

Freeze-in sterile neutrino dark matter in a class of $U(1)'$ models with inverse seesaw

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Motivations

- **Neutrino Mass:** Neutrino is massless in SM. Several neutrino oscillation experiments confirm neutrino has mass.
- **Existance of dark matter:** The galaxy rotation curve and Bullet cluster experiment by Chandra X-ray observatory → Cannot be explained by standard model (SM) of particle physics (other alternative solution: Modified Gravity).
- **511keV line:**The long standing puzzle of 511 keV line observed by the INTEGRAL satellite at the Milky way centre needs to be explained.
- **WIMPs:** The non observations of any signal in the direct detection dark matter experiments put stringent bounds on WIMPs of masses $\gtrsim 1\text{GeV}$.

The Model

	Q_{L_i}	u_{R_i}	d_{R_i}	ℓ_{L_i}	e_{R_i}	ν_{R_α}	H	Φ	S
$SU(3)_C$	3	3	3	1	1	1	1	1	1
$SU(2)_L$	2	1	1	2	1	1	2	1	1
$U(1)_Y$	1/6	2/3	-1/3	-1/2	-1	0	1/2	0	0
$U(1)'$	x_q	x_u	x_d	x_l	x_e	x_ν	$\frac{x_h}{2}$	$-x_\phi$	0

The most general Yukawa Lagrangian invariant under this new gauge group is,

$$-\mathcal{L}_{\text{Yukawa}} = Y_e \bar{\ell}_L H e_R + Y_\nu \bar{\ell}_L \tilde{H} \nu_R + Y_u \bar{Q}_L \tilde{H} u_R + Y_d \bar{Q}_L H d_R + y_{NS} \bar{\nu}_R \Phi S + \frac{1}{2} \bar{S}^c M_\mu S + \text{h.c.}$$

$$x_\ell = -x_\Phi - \frac{x_H}{2}, x_e = -x_\Phi - x_H, x_q = \frac{1}{6}(2x_\Phi + x_H), x_u = \frac{1}{3}(2x_H + x_\Phi),$$

$$x_d = \frac{1}{3}(x_\Phi - x_H)$$

Freeze-in production of dark matter:

The Boltzmann equation to solve the relic density

$$\dot{n} + 3Hn = \langle \sigma v(ab \rightarrow NN) \rangle n_a n_b.$$

The final abundance of the RHN DM is,

$$\frac{dY}{dT} = \frac{\langle \sigma v(ab \rightarrow NN) \rangle n_a n_b}{sHT}.$$

The thermal average dark matter annihilation cross section

$$\langle \sigma v \rangle = \frac{1}{(sY_{EQ})^2} g_N^2 \frac{m_N}{64\pi^2 X} \int_{4m_N^2}^{\infty} ds \times 2(s - 4m_N^2) \sigma(s) \sqrt{s} K_1\left(\frac{\sqrt{s}}{T}\right),$$

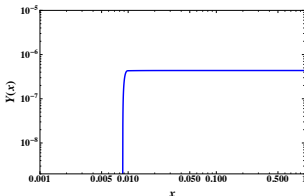
Dark matter in our $U(1)'$ model:

For the process $f\bar{f} \rightarrow NN$ in the s -channel mediated by the Z' as,

$$|\bar{\mathcal{M}}|^2 = \frac{g'^4 y_N'^2}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2} \left[c_V^2 \{s^2 + (s - 4m_f^2)(s - 4m_N^2) \cos^2 \theta - \right. \\ \left. 4s(m_N^2 - m_f^2) - 16m_N^2 m_f^2\} + c_A^2 \{s^2 + (s - 4m_f^2)(s - 4m_N^2) \cos^2 \theta - \right. \\ \left. -4s(m_N^2 + m_f^2) + 16m_N^2 m_f^2 + \frac{16m_f^2 m_N^2}{M_{Z'}^4} (s - M_{Z'}^2)^2\} \right]$$

Heavy Z' ($M_{Z'}' \gg T_R \gg m_N$):

$$\Omega_N h^2 \simeq 0.12 \times \left(\frac{m_N}{1\text{MeV}} \right) \left(\frac{10}{g_*} \right)^{\frac{3}{2}} \left(\frac{T_R}{5\text{MeV}} \right)^3 \left(\frac{9.7\text{TeV}}{M_{Z'}'/g'} \right)^4 \\ \times \frac{x_\phi^2}{3} \left[2 \left\{ \left(\frac{3}{4} x_h + x_\phi \right)^2 + \left(\frac{x_h}{4} \right)^2 \right\} \right. \\ \left. + \left\{ \left(x_\phi + \frac{x_h}{4} \right)^2 + \left(\frac{x_h}{4} \right)^2 \right\} \right].$$



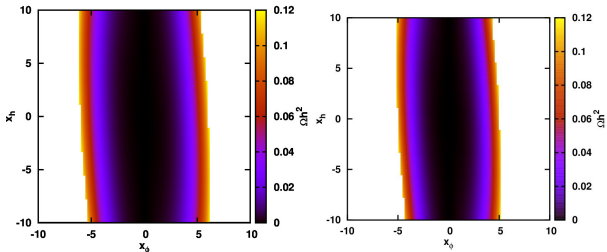


Figure: Relic density plot in the $x_h - x_\phi$ plane. In the left panel, we have fixed the values of the parameters as $m_N = 1$ MeV, $T_R = 10$ MeV, $M_{Z'} = 10$ TeV and $g' = 0.1$ whereas the right panel is for $m_N = 10$ MeV, $T_R = 50$ MeV, $M_{Z'} = 5$ TeV and $g' = 0.01$.

Light Z' ($M_{Z'} \ll m_N$):

$$\Omega_N h^2 = 0.12 \times \left(\frac{106.75}{g_*} \right)^{\frac{3}{2}} \left(\frac{g'}{3.04 \times 10^{-6}} \right)^4 x_\phi^2 (10x_h^2 + 13x_\phi^2 + 16x_H x_\phi).$$

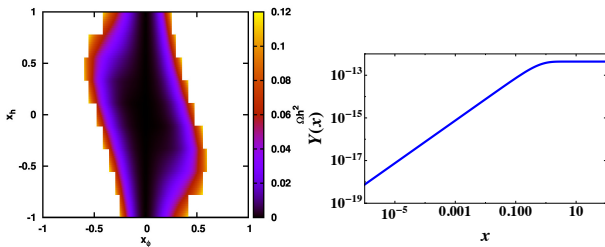


Figure: Relic density plot in $x_h - x_\phi$ plane. In the left panel, We have varied both x_h and x_ϕ from -1 to $+1$ for $g' = 3.04 \times 10^{-6}$ and in right panel, $M_{Z'} = 0.1\text{GeV}$, $m_N = 1\text{TeV}$, $x_\phi = 1$, $x_h = -1.2$ for $g' = 3.65 \times 10^{-6}$.

Intermediate Z' ($m_N \ll M_{Z'} \ll T_R$):

$$\Omega_N h^2 \simeq 0.12 \left(\frac{g'}{6.54 \times 10^{-9}} \right)^2 \left(\frac{m_N}{10 \text{keV}} \right) \left(\frac{10 \text{GeV}}{M_{Z'}} \right) \left[\frac{x_\phi^2 (10x_h^2 + 13x_\phi^2 + 16x_h x_\phi)}{\frac{1}{36} (241x_h^2 + 418x_\phi^2 + 436x_h x_\phi) + x_\phi^2} \right].$$

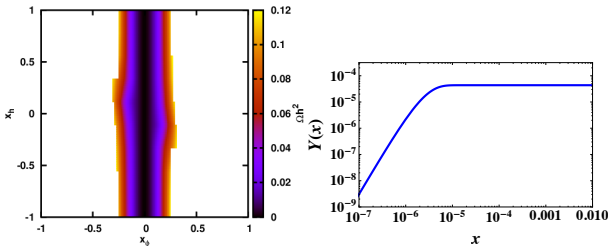
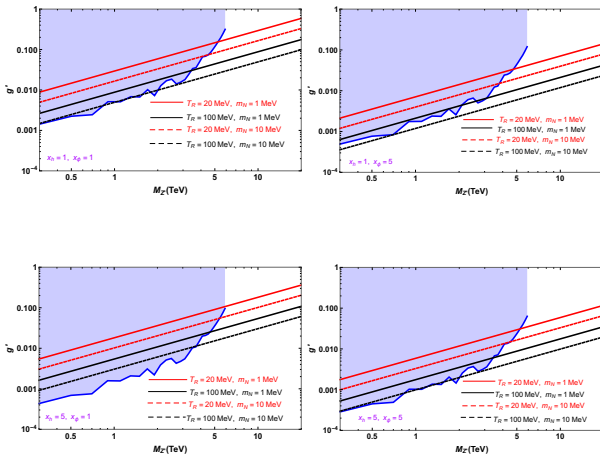
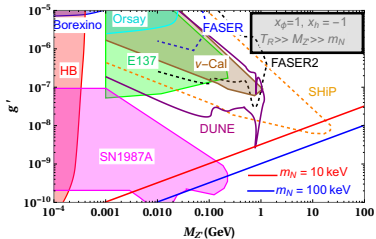
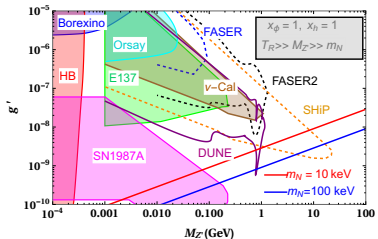
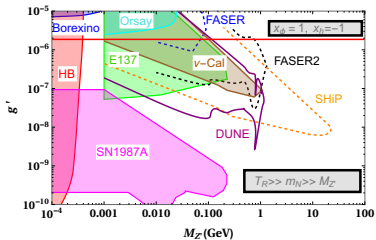
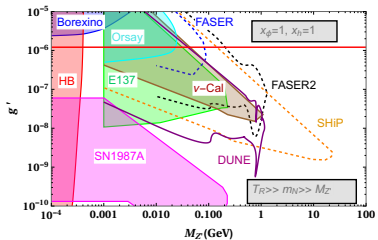


Figure: Relic density plot in $x_h - x_\phi$ plane. In the left panel, We have varied both x_h and x_ϕ from -1 to $+1$ for $g' = 6.54 \times 10^{-9}$, $m_N = 10 \text{keV}$, $M_{Z'} = 1 \text{GeV}$ and in right panel, $x_\phi = 1$, $x_h = -1.2$ for $g' = 2 \times 10^{-9}$, $m_N = 10 \text{keV}$, $M_{Z'} = 10 \text{GeV}$.

Bounds on coupling for heavy Z' from Experiments



Bounds on coupling for light Z' from Experiments



Discussions

- Detailed study of dark matter in $U(1)'$ model for a wide range of Z' mass.
- The heavy Z' can be probed in ATLAS, CMS whereas the light Z' can be probed from several lifetime frontier experiments like SHiP, FASERs etc.
- The decay of MeV scale dark matter can produce the long standing problem of galactic 511keV line.

Thank You!