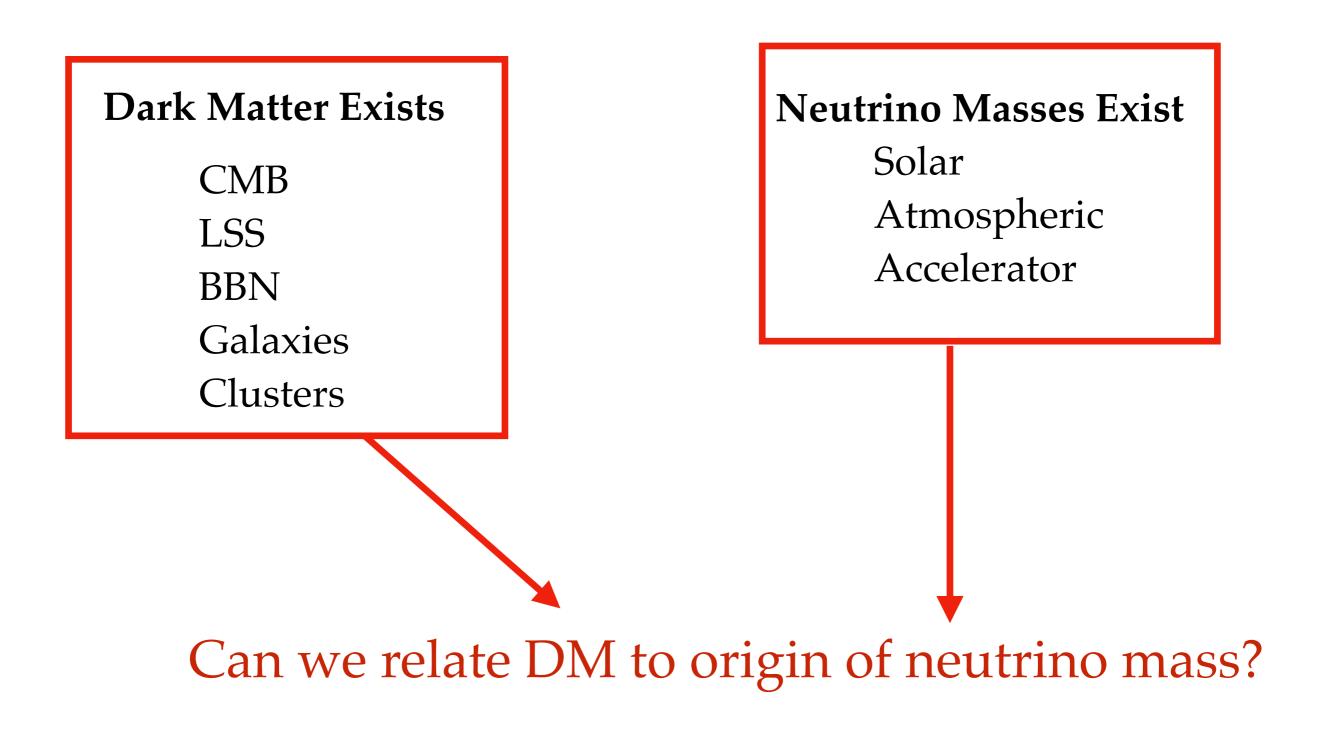
#### 



## Ultrafeeble Neutrino Interactions w/ Ultralight DM Gordan Krnjaic FNAL/UChicago

2205.06821 + Abhish Dev, Pedro Machado, Hari Ramani Mitchell 2022, May 26, 2022





Model Description

**Oscillations Regimes** 

Electron Neutrino Constraints

Mu/Tau Neutrino Variation



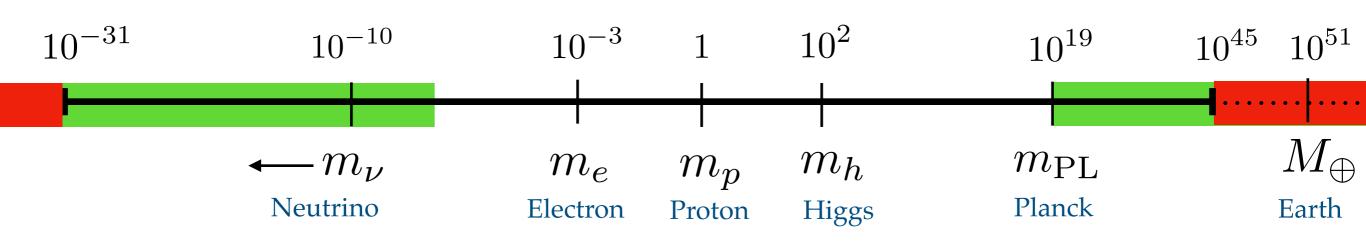
### **Model Description**

**Oscillations Regimes** 

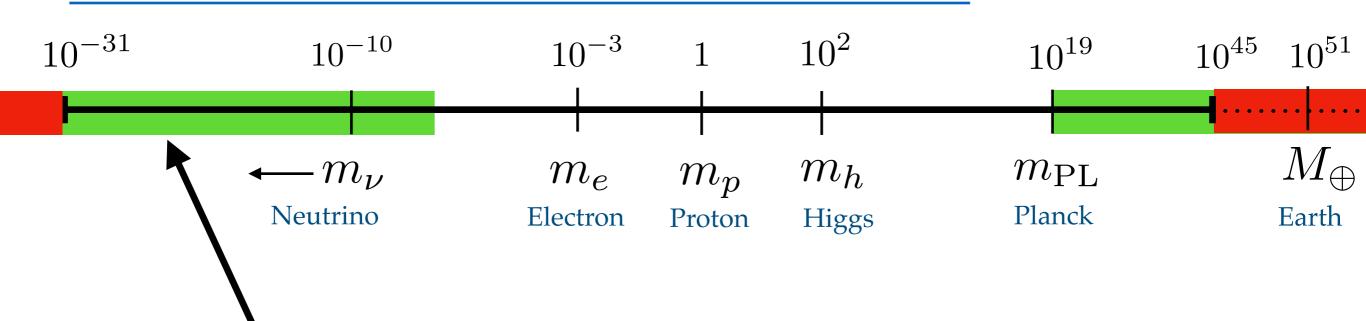
Electron Neutrino Constraints

Mu/Tau Neutrino Variation

### Huge Range of Possible DM masses

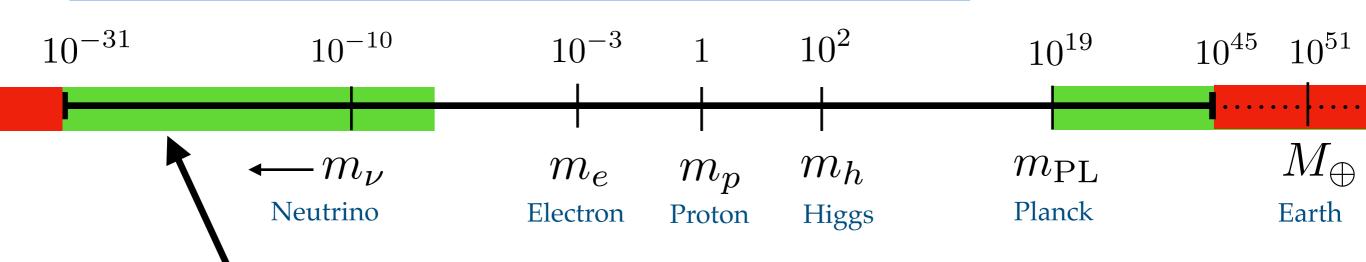


### Huge Range of Possible DM masses



Must be bosonic, feebly coupled (can't thermalize w/SM)

### Huge Range of Possible DM masses



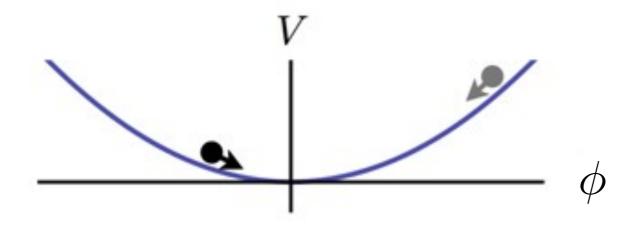
Must be bosonic, feebly coupled (can't thermalize w/SM)

de-Broglie wavelength 
$$\lambda_{\phi} = \frac{1}{m_{\phi}v_{\phi}} \approx 200 \,\mathrm{km} \left(\frac{\mathrm{neV}}{m_{\phi}}\right) \left(\frac{10^{-3}}{v_{\phi}}\right)$$

exceeds interparticle spacing ——→ classical field

### Scalar Field Cosmological Evolution

Early universe misalignment - original field value set by initial conditions



$$\ddot{\phi} + 3H\dot{\phi} + m_{\phi}^2\phi = 0$$

Equation of motion

$$\phi(\vec{r},t) = \frac{\sqrt{2\rho_{\phi}(t)}}{m_{\phi}} \cos[m_{\phi}(t+\vec{v}_{\phi}\cdot\vec{r}) + \varphi(\vec{r})]$$

Begins oscillation @ horizon crossing  $m_{\phi} \sim H$ 

Redshifts like non relativistic matter

$$\rho_{\phi} \sim m_{\phi}^2 \phi^2 \propto a^{-3}$$

### Neutrino-DM Coupling

L = 2 scalar DM induces Majorana mass for right handed neutrinos

$$\mathcal{L} \supset y_{\nu} H\ell N + \frac{y_{\phi}}{2} \phi NN + h.c.$$

Post EWSB: static Dirac mass and dynamical Majorana mass

$$m_D = \frac{y_\nu v}{\sqrt{2}} \quad , \quad m_M = \frac{y_\phi}{2}\phi(t)$$

### Neutrino-DM Coupling

L = 2 scalar DM induces Majorana mass for right handed neutrinos

$$\mathcal{L} \supset y_{\nu} H\ell N + \frac{y_{\phi}}{2} \phi NN + h.c.$$

Post EWSB: static Dirac mass and dynamical Majorana mass

$$m_D = \frac{y_{\nu}v}{\sqrt{2}} , \quad m_M = \frac{y_{\phi}}{2}\phi(t)$$

Small Majorana mass = pseudo-Dirac fermion

$$m_{h,\ell}^2 = m_D^2 \pm m_D m_M \equiv m_\nu^2 \pm \frac{1}{2} \delta m^2$$

DM density dependent splitting between heavy/light eigenstates

Active and sterile states (one pair per generation)

$$|\nu_e\rangle = \frac{1}{\sqrt{2}} (|\nu_h\rangle + |\nu_\ell\rangle)$$
$$|\nu_s\rangle = \frac{1}{\sqrt{2}} (|\nu_h\rangle - |\nu_\ell\rangle)$$

Near maximal mixing

$$\tan\left(2\theta\right) = \frac{2m_D}{m_M} \gg 1$$

Active and sterile states (one pair per generation)

$$|\nu_e\rangle = \frac{1}{\sqrt{2}} (|\nu_h\rangle + |\nu_\ell\rangle)$$
 Near maximal mixing  
$$|\nu_s\rangle = \frac{1}{\sqrt{2}} (|\nu_h\rangle - |\nu_\ell\rangle)$$
 
$$\tan(2\theta) = \frac{2m_D}{m_M} \gg 1$$

Neutrino oscillations between active/sterile  $\nu_a \rightarrow \nu_s$ 

Survival prob 
$$P_{ee}(t) = |\langle \nu(t) | \nu_e \rangle|^2 = \cos^2 \left( \frac{1}{4E_{\nu}} \int_0^t dt' \delta m^2(t') \right)$$

Analogy with 2-flavor oscillation in more familiar constant-mass case

$$P(\nu_a \to \nu_s) = \sin^2(2\theta) \sin^2\left(\frac{\delta m^2 L}{4E_\nu}\right)$$

Density dependent active-sterile oscillation

$$1 - P_{ee} = \sin^2 \left\{ \frac{m_D}{2E_\nu} \frac{y_\phi \sqrt{2\rho_\phi}}{m_\phi^2} \left( \sin\left[m_\phi t + \varphi\right] - \sin\varphi \right) \right\}$$

Defines effective mass for comparison with reported limits

$$\delta m_{\rm eff}^2 \equiv \frac{2y_{\phi}m_D}{m_{\phi}}\sqrt{2\rho_{\phi}}$$



Model Description

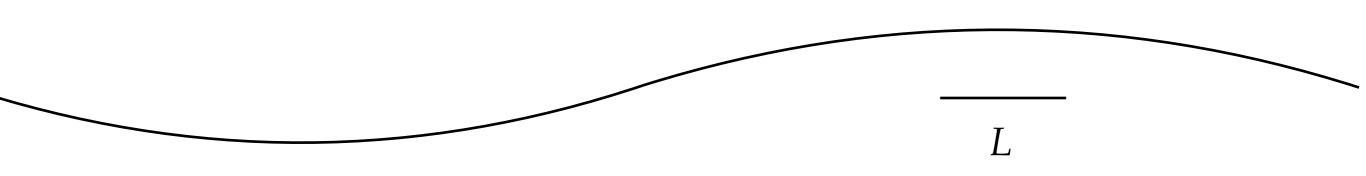
**Oscillations Regimes** 

**Electron Neutrino Constraints** 

Mu/Tau Neutrino Constraints

### Constant $\phi$ : $m_{\phi}L \lesssim 1$

Neutrino encounters same offset phase across full trajectory



### Constant $\phi$ : $m_{\phi}L \lesssim 1$

Neutrino encounters same offset phase across full trajectory

$$1 - P_{ee} \approx \sin^2 \left( \frac{L}{4E_{\nu}} \frac{2y_{\phi}m_D}{m_{\phi}} \sqrt{2\rho_{\phi}} \cos \varphi \right)$$

. **\_** 

### Constant $\phi$ : $m_{\phi}L \lesssim 1$

Neutrino encounters same offset phase across full trajectory

$$1 - P_{ee} \approx \sin^2 \left( \frac{L}{4E_{\nu}} \frac{2y_{\phi}m_D}{m_{\phi}} \sqrt{2\rho_{\phi}} \cos \varphi \right)$$

Nontrivial effect when argument is order-one

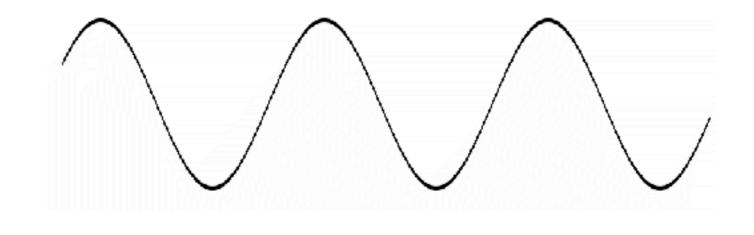
$$\delta m_{\text{eff}}^2 \equiv \frac{2y_{\phi}m_D}{m_{\phi}}\sqrt{2\rho_{\phi}} \quad \longrightarrow \quad y_{\phi} < \frac{m_{\phi}}{2m_D}\frac{\delta m_{\text{lim}}^2}{\sqrt{2\rho_{\phi}}},$$

Combination behaves like an effective mass-squared difference Can be matched with existing limits on pseudo-dirac neutrinos

### **Oscillations Regimes**

### Modulating $\phi$ : $(m_{\phi}v_{\phi}L < 1 \ll m_{\phi}L)$

# Neutrino encounters many cycles of scalar modulation ... But all in the **same phase** domain



L

### Modulating $\phi$ : $(m_{\phi}v_{\phi}L < 1 \ll m_{\phi}L)$

Neutrino encounters many cycles of scalar modulation ... But all in the **same phase** domain

Survival probability is time averaged along trajectory

$$\langle 1 - P_{ee} \rangle \approx \sin^2 \left( \frac{y_{\phi} m_D}{2E_{\nu} m_{\phi}^2} \sqrt{2\rho_{\phi}} \right)$$

Modulating 
$$\phi$$
:  $(m_{\phi}v_{\phi}L < 1 \ll m_{\phi}L)$ 

Neutrino encounters many cycles of scalar modulation ... But all in the **same phase** domain

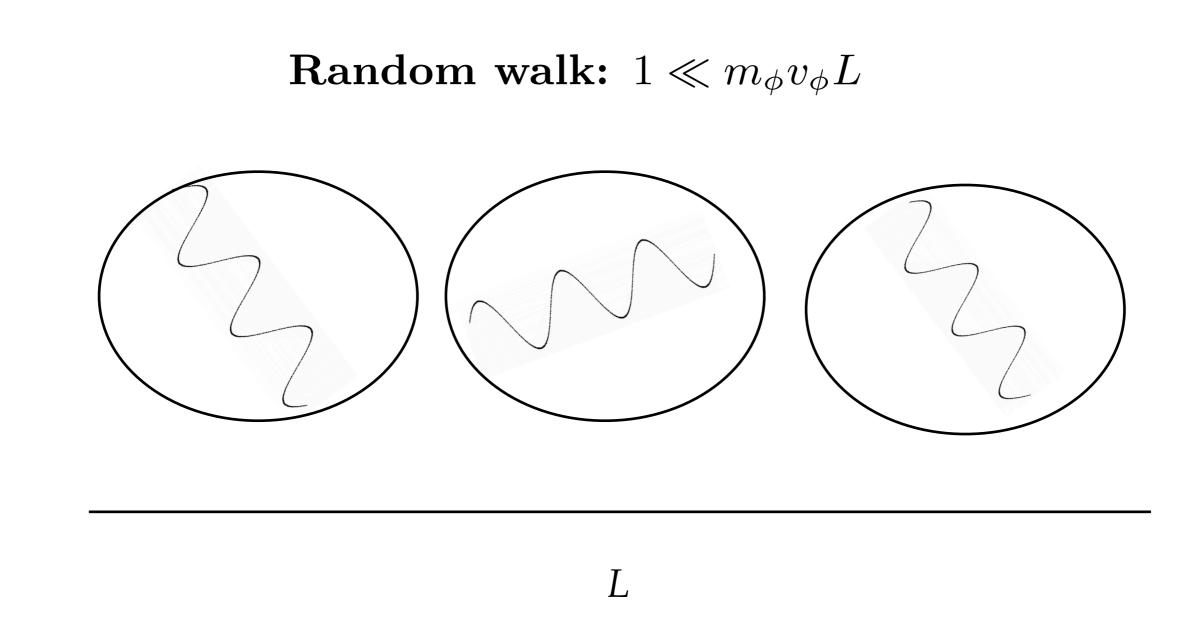
Survival probability is time averaged along trajectory

$$\langle 1 - P_{ee} \rangle \approx \sin^2 \left( \frac{y_{\phi} m_D}{2E_{\nu} m_{\phi}^2} \sqrt{2\rho_{\phi}} \right)$$

Demand order one argument for nontrivial active-sterile oscillation

$$y_{\phi}^{\lim} = \frac{\delta m_{\lim}^2 m_{\phi}^2 L}{2m_D \sqrt{2\rho_{\phi}}}$$

### **Oscillations Regimes**



Neutrino encounters many different scalar domains

— random walk across phase variation

**Random walk:**  $1 \ll m_{\phi} v_{\phi} L$ 

Effective phase along trajectory  $\varphi_{\rm eff} \sim \sqrt{m_{\phi} v_{\phi} L}$ 

**Random walk:**  $1 \ll m_{\phi} v_{\phi} L$ 

Effective phase along trajectory  $\varphi_{\rm eff} \sim \sqrt{m_{\phi} v_{\phi} L}$ 

Path averaged survival probability

$$\langle 1 - P_{ee} \rangle \approx \sin^2 \left( \frac{y_{\phi} m_D \sqrt{2\rho_{\phi} v_{\phi} L}}{2E_{\nu} m_{\phi}^{3/2}} \right)$$

Demand order one argument for nontrivial effect

$$y_{\phi}^{\rm lim} = \frac{\delta m_{\rm lim}^2}{m_D} \sqrt{\frac{m_{\phi}^3 L}{2\rho_{\phi} v_{\phi}}}.$$



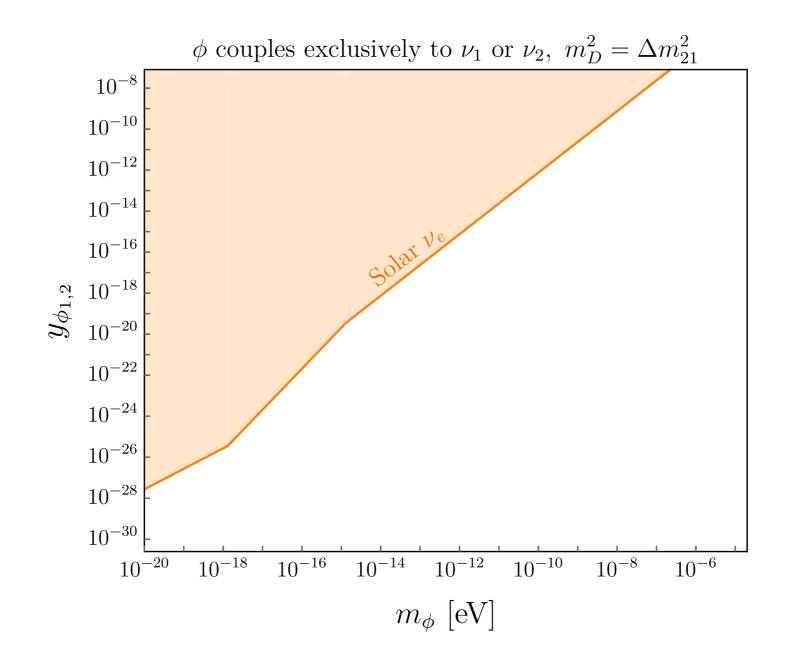
Model Description

**Oscillations Regimes** 

**Electron Neutrino Constraints** 

Mu/Tau Neutrino Constraints

 $\nu_e \rightarrow \nu_s$  oscillations on baseline  $L = 1.5 \times 10^8$  km



Impose limit based on  $\delta m_{
m lim}^2 < 10^{-12} \, {\rm eV}^2$  de Gouvea, Huang, Jenkins 0906.1611

 $\nu_e \rightarrow \nu_s$  oscillations **before** neutrino decoupling add to Neff Faster expansion rate  $\longrightarrow$  less primordial helium

 $\nu_e \rightarrow \nu_s$  oscillations after neutrino decoupling reduce  $\nu_e$ earlier n/p freeze out  $\longrightarrow$  less primordial helium

### **BBN/CMB Electron Neutrino Oscillations**

Dodelson Widrow sterile neutrino production

Solve for ratio of momentum moments

$$r_{\beta} \equiv \frac{\langle p^{\beta} \rangle_s}{\langle p^{\beta} \rangle_a}$$

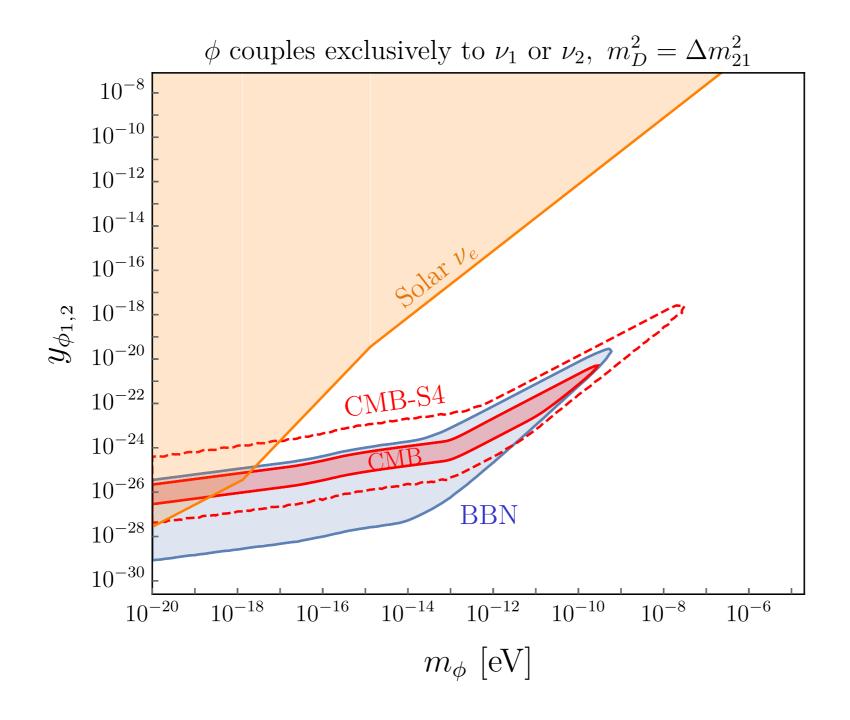
$$\frac{dr_{\beta}}{dT} = -\frac{1}{2HT\langle p^{\beta}\rangle_a} \int \frac{d^3p}{(2\pi)^3} \frac{p^{\beta}\Gamma\sin^2(2\theta_M)}{e^{p/T}+1}$$

Vacuum mixing angle  $\theta_0 = \tan^{-1} \left( \frac{y_{\phi} \sqrt{2\rho_{\phi}}}{m_D m_{\phi}} \right)$ Effective mixing angle in vacuum

$$\sin^{2}(2\theta_{M}) = \frac{\sin^{2}(2\theta_{0})}{\left[\cos(2\theta_{0}) - 2pV_{\text{eff}}/\Delta m^{2}\right]^{2} + \sin^{2}(2\theta_{0})}$$

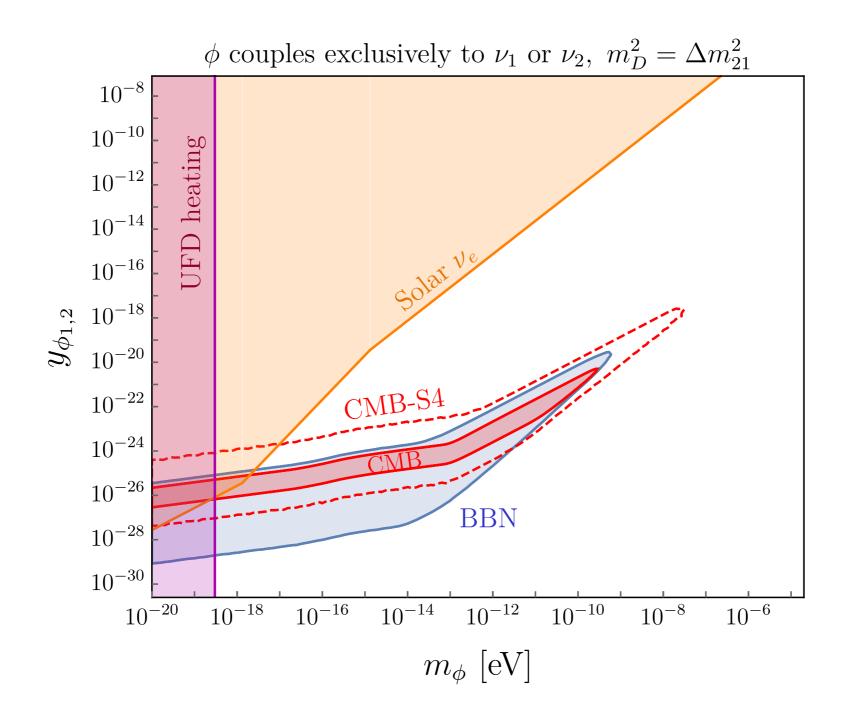
Dodelson, Widrow 9303287 Dolgov, Villante 0308083

### **BBN/CMB Electron Neutrino Oscillations**



Effective mixing angle in vacuum

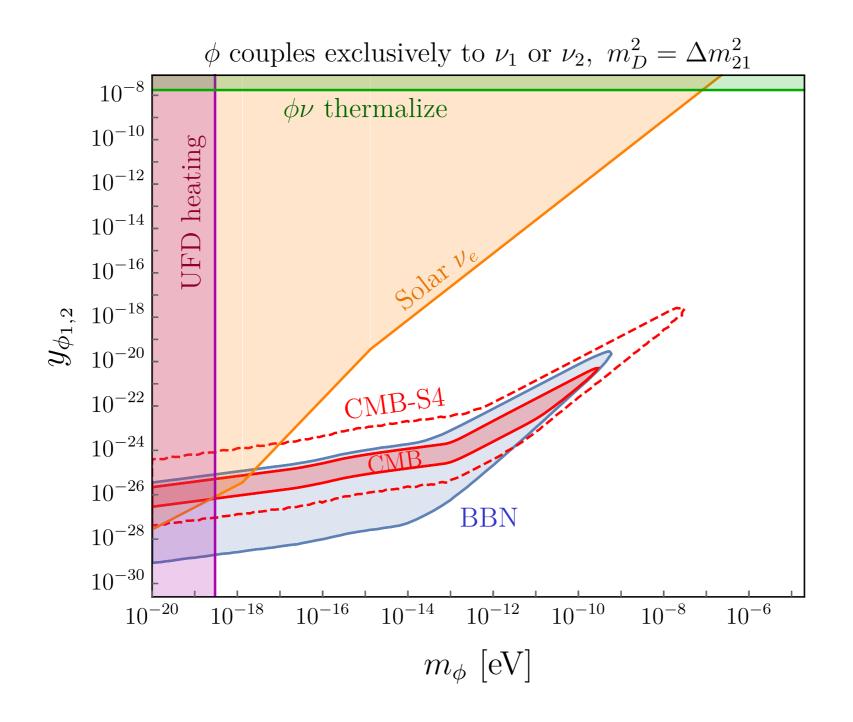
### Fuzzy Dark Matter Bounds



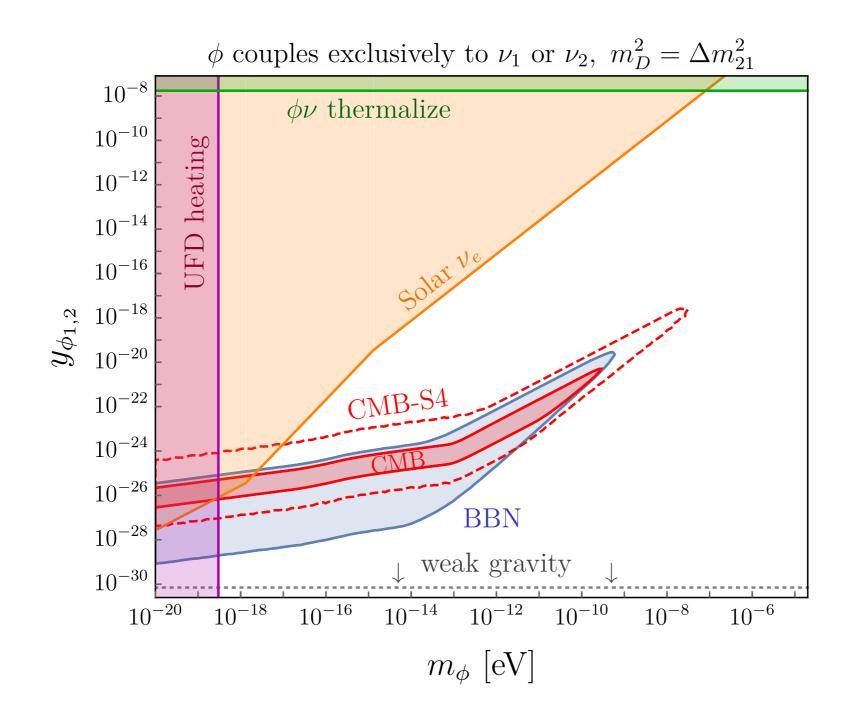
Potential fluctuations in ultra faint dwarfs affect velocity dispersion of visible stars

Hu, Barnana, Gruzinov 0003365 Dalal, Kravtsov, 2203.05750

### Thermalization w/ SM Plasma



### Compare to Gravity between Neutrinos



Low scalar mass bounds comparable to gravity between neutrinos



Model Description

**Oscillations Regimes** 

Electron Neutrino Constraints

**Mu/Tau Neutrino Variation** 

BBN/CMB simpler No electron neutrino depletion after decoupling Dolgov, Villante 0308083 **BBN/CMB simpler** 

No electron neutrino depletion after decoupling Dolgov, Villante 0308083

### **Atmospheric oscillations beat solar**

Deficit of muon neutrinos

de Gouvea, Huang, Jenkins 0906.1611

**BBN/CMB simpler** 

No electron neutrino depletion after decoupling Dolgov, Villante 0308083

### **Atmospheric oscillations beat solar**

Deficit of muon neutrinos

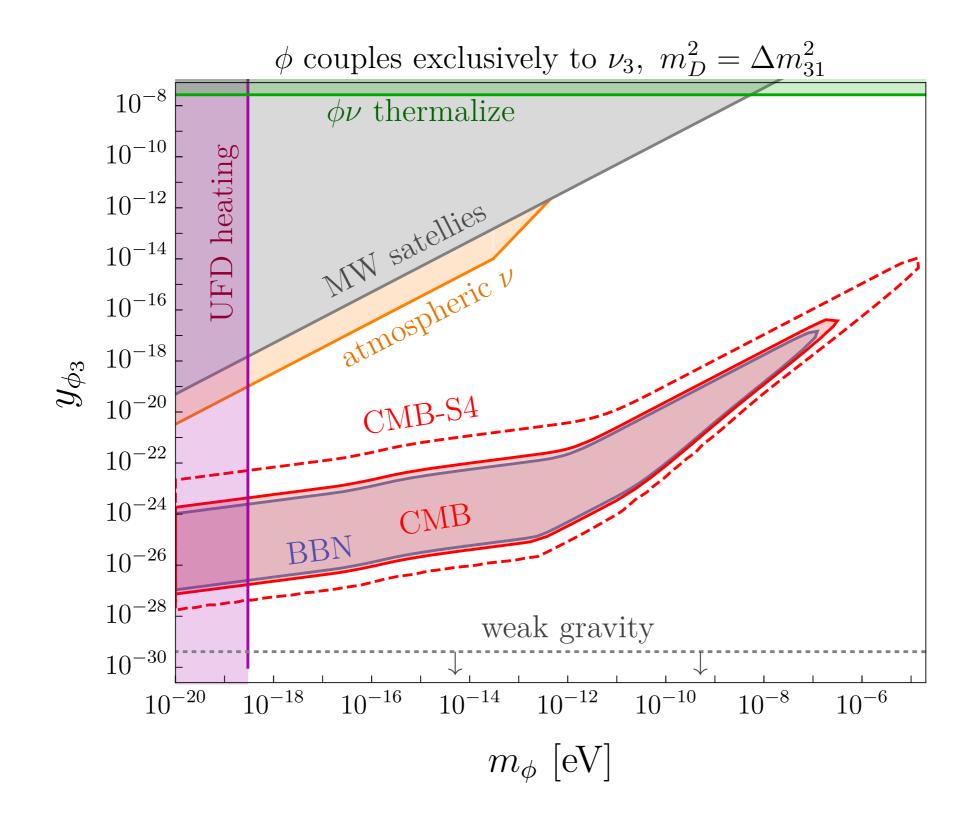
de Gouvea, Huang, Jenkins 0906.1611

### **Structure formation**

Scalar must redshift like matter after  $z \sim 10^6$ Quadratic term must dominate in potential Bound from Milky Way satellites

Das, Nadler 2010.01137

### Mu/Tau Neutrino Oscillations



Ultralight DM induced Majorana mass —> pseudo-Dirac neutrinos

Active/sterile mass splitting time/density dependent

Strong bounds from lab/astro/cosmo

Solar/atmospheric oscillations Milky Way satellites BBN/CMB/Neff Ultra Faint Dwarfs