

LHC fixed-target experiments

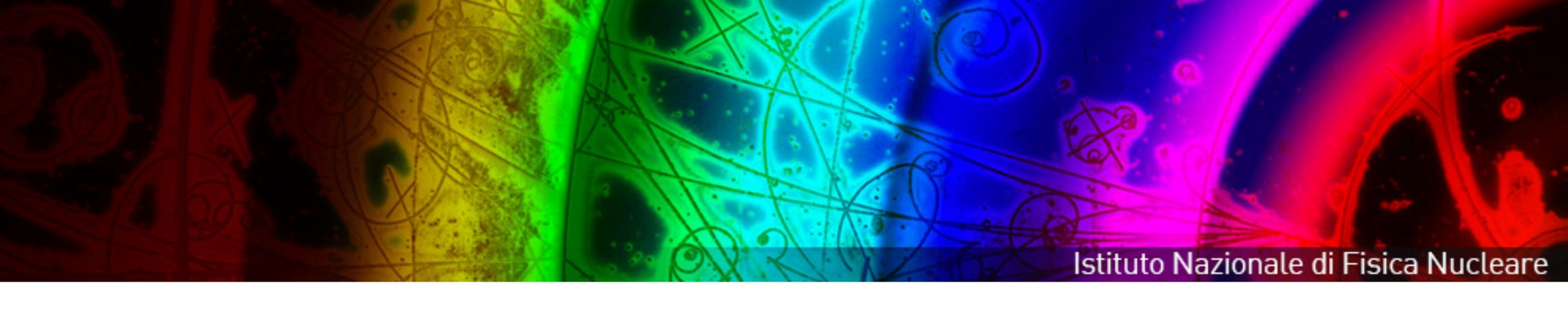
Pasquale Di Nezza







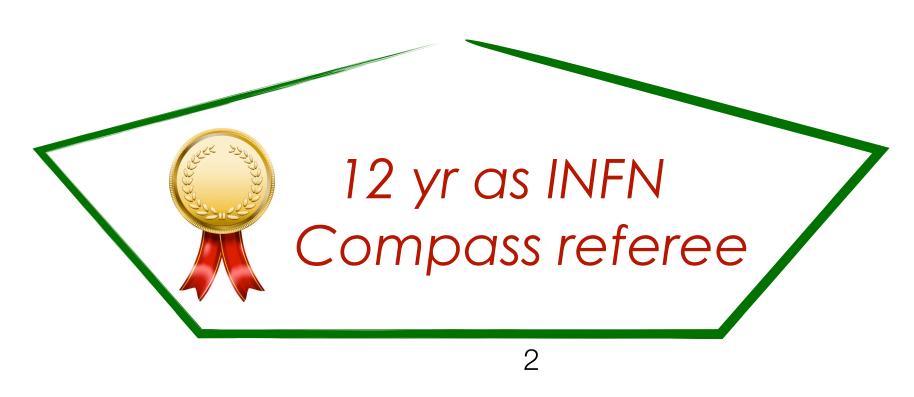




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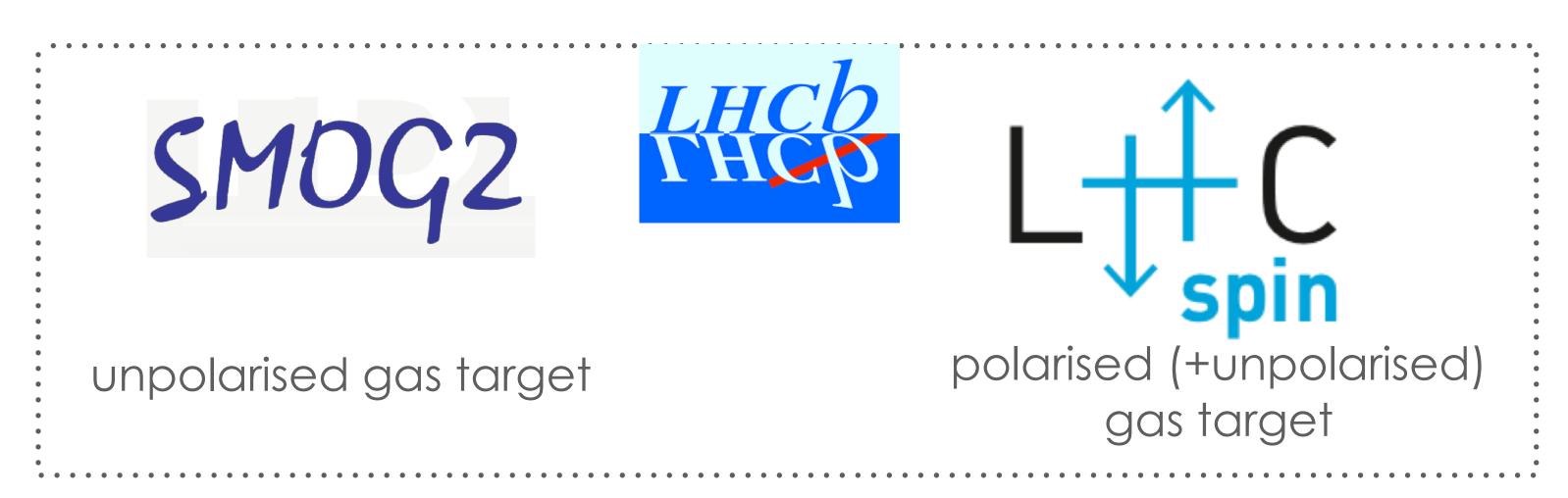




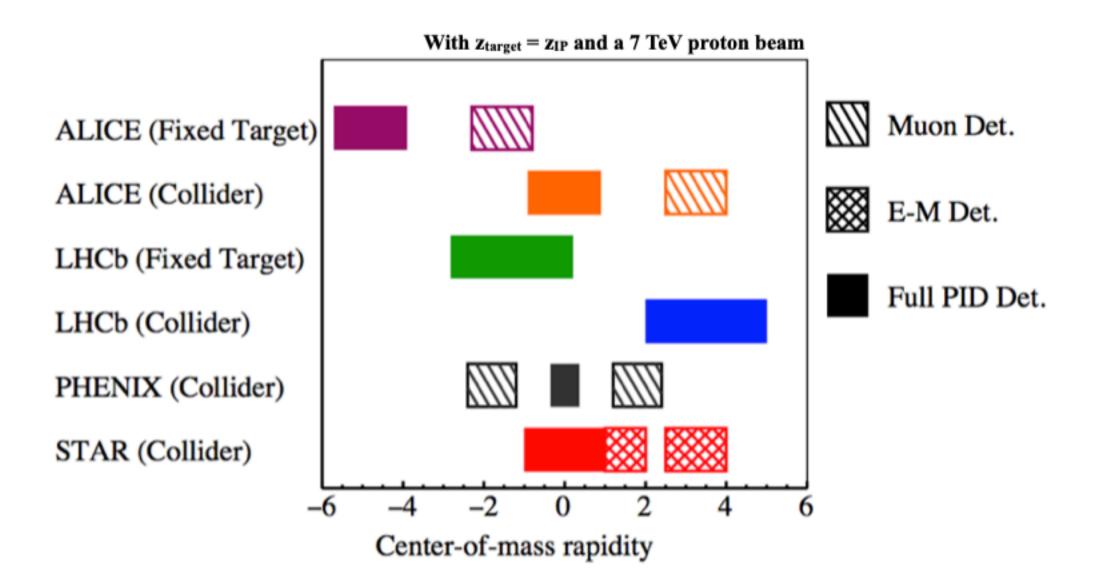
Collisions provided by a TeV-scale beam (LHC) on fixed target will exploit a unique kinematic region poorly probed. Advanced detectors make available probes never accessed before

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Acceptance in center-of-mass rapidity



gaseous targets @



The LHCb detector

- LHCb is a general-purpose forward spectrometer, fully instrumented in $2 < \eta < 5$, and optimised for c and b hadron detection
- Excellent momentum resolution with VELO + tracking stations:

$$\sigma_p/p = 0.5 - 1.0\% \ (p \in [2,200] \text{ GeV})$$

Particle identification with RICH+CALO+MUON

$$\epsilon_{\mu} \sim 98\%$$
 with $\epsilon_{\pi \to \mu} \lesssim 1\%$

• Low momentum muon trigger:

$$p_{T_u} > 1.75 \text{ GeV } (2018)$$

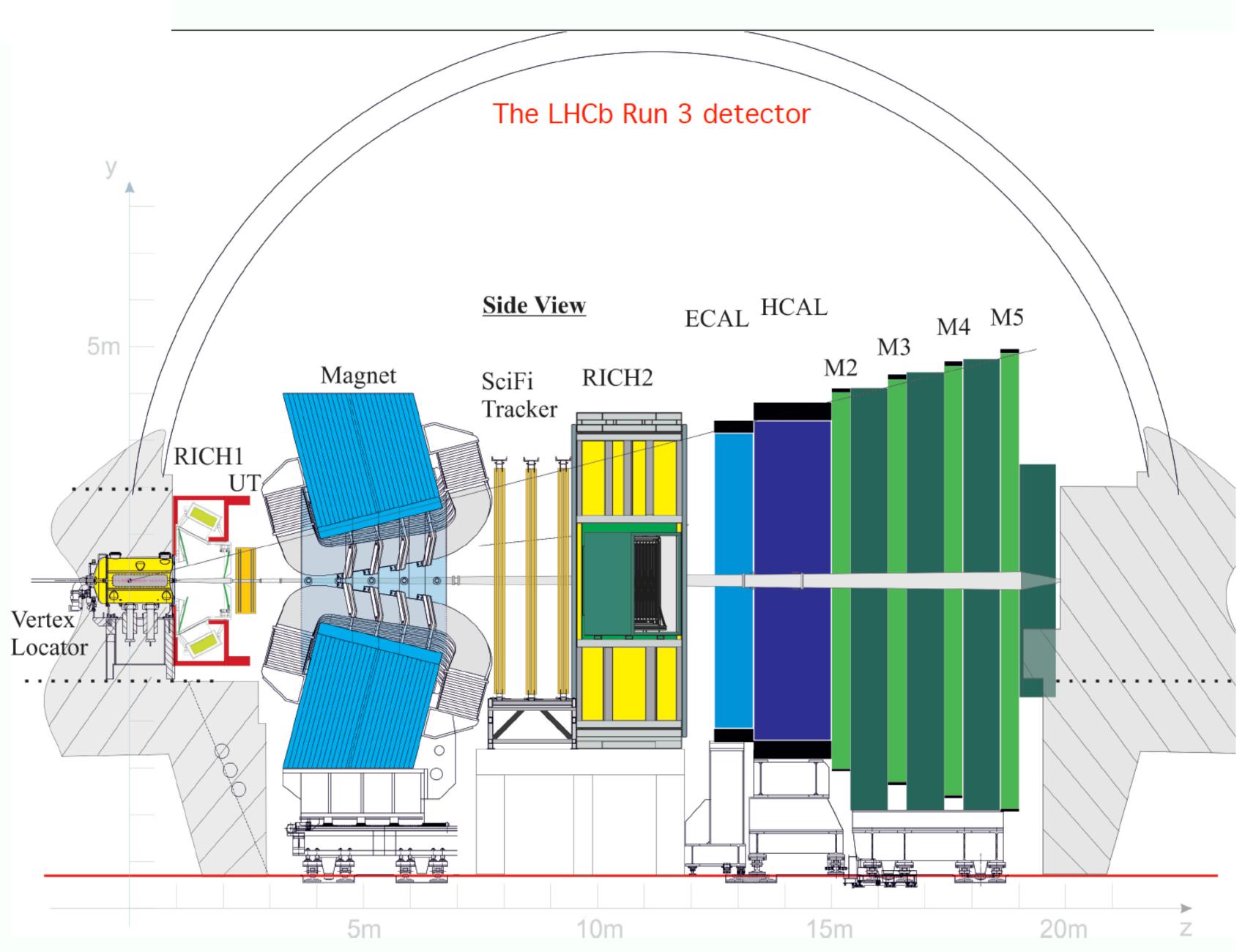
will be reduced thanks to the new fullysoftware trigger

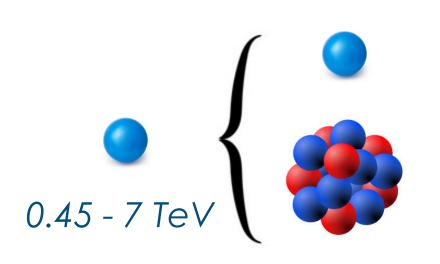
 Major detector upgrades performed during LS2 for the Run 3 (5x luminosity)

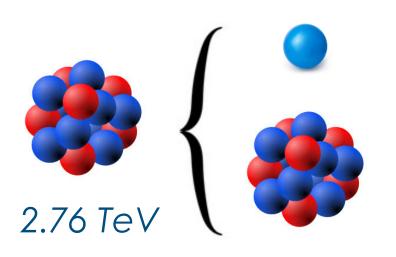
[JINST 3 (2008) S08005]

[IJMP A 30, 1530022 (2015)]

[Comput Softw Big Sci 6, 1 (2022)]







pp or pA collisions: 0.45 - 7 TeV beam on fix target

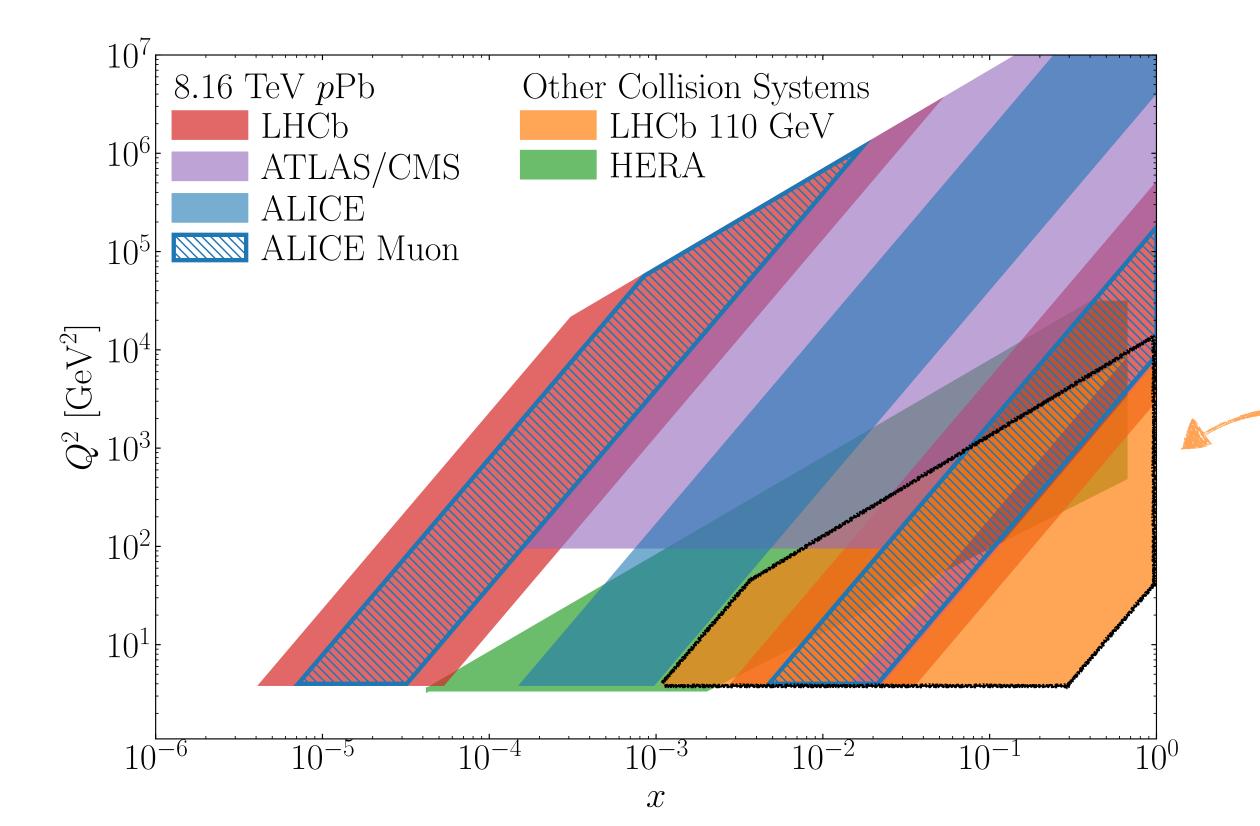
$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 41 - 115 \text{ GeV}$$

$$y_{CMS} = 0 \rightarrow y_{lab} = 4.8$$

AA collisions: 2.76 TeV beam on fix target

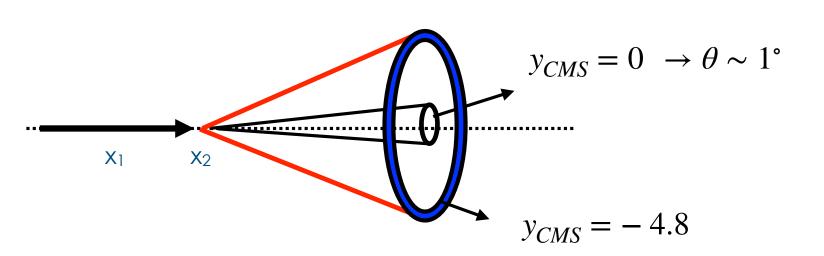
$$\sqrt{s_{NN}} \simeq 72 \; GeV$$

$$y_{CMS} = 0 \rightarrow y_{lab} = 4.3$$



1: beam; 2: target

Large CM boost, large x_2 values ($x_F < 0$) and small x_1



$$\gamma = \frac{\sqrt{s_{NN}}}{2m_p} \simeq 60$$

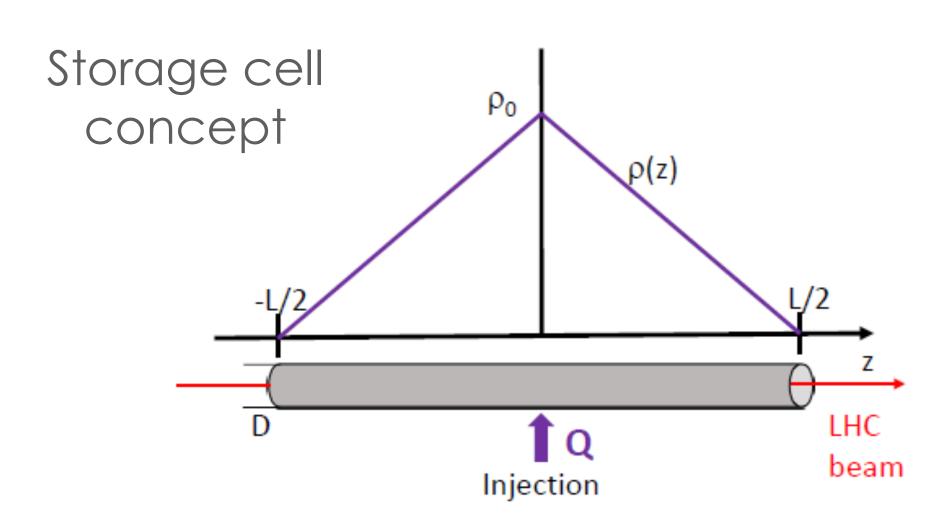
Broad and poorly explored kinematic range

SMOG2 an unpolarised target at

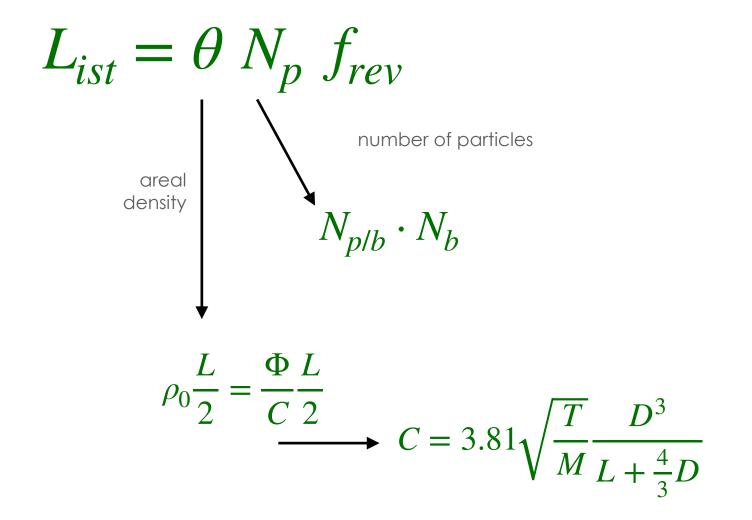


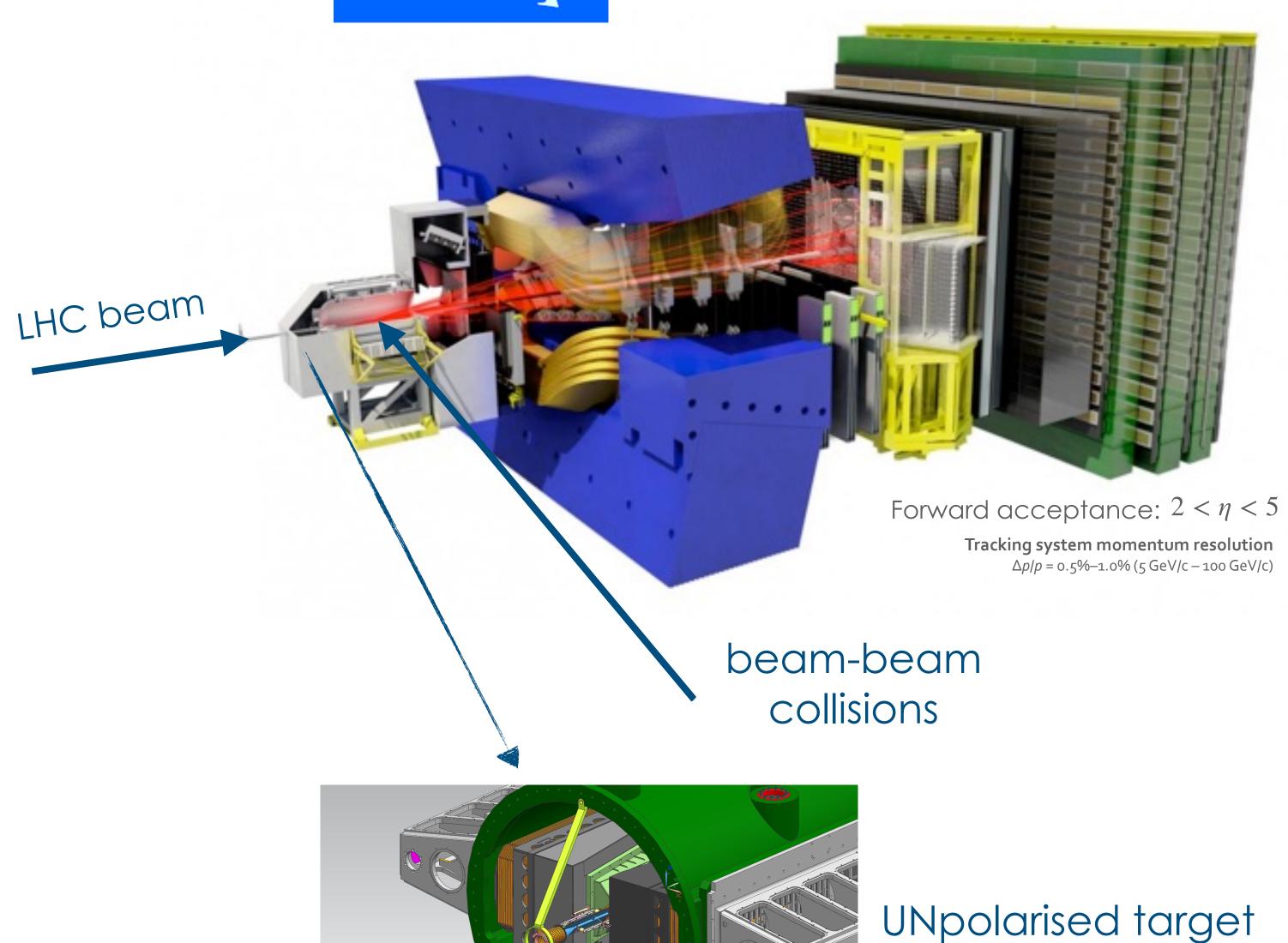
JINST 3 (2008) S08005 IJMPA 30 (2015) 1530022

(beam-gas)





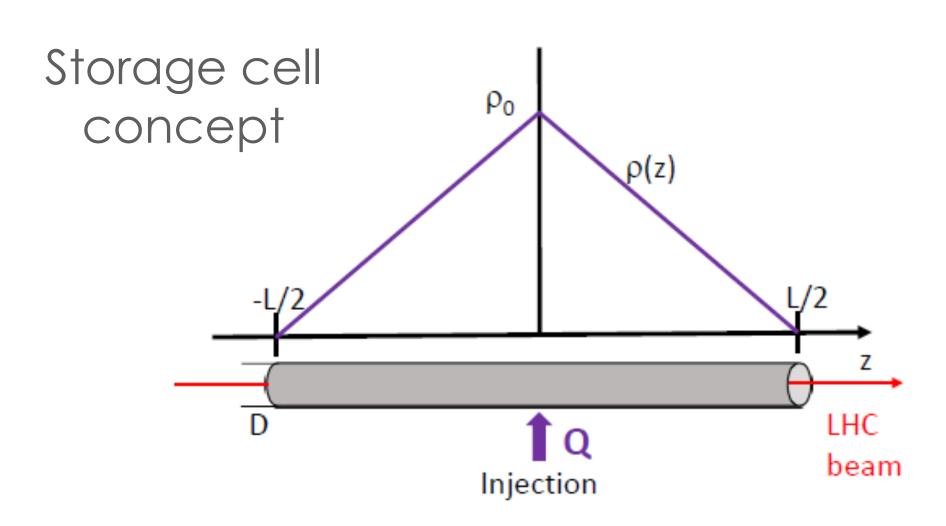




SMOG2 an unpolarised target at



JINST 3 (2008) S08005 IJMPA 30 (2015) 1530022



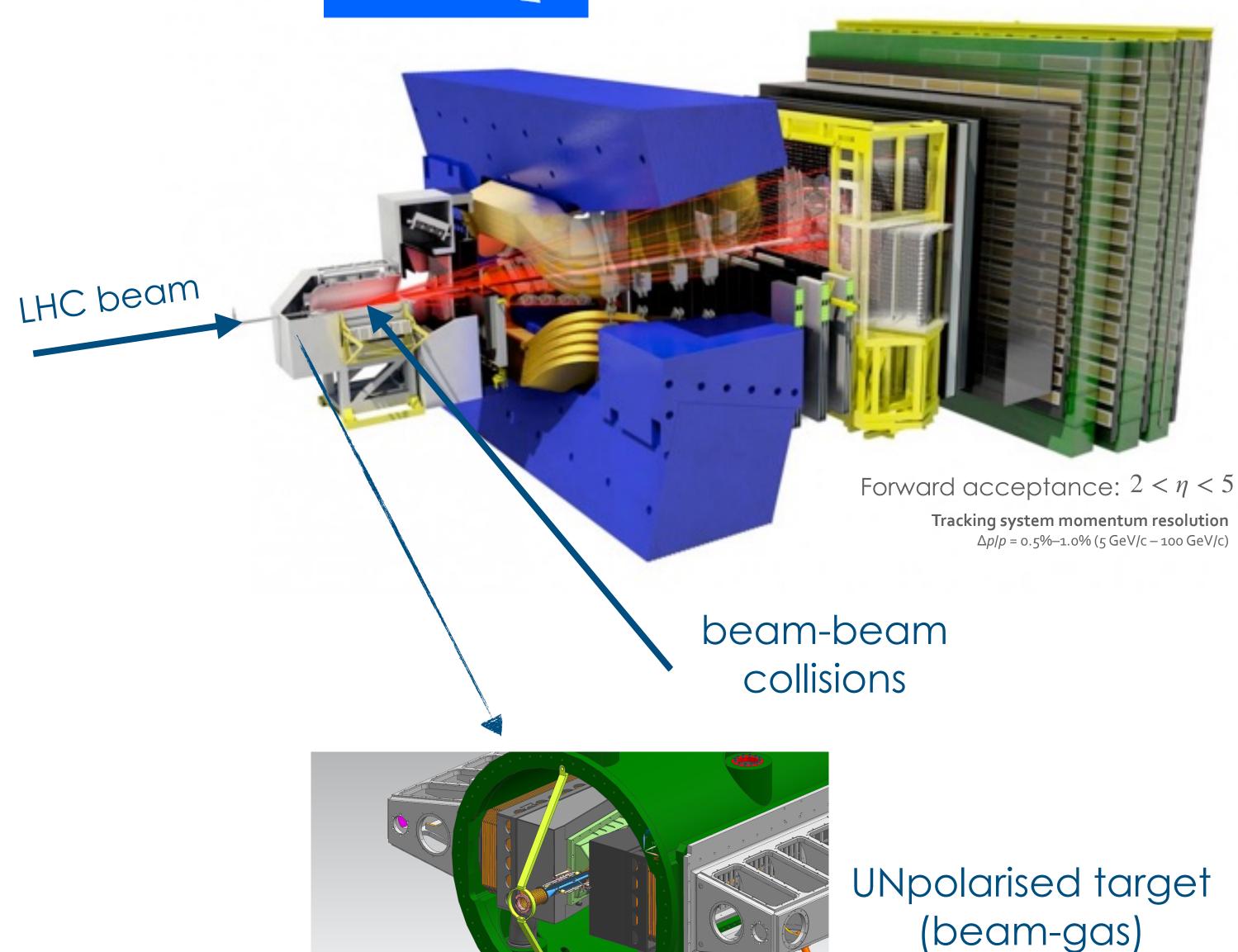
Luminosity

$$L_{ist} = \theta N_p f_{rev}$$

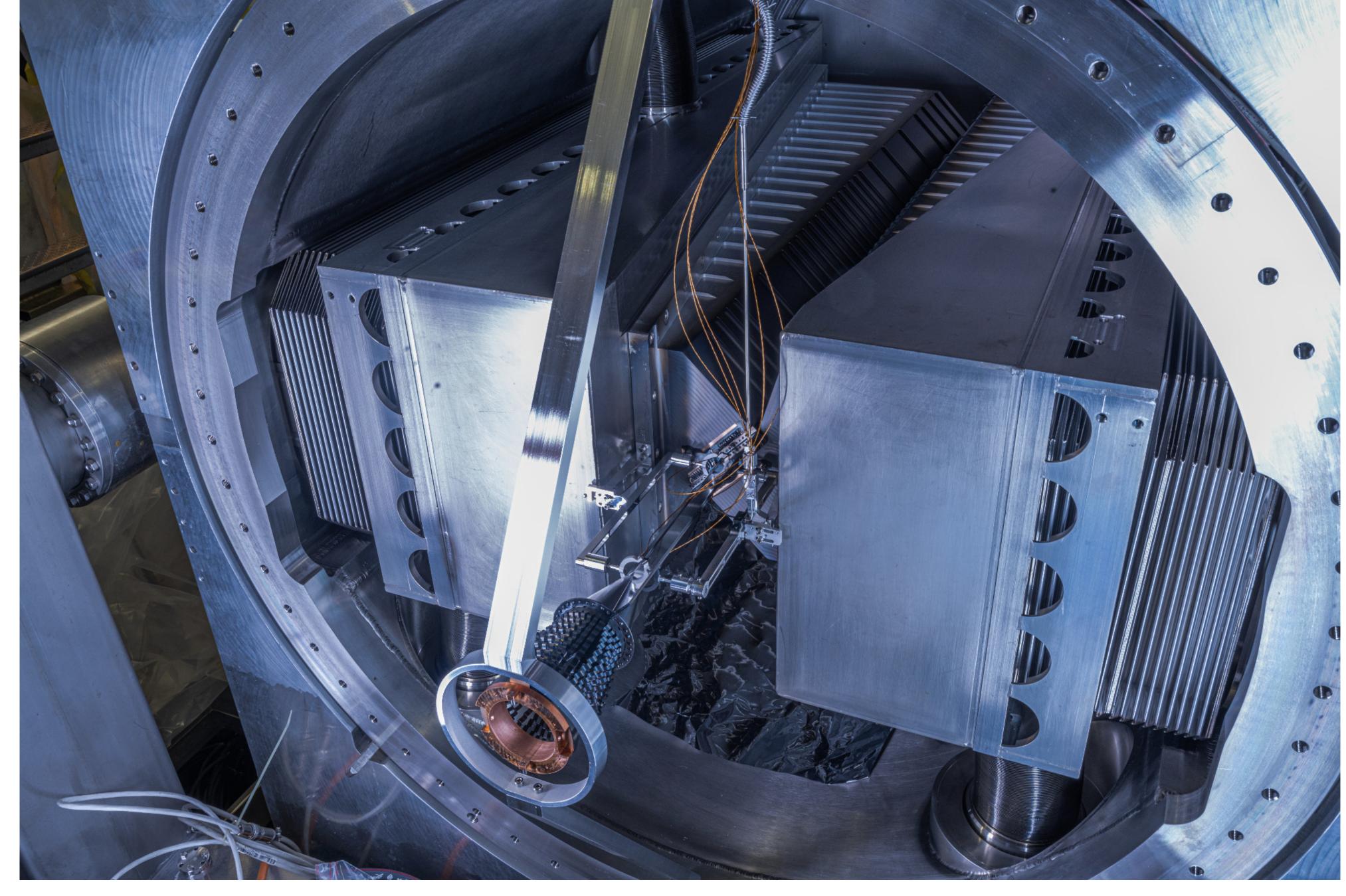
$$\begin{array}{c|c} \text{number of particles} \\ N_{p/b} \cdot N_b \end{array}$$

$$\rho_0 \frac{L}{2} = \frac{\Phi}{C} \frac{L}{2}$$

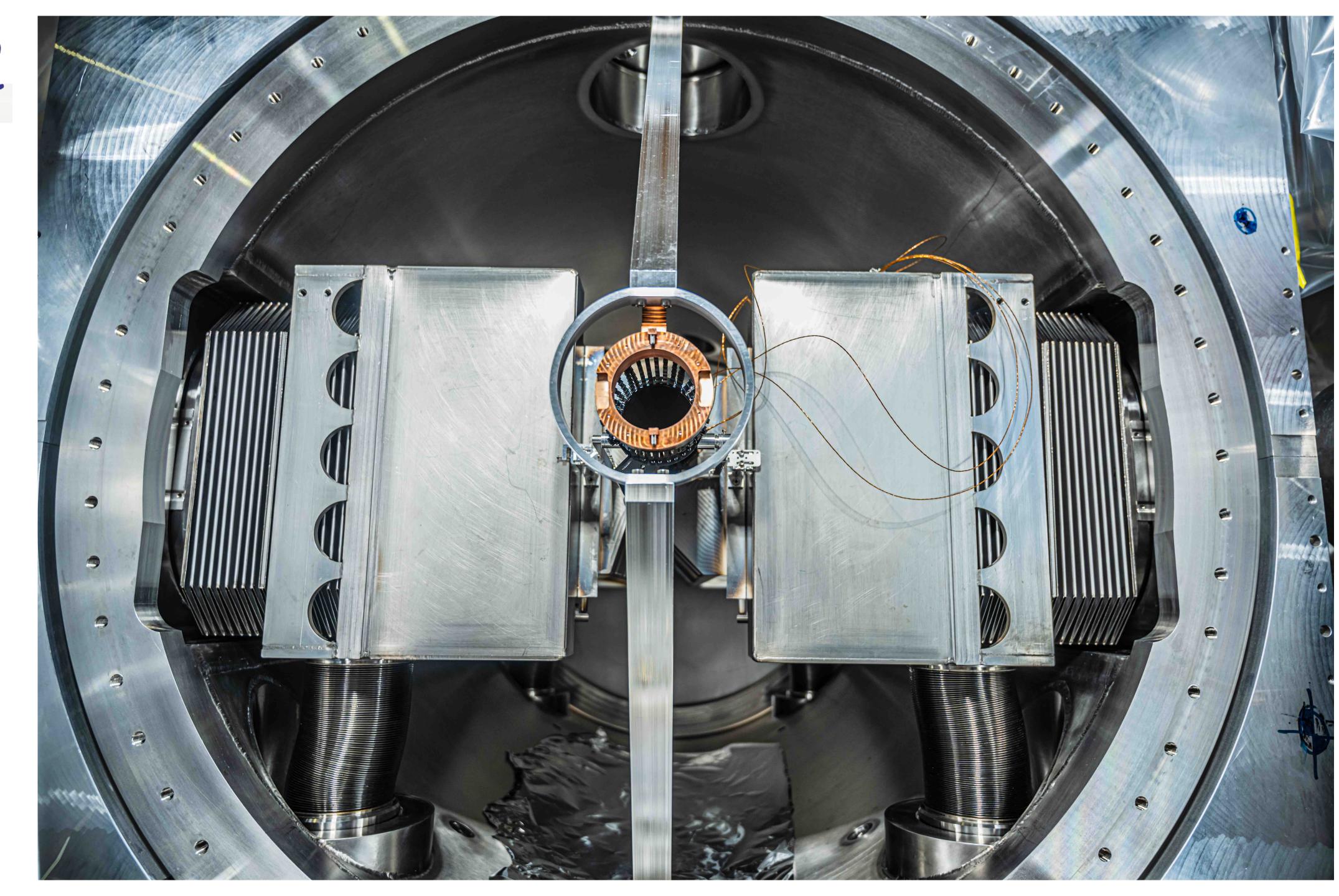
$$C = 3.81 \sqrt{\frac{T}{M}} \frac{D^3}{L + \frac{4}{3}D}$$



SMOG2

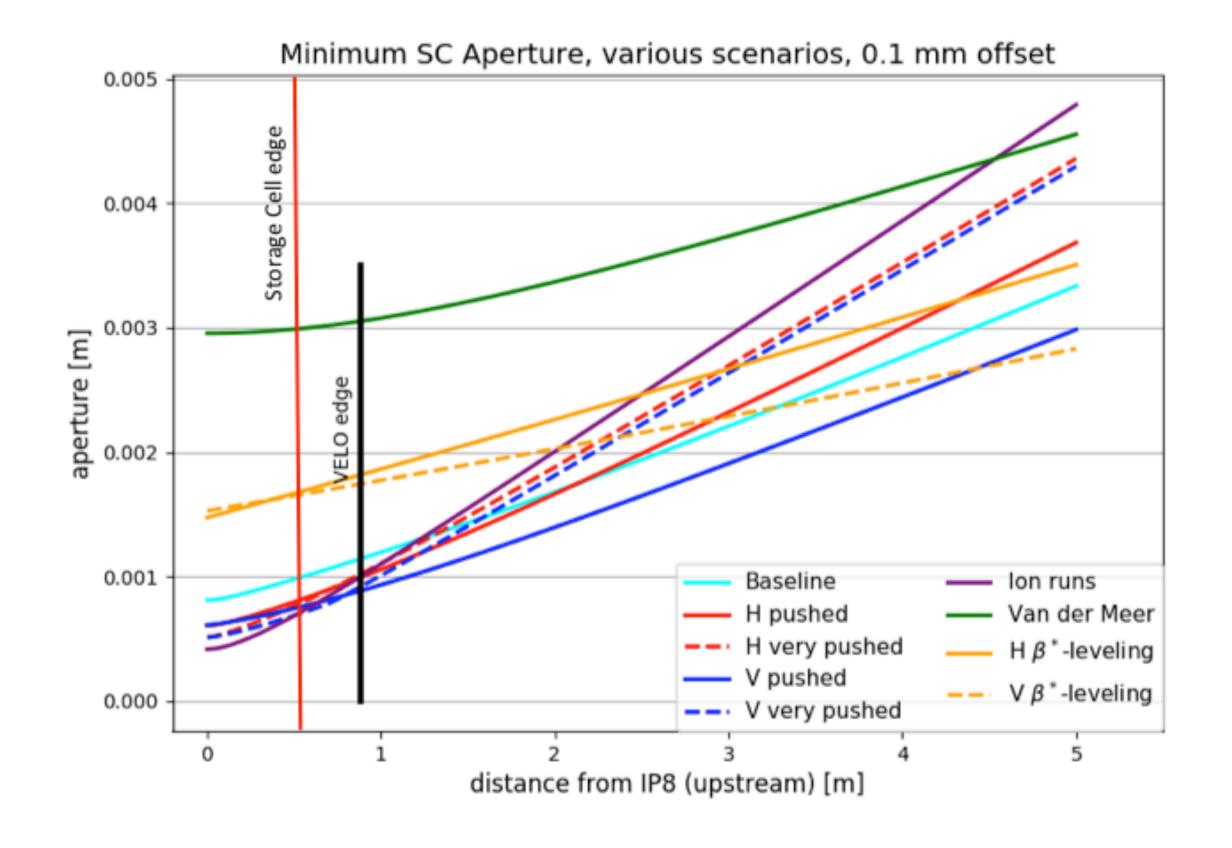


SMOG2



SMOG2 **Openable** Storage Cell dimensions: 200 mm (L) - 10 mm (D)

The physical aperture of 5 mm gives 2 mm margin to minimum allowed aperture (VdM scan)



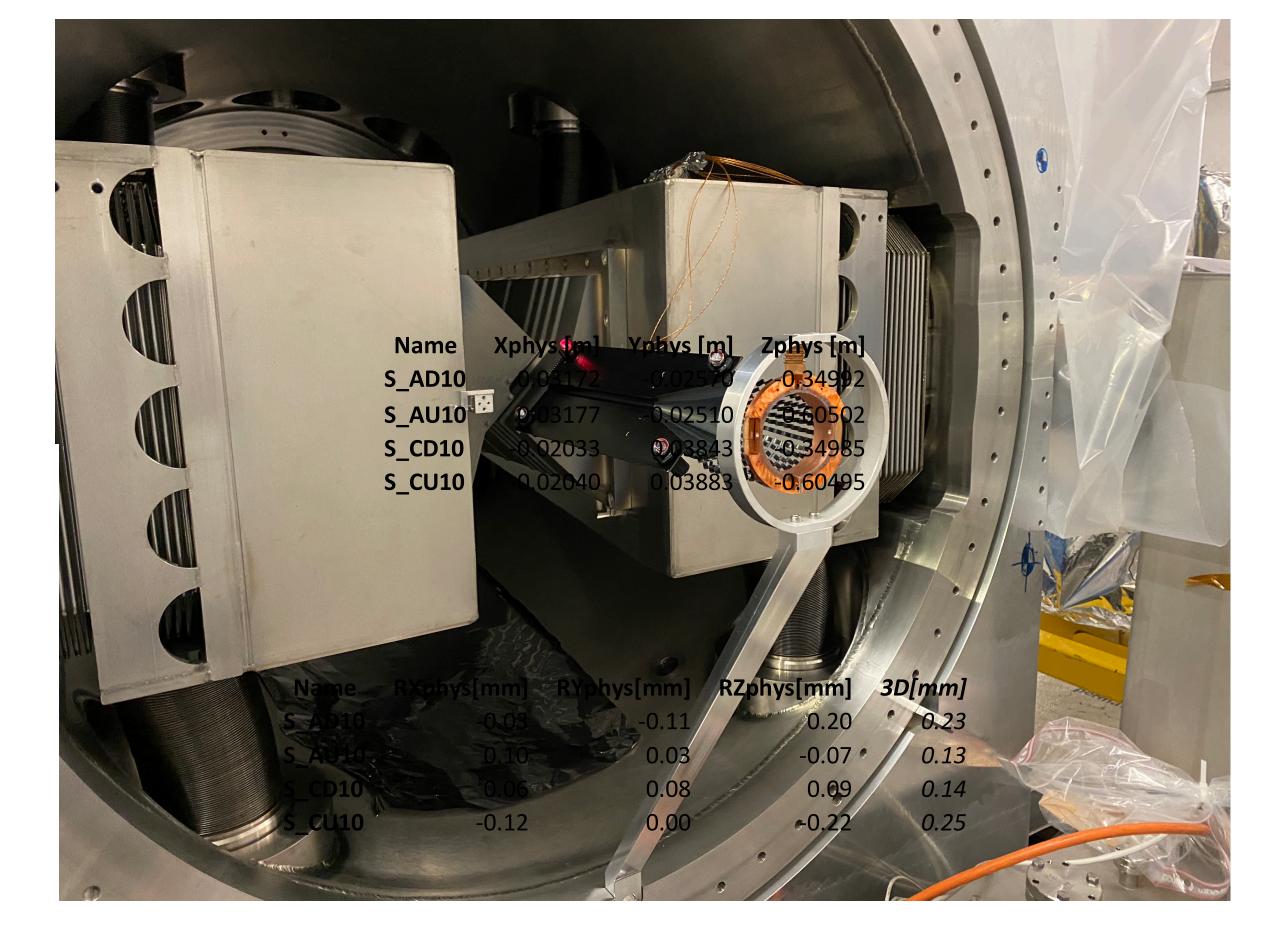


Table 9: Final position of the SMOG2 and its offset to the nominal position

Position of SMOG2				Offset to nominal		
Name	Xphys [m]	Yphys [m]	Zphys [m]	dXphys [mm]	dYphys [mm]	dZphys [mm]
S_E	-0.00142	-0.00017	-0.61739	-0.25	0.14	0.11
S_S	-0.00136	-0.00040	-0.33739	-0.19	-0.14	0.11
S_ROLL	-0.00082	0.99983	-0.61658			

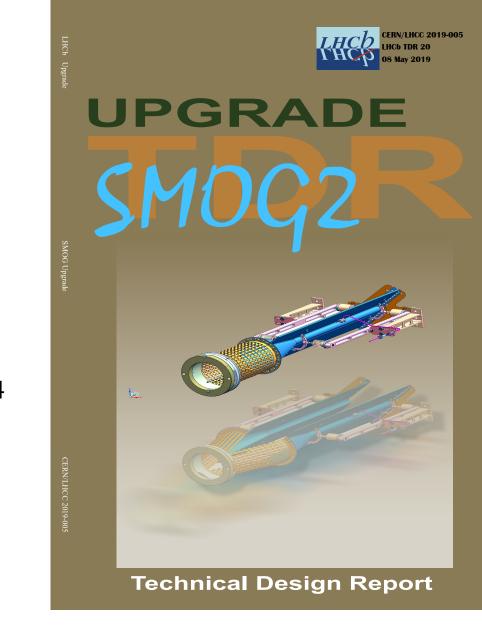
Excellent alignment reached

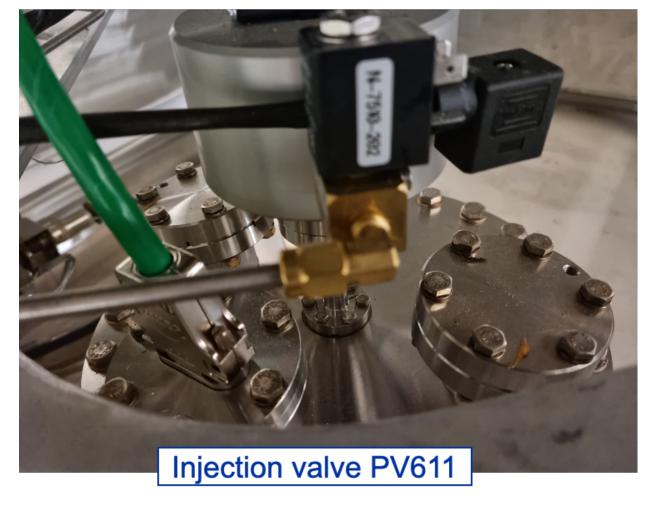
SMOG2



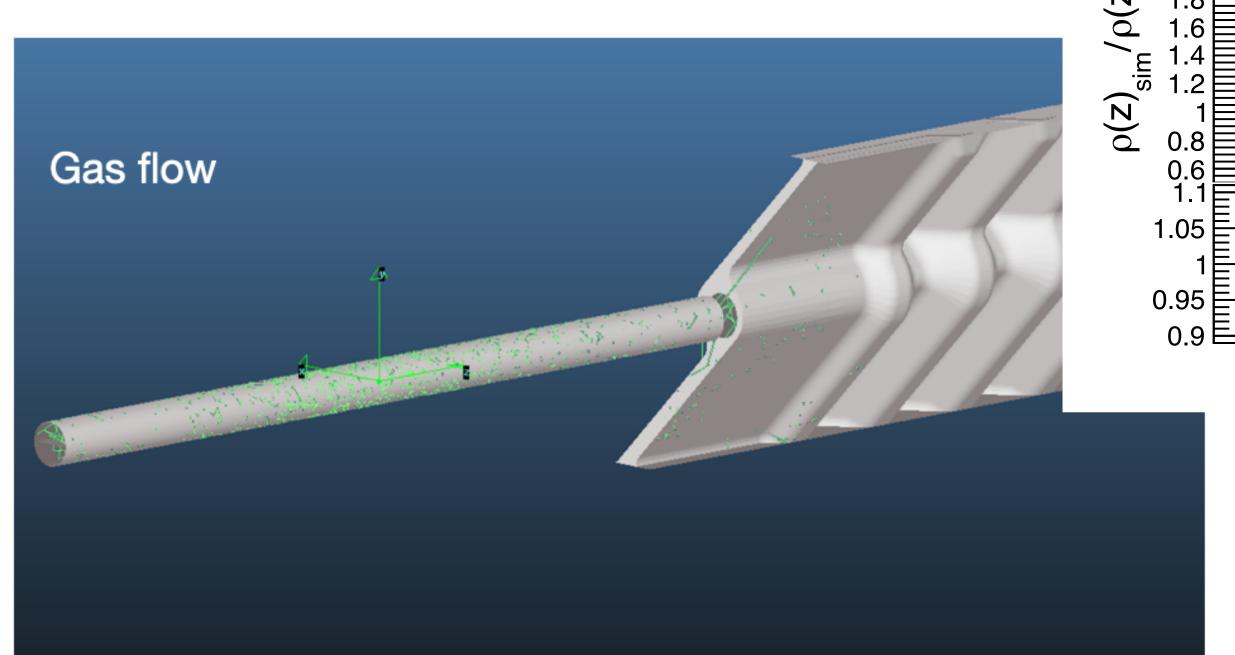
- The system is completely installed (storage cell + GFS + triggers + reconstruction)
- Negligible impact on the beam lifetime ($\tau_{beam-gas}^{p-H_2}\sim 2000$ days , $\tau_{beam-gas}^{Pb-Ar}\sim 500$ h)
- Injectable gases: He, Ne, Ar ... H₂, D₂, N₂, O₂, Kr, Xe
- Flux known with 1% precision, measured relative contamination 10-4

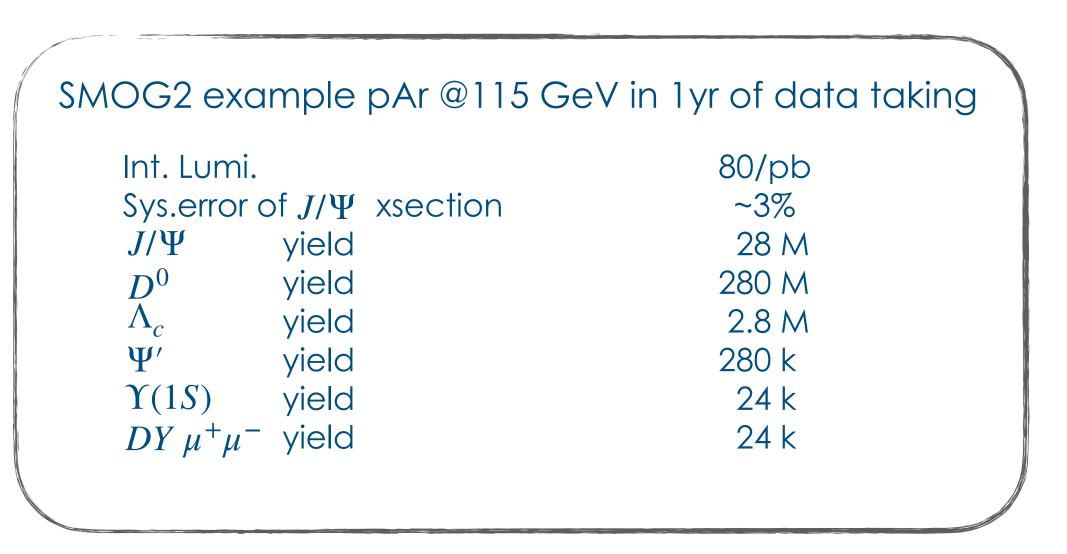


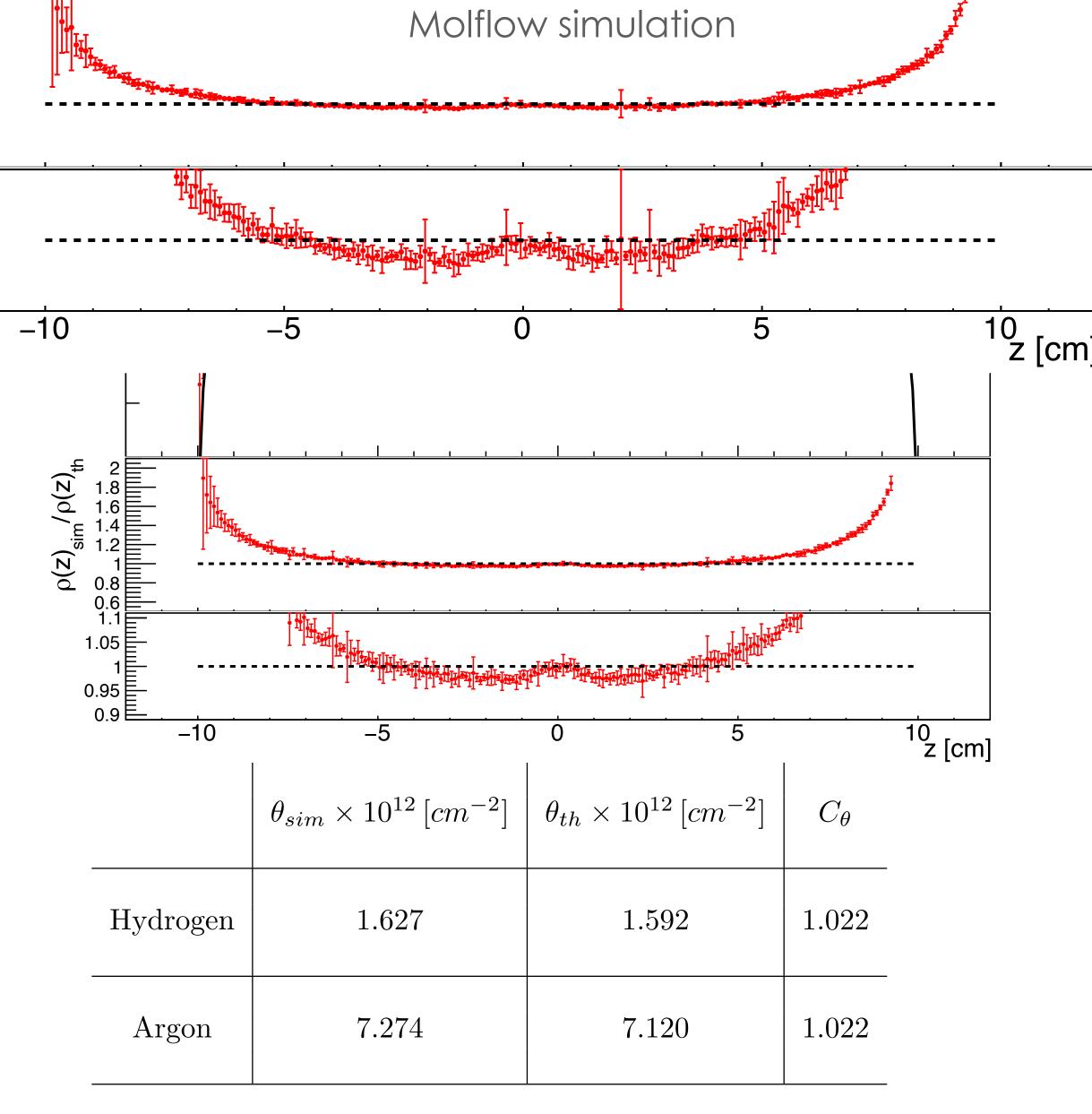




The storage cell advantage: up to x50 density wrt free flow Molflow simulation





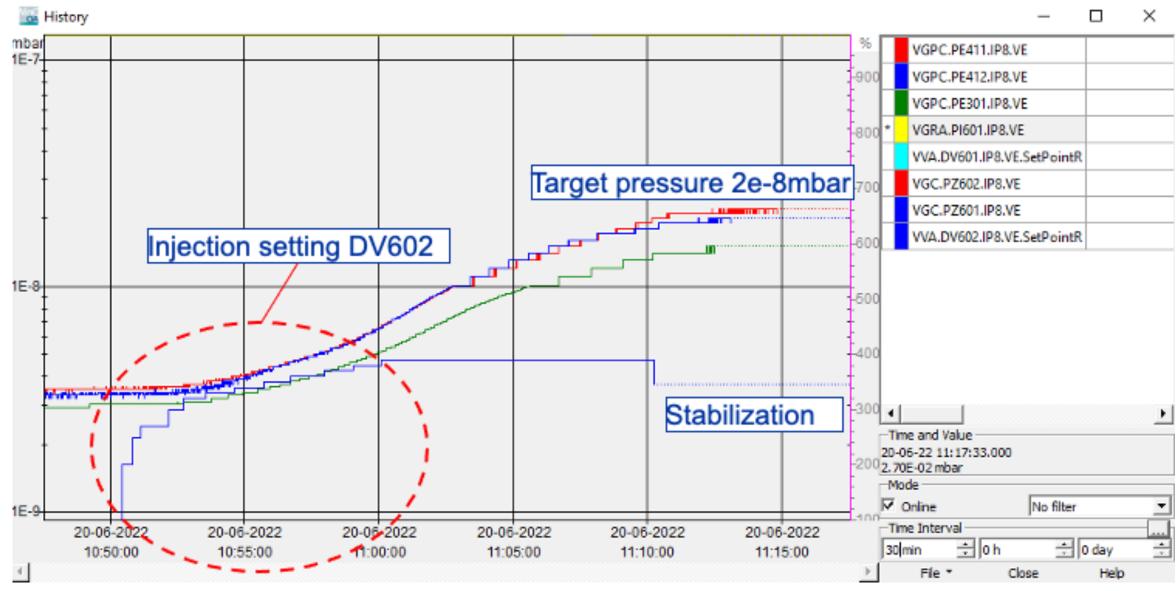


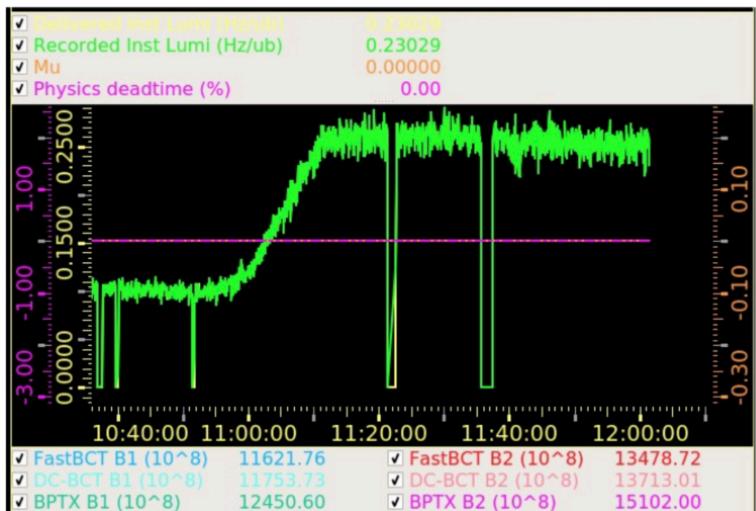
Systematic uncertainty on Luminosity: 1.5%

Very high statistics with a low gas flow

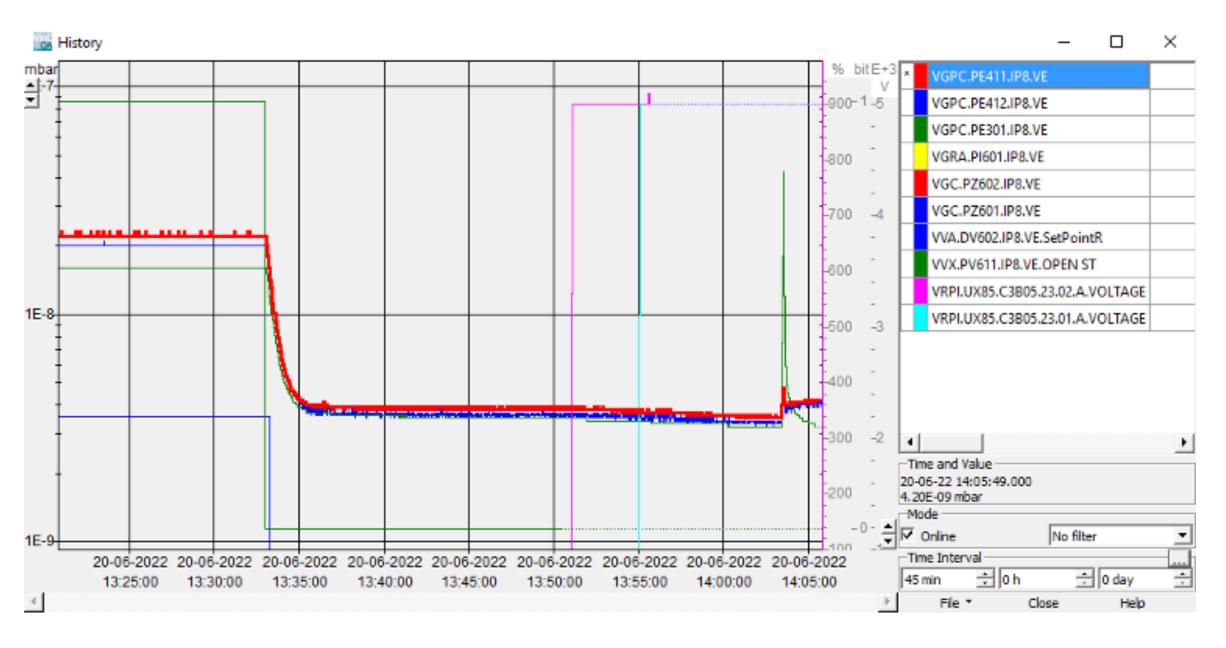
SMOG2 gas injection at LHC Run3 started in June 2022

Pressure increase into the primary vacuum





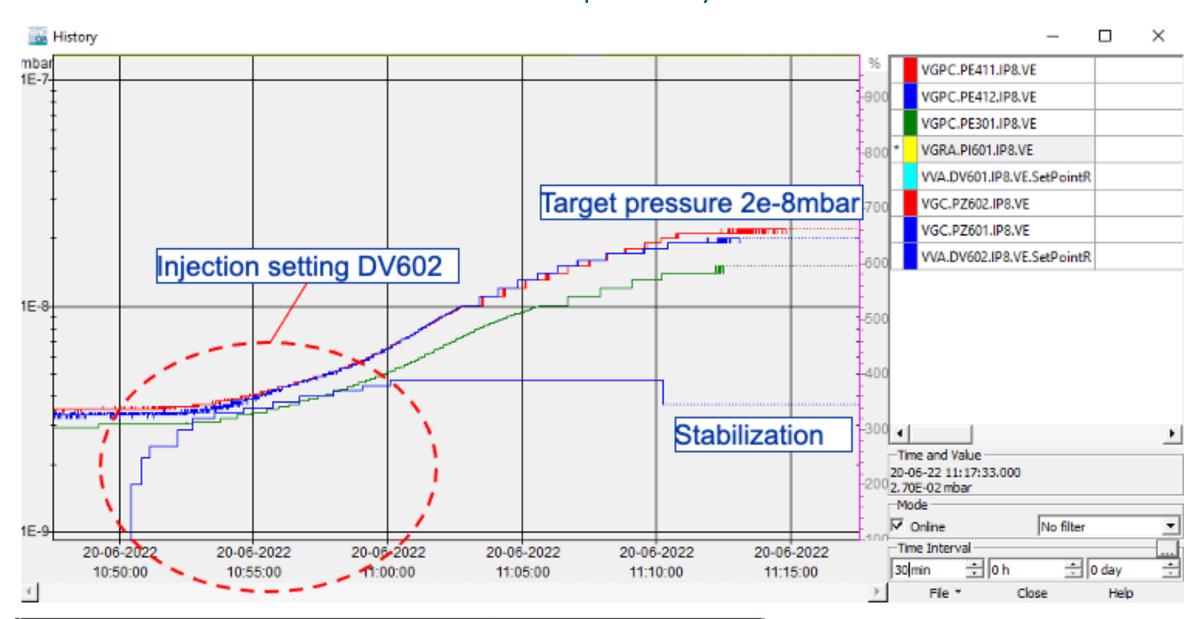
Luminosity increase seen by the LHCb luminometer (Plume) Vacuum recovery after the gas injection stop

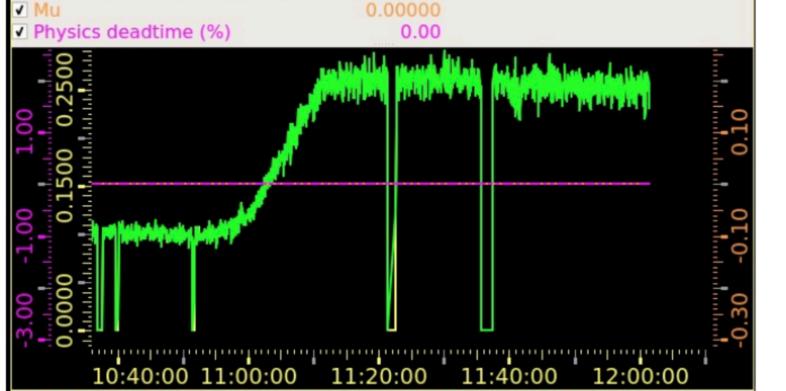


Extremely useful also for the LHCb commissioning

SMOG2 gas injection at LHC Run3 started in June 2022

Pressure increase into the primary vacuum





▼ BPTX B2 (10⁸)

15102.00

11621.76

12450.60

0.23029

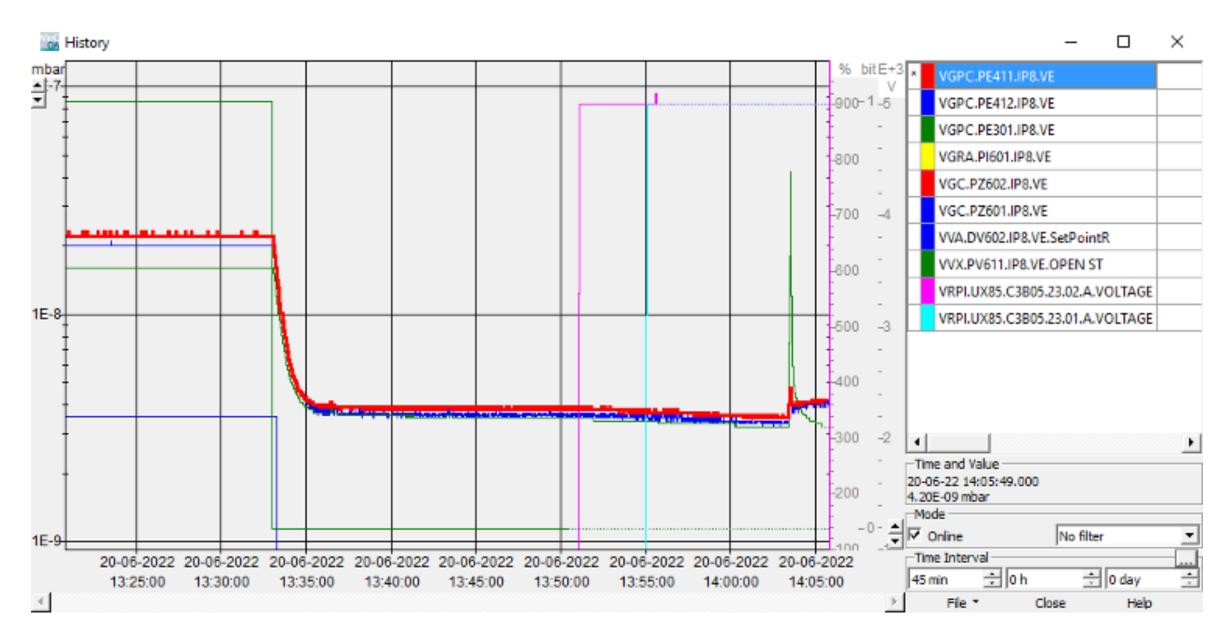
Recorded Inst Lumi (Hz/ub)

▼ BPTX B1 (10^8)

Luminosity increase seen by the LHCb luminometer (Plume)

Extremely useful also for the LHCb commissioning

Vacuum recovery after the gas injection stop



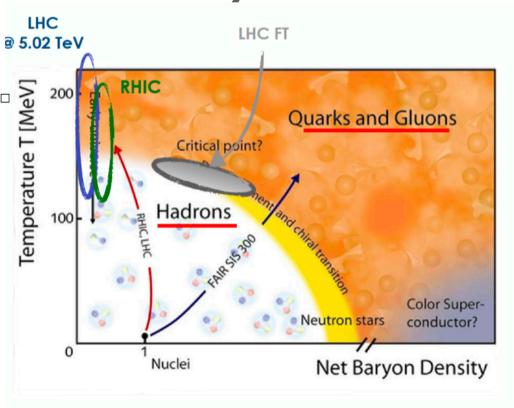
LHC official statement

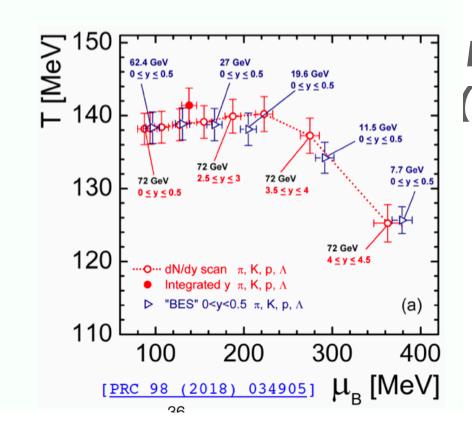
No negative feedback when there is gas injection. Green light to inject when needed

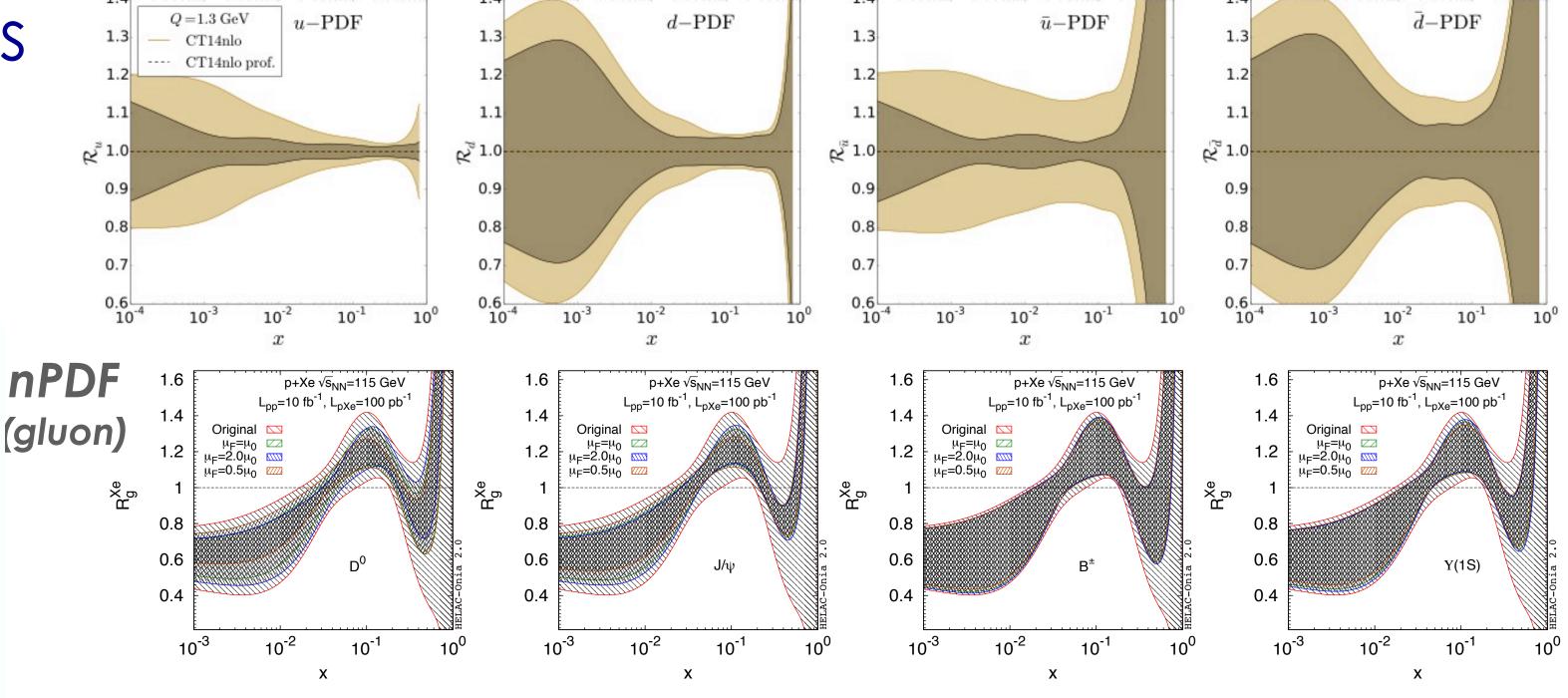
SMOC2... few highlights

http://cds.cern.ch/record/2649878/files/

Heavy-Ion and QCD phase space

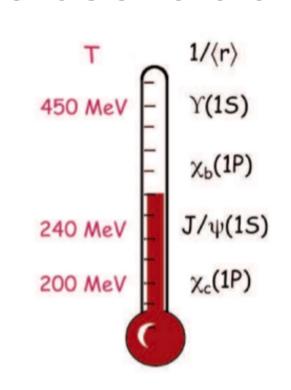




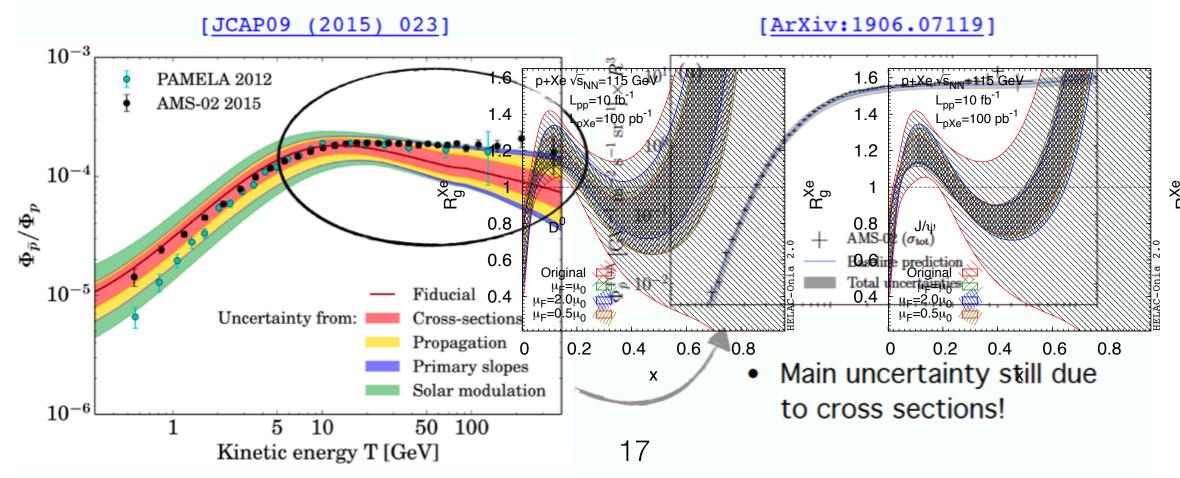


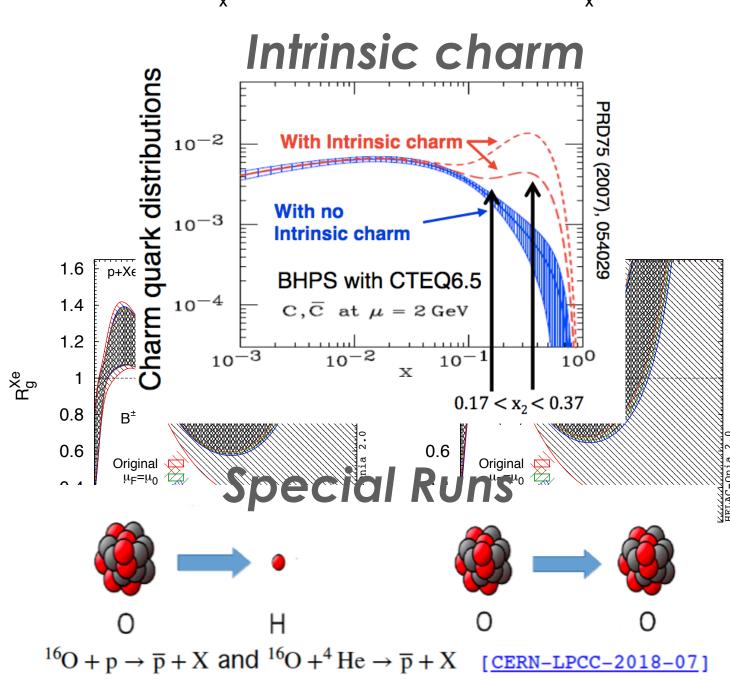
estimation with 10 fb⁻¹

$c\bar{c}$ bound states









arXiv:1807.00603

Not only standard but also special runs

Run 3

M.Lamont's talk at ICHEP22

Proton-proton

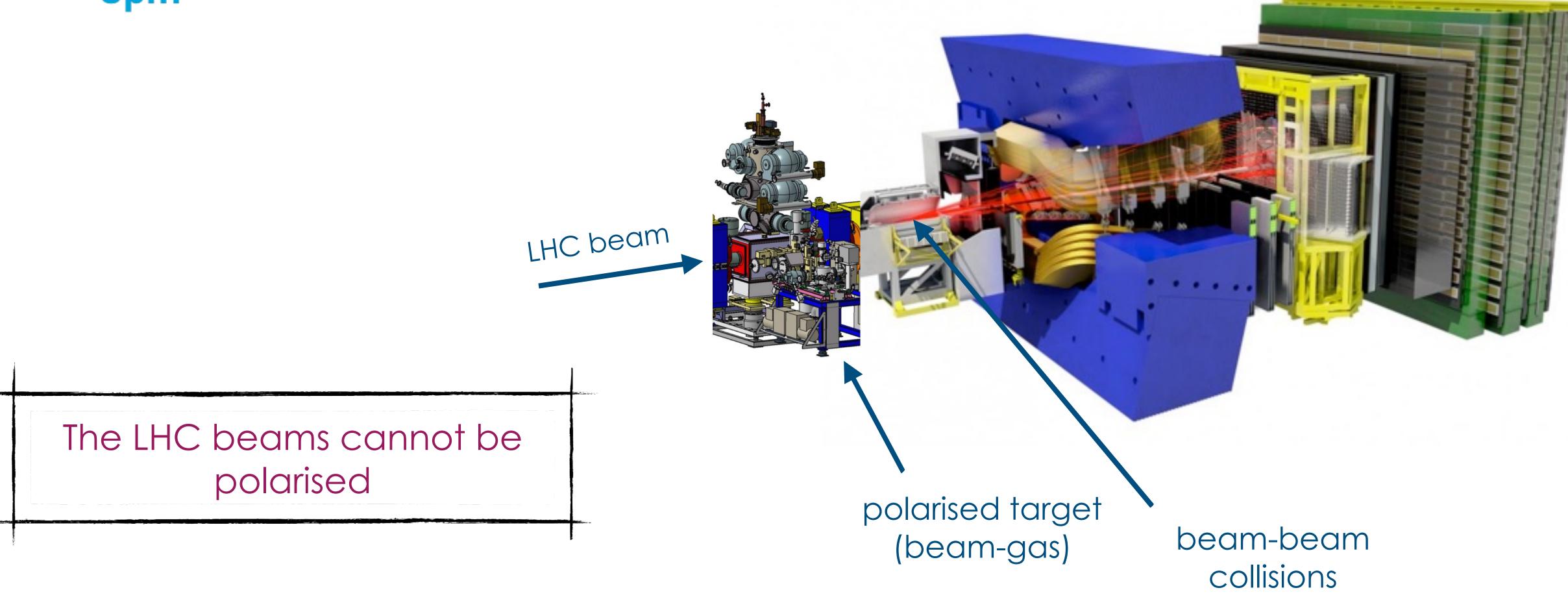
- 6.8 TeV
- Levelled to a maximum luminosity 2.05×10^{34} cm⁻²s⁻¹ in ATLAS and CMS
- Levelled to a target of $^{\sim}1.4 \times 10^{31}$ cm⁻²s⁻¹ and 2×10^{33} cm⁻²s⁻¹ in ALICE and LHCb respectively
- ~1.8×10¹¹ protons/bunch in 2023 2025 long levelling times!

Mode	ATLAS/CMS	LHCb	ALICE
proton-proton	250 - 270 fb ⁻¹	25 – 30 fb ⁻¹	200 pb ⁻¹
lead-lead	7 nb ⁻¹	1 nb ⁻¹	7 nb ⁻¹
proton-lead	0.5 pb ⁻¹	0.1 pb ⁻¹	0.25 pb ⁻¹
oxygen-oxygen	0.5 nb ⁻¹	0.5 nb ⁻¹	0.5 nb ⁻¹
proton-oxygen	LHCf 1.5 nb ⁻¹	2.0 nb ⁻¹	

Special run type	Experiment
VdM scans, etc.	all
Low-PU (<0.02) pp	LHCf
High β*(90m) pp	TOTEM
High, β*(3/6km) pp	TOTEM, ATLAS
p-He(SMOG) @ 450 GeV	LHCb

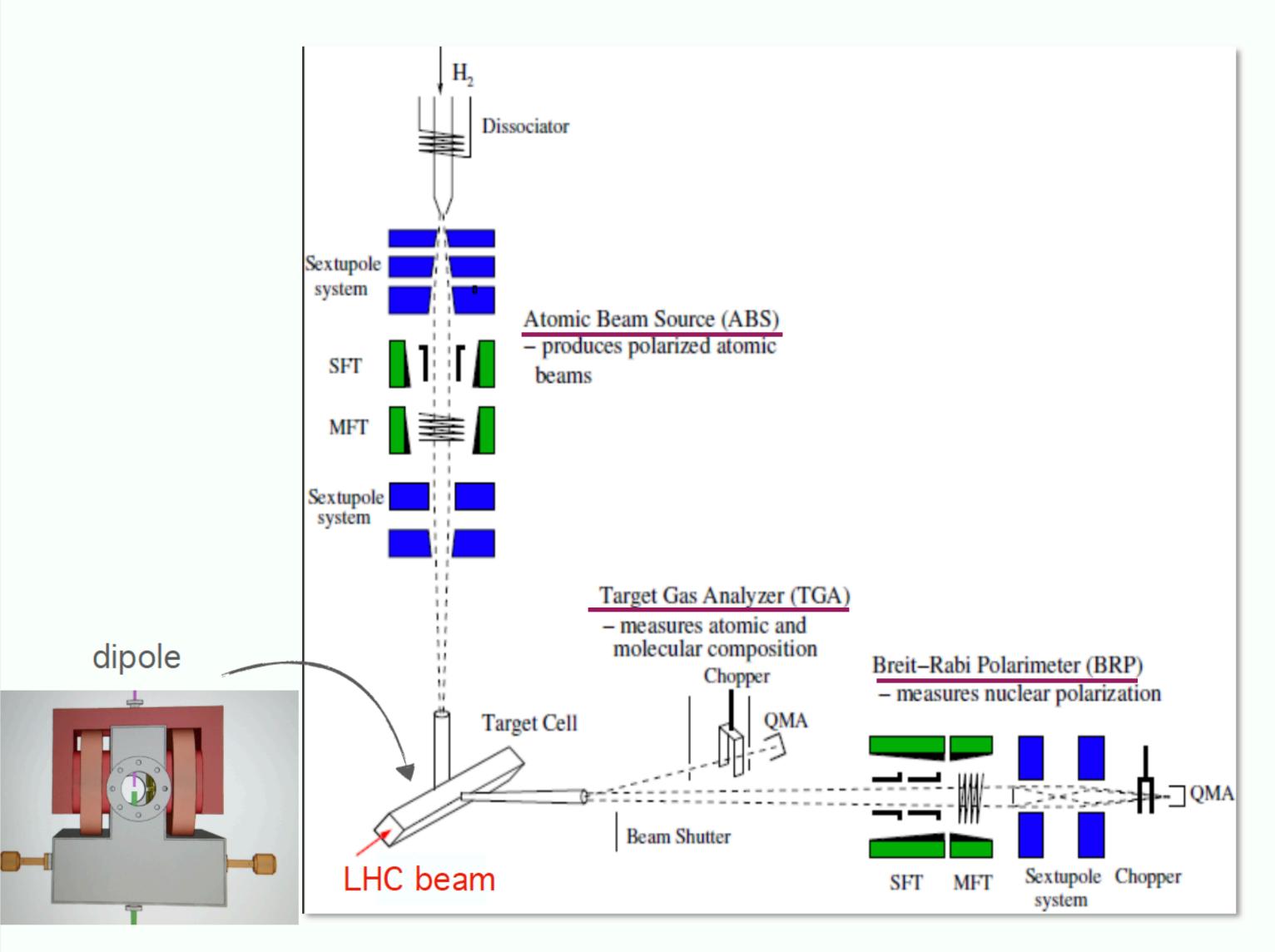




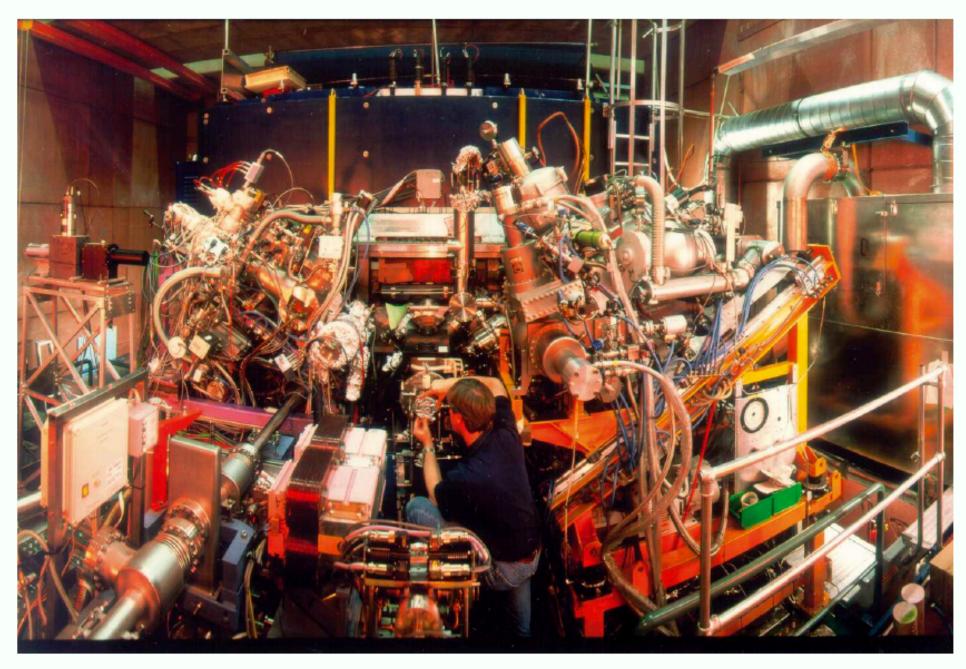


SMOG2 is not only a unique project itself, but also a great playground for L_{spin}

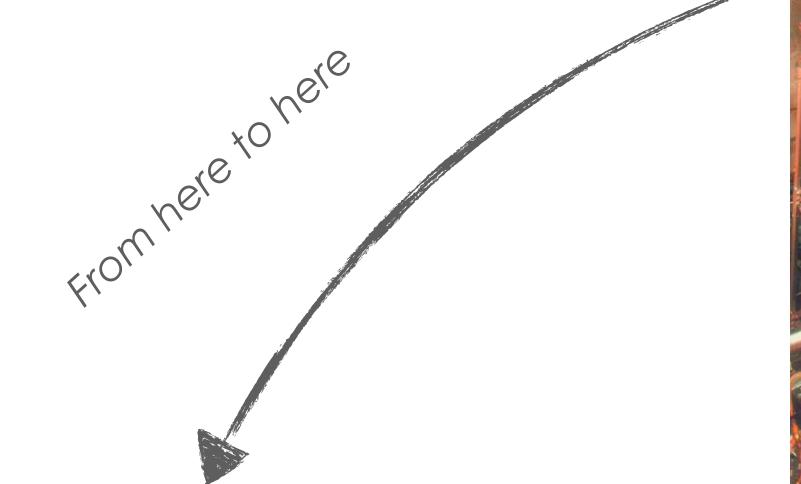
LHCspin experimental setup

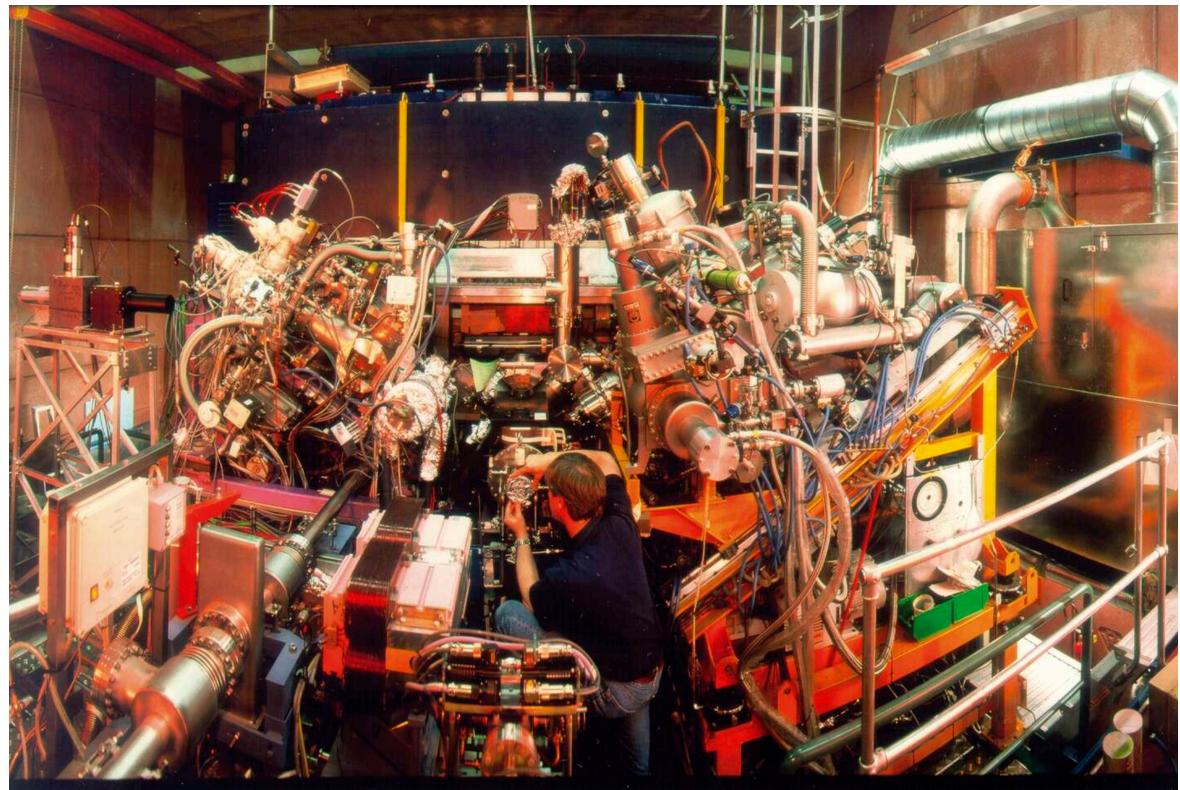


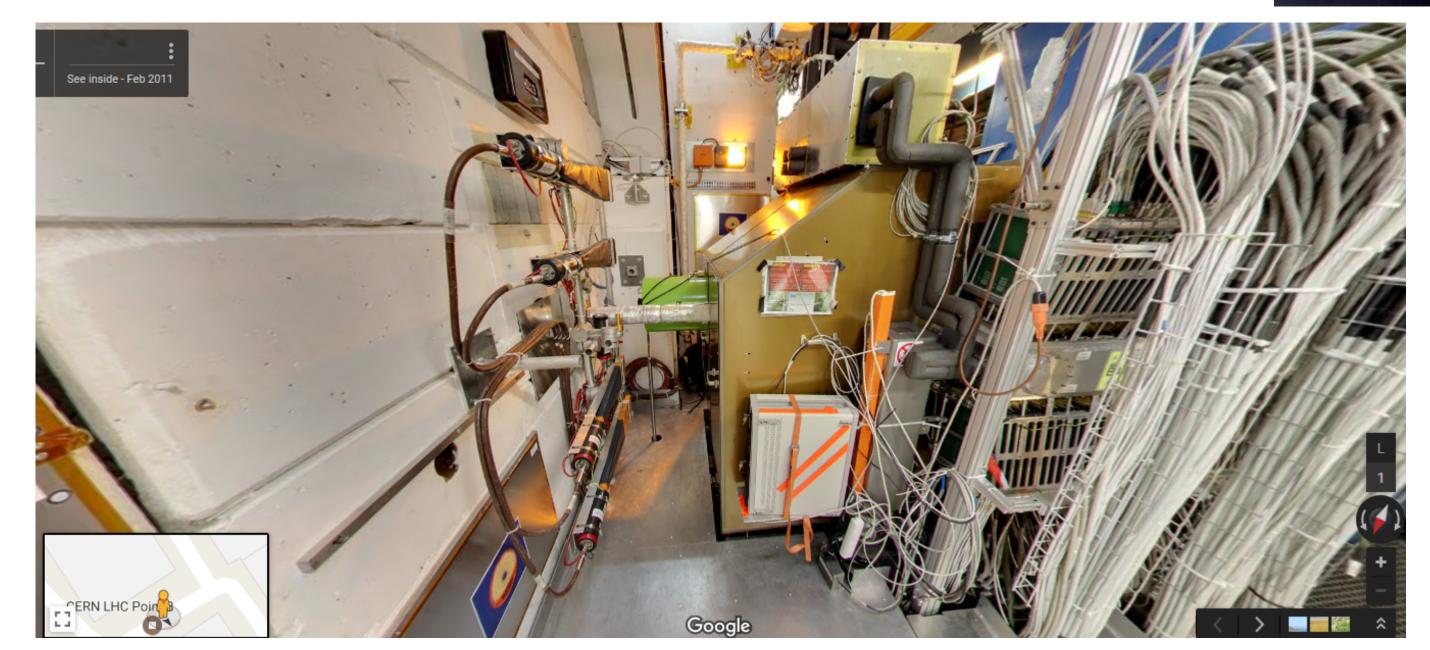
- Start from the well established HERMES setup @ DESY...
- ... to create the next generation of fixed target polarisation techniques!



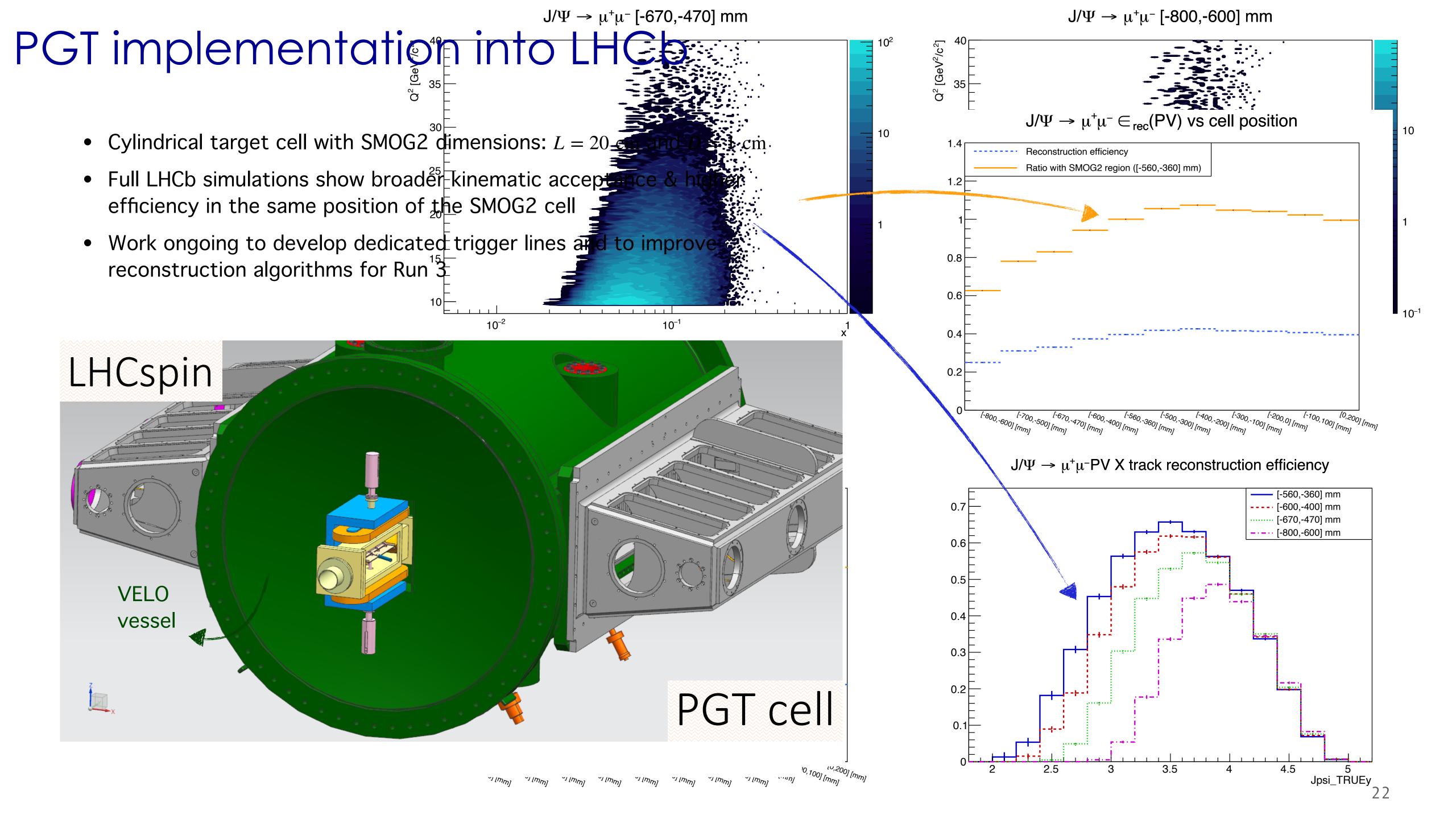
HERMES PGT





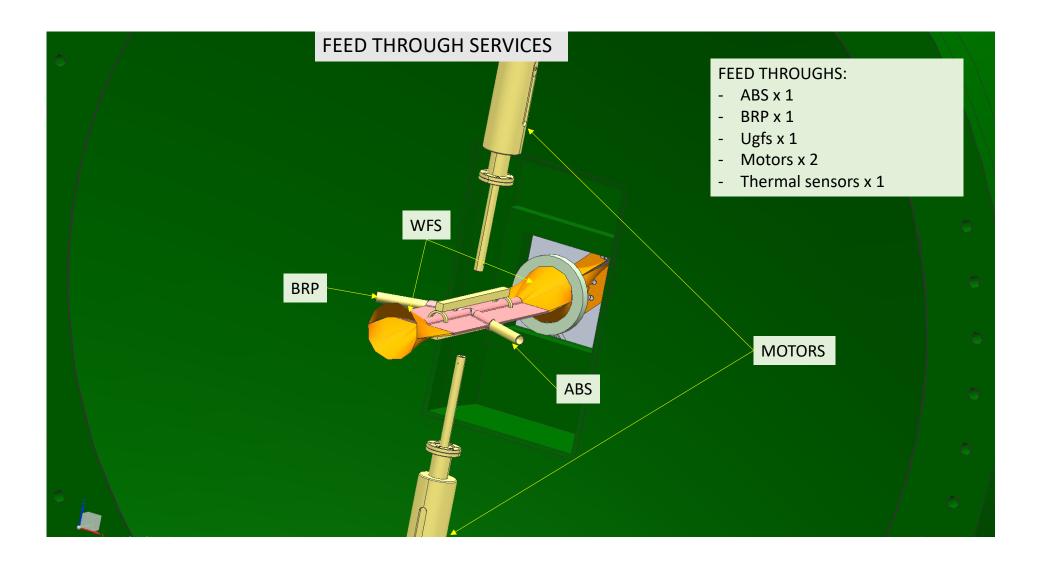


Space available in front of LHCb



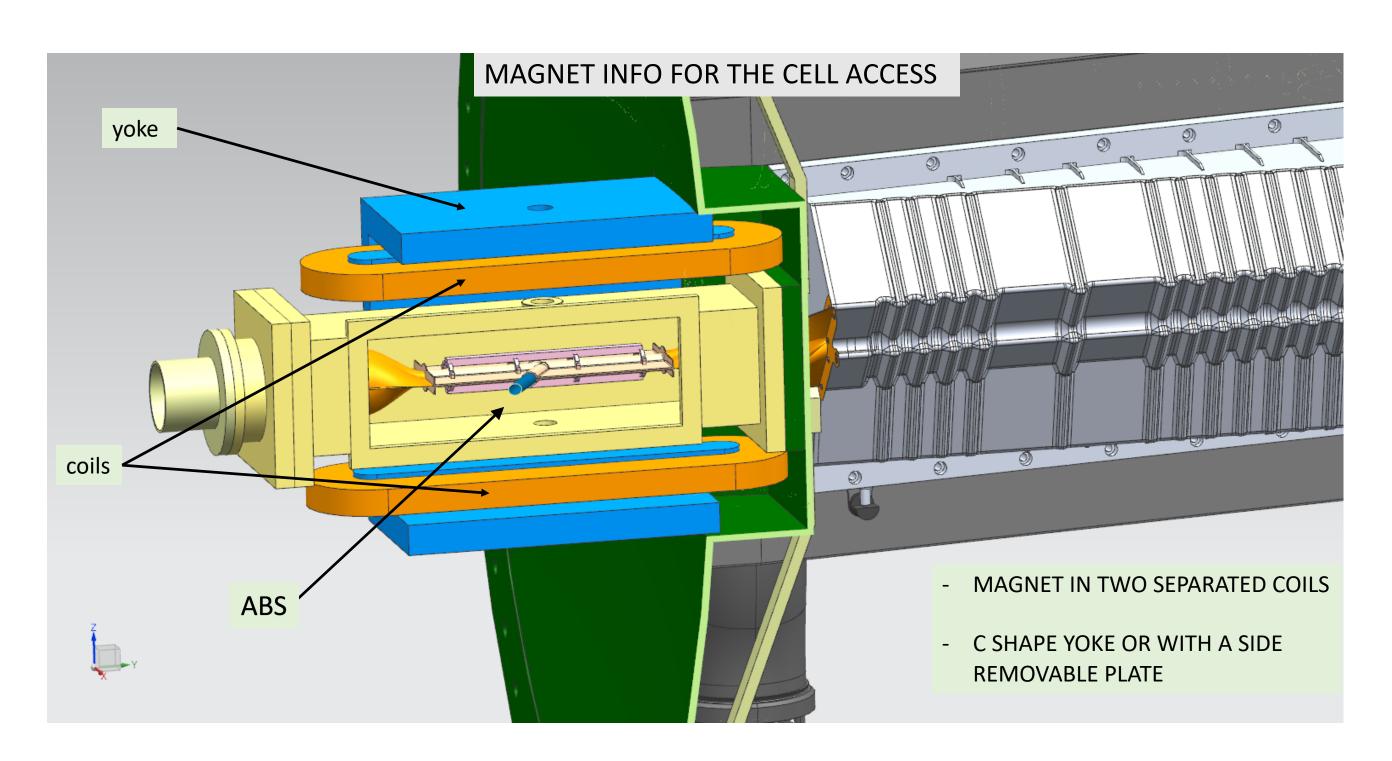
PGT implementation into LHCb

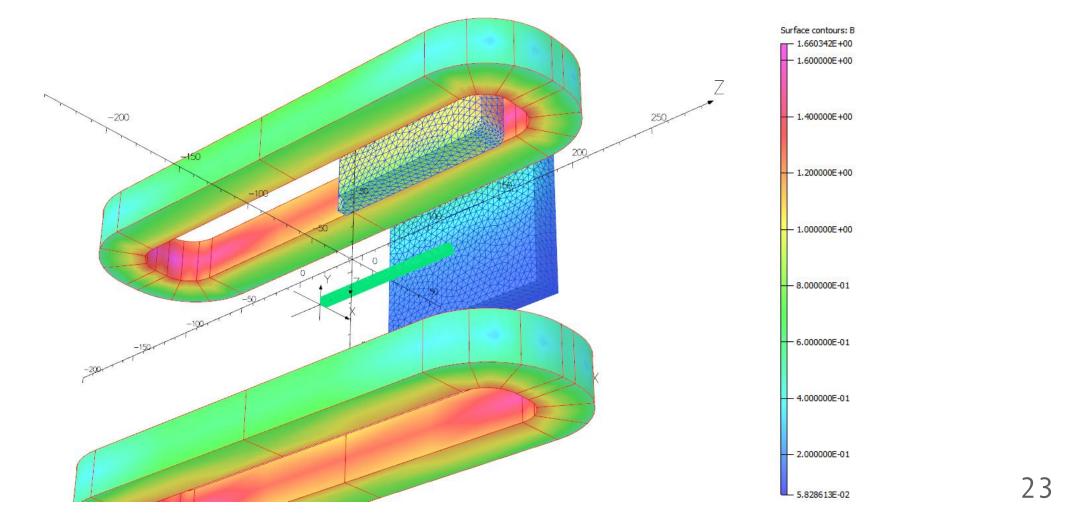
Inject both polarised and unpolarised gases via ABS and UGFS



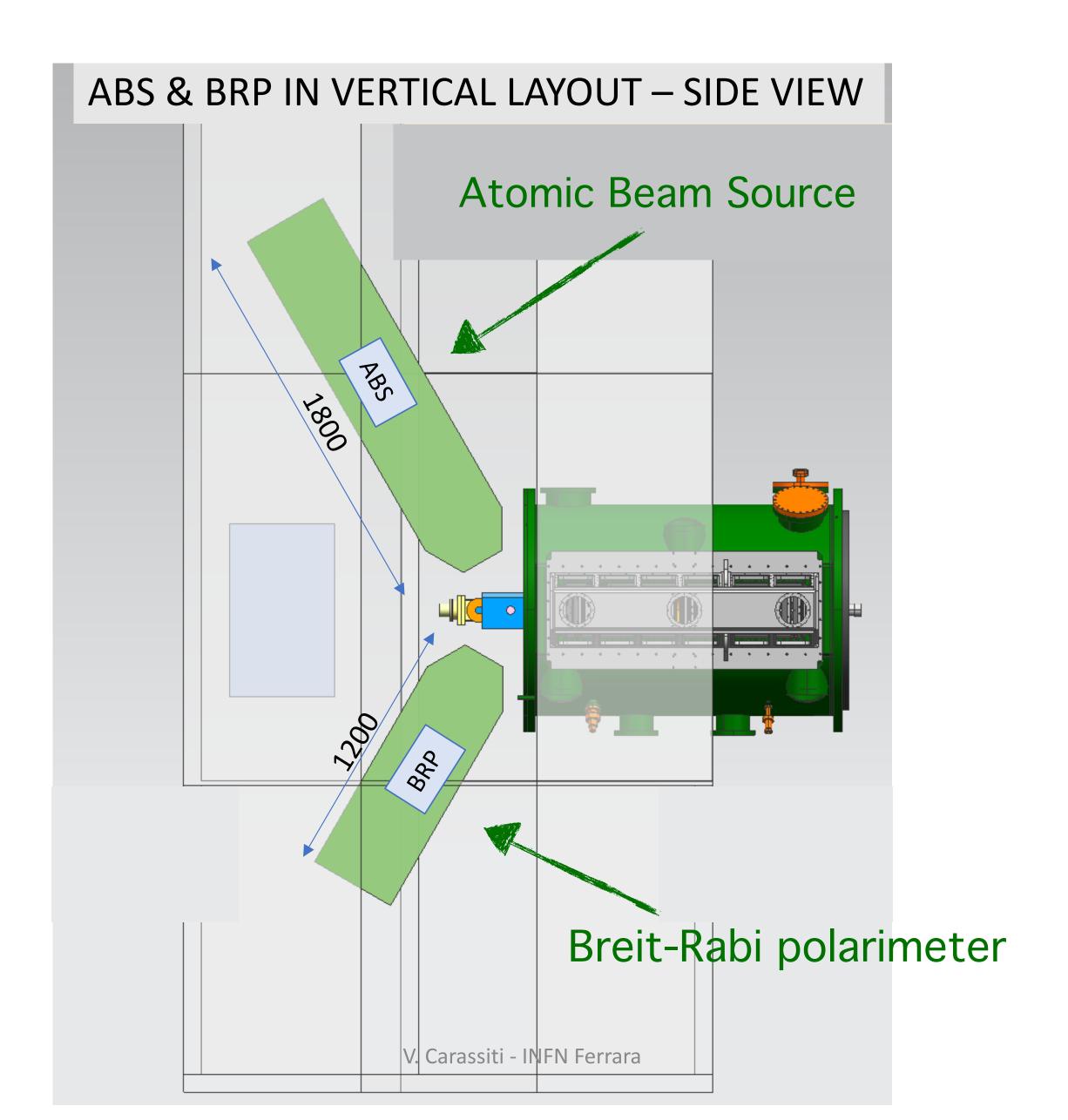
- Compact dipole magnet → static transverse field
- Superconductive coils + iron yoke configuration fits the space constraints
- $B=300~\mathrm{mT}$ with polarity inversion, $\Delta B/B \simeq 10\,\%$, suitable to avoid beam-induced depolarisation [Pos (SPIN2018)]

Possibility to switch to a solenoid and provide longitudinal polarisation (e.g. in LHC Run 5)





ABS & BRP implementation into LHCb



- Reduce the size of both ABS and BRP to fit into the available space in the LHCb cavern: a challenging R&D!
- No need for additional detectors in LHCb: only a modification of the VELO flange is needed
- $P \simeq 85\%$ achieved at HERMES

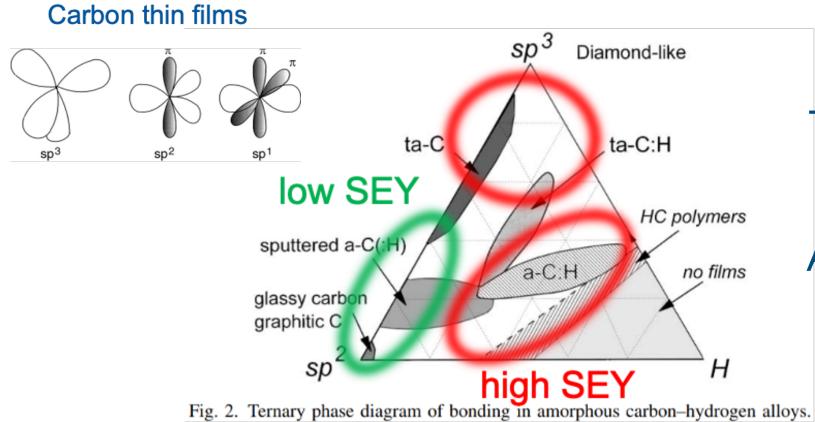
Injected intensity of H-atoms:

$$\phi = 6.5 \times 10^{16} \text{ s}^{-1}$$

Achievable Luminosity (HL-LHC):

$$\sim 8 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$$

Role of the storage cell coating

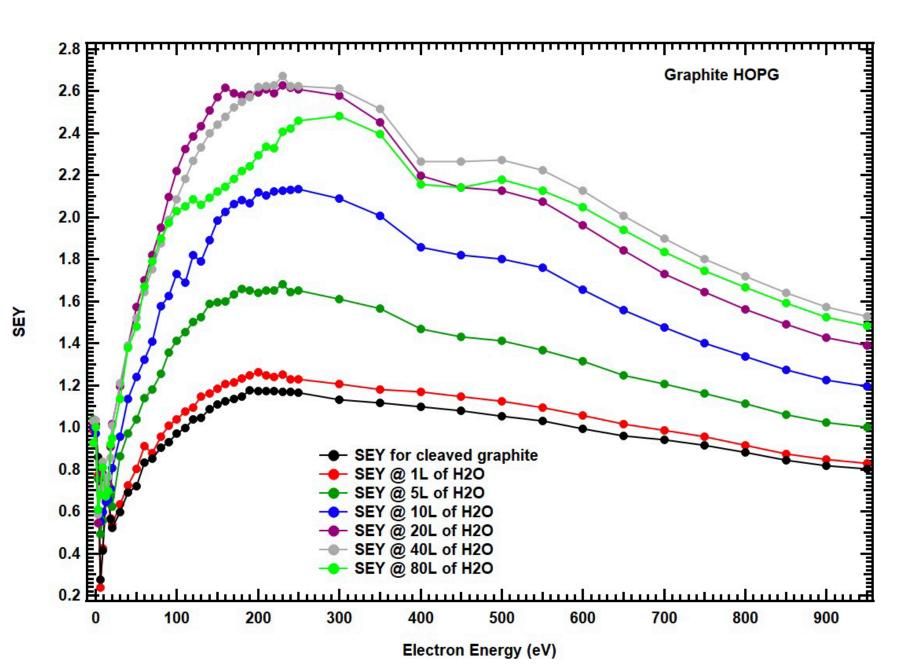


J. Robertson/Materials Science and Engineering R 37 (2002) 129–281

The material of the cell walls must have a low Secondary Electron Yield (e-cloud)

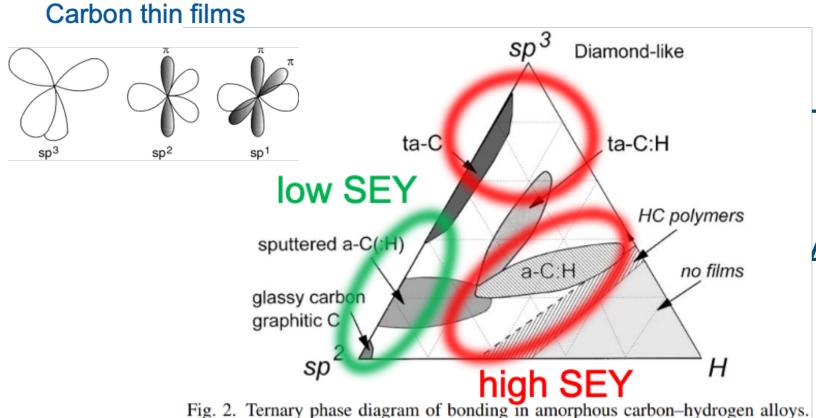
As for SMOG2, Amorphous Carbon is ok. Has it a low H recombination as well?

Studies ongoing in order to understand if <u>carbon films</u> with low secondary Electron Yield cope with the required "recombination' rate of polarized H atoms injected in the storage cell



... or follow the HERMES experience to have an ice coating (low SEY, low H recombination)

Role of the storage cell coating

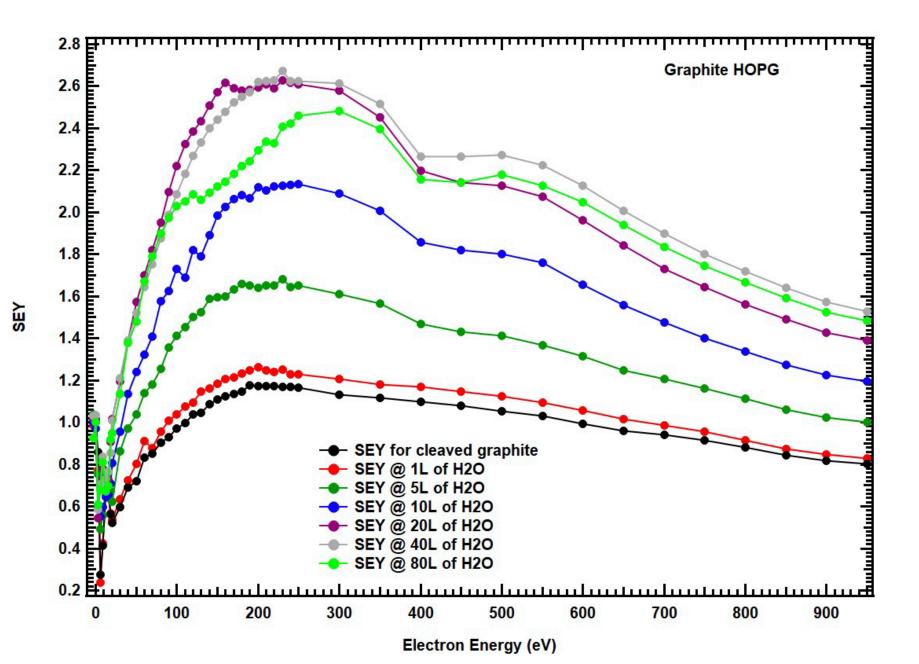


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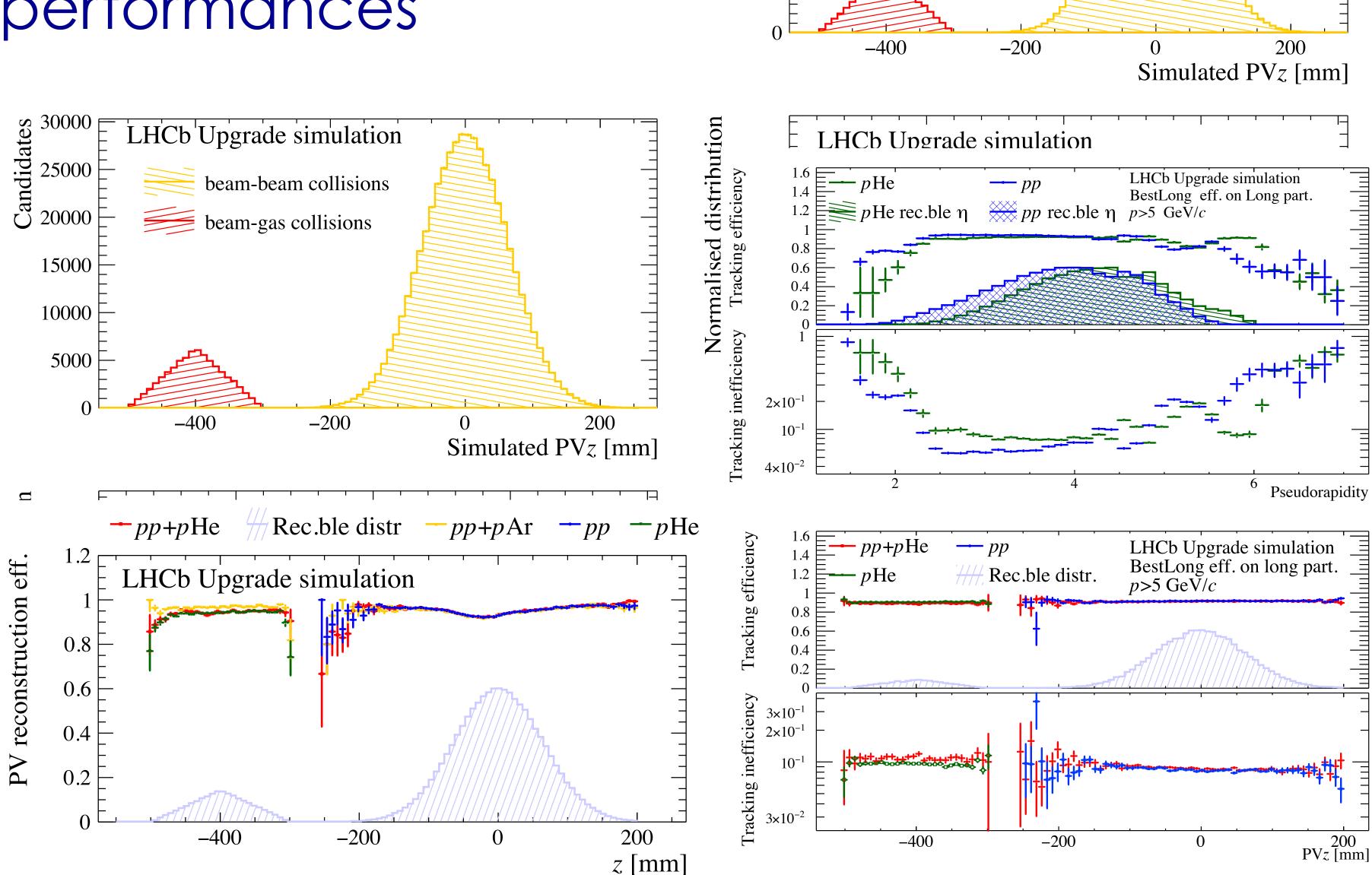
... or follow the HERMES experience to have an ice coating (low SEY, low H recombination)

Backup solution is also being investigated: a jet target that provides lower density (~10¹² atoms/cm²) but higher polarisation degree (up to 90%) and lower systematics

SMOG2/LHCspin performances

- beam-beam and beam-gas interaction regions are well detached
- Negligible increase of multiplicity:
 1 3 % throughput decrease when adding beam-gas to the LHCb event reconstruction sequence

 Full reconstruction efficiency (PV & tracks) retained in the beam-gas region



10000

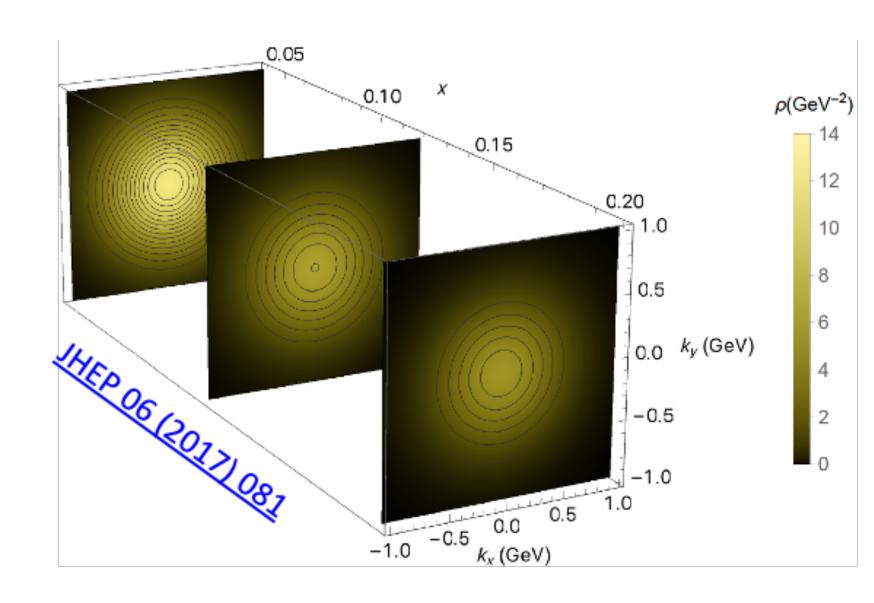
5000

LHCb is the only experiment able to run in collider and fixed-target mode simultaneously!

[LHCB=FIGURE=2022-002]

The physics goals of L++C

- Multi-dimensional nucleon structure in a poorly explored kinematic domain
- Measure experimental observables sensitive to both quarks and gluons TMDs
- Make use of new probes (charmed and beauty mesons)
- Complement present and future SIDIS results
- Test non-trivial process dependence of quarks and (especially) gluons TMDs
- Measure exclusive processes to access GPDs



quark pol.

ur e	U	L	Т
U	f_1		h_1^{\perp}
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^{\perp}	g_{1T}	h_1, h_{1T}^{\perp}

nucleon pol

Theoretically cleanest hard h-h scattering process:

• LHCb has excellent μ – ID & reconstruction for $\mu^+\mu^-$

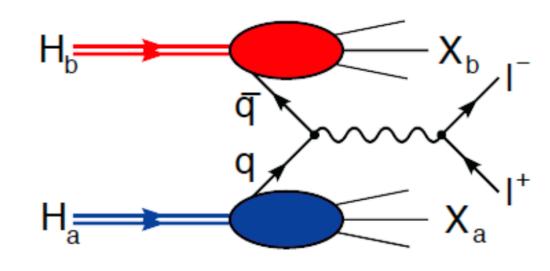
dominant:
$$\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^{+}\mu^{-}$$
 suppressed: $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^{+}\mu^{-}$

• Sensitive to unpol. and BM TMDs for $q_T \ll M_T$

$$d\sigma_{UU}^{DY} \propto f_1^{\bar{q}} \otimes f_1^{q} + \cos 2\phi \ h_1^{\perp,\bar{q}} \otimes h_1^{\perp,q}$$

- H & D targets allow to study the antiquark content of the nucleon
- SeaQuest (E906): $\bar{d}(x) > \bar{u}(x) \rightarrow \text{proton sea is not flavour symmetric}$
- intrinsic heavy quarks?

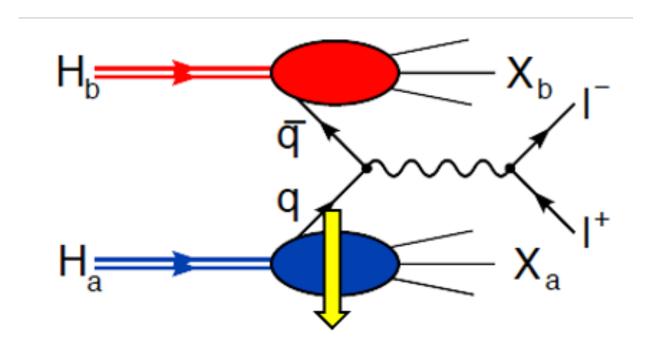
Unpolarized Drell-Yan



... still a lot to be understood and investigated

Quark TMDs

Transv. polarized Drell-Yan



Sensitive to quark TMDs through TSSAs

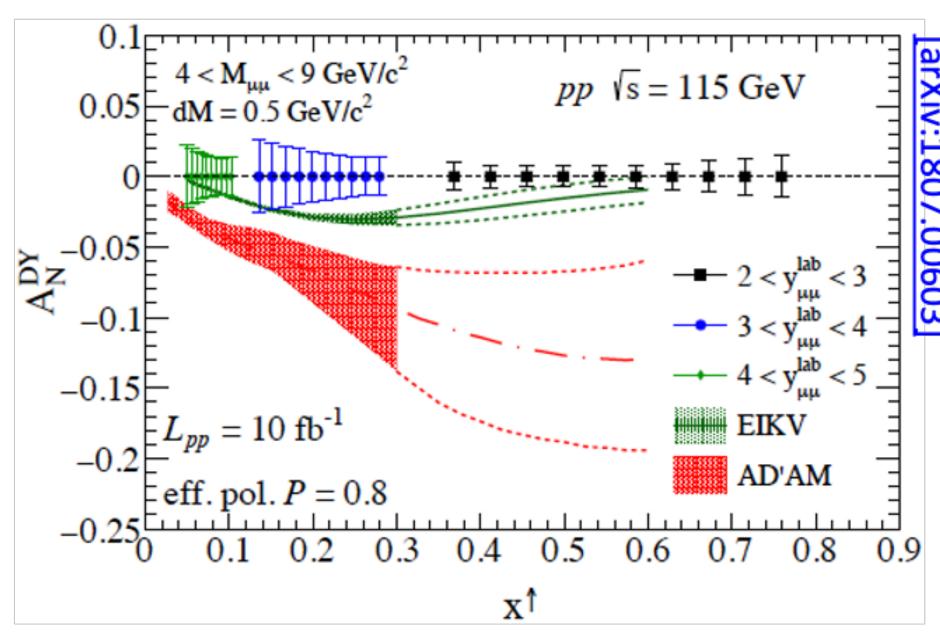
$$A_{N}^{DY} = \frac{1}{P} \frac{\sigma_{DY}^{\uparrow} - \sigma_{DY}^{\downarrow}}{\sigma_{DY}^{\uparrow} + \sigma_{DY}^{\downarrow}} \implies A_{UT}^{sin\phi_S} \sim \frac{f_1^q \otimes f_{1T}^{\downarrow q}}{f_1^q \otimes f_1^q}, \quad A_{UT}^{sin(2\phi - \phi_S)} \sim \frac{h_1^{\downarrow q} \otimes h_1^q}{f_1^q \otimes f_1^q}, \dots$$

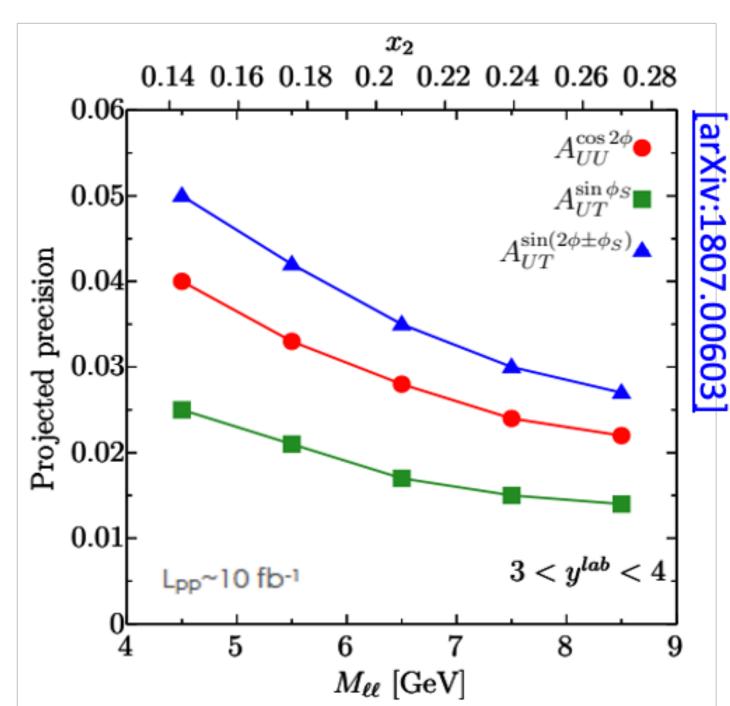
(ϕ : azimuthal orientation of lepton pair in dilepton CM)

- Extraction of qTMDs from DY does not require knowledge of FF
- Verify sign change of Sivers function wrt SIDIS

$$f_{1T}^{\perp}\big|_{DY} = -f_{1T}^{\perp}\big|_{SIDIS}$$

Test flavour sensitivity using both H and D targets





Probing the gTMDs

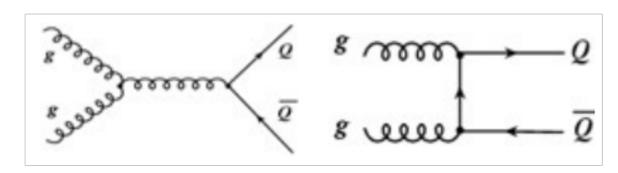
gluon pol.

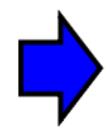
UCircularlyLinearlyU f_1^g $h_1^{\perp g}$ L g_{1L}^g $h_{1L}^{\perp g}$ T $f_{1T}^{\perp g}$ g_{1T}^g

nucleon pol

Theory framework well consolidated ...but experimental access still extremely limited!

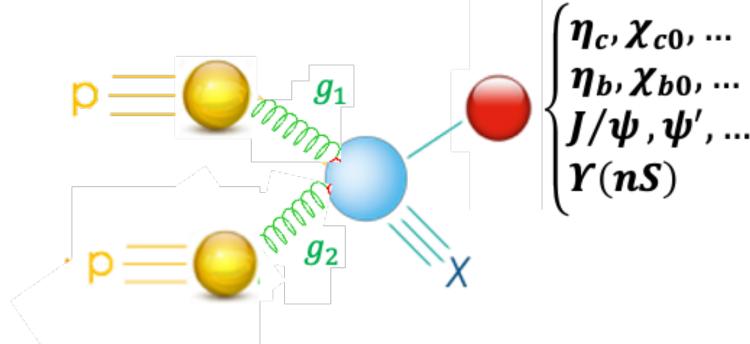
In high-energy hadron collisions, heavy quarks are dominantly produced by gg fusion:





The most efficient way to access the gluon dynamics inside the proton at LHC is to measure heavy-quark observables

• Inclusive quarkonia production in (un)polarized pp interaction $(pp^{(\uparrow)} \to [Q\bar{Q}]X)$ turns out to be an ideal observable to access gTMDs (assuming TMD factorization)

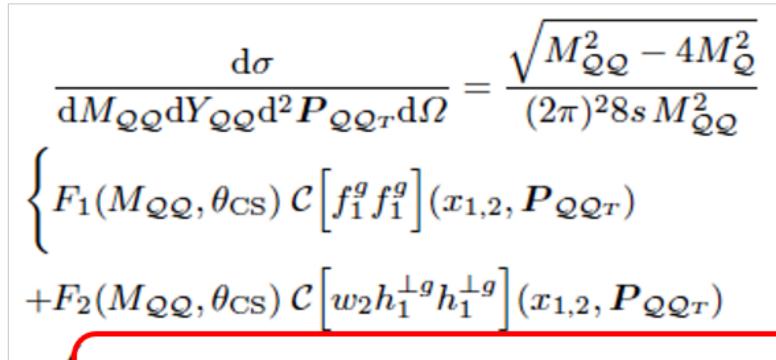


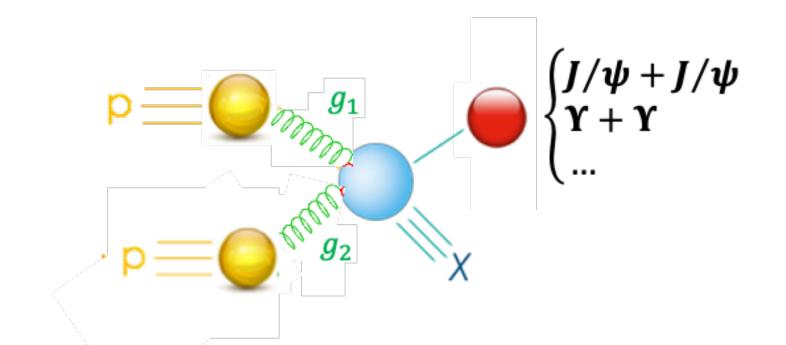
• TMD factorization requires $q_T(Q) \ll M_Q$. Can look at associate quarkonia production, where only the relative q_T needs to be small:

E.g.:
$$pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$$

•Due the larger masses this condition is more easily matched in the case of **bottomonium**, where TMD factorization can hold at larger q_T (although very challenging for experiments!)

Probing the gTMDs

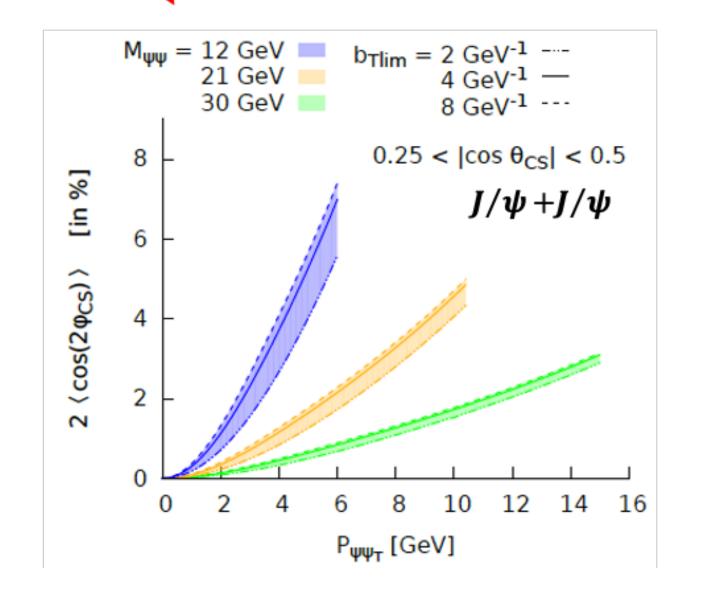




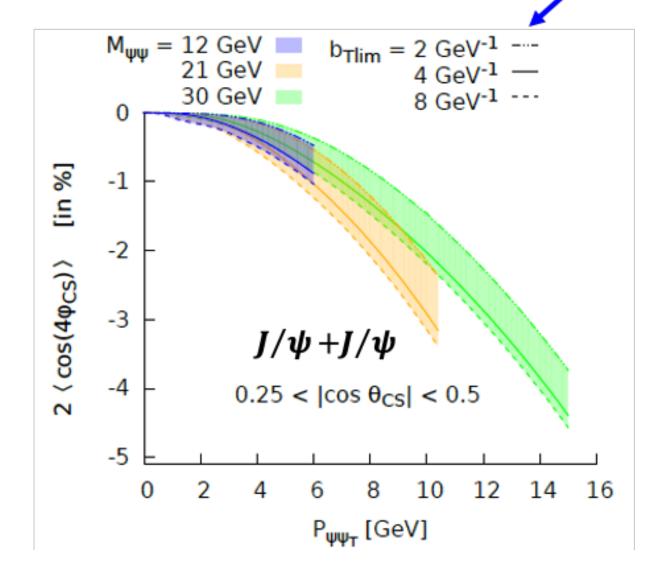
	gluon pol.				
		U	Circularly	Linearly	
pol.	U	f_1^g		$h_1^{\perp g}$	
nucleon pol	L		g_{1L}^g	$h_{1L}^{\perp g}$	
nnc	Т	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g,h_{1T}^{\perp g}$	

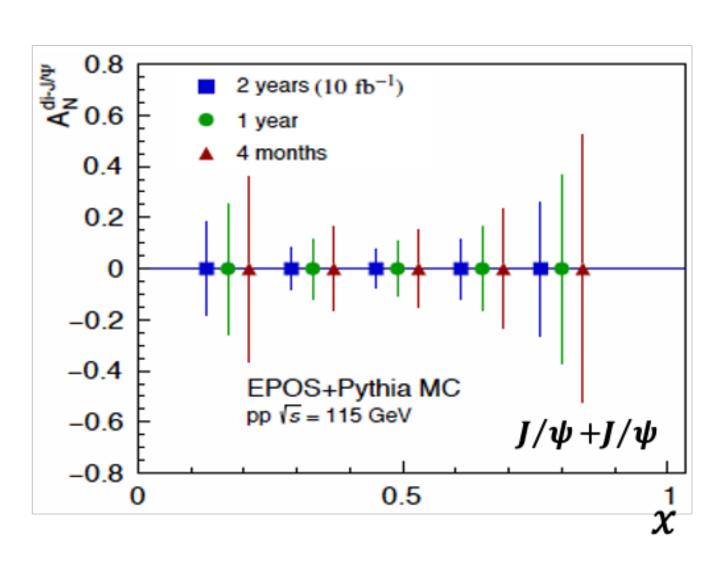
$$+\left(F_{3}(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}})\,\mathcal{C}\left[w_{3}f_{1}^{g}h_{1}^{\perp g}\right](x_{1,2},\boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_{T}})+F_{3}'(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}})\,\mathcal{C}\left[w_{3}'h_{1}^{\perp g}f_{1}^{g}\right](x_{1,2},\boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_{T}})\right)\cos2\phi_{\mathrm{CS}}\left(+F_{4}(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}})\,\mathcal{C}\left[w_{4}h_{1}^{\perp g}h_{1}^{\perp g}\right](x_{1,2},\boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_{T}})\cos4\phi_{\mathrm{CS}}\right)$$

Predictions based on CSM + TMD evolution for $x_1 \sim x_2 \sim 10^{-3}$ at forward rapidity [EPJ C 80, 87 (2020)]



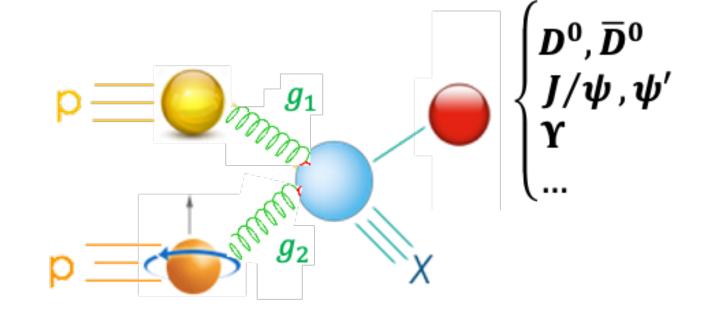
Azimuthal amplitudes ~5%!





Probing the gluon Sivers function

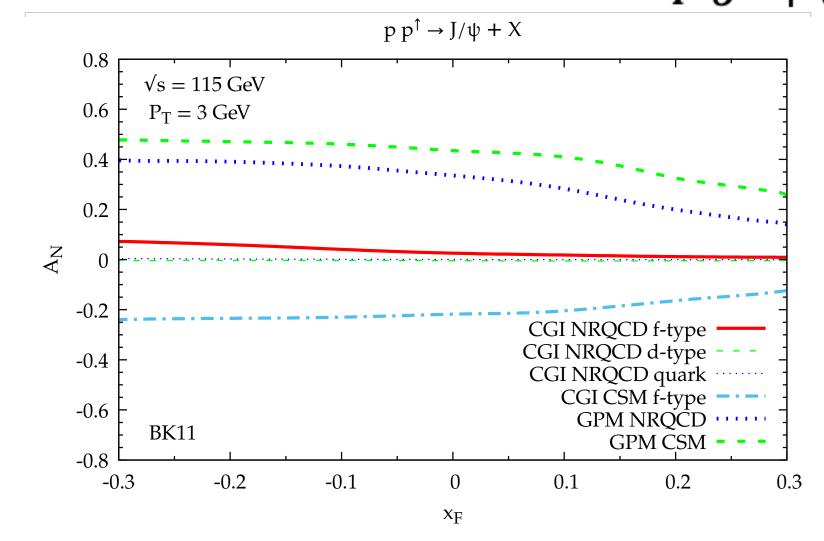
$$\Gamma_T^{\mu\nu}(x,\boldsymbol{p}_T) = \frac{x}{2} \left\{ g_T^{\mu\nu} \frac{\epsilon_T^{\rho\sigma} p_{T\rho} S_{T\sigma}}{M_p} (f_{1T}^{\perp g}(x,\boldsymbol{p}_T^2) + \dots \right\}$$



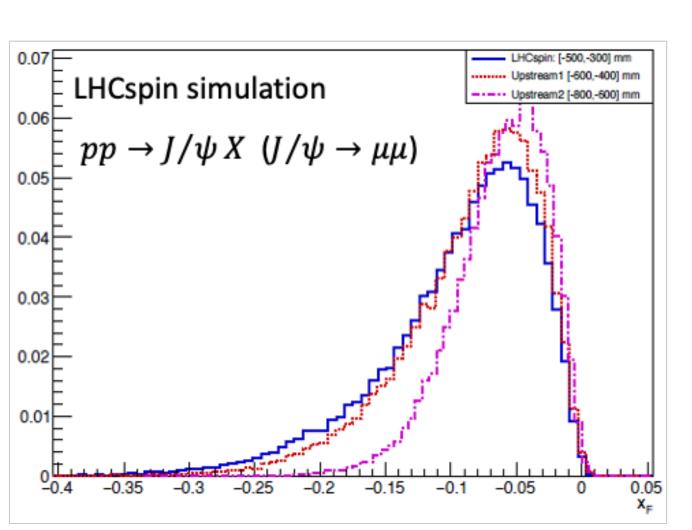
	gluon pol.				
		U	Circularly	Linearly	
pol.	U	f_1^g		$h_1^{\perp g}$	
nucleon pol.	L		g_{1L}^g	$h_{1L}^{\perp g}$	
nnc]	Т	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g,h_{1T}^{\perp g}$	

- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- sensitive to color exchange among IS and FS and gluon OAM
- expected to be small (quasi-saturation of Burkardt sum rule by $f_{1T}^{\perp q}$ and QCD predictions in large- N_c limit)
- can be accessed through the Fourier decomposition of the TSSAs for inclusive heavy meson production

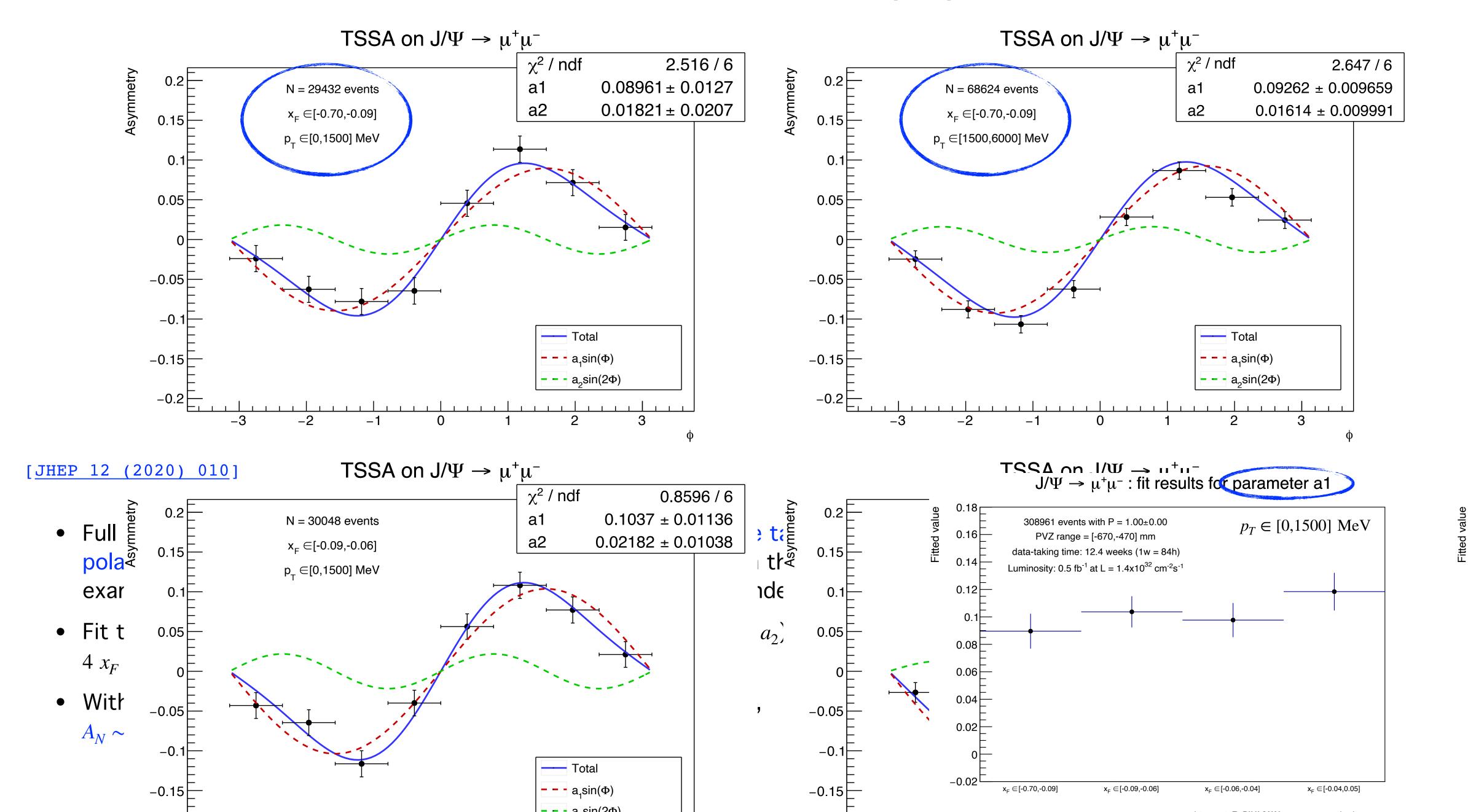
$$A_{N} = \frac{1}{P} \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} \propto \left[f_{1T}^{\perp g}(x_{a}, k_{\perp a}) \otimes f_{g}(x_{b}, k_{\perp b}) \otimes d\sigma_{gg \to QQg} \right] \sin \phi_{S} + \cdots$$



- Predictions for pol. FT meas. at LHC (LHCspin-like)
- Phys. Rev. D 102, 094011 (2020)
- $pp^{\uparrow} \rightarrow J/\psi + X$
- based on GPM & CGI-GPM
- Expected amplitudes could reach 5-10% in the $x_F < 0$ region



A TSSA analysis at LHCspin with $J/\Psi \to \mu^+\mu^-$ events



0.16

0.12

0.1

0.08

0.06

0.04

0.02

Knowledge of the polarisation deg

TSSA on $J/\Psi \rightarrow \mu^{+}\mu^{-}$

 χ^2 / ndf

a1

a2

- To estimate the systematic error due to the measurement of the polarisation degree, the analysis is repeated with different ΔP
- Very relevant for the R&D (e.g. cell vs jet target). With the shown analysis*:
- 5% error (realistic value) → negligible effect

N = 29432 events

 $x_{r} \in [-0.70, -0.09]$

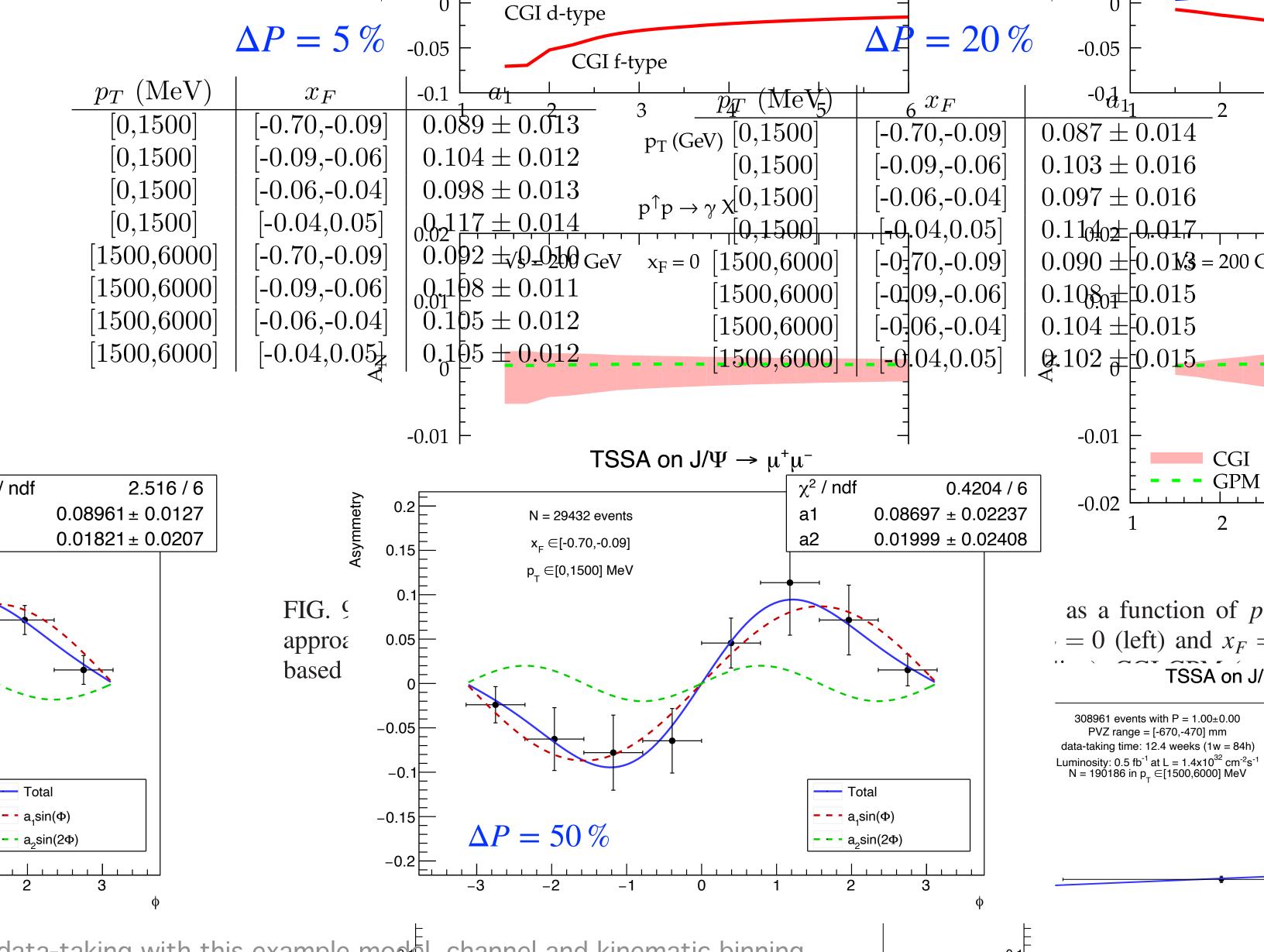
p_⊤ ∈[0,1500] MeV

- 20% error \rightarrow 30-40% of the stat. error
- 50% error → syst. dominated

0.15

0.05

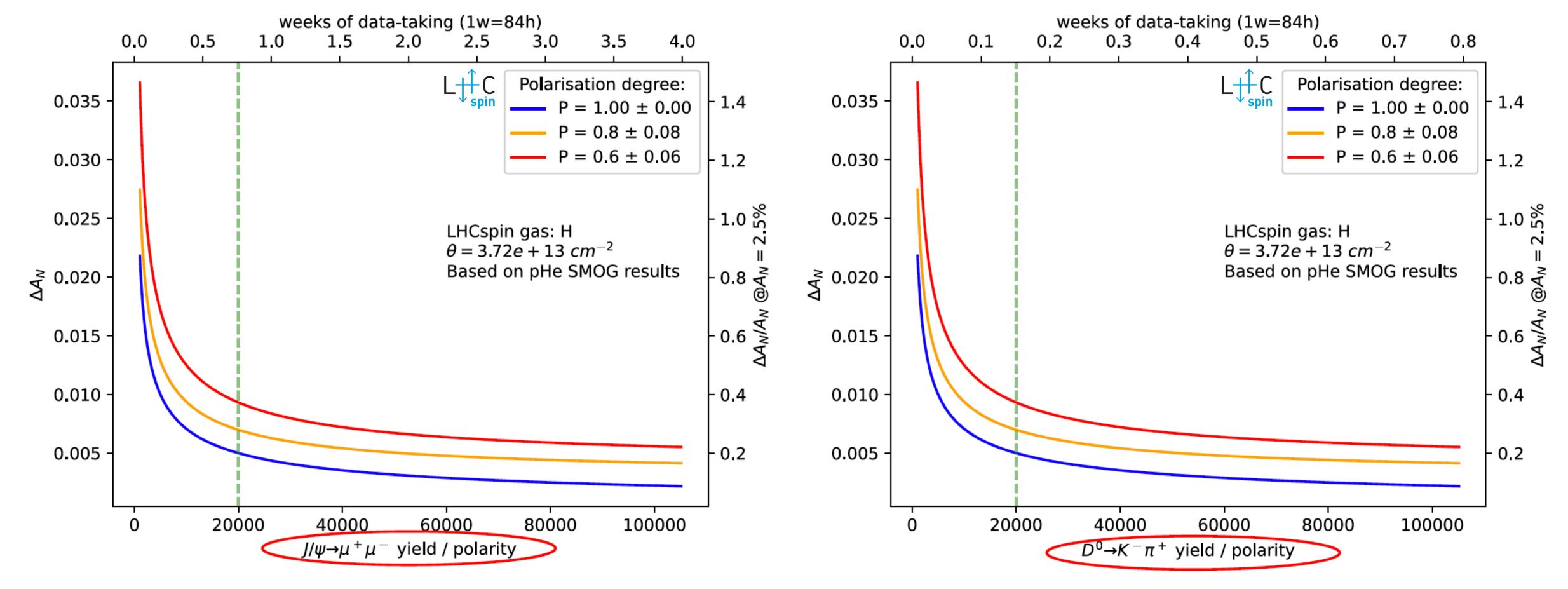
-0.05 -



* i.e. ~ 3 months of data-taking with this example model, channel and kinematic binning

LHCspin event rates

Precise spin asymmetry on $J/\Psi \to \mu^+\mu^-$ and $D^0 \to K^-\pi^+$ for pH^\uparrow collisions in just few weeks with Run3 luminosity! Statistics further enhanced by a factor 3-5 in LHCb upgrade II



reconstructed particles

UPC and gGPDs

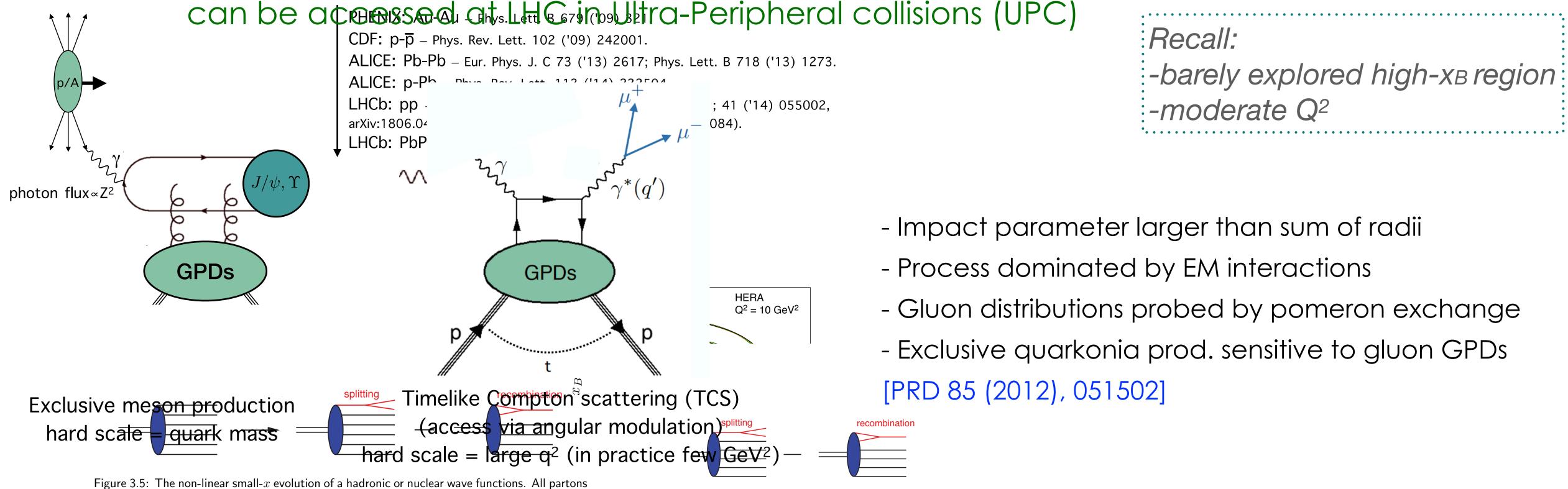
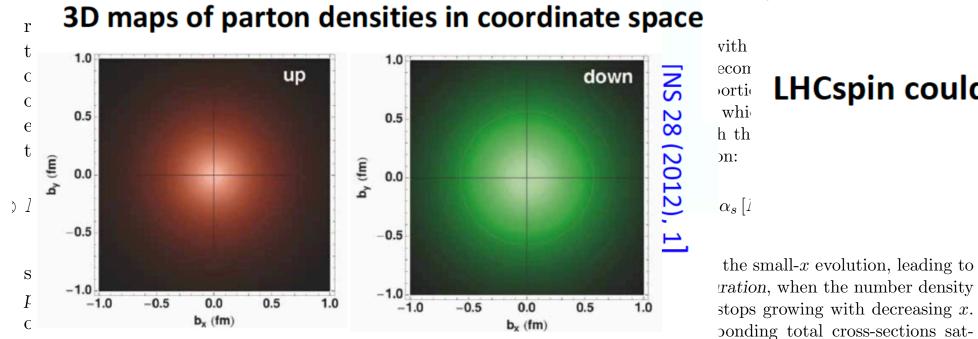


Figure 3.5: The non-linear small-x evolution of a hadronic or nuclear wave functions. All partons (quarks and gluons) are denoted by straight solid lines for simplicity.

an of a hadronic or nuclear wave functions. All partons

GPD	$oldsymbol{U}$	L	T
$oldsymbol{U}$	H		\mathcal{E}_T
L		$ ilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T,\; ilde{H}_T$

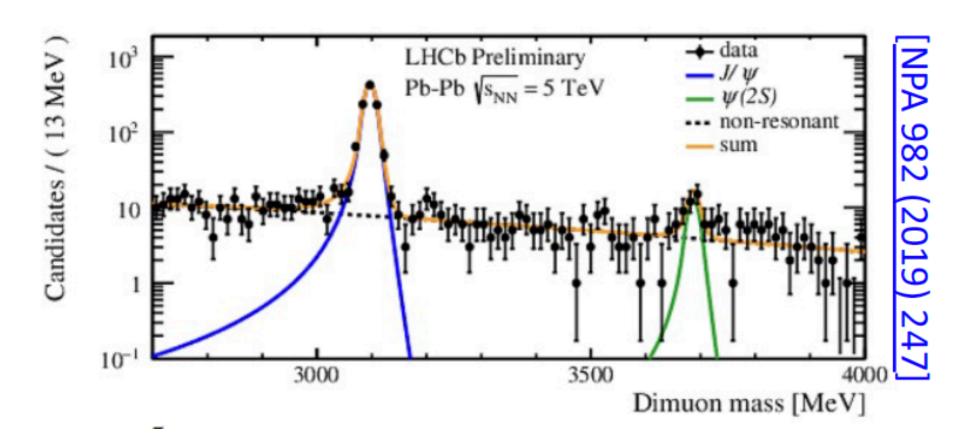


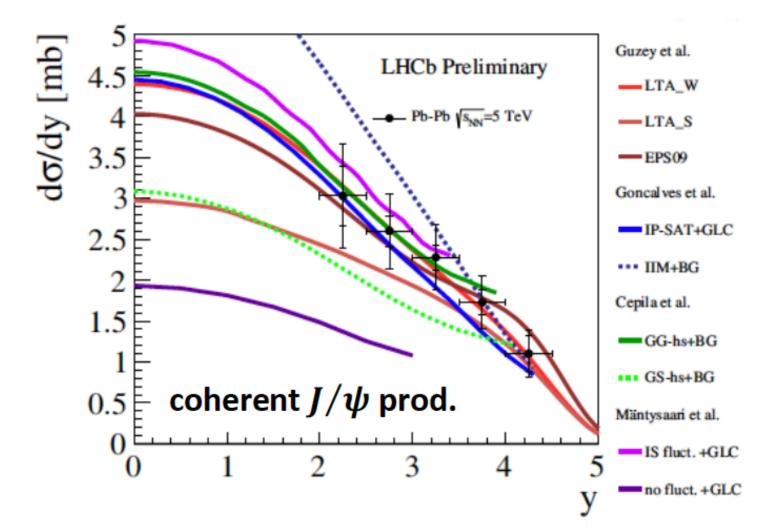
LHCspin could allow to access the GPD E^g (a key ingredient of the Ji sum rule)

$$J^{g} = \frac{1}{2} \int_{0}^{1} dx \Big(H^{g}(x, \xi, 0) + E^{g}(x, \xi, 0) \Big)$$

UPC and gGPDs

First results from LHCb in PbPb UPC





SMOG2

Continuum 2 muons, statistical uncertainty on $cos(\phi)$ modulation, p_T cut included

pp	pD	pAr	pKr	pXe
30 %	_	10 %	20 %	15 %

Continuum 2 muons, statistical uncertainty on $cos(\phi)$ modulation, p_T cut not included

Pbp	PbAr
_	30 %

Note: luminosity uncertainty does not enter and rest of systematic uncertainties expected small, since modulation.

J/ψ, total uncertainty on cross section, assuming 4% uncertainty on luminosity

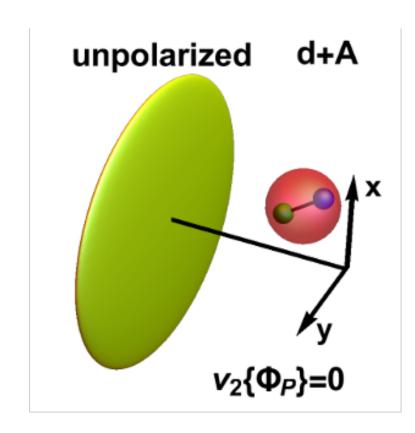
pp	pD	pAr	pKr	pXe
10 %	_	5 %	5 %	5 %

Pbp	PbAr
_	5 %

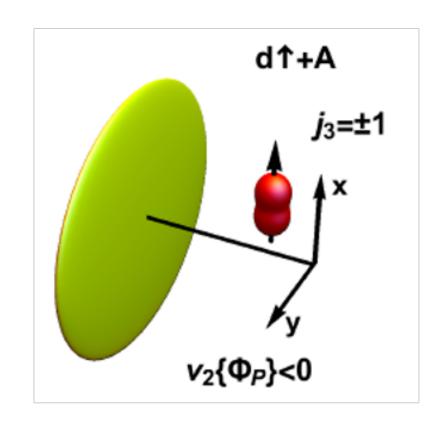
C. van Hulse's slide

Spin physics in heavy-ion collisions

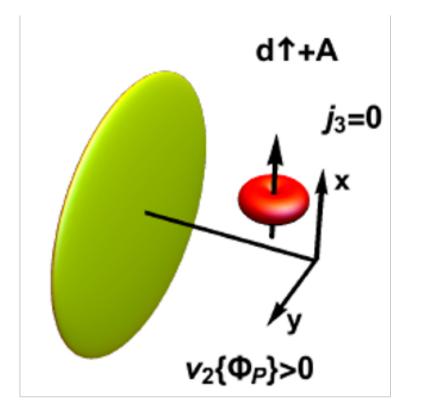
- •probe collective phenomena in heavy-light systems through ultrarelativistic collisions of heavy nuclei with trasv. pol. deuterons
- •polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the **elliptic flow** relative to the polarization axis (**ellipticity**).



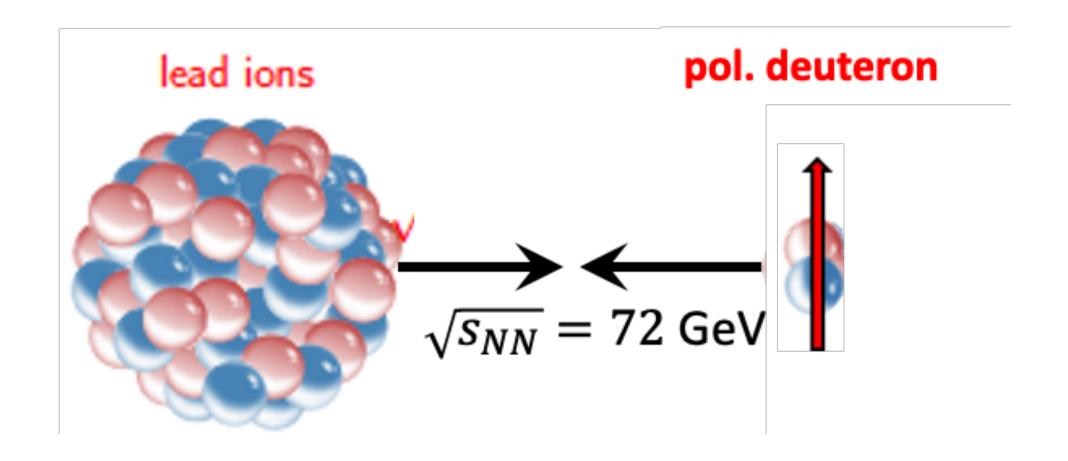
Unpol. deuterons: the fireball is azimuthally symmetric and $v_2 \approx 0$.

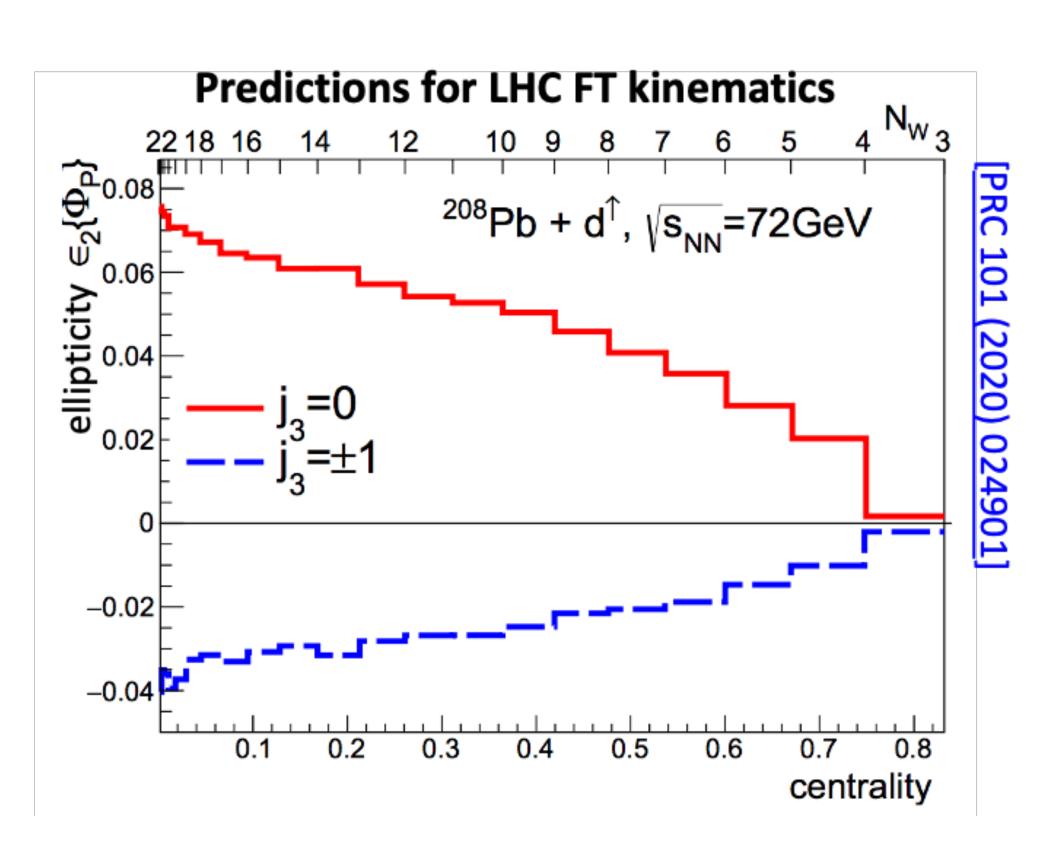


 $j_3 = \pm 1 \rightarrow \text{prolate fireball}$ stretched along the pol. axis, corresponds to $v_2 < 0$



 $j_3 = 0 \rightarrow \text{oblate fireball}$ corresponds to $v_2 > 0$

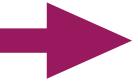




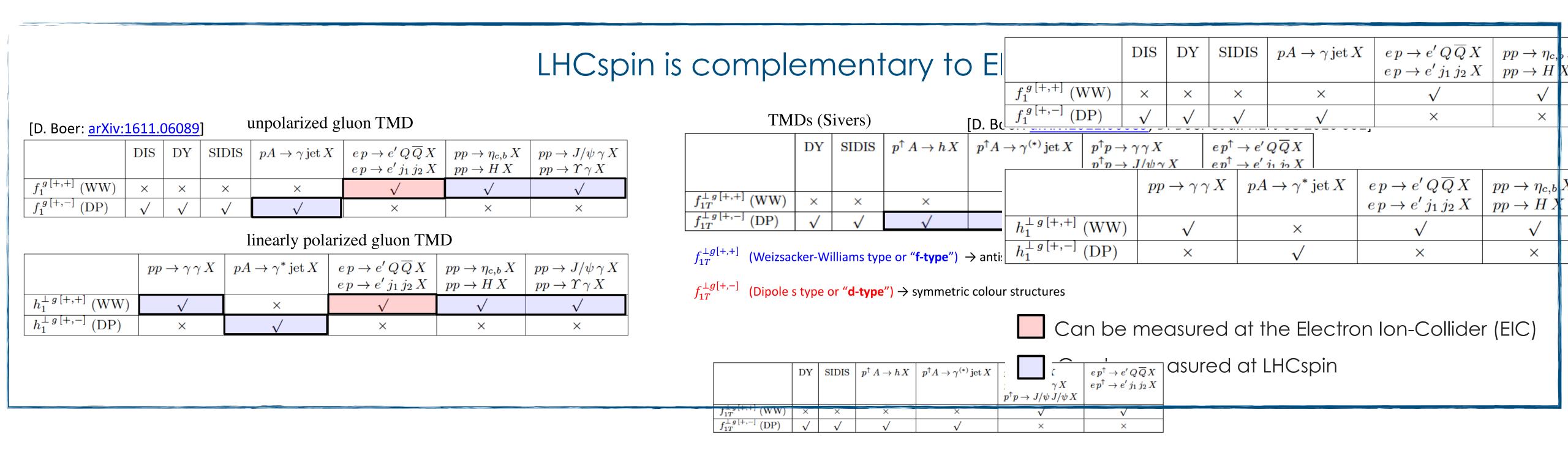
International framework and feedback

Several experiments dedicated to spin physics, but with many limitations:

very low energy, no rare probes, no ion beam, ...



LHCspin is unique in this respect



"Ambitious and long term LHC-Fixed Target research program. The efforts of the existing LHC experiments to implement such a programme, including specific R&D actions on the collider, deserve support (European Strategy for Particle Physics)

because the asymmetries in question have a process dependence between pp and lp that is predicted by theory: SERN Physics Beyond Collider)

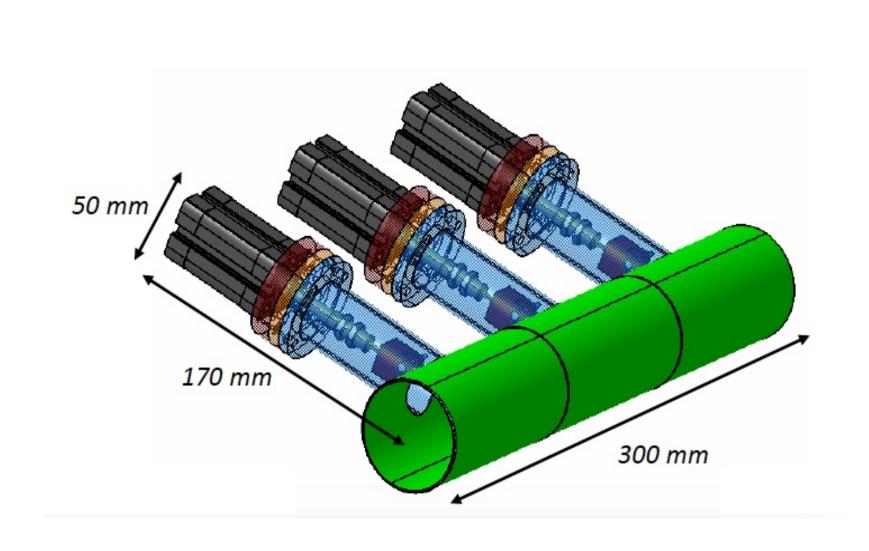
solid target @



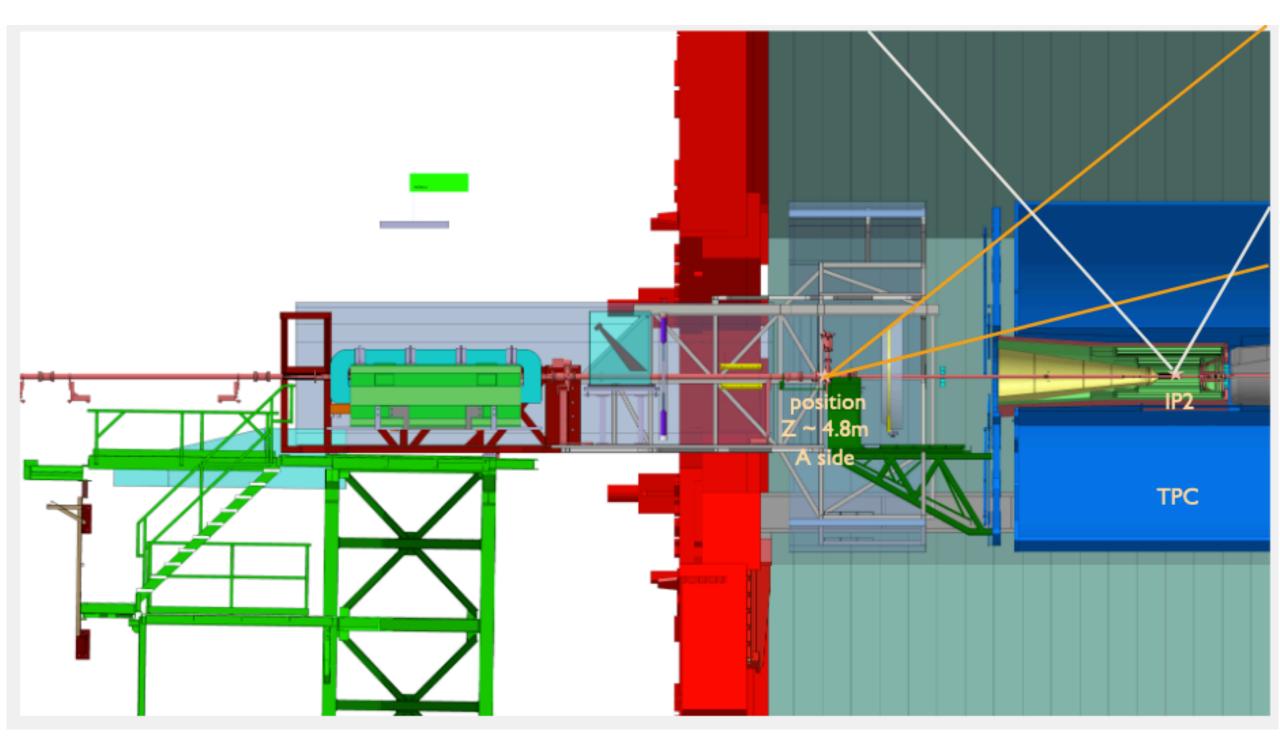
The ALICE unpolarised solid target

Two main physics goals:

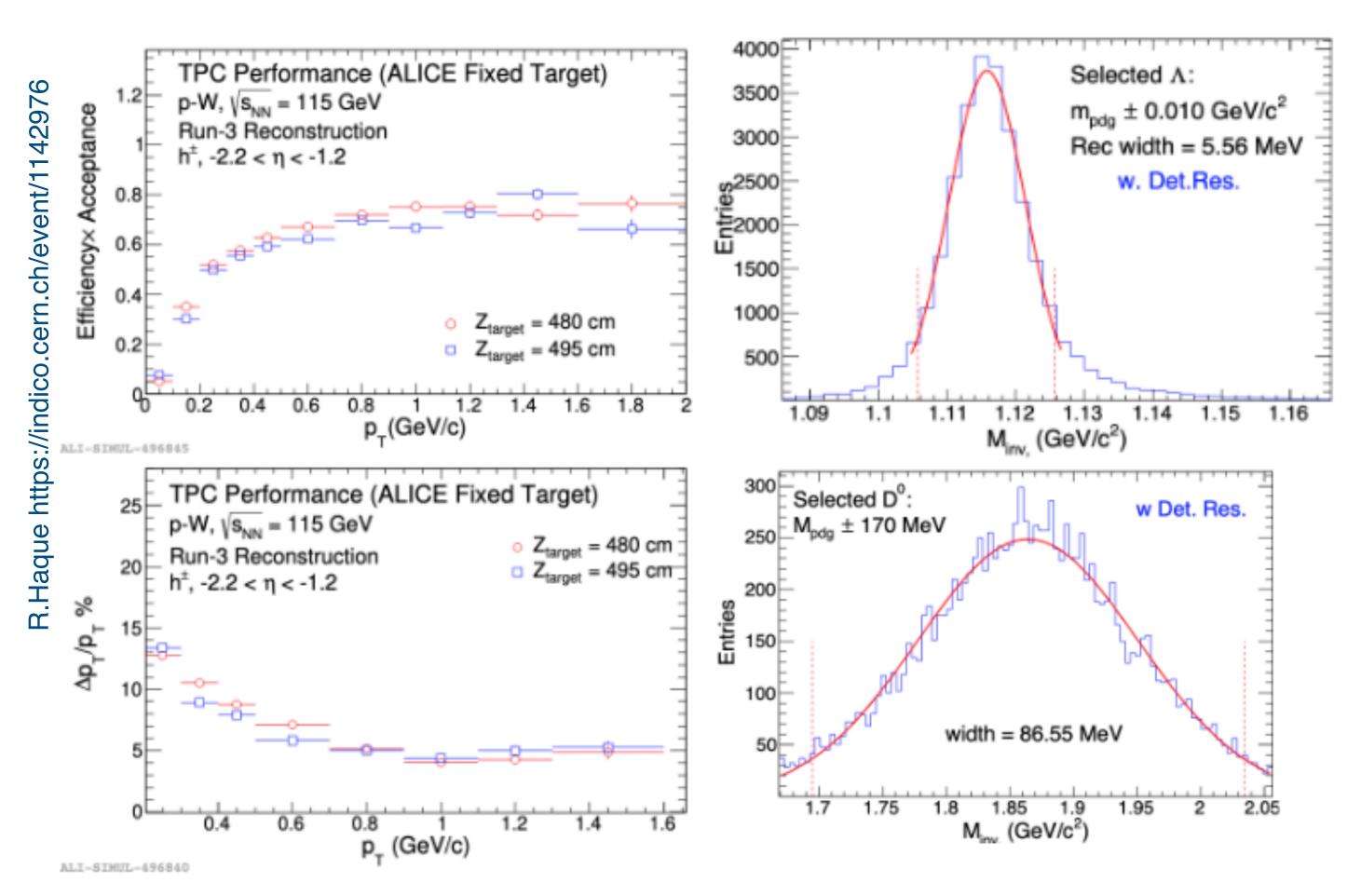
- Advance the understanding of the large-x gluon, antiquark and heavy-quark content in the nucleon and nucleus (structure of nucleon and nuclei at large-x, gluon EMC effect in nuclei, intrinsic charm in nucleon)
- Study heavy-ion collisions between SPS and RHIC energies towards large rapidities (longitudinal expansion of QGP formation, collectivity in small systems with heavy quarks, factorisation of CNM effects)
- Proton beam halo channelled with a bent crystal on a retractable solid target (C,W, Ti, ...)
- Backward cms rapidity coverage with forward detectors in the lab thanks to the boost



retractable solid target



Some of the performances



space availability at z = 3259 m



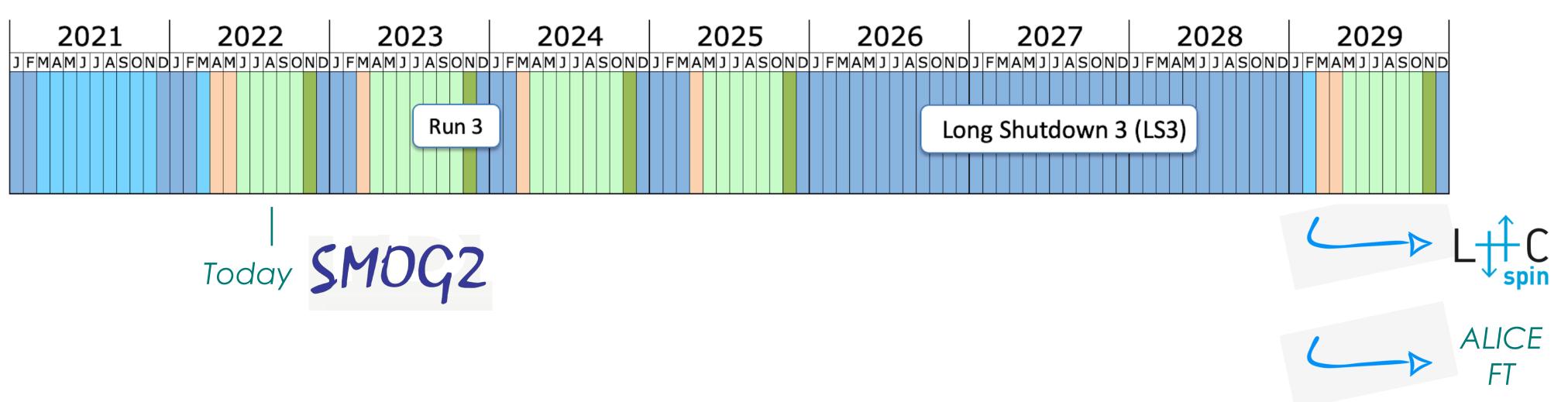
Proton beam collimation studies performed: loss maps, positioning of the crystal system and of the absorbers

LOI in ALICE (2022) —> aim for installation during LS3 (2026-2028)

Some of the results achieved

- Λ : efficiency and p_T resolution sufficient for analysis (without extra vertex detector)
- Do: TPC vertex resolution not sufficient to use secondary vertex method for analysis. Investigating combinatorial background method, reduced target size and constraints on beam spot position for tracking
- Integration solutions to comply with FOCAL and ITS motion constraints during EYETS
- Physics performance with realistic detector conditions

Conclusions



Fixed target physics at LHC is an exiting reality



has potentialities in the unpolarised case showing complementarity to LHCb



SMOC2 already operative and taking unpolarised data



is an innovative and unique project conceived to bring polarized physics at the LHC. It is extremely ambitious in terms of both physics reach and technical complexity. It could be installed in a realistic time schedule and costs





