

## Onset of deformation at $N=60$

- Sr and Zr show rapid and dramatic onset of deformation at $N=60$.
- Smooth increase for the Kr isotopic chain ${ }^{1}$.
- Low-lying intruder configurations $\rightarrow$

${ }^{3}$ P.E. Garrett, M. Zielińska, and E. Clément, Prog. Part. Nucl. Phys. 163, 103931 (2021).
 shape coexistence ${ }^{2,3}$
${ }^{1}$ M. Albers, et al., Phys. Rev. Lett. 108, 062701 (2012).
2 J.E. García-Ramos and K. Heyde, Phys. Rev. C 100, 044315 (2019).


## Deformation-driving orbitals

- Proton excitations across $Z=40$ to $\pi g_{9 / 2}$ orbital.
- Ground-state configuration at $N=60$.
- Filling neutron orbitals lowers energy of $\pi g_{9 / 2}$.
- Large overlap of $\pi g_{9 / 2}$ and $v g_{7 / 2} \ldots$ (lesser extent $v h_{11 / 2}$ )
- Tensor force ${ }^{1} \rightarrow$ Type-II shell evolution ${ }^{2}$
- Close proximity of $v s_{1 / 2}, v d_{3 / 2}, v g_{7 / 2}, v h_{11 / 2}$ orbitals.

- Enhanced quadrupole interaction from coherent contributions of configurations $\rightarrow$ deformation
${ }^{1}$ P. Federman and S. Pittel, Phys. Lett. B 69, 385 (1977).
${ }^{2}$ T. Togashi, Y. Tsunoda, T. Otsuka, and N. Shimizu, Phys. Rev. Lett. 117, 172502 (2016).


## The neutron-rich krypton isotopes

- Evolution of shape change is not as dramatic as Zr and Sr , but occurs smoothly.
- Mean-field based models ${ }^{1}$ and IBM "mapping"2,3 show configuration mixing.
- Not present in Zr and Sr...
- How do the single-particle orbitals and their occupancies evolve in the Kr isotopes?
${ }^{1}$ T.R. Rodríguez, Phys. Rev. C 90, 034306 (2014).
${ }^{2}$ K. Nomura, et al., Phys. Rev. C 96, 034310 (2017).
${ }^{3}$ R.-B. Gerst, et al., Phys. Rev. C 105, 024302 (2022).


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${ }^{1}$ T. Rząca-Urban, et al., Phys. Rev. C 95, 064302 (2017).

## Odd-mass Kr isotopes

${ }^{2}$ G. Lhersonneau, et al., Phys. Rev. C 63, 034316 (2001).
${ }^{3}$ R.-B. Gerst, et al., Phys. Rev. C 105, 024302 (2022).

- Energy gap of $v d_{3 / 2}-v g_{7 / 2}$ is reduced with increasing $N$.
- Occupancy of each orbital drives change in $\pi g_{9 / 2}$ energy.
- $v g_{7 / 2}$ and $v h_{11 / 2}$ most important.


## Use one-neutron transfer:

- Determine $\ell$ of the states.
- Relative occupancies of key orbitals.
- Identify fragmentation of SP strength.
- Determine centroid - related to ESPEs.


But we have collectivity and fragmentation! Unidentified $h_{11 / 2}$ states?


UNiversity of

## Transfer reactions at ISS $-{ }^{92,94} \mathrm{Kr}(d, p)$

- "Standard" setup with new array, plus gas recoil detector.
- Kinematics with array at 60 mm and magnetic field $=2.05$ Tesla.
- Maximum possible beam energy from HIE-ISOLDE $\sim 7.5 \mathrm{MeV} / u$.
- DWBA calculations with PToLemy + global optical model parameters ${ }^{1,2}$.
- Monitor detector for $(d, d)$ normalisation $\rightarrow$ extraction of $C^{2} S$



${ }^{1}$ H. An and C. Cai, Phys. Rev. C 73, 054605 (2006).
${ }^{2}$ A.J. Koning and J.P. Delaroche, Nucl. Phys. A 713, 231 (2003).


## Beam intensity and availability

- Experience from runs in 2009, 2010 (REX) and 2018 (HIE).
- Requires Mo-free ion source to reduce contaminants.
- ${ }^{92,94} \mathrm{Kr}$ otherwise "easy", but ${ }^{96} \mathrm{Kr}$ has short half-life ( 80 ms ).
- Synchronisation of TRAP and EBIS, plus T1 timing in analysis.
- Extra shift requested to test this for future feasibility.
- TAC comments inline with expectations and proposed plan.

| Isotope | $T_{1 / 2}$ | Primary yield | Yield at ISS |
| :---: | :---: | :---: | :---: |
| ${ }^{92} \mathrm{Kr}$ | 1.84 s | $1.0 \times 10^{8}$ ions $/ \mu \mathrm{C}$ | $5.2 \times 10^{6} \mathrm{pps}$ |
| ${ }^{94} \mathrm{Kr}$ | 212 ms | $3.3 \times 10^{6} \mathrm{ions} / \mu \mathrm{C}$ | $1.7 \times 10^{5} \mathrm{pps}$ |
| ${ }^{96} \mathrm{Kr}$ | 0 | 1.9 | $10^{5}$ |

## Realistic Simulations - ${ }^{92,94} \mathrm{Kr}(d, p)$

- Aim to be sensitive to $C^{2} S>0.1$ in $g_{7 / 2}$ state.
- Identification of unobserved $h_{11 / 2}$ state at $C^{2} S>0.1$.
- Fragmentation of strength expected in higher-lying states.


NPTool: A. Matta et al. J. Phys. G 43, 045113 (2016)
ISS implementation: M. Labiche (STFC Daresbury)

## Realistic Simulations - ${ }^{92,94} \mathrm{Kr}(d, p)$




- Varying $C^{2} S$ to see effect on doublet fitting.
- Fitted peak width $\sim 120 \mathrm{keV}$ with $0.1 \mathrm{mg} / \mathrm{cm}^{2} \mathrm{CD}_{2}$ target.


## Summary and beam time request

- ${ }^{92,94} \mathrm{Kr}(d, p)$ at ISS to map singleparticle strength towards $N=60$.
- Smooth onset of deformation.



# Transfer reactions on the neutron-rich krypton isotopes 

## Thank you!


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## Backup slides will follow...

## ISS on-axis silicon array



## Gas ionisation detector

- Built, delivered and tested with beam in 2021.
- Blocker required to reduce direct/scattered beam.
- Trigger validation from array to reduce data rate, plus fast shapers (upgrades).



## Recent results $-{ }^{95} \mathrm{Kr}$

- Most recent results for ${ }^{95} \mathrm{Kr}$.
- $\gamma$-ray spectroscopy with DALI2@RIBF ${ }^{1}$.
- Isomeric nature of $\left(7 / 2^{+}\right)$state.
- Prompt-delayed correlations



Structure of krypton isotopes calculated with symmetry-conserving configuration-mixing methods
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## Quantum Phase Transition in the Shape of $\mathbf{Z r}$ isotopes

T. Togashi, Y. Tsunoda, T. Otsuka, and N. Shimizu Phys. Rev. Lett. 117, 172502 (2016).

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## Abrupt shape transition at neutron number $N=60: B(E 2)$ values in ${ }^{94,96,98}$ Sr from fast $\boldsymbol{\gamma} \boldsymbol{-} \boldsymbol{\gamma}$ timing

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Letter of Intent to the ISOLDE and Neutron Time-of-Flight Committee

## Shape transition at $N=60$ : Development of neutron-rich $\mathbf{S r}$ beams

September 28, 2021
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Figure 2: Calculated levels compared to (selected) experimentally known states. Levels are labeled with their energy and, where known, spin and parity. For the MCSM calculations prolate (oblate) states are marked in blue (orange). States with calculated triaxial deformation are shown in green. Note that for ${ }^{97} \mathrm{Sr}$ the triaxial degree of deformation could not yet be assessed in the calculations due to computational limitations.

## Analagous mechanism

- Large-scale MCSM calculations in $n$-deficient Hg isotopes.
- Occupancy of $v i_{13 / 2}$ changes along with $\pi h_{9 / 2}$.
- Fine energy balance of orbitals and oddneutron shifts energy enough to change configuration of the ground state.
- Can be probed with 1-neutron transfer from even-mass $\rightarrow$ confirm with small spectroscopic factors for corresponding states.

${ }^{1}$ B.A. Marsh, et al., Nature Physics 14, 1163 (2018).


## DWBA

- $d_{3 / 2}$ max. $\Leftrightarrow g_{7 / 2}$ min.
- And vice versa...
- Doublet fitting made easier.
- $s_{1 / 2}$ max. lower than range.
- No problem as $2^{\text {nd }}$ max. in range.
- $h_{11 / 2}$ requires max. energy.
- Difficult to get definitive $\ell$.


## Realistic Simulations - ${ }^{92} \mathrm{Kr}(d, p)$



- Varying $C^{2} S$ to see effect on doublet fitting.



## Realistic Simulations - ${ }^{94} \mathrm{Kr}(d, p)$




- Varying $C^{2} S$ to see effect on doublet fitting.


## Combined DWBA distributions - ${ }^{94} \mathrm{Kr}$




FIG. 16. Comparison of experimental (expt) spectroscopic factors $\left(C^{2} S\right)$ to those from shell model calculations carried out in model spaces (a), (b), and (c), see text. States are labeled by the neutron single-particle orbital populated in the transfer reaction.


SINGLE-PARTICLE STRUCTURE OF NEUTRON-RICH Sr
PHYSICAL REVIEW C 100, 054321 (2019)

TABLE IV. Comparison of ${ }^{2} \mathrm{H}\left({ }^{94,96} \mathrm{Sr}, p\right)$ spectroscopic factors to shell model calculations for low-lying states. The labels SM (a), (b) and (c) denote the three proton model spaces that were investigated (see text).

| Nucleus | $J^{\pi}$ | exp. |  | SM (a) |  | SM (b) |  | SM (c) |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $E(\mathrm{keV})$ | $C^{2} S$ | $E(\mathrm{keV})$ | $C^{2} S$ | $E(\mathrm{keV})$ | $C^{2} S$ | $E(\mathrm{keV})$ | $C^{2} S$ |
| ${ }^{95} \mathrm{Sr}$ | $\frac{1}{2}^{+}$ | 0 | 0.41(9) | 0 | 0.553 | 0 | 0.449 | 0 | 0.413 |
|  | $\frac{3}{2}+$ | 352 | 0.53(8) | 766 | 0.865 | 412 | 0.767 | 375 | 0.744 |
|  | $\frac{5}{2}+$ | 681 | 0.16(3) | 691 | 0.146 | 585 | 0.180 | 523 | 0.201 |
|  | $\frac{7}{2}^{+}$ | 556 |  | 1086 | 0.959 | 602 | 0.828 | 205 | 0.757 |
| ${ }^{97} \mathrm{Sr}$ | $\frac{1}{2}^{+}$ | 0 | 0.10(5) | 1631 | 0.013 | 1279 | 0.024 | 417 | 0.002 |
|  | $\frac{3}{2}^{+}$ | 167 | 0.25(5) | 0 | 0.881 | 0 | 0.804 | 117 | 0.713 |
|  | $\frac{7}{2}^{+}$ | 308 |  | 270 | 0.979 | 149 | 0.931 | 0 | 0.819 |
|  | $\frac{5}{2}^{+}$ | 522 | 0.13(5) | 1714 | 0.025 | 1336 | 0.042 | 57 | 0.000 |



FIG. 10. Doppler-corrected $\gamma$-ray energy spectra measured with DALI2 for the two-nucleon removal and one-neutron knockout reaction channels ${ }^{97} \mathrm{Rb}(p, 2 p n){ }^{95} \mathrm{Kr}$ and ${ }^{96} \mathrm{Kr}(p, p n){ }^{95} \mathrm{Kr}$. The spectra were fitted with simulated response functions (red) and a twocomponent exponential background (blue dashed curve).

## $\gamma$-ray spectroscopy of low-lying yrast and non-yrast states in neutron-rich ${ }^{94,95,96} \mathrm{Kr}$

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FIG. 9. Background-subtracted EURICA energy spectrum with a gate on ${ }^{95} \mathrm{Kr}$ in the ZeroDegree PID. The two peaks correspond to the two delayed transitions in ${ }^{95} \mathrm{Kr}$ at 81.7(2) and 113.8(2) keV depopulating the known $(7 / 2)^{+}$isomer $[29,40]$ as shown in the inset of the figure (see text for details).

