The current state of solar modelling (and how we got there)

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Duvall-law inversion

$$F(w) = \int_{r_{t}}^{R_{o}} \left(\frac{r^{2}}{c^{2}} - \frac{1}{w^{2}}\right)^{1/2} \frac{dr}{r} = \frac{\pi(n+\alpha)}{\omega}$$

$$\frac{c_{0}^{2} - c_{m}^{2}}{c_{m}^{2}} = \frac{c_{0}^{2} - c_{m}^{2} - c_{m}^{2}}{c_{m}^{2}} = \frac{c_{$$

Christensen-Dalsgaard, Duvall, **Gough**, Harvey & Rhodes (1985; Nature, **315**, 378 – 382) A possible explanation is that the opacity used to construct the theoretical solar model is too low in the outer part of the radiative interior. An increase of $\sim 20\%$ in opacity at temperatures between $\sim 10^5$ and $\sim 4 \times 10^6$ K, for example, would reduce the discrepancy to an insignificant level.

Differential Duvall law

$$S(w)\frac{\delta\omega}{\omega} = \int_{r_{t}}^{R} \left(1 - \frac{c^{2}}{w^{2}r^{2}}\right)^{-1/2} \frac{\delta c}{c} \frac{dr}{c} + \pi \frac{\delta\alpha}{\omega}$$
$$S(w) \equiv \int_{r_{t}}^{R} \left(1 - \frac{c^{2}}{w^{2}r^{2}}\right)^{-1/2} \frac{dr}{c}$$

Linearized frequency differences

$$\frac{\delta\omega_{nl}}{\omega_{nl}} = \int_0^R \left[K_{c^2,\rho}^{nl}(r) \frac{\delta_r c^2}{c^2}(r) + K_{\rho,c^2}^{nl}(r) \frac{\delta_r \rho}{\rho}(r) \right] \mathrm{d}r + Q_{nl}^{-1} \mathcal{G}(\omega_{nl}) + \sigma_{nl} ,$$

$$\delta M = 4\pi \int_0^R \frac{\delta_r \rho(r)}{\rho(r)} \rho(r) r^2 \mathrm{d}r = 0 \; .$$



Small frequency separations

$$\nu_{nl} \sim \Delta \nu \left(n + \frac{l}{2} + \alpha \right) + \epsilon_{nl}$$

where

$$\epsilon_{nl} \simeq l(l+1) \frac{\Delta \nu}{4\pi^2 \nu_{nl}} \int_0^R \frac{\mathrm{d}c}{\mathrm{d}r} \frac{\mathrm{d}r}{r}$$

Frequency separations: $\Delta v_{nl} = v_{nl} - v_{n-1l} \approx \left(2 \int dr/c\right)^{-1}$

$$\delta\nu_{nl} = \nu_{nl} - \nu_{n-1\,l+2} \simeq -(4l+6)\frac{\Delta\nu}{4\pi^2\nu_{nl}} \int_0^R \frac{\mathrm{d}c}{\mathrm{d}r} \frac{\mathrm{d}r}{r}$$

Low-degree helioseismology



Elsworth et al. (1990; Nature 347, 536)

Oscillations and neutrinos



Elsworth et al. (1990; Nature, 347, 535)

Neutrinos and low-degree helioseismology



JC-D (1991; Geophys. Astrophys. Fluid Dyn. 62,123)





Neutrinos and inferred solar structure

0.1 Mixed $S_{34} = 0$ OPAL EOS /Sun 0.05 (Model-Sun) 0 -0.050.2 0.4 0.6 0.8 0 R/R_{o}

Bahcall et al. (1997; Phys. Rev. Lett. 78, 171)

Effects of He diffusion and settling



Model S



Model:

- Early OPAL EOS and opacity
- Simplified treatment of diffusion and settling
- Grevesse & Noels (1993) composition
- Bahcall & Pinsonneault (1995) nuclear
 parameters

Christensen-Dalsgaard et al. (1996; Science 272, 1286)

Basu et al. (1997; MNRAS 292, 243)

Global seismic parameters

$d_{\rm cz} = 0.287 \pm 0.001 R$

Basu & Antia (1997; MNRAS, 287, 189)

 $Y_{\rm env} = 0.2485 \pm 0.0034$

Basu & Antia (2004; Phys. Rep. 457, 217)

Detecting effects of relativistic electrons



Elliott & Kosovichev (1998; ApJ 500, L199)



Rotational mixing beneath convection zone



Overshoot as acoustic glitch



Christensen-Dalsgaard, Monteiro, Rempel and Thompson (2011; MNRAS 414, 1158)

Glitch parameters, observations and models



- 72-d data sets
- 1-yr data sets

Christensen-Dalsgaard, Monteiro, Rempel and Thompson (2011; MNRAS 414, 1158)

Overshoot models



Christensen-Dalsgaard, Monteiro, Rempel and Thompson (2011; MNRAS 414, 1158)

Glitch parameters, observations and models



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Modified surface composition



- GN93: $(Z/X)_s = 0.0245$
- AGS05: $(Z/X)_s = 0.0165$
- AGSS09: $(Z/X)_s = 0.0181$





Christensen-Dalsgaard & Houdek (2010; Astrophys. Space Sci., 328, 51)

Restored model S for AGSS09



Opacity 'kernels'



$$\delta \log \kappa = A_{\kappa} \exp[-(\log T - \log T_{\kappa})^2 / \Delta_{\kappa}^2]$$

$$\Delta_{\kappa} = 0.02$$

$$\log T_{\kappa} = 7.1 \qquad -----$$

$$\log T_{\kappa} = 6.9 \qquad ------$$

$$\log T_{\kappa} = 6.65 \qquad ------$$

$$\log T_{\kappa} = 6.4 \qquad ------$$

A somewhat desperate attempt?

