## MiNLO

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#### Outline

- MiNLO
  - Case for next-to-leading order calculations
  - Renormalization and factorization scales
  - Motivations for MiNLO
  - MiNLO example
  - Applications
  - MiNLO'

#### Case for next-to-leading order calculations

- Help looking for / setting limits on SUSY
- Take  $\tilde{g}\tilde{g}$  product<sup>n</sup>  $\rightarrow$  4 jets + MET
- Main bkg Z+4 jets [ 4 jet = 4  $\alpha_s$ 's ]





We prefer to know this bkg at NLO

#### Case for next-to-leading order calculations

- Precision Higgs physics
- $pp \rightarrow t\bar{t}H$  probes top Yukawa at tree level
- Has significant irreducible background from  $pp \rightarrow t\bar{t}b\bar{b}$





#### • We prefer to know this bkg at NLO

[Bredenstein, Denner, Dittmaier, Pozzorini PRL 103 (2009)]
[Bevilacqua, Czakon, Papadopoulos, Pittau, Worek JHEP 0909 (2009)]

#### Case for next-to-leading order calculations

• Recent years have seen amazing progress in NLO calculat<sup>n</sup>s:



- Increasing complexity brings increasing powers of  $\alpha_{S}(\mu)$
- More emphasis on choosing  $\mu$ 's carefully

 'Good scales' commonly considered to be so retrospectively on seeing higher order corr<sup>n</sup>s and residual scale sensitivity are small

 'Bad scales' commonly declared on finding large higher order corr<sup>n</sup>s & res. scale sensitivity : typically diagnosed as large unphysical scale logs

#### Q1: are large higher order corr<sup>n</sup>s all down to large $\mu_{R/F}$ logs?

 Big higher order corr<sup>n</sup>s can have real physical origins: new prod<sup>n</sup> channels, big colour factors, large gluon flux, I.R. logs ...



 Adjusting scale to make corr<sup>n</sup>s / sensitivity small can effectively 'eat' unrelated physics in scale choice

• In single/few scale processes it's harder to make a bad scale choice



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• In procs with more jets, i.e. more scales, it's hard to know what to do



- BSM background: W+3 jets [3 jets = 3  $\alpha$  s's]
- BlackHat paper points out physical dist<sup>n</sup>s can go -ve for  $\mu_R = \mu_F = E_{T,W}$



 For sufficiently poor choices (of scales) large logs can appear in some distributions, invalidating even an NLO prediction —— BlackHat collaboration



- Residual scale dependence of NLO calc<sup>n</sup>s goes as ~  $\alpha_{\rm S}^n \log^n \frac{\mu_R^2}{O^2}$  [n≥2]
- But often have situations were NNLO etc corr<sup>n</sup>s go as ~  $\alpha_{\rm S}^n \log^{2n} \frac{Q^2}{p_T^2}$  [n≥2]



 Is it reasonable to worry about single log terms beyond NLO & not also worry about IR Sudakov double logs?

#### MiNLO: Multi-scale improved Next-to-leading Order

Nason, Zanderighi, KH

- <u>In a nutshell:</u>
  - Determine the parton shower branching history associated with the kinematic config<sup>n</sup>s in the NLO events in the x-sec<sup>n</sup> integrals.
  - Take **all-orders rad corr<sup>n</sup>s** that the shower would associate to such config<sup>n</sup>s, as in CKKW<sup>\*</sup>, and **match** them with the exact NLO corr<sup>n</sup>s
- Addressing questions / objections of last few slides:
  - A1: small/moderate NLO corr<sup>n</sup>s/scale sensitivity isn't a consideration
  - A2: PS have natural uniquely defined scale setting for multi-scale probs
  - A3: PS resum large IR double logs as well as single scale logs
- Catani, Krauss, Kuhn, Webber

Ask how a parton shower would have generated this event:



- Emissions strongly ordered in hardness factorise from one another
- In PS each branching is like its own <u>simple</u> process with own scale
- Evaluating each  $\alpha_{S}(\mu)$  associated to a branching vertex at branching's own pT sums large class of higher order corr<sup>n</sup>s



extra coupling constant ratios:

 $\frac{\alpha s(q_2)}{\alpha s(\mu)}$ 



extra coupling constant ratios:

 $\frac{\alpha s(q_2)}{\alpha s(\mu)} \quad \frac{\alpha s(q_1)}{\alpha s(\mu)}$ 





 Showers also account for the probability that the gluons exiting gg→H vertex evolve from qm to q1 and q2 without emitting any resolvable radn (q>q2) : Sudakov form factors



associated with ratios of self-energies [ratios of self-energies in this cartoon]



factors are also needed to account for any external legs produced above q2 evolving down to q2



 Parton shower also effectively sets the factorisation scale for such a configuration to the scale beneath which all activity is integrated out, q2...

$$d \sigma_{H_{JJ}}^{LO} = d x_1 d x_2 d \Phi_{H_{JJ}} f_{h_1}(x_1, \mu_F) f_{h_2}(x_2, \mu_F) \frac{1}{2s} \mathcal{M}(\Phi_{H_{JJ}}, \mu_F)$$

$$LO \text{ terms} = O(\alpha_s^4)$$

$$d \sigma_{H_{JJ}}^{P.5.} = d x_1 d x_2 d \Phi_{H_{JJ}} f_{h_1}(x_1, \mu_F) f_{h_2}(x_2, \mu_F) \frac{1}{2s} \mathcal{M}^{P.5}(\Phi_{H_{JJ}})$$

$$\times \frac{\alpha s(q_2)}{\alpha s(\mu)} \cdot \frac{\alpha s(q_1)}{\alpha s(\mu)} \cdot \frac{\alpha^2 s(q_m)}{\alpha^2 s(\mu)}$$

$$\times \frac{\Delta_g(q_m; q_2)}{\Delta_g(q_2; q_2)} \cdots \frac{\Delta_g(q_1; q_2)}{\Delta_g(q_2; q_2)}$$

$$Beyond LO corrections$$

$$\sim 1 + O(\alpha s) + ...$$

$$\times \frac{f(x_1, q_2)}{f(x_1, \mu_F)} \cdot \frac{f(x_2, q_2)}{f(x_2, \mu_F)}$$

$$\begin{aligned}
\begin{aligned}
d & \bigcup_{H_{JJ}}^{LO} = \underbrace{dx_1 dx_2 d\Phi_{HJJ}}_{H_{JJ}} \underbrace{f_{u_1}(x_1, \mu_F)}_{u_1} \underbrace{f_{u_1}(x_2, \mu_F)}_{2S} \underbrace{\frac{1}{2S}}_{M(\Phi_{HJJ}, \mu_A)} \\
& LO terms = O(\alpha_S^4)
\end{aligned}$$

$$\begin{aligned}
& \bigcup_{H_{JJ}}^{P.S.} = \underbrace{dx_1 dx_2 d\Phi_{HJJ}}_{M_J} \underbrace{f_{u_1}(x_1, \mu_F)}_{u_1} \underbrace{f_{u_1}(x_2, \mu_F)}_{2S} \underbrace{\frac{1}{2S}}_{2S} \underbrace{M(\Phi_{HJJ}, \mu_A)}_{2S} \\
& \times \frac{\alpha S(q_2)}{\alpha S(\mu)} \cdot \frac{\alpha S(q_1)}{\alpha S(\mu)} \cdot \frac{\alpha^2 S(q_m)}{\alpha^2 S(\mu)} \\
& \times \frac{\Delta_g(q_m; q_2)}{\Delta_g(q_2; q_2)} \cdots \cdot \frac{\Delta_g(q_1; q_2)}{\Delta_g(q_2; q_2)} \\
& \times \frac{f(x_1, q_2)}{f(x_1, \mu_F)} \cdot \frac{f(x_2, q_2)}{f(x_2, \mu_F)}
\end{aligned}$$

 $d\sigma_{H_{JJ}} = dx_1 dx_2 d\Phi_{H_{JJ}} f_{h_1}(x_1, \mu_F) f_{h_2}(x_2, \mu_F) \frac{1}{2S} M(\Phi_{H_{JJ}}, \mu_R)$  $d\sigma_{H_{JJ}}^{\text{MiLO}} \equiv d\sigma_{H_{JJ}}^{\text{LO}} \times \frac{d\sigma_{H_{JJ}}^{\text{P.S.}}}{\left[d\sigma_{H_{JJ}}^{\text{P.S.}}\right]_{\text{LO}}}$  $\alpha_{\rm S}(q_2) \quad \alpha_{\rm S}(q_1) \quad \alpha_{\rm S}^2(q_m)$  $\overline{\alpha_{\rm S}(\mu)}$   $\alpha_{\rm S}(\mu)$   $\alpha_{\rm S}^2(\mu)$  $= \int \Delta_g(q_m; q_2) \qquad \Delta_g(q_1; q_2)$  $\Delta_g(q_2; q_2)$   $\Delta_g(q_2; q_2)$  $f(x_1,q_2) f(x_2,q_2)$  $f(x_1, \mu_F) f(x_2, \mu_F)$ 

#### Example: H+2 jets MiNLO at next-to-leading order

 To extend from LO `MiLO' example to NLO apply the same all orders shower corr<sup>n</sup>s to the conventional NLO HJJ computation

![](_page_25_Figure_2.jpeg)

• And subtract a term to render the expansion in  $\alpha$  s unchanged to NLO .....

#### Application: H+1 jet MiNLO at next-to-leading order

![](_page_26_Figure_1.jpeg)

- Resummation matched to NLO inclusive  $gg \rightarrow H xsec^n$  [ = 1 in ratios ]
- HJ RUN: NLO H + 1 jet with  $\mu_R = \mu_F = p_{T,H}$
- HJ FXD: NLO H + 1 jet with  $\mu_R = \mu_F = M_H$
- HJ-MiNLO<sub>Ratio</sub> conventional NLO H + 1 jet at high pT
- HJ-MiNLO,→, resummation result at low p⊤
- $H_J-M_HO \rightarrow sensible scale unc. band (doesn't shrink as <math>p_T \rightarrow 0$ )

 $\varphi_{j1j2}$ 

#### Application: Z+2 jet MiNLO ⊕ Pythia vs ATLAS

![](_page_27_Figure_2.jpeg)

- Left: improves Z + 2 NLO s.t. gives even predict<sup>n</sup> for  $\geq 0$  jet evts!
- Right: NLO accuracy retained [& improved] for ≥ 2 jet events
- Equally nice improvement & agreement for ATLAS W+jets data

[ Campbell, Ellis, Nason, Zanderighi]

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![](_page_28_Picture_5.jpeg)

#### MiNLO' for H+1-jet

- MiNLO matches fully differential NLO to LL [NLL<sub>o</sub>] resummation
- MiNLO finite in all ph.space: no need of gen. cuts

![](_page_29_Figure_3.jpeg)

Question: what's MiNLO accuracy for inclusive quantities?

- H+1-jet spectrum known analytically to high accuracy
- Allows to determine Sudakov s.t. H+1-jet is NLO for H+0-jet

$$\begin{split} & \mathsf{MiNLO} \to \mathsf{MiNLO'} \\ & \Delta(Q, p_{\mathrm{T}}) \to \Delta'(Q, p_{\mathrm{T}}) = \Delta(Q, p_{\mathrm{T}}) \ \delta\Delta(Q, p_{\mathrm{T}}) \\ & \delta\Delta(Q, p_{\mathrm{T}}) = \exp\left[\int_{0}^{L} dL' \ \bar{\alpha}_{\mathrm{S}}^{2}(p_{\mathrm{T}}') \left[\widetilde{R}_{21} L' + \widetilde{R}_{20} - \bar{\beta}_{0} \mathcal{H}_{1}\right]\right] \\ & \frac{d\sigma'_{\mathcal{M}}}{d\Phi} = \int_{L} \frac{d\sigma'_{\mathcal{R}}}{d\Phi dL} + \int_{L} \frac{d\sigma_{\mathcal{F}}}{d\Phi dL} = \frac{d\sigma_{\mathrm{NLO}}}{d\Phi} \end{split}$$

MiNLO' simultaneously NLO for H & H+1-jet prod<sup>n</sup>

Higgs rapidity

![](_page_31_Figure_2.jpeg)

- Conventional NLO H prod<sup>n</sup>: red
- MiNLO' H+1-jet+parton shower: green
- Agree with each other ~ to within the line thickness

- Conventional scale setting in complex processes subject to large ambiguities, and lacking physical justifications
- MiNLO is a physically well motivated scale setting for processes with jets: identifies all important scales, treats all coherently
- MiNLO taken up by independent groups: applied in complex processes, used in new NLO merging schemes [e.g. FxFx]
- MiNLO' = NLO x NLO calc<sup>n</sup>s w.o. merging scales
- MiNLO' extended to (N)NLO x NLO x NLO accuracy in H+2 jets

![](_page_33_Picture_1.jpeg)

MiNLO pgs. 394 - 398

#### Thank you Paolo

![](_page_34_Picture_1.jpeg)