Probing the Supersymmetric Grand Unified Theories at the Future Proton-Proton Colliders and Hyper-Kamiokande Experiment

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Waqas Ahmed, TL, Shabbar Raza and Fang-Zhou Xu, Phys. Lett. B **819**, 136378 (2021) [arXiv:2007.15059 [hep-ph]]; in preparation.

- ► A natural solution to the gauge hierarchy problem in the SM.
- Gauge coupling unification can be achieved.
- The Lightest Supersymmetric Particle (LSP) such as the LSP neutralino etc can be a dark matter candidate.
- The electroweak gauge symmetry can be broken radiatively due to the large top quark Yukawa coupling.

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The Grand Unified Theories: SU(5) and SO(10)

- Gauge interaction unification.
- ► In SO(10) model, one family of the SM fermion forms a spinor 16 representation.
- Yukawa coupling unification.
- Charge quantization.
- Weak mixing angle at weak scale M_Z .
- ► Neutrino masses and mixings by seesaw mechanism.
- Prediction: dimension-six proton decay via heavy gauge boson exchange.

- Gauge symmetry breaking.
- Doublet-triplet splitting problem.
- Proton decay problem.
- Fermion mass problem.

The wrong prediction on the fermion mass ratios: $m_e/m_\mu = m_d/m_s.$

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- Calabi-Yau compactification of heterotic string theory
- Orbifold compactification of heterotic string theory
- D-brane models on Type II orientifolds
- Free fermionic string model builing
- *F*-Theory Model Building

Supersymmetry is a bridge between the promising low energy phenomenology and high-energy fundamental physics.

String Theory \rightarrow String Models \rightarrow SUSY GUTs \rightarrow SSMs \rightarrow SM

Can we probe this grand particle physics paradiagm via the future pp colliders and other experiments? If yes, what is the center-of-mass energy needed?

We can probe the SUSY GUTs at the future experiments!

- ► Lepton colliders: CEPC, CLIC, *FCC*_{ee}, and ILC.
- ► Hadron colliders: HL-LHC, HE-LHC, *FCC*_{hh}, SppC, and VLHC.

To probe the new physics beyond the SM, we do need future proton-proton colliders.

- Question: what is the concrete scientific goal for the future pp colliders?
- Question: what is the center-of-mass energy needed for this scientific goal?

- Supersymmetry cannot be the scientific goal since the sparticles can be very heavy and then we cannot probe supersymmetry.
- The supersymmetric GUTs with grand desert hypothesis can be the scientific goal ¹.

¹Waqas Ahmed, TL, Shabbar Raza and Fang-Zhou Xu, Phys. Lett. B **819**, 136378 (2021) [arXiv:2007.15059 [hep-ph]]; in preparation.

- Grand desert hypothesis: no new physics between the sparticle mass scale and the GUT scale.
- ▶ For the GUTs with the GUT scale $M_{GUT} \le 1.0 \times 10^{16}$ GeV, we can probe the dimension-six proton decay via heavy gauge boson exchange at the Hyper-Kamiokande experiment.
- ▶ For the GUTs with $M_{GUT} \ge 1.0 \times 10^{16}$ GeV, we can probe the gluino and/or squarks at the future pp colliders.
- Providing the "Concrete Scienctific Goal" for the future pp colliders and Hyper-Kamiokande experiment.

For the GUTs with the GUT scale $M_{GUT} \leq 1.0 \times 10^{16}$ GeV, we can probe the dimension-six proton decay via heavy gauge boson exchange at the Hyper-Kamiokande experiment.

The Predictions of the SUSY GUTs: Proton Decays

► The dimension-six proton decay via superheavy (X_µ, Y_µ) gauge boson exchanges

$$SU(5) = \begin{pmatrix} SU(3)_C & (\overline{X}_{\mu}, \overline{Y}_{\mu}) \\ (X_{\mu}, Y_{\mu}) & SU(2)_L \end{pmatrix}$$

- The dimension-five proton decays via colored Higgsino exchanges in the supersymmetric GUTs.
- ▶ We do not consider the dimension-five proton decays.

We can have the dimension-five proton decays at renormalizable level by introducing vector-like particles, and can nelect the dimension-five proton decays in some models such as the flipped SU(5) models.

The Dimension-Six Proton Decay via (X_{μ}, Y_{μ}) Exchanges



The GUTs with $M_{GUT} \leq 1.0 \times 10^{16}$ GeV

• The proton lifetime from the dimension-six proton decay $p \rightarrow e^+ \pi^0$ via heavy gauge boson exchange is

$$egin{aligned} \pi_p &\simeq & 1.0 imes 10^{34} imes \left(rac{2.5}{A_R}
ight)^2 imes \left(rac{0.04}{lpha_{
m GUT}}
ight)^2 \ & imes \left(rac{M_{
m GUT}}{1.0 imes 10^{16} \ {
m GeV}}
ight)^4 \ {
m years} \; . \end{aligned}$$

- ► The current lower limit from the Super-Kamiokande experiment gives τ_p > 2.4 × 10³⁴ years ², and thus we obtain M_{GUT} ≥ 1.0 × 10¹⁶ GeV.
- At the future Hyper-Kamiokande experiment, we can probe the proton lifetime at least above 1.0 × 10³⁵ years ³. Therefore, the GUTs with M_{GUT} ≤ 1.0 × 10¹⁶ GeV is definitely within its reach.
- ²K. Abe *et al.* [Super-Kamiokande], Phys. Rev. D **95**, no.1, 012004 (2017) [arXiv:1610.03597 [hep-ex]]; A. Takenaka *et al.* [Super-Kamiokande], Phys. Rev. D **102**, 112011 (2020) [arXiv:2010.16098 [hep-ex]]. ³K. Abe *et al.* [Hyper-Kamiokande], [arXiv:1805.04163 [physics.ins-det]]. □ → (¬→) (≥→) (≥→) (≥→)

- In fact, if we did not observe the dimension-six proton decay from the Super-Kamiokande experiment and future Hyper-Kamiokande experiment, we obtain $M_{X\mu/Y\mu} \geq 1.0 \times 10^{16} \text{ GeV}$, not $M_{\text{GUT}} \geq 1.0 \times 10^{16} \text{ GeV}$.
- There indeed exists some subtleties to define the GUT scale due to threshold corrections.
- ► Because $M_{X_{\mu}/Y_{\mu}} \leq M_{GUT}$, our study is not affected by these subtleties.

For the GUTs with $M_{GUT} \ge 1.0 \times 10^{16}$ GeV, we can probe the gluino and/or squarks at the future pp colliders.

The supersymmetry breaking mediation mechanisms

- Gravity mediated supersymmetry breaking.
- Anomaly mediated supersymmetry breaking.
- Gauge mediated supersymmetry breaking.

- ► At the future 100 TeV pp Colliders such as FCC_{hh} and SppC, gluino *g̃* via heavy flavor decay, gluino via light flavor decay, and the first-two generation squarks *q̃* can be discovered for their masses up to about 11 TeV, 17 TeV, and 14 TeV, respectively. If the gluino and first-two generation squark masses are similar, they can be probed up to 20 TeV.
- ► To discover the gluino \tilde{g} with mass around 15 TeV via heavy flavor decay, we need the 160 TeV pp collider such as the VLHC.

- ► The SM-like Higgs boson mass m_h ⊂ [123, 127] GeV, and gluino mass m_ğ ≥ 2.2 TeV.
- ► The constraints from rare decay processes $B_s \rightarrow \mu^+ \mu^-$, $b \rightarrow s\gamma$, and $B_u \rightarrow \tau \nu_{\tau}$.
- ► To be general, we do not require the relic abundance of the LSP neutralino to satisfy the Planck bound within 5σ $0.114 \leq \Omega_{\rm CDM} h^2 \leq 0.126$.

We perform the random scans for the following mSUGRA/CMSSM parameter space $% \mathcal{M} = \mathcal{M} = \mathcal{M} + \mathcal{M} +$

$$\begin{array}{l} 0 \leq M_0 \leq 90 \ {\rm TeV}, \\ 0 \leq M_{1/2} \leq 30 \ {\rm TeV}, \\ -3 \leq A_0/M_0 \leq 3, \\ 2 \leq \tan \beta \leq 60 \end{array}$$

with $\mu > 0$ and $m_t = 173.2 \,\mathrm{GeV}$.



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- The upper bound on the universal gaugino mass M_{1/2} is about 7 TeV, and then the upper bound on gluino mass is 15 TeV.
- In our viable parameter space, the lightest squark is generically to be light stop, and thus we do have gluino via heavy flavor decay. To probe such gluino with mass up to 15 TeV, we find that the center-of-mass energy of the future pp collider needs to be about 160 TeV.
- ► The needed center-of-mass energy is 160 TeV.
- In the viable parameter space with correct dark matter relic density and tan β > 7.5, gluino mass is less than 10 TeV. Thus, we can probe it at the FCC_{hh} and SppC.

We perform the random scans over the following parameter space of the minimal AMSB

$$\begin{split} 1\,\mathrm{TeV} &\leq M_0 \leq 7.5\,\mathrm{TeV},\\ 100\,\mathrm{TeV} &\leq M_{3/2} \leq 3000\,\mathrm{TeV},\\ 2 \leq \tan\beta \,\leq 60 \end{split}$$

with $\mu > 0$ and $m_t = 173.2 \,\mathrm{GeV}$.

With one pair of the messenger fields in the **5** and $\overline{5}$ representations of SU(5), we perform random scans over the following parameter space

$$\begin{split} 5\times 10^5\,{\rm GeV} \, \leq & \Lambda \equiv F/M_{\rm mess} \, \leq 10^{10}\,{\rm GeV}, \\ 2\times \Lambda \, \leq & M_{\rm mess} \leq 10^{15}\,{\rm GeV}, \\ 2 \leq & \tan\beta \, \leq 60 \end{split}$$

with $\mu > 0$ and $m_t = 173.2 \,\mathrm{GeV}$.

Anomaly and Gauge Mediated Supersymmetry Breakings



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- For anomaly mediation, the squark and gluino masses are less than about 5 TeV.
- For anomaly mediation, the squark and gluino masses are less than about 8 TeV and 6 TeV, respectively.
- ▶ The GUTs with anomaly and gauge mediated supersymmetry breakings are well within the reaches of the future 100 TeV pp colliders such as the FCC_{hh} and SppC.

We cannot probe the Pati-Salam models and SM-like models from four-dimensional free-fermionic string constructions as well as the Pati-Salam models from D-brane models on Type II orientifolds.

No dimension-six proton decay.

- All the other string models might be probed except the models with one-step string-scale gauge coupling unification.
- The $\mathcal{F} SU(5)$ can be probed and is under investigation.

The flipped SU(5) model with vector-like particles and string-scale gauge coupling unification from F-theory model building or free-fermionic string constructions.

- Supersymmetry is a bridge between the low energy phenomenology and high-energy fundamental physics, and thus is the promising new physics beyond the SM.
- Gauge coupling unification in the supersymmetric SM strongly implies the GUTs.
- With the grand desert hypothesis, we show that the supersymmetric GUTs and some string models can be probed at the future pp colliders and Hyper-Kamiokande experiment.

Thank You Very Much for Your Attention!

