

# Low-energy Neutrino-nucleus Interactions

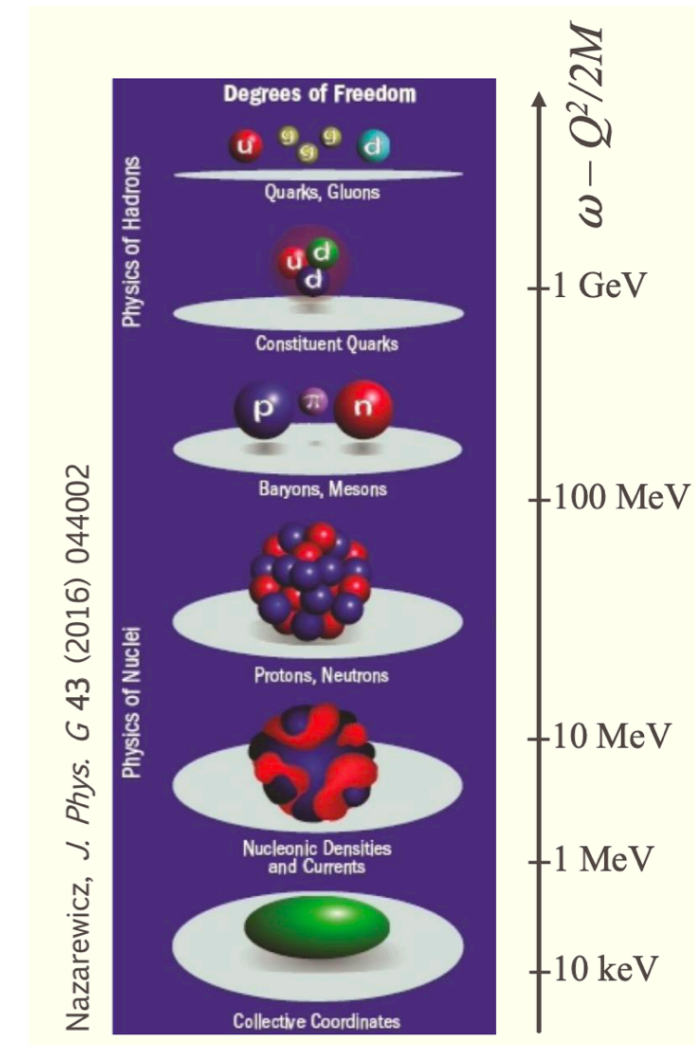
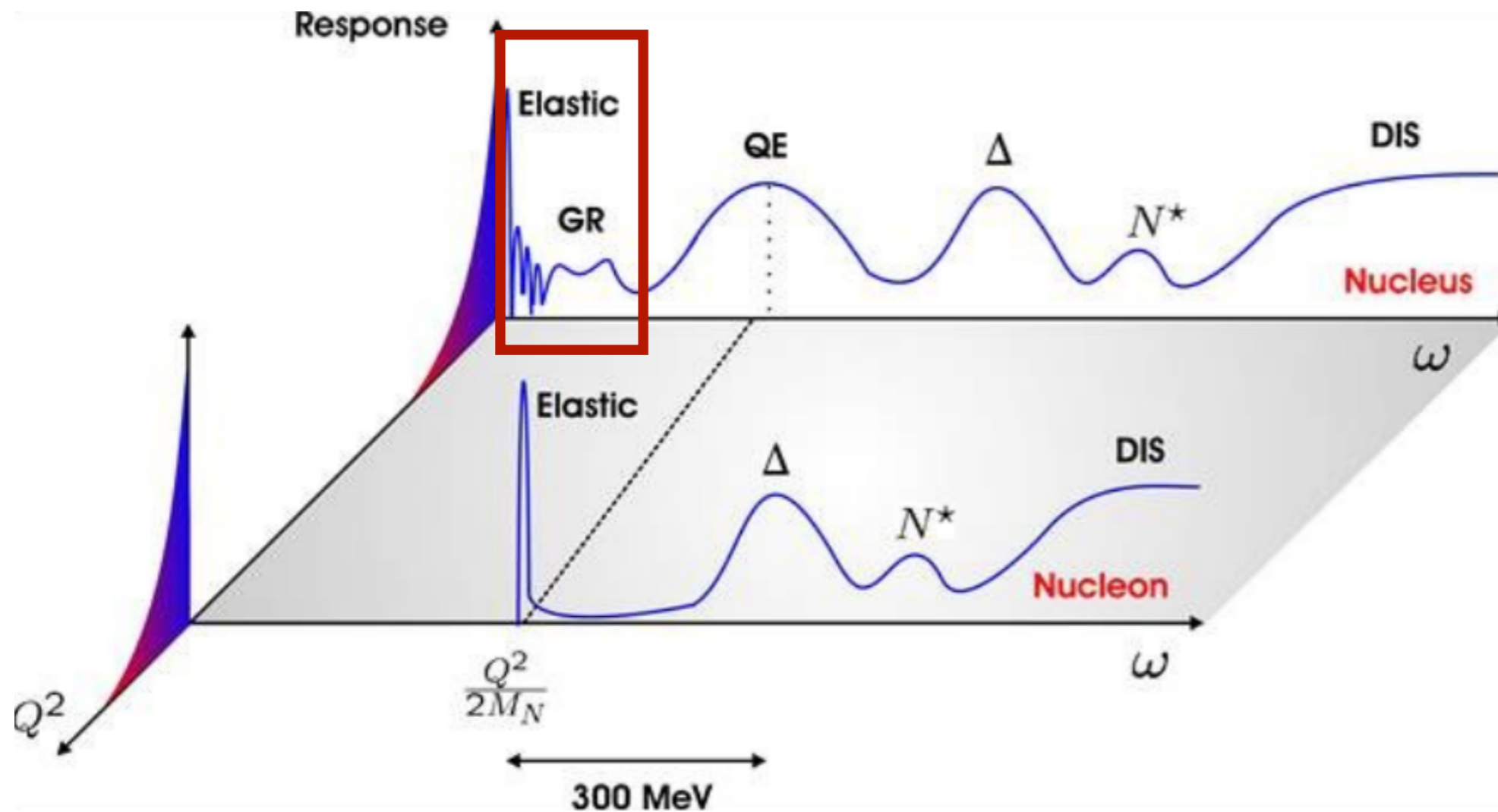
**Vishvas Pandey**



The Mitchell Conference on Collider, Dark Matter, and Neutrino Physics  
May 24 - 27, 2022, College Station, Texas

# Low-energy Neutrino Interactions

◆ Low-energy  $\approx 10$ s of MeV ( $E_\nu$  and/or  $\omega$ )

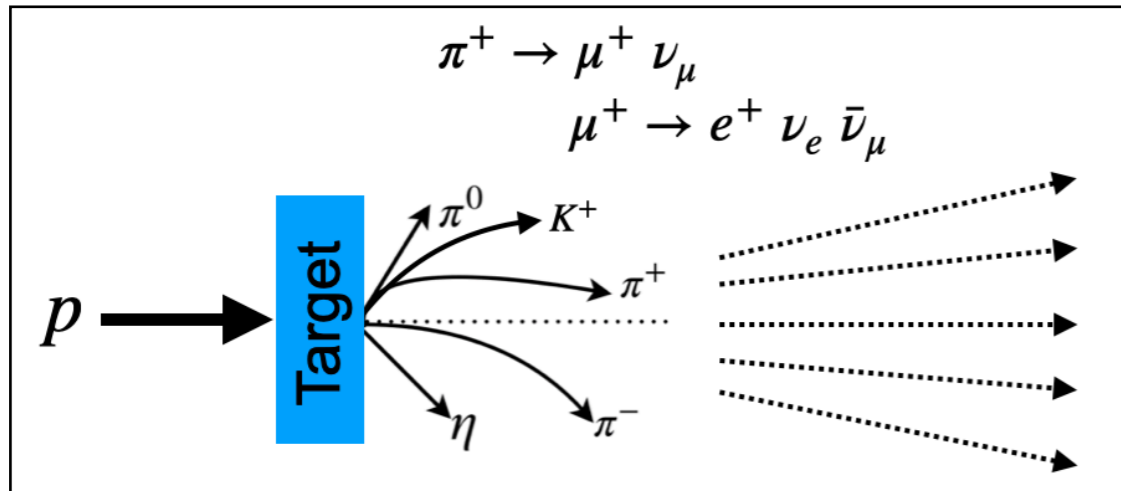


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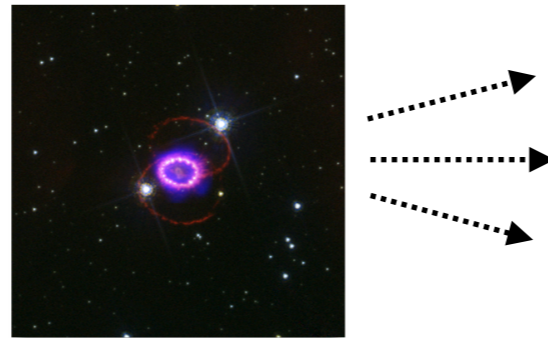
### ■ Pion decay-at-rest (piDAR) Neutrinos

(SNS at ORNL, LANSCE at LANL, MLF at JPARC, FNAL, ...)

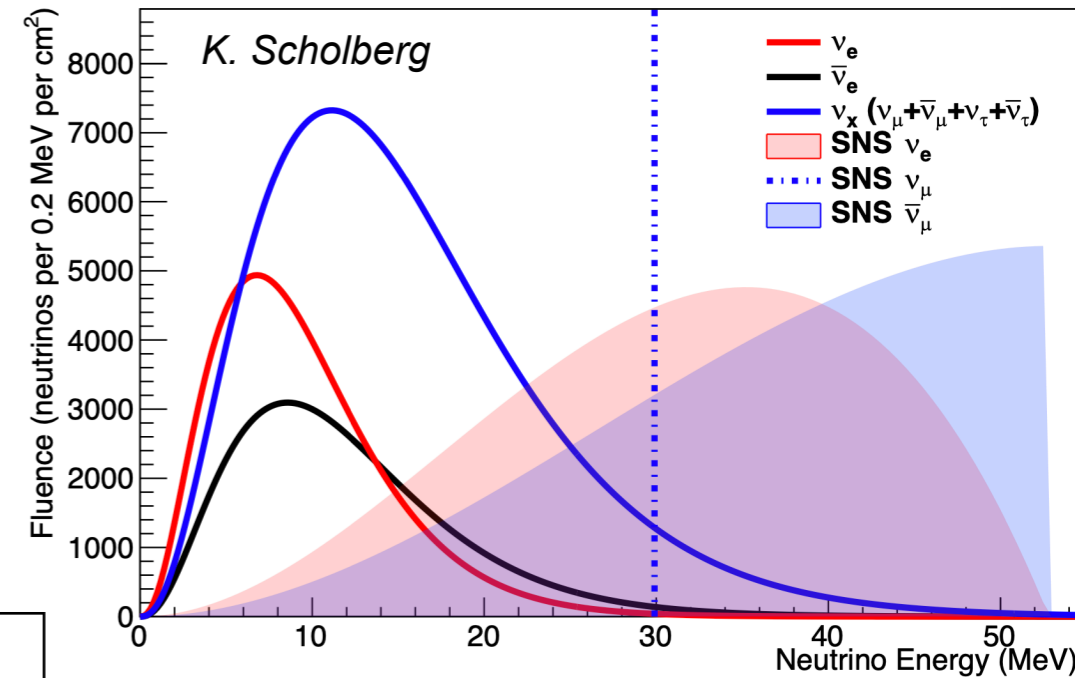


### ■ Core-collapse Supernova Neutrinos

A short, sharp “neutronization” (or “breakout”) burst primarily composed of  $\nu_e$  from  $e^- + p \rightarrow \nu_e + n$ .



piDAR and Supernova Neutrinos

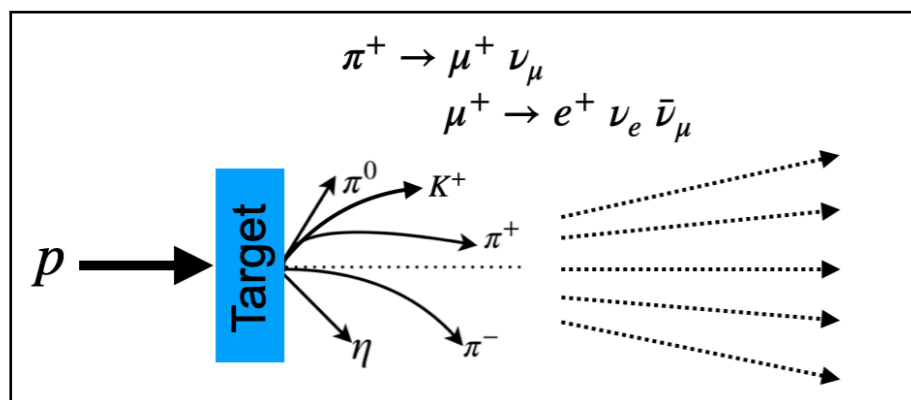


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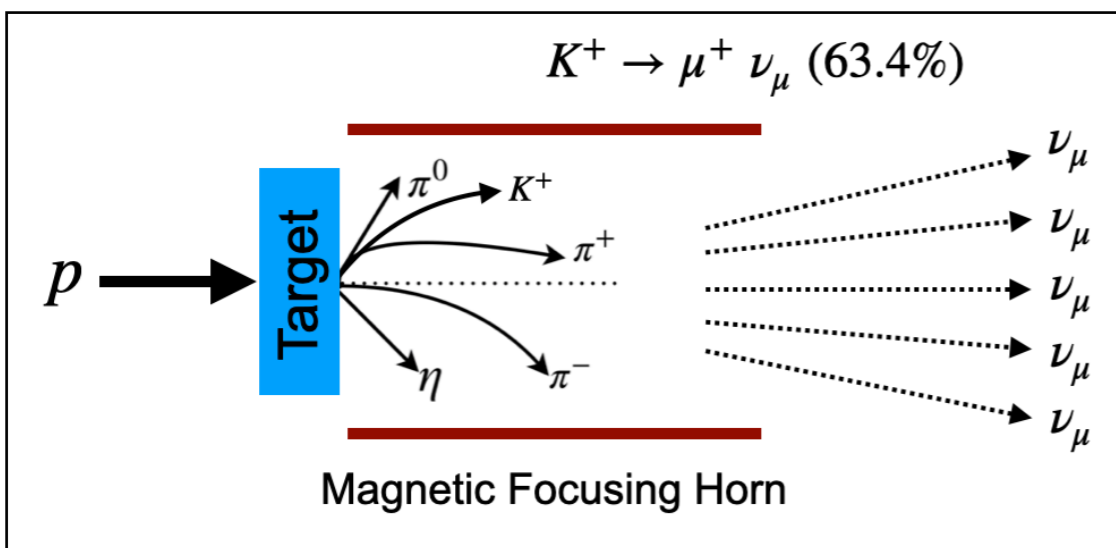


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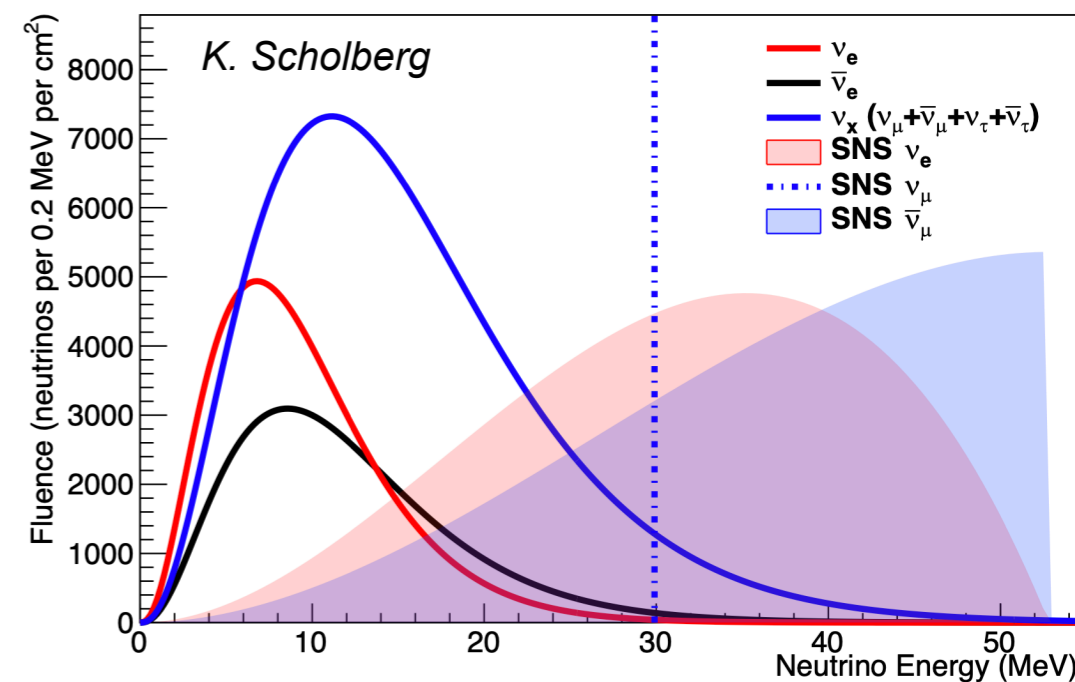
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### ■ Kaon decay-at-rest (KDAR) Neutrinos

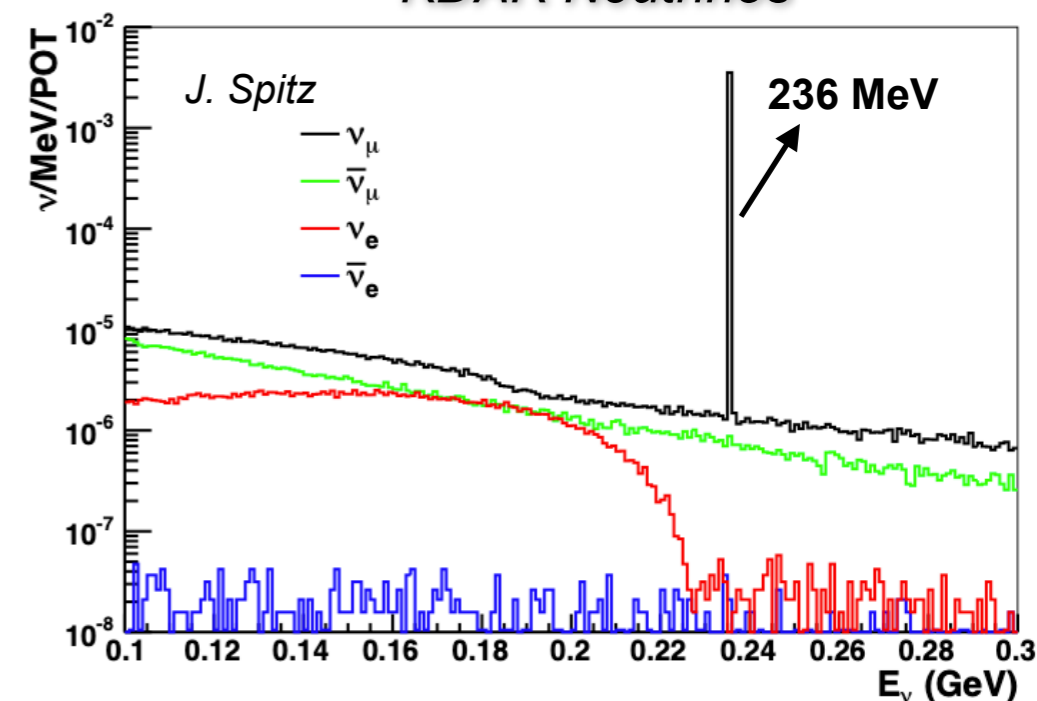
(NuMI at FNAL, MLF at JPARC, ...)



piDAR and Supernova Neutrinos



KDAR Neutrinos



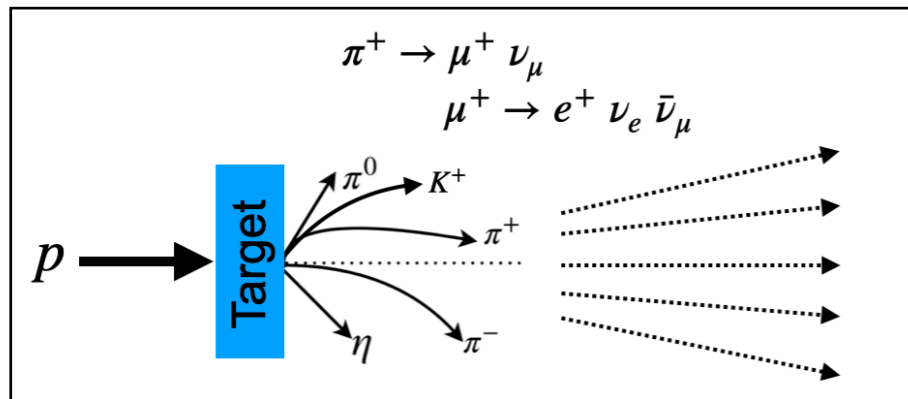


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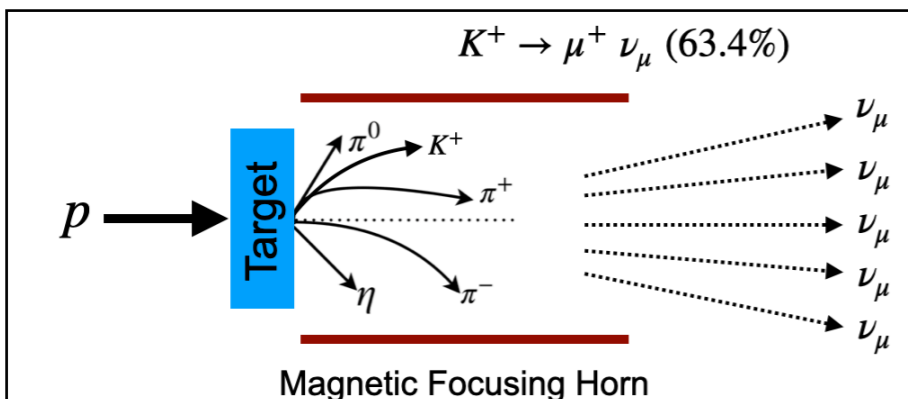


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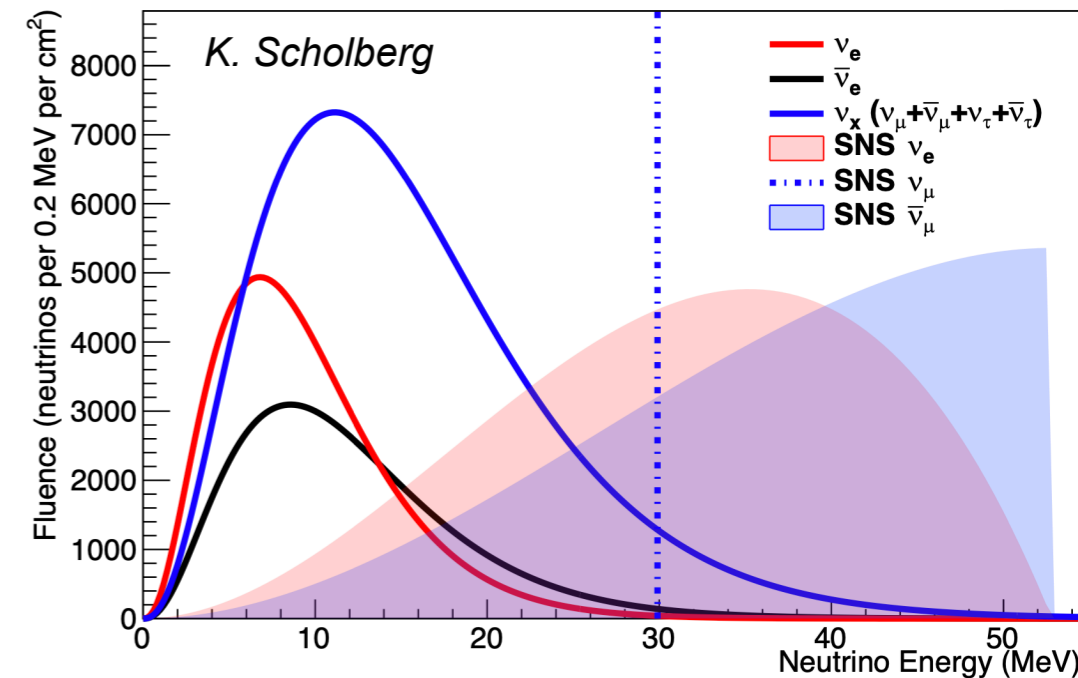
(NuMI at FNAL, MLF at JPARC, ...)



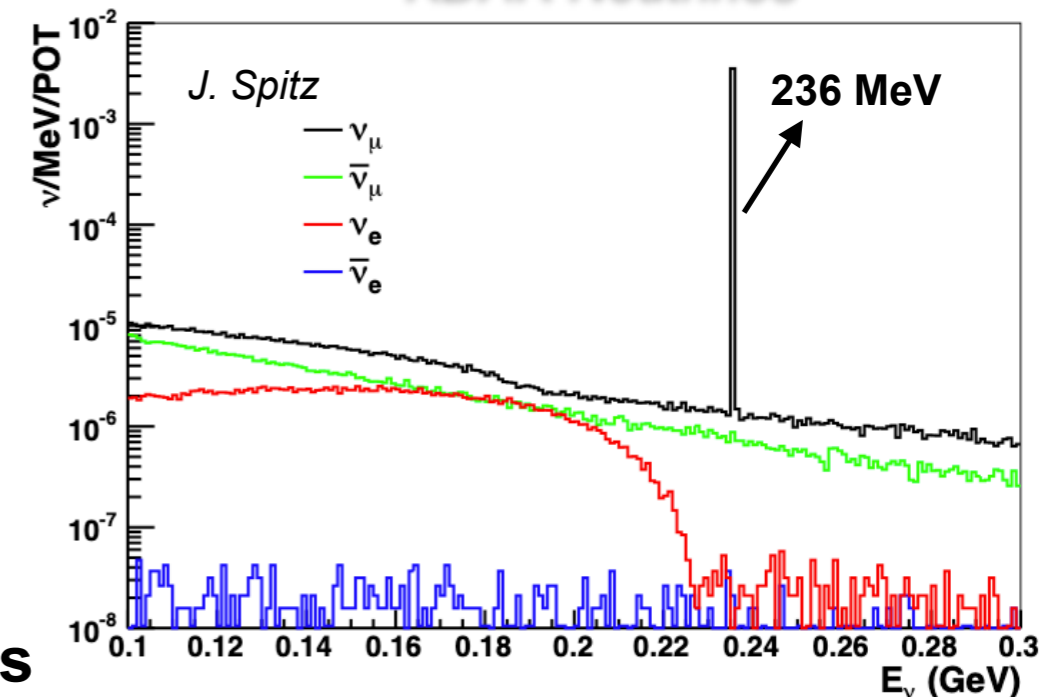
### ■ Forward Scattering of decay-in-flight (DIF) Neutrinos

(BNB/NuMI at FNAL, JPARC, ...)

piDAR and Supernova Neutrinos

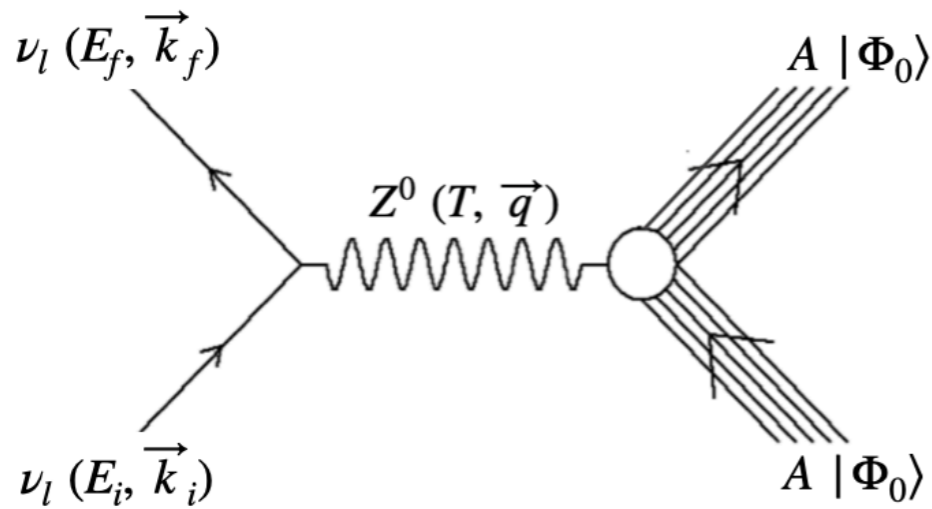


KDAR Neutrinos



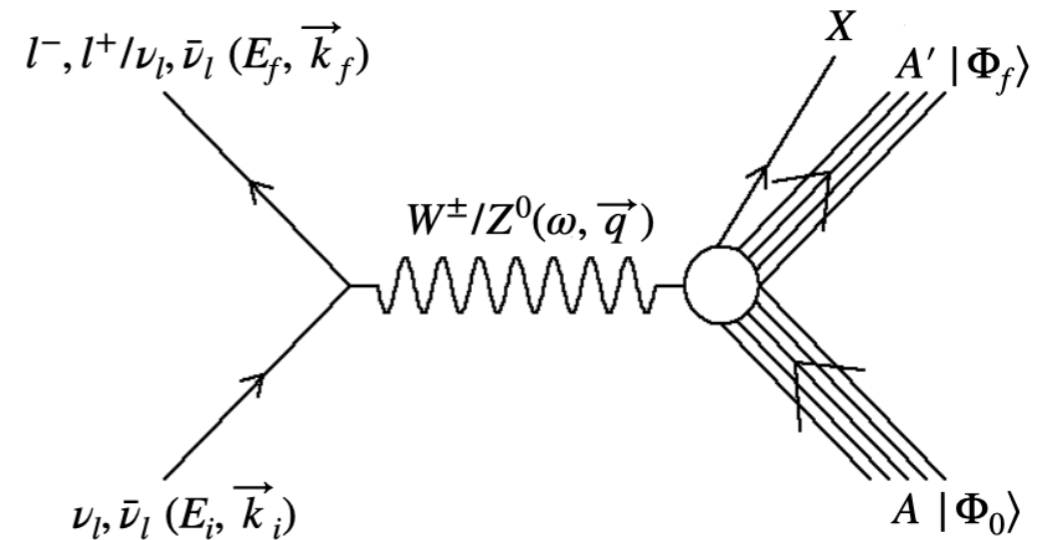
# Coherent Elastic and Inelastic Neutrino-Nucleus Scattering

## Coherent elastic [CEvNS]

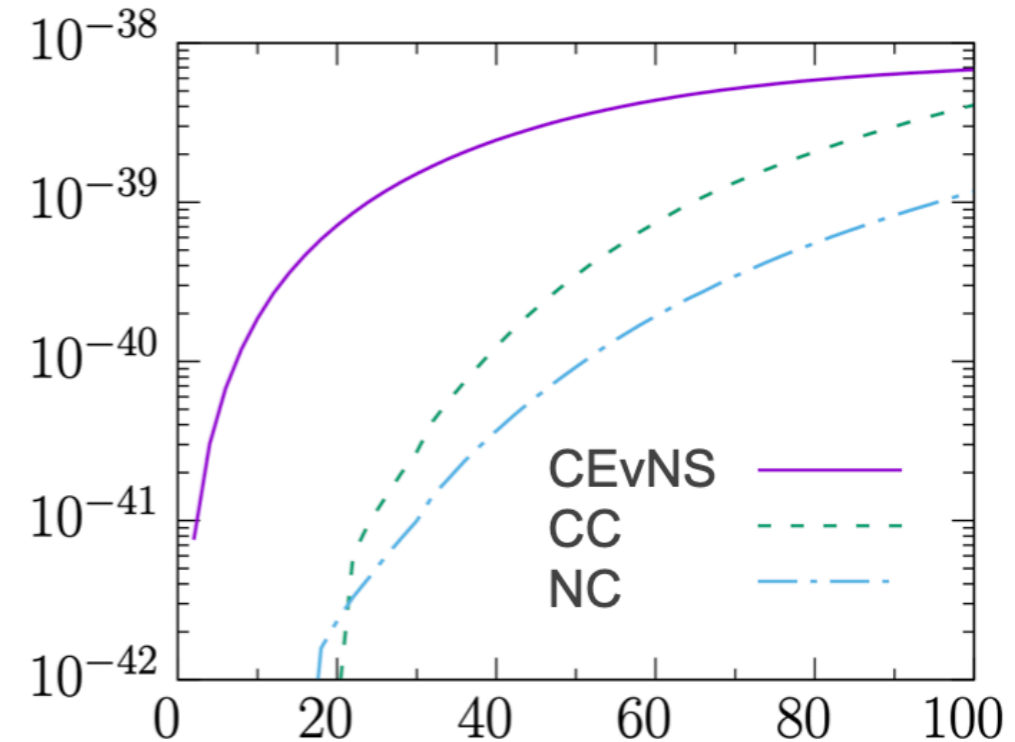
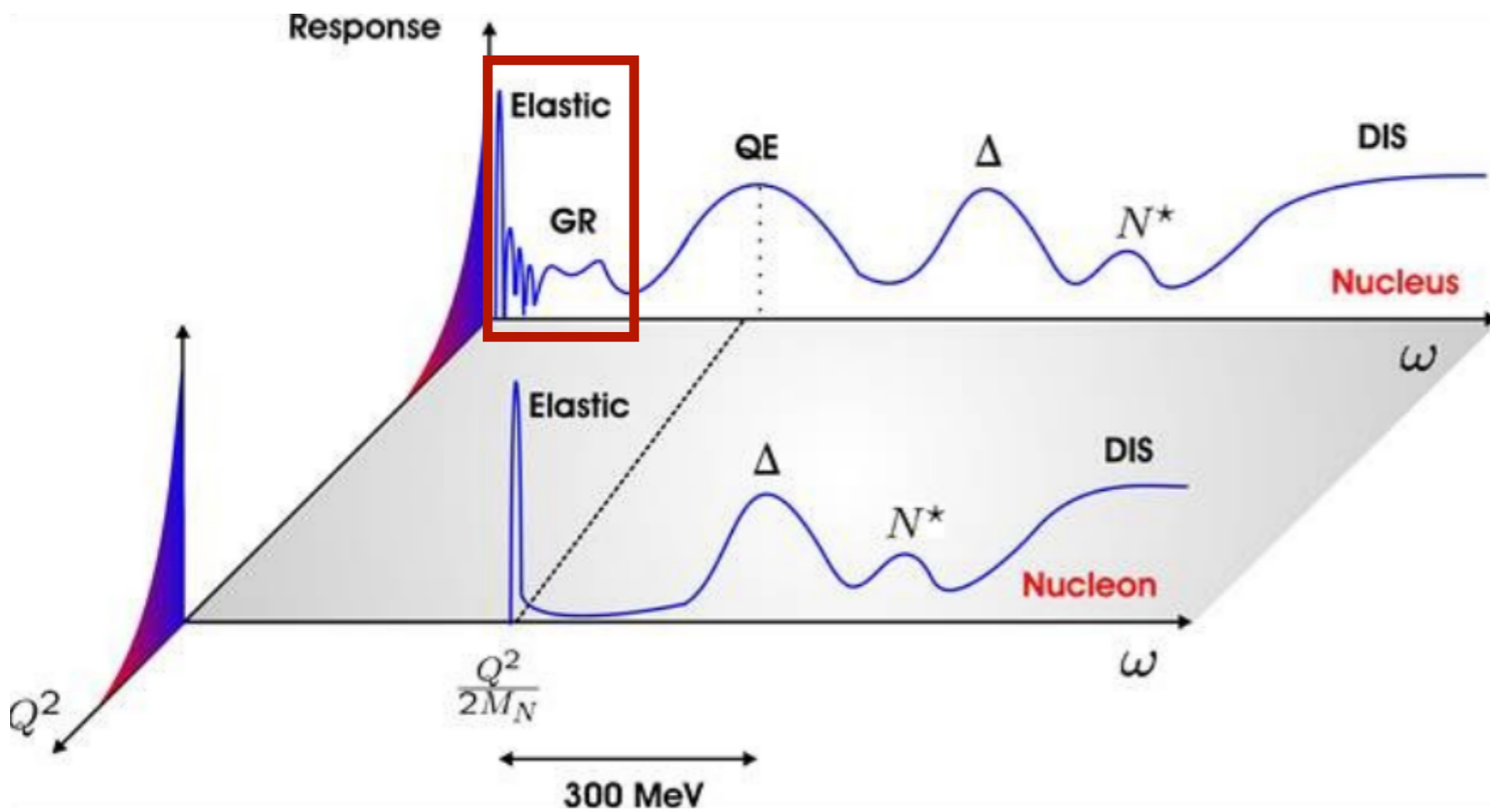


- Tiny recoil energy, large cross section
- Final state nucleus stays in its ground state
- Signal: keV energy nuclear recoil

## Inelastic CC/NC

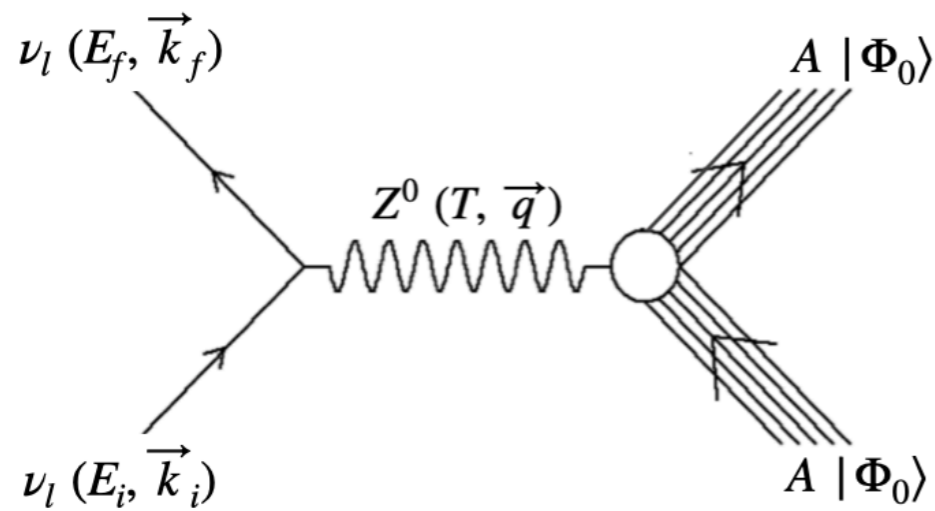


- Nucleus excites to states with well-defined excitation energy, spin and parity ( $J^\pi$ )
- Followed by nuclear de-excitation into gammas, n, p, and nuclear fragmentations.



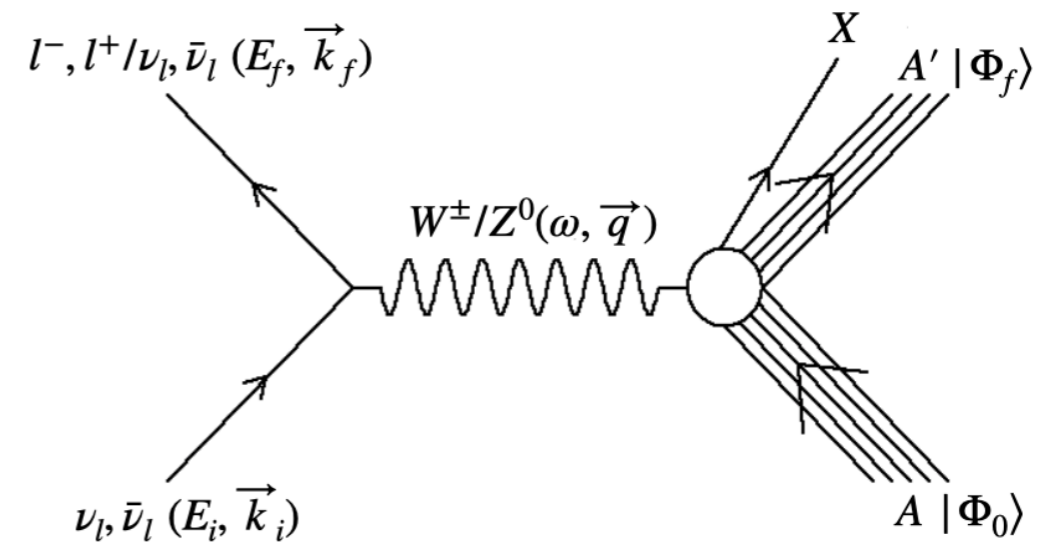
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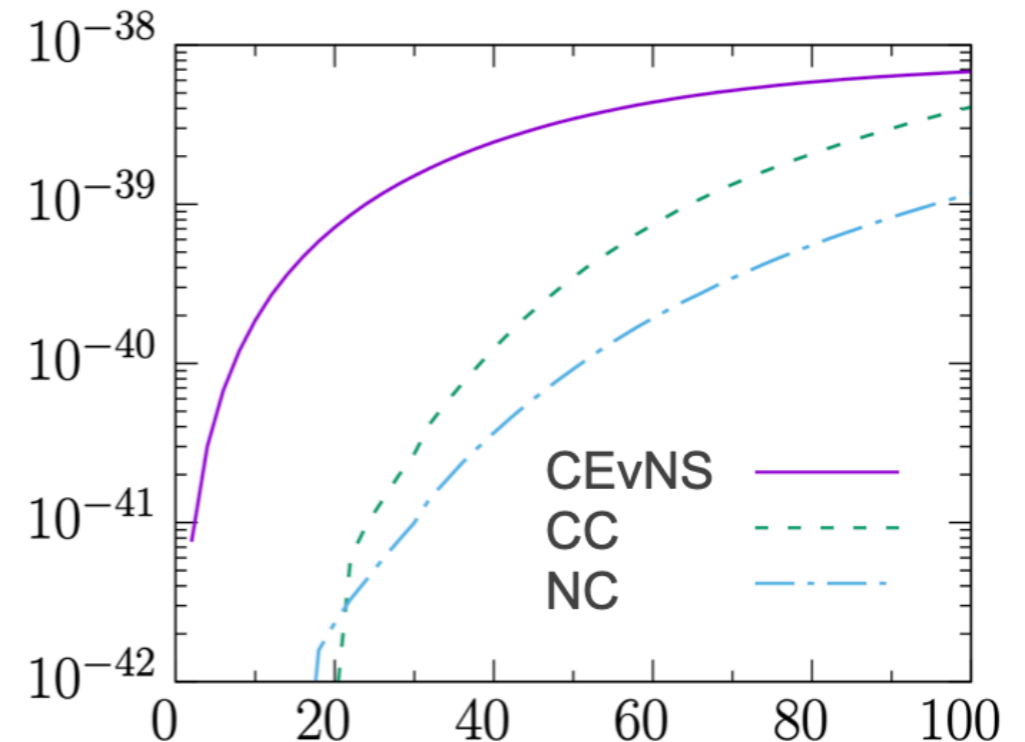


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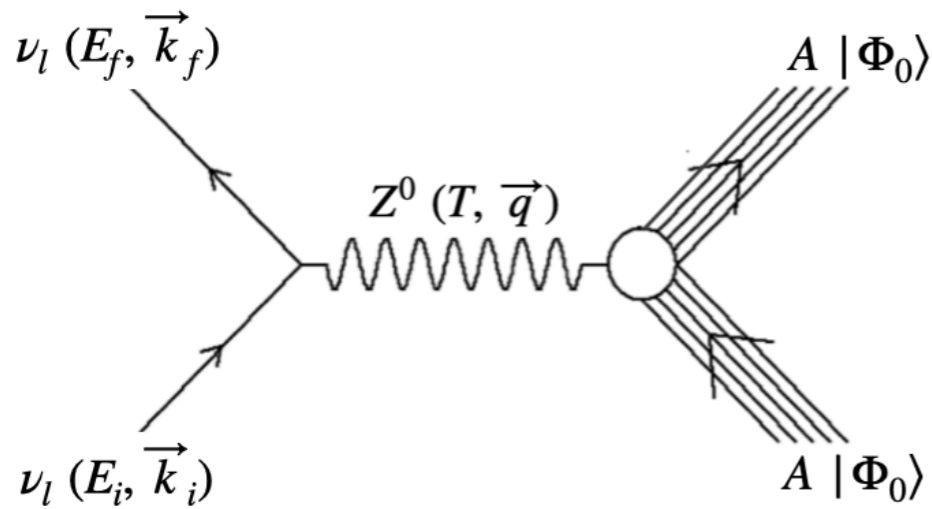
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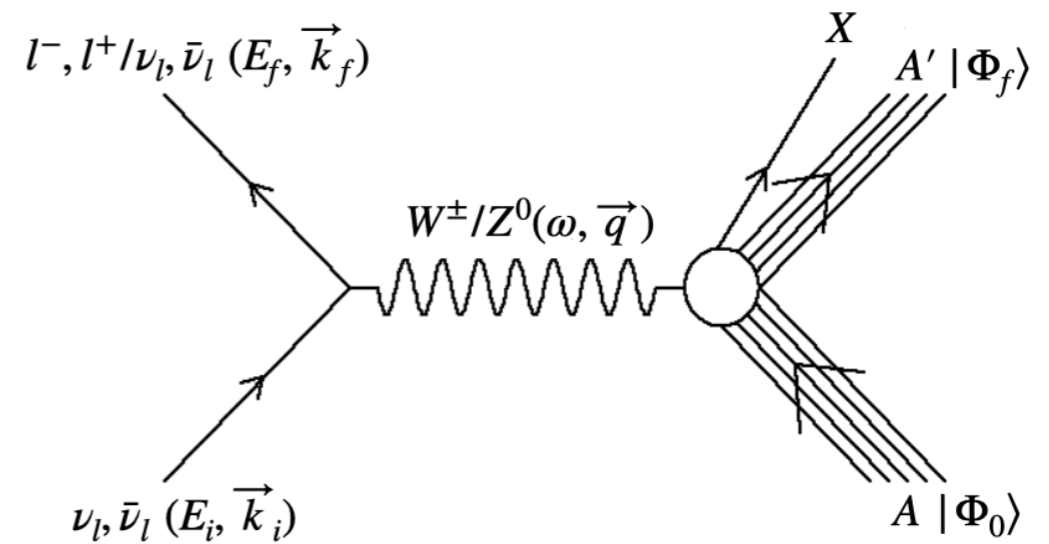
N. Van Dessel, V. Pandey, H. Ray, N. Jachowicz, [arXiv:2007.03658 \[nucl-th\]](https://arxiv.org/abs/2007.03658)

# Coherent Elastic and Inelastic Neutrino-Nucleus Scattering

## Coherent elastic [CEvNS]



## Inelastic CC/NC



$$\sum_{fi} |\mathcal{M}|^2 \propto \frac{G_F^2}{2} L_{\mu\nu} W^{\mu\nu}$$

$$\text{Leptonic Tensor: } L_{\mu\nu} = \sum_{fi} (\mathcal{J}_{l,\mu})^\dagger \mathcal{J}_{l,\nu}$$

$$\text{Hadronic Tensor: } W^{\mu\nu} = \sum_{fi} (\mathcal{J}_n^\mu)^\dagger \mathcal{J}_n^\nu$$

$$\text{Transition Amplitude: } \mathcal{J}_n^\mu = \langle \Phi_0 | \hat{J}_n^\mu(q) | \Phi_0 \rangle$$

$$\text{Transition Amplitude: } \mathcal{J}_n^\mu = \langle \Phi_f | \hat{J}_n^\mu(q) | \Phi_0 \rangle$$

Cross Section:

$$d\sigma \propto \frac{G_F^2}{4\pi} Q_W^2 F_W^2(q)$$

Cross Section:

$$d\sigma \propto \frac{G_F^2}{4\pi} \sum_{J^\pi} [v_{CC} W_{CC} + v_{CL} W_{CL} + v_{LL} W_{LL} + v_T W_T \pm v_{T'} W_{T'}]$$



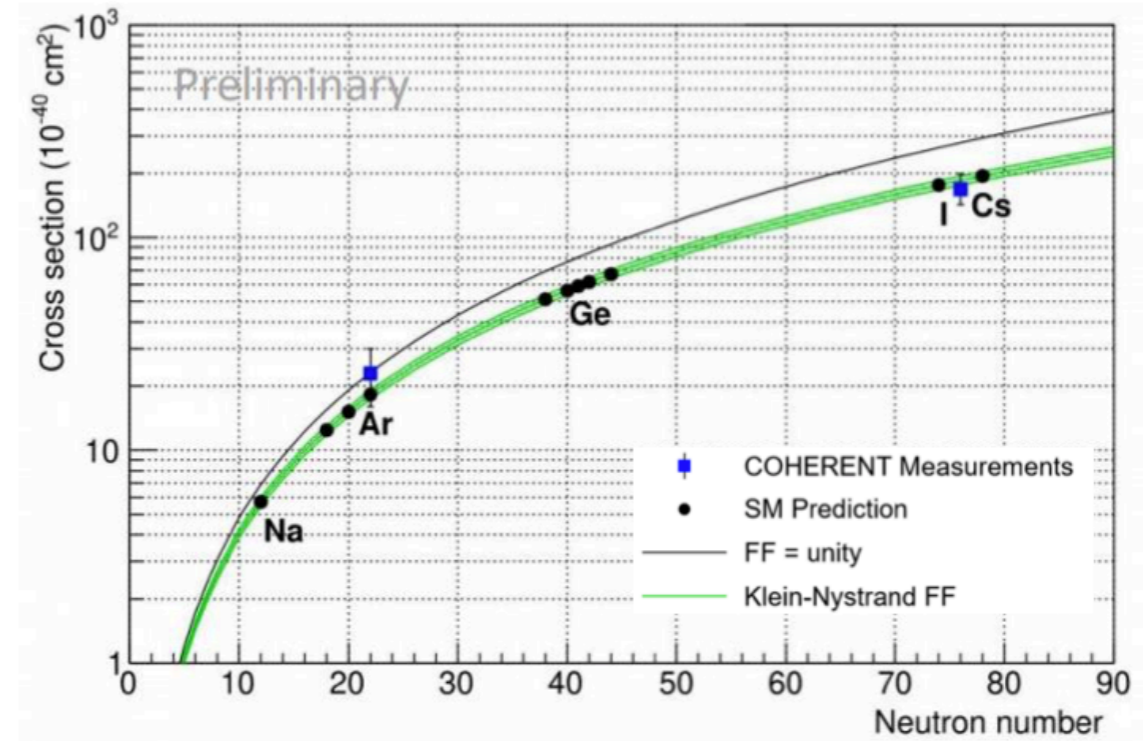
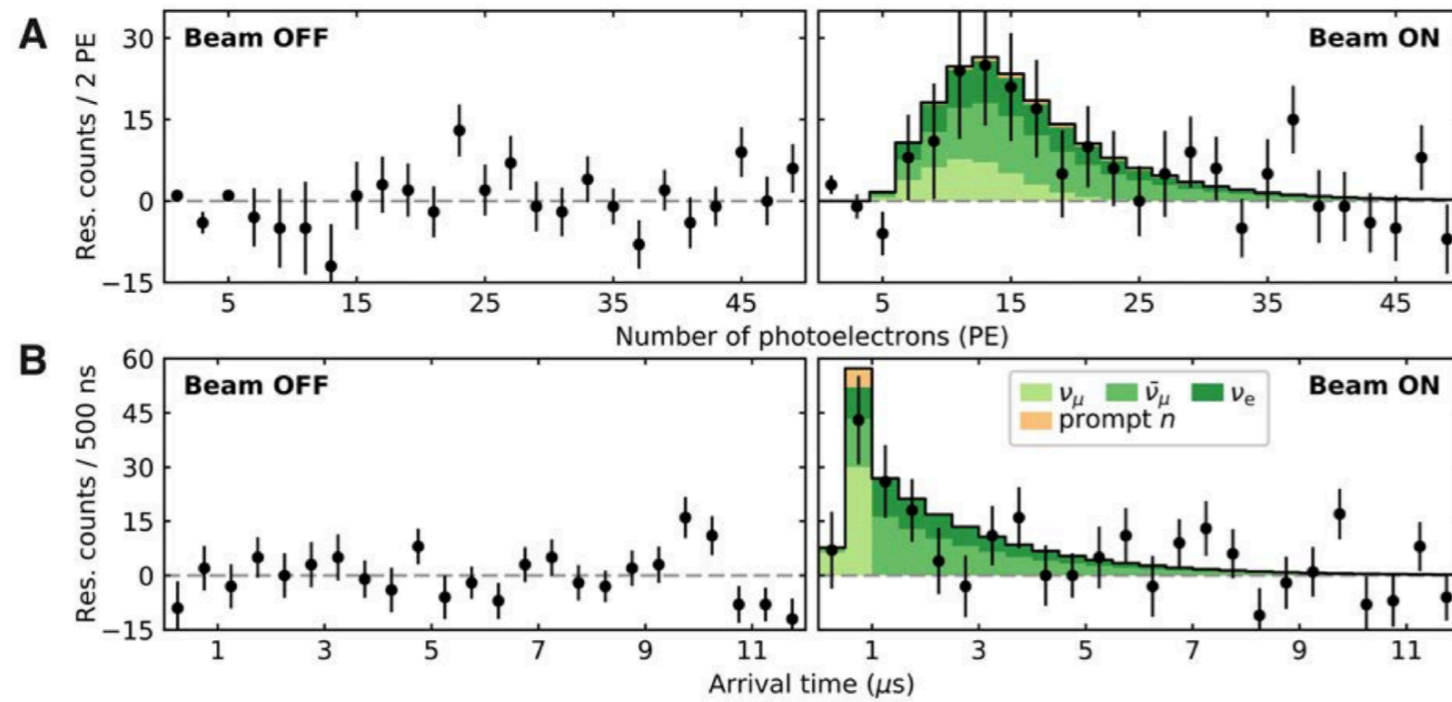
# Coherent Elastic Neutrino-Nucleus Scattering

## COHERENT Collaboration at SNS at ORNL

Monday Morning's Session

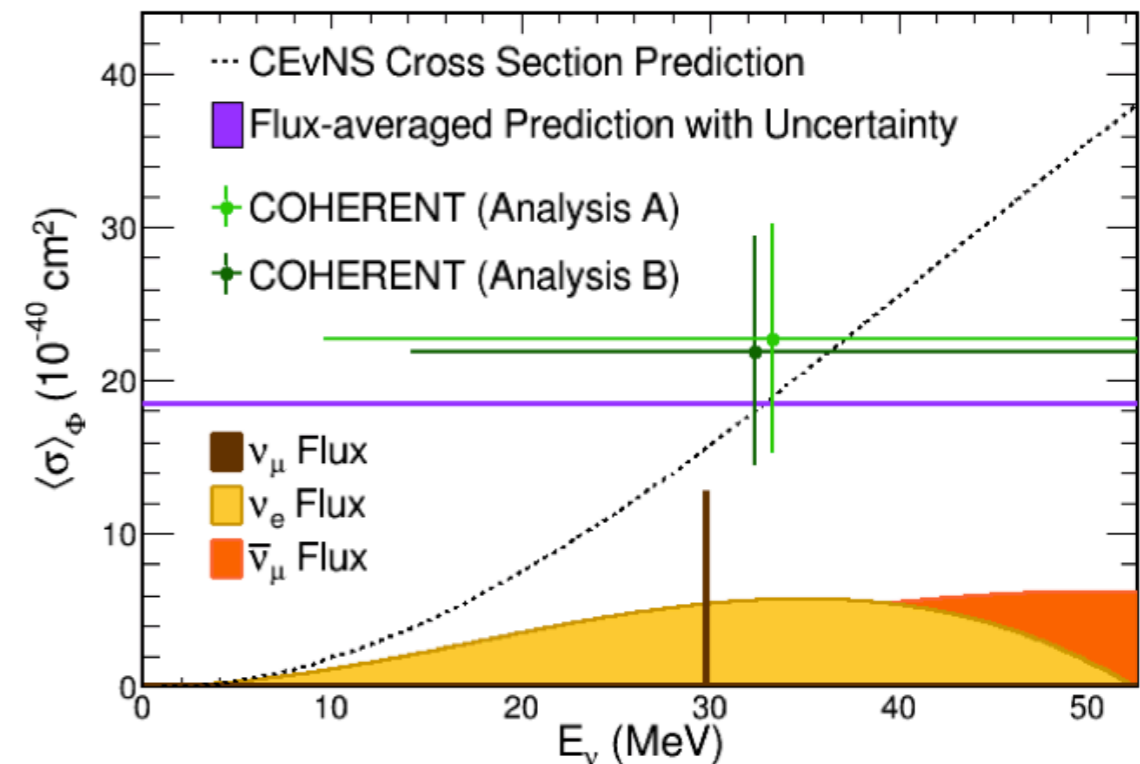
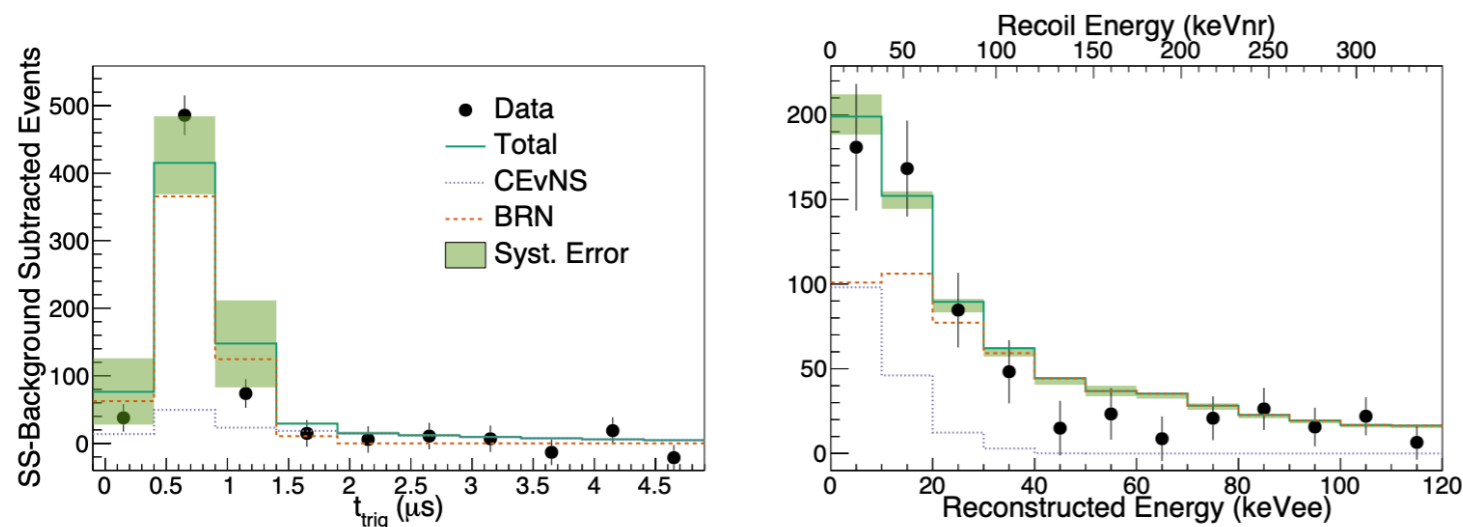
### 14 kg CSI detector

COHERENT Collaboration, *Science* 357, 6356, 1123–1126 (2017)



### 24 kg LAr (CENNS-10) detector

COHERENT Collaboration, *Phys. Rev. Lett.* 126, 012002 (2021)



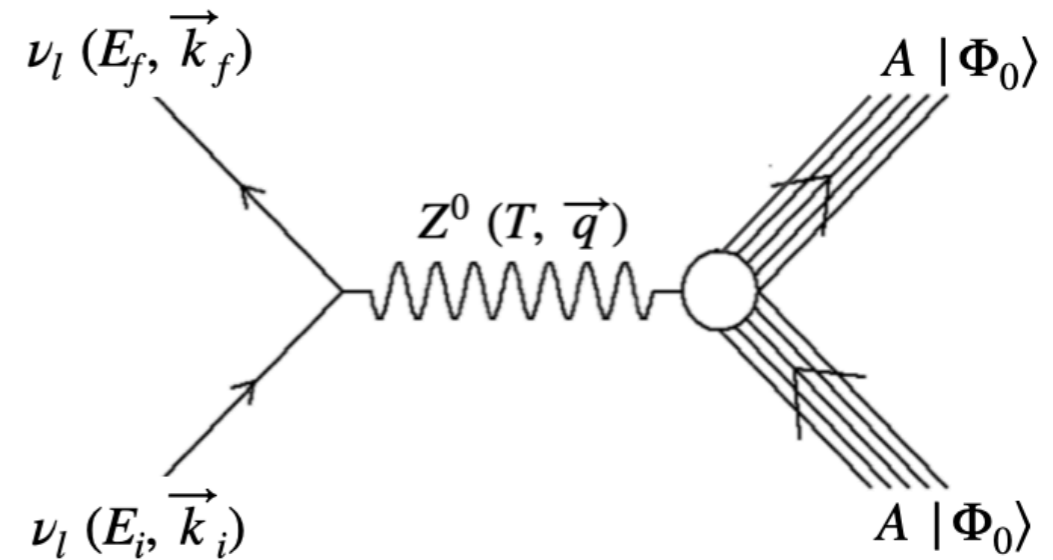


# CEvNS Cross Section and Form Factors

## ■ Cross section\*:

$$\frac{d\sigma}{dT} = \frac{G_F^2}{\pi} M_A \left[ 1 - \frac{T}{E_i} - \frac{M_A T}{2E_i^2} \right] \frac{Q_W^2}{4} F_W^2(q)$$

$$\frac{d\sigma}{d\cos\theta_f} = \frac{G_F^2}{2\pi} E_i^2 (1 + \cos\theta_f) \frac{Q_W^2}{4} F_W^2(q)$$



$$T \in \left[ 0, \frac{2E_i^2}{(M_A + 2E_i)} \right]$$

$$Q_W^2 = [g_n^V N + g_p^V Z]^2$$

## ■ Weak Form Factor:

$$\begin{aligned} Q_W F_W(q) &\approx \langle \Phi_0 | \hat{J}_0(q) | \Phi_0 \rangle \\ &\approx (1 - 4 \sin^2 \theta_W) Z F_p(q) - N F_n(q) \\ &\approx 2\pi \int d^3r \left[ (1 - 4 \sin^2 \theta_W) \rho_p(r) - \rho_n(r) \right] j_0(qr) \end{aligned}$$

*N. Van Dessel, V. Pandey, H. Ray, N. Jachowicz, arXiv:2007.03658 [nucl-th]*

Charge density and charge form factor: proton densities and charge form factors are well known through decades of elastic electron scattering experiments.

Neutron densities and neutron form factor: neutron densities and form factors are poorly known. Note that CEvNS is primarily sensitive to neutron density distributions.

\*barring radiative corrections, for radiative corrections, see:

*O. Tomalak, P. Machado, V. Pandey, R. Plestid, JHEP 02, 097 (2021)*

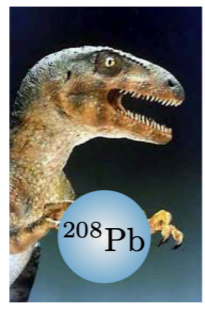
# CEvNS and PVES

- **Electroweak probes** such as parity-violating electron scattering (**PVES**) and **CEvNS** provide relatively model-independent ways of determining weak form factor and neutron distributions.
- The parity violating asymmetry for elastic electron scattering is the fractional difference in cross section for positive helicity and negative helicity electrons.

$$A_{pv} = \frac{d\sigma/d\Omega_+ - d\sigma/d\Omega_-}{d\sigma/d\Omega_+ + d\sigma/d\Omega_-} = \frac{G_F q^2 |Q_W| F_W(q^2)}{4\pi\alpha\sqrt{2}Z F_{ch}(q^2)}$$

Experiment	Target	$q^2$ (GeV <sup>2</sup> )	$A_{pv}$ (ppm)	$\pm\delta R_n$ (%)
PREX	<sup>208</sup> Pb	0.00616	0.550 ± 0.018	1.3
CREX	<sup>48</sup> Ca	0.0297		0.7
Qweak	<sup>27</sup> Al	0.0236	2.16 ± 0.19	4
MREX	<sup>208</sup> Pb	0.0073		0.52

[arXiv:2203.06853 \[hep-ex\]](https://arxiv.org/abs/2203.06853)



Pb Radius Experiment (PREX)



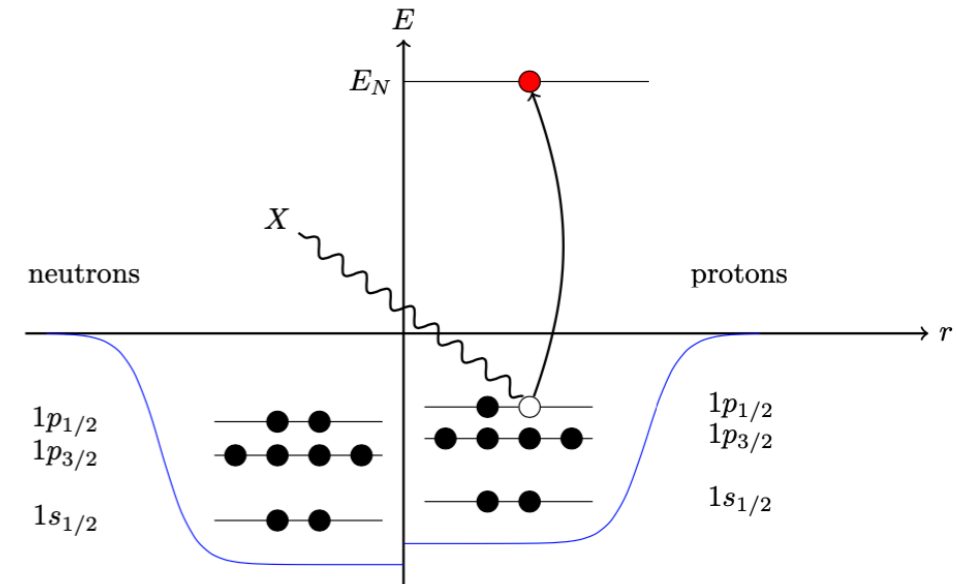
Calcium Radius Experiment (CREX)



Mainz Radius Experiment (MREX)  
At P2 experimental hall with <sup>208</sup>Pb

# CEvNS Cross Section Calculations

- Nuclear ground state described as a many-body quantum mechanical system where nucleons are bound in an effective nuclear potential.
- Solve Hartree-Fock (**HF**) equation with a Skyrme (**SkE2**) nuclear potential to obtain single-nucleon wave functions for the bound nucleons in the nuclear ground state.
- Evaluate proton and neutron density distributions and form factors



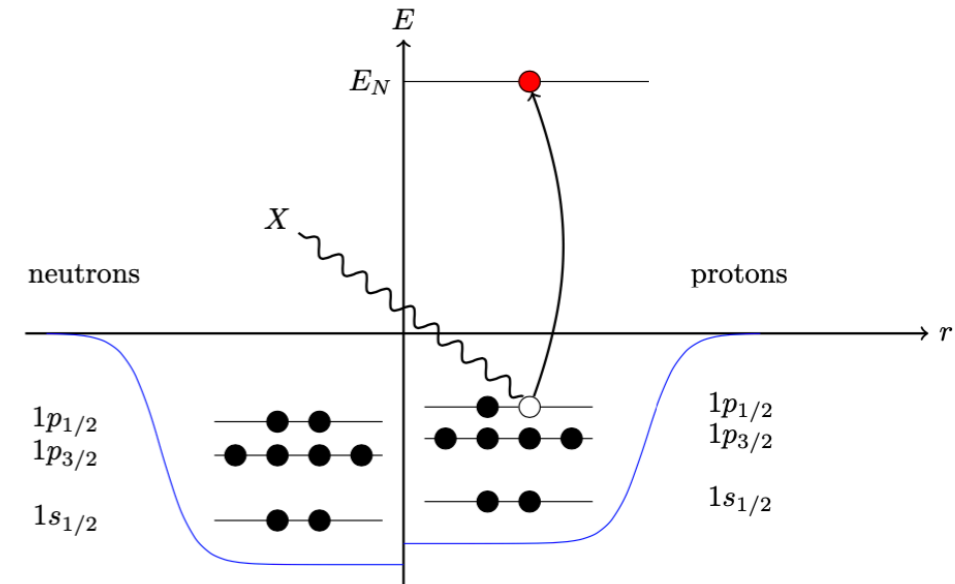
$$\rho_{\tau}(r) = \frac{1}{4\pi r^2} \sum_{\alpha} v_{\alpha,\tau}^2 (2j_{\alpha} + 1) |\phi_{\alpha,\tau}(r)|^2$$

$$F_{\tau}(q) = \frac{1}{N} \int d^3r j_0(qr) \rho_{\tau}(r) \quad (\alpha \in n_{\alpha}, l_{\alpha}, j_{\alpha})$$

$$(\tau = p, n)$$

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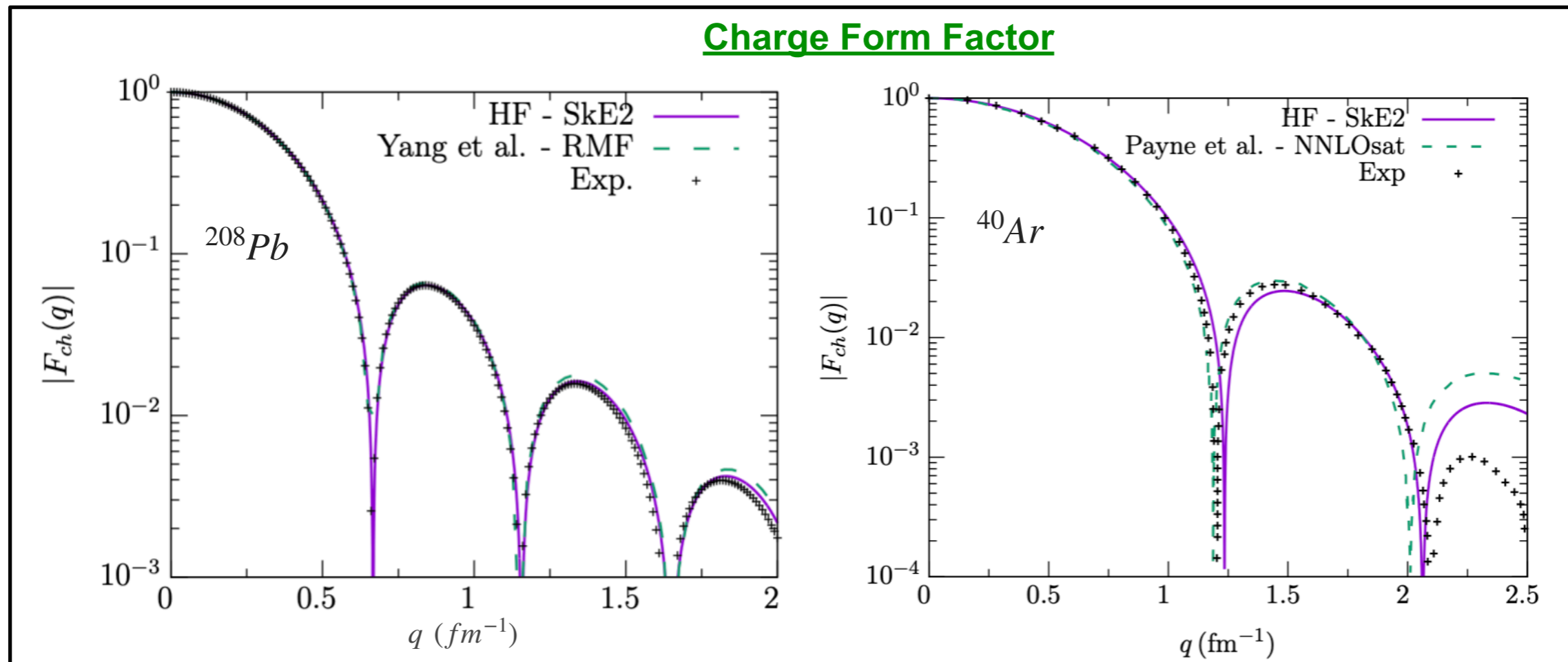
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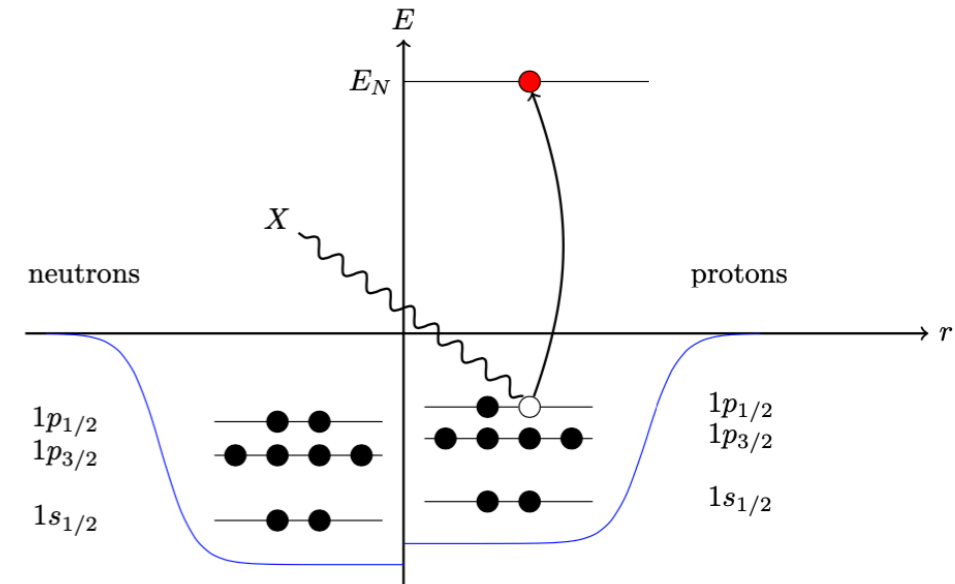
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*N. Van Dessel, V. Pandey, H. Ray, N. Jachowicz, arXiv:2007.03658 [nucl-th]*

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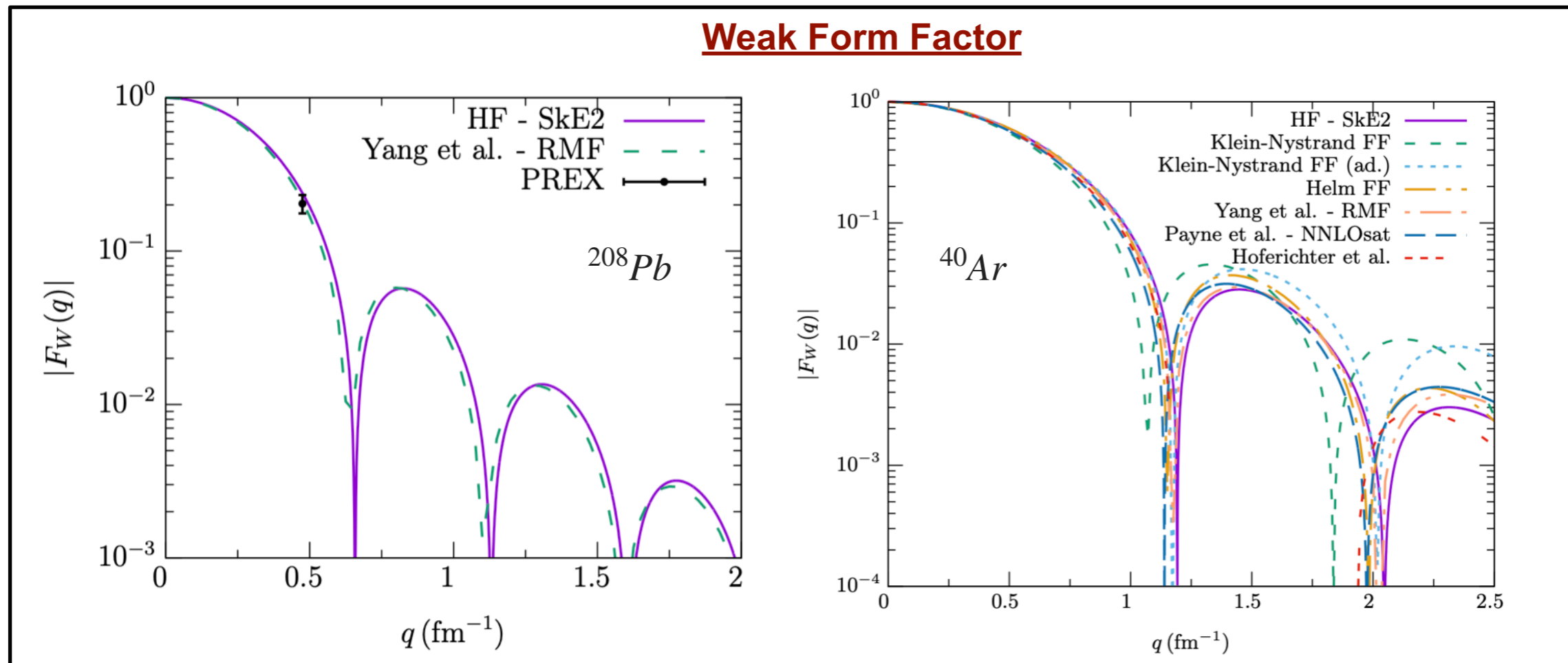


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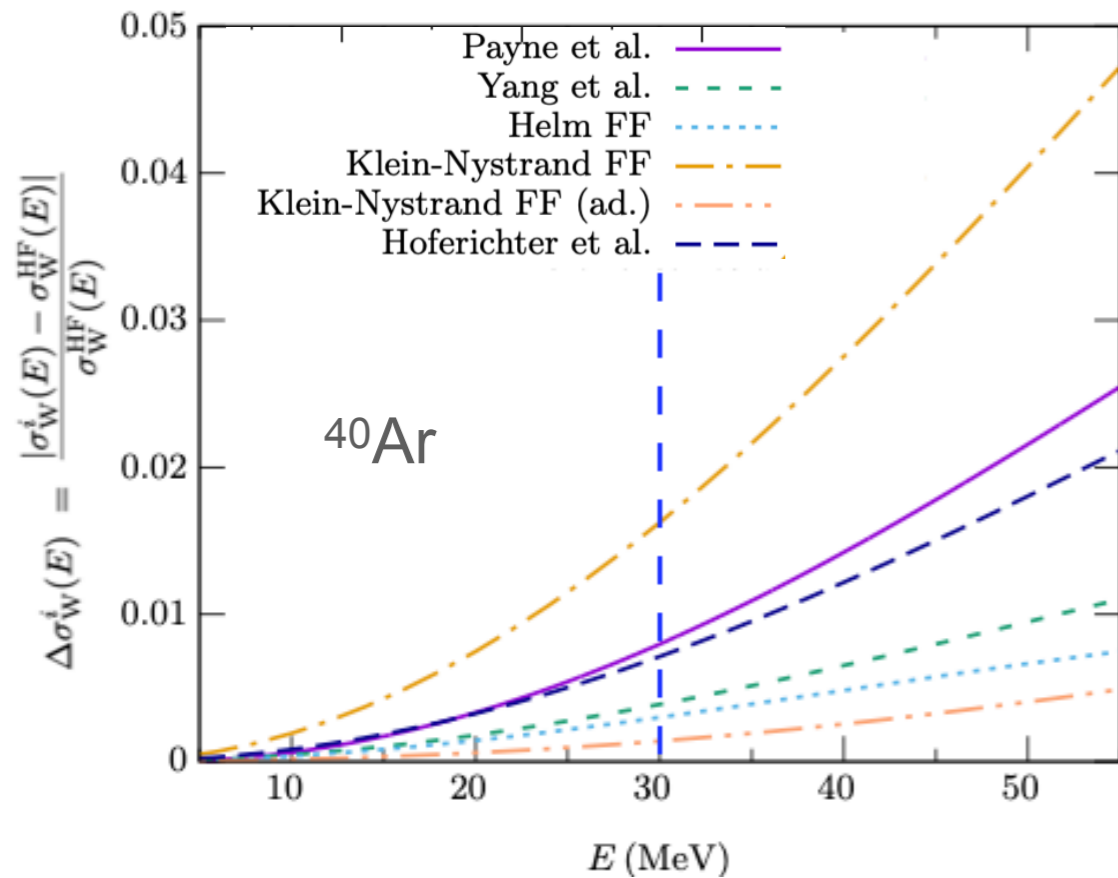
*N. Van Dessel, V. Pandey, H. Ray, N. Jachowicz, arXiv:2007.03658 [nucl-th]*

*Data: S. Abrahamyan et al., Phys. Rev. Lett. 108, 112502 (2012)*



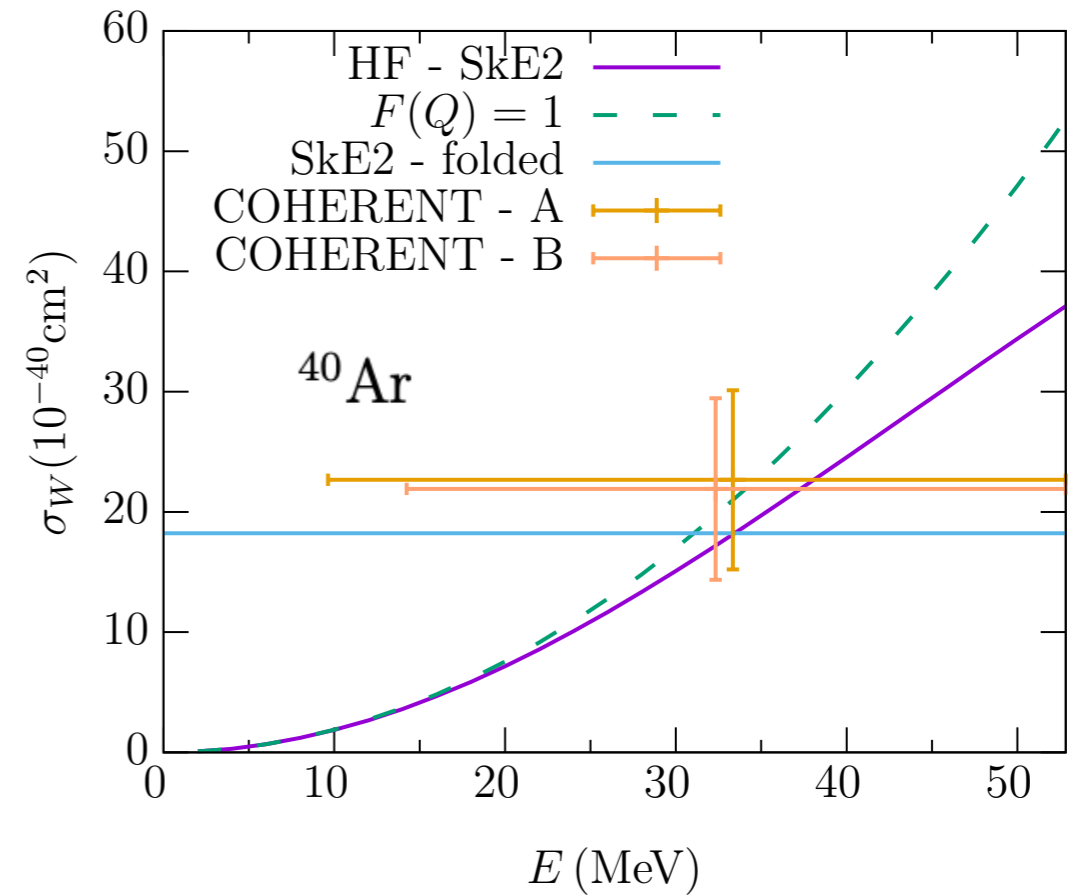
# CEvNS Cross Section Calculations

- Relative CEvNS cross section differences between the results of different calculations:



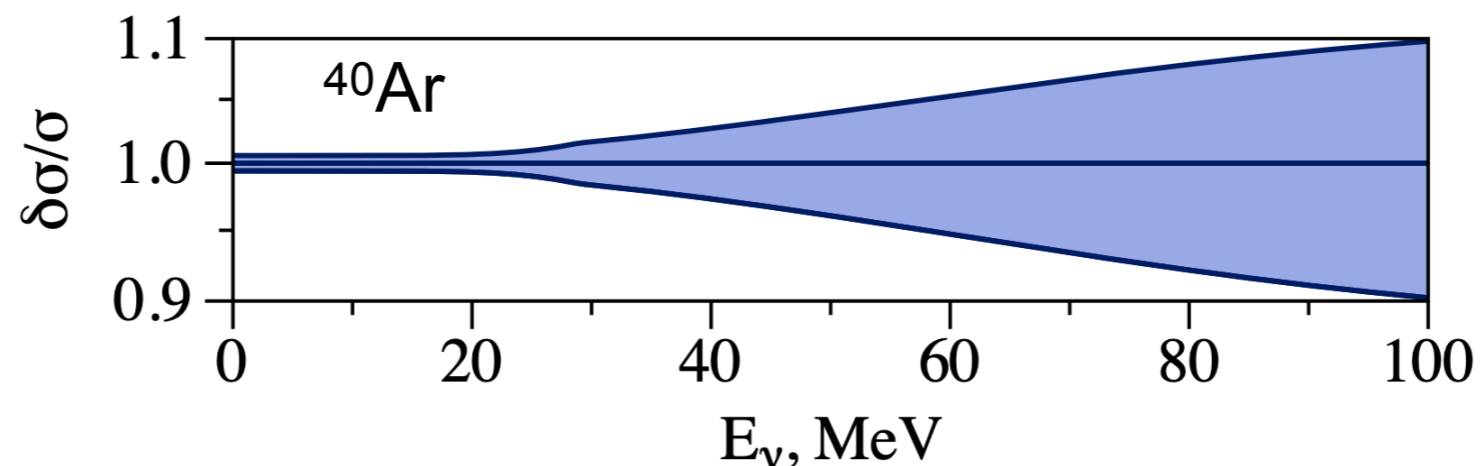
*N. Van Dessel, V. Pandey, H. Ray, N. Jachowicz, arXiv:2007.03658 [nucl-th]*

- Comparison with COHERENT data



*COHERENT data: arXiv:2003.10630 [nucl-ex].*

- Relative CEvNS cross section theoretical uncertainty on  $^{40}\text{Ar}$  (includes nuclear, nucleonic, hadronic, quark levels as well as perturbative errors):



*O. Tomalak, P. Machado, V. Pandey, R. Plestid, JHEP 02, 097 (2021)*

# 10s of MeV Inelastic Neutrino-Nucleus Scattering

- CEvNS experiments at stopped-pion sources are also powerful avenues to measure 10s of MeV inelastic CC and NC cross sections subject to detailed underlying nuclear structure and dynamics.
  - These are vital in the detection of core-collapse supernovae, but are almost completely unexplored experimentally so far.

Past measurements on Carbon

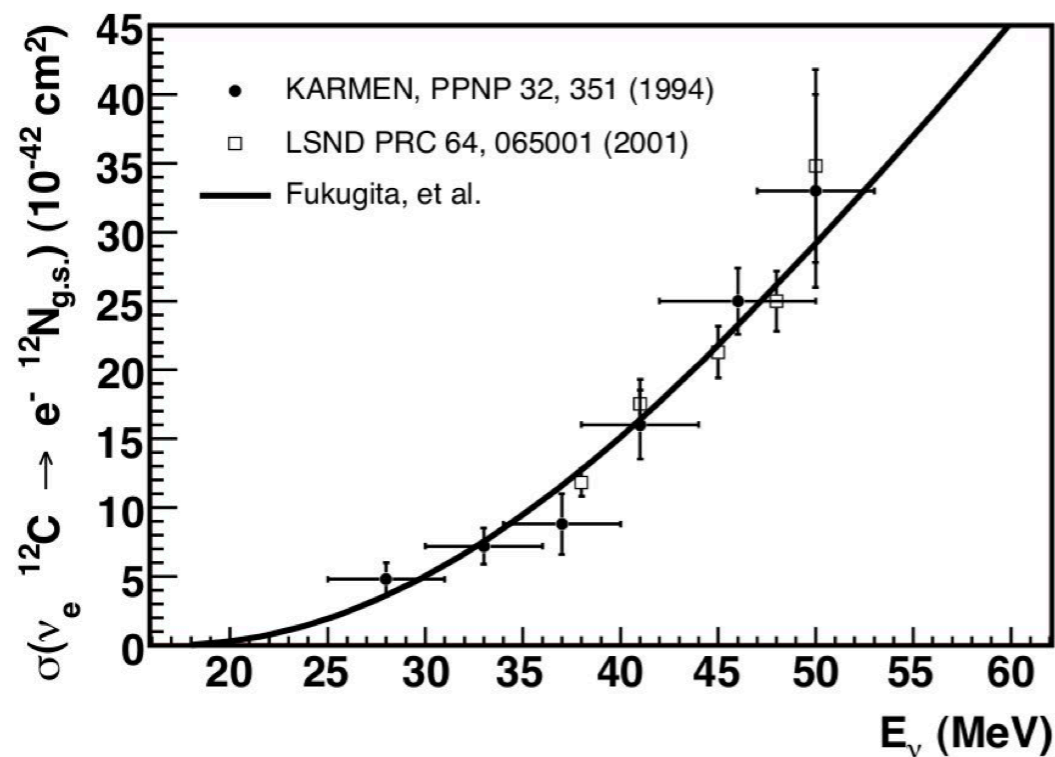
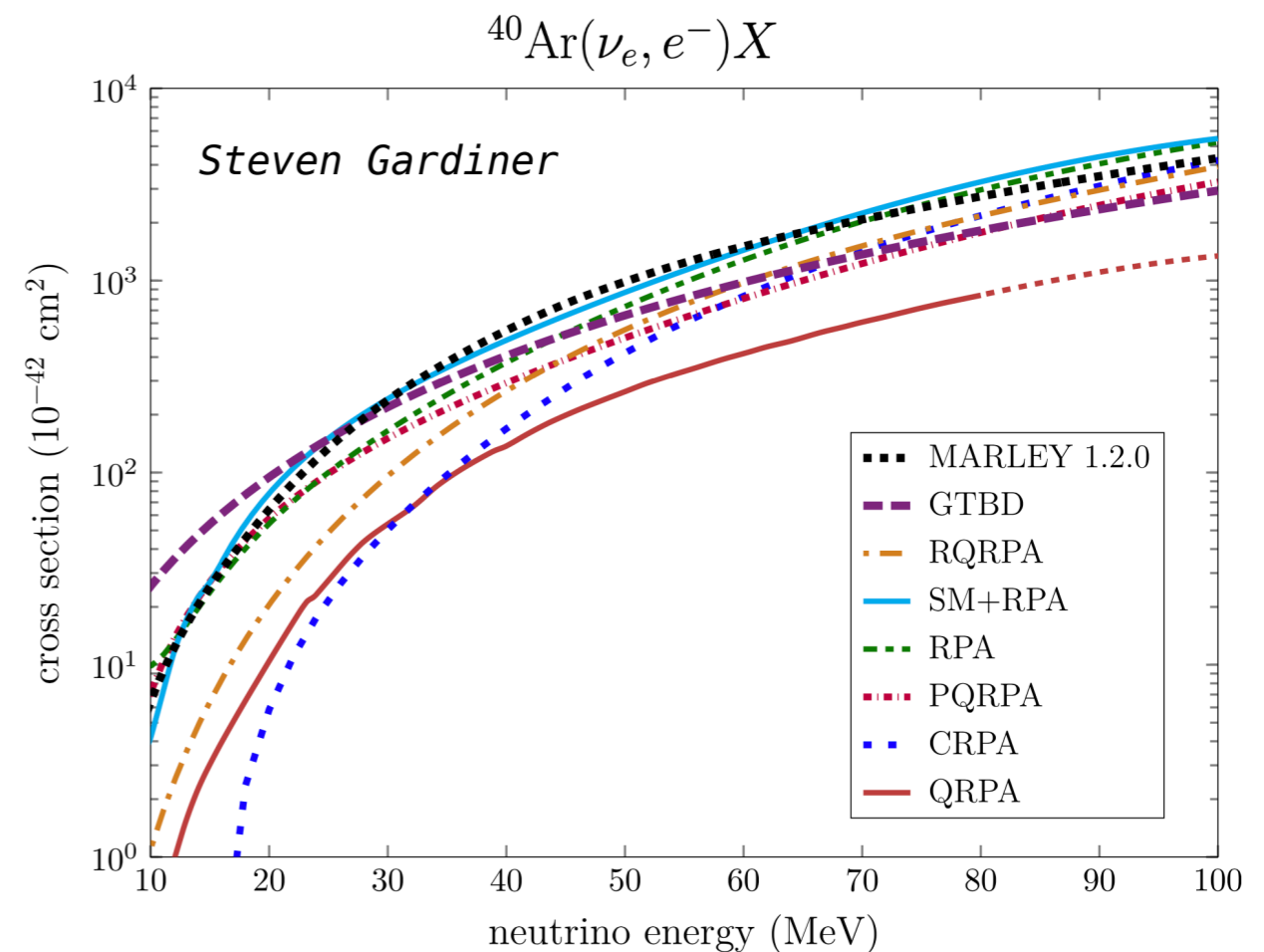


FIG. 6 Cross-section as a function of neutrino energy for the exclusive reaction  $^{12}\text{C}(\nu_e, e^-)^{12}\text{N}$  from  $\mu^-$  decay-at-rest neutrinos. Experimental data measured by the KARMEN (Zeitnitz *et al.*, 1994) and LSND (Athanasopoulos *et al.*, 1997; Auerbach *et al.*, 2001) experiments. Theoretical prediction taken from Fukugita *et al.* (Fukugita *et al.*, 1988).

*Rev. Mod. Phys.* 84,1307 (2012)

No measurements on Argon yet

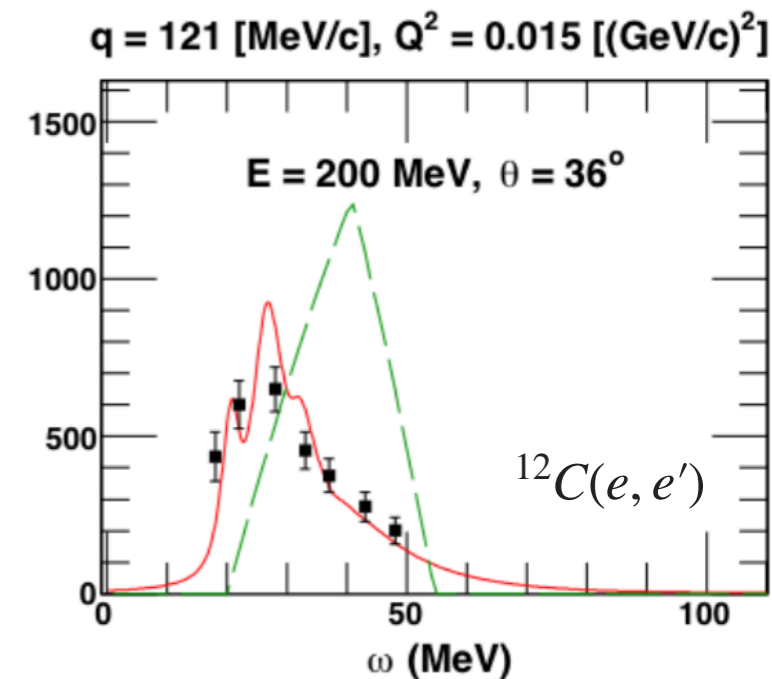
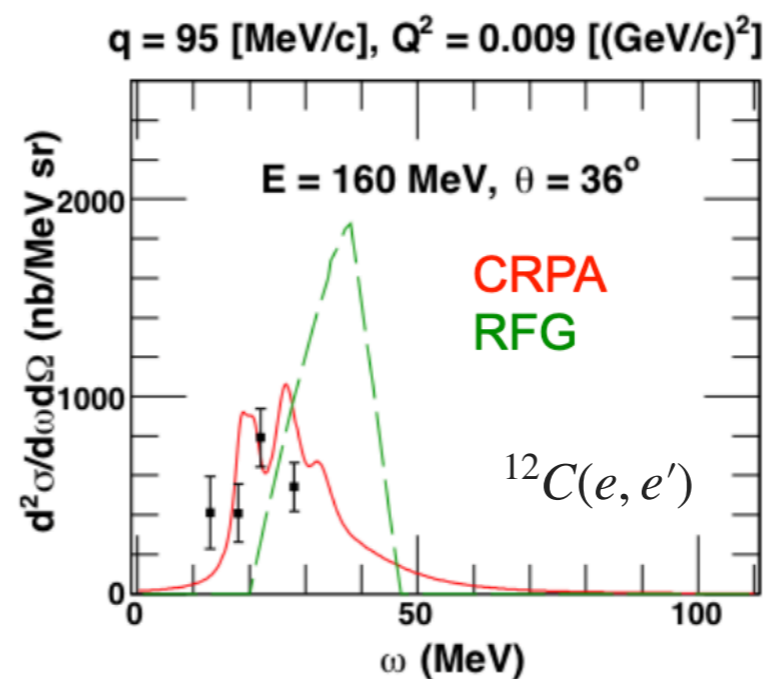
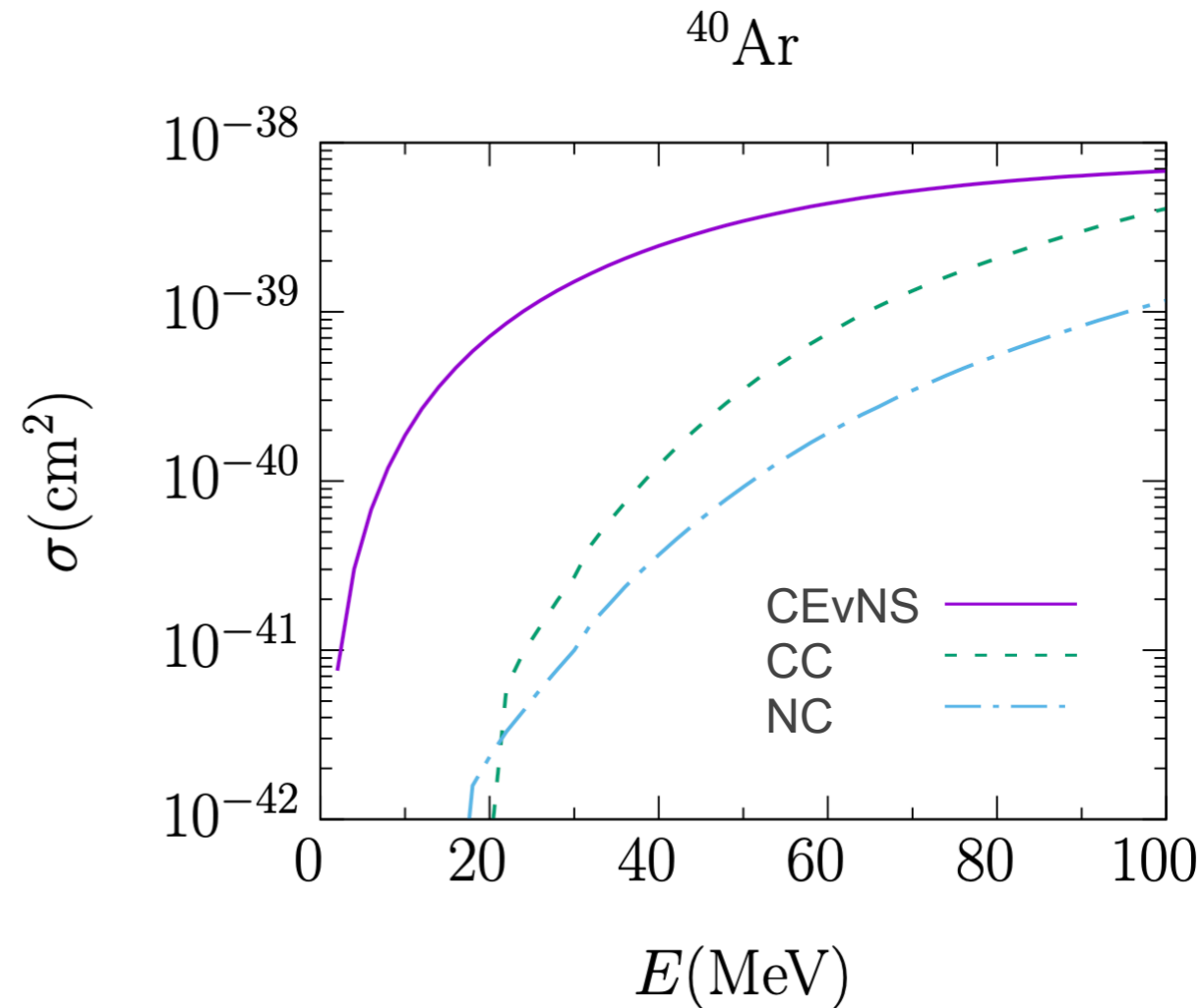
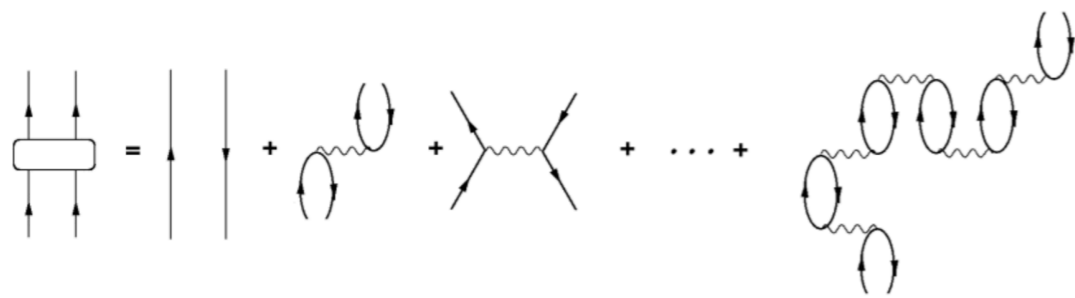


◆ Need inelastic neutrino-nucleus cross section measurements at stopped-pion sources.

# 10s of MeV Inelastic Neutrino-Nucleus Scattering Calculations

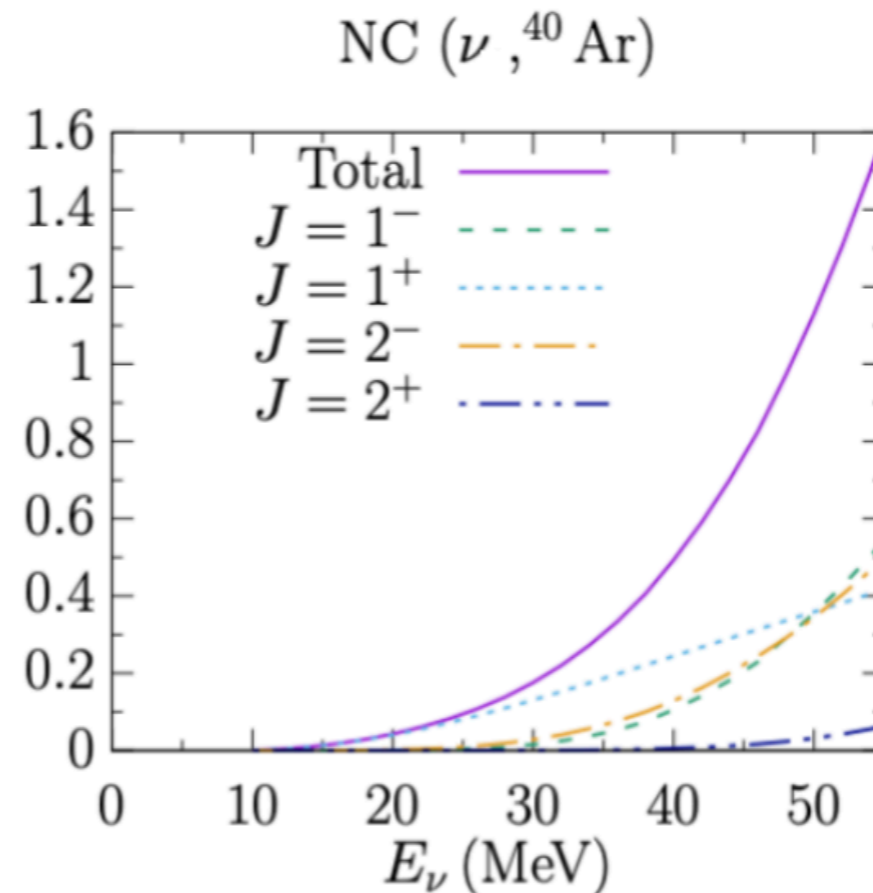
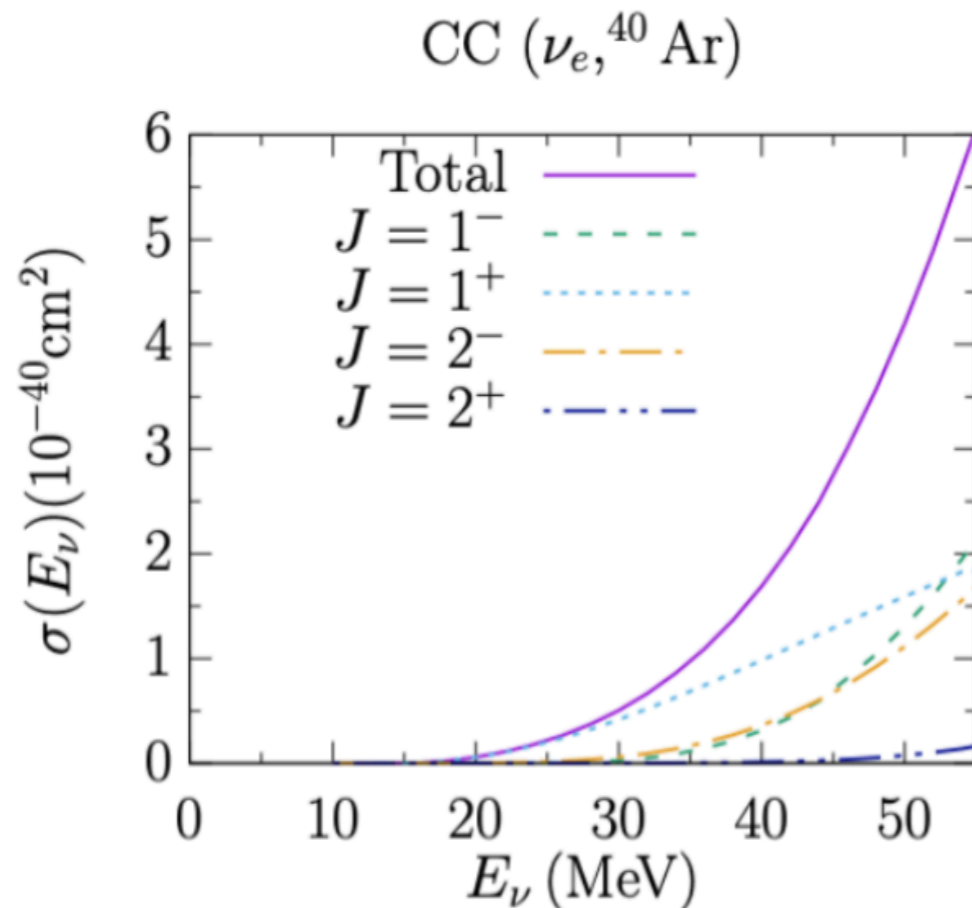
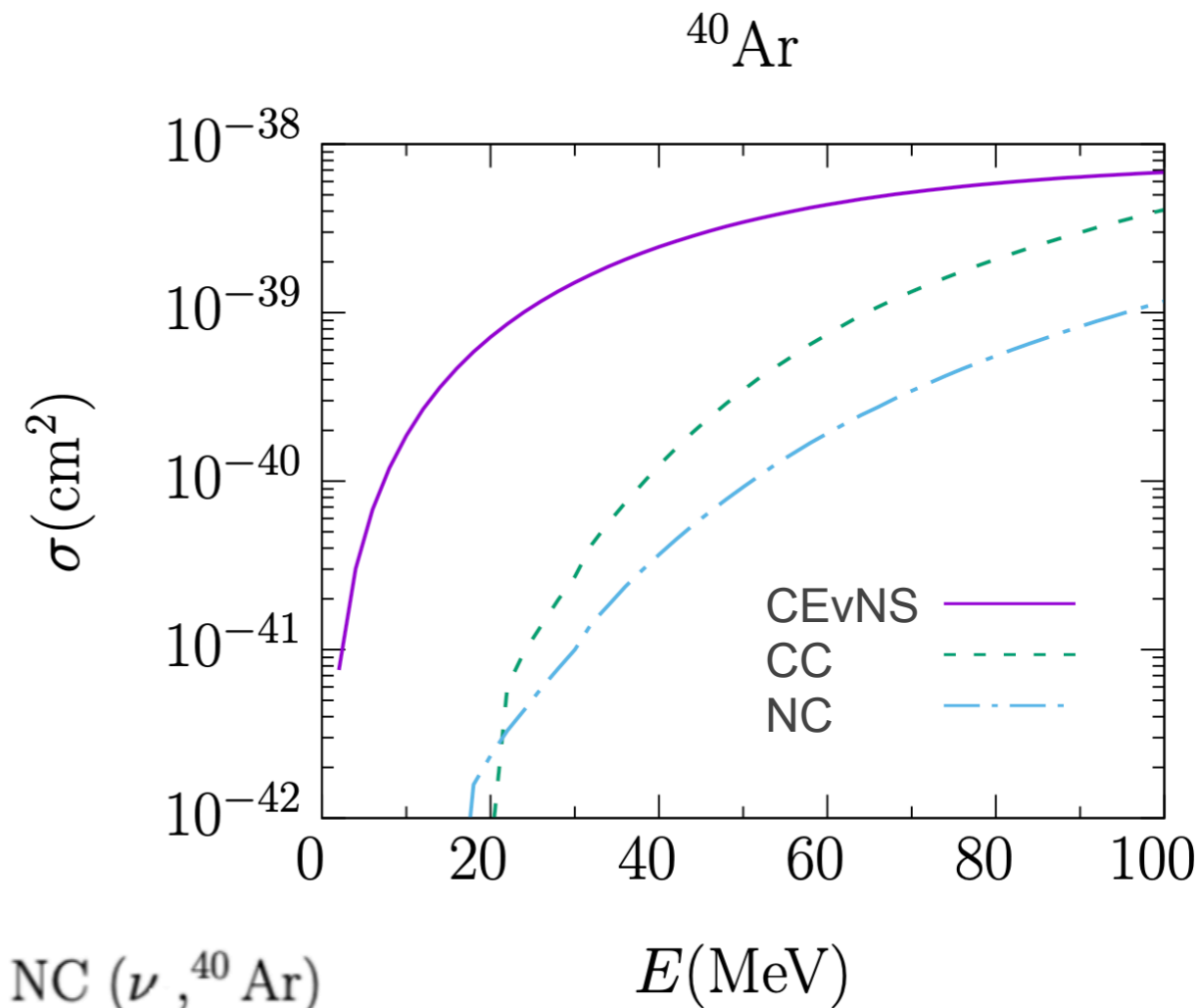
- In the inelastic cross section calculations, the influence of long-range correlations between the nucleons is introduced through the **continuum Random Phase Approximation (CRPA)** on top of the HF-SkE2 approach.
- CRPA effects are vital to describe the process where the nucleus can be excited to low-lying collective nuclear states.
- The local RPA-polarization propagator is obtained by an iteration to all orders of the first order contribution to the particle-hole Green's function.

$$\Pi^{(RPA)}(x_1, x_2; E_x) = \Pi^{(0)}(x_1, x_2; E_x) + \frac{1}{\hbar} \int dx dx' \Pi^0(x_1, x; E_x) \times \tilde{V}(x, x') \Pi^{(RPA)}(x', x_2; E_x)$$



# 10s of MeV Inelastic Neutrino-Nucleus Scattering Calculations

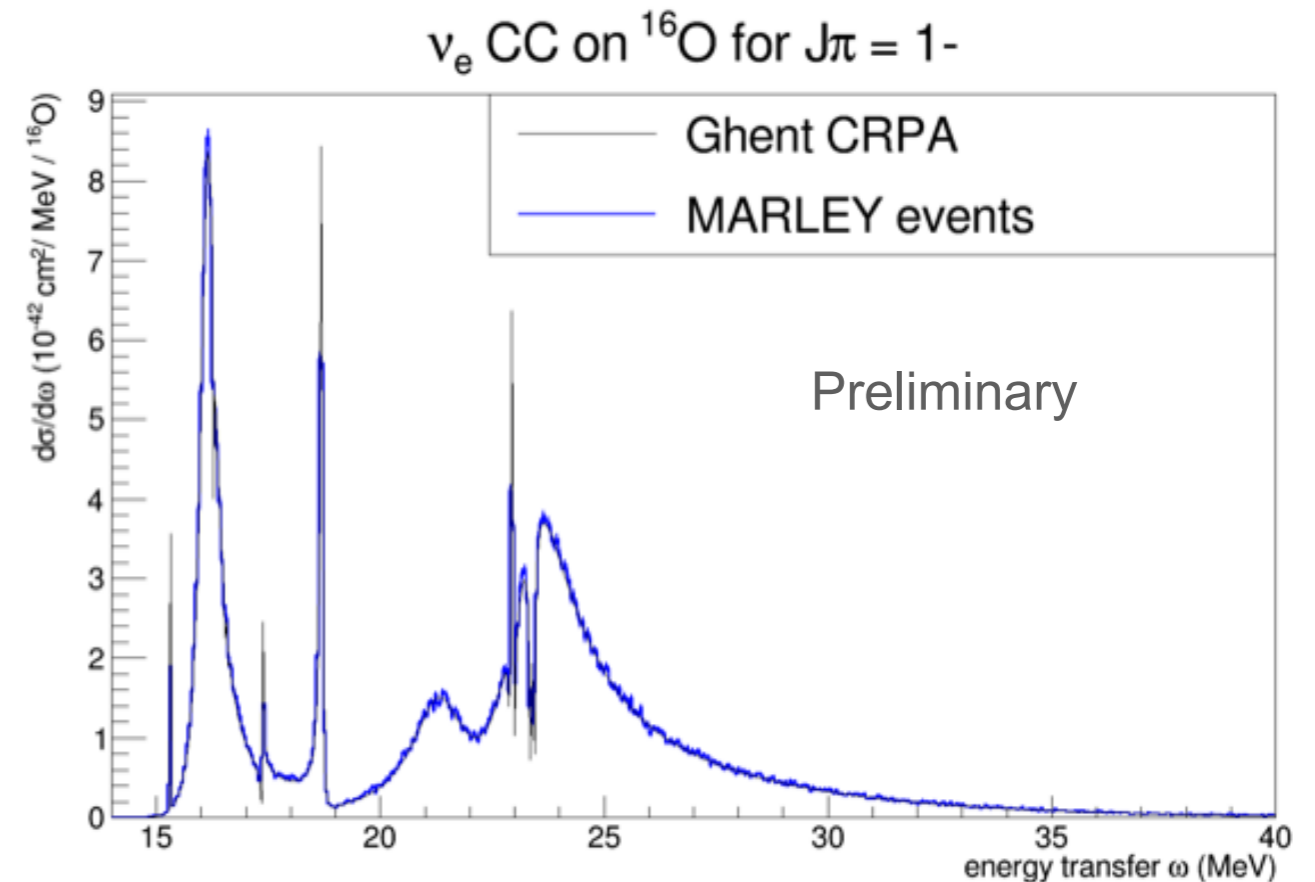
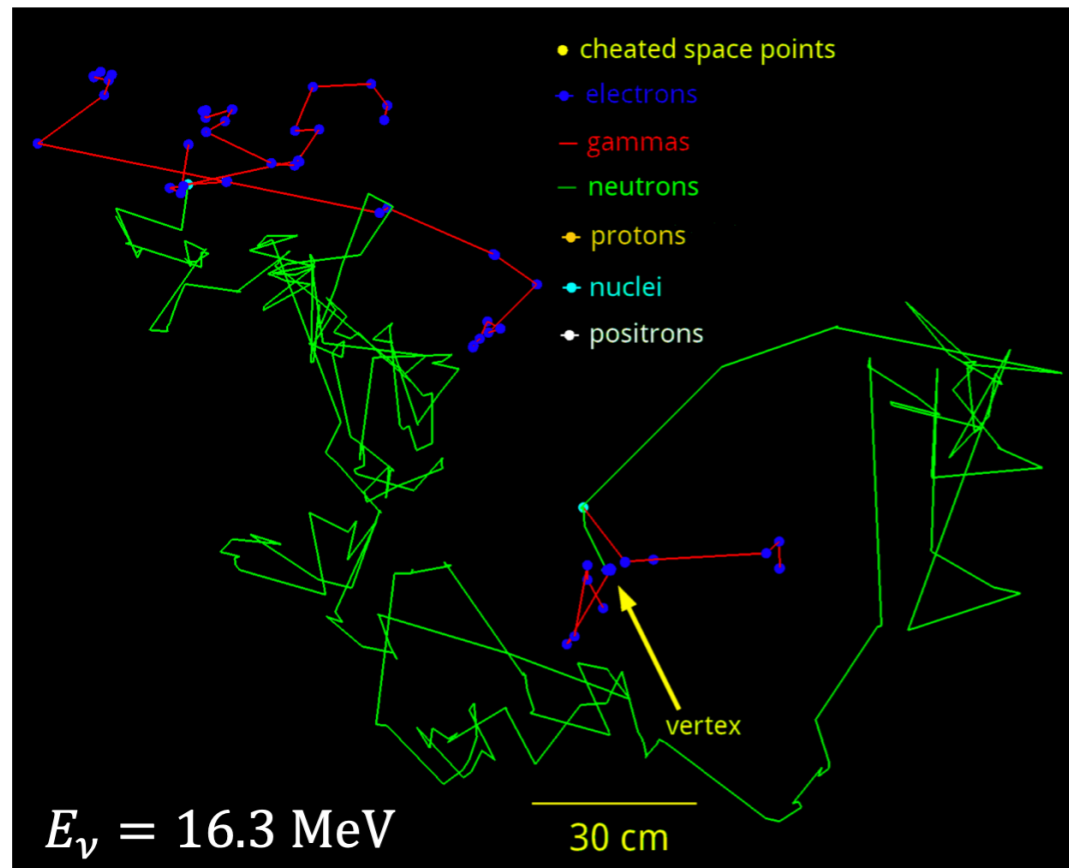
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# 10s of MeV Inelastic Neutrino-Nucleus Scattering: Model in Generators

- Model implementation in MARLEY (low-energy generator used by DUNE) is currently on-going, in collaboration with Steven Gardiner (FNAL) (sole author of MARLEY).



[arXiv:2008.06647 \[hep-ex\] \[DUNE Collaboration\]](https://arxiv.org/abs/2008.06647)

- MARLEY includes allowed approximation (long-wavelength ( $q \rightarrow 0$ ) and slow nucleons ( $p_N/m_N \rightarrow 0$ ) limit), Fermi and Gamow-Teller matrix elements. CRPA includes full expansion of nuclear matrix element (allowed as well as forbidden transition).

- Model recently implemented in GENIE.

[S. Dolan, A. Nikolakopoulos, O. Page, S. Gardiner, N. Jachowicz and V. Pandey, arXiv:2110.14601 \[hep-ex\].](https://arxiv.org/abs/2110.14601)



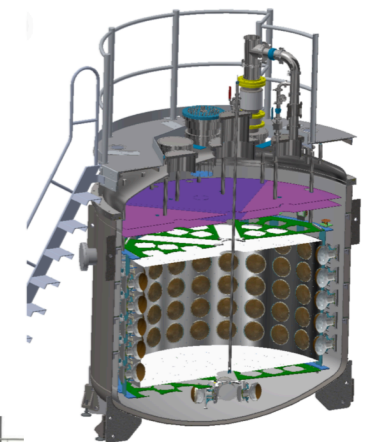
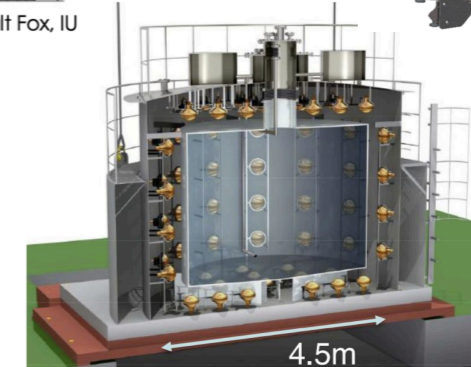
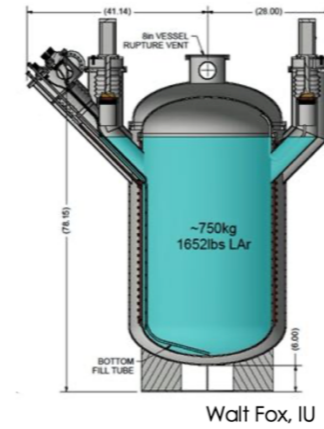
# Low-energy Neutrinos: Near-Future Measurements

- **COHERENT at SNS:** COH-Ar-10 (24kg) LAr detector. COH-Ar-750 (750 kg) LAr detector is underway.

Iodine (NalvE) and Pb, Fe, Cu (NIN cubes) detectors.

- **Coherent CAPTAIN Mills at LANL:** 10 ton LAr detector at Lujan center at LANL. Collected data in 2019 and 2021.

- **JSNS<sup>2</sup> at JPARC-MLF:** 50 ton gd-loaded LS detector.



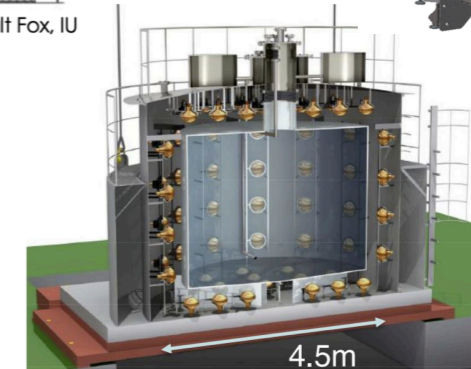
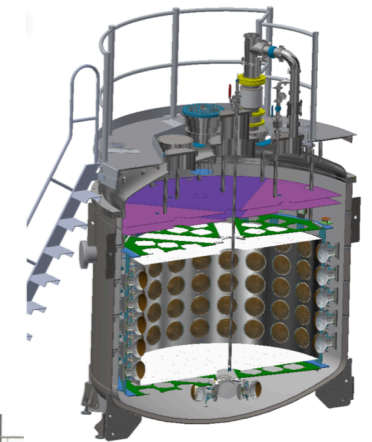
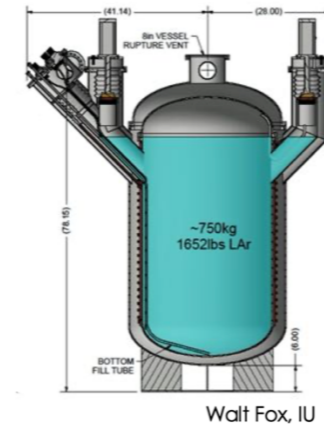
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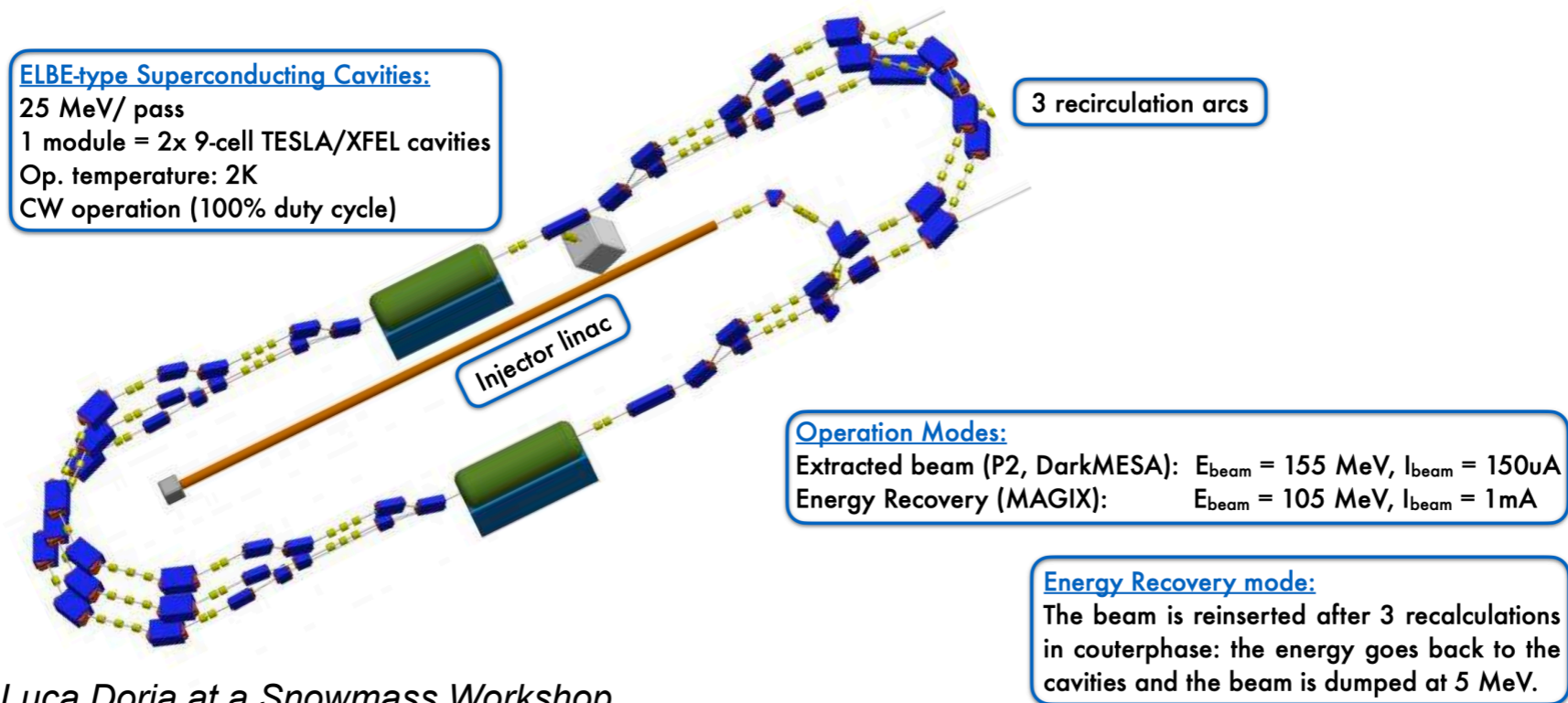
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- **10s of MeV electron scattering experiment is planned at MESA, Mainz:**

## MESA: Mainz Energy-Recovery Superconducting Accelerator



Luca Doria at a Snowmass Workshop

# Mono-energetic KDAR Neutrinos

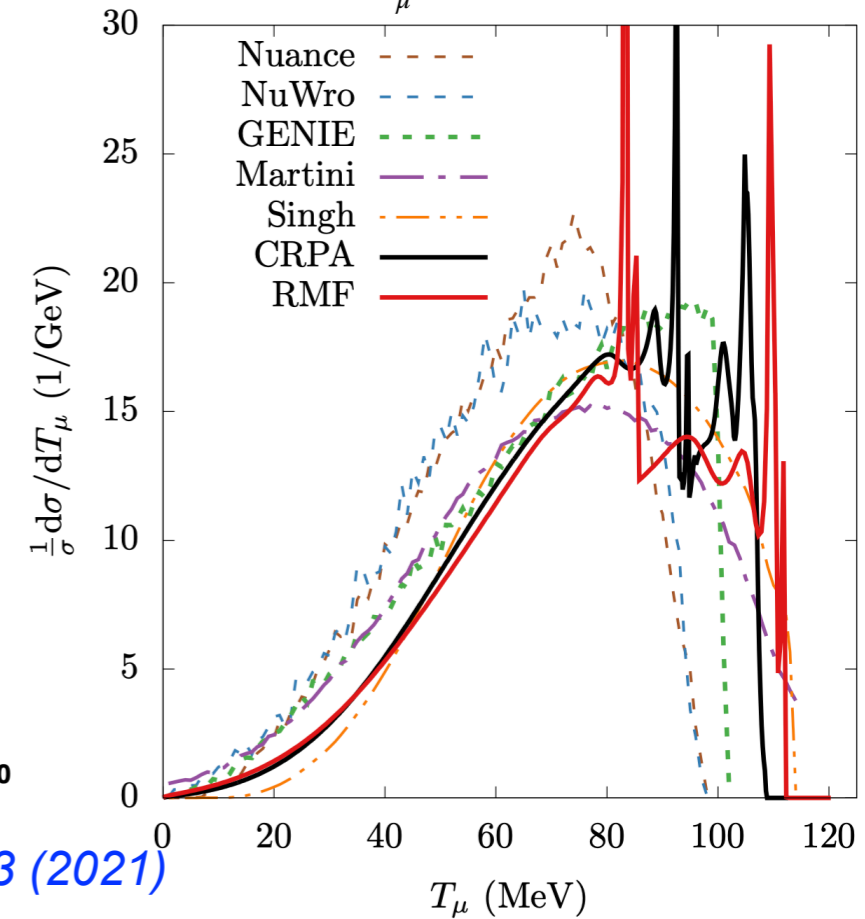
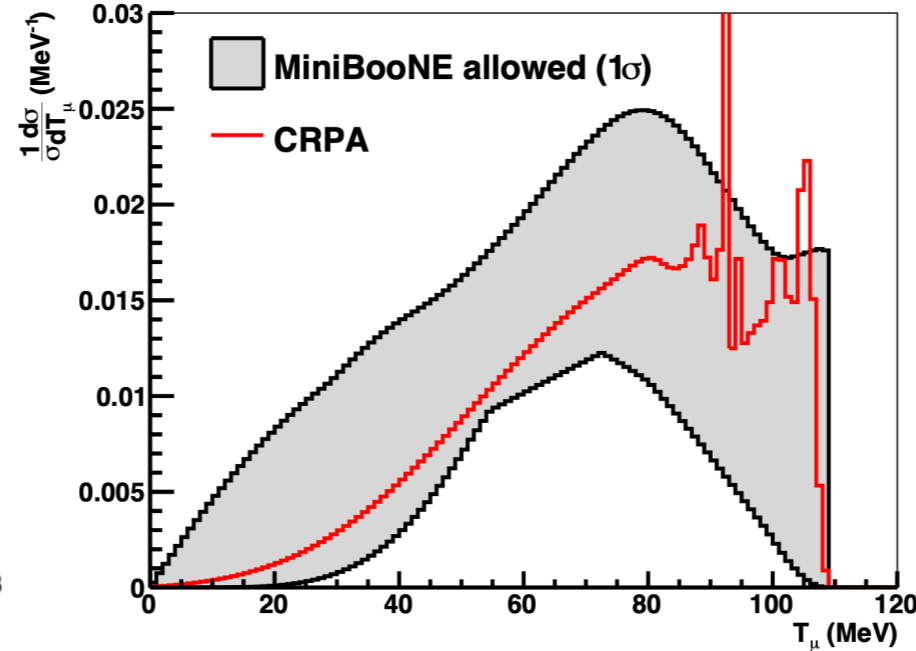
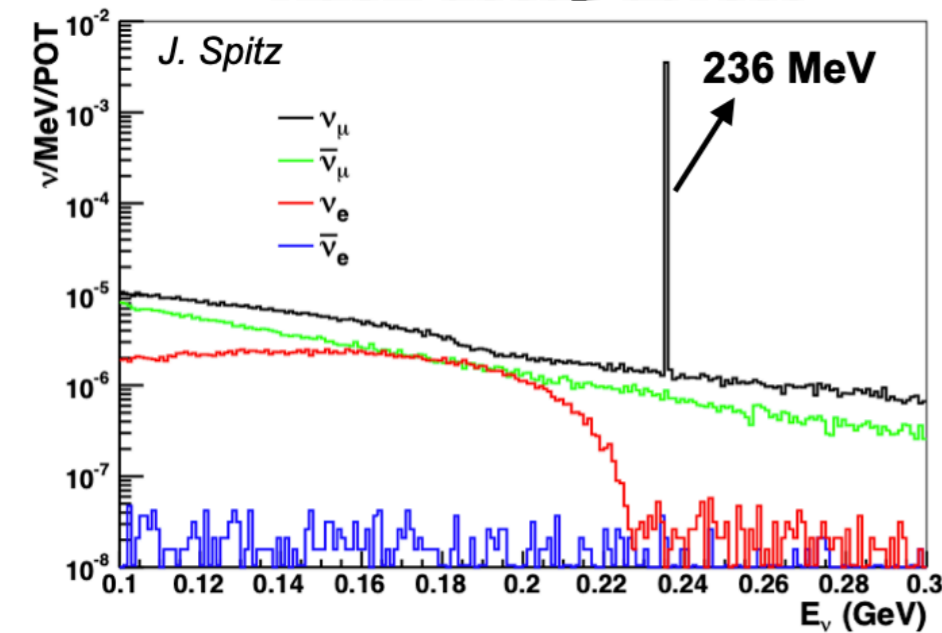
- Mono-energetic KDAR neutrinos at NuMI beam dump (FNAL) and at MLF (JPARC).

$$K^+ \rightarrow \mu^+ \nu_\mu, E_{\nu_\mu} = 236 \text{ MeV}$$

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## Kaon decay at rest

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*A. Nikolakopoulos, V. Pandey, J. Spitz and N. Jachowicz, Phys. Rev. C 103, 064603 (2021)*

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- Exciting near future measurements: MicroBooNE and ICARUS (argon), JSNS<sup>2</sup> at J-PARC (carbon)

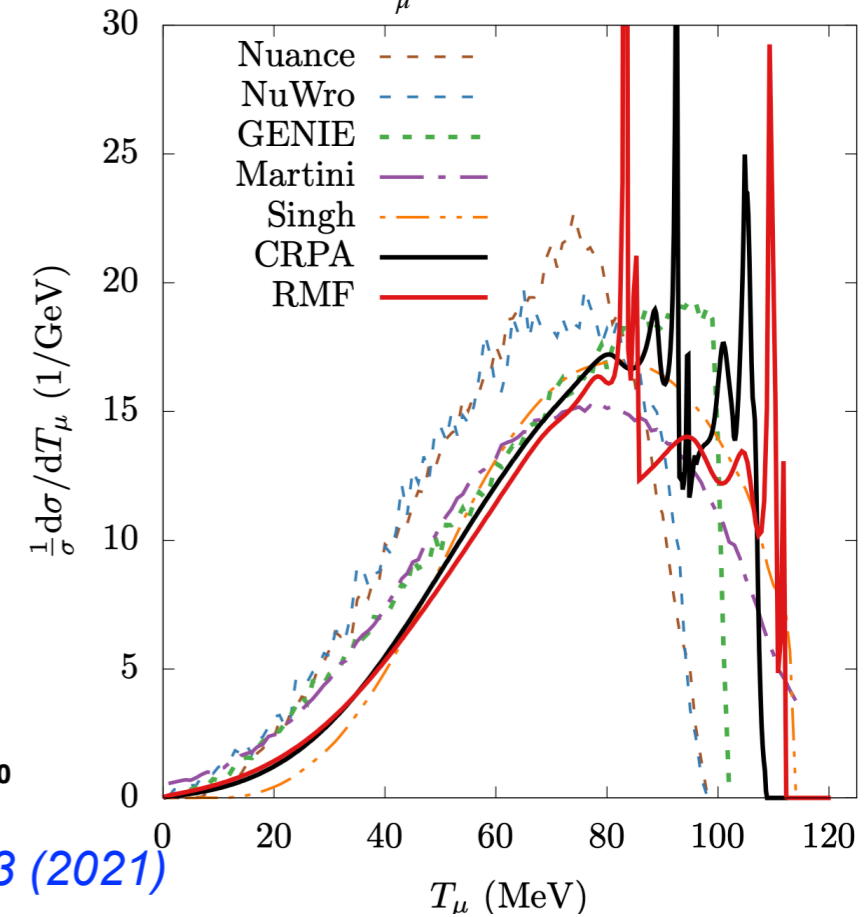
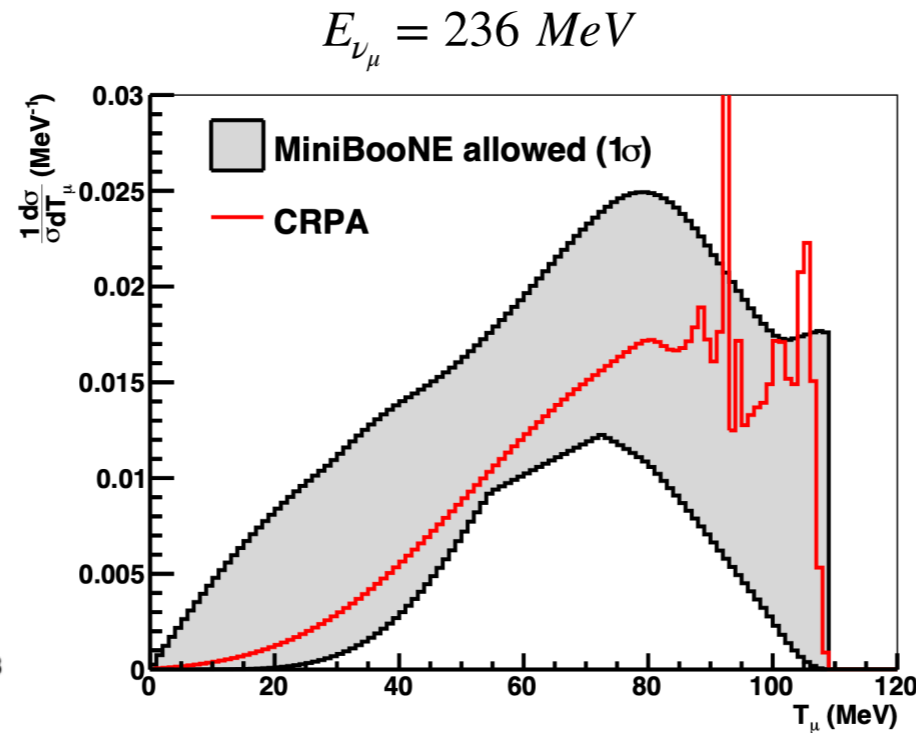
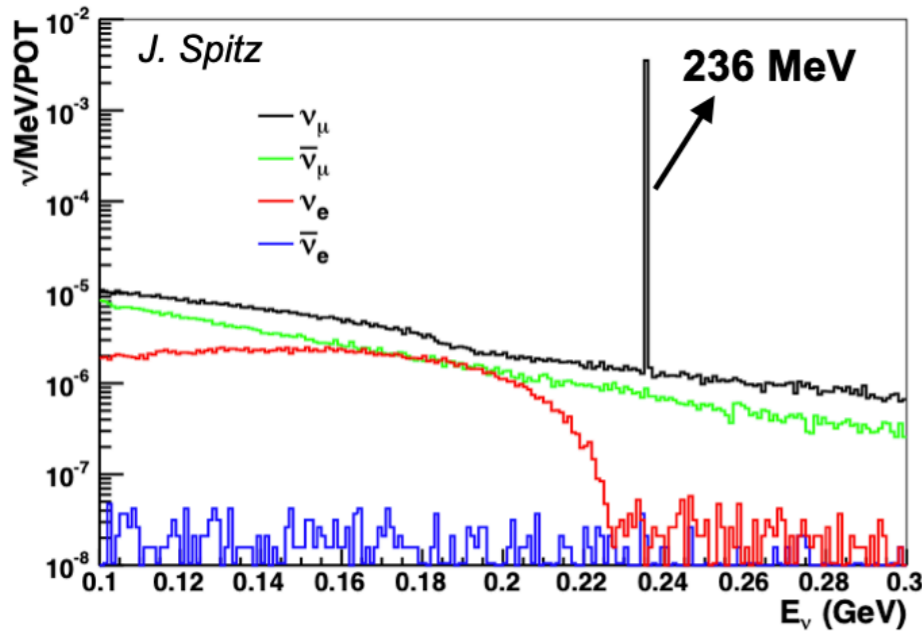
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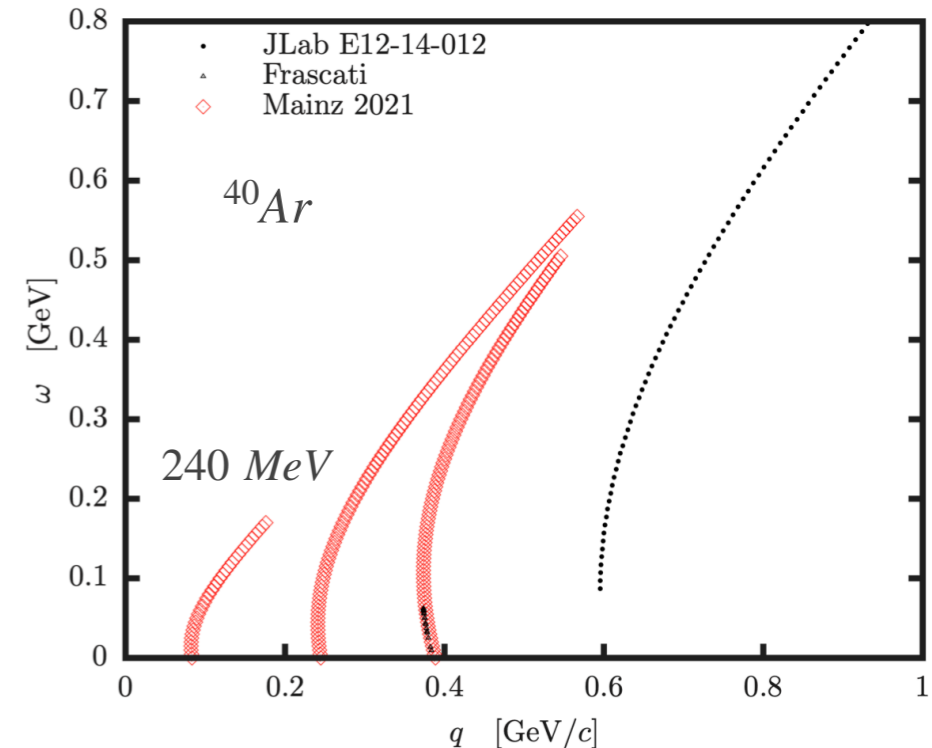
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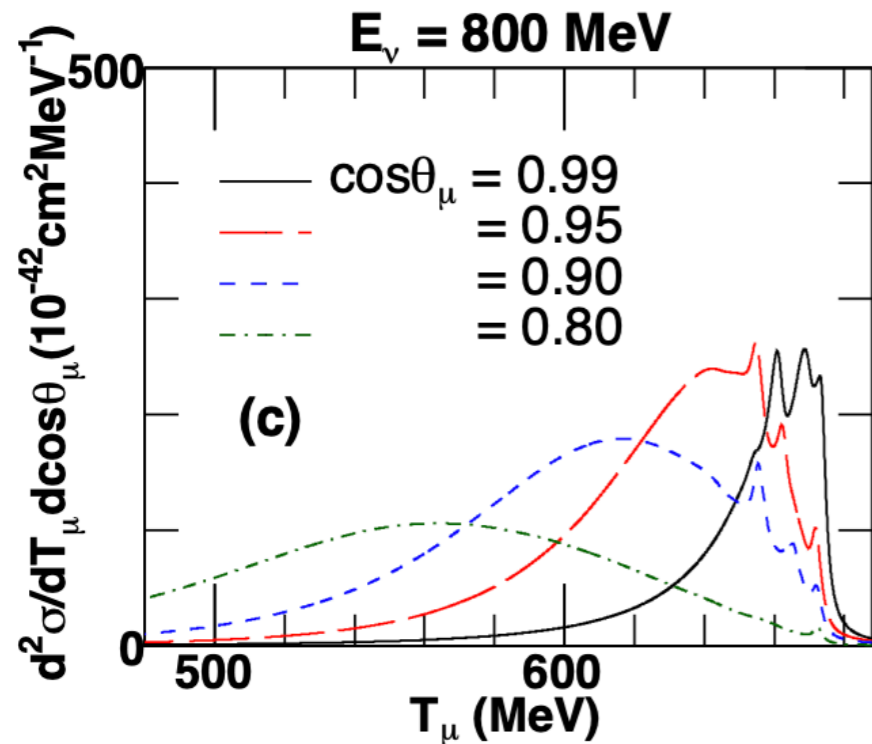
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- Exciting near future measurements: MicroBooNE and ICARUS (argon), JSNS<sup>2</sup> at J-PARC (carbon)
- 240 MeV electron scattering measurement planned at Mainz.
- Combined analysis of mono energetic electron and  $\nu_\mu$  cross sections will give great opportunity to constrain axial response at fixed energy.

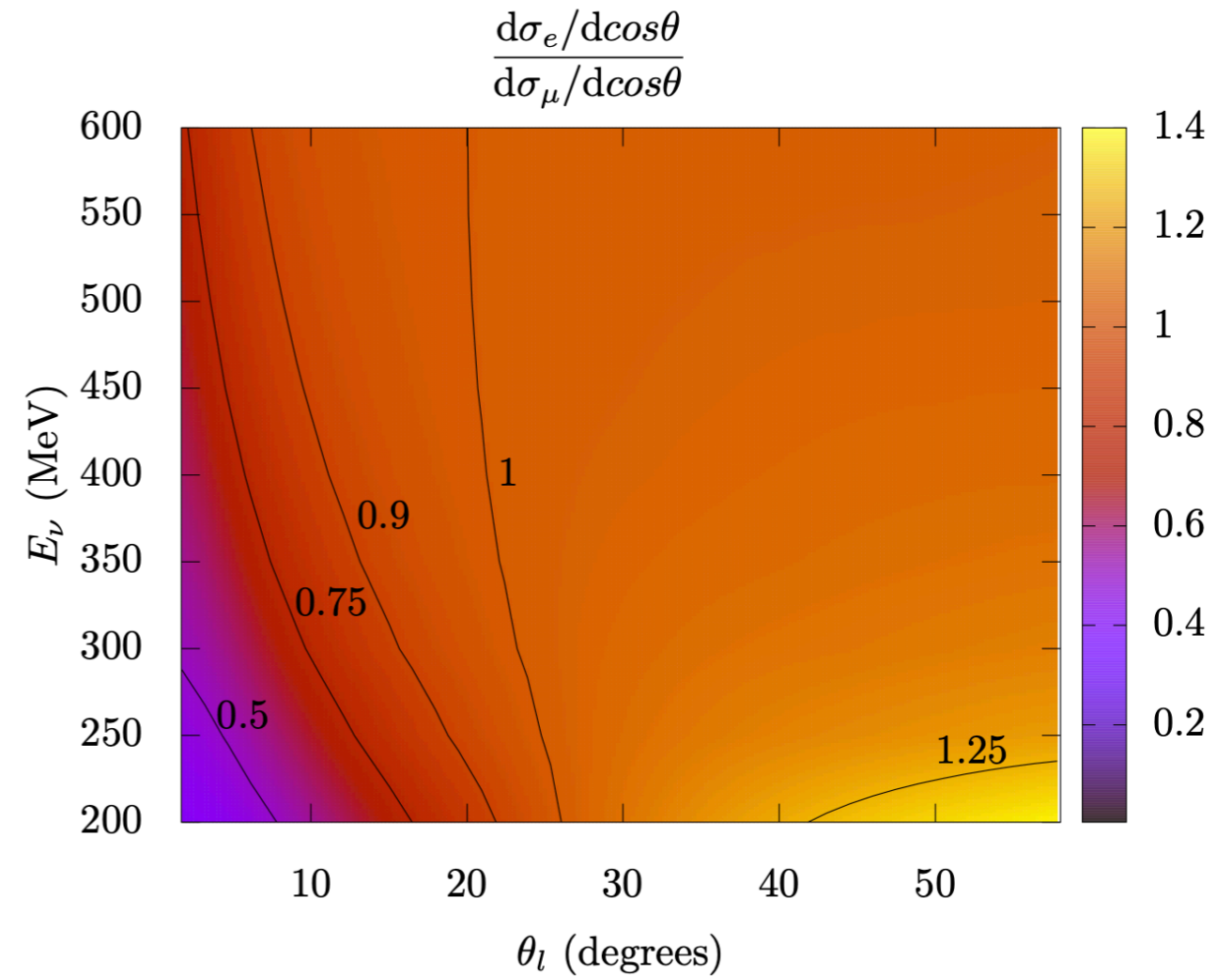




# 10s of MeV Physics in GeV-scale Beams

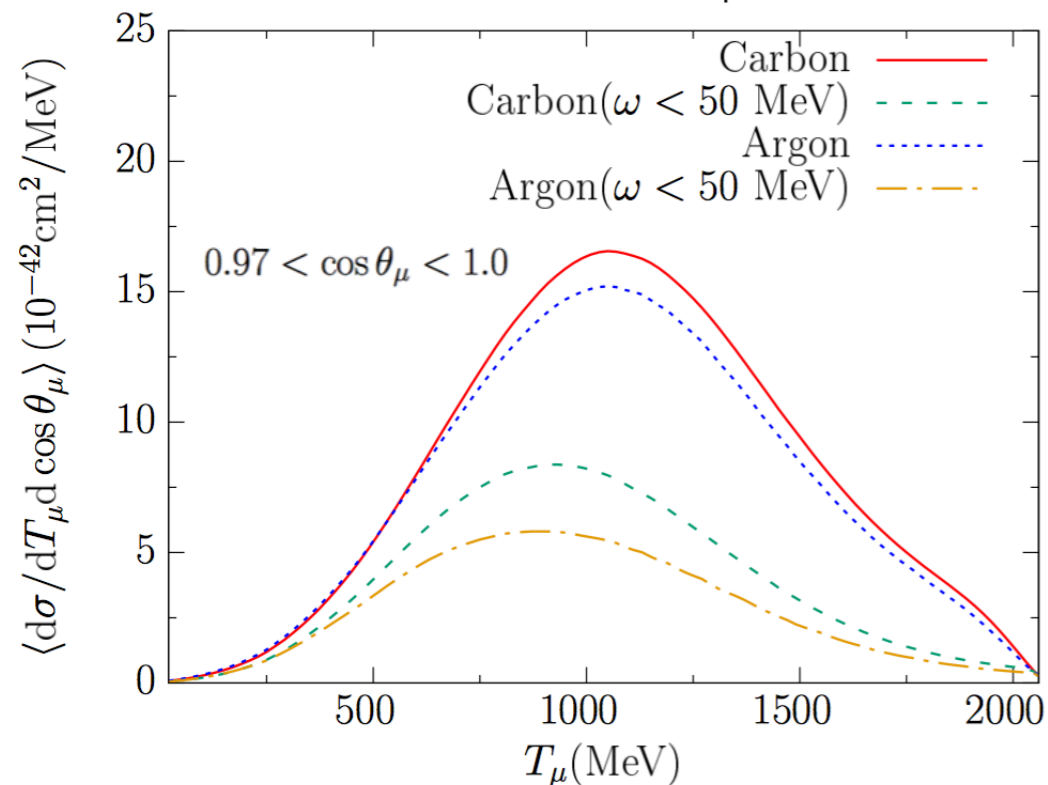


V. Pandey, N. Jachowicz, T. Van Cuyck, J. Ryckebusch, M. Martini, *Phys. Rev. C* **92**, 024606 (2015)



A. Nikolakopoulos, N. Jachowicz, N. Van Dessel, K. Niewczas, R. González-Jiménez, J. M. Udías, V. Pandey, *Phys. Rev. Lett.* **123**, 052501 (2019).

Folded with BNB  $\nu_\mu$  flux



N. Van Dessel, N. Jachowicz, R. González-Jiménez, V. Pandey, T. Van Cuyck, *Phys. Rev. C* **97**, 044616 (2018).

- At forward scattering angles (low momentum transfer), the neutrino-nucleus cross section at GeV-scale energies is impacted by the same nuclear physics effects that are important for the low-energy case more generally.
- At these kinematics, differences between final-state lepton masses become vital and affect the ratio of the charged-current  $\nu_e$  to  $\nu_\mu$  cross sections.



# Neutrino-nucleus Scattering => DM-nucleus Scattering

- Boosted Dark Matter  $\mathcal{L} \supset g_D A'_\mu \bar{\chi} \gamma^\mu \chi + e \epsilon Q_q A'_\mu \bar{q} \gamma^\mu q$   
*B. Dutta, et al., arXiv:2006.09386 [hep-ph]*

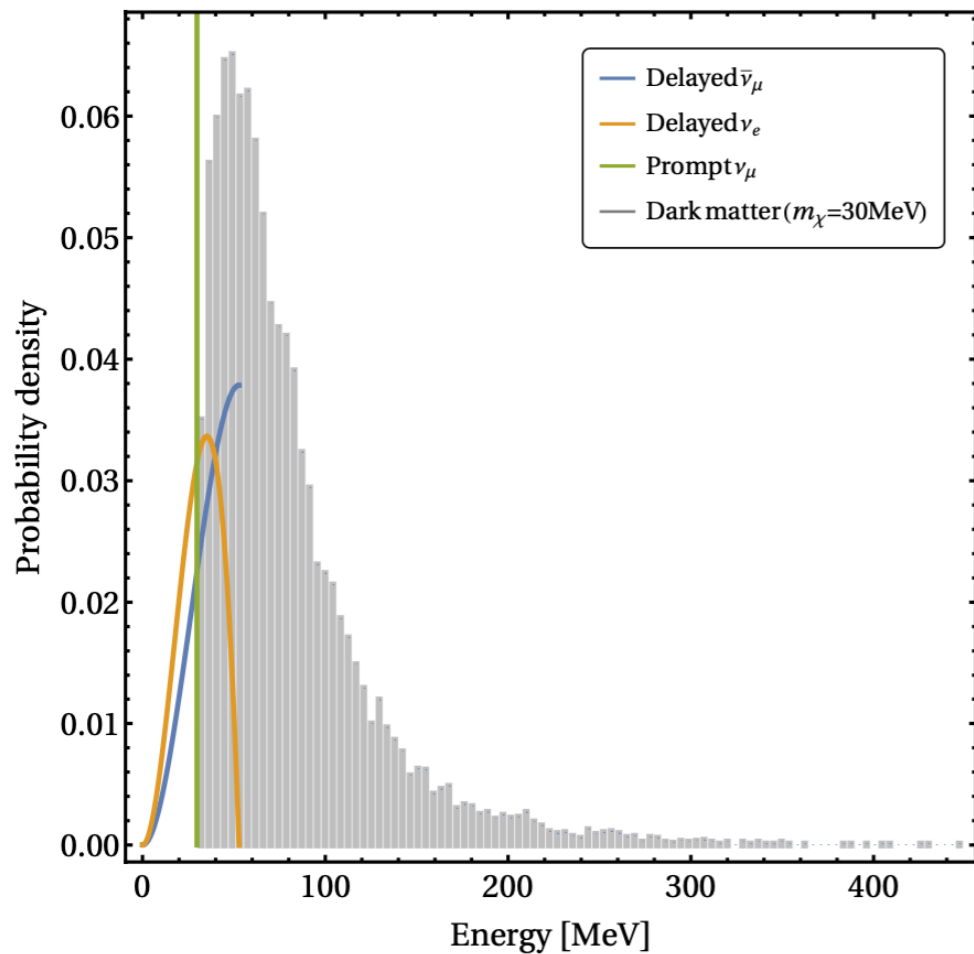
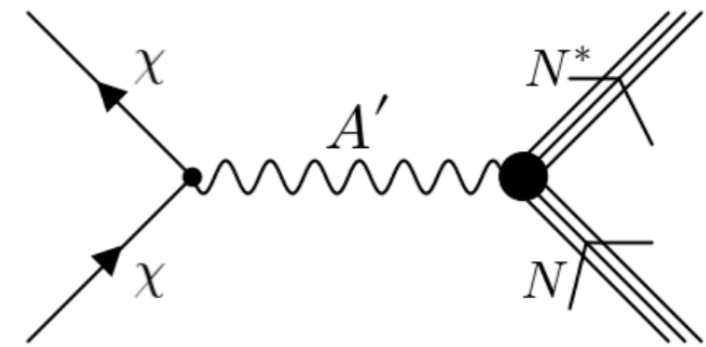
- Dark photon produced in pion decay (e.g. at SNS or at LANL)

$$\pi^- + p \rightarrow n + A'$$

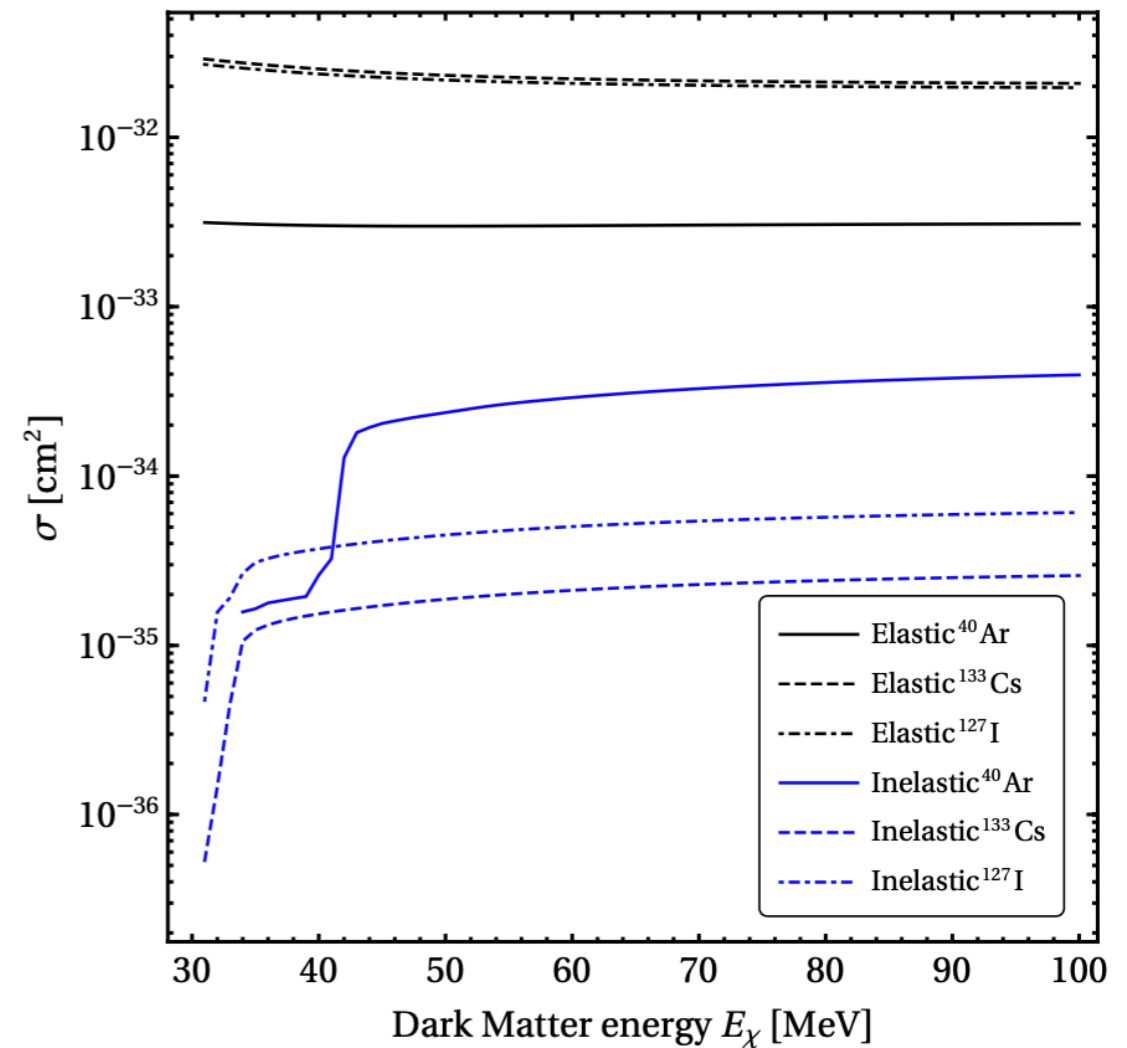
$$\pi^+ + n \rightarrow p + A'$$

$$\pi^0 \rightarrow \gamma + A'$$

- Performing a similar DM-nucleus scattering calculations (dark matter interacting through an  $A'$ ) as for neutrino-nucleus case.



Energy spectra of  $\pi$ -DAR neutrinos and a sample DM spectrum assuming  $m_{A'} = 3m_\chi$



# Summary

- Interactions of low energy (10s of MeV) neutrinos - elastic (CEvNS) and inelastic - are interesting for studies of various SM and BSM processes.
- Neutrino-nucleus interactions at these energies are sensitive to neutron radius and weak elastic form factor (CEvNS), and underlying nuclear structure (inelastic).
- Various neutrino facilities and experiments are sensitive to these processes.

