

IWHSS-2022

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CERN



**International Workshop on Hadron Structure  
and Spectroscopy - 2022**

# **Spin effects in unpolarized SIDIS using the string + $^3P_0$ model**

Albi Kerbizi

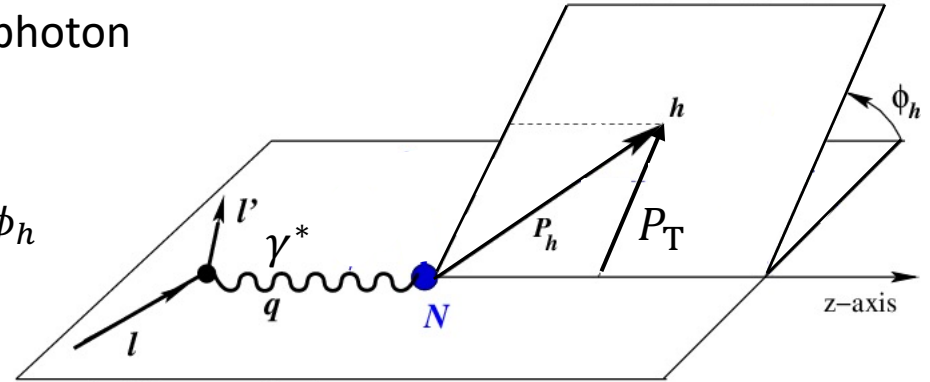
*in collaboration with Anna Martin*



# Spin effects in the azimuthal asymmetries in unpolarized SIDIS

Unpolarized SIDIS cross section in the one-photon exchange approximation:

$$\frac{d\sigma}{dx dz dy dP_T^2 d\phi_h} \propto 1 + \epsilon_1(y) A_{UU}^{\cos \phi_h} \cos \phi_h + \epsilon_2(y) A_{UU}^{\cos 2\phi_h} \cos 2\phi_h$$



Partonic content of the azimuthal asymmetries:

$$A_{UU}^{\cos \phi_h} = \frac{2M_N}{Q} C \left[ \overset{\text{Cahn effect}}{\overset{\text{twist 3}}{-\frac{\hat{h} \cdot \vec{k}_T}{M_N} f_1^q D_{1q}^h}} - \overset{\text{Boer-Mulders}}{\overset{\text{twist 3}}{\frac{(\hat{h} \cdot \vec{p}_\perp) k_T^2}{zM_N^2 M_h} h_1^{\perp q} H_{1q}^{\perp h} + \dots}} \right] / F_{UU}$$

$$A_{UU}^{\cos 2\phi_h} = C \left[ \overset{\text{Cahn effect}}{\overset{\text{twist 4}}{2 \frac{2(\vec{k}_T \cdot \hat{h})^2 - k_T^2}{Q^2} f_1^q D_{1q}^h}} - \overset{\text{Boer-Mulders}}{\overset{\text{twist 2}}{\frac{2(\vec{k}_T \cdot \hat{h})(\vec{p}_\perp \cdot \hat{h}) - \vec{k}_T \cdot \vec{p}_\perp}{zM_N M_h} h_1^{\perp q} H_{1q}^{\perp h}}} \right] / F_{UU}$$

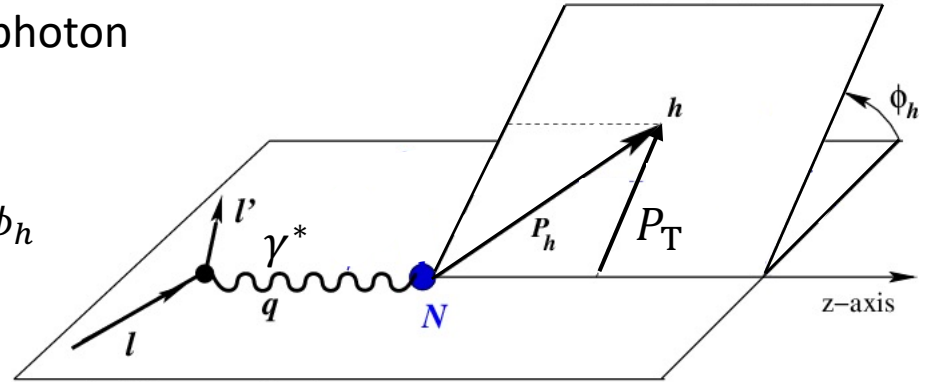
(the only know term..)

$$F_{UU} = C [f_1^q D_1^q]$$

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(the only known term..)

$$F_{UU} = C [f_1^q D_{11}^q]$$

- Important observables, enable access to

- i. intrinsic  $\vec{k}_T$
- ii. Boer-Mulders TMD

Many data exist:

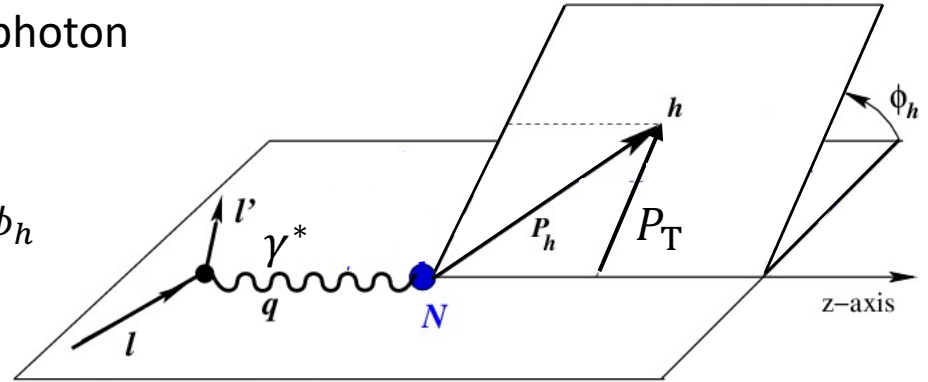
HERMES	p, d
COMPASS	d, recently p
JLAB	p

large azimuthal asymmetries observed

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*Cahn effect*  
twist 4  
(the only known term..)

*Boer-Mulders*  
twist 2

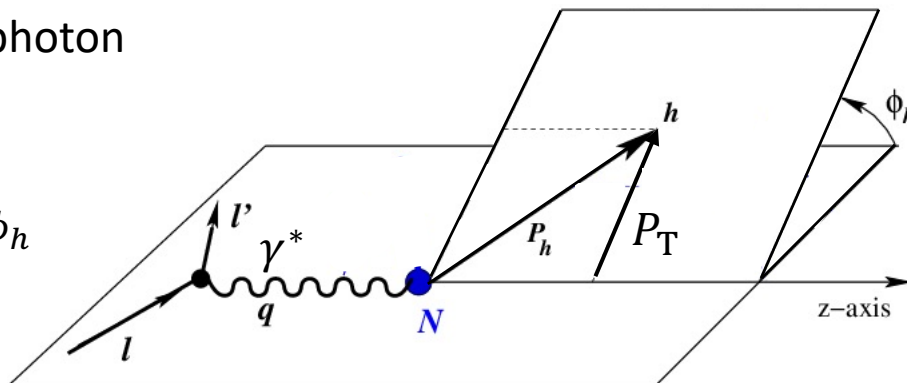
*Boer-Mulders TMD:*  
describes transversely polarized quarks in an unpolarized nucleon

*Collins FF:*  
fragmentation of transversely polarized quarks in unpolarized hadrons

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Boer-Mulders

twist 2

Collins FF:  
fragmentation of transversely polarized quarks in unpolarized hadrons

$$F_{UU} = C [f_1^q D_1^q]$$

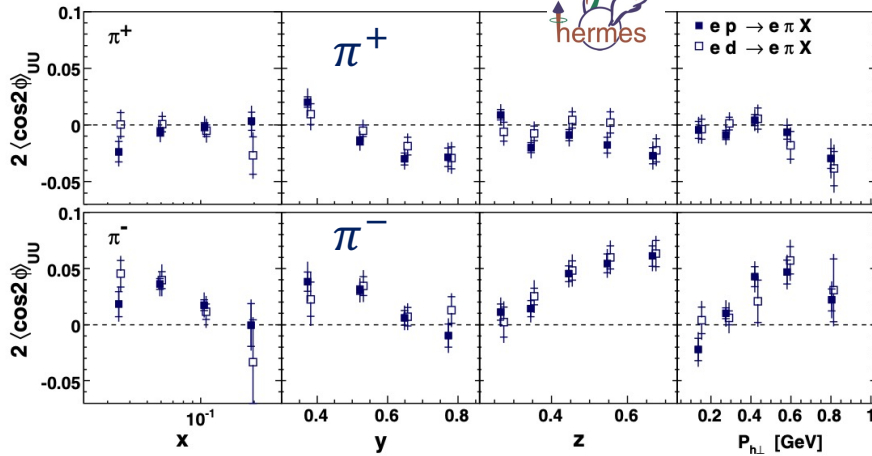
Focus on  $A_{UU}^{\cos 2\phi_h}$  in the next slides

# Data on the $A_{UU}^{\cos 2\phi_h}$ asymmetry

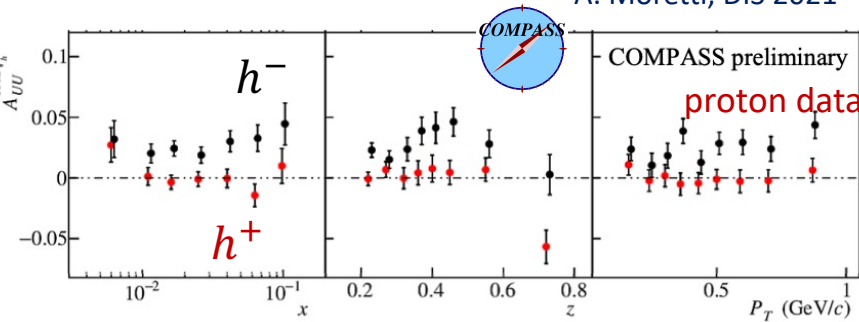
NPB 956 (2020) 115039



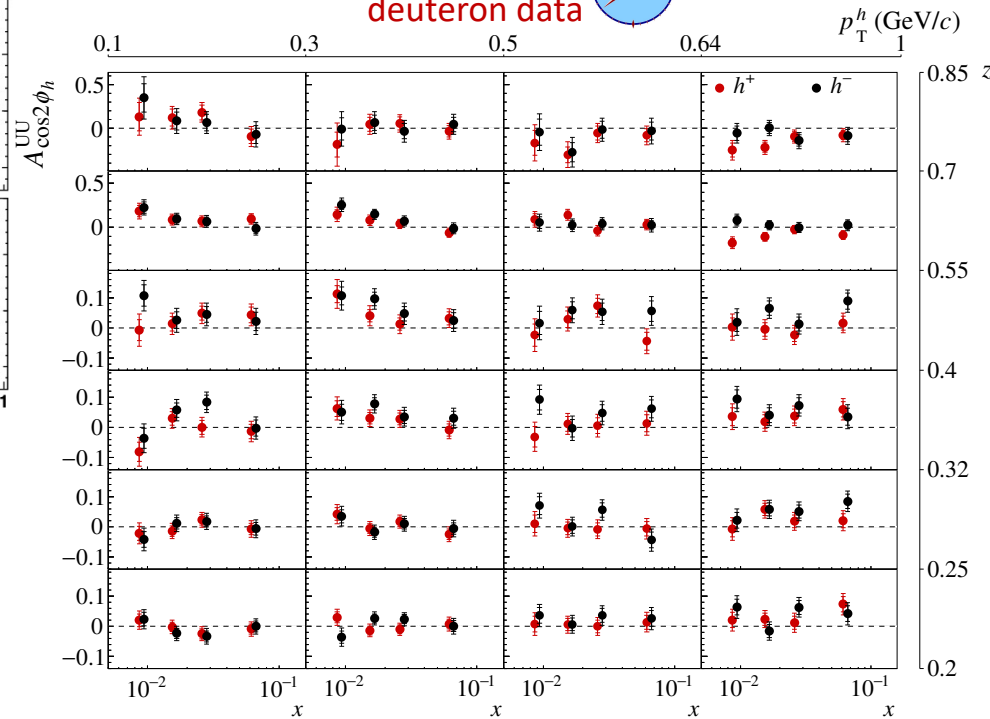
PRD 87, 012010 (2013)



A. Moretti, DIS 2021



deuteron data

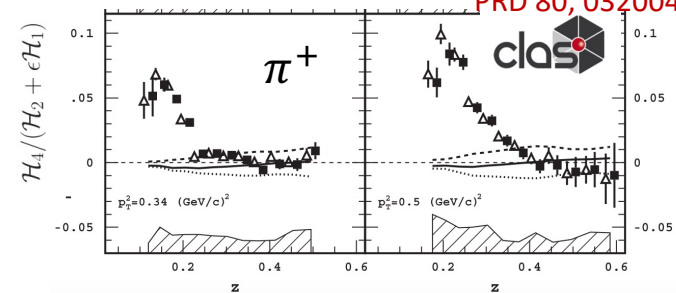


- phenomenological analyses have been performed  
e.g., Barone et al., PRD91 (2015) 7, 074019

- still no clear interpretation of these asymmetries

*Boer-Mulders function still unknown..*

PRD 80, 032004 (2009)



## Can we simulate the unpolarized azimuthal asymmetries using event generators?

Cahn effect already implemented in LEPTO → simulation of  $A_{UU}^{\cos \phi_h}$   
[A. Kotzinian, arXiv:0510359]

Boer-Mulders effect never implemented in full event generators  
(requires a model for polarized hadronization)

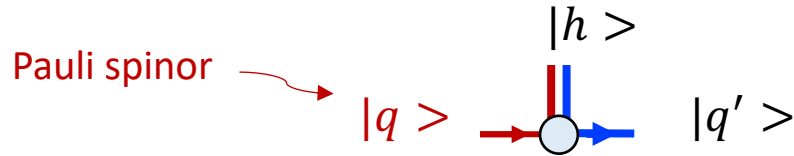
this work, simulations using :

- the  $\text{string}+^3\text{P}_0$  model of hadronization
- the **PYTHIA 8** generator

→ simulation of  $A_{UU}^{\cos 2\phi_h}$

# The string+<sup>3</sup>P<sub>0</sub> model of hadronization

- A quantum mechanical model for the **elementary splitting**



AK, X. Artru, A. Martin, *PRD104* (2021) 11, 114038  
 AK, X. Artru, Z. Belghobsi, A. Martin, *PRD* 100 (2019) 1, 014003  
 AK, X. Artru, Z. Belghobsi, F. Bradamante, A. Martin, *PRD* 97 (2018) 7, 074010

described by the  $2 \times 2$  splitting amplitude  $T$

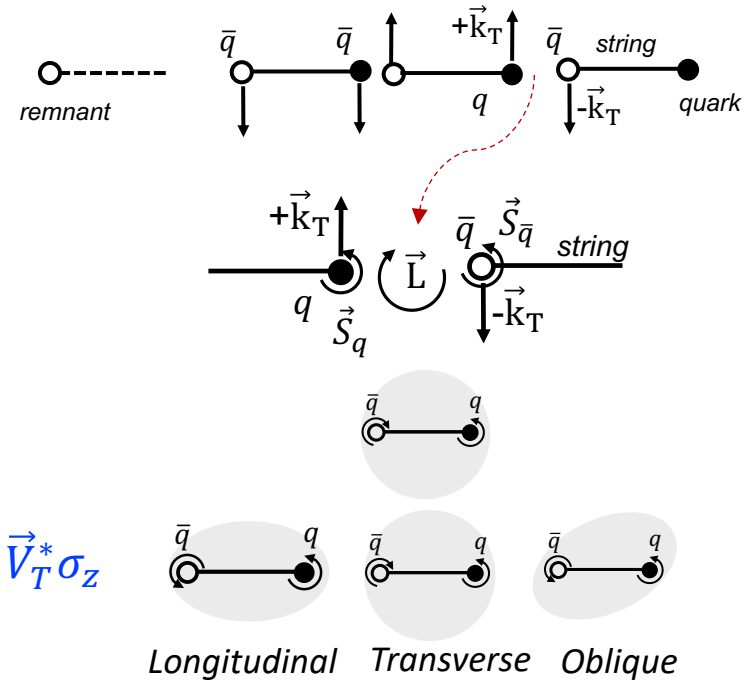
$$T = (\text{Lund Splitting Function})^{1/2} \times$$

$\times$  <sup>3</sup>P<sub>0</sub> mechanism

$$\mu + \sigma_z \vec{\sigma}_T \cdot \vec{k}'_T$$

$\times$  coupling

$$\left\{ \begin{array}{l} \text{PS meson:} \\ \text{VM with pol. } \vec{V} \end{array} \right. \quad \begin{array}{l} \sigma_z \\ G_L V_L^* \mathbf{1} + G_T \vec{\sigma}_T \cdot \vec{V}_T^* \sigma_z \end{array}$$





# The string+ $^3P_0$ model of hadronization: parameters

- A quantum mechanical model for the splitting

AK, X. Artru, A. Martin, PRD104 (2021) 11, 114038  
 u, Z. Belghobsi, A. Martin, PRD 100 (2019) 1, 014003  
 u, Z. Belghobsi, F. Bradamante, A. Martin, PRD 97  
 174010

Complex mass  $\mu$ :

$\text{Re } \mu \rightarrow$  longitudinal spin effects (jet handedness)

$\text{Im } \mu \rightarrow$  transverse spin effects (Collins, dihadron)

desc

$T = (\text{Lund Splitting Function})^{1/2} \times$

$\times$   $^3P_0$  mechanism

$$\mu + \sigma_z \vec{\sigma}_T \cdot \vec{k}'_T$$

$\times$  coupling

PS meson:

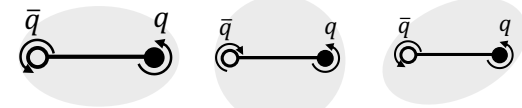
VM with pol.  $\vec{V}$

$\sigma_z$

$$G_L V_L^* \mathbf{1} + G_T \vec{\sigma}_T \cdot \vec{V}_T^* \sigma_z$$

$$f_L = \frac{|G_L/G_T|^2}{2 + |G_L/G_T|^2} \rightarrow \text{fraction of L pol. VMs}$$

$$\theta_{LT} = \arg\left(\frac{G_L}{G_T}\right) \rightarrow \text{oblique polarization}$$



Longitudinal Transverse Oblique

# The interface with PYTHIA 8: StringSpinner

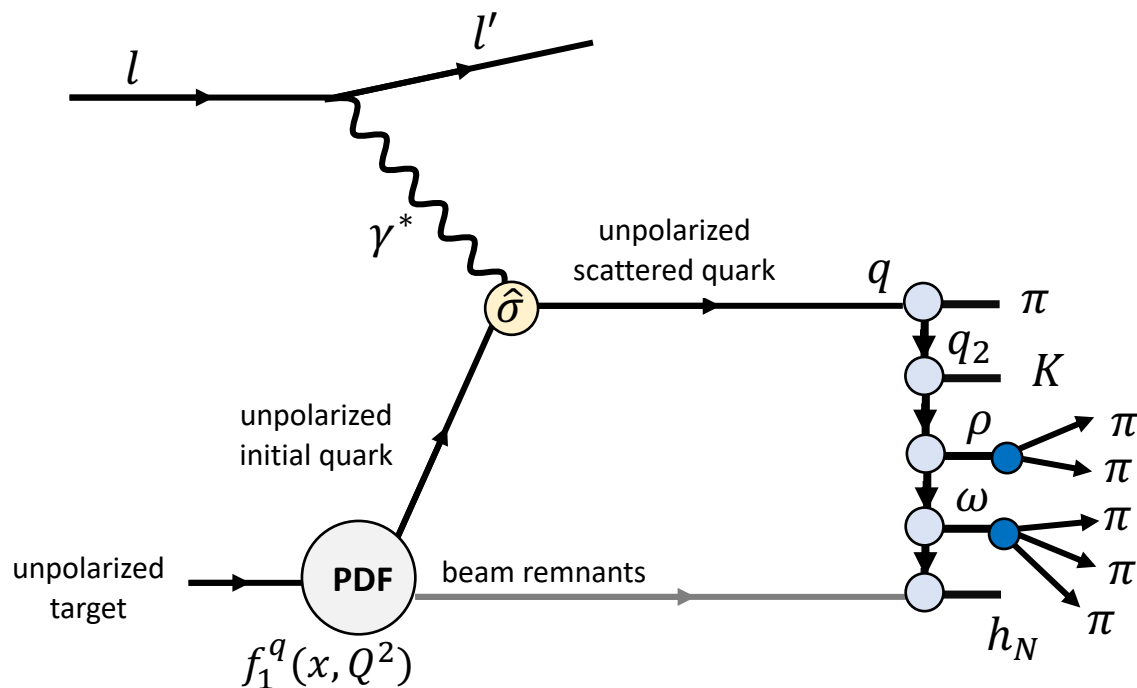
String+ $^3P_0$  interfaced to PYTHIA 8 for DIS as an external package → **StringSpinner**

public version → only PS meson production

here PS and VM production

[AK, L. Lönnblad, CPC 272 (2022) 108234]

[AK, talk at ICHEP-2022]



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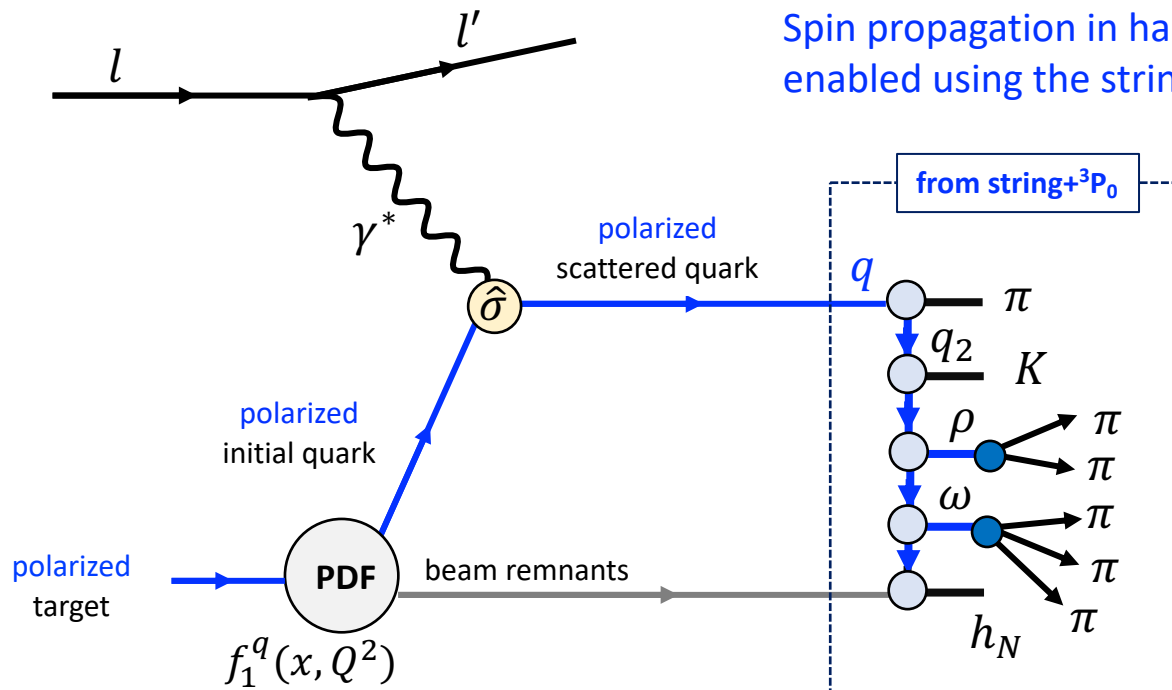
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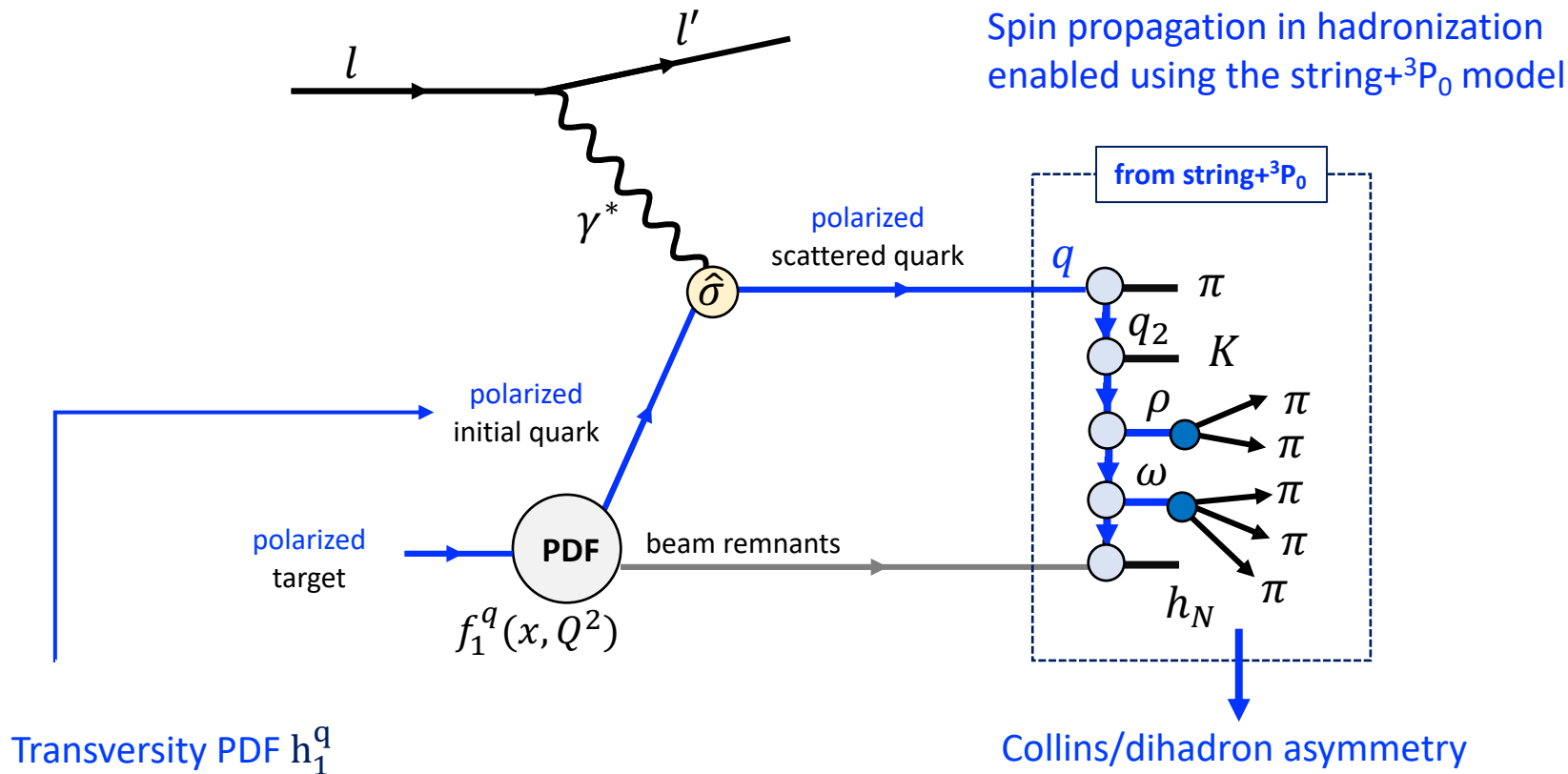
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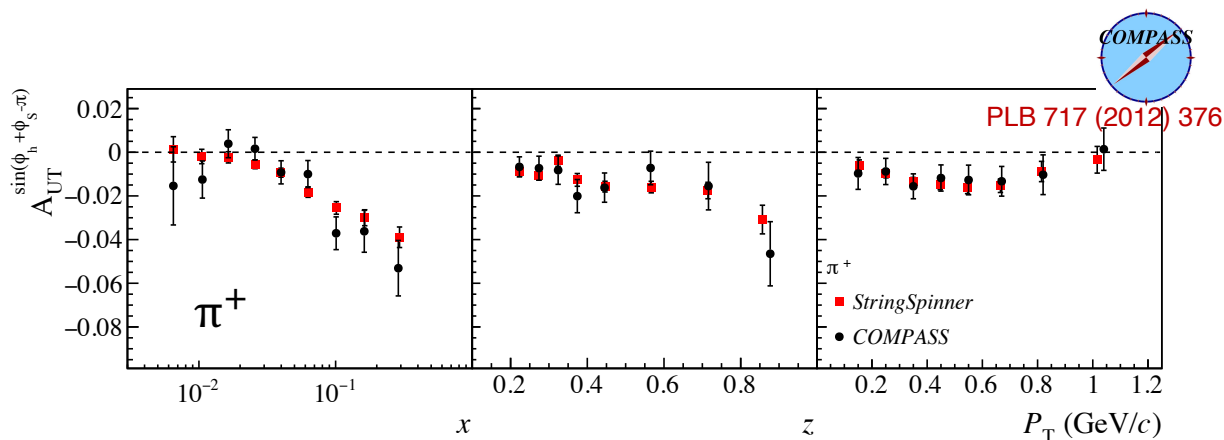
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[AK, talk at ICHEP-2022]

## Example: Comparison with Collins asymmetries in SIDIS

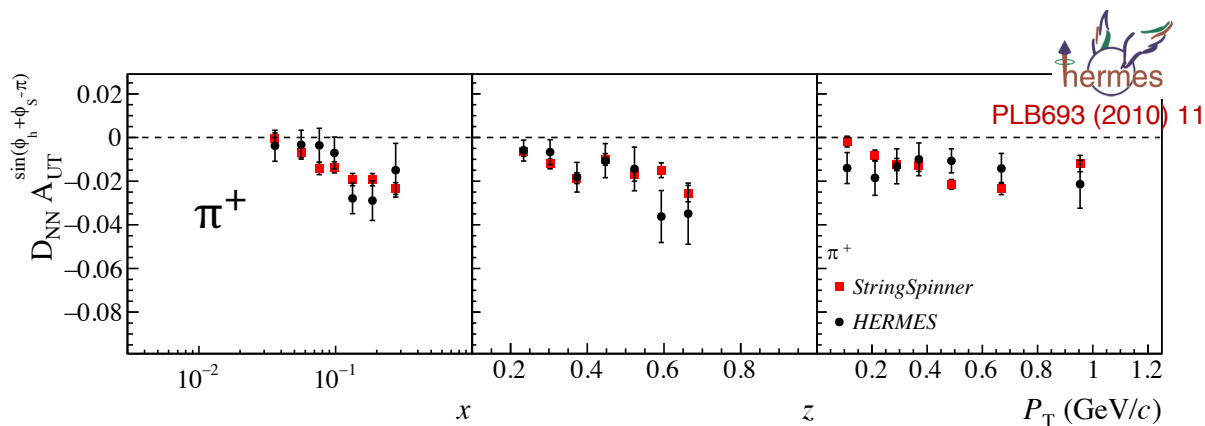


Parameters for spin effects

$$\mu = (0.42 + i 0.76) \text{ GeV}/c^2$$

$$f_L = 0.93, \theta_{LT} = 0$$

Satisfactory description  
also for  $\pi^-$



Other sets of parameters  
also possible

# The interface with PYTHIA 8: StringSpinner

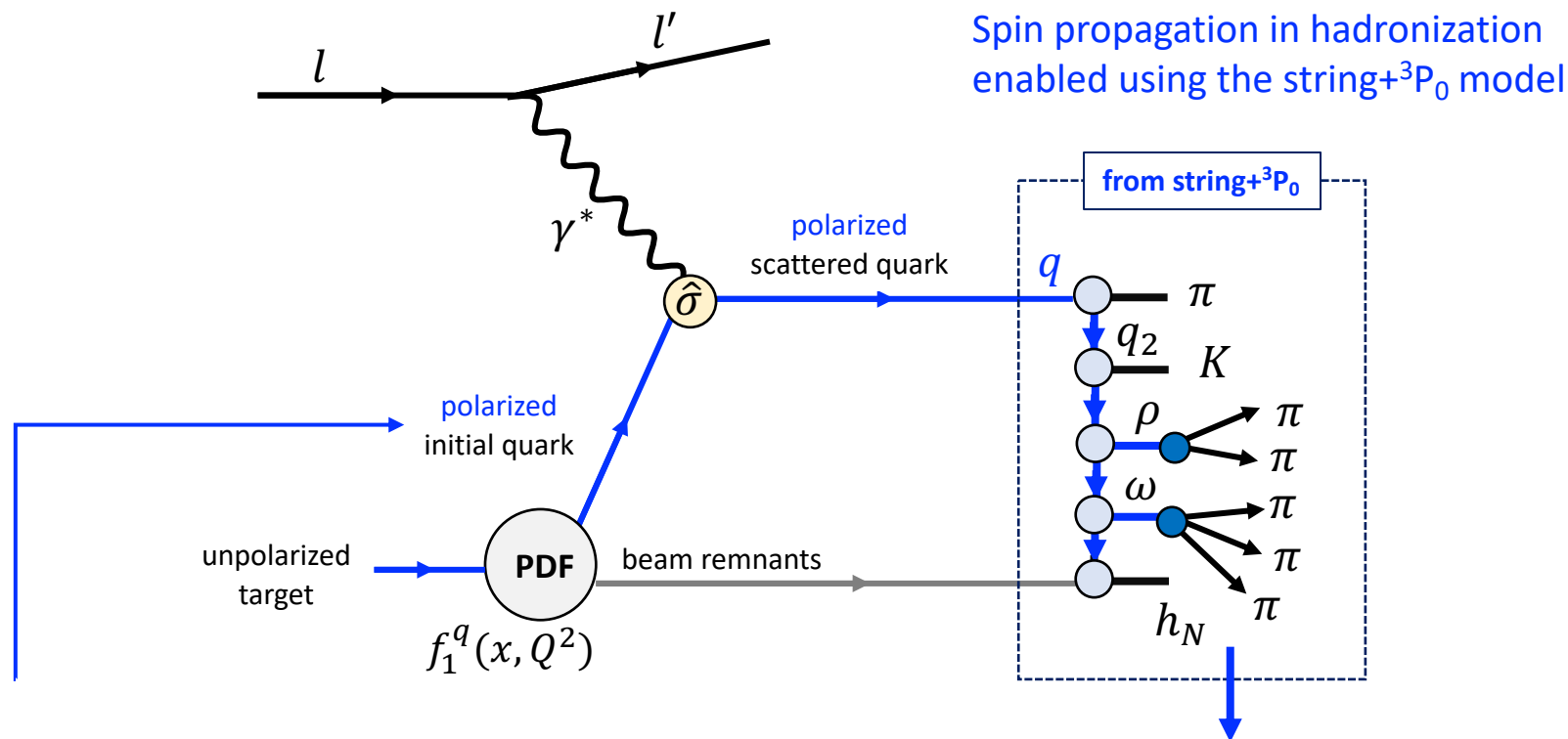
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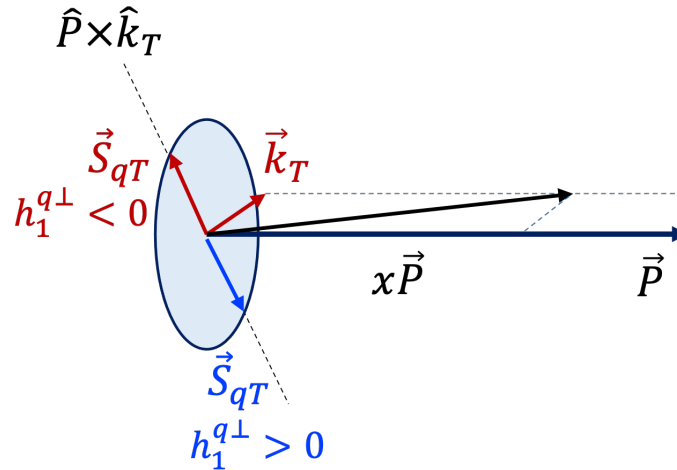
Transversity PDF  $h_1^q$

Boer-Mulders TMD PDF  $h_1^{\perp q}$  (new)

Collins/dihadron asymmetry

$A_{UU}^{\cos 2\phi_h}$  asymmetry (new)

## Recall on the Boer-Mulders effect



Quark transverse polarization due to the Boer-Mulders function

[Boer,Mulders, PRD57 (1998) 5780]

$$\vec{S}_{qT} = \frac{k_T}{M_N} \frac{h_1^{\perp q}}{f_1^q} (-\hat{P} \times \hat{k}_T)$$

Positivity condition, assuming  $h_{1L}^{\perp} = 0$

[Bacchetta et al, PRL 85 (2000) 712-715]

$$\frac{k_T}{M_N} \frac{|h_1^{\perp q}|}{f_1^q} \leq 1$$

# Inclusion of the Boer-Mulders effect in StringSpinner

Consider fully polarized quarks, i.e.

$$\vec{S}_{qT} = \text{sign}(h_1^{\perp q}) \hat{q} \times \hat{k}_T$$

saturation of positivity condition,  
only the sign of  $h_1^{\perp q}$  is relevant

Scattered quark depolarized and reflected

$$\vec{S}'_{qT} = D_{NN} \times (\vec{S}_{qT} - 2 \vec{S}_{qT} \cdot \hat{l}_T)$$

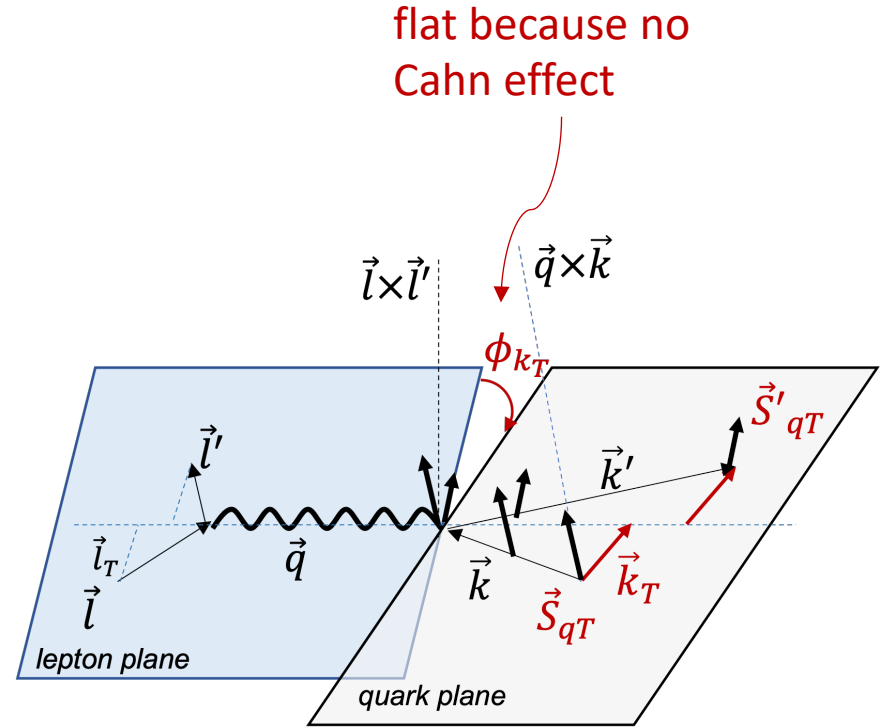
Intrinsic  $\vec{k}_T$  generated from Pythia according to the distribution

$$dk_T^2 e^{-k_T^2/\langle k_T^2 \rangle} \times d\phi_{k_T}/2\pi$$

We take

$$\langle k_T^2 \rangle = 0.1 \left( \frac{\text{GeV}}{c} \right)^2$$

no flavor dependence  
no kinematic dependence



Lepton-quark hard scattering in the GNS.



(selection of) Preliminary results on  $A_{UU}^{\cos 2\phi_h}$

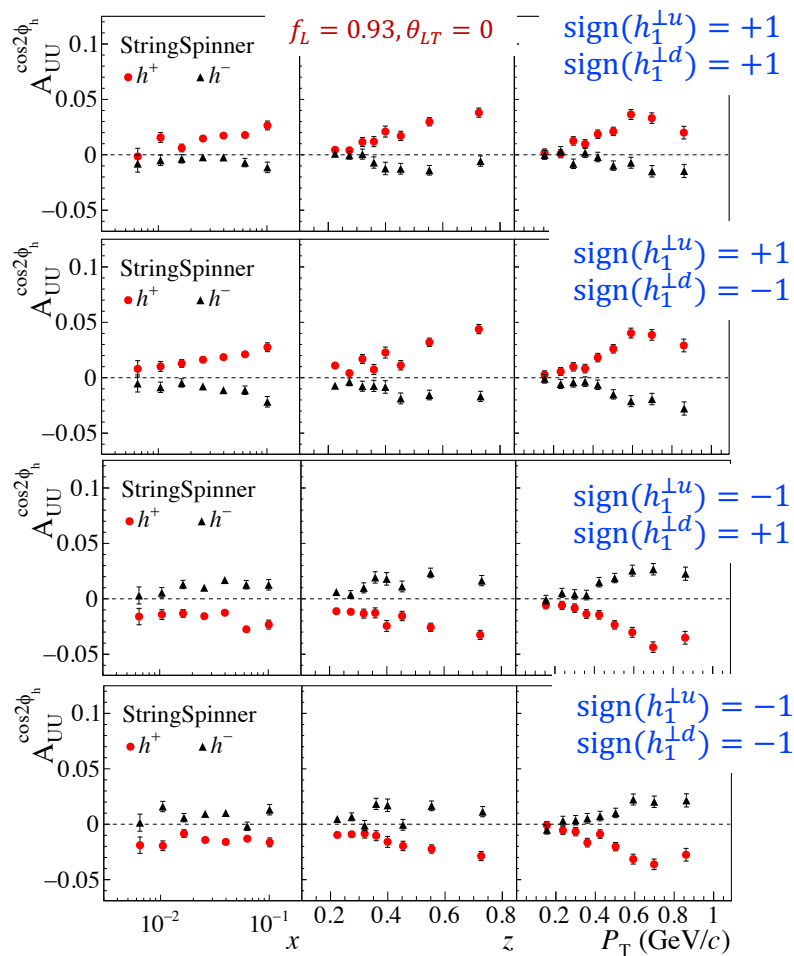
simulations of SIDIS off unpolarized protons in the COMPASS kinematics

# Contribution of the Boer-Mulders TMD to $A_{UU}^{\cos 2\phi_h}$

DIS in the COMPASS kinematics with a proton target

valence  $u$  and  $d$  quarks polarized

sea quarks unpolarized



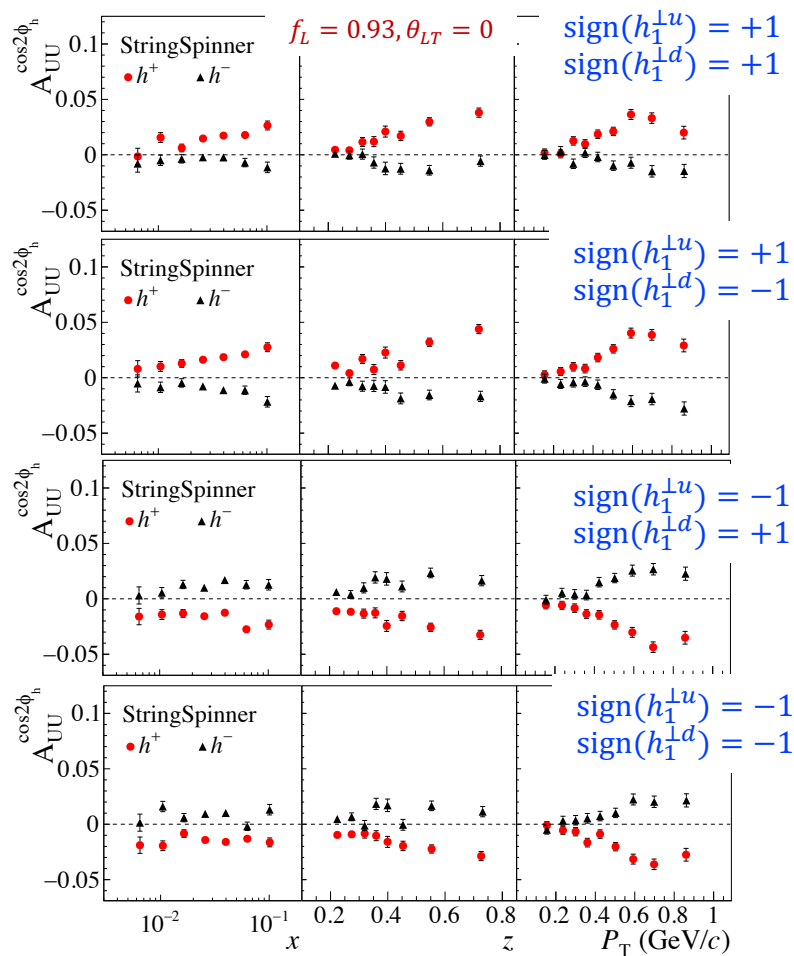
kinematic cuts as in data

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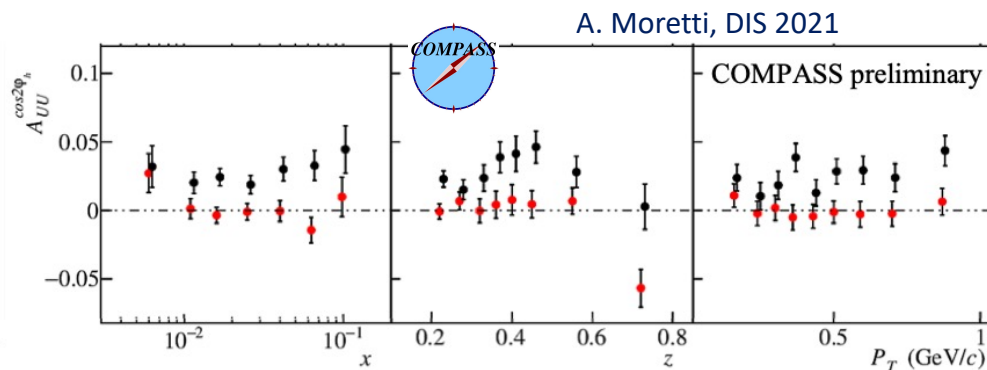
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sea quarks unpolarized



kinematic cuts as in data



difficult to describe the data taking into account only valence quarks

→ compensate by choosing appropriately the sign of the Boer-Mulders TMD for sea quarks?

Note:

- same sign of the asymmetry for leading hadrons as the Boer-Mulders TMD
- smaller effect for subleading hadrons

# Contribution of the Boer-Mulders TMD to $A_{UU}^{\cos 2\phi_h}$

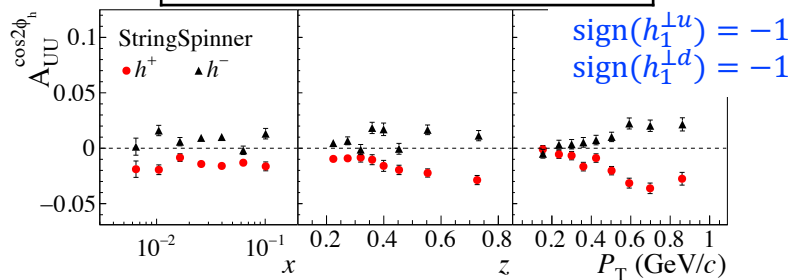
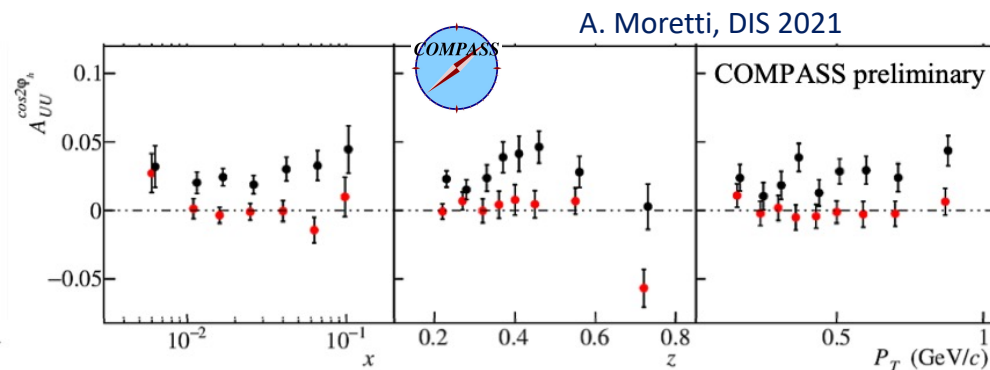
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valence  $u$  and  $d$  quarks polarized

$u/d$  sea quarks polarized

$\pi^+$	$A_{UU}^{\cos 2\phi_h}$	$\text{sign}(h_1^{\perp q})$
Data	$\sim 0$	
$u^{val}$	-	-
$\bar{d}$	+	+
$u^{sea}$	+	+

$\pi^-$	$A_{UU}^{\cos 2\phi_h}$	$\text{sign}(h_1^{\perp q})$
Data	+	
$u^{val}$	$\sim 0$	-
$\bar{u}$	+	+
$d^{sea}$	+	+



kinematic cuts as in data

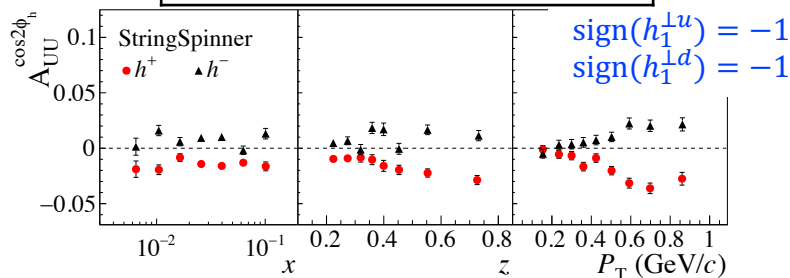
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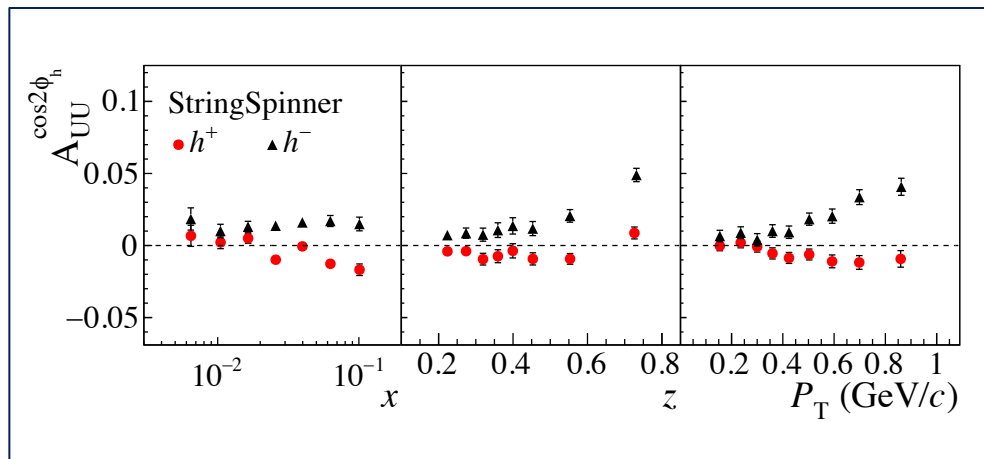
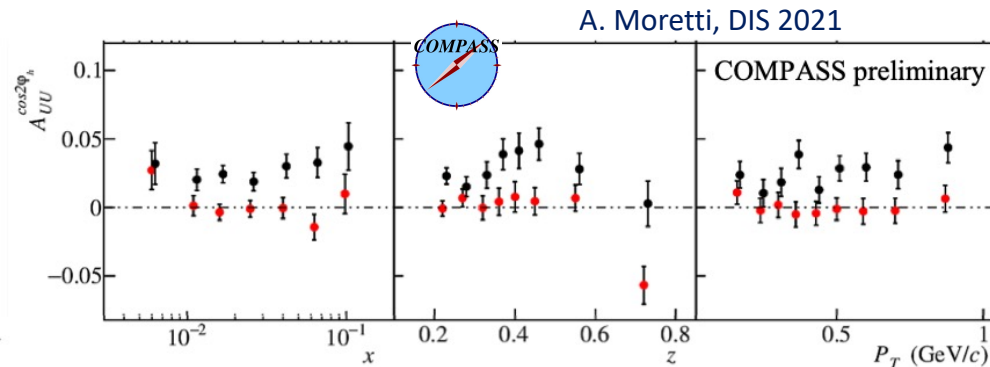
valence  $u$  and  $d$  quarks polarized,  
 $u/d$  sea quarks polarized

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$u^{val}$	-	-
$\bar{d}$	+	+
$u^{sea}$	+	+

$\pi^-$	$A_{UU}^{\cos 2\phi_h}$	$\text{sign}(h_1^{\perp q})$
Data	+	
$u^{val}$	$\sim 0$	-
$\bar{u}$	+	+
$d^{sea}$	+	+



kinematic cuts as in data



- Better comparison with data
- still non vanishing asymmetry for  $h^+$  (large  $x_B$ )
- different trends at large  $z_h$

## Conclusions

- We have introduced spin effects in PYTHIA 8 hadronization for SIDIS with PS and VM production by using the *string+<sup>3</sup>P<sub>0</sub>* model → *StringSpinner*  
*the most recent version to be published soon*
- Good description of available data on TSA
- Introduction of *spin effects for unpolarised SIDIS in StringSpinner* ongoing  
Boer-Mulders effect included → preliminary results on  $A_{UU}^{\cos 2\phi_h}$   
next step, addition of the Cahn effect

# Backup

# Relevant free parameters for string fragmentation

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see Kerbizi, Artru, Martin, PRD104 (2021) 11, 114038

## Pythia parameters:

StringZ:aLund	0.90
StringZ:bLund	0.50 (GeV/c <sup>2</sup> ) <sup>-2</sup>
StringPT:sigma	0.37 GeV/c
StringPT:enhancedFraction	0.00
StringPT:enhancedWidth	0.00 GeV/c
BeamRemnants:primordialKTremnant	0.00
BeamRemnants:halfScaleForKT	0.00
BeamRemnants:halfMassForKT	0.00
BeamRemnants:primordialKThard	0.50

## String+<sup>3</sup>P<sub>0</sub> parameters

Re( $\mu$ )	0.42 GeV/c <sup>2</sup>
Im( $\mu$ )	0.76 GeV/c <sup>2</sup>
$f_L$	0.93
$\theta_{LT}$	0.00



# Comparison with Collins asymmetries in SIDIS for $\pi^-$

Parameters for spin effects

$$\mu = (0.42 + i 0.76) \text{ GeV}/c^2$$

$$f_L = 0.93, \theta_{LT} = 0$$

