

# Tackling critical slowing down using global correction steps with equivariant flows within the 2D Schwinger model

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Meinerzhagen  
16.08.2022

## Simulations at the precision frontier

- critical slowing down via Dirac's index
- global corrections within Monte Carlo simulations

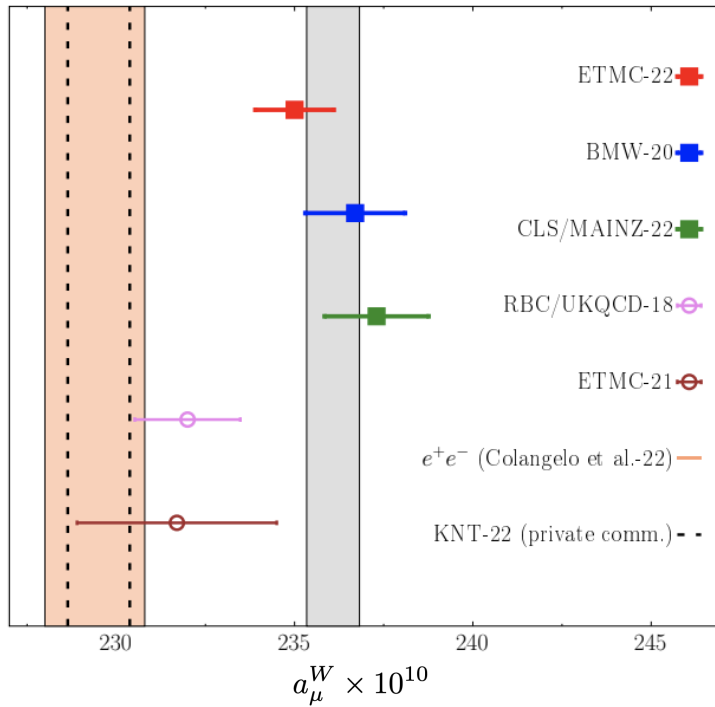
## Generative models for $U(1)$

- insides into gauge invariant flows
- scalability via domain decomposition

## Global corrections with the fermion determinant

- towards high acceptance rate
- towards low autocorrelation

## Comments on steps towards 4D-QCD ...



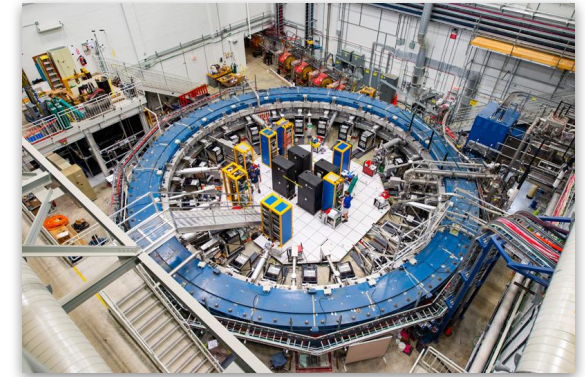
## Exciting times for Lattice Quantum Chromodynamics

Muon and Flavor Physics are indicating New Physics; ab initio LQCD calculations are needed

Search for new physics in the precision frontier by

- high precision measurements
- theoretical prediction

deviations are signs for new physics



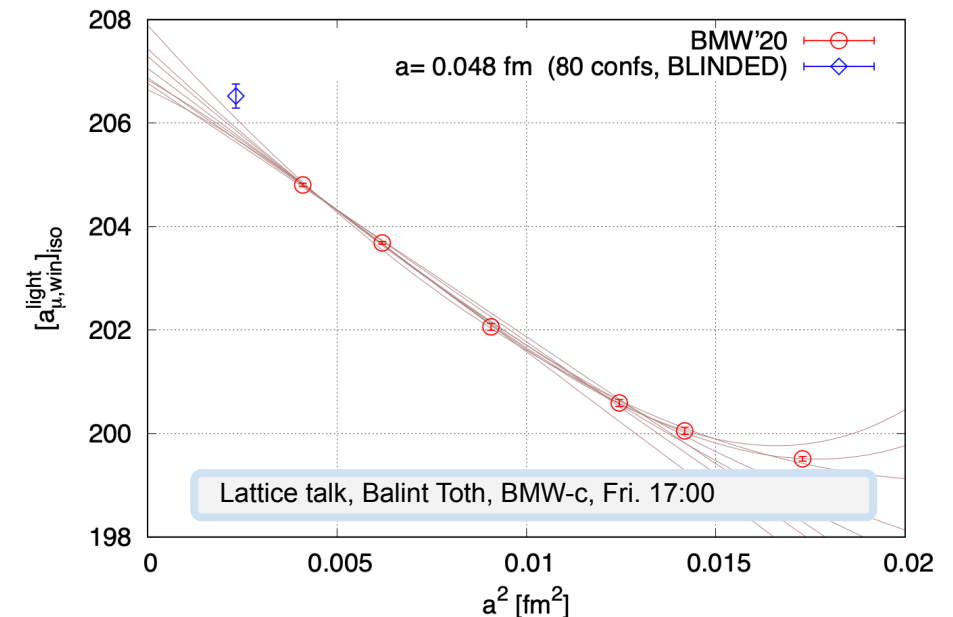
### Anomalous magnetic moment of muon:

Muon g-2 Experiment at FermiLab confirmed results

- $4\sigma$  deviation between experiment and data-driven approach
- $4\sigma$  deviation between lattice and data-driven approach

To resolve this puzzle:

**Precision Measurement of Lattice QCD are needed**





## Simulation at the Precision Frontier:

Very fine lattice spacing needed to match future experiments precision

Standard large scale MCMC method:

- Hybrid Monte Carlo (HMC) algorithm
  - based on molecular dynamics

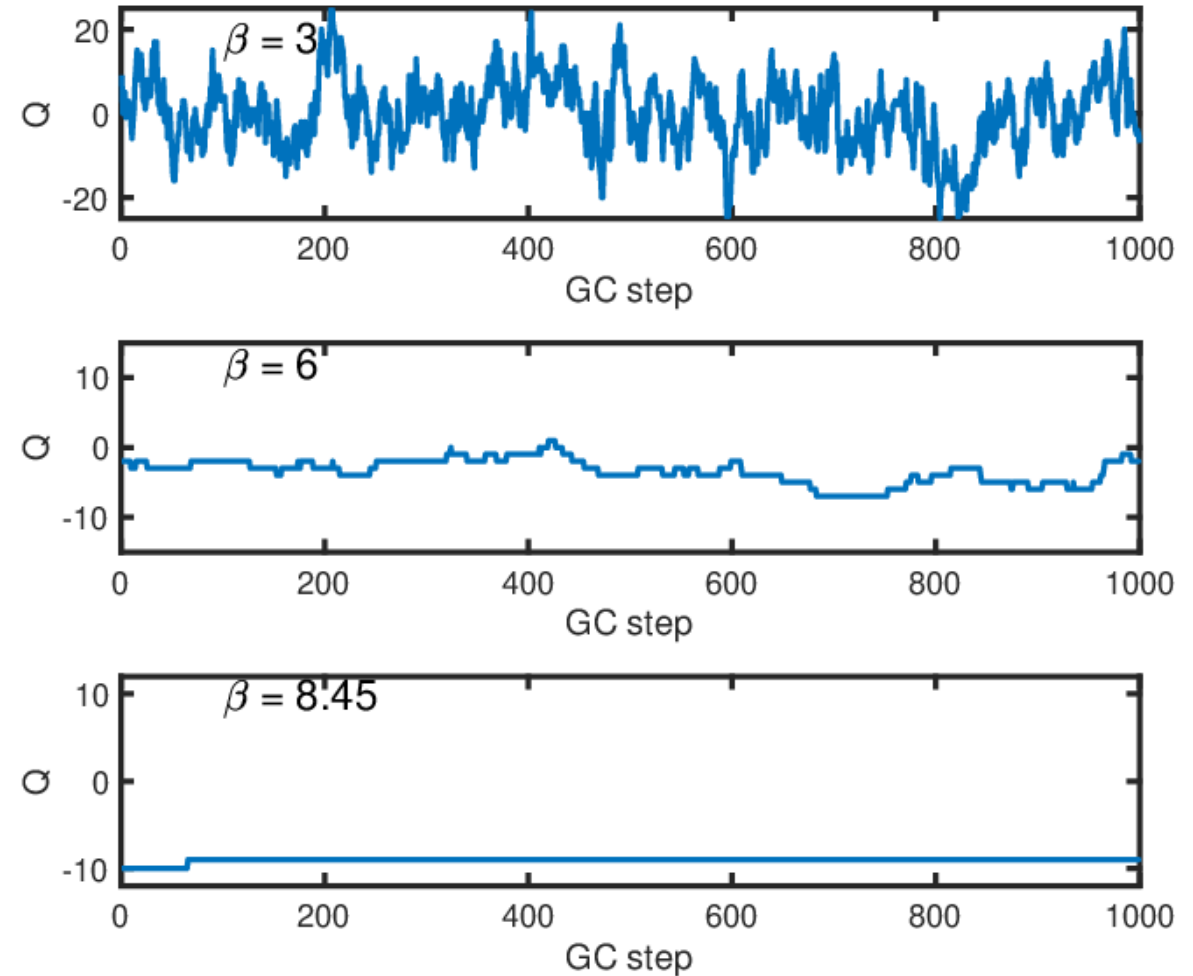
$$\dot{P} = -\frac{\partial H}{\partial U} \quad \text{and} \quad \dot{U} = \frac{\partial H}{\partial P}$$

for very fine lattice spacings  $a < 0.05$  fm  
the HMC algorithm freezes out a topological sector

S. Schaefer et al., Null. Phys. B 845 (2011) 93-119

severe critical slowing down

- Efficient algorithm in QCD missing (openBC would be a possibility)



here, we will take a look to a simpler case: 2D Schwinger Model

The Index theorem gives some illustrative insides:

$$N_R - N_L = Q^{geo}$$

Ativah and Singer. 1963

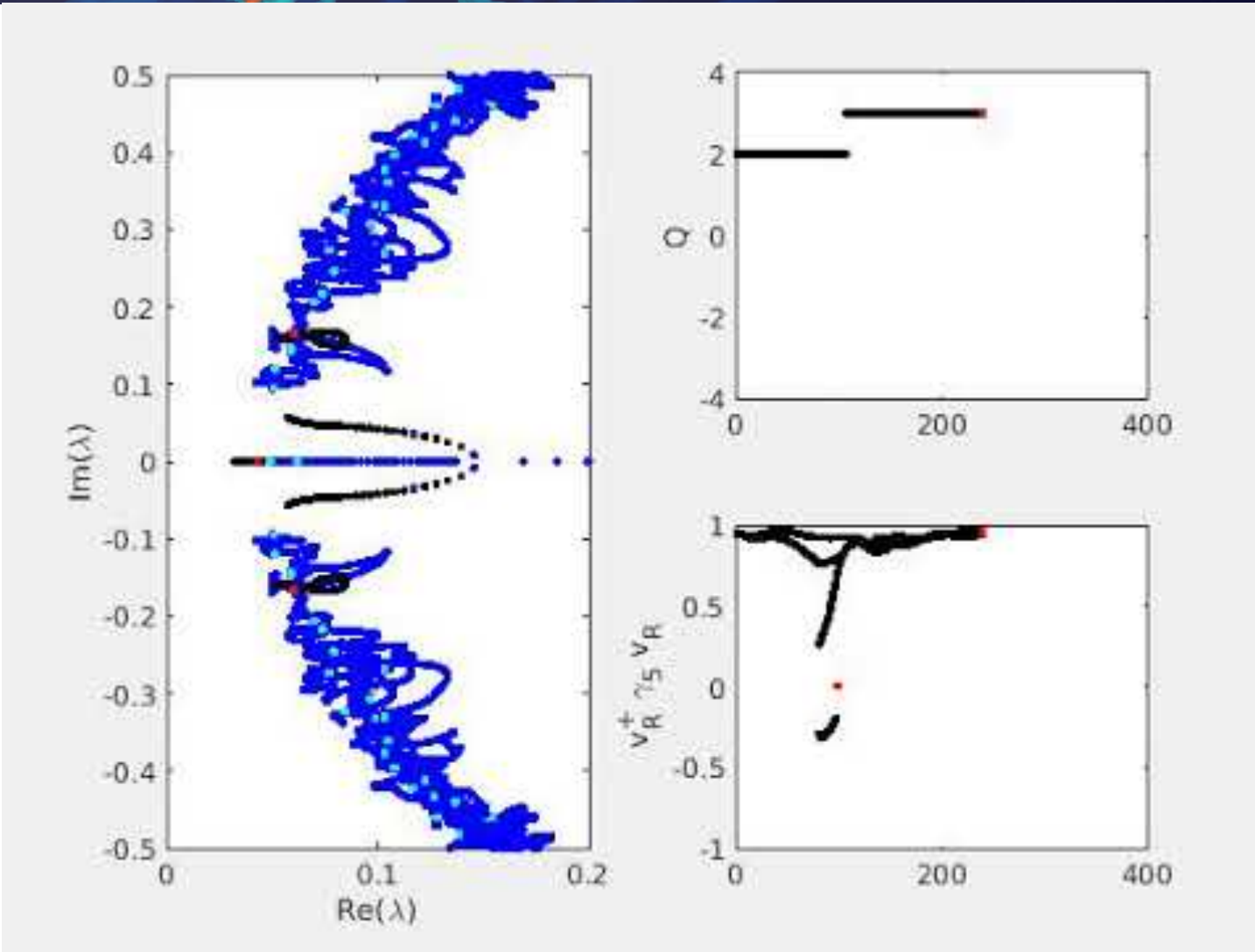
with the geometric definition:  $Q^{geo} = \frac{1}{2\pi} \sum_x \theta_{12}(x)$

and  $\text{Index}(D) = N_R - N_L = \sum_i \chi_i|_{\lambda(D)=0}$

$\chi_i \equiv \text{sign}(v_{i,R}^\dagger \gamma_5 v_{i,R})$

Microcanonical simulations suppresses Q transition

$$U \longrightarrow U'$$
$$H(U) \equiv H(U') = P^2 + \beta S(U') - 2 \sum_i \ln \lambda_i$$



## General structure:

1. Propose  $U'$  according to  $T_0(U \rightarrow U')$
2. Correct with  $P_{acc}(U \rightarrow U') = \min \left[ 1, \frac{\tilde{\rho}(U)\rho(U')}{\rho(U)\tilde{\rho}(U')} \right]$

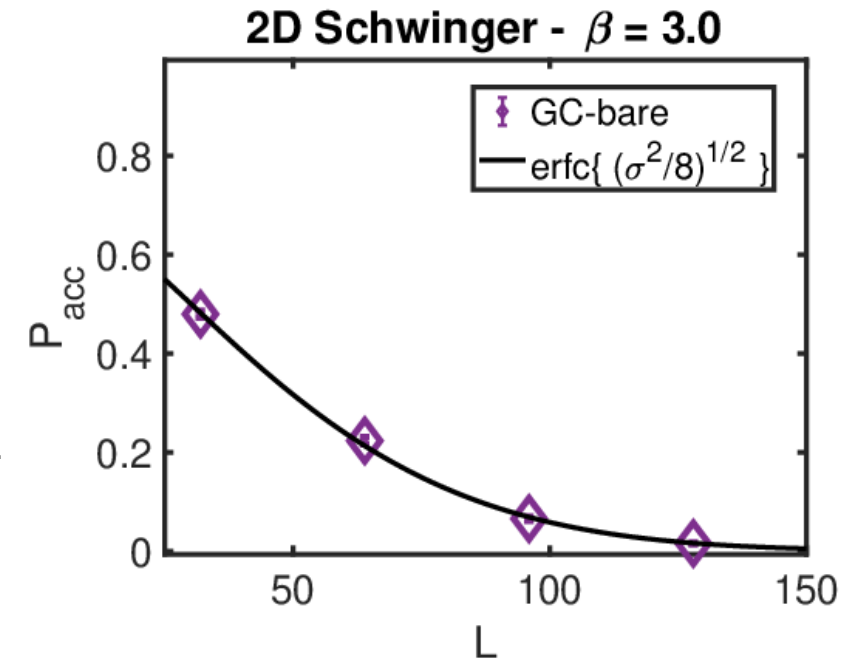
In case ratio of distributions  $(\tilde{\rho}(U)\rho(U'))/(\rho(U)\tilde{\rho}(U'))$  is log-normal distributed.

- for the acceptance rate follow [Creutz, Phys. Rev. D38 \(1988\) 1228–1238](#)

$$P_{acc} = \text{erfc}\left\{ \sqrt{\sigma^2(\Delta S)/8} \right\}$$

with  $\Delta S = \ln\{\rho(U')\} - \ln\{\rho(U)\} + \ln\{\tilde{\rho}(U)\} - \ln\{\tilde{\rho}(U')\}$

where  $\ln(\rho(U))$  is an extensive quantity, thus  $\sigma^2(\Delta S) \propto V$



- GC-step is very fast ineffective :

$$P_{acc} \rightarrow e^{-V}$$

How to control  $\sigma^2(\Delta S)$

1. by using correlations between  $\rho$  and  $\tilde{\rho}$
2. by reduction of degrees of freedom of  $\rho$  and  $\tilde{\rho}$

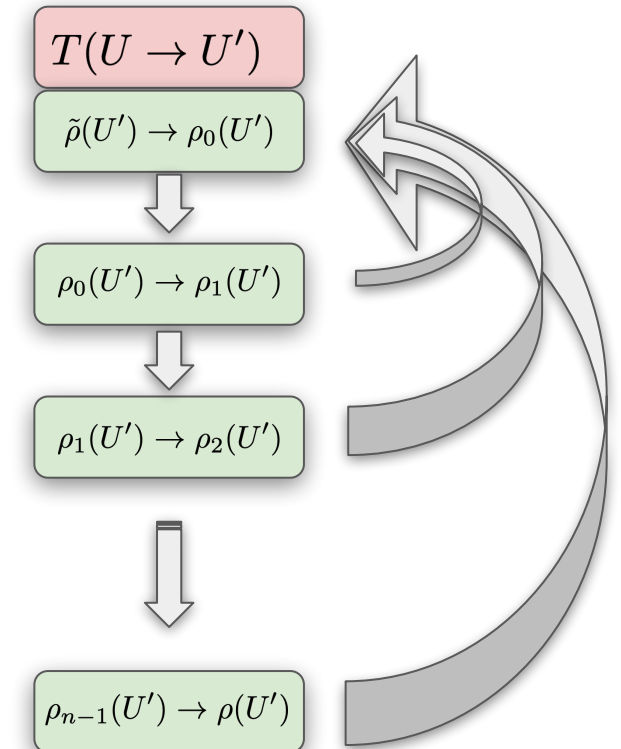
Generalization leads to factorization with parametrization of  $\rho$  via

$$\rho_n(U) = P_0(U, \alpha_i^{(0)}) P_1(U, \alpha_i^{(1)}) \dots P_n(U, \alpha_i^{(n)})$$

and GC step is splitting up into  $n$  successive steps

$$P_{acc}^i(U \rightarrow U') = \min \left[ 1, \frac{\rho_{j-1}(U, \alpha_i^{(j-1)}) \rho_j(U', \alpha_i^{(j)})}{\rho_j(U, \alpha_i^{(j)}) \rho_{j-1}(U', \alpha_i^{(j-1)})} \right]$$

- Iterate each step to filter out local fluctuations

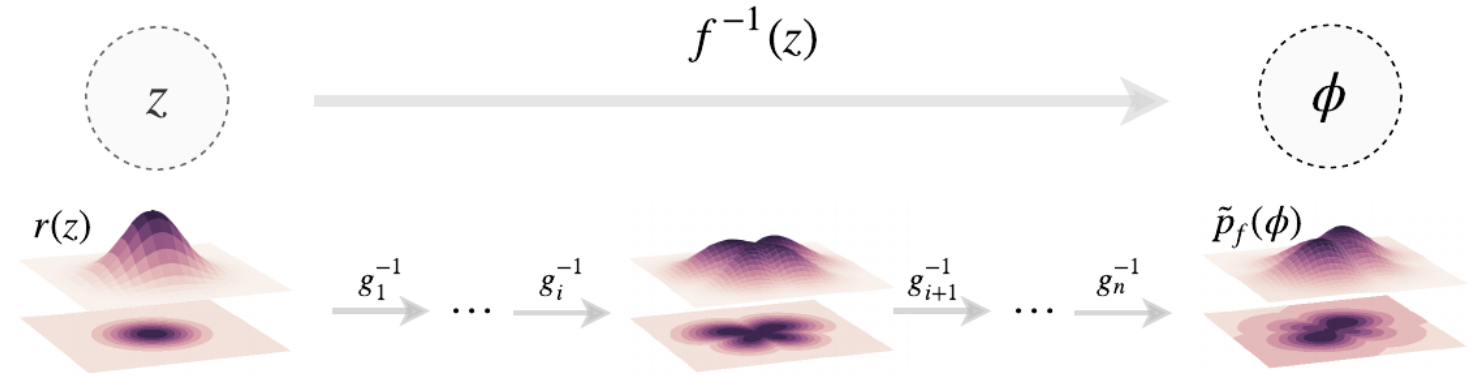


## An example: Generative model in U(1) with gauge invariant flow

### Idee:

Use a flow map  $f^{-1}(z)$  to propose new configurations with known distribution

$$\tilde{p}(\phi) = r(f(\phi)) \cdot \left| \det \frac{\partial f(\phi)}{\partial \phi} \right|$$



- introduce coupling layers with

$$g_i^{-1}(z) := \begin{cases} \phi_a = z_a \\ \phi_b = (z_b - t_i(z_a)) \odot e^{-s_i(z_a)}. \end{cases}$$

- train the coupling layers (s,t) by minimizing the loss-function

$$\begin{aligned} L(\tilde{P}) &:= D_{KL}(\tilde{P}||p) - \log Z \\ &= \int \prod_j d\phi_j \tilde{P}(\phi) (\log \tilde{P}(\phi) + S(\phi)). \end{aligned}$$

successfully applied to ultra local 2D discrete lattice models by

- $\phi^4$  Albergo et al., Phys.Rev.D 100 (2019) 3, 034515
- U(1), Kanwar et al., Phys.Rev.Lett.125 (2020) 12, 121601
- SU(2), SU(3) Boyda et al., Phys.Rev.D 103 (2021) 7, 074504

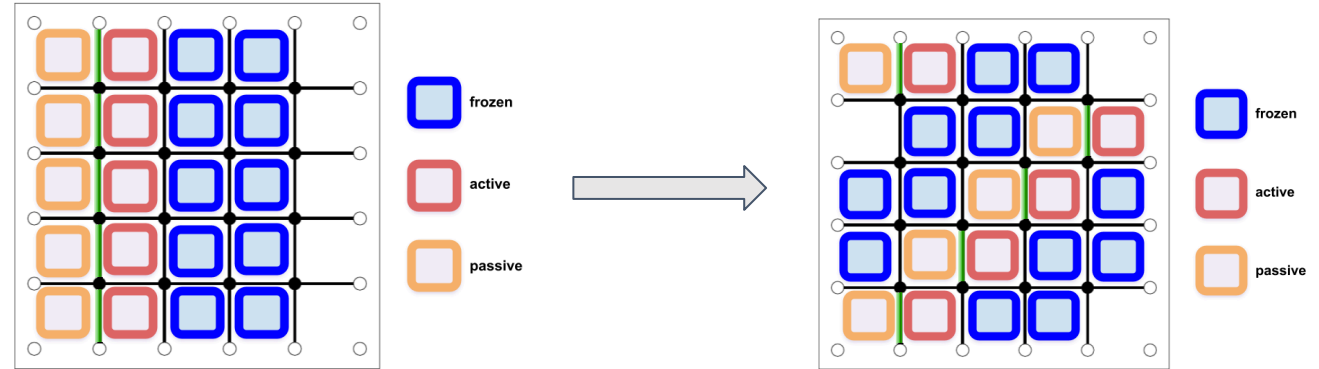
- can overcome critical slowing down

Albergo et al., arXiv:2101.08176



## How to design coupling layers:

$$g_i^{-1}(z) := \begin{cases} \phi_a = z_a \\ \phi_b = (z_b - t_i(z_a)) \odot e^{-s_i(z_a)}. \end{cases}$$

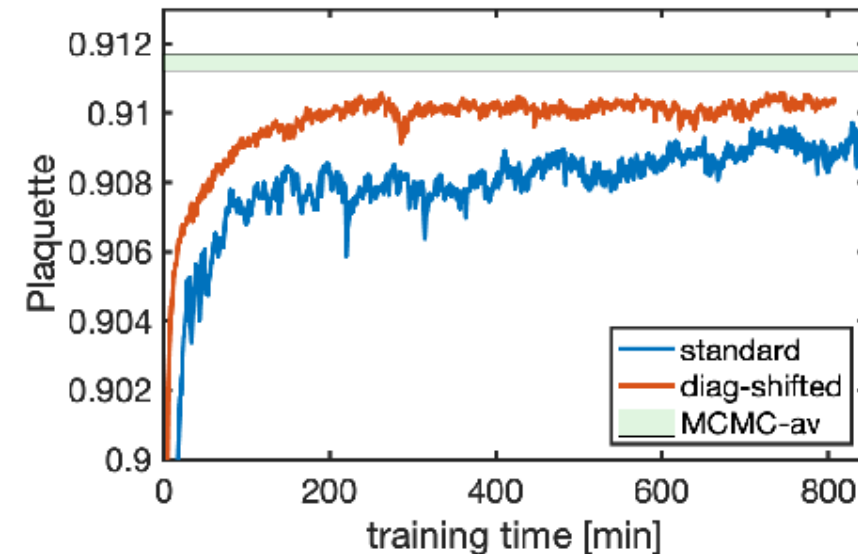


$t_i$  and  $s_i$  consists of neural networks

- Can be design to contain symmetries
  - Gauge invariant by masks and proposing plaquettes
  - Partially translation invariant by convolutional networks

### Structure of networks

- convolutional kernels with size 3
  - note that only frozen plaquettes are used as input values
- with hidden layers (here default 2 with 8 nodes)
- 8 coupling layers corresponds to a full update



Updating masks improve convergence rate and acceptance from 30% to 50%

Let's defined our minimization condition:

The loss function:

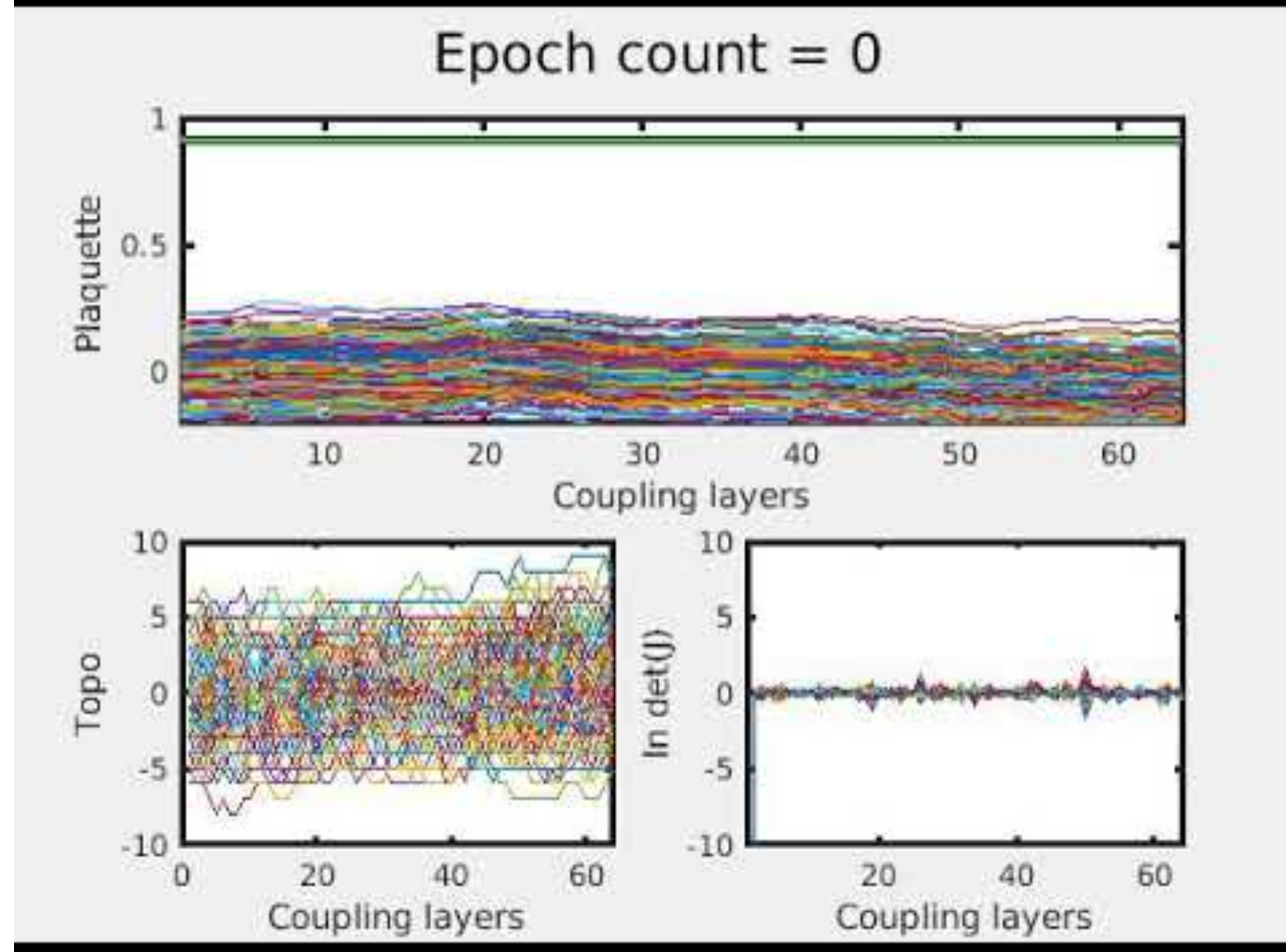
$$\begin{aligned} L(\tilde{P}) &:= D_{KL}(\tilde{P}||p) - \log Z \\ &= \int \prod_j d\phi_j \tilde{P}(\phi) (\log \tilde{P}(\phi) + S(\phi)). \end{aligned}$$

- with ultra-local plaquette action:

$$\ln(\rho(U)) = -\beta \sum_x P_{12}(U)$$

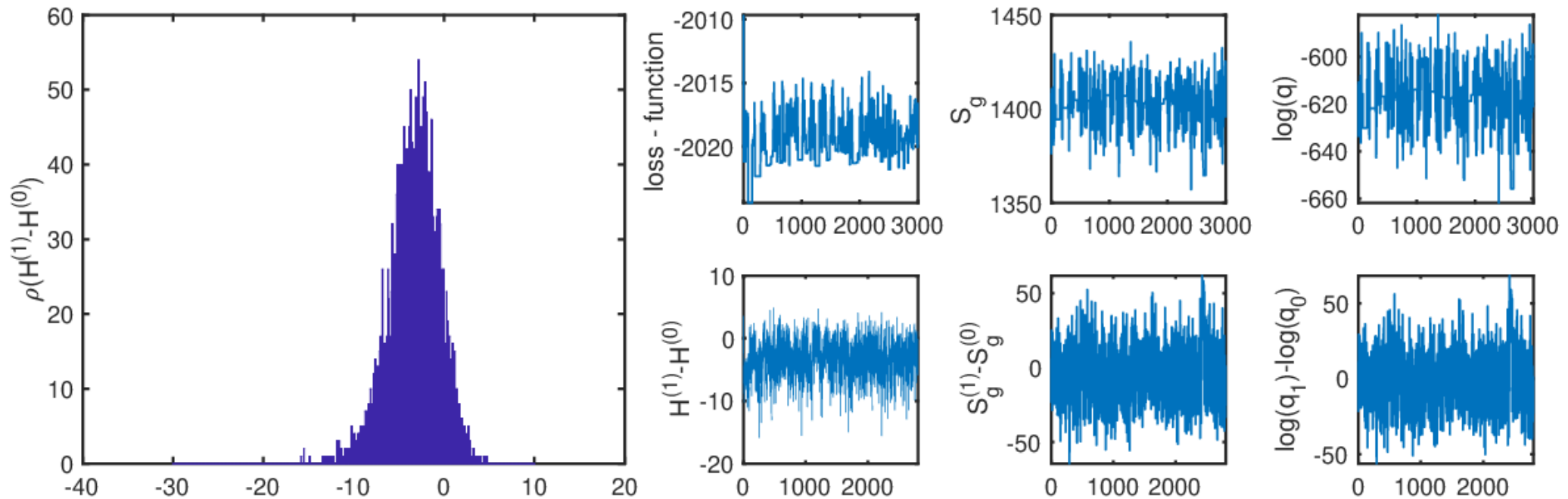
- and flow distribution:

$$\tilde{\rho}(U) = \rho_{trival}(m^{-1}(U)) \prod_j \det J(g_j^{-1}(\alpha_{i,j}^{(0)}))$$



## Correlations of distribution $\tilde{\rho}$ and $\rho$

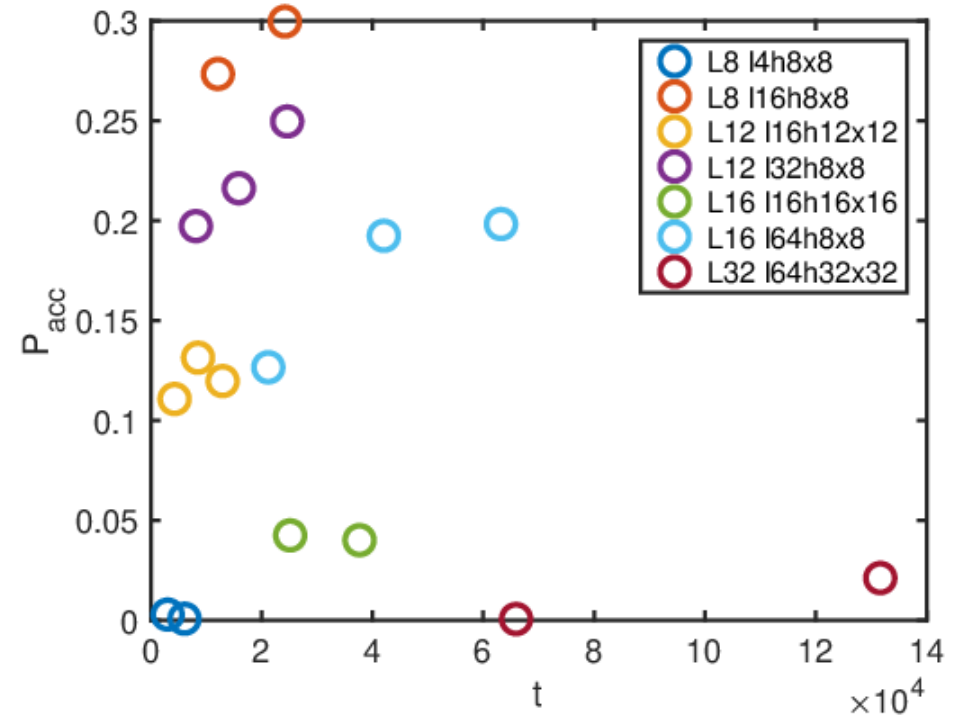
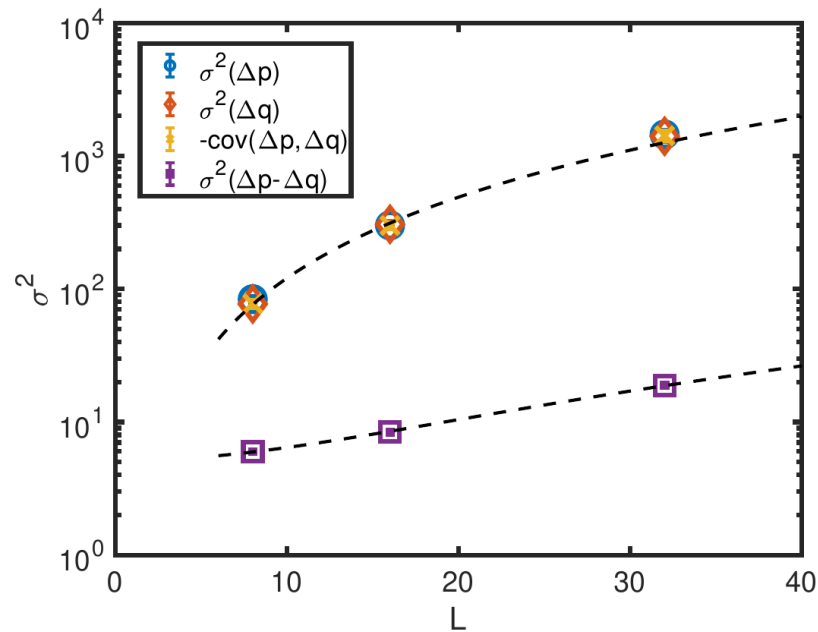
- covariance need to be of  $\text{cov}(\tilde{\rho}, \rho) \propto \mathcal{O}(V)$  to compensate extensive variances  $\sigma^2 \propto \mathcal{O}(V)$



Works for  $L=8 \rightarrow L=16$

## Volume scaling of gauge invariant flow:

- coupling layer dof are scaled with volume
  - $l$ : coupling layers
  - $h$ : hidden layers
- scaled  $l$  and  $h$  with  $V=L^2$  while decreasing minimization rate



## Fine tuning problem:

Covariances of distributions scales like variances

$$\text{var}(\Delta p) + \text{var}(\Delta q) \approx 2\text{cov}(\Delta p, \Delta q)$$

But  $\sigma^2 = \text{var}(\Delta p) + \text{var}(\Delta q) + 2 \cdot \text{cov}(\Delta p, \Delta q)$   
still grows with the volume



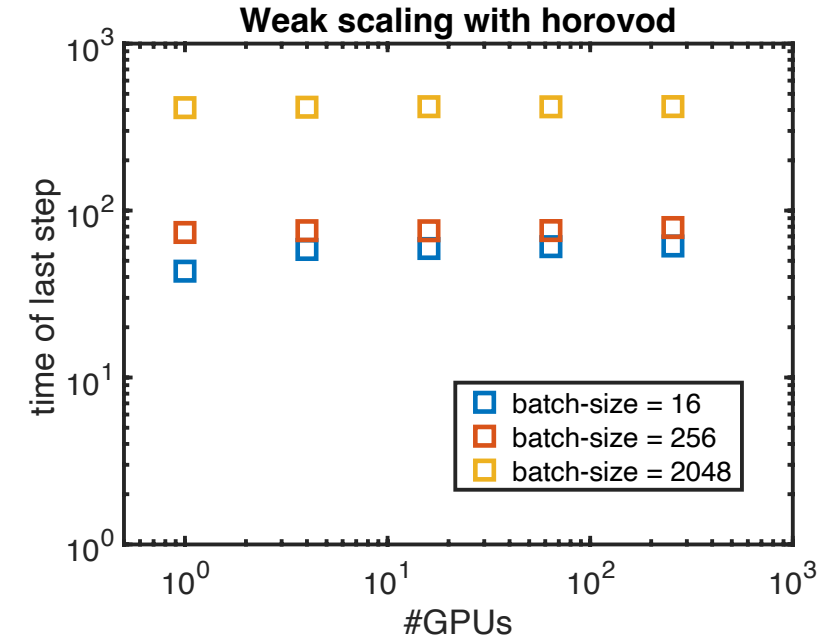
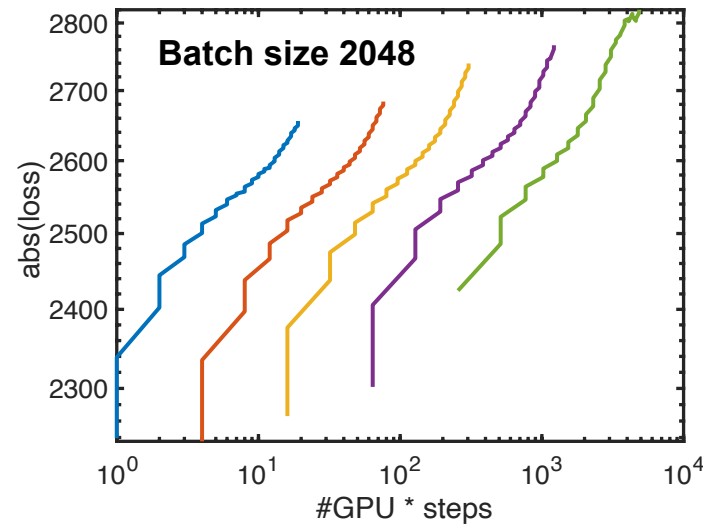
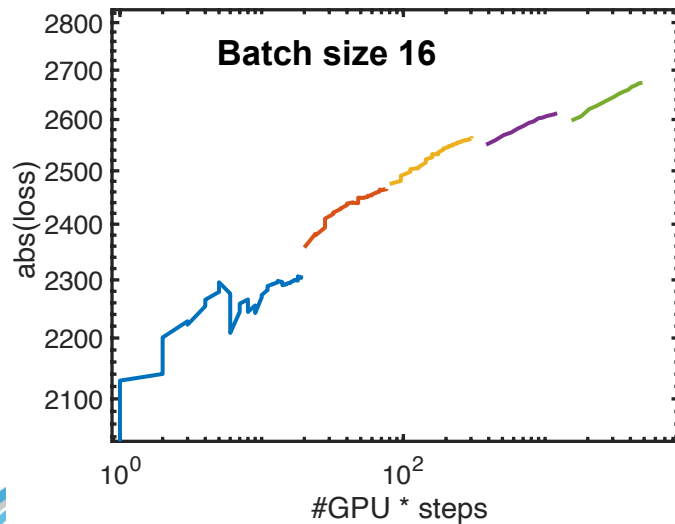
## How to scale up:

### Exercise with horovod

- Simple to implement but needs fine tuning
- adds new batch to each additional GPU
  - Total batch-size = #GPUs x local batch-size

#### Modifications:

- Switch to double precision
- Use Ada..
- Use stepsize decay



### Benchmark runs on JUWELS-BOOSTER

- Loosely coupled scales weakly perfect
- For smaller batch-sizes works fine
- For larger batch-sizes convergence deteriorates

## Lattice action are local

every highly optimized lattice algorithm are based on it

- multigrid, multi level, hierarchical probing, low-mode averaging, etc.

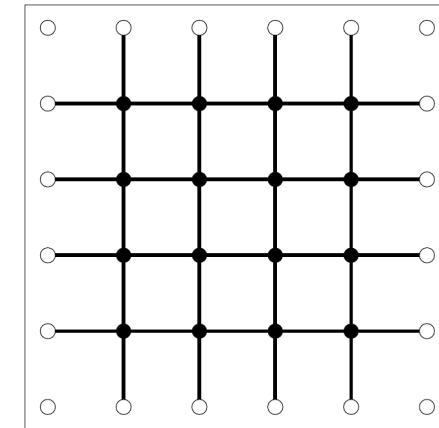
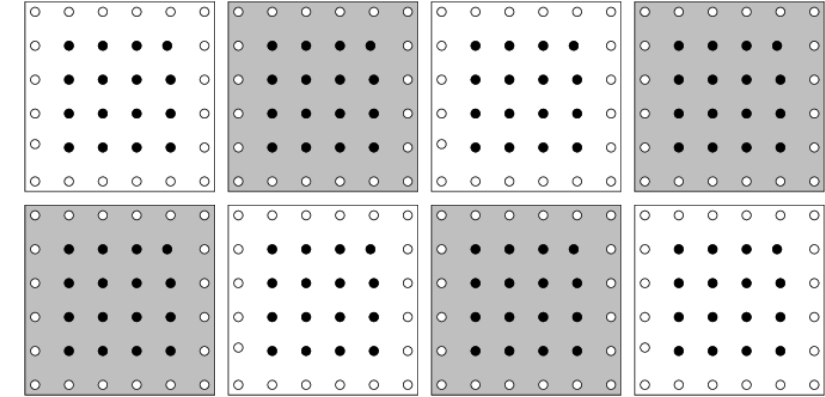
mainly based on Domain Decomposition of the lattice  
then the ultra local plaquette action splits up into

$$S_g = -\beta \left\{ \sum_b \sum_{x \in b} P_{12}(x) + \sum_{\delta b} \sum_{x \in \delta b} P_{12}(x) \right\}$$

Plaquettes inside of the blocks      Plaquettes between blocks

## For the gauge invariant flow

- update only links/plaquettes inside blocks
- create maps of active links within each block

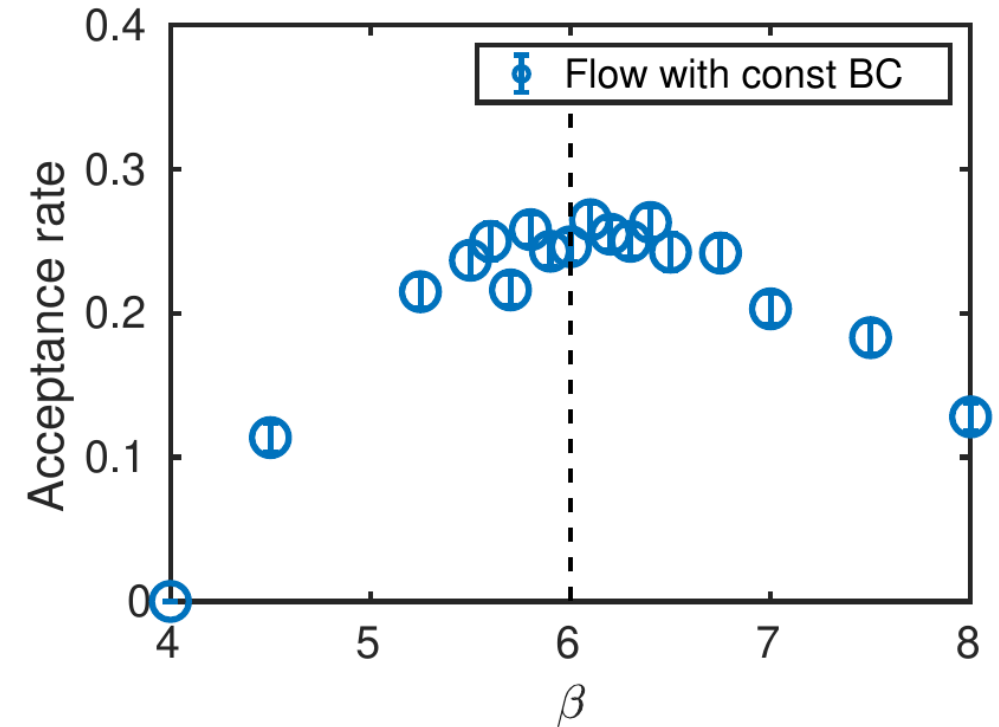


Taken from: M. Luscher, CPC 165 (2005) 199-220

## Adaptation of training procedure

By:

- Using the periodic trained model to generate boundaries or starting from random and shift lattice after each epoch
- Using different boundaries for each batch with total batch size 4096
- Increase iteration before boundaries updated to 1000
- Using diagonal masks to increase overlap with frozen plaquettes (faster convergence)



Acceptance rate of fixed boundaries drops down to  $\sim 25\%$  with  $L = 8$  (from 50% periodic case)

- due to the ultra locality of gauge action:  
larger volumes are trivial to generate

**Action with fermions:**

$$P(U) = Z^{-1} \left( \prod_j^{N_f} \det D_j(U) \right) e^{-\beta S_g(U)}$$

with  $\det D(U)$  is a *localised* action

- distance interaction decays with

$$cov(x, y) \propto \exp\{-m_{PS}|x - y|\}$$

Idea: using exact decomposition of fermion action:

$$\det D = \det S_{red} \cdot \det S_{pink} \cdot \det D_{blue}$$

effective long range decomposition of the fermion determinant

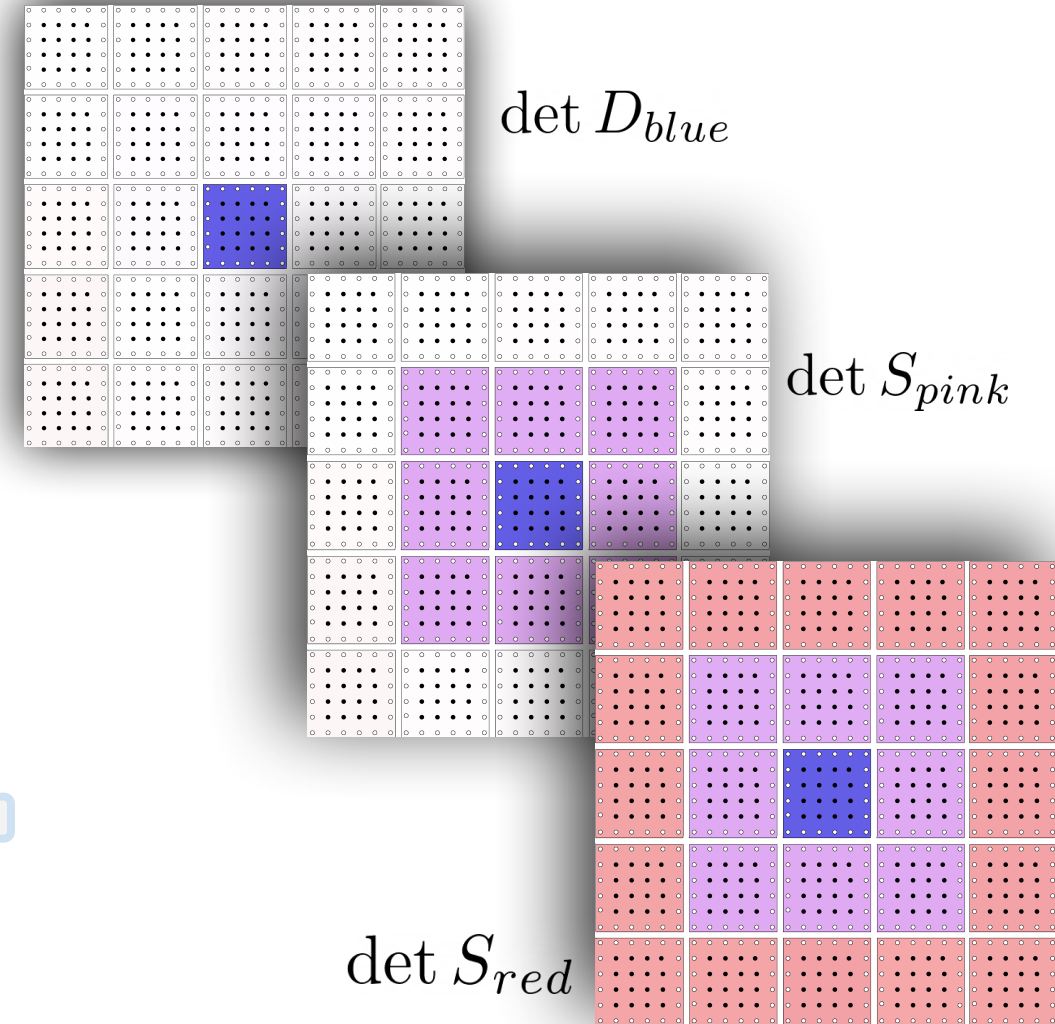
M. Luscher, CPC 165 (2005) 199-220

J. F. et al., CPC 184 (2013) 1522-1534

M. Cè et al., Phys.Rev.D 93 (2016) 9, 094507

M. Cè et al., Phys.Rev.D 95 (2017) 3, 034503

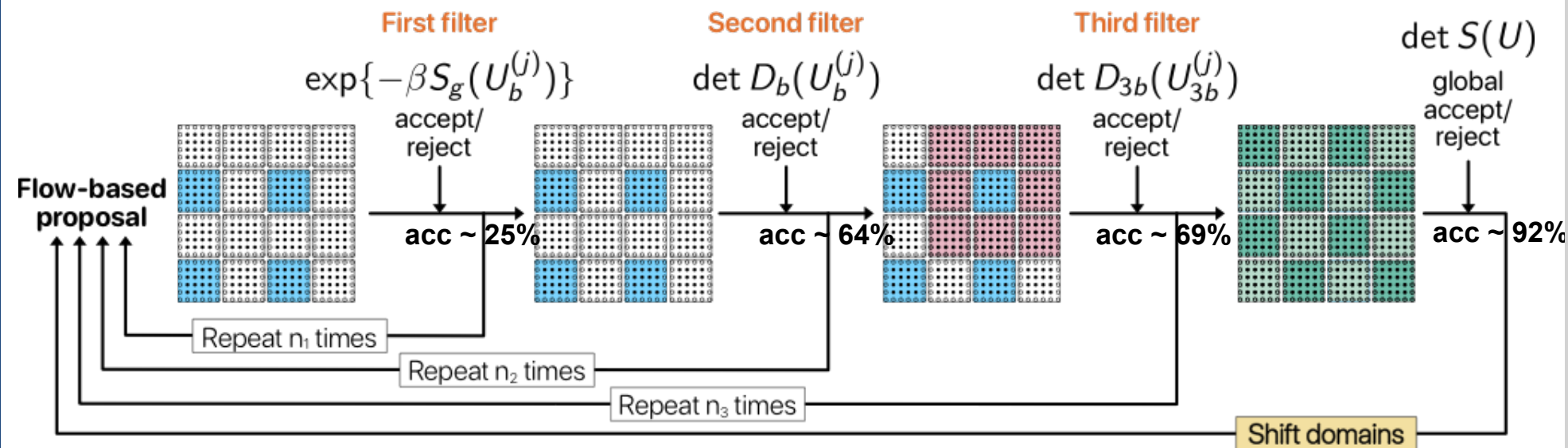
## Recursive Domain Decomposition





## Global Correction Monte Carlo algorithms with equivariant flows:

J. F., arxiv:2201.02216



## Multilevel hierarchical filter steps with 4 levels

Enhancing acceptance rate by

- within level 1, 2, 3 each active block can be updated independently from each other
- use correlation between actions via parameterization,
  - e.g. for the gauge coupling

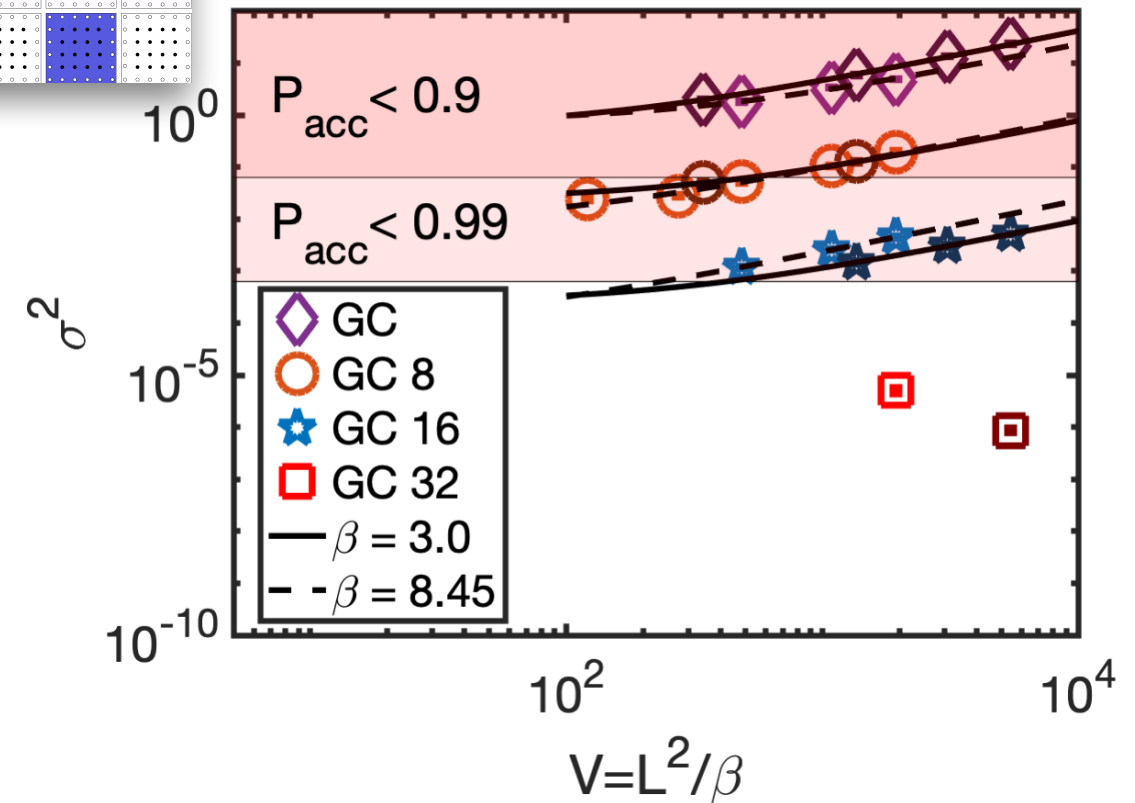
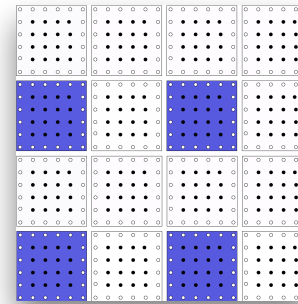
$\beta$	3.0	6.0	8.45
5 level flowGC with $d = 16$ :			
Level 4			
with $\sigma^2$	0.0052	0.0369	0.0046
and $P_{acc}$	0.9713	0.9235	0.9727
$\delta\beta_4^{(3)}$	-2.0037	-2.0182	-2.0087
$\delta\beta_4^{(2)}$	1.0027	1.0061	1.0083
$\delta\beta_4^{(1)}$	-0.0003	0.0008	0.0004
Level 3	$n_1 = 2$		
with $\sigma^2$	0.6688	0.6190	0.1546
and $P_{acc}$	0.6826	0.6940	0.8441
$\delta\beta_3^{(2)}$	-1.1730	-1.3635	-1.3534
$\delta\beta_3^{(1)}$	-0.0006	0.0149	0.0125
Level 2	$n_2 = 4$		
with $\sigma^2$	1.4384	0.8325	0.1857
and $P_{acc}$	0.5487	0.6482	0.8294
$\delta\beta_2^{(1)}$	-0.2482	-0.3082	-0.2863
Level 1	$n_1 = 100$		
with $P_{acc}$	0.5669	0.2501	0.2794
2 level GC:			
with $\sigma^2$	12.3774	9.7119	3.7260
and $P_{acc}$	0.0786	0.1192	0.3345

## Acceptance rate:

- select  $L=8$  flow proposals
- updating every 4th block, which introduces a distance between active blocks by  $d = Lbs$  which results into 16% of links updated per step (independent of global volume!)

- runs for different  $Lbs = 8, 16, 32$  with 4 lvl filter steps
- variance is very efficient reduced for larger  $Lbs$
  - volume scaling remains

How a change of 16% influence sampling rates ?



## Topological charge:

Usually we are using the autocorrelation time for comparison, but HMC freezes and  $\tau_{int}$  is not measurable

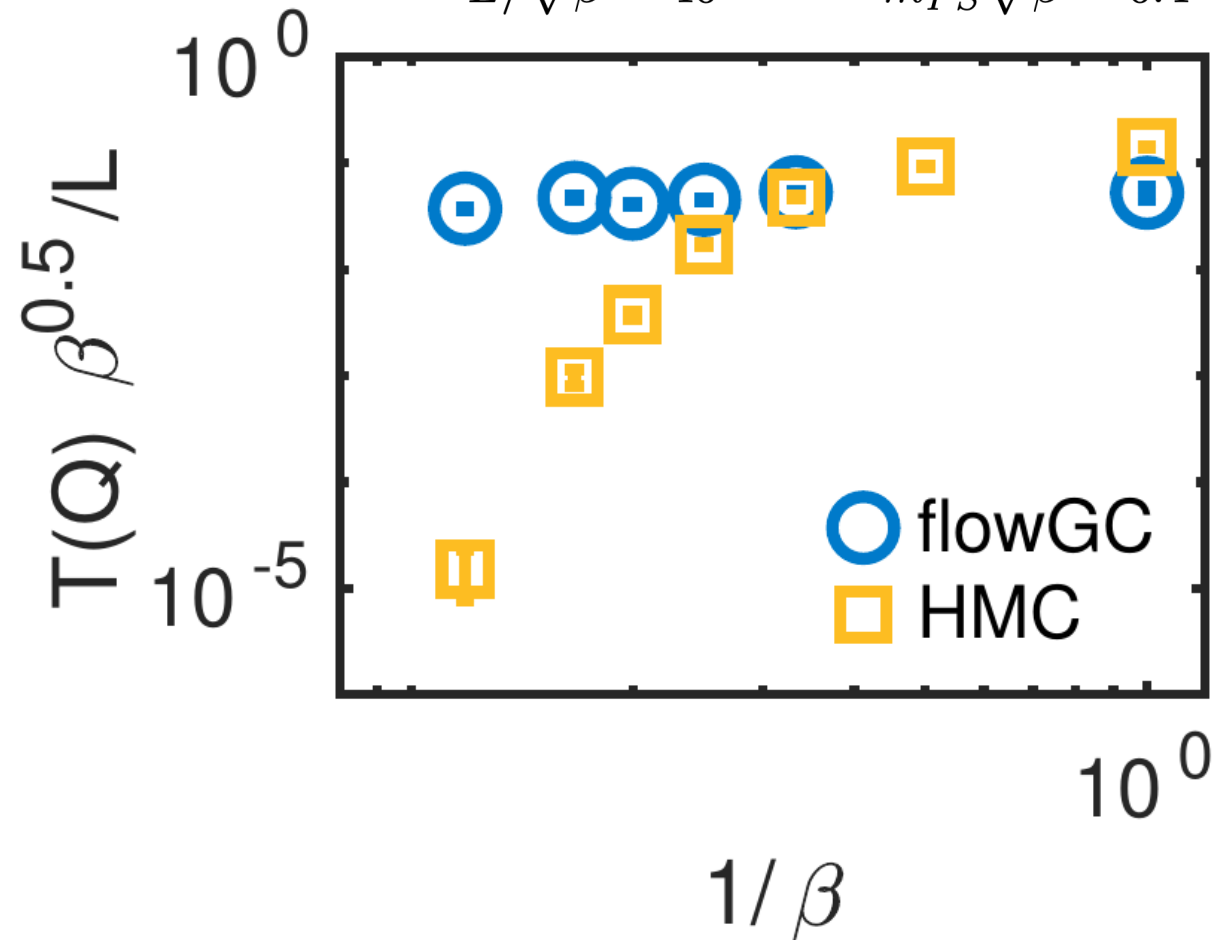
Instead one can define a tunneling rate:

$$T(Q) = \langle |Q_i - Q_{i+1}| \rangle$$

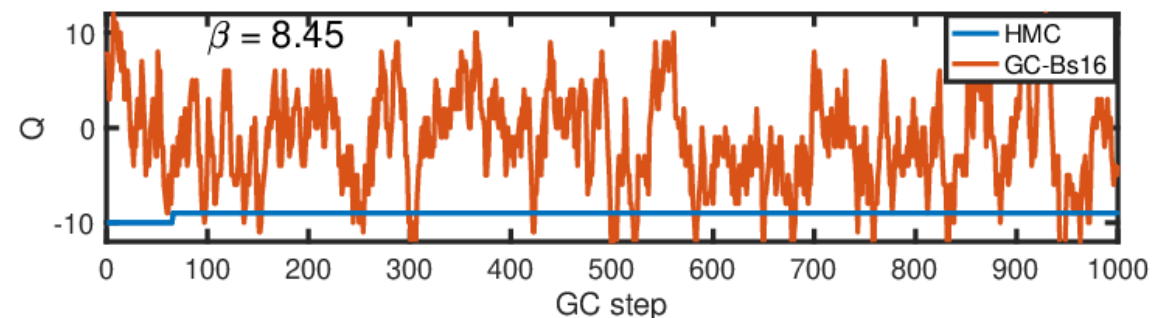
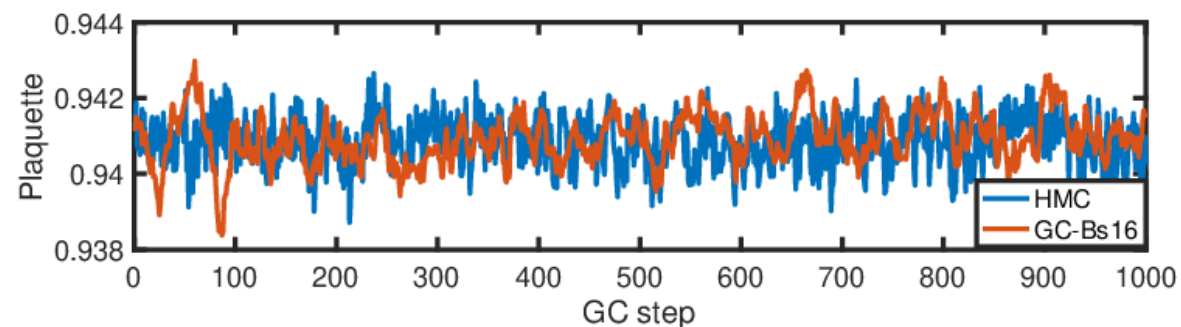
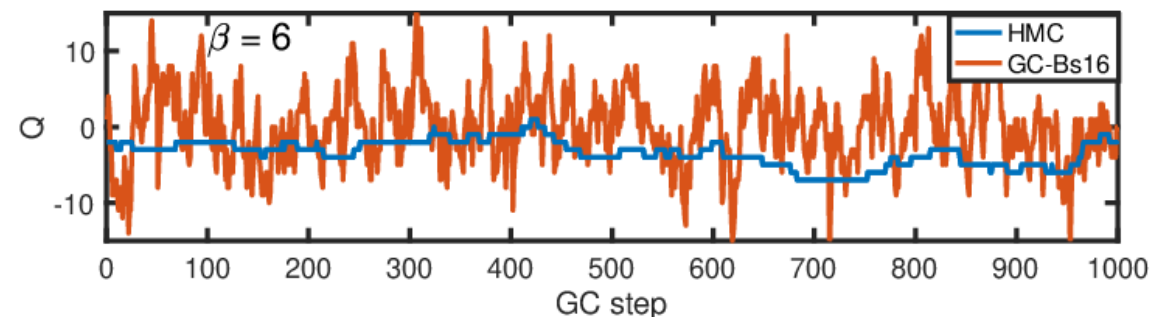
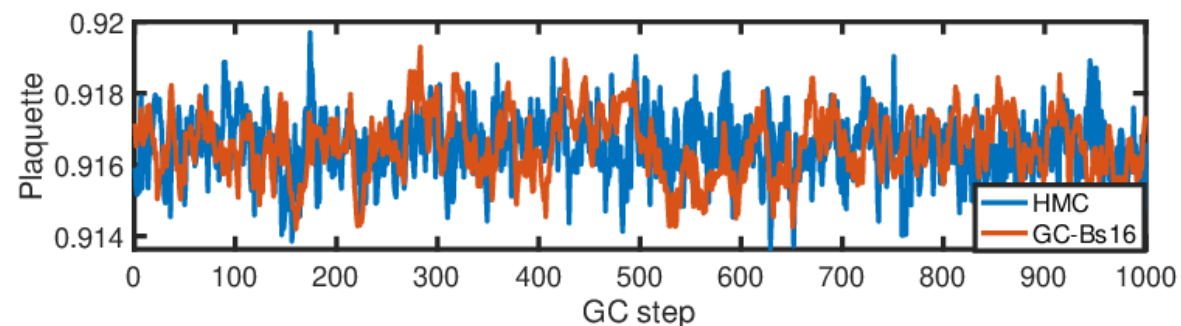
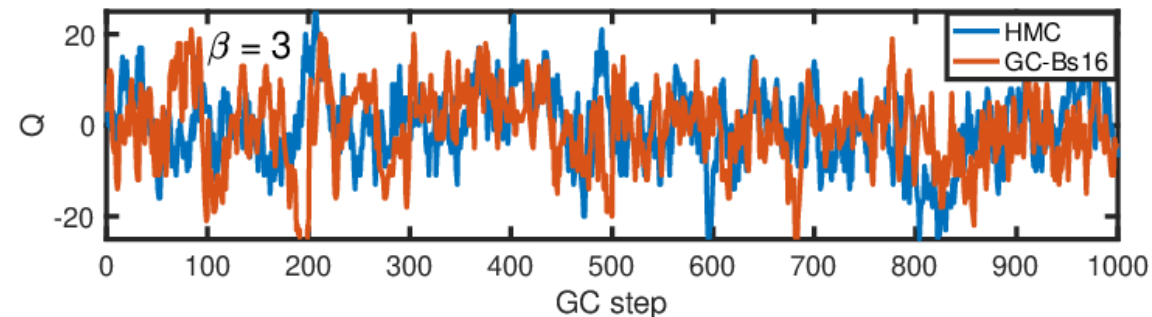
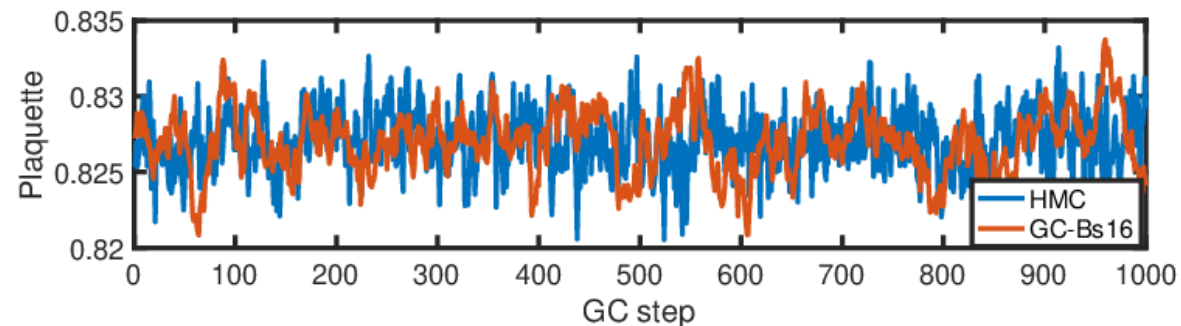
GC shows no critical slowing down and topological tunneling scales

At constant line of physics:

$$L/\sqrt{\beta} \sim 40 \quad \text{and} \quad m_{PS}\sqrt{\beta} \sim 0.4$$



## Runs at L=128

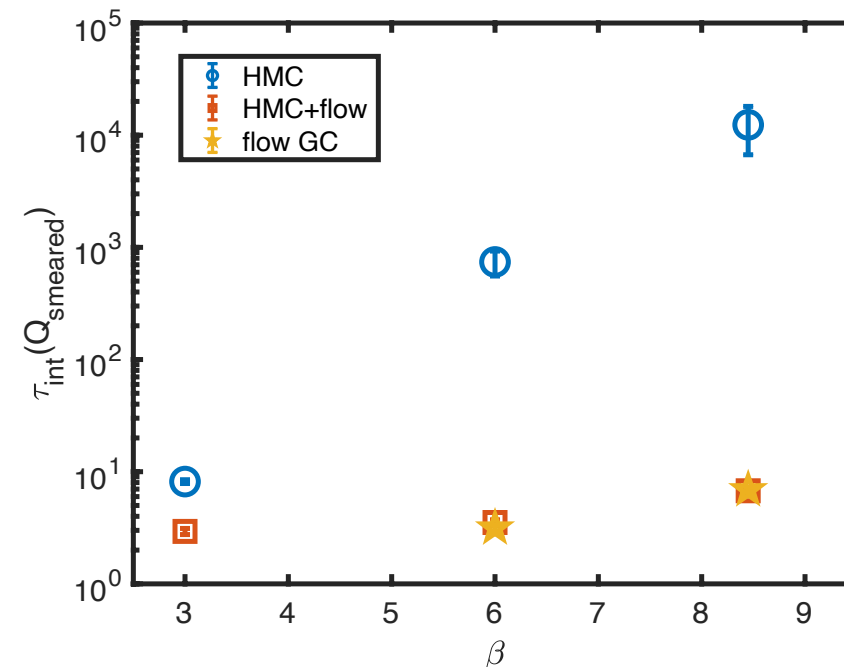
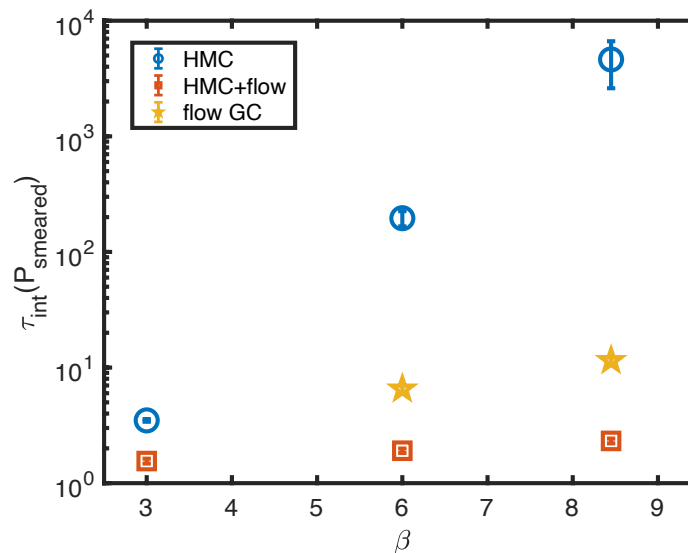
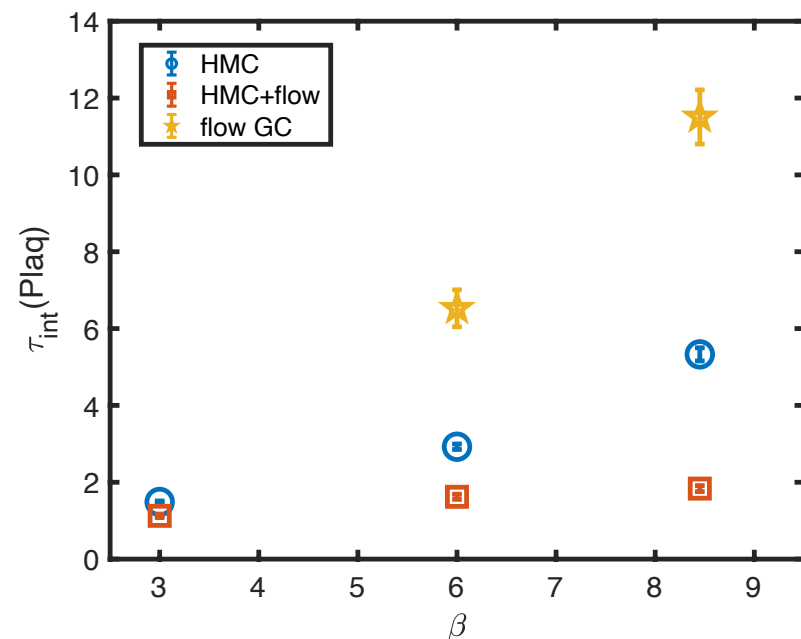




**Idea:** combination with HMC and high statistic runs

Similar to

D. Albandea et al., Eur.Phys.J.C 81 (2021) 10, 873



with  $\tau_{int}(P_0) \propto \beta^c$

- HMC :  $c = -0.5$
- HMC+flow:  $c = -0.5$
- flow GC :  $c = -1.5$

**HMC+flow:**  
Outperforming other  
Methods

**Runs done on L=32**  
**No constant line**  
**of physics**

with  $\tau_{int}(Q) \propto \beta^d$

- HMC :  $d = -7.0$
- HMC+flow:  $d = -0.8$
- flow GC :  $d = -2.0$

## GC+flow proposal can solve critical slowing down in the 2D Schwinger Model

### Major challenges addressed

J. F., arXiv:2201.02216

- very high acceptance rate by keeping 16% of links active towards large volumes
- Tunneling rate of topological charge relative constant towards finer lattice spacings

### Combination with HMC promising towards more complex and larger models

Which depends on:

- Flow proposals within 4D with SU(3)
- Block acceptance can break down (so far  $6^4$  are reached)  
flow proposals with fermions should help

J. F. et al., CPC 184 (2013) 1522-1534

M. Albergo et al., Phys.Rev.D 104 (2021) 11,11450

R. Abbott et al., arXiv:2207.08945

### Normalising flows: Volume scaling needs to be addressed

- Parameter/function/method space is large
  - A lot of possibilities/potential : training procedure, mapping, factorizations ...
  - ... but there is the danger of the parameter/methods desert

P. Shanahan, Talk, 16.08, 10:40

## GC - steps - Status

**Techniques introduced in** J. F. et al., CPC 184 (2013) 1522-1534

### Factorisation of determinant and its computation

- Use LU until  $L=4$
- Use Stochastic estimators for  $L>4$ 
  - Only one source per ratio (need for rel. gauge fixing)

### New developments (so far not implemented):

- increase distances between active domains M. Cè et al., Phys.Rev.D 95 (2017) 3, 034503
- Use GC-steps as topological tunnelling steps and not as full MCMC method L. Guisti et al., Phys. Let. B 829 (2022) 137103

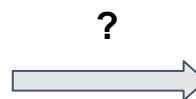
### New implementation for an efficient steps

This should/could include:

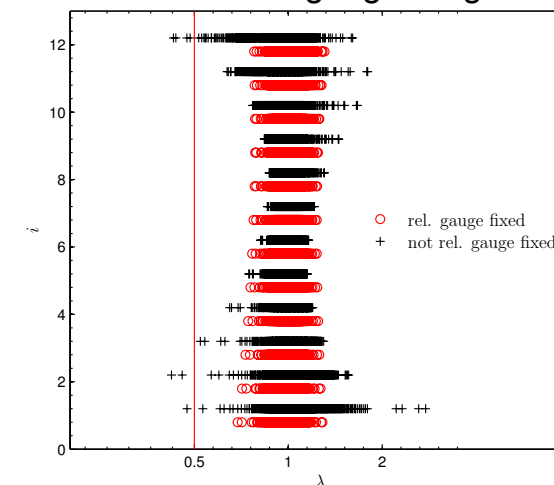
- Flexible parallelisation techniques
  - Decomposition is not equally distribute computing
  - Active domains are computational hot spots
- Modularity
  - LU-decomposition requires thick nodes
  - Sparse matrix inversions more efficient on GPUs

Included in design of lyncs

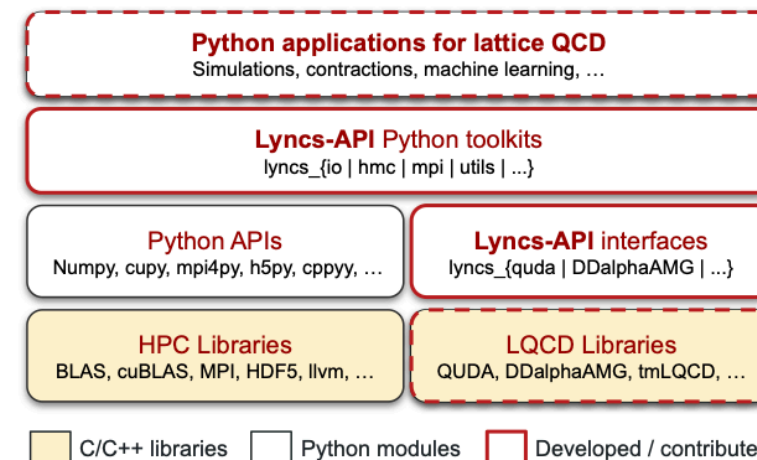
- requires lyncs-GC



Spectrum of Dirac Operator under relative gaugefixing



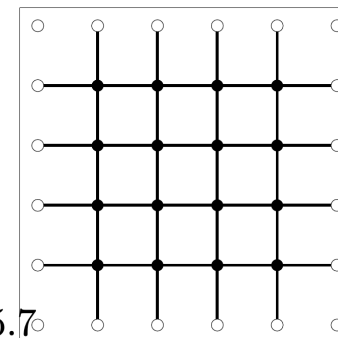
### Python ecosystem for Lattice QCD



## SU(3) - domain size

How large has to be the block ? Roughly  $L > 0.4 \text{ fm}$ , which is  $\sim 10^4$  at  $a=0.04 \text{ fm}$

- HB-Overrelaxation study seems to confirm that (here  $8^4$  within  $16^4$ )



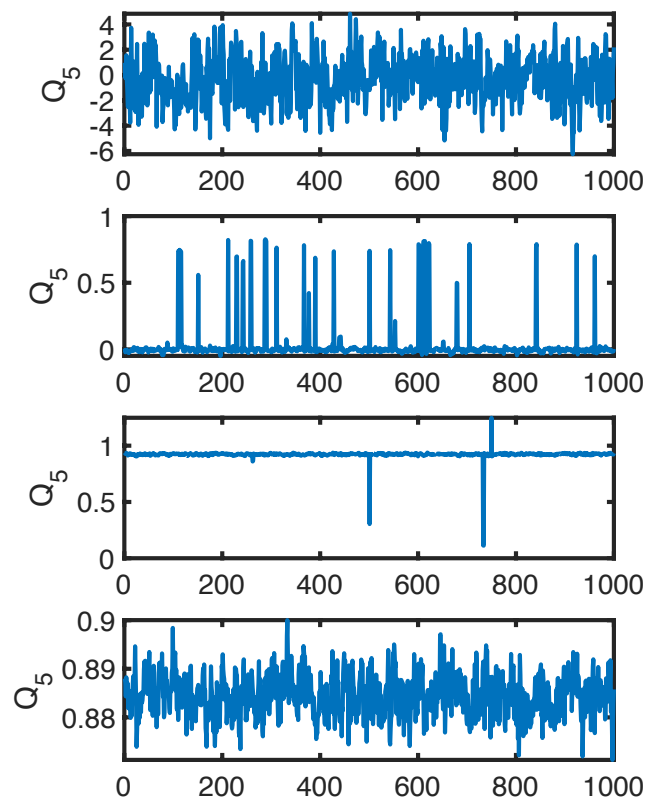
$\beta = 5.7$

$\beta = 6.0$

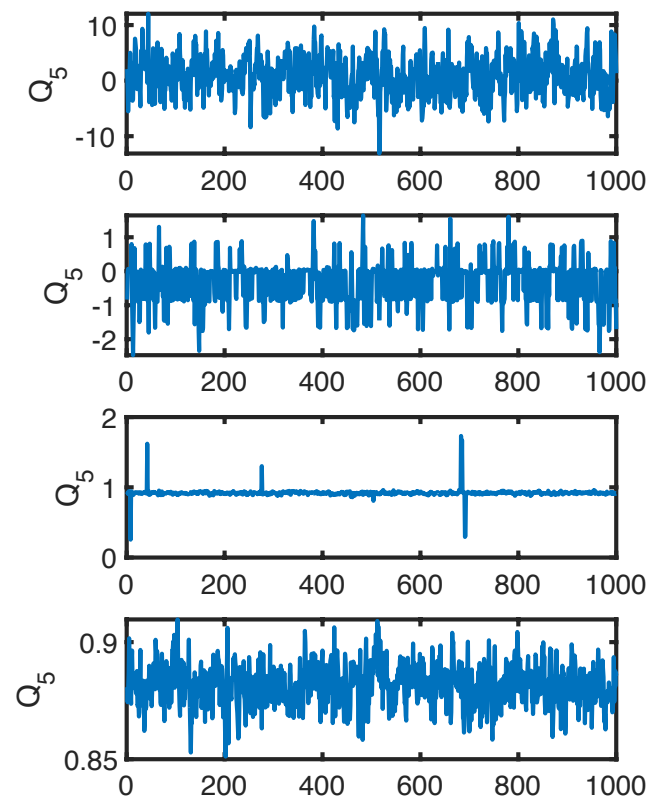
$\beta = 6.18$

$\beta = 6.4$

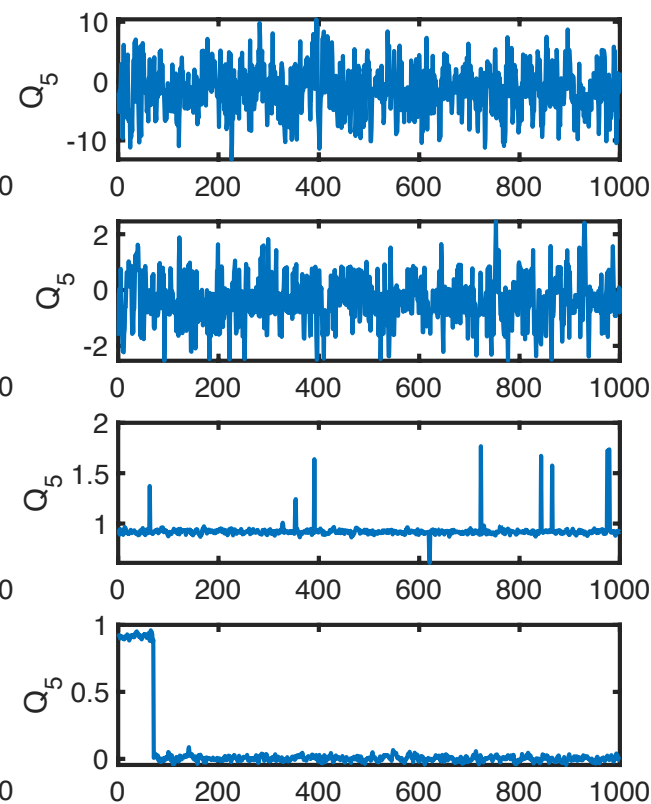
**L=6**



**L=8 - c2bn =1**



**L=8**





## SU(3) - updates

**Need for an update procedure which can (ideally guarantee) tunnelling of topology**

- Generative models
- Continuous flows
- Instanton-updates (seems not to work)
  - maybe in combination with flows
- Re-thermalization (brute force)
- Local HMC (brute force)

- **at the physical point**

- require at least 1 fm distance between active domains
  - Within a 5 fm box
    - 162 blocks of size 2.5 fm possible
    - should be okay (if acceptance rate is fine)

M. Dalla Brida et al., Phys.Lett.B 816 (2021) 136191

## How large are the costs ?

**GC:** nested accept-reject steps will scale with the most expensive step

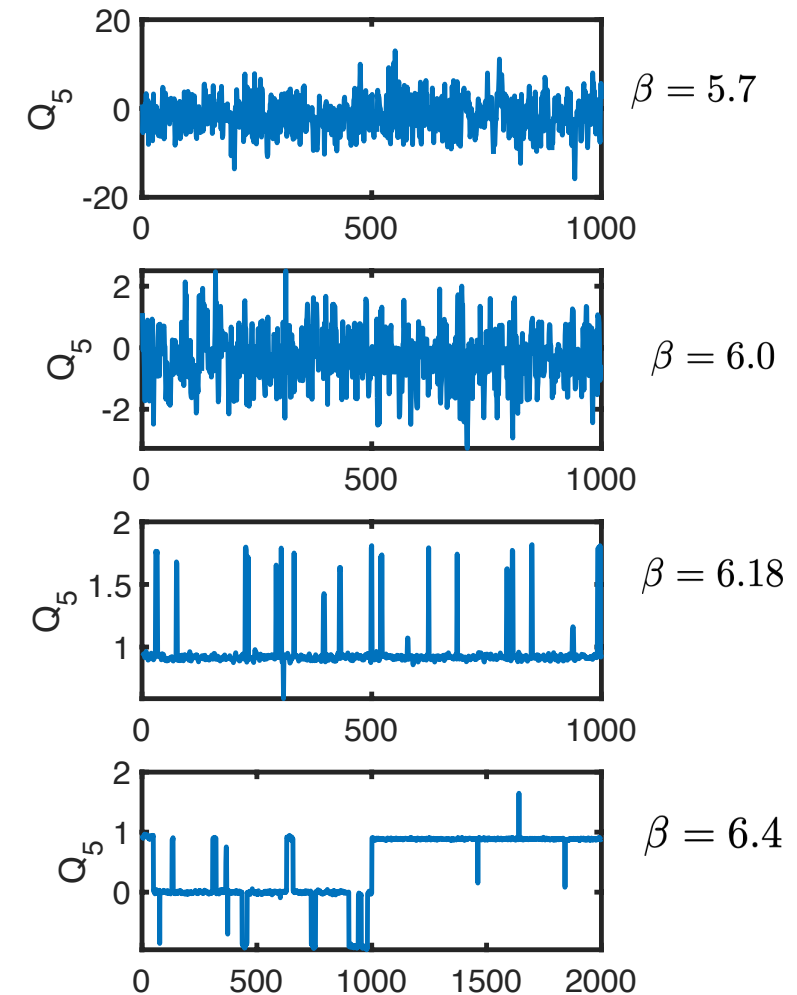
- PG step or local determinant (potential  $V^2$  scaling)

**Multi-level/HMC-updates:** of the larger domains at 0.04 fm

- Scales with  $162 \cdot 60^4$  ( $\sim 8 \times 128^4$ )

**Not clear in the moment  
which method will work**

## HB+OR re-thermalization on L=8



**Note that with computing at the exascale**

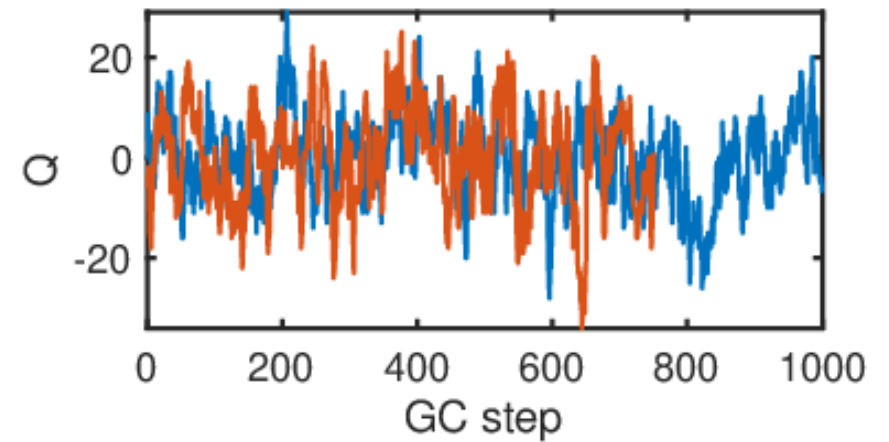
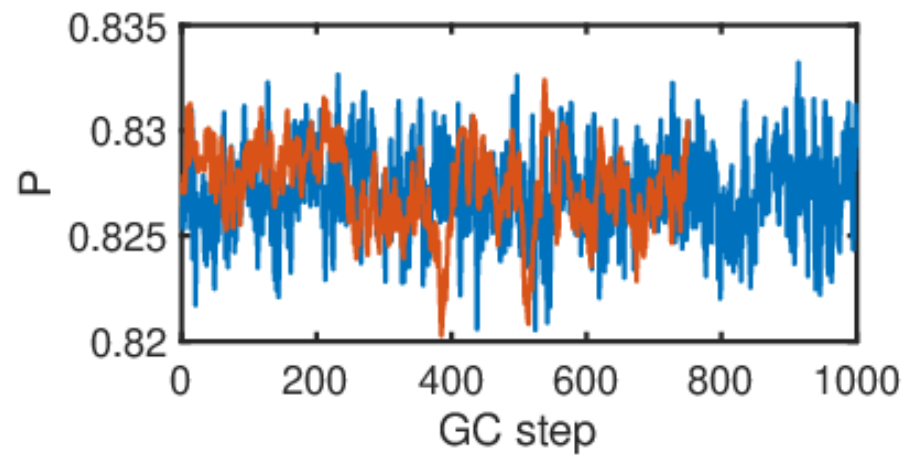
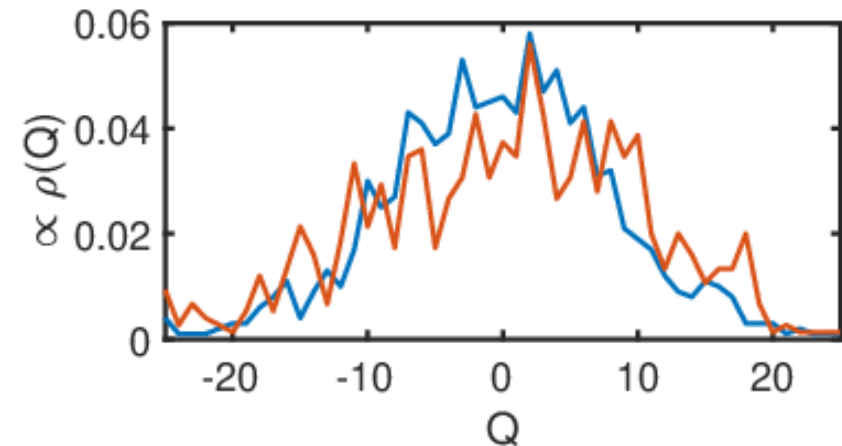
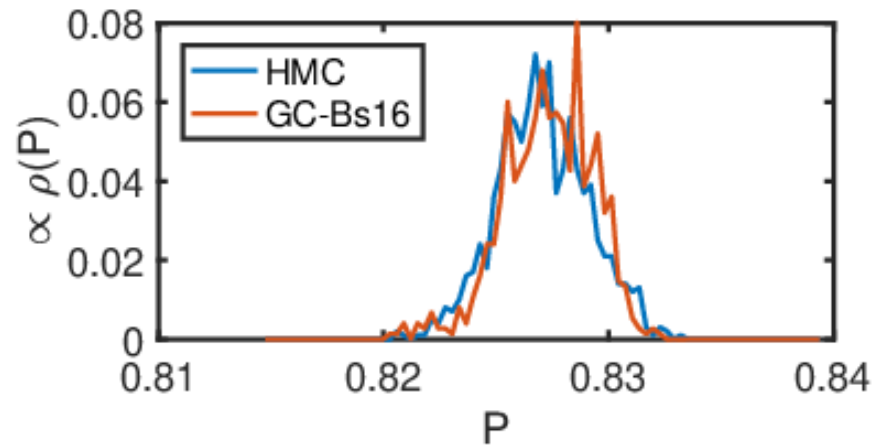
- Computational resources available to run 10k MDU for  $L=128$  at physical pion masses reaching  $a < 0.05$  fm in reach if we can mild down topological freezing  
... novel idea's and implementation are needed.

**Thank you**

## Appendix

2D Schwinger -  $\beta = 3$  -  $L = 128$ 

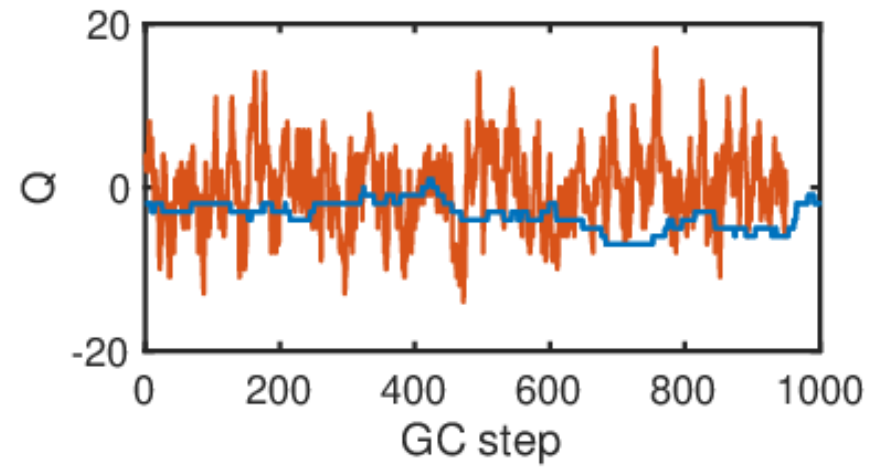
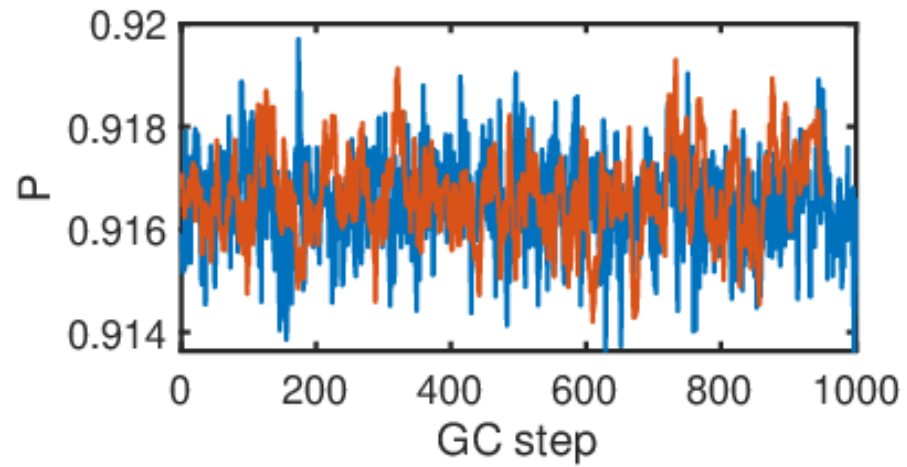
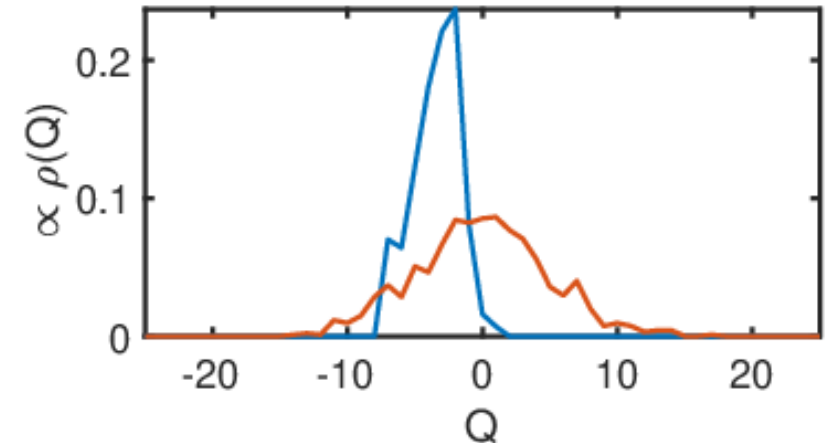
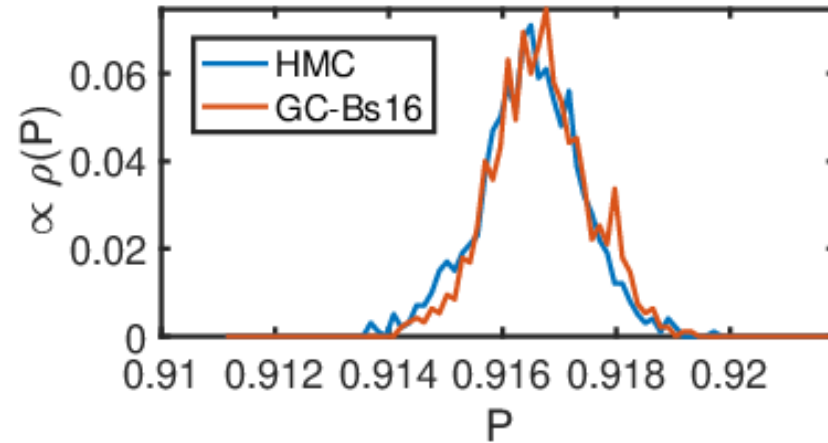
- $L = 128$
- $\beta = 3.0$
- $m = -0.082626$





2D Schwinger -  $\beta = 6$  -  $L = 128$ 

- $L = 128$
- $\beta = 6.0$
- $m = -0.0342$



2D Schwinger -  $\beta = 8.45$  -  $L = 128$ 

- $L = 128$
- $\beta = 8.45$
- $m = 0.0$

