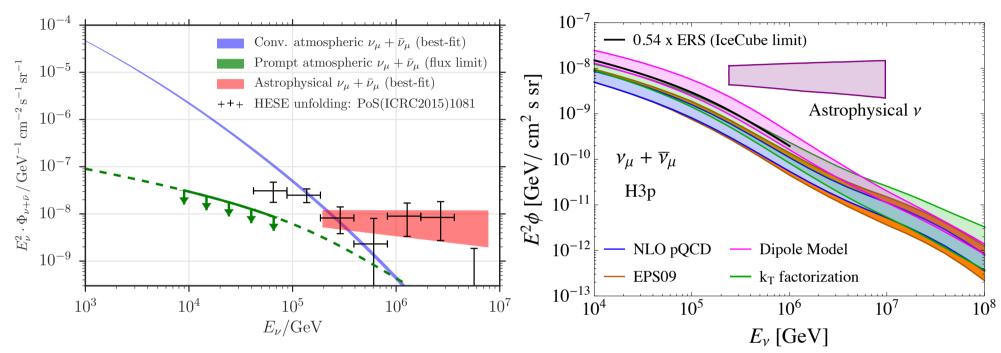


Figure from the FPF Snowmass White Paper draft

What rapidity ranges?



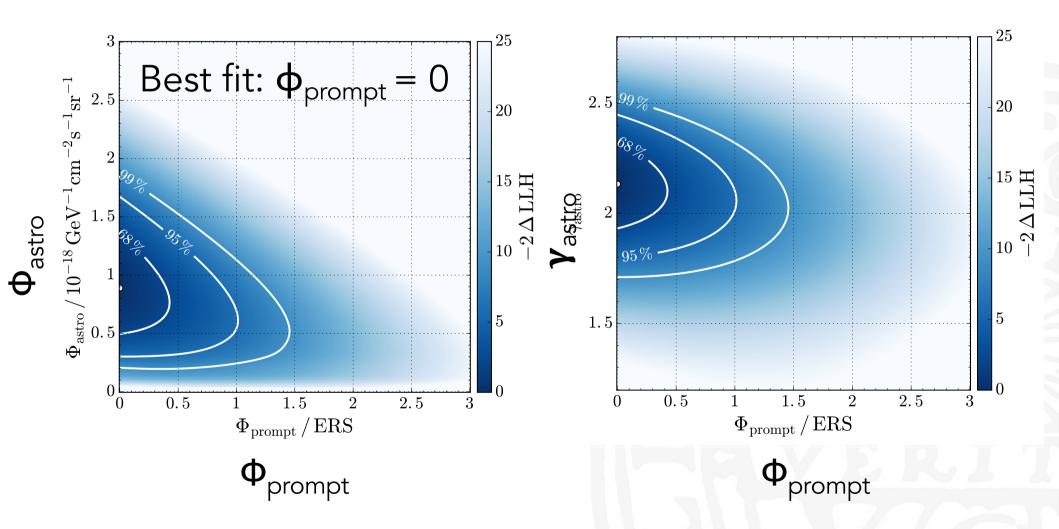
IceCube vs prompt flux



3-year IceCube limit on the prompt flux (at 90% CL): 0.54 x (*ERS modified with H3p CR's*) Best fit is φ_{prompt} = 0

L. Rädel & S. Schoenen (IceCube), PoS ICRC2015, 1079

IceCube fits to Φ_{astro} and Φ_{prompt}



IceCube, arXiv:1607.08006

What can FPF do?

Generally: constrain PDFs at very small **and** very large x

Neutrinos: (see also talk by Hallsie Reno)

✓ FASER ν 2: neutrinos at η > 8.6 or 8.8

✓ Advanced SND:

- Neutrinos in far detector at 7.2 < η < 8.4
- Near detector $4 < \eta < 5$

− Can measure neutrinos from charm here,
 learn about charm → neutrino fragmentation

✓ FLArE: v_{τ} flux

Basically, predict neutrinos at FPF and in atmosphere with same underlying QCD calculations