

Hard Exclusive Reactions at COMPASS at CERN

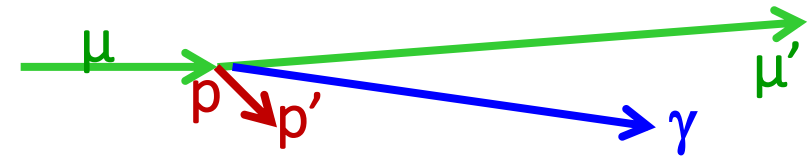
Exclusive photon (DVCS) and meson (HEMP) production at small transfer for GPD studies



Nicole d'Hose – CEA – Université Paris-Saclay
for the COMPASS Collaboration



$$\text{DVCS} : \mu p \rightarrow \mu' p' \gamma$$



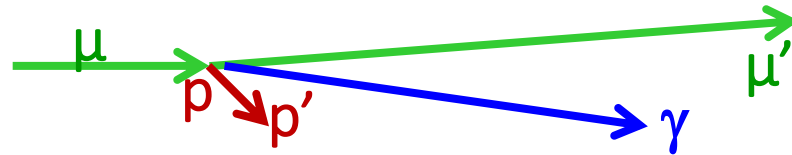
$$\text{Pseudo-Scalar Meson} : \mu p \rightarrow \mu' p' \pi^0$$

Markéta Pešková's talk today at 12:05

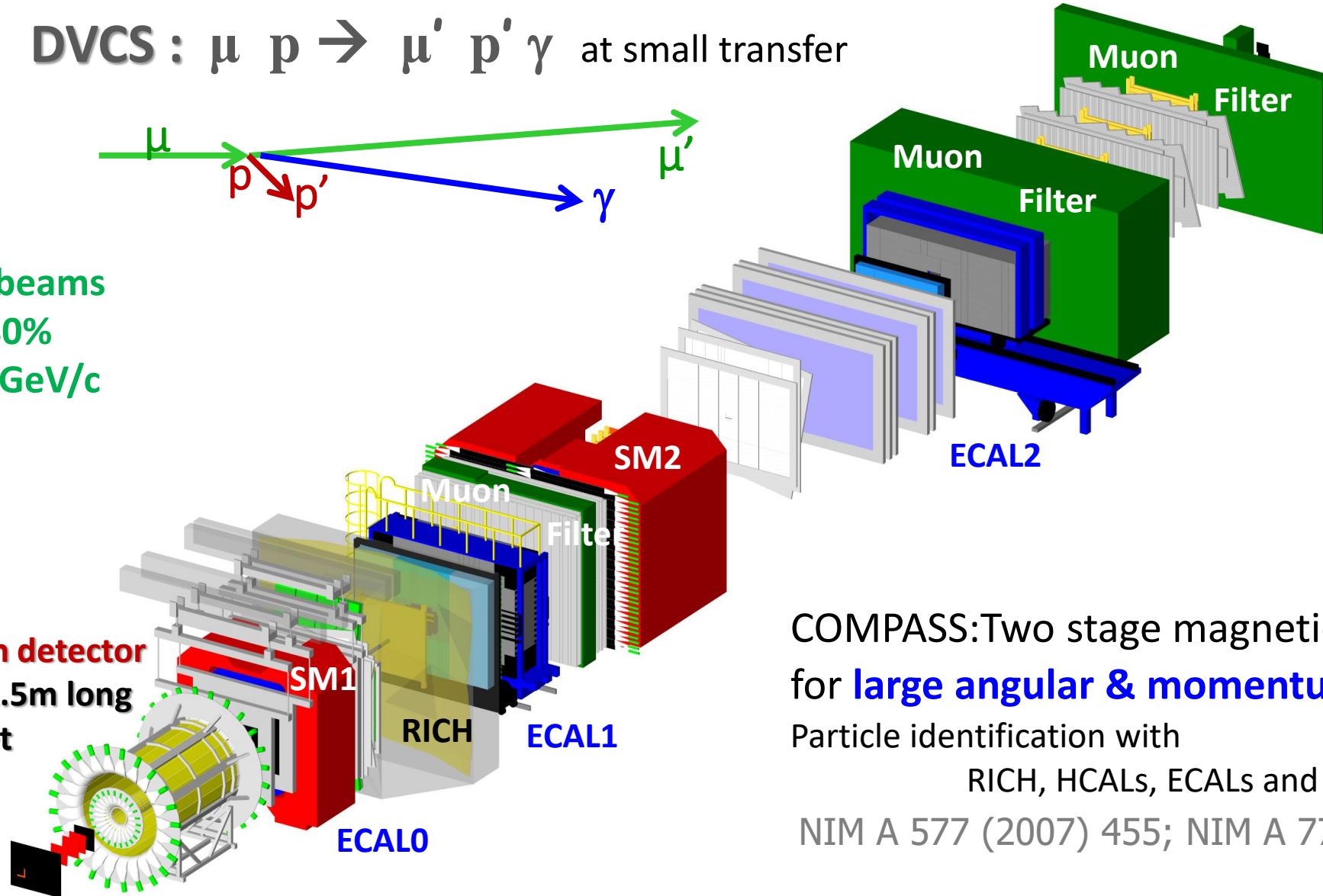
$$\text{Vector Meson} : \mu p \rightarrow \mu' p' \rho \text{ or } \omega \text{ or } \dots$$

Measurement of exclusive cross sections at COMPASS

DVCS : $\mu p \rightarrow \mu' p' \gamma$ at small transfer



Both μ^+ and μ^- beams
Polarisation $\sim \pm 80\%$
Momentum 160 GeV/c

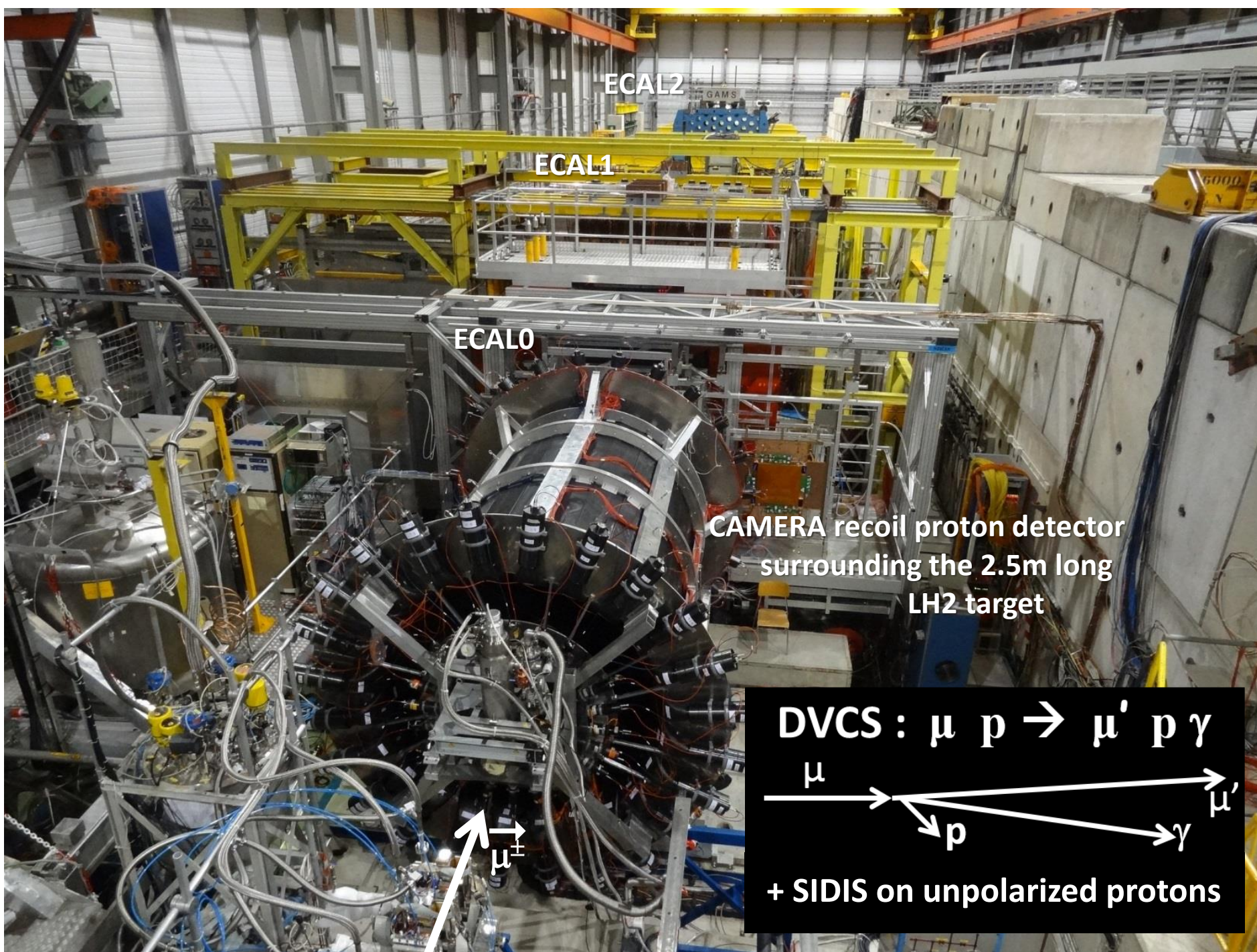


COMPASS: Two stage magnetic spectrometer for **large angular & momentum acceptance**

Particle identification with

RICH, HCALs, ECALs and muon filters

NIM A 577 (2007) 455; NIM A 779 (2015) 69



ECAL2

ECAL1

ECAL0

CAMERA recoil proton detector
surrounding the 2.5m long
LH2 target

DVCS : $\mu p \rightarrow \mu' p \gamma$

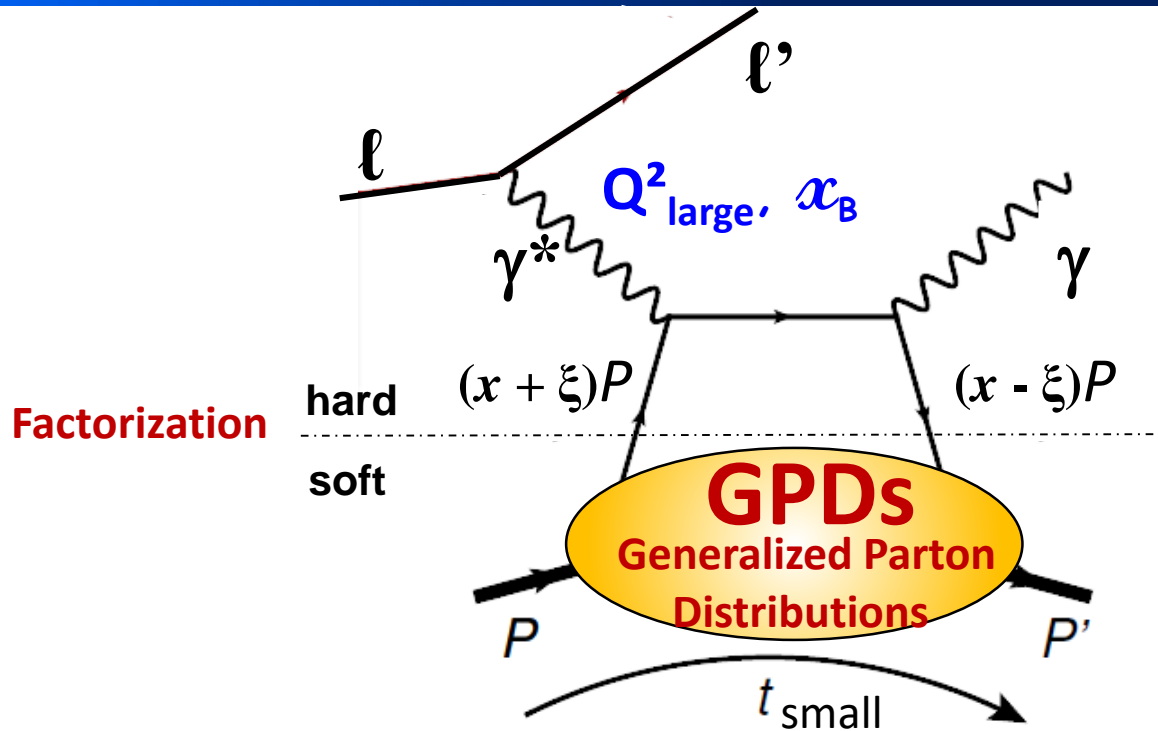


+ SIDIS on unpolarized protons

2012:
1 month pilot run

2016 -17:
2 x 6 month
data taking

Deeply virtual Compton scattering (DVCS)



D. Mueller *et al*, Fortsch. Phys. 42 (1994)

X.D. Ji, PRL 78 (1997), PRD 55 (1997)

A. V. Radyushkin, PLB 385 (1996), PRD 56 (1997)

DVCS: $\ell p \rightarrow \ell' p' \gamma$

the golden channel

because it interferes with
the Bethe-Heitler process

also meson production

$\ell p \rightarrow \ell' p' \pi, \rho, \omega$ or ϕ or $J/\psi \dots$

The GPDs depend on the following variables:

x : average
 ξ : transferred } quark longitudinal
momentum fraction

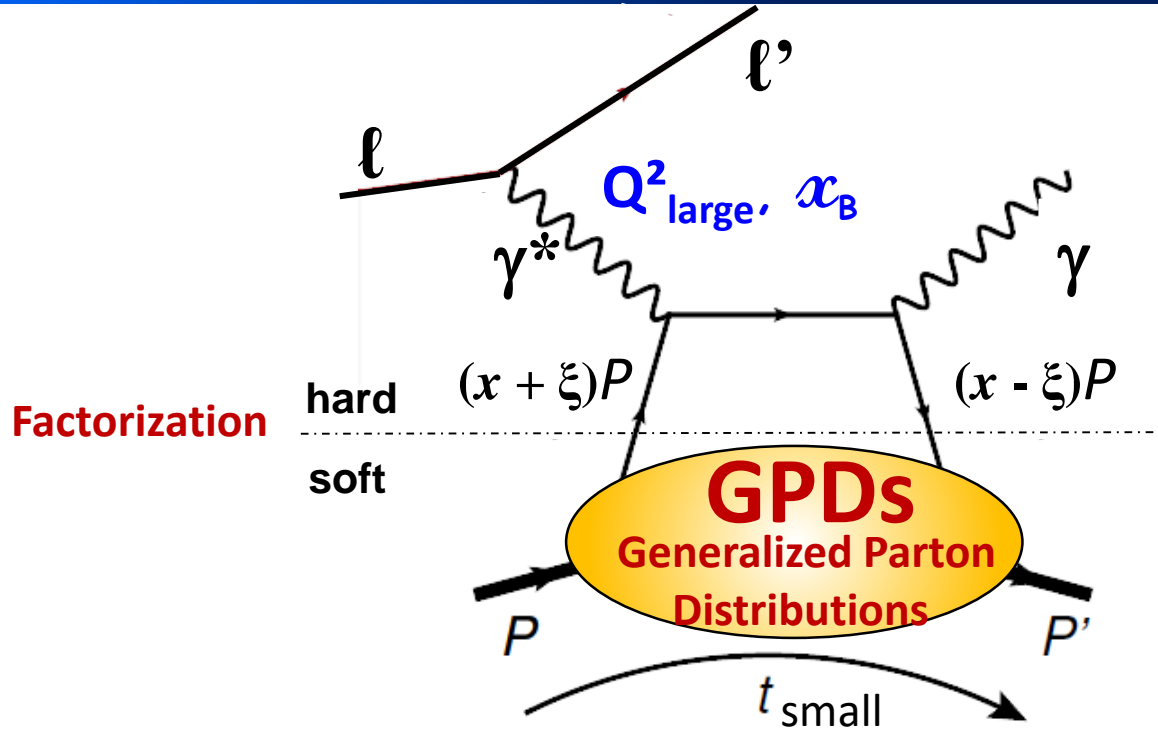
t : proton momentum transfer squared
related to b_{\perp} via Fourier transform

The variables measured in the experiment:

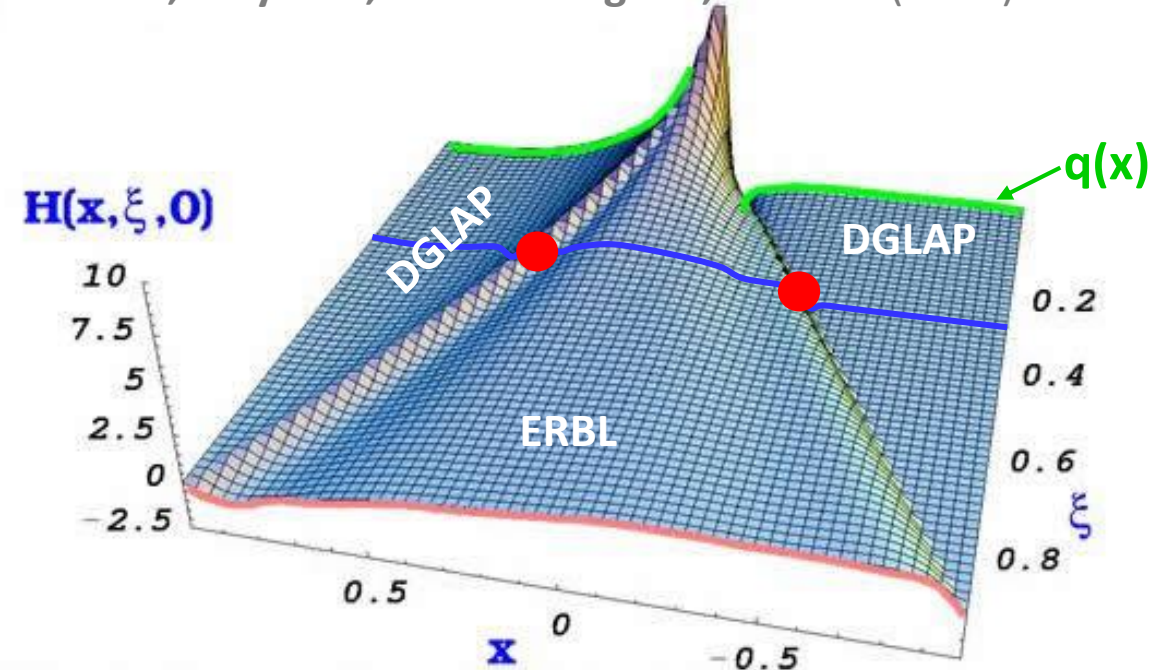
$E_{\ell}, Q^2, x_B \sim 2\xi / (1 + \xi),$

t (or $\theta_{\gamma^* \gamma}$) and ϕ ($\ell \ell'$ plane / $\gamma \gamma^*$ plane)

Deeply virtual Compton scattering (DVCS)



Goeke, Polyakov, Vanderhaeghen, PPNP47 (2001)



The amplitude DVCS at LT & LO in α_s (GPD \mathcal{H}):

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i\pi H(x = \pm \xi, \xi, t)$$

Real part Imaginary part

In an experiment we measure
Compton Form Factor \mathcal{H}

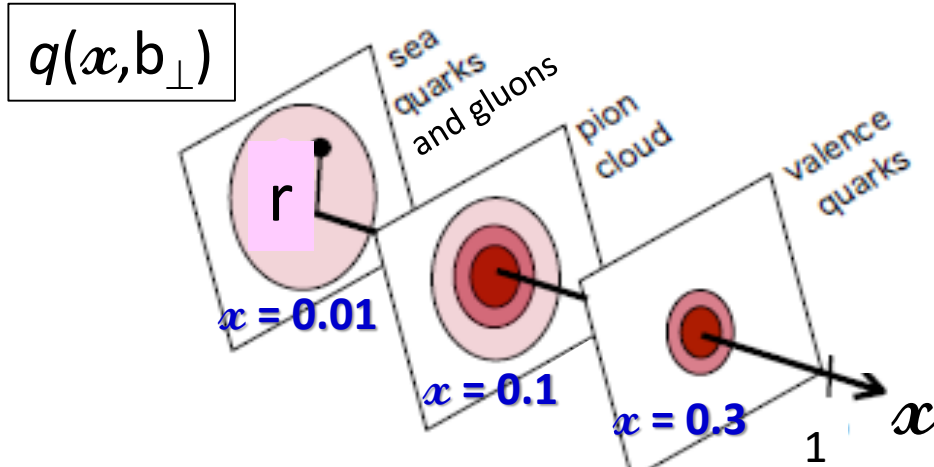
$$\text{Re}\mathcal{H}(\xi, t) = \pi^{-1} \int_0^1 dx \frac{2x \text{Im}\mathcal{H}(x, t)}{x^2 - \xi^2} + \Delta(t)$$

Deeply virtual Compton scattering (DVCS)

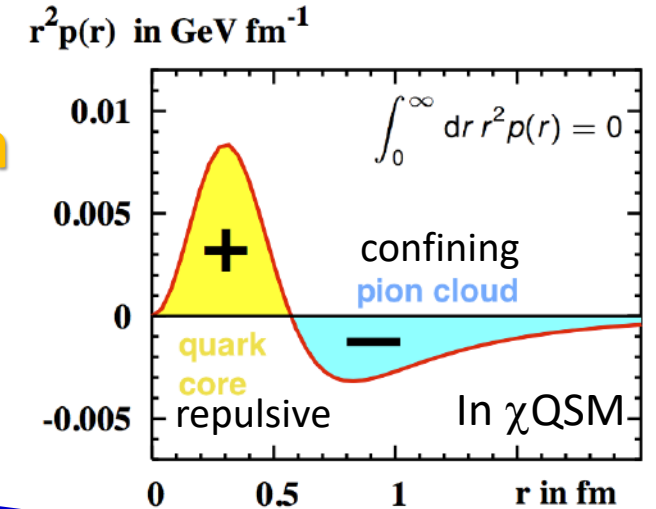
M. Burkardt, PRD66(2002)

M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)

Mapping in the transverse plane



Pressure Distribution



FT of $H(x, \xi=0, t)$

The amplitude DVCS at LT & LO in α_s (GPD \mathcal{H}):

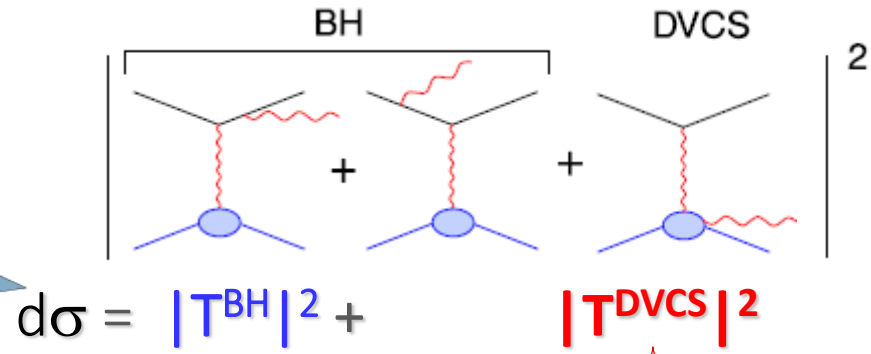
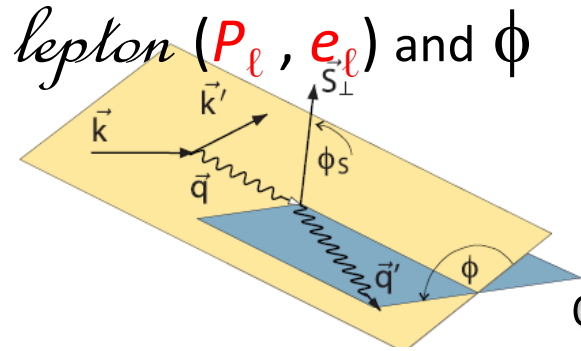
$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i \pi H(x = \pm \xi, \xi, t)$$

In an experiment we measure
Compton Form Factor \mathcal{H}

$$\text{Re}\mathcal{H}(\xi, t) = \pi^{-1} \int_0^1 dx \frac{2x \text{Im}\mathcal{H}(x, t)}{x^2 - \xi^2} + \Delta(t)$$

$d_1(t)$
D-term

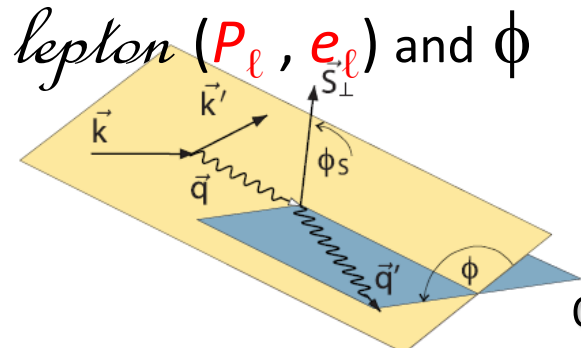
Deeply virtual Compton scattering (DVCS)



$$d\sigma = |T^{BH}|^2 + |T^{DVCS}|^2 + \text{Interference Term}$$

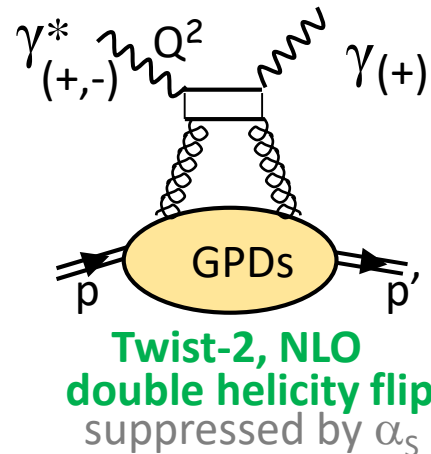
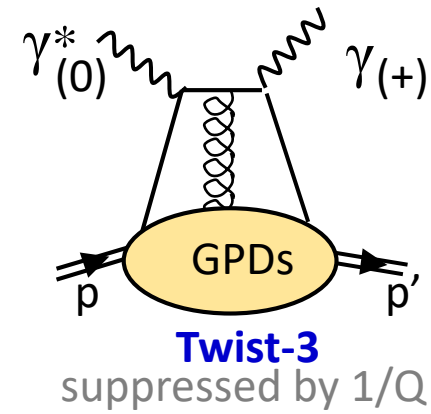
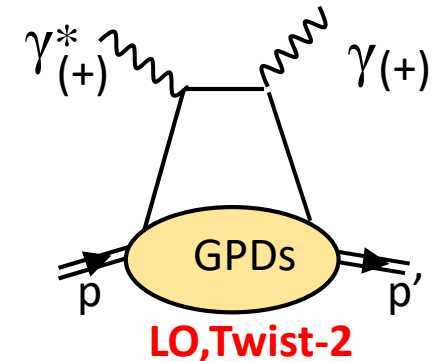
$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Deeply virtual Compton scattering (DVCS)



$$d\sigma = |T^{BH}|^2 + |T^{DVCS}|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



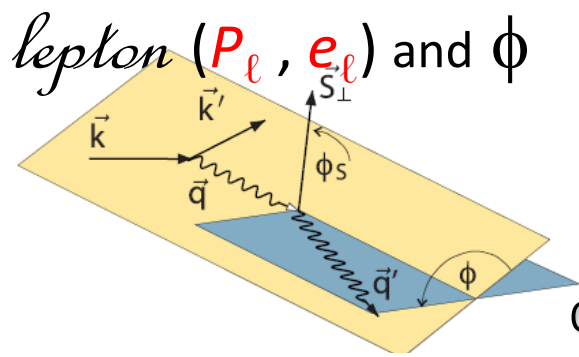
With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

α_s suppressed

Deeply virtual Compton scattering (DVCS)



$$d\sigma = \left| T^{BH} \right|^2 + \left| T^{DVCS} \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

With both μ^+ and μ^- beams we can build:

① beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$$

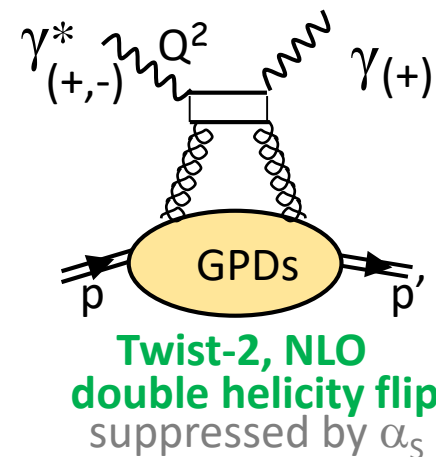
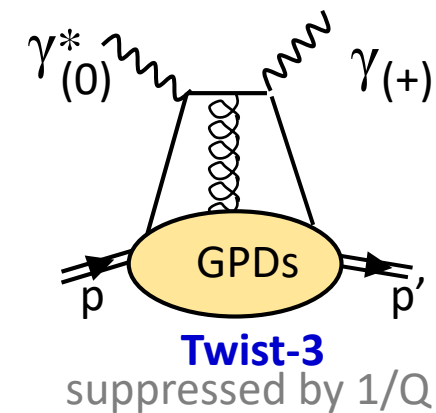
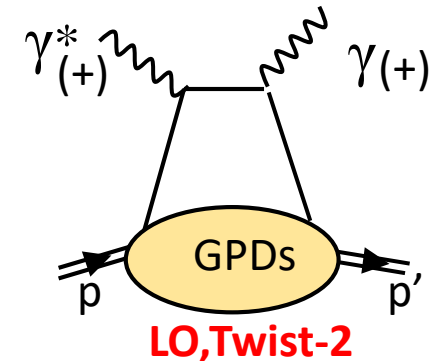
$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

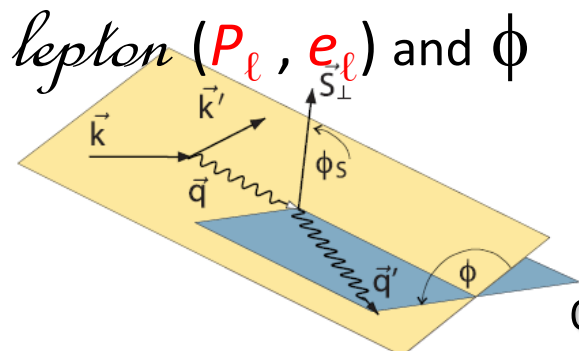
$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

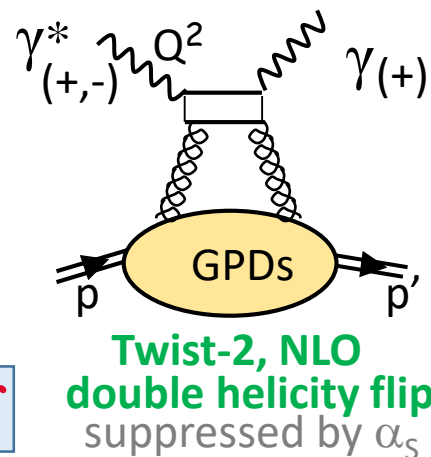
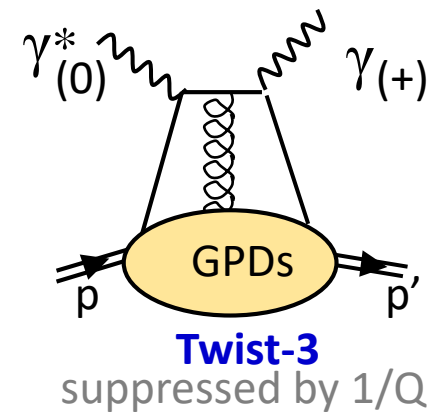
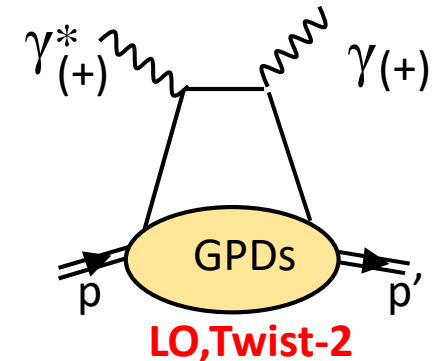


Deeply virtual Compton scattering (DVCS)



$$d\sigma = |T^{BH}|^2 + |T^{DVCS}|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With both μ^+ and μ^- beams we can build:

① beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$$

② difference

$$\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

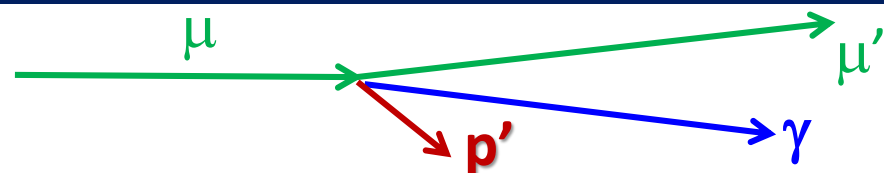
$$s_1^I \propto \text{Im } \mathcal{F}$$

$$c_1^I \propto \text{Re } \mathcal{F}$$

$$\mathcal{F} = F_1 \mathcal{H} + \xi (F_1 + F_2) \mathcal{H} - t/4m^2 F_2 \mathcal{E} \xrightarrow{\text{for proton}} F_1 \mathcal{H}$$

COMPASS 2016 data Selection of exclusive single photon production

Comparison between the observables given by the spectro or by CAMERA



DVCS: $\mu p \rightarrow \mu' p \gamma$

1) $\Delta\varphi = \varphi^{\text{cam}} - \varphi^{\text{spec}}$

2) $\Delta p_T = p_T^{\text{cam}} - p_T^{\text{spec}}$

3) $\Delta z_A = z_A^{\text{cam}} - z_A^{\text{inter}}$ and vertex

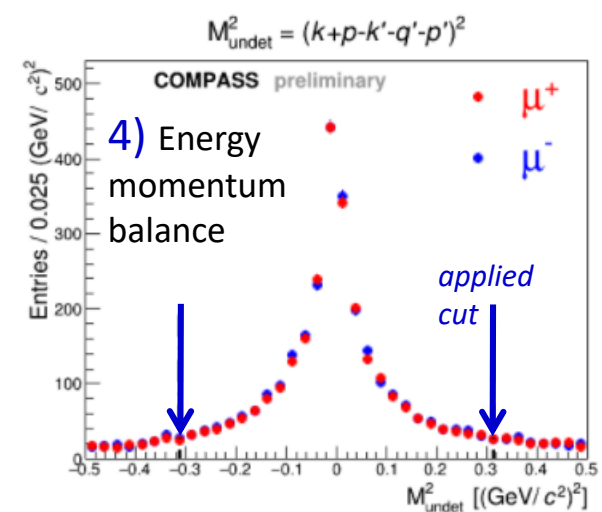
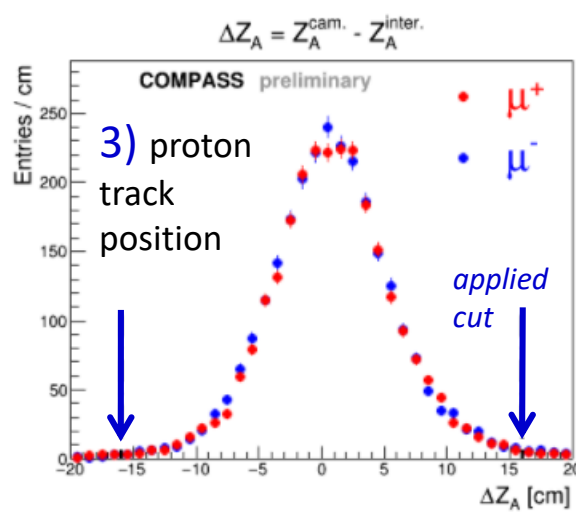
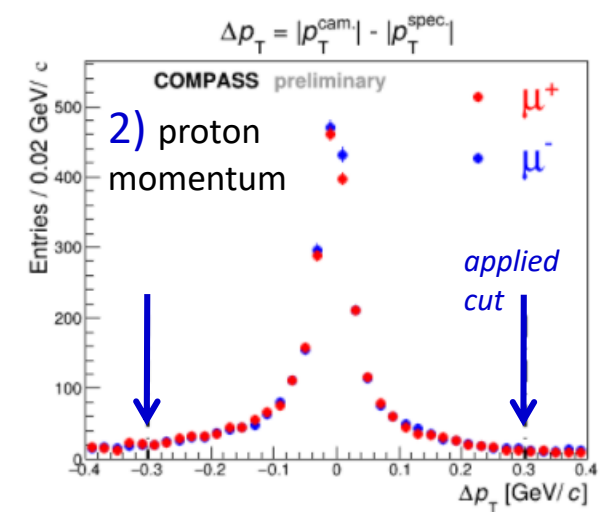
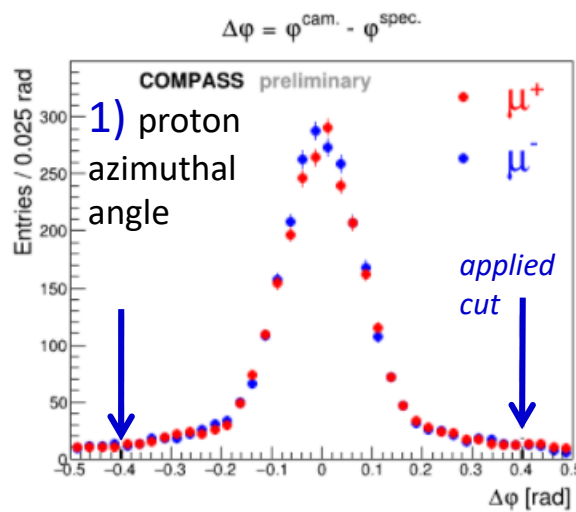
4) $M_{X=0}^2 = (p_{\mu_{\text{in}}} + p_{p_{\text{in}}} - p_{\mu_{\text{out}}} - p_{p_{\text{out}}} - p_{\gamma})^2$

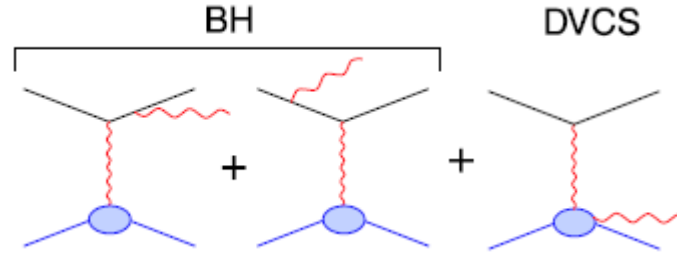
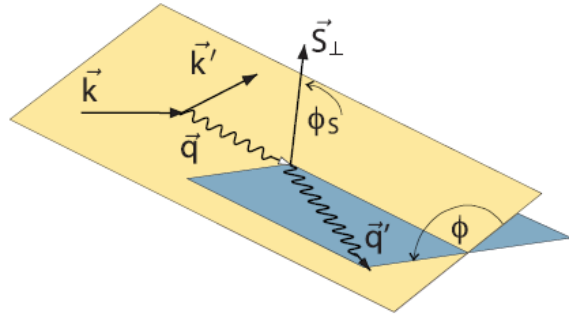
Good agreement between $\vec{\mu}^+$ and $\vec{\mu}^-$ yields

Important achievement for:

① $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$ **Easier, done first**

② $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$ **Challenging, but promising**





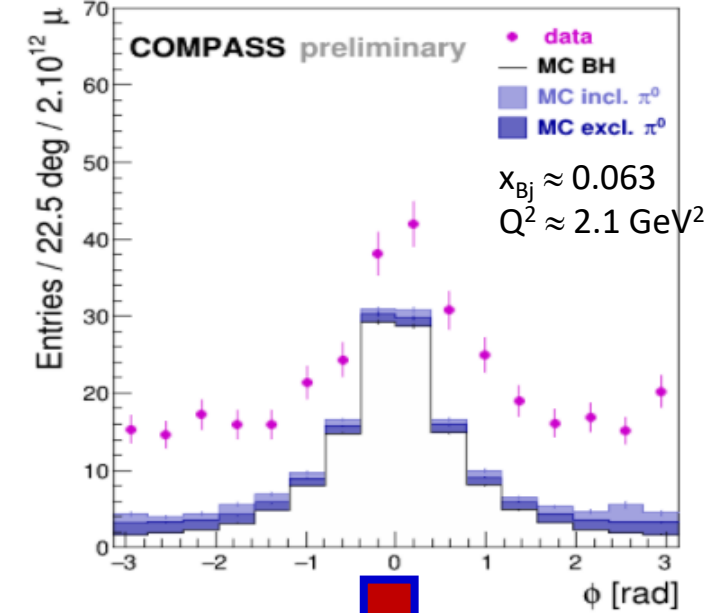
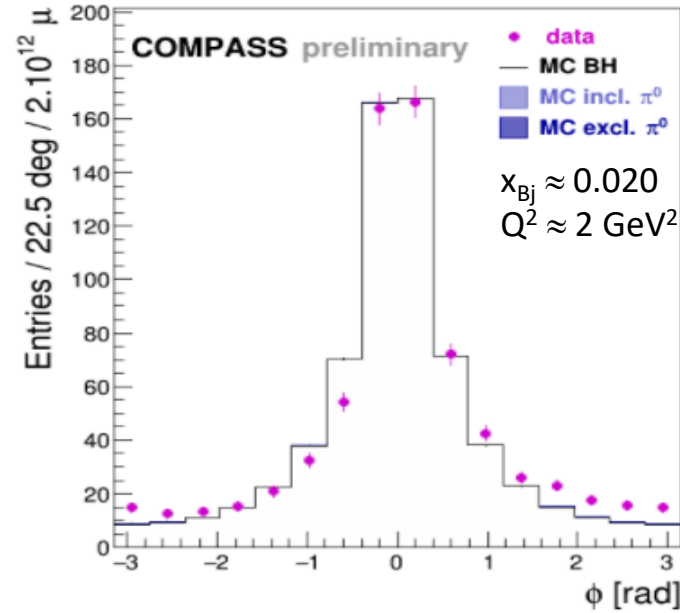
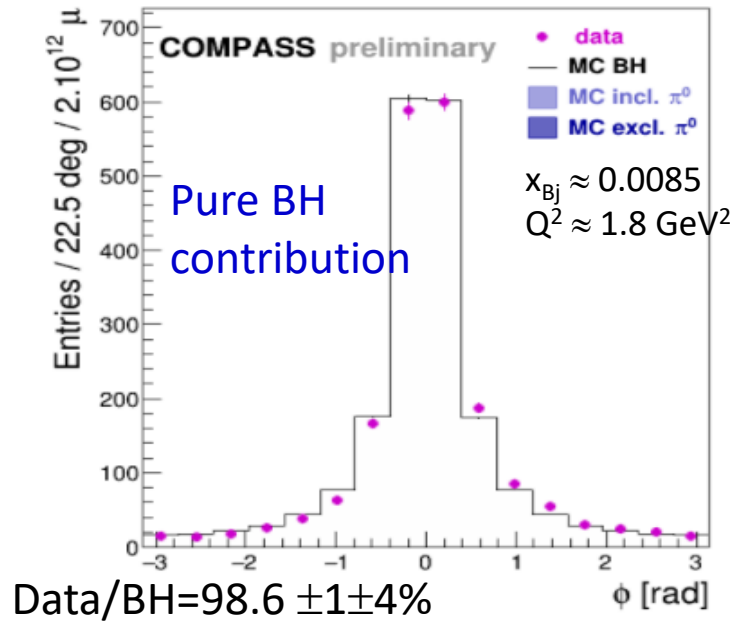
$$\Sigma = d\sigma(\mu^+) + d\sigma(\mu^-)$$

$$d\sigma \propto |T^{BH}|^2 + \text{Interference Term} + |T^{DVCS}|^2$$

$80 < v \text{ [GeV]} < 144$

$32 < v \text{ [GeV]} < 80$

$10 < v \text{ [GeV]} < 32$



DVCS above the **BH** contrib.

MC: BH contribution evaluated for the integrated luminosity
 π^0 background contribution from SIDIS (LEPTO) + exclusive production (HEPGEN)

At COMPASS using polarized positive and negative muon beams:

$$S_{CS,U} \equiv d\sigma^{\leftarrow +} + d\sigma^{\rightarrow -} = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

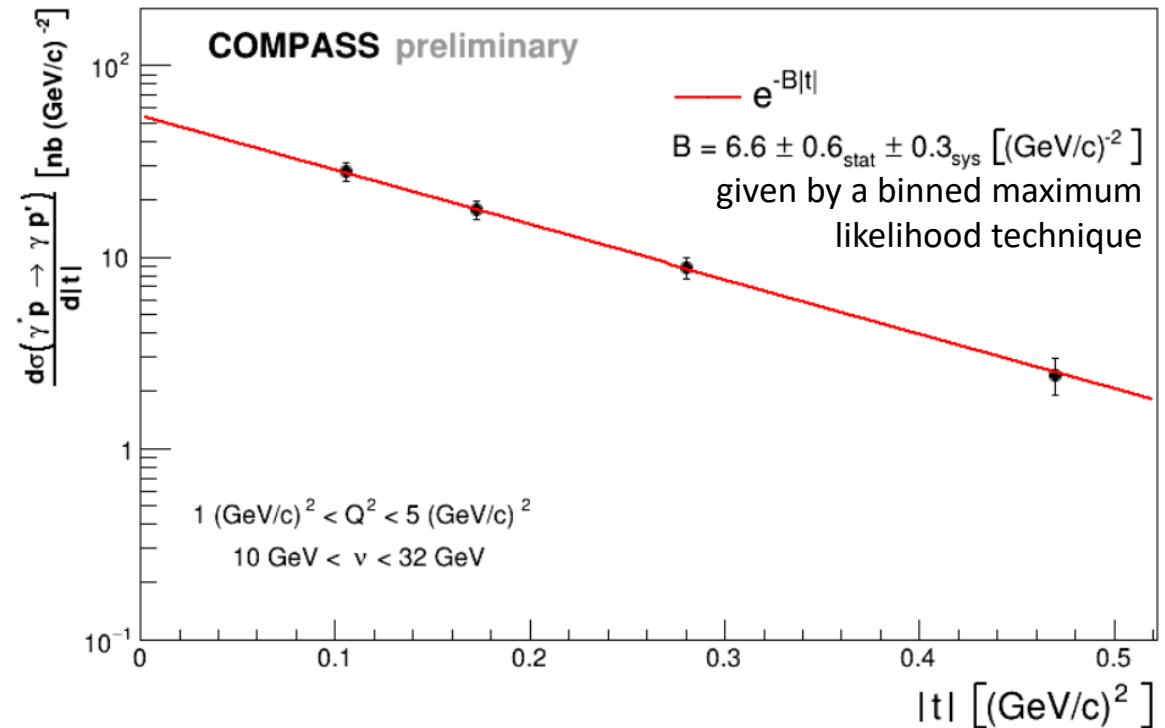
calculable
can be subtracted

All the other terms are cancelled in the integration over ϕ

$$\frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt} = \int_{-\pi}^{\pi} d\phi (d\sigma - d\sigma^{BH}) \propto c_0^{DVCS}$$

$$\frac{d\sigma^{\gamma^* p}}{dt} = \frac{1}{\Gamma(Q^2, \nu, E_\mu)} \frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt}$$

Flux for transverse virtual photons



COMPASS 2016 Transverse extension of partons in the sea quark range

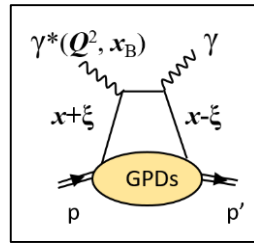
$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (\text{Im}\mathcal{H})^2$$

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2}\mathcal{E}\mathcal{E}^*$$

In the COMPASS kinematics, $x_B \approx 0.06$, dominance of $\text{Im}\mathcal{H}$
 97% (GK model) 94% (KM model)

$$\text{Im}\mathcal{H} = H(x=\xi, \xi, t)$$

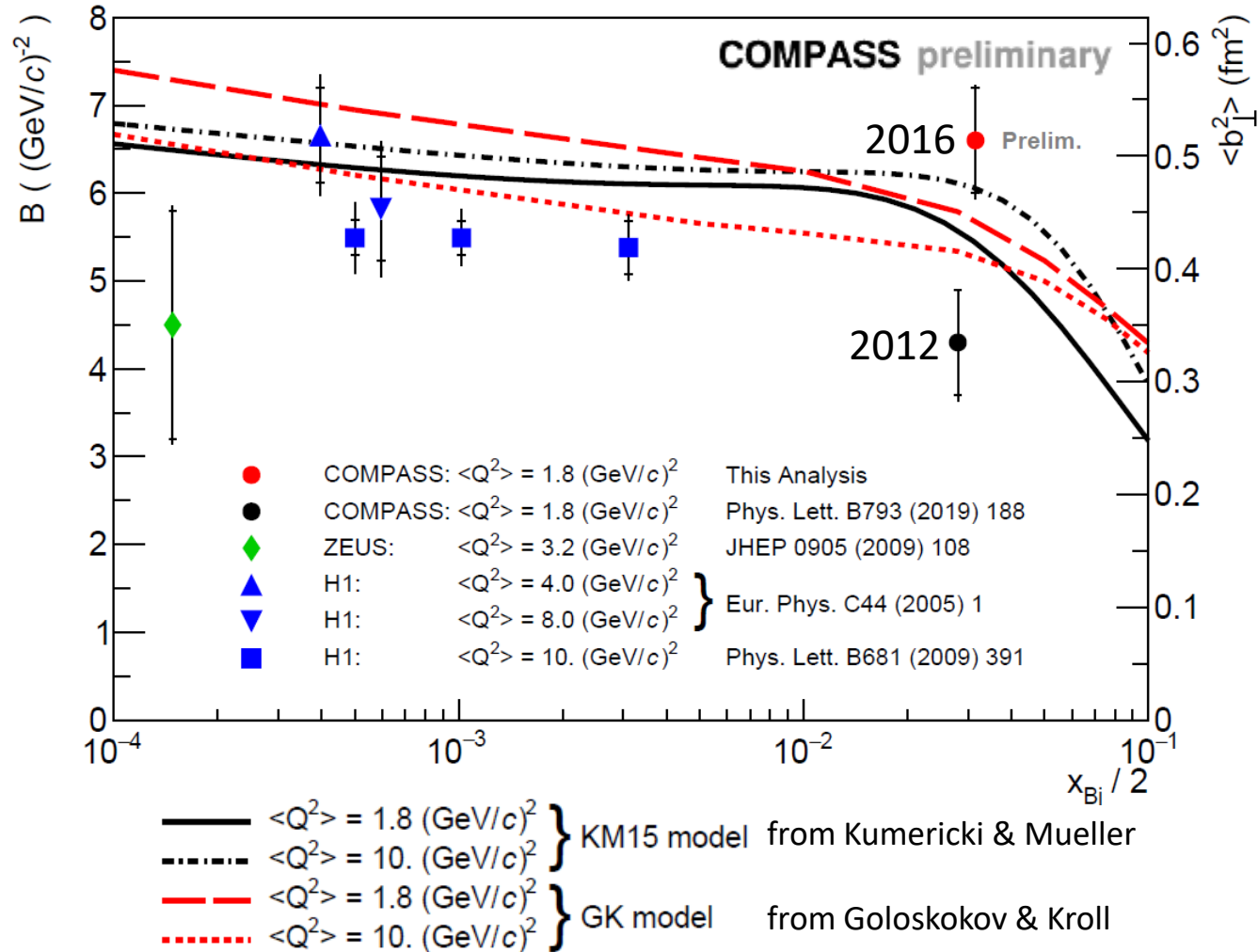
$$x = \xi \approx x_B/2 \text{ close to } 0$$



$$q(x, \mathbf{b}_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} H^q(x, 0, -\Delta_\perp^2).$$

$$\langle b_\perp^2 \rangle_x^f = \frac{\int d^2b_\perp b_\perp^2 q_f(x, b_\perp)}{\int d^2b_\perp q_f(x, b_\perp)} = -4 \frac{\partial}{\partial t} \log H^f(x, \xi=0, t) \Big|_{t=0}$$

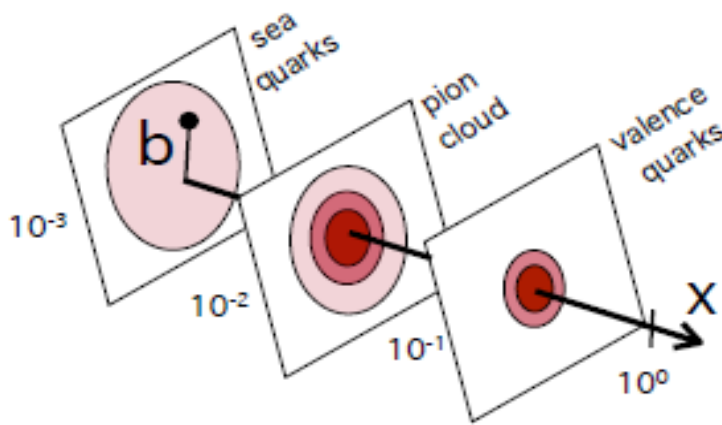
$$\langle b_\perp^2(x) \rangle \approx 2B(\xi)$$



COMPASS 2016 Transverse extension of partons in the sea quark range

$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (Im\mathcal{H})^2$$

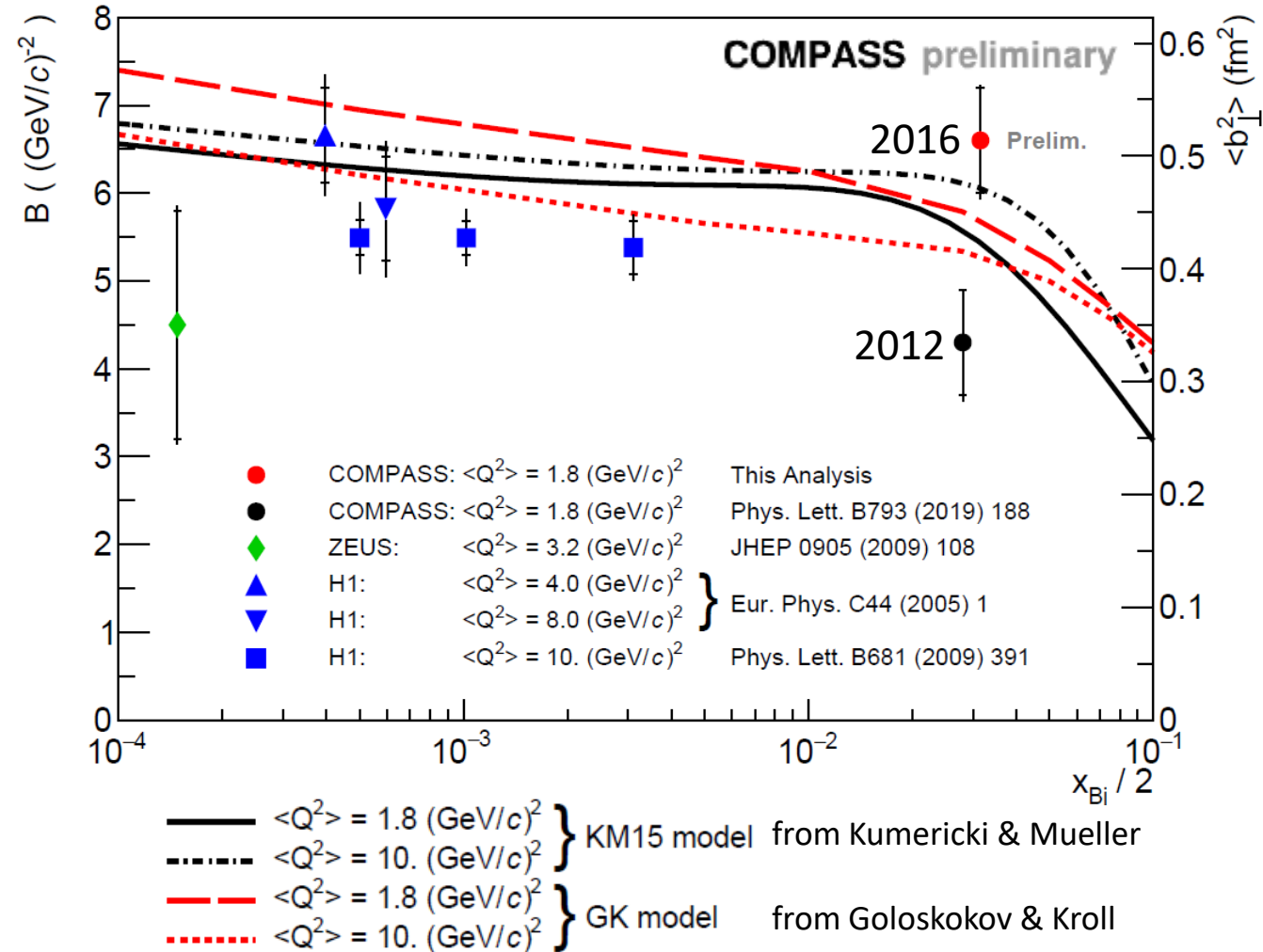
$$\langle b_{\perp}^2(x) \rangle \approx 2B(\xi)$$



3 σ difference between 2012 and 2016 data

- more advanced analysis with 2016 data
- π^0 contamination with different thresholds
- binning with 3 variables (t, Q², v) or 4 variables (t, ϕ , Q², v)

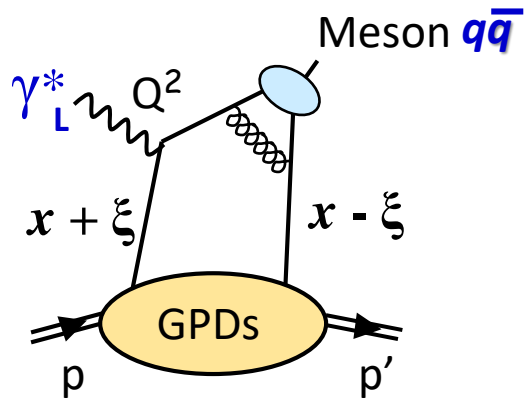
2012 statistics = Ref
 2016 analysed statistics = 2.3 \times Ref
 2016+2017 expected statistics = 10 \times Ref



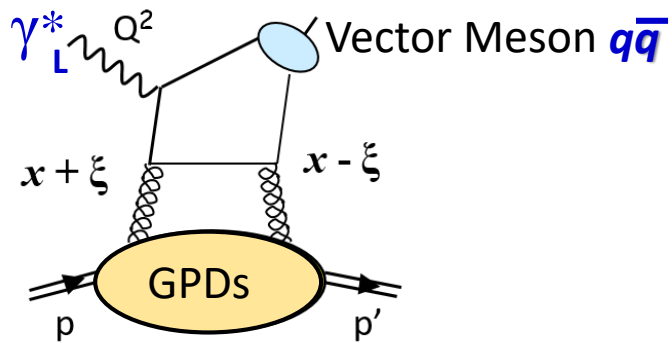
GPDs and Hard Exclusive Meson Production

Factorisation proven only for σ_L

Quark contribution



Gluon contribution at the same order in α_s



The meson wave function

Is an additional non-perturbative term

4 chiral-even GPDs: helicity of parton unchanged

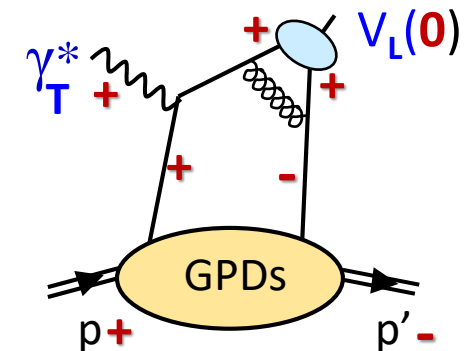
$H^q(x, \xi, t)$	$E^q(x, \xi, t)$	(as Sivers with OAM)	For Vector Meson
$\tilde{H}^q(x, \xi, t)$	$\tilde{E}^q(x, \xi, t)$		For Pseudo-Scalar Meson

+ 4 chiral-odd or transversity GPDs: helicity of parton changed
(not possible in DVCS)

$H_T^q(x, \xi, t)$	(as trans- versity)	$E_T^q(x, \xi, t)$	
$\tilde{H}_T^q(x, \xi, t)$		$\tilde{E}_T^q(x, \xi, t)$	$\bar{E}_T^q = 2\tilde{H}_T^q + E_T^q$ (as Boer-Mulders)

σ_T is asymptotically suppressed by $1/Q^2$ but large contribution observed
GK model: k_T of q and \bar{q} and Sudakov suppression factor are considered

$\mathcal{M}_{0-, ++}$ $\gamma_T^* \rightarrow V_L$ sensitive to H_T^q
and to a twist-3 meson wave function



HEMP with Transversely Polarized Target without RPD

Gparity: $G(\pi)=-1$; $G(\rho)=+1$; $G(\omega)=-1$

$$\rho^0 \rightarrow \pi^+\pi^-$$

$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u \oplus \frac{1}{3} E^d + \frac{3}{4} \frac{E_g}{x} \right)$$

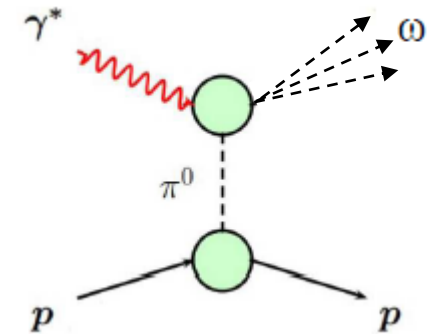
$$\omega \rightarrow \pi^+\pi^- \pi^0$$

$$E_{\omega} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u \ominus \frac{1}{3} E^d + \frac{1}{4} \frac{E_g}{x} \right)$$

E^u and E^d of opposite sign

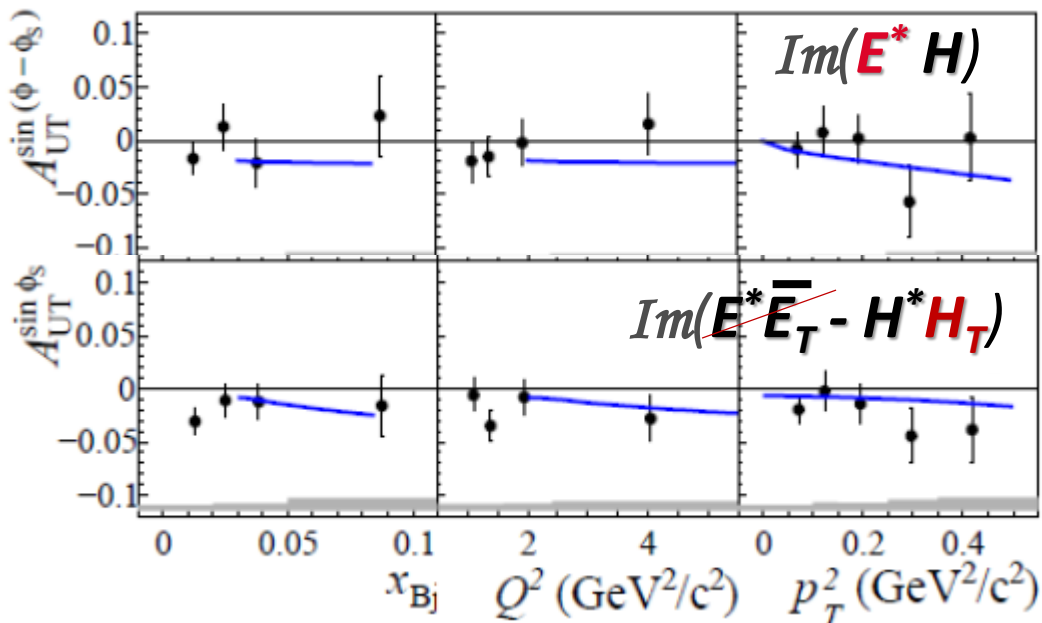
ω is more promising
(see the larger scale)
but there is the inherent
pion pole contribution

$$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma)$$

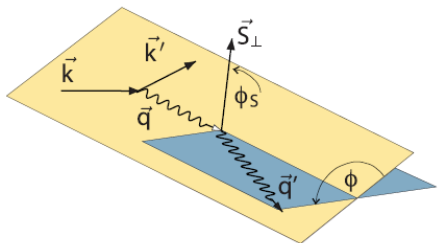
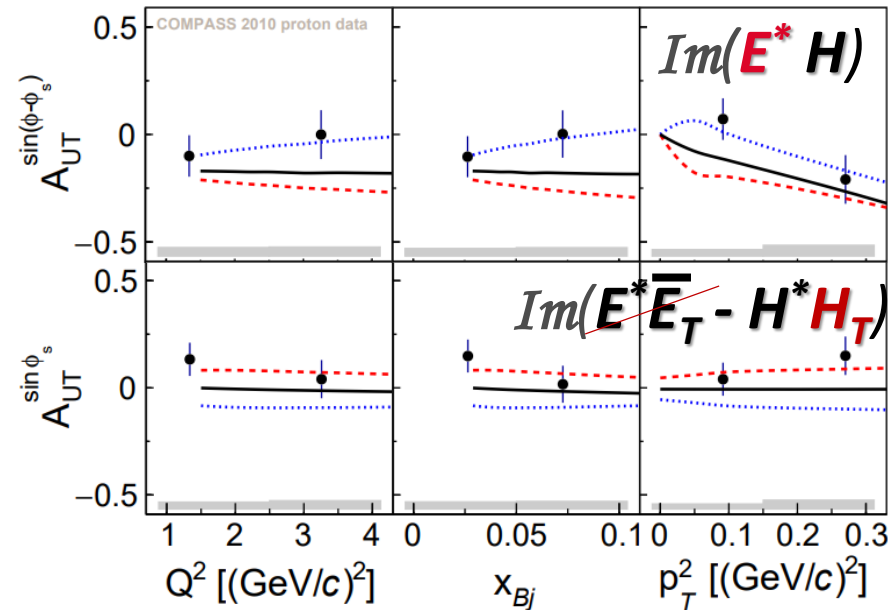


- ▶ positive $\pi\omega$ form factor
- ▶ no pion pole
- ▶ negative $\pi\omega$ form factor

COMPASS, NPB 865 (2012) 1-20, PLB731 (2014) 19



COMPASS, NPB 915 (2017)



GK Goloskokov, Kroll, EPJC42,50,53,59,65,74 GPD model constrained by HEMP at small x_B
 longitudinal $\gamma_L^* p \rightarrow M p$ and transv. polar. $\gamma_T^* p \rightarrow M p$
 quark and gluon contributions (GPDs H, E, H_T, E_T) and beyond leading twist

exclusive VM production with Unpolarised Target and SDME

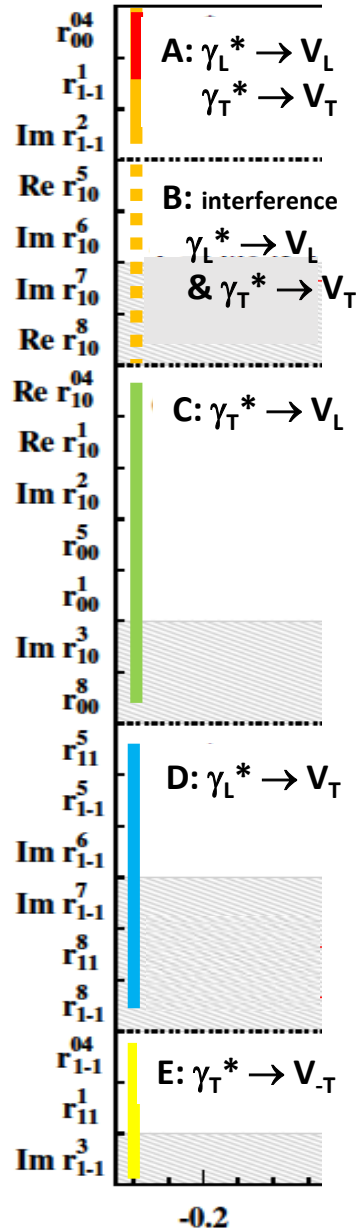
experimental angular distributions:

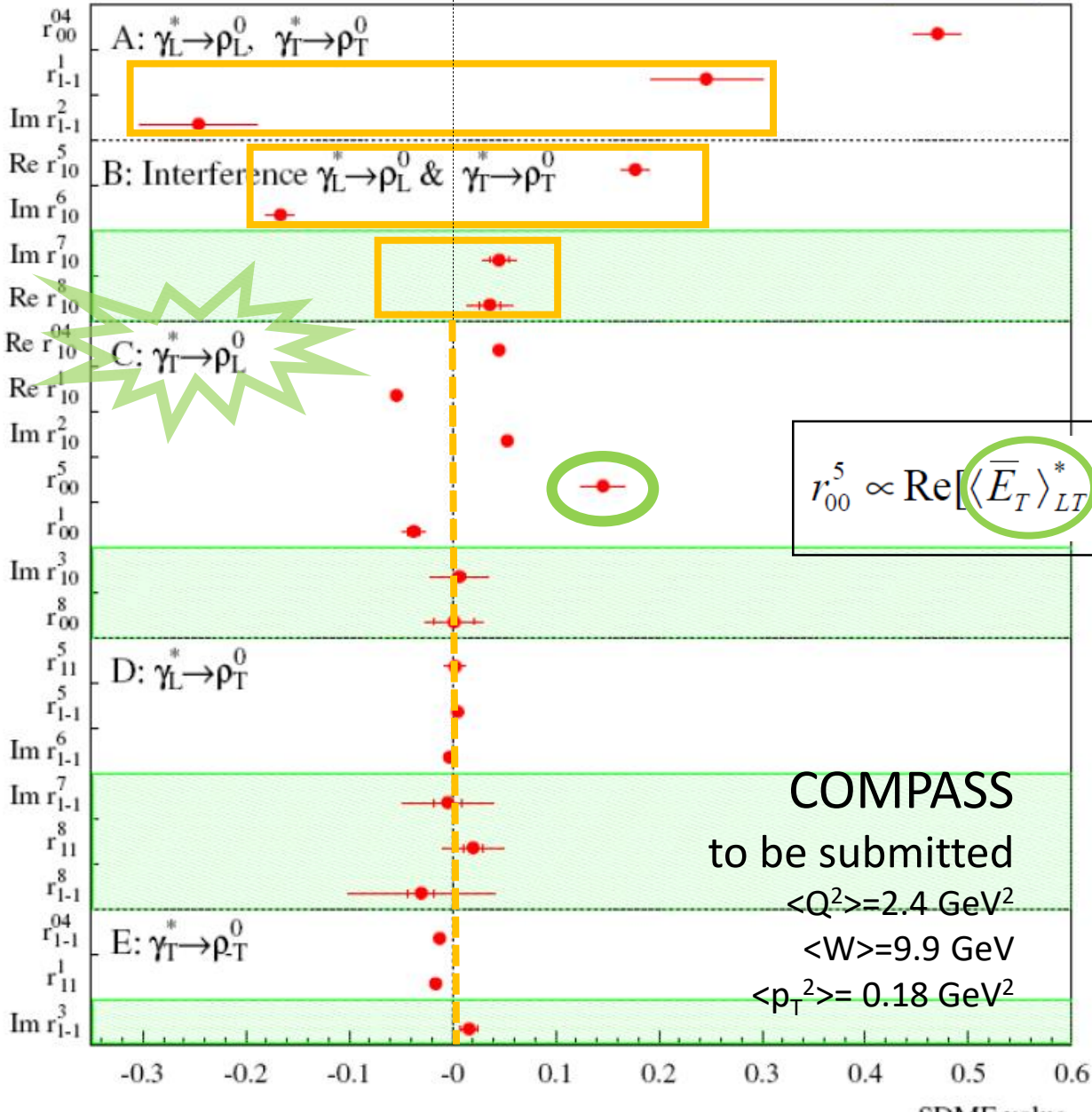
$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta)$$

15 'unpolarized' and 8 'polarized' SDMEs

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi \left(r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi \right) \\ & \left. - \epsilon \sin 2\Phi \left(\sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos \Phi \left(r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin \Phi \left(\sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi \right) \right], \\ \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\sqrt{1-\epsilon^2} \left(\sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos \Phi \left(\sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \sin \Phi \left(r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi \right) \right] \end{aligned}$$

ϵ close to 1,
small \mathcal{W}^L
no L/T separation





If SCHC ($\lambda_\gamma = \lambda_\nu$)

$$r_{1-1}^1 + \text{Im}\{r_{1-1}^2\} = 0 \quad = 0.000 \pm 0.005 \pm 0.003,$$

$$\text{Re}\{r_{10}^5\} + \text{Im}\{r_{10}^6\} = 0 \quad = 0.011 \pm 0.002 \pm 0.002,$$

$$\text{Im}\{r_{10}^7\} - \text{Re}\{r_{10}^8\} = 0 \quad = 0.009 \pm 0.014 \pm 0.028.$$

measurements:

All the other SDME in classes C, D, E should be 0
 not observed for class C

$$r_{00}^5 \propto \text{Re} \left[\langle \bar{E}_T \rangle_{LT}^* \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle_{LT}^* \langle E \rangle_{LL} \right]$$

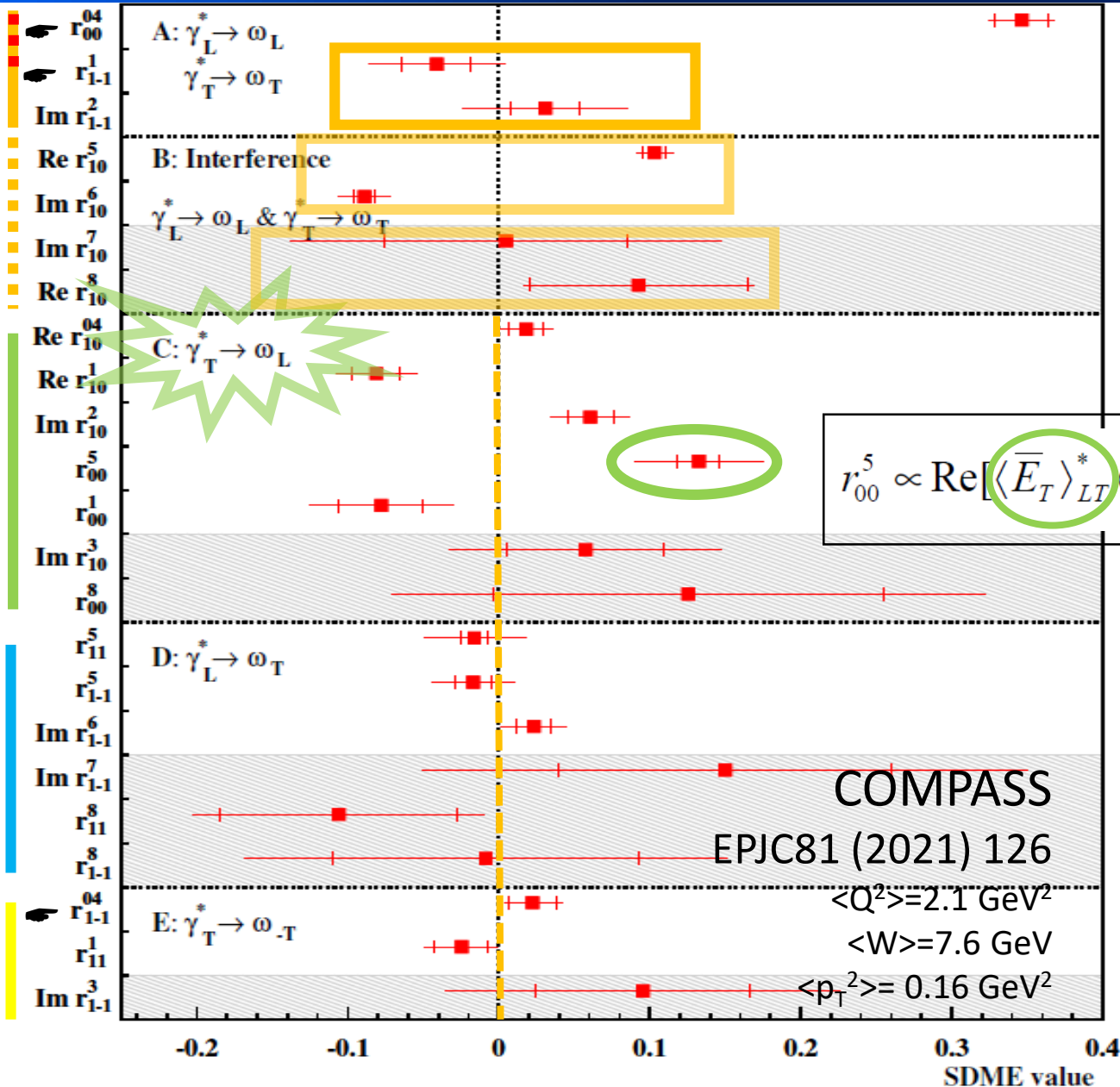
From Goloskokov and Kroll, EPJC74 (2014) 2725

$$F_{\rho^0} = 2/3 F^u + 1/3 F^d$$

→ The first term dominates and r_{00}^5 probes \bar{E}_T

COMPASS
 to be submitted
 $\langle Q^2 \rangle = 2.4 \text{ GeV}^2$
 $\langle W \rangle = 9.9 \text{ GeV}$
 $\langle p_T^2 \rangle = 0.18 \text{ GeV}^2$

(H^u, H^d) of same sign
 $(\bar{E}_T^u, \bar{E}_T^d)$ of same sign
 (H_T^u, H_T^d) of opposite sign
 (E^u, E^d) of opposite sign



If SCHC ($\lambda_\gamma = \lambda_V$)

$$r_{1-1}^1 + \text{Im}\{r_{1-1}^2\} = 0 \quad = -0.010 \pm 0.032 \pm 0.047$$

$$\text{Re}\{r_{10}^5\} + \text{Im}\{r_{10}^6\} = 0 \quad = 0.014 \pm 0.011 \pm 0.013$$

$$\text{Im}\{r_{10}^7\} - \text{Re}\{r_{10}^8\} = 0 \quad = -0.088 \pm 0.110 \pm 0.196$$

measurements:

All the other SDME in classes C, D, E should be 0
 not observed for class C

$$r_{00}^5 \propto \text{Re} \left[\langle \bar{E}_T \rangle_{LT}^* \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle_{LT}^* \langle E \rangle_{LL} \right]$$

From Goloskokov and Kroll,
 EPJC74 (2014) 2725

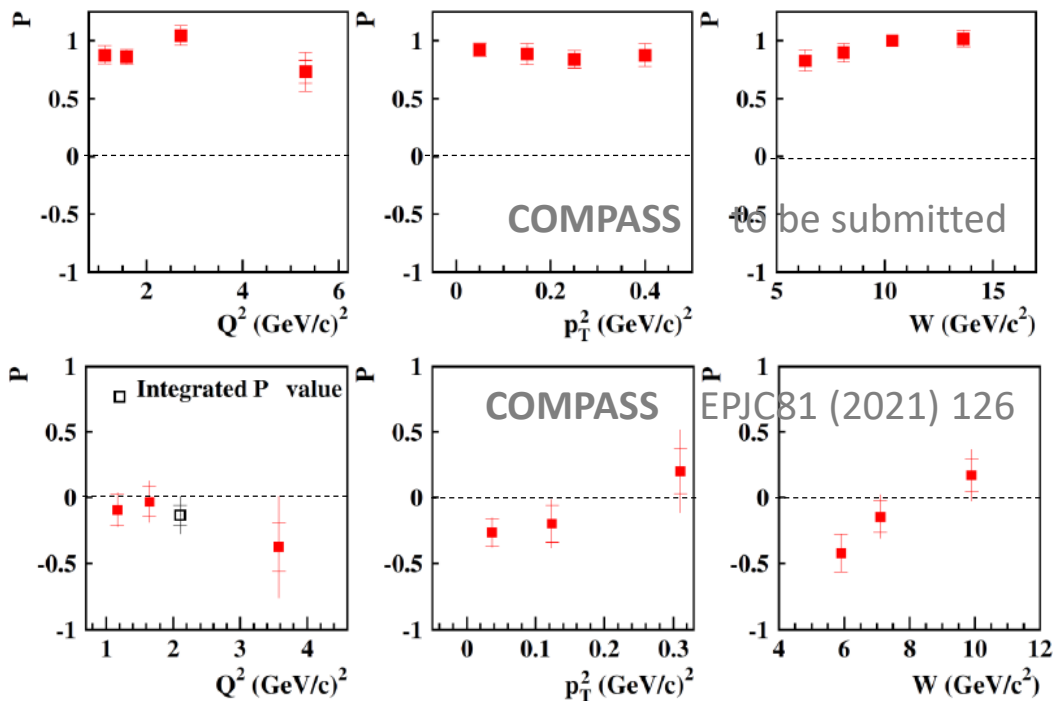
$$F_\omega = 2/3 F^u - 1/3 F^d$$

→ Both 2 terms are important

(H^u, H^d) of same sign
 $(\bar{E}_T^u, \bar{E}_T^d)$ of same sign
 (H_T^u, H_T^d) of opposite sign
 (E^u, E^d) of opposite sign

Natural (N) to Unnatural (U) Parity Exchange for $\gamma_T^* \rightarrow V_T$

$$P = \frac{2r_{1-1}^1}{1 - r_{00}^{04} - 2r_{1-1}^{04}} \approx \frac{d\sigma_T^N(\gamma_T^* \rightarrow V_T) - d\sigma_T^U(\gamma_T^* \rightarrow V_T)}{d\sigma_T^N(\gamma_T^* \rightarrow V_T) + d\sigma_T^U(\gamma_T^* \rightarrow V_T)}$$

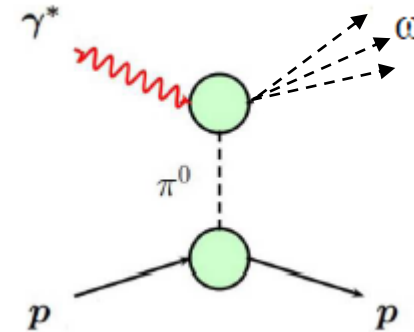


ρ^0 : $P \sim 1 \rightarrow$ NPE dominance $P \sim 1$
NPE with GPDs H, E

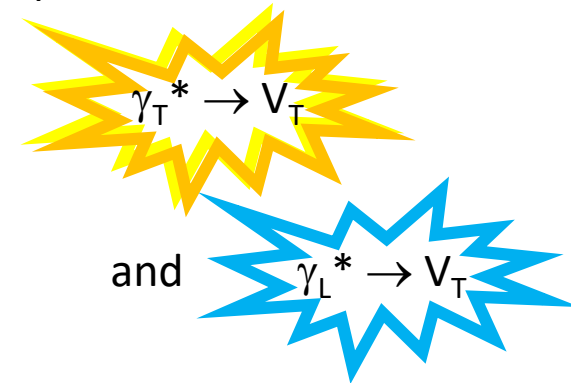
ω : $P \sim 0 \rightarrow$ NPE \sim UPE
 UPE dominance at small W and p_T^2
UPE with GPDs \tilde{H}, \tilde{E} and the dominant pion pole

The pion pole exchange (UPE) is large for ω compared to ρ^0

$$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma) \quad \text{as for the transition } \pi^0 V \text{ FF}$$



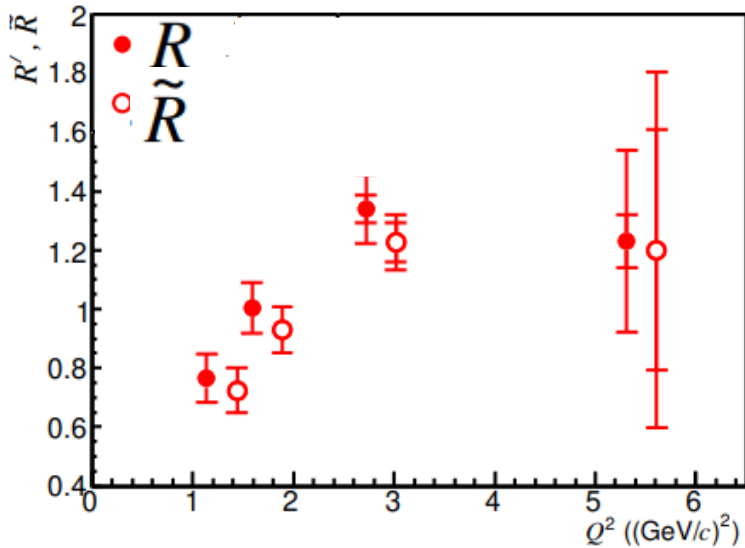
It plays an important role in ω production for:



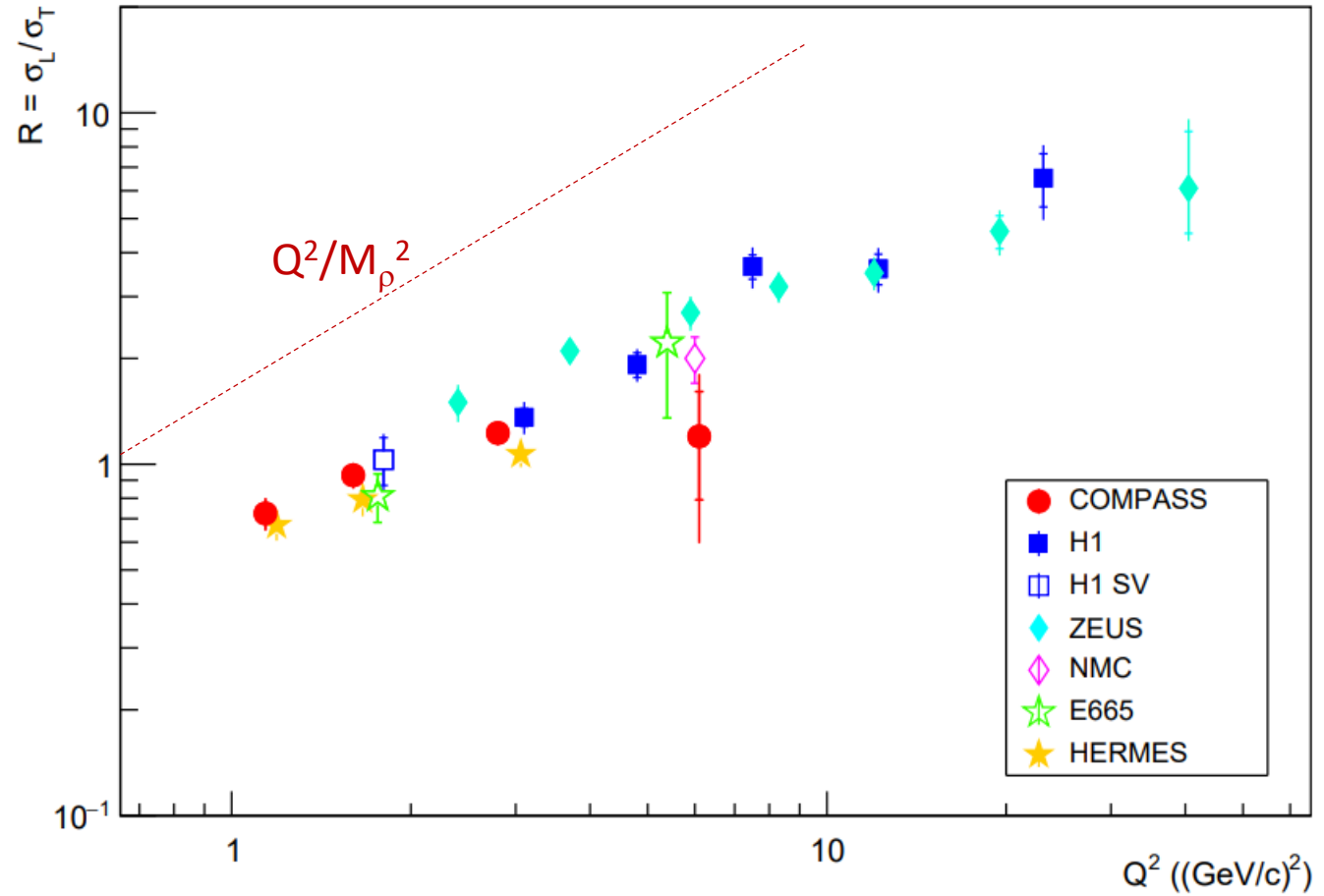
$$R = \frac{\sigma_L(\gamma_L^* \rightarrow V)}{\sigma_T(\gamma_T^* \rightarrow V)}$$

$$R = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}} \quad \text{only if SCHC}$$

In COMPASS domain evaluation of R and \tilde{R} considering violation of SCHC (and only NPE)



for all the experiments with $Q^2 > 1 \text{ GeV}^2$



Deviation from the pQCD LO prediction in Q^2/M_ρ^2 : QCD evolution and q_T Transverse size effects of the meson smaller for σ_L than for σ_T

- ✓ DVCS and the sum $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$
 - c_0 and s_1 and constrain on $\text{Im}\mathcal{H}$ and Transverse extension of partons
- ✓ DVCS and the difference $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$
 - c_1 and constrain on $\text{Re}\mathcal{H}$ and D-term and pressure distribution
- ✓ On-going analysis (Cross section, SDME) for HEMP of $\pi^0, \rho^0, \omega, \phi, J/\psi$
 - ✓ Transversity GPDs
 - ✓ Gluon GPDs
 - ✓ Flavor decomposition



Exclusive single photon events selection at high ν

$$80 < \nu \text{ [GeV]} < 144$$

$$\Delta\phi = \phi^{\text{cam.}} - \phi^{\text{spec.}}$$

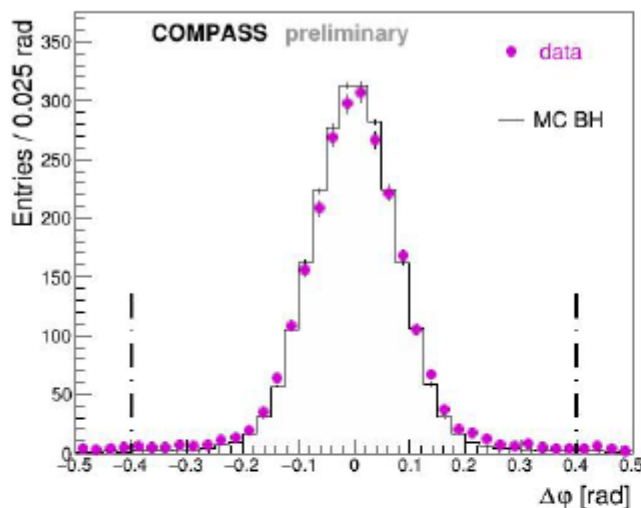
$$\Delta p_T = |p_T^{\text{cam.}}| - |p_T^{\text{spec.}}|$$

$$\Delta Z_A = z_A^{\text{cam.}} - z_A^{\text{spec.}}$$

$$M_{\text{undet}}^2 = (k + p - k' - q' - p')^2$$

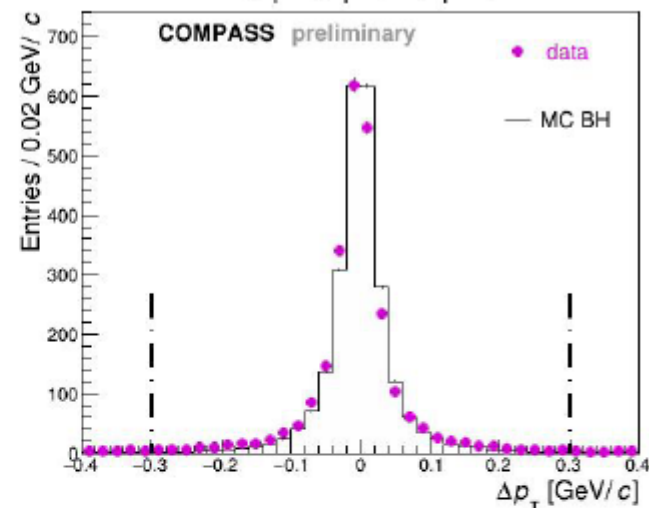
Each variable is plotted with the cuts on the 3 others applied

$$\Delta\phi = \phi^{\text{cam.}} - \phi^{\text{spec.}}$$

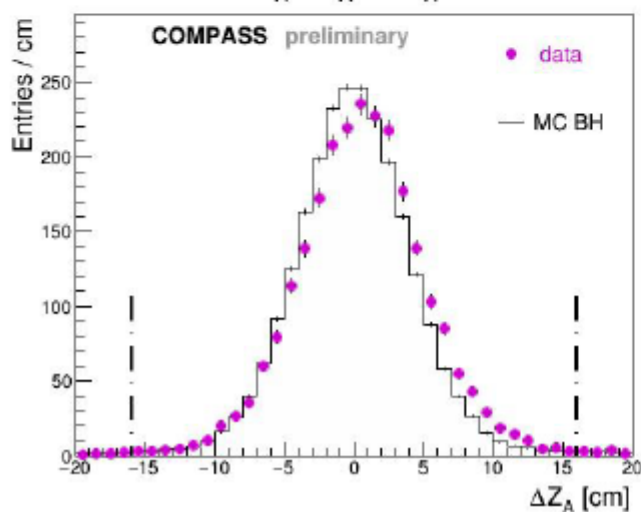


μ^+ and μ^- yields
renormalised to 10^{12} muons each

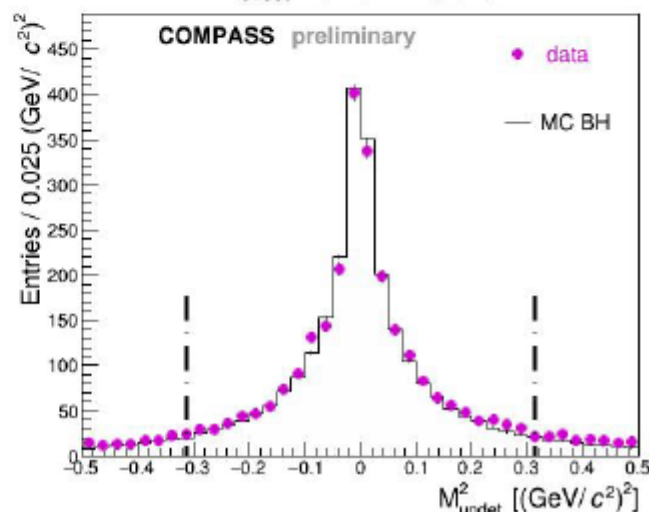
$$\Delta p_T = |p_T^{\text{cam.}}| - |p_T^{\text{spec.}}|$$



$$\Delta Z_A = z_A^{\text{cam.}} - z_A^{\text{inter.}}$$

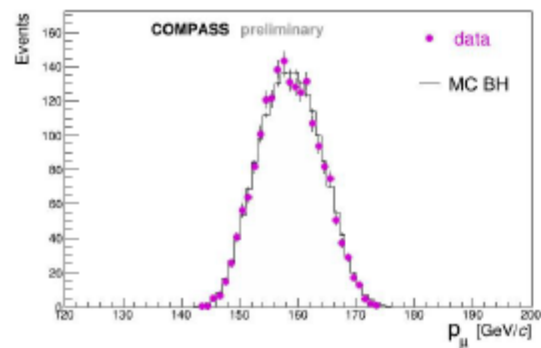


$$M_{\text{undet}}^2 = (k + p - k' - q' - p')^2$$

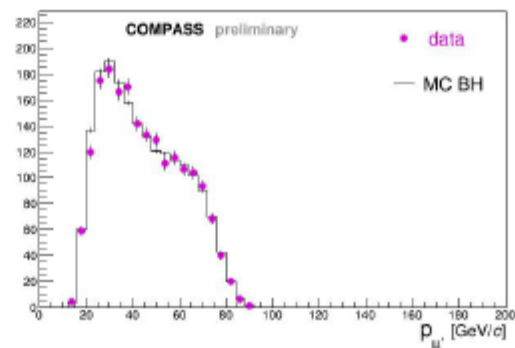


Muons

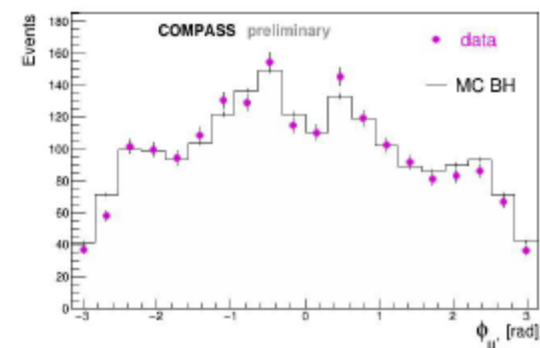
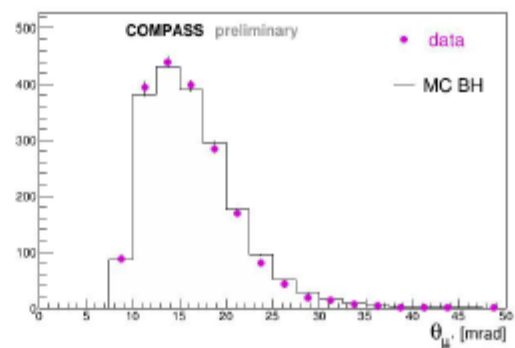
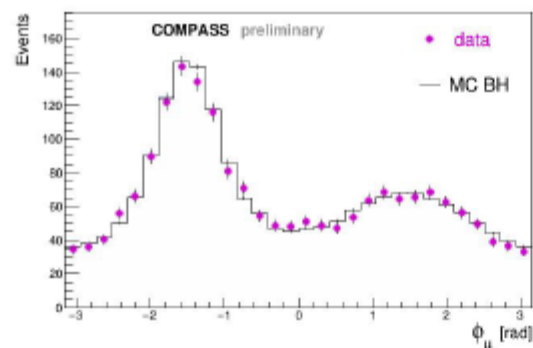
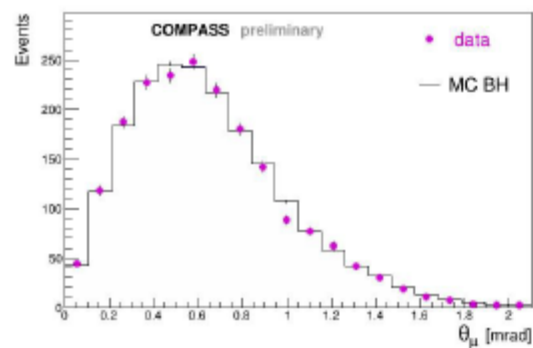
$80 < \nu \text{ [GeV]} < 144$



Beam muon

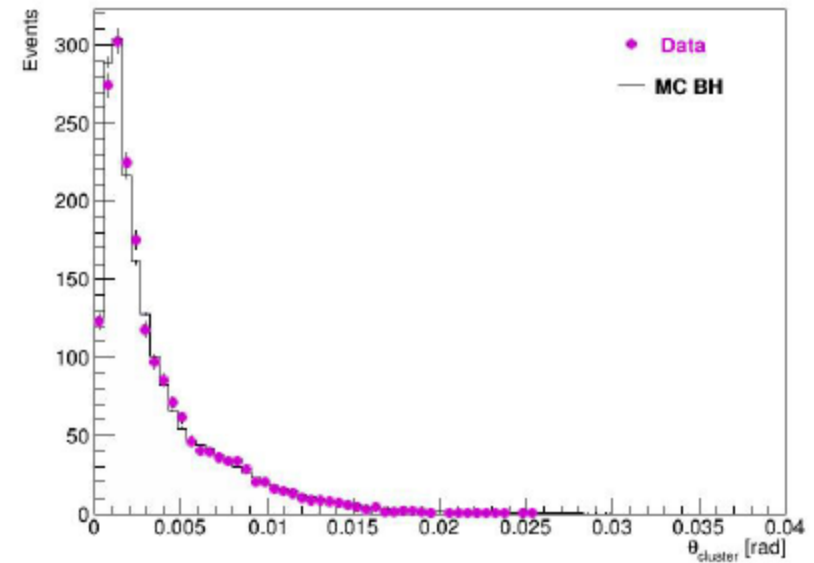
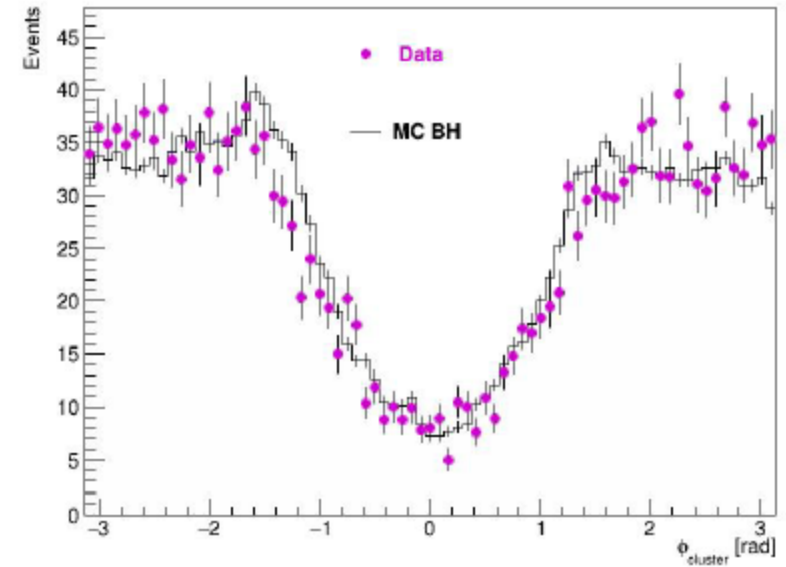
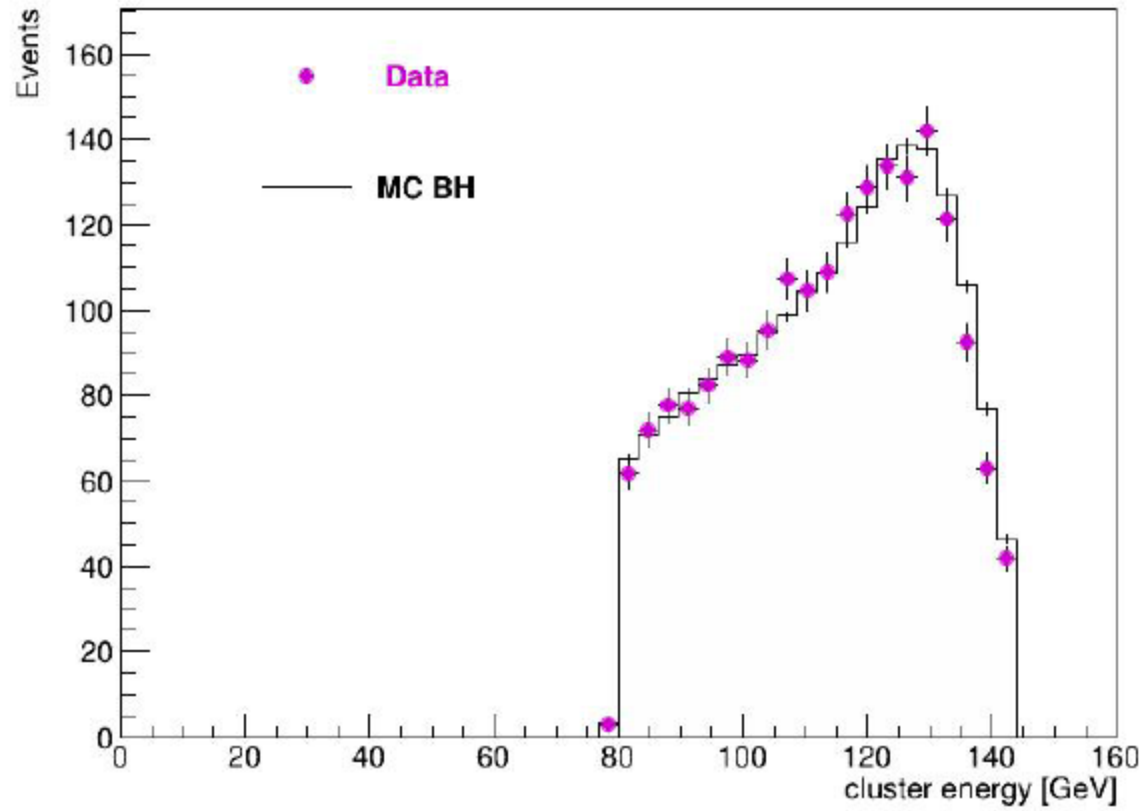


Scattered muon



Photon - Ecal2 (Bethe-Heitler)

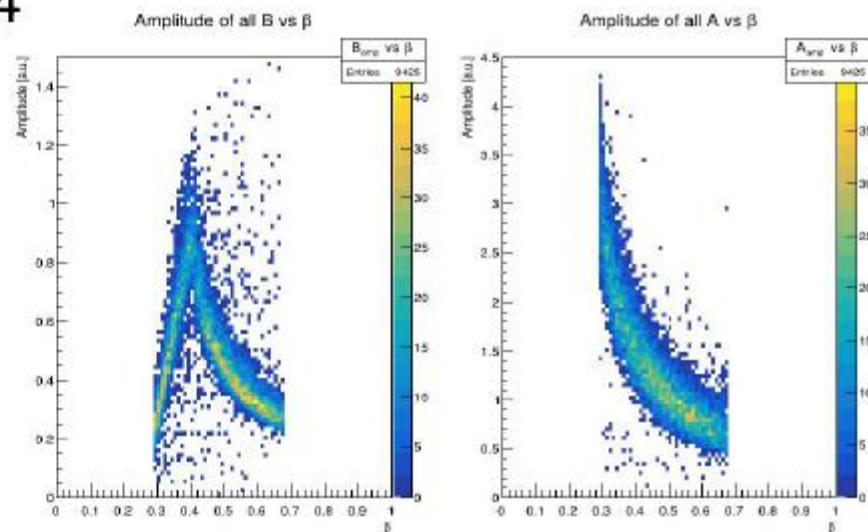
$80 < \nu \text{ [GeV]} < 144$



Proton

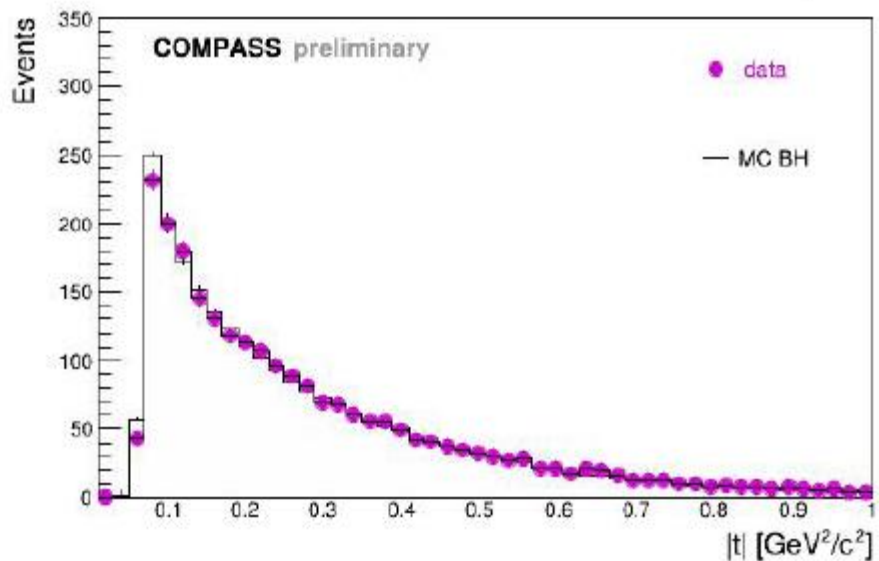
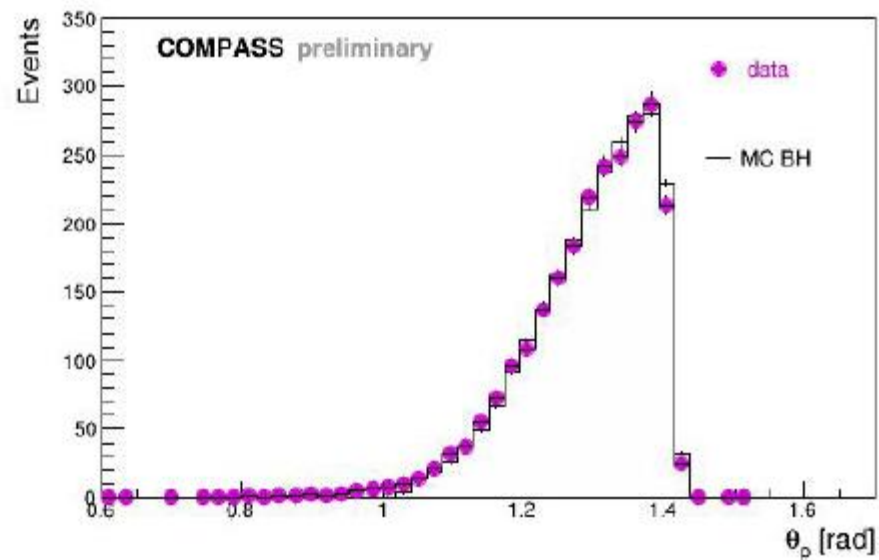
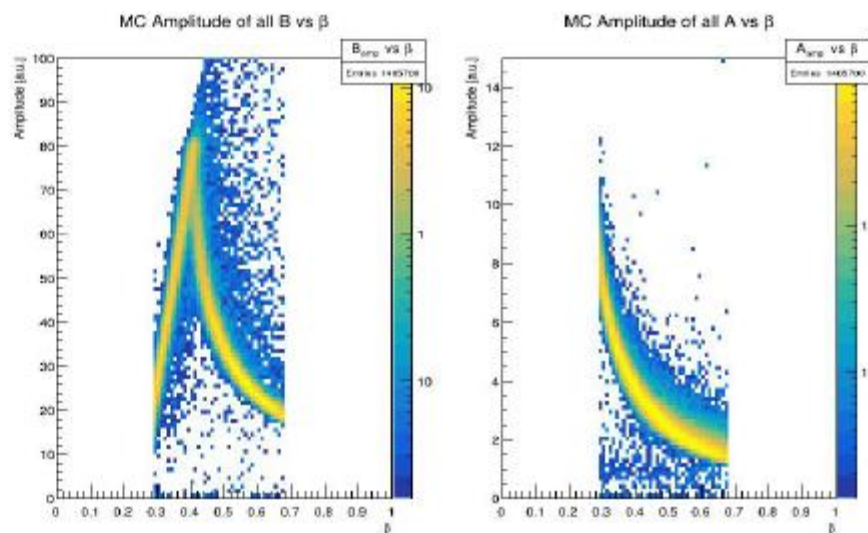
$80 < \nu \text{ [GeV]} < 144$

DATA



(t cut applied)

MC BH

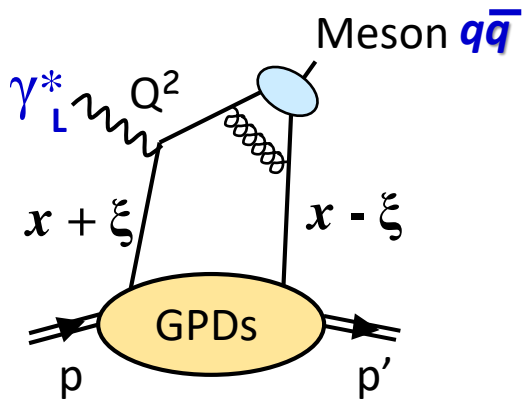


GPDs and Hard Exclusive Meson Production

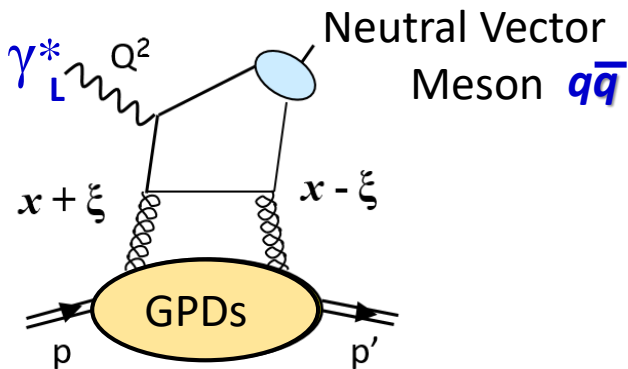
Factorisation proven only for σ_L

The meson wave function is an additional non-perturbative term

Quark contribution



Gluon contribution at the same order in α_s



4 chiral-even GPDs: helicity of parton unchanged

$H^q(x, \xi, t)$	$E^q(x, \xi, t)$	(as Sivers with OAM)	For Vector Meson
$\tilde{H}^q(x, \xi, t)$	$\tilde{E}^q(x, \xi, t)$		For Pseudo-Scalar Meson

Flavor decomposition (val and sea quarks and gluons) Diehl, Vinnikov

PLB 609 (2005)

$$F_{\rho^0} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} F^u + \frac{1}{3} F^d + \frac{3}{8} \frac{F_g}{x} \right)$$

$$F_{\omega} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} F^u - \frac{1}{3} F^d + \frac{1}{8} \frac{F_g}{x} \right)$$

$$F_{\phi} = -\frac{1}{3} F^s - \frac{1}{8} \frac{F_g}{x}$$

F for $H, E \dots$

✓ H^u, H^d of same sign

$\sigma_L(\rho^0) \sim 9 \times \sigma_L(\omega)$ with Unpol Target

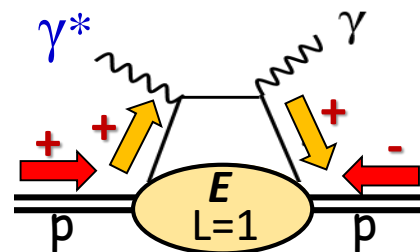
✓ E^u, E^d of opposite sign

$$A_{UT}^{\sin(\phi-\phi_s)} \sim \text{Im}[\langle E \rangle^* \langle H \rangle]$$

$$A_{UT}^{\sin(\phi-\phi_s)}(\omega) > A_{UT}^{\sin(\phi-\phi_s)}(\rho^0) \text{ with Trans Pol Target}$$

Access to the GPD E with transversely polarized target

with DVCS on the proton or Vect Meson - OR - DVCS off the neutron



the Holy Grail with E : to reveal OAM

$$J_i: 2J^q = \int x (H^q(x, \xi, 0) + E^q(x, \xi, 0)) dx$$

π^0 are one of the main background sources for excl. photon events.

Two possible case:

- **Visible** (both γ detected \rightarrow subtracted)

the exclusive single photon candidate is combined with all detected photons below the DVCS threshold: 4,5 GeV in ECAL0, 1

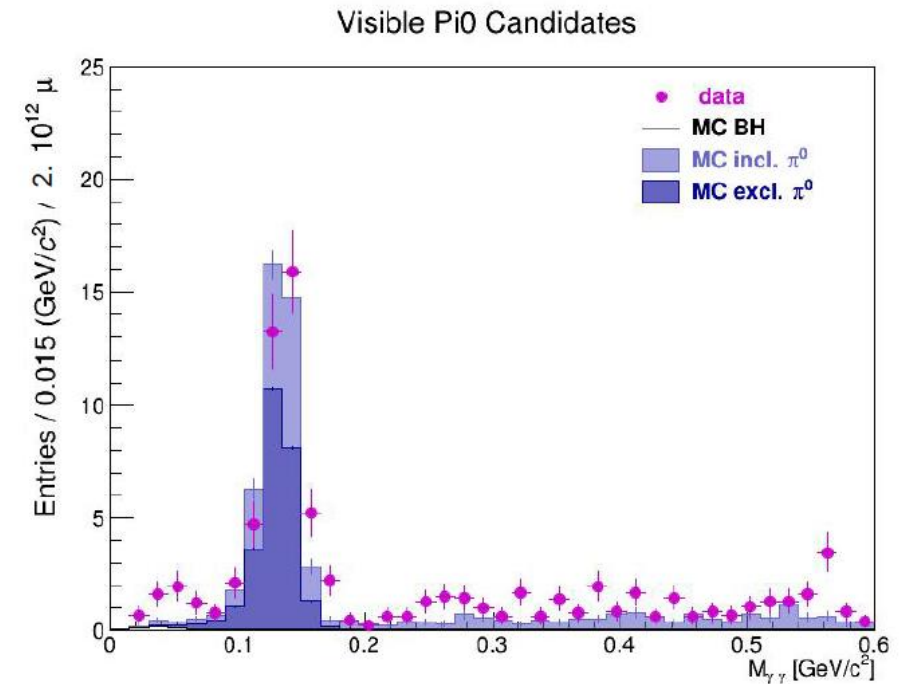
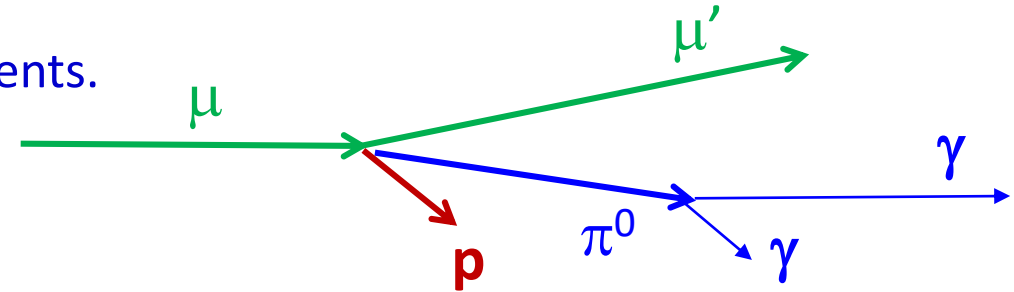
- **Invisible** (one γ lost \rightarrow estimated by MC)

- Semi-inclusive LEPTO 6.1
- Exclusive HEPGEN π^0
- (Goloskokov-Kroll model)

Comparing the two components to the data allows the determination of their relative normalisation.

$$r_{\text{LEPTO}} = 40 \pm 10\%$$

The sum of the 2 components is normalized to the visible π^0 contamination in the $M_{\gamma\gamma}$ peak



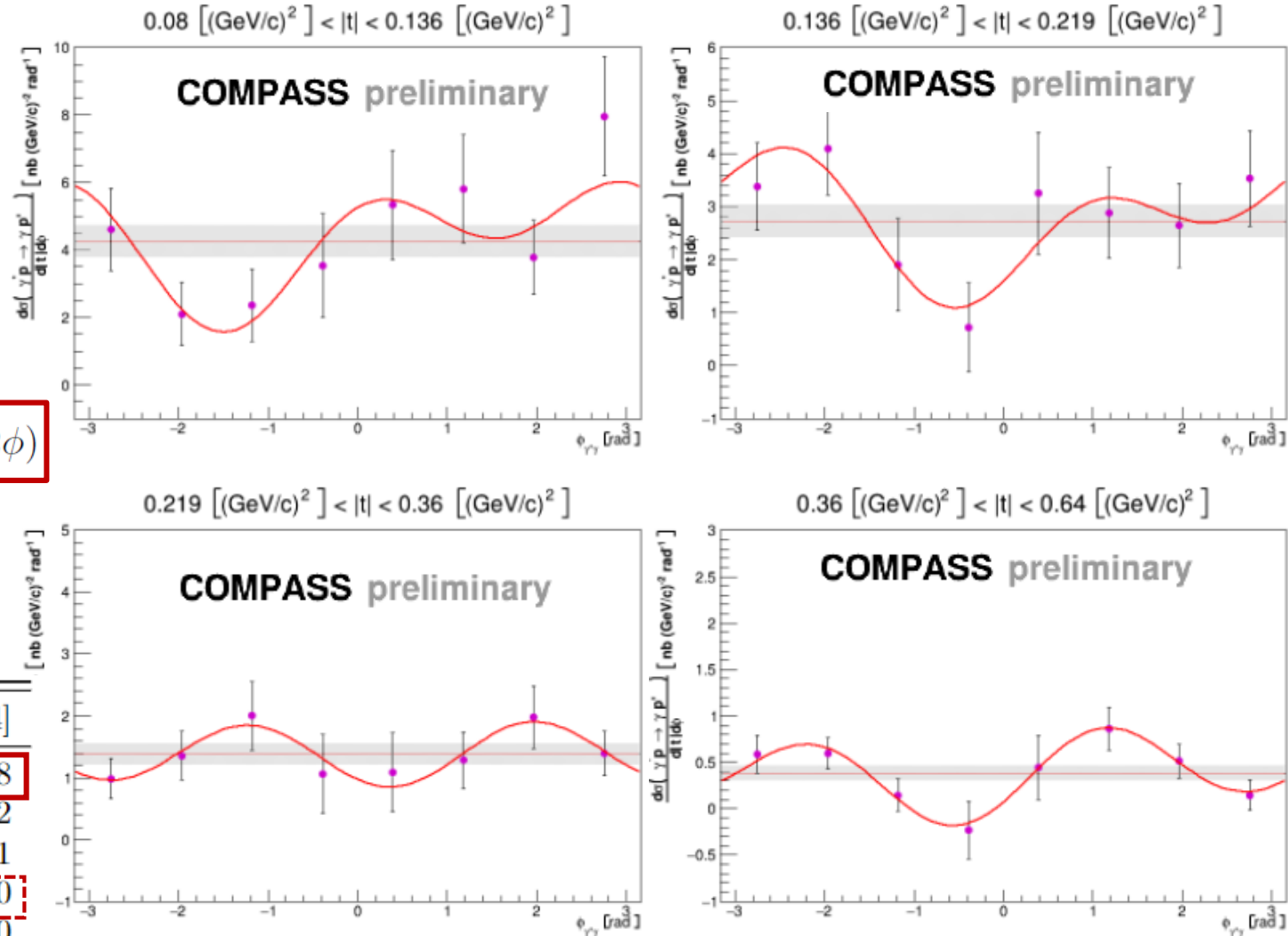
$$d\sigma_{\text{DVCS}}^{\gamma^* p \rightarrow \gamma p'} / d|t|d\phi$$

$10 \text{ GeV} < \nu < 32 \text{ GeV}$

$1 \text{ (GeV/c)}^2 < Q^2 < 5 \text{ (GeV/c)}^2$

$0.08 \text{ (GeV/c)}^2 < |t| < 0.64 \text{ (GeV/c)}^2$

In 4 $|t|$ bins with similar statistics



$$f(\phi) = c_0 + c_1 \cos(\phi) + c_2 \cos(2\phi) + s_1 \sin(\phi) + s_2 \sin(2\phi)$$

$$c_0^{\text{DVCS}} \propto (\text{Im}\mathcal{H})^2 \quad s_1^I \propto (\text{Im}\mathcal{H})$$

Results of the fit of $d\sigma_{\text{DVCS}}^{\gamma^* p \rightarrow \gamma p'} / d|t|d\phi$ [nb (GeV/c)⁻² rad⁻¹]

$ t $ -bin	[0.08,0.136]	[0.136,0.219]	[0.219,0.36]	[0.36, 0.64]
c_0	4.26 ± 0.48	2.73 ± 0.31	1.38 ± 0.17	0.38 ± 0.08
c_1	-0.32 ± 0.73	-0.94 ± 0.45	-0.06 ± 0.25	-0.12 ± 0.12
c_2	1.30 ± 0.67	-0.20 ± 0.44	-0.35 ± 0.24	-0.20 ± 0.11
s_1	1.38 ± 0.61	1.09 ± 0.43	0.01 ± 0.24	0.17 ± 0.10
s_2	0.09 ± 0.78	0.44 ± 0.31	-0.33 ± 0.23	0.33 ± 0.10

At COMPASS using polarized positive and negative muon beams:

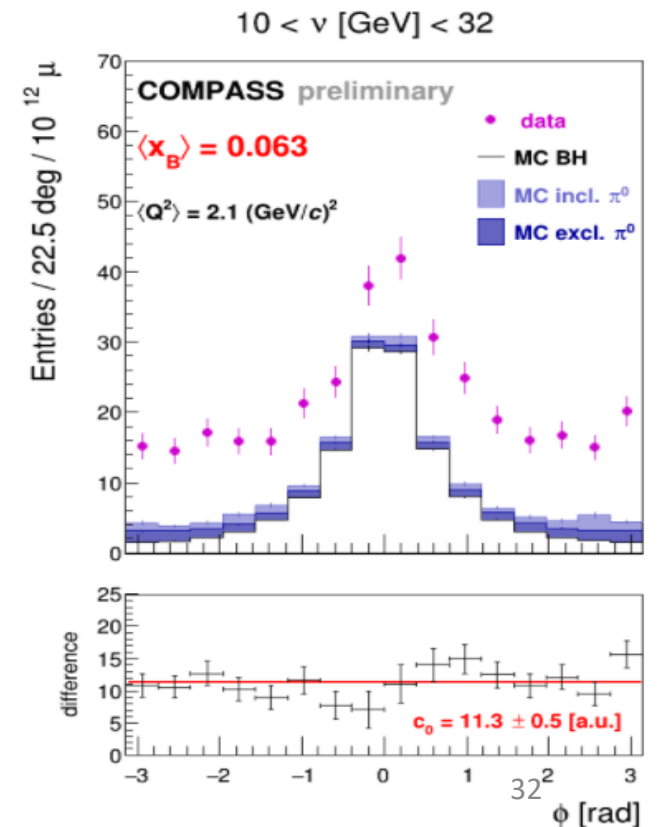
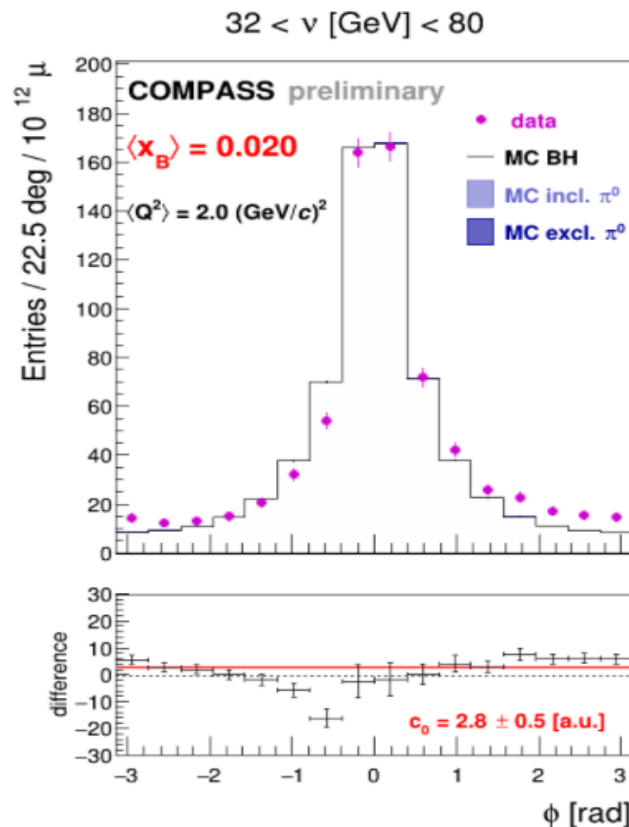
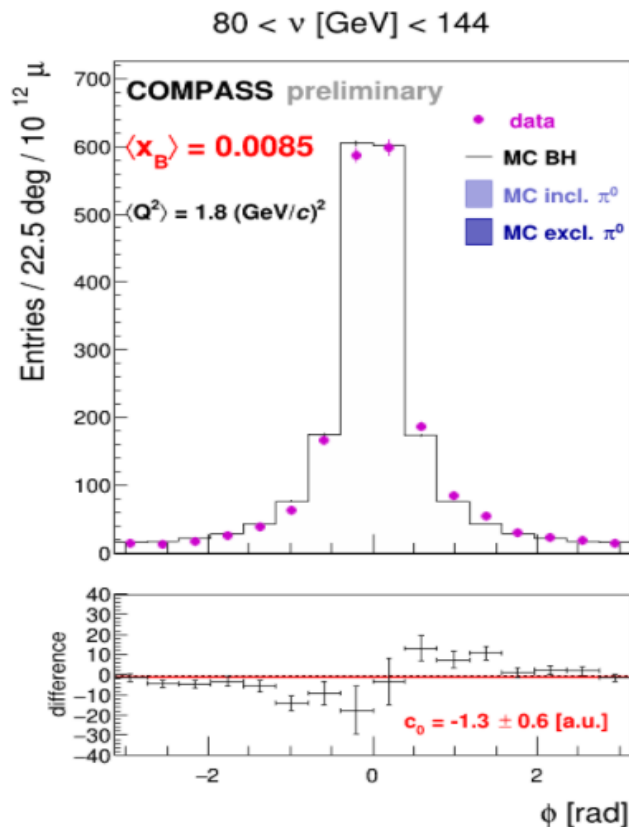
$$S_{CS,U} \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

At small x_B

$$c_0^{DVCS} \propto (\text{Im } \mathcal{H})^2$$

$$s_1^I \propto (\text{Im } \mathcal{H})$$



8 GPDs in parallel of 8 TMDs

		Quark Polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	H		$\bar{E}_T = 2\tilde{H}_T + E_T$
	L		\tilde{H}	\tilde{E}_T
	T	E	\tilde{E}	H_T, \tilde{H}_T

		Quark Polarization		
		U	L	T
Nucleon Polarization	U	f_1 unpolarized 		h_1^\perp Boer-Mulders
	L		g_{1L} helicity 	h_{1L}^\perp longi-transversity (worm-gear)
	T	f_{1T}^\perp Sivers 	g_{1T} trans-helicity (worm-gear) 	h_1 transversity h_{1T}^\perp pretzelosity

Nucleon spin Quark spin

For valence contributions:

H^u, H^d of **same** sign

E^u, E^d of **opposite** sign

\tilde{H}^u, \tilde{H}^d of **opposite** sign

\tilde{E}^u, \tilde{E}^d of **opposite** sign

H_T^u, H_T^d of **opposite** sign

\bar{E}_T^u, \bar{E}_T^d of **same** sign

GPDs and TMDs

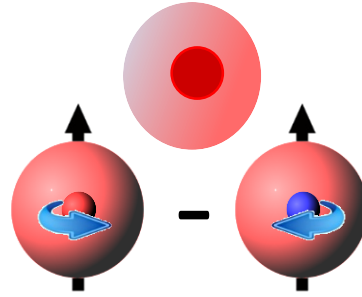
Chiral-even

$$H \leftrightarrow q$$

$$E \leftrightarrow f_{1T}^\perp$$

the Holy Grail: to reveal OAM

$$J_i: 2J^q = \int \mathbf{x} (H^q(\mathbf{x}, \xi, 0) + E^q(\mathbf{x}, \xi, 0)) d\mathbf{x}$$



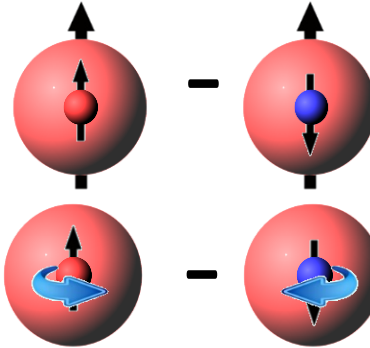
Sivers: quark k_T
and nucleon transv. Spin

T-odd

Chiral-odd

$$H_T \leftrightarrow h_1$$

$$\bar{E}_T = 2\tilde{H}_T + E_T \leftrightarrow h_1^\perp$$

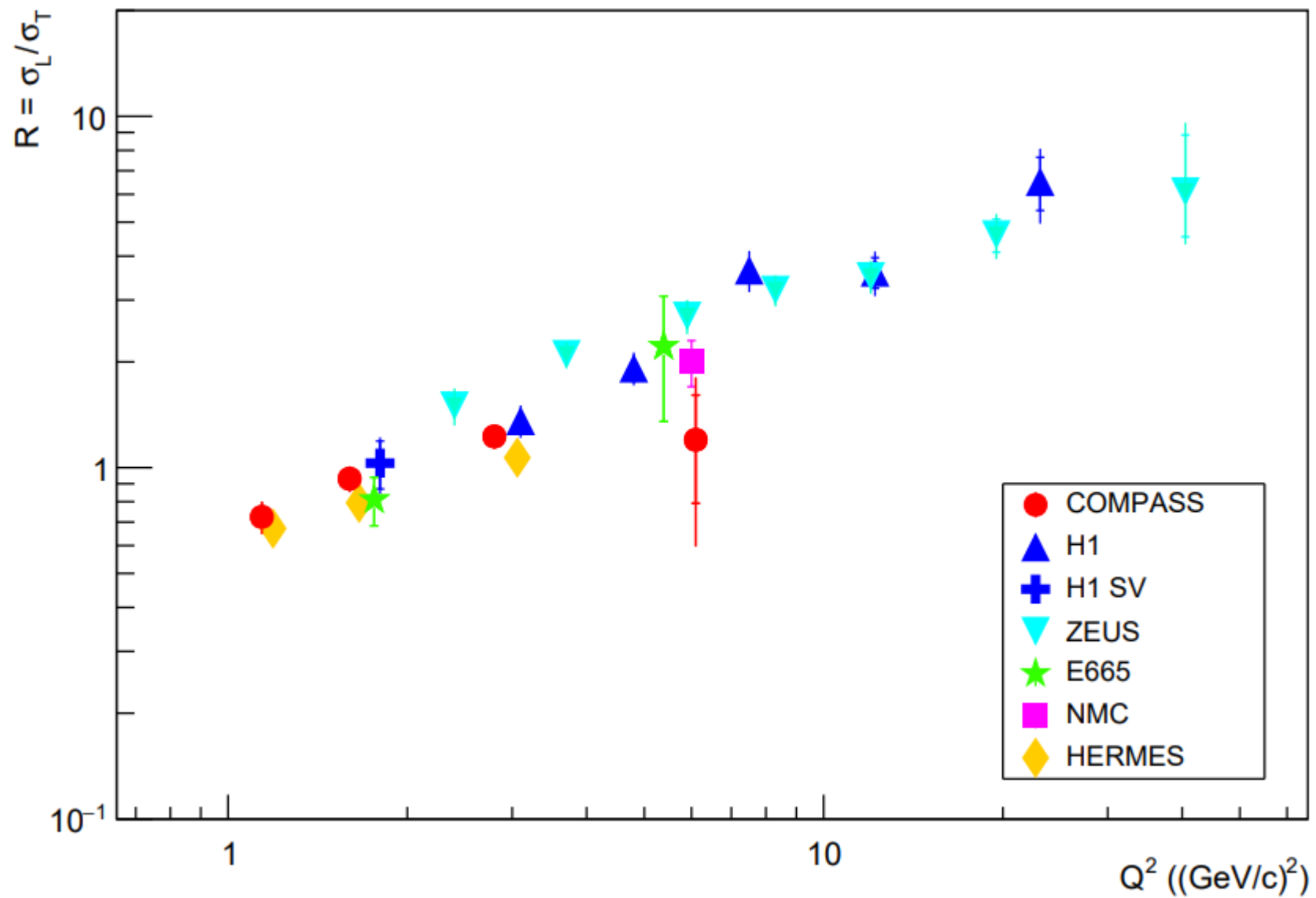


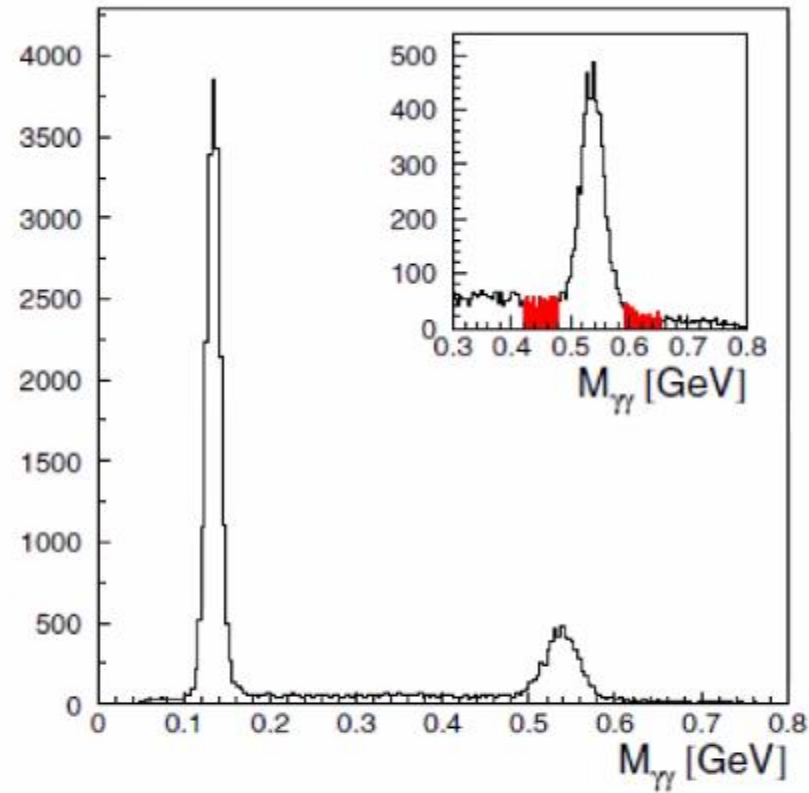
Transversity: quark spin
and nucleon transv. spin

Boer-Mulders: quark k_T
and quark transverse spin

T-odd

$$F_{\rho^0} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} F^u + \frac{1}{3} F^d + \frac{3}{8} \frac{F_g}{x} \right)$$
$$F_{\omega} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} F^u - \frac{1}{3} F^d + \frac{1}{8} \frac{F_g}{x} \right)$$
$$F_{\phi} = -\frac{1}{3} F^s - \frac{1}{8} \frac{F_g}{x}.$$

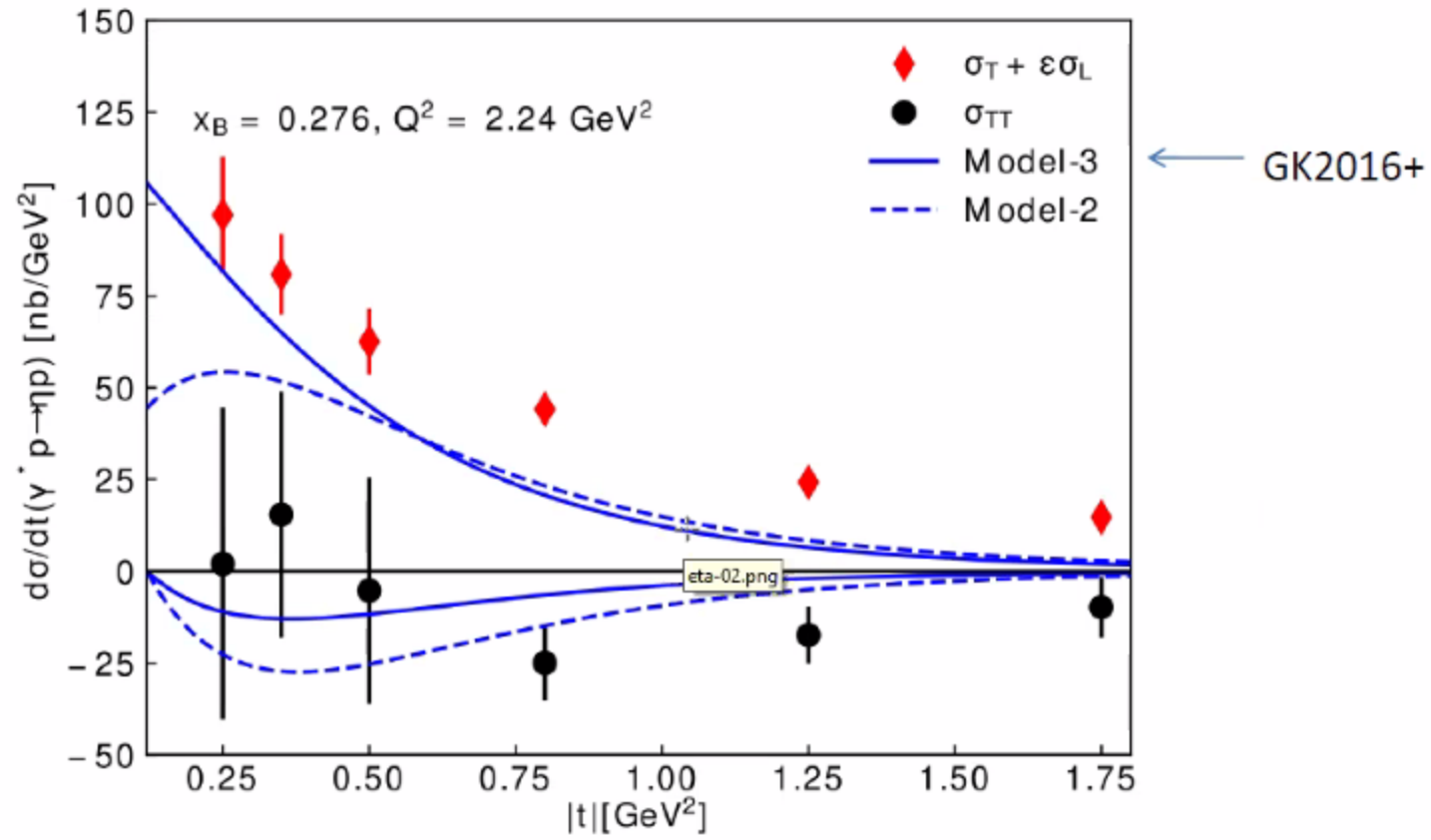




Eta and pi0
At JLab

FIG. 8. The two-photon invariant-mass distribution $M_{\gamma\gamma}$ after all exclusivity cuts have been applied, for the case where the two photons are detected by the IC. The large peak at lower $M_{\gamma\gamma}$ is due to π^0 electroproduction and the smaller peak at higher $M_{\gamma\gamma}$ is due to η electroproduction. The inset magnifies the region around the η peak.

CLAS η data & GK fits preliminary



π GK2016+ prediction η

For COMPASS

