



## **QCD** at LHC

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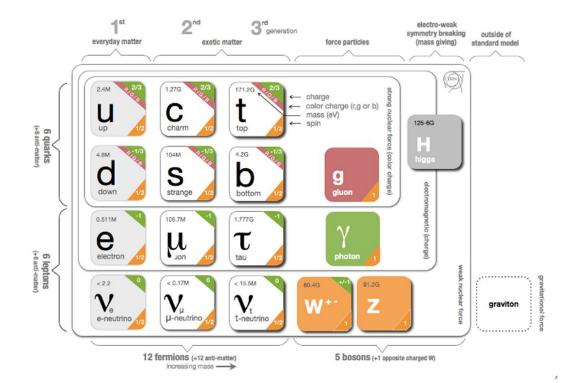
Sarajevo School of High Energy Physics October 2022



### Standard Model



➤ Theory developed in the 1960s that describes the fundamental particles and the interactions between them



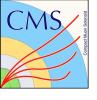


# Quantum Chromodynamics (QCD)



- The strong interaction described by the Quantum
   Chromodynamics (QCD) theory which is based on the gauge symmetry group SU(3)
- The particles that interact with the strong force (quarks and gluons) have an additional quantum number called color
- The quarks one of the three colors: red, blue or green (antiquarks - anticolors)
- The eight gluons the combination of color and anticolor

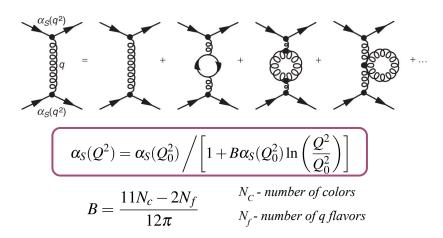




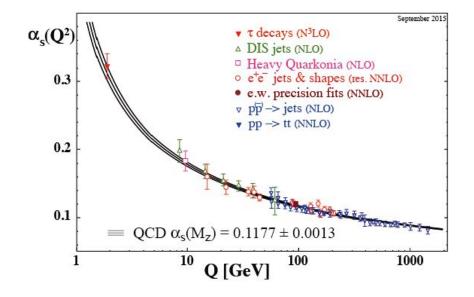
# Running of the strong coupling constant $\alpha_s$



All couplings run (QED, QCD, EW), i.e. they depend on the momentum scale (Q<sup>2</sup>)
of process



- ⇒ At high scales Q, coupling becomes small quarks and gluons are almost free, interactions are weak
- ⇒ At low scales, coupling becomes strong

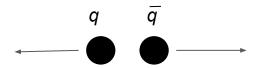


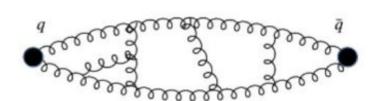


### Color confinement



### Can we detect free quarks?







- The color field between the quarks is squeezed into a tube
- The energy stored in the field is proportional the separation of the quarks



$$V(r) \sim \lambda r$$

⇒ infinite amount of energy to separate two quarks to infinity

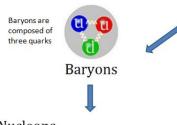


### Color confinement



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- Quark arrange themselves into bound hadronic states that are colourless combinations with no colour field between them
- Hadrons can be mesons and baryons



#### **Nucleons**

Particle	Mass $(MeV/c^2)$	τ (sec)
P	938.2	$> 10^{11}$
n	939.5	$10^{3}$

#### Hyperons

Particle	Mass (MeV/c <sup>2</sup> )	τ (sec)
Λ	1115	$2.6 \times 10^{-10}$
$\Sigma^+$	1189	$0.8 \times 10^{-10}$
$\Sigma^0$	1192	$10^{-14}$
$\Sigma^-$	1197	$1.6 \times 10^{-10}$
$\mathcal{\Xi}^0$	1314	$3 \times 10^{-10}$
$\mathcal{E}^-$	1321	$1.8 \times 10^{-10}$
$\Omega^-$	1675	$1.3\times10^{-10}$

#### Hadrons





#### Mesons



#### **Pions**

Particle	Mass $(MeV/c^2)$	τ (sec)
$\pi^-, \pi^+$	139	$2.5 \times 10^{-8}$
$\pi^0$	135	$1.8 \times 10^{-1}$

#### Kaons

Particle	Mass ( $MeV/c^2$ )	$\tau$ (sec)
$K^-, K^+$	494	$1.2 \times 10^{-}$
$K^0$	498	
η	550	$10^{-18}$











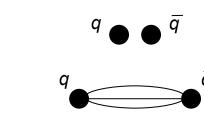


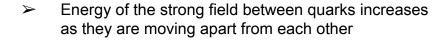


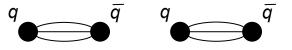
Energy of the strong field between quarks increases as they are moving apart from each other









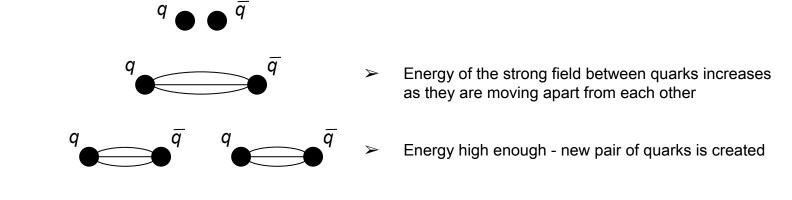


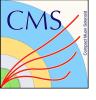
Energy high enough - new pair of quarks is created



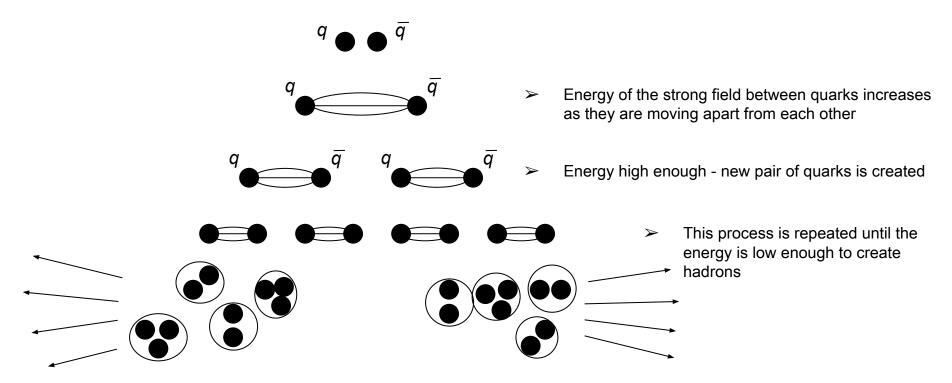


This process is repeated until...



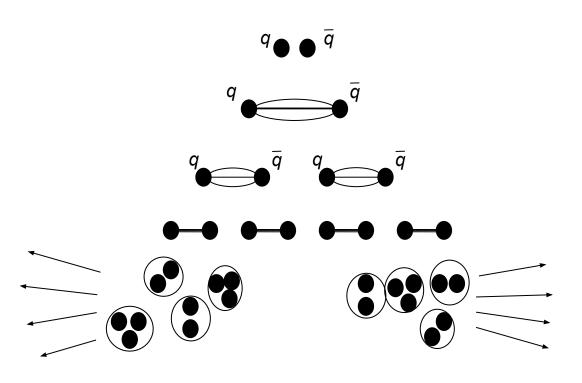




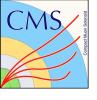








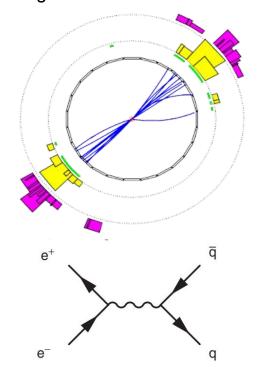
- This process is called hadronization
- The produced hadrons are often the results of boosted interactions, which makes the particles to be collinear and form what is called a jet

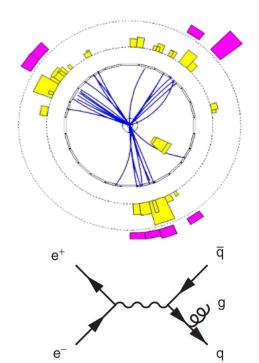


### Jet production in e<sup>+</sup>e<sup>-</sup> collisions



• Jet production in high-energy electron—positron collisions also provides direct evidence for the existence of gluons







### Jets at LHC



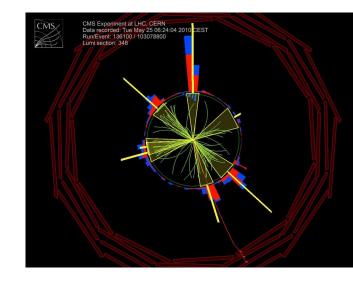
 Jets - less well defined than leptons or muons, but important for understanding LHC physics

#### Standard Model physics:

- Many standard model processes produce jets, are sensitive to  $\alpha_{_{\rm S}}$
- Multijet cross section measurements test QCD
- Hadronic decays of heavy particles

#### New physics searches:

 Many searches looking for final states with jets, or in regions of phase space with high jet multiplicities

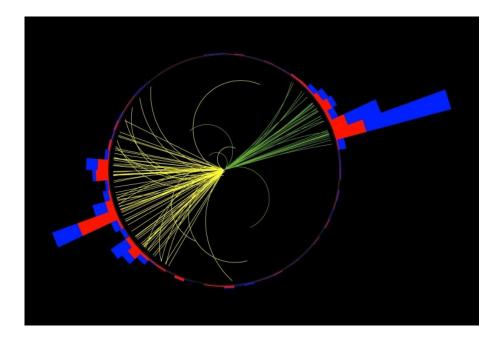




### Jets at LHC



 Jets of particles leave signals in components such as the tracker and the electromagnetic and hadronic calorimeters

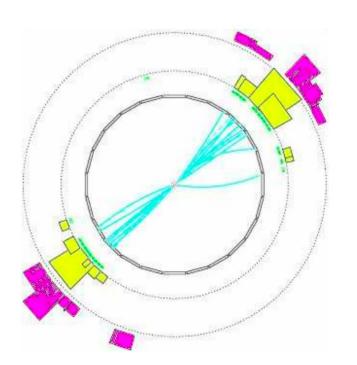


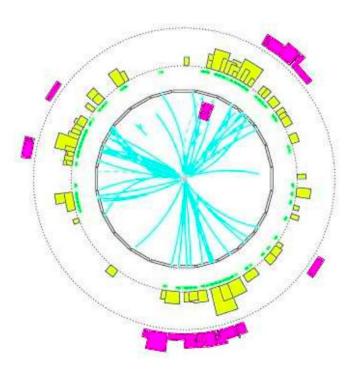


# Jet finding



### How many jets do you see?



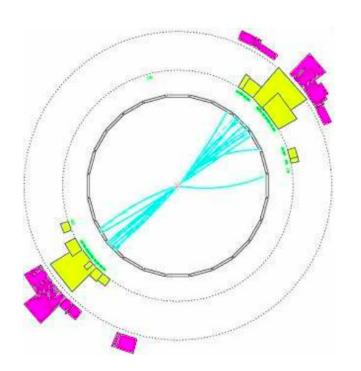


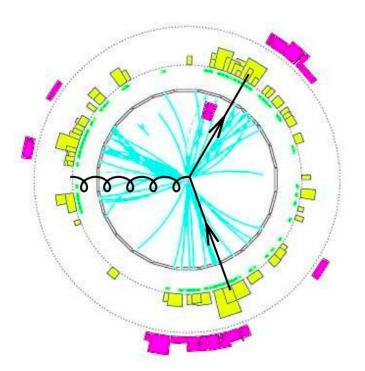


# Jet finding



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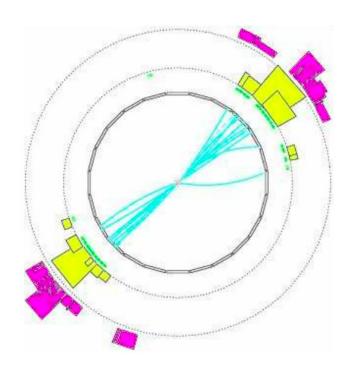


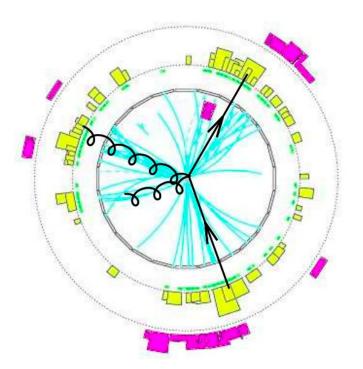


# Jet finding



### How many jets do you see?







### Jets algorithms



- To decide unambiguously whether a measured object belongs to a jet or not, a mathematical prescription is required - jet algorithm
- Algorithm has to be applicable to theoretical calculations as well as to measurements from different experiments

Jets should be invariant with respect to certain modifications of the event:

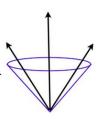
#### collinear splitting:

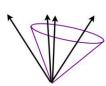
Collinear splittings should not bias jet finding

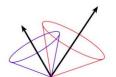
#### infrared emission :

- soft radiation should not affect jet configuration
- Only observables that are IR safe can be calculated in pQCD













## Jets algorithms



### Cone algorithms

- o particles in jets will show up in conical regions and thus they cluster based on (η-φ) space, resulting in jets with rigid circular boundaries
- select the most energetic particle as a seed with the
- o constituents within cone of radius R are considered part of the jet

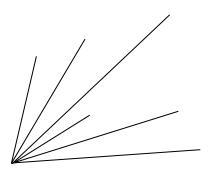
### Sequential Clustering Algorithms

 $\circ$  particles within jets will have small differences in transverse momenta and thus groups particles based on momentum space, resulting in jets that have fluctuating areas in (η-  $\phi$ ) space.





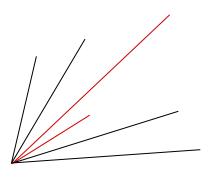
- Work their way backwards through this branching by combining pairs of particles into a single one
- A distance measure based on angular separation and energy/pT of constituent has to be determined
- The particles that are closest are combined







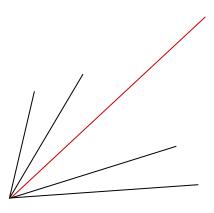
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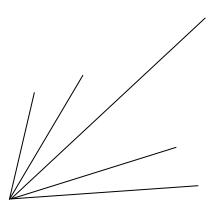
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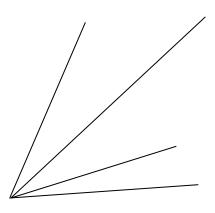
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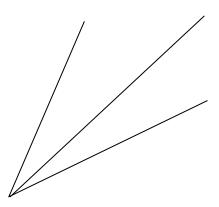
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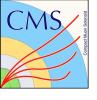






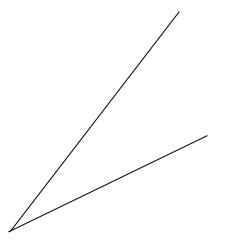
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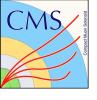


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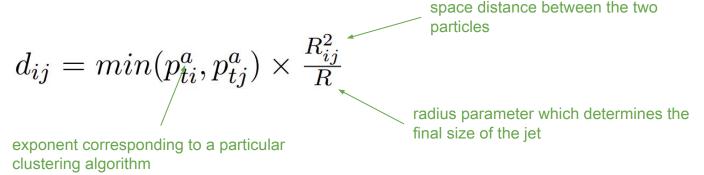
How do we decide which particles should we combine?

When do we stop combining particles?



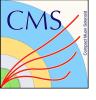


Distance between two particles:



 Distance between the beam axis and the detected particle:

$$d_{iB} = p_{ti}^a$$





Distance between two particles:

$$d_{ij} = min(p_{ti}^a, p_{tj}^a) \times \frac{R_{ij}^2}{R}$$

 Distance between the beam axis and the detected particle:

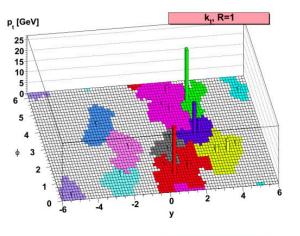
$$d_{iB} = p_{ti}^a$$

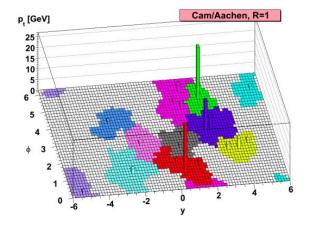
#### **Algorithm flow:**

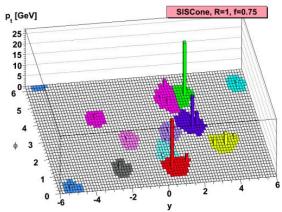
- Find the minimum distance
- If d<sub>ij</sub> is the minimum: particles i and j are combined into one particle (ij), i and j are removed from the list of particles
- If  $d_{iB}$  is the minimum: i is labelled a final jet and removed from the list of particles
- This process is repeated until:
  - o all particles are part of a jet with the distance between the jet axes  $R_{ij}$  greater than R inclusive clustering
  - desired amount of jets have been found - exclusive clustering

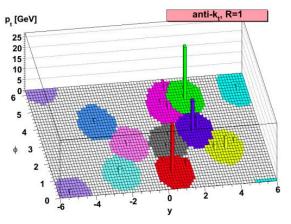












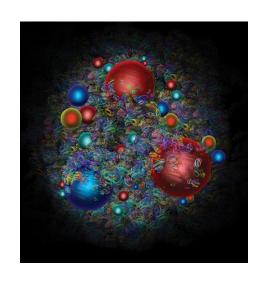


## Proton-proton collisions



- Protons are baryons composed of two u quarks and one d quark ("valence" quarks) and "sea" quarks
- Quarks and gluons within protons are referred to as partons

- Structure of proton described by the Parton model
  - $\Rightarrow$  partons carry a fraction of the total proton momentum, which gives the probability that parton has a fraction x of the total proton momentum  $P(p_i = xP)$



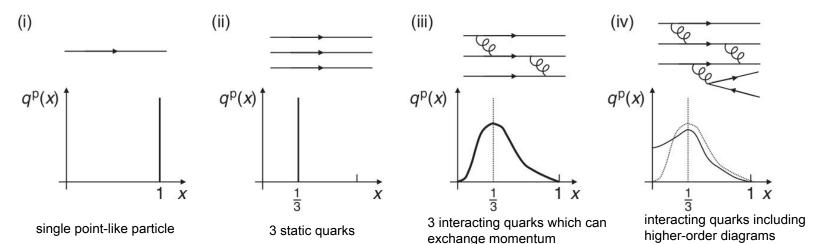


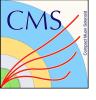
### Parton distribution function



- The quarks inside the proton will interact with each other through the exchange of gluons
- The dynamics of this interacting system will result in a distribution of quark momenta within the proton
- These distributions are expressed in terms of Parton Distribution Functions (PDFs)

#### Forms of the quark PDFs:





### Parton distribution function



- Functional forms of the PDFs not a priori known and have to be obtained from experiment
  - deep inelastic scattering (DIS) experiments such as lepton-hadron collider HERA
  - hadron colliders such as LHC
  - the fixed-target experiments
- The PDFs depend on the scale at which the hadron is probed
- Evolution of PDFs with the scale  $\mu_F$  described by the **Dokshitzer-Gribov-Lipatov- Altarelli-Parisi** (**DGLAP**) equations:

$$\mu_F \frac{\delta f_a(x, \mu_F)}{\delta \mu_F^2} = \sum \frac{\alpha_s(\mu_F^2)}{2\pi} \int_x^1 \frac{d\xi}{\xi} P_{a \to bc}(\frac{x}{\xi}, \mu_F) f_a(\xi, \mu_F)$$

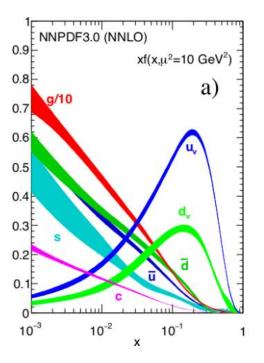
- $P_{a o bc}$  Altarelli-Parisi splitting function, that gives the probability for a parton a to split into two partons bc.
- $\xi$  resulting particle momentum fraction of the quark with momentum  $p_a$

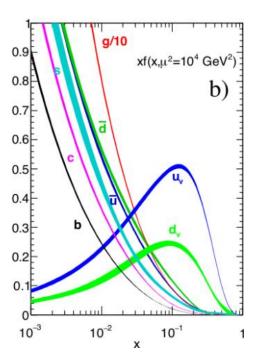


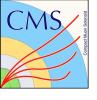
### Parton distribution function



 Because of the universality of the PDFs, it is possible to use PDFs extracted from well-known processes to obtain predictions corresponding to different scales or different final states







### Cross section



- > The **cross section** is a measure of quantum mechanical probability for the interaction
- > It depends on the fundamental physics involved in the Feynman diagrams that contribute to the process
- The cross section of a proton-proton interaction cannot be computed easily, due to the complex structure of protons

$$\sigma_{pp\to X} = \sum_{a,b} \int \int dx_a dx_b f_a(x_a, \mu_F) f_b(x_b, \mu_F) \underline{\sigma_{ab\to X}(x_a x_b, \mu_R, \mu_F)}$$

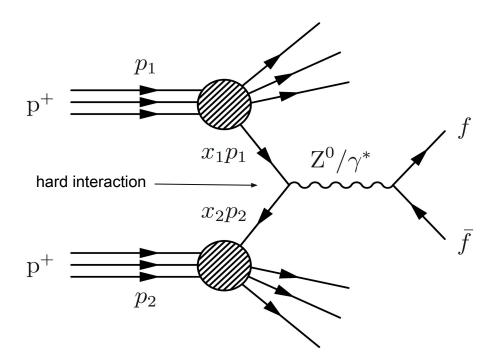
cross section at the partonic level.



# Drell-Yan process



Annihilation of quark-antiquark pairs from hadrons with the creation of a Z boson or a virtual photon, which decays into a lepton-antilepton pair

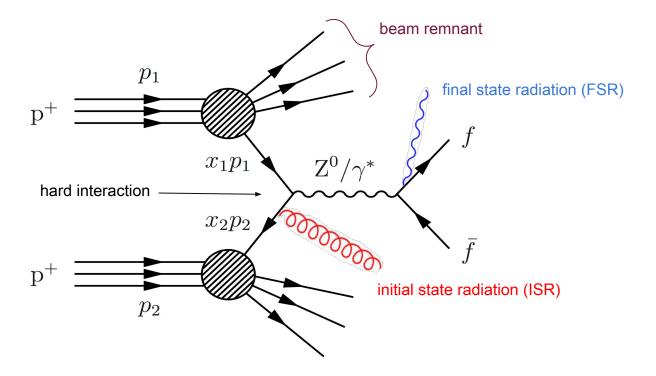




## Drell-Yan process



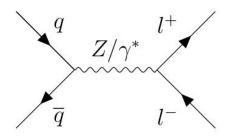
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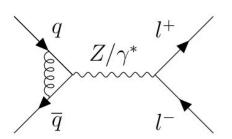


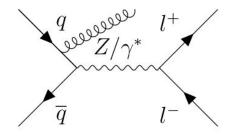


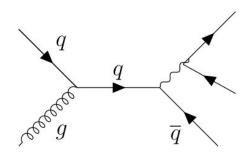
# Drell-Yan process higher orders













## Drell-Yan process cross section



The cross section of the Drell-Yan process:

$$\sigma(pp \to l^+ l^-) = \sum_q \int \int dx_q dx_{\overline{q}} f_q(x_q, \mu_F^2) f_{\overline{q}}(x_{\overline{q}}, \mu_F^2) \sigma(q\overline{q} \to l^+ l^-)$$

• Following the perturbative QCD, the partonic cross section can be expanded in series with respect to the coupling constant  $\alpha_s$ :

$$\sigma(q\overline{q} \to l^+l^-) = \sigma_{LO} + \alpha_s \sigma_{NLO} + \dots$$

• The partonic cross section can be calculated using the Matrix Element of the Feynman diagram



## Theoretical predictions



 To establish connection between theory and experiment it is necessary to simulate the proton-proton interactions

- The evolution of an event in simulation starts with the two beam particles that are colliding
- The partons from beams start irradiating initial state shower is simulated
- The incoming partons enter the hard interaction and the outgoing particles are produced
- In the hard process, **short-lived resonances** can be created and their decay is considered in this step
- The outgoing particles undergo radiation final state radiation is simulated.
- The simulation of underlying events.
- The process of hadronization is simulated.
- The decay of long-life particles, such as τ leptons or B-hadrons.

