Stasis in an Expanding Universe: A Recipe for Stable Mixed-Component Cosmological Eras

Fei Huang

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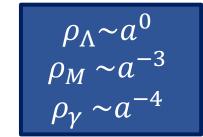
in collaboration with

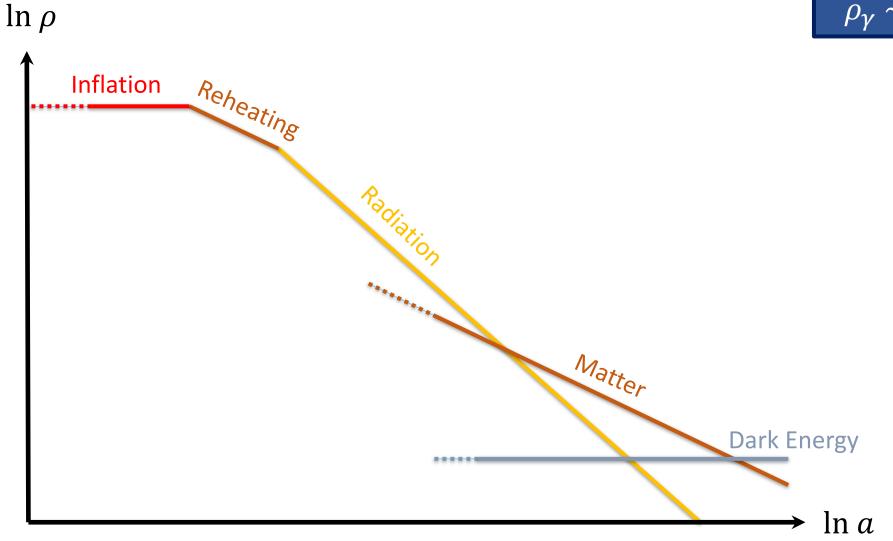
Keith Dienes, Lucien Heurtier, Doojin Kim, Tim Tait, Brooks Thomas





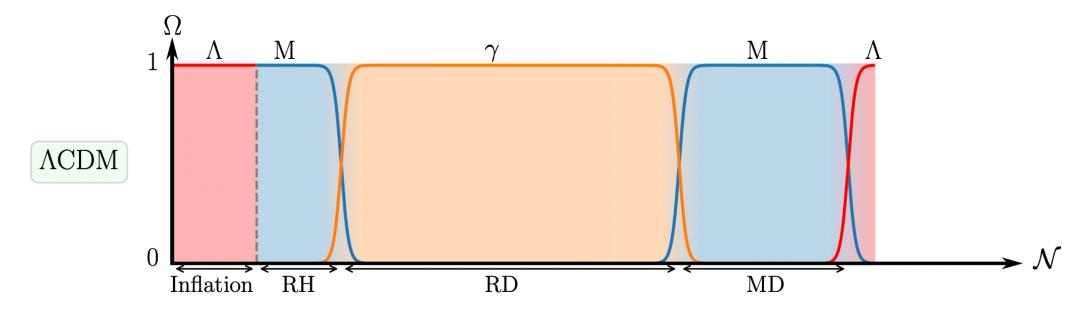
The Standard Lore: ΛCDM





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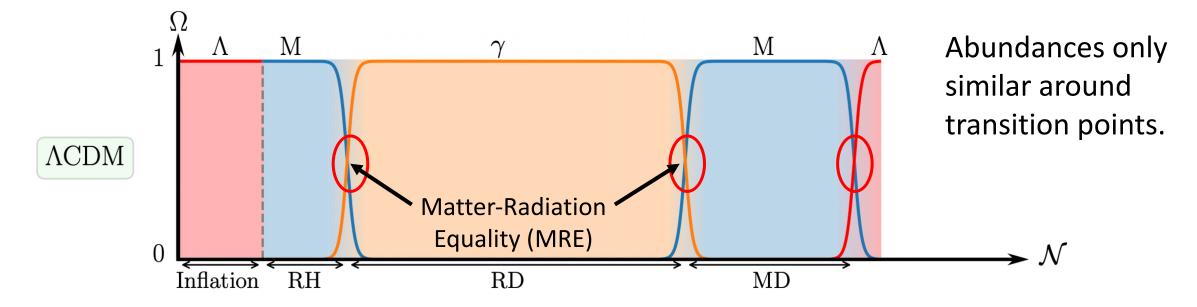
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 What if the MRE is not just a point, but a sustained era in the cosmological timeline?

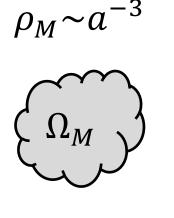
• More generally, is it possible to maintain a constant Ω_M and Ω_γ over a sizable period of time?

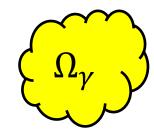
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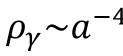
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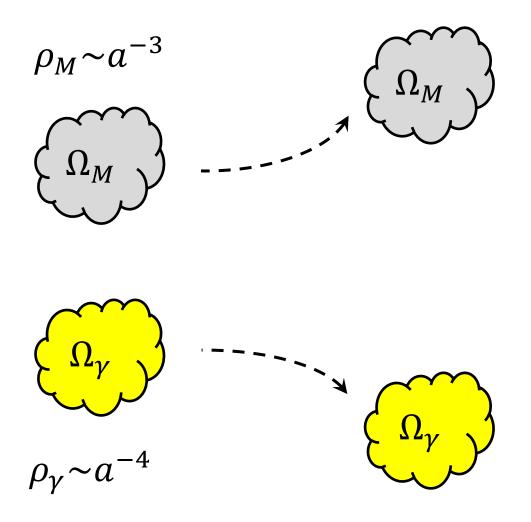


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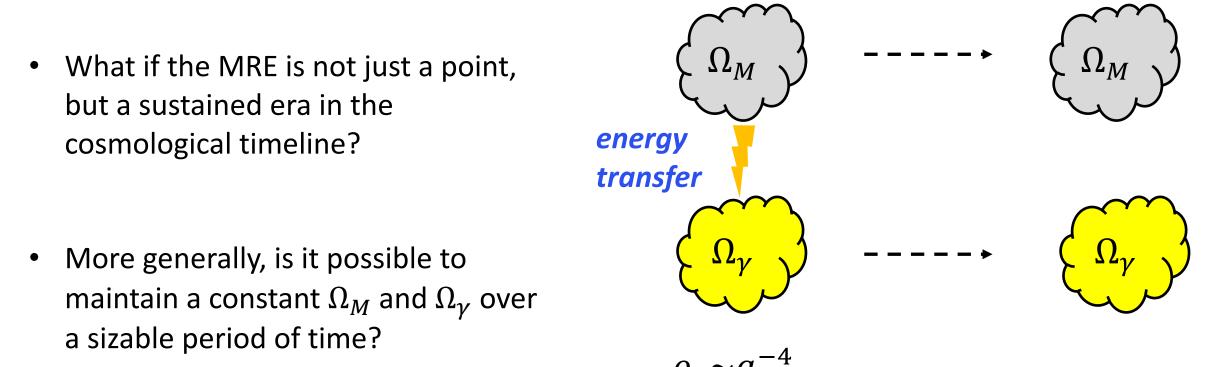
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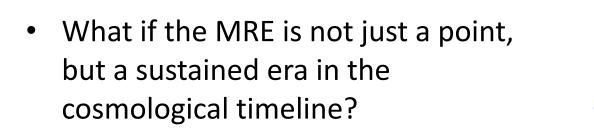
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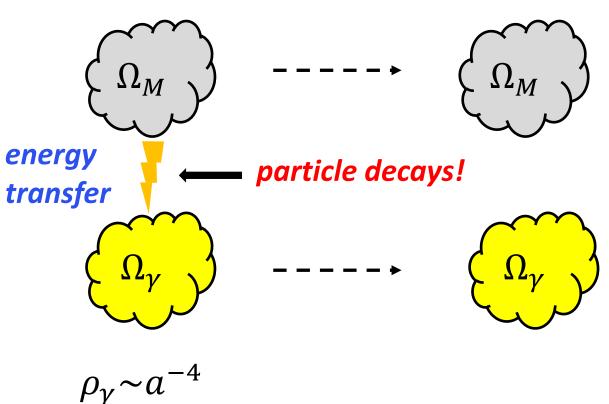
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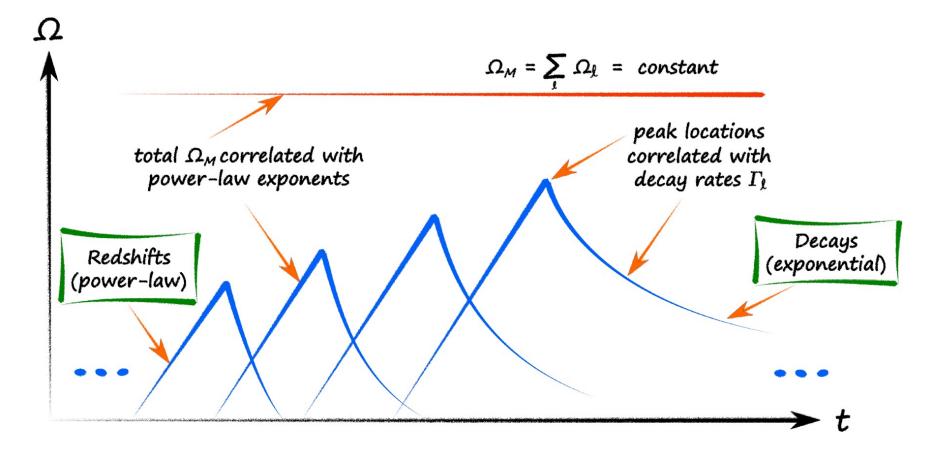


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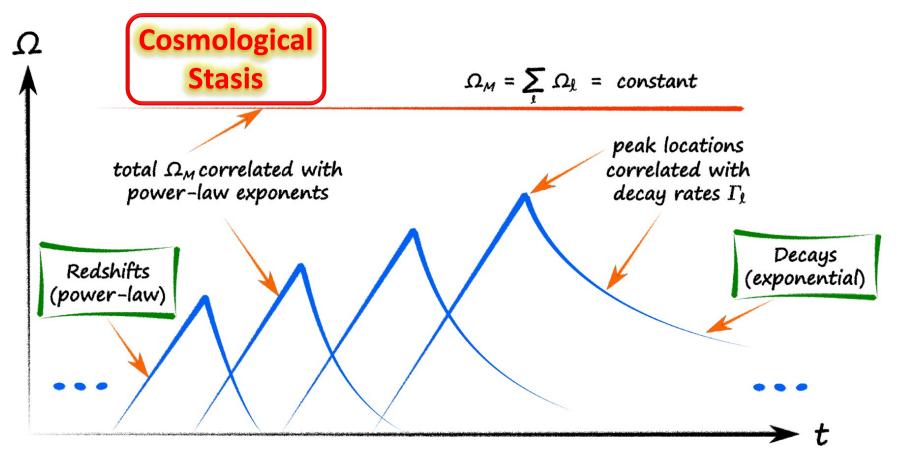


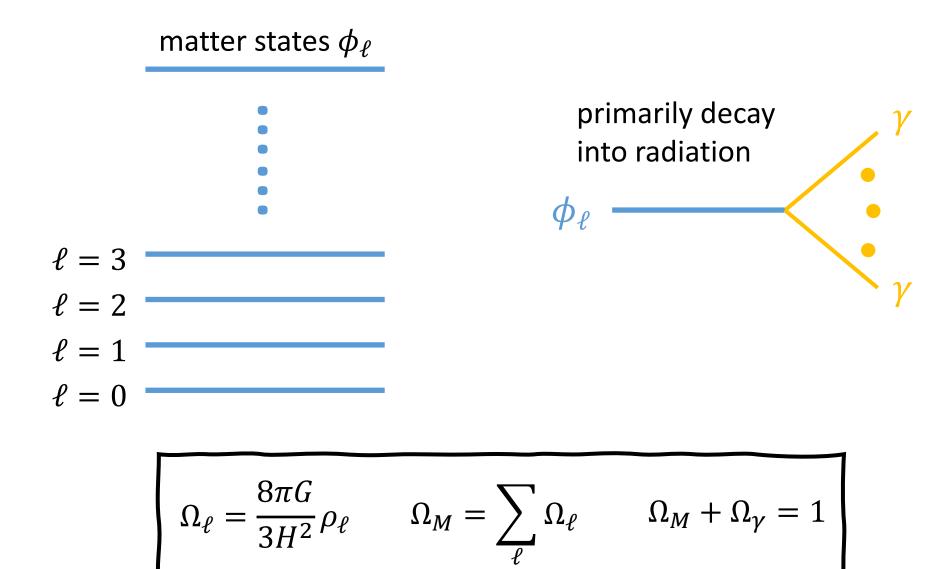
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- How about a series of decays from many matter states?



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Boltzmann equations

$$\frac{d\rho_{\ell}}{dt} = -3H\rho_{\ell} - \Gamma_{\ell}\rho_{\ell} \qquad \qquad \frac{d\rho_{\gamma}}{dt} = -4H\rho_{\gamma} + \sum_{\ell}\Gamma_{\ell}\rho_{\ell}$$

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Friedmann equation

$$H^2 = \frac{8\pi G}{3} \left(\rho_M + \rho_\gamma \right)$$

Boltzmann equations

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$$H^2 = \frac{8\pi G}{3} \left(\rho_M + \rho_\gamma \right)$$

$$\frac{d\Omega_M}{dt} = -\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} + H(\Omega_M - \Omega_M^2)$$

$$\frac{d\Omega_M}{dt} = 0$$

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} = H(\Omega_M - \Omega_M^2)$$

<u>Necessary</u> condition for Stasis, but <u>NOT</u> <u>sufficient</u>

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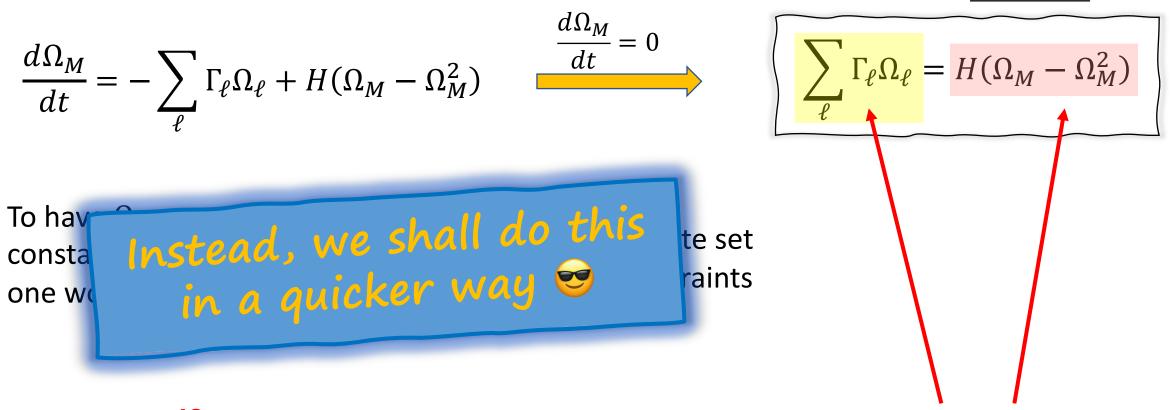
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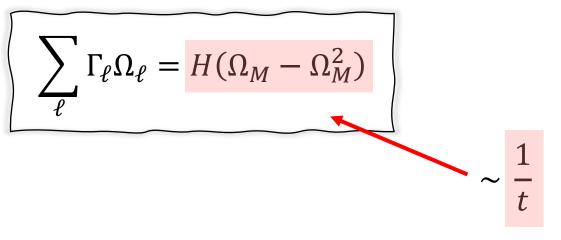
To have Ω_M remain constant in time, one would require

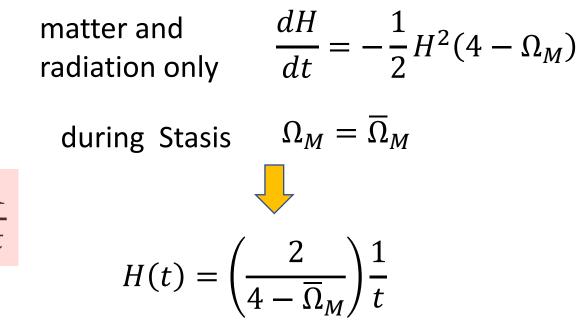
$$\left| \frac{d^n \Omega_M}{dt^n} = 0 \right| \longrightarrow \begin{array}{l} \text{an infinite set} \\ \text{of constraints} \end{array}$$

<u>Necessary</u> condition for Stasis, but not sufficient



Assuming $\frac{d\Omega_M}{dt} = 0$ is satisfied at some *fiducial time* t_* , demand that <u>both sides</u> evolve with time in the same manner.





$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} = H(\Omega_{M} - \Omega_{M}^{2})$$

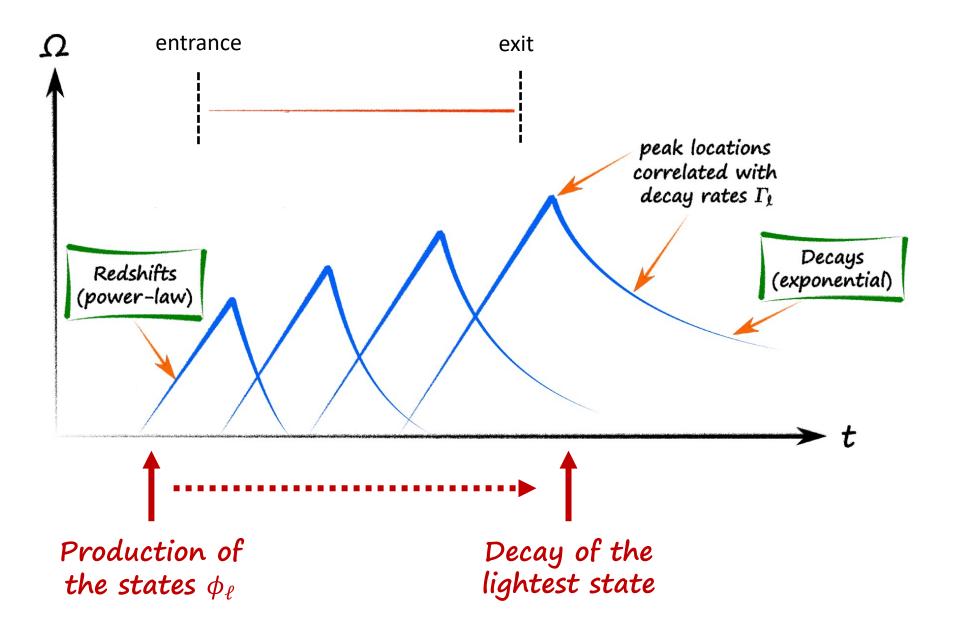
$$\sim \frac{1}{t}$$
matter and $\frac{dH}{dt} = -\frac{1}{2}H^{2}(4 - \Omega_{M})$
during Stasis $\Omega_{M} = \overline{\Omega}_{M}$

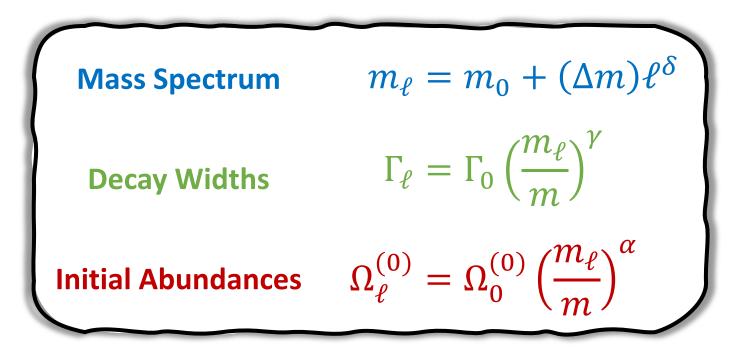
$$H(t) = \left(\frac{2}{4 - \overline{\Omega}_{M}}\right)\frac{1}{t}$$

$$\sum_{\ell} \frac{\Gamma_{\ell}}{\Omega_{\ell}(t)} = \frac{2\overline{\Omega}_{M}(1-\overline{\Omega}_{M})}{4-\overline{\Omega}_{M}} \frac{1}{t}$$
$$\sum_{\ell} \Omega_{\ell}(t) = \overline{\Omega}_{M}$$

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \overline{\Omega}_{M})}{4 - \overline{\Omega}_{M}} \frac{1}{t}$$

Stasis (Eternal)





Mass Spectrum $m_{\ell} = m_0 + (\Delta m)\ell^{\delta}$ Decay Widths $\Gamma_{\ell} = \Gamma_0 \left(\frac{m_{\ell}}{m}\right)^{\gamma}$ Initial Abundances $\Omega_{\ell}^{(0)} = \Omega_0^{(0)} \left(\frac{m_{\ell}}{m}\right)^{\alpha}$

Depends on particle physics model • KK excitations of a 5-d scalar field compactified on a circle of radius R- $\delta \sim 1$ for $mR \ll 1$

- $\delta \sim 2$ for $mR \gg 1$
- Bound states of strongly-coupled gauge theory
 - δ~1/2

Depends on decay mode

• if ϕ_{ℓ} decays to photons through contact operator $\mathcal{O}_{\ell} \sim c_{\ell} \phi_{\ell} \mathcal{F} / \Lambda^{d-4}$, $\gamma = 2d - 7$, e.g., $\gamma \sim \{3, 5, 7\}$

Depends on the production mechanism

- $\alpha < 0$ for misalignment production
- both $\alpha > 0$ or $\alpha < 0$ for thermal freeze-out
- $\alpha = 1$ for universal inflaton decay

 $m_{\ell} = m_0 + (\Delta m) \ell^{\delta}$ **Mass Spectrum** $\Gamma_{\ell} = \Gamma_0 \left(\frac{m_{\ell}}{m}\right)^{\gamma}$ **Decay Widths** $\Omega_{\ell}^{(0)} = \Omega_{0}^{(0)} \left(\frac{m_{\ell}}{-}\right)^{\alpha}$ **Initial Abundances** Free parameters $\left\{ \alpha, \gamma, \delta, m_0, \Delta m, \Gamma_0, \Omega_0^{(0)}, t^{(0)} \right\}$ initial conditions

Depends on particle physics model
KK excitations of a 5-d scalar field compactified on a circle of radius R

δ~1 for mR ≪ 1
δ~2 for mR ≫ 1

Bound states of strongly-coupled gauge theory

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Recall the condition for stasis

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \overline{\Omega}_{M})}{4 - \overline{\Omega}_{M}} \frac{1}{t}$$

With our particular model, we can work out the LHS

Recall the condition for stasis

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 $\Omega_\ell(t)$

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Production at $t^{(0)}$ (before stasis)

 $\Omega_{\ell}(t)$

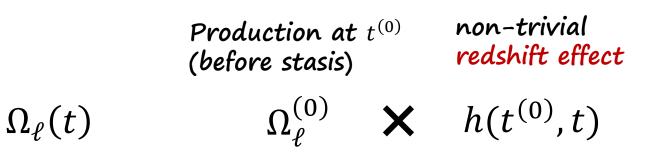
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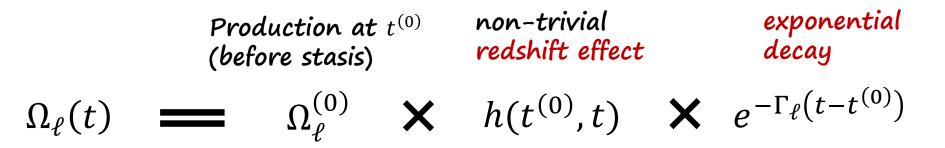
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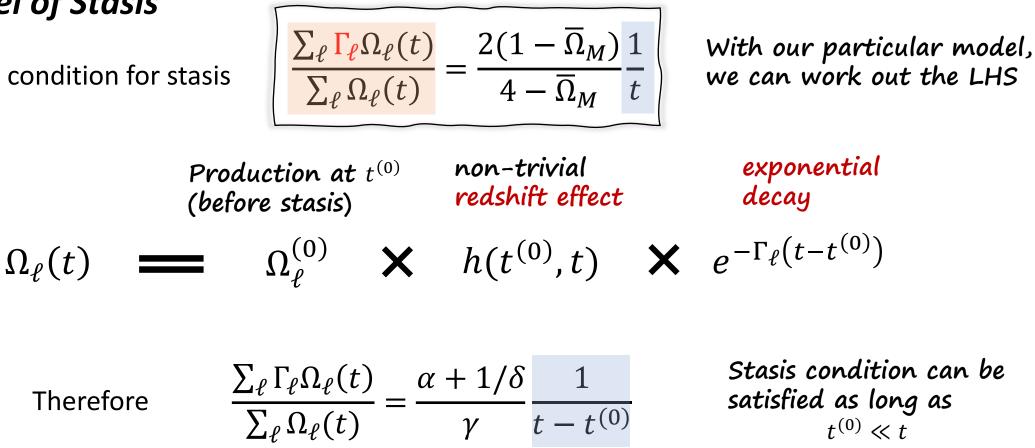
$$\Omega_{\ell}(t) = \Omega_{\ell}^{(0)} \times h(t^{(0)}, t) \times e^{-\Gamma_{\ell}(t-t^{(0)})}$$
$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t) = \Gamma_{0} \Omega_{0}^{(0)} h(t^{(0)}, t) \sum_{\ell} \left(\frac{m_{\ell}}{m_{0}}\right)^{\alpha+\gamma} e^{-\Gamma_{0}\left(\frac{m_{\ell}}{m_{0}}\right)^{\gamma} (t-t^{(0)})}$$

Continuous

Recall the

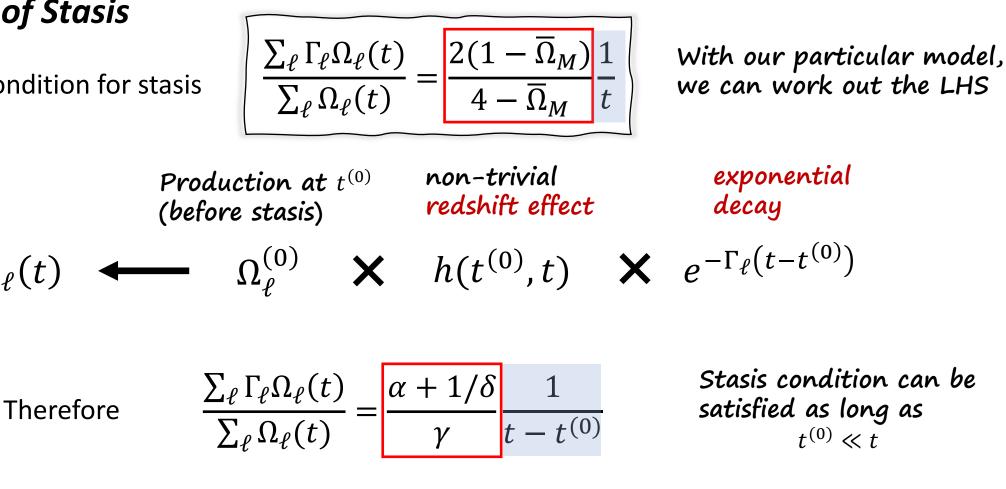
$$\begin{split} & \frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \overline{\Omega}_{M})}{4 - \overline{\Omega}_{M}} \frac{1}{t} \end{split} \qquad \text{With our particular model,} \\ & \text{we can work out the LHS} \end{aligned} \\ & \frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \overline{\Omega}_{M})}{4 - \overline{\Omega}_{M}} \frac{1}{t} \end{aligned} \qquad \text{With our particular model,} \\ & \text{we can work out the LHS} \end{aligned} \\ & \Omega_{\ell}(t) = \Pr_{0} \Omega_{0}^{(0)} h(t^{(0)}, t) \sum_{\ell} (m_{\ell})^{\alpha + \gamma} e^{-\Gamma_{0}(m_{\ell})^{\gamma}(t - t^{(0)})} \\ & \Pi_{\ell} \Omega_{\ell}(t) = \Gamma_{0} \Omega_{0}^{(0)} h(t^{(0)}, t) \sum_{\ell} \left(\frac{m_{\ell}}{m_{0}}\right)^{\alpha + \gamma} e^{-\Gamma_{0}\left(\frac{m_{\ell}}{m_{0}}\right)^{\gamma}(t - t^{(0)})} \\ & \Pi_{\ell} \Pi_{\ell}$$

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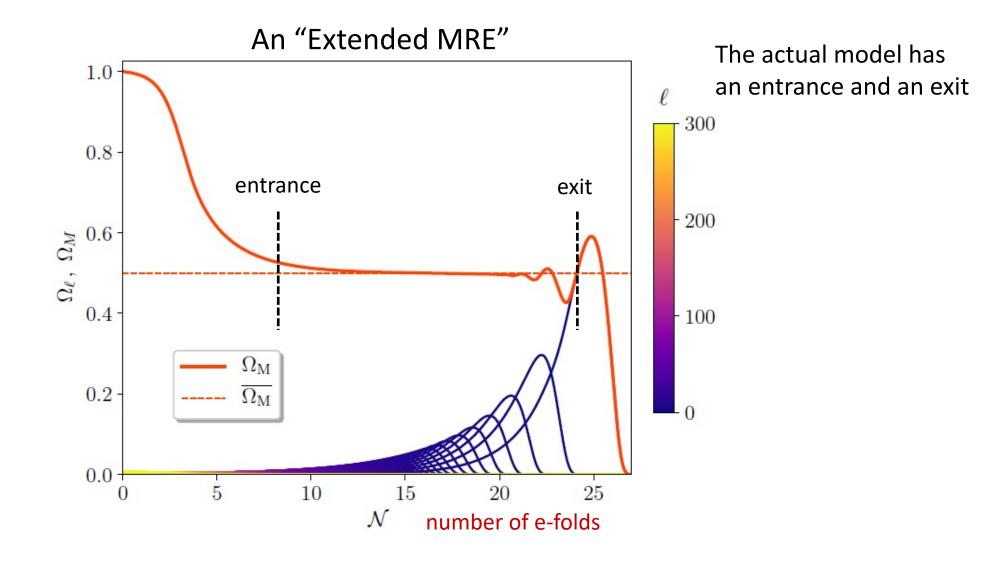


Matter abundance during stasis is determined by model parameters

$$\overline{\Omega}_M = \frac{2\gamma\delta - 4(1 + \alpha\delta)}{2\gamma\delta - (1 + \alpha\delta)}$$

It's time to test it numerically...

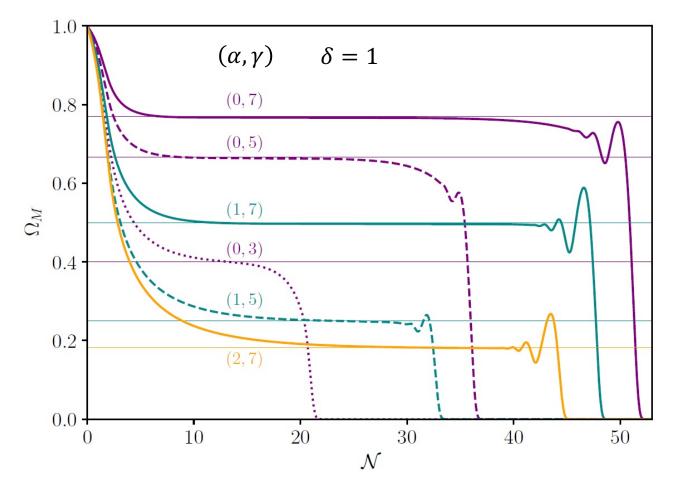
Examples



 $(\alpha, \gamma, \delta) = (1,7,1) \rightarrow \overline{\Omega}_M = 1/2, \Delta m = m_0, \Gamma_{N-1}/H^{(0)} = 0.01, N = 300$

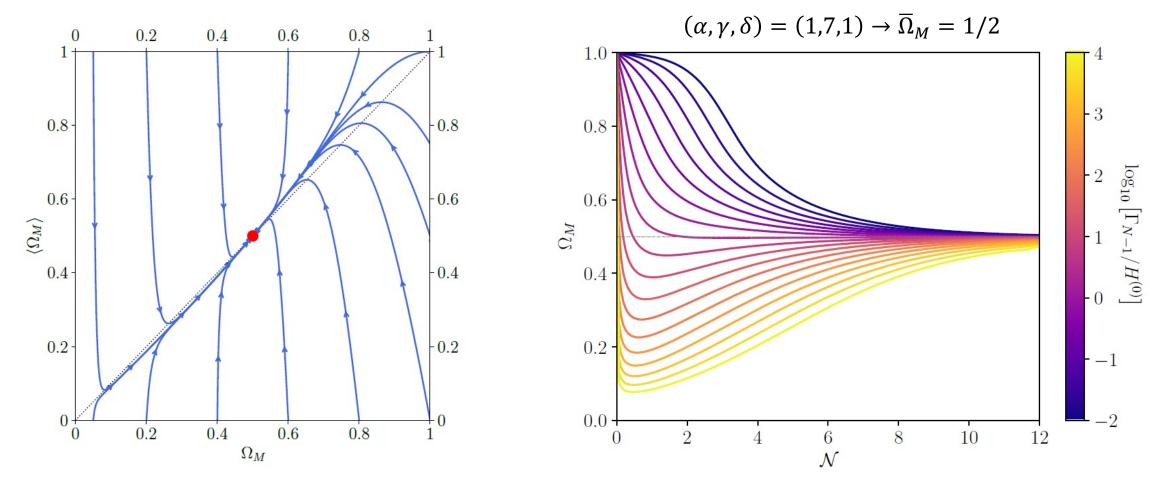
Examples

Fine tuned?

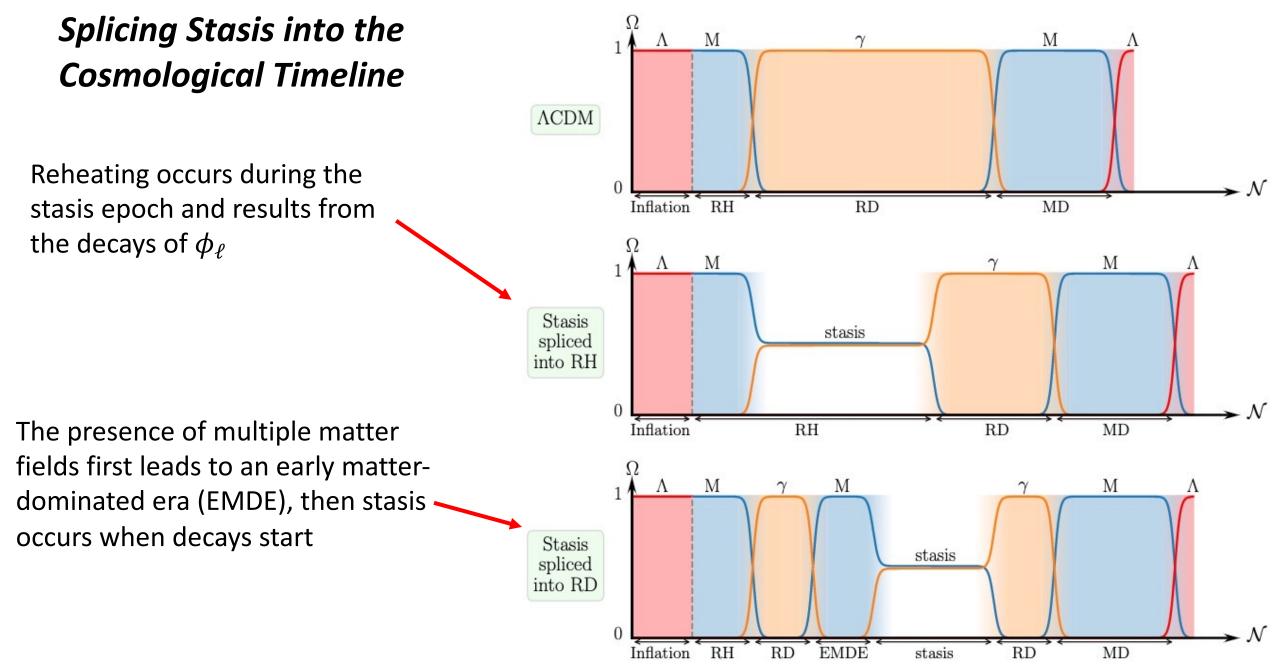


Prolonged epochs in which $\Omega_M \sim \overline{\Omega}_M$ can be achieved in multiple cases!

Stasis as a Global Attractor



The attractor is GLOBAL!



Cosmological Implications

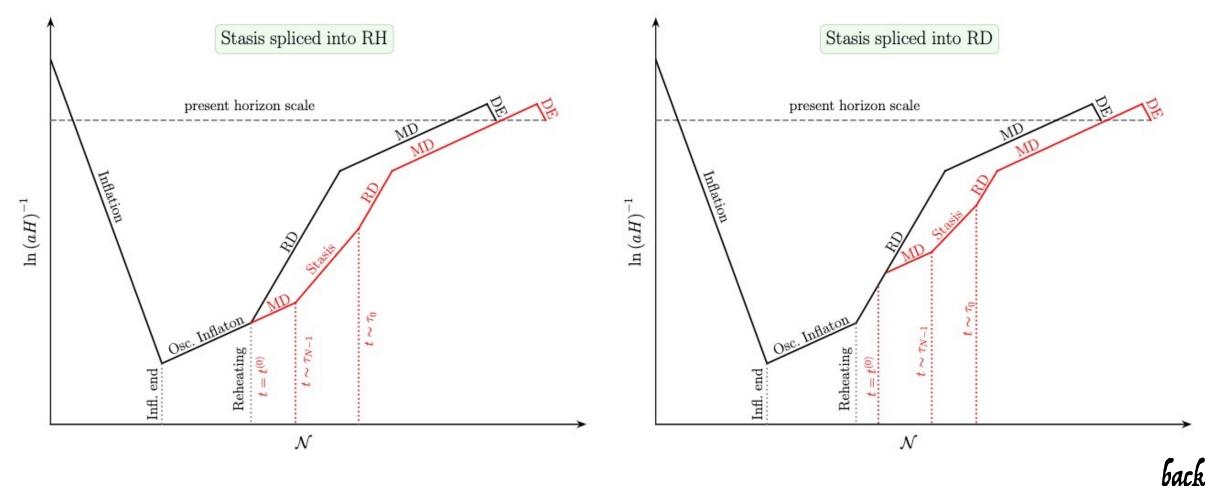
- Since the comoving Hubble radius grows more slowly if stasis exists, perturbation modes re-enter the horizon at a later time, can affect predictions for *inflationary observables*
- **Density perturbations** grow faster during stasis than in Λ CDM, can result in formation of compact objects such as PBH or compact minihalos
- **DM relic abundance** would be affected if produced prior to or during Stasis due to different expansion history and injection of entropy. DM can also be the decay products of ϕ_{ℓ} or the lightest ϕ_{ℓ} with smaller decay width and parametrically smaller initial abundance.
- Generation of lepton or baryon asymmetry would also be affected if it occurs prior to or during Stasis, or if the asymmetry is produced through ϕ_{ℓ} decays
- Numerous directions to be explored...

Summary

- Decays of a tower of states could result in an epoch in which the matter and radiation abundances are constant and can take any value.
- Such a scenario could occur naturally in well motivated models of particle physics, like KK theory and strongly-coupled gauge theories.
- Stability analysis shows the stasis solution is a global attractor, No fine-tuning needed.
- Stasis can be spliced into the cosmological timeline in various ways, leading to new possibilities in cosmology.

Cosmological Implications

The insertion of a stasis period delays the subsequent timeline relative to Λ CDM expectations



Cosmological Implications

Stasis and EMDE can place the universe on the same cosmological timeline

