

Extreme plasma physics in the laboratory and in astrophysics

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Acknowledgements

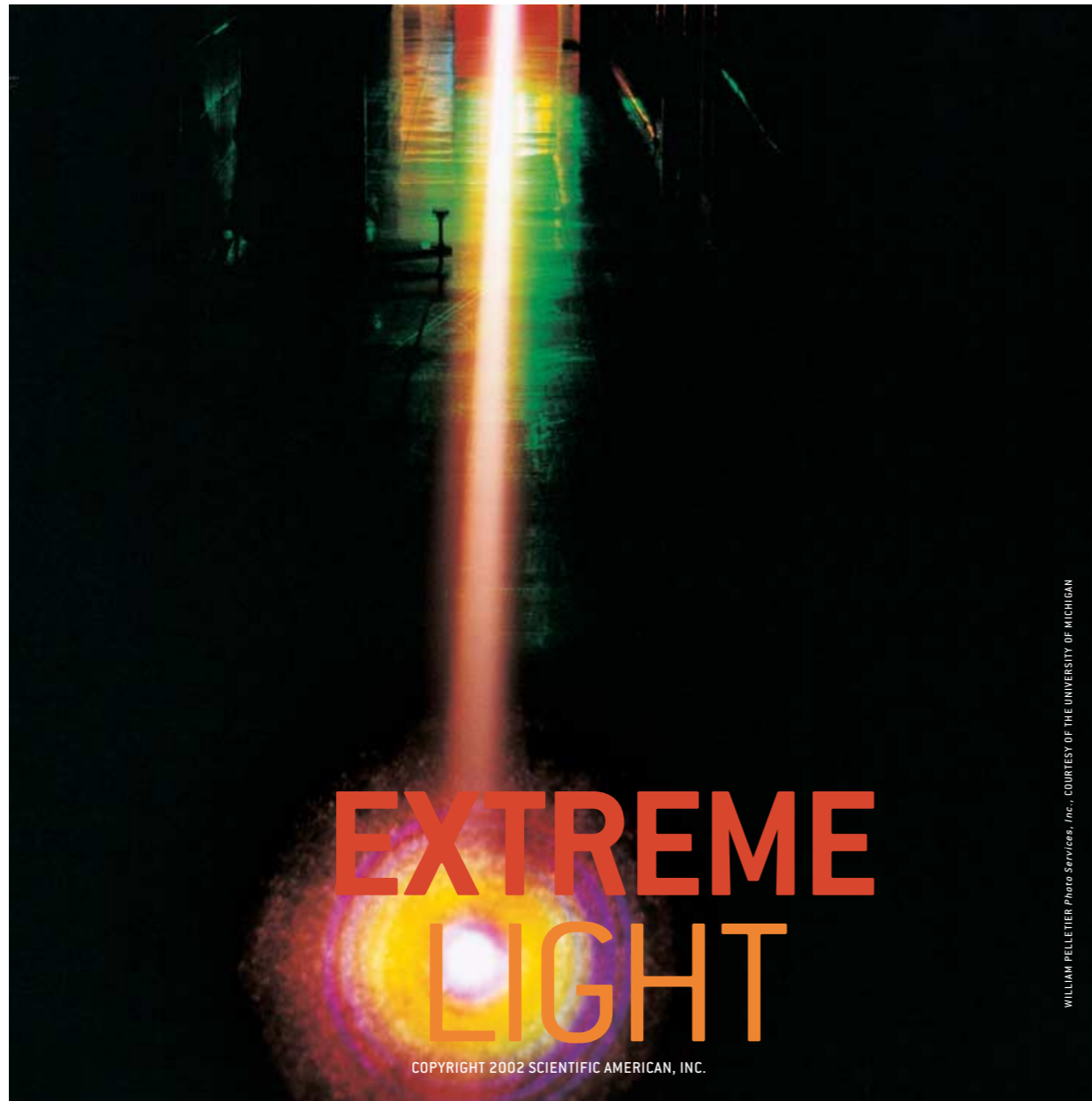
**R. A. Fonseca, T. Grismayer, J. Vieira, K. Schoeffler,
M. Vranic, B. Martinez, F. Cruz, F. Del Gaudio**

**W.B.Mori, R. Bingham, D. Uzdensky, A. Spitkovsky,
G. Gregori**

Simulations performed at **Accelerates (IST), Dawson2 (UCLA), Jugene/Juqueen (FZ Jülich), SuperMUC (Münich), MareNostrum (BSC), Sequoia (LLNL)**

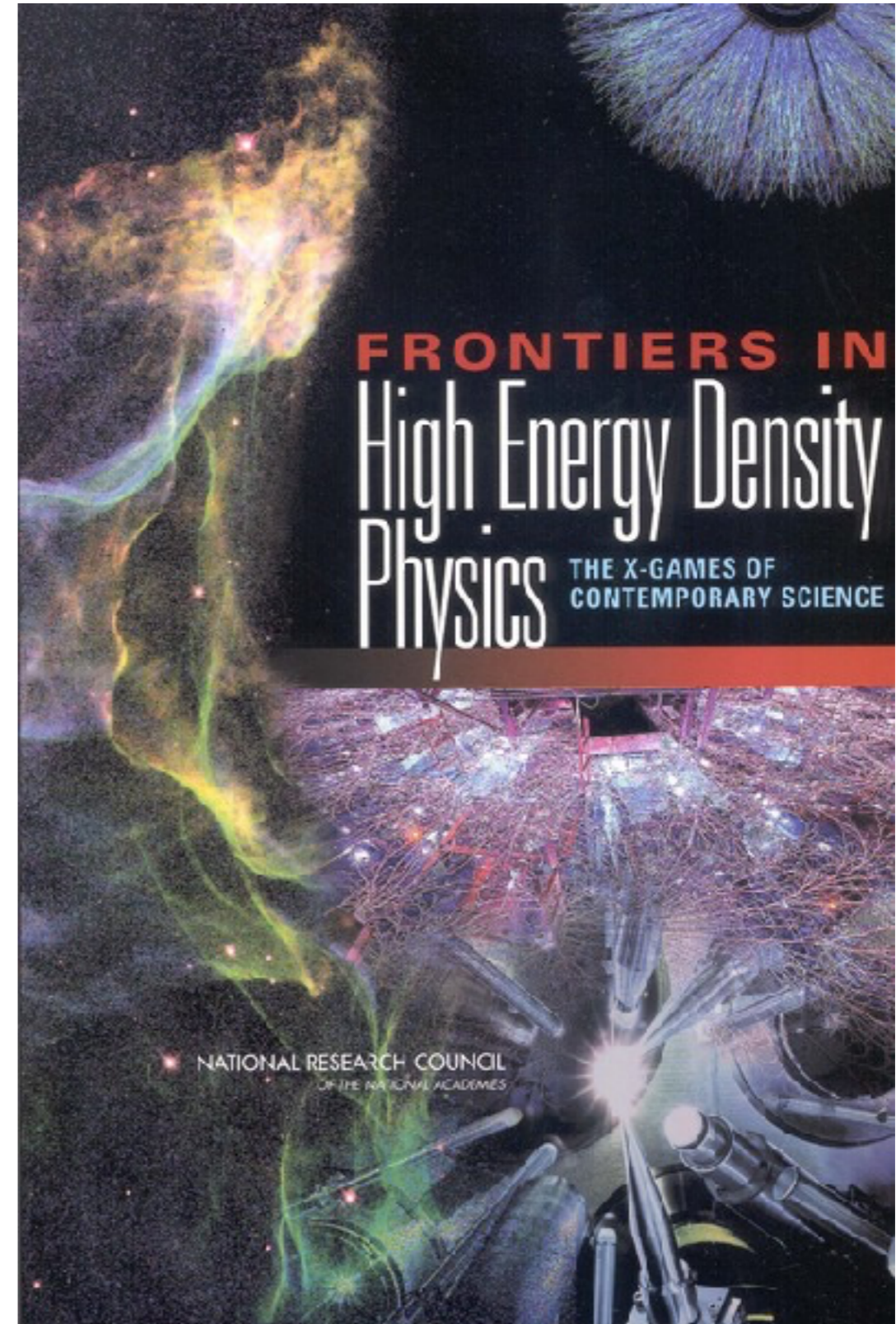


GoLP | golp.tecnico.ulisboa.pt
Committed to continuously raising the bar



Focusing light with the power of
1,000 HOOVER DAMS onto a point the size of
A CELL NUCLEUS accelerates electrons to the speed
of light in a femtosecond

By Gérard A. Mourou and Donald Umstadter

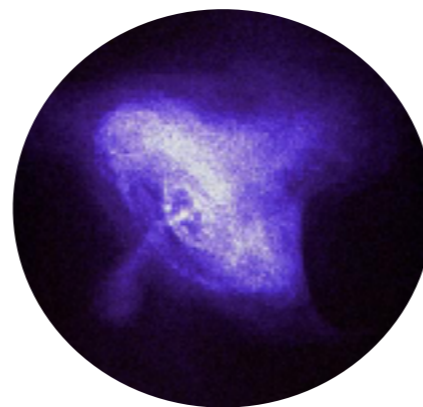


Plasma physics supplemented by several additional physical effects that
“...would be considered “exotic” in traditional plasma physics e.g.
special-relativistic effects relativistically hot plasmas and relativistic bulk motions
radiation-reaction effects e.g., synchrotron or inverse-Compton radiative cooling
electron-positron pair creation
ultra-strong magnetic fields QED effects such as 1-photon pair creation
general-relativistic effects.”

D. Uzdensky *et al.*, Extreme Plasma Astrophysics, arXiv:1903.05328



**Ultra intense
laser and
particle beams**



**Relativistic
Astrophysics**

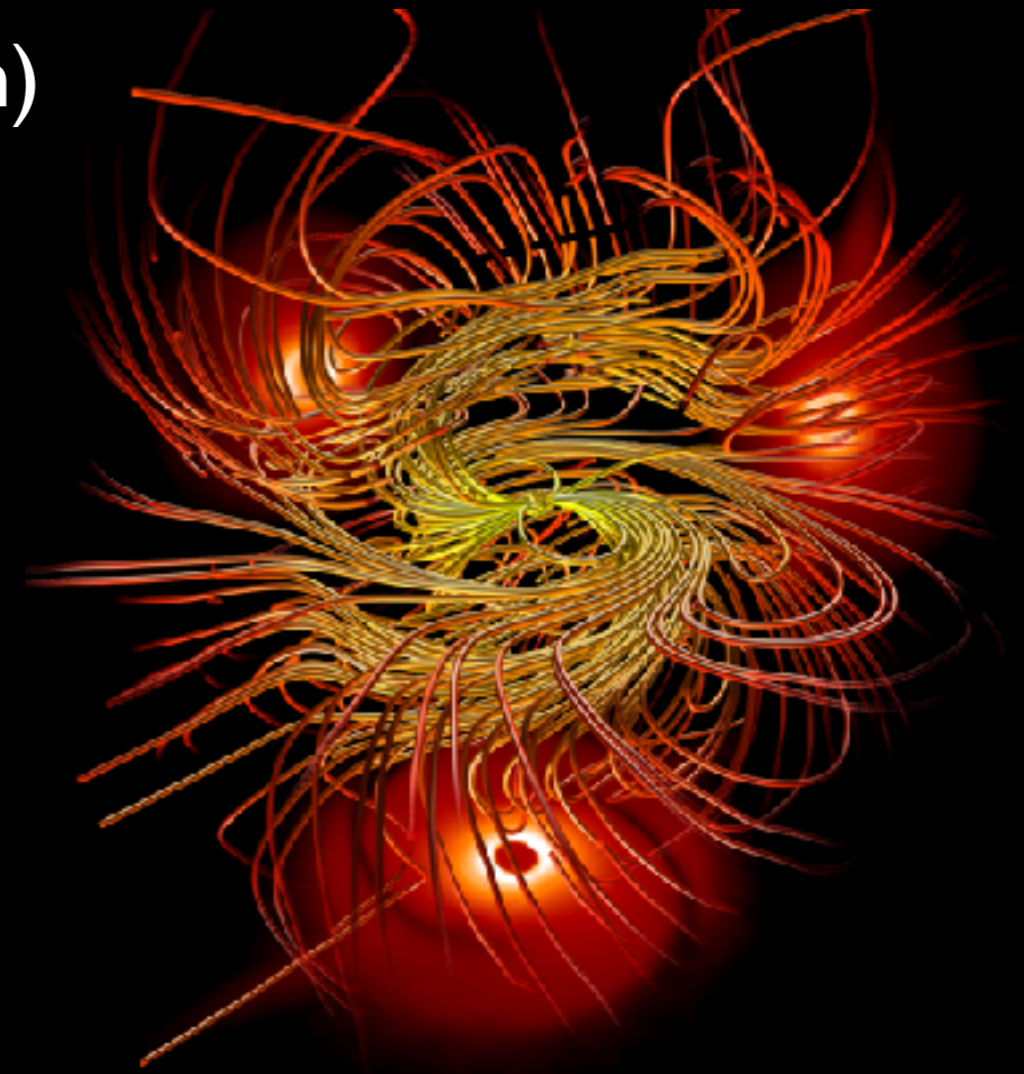


Sci-fi

What is the behavior of matter (and vacuum) at the intensity frontier ($> 10^{23} \text{ W/cm}^2$) in lab and in astro?

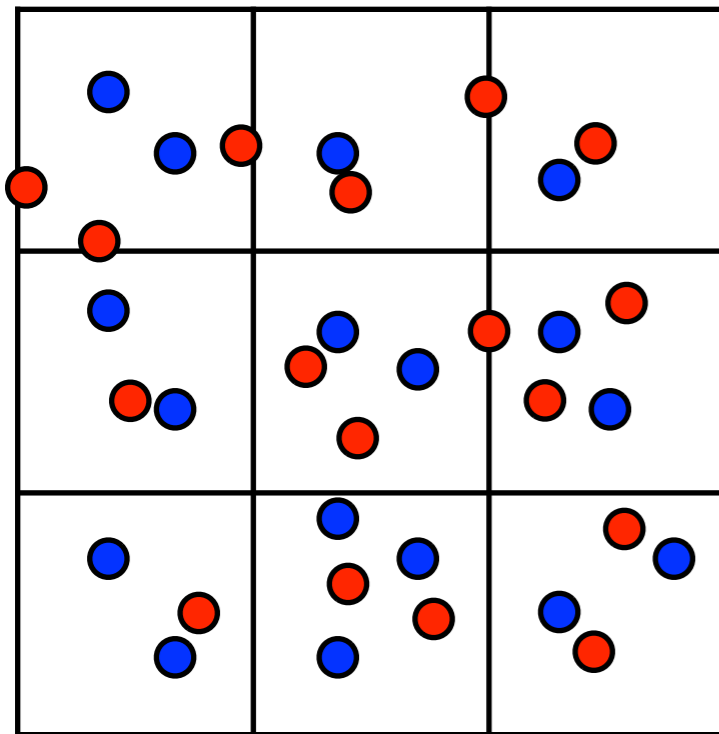
How can we understand and explore the complex and nonlinear behavior with a combination of simulations + theory + experiments with lasers and beams?

epp.tecnico.ulisboa.pt



Particle-in-cell simulations

Solving Maxwell's equations on a grid with self-consistent charges and currents due to charged particle dynamics



State-of-the-art

$\sim 10^{12}$ particles
 $\sim (12000)^3$ cells

RAM \sim 1 Gbyte - 100 TByte

Run time: hours to months

Data/run \sim few MB - 100s TByte

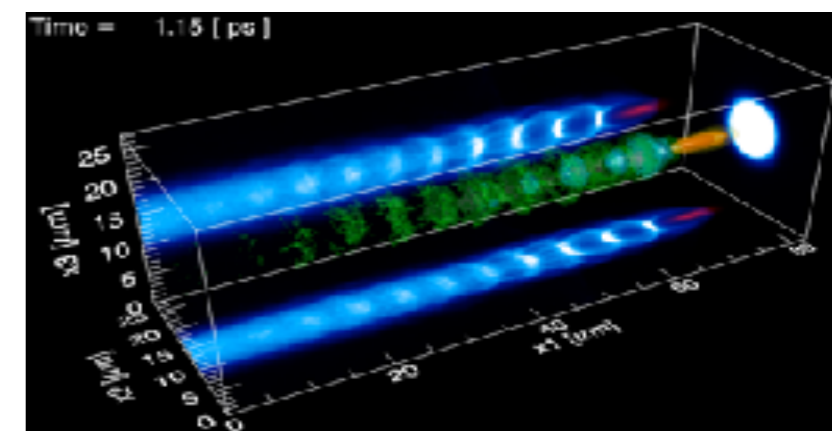
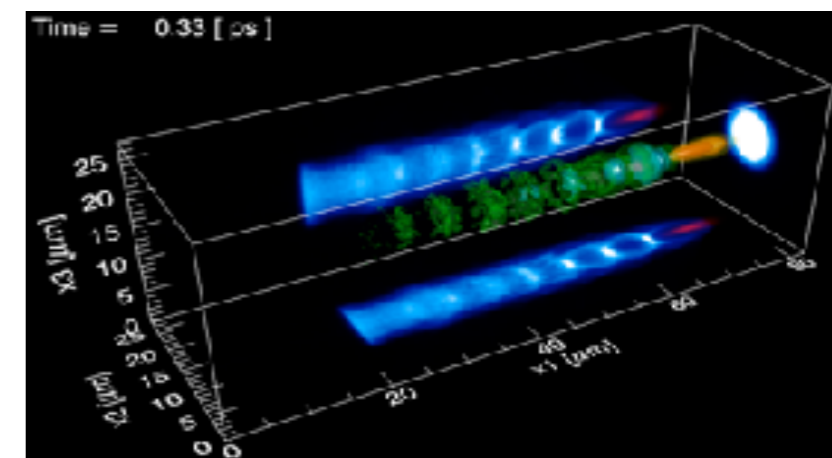
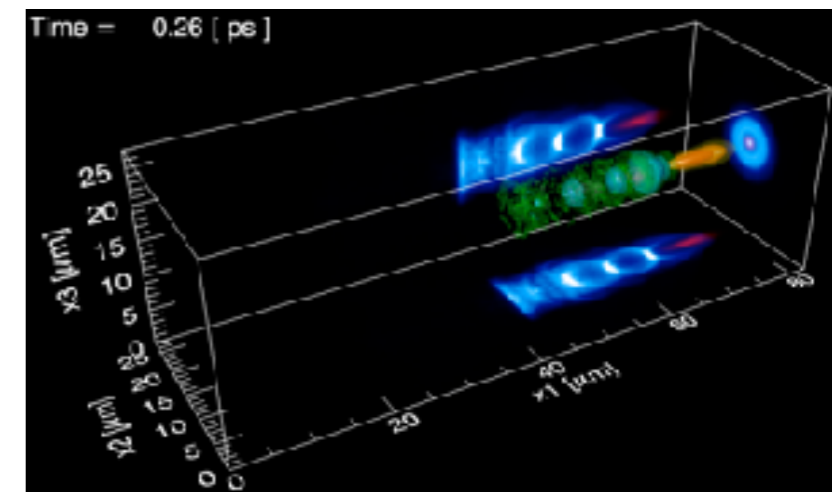
One-to-one simulations of plasma based accelerators & cluster dynamics

Weibel/two stream instability in astrophysics, relativistic shocks, fast igniton/inertial fusion energy, low temperature plasmas

Particle-in-cell (PIC) - (Dawson, Buneman, 1960's)

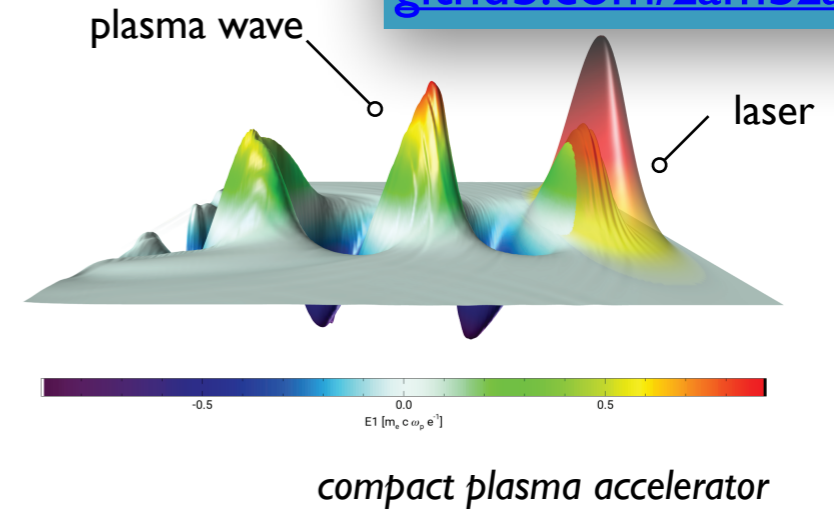
Maxwell's equation solved on simulation grid

Particles pushed with Lorentz force



- **Particle-in-cell Code suite.** Fully relativistic electro-magnetic 1D and 2D (FDTD/spectral) and 1D Electrostatic.
- **Python interface.** All simulation codes/parameters can be controlled via Python interface.
- **Educational examples.** Set of Python Jupyter notebooks with detailed physics problem description and simulation setup.

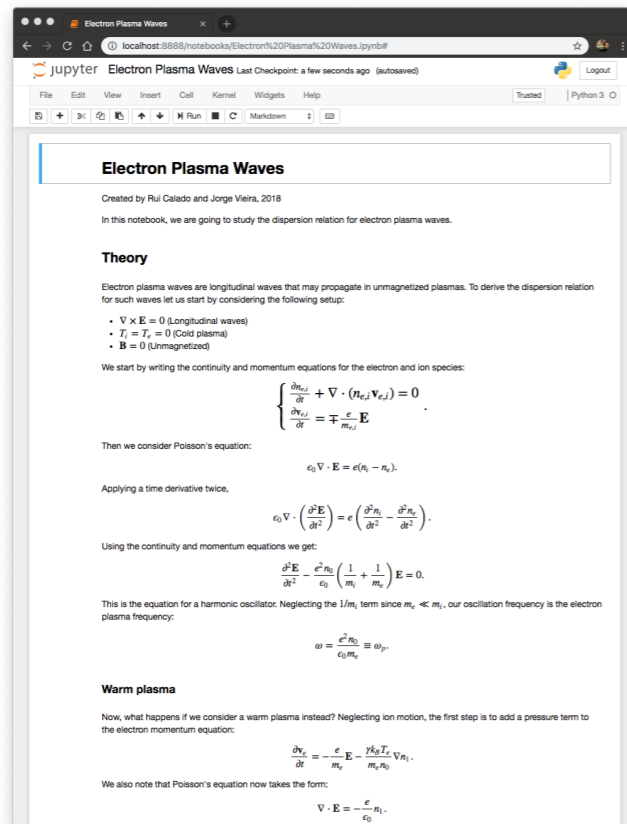
Come find us on GitHub
github.com/zambzamb/zpic



Example notebook



Theoretical introduction



Electron Plasma Waves

Created by Rui Calado and Jorge Vieira, 2018

In this notebook, we are going to study the dispersion relation for electron plasma waves.

Theory

Electron plasma waves are longitudinal waves that may propagate in unmagnetized plasmas. To derive the dispersion relation for such waves let us start by considering the following setup:

- $\nabla \times \mathbf{E} = 0$ (Longitudinal waves)
- $T_i = T_e = 0$ (Cold plasma)
- $\mathbf{B} = 0$ (Unmagnetized)

We start by writing the continuity and momentum equations for the electron and ion species:

$$\begin{cases} \frac{\partial n_{e,i}}{\partial t} + \nabla \cdot (n_{e,i} \mathbf{v}_{e,i}) = 0 \\ \frac{\partial \mathbf{v}_{e,i}}{\partial t} = \frac{e}{m_{e,i}} \mathbf{E} \end{cases}$$

Then we consider Poisson's equation:

$$\epsilon_0 \nabla \cdot \mathbf{E} = e(n_e - n_i)$$

Applying a time derivative twice,

$$\epsilon_0 \nabla \cdot \left(\frac{\partial^2 \mathbf{E}}{\partial t^2} \right) = e \left(\frac{\partial^2 n_e}{\partial t^2} - \frac{\partial^2 n_i}{\partial t^2} \right)$$

Using the continuity and momentum equations we get:

$$\frac{\partial^2 \mathbf{E}}{\partial t^2} - \frac{e^2 n_0}{\epsilon_0} \left(\frac{1}{m_e} + \frac{1}{m_i} \right) \mathbf{E} = 0$$

This is the equation for a harmonic oscillator. Neglecting the $1/m_i$ term since $m_e \ll m_i$, our oscillation frequency is the electron plasma frequency:

$$\omega = \frac{e^2 n_0}{\epsilon_0 m_e} \equiv \omega_p$$

Warm plasma

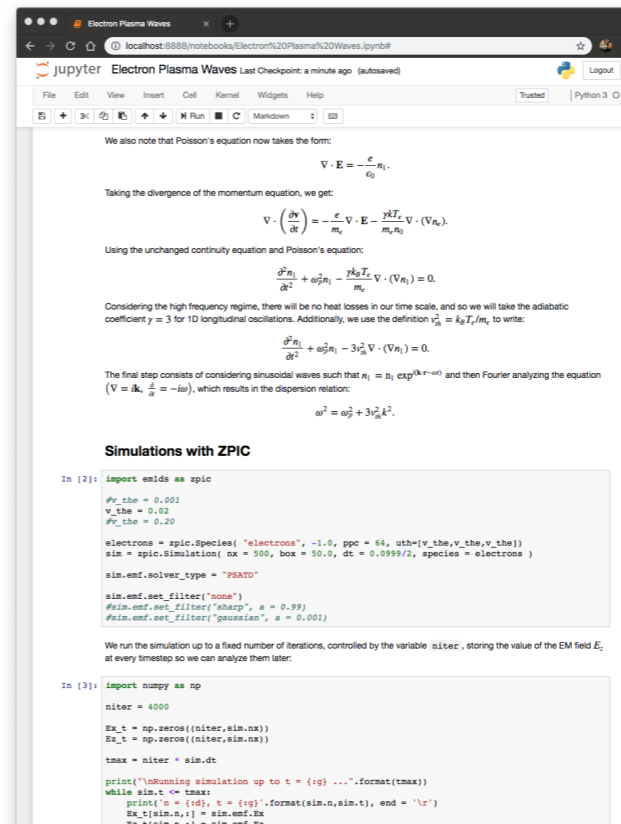
Now, what happens if we consider a warm plasma instead? Neglecting ion motion, the first step is to add a pressure term to the electron momentum equation:

$$\frac{\partial \mathbf{v}_e}{\partial t} = -\frac{e}{m_e} \mathbf{E} - \frac{\gamma k_B T_e}{m_e n_0} \nabla n_e$$

We also note that Poisson's equation now takes the form:

$$\nabla \cdot \mathbf{E} = -\frac{e}{\epsilon_0} n_e$$

simulation initialisation



We also note that Poisson's equation now takes the form:

$$\nabla \cdot \mathbf{E} = -\frac{e}{\epsilon_0} n_e$$

Taking the divergence of the momentum equation, we get:

$$\nabla \cdot \left(\frac{\partial \mathbf{v}}{\partial t} \right) = -\frac{e}{m_e} \nabla \cdot \mathbf{E} - \frac{\gamma k_B T_e}{m_e n_0} \nabla \cdot (\nabla n_e)$$

Using the unchanged continuity equation and Poisson's equation:

$$\frac{\partial^2 n_e}{\partial t^2} + \omega_p^2 n_e - \frac{\gamma k_B T_e}{m_e} \nabla^2 (\nabla n_e) = 0$$

Considering the high frequency regime, there will be no heat losses in our time scale, and so we will take the adiabatic coefficient $\gamma = 3$ for 1D longitudinal oscillations. Additionally, we use the definition $v_{th}^2 = k_B T_e / m_e$ to write:

$$\frac{\partial^2 n_e}{\partial t^2} + \omega_p^2 n_e - 3v_{th}^2 \nabla^2 (\nabla n_e) = 0$$

The final step consists of considering sinusoidal waves such that $n_e = n_1 \exp^{ikx - i\omega t}$ and then Fourier analyzing the equation ($\nabla = ik$, $\frac{\partial}{\partial t} = -i\omega$), which results in the dispersion relation:

$$\omega^2 = \omega_p^2 + 3v_{th}^2 k^2$$

Simulations with ZPIC

```
In [2]: import emids as spic
p_r_0 = 0.001
v_the = 0.02
p_r_0 = 0.20

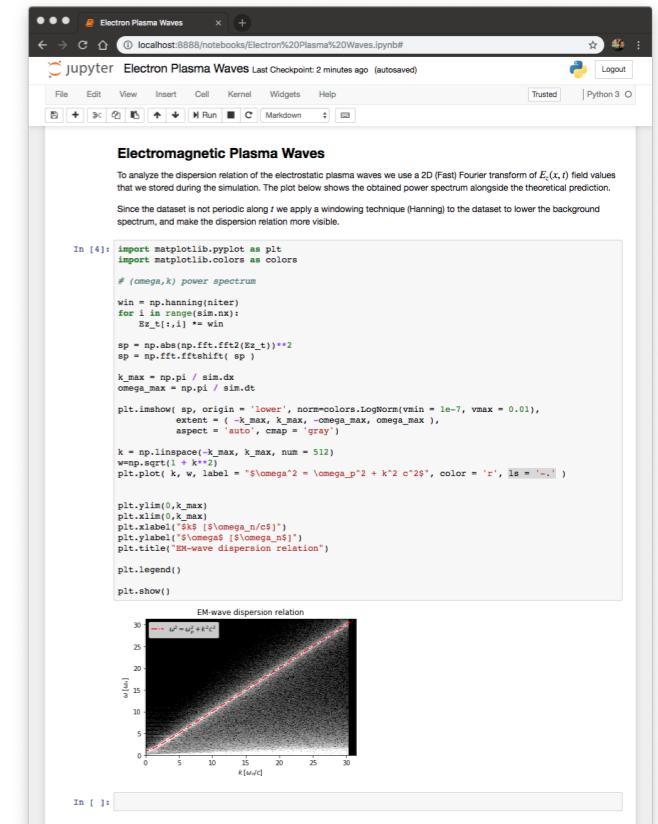
electrons = spic.Species("electrons", -1.0, ppc = 64, utb=[v_the, v_the, v_the])
sim = spic.Simulation(nx = 500, box = 50.0, dt = 0.0999/2, species = electrons)
sim.emf_solver_type = "PSATD"
sim.emf_set_filters("none")
sim.emf_set_filter("sharp", a = 0.99)
sim.emf_set_filter("gaussian", a = 0.001)

We run the simulation up to a fixed number of iterations, controlled by the variable niter, storing the value of the EM field E_e at every timestep so we can analyze them later.

In [3]: import numpy as np
niter = 4000
E_x_t = np.zeros((niter, sim.nx))
E_z_t = np.zeros((niter, sim.nx))
tmax = niter * sim.dt

print("\nRunning simulation up to t = {}s ...".format(tmax))
while sim.t <= tmax:
    print("t = {}s".format(sim.t), end = '\r')
    E_x_t[sim.n, :] = sim.emf.E_x
    E_z_t[sim.n, :] = sim.emf.E_z
```

analysis and questions for discussion



Electromagnetic Plasma Waves

To analyze the dispersion relation of the electrostatic plasma waves we use a 2D (Fast) Fourier transform of $E_e(x, t)$ field values that we stored during the simulation. The plot below shows the obtained power spectrum alongside the theoretical prediction.

Since the dataset is not periodic along t we apply a windowing technique (Hanning) to the dataset to lower the background spectrum, and make the dispersion relation more visible.

```
In [4]: import matplotlib.pyplot as plt
import matplotlib.colors as colors
# (omega, k) power spectrum
win = np.hanning(niter)
for i in range(sim.nx):
    E_x_t[i, :] *= win
sp = np.abs(np.fft.fft2(E_x_t))**2
sp = np.fft.fftshift(sp)

k_max = np.pi / sim.dx
omega_max = np.pi / sim.dt

plt.imshow(sp, origin = 'lower', norm=colors.LogNorm(vmin = 1e-7, vmax = 0.01),
          extent = (-k_max, k_max, -omega_max, omega_max),
          aspect = 'auto', cmap = 'gray')

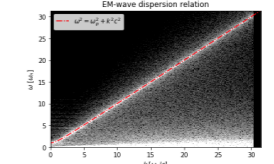
k = np.linspace(-k_max, k_max, num = 512)
wpsqrt = np.sqrt(1 + k**2)
plt.plot(k, wpsqrt, label = "$\omega_{\text{the}} = \sqrt{\omega_p^2 + k^2 c^2}$", color = 'r', ls = '--')

plt.ylim(0, k_max)
plt.xlim(0, omega_max)
plt.xlabel("$k$ [$\omega_{\text{pe}}/c$]")
plt.ylabel("$\omega$ [$\omega_{\text{pe}}$]")
plt.title("EM-wave dispersion relation")
plt.legend()

plt.show()

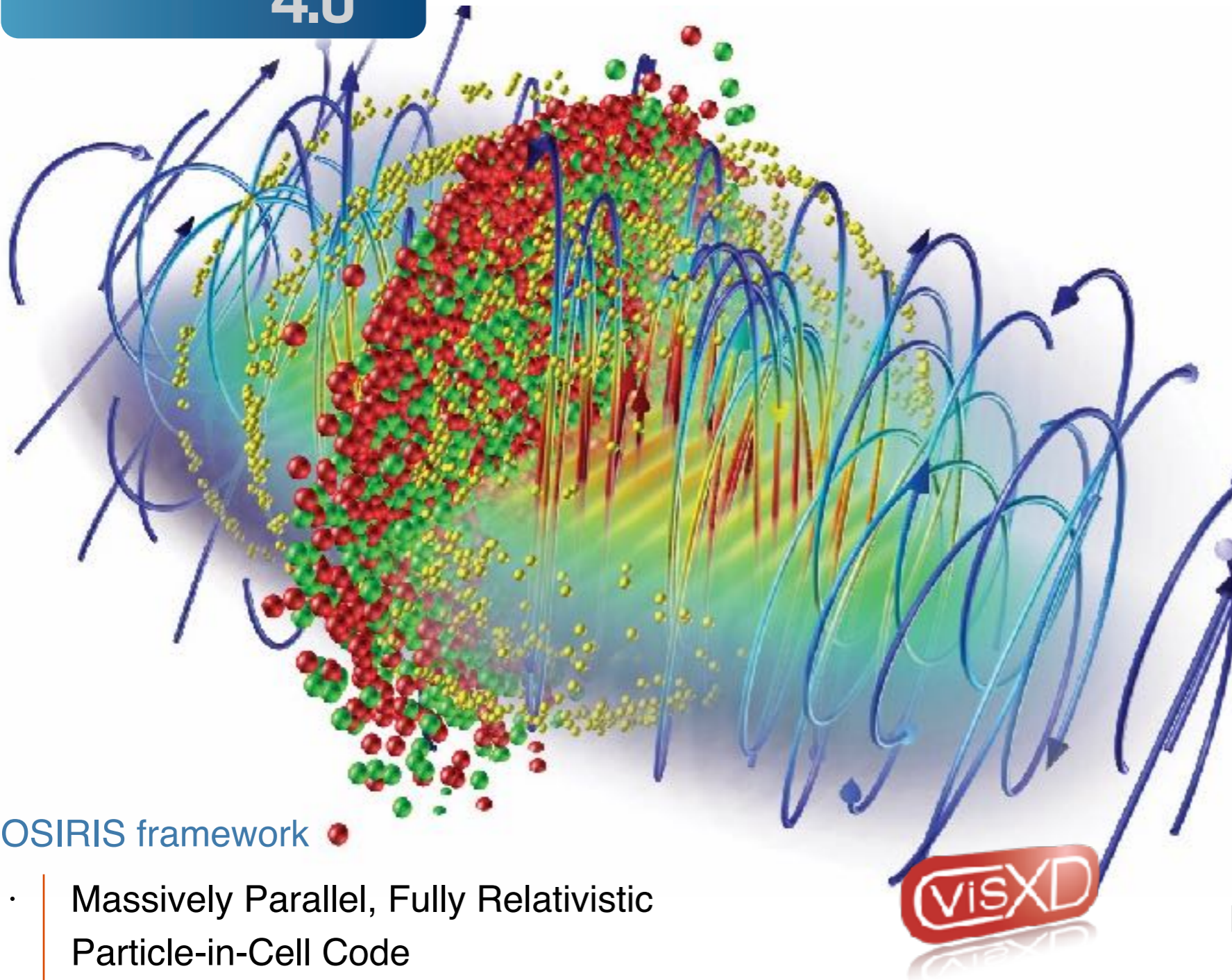
In [ ]:
```

EM-wave dispersion relation





Committed to open science



Open-access model

- 40+ research groups worldwide are using OSIRIS
- 250+ publications in leading scientific journals
- Large developer and user community
- Detailed documentation and sample inputs files available

Using OSIRIS 4.0

- The code can be used freely by research institutions after signing an MoU
- Find out more at:

<http://epp.tecnico.ulisboa.pt/osiris>

OSIRIS framework

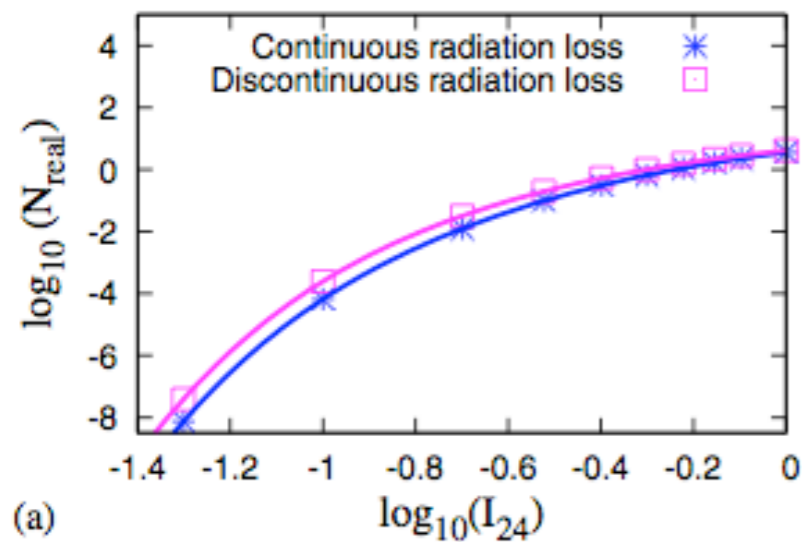
- Massively Parallel, Fully Relativistic Particle-in-Cell Code
- Parallel scalability to 2 M cores
- Explicit SSE / AVX / QPX / Xeon Phi / CUDA support
- Extended simulation/physics models



Ricardo Fonseca: ricardo.fonseca@tecnico.ulisboa.pt

Monte Carlo simulations demonstrating pair production via real photons per electron

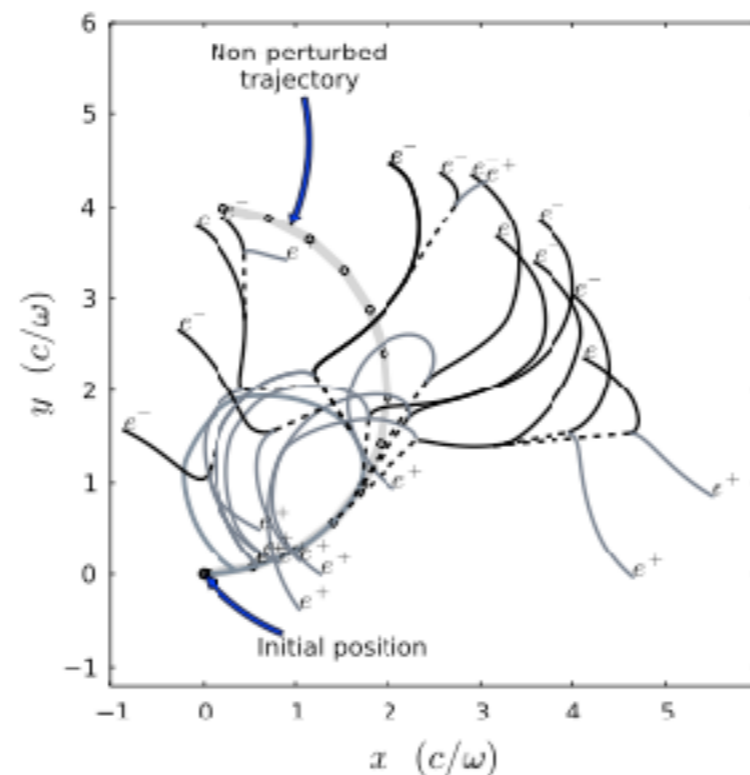
J. G. Kirk, A. R. Bell, and I. Arka, PPCF (2009)
 R. Ducloux, J. G. Kirk & A. R. Bell, PPCF (2010)



Number of pairs produced

PIC simulations of QED cascade in various configuration (counter propagating laser, rotating field)

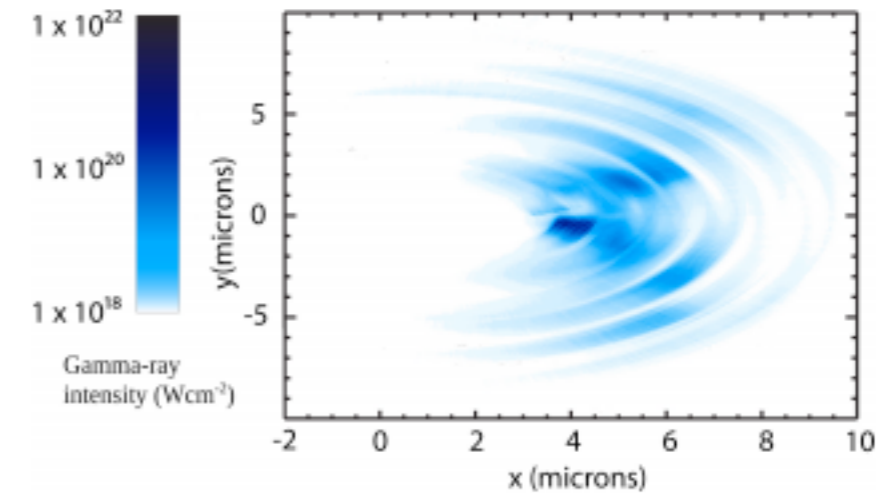
N.V. Elkina et al, Phys. Rev. ST.AB (2011)
 E.N. Nerush, et al, Phys. Rev. Lett. (2011)



Cascade in rotating field

Dense pair Plasmas and Ultra-Intense Bursts of Gamma-Rays from Laser-Irradiated Solids

C.P. Ridgers, et al Phys. Rev.Lett., (2012)



Gamma rays from laser-irradiated solid

QED-PIC loop + Particle Merging

$$\frac{d\mathbf{p}}{dt} = \mathbf{F}_L + \frac{d\mathcal{P}_\gamma}{dt d\chi} \quad \text{Probabilistic}$$



Integration of equations of motion:
moving particles
 $\mathbf{F}_p \rightarrow \mathbf{u}_p \rightarrow \mathbf{x}_p$

Emission of photons
Probability of pair creation
→ new particles

Interpolation:
evaluating force on particles
 $(\mathbf{E}, \mathbf{B})_i \rightarrow \mathbf{F}_p$



Deposition:
calculating current on grid
 $(\mathbf{x}, \mathbf{u})_p \rightarrow \mathbf{j}_i$

Particle Merging

Integration of field equations:
updating fields
 $(\mathbf{E}, \mathbf{B})_i \leftarrow \mathbf{J}_i$

M.Vranic, et al.,
Comput. Phys. Commun. (2015)

$$\frac{\partial \mathbf{E}}{\partial t} = c \nabla \times \mathbf{B} - 4\pi \mathbf{j}$$

$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E}$$

For a review
A. Gonoskov et al., PRE (2015)

The emergence of (relativistic) quantum behaviour with intense fields

Schwinger field

$$E_s = \frac{m^2 c^3}{e \hbar} \quad E_s \simeq 1.32 \times 10^{18} \text{ V/cm}$$

Pair creation probability

$$W \propto \exp(-\pi E_s / E)$$

Normalized electric field

$$\chi = \frac{E}{E_s}$$



Julian Schwinger, 1918-1994

Generalization for any Lorentz frame

$$\chi = \frac{1}{E_s} \sqrt{(\gamma \mathbf{E} + \frac{\mathbf{p}}{mc} \times \mathbf{B})^2 - (\frac{\mathbf{p}}{mc} \cdot \mathbf{E})^2}$$

Lower E still leads to pair creation due to Lorentz boost

$$\chi \simeq \frac{\gamma E_{\perp}}{E_s}$$

For reviews:

GA Mourou, T Tajima, SV Bulanov, RMP (2002); M. Marklund and P. K. Shukla, RMP (2006);

A Di Piazza, C Müller, KZ Hatsagortsyan, CH Keitel RMP (2012)

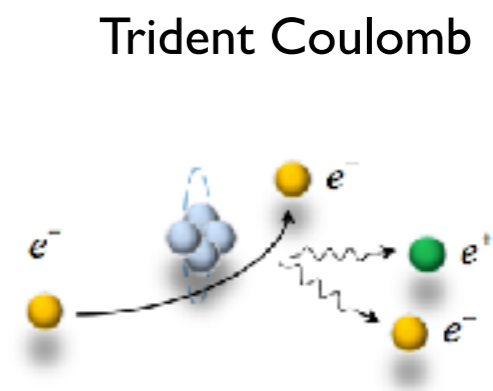
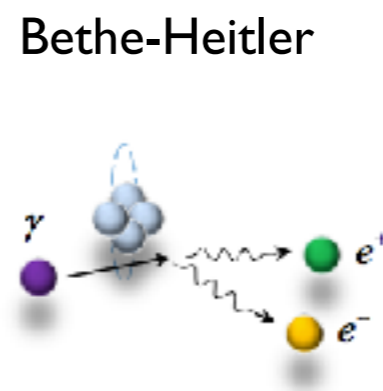
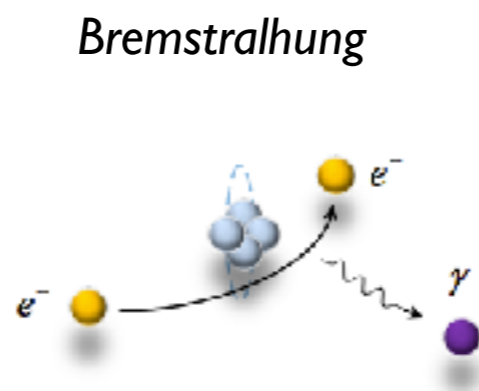
Plethora of QED processes

Incoherent processes

- Pair production
- Photon production
- Photon annihilation
- “Comptonisation”

$$\sigma \sim \alpha^n \sigma_T \times f(\mathcal{E})$$

Cross section

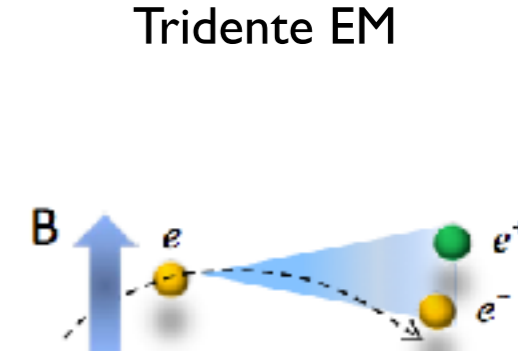
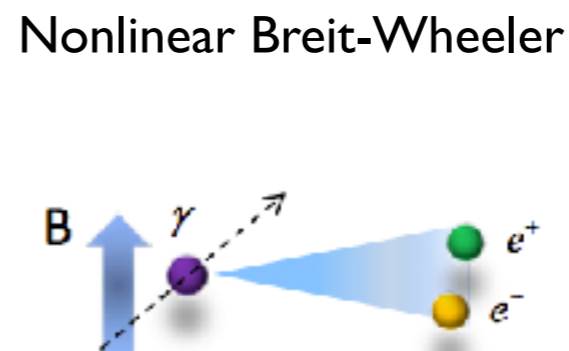
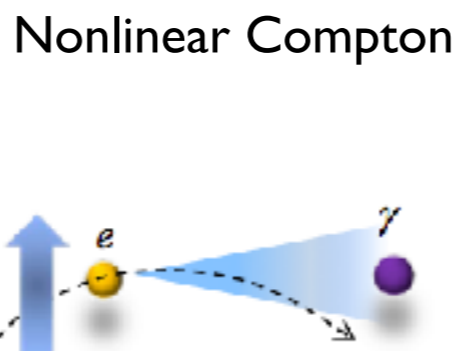


Coherent processes

- Nonlinear Breit-Wheeler
- Nonlinear Compton
- Schwinger mechanism
- Photon splitting

$$d\mathcal{P}/dt \sim (\alpha c/\lambda_C) f(E/E_S, \mathcal{E})$$

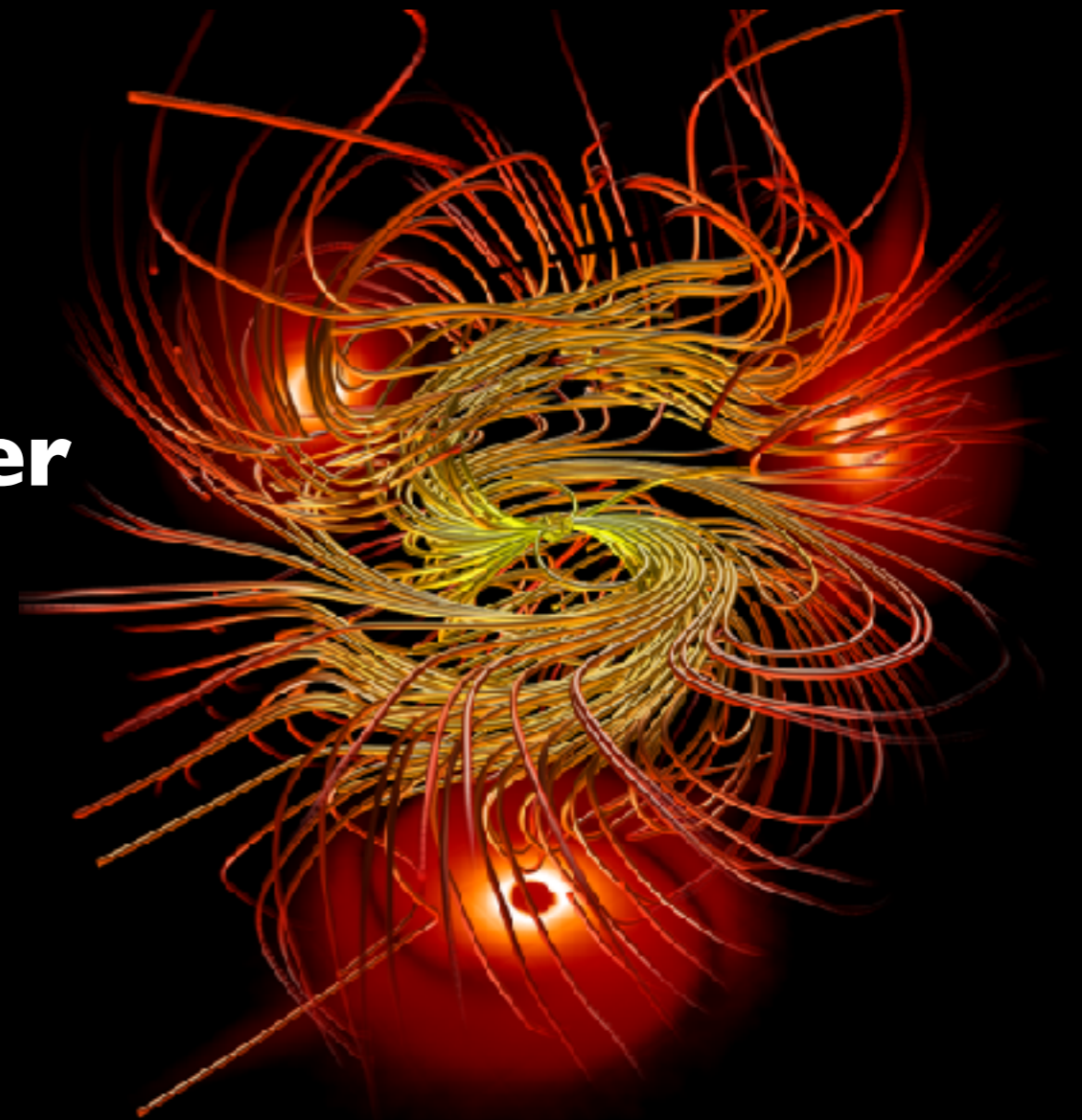
Probability of the process



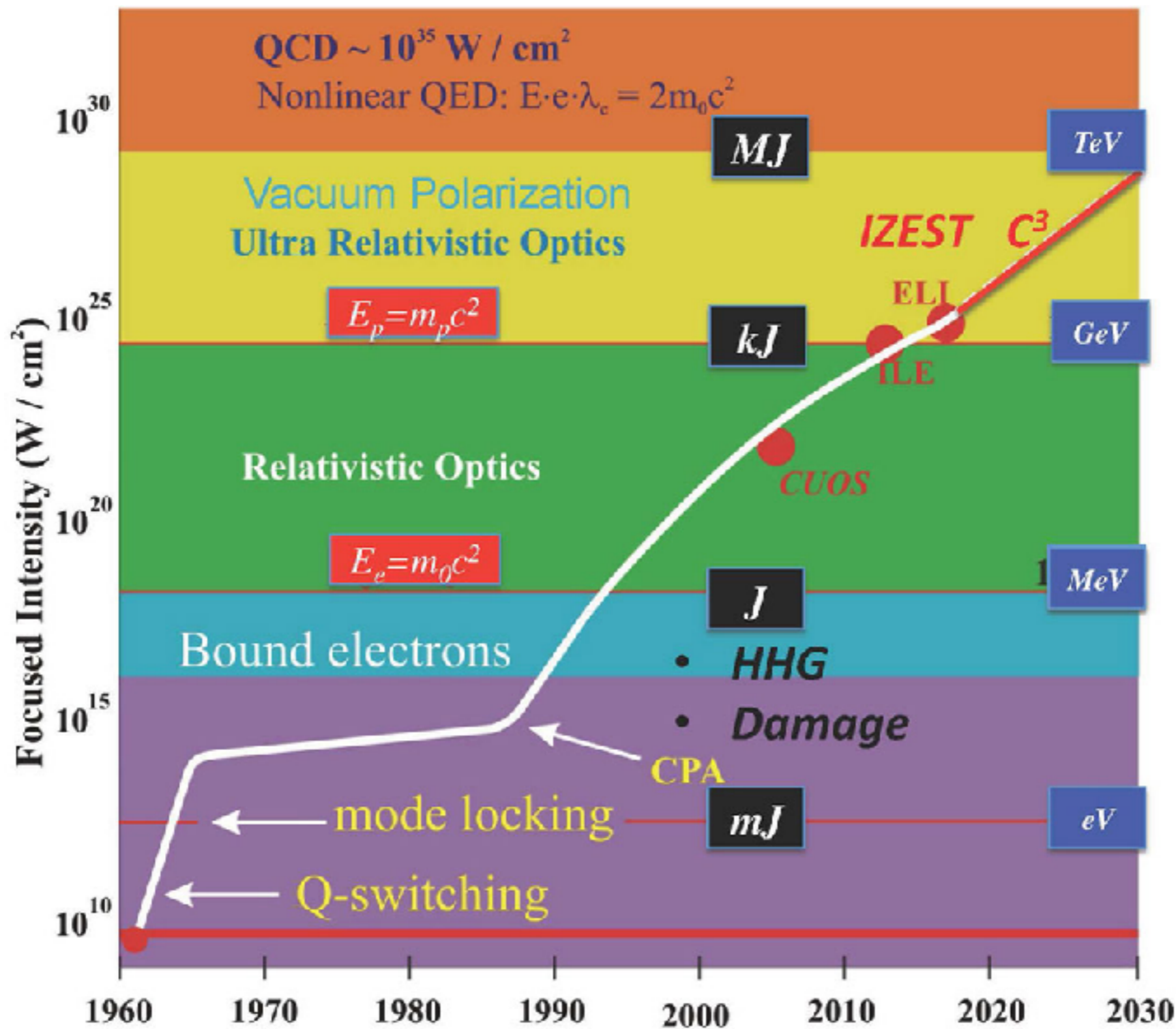
Courtesy: T. Grismayer and B. Martinez

Ultra intense lasers & particle beams and the intensity frontier

QED plasma processes in relativistic astrophysics (pulsars)



At the focus of intense lasers



Adapted from Mourou, Tajima, Bulanov, Rev. Mod. Phys. (2006)

Existing or planned particle beams

LHC @ CERN $I \sim 2.5 \times 10^{19} \text{ W/cm}^2$

100 kJ, 7 TeV per proton, 10^{11} protons per beam; 10 cm long bunch; 200 microns spot

SPS @ CERN $I \sim 1.5 \times 10^{18} \text{ W/cm}^2$

~7 kJ, 0.5 TeV per proton, 10^{11} protons per beam; 10 cm long bunch; 200 microns spot

ILC $I \sim 1.5 \times 10^{24} \text{ W/cm}^2$

1.6 kJ, 0.5 TeV per electron/positron, 2×10^{10} electrons/positrons; < 10 nm width in x; < ~100 nm width in y; 6 mm long

SLAC $I \sim 1.2 \times 10^{19} \text{ W/cm}^2$

160 J, 50 GeV per electron/positron, 2×10^{10} electrons/positrons; ~50 microns long; ~50 microns spot

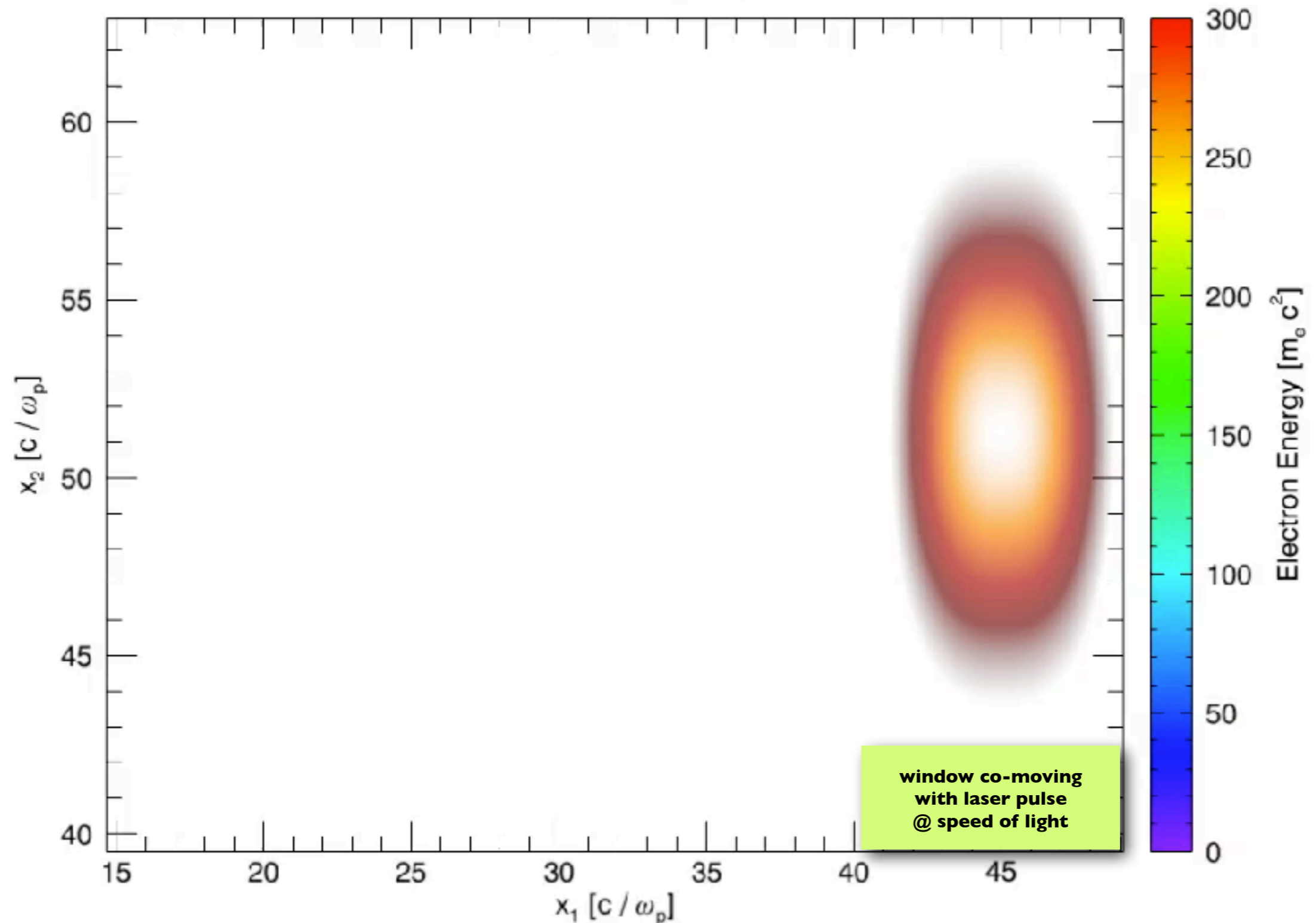
Plasma accelerators are an example of extreme plasma physics at the forefront of Science

Simulations + lasers + sources directly impacted this progress

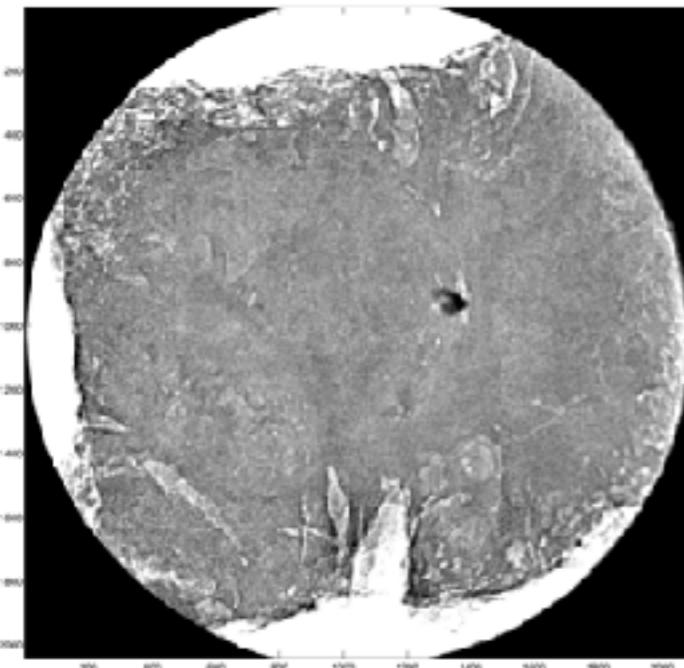


Self-injection, Dephasing, and Depletion

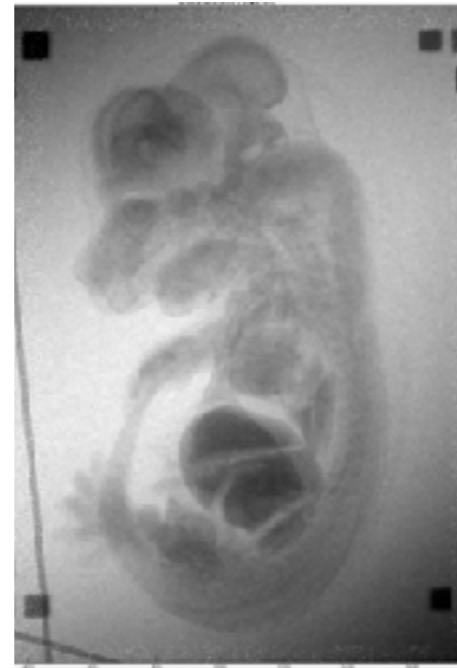
Time = 0.00 [1 / ω_p]



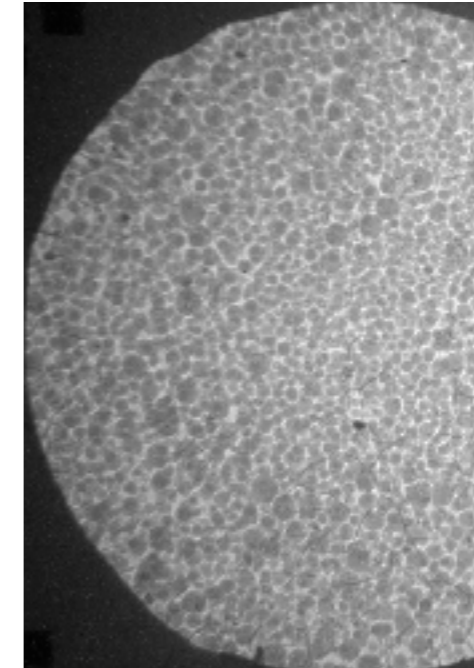
Application to bone micro-structure* imaging (above)
biological, medical, material and shock imaging (below)



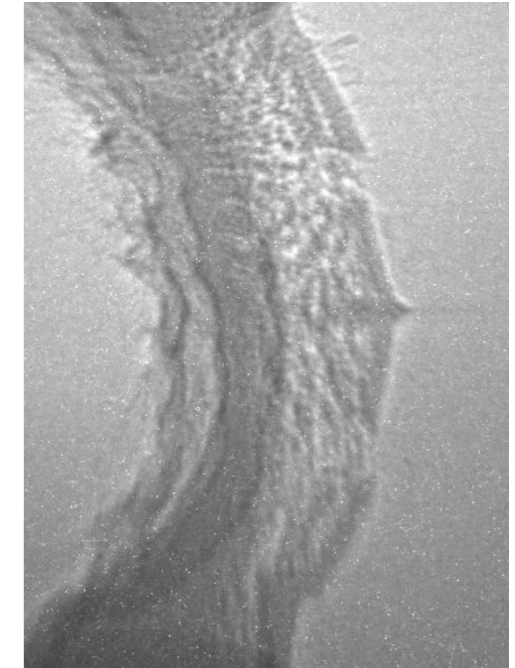
5 mm thick breast sample



mouse embryo raw image and HQ
tomographic reconstruction



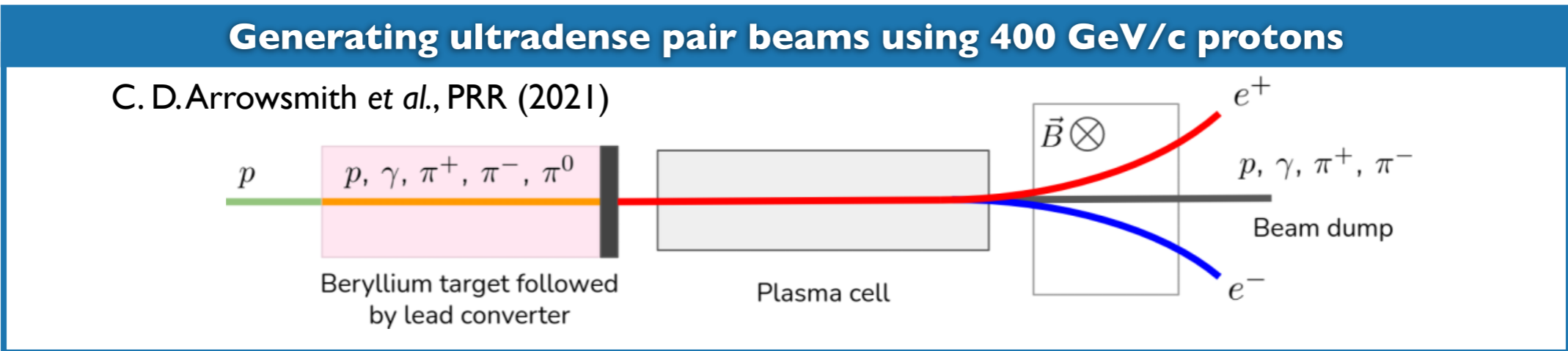
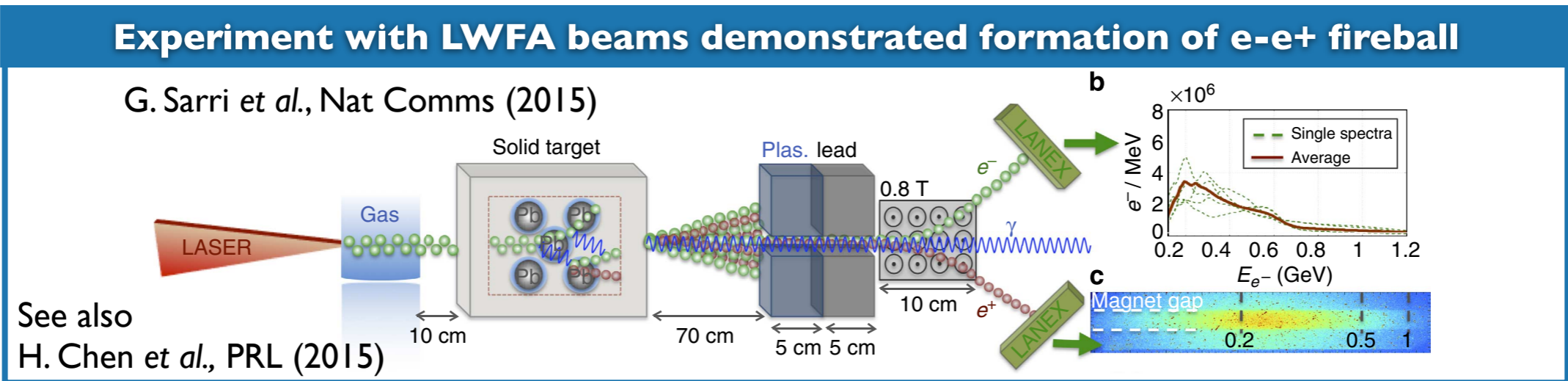
imaging of sintering
powders



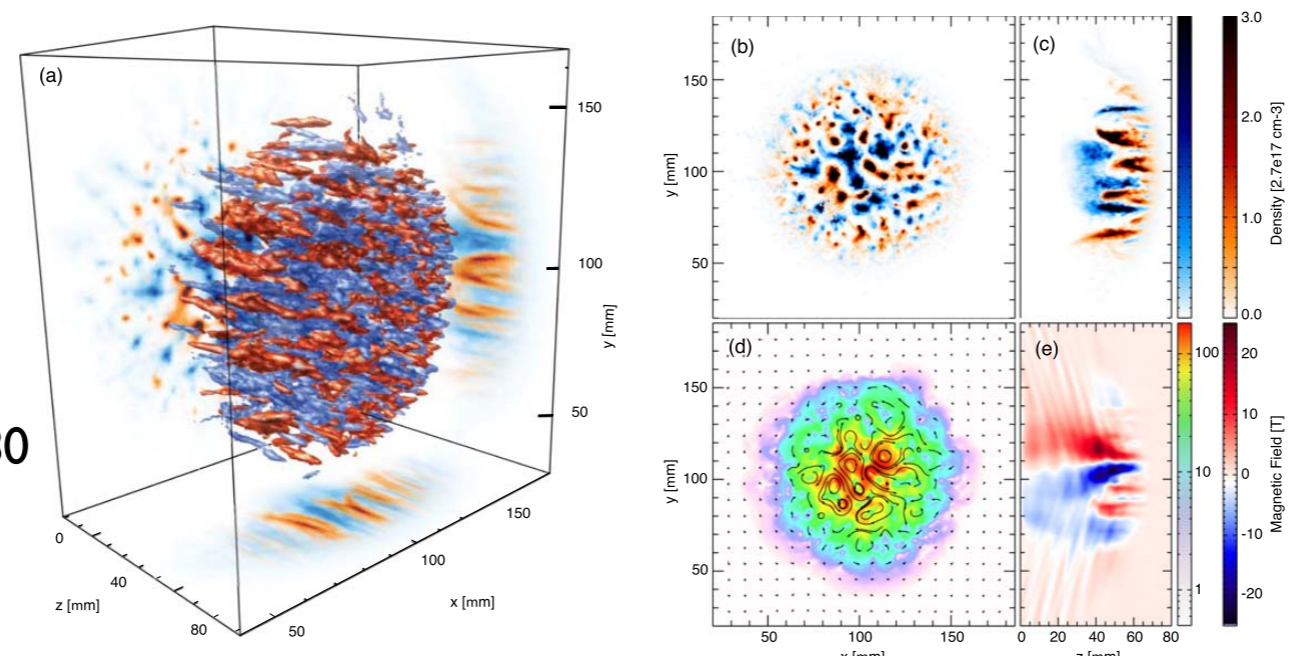
imaging of laser
induced shocks

J. M. Cole, J. C. Wood, N. C. Lopes *et al.*, *Sci. Rep.* 5, 13244(2015)

e-e+ fireballs from lasers and beams



P. Muggli *et al.*, arXiv:1306.4380
 N. Shukla *et al.*, JPP (2018)
 N. Shukla *et al.*, NJP (2020)

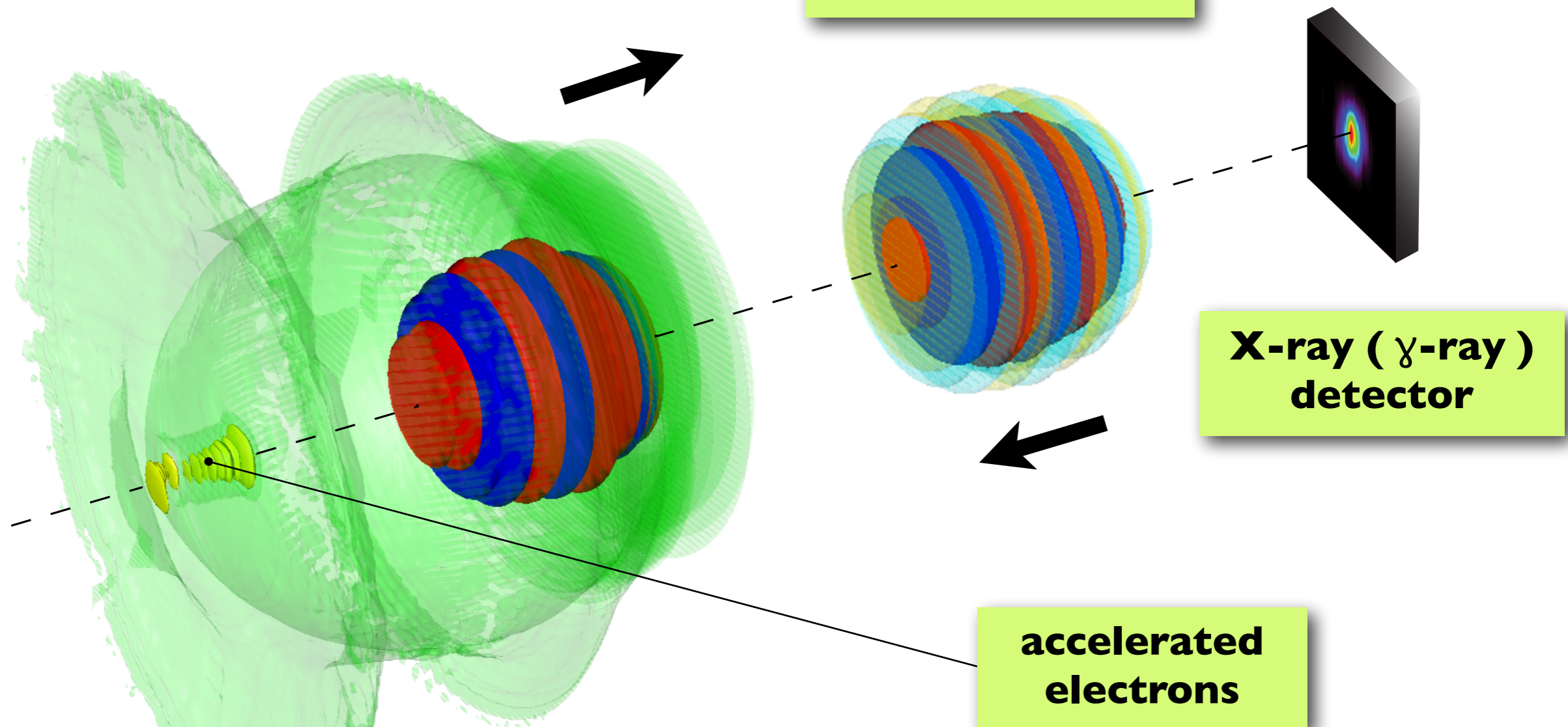


All-optical radiation reaction configuration

Identifying radiation reaction signatures and the emergence of quantum behaviour

**laser wakefield accelerator in
bubble regime**

**second laser
 $I \sim 10^{21} \text{ W/cm}^2$**



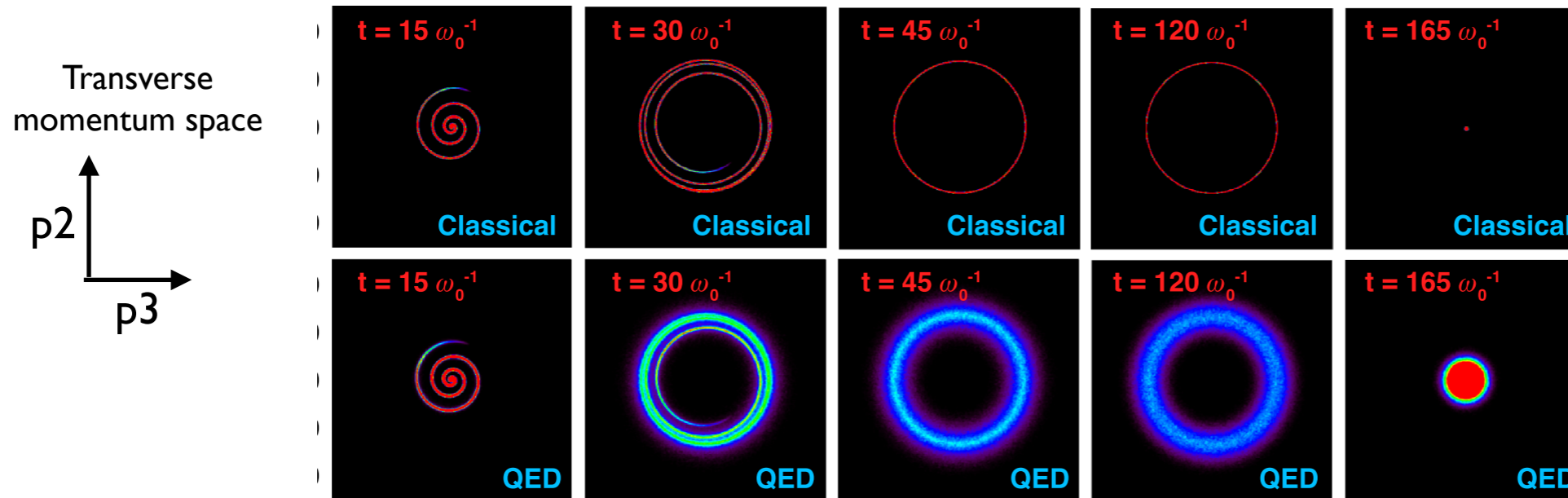
**accelerated
electrons**

**X-ray (γ -ray)
detector**

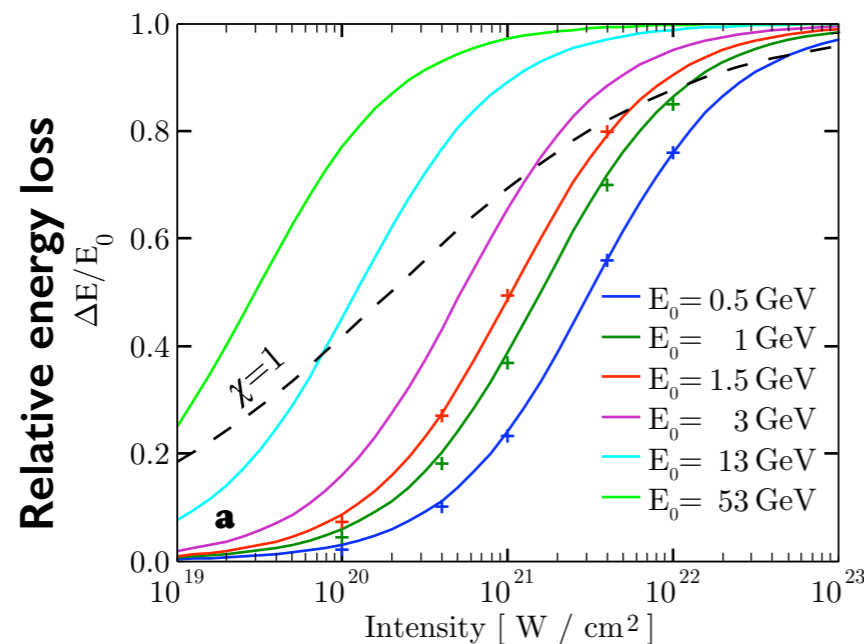
S.V. Bulanov *et al.*, NIM A (2011)
A. G. R. Thomas *et al.*, PRX (2012)
M. Vranic *et al.*, PRL (2014)

[also connection with S. Cippicia *et al.*, Nature Physics (2011)]

The emergence of QED behaviour: classical to quantum transition in the relativistic regime

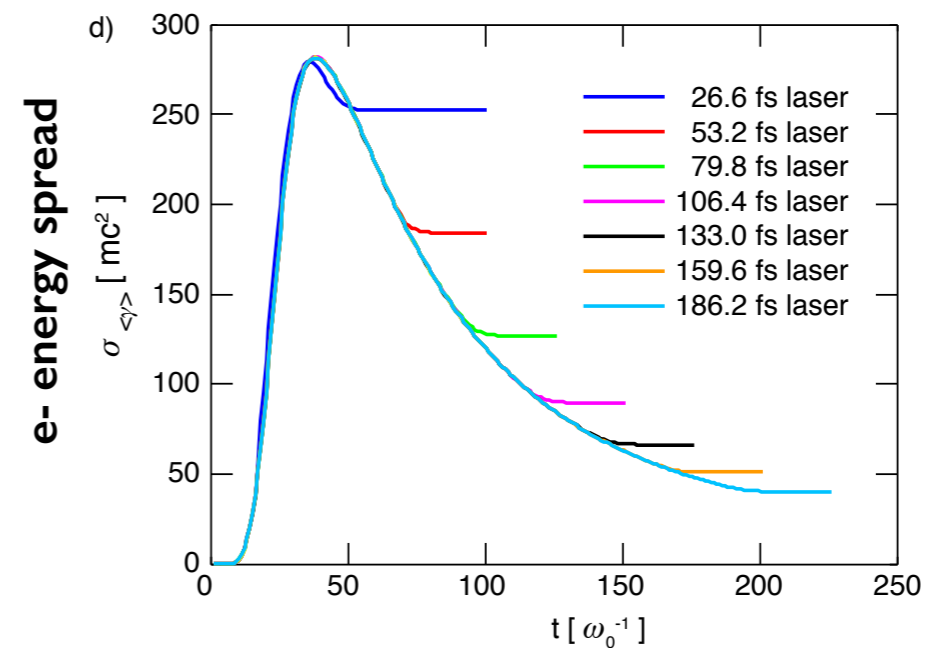


Energy loss for electron beams (classical)



M.Vranic *et al.*, PRL (2014)
M.Vranic *et al.*, Comput. Phys. Comm. (2016)

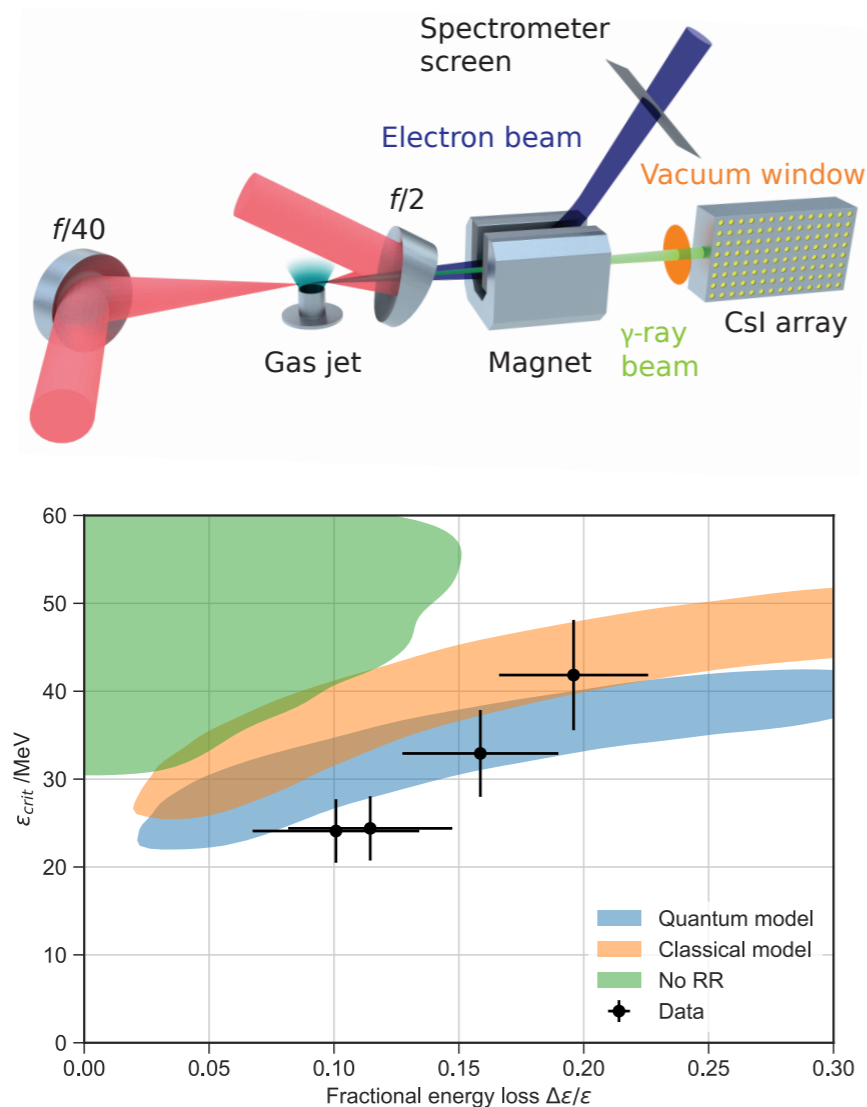
Competition: drift (classical) vs diffusion (quantum)



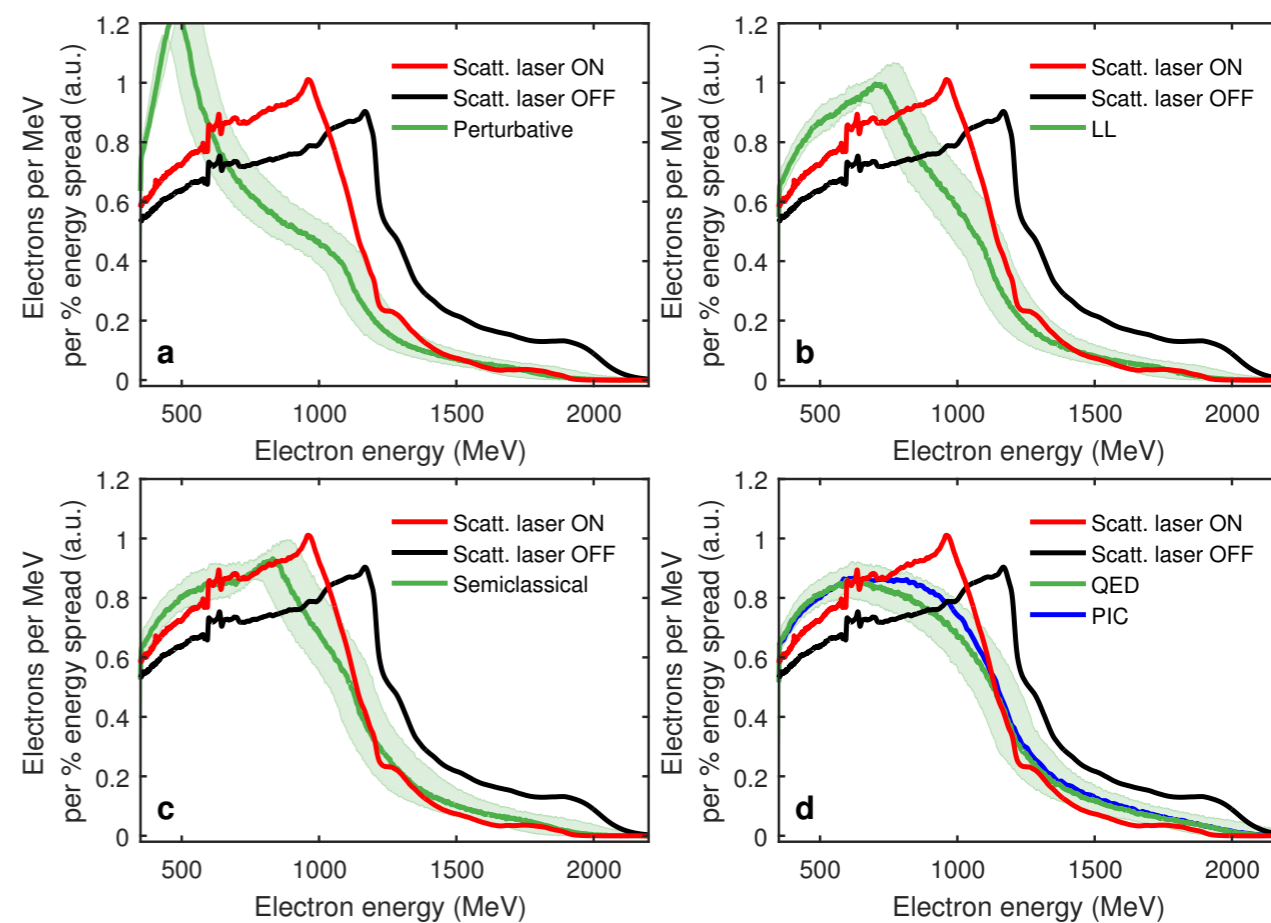
M.Vranic *et al.*, New J. Phys. (2016)
Also N. Neitz and A. Di Piazza, PRL (2013), S.Yoffe *et al.*, NJP (2015)

Recent experimental results provide evidence and open new questions

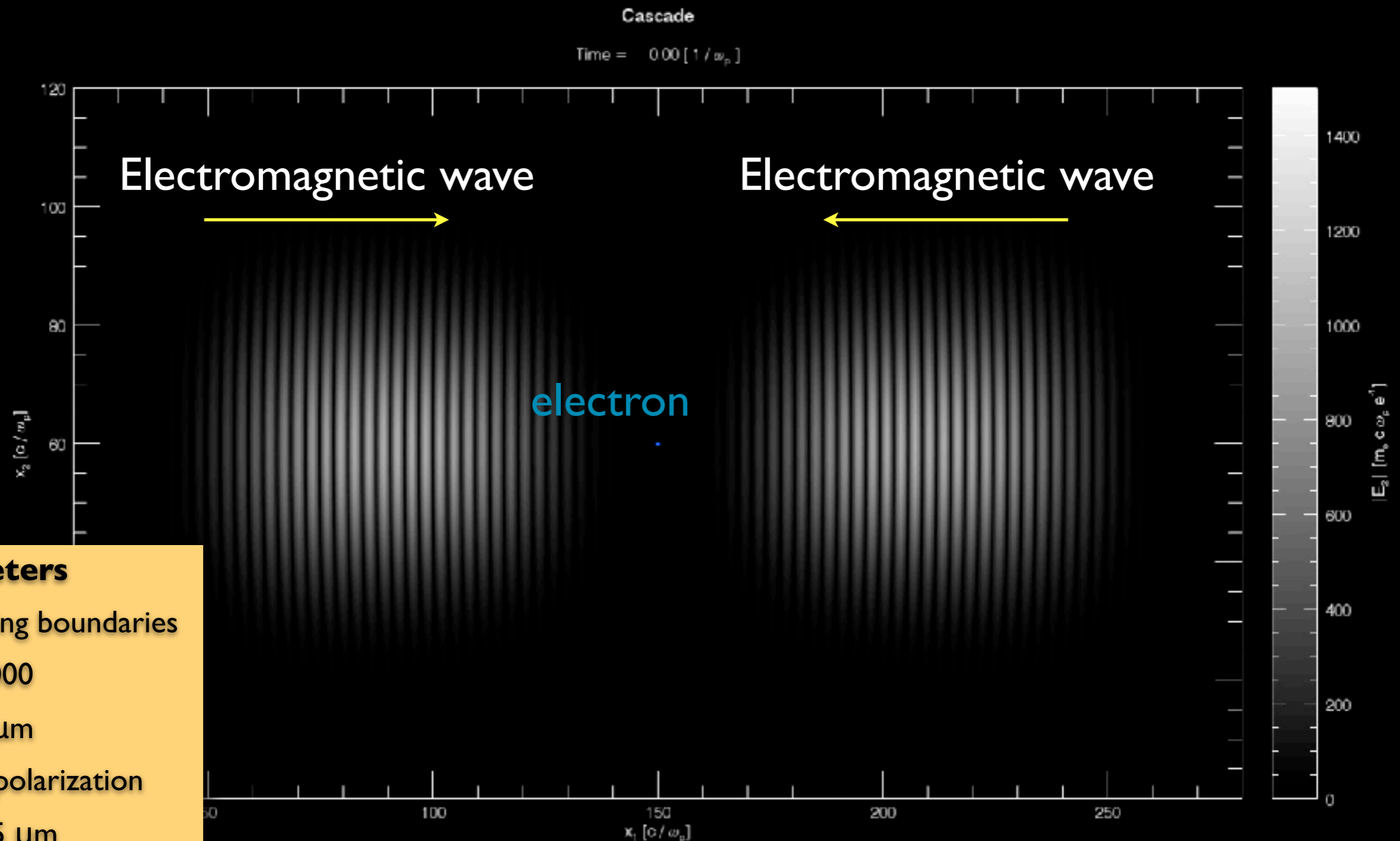
J. M. Cole et al., PRX (2018)



K. Poder et al., PRX (2018)



T. Grismayer *et al.*

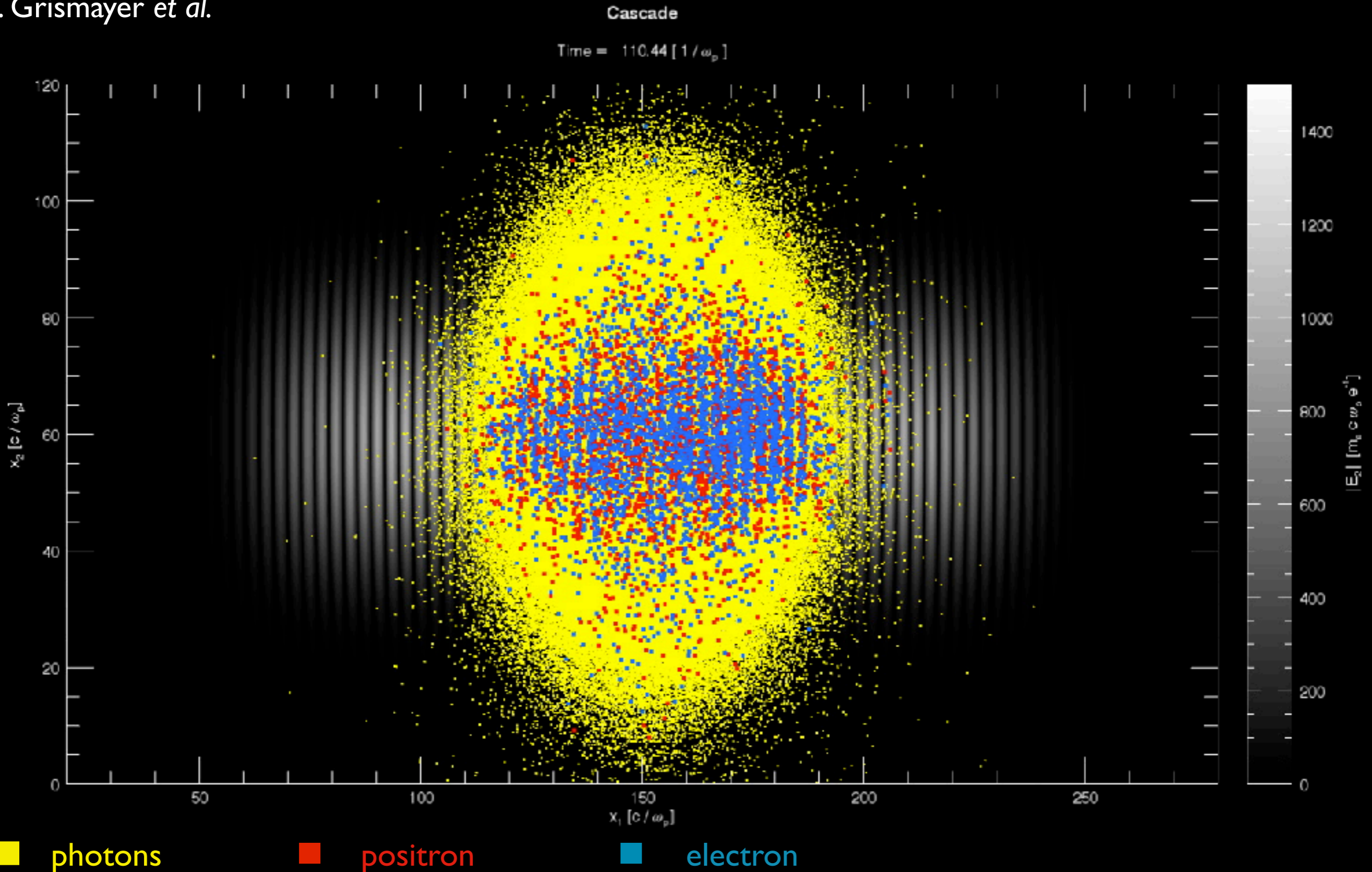


Parameters

- absorbing boundaries
- $a_0 = 1000$
- $\lambda_0 = 1 \mu\text{m}$
- Linear polarization
- $W_0 = 5 \mu\text{m}$
- $\tau = 30 \text{ fs}$

QED cascades in counter propagating electromagnetic fields

T. Grismayer *et al.*



This requires finding where the parameter $\frac{1}{E_S} \sqrt{\left(\gamma \mathbf{E} + \frac{\mathbf{p}}{mc} \times \mathbf{B}\right)^2 - \left(\frac{\mathbf{p}}{mc} \cdot \mathbf{E}\right)^2}$ reaches high values

Linear

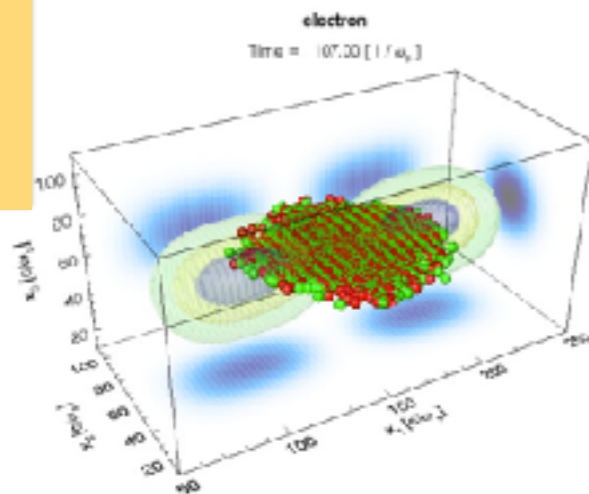
Double clockwise

Clockwise-anti clockwise

Parameters

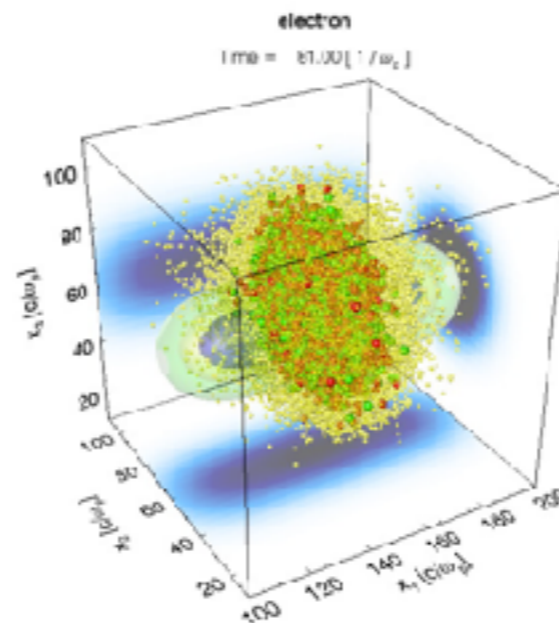
- absorbing boundaries
- $a_0 = 1000$
- $\lambda_0 = 1 \mu\text{m}$
- $W_0 = 5 \mu\text{m}$
- $\tau = 30 \text{ fs}$

- electron
- positron



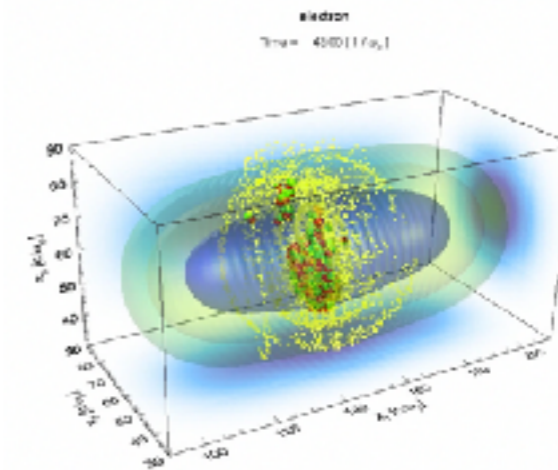
Particles remain in the x_1-x_2 plane

- electron
- positron
- photon



Particles explore the whole space

- electron
- positron
- photon



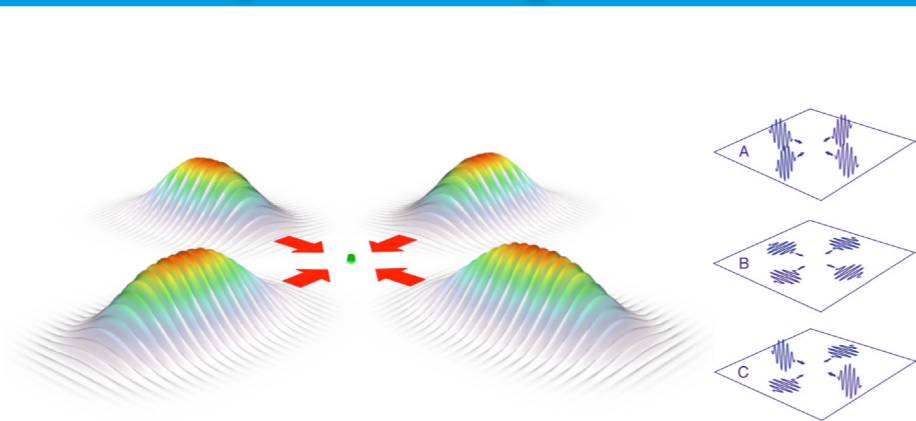
Particles rotate mainly in the x_2-x_3 plane

T. Grismayer et al., *PRE* 95, 023210 (2017)

T. Grismayer et al., *APS DPP* (2012)

$$\chi = \frac{1}{E_S} \sqrt{\left(\gamma \mathbf{E} + \frac{\mathbf{p}}{mc} \times \mathbf{B}\right)^2 - \left(\frac{\mathbf{p}}{mc} \cdot \mathbf{E}\right)^2}$$

4 pulse configuration



$$\Gamma \sim \frac{8}{15\pi} \left(\frac{2\pi}{3}\right)^{\frac{1}{4}} \frac{\alpha}{\tau_c \bar{\gamma}} K_{1/3}^2 \left(\frac{4}{3\bar{\chi}_e}\right)$$

$$\chi_e \approx \frac{1}{E_S} \left\| \frac{\vec{p}}{mc} \times \vec{B} \right\|$$

$$\bar{\gamma} \approx 6a_0$$

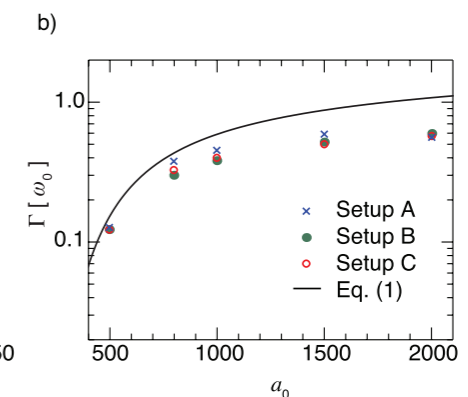
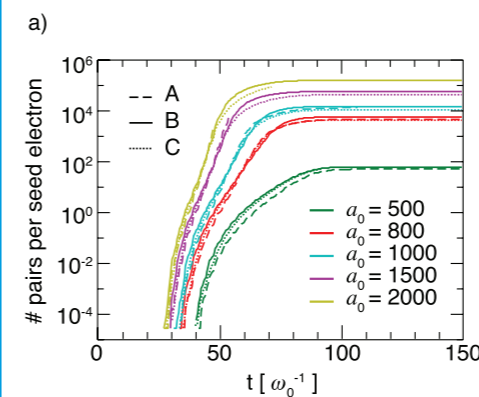
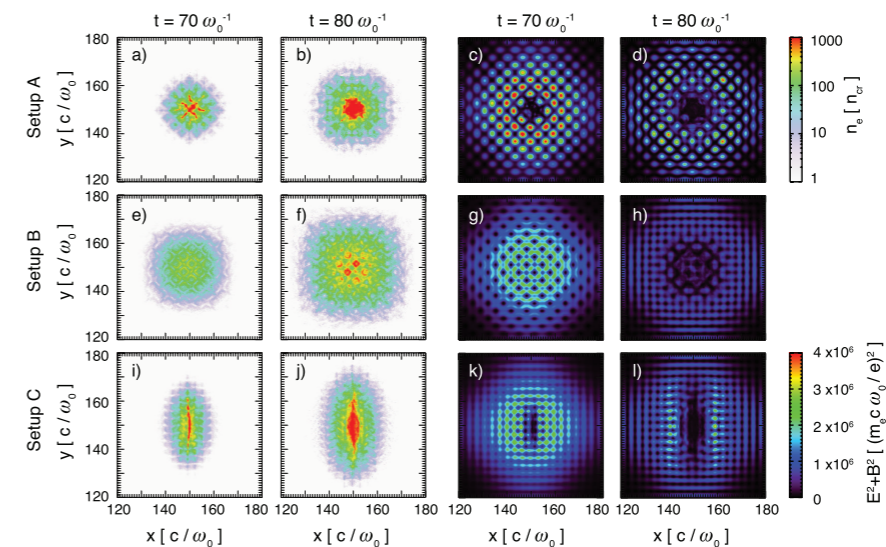
$$\bar{\chi}_e \approx 12a_0^2 / (\pi a_S)$$

$$a_S = mc^2 / (\hbar \omega_0)$$

M.Vranic *et al.*, PPCF (2017), arXiv:1609.08081

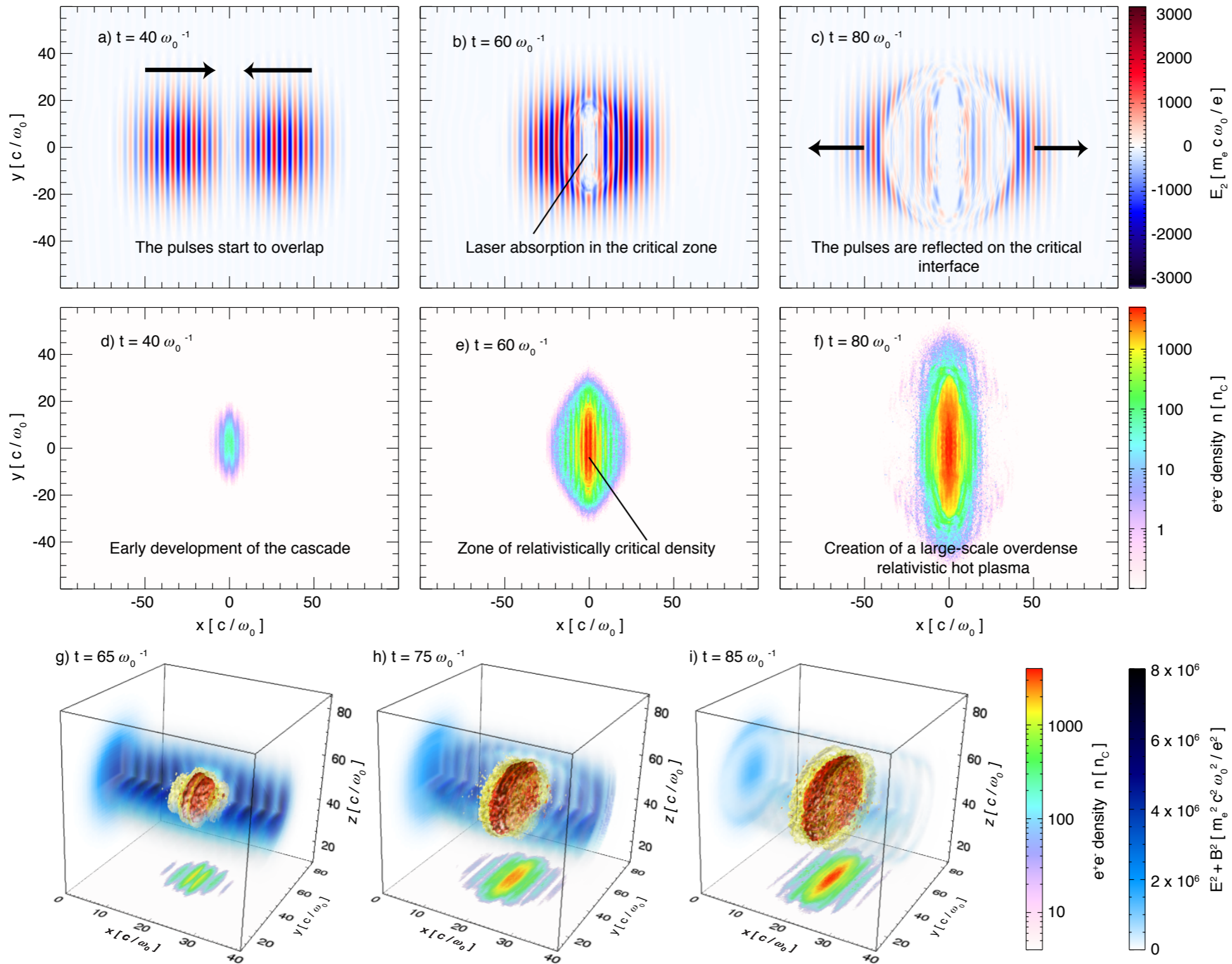
Z. Gong *et al.*, PRE (2017), arXiv:1609.09567

Standing wave pattern enhances e-e+ plasma



Ideal configuration has been first explored in I. Gonoskov *et al.*, Phys Rev A (2012)

Self-consistent dynamics of e-e+ plasma impacts field dynamics

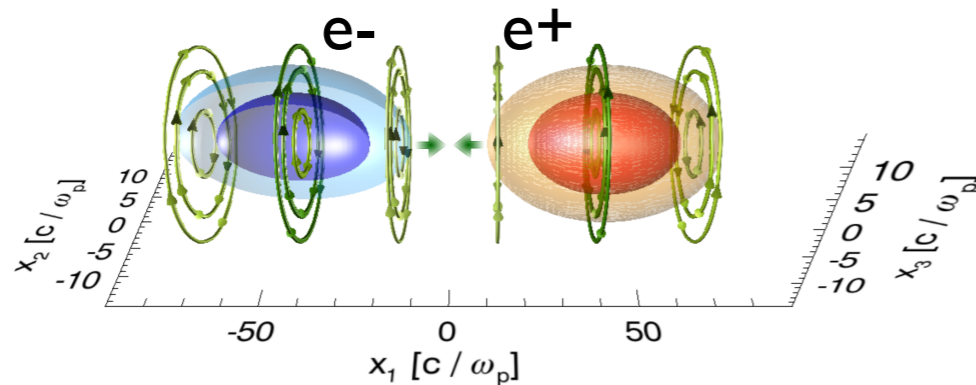


T. Grismayer *et al.*, PRE (2017)

T. Grismayer *et al.*, PoP (2016)

See also A. Fedotov *et al.*, PRL (2010); E. Nerush *et al.*, PRL (2011); V. Bashmakov *et al.*, PoP (2014)

Beam-beam collider



At the interaction point, the particles in one beam will feel:

$$E_{\perp} \simeq B_{\perp} \sim en_0\sigma_0 \quad \chi(r, z) = 2\gamma B_{\perp}(r, z)/E_s$$

$$E_{\parallel} \sim E_{\perp}/\gamma$$

Beams will then pinch (in e- e- collider will diverge). Number of pinching points during crossing time is

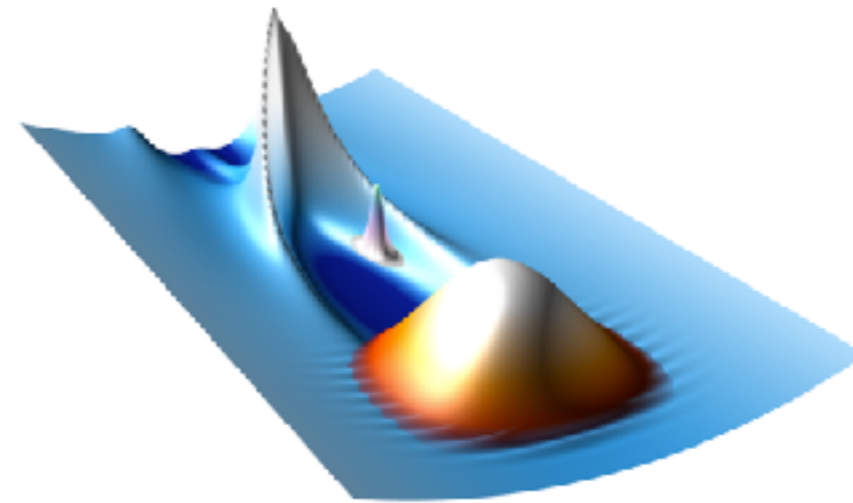
$$D = \frac{r_e N \sigma_z}{\gamma \sigma_0^2} \quad \text{Disruption parameter}$$

P. Chen, S. Rajagopalan and J. Rosenzweig, PRE (1989)

P. Chen and K. Yokoya *et al.*, PRD (1988)

T. Katsouleas *et al.*, PoP (1990)

Betatron in ion channels



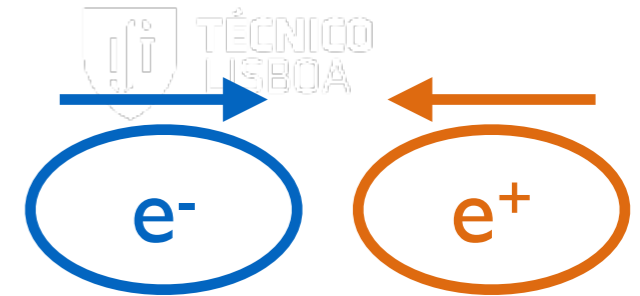
$$E_{\perp} = 4\pi en_{e0} \frac{r}{2} \quad \chi(r) = \gamma \frac{E_{\perp}}{E_s}$$

In the blowout regime $n_{\text{beam}} > n_{e0}$

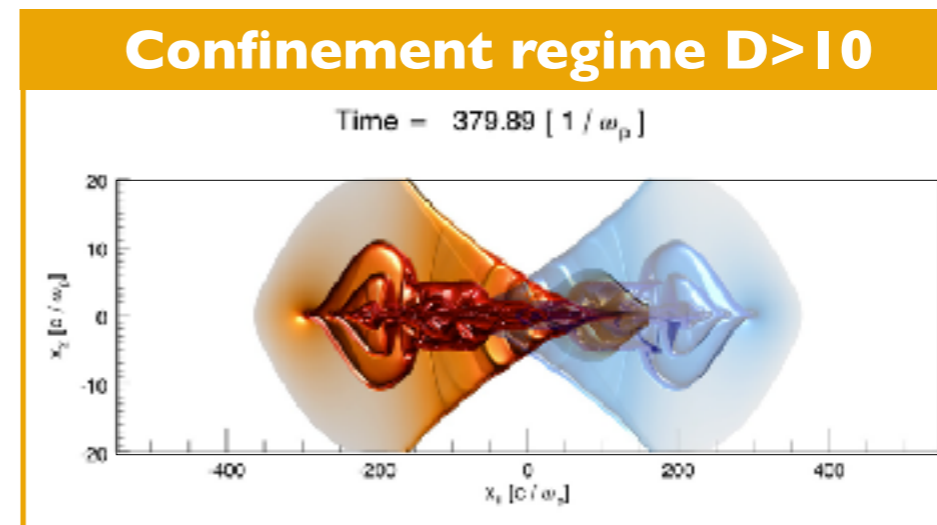
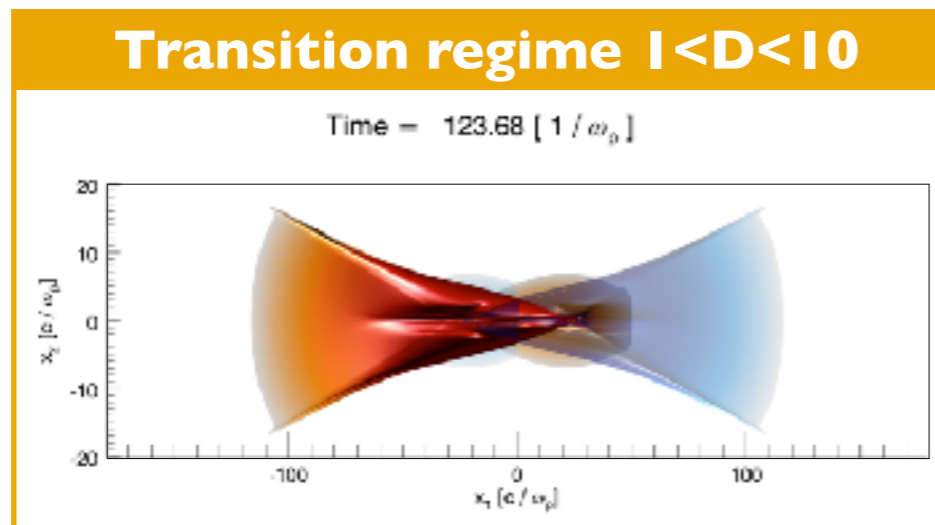
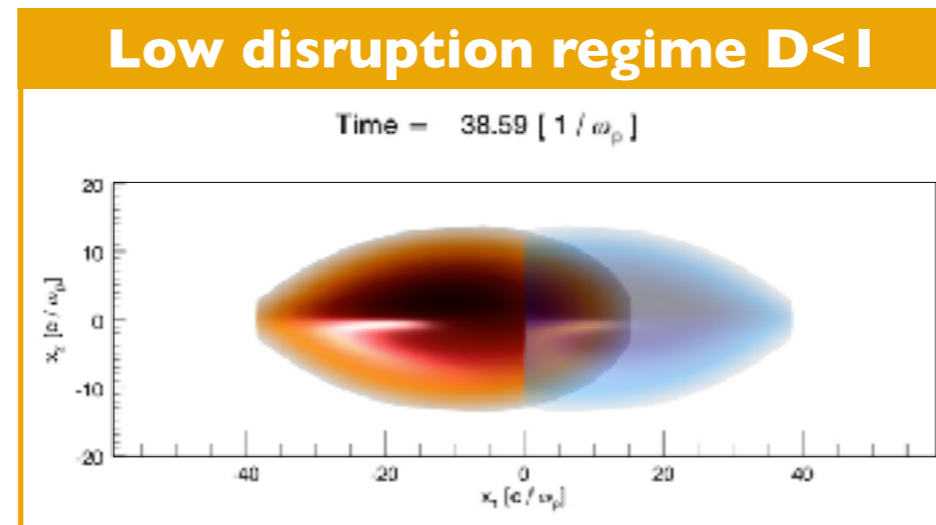
Beam can drive ion channel

For small emittance, long propagation in ion channel @ 30 GeV

$$\sigma_0 \approx \mu\text{m} \quad \chi \approx 0.1 - 1$$



$$D = \frac{r_e N \sigma_z}{\gamma \sigma_0^2}$$

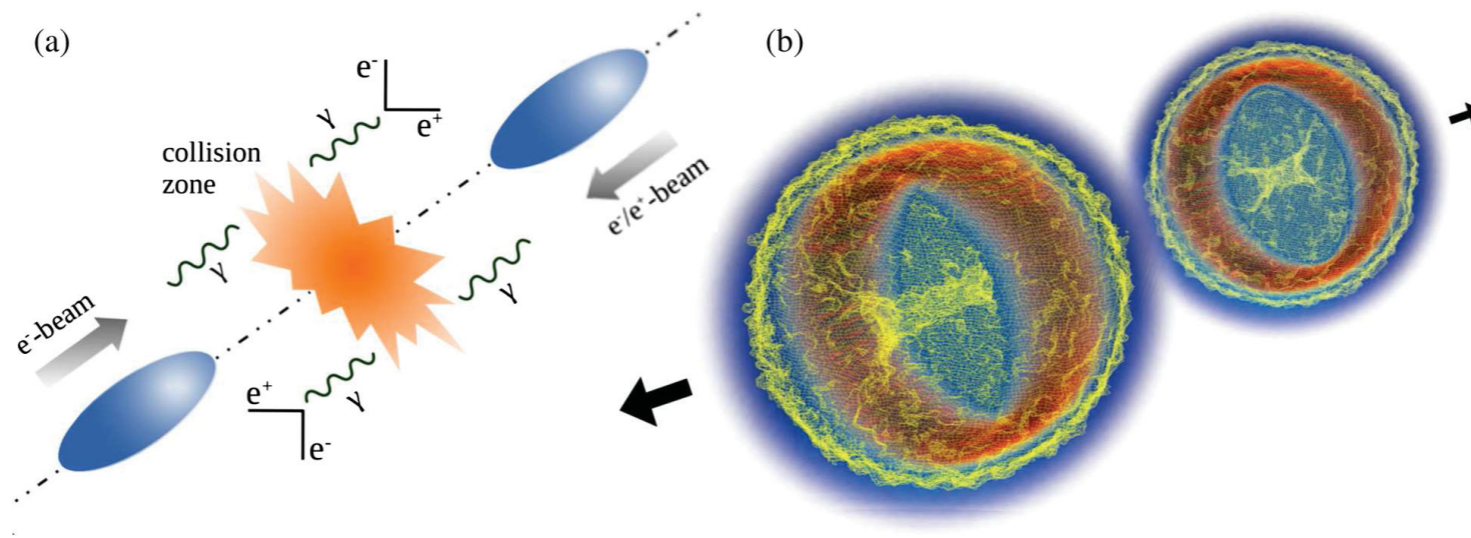


- Electron beam density
- Positron beam density

Exploring LWFA like beams for radiation in the low D regime
 F. Del Gaudio, *et al.*, PRAB 2019

Medium to high D regime - W. L. Zhang *et al.*, in preparation

SLAC PWFA @ 500 GeV can reach the QED regime (and beyond)



Energy : 500 GeV
N : 10^{10}
Length : 15 μm
Spot size : 15.7 nm

@ Interaction Point

$$D > 1 \quad \chi > 30$$

Abundant secondary pairs are created in a stochastic manner, leading to the kink instability

Additional note: also possible to reach conditions where QED might “breakdown”

Ritus-Narozhny conjecture: the expansion parameter of QED in the strong field sector is

V. Ritus, AP (1972); N. Narozhny, PRD (1980)

A. Fedotov, JP Conf Series (2017)

$$\alpha \chi^{2/3}$$

i.e. QED becomes a strongly coupled theory if $\alpha \chi^{2/3} \gtrsim 1$

V. Yakimenko *et al.*, PRL (2019)

also C. Baumann *et al.*, Sci Reports (2019), T. Blackburn *et al.*, NJP (2019)

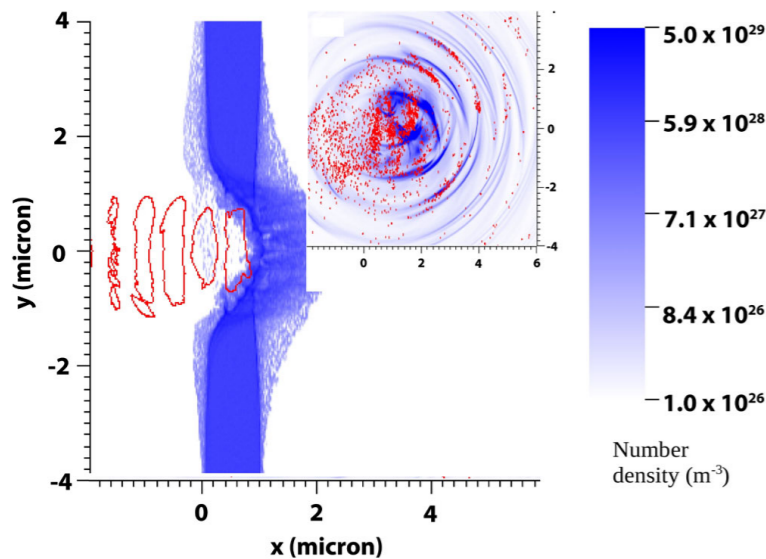
for electron beam - laser configuration

How much plasma physics in the intense laser/beam scenarios or what is the role/interplay of collective effects/QED?

System (plasma) size $W_0^3 \sim (\text{few } \mu\text{m})^3 \gtrsim (\text{few } \lambda_D)^3$

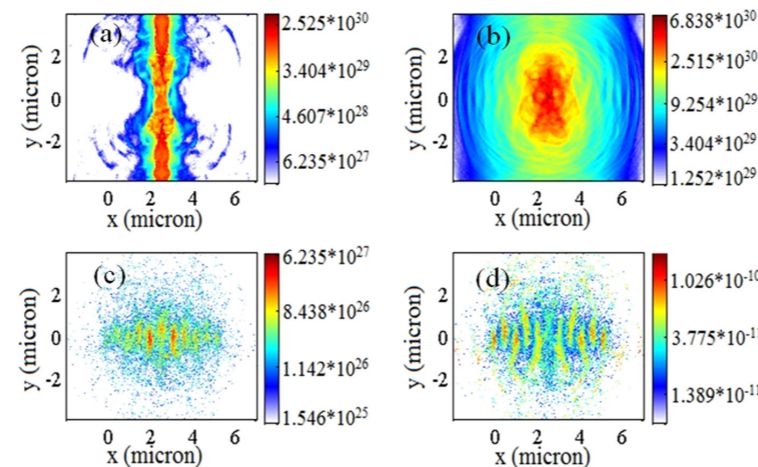
$n_{e^-e^+} \gtrsim 10^{20} \text{cm}^{-3}$ LOS *et al.*, in preparation

QED in hole boring



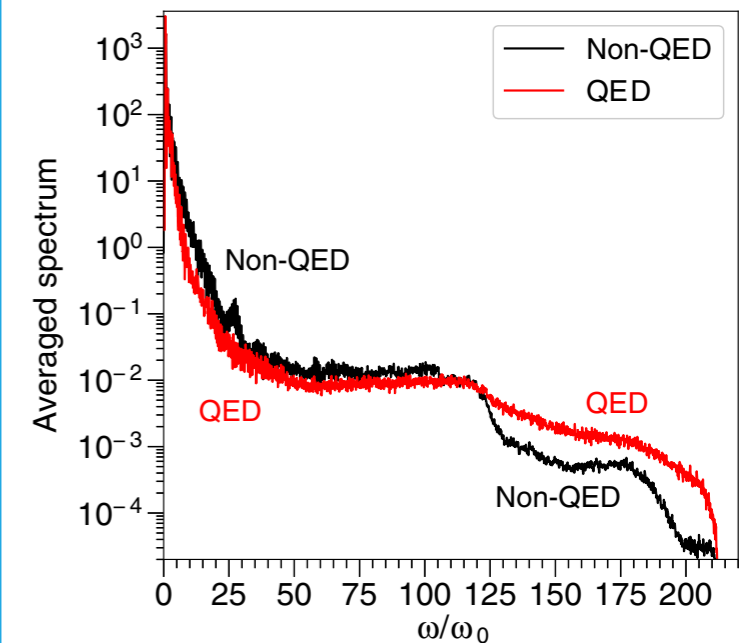
C. Ridgers *et al.*, PRL (2012)

QED in $\sim n_c$ targets with 2 lasers



W. Luo *et al.*, PoP (2015)
T. Grismayer *et al.*, PoP (2016)

QED in HHG

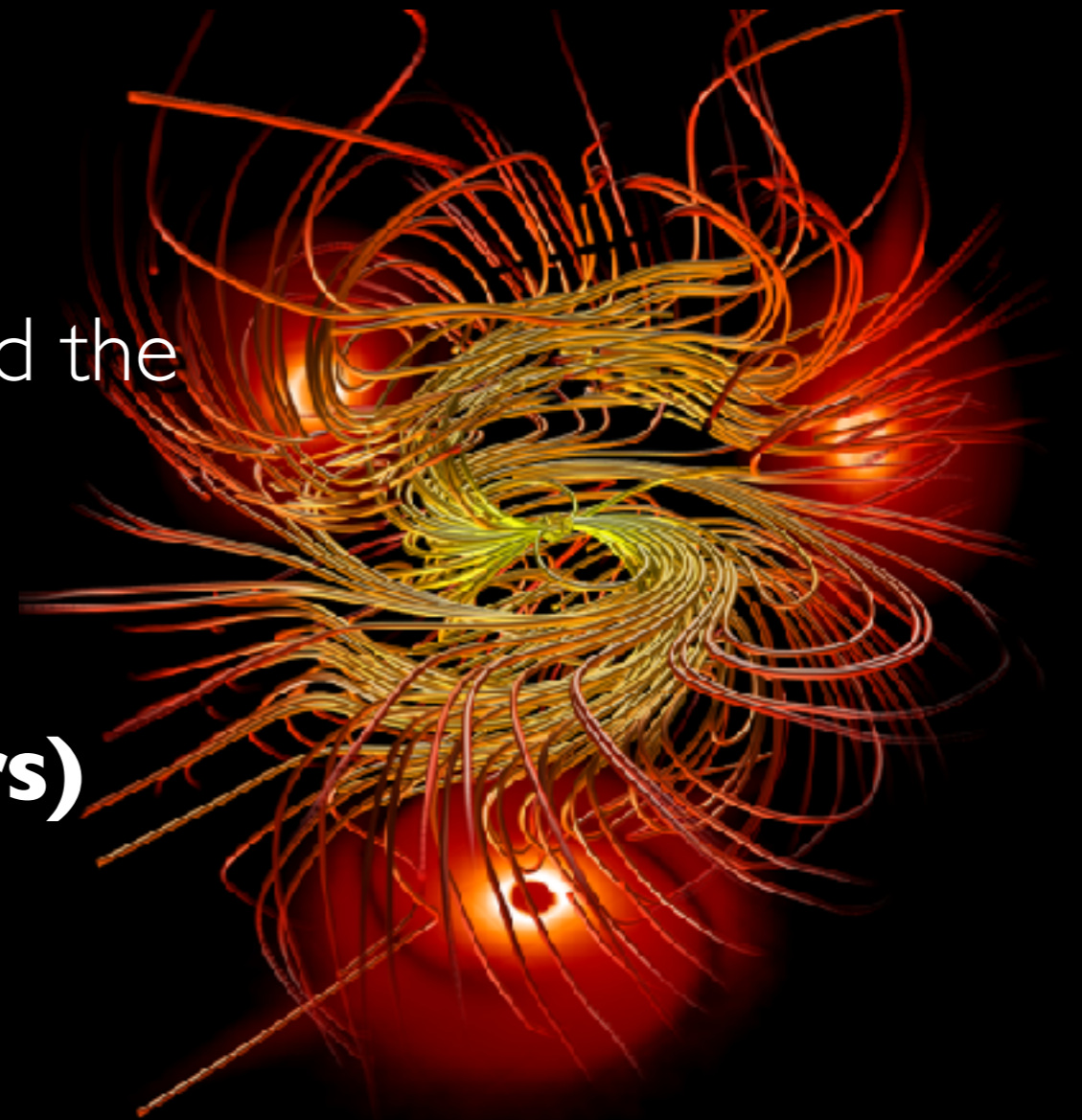


W. L. Zhang *et al.*, PRE (2021)

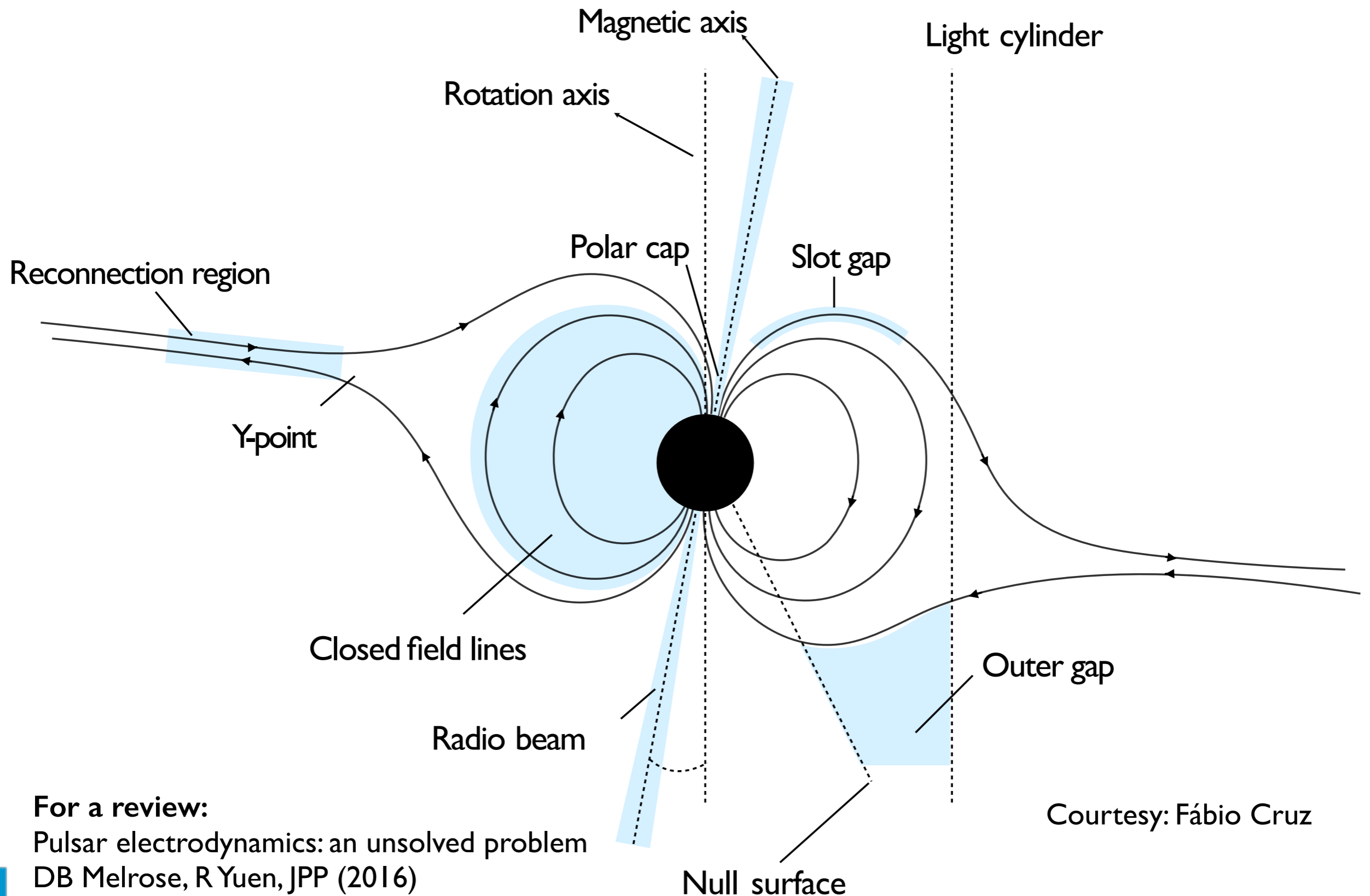
For beam-beam collisions, collective effects when $D \gtrsim 1$

Ultra intense lasers & particle beams and the intensity frontier

QED plasma processes in relativistic astrophysics (pulsars)



Understanding the properties of the magnetospheres of neutron stars/pulsars from first principles

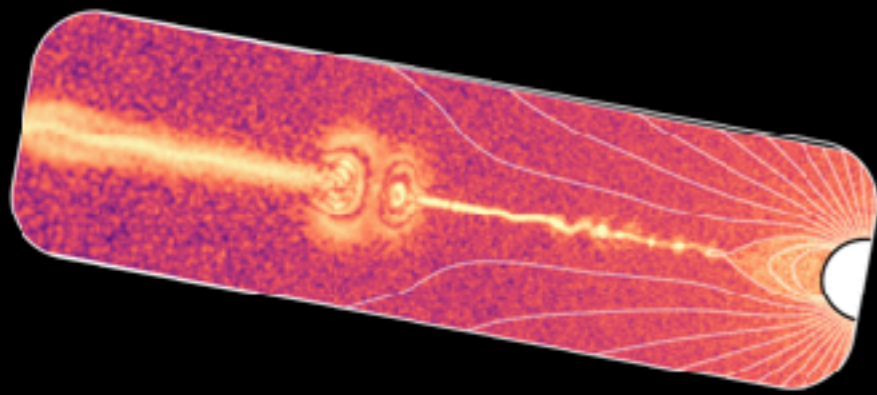


For a review:
Pulsar electrodynamics: an unsolved problem
DB Melrose, R Yuen, JPP (2016)

Courtesy: Fábio Cruz



Reconnection



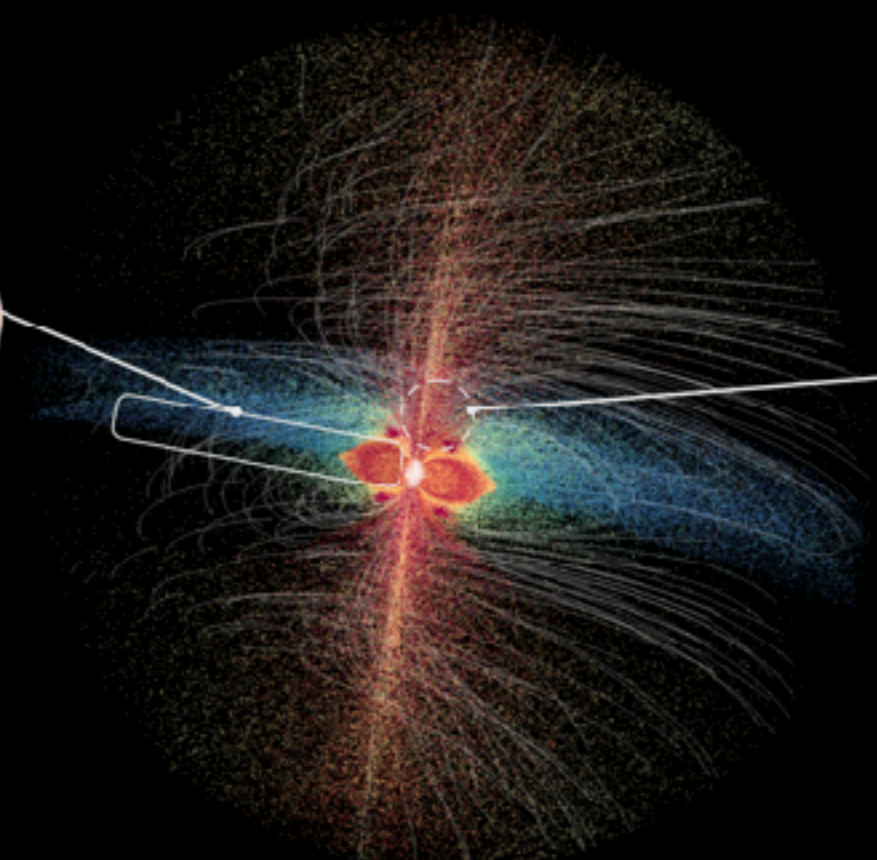
Relativistic

Uzdenski *et al.*, Spitkovsky *et al.*,
Sironi *et al.*, Philippov *et al.*

QED

Schoeffler *et al.*

Global

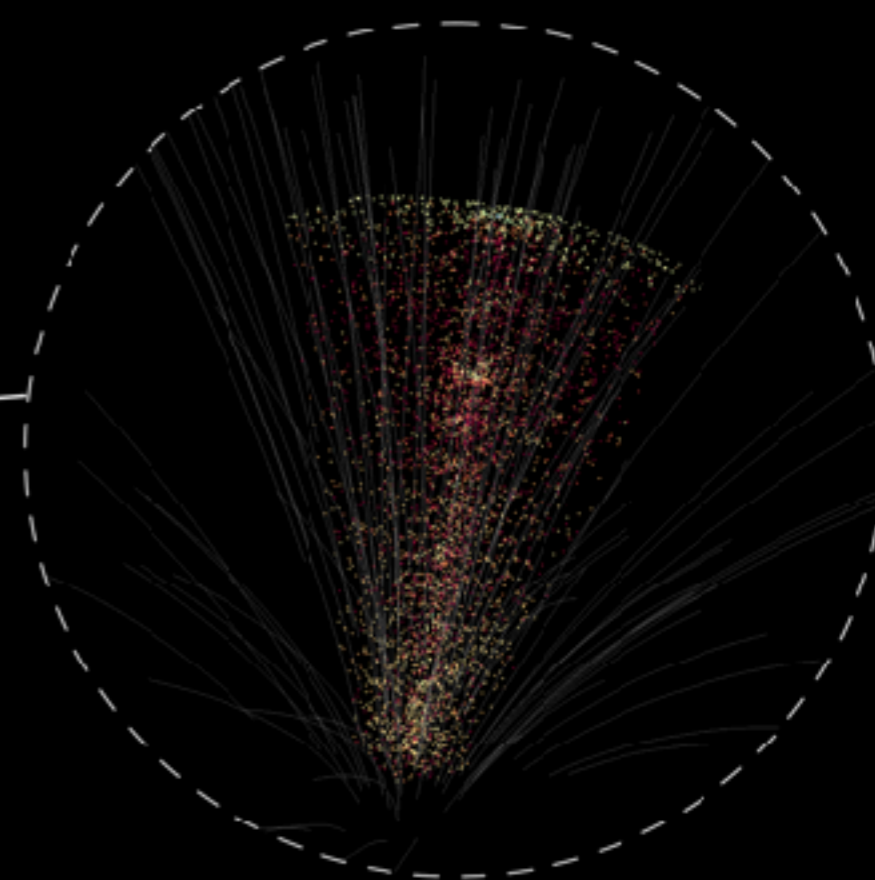


Spitkovsky *et al.*, Cerruti *et al.*,
Philippov *et al.*

GR

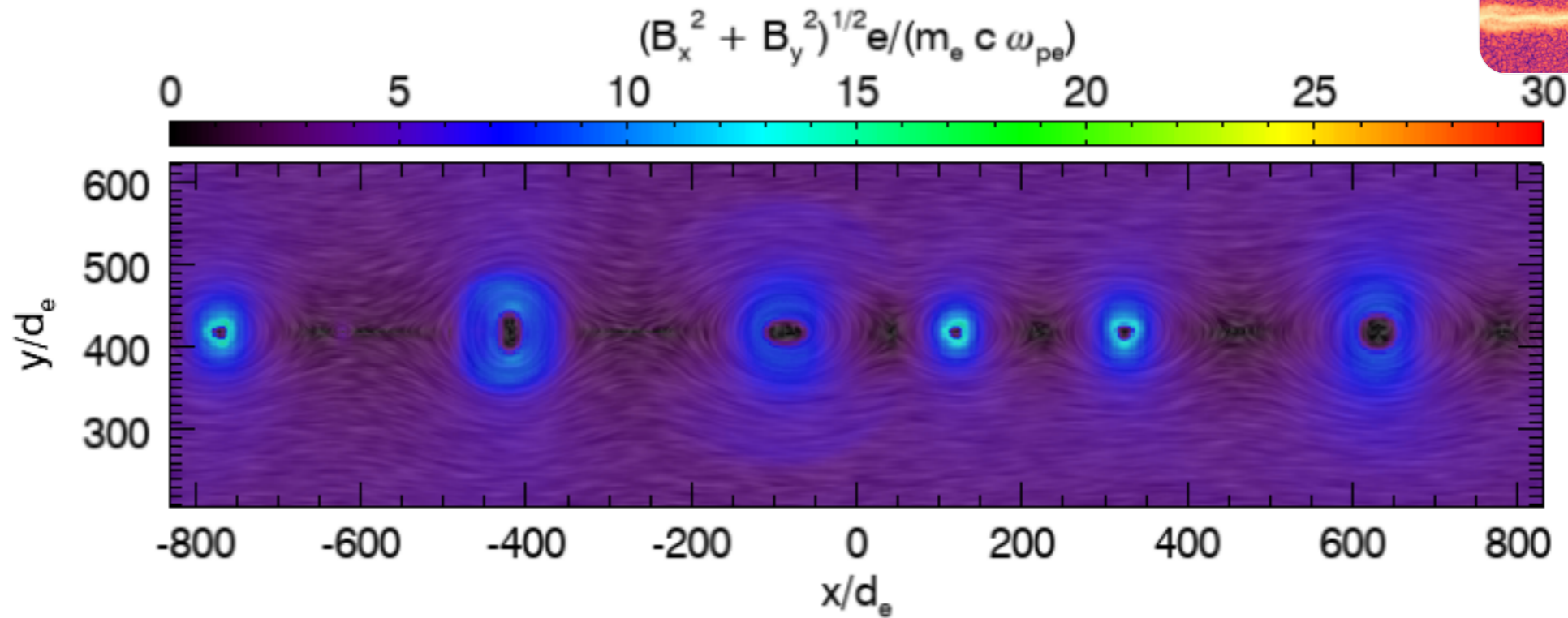
Spitkovsky *et al.*
(BH) Parfrey *et al.*, Torres *et al.*

Polar Cap

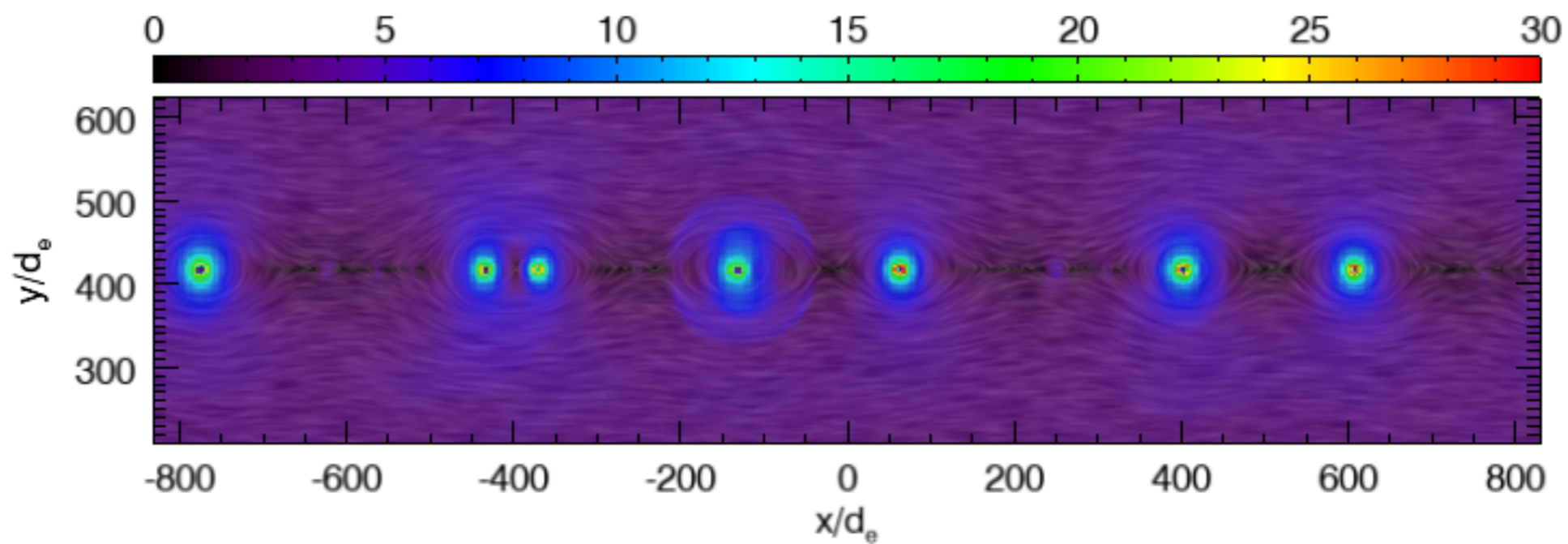


Timokhin & Arons, Philippov *et al.*,
Cruz *et al.*

**Classical relativistic
reconnection**

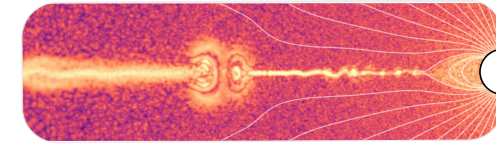
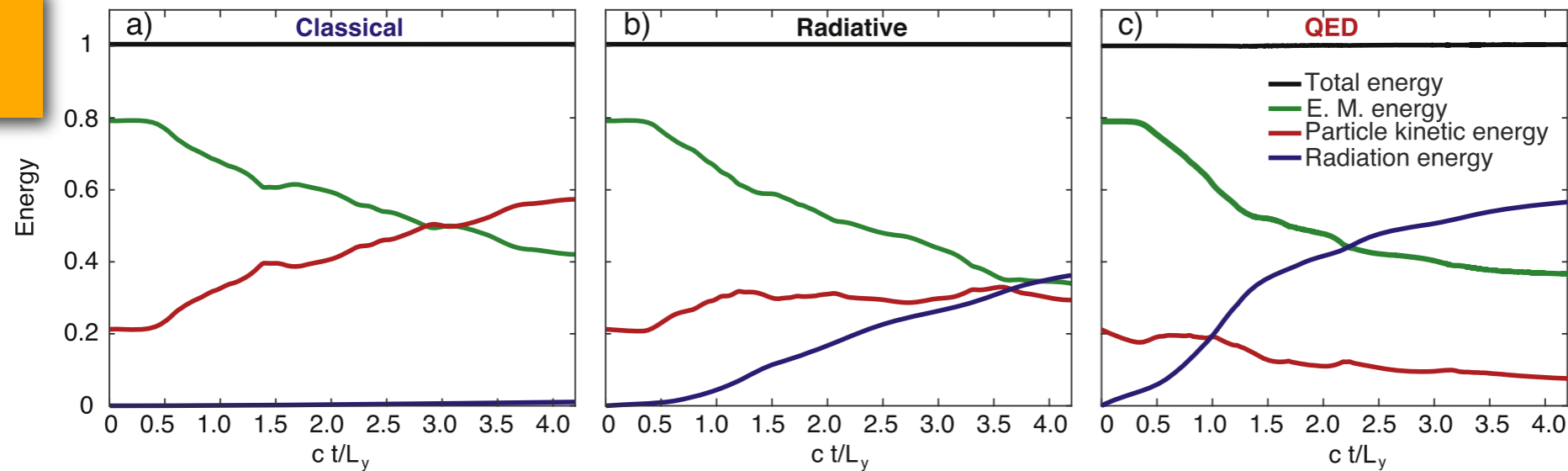


**QED reconnection
with pairs**

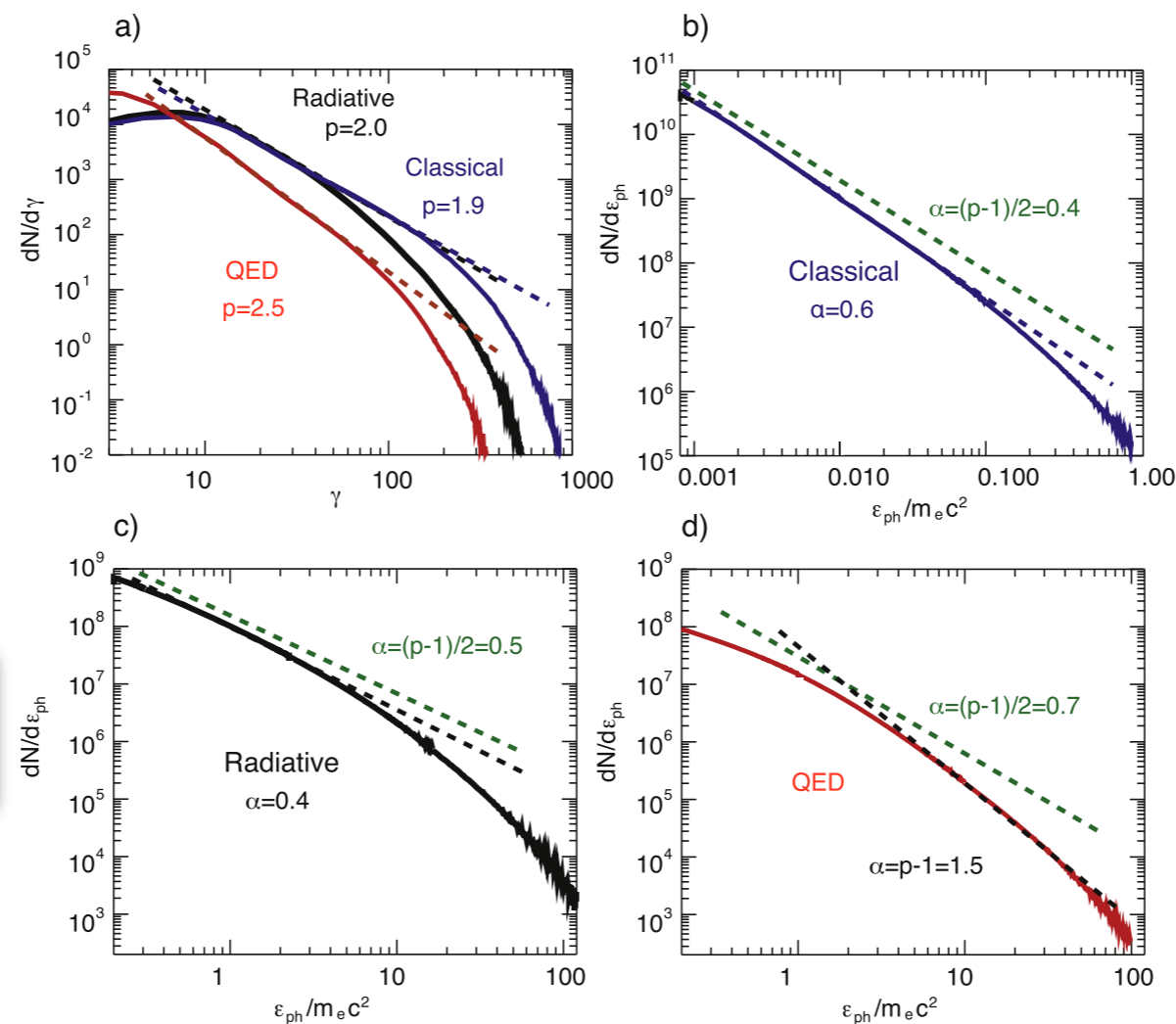


Significant differences depending on regime

Energy



e- spectrum



photon spectrum

Cyclic outflow of e-e+ bursts accompanied by kinetic instabilities

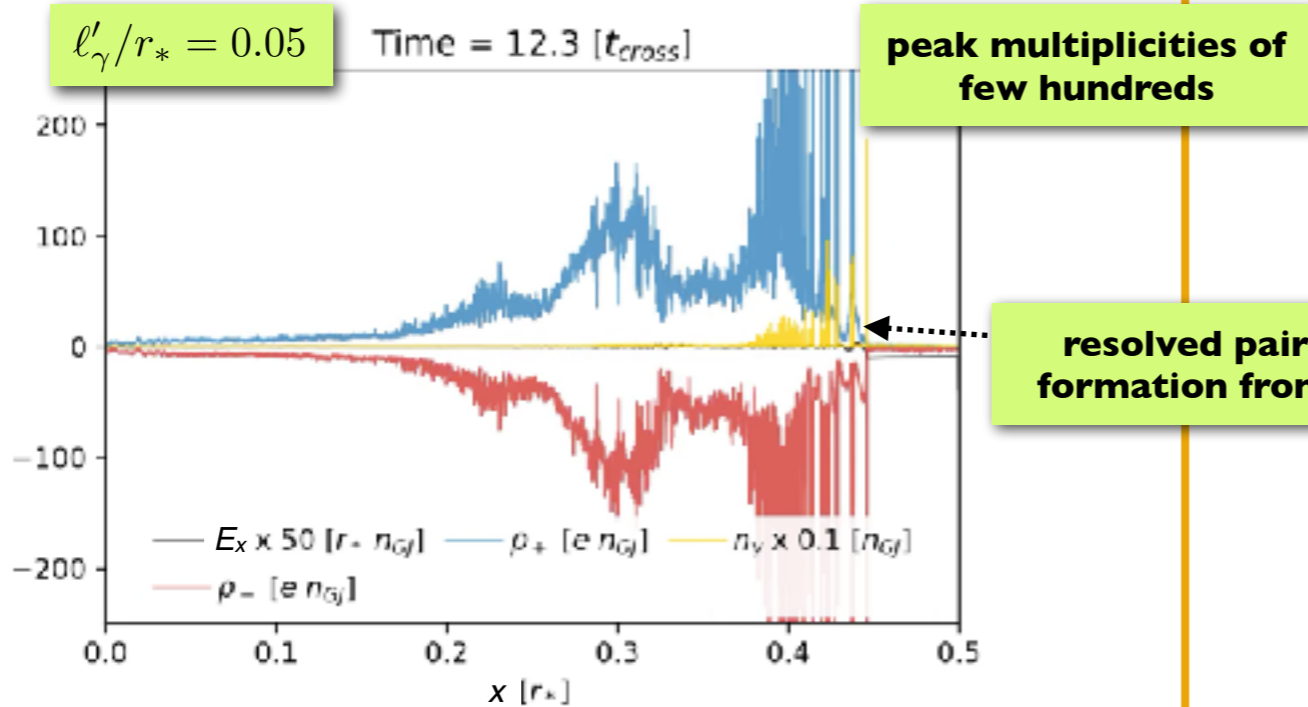
Modelling pulsar polar vacuum gaps

ID cascades with constant curvature **B** show cascade-shielding cycles with:

Co-rotation frame $\nabla \cdot \mathbf{E} = 4\pi(\rho - \rho_{GJ})$

An imposed current component \mathbf{J}_m required by the magnetosphere

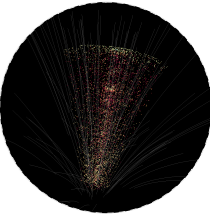
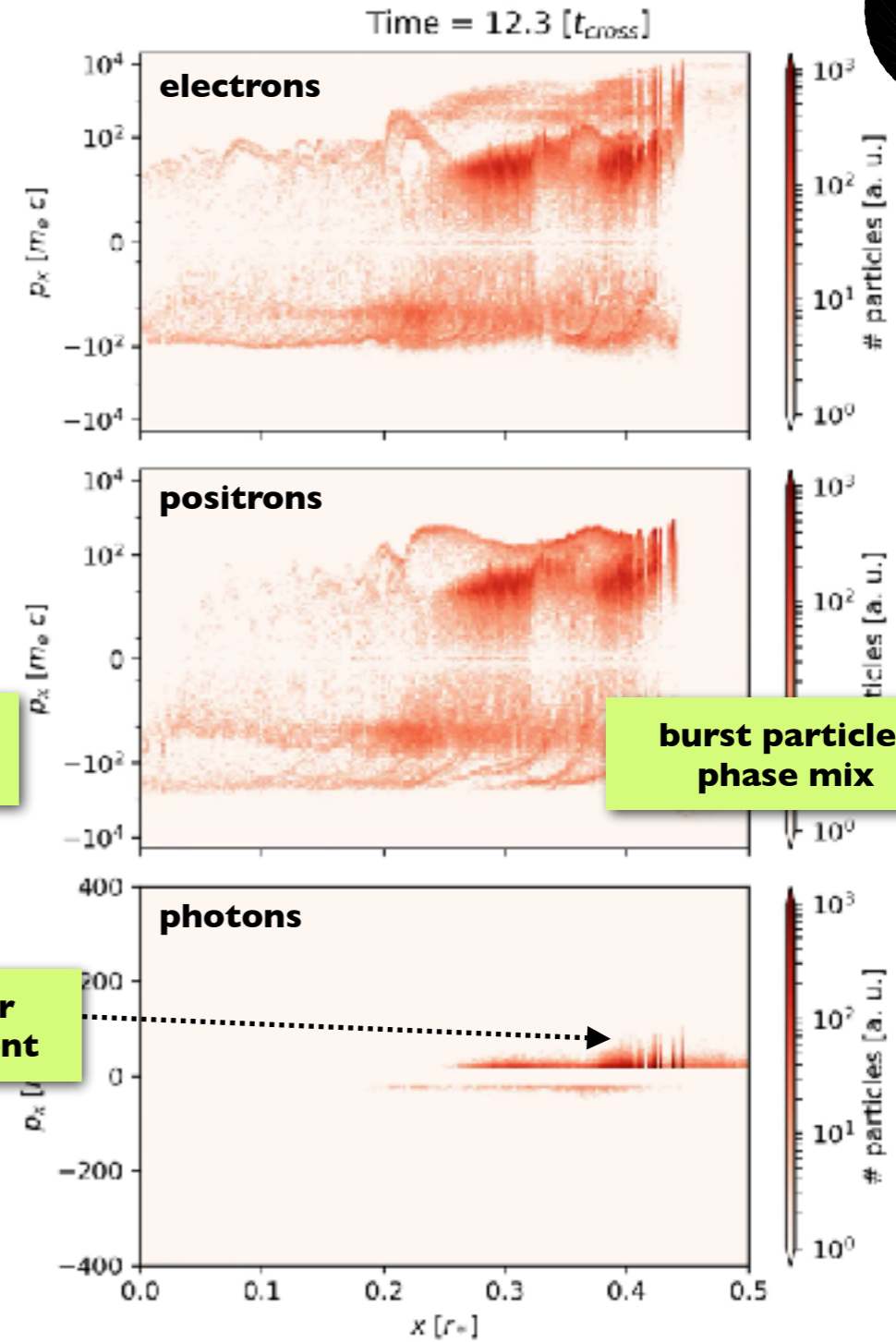
$$\frac{\partial \mathbf{E}}{\partial t} = c \nabla \times \mathbf{B} - 4\pi(\mathbf{J} - \mathbf{J}_m) \quad \mathbf{J}_m = 1.5 \rho_{GJ} c \hat{\mathbf{x}}$$



peak multiplicities of few hundreds

resolved pair formation front

burst particles phase mix



F. Cruz *et al.*, *ApJ* (2021); *idem*, *PRE* (2021)

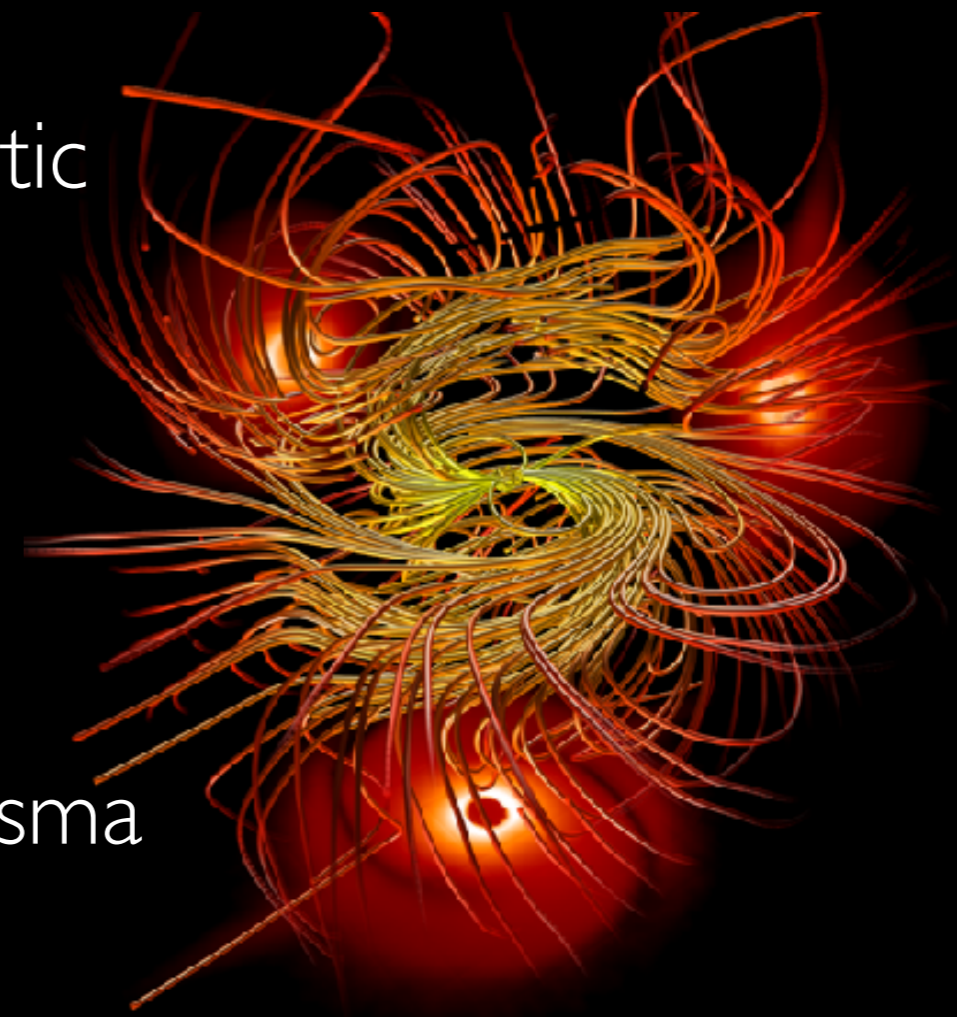
Set-up similar to A. Timokhin & Arons, see also A. Philipov *et al.*, *PRL* (2020)

Several open questions in fundamental QED processes in intense fields + *expressing/* benchmarking those processes in plasma kinetic codes

Optimization/configurations for secondary sources of gamma rays and e^+

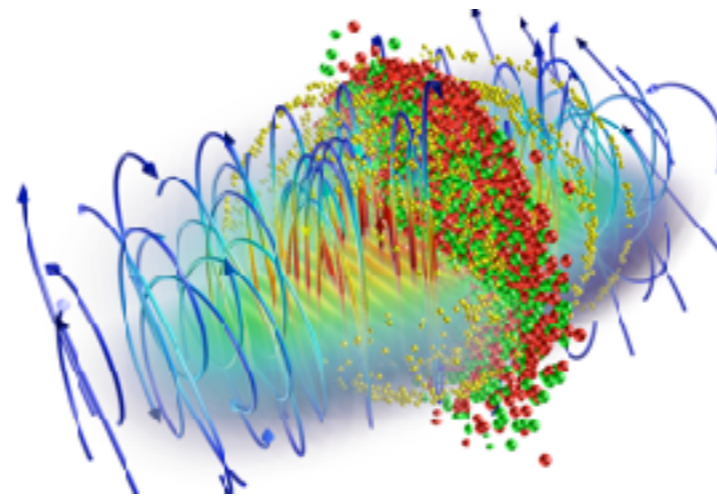
Radiation signatures of collective extreme plasma processes in laboratory and in astrophysics

Coupling with GR



Revolutionary computational power is reshaping our understanding of plasmas in extreme conditions

We now have the ability to explore multi-scale processes from first principles with unprecedented detail and explore unique highly nonlinear scenarios - some of these to be explored soon in the laboratory, others with astrophysical consequences



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Advanced Program in Plasma Science and Engineering

A new call is open to grant 3 PhD fellowships within the framework of the APPLAuSE Doctoral Programme.

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Application deadline: 28 January 2022

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