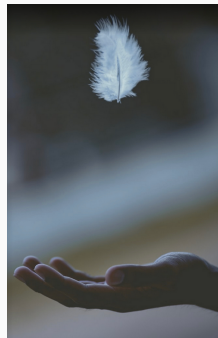


Heavy Neutral Leptons:

the FIP Physics Centre approach



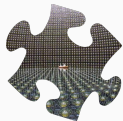
Juraj Klarić

FIPs workshop, October 21st 2022



Some puzzles for physics beyond the Standard Model

Neutrino masses



The Baryon Asymmetry of the Universe

$$n_B/n_\gamma = 6.05(7) \times 10^{-10}$$

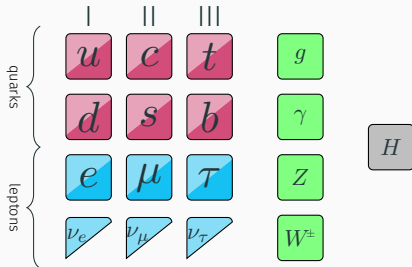
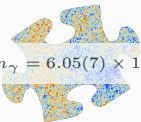
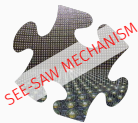


Image credits: Kamioka Observatory, ICRR, U. Tokyo; ESA and the Planck Collaboration

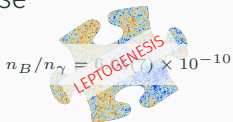
Some puzzles for physics beyond the Standard Model

Neutrino masses



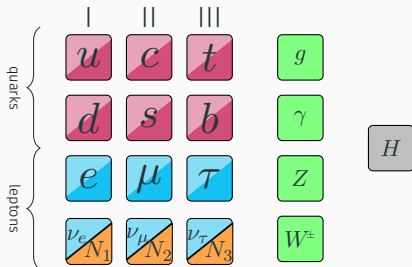
[Minkowski 1977...]

The Baryon Asymmetry of the Universe



[Fukugita/Yanagida '86...]

Image credits: Kamioka Observatory, ICRR, U. Tokyo; ESA and the Planck Collaboration



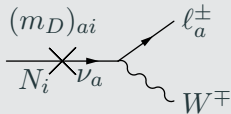
Discovering heavy neutral leptons

The seesaw mechanism and the Majorana neutrino mixing

Active neutrino masses

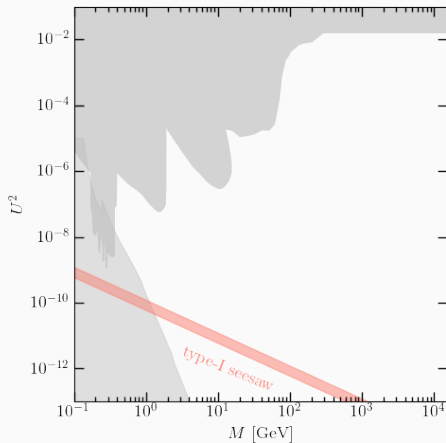
$$m_\nu = -m_D M_M^{-1} m_D^T$$

HNL mixing



$$U_{ai}^2 \equiv |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2$$



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

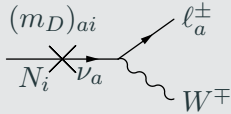
[see talks by A. Marocco, B. Dey, A. Golutvin, E. Goudzovski, A. Paolini,
J. Kopp, S. C. Middleton, JL Tastet, Giulia Ripellino]

The seesaw mechanism and the Majorana neutrino mixing

Active neutrino masses

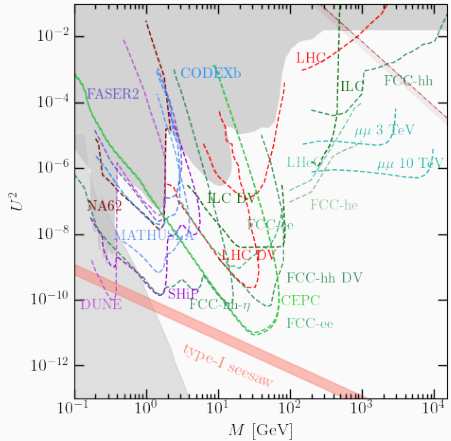
$$m_\nu = -m_D M_M^{-1} m_D^T$$

HNL mixing



$$U_{ai}^2 \equiv |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2$$



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

[see talks by A. Marocco, B. Dey, A. Golutvin, E. Goudzovski, A. Paolini,

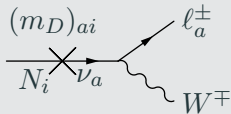
J. Kopp, S. C. Middleton, JL Tastet, Giulia Ripellino]

The seesaw mechanism and the Majorana neutrino mixing

Active neutrino masses

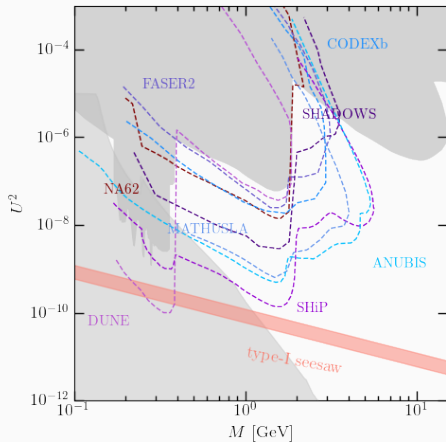
$$m_\nu = -m_D M_M^{-1} m_D^T$$

HNL mixing



$$U_{ai}^2 \equiv |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2$$



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

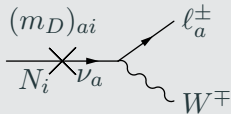
[see talks by A. Marocco, B. Dey, A. Golutvin, E. Goudzovski, A. Paolini,
J. Kopp, S. C. Middleton, JL Tastet, Giulia Ripellino]

The seesaw mechanism and the Majorana neutrino mixing

Active neutrino masses

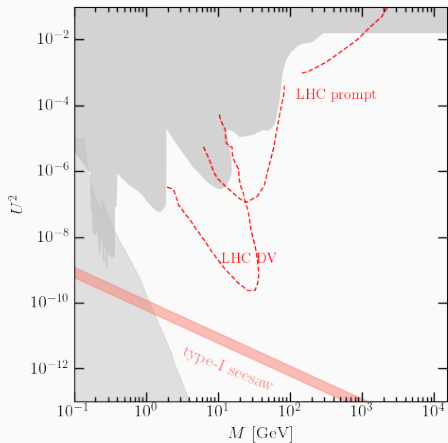
$$m_\nu = -m_D M_M^{-1} m_D^T$$

HNL mixing



$$U_{ai}^2 \equiv |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2$$



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

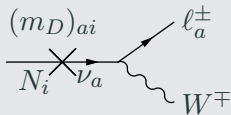
[see talks by A. Marocco, B. Dey, A. Golutvin, E. Goudzovski, A. Paolini,
J. Kopp, S. C. Middleton, JL Tastet, Giulia Ripellino]

The seesaw mechanism and the Majorana neutrino mixing

Active neutrino masses

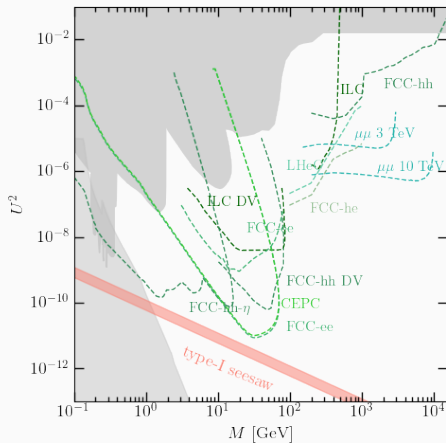
$$m_\nu = -m_D M_M^{-1} m_D^T$$

HNL mixing



$$U_{ai}^2 \equiv |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2$$



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

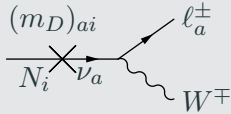
[see talks by A. Marocco, B. Dey, A. Golutvin, E. Goudzovski, A. Paolini, J. Kopp, S. C. Middleton, JL Tasset, Giulia Ripellino]

The seesaw mechanism and the Majorana neutrino mixing

Active neutrino masses

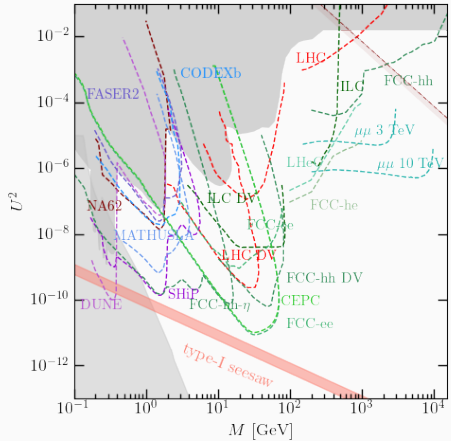
$$m_\nu = -m_D M_M^{-1} m_D^T$$

HNL mixing



$$U_{ai}^2 \equiv |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2$$



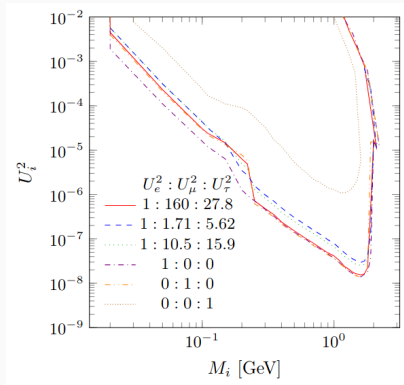
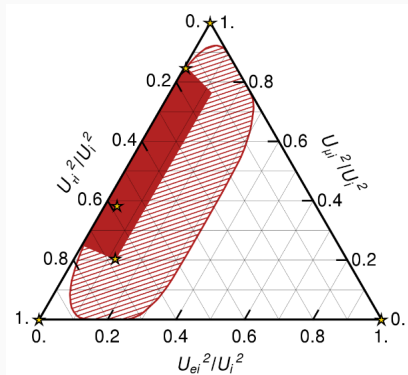
[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

[see talks by A. Marocco, B. Dey, A. Golutvin, E. Goudzovski, A. Paolini,

J. Kopp, S. C. Middleton, JL Tasset, Giulia Ripellino]

Realistic benchmarks

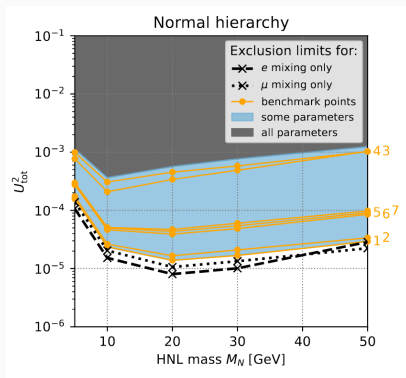
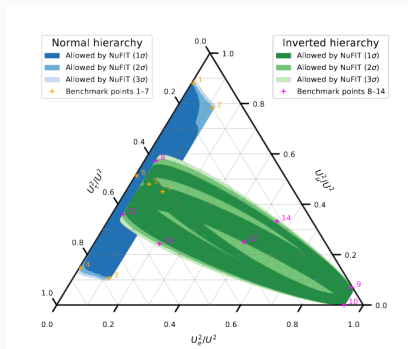
Sensitivity of experiments highly depends on mixing ratios: NA62 in beam dump mode



[Drewes/Hajer/JK/Lanfranchi 1801.04207]

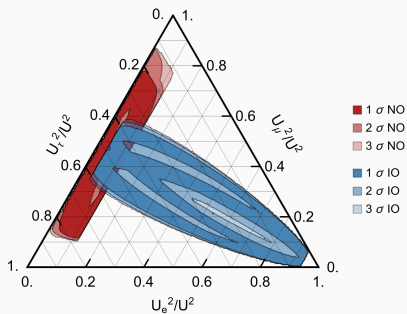
Sensitivity of experiments highly depends on mixing ratios:

ATLAS



[Tastet/Ruchayskiy/Timiryasov 2107.12980] [see also talk by JL Tastet]

Constraints from the seesaw mechanism

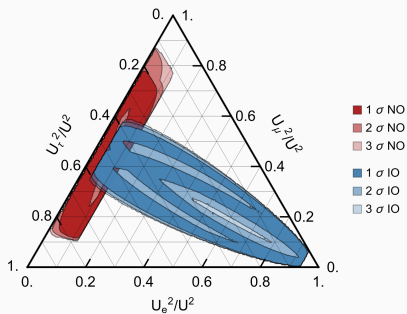


[Drewes/JK/Lopez-Pavon 2207.02742]

[using nuFIT 5.1 2007.14792]

- in the minimal seesaw model the ratios are completely determined by U_{PMNS}
- ratios dominated by Majorana phase η , Dirac phase δ and θ_{23}
- allowed ratios become smaller as we pin down the PMNS parameters
- How to choose future-proof benchmarks?

Constraints from the seesaw mechanism



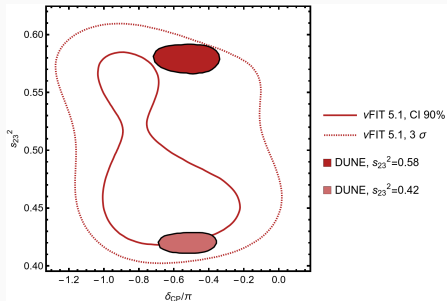
[Drewes/JK/Lopez-Pavon 2207.02742]

[using nuFIT 5.1 2007.14792]

- in the minimal seesaw model the ratios are completely determined by U_{PMNS}
- ratios dominated by Majorana phase η , Dirac phase δ and θ_{23}
- allowed ratios become smaller as we pin down the PMNS parameters
- How to choose future-proof benchmarks?

Future sensitivity?

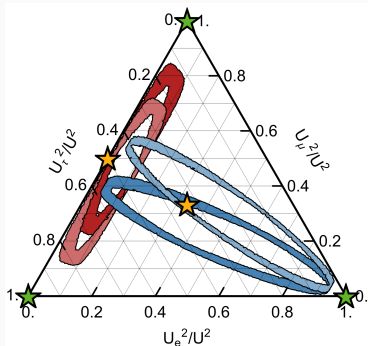
- significant improvement expected with DUNE and HyperK
- we can use the sensitivity estimates to estimate how the allowed flavor ratios change



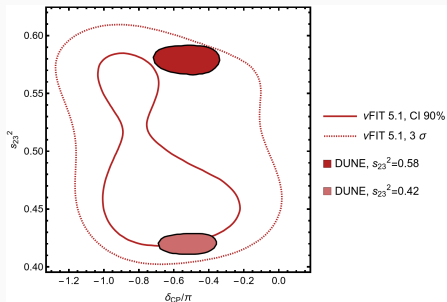
[nuFIT 5.1 2007.14792]

[DUNE TDR 2002.03005]

Future sensitivity?



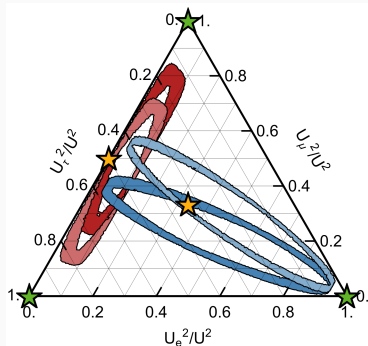
[Figure from 2207.02742]



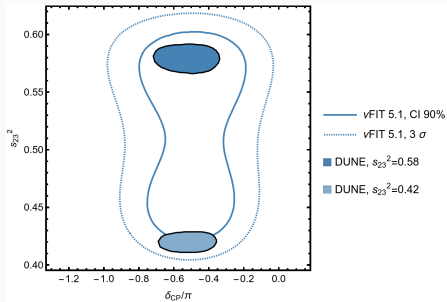
[nuFIT 5.1 2007.14792]

[DUNE TDR 2002.03005]

Future sensitivity?



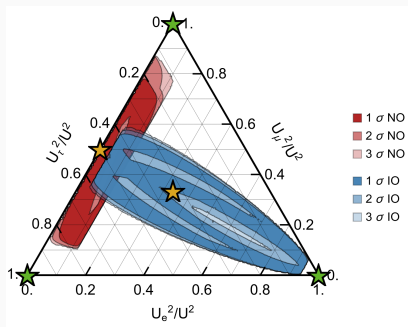
[Figure from 2207.02742]



[nuFIT 5.1 2007.14792]

[DUNE TDR 2002.03005]

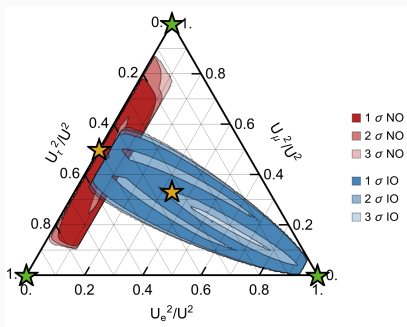
New Benchmark Points



[Figure from 2207.02742]

- selection criteria:
 1. consistency with ν -osc. data
 2. added value
 3. symmetry considerations
 4. simplicity
 5. leptogenesis
- in addition to the single flavor benchmarks, we propose the new points:
 - $U_e^2 : U_\mu^2 : U_\tau^2 = 0 : 1 : 1$
 - $U_e^2 : U_\mu^2 : U_\tau^2 = 1 : 1 : 1$
- Common benchmarks can be used to compare the reach of different searches
- Reinterpretations still highly desirable [see talk by JL Tastet]

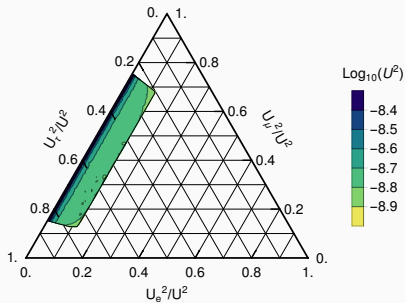
New Benchmark Points



[Figure from 2207.02742]

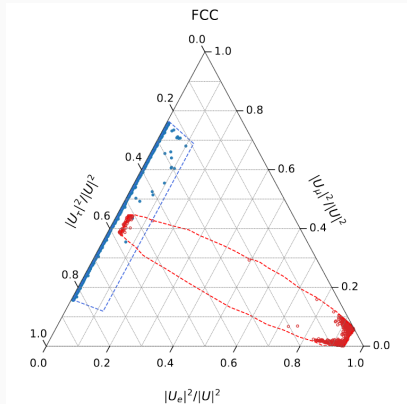
- selection criteria:
 1. consistency with ν -osc. data
 2. added value
 3. symmetry considerations
 4. simplicity
 5. leptogenesis
- in addition to the single flavor benchmarks, we propose the new points:
 - $U_e^2 : U_\mu^2 : U_\tau^2 = 0 : 1 : 1$
 - $U_e^2 : U_\mu^2 : U_\tau^2 = 1 : 1 : 1$
- Common benchmarks can be used to compare the reach of different searches
- Reinterpretations still highly desirable [see talk by JL Tastet]

Flavor constraints from leptogenesis



[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

1710.03744]

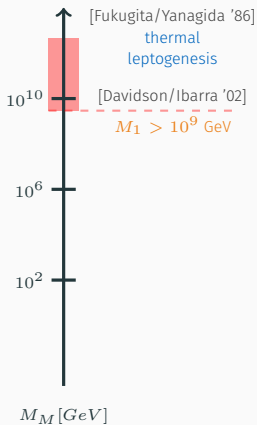


$$\Delta M/M = 10^{-2}$$

[Hernandez/Lopez-Pavon/Rius/Sandner 2207.01651]

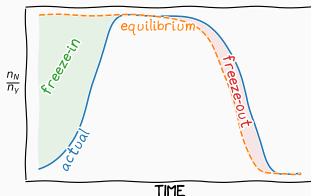
Input from cosmology: low-scale leptogenesis

Leptogenesis mechanisms



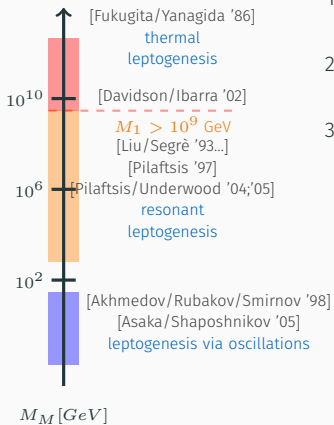
Sakharov conditions

1. Baryon number violation
sphaleron processes
2. C and CP violation
RHN decays and oscillations
3. Deviation from thermal equilibrium
freeze-in and freeze-out of RHN



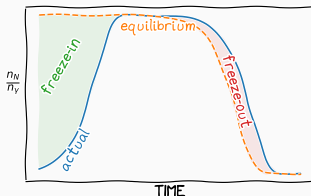
- for hierarchical RHN $M_1 \gtrsim 10^9 \text{ GeV}$

Leptogenesis mechanisms



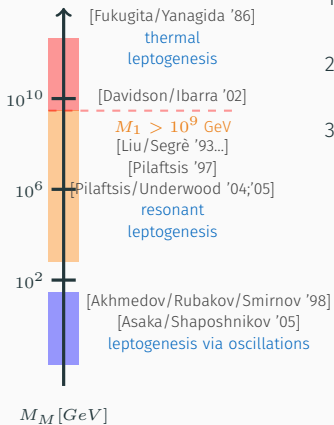
Sakharov conditions

1. Baryon number violation
sphaleron processes
2. C and CP violation
RHN decays and oscillations
3. Deviation from thermal equilibrium
freeze-in and freeze-out of RHN



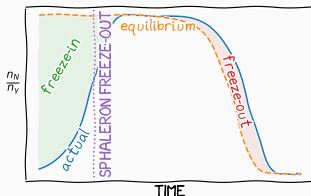
- for hierarchical RHN $M_1 \gtrsim 10^9 \text{ GeV}$
- leptogenesis works in a wide range of RHN masses

Leptogenesis mechanisms



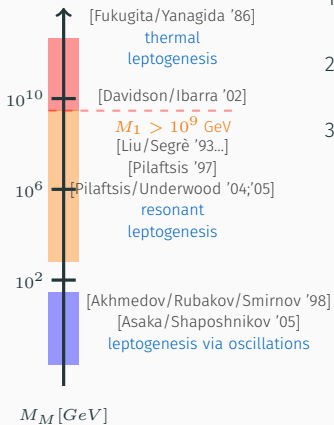
Sakharov conditions

1. Baryon number violation
sphaleron processes
2. C and CP violation
RHN decays and oscillations
3. Deviation from thermal equilibrium
freeze-in and freeze-out of RHN



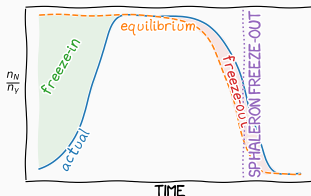
- for hierarchical RHN $M_1 \gtrsim 10^9$ GeV
- leptogenesis works in a wide range of RHN masses

Leptogenesis mechanisms



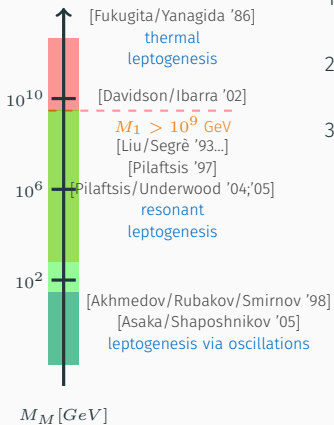
Sakharov conditions

1. Baryon number violation
sphaleron processes
2. C and CP violation
RHN decays and oscillations
3. Deviation from thermal equilibrium
freeze-in and freeze-out of RHN



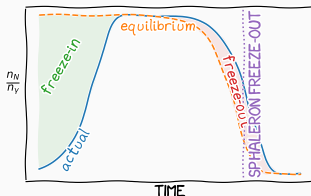
- for hierarchical RHN $M_1 \gtrsim 10^9$ GeV
- leptogenesis works in a wide range of RHN masses

Leptogenesis mechanisms



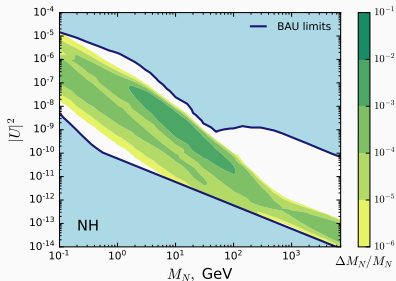
Sakharov conditions

1. Baryon number violation
sphaleron processes
2. C and CP violation
RHN decays and oscillations
3. Deviation from thermal equilibrium
freeze-in and freeze-out of RHN



- for hierarchical RHN $M_1 \gtrsim 10^9$ GeV
- leptogenesis works in a wide range of RHN masses
- how are the low-scale mechanisms connected?

Results: The minimal model with 2 RHNs

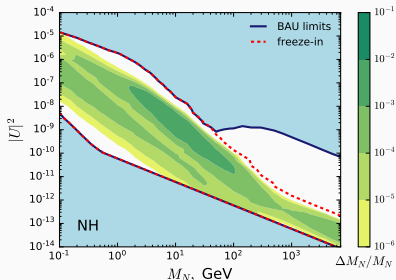


[JK/Timiryasov/Shaposhnikov 2103.16545]

- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$
- overall agreement with analytic approximations from [2207.01651] [see talk by S. Sandner]

Results: The minimal model with 2 RHNs

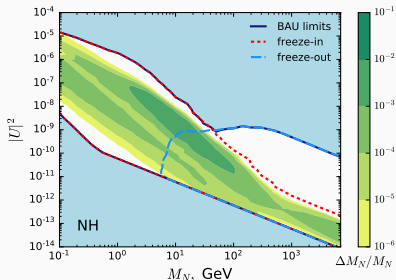


[JK/Timiryasov/Shaposhnikov 2103.16545]

- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$
- overall agreement with analytic approximations from [2207.01651] [see talk by S. Sandner]

Results: The minimal model with 2 RHNs

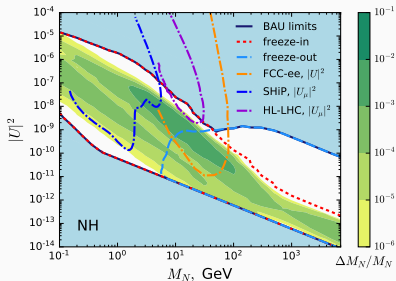


[JK/Timiryasov/Shaposhnikov 2103.16545]

- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$
- overall agreement with analytic approximations from [2207.01651] [see talk by S. Sandner]

Results: The minimal model with 2 RHNs

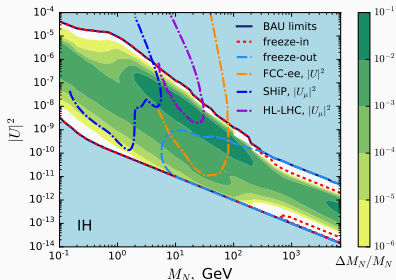


[JK/Timiryasov/Shaposhnikov 2103.16545]

- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$
- overall agreement with analytic approximations from [2207.01651] [see talk by S. Sandner]

Results: The minimal model with 2 RHNs

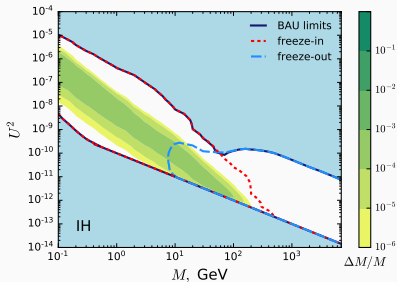


[JK/Timiryasov/Shaposhnikov 2103.16545]

- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$
- overall agreement with analytic approximations from [2207.01651] [see talk by S. Sandner]

Results: The minimal model with 2 RHNs

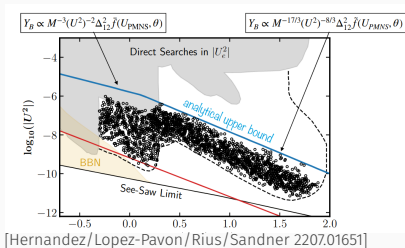


[JK/Timiryasov/Shaposhnikov 2103.16545]

- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$
- overall agreement with analytic approximations from [2207.01651] [see talk by S. Sandner]

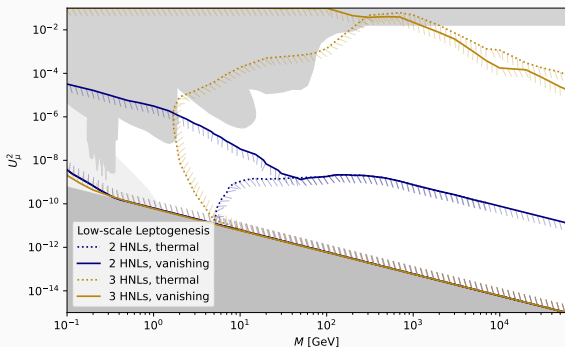
Results: The minimal model with 2 RHNs



- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$
- overall agreement with analytic approximations from [2207.01651] [see talk by S. Sandner]

Results: Leptogenesis with 3 RHNs

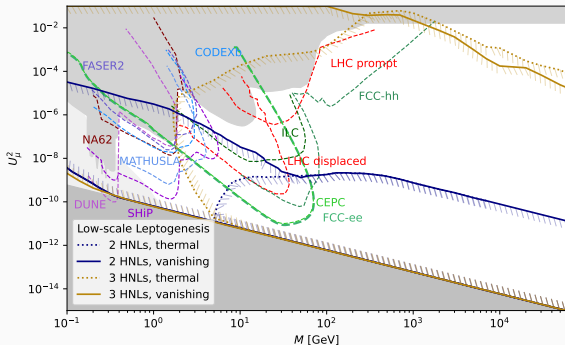


[Snowmass White Paper 2203.08039]

leptogenesis bounds from [Drewes/Georis/JK 2106.16226]

- for experimentally accessible heavy neutrino masses, all U^2 are allowed
- both **freeze-in** and **freeze-out** leptogeneses already testable at existing experiments
- the maximal value of U^2 depends on m_1

Results: Leptogenesis with 3 RHNs

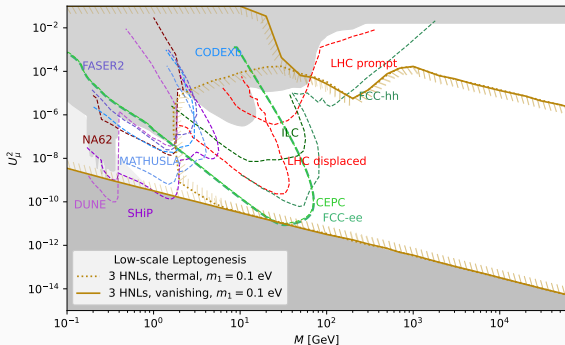


[Snowmass White Paper 2203.08039]

leptogenesis bounds from [Drewes/Georis/JK 2106.16226]

- for experimentally accessible heavy neutrino masses, all U^2 are allowed
- both freeze-in and freeze-out leptogeneses already testable at existing experiments
- the maximal value of U^2 depends on m_1

Results: Leptogenesis with 3 RHNs



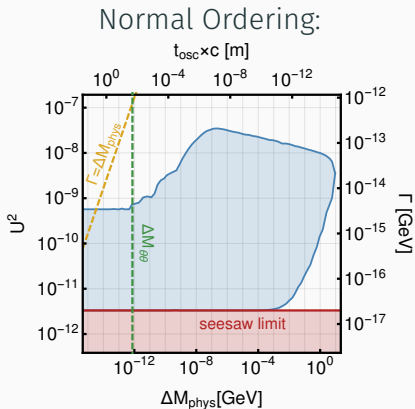
[Snowmass White Paper 2203.08039]

leptogenesis bounds from [Drewes/Georis/JK 2106.16226]

- for experimentally accessible heavy neutrino masses, all U^2 are allowed
- both freeze-in and freeze-out leptogeneses already testable at existing experiments
- the maximal value of U^2 depends on m_1

From discovery to tests

Measuring the mass splitting in model with 2 HNLs



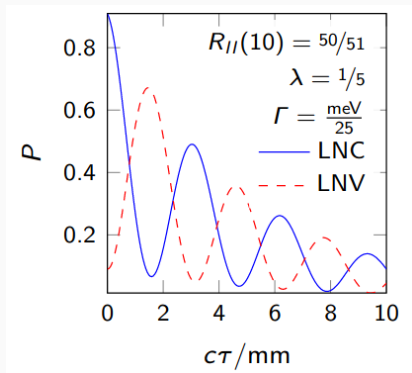
$$M = 30 \text{ GeV}$$

[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

1710.03744]

- large range of ΔM consistent with leptogenesis
- energy resolution of planned experiments - $\Delta M/M \sim \mathcal{O}(\text{few}\%)$
- Higgs vev contribution to RHN mass difference $\Delta M_{\theta\theta}$ practically implies lower limit on the mass splitting

Measuring the mass splitting in model with 2 HNLs

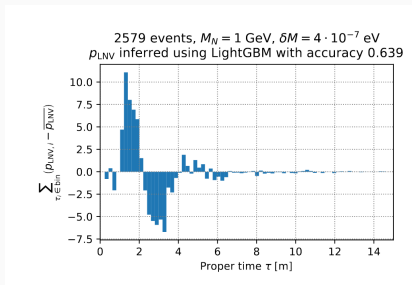


[Antusch/Hajer/Roskopp 2210.10738]

[see talk by J. Hajer]

- large range of ΔM consistent with leptogenesis
- energy resolution of planned experiments - $\Delta M/M \sim \mathcal{O}(\text{few}\%)$
- Higgs vev contribution to RHN mass difference $\Delta M_{\theta\theta}$ practically implies lower limit on the mass splitting

Measuring the mass splitting in model with 2 HNLs

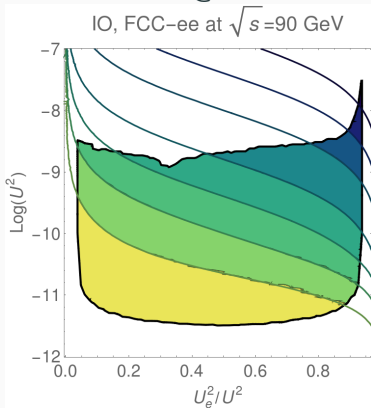


[Tastet/Timiryasov 1912.05520]

- large range of ΔM consistent with leptogenesis
- energy resolution of planned experiments - $\Delta M/M \sim \mathcal{O}(\text{few}\%)$
- Higgs vev contribution to RHN mass difference $\Delta M_{\theta\theta}$ practically implies lower limit on the mass splitting

Measuring flavor ratios at experiments

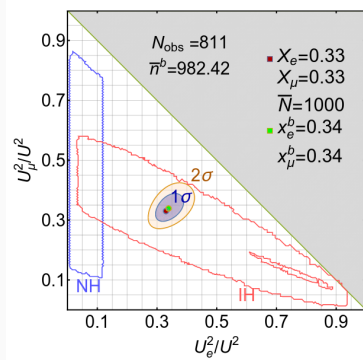
$M_N = 30$ GeV @ FCC-ee



[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

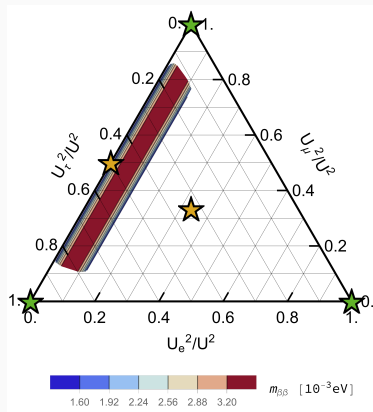
1710.03744]

$M_N = 1$ GeV @ SHiP



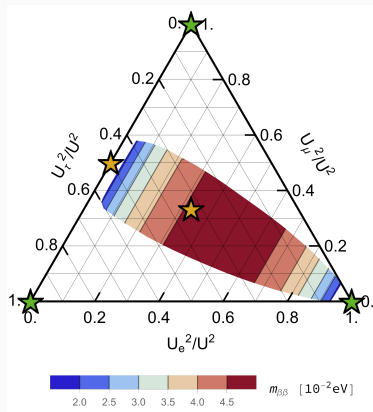
[Snowmass HNL WP 2203.08039]

Complementarity with neutrinoless double beta decay



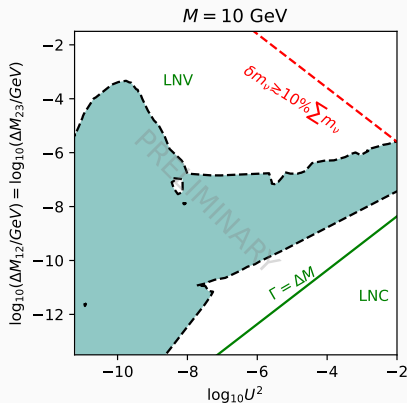
- $m_{\beta\beta}$ is a complementary probe of the flavor mixing ratios for $M_N \gg 100 \text{MeV}$
- “NH” benchmark minimizes $m_{\beta\beta}$ for normal hierarchy
- “IH” benchmark close to maximal $m_{\beta\beta}$ for inverted hierarchy

Complementarity with neutrinoless double beta decay



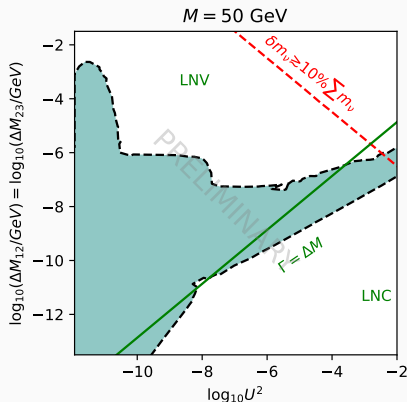
- $m_{\beta\beta}$ is a complementary probe of the flavor mixing ratios for $M_N \gg 100\text{MeV}$
- “NH” benchmark minimizes $m_{\beta\beta}$ for normal hierarchy
- “IH” benchmark close to maximal $m_{\beta\beta}$ for inverted hierarchy

3 HNLs: range of mixing angles U^2 and mass splittings ΔM



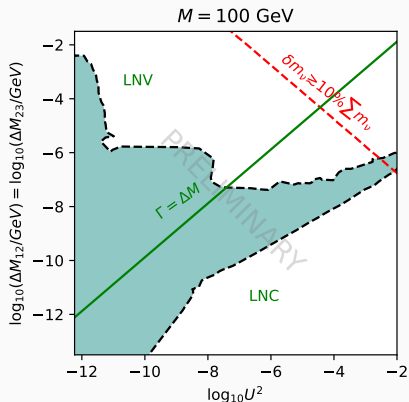
- benchmark with fixed $U_{\alpha I}^2/U^2$
- upper bound on U^2 arises through a combination of baryogenesis + fine tuning constraints
- leptogenesis consistent with both LNV and LNC RHN decays
- nontrivial LNV/LNC ratios can further constrain the RHN parameters

3 HNLs: range of mixing angles U^2 and mass splittings ΔM



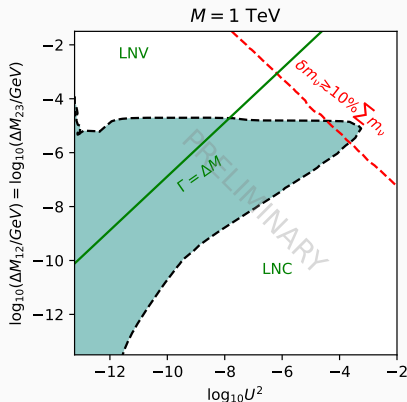
- benchmark with fixed $U_{\alpha I}^2/U^2$
- upper bound on U^2 arises through a combination of baryogenesis + fine tuning constraints
- leptogenesis consistent with both LNV and LNC RHN decays
- nontrivial LNV/LNC ratios can further constrain the RHN parameters

3 HNLs: range of mixing angles U^2 and mass splittings ΔM



- benchmark with fixed $U_{\alpha I}^2/U^2$
- upper bound on U^2 arises through a combination of baryogenesis + fine tuning constraints
- leptogenesis consistent with both LNV and LNC RHN decays
- nontrivial LNV/LNC ratios can further constrain the RHN parameters

3 HNLs: range of mixing angles U^2 and mass splittings ΔM



- benchmark with fixed $U_{\alpha I}^2/U^2$
- upper bound on U^2 arises through a combination of baryogenesis + fine tuning constraints
- leptogenesis consistent with both LNV and LNC RHN decays
- nontrivial LNV/LNC ratios can further constrain the RHN parameters

[Drewes/Georis/JK 22Xx.xxxx]

Conclusions

- experimental sensitivities differ depending on flavor ratios
- tools for **reinterpretation** are highly desirable
- proposed new benchmarks for HNL searches
- when accessible **LNV** is an important observable $R_{\ell\ell}$
- leptogenesis is a viable baryogenesis mechanism for **all heavy neutrino masses** above the $\mathcal{O}(100)$ MeV scale
- leptogenesis is within reach at planned future experiments
 - synergy between **high-energy** and **high-intensity** frontiers!
 - together they can cover a large portion of the low-scale leptogenesis parameter space

Thank you!

Large mixing angles and approximate B-L symmetry

- large U^2 require cancellations between different entries of the Yukawa matrices F
- this cancellation can be associated with an approximate lepton number symmetry

[Shaposhnikov hep-ph/0605047, Kersten Smirnov

0705.3221, Moffat Pascoli Weiland 1712.07611]

- symmetry broken by small parameters $\epsilon, \epsilon', \mu, \mu'$

Pseudo-Dirac pairs

$$N_s = \frac{N_1 + iN_2}{\sqrt{2}}, N_w = \frac{N_1 - iN_2}{\sqrt{2}}$$

B-L parametrisation

$$M_M = \bar{M} \begin{pmatrix} 1 - \mu & 0 & 0 \\ 0 & 1 + \mu & 0 \\ 0 & 0 & \mu' \end{pmatrix}$$

$$F = \frac{1}{\sqrt{2}} \begin{pmatrix} F_e(1 + \epsilon_e) & iF_e(1 - \epsilon_e) & F_e\epsilon'_e \\ F_\mu(1 + \epsilon_\mu) & iF_\mu(1 - \epsilon_\mu) & F_\mu\epsilon'_\mu \\ F_\tau(1 + \epsilon_\tau) & iF_\tau(1 - \epsilon_\tau) & F_\tau\epsilon'_\tau \end{pmatrix}$$

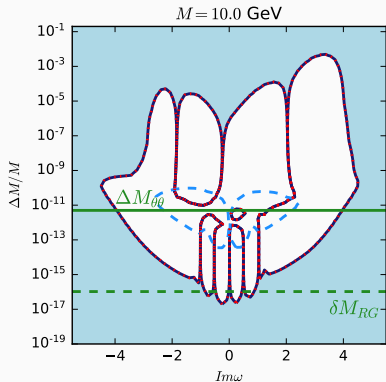
Fine tuning

- if present, symmetries are manifest to all orders in p.t.
- in the case of a large B-L breaking, radiative corrections can cause large neutrino masses
- we can use the size of radiative corrections to the light neutrino masses to quantify tuning

Fine Tuning

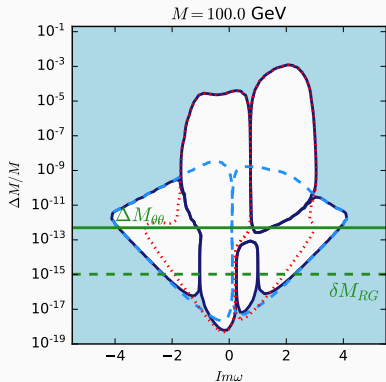
$$f.t.(m_\nu) = \sqrt{\sum_{i=1}^3 \left(\frac{m_i^{\text{loop}} - m_i^{\text{tree}}}{m_i^{\text{loop}}} \right)^2}$$

Slices of the parameter space



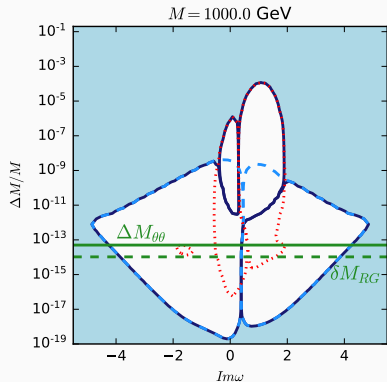
- two characteristic mass splittings
- mass splitting induced by the Higgs $\Delta M_{\theta\theta}$
- mass splitting induced by RG running δM_{RG}

Slices of the parameter space



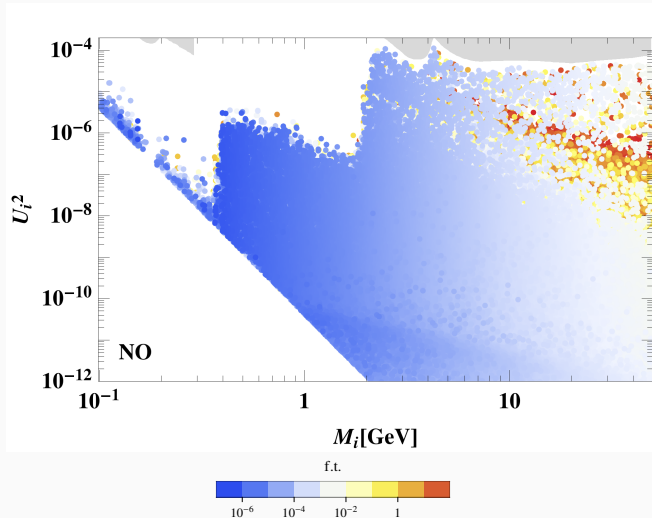
- two characteristic mass splittings
- mass splitting induced by the Higgs $\Delta M_{\theta\theta}$
- mass splitting induced by RG running δM_{RG}

Slices of the parameter space



- two characteristic mass splittings
- mass splitting induced by the Higgs $\Delta M_{\theta\theta}$
- mass splitting induced by RG running δM_{RG}

Results: Leptogenesis with 3 RHN (Normal Ordering)



[Abada/Arcadi/Domcke/Drewes/JK/Lucente 1810.12463]

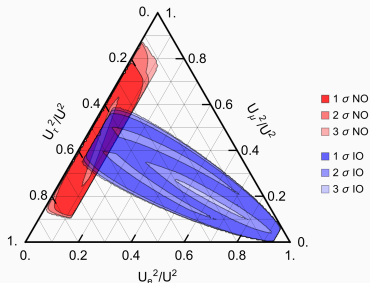
Hierarchy in the washout

- lepton asymmetry can survive washout if hidden in a particular flavor
- washout suppression

$$f \equiv \frac{\Gamma_a}{\Gamma} \sim \frac{U_a^2}{U^2}$$

- for 2 RHN $f > 5 \times 10^{-3}$
- for 3 RHN $f \ll 1$ possible

2 RHNs:



[Snowmass White Paper 2203.08039]

[Drewes/Garbrecht/Gueter/JK 1609.09069]

[Caputo/Hernandez/Lopez-Pavon/Salvado 1704.08721]

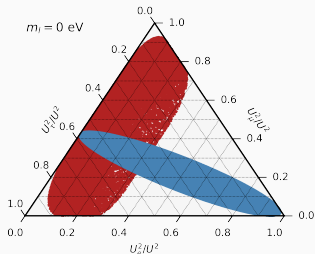
Hierarchy in the washout

- lepton asymmetry can survive washout if hidden in a particular flavor
- washout suppression

$$f \equiv \frac{\Gamma_a}{\Gamma} \sim \frac{U_a^2}{U^2}$$

- for 2 RHN $f > 5 \times 10^{-3}$
- for 3 RHN $f \ll 1$ possible

3 RHNs:



[Drewes/Georis/JK 220x.xxxx]

[Chrzaszcz/Drewes/Gonzalo/Harz/Krishnamurthy/Weniger 1908.02302]

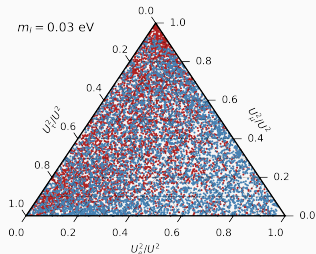
Hierarchy in the washout

- lepton asymmetry can survive washout if hidden in a particular flavor
- washout suppression

$$f \equiv \frac{\Gamma_a}{\Gamma} \sim \frac{U_a^2}{U^2}$$

- for 2 RHN $f > 5 \times 10^{-3}$
- for 3 RHN $f \ll 1$ possible

3 RHNs:



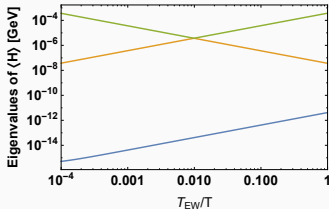
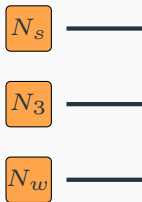
[Drewes/Georis/JK 220x.xxxx]

[Chrzaszcz/Drewes/Gonzalo/Harz/Krishnamurthy/Weniger 1908.02302]

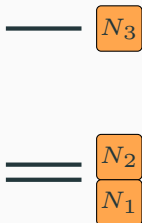
Enhancement due to level crossing

- in the $B - L$ symmetric limit two heavy neutrinos form a pseudo-Dirac pair
- the “3rd” heavy neutrino can be heavier than the pseudo-Dirac pair
- for $T \gg T_{EW}$, the pseudo-Dirac pair also has a thermal mass

$T \gg T_{EW}$

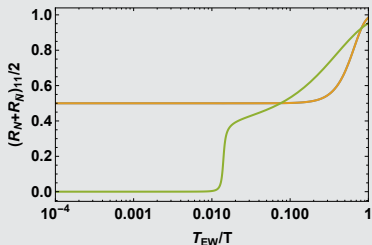


$T \ll T_{EW}$

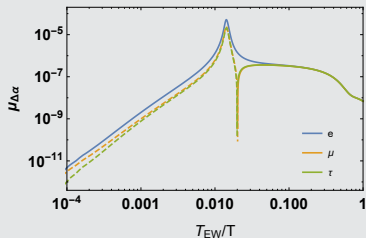


Enhancement due to level crossing

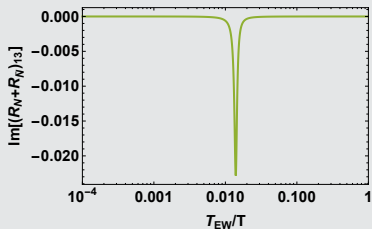
Heavy Neutrino Densities



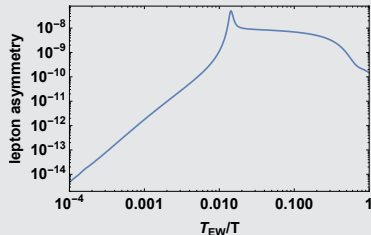
Lepton flavour asymmetries



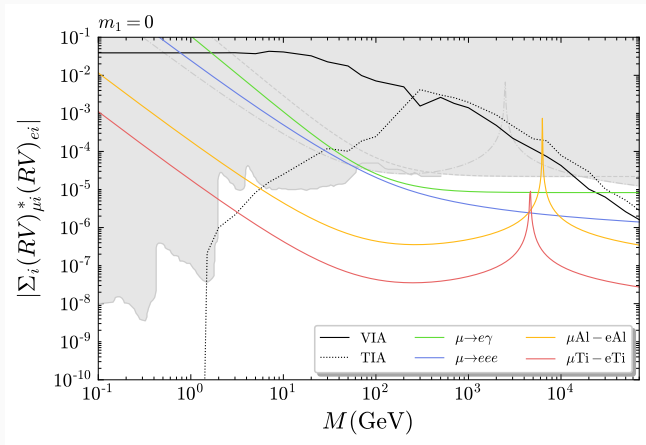
Heavy Neutrino correlations



Lepton number asymmetry



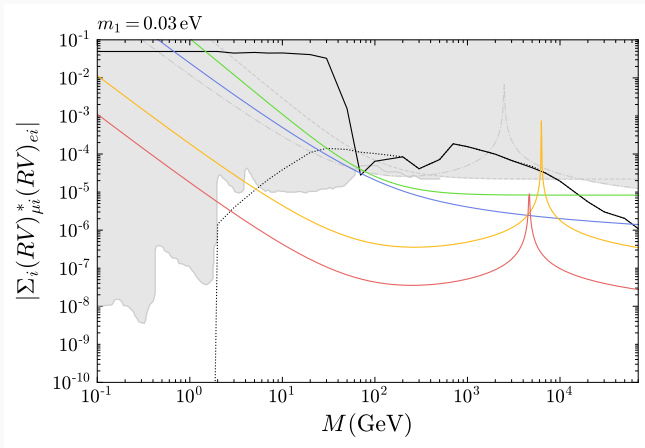
Indirect probes: Charged LFV



[Granelli/JK/Petcov 2206.04342]

- parameters space in the TeV region already severely constrained by cLFV observables
- future $\mu \rightarrow e$ conversion experiments can probe a large part of the $N = 3$ parameter space

Indirect probes: Charged LFV



[Granelli/JK/Petcov 2206.04342]

- parameters space in the TeV region already severely constrained by cLFV observables
- future $\mu \rightarrow e$ conversion experiments can probe a large part of the $N = 3$ parameter space