# Evolution of $\mathbf{N}=50$ shell and neutron single-particle states towards ${ }^{78} \mathrm{Ni}$ : ${ }^{79 \mathrm{Zn}(\mathrm{d}, \mathrm{py})}{ }^{80 \mathrm{Zn}}$ 


$79 \mathrm{Zn}(\mathrm{n}, \gamma)$ cross section and r-process
around $A=80$ mass region

Spokesperson: E. Sahin, University of Oslo, Oslo, Norway

- Location of the neutron single particle orbitals
- Size of the $\mathrm{N}=50$ gap at $\mathrm{Z}=30$ near ${ }^{78} \mathrm{Ni}(Z=28)$

$$
\nu 1 g_{9 / 2}^{-1} \otimes \nu 2 d_{5 / 2}^{1} \quad \text { States }
$$

I=2 transfer $\left(2^{+}, 3^{+}, \ldots, 7^{+}\right)$from the g.s. ${ }^{79} \mathrm{Zn}, 9 / 2^{+}$
Sensitive to changes in the $\mathrm{N}=50$ shell gap due to monopole interaction
SM calculations exist only for $5^{+}$and 6+ from $\mathrm{Z}=28$ to $\mathrm{Z}=36$

1p-1h excitations


## States from 1p-1h excitations



$$
\nu 1 g_{9 / 2}^{-1} \otimes \nu 2 d_{5 / 2}^{1} \quad \text { States }
$$

B $\mathrm{I}=2$ transfer $\left(2^{+}, 3^{+}, \ldots, 7^{+}\right)$from the g.s. $79 \mathrm{Zn}, 9 / 2^{+}$

- Sensitive to the $\mathrm{N}=50$ shell gap due to monopole interaction
SM calculations exist only for $5^{+}$and $6^{+}$

$$
\nu 1 g_{9 / 2}^{-1} \otimes \nu 3 s_{1 / 2}^{1} \quad \text { States }
$$

- $1=0$ transfer $\left(4^{+}, 5^{+}\right)$from the g.s.
- Location information, $3 s_{1 / 2}$
* No SM calculations
* Estimated to be few hundreds keV higher than the I=2 states
J.S. Thomas et al., Phys. Rev. C 71, 021302R (2005)
${ }^{80} \mathrm{Zn} \nu(1 p-1 h)$


Location of the $\mathrm{s}_{1 / 2}$ and $\mathrm{d}_{5 / 2}$ states in the $\mathrm{N}=51$ isotones


$$
\nu 1 g_{9 / 2}^{-1} \otimes \nu 2 d_{5 / 2}^{1} \quad \text { States }
$$

- $\mathrm{I}=2$ transfer $\left(2^{+}, 3^{+}, \ldots, 7^{+}\right)$from the g.s. $79 \mathrm{Zn}, 9 / 2^{+}$
- Sensitive to the $\mathrm{N}=50$ shell gap due to monopole interaction
SM calculations exist only for $5^{+}$and $6^{+}$

$$
\nu 1 g_{9 / 2}^{-1} \otimes \nu 3 s_{1 / 2}^{1} \quad \text { States }
$$

8 I=0 transfer ( $4^{+}, 5^{+}$) from the g.s.

* Location i calculations nformation, $3 \mathrm{~s}_{1 / 2}$
- No SM
- Estimated to be few hundreds keV higher than the $\mathrm{I}=2$ states

$$
\nu 1 g_{9 / 2}^{-2} \otimes \nu 3 s_{1 / 2}^{2} \quad \text { States }
$$

- I=0 transfer from the isomeric state in 79 Zn
- $\mathrm{E}_{1 / 2+}=1.05 \mathrm{MeV} \& \mathrm{t}_{1 / 2} \geq 200 \mathrm{~ms}$
- Possible assignment as (1p-2h) excitations across $\mathrm{N}=50$ (PFSDG-U int.)
- Evidence for a $0^{+}$intruder state


ISOLDE: X. F. Yang et al., PRL 116, 182502 (2016).

79Zn $\boldsymbol{v}(1 p-2 h)$

$\nu 1 g_{9 / 2}^{-1} \otimes \nu 2 d_{5 / 2}^{1} \quad$ States

- $1=2$ transfer $\left(2^{+}, 3^{+}, \ldots, 7^{+}\right)$from the g.s., $9 / 2^{+}$
- Sensitive to the $\mathrm{N}=50$ shell gap due to monopole interaction
SM calculations exist only for $5^{+}$and $6^{+}$

$$
\nu 1 g_{9 / 2}^{-1} \otimes \nu 3 s_{1 / 2}^{1} \quad \text { States }
$$

- $\mathrm{I}=0$ transfer $\left(4^{+}, 5^{+}\right)$from the g.s.
* Location i calculations nformation, $3 \mathrm{~s}_{1 / 2}$
- No SM
- Estimated to be few hundreds keV higher than the $\mathrm{I}=2$ states
$\mathrm{E}\left(0^{+}\right) \cong 2 \mathrm{MeV} \boldsymbol{\nu}(2 \mathrm{p}-2 \mathrm{~h})$

80Zn $\boldsymbol{\nu}(2 p-2 h)$


ISOLDE: X. F. Yang et al., PRL 116, 182502 (2016).

$$
\nu 1 g_{9 / 2}^{-1} \otimes \nu 2 d_{5 / 2}^{1} \quad \text { States }
$$

- $1=2$ transfer $\left(2^{+}, 3^{+}, \ldots, 7^{+}\right)$from the g.s., $9 / 2^{+}$
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SM calculations exist only for $5^{+}$and $6^{+}$

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\nu 1 g_{9 / 2}^{-1} \otimes \nu 3 s_{1 / 2}^{1} \quad \text { States }
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\nu 1 g_{9 / 2}^{-2} \otimes \nu 3 s_{1 / 2}^{2} \quad \text { States }
$$

- $\mathrm{I}=0$ transfer from the isomeric state in 79 Zn
- $E_{1 / 2+}=1.05 \mathrm{MeV} \& \mathrm{t}_{1 / 2} \geq 200 \mathrm{~ms}$
- Possible assignment as (1p-2h) excitations across N=50 (PFSDG-U int.)
- Evidence for a $0^{+}$intruder state

$2^{+}$state at ISOLDE Coulomb excitation : J. Van de Walle et al., PRL 99, 142501 (2007).

RIKEN inelastic scattering and proton removal: ${ }^{9} \mathrm{Be}(80 \mathrm{Zn}, 80 \mathrm{Zn})$ and ${ }^{9} \mathrm{Be}\left({ }^{81} \mathrm{Ga}, 80 \mathrm{Zn}\right)$ :
Y. Shiga et al., PRC 93024320 (2016).
$5^{+}, 6^{+}$: K. Sieja and F. Nowacki , PRC85 051301(R) (2012)

## DWBA Calculations via FRESCO



## Setup and Experiment



## ${ }^{79} \mathrm{Zn}(\mathrm{d}, \mathrm{p} \gamma)^{80 \mathrm{Zn}}$ inverse kinematics

Beam energy (79Zn)
Beam intensity on MINIBALL
Target thickness ( $\mathrm{CD}_{2}$ )
Cross sections

395 MeV ( $5 \mathrm{MeV} /$ nuc)
$4 \lambda 10 / \mathrm{pps} \rightarrow 4 \times 10^{3} \mathrm{pps}$
1 mg $\mathrm{cm}^{2} \rightarrow 2 \mathrm{mg} / \mathrm{cm}^{2}$
DWBA via FRESCO

Isolde Database: $10^{6} \mathrm{pps} / \mu \mathrm{C}$ UCx
Suggested value from the earlier records: $\sim 5.10^{5} \mathrm{pps} / \mu \mathrm{\mu C}$ UCx
$1.6 \mu \mathrm{C}$ total proton intensity $+5 \%$ transmission eff. $==>\mathbf{4} \mathbf{x 1 0 ^ { 4 }} \mathbf{p p s}$ at MINIBALL
$\mathrm{Rb} / \mathrm{Ga}$ contamination: $\mathbf{4 \times 1 0 ^ { 3 }} \mathbf{~ p p s}$ recommended by TAC .



-TREX angular coverage: $60 \%$

- MINIBALL efficiency on average $6 \%$ at 1 MeV
- Average cross section for $\nu 1 g_{9 / 2}^{-1} \otimes \nu 2 d_{5 / 2}^{1}$ states : 20 mb
- Average cross section for $\nu 1 g_{9 / 2}^{-1} \otimes \nu 3 s_{1 / 2}^{1}$ states : 28 mb
(particle- $\boldsymbol{\gamma}$ )
(particle- $\boldsymbol{\gamma} \boldsymbol{\gamma}$ )



## Identification of states in the worst scenario:

Low-spin states have lower cross section and energy behaviour also follows a certain trend (parabola) .

One can use gamma-tagged proton angular distributions and see how this compares with the expected trends of the excitation energy and the s.p. cross sections.

Plus shell model calculations might be helpful.

|  | This proposal 79Zn(d,py)80Zn | Accepted proposal Spokesperson: R. Orlandi $\text { IS556 } \quad 80 Z n(d, p y)^{81 Z n}$ |
| :---: | :---: | :---: |
| Beam intensity | $5 \times 10^{4} \mathrm{pps} / \mu \mathrm{C}$ at UCx <br> $4 \times 10^{3} \mathrm{pps}$ at MINIBALL | $3 \times 10^{4} \mathrm{pps} / \mu \mathrm{C}$ at UCx <br> $2.3 \times 10^{3} \mathrm{pps}$ at MINIBALL |
| Target thickness | $2 \mathrm{mg} / \mathrm{cm}^{2}$ | $2 \mathrm{mg} / \mathrm{cm}^{2}$ |
| Required beam time | 7 days | 7 days |
| Expected yield | $300 \mathrm{p} \gamma / 30 \mathrm{p} \gamma \gamma$ events for one $\mathrm{I}=2$ state $400 \mathrm{p} \gamma / 40 \mathrm{p} \gamma \gamma$ events for one $\mathrm{I}=0$ state | $350 \mathrm{p} \gamma$ events for the I=0 state, $1 / \mathbf{2}^{+}$ |

[^0]Alternative solution could be dropping the neutron converter (UCx + neutron converter + quartz line + RILIS)

## Thank you

## Additional slides



Shell gap from mass is correlated.
SM approach:
K. Sieja and F. Nowacki , PRC85 051301(R) (2012)
F. Nowacki et al., PRL 117, 272501 (2016)
A. Welker et al., PRL 119, 192502 (2017).


Correlation effects can be explored and help theory

Mean-field approach:
M. Bender et al. PRC78, 054312 (2008)

## Momentum matching in transfer reactions

${ }^{60} \mathrm{Ni}\left(\alpha,{ }^{3} \mathrm{He}\right): \mathrm{Q}_{\text {g.s. }}=-12.8 \mathrm{MeV}->$ high momentum transfers ${ }^{60} \mathrm{Ni}(\mathrm{d}, \mathrm{p}): \mathrm{Q}_{\mathrm{g} . \mathrm{s} .}=5.6 \mathrm{MeV}$-> low momentum transfers



FIG. 8. Measured differential cross sections and DWBA fits for $l=2$ transitions. All fits are based on NLFR calculations using L/B parameters.

TABLE V. Summary of ( $d, p$ ) results for levels in ${ }^{88} \mathrm{Sr}$.

| Level No. ${ }^{\text {a }}$ | E* (keV) | Cosman and Slater ${ }^{\text {a }}$ |  |  | This experiment ${ }^{\text {b }}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  | $J^{\pi}$ |  |
|  |  | $l$ | $G_{l J}{ }^{\text {c }}$ | $l$ | $G_{l j}^{88}$ | $G_{i j}^{89}$ | (assumed) ${ }^{\text {d }}$ | $S_{l j}^{88}$ |
| 1 | (1836) | 2 | 0.126 | 2 | (0.13) |  | $2^{+}$ | 0.25 |
| 5 | 4032 | 2 | 0.279 | 2 | 0.35 |  | $2^{+}$ | 0.71 |
| 6 | 4294 | 2 | 0.376 | 2 | 0.53 |  | $4^{+}$ | 0.59 |
| 7 | 4413 | 2 | 0.875 | 2 | 1.18 |  | [5] ${ }^{+}$ | 1.07 |
| 8 | (4450) | 2 | 0.083 | (2) | ( $\sim 0.10$ ) |  | $[4]^{+}$ | (0.11) |
| 10 | 4514 | 2 | 1.080 | 2 | 1.31 |  | $[6]^{+}$ | 1.00 |
| 12 | 4633 | 2 | 0.564 | 2 | 0.68 |  | [3] ${ }^{+}$ | 0.97 |
| 13 | 4744 | 2 | 0.805 | 2 | 0.14 | 5.86 | $[4]^{+}$ | 0.28(0.16) |
| 17 | 5094 | 2 | 1.040 | 2 | 1.33 |  | $[7]^{+}$ | 0.89 |
|  |  |  | 5.228 |  | 5.75 | $\overline{5.86}$ |  |  |
| 15 | $4873{ }^{\text {e }}$ | 0 | 0.230 | 0 | $0.24{ }^{\text {e }}$ |  | [4] ${ }^{+}$ | $0.26{ }^{\text {e }}$ |
| 21 | 5416 | 0 | 0.105 | 0 | 0.13 |  | [5] ${ }^{+}$ | 0.12 |
| 22 | 5466 | 0 | 0.563 | 0 | 0.61 |  | $4^{+}$ | 0.67 |
| 23 | (5506) | 0 | 0.027 | (0) | <0.01 |  |  |  |
| 25 | 5729 | 0 | 0.789 | 0 | 0.94 |  | $[5]^{+}$ | 0.85 |
| 26 | (5780) | 0 | 0.405 | 0 | $0( \pm 0.03)$ | 1.92 |  |  |
| 32 | (6214) | 0 | 0.031 | (0) | $(\sim 0.03)$ |  |  |  |
|  |  |  | 2.150 |  | 1.95 | 1.92 |  |  |

Equivalent of $5^{+}, 6^{+}$states in ${ }^{82} \mathrm{Ge}$ is found to be $13 / 2^{-}, 15 / 2^{-}$in ${ }^{83} \mathrm{As}$


## SM Calculations:

A.F. Lisetskiy, B.A. Brown, M. Horoi, H. Grawe,

Phys. Rev. C 70 (2004) 044314.
Interaction: JJ4B + SDI
Model spaces: pfg9+sdg
Inert Core nucleus: ${ }^{56} \mathrm{Ni}$
Tensor interactions are included

The SPEs relative to the ${ }^{56} \mathrm{Ni}$ core have been derived from the SPEs with respect to the doubly-magic ${ }^{78} \mathrm{Ni}$ core.

| Model Space | Single-Particle Energy_ |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| pfg | $\mathrm{E}\left(1 \mathrm{f}_{5 / 2}\right)$ | $\mathrm{E}\left(2 \mathrm{p}_{3 / 2}\right)$ | $\mathrm{E}\left(2 \mathrm{p}_{1 / 2}\right)$ | $\mathrm{E}\left(1 \mathrm{~g}_{9 / 2}\right)$ |
|  | -9.28590 | -9.65660 | -8.26950 | -5.89440 |
| sdg | $\mathrm{E}\left(2 \mathrm{~d}_{5 / 2}\right)$ | $\mathrm{E}\left(3 \mathrm{~s}_{1 / 2}\right)$ | $\mathrm{E}\left(1 \mathrm{~g}_{7 / 2}\right)$ |  |
|  | -1.19440 | -0.16800 | 0.2700 |  |

$$
E\left(v d_{5 / 2}-v g_{9 / 2}\right)=\text { parameter }
$$

## Weak r-process

Sensitivity study to the n-capture process in the context of neutron-rich supernova and collapsar accretion disk winds.

weak $r$ process-a rapid neutron capture process that forms a solar-type $A \sim 80 r$-process peak and potentially nuclei up to the $A \sim 130$ peak.

Nuclear data relevant to r-process calculations are (pinparity assignments, excitation energies, spectroscopic factors and can be extracted from transfer reactions, such as (d, p).

## Oslo method in inverse kinematics

${ }^{86} \mathrm{Kr}(\mathrm{d}, \mathrm{p}){ }^{87} \mathrm{Kr}$ iThemba
V. W. Ingeberg, Master thesis 2016
V.W.Ingeberg et al. Eur. Phys. J. A (2020) 56:68.

## 

## MINIBALL


+6 $\mathrm{LaBr}_{3}$




V. Ingeberg, to be submitted to PRC, Feb 2022. V. Ingeberg, PhD thesis to be submitted, Feb 2022


Calculations by A.C. Larsen, Univ. of Oslo Temperature ( $10^{9} \mathrm{~K}$ )

A.C. Larsen et al., PRC 97, 054329 (2018)



[^0]:    Future perspective: New neutron converter will increase expected yield as indicated in IS556. Test has been done Ref: J.P. Ramos et al., NIMB 463, 357 (2020).

