What is the $i\varepsilon$ for the S-matrix?

Sebastian Mizera (IAS Princeton)

based on **hep-th/2204.02988**
with Hofie Sigríðar Hannesdóttir

There's a long history in understanding imprints of *causality* on scattering amplitudes

(microcausality, macrocausality, Bogoliubov causality, no Shapiro time advances, …)

[Bogoliubov, Schutzer, Tiomno, van Kampen, Gell-Mann, Goldberger, Thirring, Wanders, Iagolnitzer, Eden, Landshoff, Peres, Branson, Omnes, Chandler, Pham, Stapp, Adams, Arkani-Hamed, Dubovsky, Grinstein, O'Connell, Wise, Giddings, Porto, Camanho, Edelstein, Maldacena, Zhiboedov, Tomboulis, Minwalla, Caron-Huot, Mazac, Rastelli, Simmons-Duffin, Riva, Capatti, Hirschi, Kermanschah, Pelloni, Ruijl, …]

Although never made precise, it is generally believed that causality is encoded in *complex-analytic* properties of the S-matrix

Over the years, complexification become bread and butter in scattering amplitude computations. Multiple practical reasons:

- Theory of complex angular momenta, dispersion relations, on-shell recursion relations, …

- Crossing symmetry

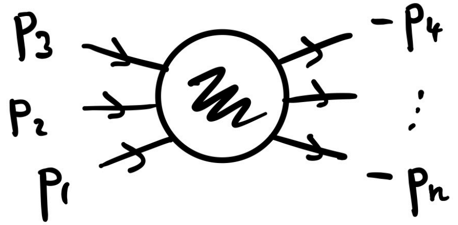
$$\underbrace{e^+e^- \to \gamma\gamma}_{s>0} \qquad \underbrace{\gamma e^- \to \gamma e^-}_{s<0}$$

can we get it "for free"?

Before we start… conventions

- S-matrix operator: $S = \mathbb{1} + iT$

$$\mathbf{T} = \langle \text{out} | T | \text{in} \rangle =$$



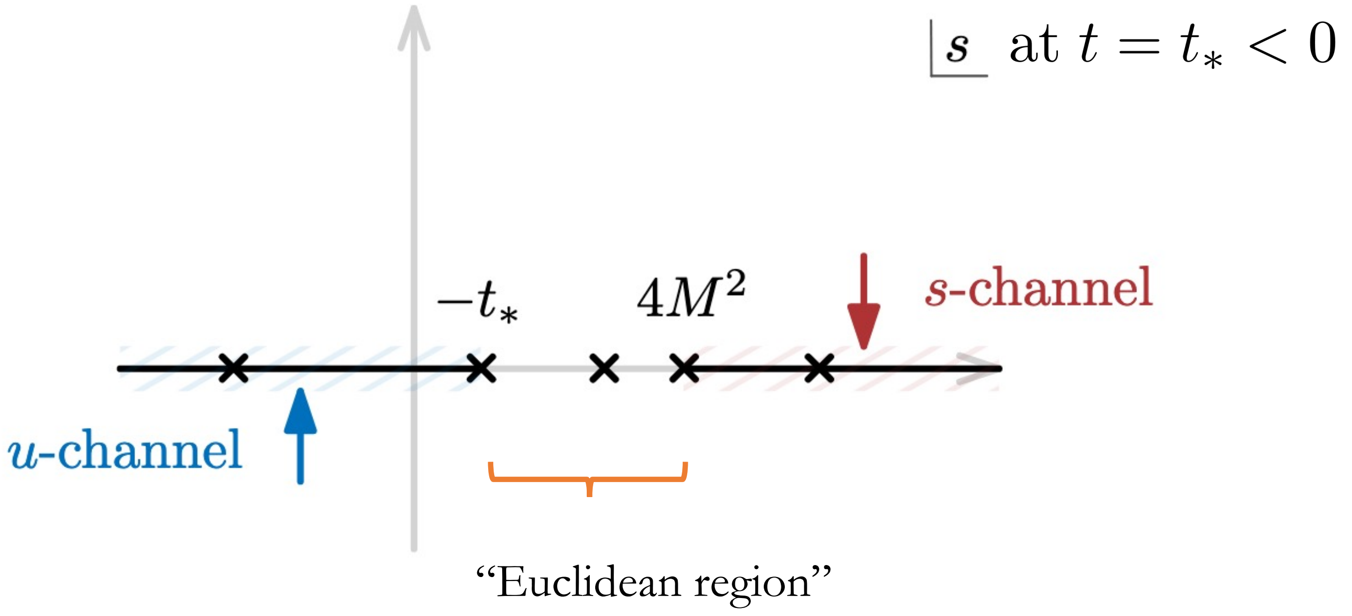
- For $2 \to 2$ scattering

$$s = (p_1 + p_2)^2 \qquad t = (p_2 + p_3)^2 \qquad u = (p_1 + p_3)^2$$

subject to momentum conservation

$$s + t + u = \sum_{i=1}^{4} M_i^2$$

Analyticity is best understood for $2 \to 2$ scattering of the lightest state in theories with a mass gap M for low momentum transfer:



$$\mathbf{T}(s, t_*) = \lim_{\varepsilon \to 0^+} \mathbf{T}_{\mathbb{C}}(s + i\varepsilon, t_*)$$

causality?

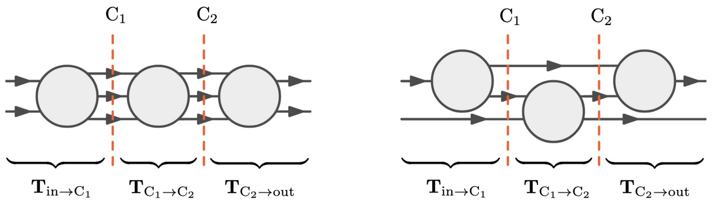
How does this picture extend to more realistic processes?

Standard Model $\supset$ massless and unstable particles, UV/IR divergences, higher-point processes

Outline

- Unitarity constraints

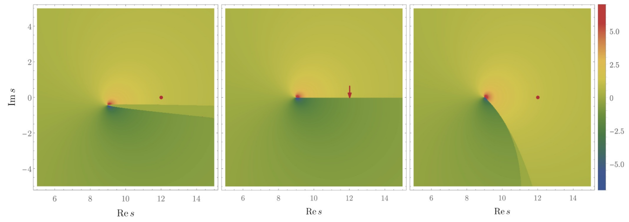
Holomorphic cutting rules



Discontinuities beyond normal thresholds

- Causality constraints

Different ways of implementing causality



Deforming branch cuts in the kinematic space

Unitarity constraints

Unitarity, $SS^\dagger = \mathbb{1}$, implies that

Sum-integral over all the intermediate states

$$\frac{1}{2i}(T - T^\dagger) = \frac{1}{2} T T^\dagger$$

RHS is non-holomorphic and doesn't manifest all singularities

Eliminate $T^\dagger = T(\mathbb{1} + iT)^{-1}$ and expand the geometric series

[related work Coster, Stapp, Blazek, Matak]

This results in *holomorphic cutting rules*

positivity not manifest

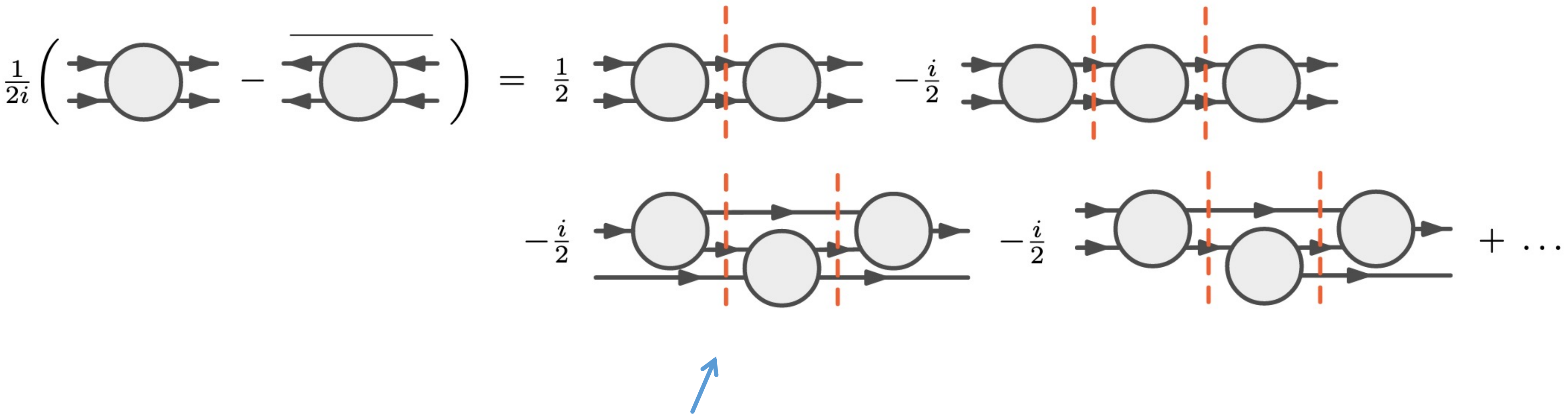
$$\frac{1}{2i}(T - T^\dagger) = -\frac{1}{2}\sum_{c=1}^{\infty}(-iT)^{c+1}$$

number of unitarity cuts

- The place where a new term on the RHS starts contributing is called a *threshold:* a potentially violent event that could give rise to *singularities* or *branch cuts*
- The phase-space is so small, it only allows for *classical* scattering

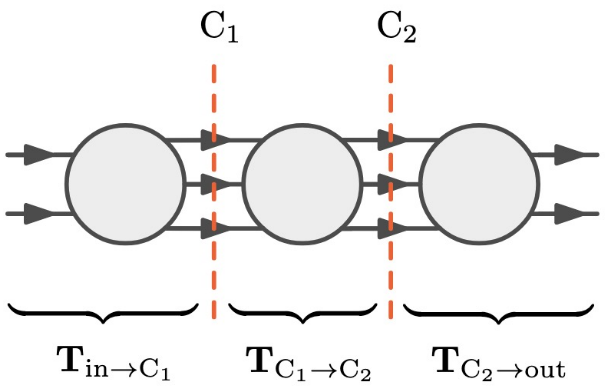
[Coleman, Norton]

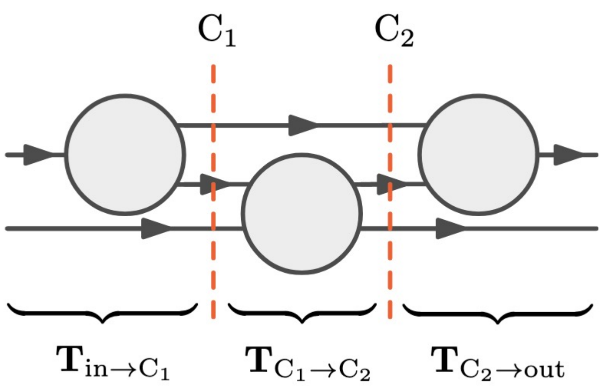
Diagrammatically



Putting propagators on shell: $\dfrac{-1}{q^2 - m^2 + i\varepsilon} \to 2\pi\delta^+(q^2 - m^2)$

There are two types of thresholds on the RHS:





Normal thresholds

(purely temporal)

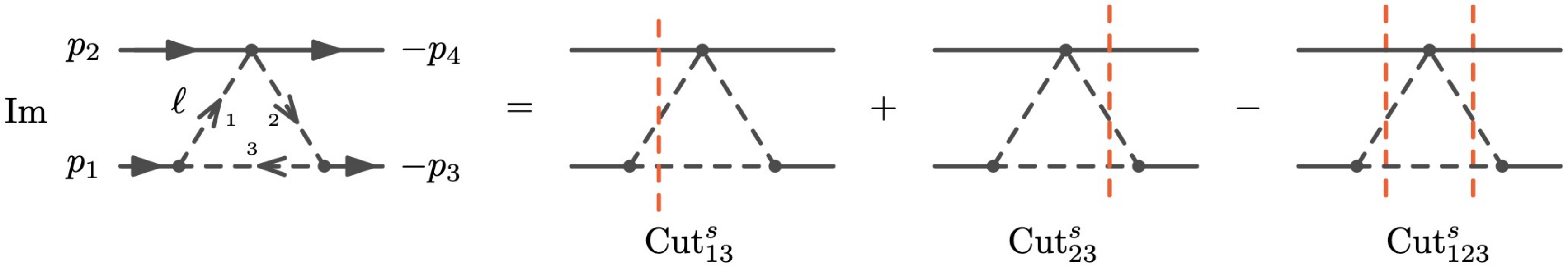
"Anomalous" thresholds
or Landau singularities

(spatially spread out)

$\mathbf{T}_{\mathbb{C}}$ doesn't have any new types of singularities in perturbation theory

[Bjorken, Landau, Nakanishi]

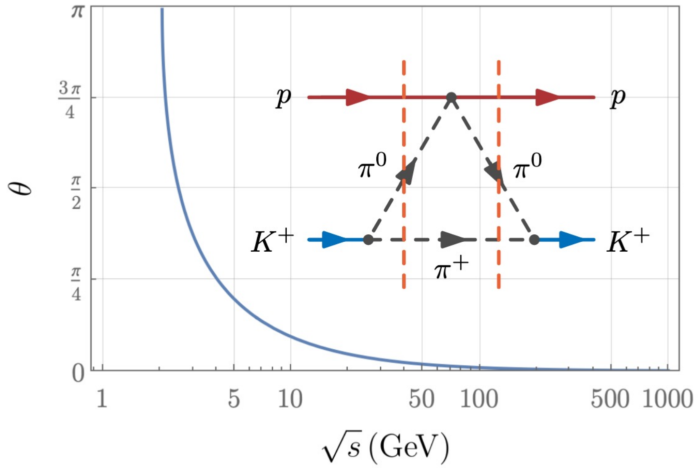
Simplest example

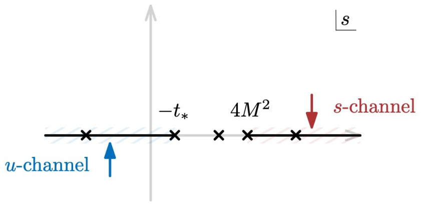


When can we build a triangle diagram with 3 momenta?

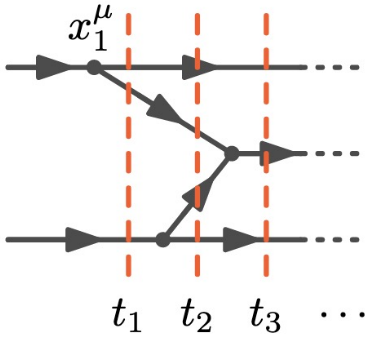
$$\cos\theta = 1 - \frac{2s\left(m_{K^+}^2 - (m_{\pi^0} + m_{\pi^+})^2\right)\left(m_{K^+}^2 - (m_{\pi^0} - m_{\pi^+})^2\right)}{m_{\pi^+}^2\left(s - (m_{K^+} + m_p)^2\right)\left(s - (m_{K^+} - m_p)^2\right)}$$

- Heavily suppressed compared to tree-level processes
- Widths move the peak to a complex plane: Breit-Wigner-like distribution



How is this consistent with  ?

At a threshold, we can time order the interaction vertices:



If all external particles are stable, must have *at least* 2 incoming particles at x_1
For $2 \to 2$ this implies only normal thresholds for physical kinematics

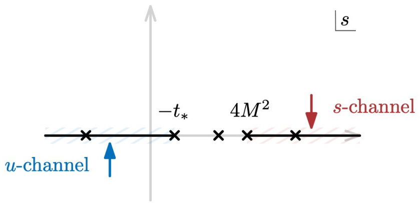
$$s = (m_1 + m_2 + \ldots)^2$$

We need to worry about anomalous thresholds for

- Higher-point scattering
- $2 \to 2$ processes with unstable particles
- Discontinuities of amplitudes
- Landau singularities in analytic expressions

Recent pheno-oriented work includes hadron spectroscopy, $\mathrm{b\bar{b}H}$ production, $\mathrm{ZZ} \to \mathrm{ZZ}$ scattering, …

[Liu, Oka, Zhao, Meissner, Guo, Denner, Dittmaier, Hahn, Boudjema, Ninh, Passarino, …]

It is reasonable to ask how much of the 

intuition survives

In particular, in the absence of the Euclidean region:

- Can we always uplift the S-matrix to a complex-analytic function in a way consistent with causality?

$$\mathbf{T}(s,t_*) \stackrel{?}{=} \lim_{\varepsilon\to 0^+} \mathbf{T}_{\mathbb{C}}(s+i\varepsilon, t_*)$$

- Is the imaginary (absorptive) part

$$\operatorname{Im}\mathbf{T}(s,t_*) = \tfrac{1}{2i}\Big(\mathbf{T}(s,t_*) - \overline{\mathbf{T}(s,t_*)}\Big)$$

always equal to the discontinuity

$$\operatorname{Disc}_s \mathbf{T}_{\mathbb{C}}(s,t_*) = \lim_{\varepsilon\to 0^+} \tfrac{1}{2i}\Big(\mathbf{T}_{\mathbb{C}}(s+i\varepsilon,t_*) - \mathbf{T}_{\mathbb{C}}(s-i\varepsilon,t_*)\Big) \quad ?$$

Where do we even start?

Convert into algebraic problems for every Feynman diagram:

We'll explain these conditions on the next slides

$$
\begin{aligned}
\mathcal{V} &= 0 \quad \text{for any } \alpha\text{'s} \quad &\Leftrightarrow \quad \text{branch cut} \\
\partial_{\alpha_e}\mathcal{V} &= 0 \quad \text{for any } \alpha\text{'s} \quad &\Leftrightarrow \quad \text{branch point} \\
\operatorname{Im}\mathcal{V} &> 0 \quad \text{for all } \alpha\text{'s} \quad &\Leftrightarrow \quad \text{causal branch}
\end{aligned}
$$

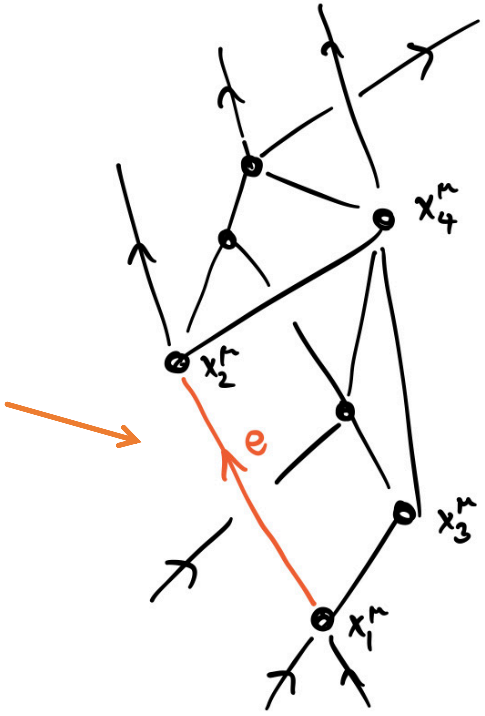
We already know *branch points* are classical scattering configurations:

Momentum q_e^μ

Mass m_e

Schwinger proper time $\alpha_e \geqslant 0$

Space-time displacement $\Delta x_e^\mu = \alpha_e q_e^\mu$



Momentum conservation at every vertex:

$$\sum_{e \ni v} q_e^\mu + \sum_{i \ni v} p_i^\mu = 0$$

Local interactions at vertices:

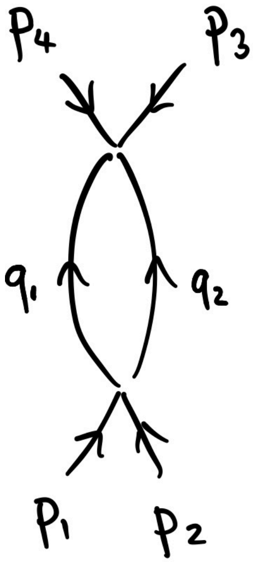
$$x_j^\mu - x_i^\mu = \sum_{e: i \to j} \Delta x_e^\mu$$

On-shell conditions for every edge:

$$q_e^2 - m_e^2 = 0$$

Landau equations [Bjorken, Landau, Nakanishi]

Bubble diagram



$s = (p_1 + p_2)^2$

momentum conservation

$$p_1^\mu + p_2^\mu = q_1^\mu + q_2^\mu = -p_3^\mu - p_4^\mu$$

locality

$$\alpha_1 q_1^\mu = \alpha_2 q_2^\mu$$

on-shellness

$$q_1^2 - m_1^2 = 0$$
$$q_2^2 - m_2^2 = 0$$

$$q_1^\mu = \ell^\mu, \qquad q_2^\mu = p_1^\mu + p_2^\mu - \ell^\mu$$

$$\ell^\mu = (p_1^\mu + p_2^\mu)\frac{\alpha_2}{\alpha_1+\alpha_2}$$

$$s\left(\frac{\alpha_2}{\alpha_1+\alpha_2}\right)^2 - m_1^2 = 0, \qquad s\left(\frac{\alpha_1}{\alpha_1+\alpha_2}\right)^2 - m_2^2 = 0$$

Lorentz invariant

Can be concisely summarized as:

$$\begin{cases} s\left(\frac{\alpha_2}{\alpha_1+\alpha_2}\right)^2 - m_1^2 = 0, \\ s\left(\frac{\alpha_1}{\alpha_1+\alpha_2}\right)^2 - m_2^2 = 0 \end{cases} \iff \begin{cases} \mathcal{V} = s\frac{\alpha_1\alpha_2}{\alpha_1+\alpha_2} - m_1^2\alpha_1 - m_2^2\alpha_2 \\ \partial_{\alpha_1}\mathcal{V} = 0 \\ \partial_{\alpha_2}\mathcal{V} = 0 \end{cases}$$

The solutions are

$$(\alpha_1 : \alpha_2) = \left(\frac{1}{m_1} : \pm\frac{1}{m_2}\right) \qquad s = (m_1 \pm m_2)^2$$

Projective invariance in Schwinger parameters and kinematic variables separately

$+$ normal threshold

$-$ pseudo-normal threshold

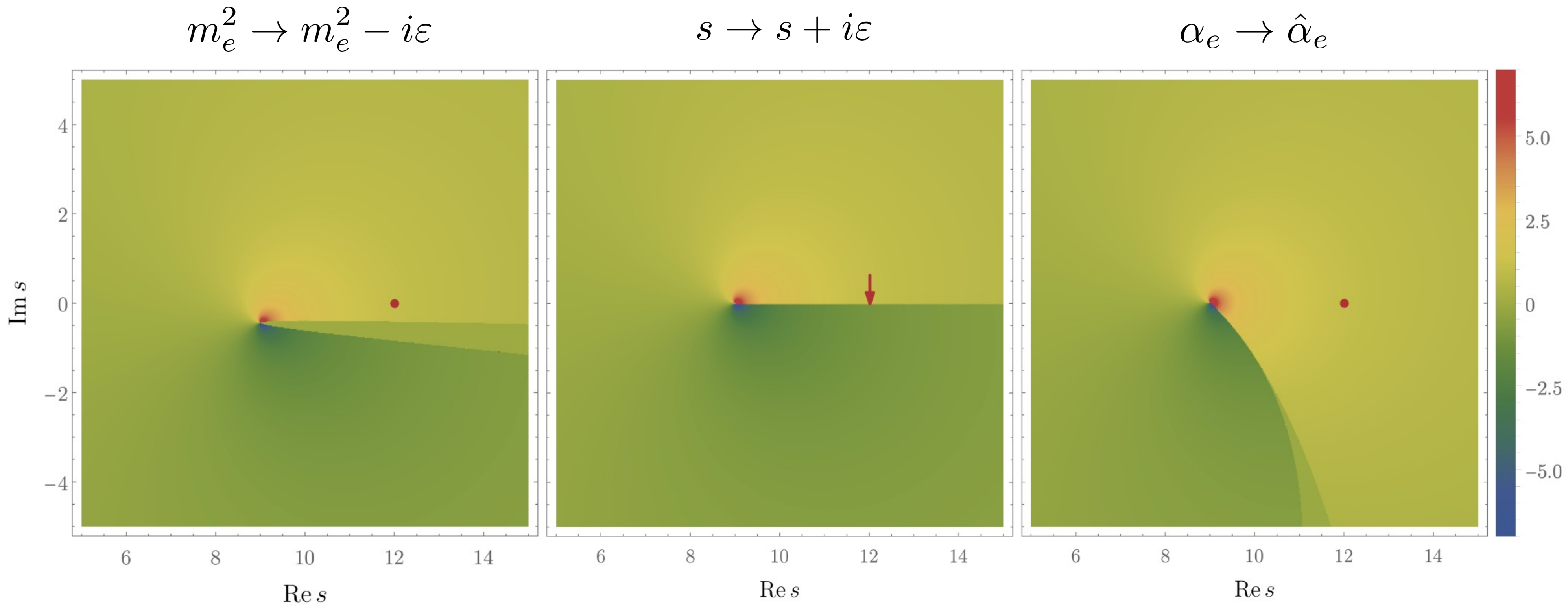
In practice, Schwinger parametrization of the bubble diagram gives:

$$\int_0^\infty \frac{\mathrm{d}\alpha_1\,\mathrm{d}\alpha_2}{\mathcal{V}^{2-\mathrm{D}/2}}\delta(\alpha_1+\alpha_2-1) \qquad \text{with} \qquad \mathcal{V} = s\frac{\alpha_1\alpha_2}{\alpha_1+\alpha_2} - m_1^2\alpha_1 - m_2^2\alpha_2$$

- When $\mathcal{V}=0$, we have to make a decision how to deform away from it (branch cut)

- Causal branch determined by $\operatorname{Im}\mathcal{V} > 0$

There are three options for implementing $\mathrm{Im}\,\mathcal{V} > 0$ ($\mathcal{V} = s\frac{\alpha_1\alpha_2}{\alpha_1+\alpha_2} - m_1^2\alpha_1 - m_2^2\alpha_2$):



Feynman $i\varepsilon$
(moves branch points, unphysical)

Kinematic $s+i\varepsilon$
(doesn't work in general)

Branch cut deformations
(next slides)

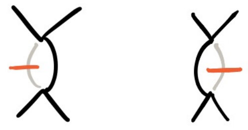
This structure is not a coincidence!
For any Feynman diagram we can define the *worldline action*

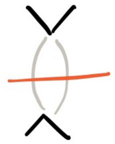
$$\mathcal{V}(\alpha_e; s_{ij}, m_e) = \frac{\mathcal{F}}{\mathcal{U}},$$

where the two *Symanzik polynomials* are given by

for the experts

$$\mathcal{U} = \sum_{\substack{\text{spanning} \\ \text{trees } T}} \prod_{e \notin T} \alpha_e, \qquad \mathcal{F} = \sum_{\substack{\text{spanning} \\ \text{2-trees } T_L \sqcup T_R}} p_L^2 \prod_{e \notin T_L, T_R} \alpha_e - \mathcal{U} \sum_{e=1}^{\mathrm{E}} m_e^2 \, \alpha_e$$

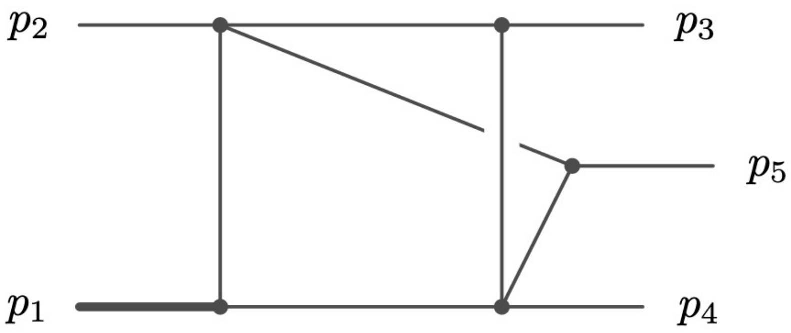




$$\Rightarrow \; \mathcal{V} = s\frac{\alpha_1\alpha_2}{\alpha_1+\alpha_2} - m_1^2\alpha_1 - m_2^2\alpha_2$$

Nowadays we have powerful algebraic geometry tools to address such questions

[SM, Telen '21]



$$\Sigma_5 = (s_{12}s_{15} - s_{12}s_{23} - s_{15}s_{45} + s_{34}s_{45} + s_{23}s_{34})^2 - 4s_{23}s_{34}s_{45}(s_{34} - s_{12} - s_{15})$$

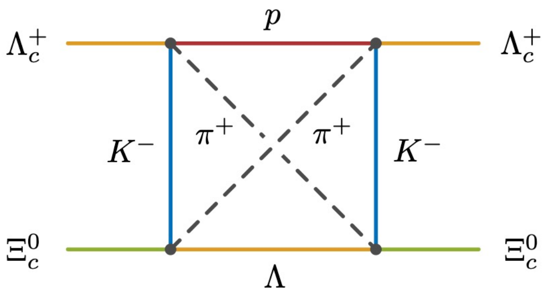
Two-Loop Hexa-Box Integrals for Non-Planar Five-Point One-Mass Processes

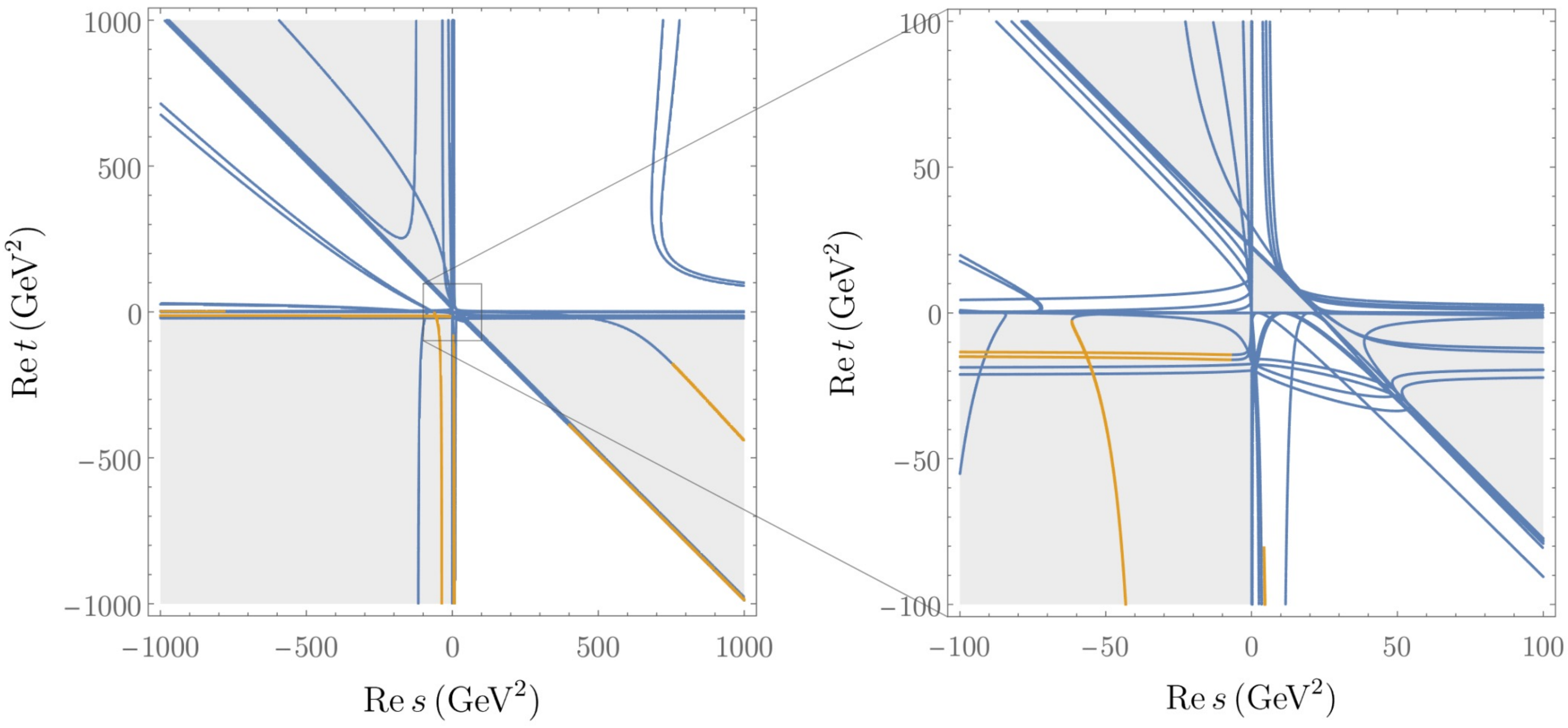
Samuel Abreu,[1,2,3] **Harald Ita,**[4] **Ben Page,**[1] **Wladimir Tschernow**[4]

[1] *Theoretical Physics Department, CERN, Geneva, Switzerland*
[2] *Mani L. Bhaumik Institute for Theoretical Physics, Department of Physics and Astronomy, UCLA, Los Angeles, CA 90095, USA*
[3] *Higgs Centre for Theoretical Physics, School of Physics and Astronomy, The University of Edinburgh, Edinburgh EH9 3FD, Scotland, UK*
[4] *Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, D-79104 Freiburg, Germany*

[hep-ph/2107.14180]

They very quickly get out of hand, e.g.,





Summary thus far

Worldline action

$$\mathcal{V} = 0 \quad \text{for any } \alpha\text{'s} \quad \Leftrightarrow \quad \text{branch cut}$$

Extremizing gives a classical saddle point

$$\partial_{\alpha_e}\mathcal{V} = 0 \quad \text{for any } \alpha\text{'s} \quad \Leftrightarrow \quad \text{branch point}$$

$$\operatorname{Im}\mathcal{V} > 0 \quad \text{for all } \alpha\text{'s} \quad \Leftrightarrow \quad \text{causal branch}$$

How to implement consistently?

Analytic properties can be determined without explicit computations

Causality: giving worldlines a small phase

$$\alpha_e \to \alpha_e \exp(i\varepsilon \partial_{\alpha_e} \mathcal{V})$$
$$= \alpha_e \left(1 + i\varepsilon \partial_{\alpha_e} \mathcal{V} + \ldots\right)$$

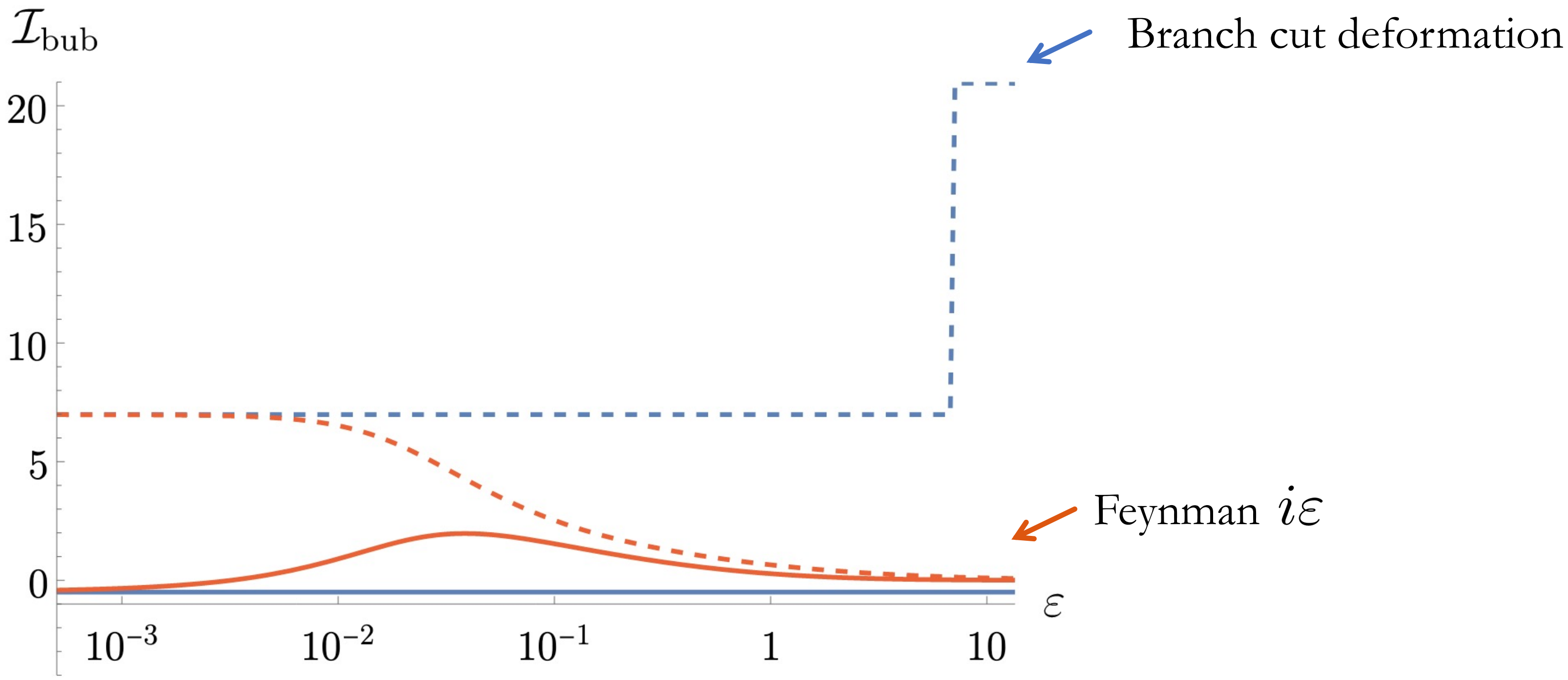
[related work Chandler, Nagy, Soper, SM, …]

At the level of the action:

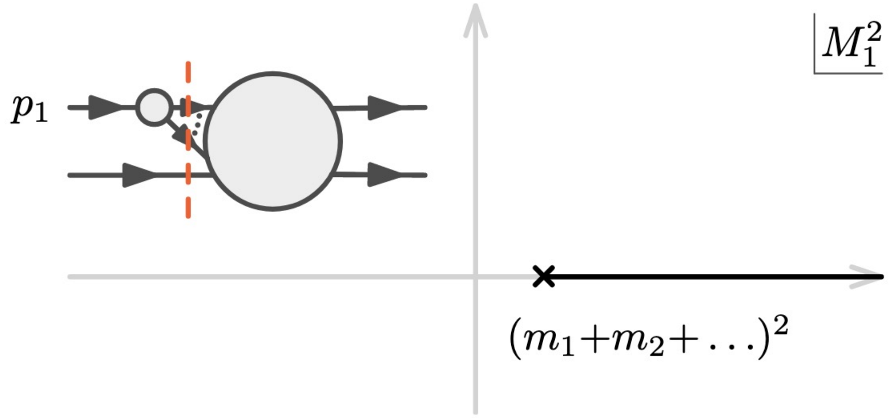
$$\mathcal{V} \to \mathcal{V} + i\varepsilon \underbrace{\sum_{e=1}^{\mathrm{E}} \alpha_e (\partial_{\alpha_e} \mathcal{V})^2}_{\geqslant 0} + \ldots$$

← Breaks down directly at the branch points

In practice, we only need a *sufficiently small* ε
(as opposed to infinitesimal)



Why couldn't we just use $s + i\varepsilon$? First sign of problems:



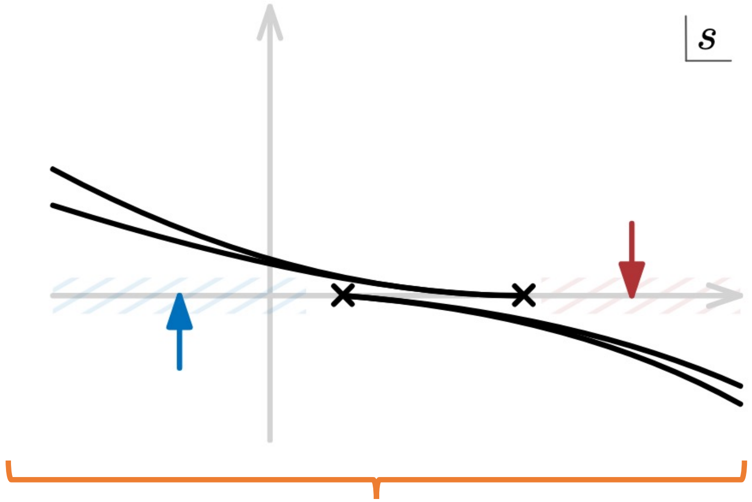
Off-shell:
branch cut between
$M_1^2 \pm i\varepsilon$

By momentum conservation

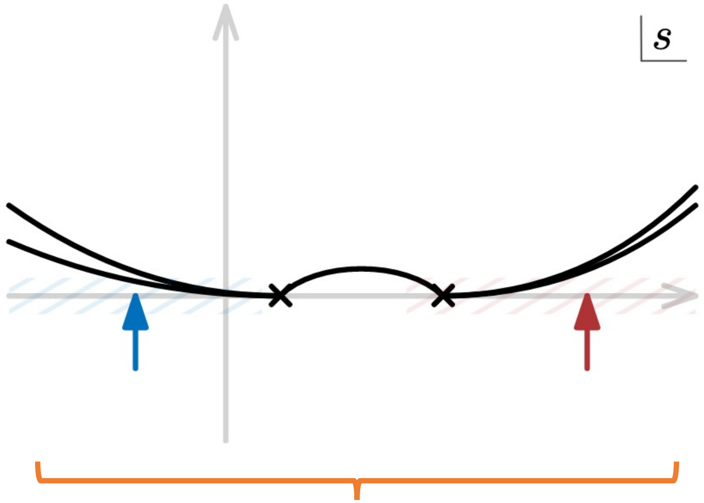
$$(s \mp i\varepsilon) + t + u = \sum_{i=1}^{4} M_i^2$$

On-shell:
branch cut between
$s \mp i\varepsilon$

Once we encounter a branch cuts for all s, there are two possibilities:



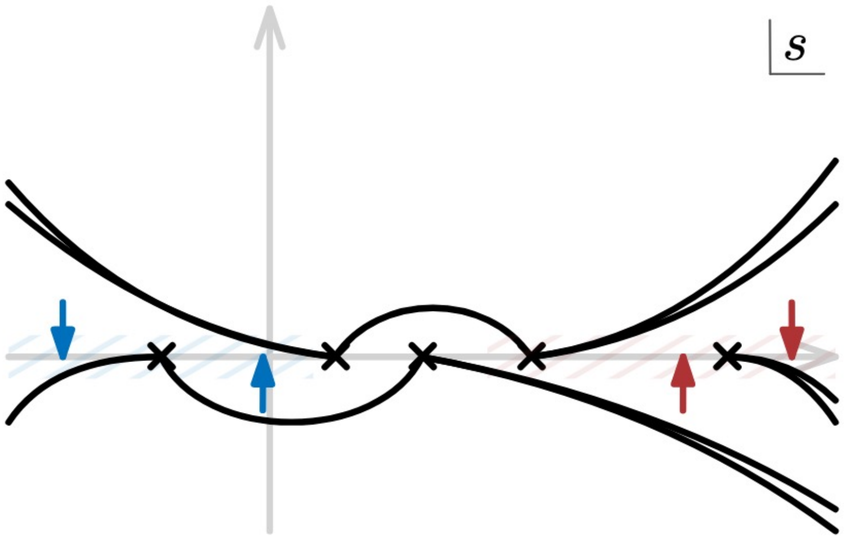
Can connect upper- and lower-half planes



Cannot connect
(two distinct analytic functions)

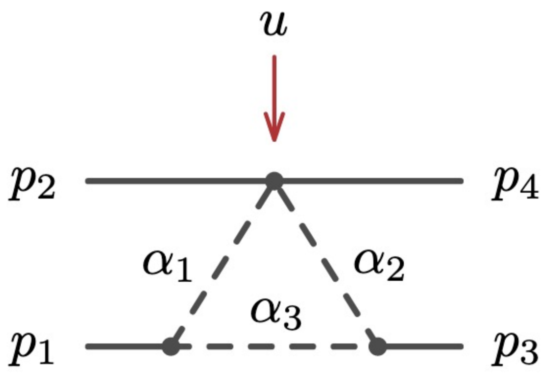
In general, there's no unique way to approach physical regions!

We are forced to perform branch cut deformations:



physical sheet

Problems with analyticity: simplest example

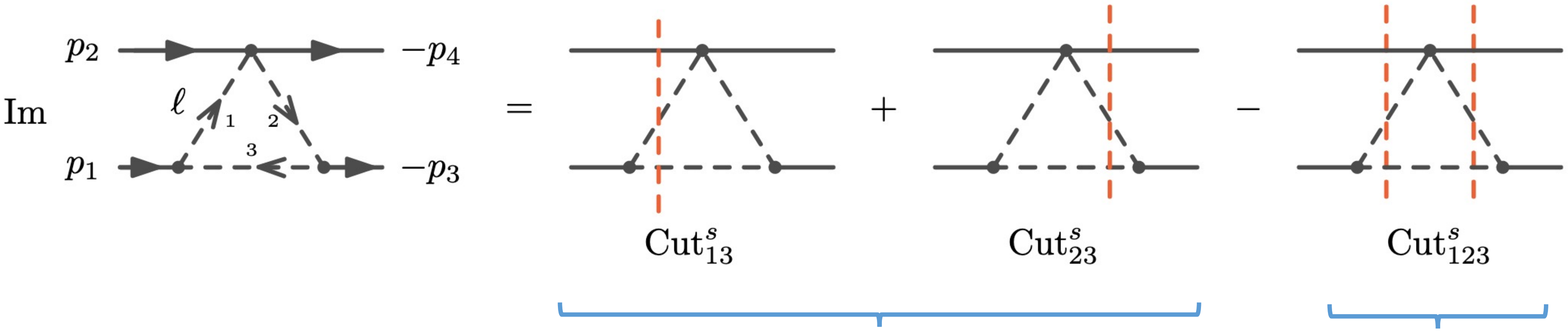


Internal masses m

External masses $M > 2m$

Action: $$\mathcal{V} = \frac{u\alpha_1\alpha_2 + M^2\alpha_3(\alpha_1+\alpha_2)}{\alpha_1+\alpha_2+\alpha_3} - m^2(\alpha_1+\alpha_2+\alpha_3)$$

Unitarity in the s-channel:

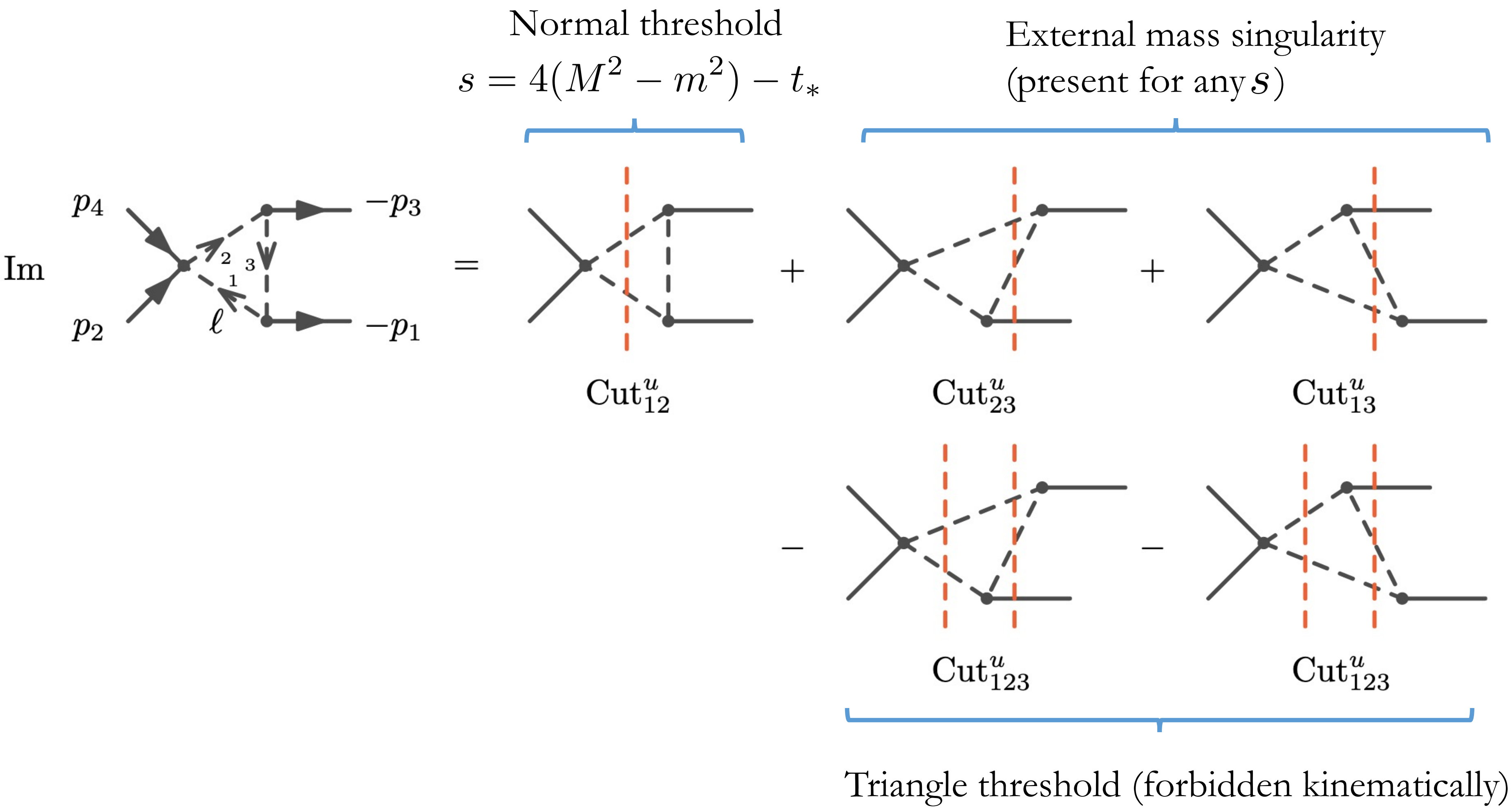


External mass singularity
(present for any s)

Triangle threshold

$$s = \frac{M^4}{m^2} - t_*$$

Unitarity in the u-channel:



Two distinct analytic functions in the UHP and LHP:

$$\mathcal{I}_{\text{tri}}^{\text{UHP}}(s,t) = \frac{z}{4M^2\beta_z} \sum_{\zeta\in\{-1,1\}} \Bigg\{ \zeta\,\mathrm{Li}_2\left(\frac{1+\frac{z}{2}-\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right) + \zeta\,\mathrm{Li}_2\left(1-\frac{1+\frac{z}{2}+\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right)$$
$$+ 2\,\mathrm{Li}_2\left(\frac{1+\beta_z}{1+\zeta\beta_z\beta_{yz}}\right) - 2\,\mathrm{Li}_2\left(\frac{1-\beta_z}{1+\zeta\beta_z\beta_{yz}}\right) + 2\pi i \log\left(\frac{1+\beta_z}{1+\beta_z\beta_{yz}}\right)$$
$$+ \zeta \log\left(\frac{1+\frac{z}{2}+\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right)\left[-\pi i + \log\left(-1+\frac{1+\frac{z}{2}+\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right)\right]\Bigg\},$$

$$\mathcal{I}_{\text{tri}}^{\text{LHP}}(s,t) = \frac{z}{4M^2\beta_z} \sum_{\zeta\in\{-1,1\}} \Bigg\{ \zeta\,\mathrm{Li}_2\left(\frac{1+\frac{z}{2}-\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right) + \zeta\,\mathrm{Li}_2\left(1-\frac{1+\frac{z}{2}+\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right)$$
$$+ 2\,\mathrm{Li}_2\left(\frac{1+\beta_z}{1+\zeta\beta_z\beta_{yz}}\right) - 2\,\mathrm{Li}_2\left(\frac{1-\beta_z}{1+\zeta\beta_z\beta_{yz}}\right) - 2\pi i \log\left(\frac{1+\beta_z}{1+\beta_z\beta_{yz}}\right)$$
$$+ \zeta \log\left(\frac{1+\frac{z}{2}+\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right)\left[\pi i + \log\left(-1+\frac{1+\frac{z}{2}+\zeta\beta_z}{1+\frac{z}{2}+\zeta\beta_y\beta_z}\right)\right]\Bigg\}.$$

where

$$y = -\frac{4m^2}{u}, \qquad z = -\frac{4M^2}{u},$$
$$\beta_y = \sqrt{1+y}, \qquad \beta_z = \sqrt{1+z}, \qquad \beta_{yz} = -i\sqrt{-1+\frac{4y}{z}}.$$

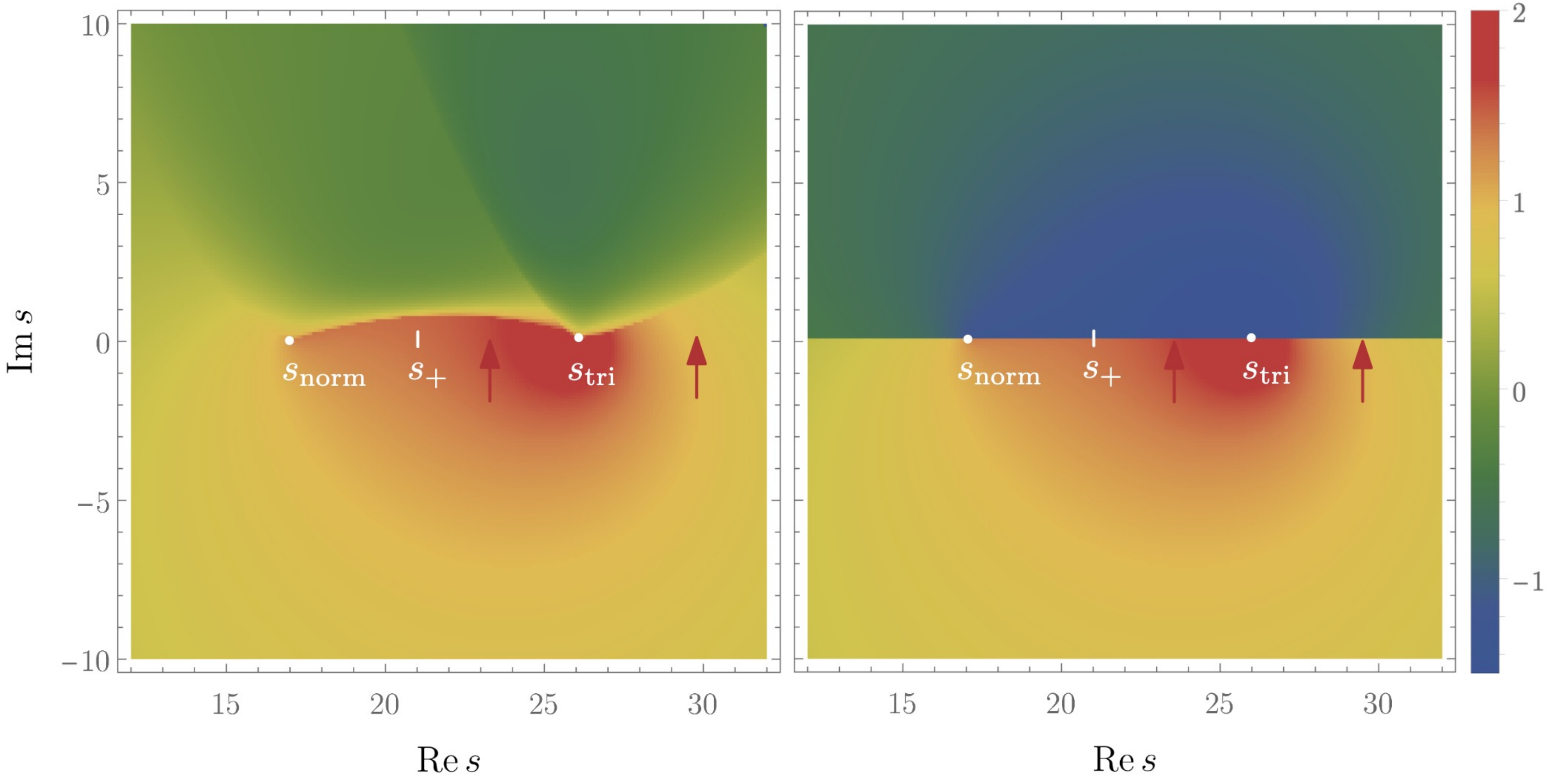
Causality requires $\mathrm{Im}\,\mathcal{V} > 0$ ($\mathcal{V} = \frac{u\alpha_1\alpha_2 + M^2\alpha_3(\alpha_1+\alpha_2)}{\alpha_1+\alpha_2+\alpha_3} - m^2(\alpha_1+\alpha_2+\alpha_3)$):

$$\mathrm{Im}\,\mathcal{V} = \mathrm{Im}\,u \frac{\alpha_1\alpha_2}{\alpha_1+\alpha_2+\alpha_3}$$

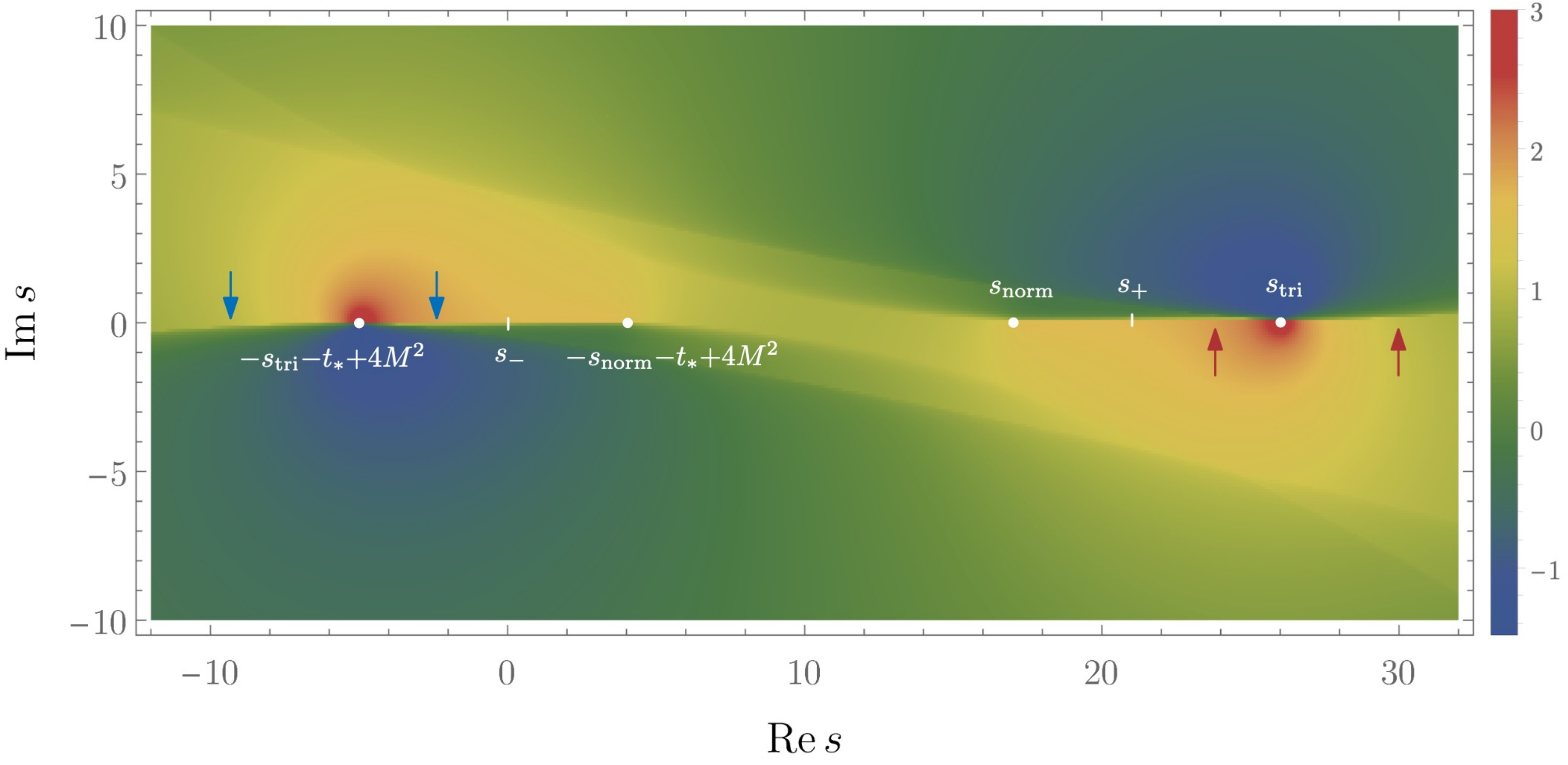
$$= -\underbrace{\mathrm{Im}\,s}_{<0} \underbrace{\frac{\alpha_1\alpha_2}{\alpha_1+\alpha_2+\alpha_3}}_{>0} > 0$$

Approach both s- and u-channel physical regions from LHP

Comparing numerical and analytic expressions:



Finally, summing over multiple Feynman diagrams



Two general results in the s-channel

- $2 \to 2$ scattering with no unstable external particles:

$$\mathbf{T}(s,t_*) = \lim_{\varepsilon\to 0^+} \mathbf{T}_{\mathbb{C}}(s+i\varepsilon,t_*) \qquad \mathrm{Im}\,\mathbf{T} = \mathrm{Disc}_s\mathbf{T}_{\mathbb{C}}$$

(previously only established when the Euclidean region exists)

- $2 \to 2$ scattering with unstable external particles:

$$\mathbf{T}(s,t_*) \neq \lim_{\varepsilon\to 0^+} \mathbf{T}_{\mathbb{C}}(s\pm i\varepsilon,t_*) \qquad \mathrm{Im}\,\mathbf{T} \neq \mathrm{Disc}_s\mathbf{T}_{\mathbb{C}}$$

(alternatively, probing higher-point analyticity)

Many open questions, for example:

- How big of a mistake we'd make by always approaching the s-channel from the UHP?

$$\propto \left(\frac{\Gamma}{M}\right)^{\#}$$

- Effect on practical Standard Model computations?

e.g., ZZ $\to$ ZZ

- What is the analogue for $2 \to 3$ scattering?

If there's time:

Fluctuations around classical saddle points

Since singularities are already determined by saddle points

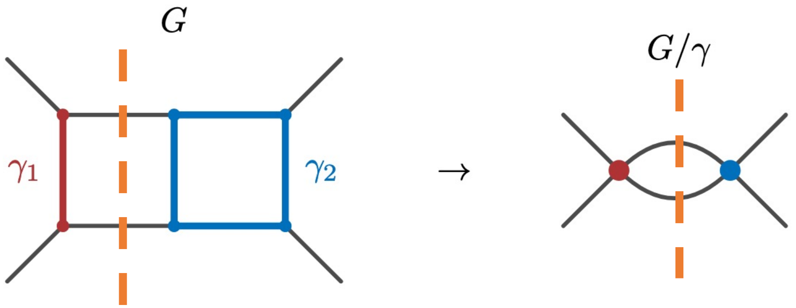
$$\alpha_e = \alpha_e^*, \qquad \Delta(s,t,M,m) = 0$$

Why don't we just study fluctuations around such saddles:

$$\alpha_e = \alpha_e^* + \delta\alpha_e + \ldots, \qquad \Delta(s,t,M,m) = 0 + \delta\Delta + \ldots$$

Local behavior around the threshold

So far limited to isolated and non-degenerate saddles
(excludes massless Feynman integrals)



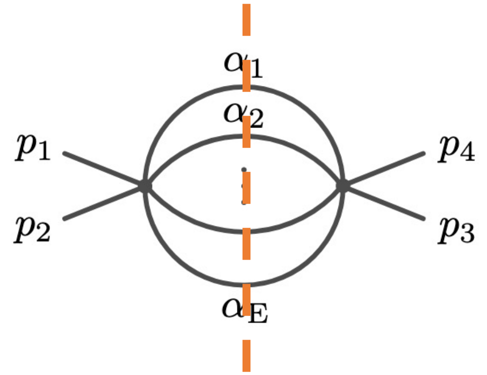
$$\mathcal{I}_G \approx \# \prod_i \mathcal{I}^*_{\gamma_i} \begin{cases} \Delta^\rho & \text{if } \rho < 0 \\ \log \Delta & \text{if } \rho = 0 \end{cases}$$

where $\rho = \frac{\mathrm{L}_{G/\gamma}\mathrm{D} - \mathrm{E}_{G/\gamma} - 2d_{\mathcal{N}_{G/\gamma}} - 1}{2}$

#loops dimension #edges numerator degree

[related work Landau, Polkinghorne, Screaton, Greenman, Kinoshita, …]

For example, near every normal threshold



$$\mathcal{I}_G \approx \# \mathcal{I}^*_{\gamma_L} \mathcal{I}^*_{\gamma_R} \begin{cases} \Gamma(-\rho) \left[\left(\sum_{e=1}^{\mathrm{E}} m_e \right)^2 - s \right]^{\rho} & \text{if } \mathrm{D} < \frac{\mathrm{E}+1}{\mathrm{E}-1}, \\ -\log \left[\left(\sum_{e=1}^{\mathrm{E}} m_e \right)^2 - s \right] & \text{if } \mathrm{D} = \frac{\mathrm{E}+1}{\mathrm{E}-1}, \end{cases}$$

$$\text{where} \quad \rho = \frac{(\mathrm{E}-1)\mathrm{D} - \mathrm{E} - 1}{2}$$

Naively, Δ^ρ would suggest that the S-matrix
can have arbitrarily-singular behavior…

We're rescued if we assume analyticity
(at most codim-1 singularities): $\mathrm{E}_{G/\gamma} - \mathrm{L}_{G/\gamma}\mathrm{D} \leqslant 1$

$$\rho = \frac{\mathrm{L}_{G/\gamma}\mathrm{D} - \mathrm{E}_{G/\gamma} - 2d_{\mathcal{N}_{G/\gamma}} - 1}{2} \geqslant -1$$

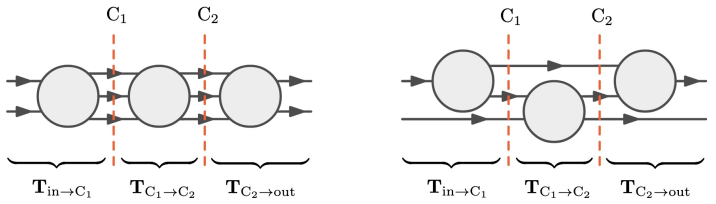
Every 1VI component can only lead to singularities of the type

$$\frac{1}{\Delta}, \qquad \frac{1}{\sqrt{\Delta}}, \qquad \log\Delta$$

Summary

- Unitarity constraints

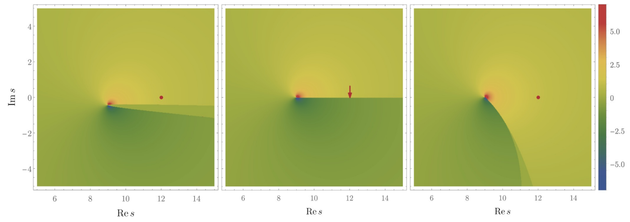
Holomorphic cutting rules



Discontinuities beyond normal thresholds

- Causality constraints

Different ways of implementing causality



Deforming branch cuts in the kinematic space

Imprints of causality more complicated than previously assumed:

- $2 \to 2$ scattering with no unstable external particles:

$$\mathbf{T}(s, t_*) = \lim_{\varepsilon \to 0^+} \mathbf{T}_{\mathbb{C}}(s + i\varepsilon, t_*) \qquad \mathrm{Im}\,\mathbf{T} = \mathrm{Disc}_s \mathbf{T}_{\mathbb{C}}$$

(previously only established when the Euclidean region exists)

- $2 \to 2$ scattering with unstable external particles:

$$\mathbf{T}(s, t_*) \neq \lim_{\varepsilon \to 0^+} \mathbf{T}_{\mathbb{C}}(s \pm i\varepsilon, t_*) \qquad \mathrm{Im}\,\mathbf{T} \neq \mathrm{Disc}_s \mathbf{T}_{\mathbb{C}}$$

(alternatively, probing higher-point analyticity)

Thank you