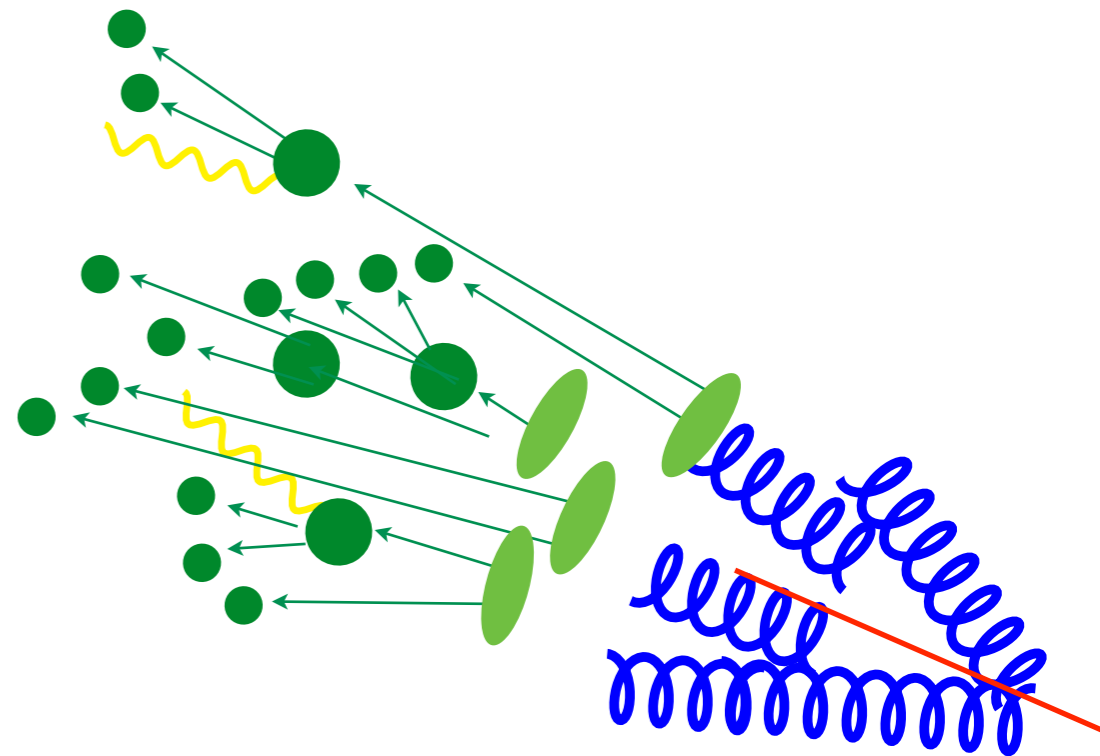


# QCD physics at the LHC

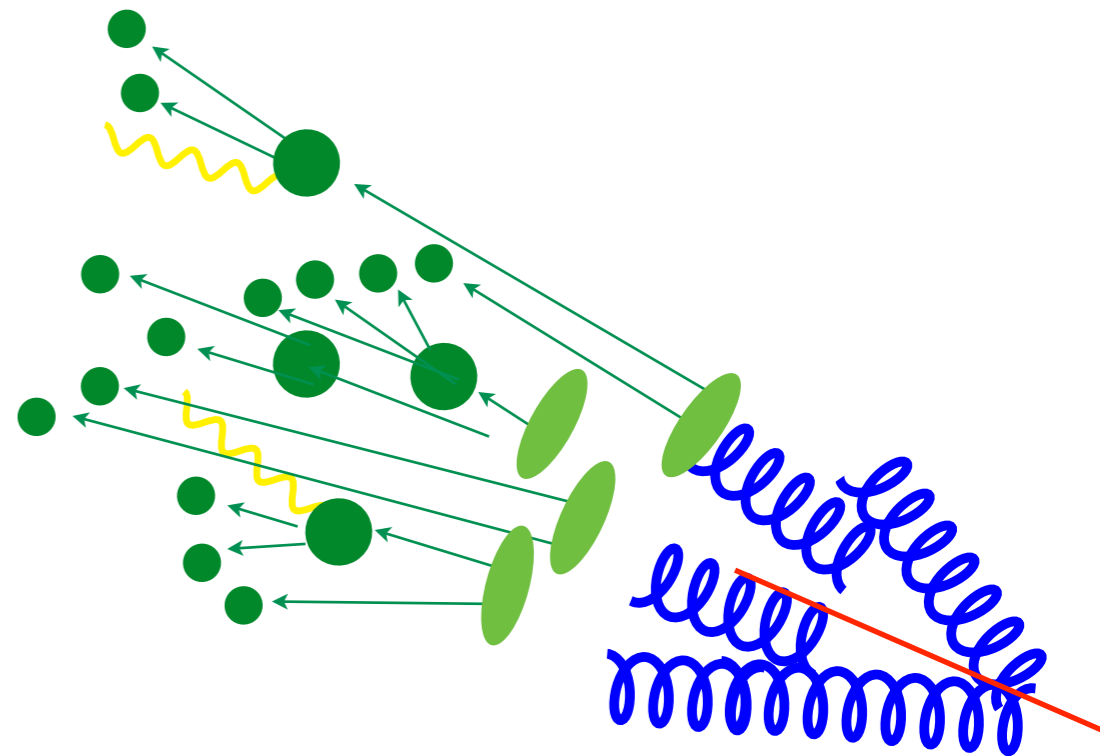
Benjamin Nachman

*Lawrence Berkeley National Laboratory*

*HEP Graduate Workshop  
Batna, Algeria, 2022*



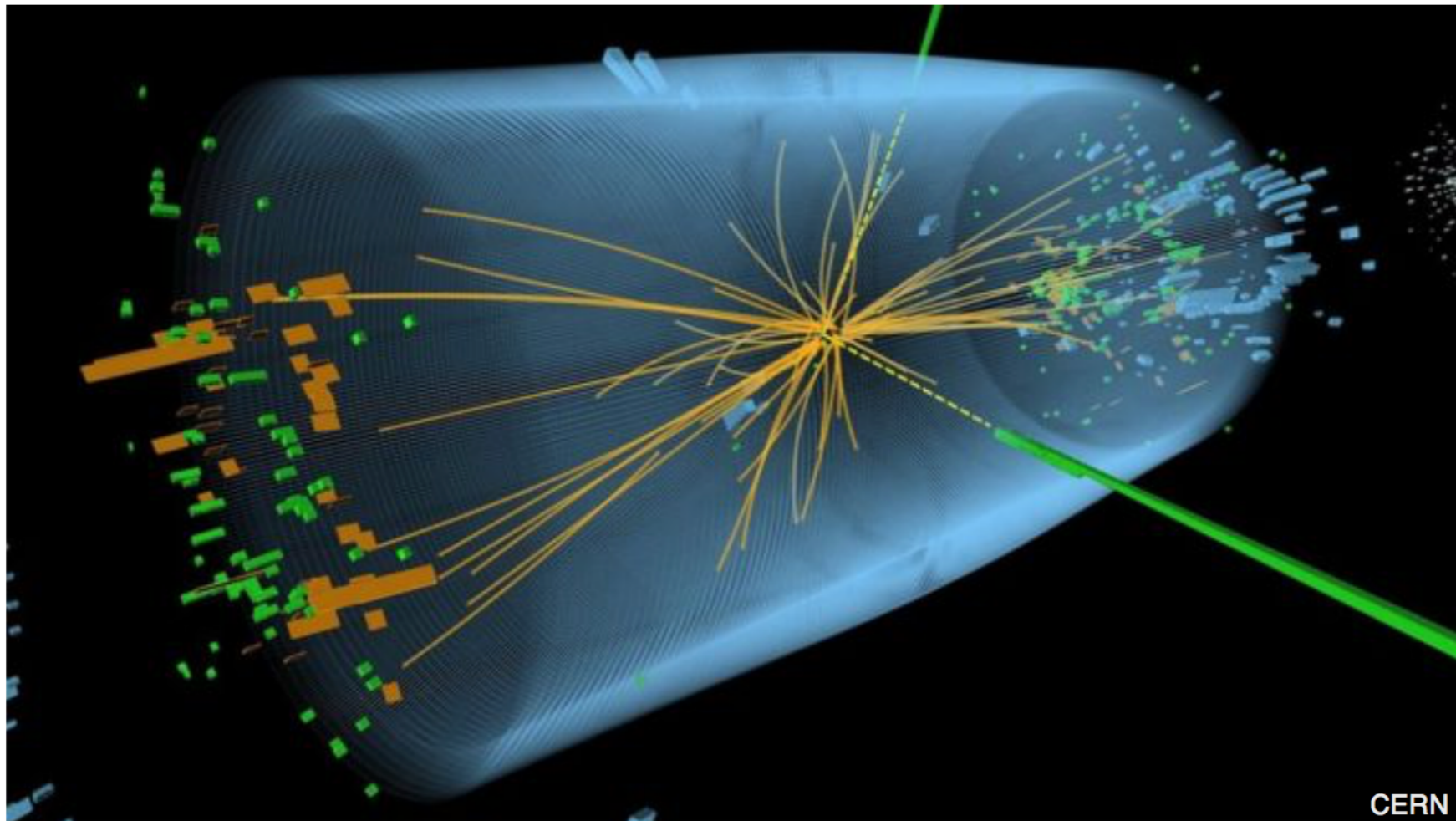
1. The Large Hadron Collider as a gluon factory
2. Overview of the QCD physics program at the LHC
3. Experimental considerations
  1. Calibration
  2. Unfolding
4. Physics topics
  1. PDFs
  2.  $\alpha_s$
  3. jet substructure
  4. other



# What is the goal of the LHC?



BBC, March 2018



A candidate Higgs boson collision event at CMS, one of the experiments at the Large Hadron Collider

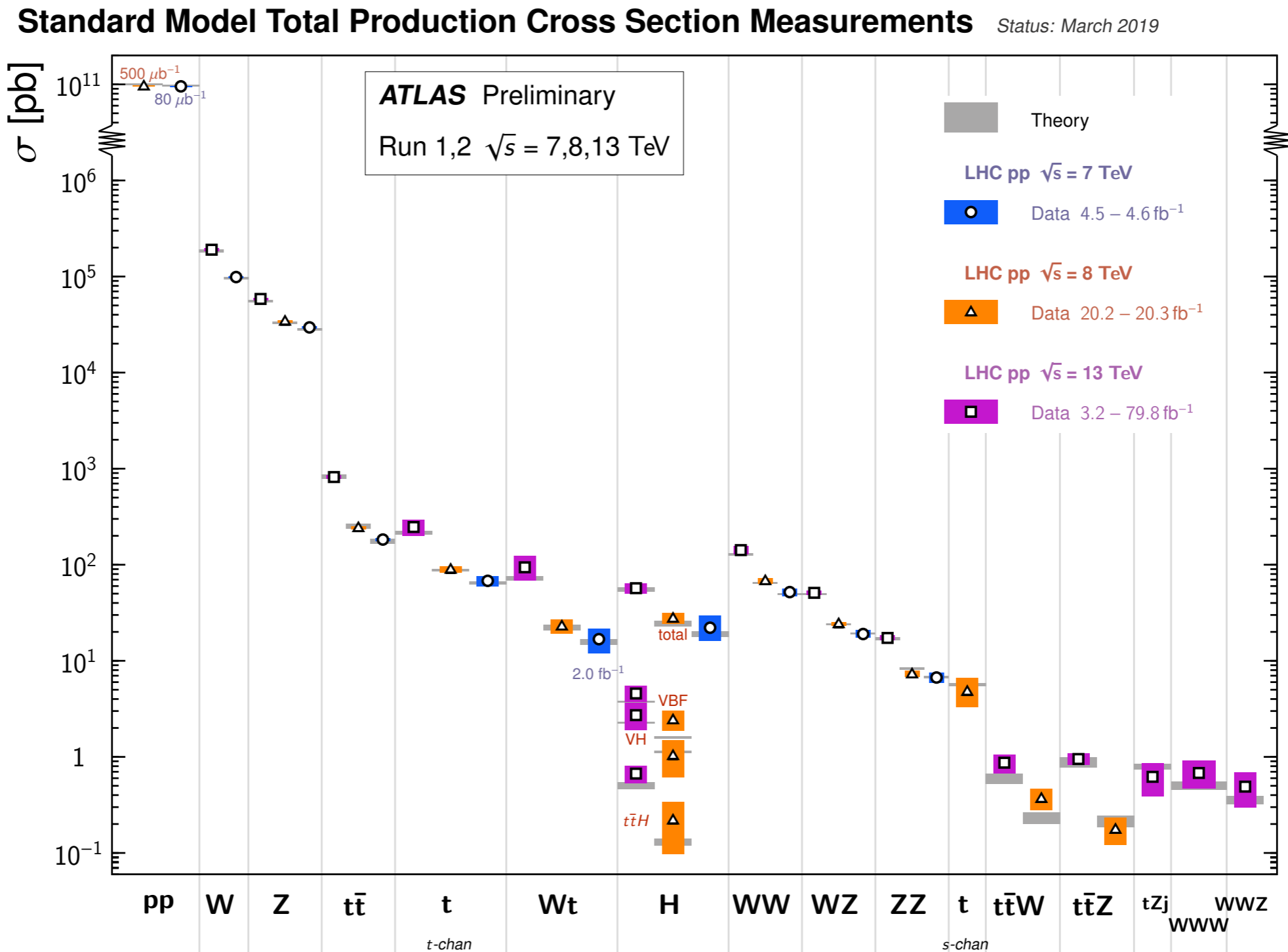
A top physicist says the construction of a "factory" to produce Higgs boson particles is a priority for the science community.

In an exclusive interview, Nigel Lockyer, head of America's premier particle physics lab, said studying the Higgs could hasten major discoveries.

# Is the LHC a Higgs factory?



Rate at which  
processes happen

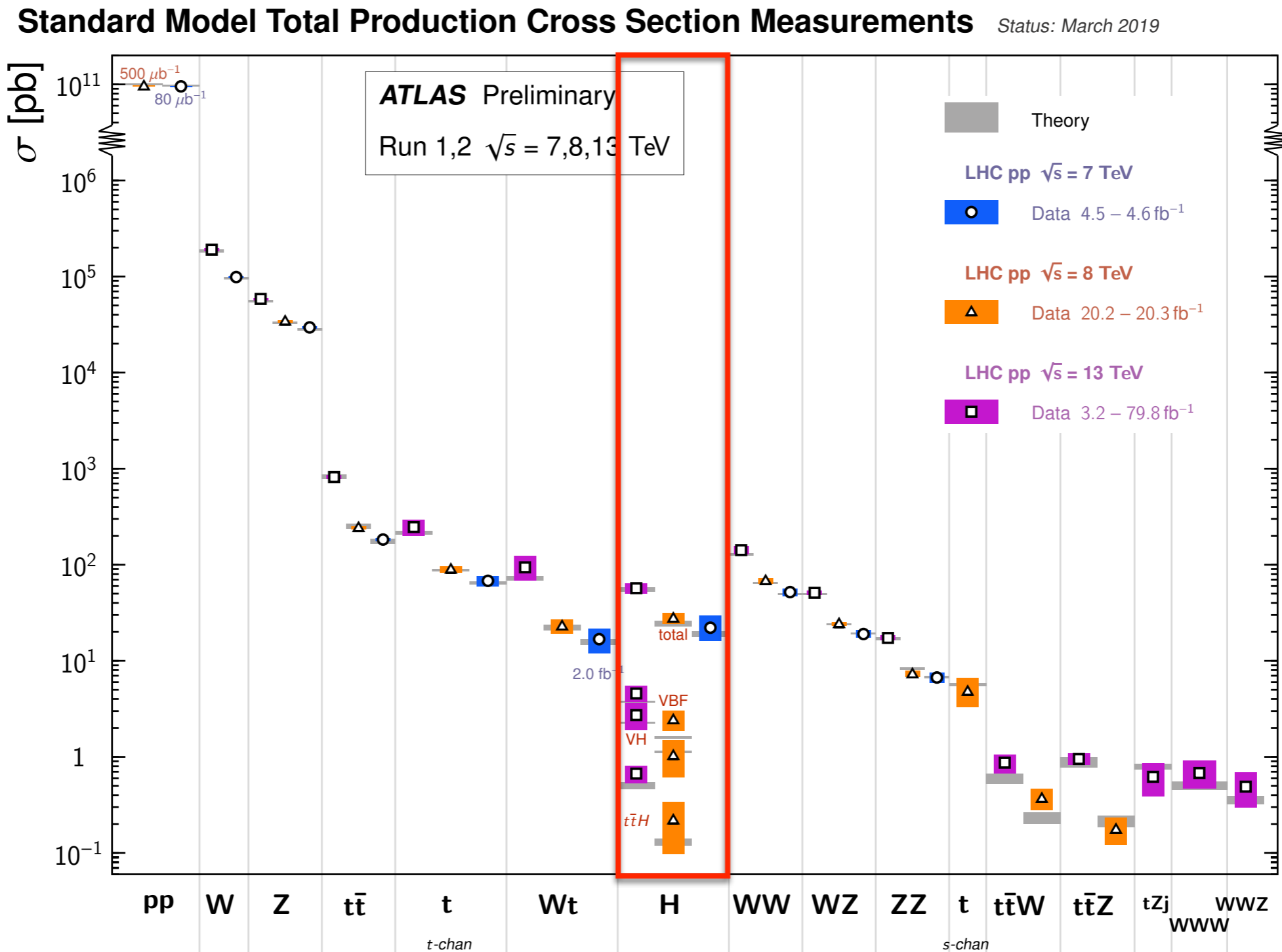


Process name: proton + proton goes to X

# Is the LHC a Higgs factory?



Rate at which  
processes happen

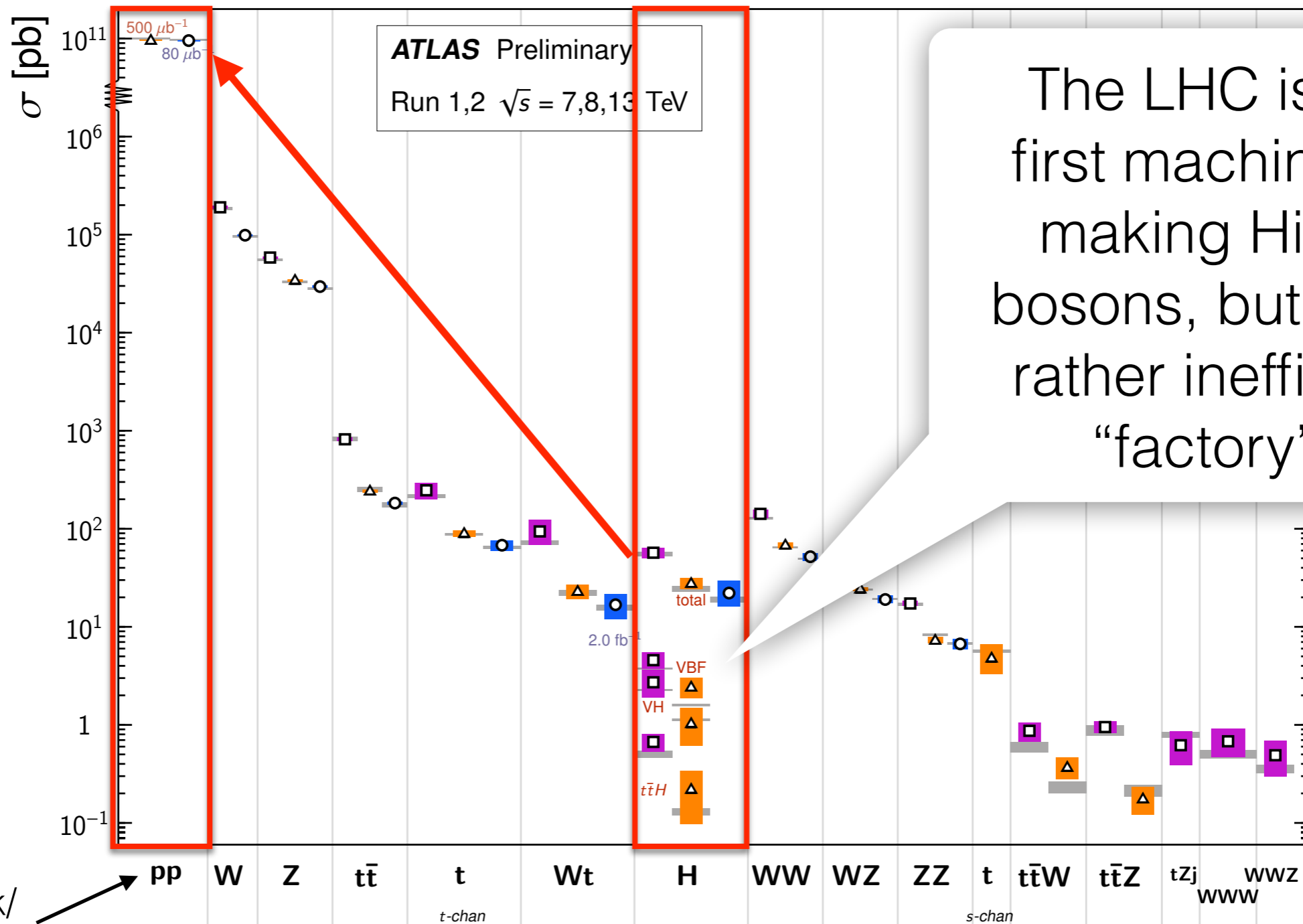


Process name: proton + proton goes to X

# Is the LHC a Higgs factory?



Standard Model Total Production Cross Section Measurements *Status: March 2019*



Rate at which processes happen

The LHC is the first machine for making Higgs bosons, but it is a rather inefficient "factory" !

Process name: proton + proton goes to X

# *Higgs versus gluons*



nature. One major ingredient in this model is a hypothetical, ubiquitous quantum field that is supposed to be responsible for giving particles their masses (this field would answer the basic question of why particles have the masses they do--or indeed, why they have any mass at all).

This field is called the **Higgs field**. As a consequence of wave-particle

- Scientific American

# Higgs versus gluons

8

nature. One major ingredient in this model is a hypothetical, ubiquitous quantum field that is supposed to be responsible for giving particles their masses (this field would answer the basic question of why particles have the masses they do--or indeed, why they have any mass at all).

This field is called the **Higgs field**. As a consequence of wave-particle

- Scientific American



This is true, but it turns out that the gluon is responsible for most of our mass - without them, the proton mass would be many orders of magnitude lighter...

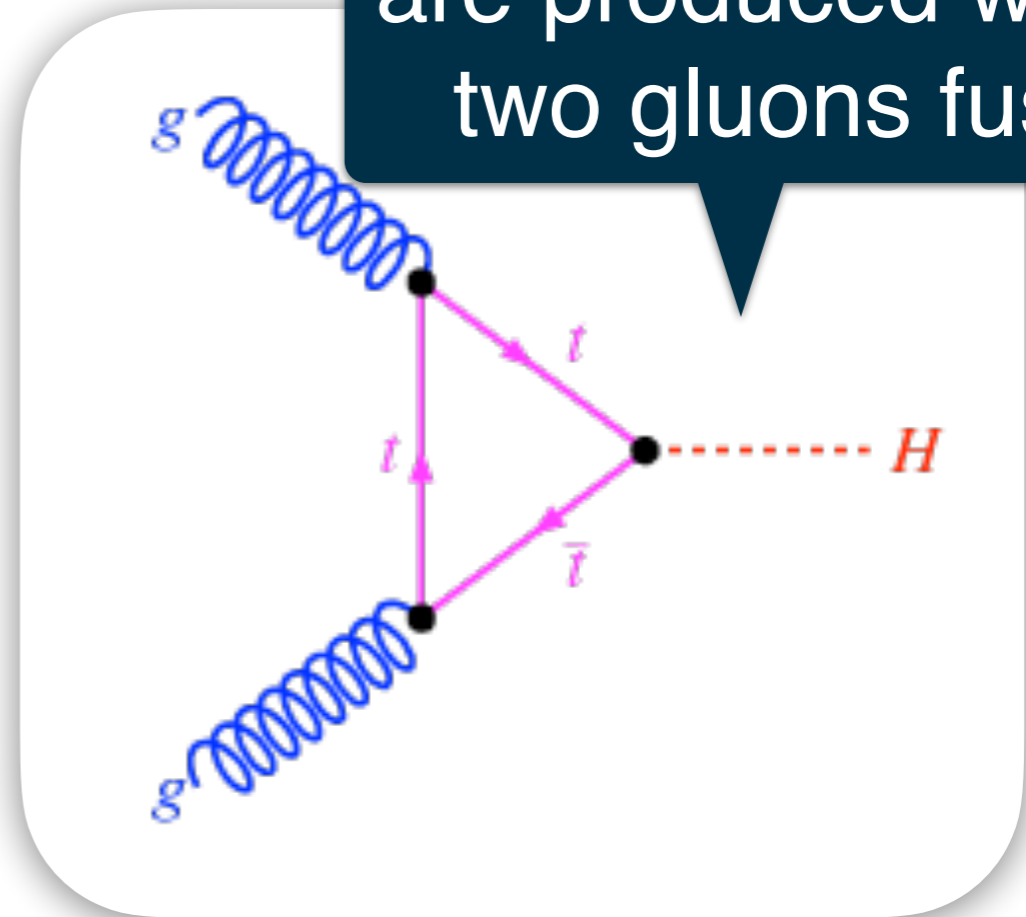
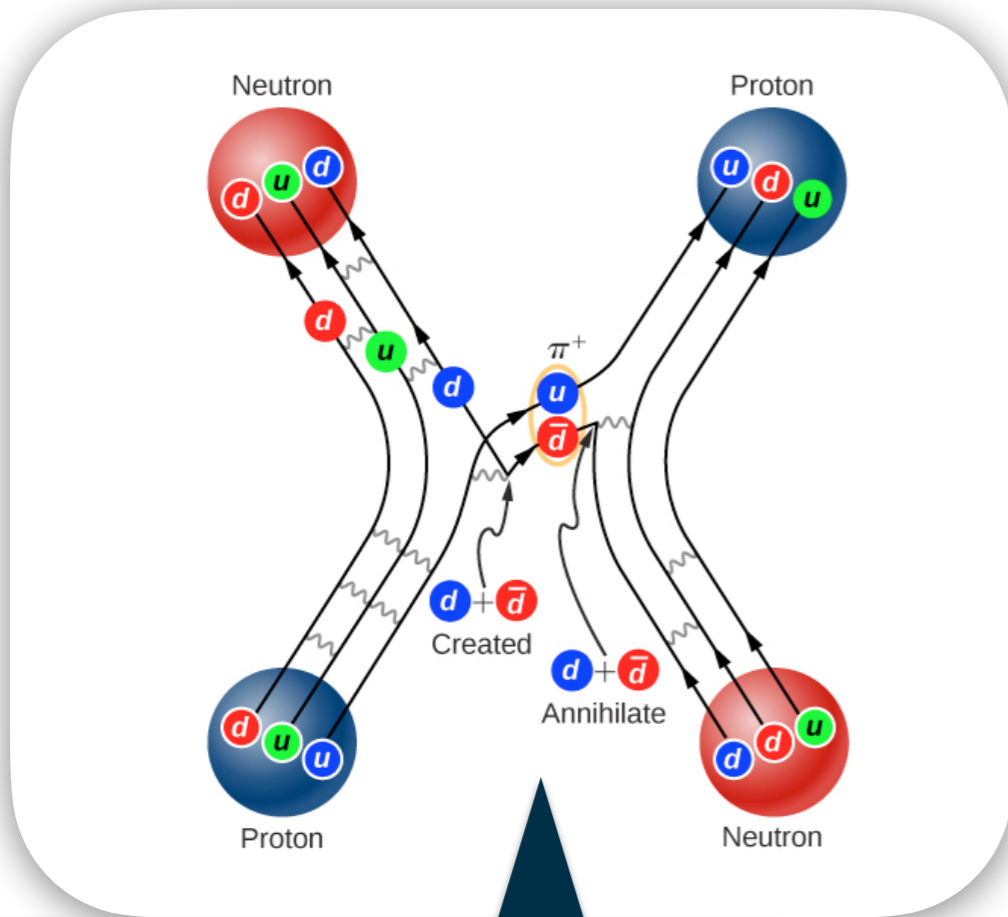


# The Higgs-gluon connection



There is a **strong** connection between the Higgs boson and gluons.

Most Higgs bosons are produced when two gluons fuse



Due to the Higgs boson, the strong force is short-ranged

...okay, enough about the Higgs boson for now - more about the strong force!

# QCD physics program at the LHC

10

1. Photon/W/Z+jets
2. Jet physics
3. Jet substructure
4. “Soft QCD” & Heavy Ions
5. Quarkonia

Parton Distribution  
Functions (PDFs)

Perturbative QCD /  
strong coupling

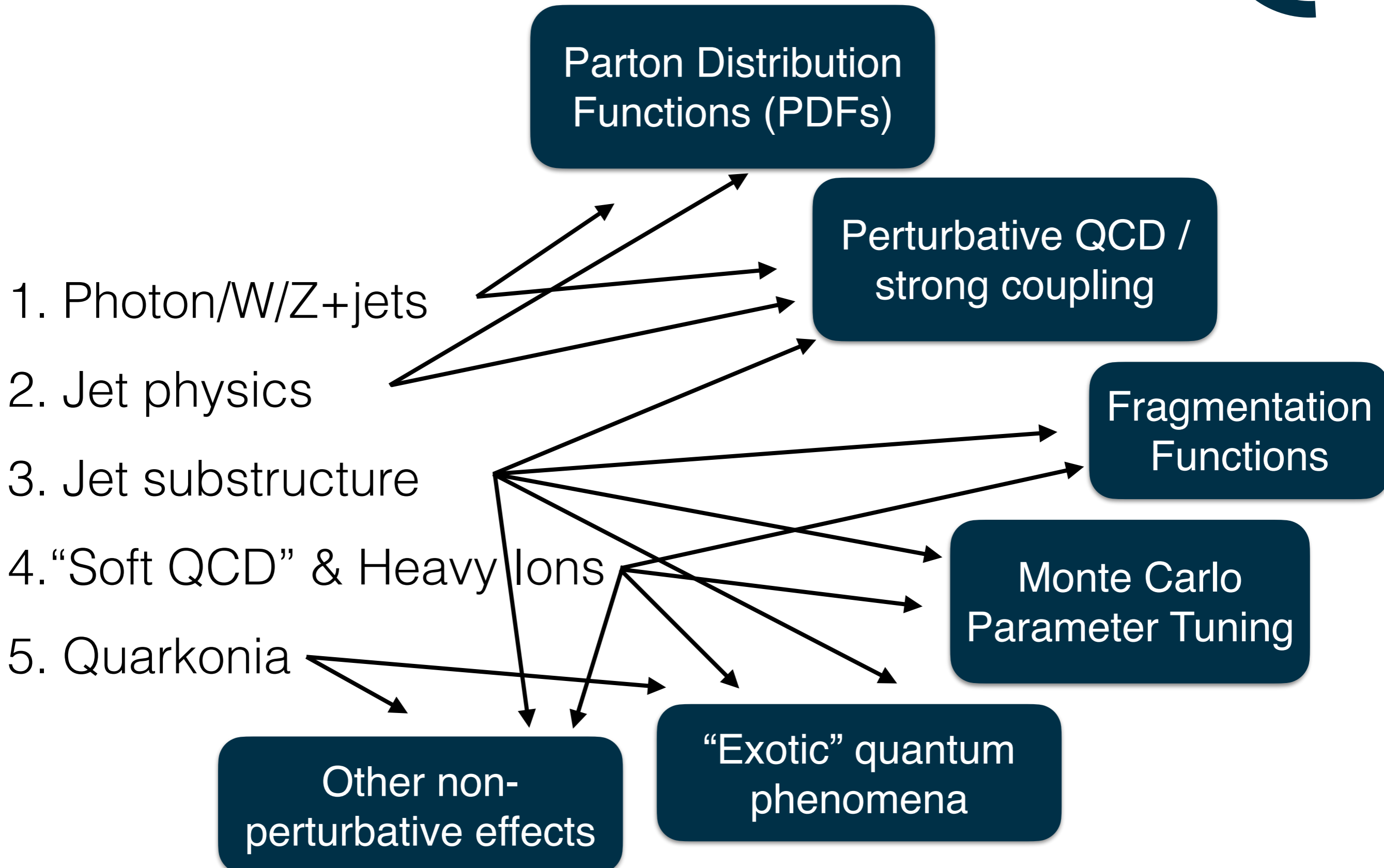
Fragmentation  
Functions

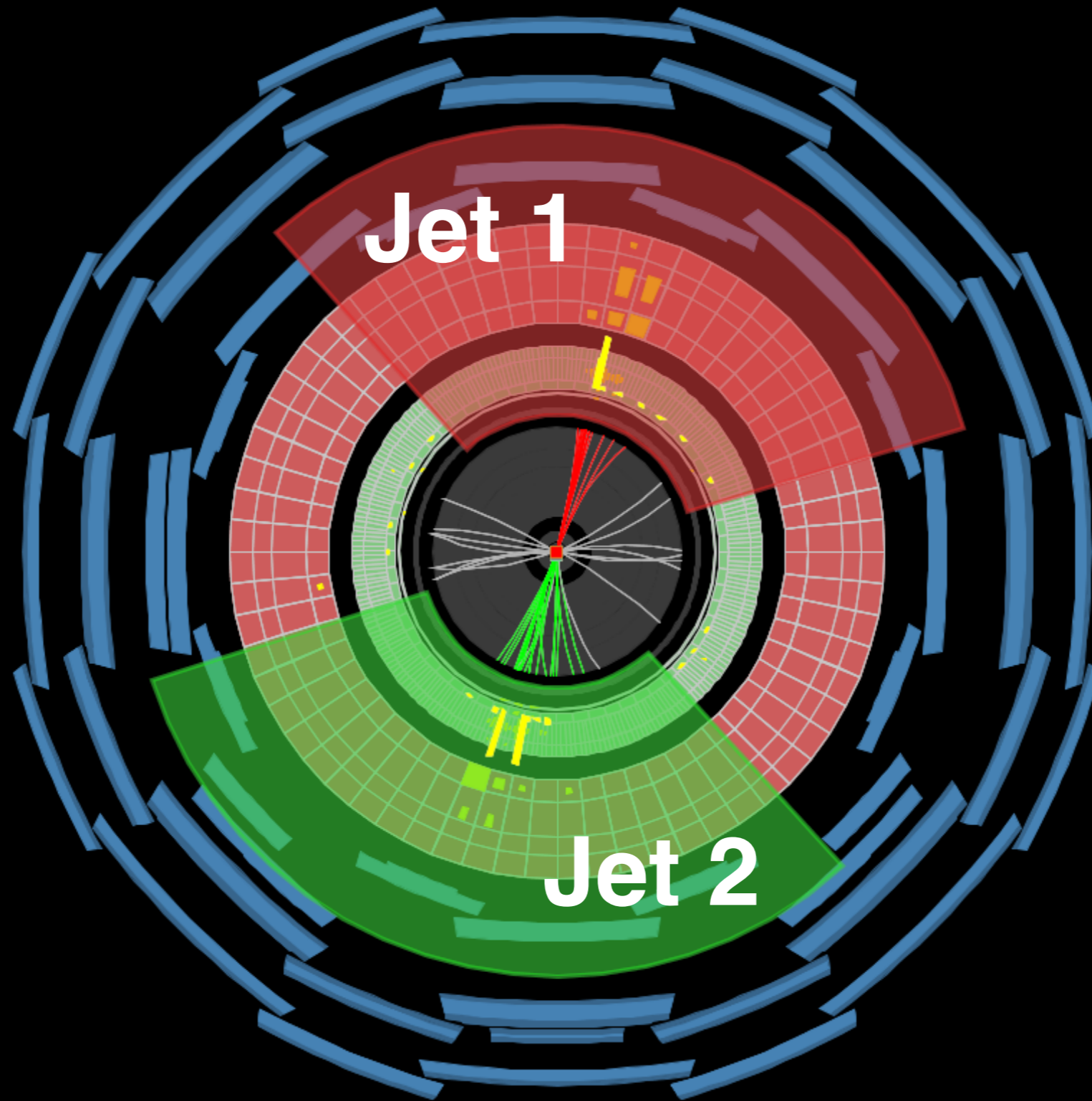
Monte Carlo  
Parameter Tuning

Other non-  
perturbative effects

“Exotic” quantum  
phenomena

# QCD physics program at the LHC

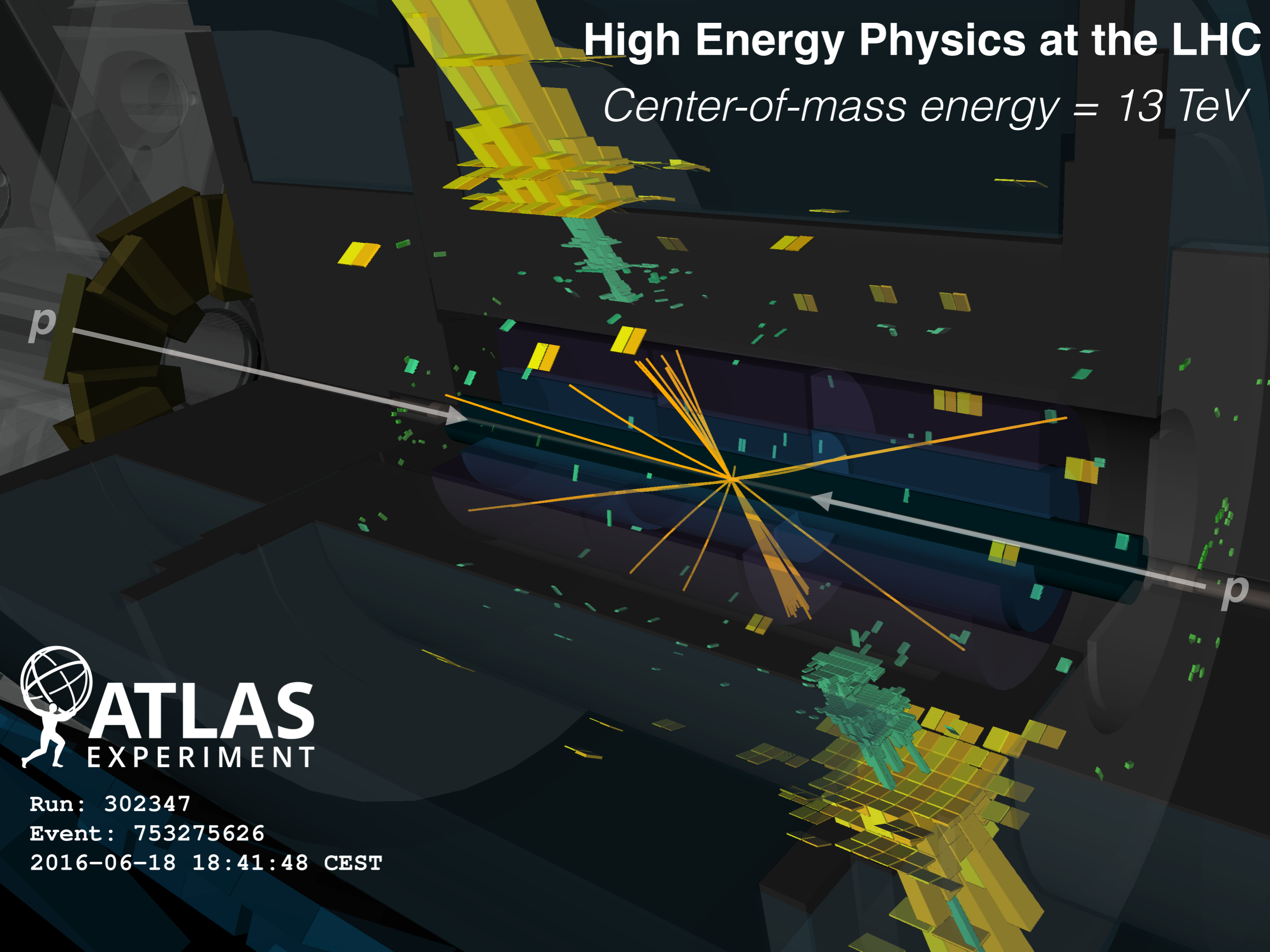




*collisions in/out of page*

# High Energy Physics at the LHC

*Center-of-mass energy = 13 TeV*



 **ATLAS**  
EXPERIMENT

Run: 302347

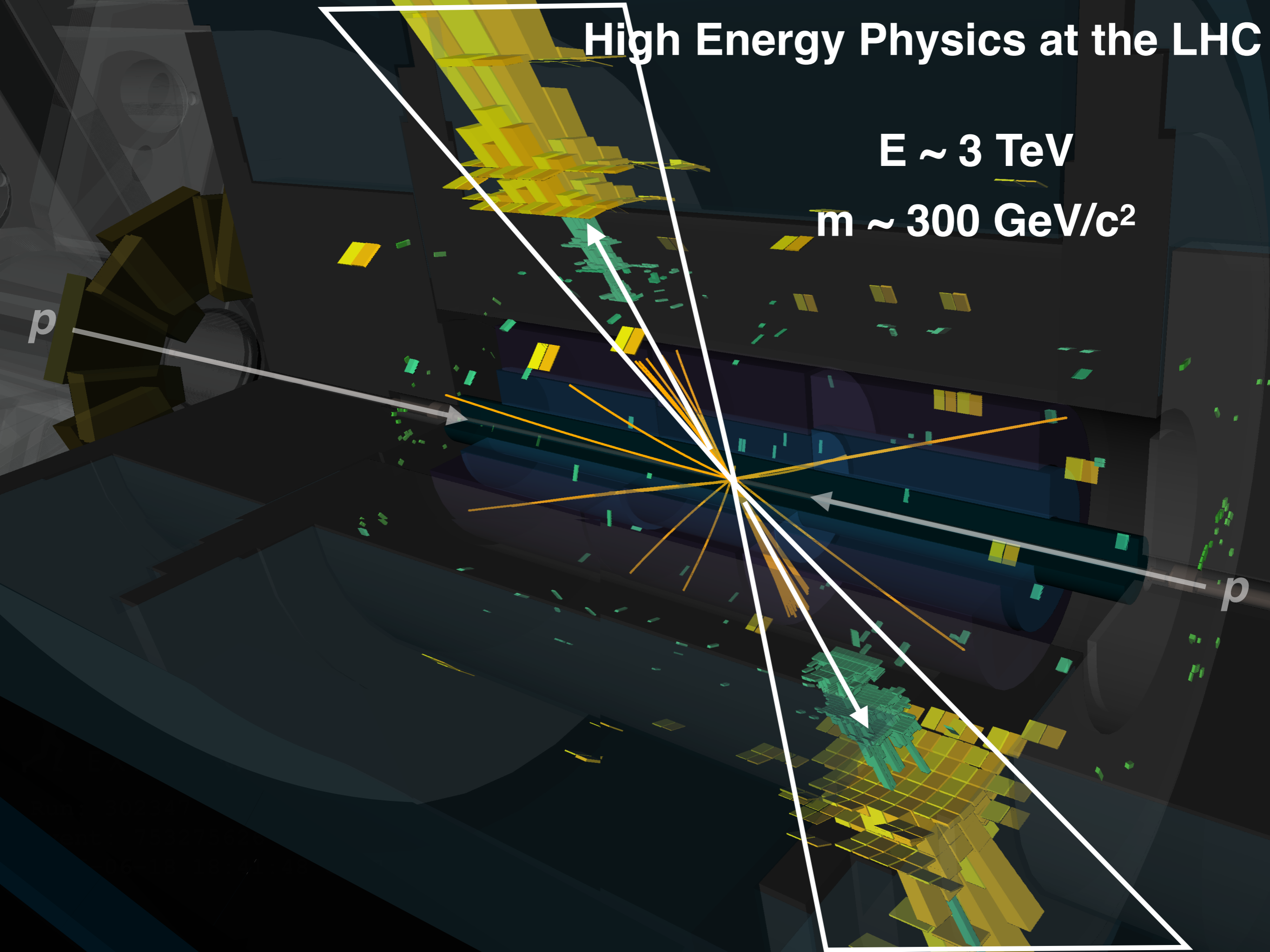
Event: 753275626

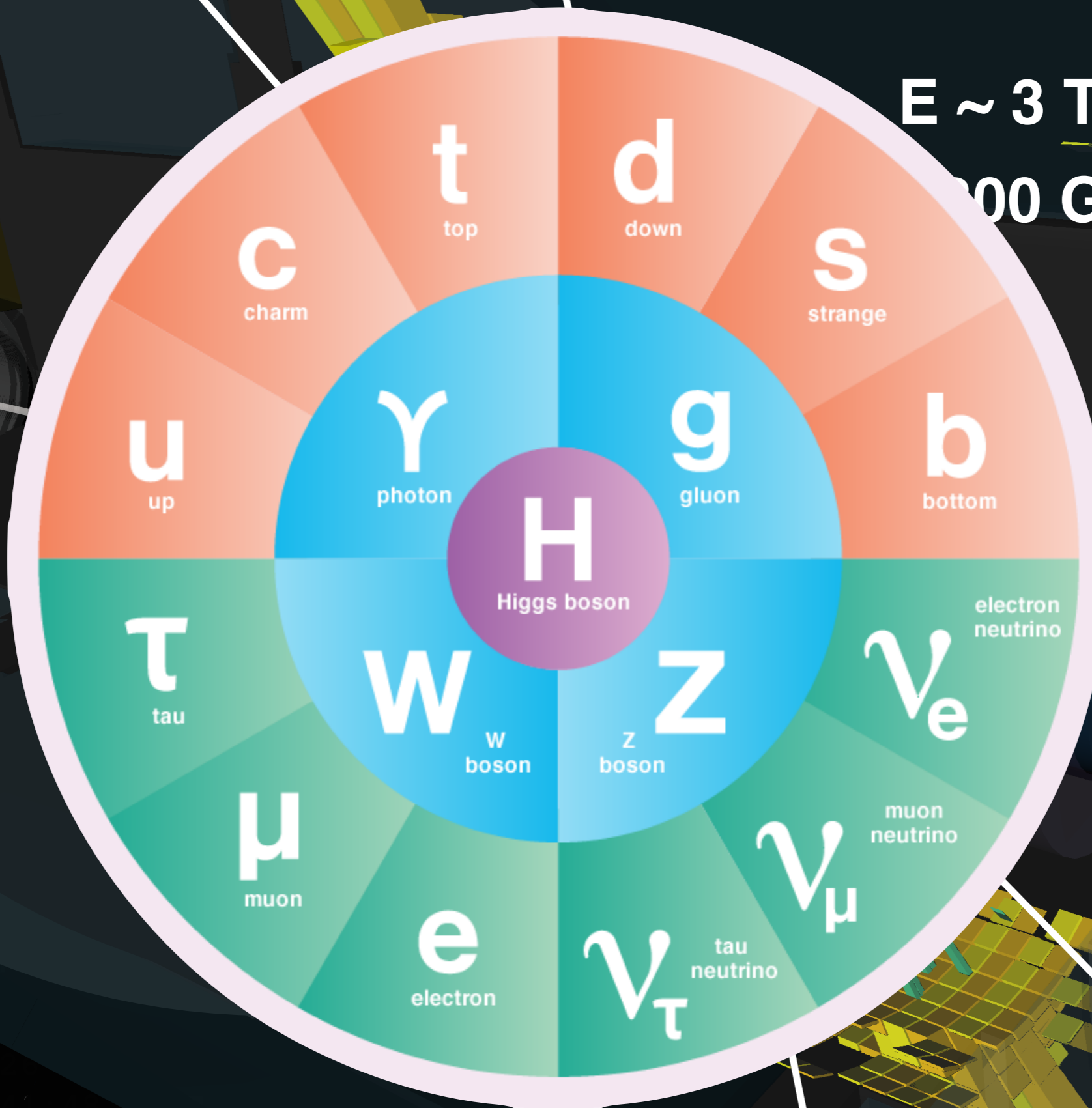
2016-06-18 18:41:48 CEST

# High Energy Physics at the LHC

$E \sim 3 \text{ TeV}$

$m \sim 300 \text{ GeV}/c^2$





$E \sim 3 \text{ TeV}$

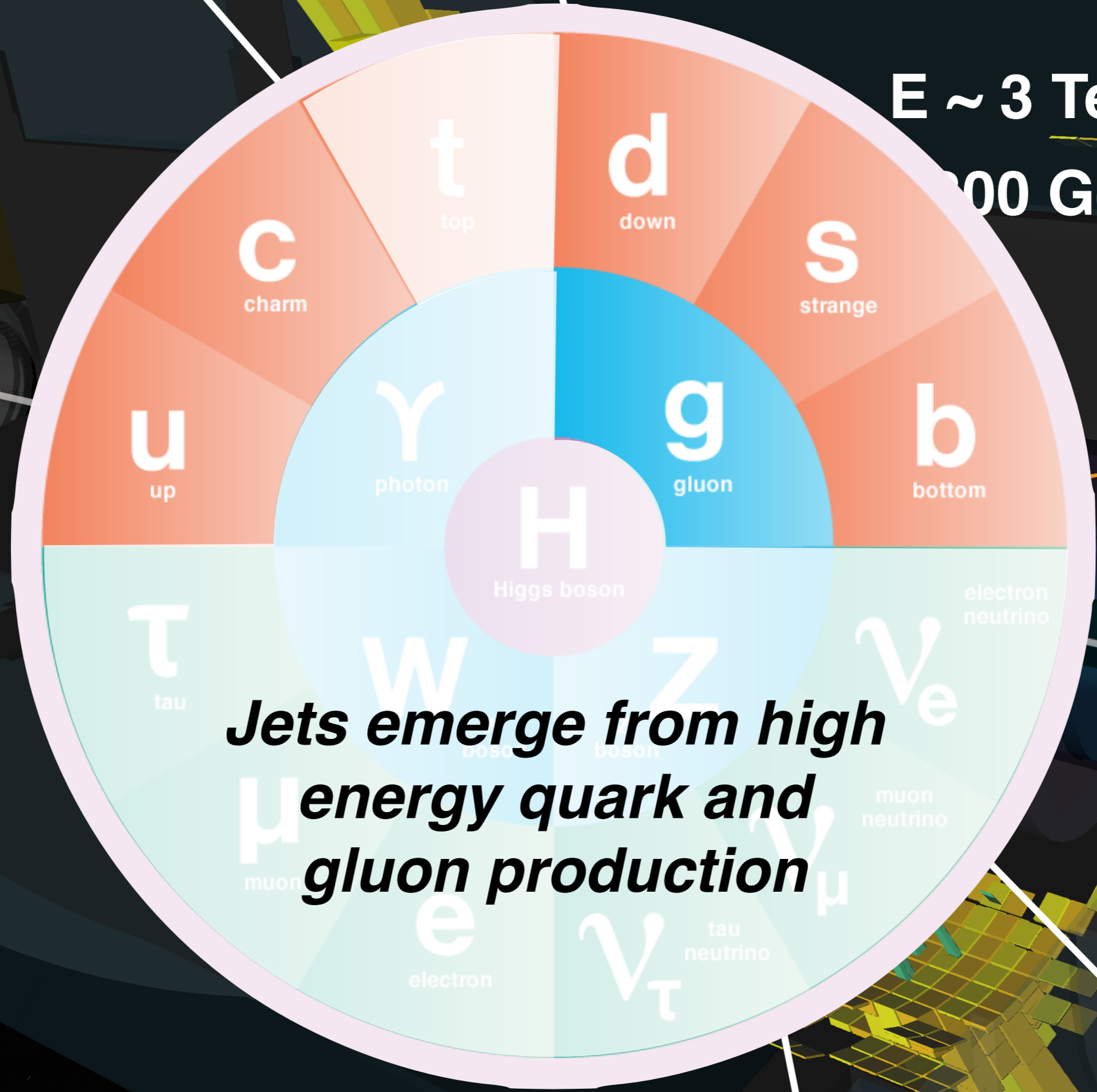
$100 \text{ GeV}/c^2$

p

p

$E \sim 3 \text{ TeV}$

$100 \text{ GeV}/c^2$



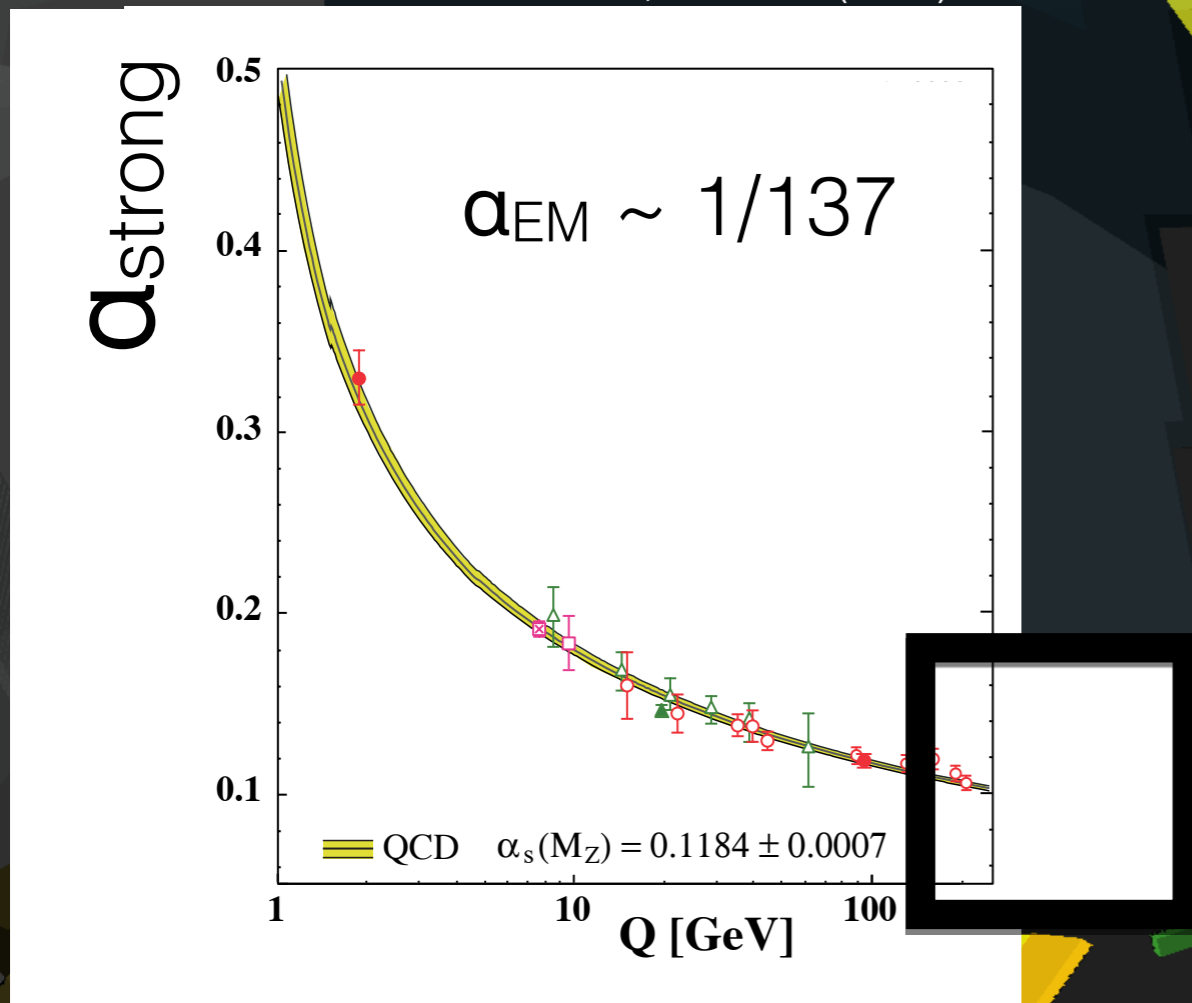
***Jets emerge from high energy quark and gluon production***

p

p



# The origin of jets



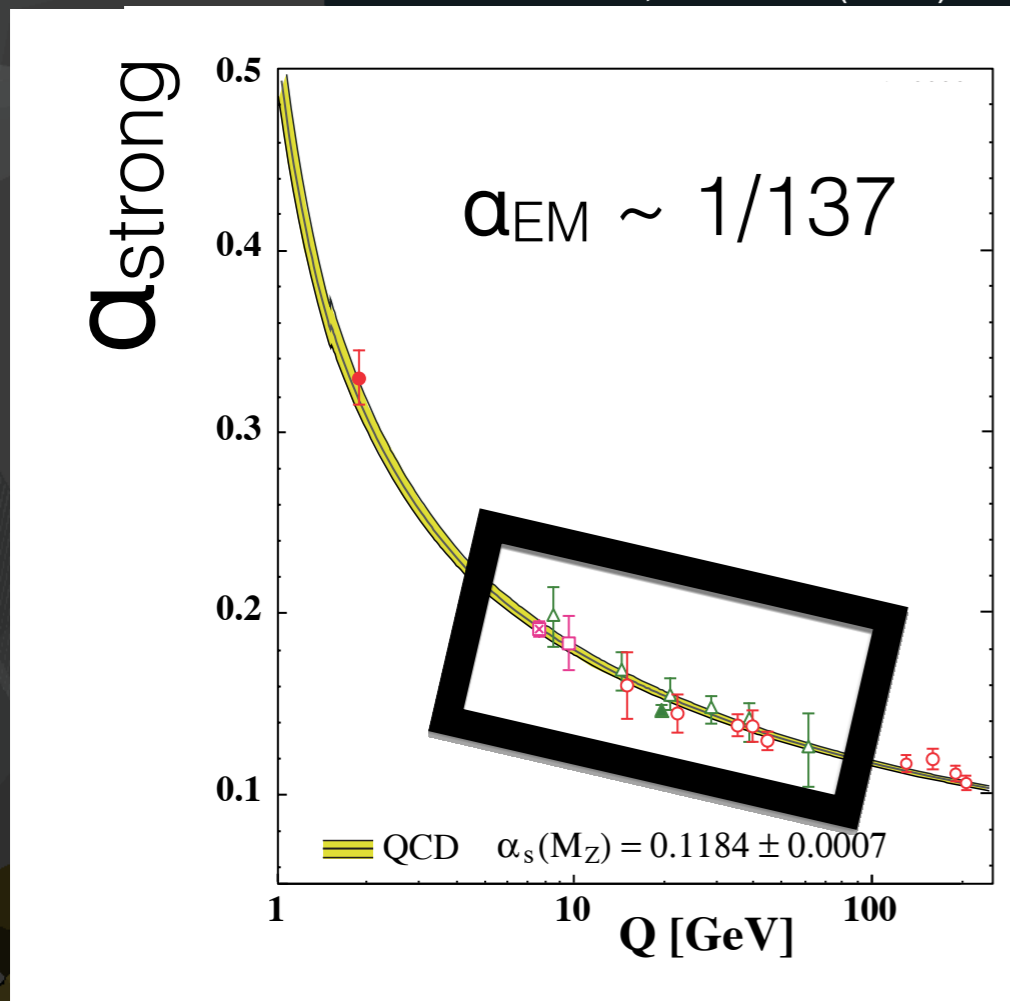
At small distance scales,  
there are quarks and gluons

They radiate gluons and  
quarks as they fly away

Quarks and gluons form  
bound states (like the proton)



# The origin of jets



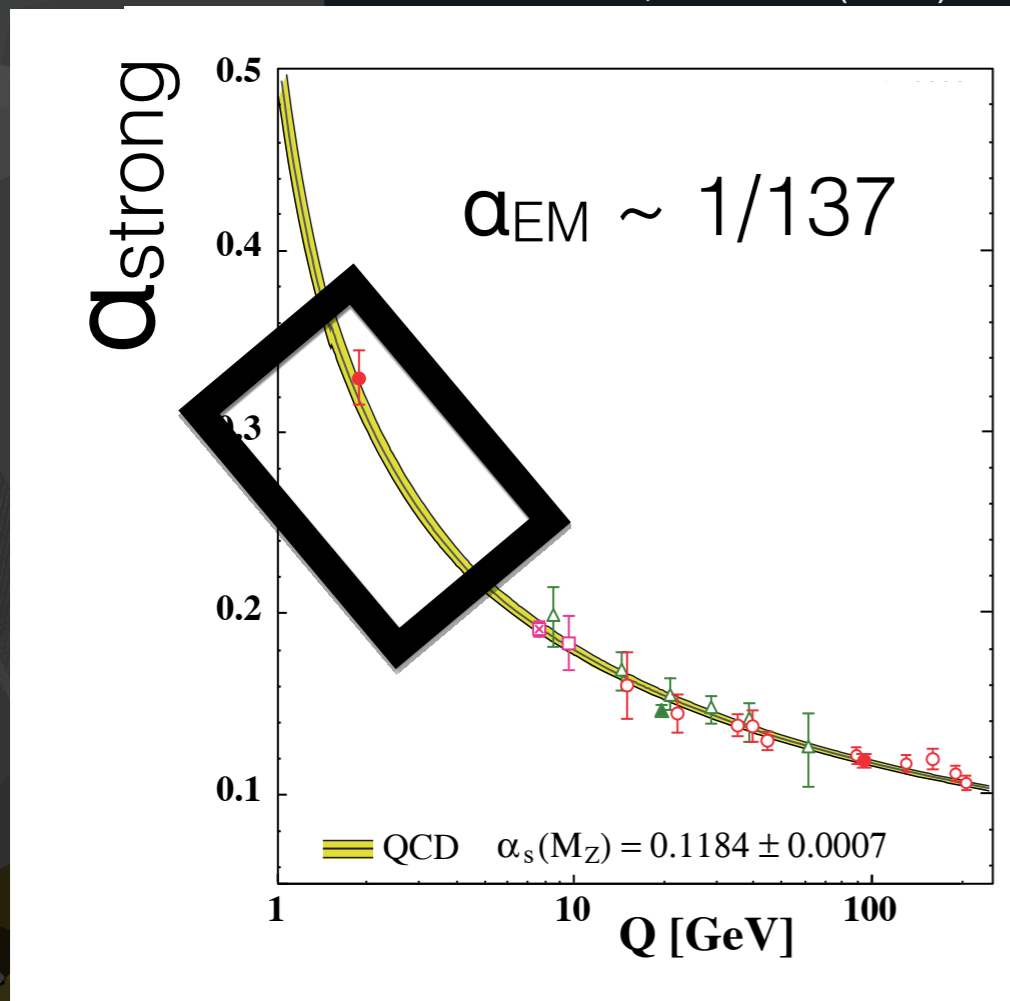
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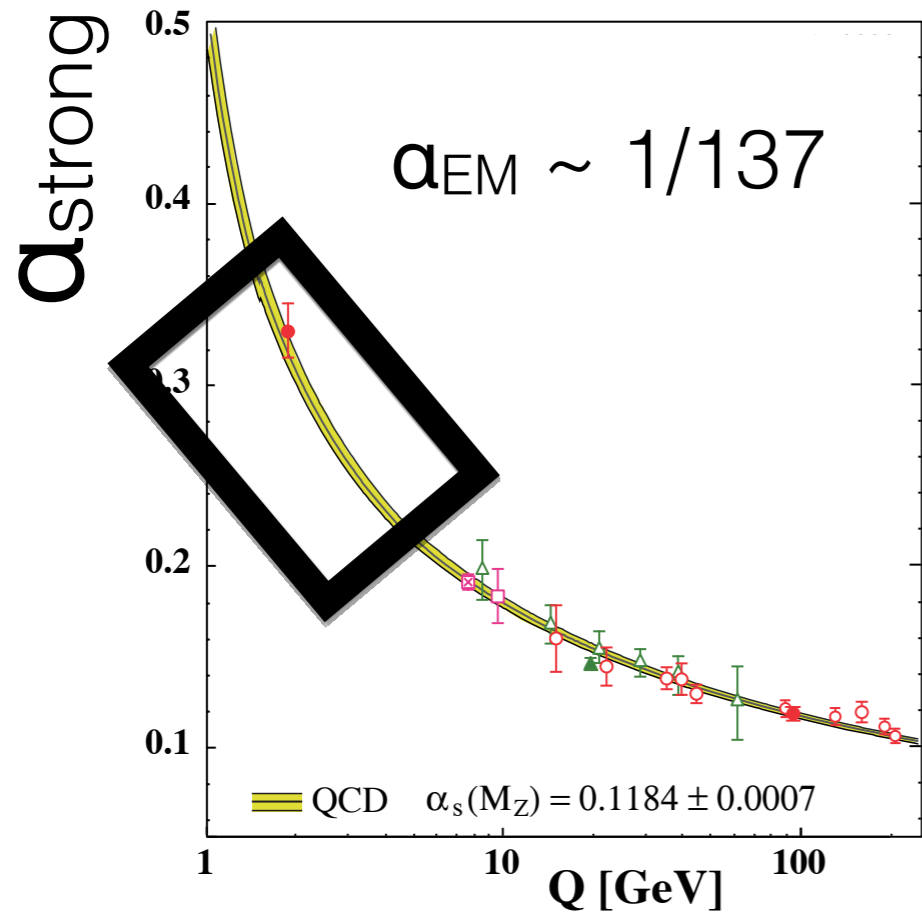


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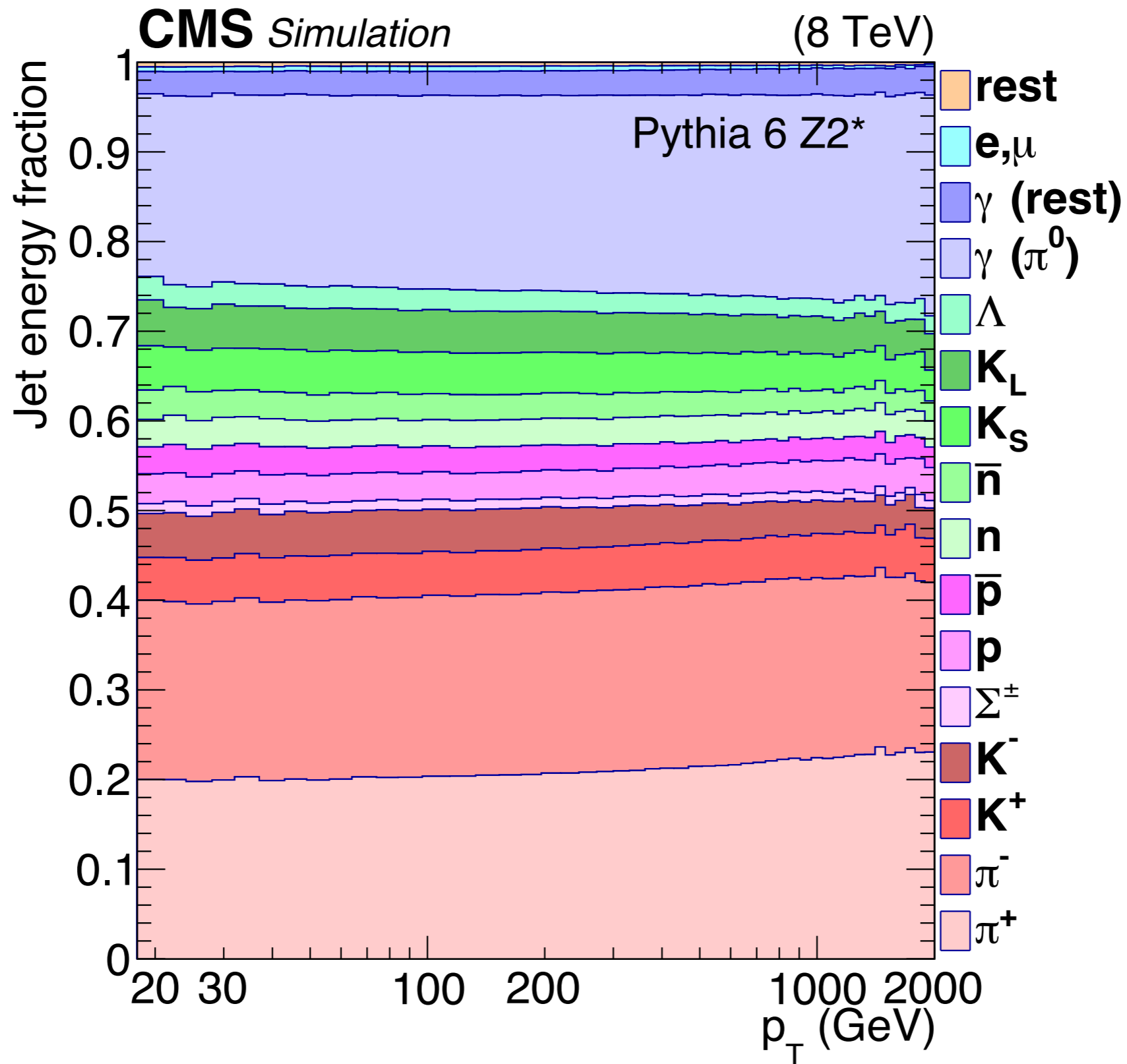
The origin  
of jets

At small distance scales,  
there are quarks and gluons  
They radiate gluons and  
quarks as they fly away  
Quarks and gluons form  
bound states (like the proton)

We measure  
these bound-  
states flying into  
our detector

# Jet constituents

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# Experimental considerations

We don't measure particles, we measure energy deposits and then infer particle properties.

Single-particle objects  
(e.g. muons)

Composite objects  
(e.g. jets)

Energy and angular biases

Energy and angular biases  
Particles in inactive material  
Particles bent out of cone  
Secondary particles  
Punchthrough

We don't measure particles, we measure energy deposits and then infer particle properties.

Single-particle objects (e.g. muons)	Composite objects (e.g. jets)
Energy and angular biases	Energy and angular biases
Usually one "bulk" correction	Particles in inactive material
	Particles bent out of cone
	Secondary particles
	Punchthrough



We don't measure particles, we measure energy deposits and then infer particle properties.

Single-particle objects  
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Usually one  
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Particles in inactive material

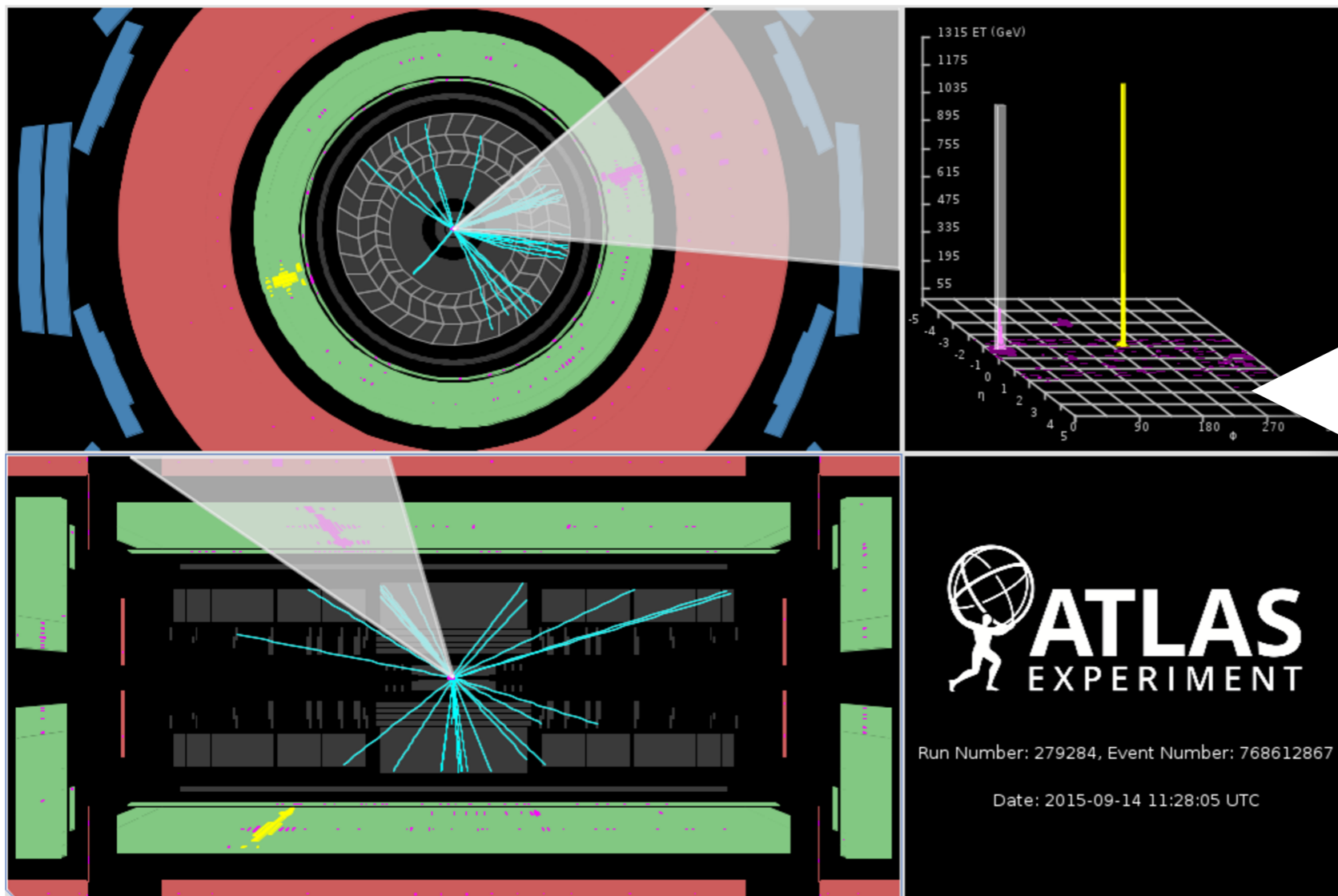
Particles bent out of cone

*Food for thought: how would you do this without depending on the prior spectrum?*

Secondary particles

Punchthrough

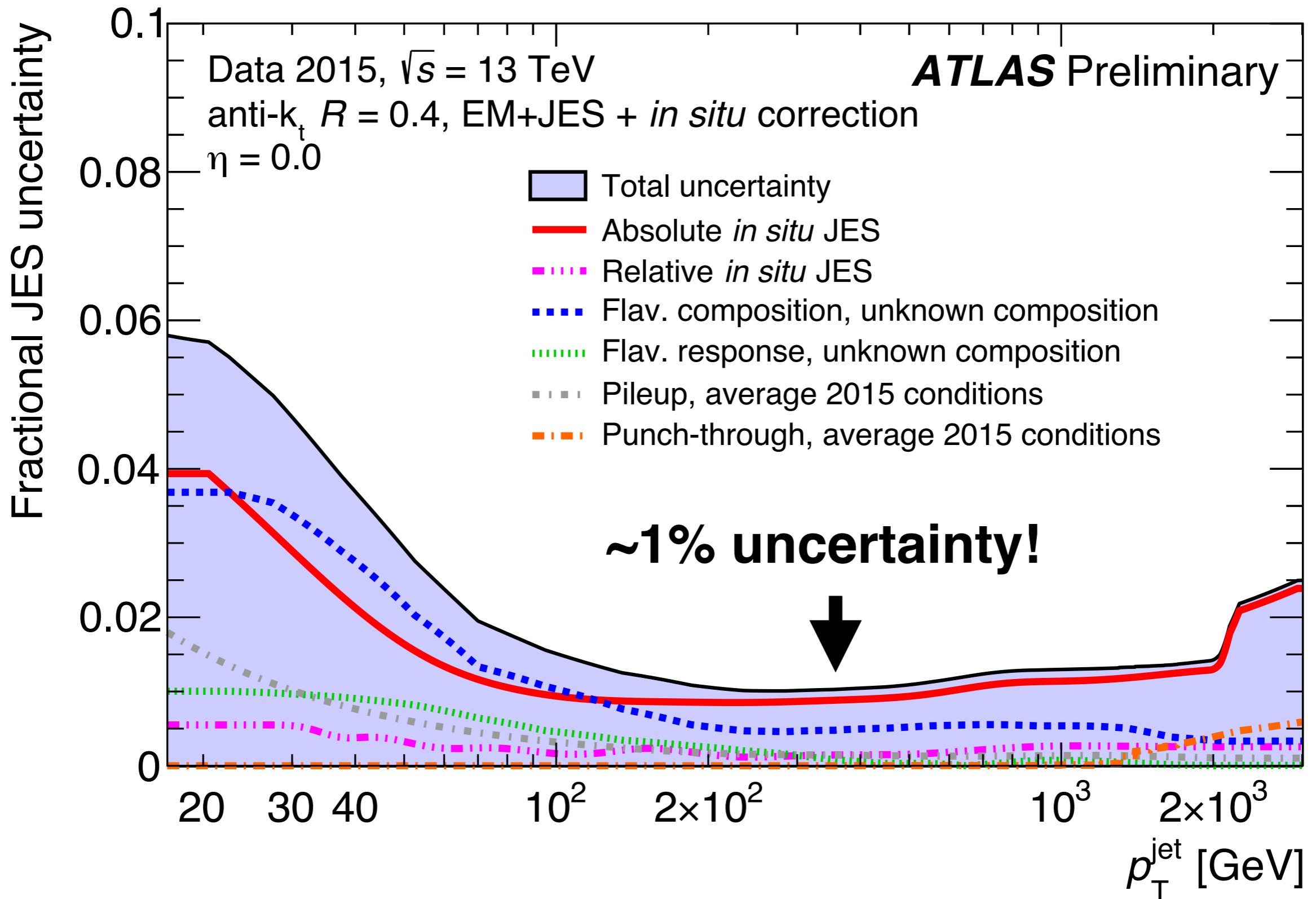
The nominal calibration is derived using simulation and then a residual calibration accounts for differences between data and simulation (derived using data).



Can use the balance of well-measured objects (e.g. photons) with jets to study the bias in data.

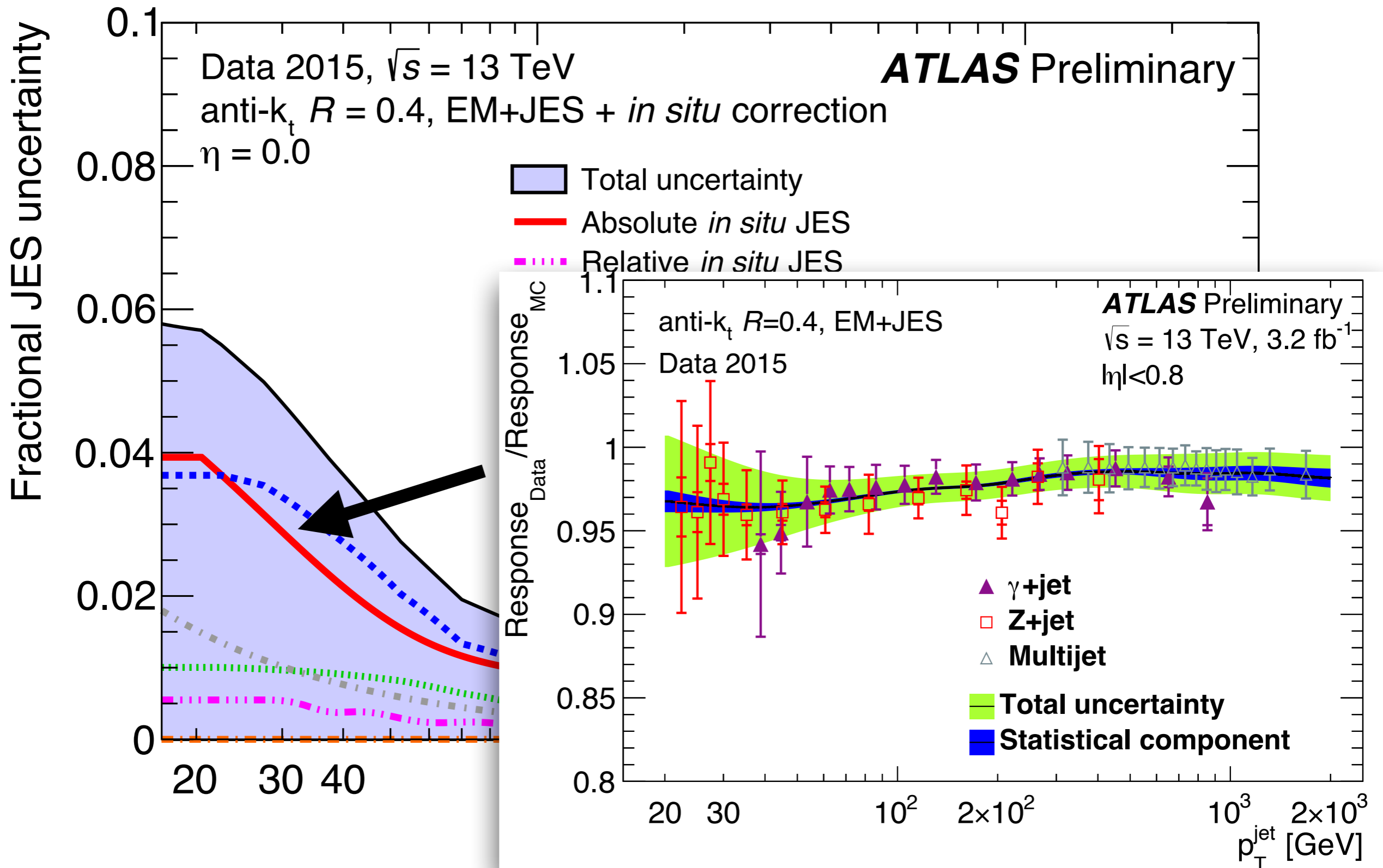
# Jet energy bias uncertainty

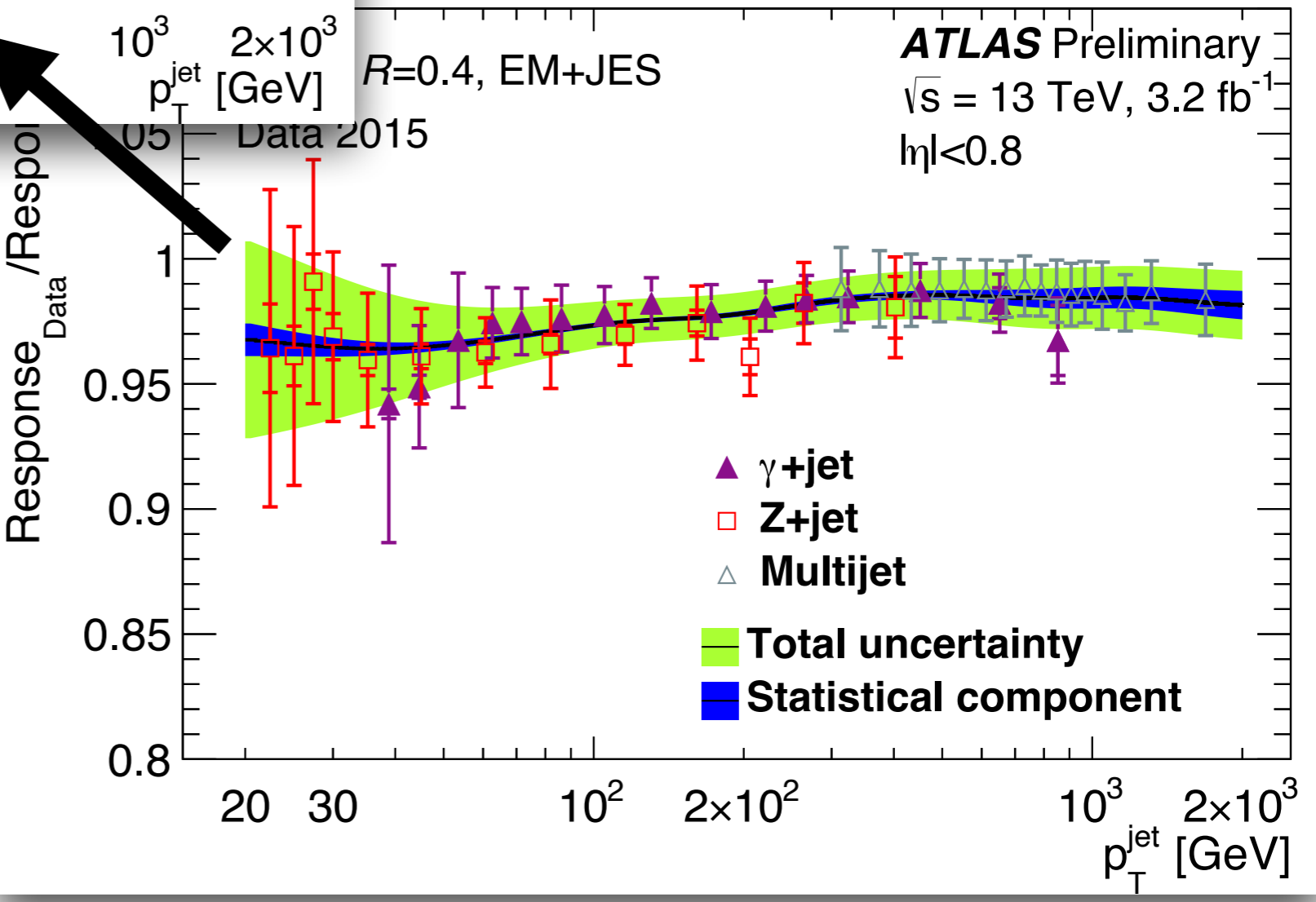
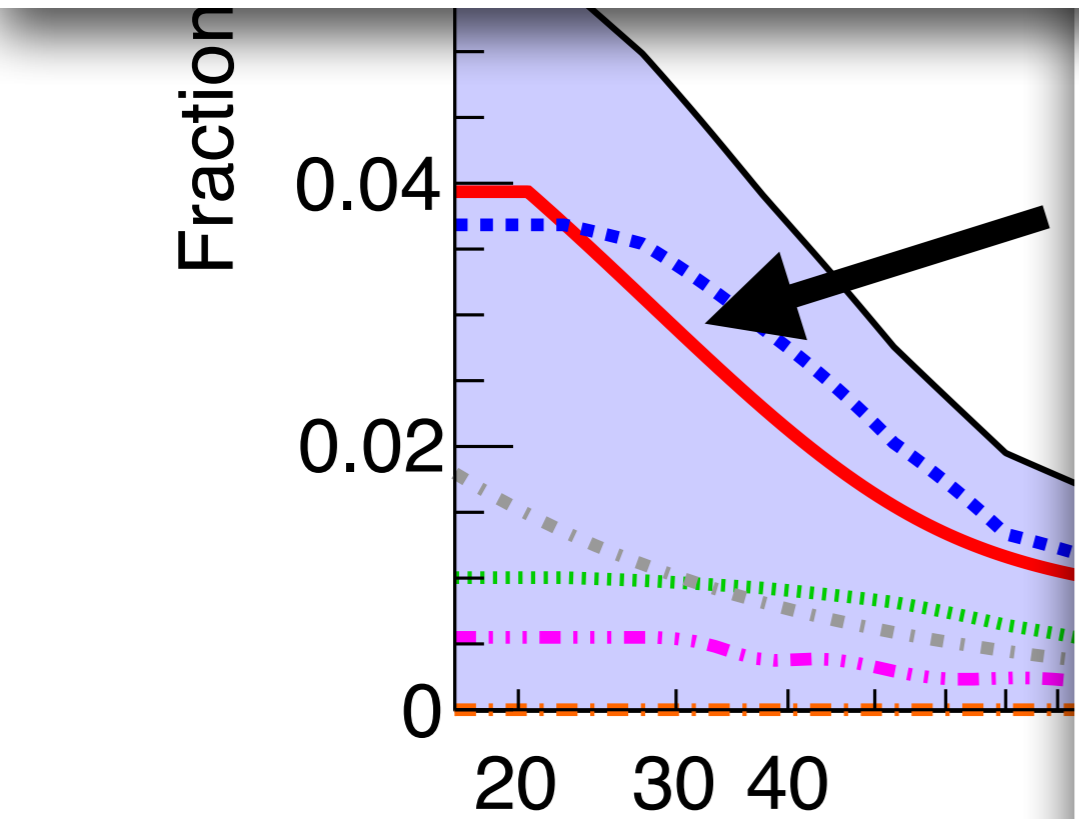
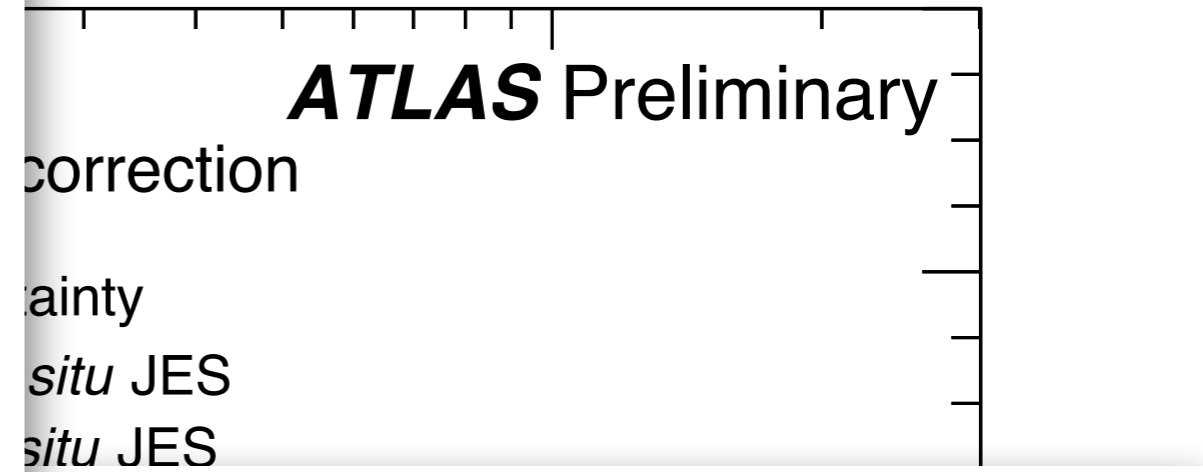
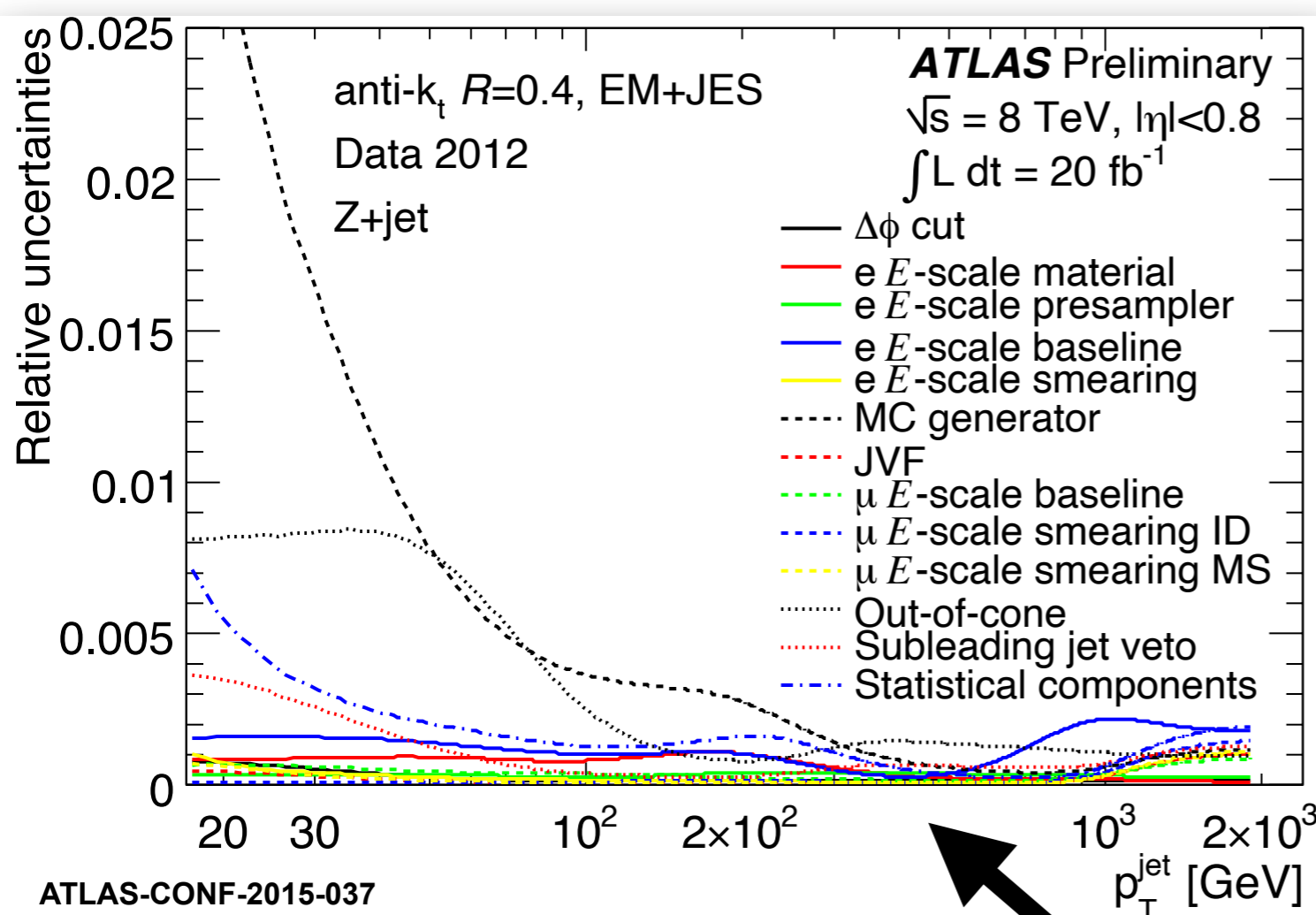
27

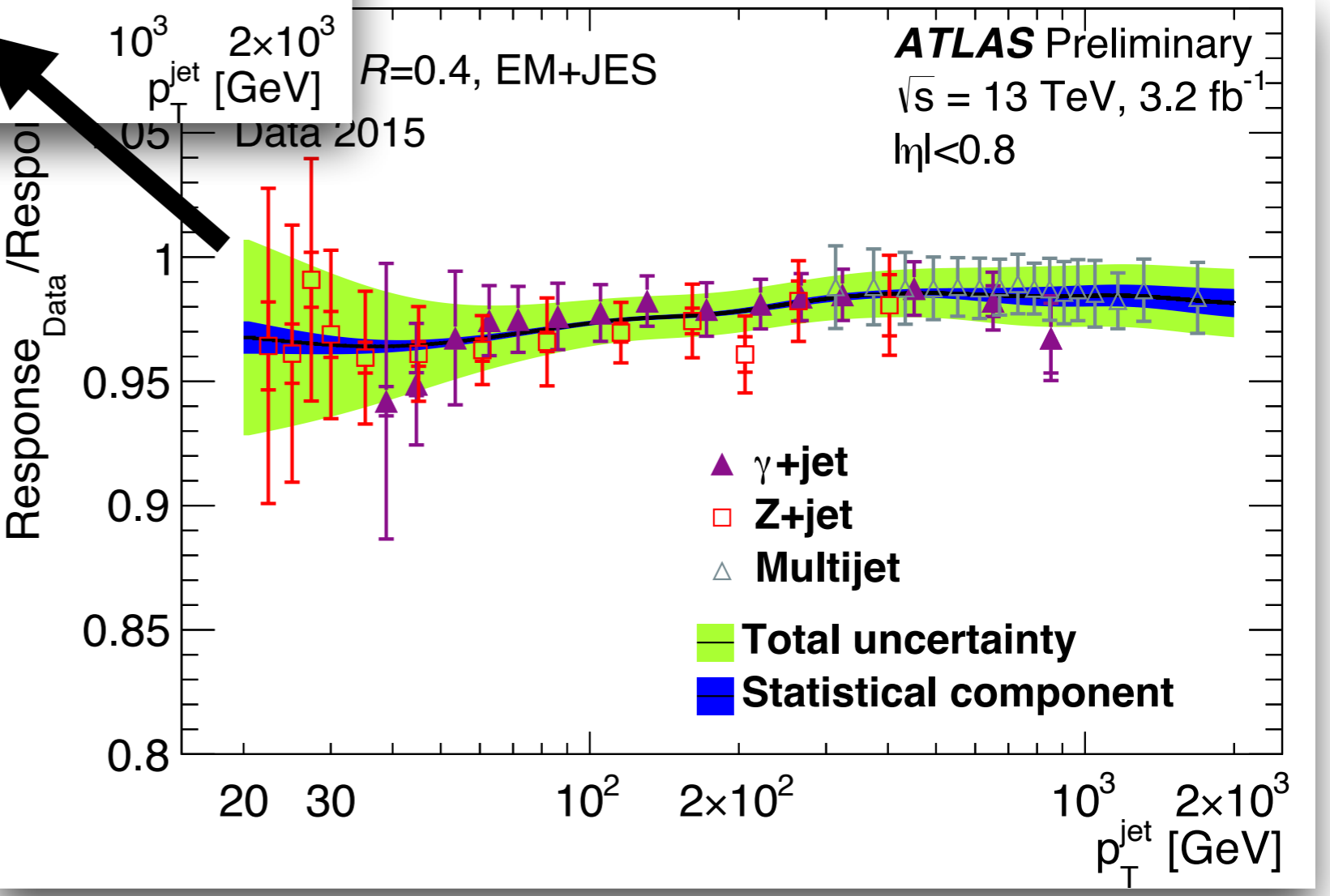
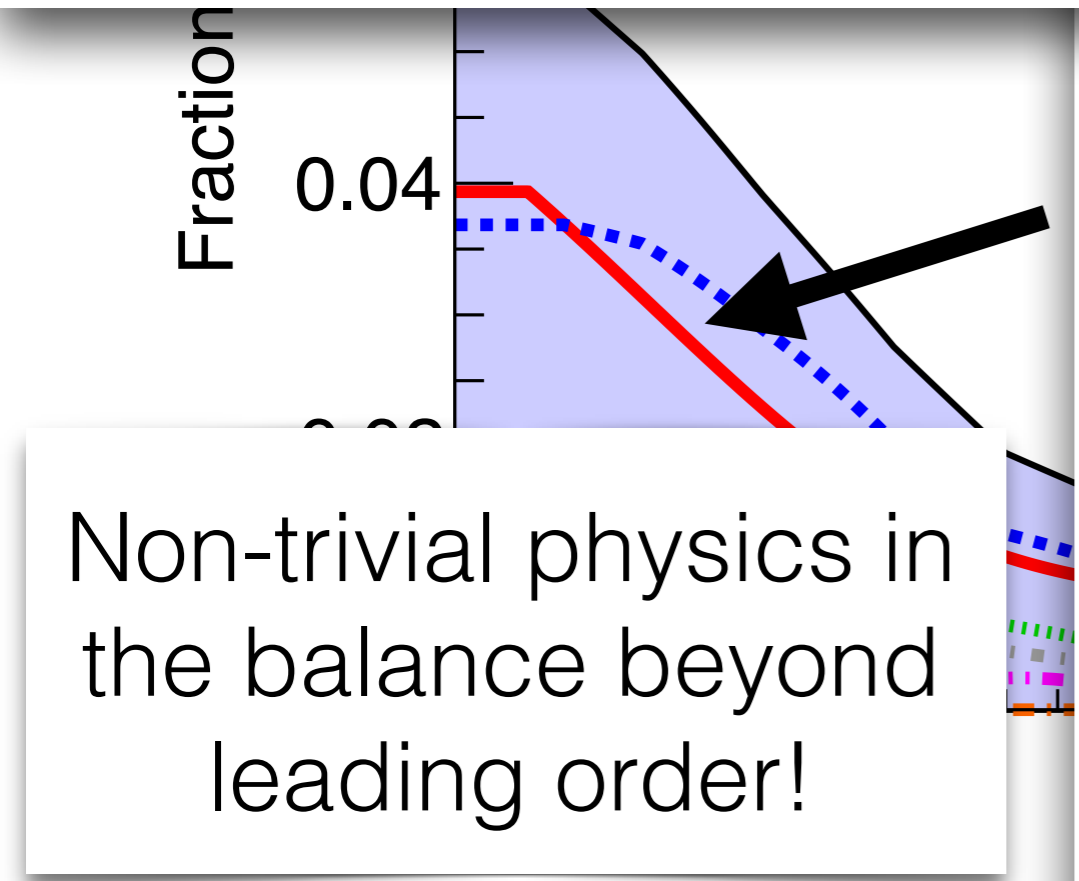
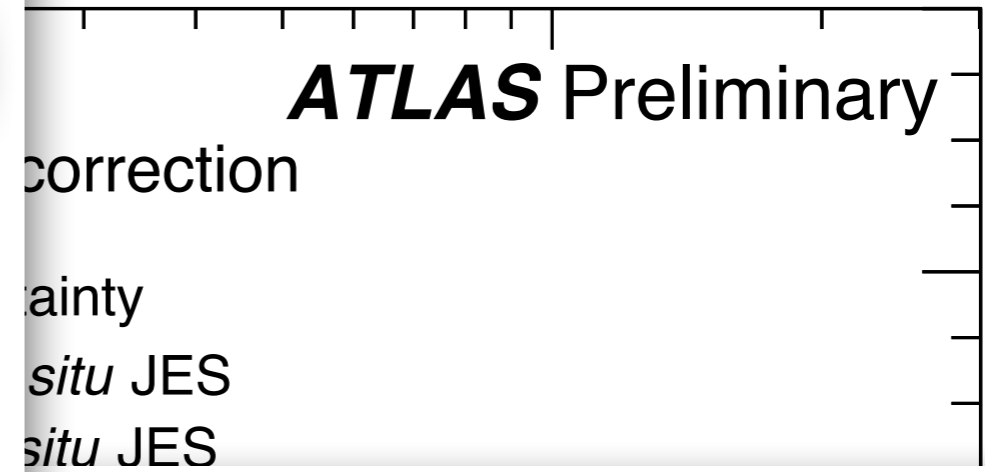
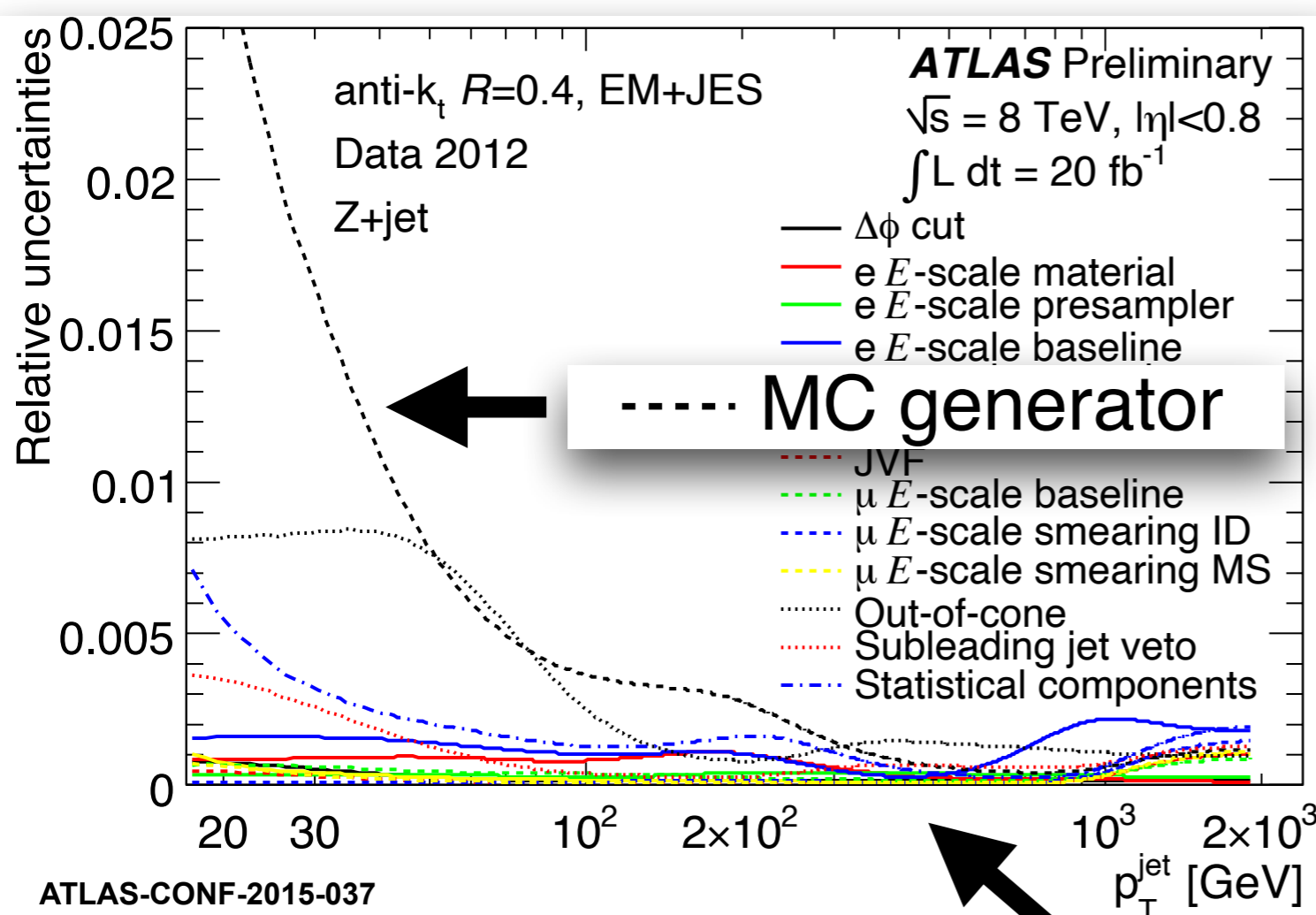


# Jet energy bias uncertainty

28







When someone says they have measured a differential cross-section, they mean that it has been **unfolded**.

Unfolding corrects for detector effects so that our data can be directly compared with theory predictions.

**It is not valid to directly compare theory predictions to detector-level data !!**

Either the theory or the data must be corrected. Correcting the data is more general and can allow for multiple theory groups to reuse the same measurement.

# What does unfolding do?

In general, unfolding needs to correct for interrelated effects:

- Acceptance and efficiency
  - ➔ Particles produced may not be measured
- Detector noise
  - ➔ Particles measured may not be from real particles
- Background processes
  - ➔ If you want to measure process  $X$ , need to remove  $Y$
- Combinatorics
  - ➔ If  $N$  particles, chance that detector can change order
- Detector distortions
  - ➔ Bias and resolution effects



# What does unfolding do?

33

In general, unfolding needs to correct for interrelated effects:

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I'll briefly  
illustrate this



# Illustrative toy example

34

$$m = Rt \quad m = \text{measured}; t = \text{true}$$

We usually call  $R$  the “response matrix” because  $m$  and  $t$  are binned (and thus vectors).

We (usually) get  $R$  from detailed detector simulations.

e.g. Geant4 (particle propagation and energy deposition) + custom code for analog and digital signal emulation.

# Illustrative toy example

35

$$m = Rt \quad m = \text{measured}; t = \text{true}$$

I know what you want to do here is  $t = R^{-1} m$ .

# Illustrative toy example

36

$$m = Rt \quad m = \text{measured}; t = \text{true}$$

I know what you want to do here is  $t = R^{-1} m$ .

*In the next slides, I hope to convince you that this is not usually a good idea.*

# Illustrative toy example

37

$$m = Rt \quad m = \text{measured}; t = \text{true}$$

$$R = \begin{pmatrix} 1 - \epsilon & \epsilon \\ \epsilon & 1 - \epsilon \end{pmatrix}$$

Consider this case, where  $0 \leq \epsilon \leq 0.5$

# Illustrative toy example

38

$$m = Rt \quad m = \text{measured}; t = \text{true}$$

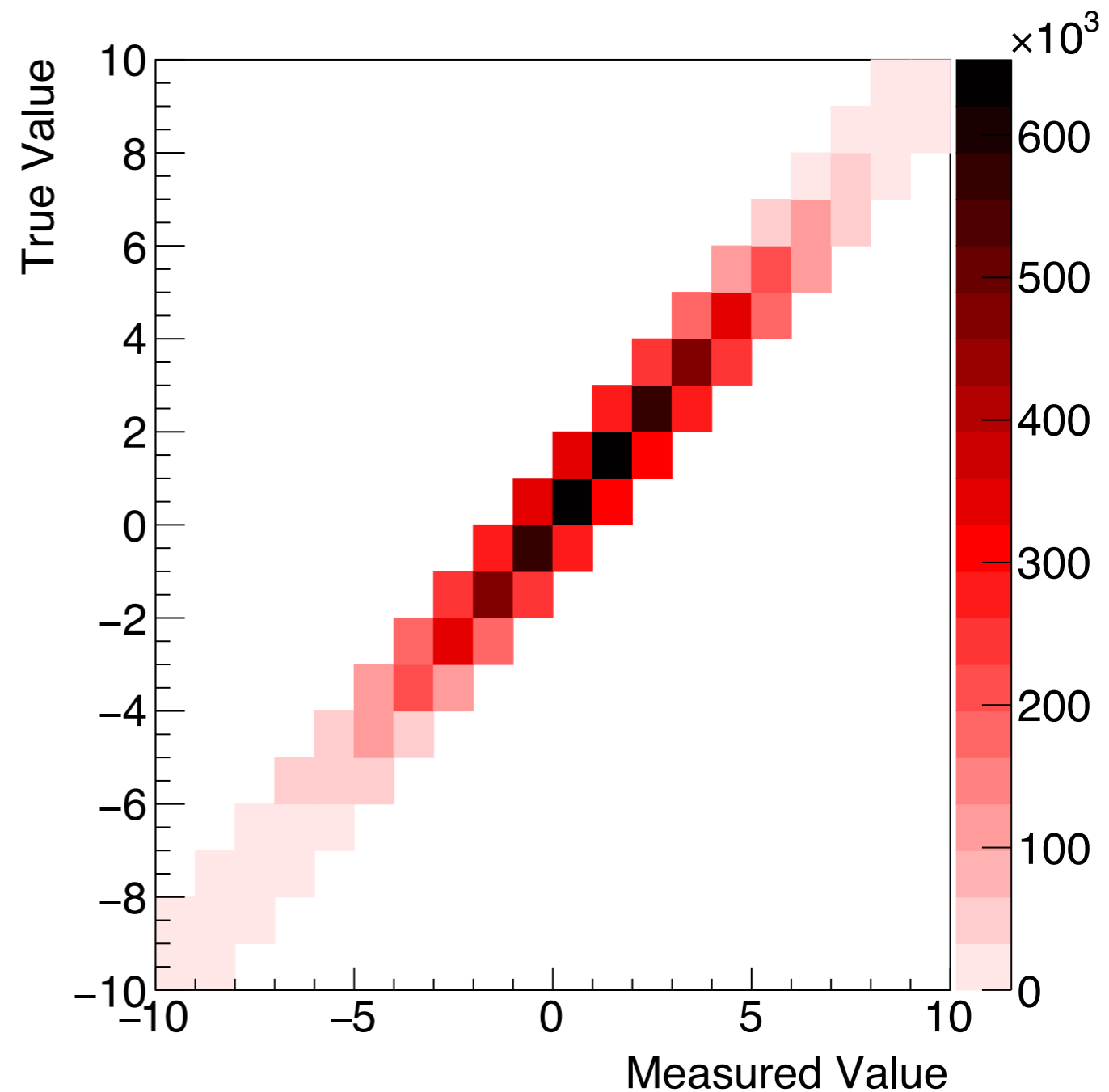
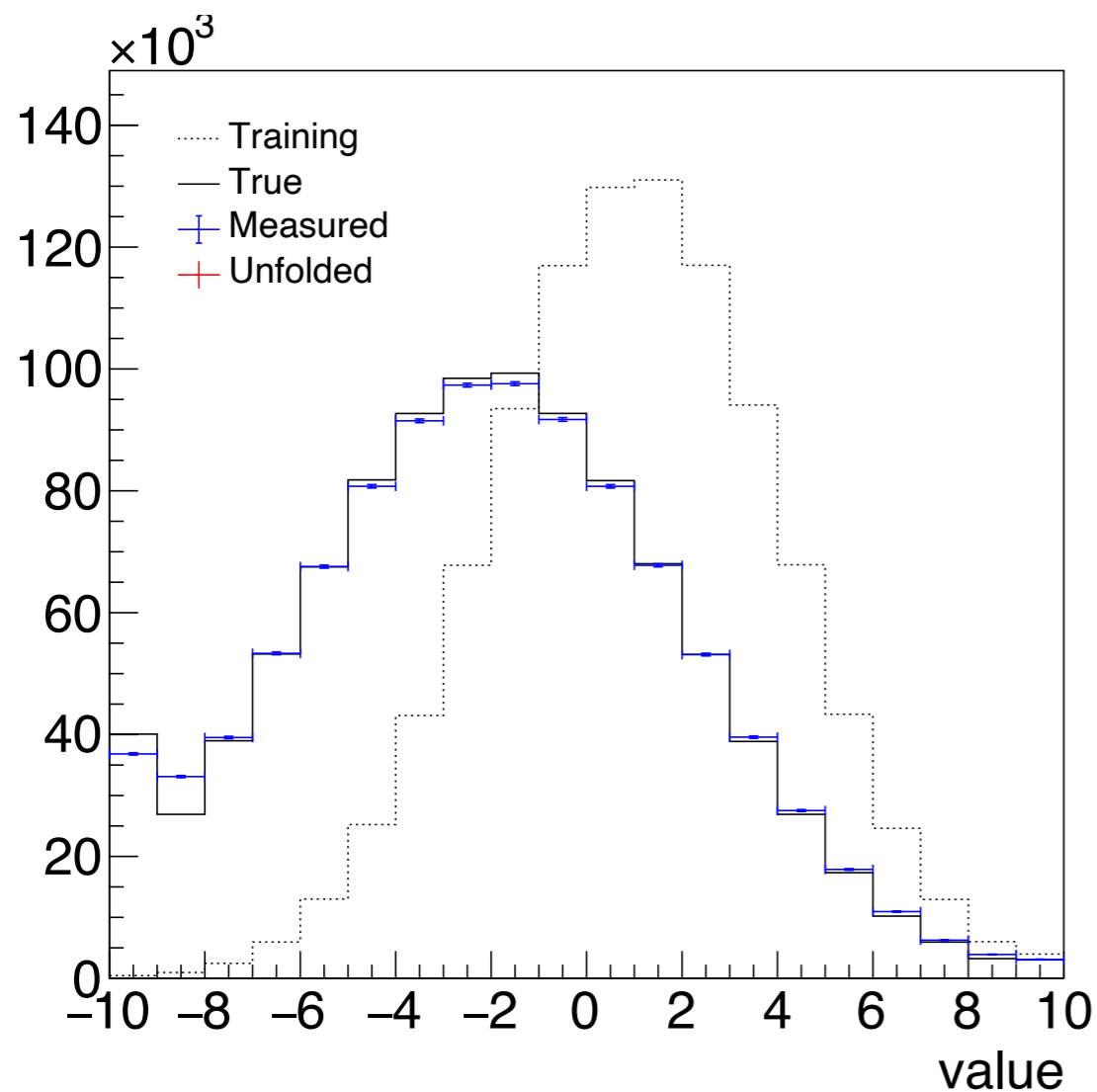
$$R = \begin{pmatrix} 1 - \epsilon & \epsilon \\ \epsilon & 1 - \epsilon \end{pmatrix}$$

$$\text{Var}(R^{-1}m) \propto 1/\text{Det}(R) = 1 - 2\epsilon$$

Statistical uncertainty blows up as  $\epsilon \rightarrow 0.5$

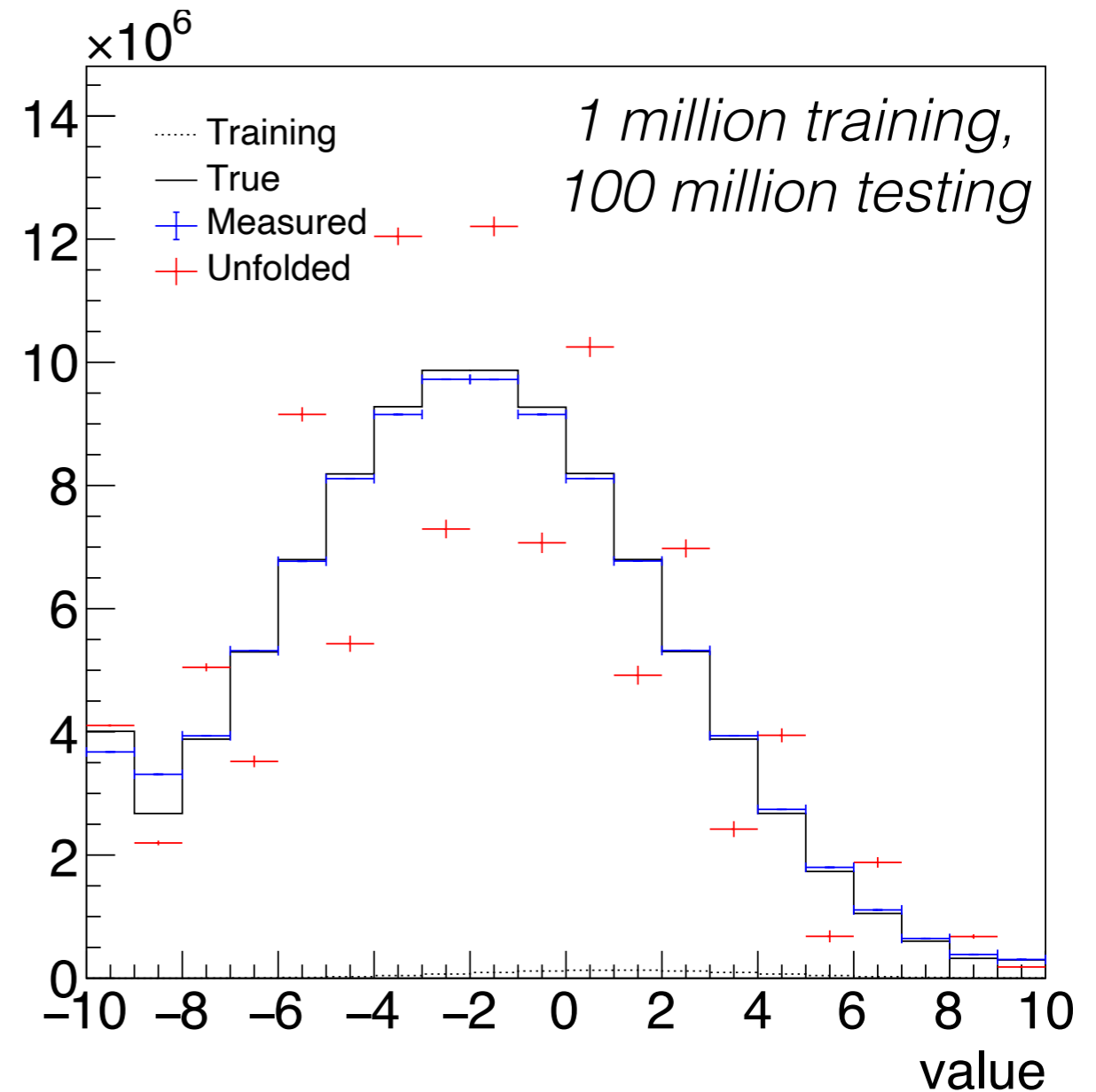
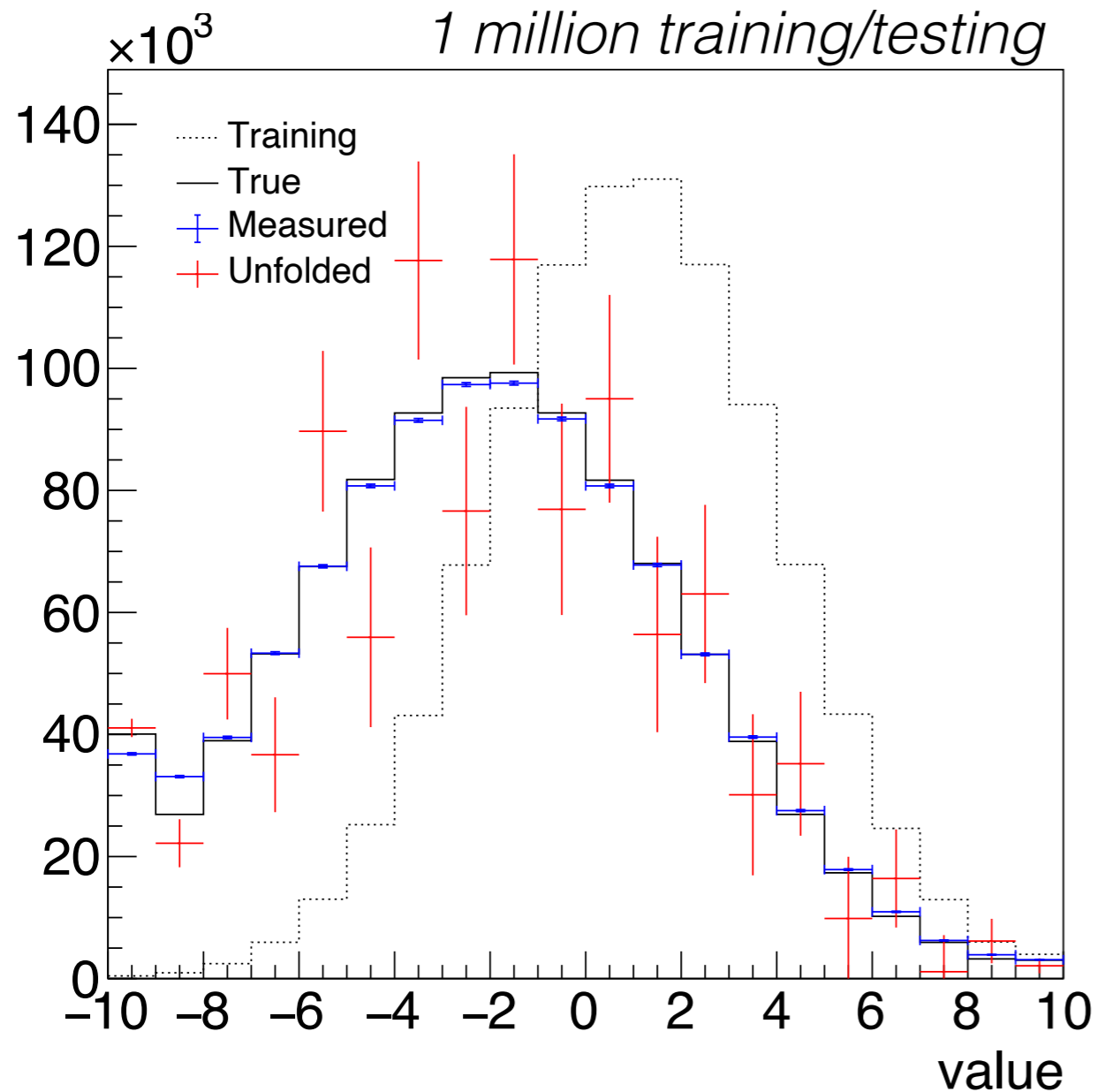
# A more realistic example

$$R_{\text{norm}} = \begin{pmatrix} 0.75 & 0.25 & 0 & & & & \\ 0.25 & 0.50 & 0.25 & 0 & & & \\ 0 & 0.25 & 0.50 & 0.25 & 0 & & \\ & 0 & 0.25 & 0.50 & 0.25 & & \\ & & 0 & 0.25 & 0.50 & & \\ & & & & & \ddots & \end{pmatrix}$$



# Unfolding by Matrix Inversion

40



Stat. uncertainty is large and there is a bias when training dataset is too small.



# The HEP solution

41

Our solution is to do regularized matrix inversion.

There are two main techniques that we use:

“Iterative Bayesian Unfolding”

$$\theta_{ij} = \frac{\Pr(m_j|t_i) \cdot \Pr(t_i)}{\sum_i \Pr(m_j|t_i) \cdot \Pr(t_i)}$$

response  
matrix

regularization  
= number of iterations

$$\Pr_{k+1}(t_i) = \sum_j \theta_{ij} \Pr_k(t_i)$$

Nucl. Inst. Meth. A 362 (1995) 487

“Singular Value  
Decomposition (SVD) Unfolding”

$$R = USV^T$$

$U, V$ , orthogonal,  $S$  diagonal & non-negative

$$d = U^T m \quad z_i(\tau) = \frac{d_i}{s_i} \cdot \frac{s_i^2}{s_i^2 + \tau}$$

$$t = Vz$$

regularization  
parameter

Nucl. Inst. Meth. A 372 (1995) 469

Main tool: RooUnfold (ROOT-based C++ code)

Our solution is to do regularized matrix inversion.

Note: regularized matrix inversion depends on unphysical regularization parameters

One chooses parameters to tradeoff bias and uncertainty.

$$\theta_{ij} = \frac{\Pr(m_j|t_i) \cdot \Pr(t_i)}{\sum_k \Pr(m_j|t_k) \cdot \Pr(t_k)}$$

**IBU Unfolding**

- depend on prior

- depends on # of iterations

Nucl. Inst. Meth. A 362 (1995) 487

$$B = U S V^T$$

**SVD Unfolding**

$$d = U^T t \cdot \frac{s_i^2}{s_i^2 + \tau}$$

$$t = V z$$

regularization parameter

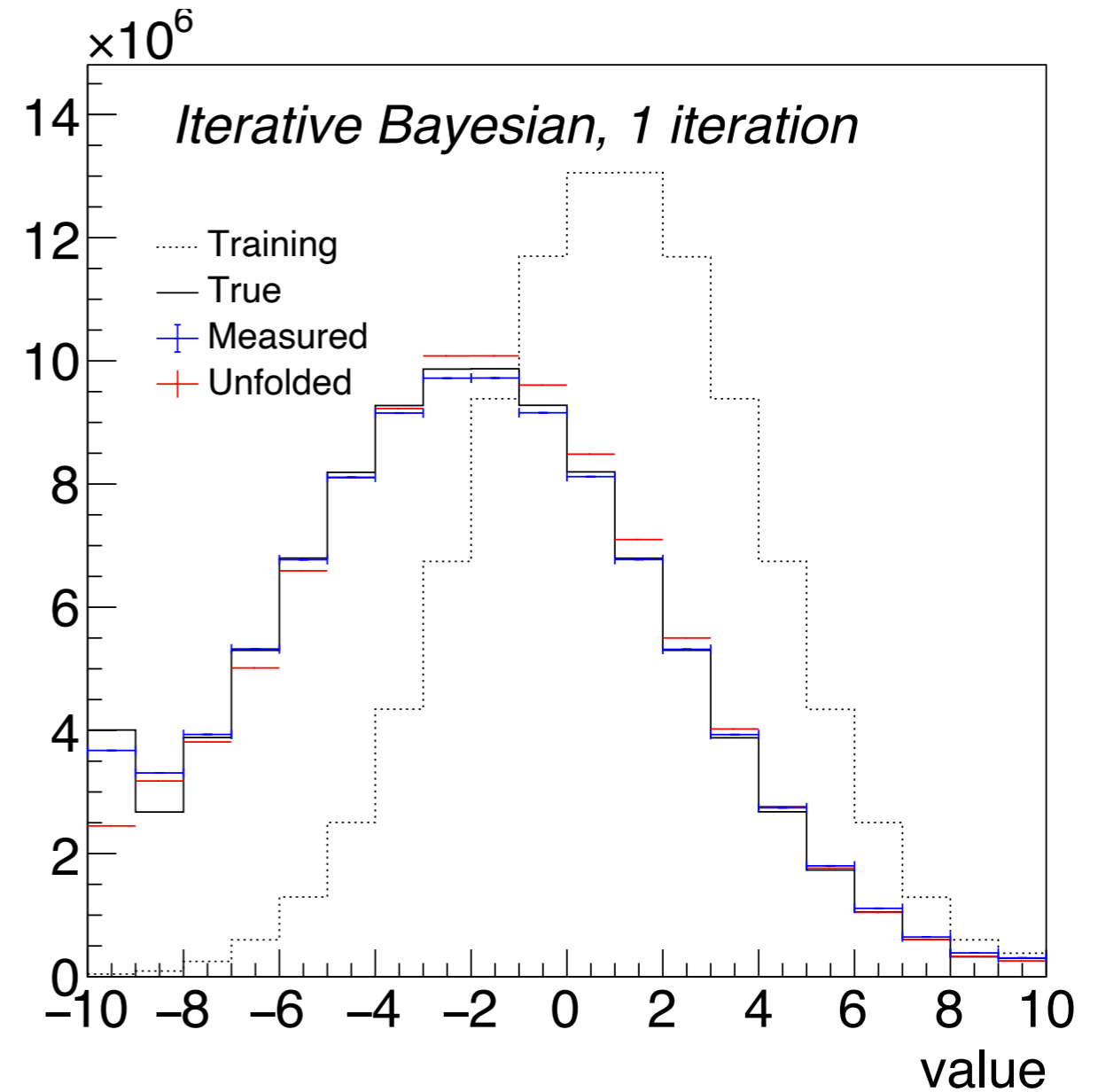
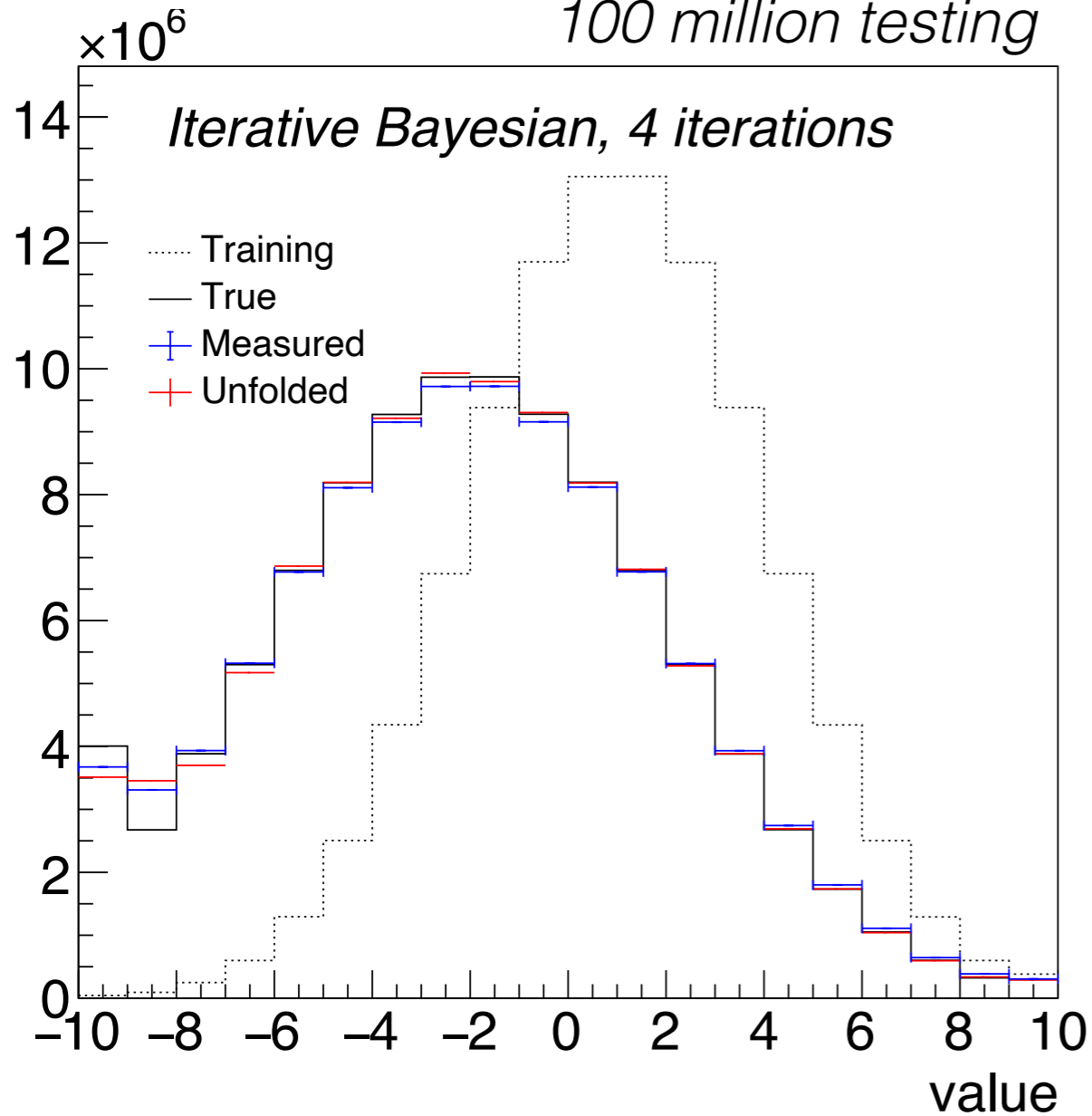
Nucl. Inst. Meth. A 372 (1995) 469

Main tool: RooUnfold (ROOT-based C++ code)

# Example: Iterative Bayesian Unfolding

43

1 million training,  
100 million testing



# One last comment: phase space def.

44

Need to decide what to unfold to, called the “fiducial volume”

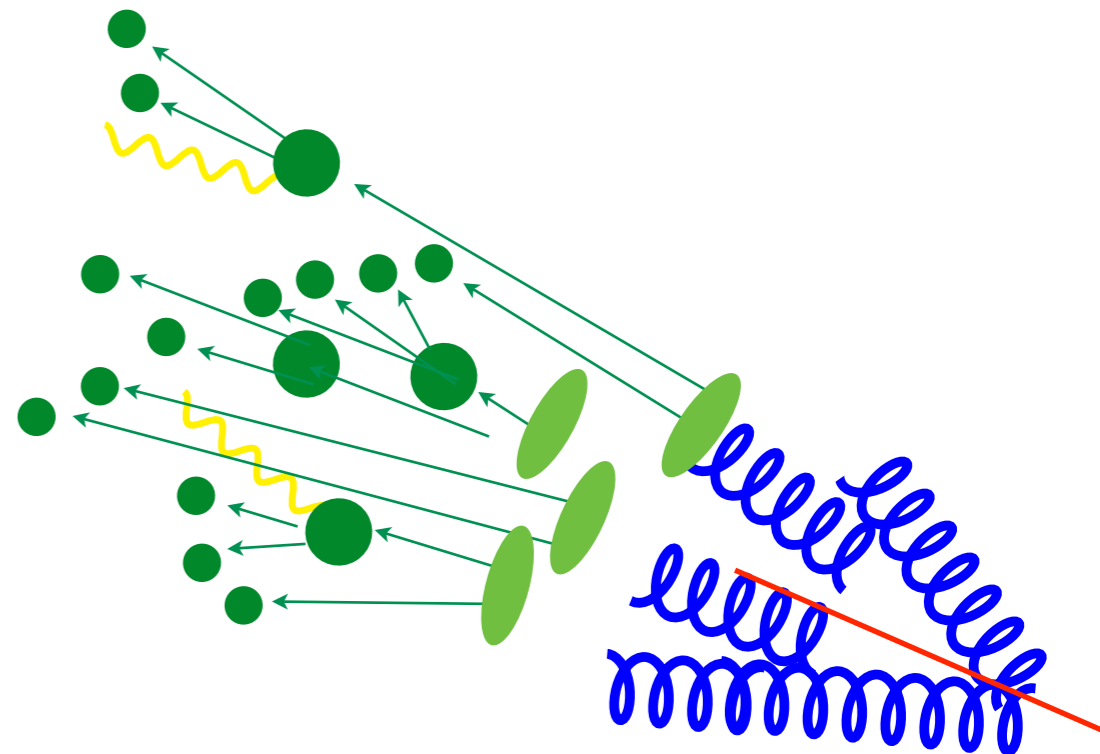
Calculations are often at the level of “born-level partons”

Measurements are at the level of hadrons.

+Non-perturbative corrections

+Resummation (“dressed”)

Born-level parton



# One last comment: phase space def.

45

Need to decide what to unfold to, called the “fiducial volume”

Calculations are often at the level of “born-level partons”

Measurements are at the level of hadrons.

The closer the target is to the observable, the less the unfolding has to do and the smaller the modeling uncertainty.

**Please don't unfold to “born-level partons”**. Better to “dress” the calculations with resummation and hadronization.

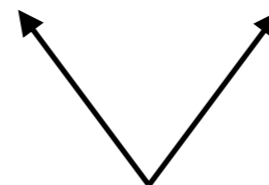
(fortunately, born measurements is becoming less fashionable these days)

# Physics topics

46



$$\sigma = \sum_{ij \in \{q, \bar{q}, g\}} \int dx_1 dx_2 f_i(x_1, \mu) f_j(x_2, \mu) \hat{\sigma}_{ij}(x_1 E, x_2 E)$$



PDFs

Every\* LHC measurement is sensitive to PDFs - the ones that are most useful for constraining them can be (1) measured precisely, (2) predicted precisely (see last term), and (3) mostly sensitive to ~one partonic channel

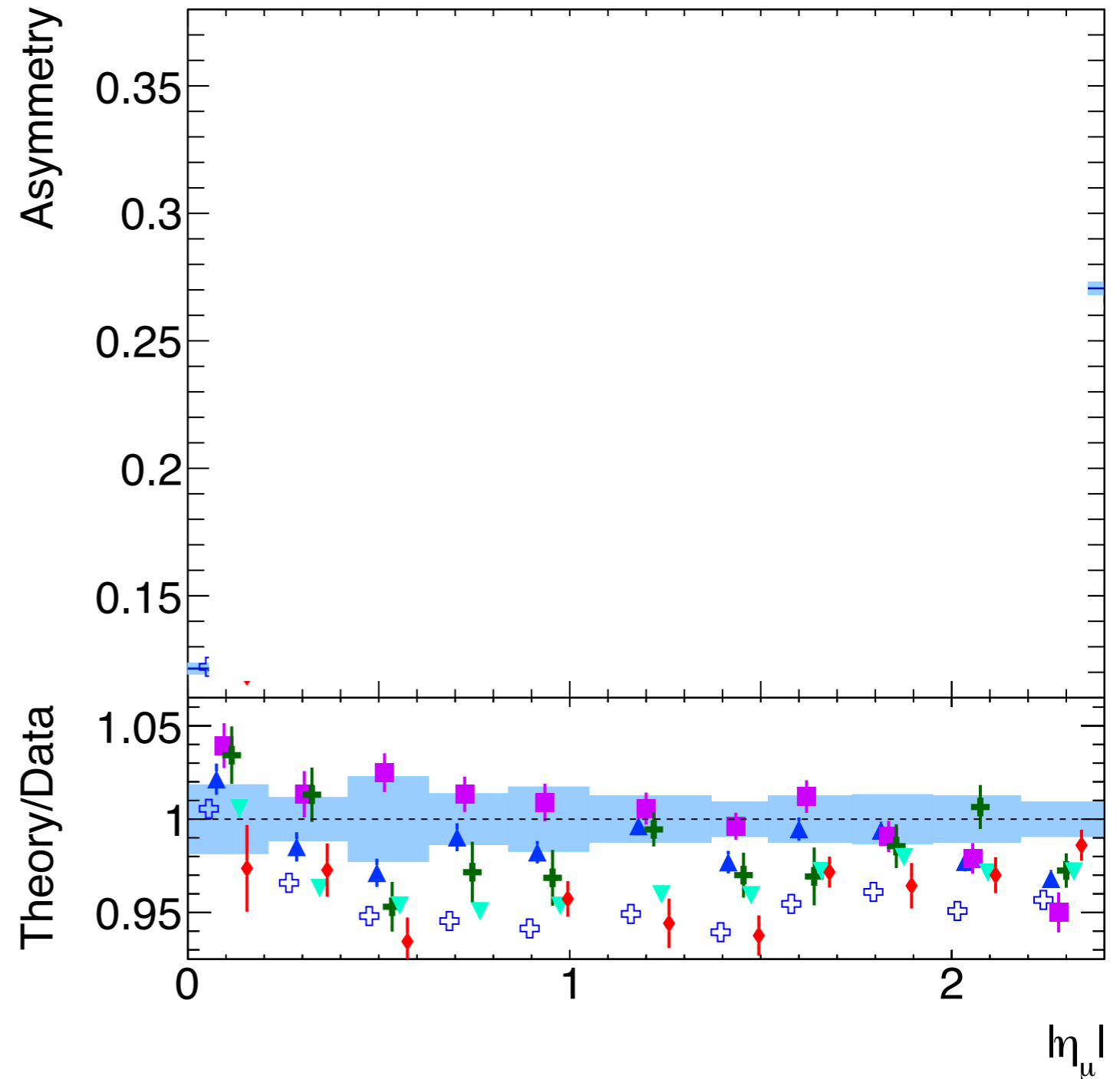
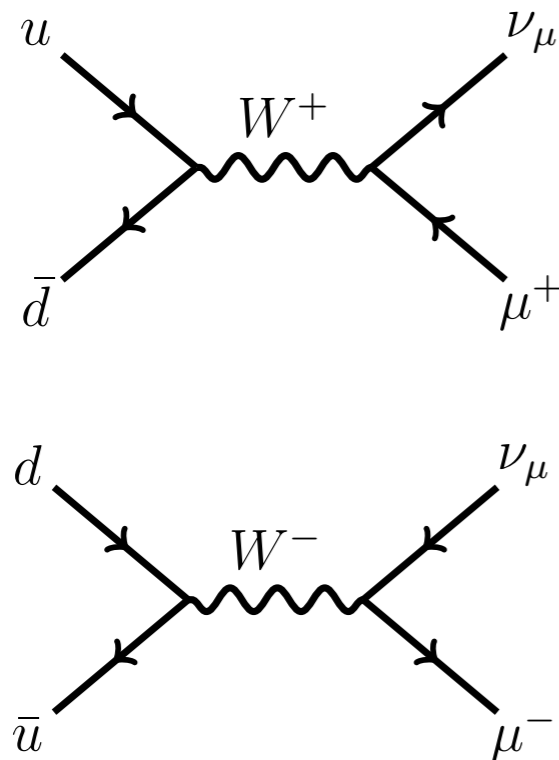
\*Except when the protons interact coherently.

# Valence quarks at moderate $x$

48

$W$ +jets is a precise probe of the  $u/d$  content of the proton.

At a  $pp$  collide, more  $W^+$  than  $W^-$





# Valence quarks at moderate $x$

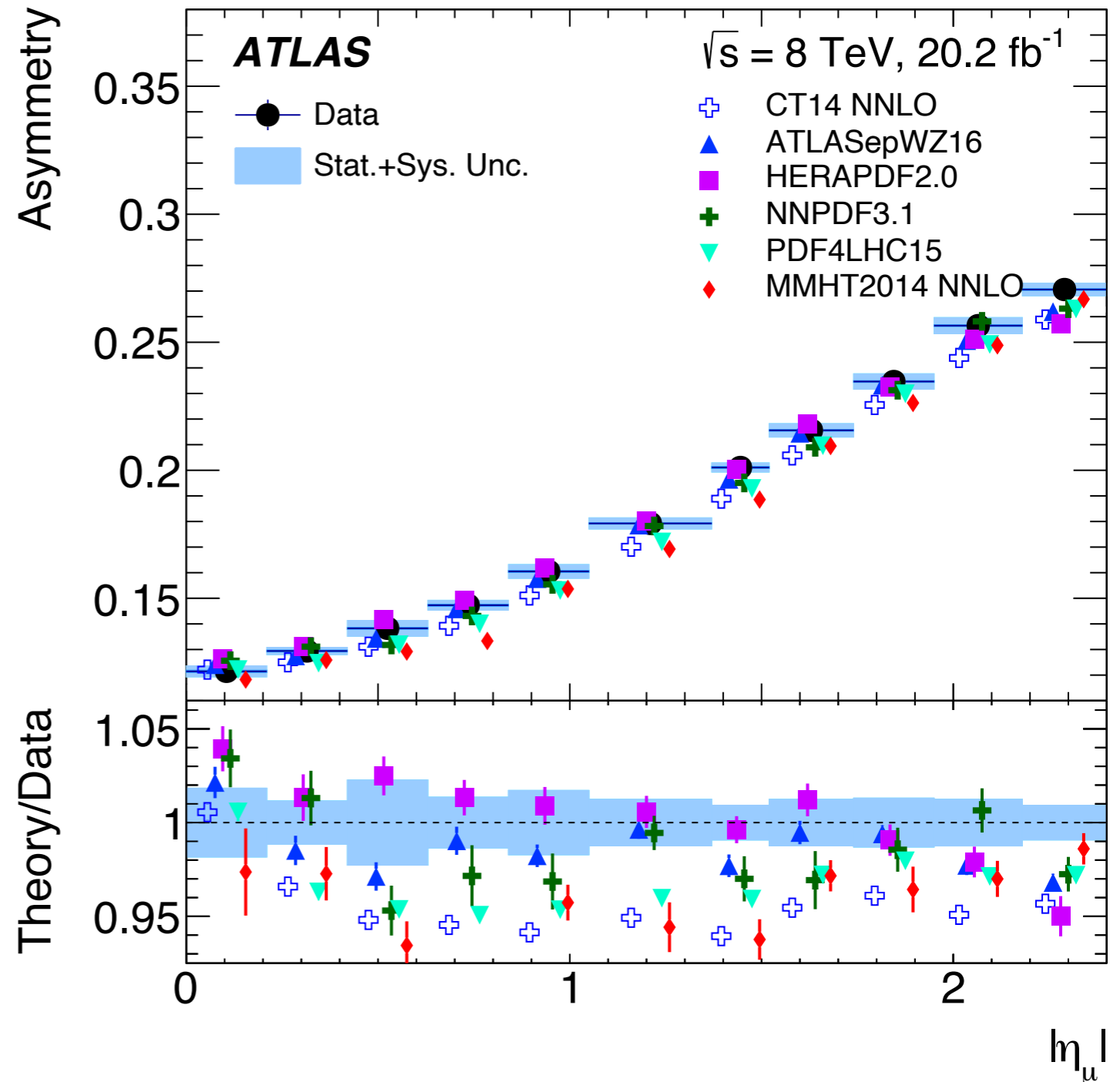
49

$W$ +jets is a precise probe of the  $u/d$  content of the proton.

At a  $pp$  collide, more  $W^+$  than  $W^-$

$$A_\mu = \frac{d\sigma_{W_{\mu^+}}/d\eta_\mu - d\sigma_{W_{\mu^-}}/d\eta_\mu}{d\sigma_{W_{\mu^+}}/d\eta_\mu + d\sigma_{W_{\mu^-}}/d\eta_\mu}$$

Why not 1/3?



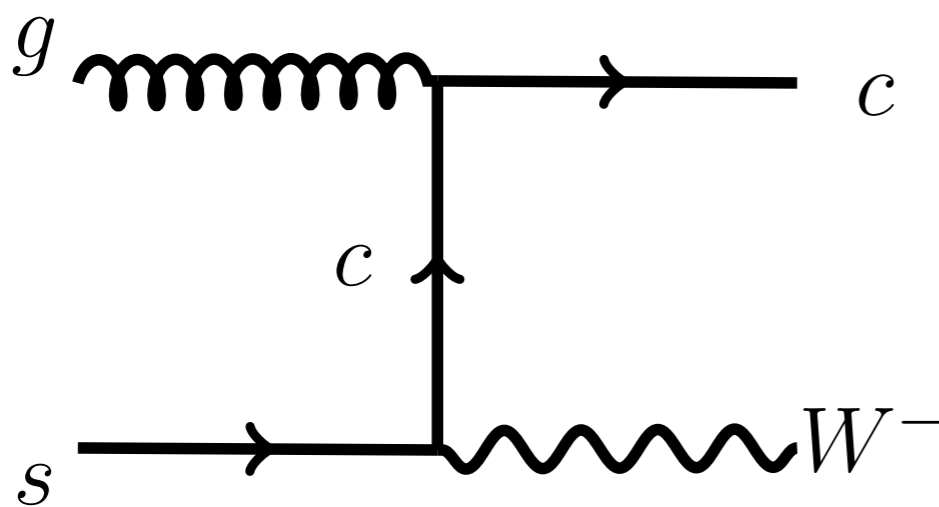
# Heavier quarks

50

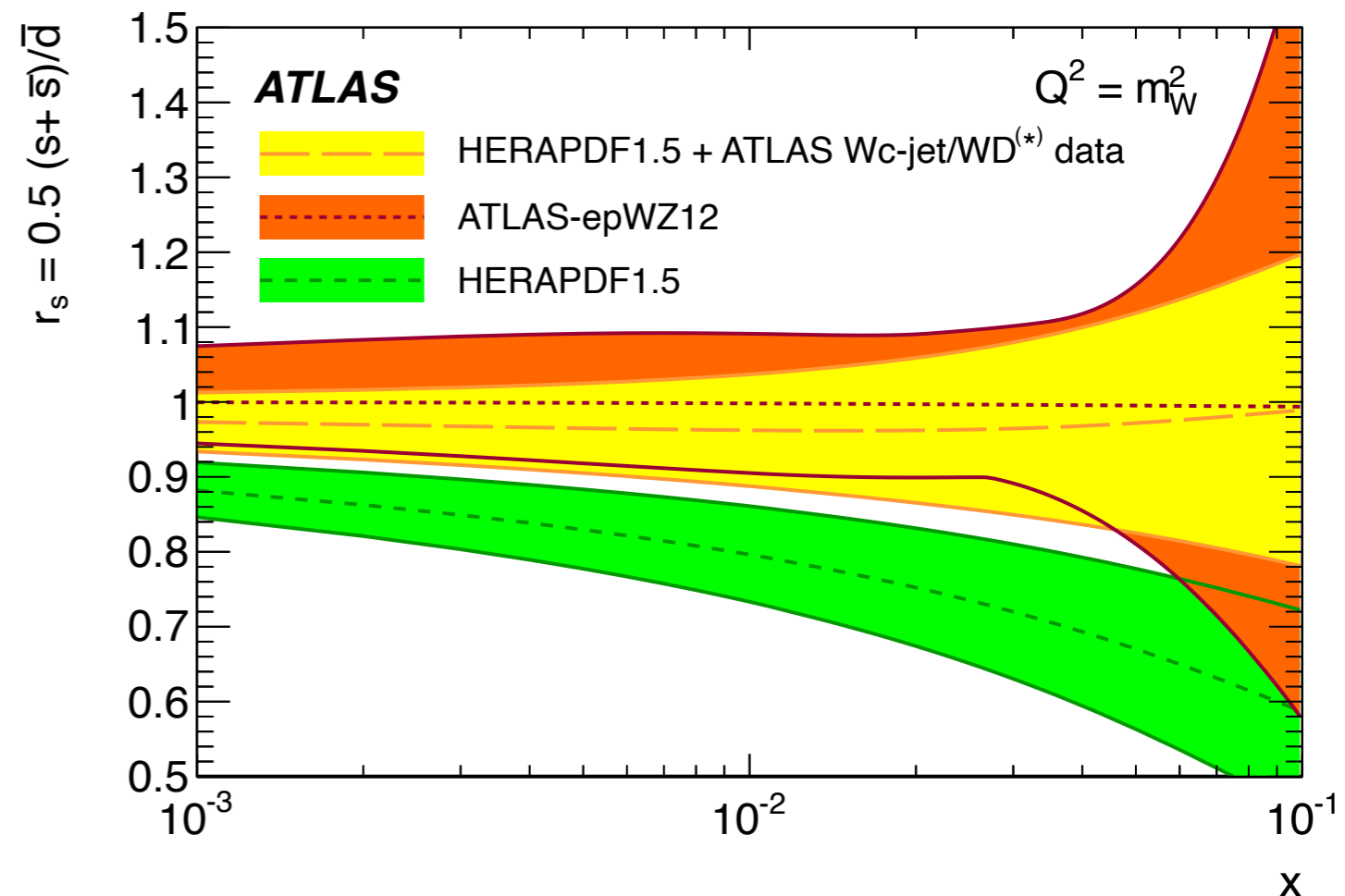
Various measurements are constructed to be sensitive to the s-, c-, and even b-component of the proton.

For example,  $W^+$ -charm measurement can constrain the sea strangeness.

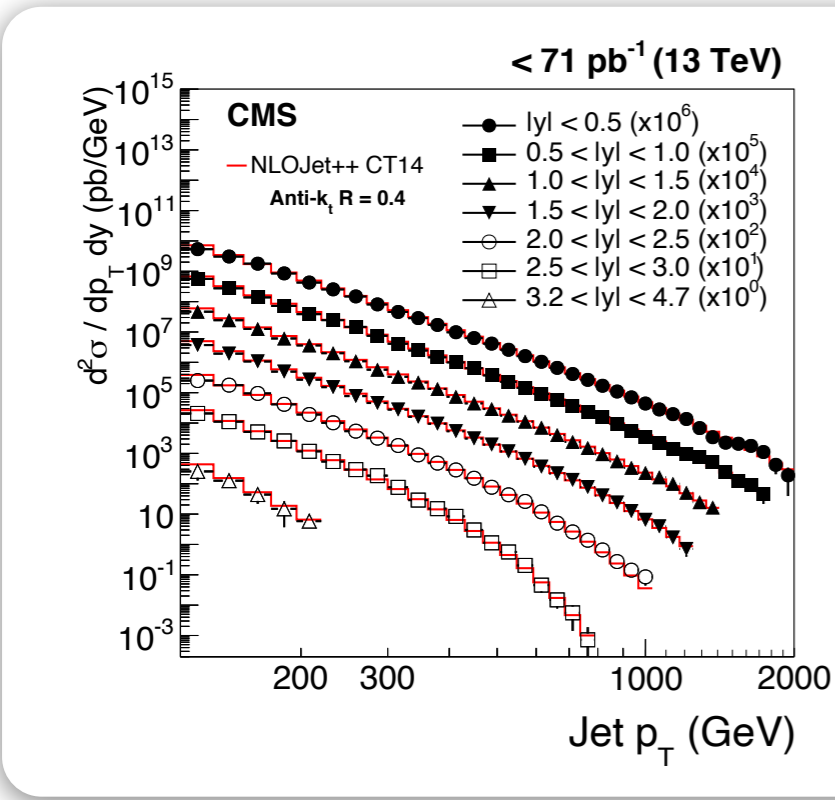
(ATLAS and to a lesser extent CMS find  $s \sim$  other sea contributions)



Question: how would you identify the  $c$ -quark?



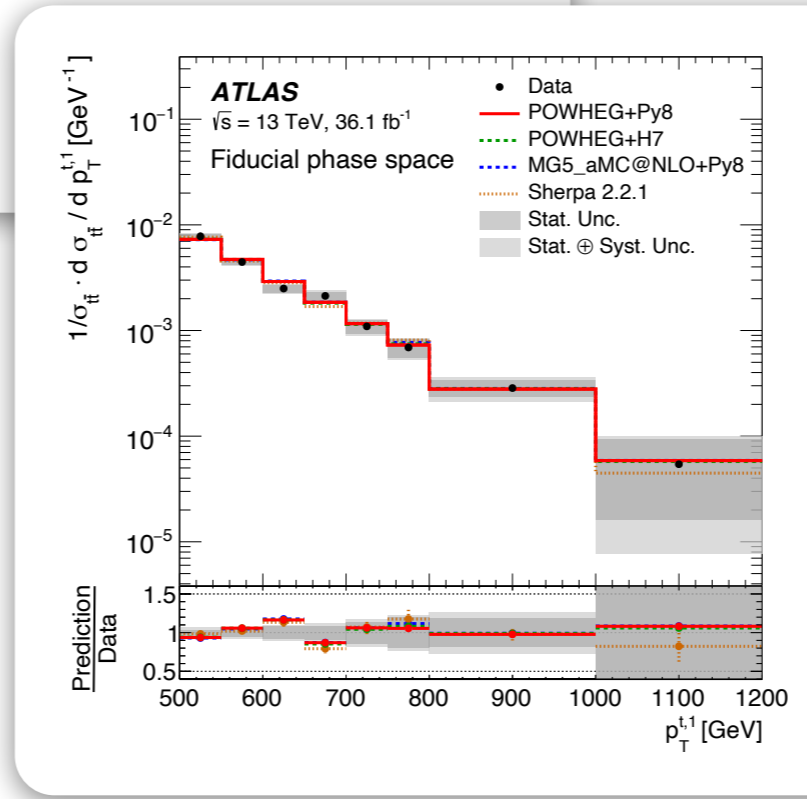
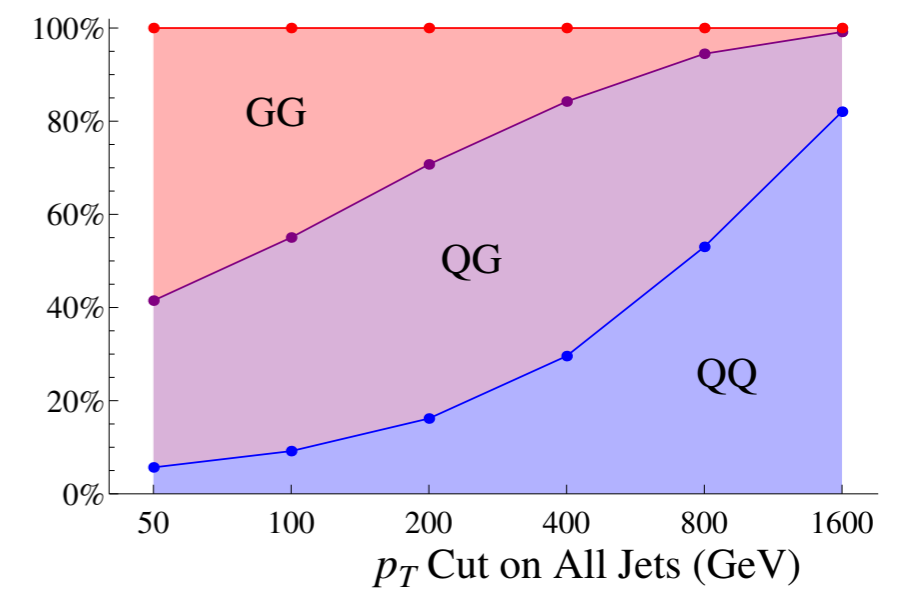
Gluon PDF at moderate - high  $x$  is constrained by two sources:



## Dijets

High stats,  $\sim 10\%$  precision across a wide kinematic range

small gg/qg fraction at high jet  $p_T$



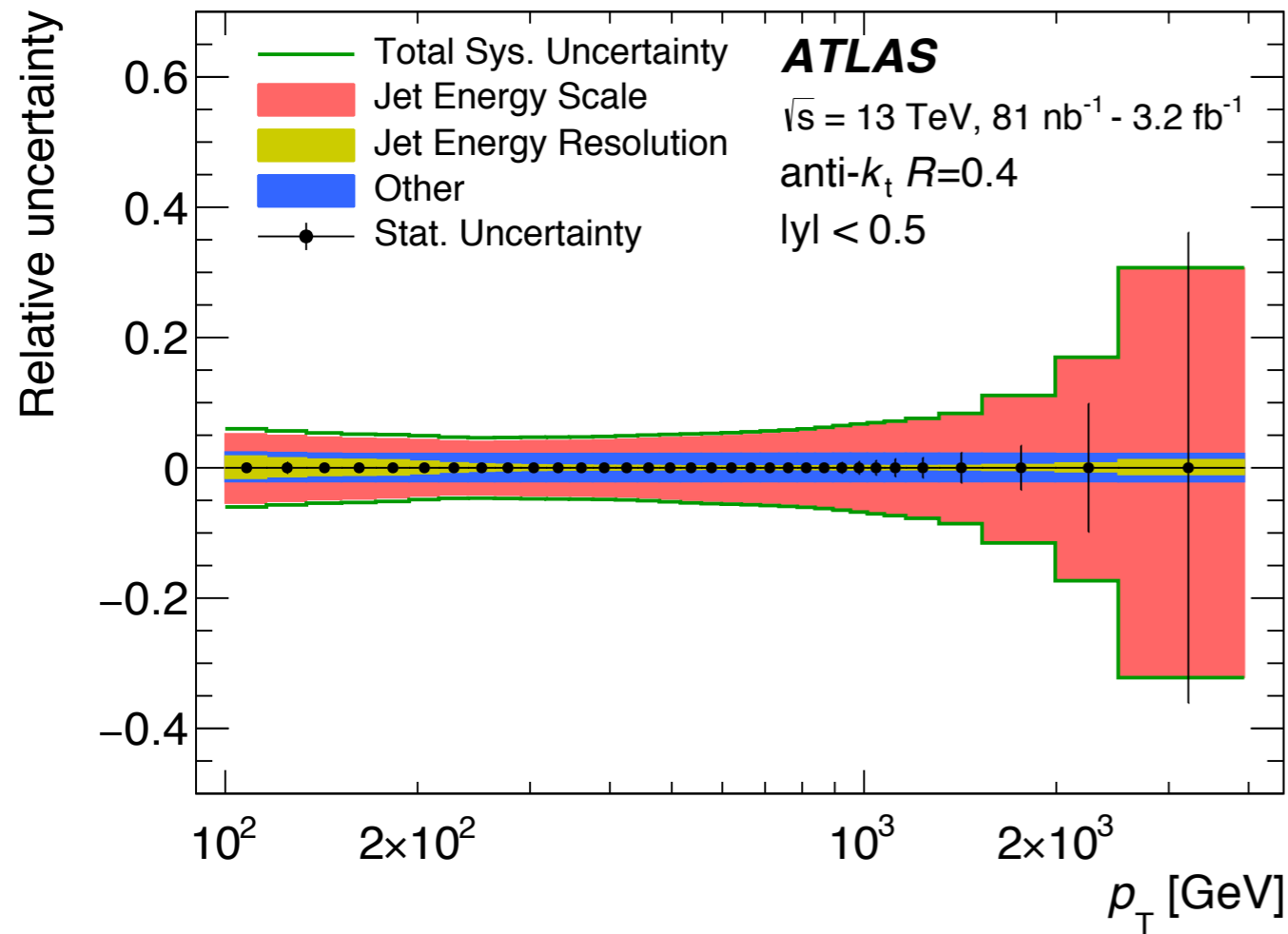
## Top quark pairs

Dominated by gg initial state

run out of statistics much earlier than inclusive jets

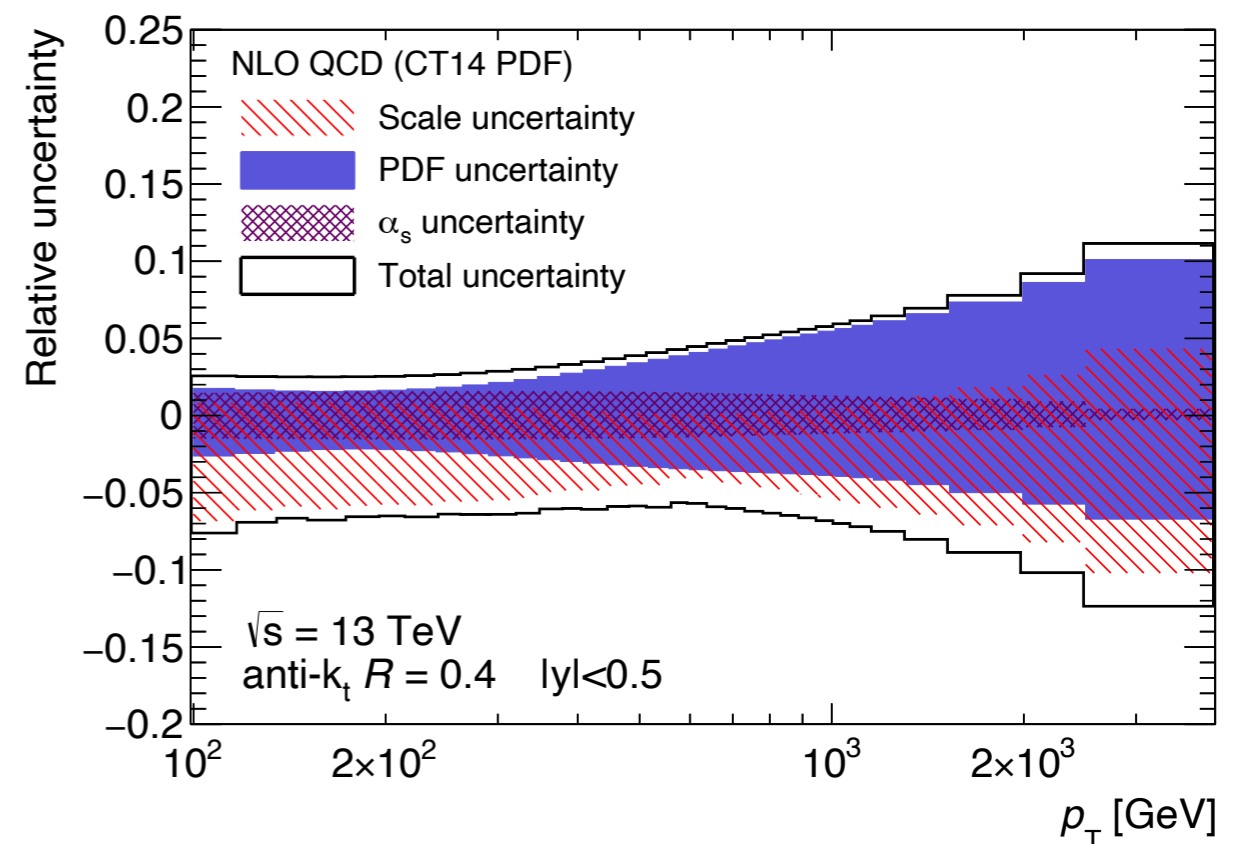
# Dijets (gluons at low $p_T$ , quarks at high $p_T$ )

52

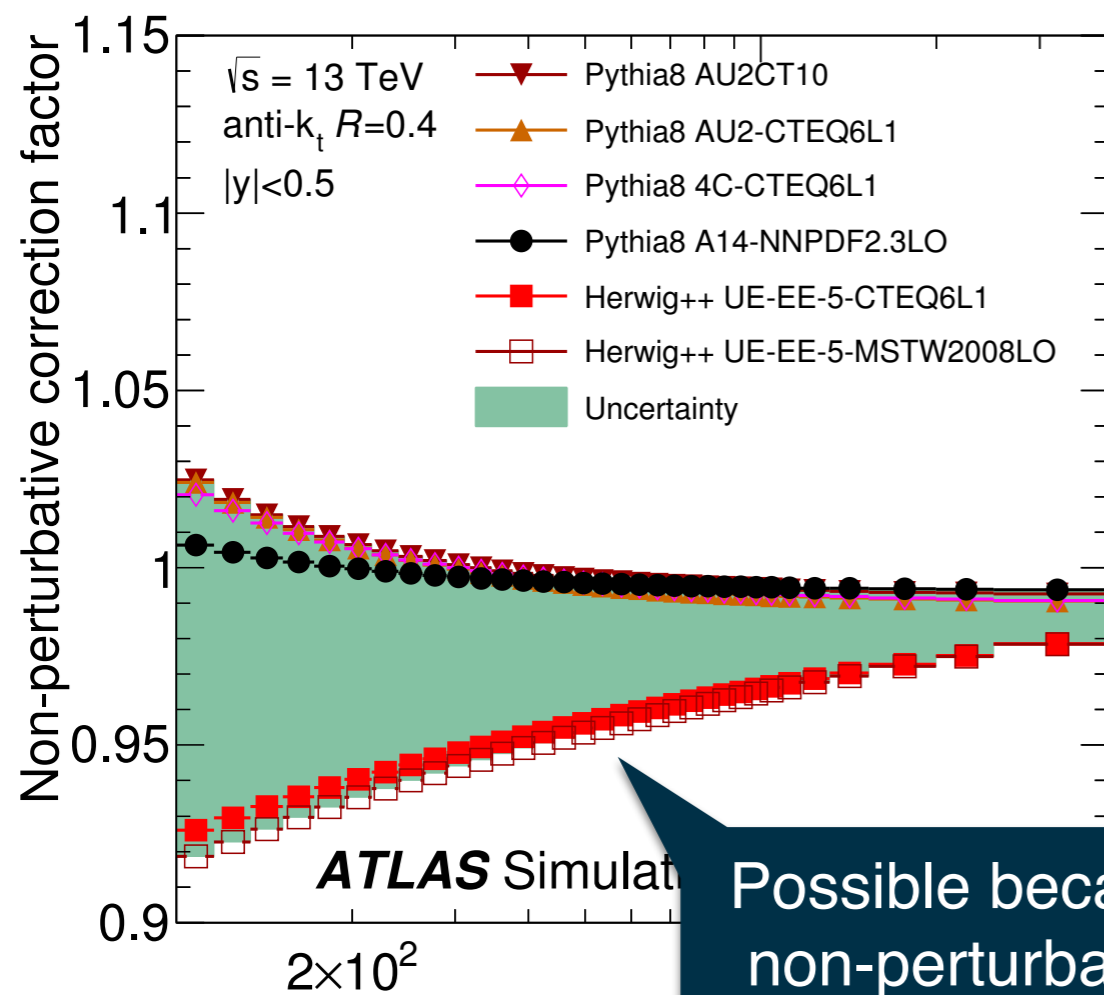


The PDF uncertainty is comparable to the measurement uncertainty at  $\sim O(100)$  GeV  
→ constraining power!

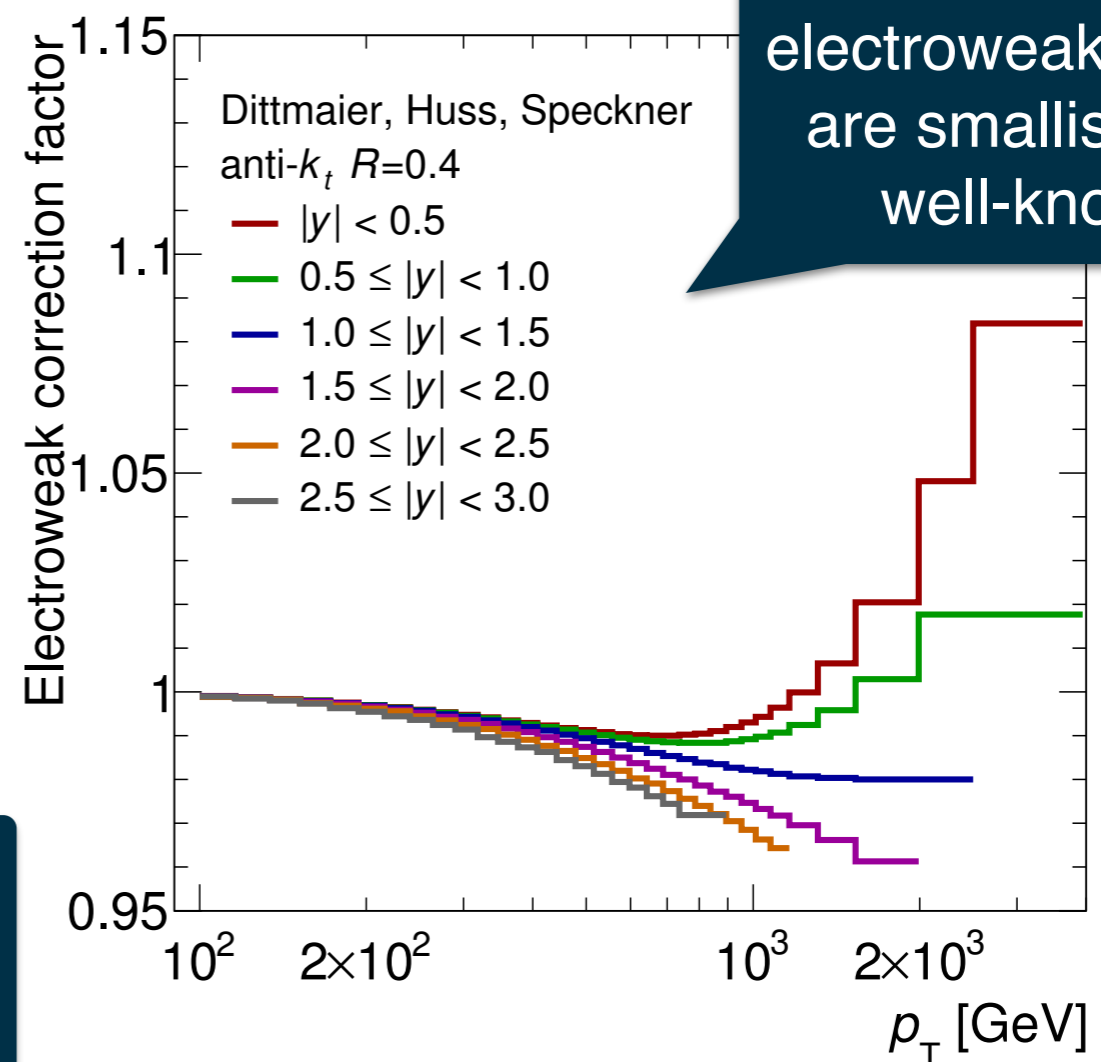
*If the energy uncertainty is 1%, why is the cross-section uncertainty much larger?*



Precise measurements and predictions can do more than provide input to PDFs - they also provide a powerful consistency test of QCD in new energy regimes.



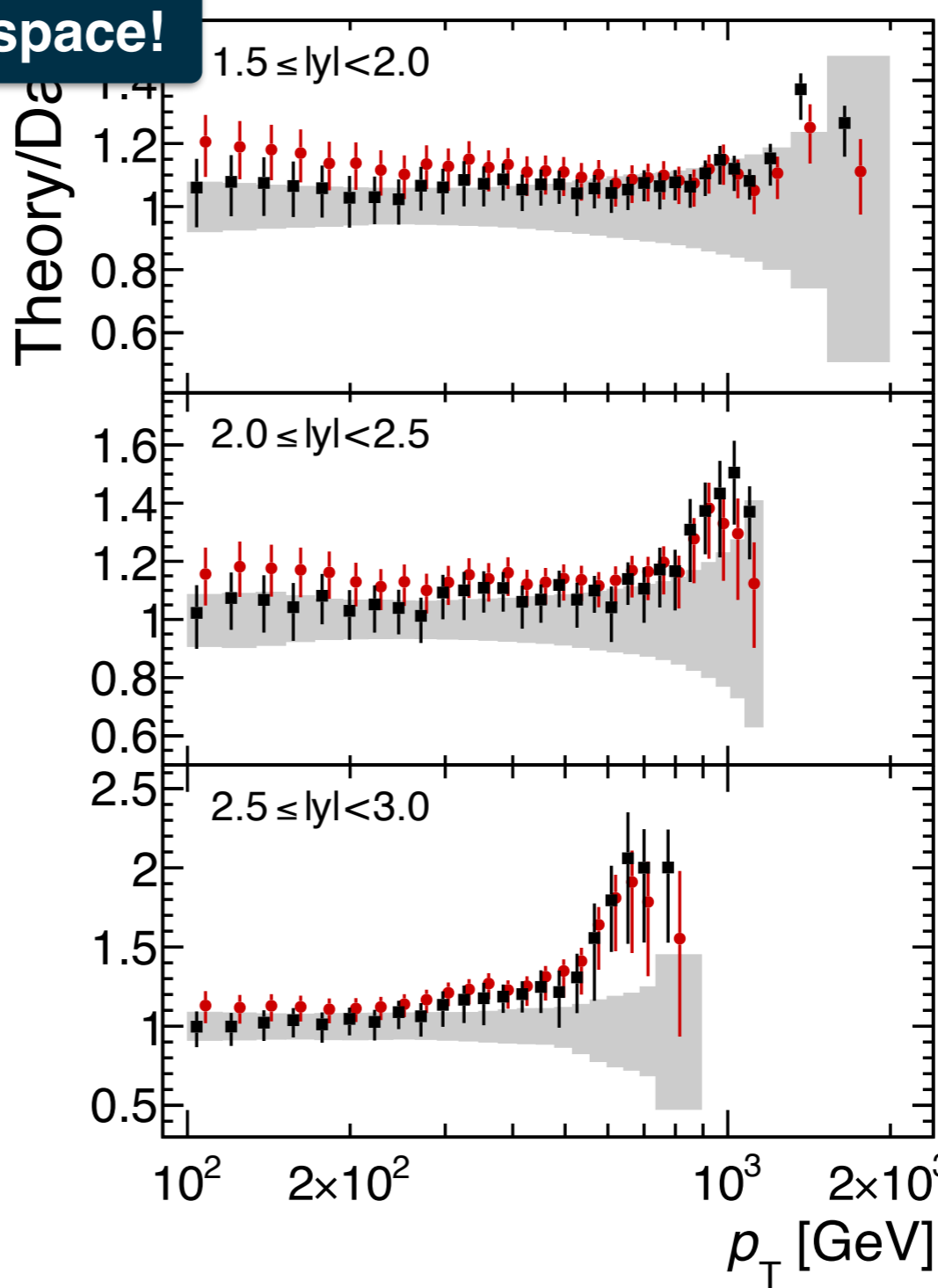
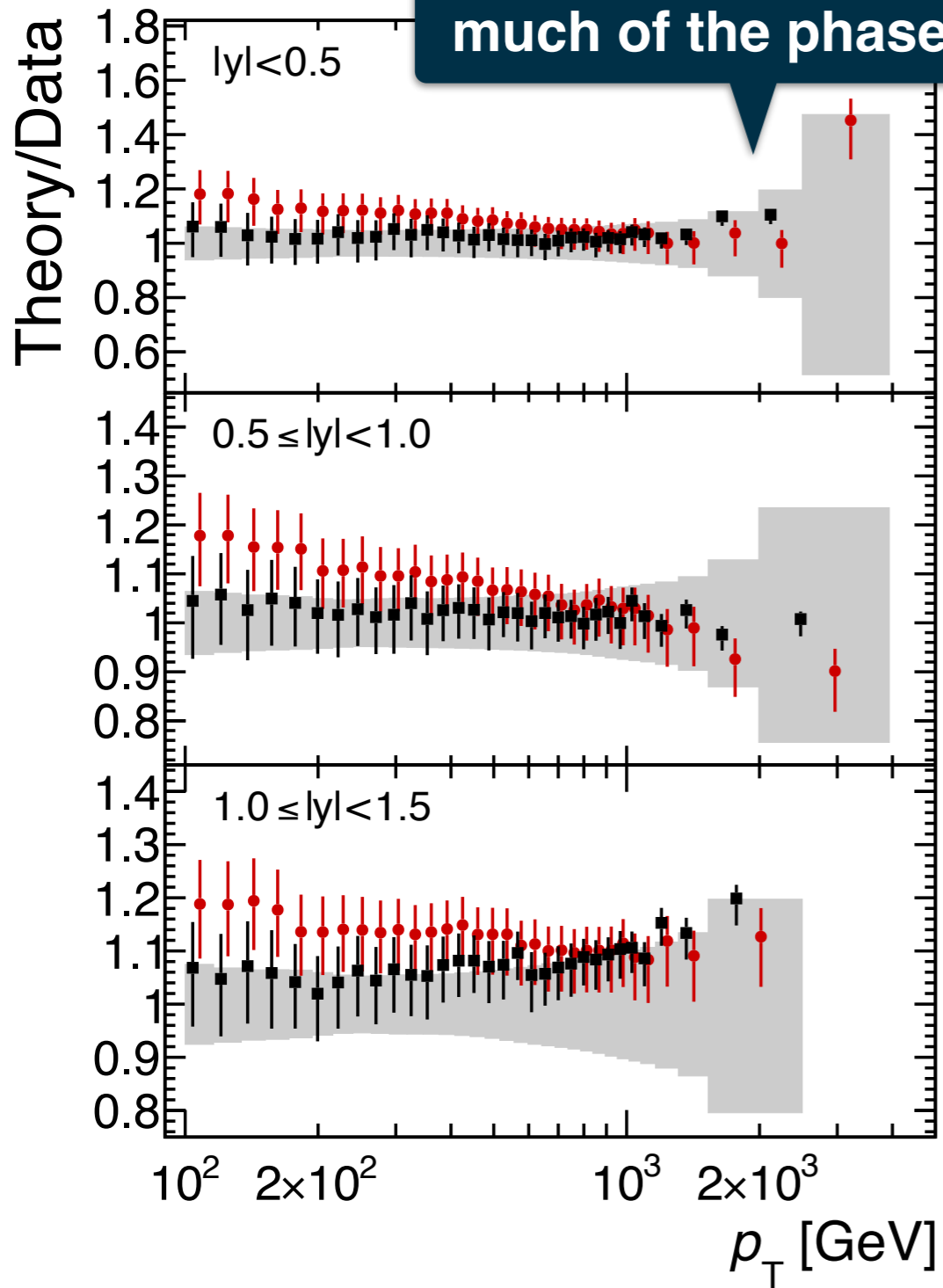
Possible because non-perturbative effects are small



Possible because electroweak effects are smallish and well-known

# Fixed-order pQCD

Impressive agreement over much of the phase space!



**ATLAS**  
 $L = 81 \text{ nb}^{-1} - 3.2 \text{ fb}^{-1}$   
 $\sqrt{s} = 13 \text{ TeV}$   
 anti- $k_t$   $R=0.4$

■ Data

□ NLO QCD

⊗  $k_{EW}$  ⊗  $k_{NP}$

□ NNLO QCD

⊗  $k_{EW}$  ⊗  $k_{NP}$

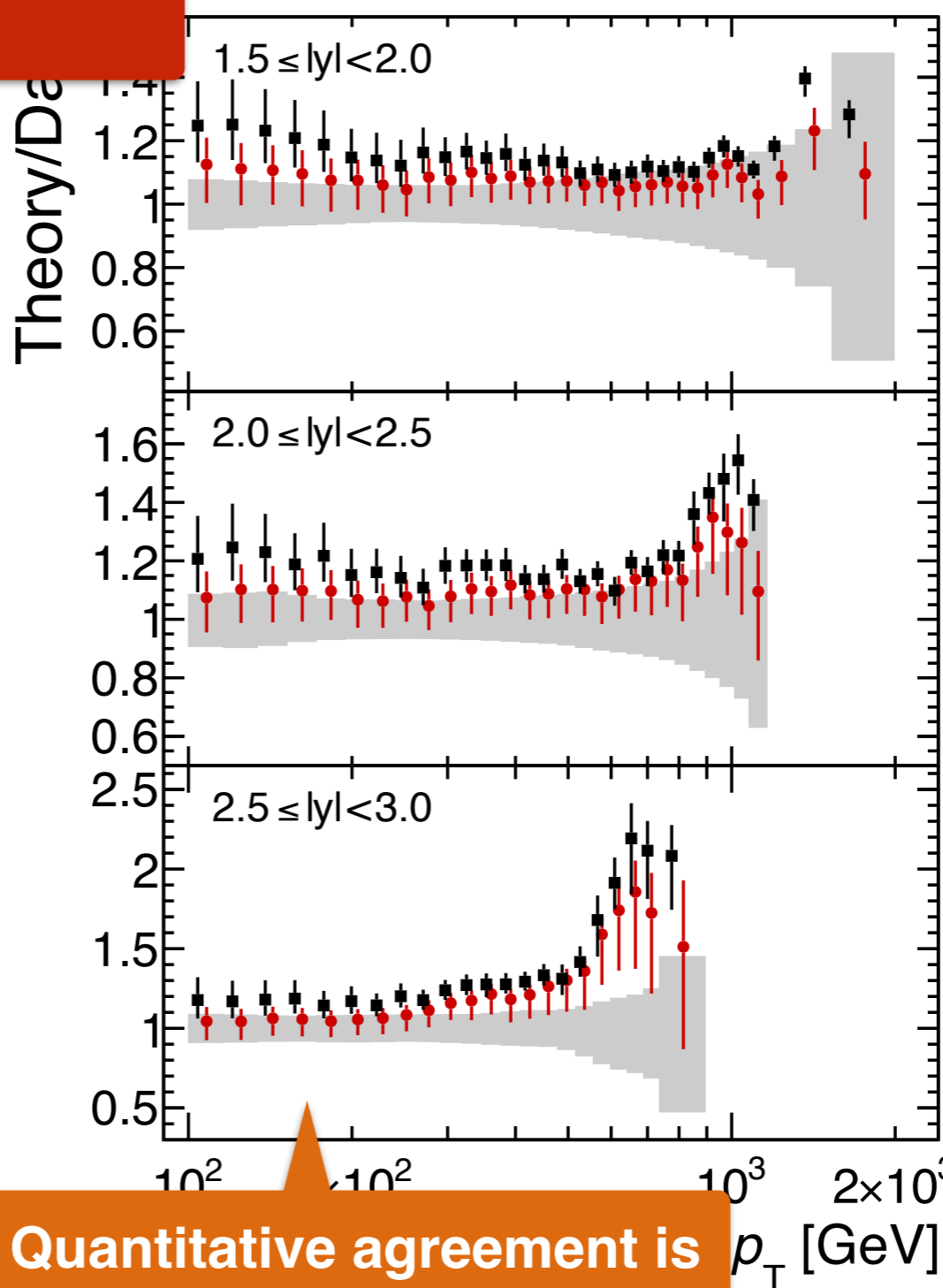
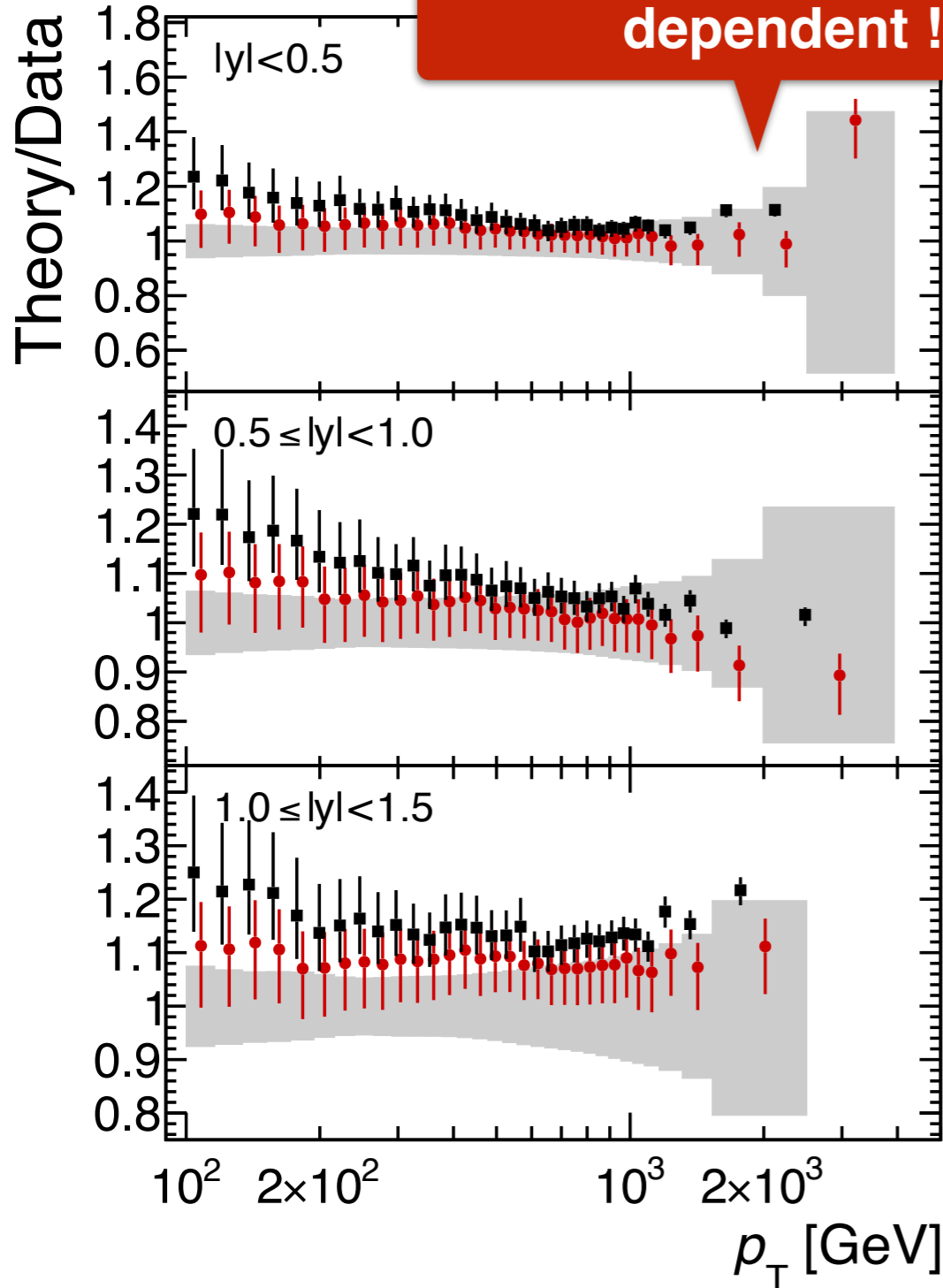
$\mu_R = \mu_F = p_T^{\text{jet}}$

● NLO  
MMHT 2014 NLO

■ NNLO  
MMHT 2014 NNLO

# Fixed-order pQCD

**Caution: story is scale-dependent !!**

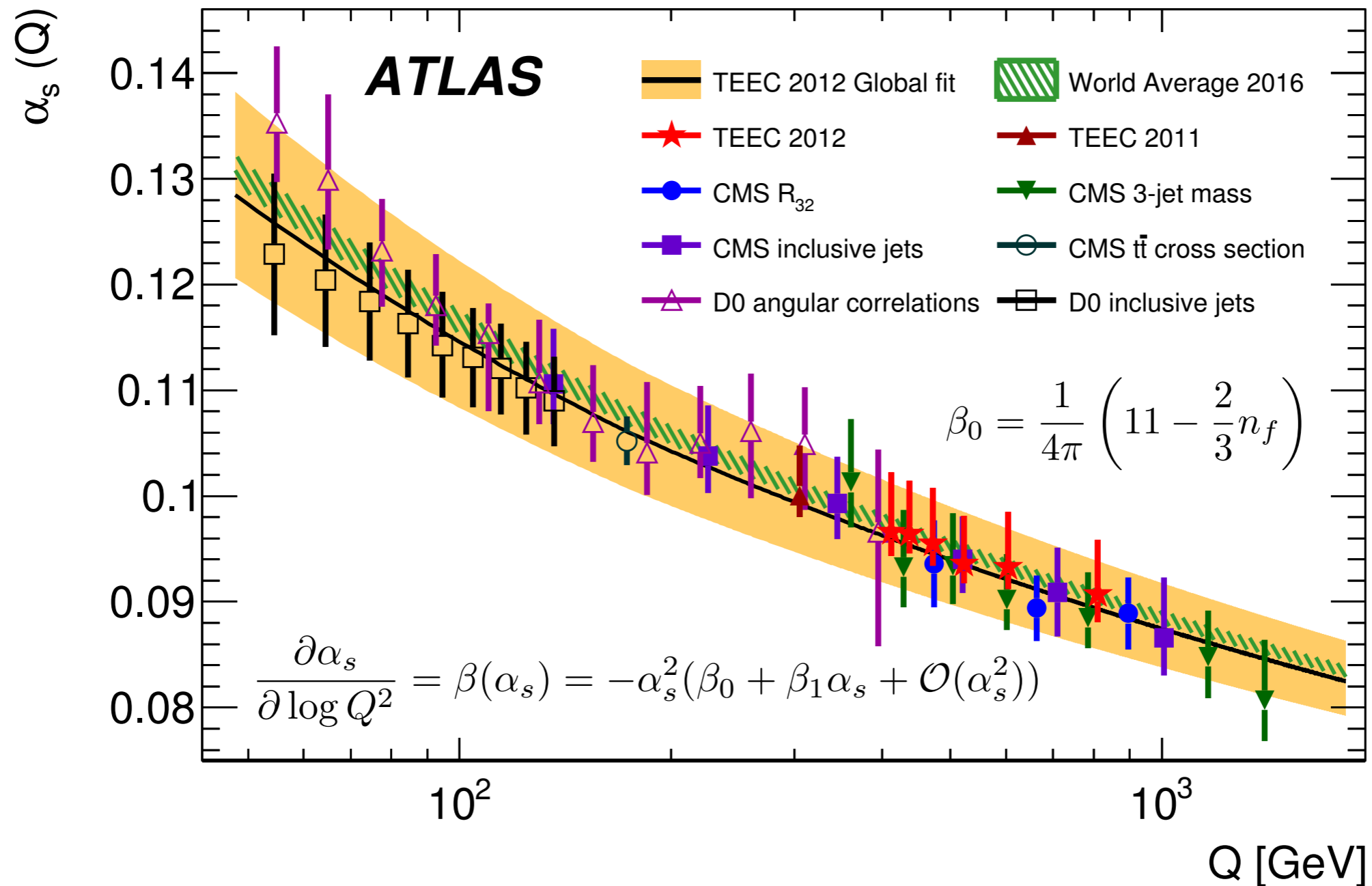


**Quantitative agreement is less good for all y-bins.**

**ATLAS**  
 $L = 81 \text{ nb}^{-1} - 3.2 \text{ fb}^{-1}$   
 $\sqrt{s} = 13 \text{ TeV}$   
 anti- $k_t$   $R=0.4$   
 ■ Data  
 ○ NLO QCD  
 ⊗  $k_{EW}$  ⊗  $k_{NP}$   
 ○ NNLO QCD  
 ⊗  $k_{EW}$  ⊗  $k_{NP}$   
 $\mu_R = \mu_F = p_T^{\text{max}}$   
 ● NLO  
 MMHT 2014 NLO  
 ■ NNLO  
 MMHT 2014 NNLO

# Strong coupling constant

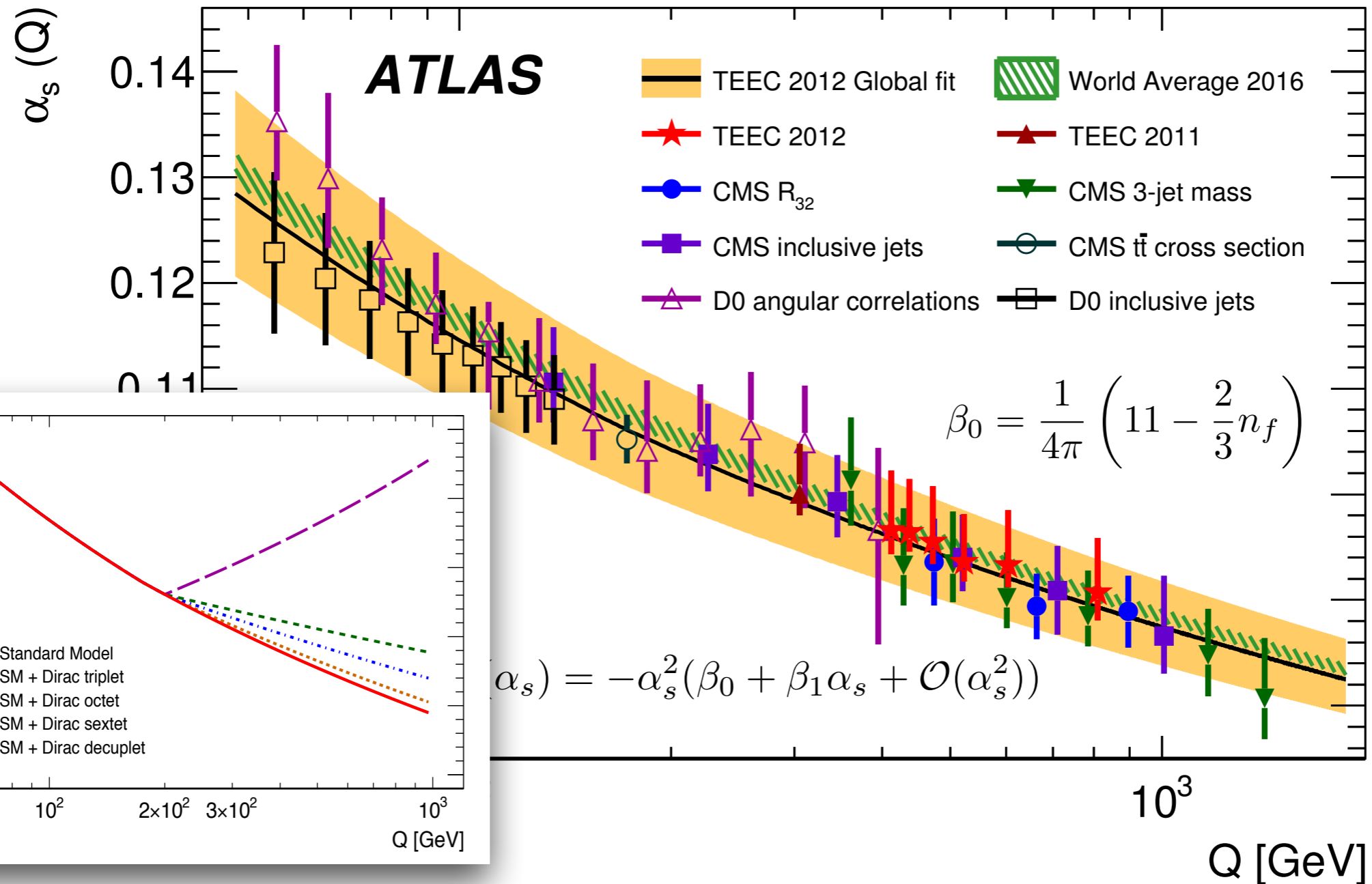
Scale challenges aside, one can use these data to test the running of the strong coupling at the highest energies.





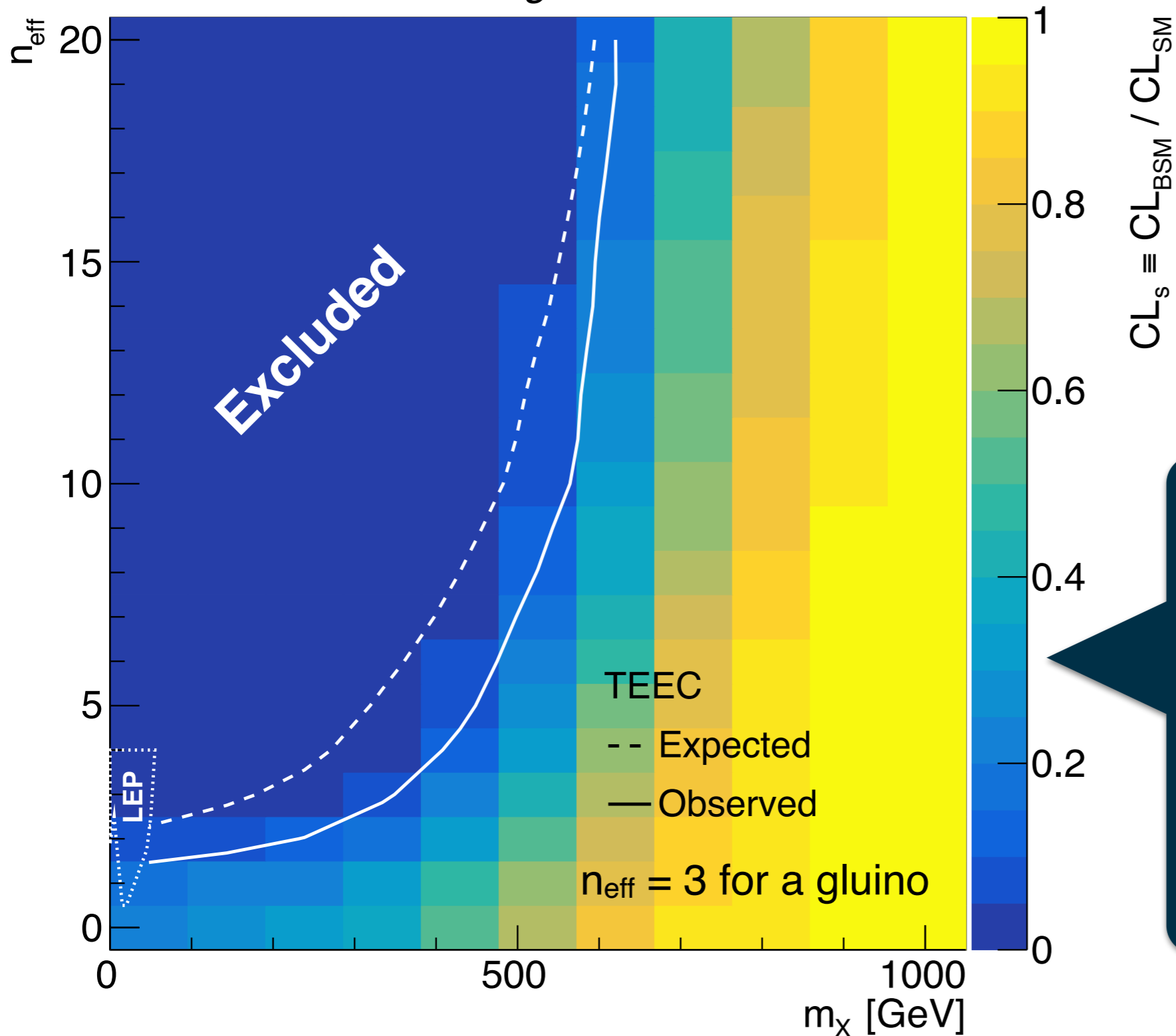
# Strong coupling constant

Scale challenges aside, one can use these data to test the running of the strong coupling at the highest energies.



# BSM from the running of $\alpha_s$

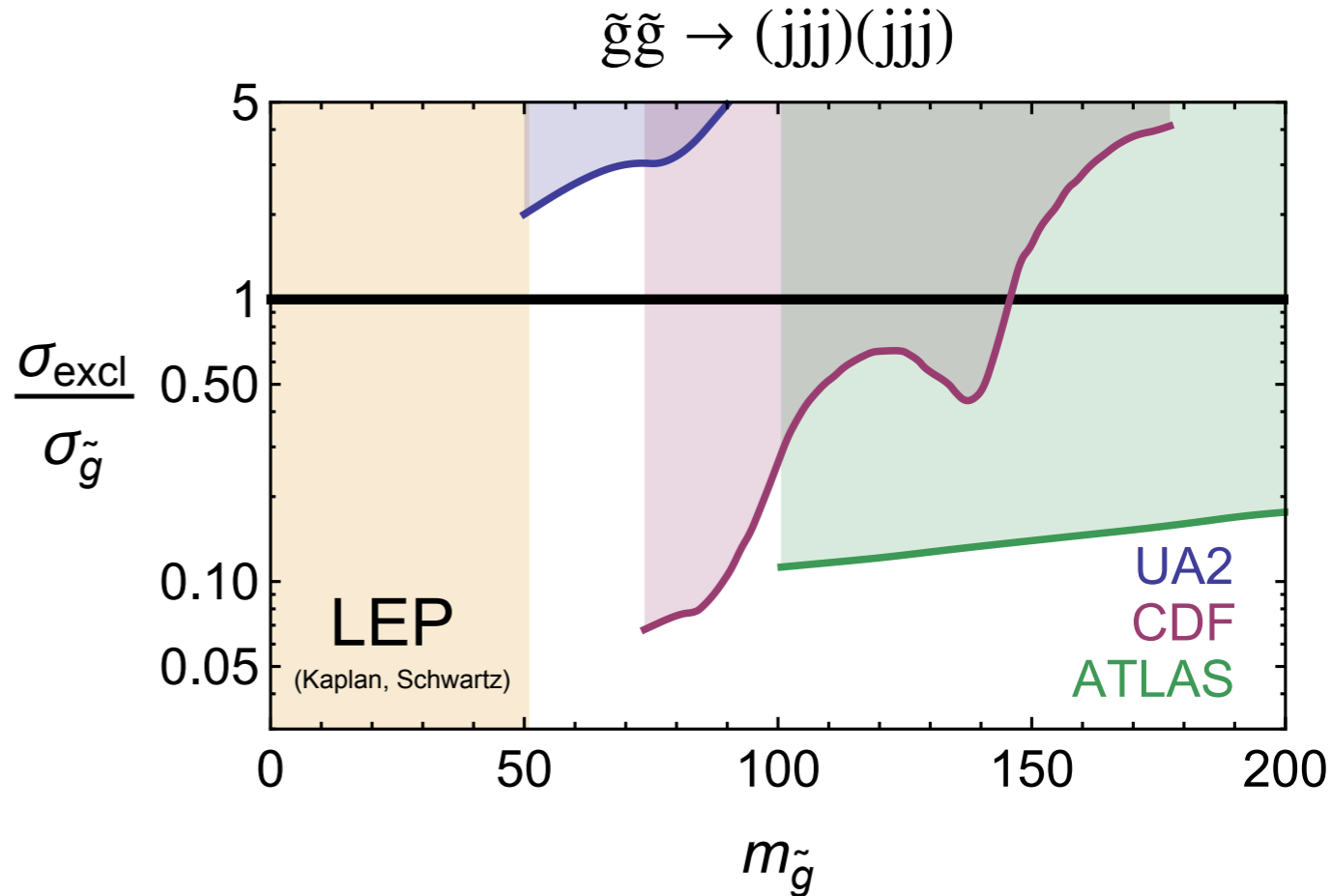
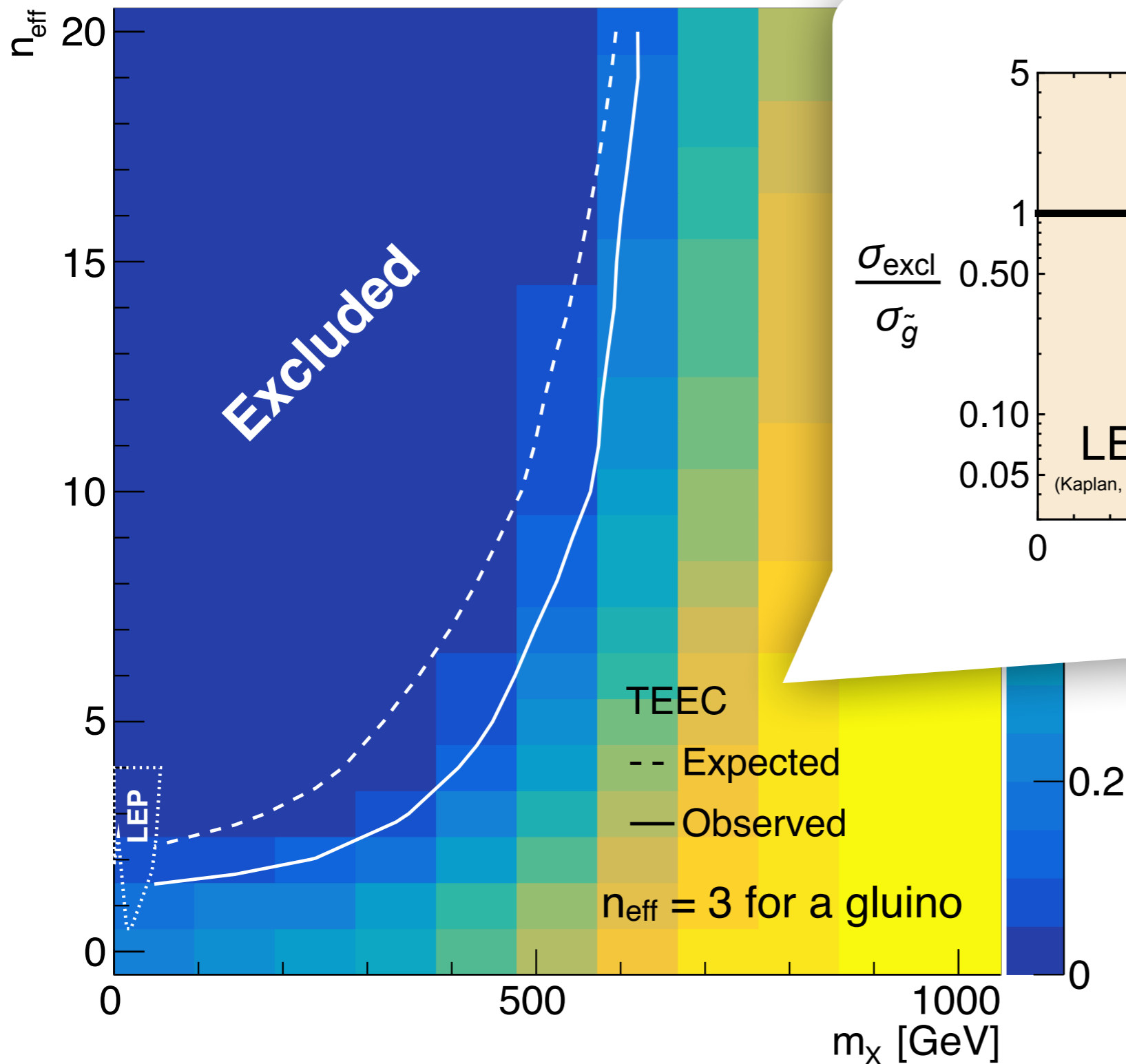
New fermion limits using NLOJet++ & ATLAS data



Complimentary to direct searches because ~agnostic to the decay of the new particles.

# BSM from the running of $a_s$

New fermion limits using NLOJet++ & ATLAS data

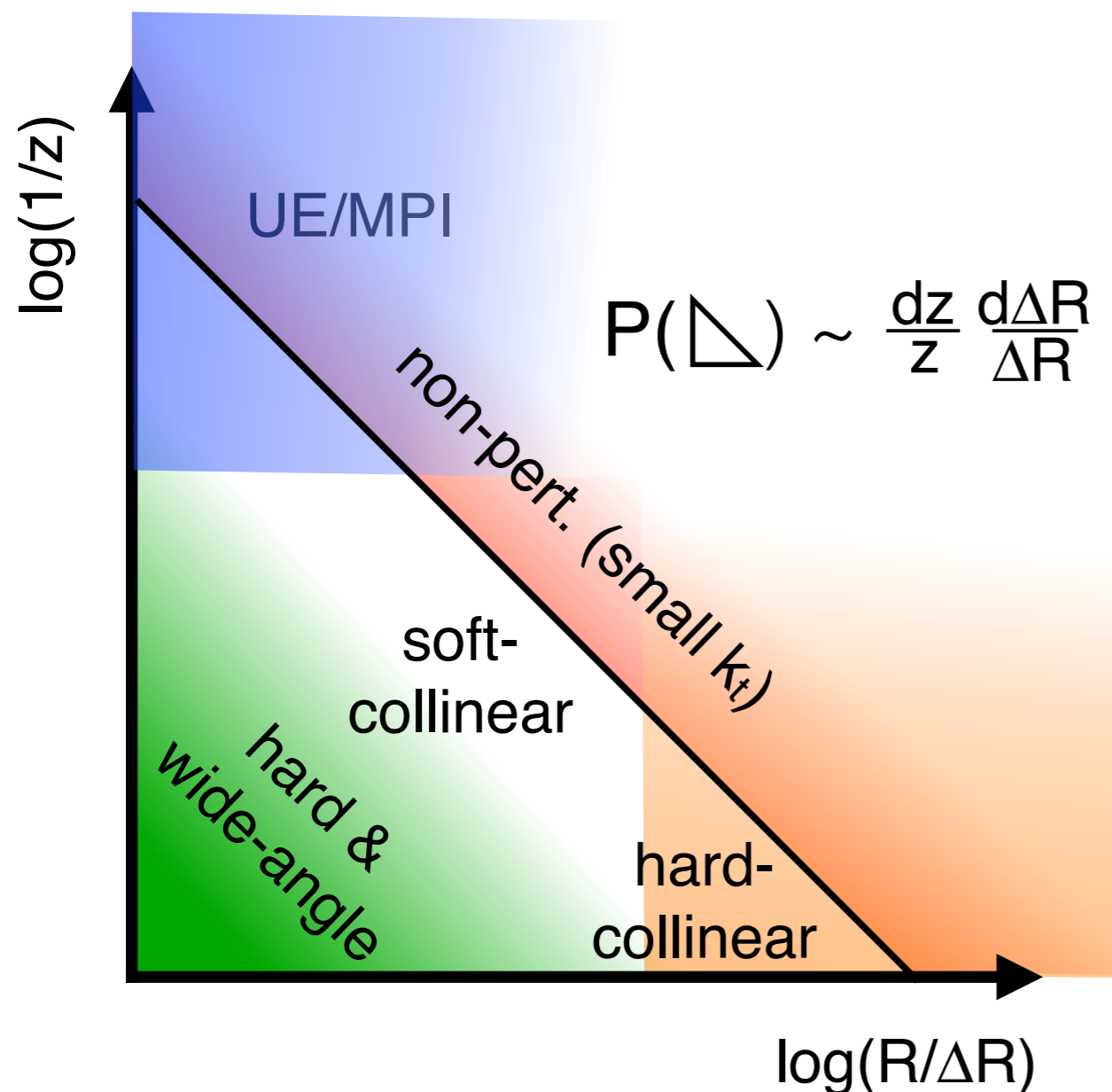


*Even though direct limits are  $O(\text{TeV})$ , there may be some gaps at low mass that techniques like this can probe.*

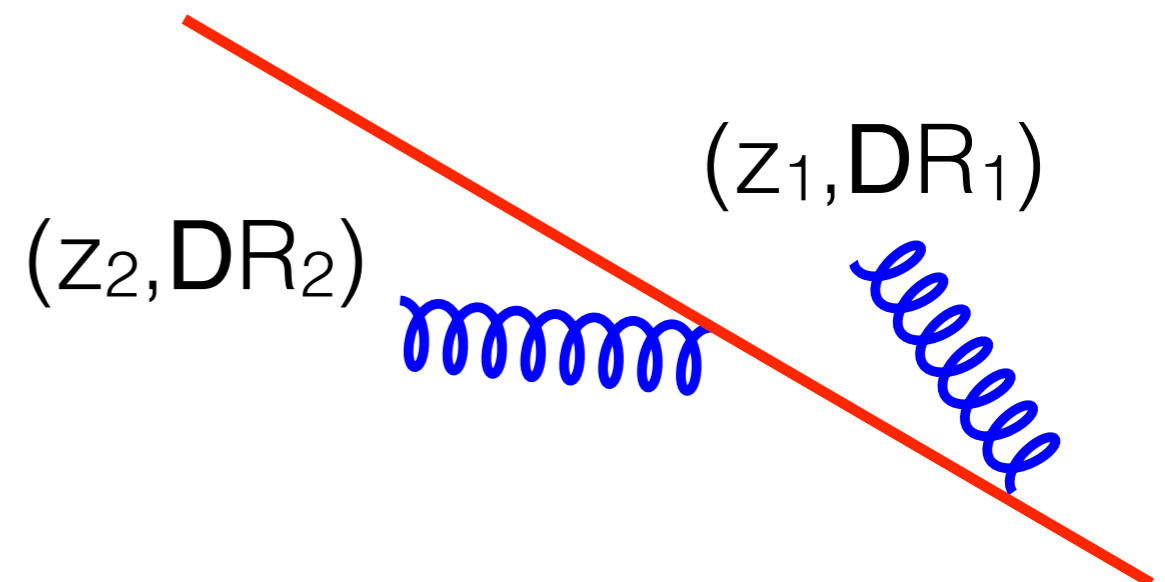
# Resummed pQCD and jet substructure

60

Well-separated jets can be described by fixed-order pQCD but the radiation pattern **inside jets** requires resummation.



A useful way of thinking about radiation inside the jet is in the context of the *Lund plane*



Well-separated jets can be described by fixed-order pQCD but the radiation pattern inside jets requires resummation.

For example, the jet mass is approximately

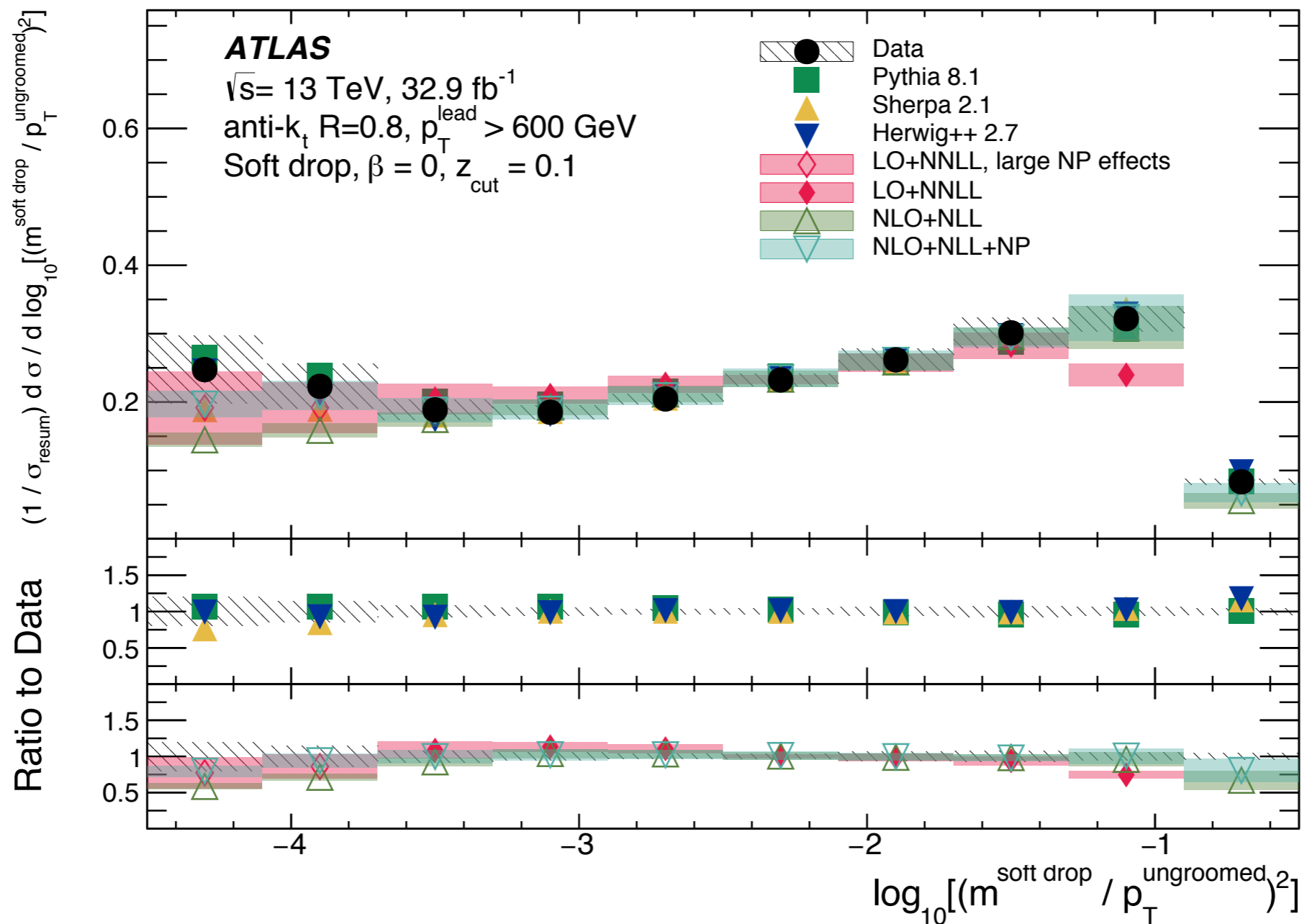
$m^2/p_T^2 \sim z\Delta R^2$  which is a line in the Lund plane:

(where did this come from?)

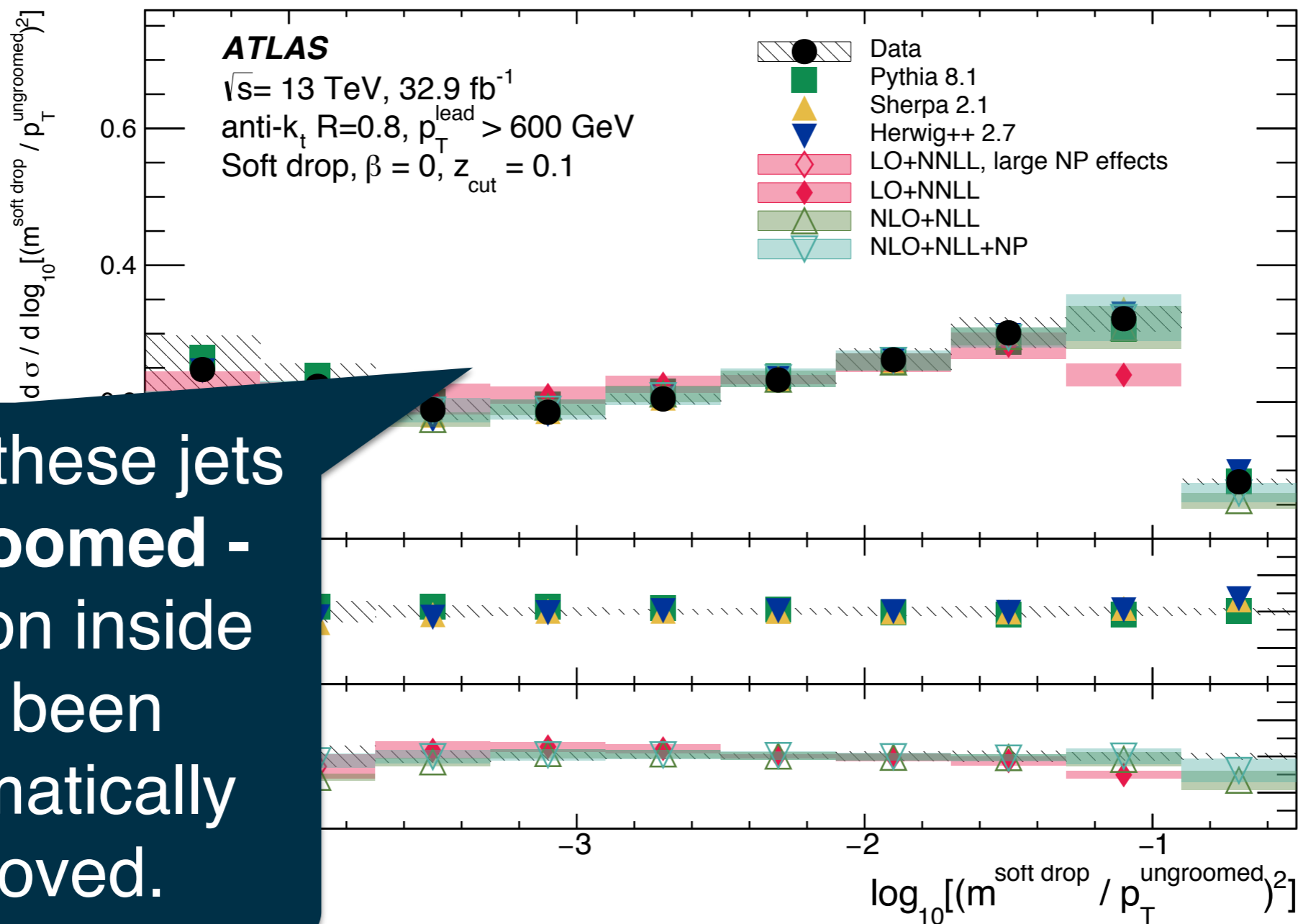
$$\log(1/z) = \text{constant} - 2 \log(R/\Delta R)$$

*One can use this to derive the probability distribution at leading logarithm accuracy.*

In fact, the jet mass has been calculated to higher accuracy and precisely measured!



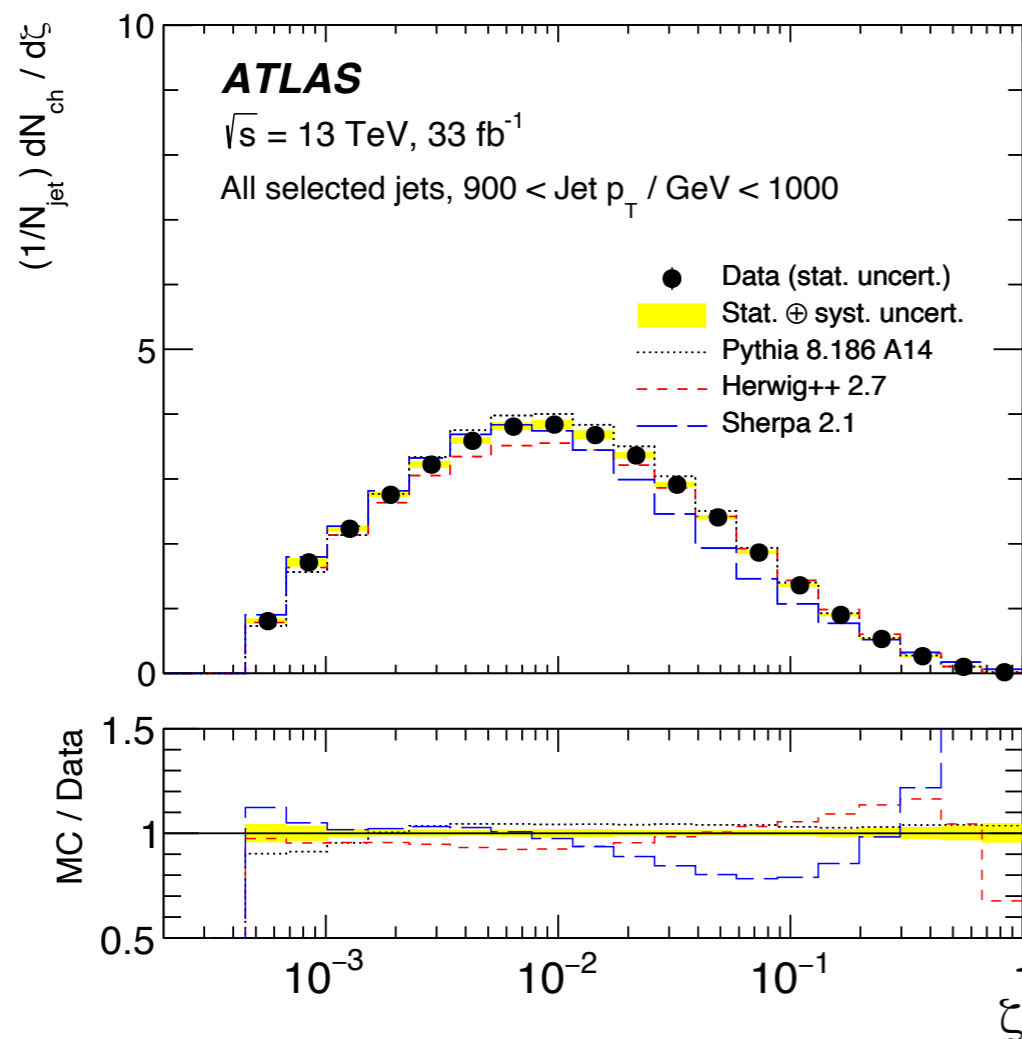
In fact, the jet mass has been calculated to higher accuracy and precisely measured!



Subtly: these jets are groomed - radiation inside has been systematically removed.

While jet substructure nowadays typically refers to observables like the jet mass and Lund plane coordinates,

there is a long history of more traditional “fragmentation function” measurements which probe similar physical regimes.



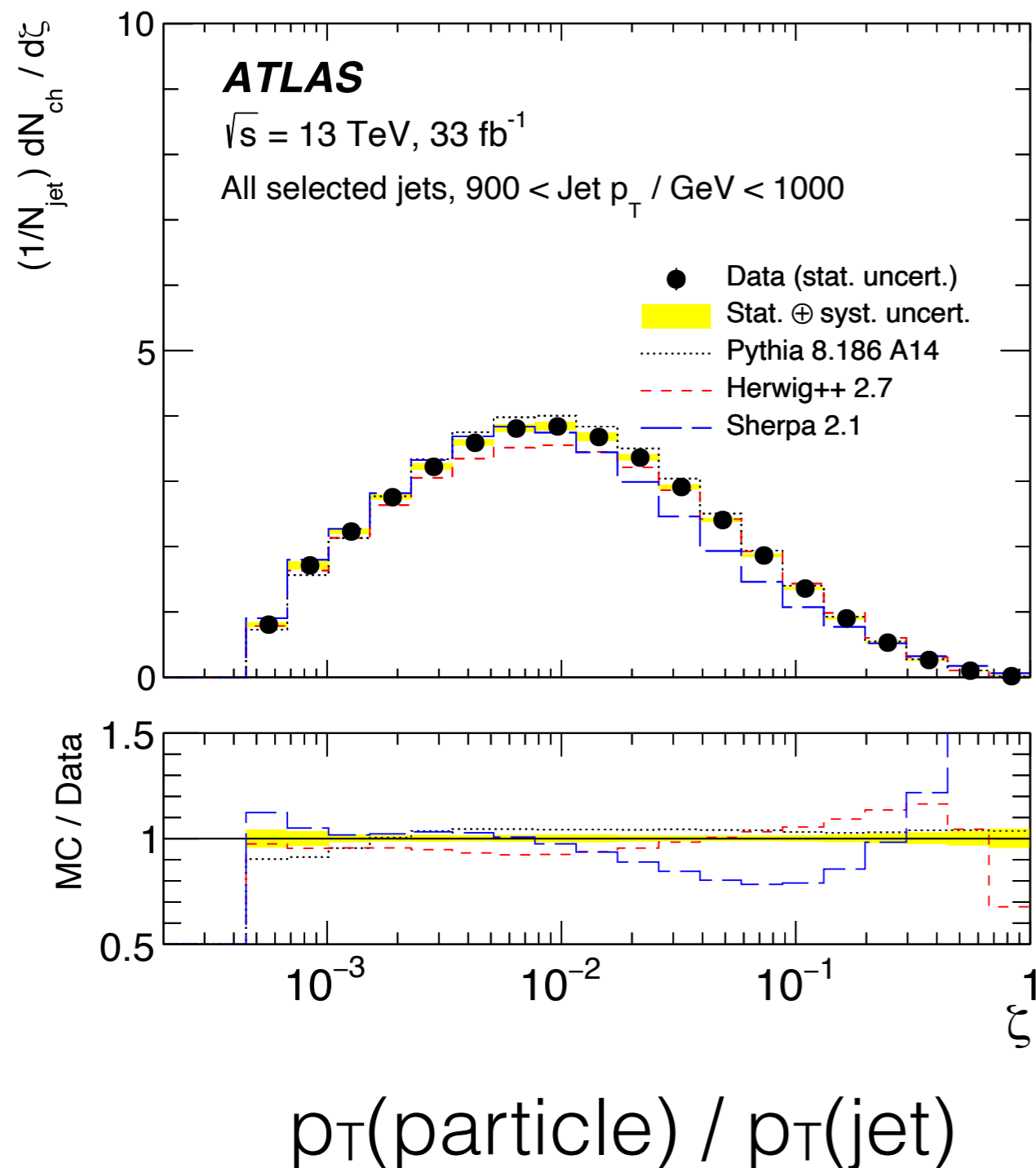
*These studies are crucial inputs to tuning Parton Shower Monte Carlo algorithms, but can also be used to study pQCD.*

←  $p_T(\text{particle}) / p_T(\text{jet})$



# Fragmentation: generic q/g jets

65



We can't generally measure individual particles, but we can for charged particles.

With charged-only, we can't predict this distribution from first principles. But for a given jet type (q/g) we can predict how it depend on jet  $p_T$ .

*How can we extract just the gluon (or quark) contribution?*

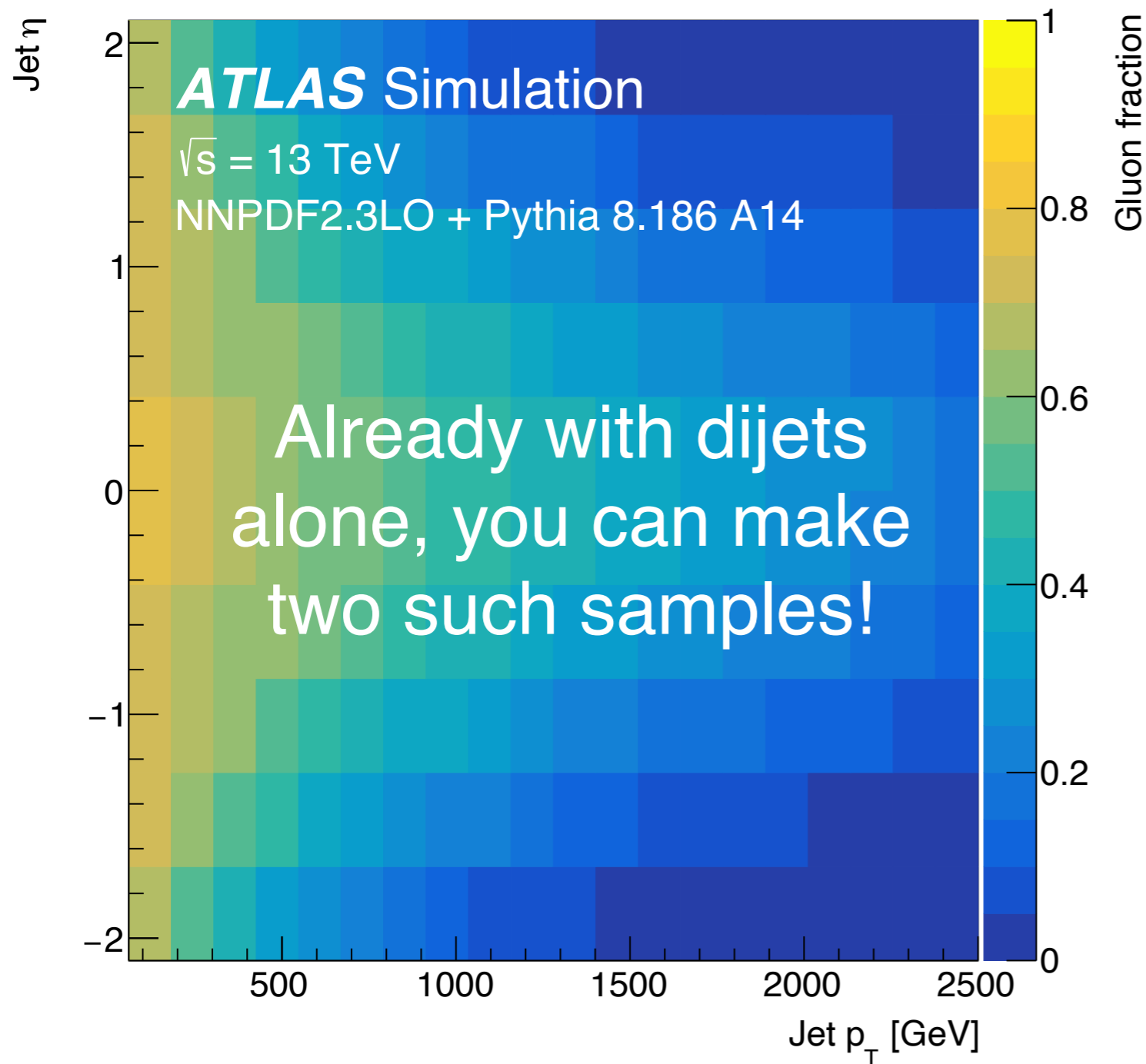
One way is to find two event samples ( $f$  and  $c$ ) with different q/g compositions.

Given the fractions  $f_q$ , you can extract the q/g bin contents  $h$ :

$$h_i^f = f_q^f h_i^q + (1 - f_q^f) h_i^g$$
$$h_i^c = f_q^c h_i^q + (1 - f_q^c) h_i^g$$

# The rapidity trick

67



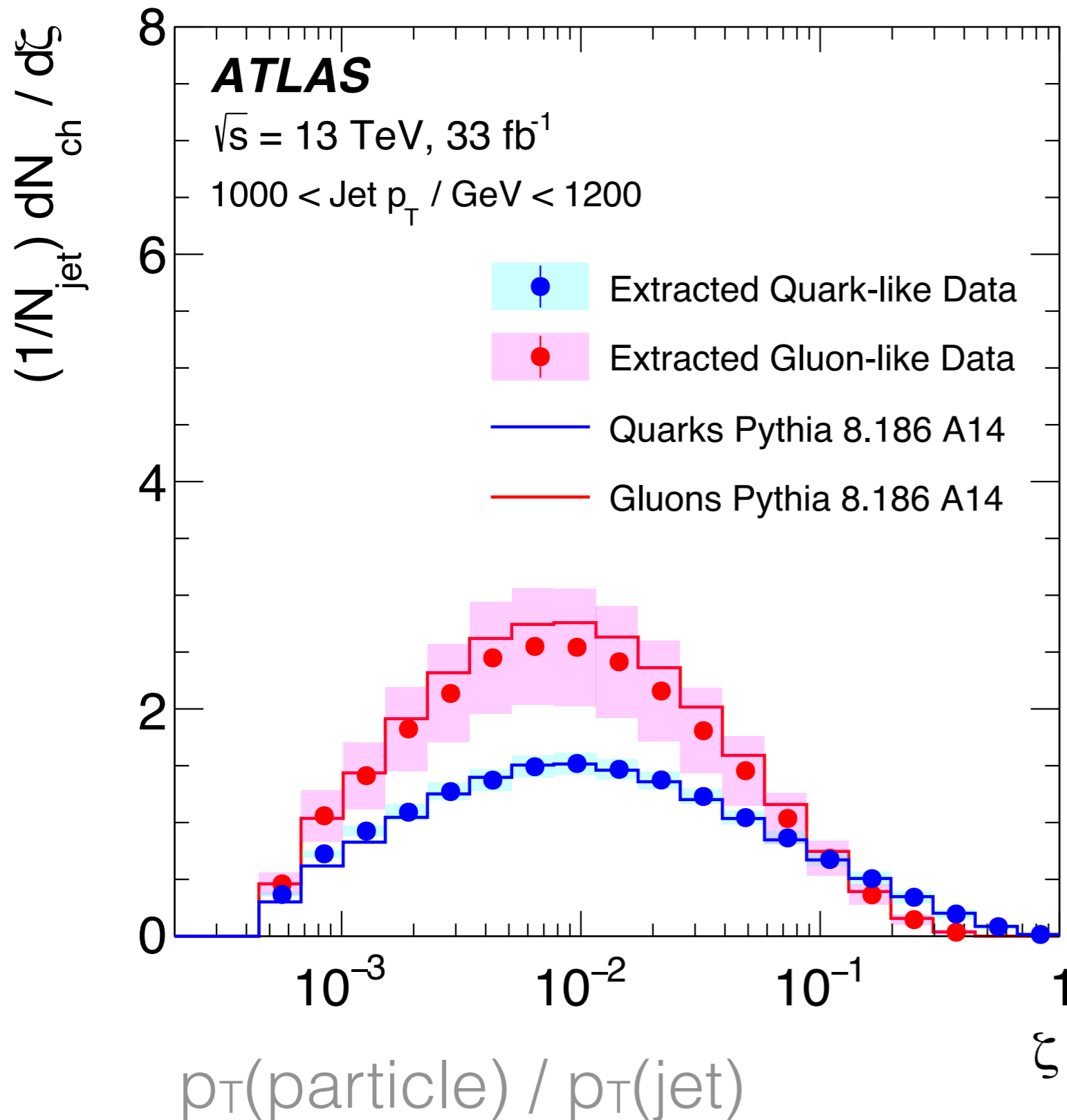
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# The rapidity trick

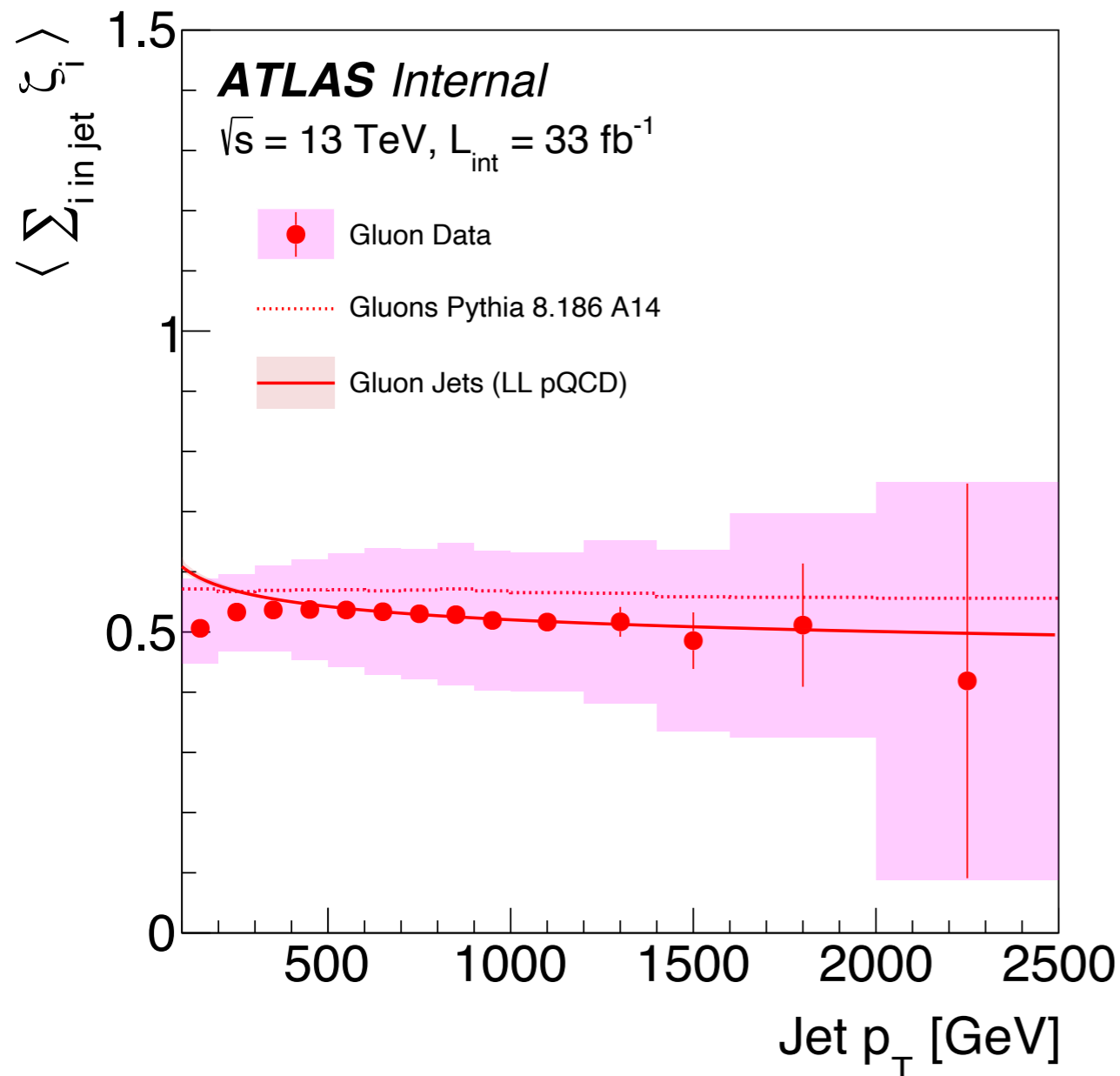


Gluons have more, softer particles than quarks.

*roughly gluons have 2x the particles as quarks since they have ~twice as much color charge*

# The $p_T$ dependence

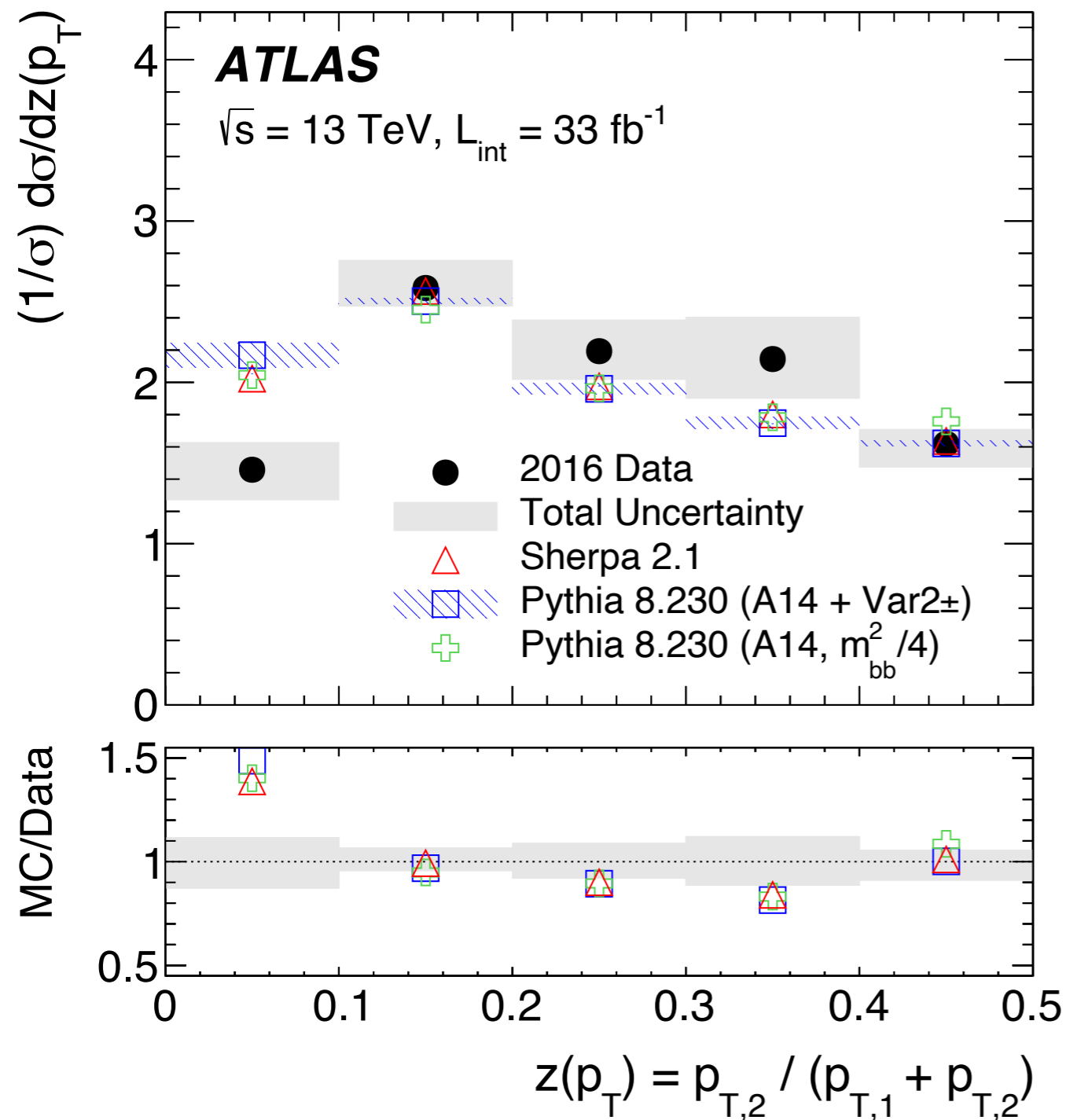
One can predict the  $p_T$  dependence of these observables, but let's just do a quick sanity check:



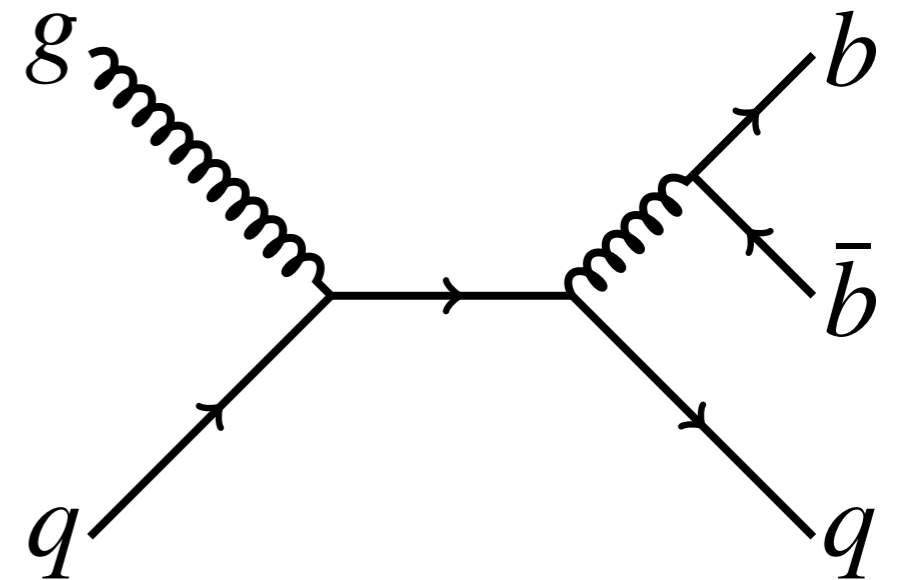
Why is the value  $\sim$ constant and a bit above 1/2?

# One last note on gluon fragmentation

70



Gluon splitting to  $b$ -quarks offers a unique opportunity to directly probe gluon fragmentation.



Question: why did I pick this diagram?

# QCD has much more to offer

71

I've offered a (biased) collection of results,  
but there are many more!

More on PDFs, fixed-order effects, ...

...Resummation, non-global effects ("entanglement"), ...

...Quark and gluon properties, W/Z/H hadronic decays, ...

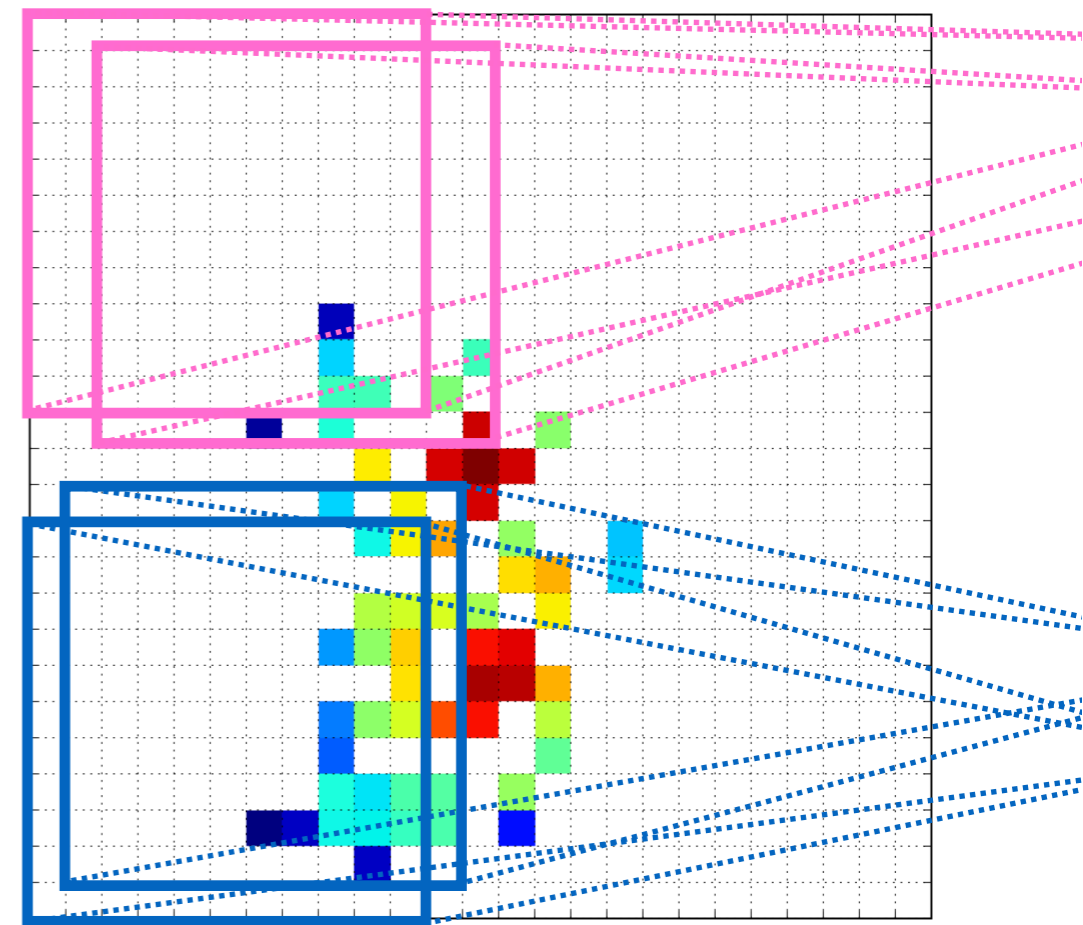
...Collective effects, connections with heavy ions, ...

...However, I stand between you and  
the end of the day, so let's wrap up!

Even though QCD has only  $\sim 1$  free parameter, it is a **rich theory with various regimes** that we can probe at the LHC.

Studying QCD is **inherently interesting** as a quantum theory of nature. Understanding it is also critical for direct and indirect **new particles searches**.

*...there are also many exciting connections to modern machine learning that I did not have time to discuss - consider attending [ML4Jets2022!](#)*



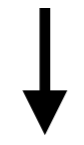


Questions?

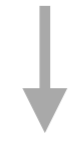


# Soft drop procedure

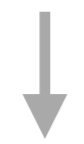
Take a jet clustered with e.g. anti- $k_t$



Re-cluster it with C/A



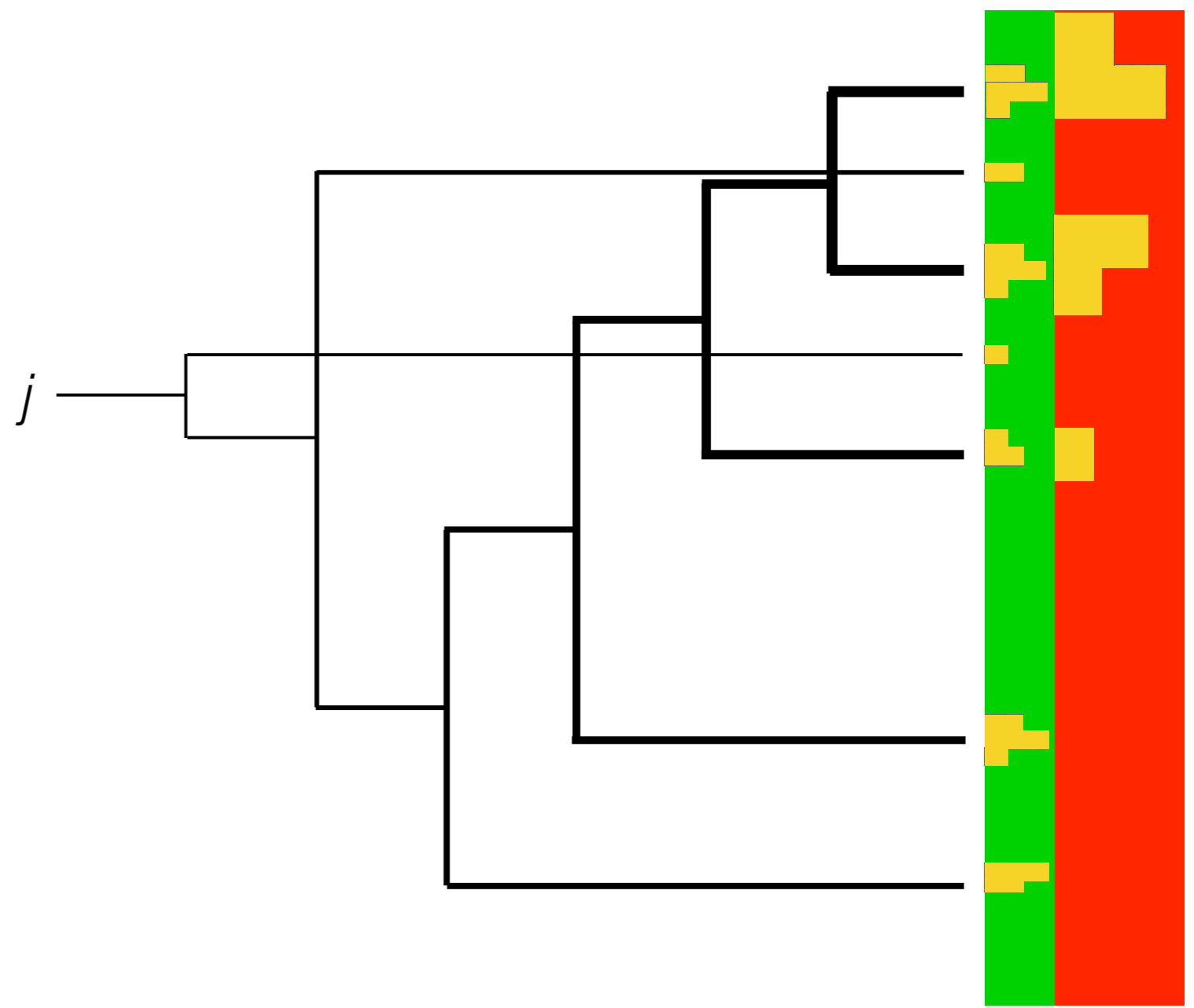
Traverse the clustering tree backwards



If a branch point satisfies the soft drop condition, stop.



Otherwise remove the softer branch and continue down the harder branch.



**clusters hardest radiation first**

# Soft drop procedure

Take a jet clustered with e.g. anti- $k_t$



Re-cluster it with C/A



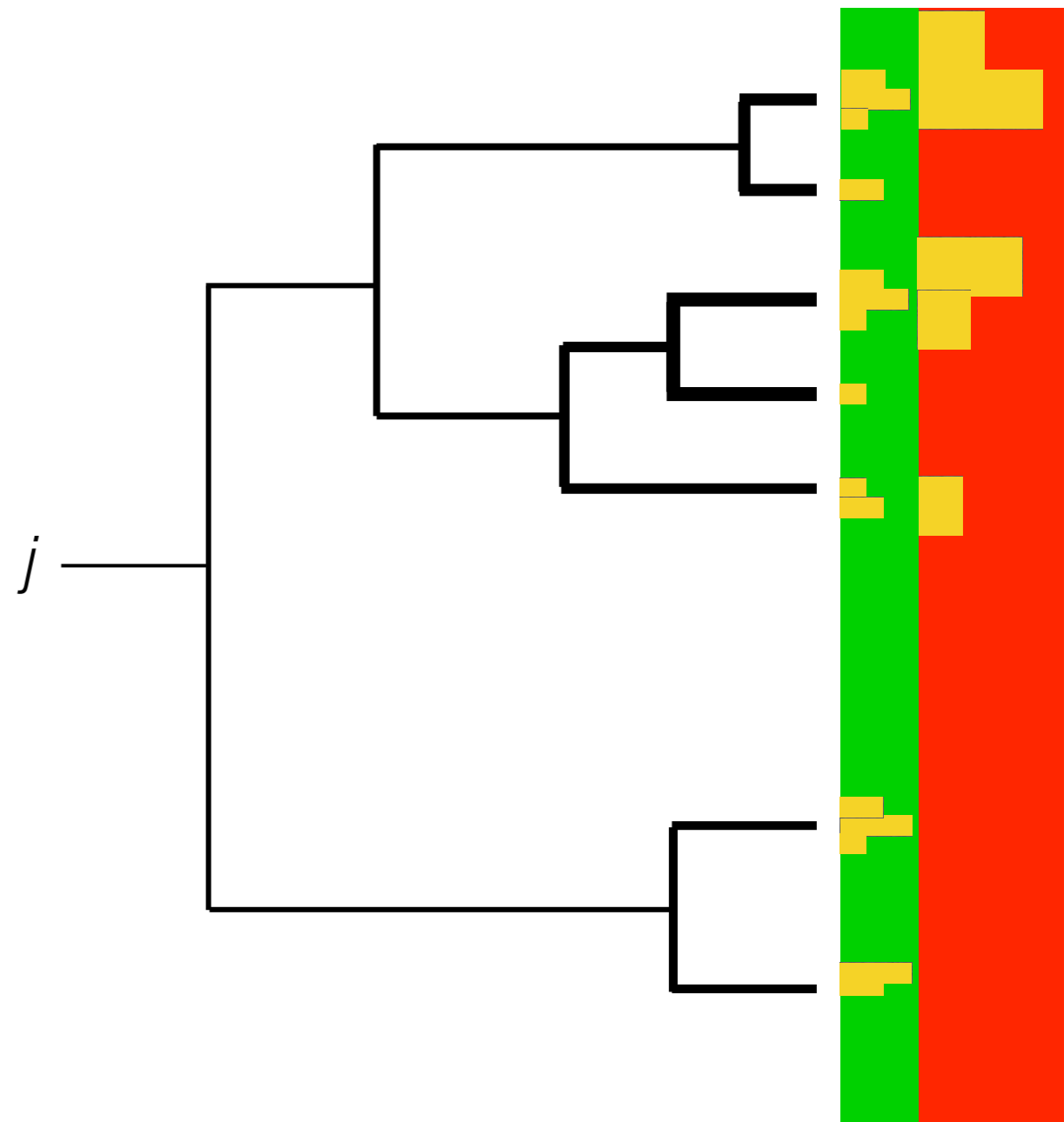
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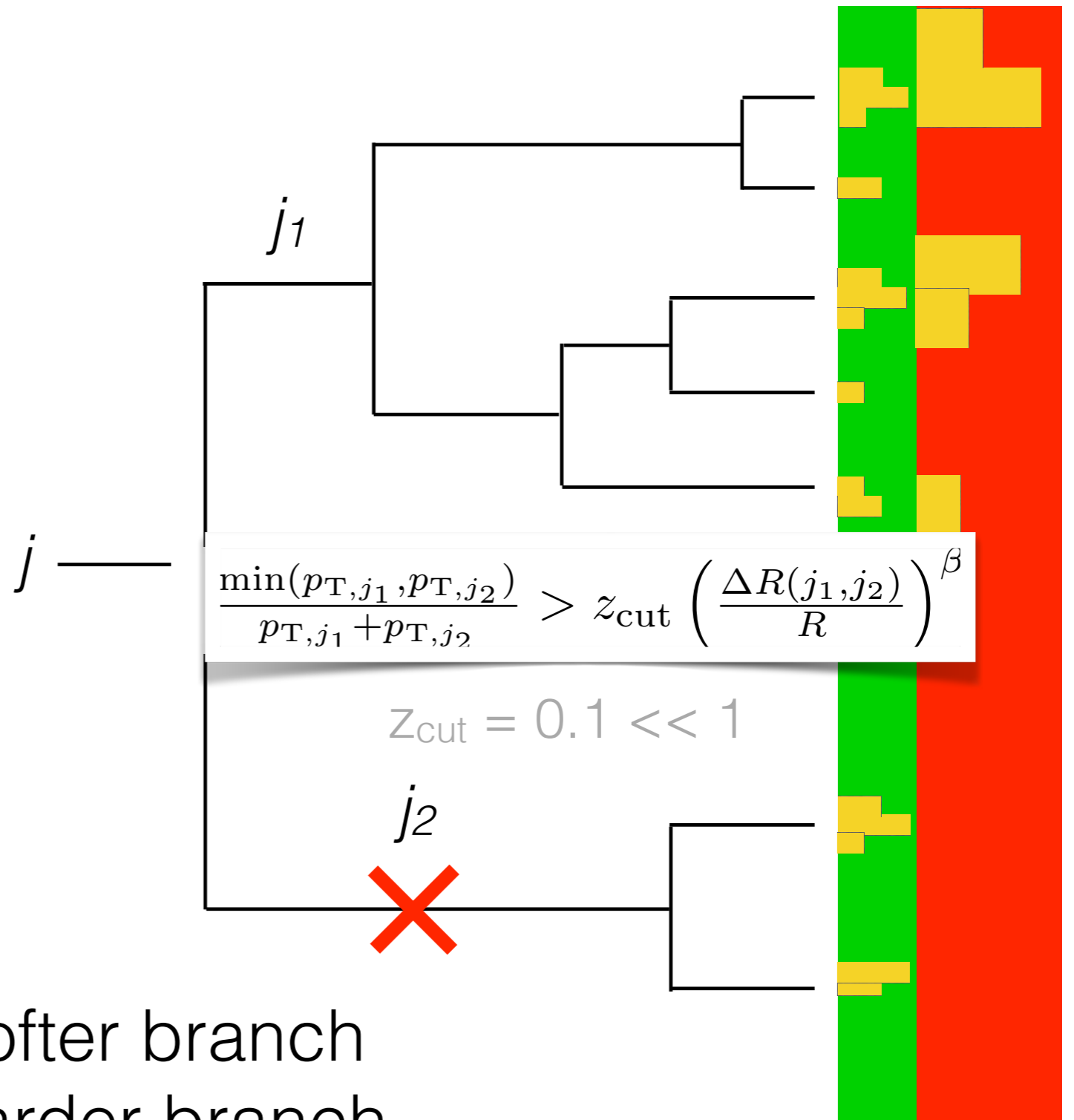
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