

CDF II W -mass anomaly faces first-order electroweak phase transition

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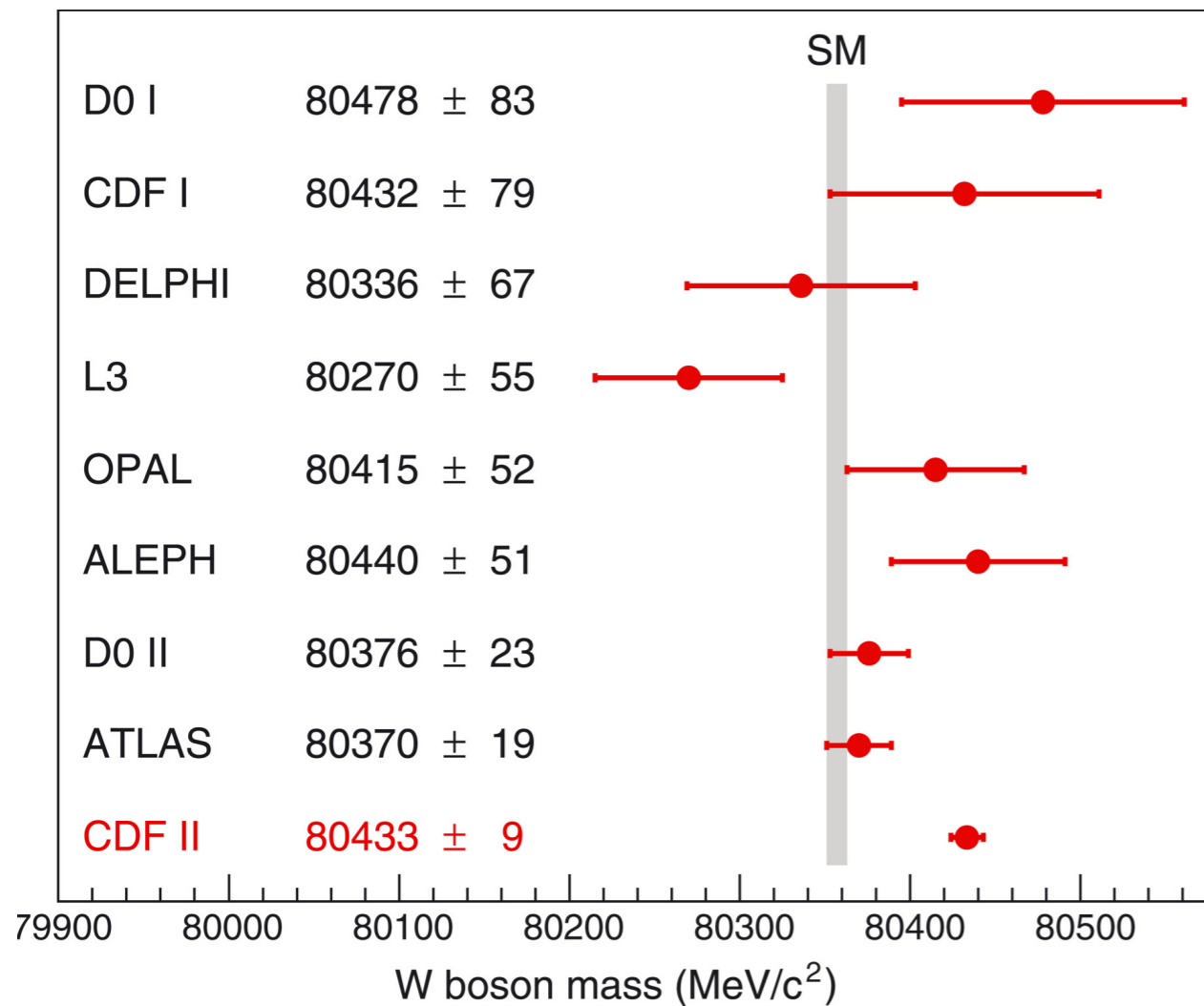
Based on:

A. Addazi, A. Marciano, A. Morais, RP, H. Yang
arXiv: 2204.10315

Fysikdagarna, June 15th 2022

W mass anomaly

**CDF II Collaboration,
Science 376, 170 (2022)**



- Surprising CDF II measurement of W mass lies 7.6σ away from the Standard Model
- Many scenarios beyond the SM have been deployed in the literature to explain this measurement (over 100 publications so far!)
- A large class of BSM scenarios offering such an explanation features the existence of a new SU(2) adjoint (triplet) scalar which provides a tree-level corrections to the SM W mass value
- Existence of such scalars may impact the Electro Weak phase transition in early Universe, possibly rendering such models testable in future gravitational-wave detectors

A minimal scalar SU(2) triplet extension

Let us focus on a simplified framework that relates the characteristics of EW phase transitions to a possible explanation of the W mass anomaly

The model by: L. Di Luzio, R. Gröber and P. Paradisi,
"Higgs physics confronts the m_W anomaly"
[arXiv:2204.05284 [hep-ph]]

W mass anomaly



Anomaly in T-parameter
(assuming $U=0$)

$$\hat{T} \simeq (0.84 \pm 0.14) \times 10^{-3}$$

A. Strumia, [arXiv:2204.04191 [hep-ph]]

Effective Field Theory
d=6 operator generates
this anomaly

$$c_{\text{HD}} \mathcal{O}_{\text{HD}} = c_{\text{HD}} (H^\dagger D_\mu H) ((D_\mu H)^\dagger H)$$

$$\hat{T} = -\frac{v^2}{2} c_{\text{HD}}$$

A new heavy state that
generates this operator

$$\Delta = (1, 3, 0) \quad \mathcal{L}_{\text{T}} = -k_\Delta H^\dagger \Delta \cdot \sigma H + \text{h.c.}$$

$$c_{\text{HD}} = -2 \frac{k_\Delta^2}{M_\Delta^4}$$

negative effective
coupling!

Scale of New Physics

Integrating out heavy new scalar triplet state yields both:
a positive contribution to the T-parameter and a modification of the Higgs potential

$$\hat{T} = \frac{k_{\Delta}^2 v^2}{M_{\Delta}^4} = 0.84 \times 10^{-3} \left(\frac{|k_{\Delta}|}{M_{\Delta}} \right)^2 \left(\frac{8.5 \text{ TeV}}{M_{\Delta}} \right)^2$$

consistent with the observed shift in W mass for a TeV-scale scalar triplet state!

Saturating the perturbativity bound $|k_{\Delta}|/M_{\Delta} \leq 4\pi$ the mass scale cannot exceed 100 TeV

Higgs quartic couplings receives a tree-level correction

$$\lambda = \lambda_{\text{bare}} + (k_{\Delta}/m_{\Delta})^2$$

$$\lambda = m^2/2v^2 \quad \begin{array}{l} m^2 \text{ Higgs mass parameter} \\ v \simeq 246 \text{ GeV Higgs VEV} \end{array}$$

d=6 Higgs self-interaction term is induced as well

$$\frac{\mu_{\Delta}}{3} \Delta^3 + \text{h.c.} \quad \Delta^3 \equiv (\Delta \cdot \sigma)(\Delta \cdot \sigma)(\Delta \cdot \sigma)$$



crucial contribution to the Higgs potential that determines the nature and the strength of the EW phase transition



$$\frac{\kappa}{\Lambda^2} (H^{\dagger} H)^3 + \text{h.c.}$$

cutoff scale

$$\frac{\Lambda}{\sqrt{\kappa}} = \frac{\sqrt{3} M_{\Delta}^3}{\sqrt{\mu_{\Delta}} k_{\Delta}^{3/2}}$$

$$\kappa \lesssim 4\pi$$

Dynamics of EWPTs

- High $T \rightarrow$ classical motion in Euclidean space described by action \hat{S}_3

$$\hat{S}_3 = 4\pi \int_0^\infty dr r^2 \left\{ \frac{1}{2} \left(\frac{d\hat{\phi}}{dr} \right)^2 + V_{\text{eff}}(\hat{\phi}) \right\},$$

- Effective potential: loop and thermal corrections

$$V_{\text{eff}}^{(1)}(\hat{\phi}) = V_{\text{tree}} + V_{\text{CW}} + \Delta V^{(1)}(T)$$

$$V_{\text{CW}} = \sum_i (-1)^F n_i \frac{m_i^4}{64\pi^2} \left(\log \left[\frac{m_i^2(\hat{\phi}_\alpha)}{\Lambda^2} \right] - c_i \right)$$

$$\Delta V^{(1)}(T) = \frac{T^4}{2\pi^2} \left\{ \sum_b n_b J_B \left[\frac{m_b^2(\hat{\phi}_\alpha)}{T^2} \right] - \sum_f n_f J_F \left[\frac{m_f^2(\hat{\phi}_\alpha)}{T^2} \right] \right\},$$

- $\hat{\phi} \rightarrow$ solution of the e.o.m. found by the path that minimizes the energy.

$$\Delta V^{(1)}(T)|_{\text{L.O.}} = \frac{T^2}{24} \left\{ \text{Tr} [M_{\alpha\beta}^2(\phi_\alpha)] + \sum_{i=W,Z,\gamma} n_i m_i^2(\phi_\alpha) + \sum_{i=t,b,\tau} \frac{n_i}{2} m_i^2(\phi_\alpha) \right\}$$

Finite-T effective potential & EW FOPTs

In unitary gauge, one-loop effective Higgs potential:

$$V_{\text{eff}}(T, h) = V_{\text{tree}}(h) + V_{T=0}^{(1)}(h) + \Delta V_T(h, T)$$

$$V_{\text{tree}}(h) = \frac{1}{2}m^2 h^2 + \frac{\lambda}{4}h^4 + \frac{\kappa}{8\Lambda^2}h^6$$

The dominant thermal correction to the Higgs mass:

$$CT^2/2$$

$$C \simeq \frac{1}{16} \left(g'^2 + 3g^2 + 4y_t^2 + 4\frac{m_h^2}{v^2} + 36\frac{\kappa v^2}{\Lambda^2} \right)$$

$$m_h^2 = 2\lambda v^2 + 3v^4\kappa/\Lambda^2$$

zero-T Higgs mass
 $m_h = 125 \text{ GeV}$

Limit on the cutoff scale imposed by the strongly 1st order EW phase transition requirement yields:

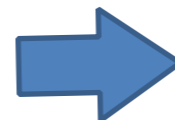
$$480 \text{ GeV} \lesssim \frac{\sqrt{3}M_\Delta^3}{\sqrt{\mu_\Delta}k_\Delta^{3/2}} \lesssim 840 \text{ GeV}$$

F. Huang et al, Phys. Rev. D94 (2016) 041702
 [arXiv:1601.01640 [hep-ph]]

$$v(T_c)/T_c > 1$$

For saturated perturbativity

$$|k_\Delta|/M_\Delta \simeq 4\pi \text{ and } |\mu_\Delta|/M_\Delta \simeq 4\pi$$



The EW FOPT bound:

$$M_\Delta \simeq 5 \div 10 \text{ TeV}$$

EW FOPTs & GWs characteristics

Inverse time-scale of the PTs:

$$\frac{\beta}{H} = T_* \left. \frac{\partial}{\partial T} \left(\frac{\hat{S}_3}{T} \right) \right|_{T_*}$$

Relative latent heat (PT strength):

$$\alpha = \frac{1}{\rho_\gamma} \left[V_i - V_f - \frac{T_*}{4} \left(\frac{\partial V_i}{\partial T} - \frac{\partial V_f}{\partial T} \right) \right]$$

$$\rho_\gamma = g_* \frac{\pi^2}{30} T_*^4$$

Probability to find a point in the false vacuum:

$$P(T) = e^{-I(T)},$$

$$I(T) = \frac{4\pi v_b^3}{3} \int_T^{T_c} \frac{\Gamma(T') dT'}{T'^4 H(T')} \left(\int_T^{T'} \frac{d\tilde{T}}{H(\tilde{T})} \right)^3$$

Percolation temperature
(temperature at which at least 34%
of the false vacuum has tunnelled
into the true vacuum)

$$I(T_*) = 0.34$$

J. Ellis, M. Lewicki, and V. Vaskonen, *Journal of Cosmology and Astroparticle Physics* **2020**, 020–020 (2020).

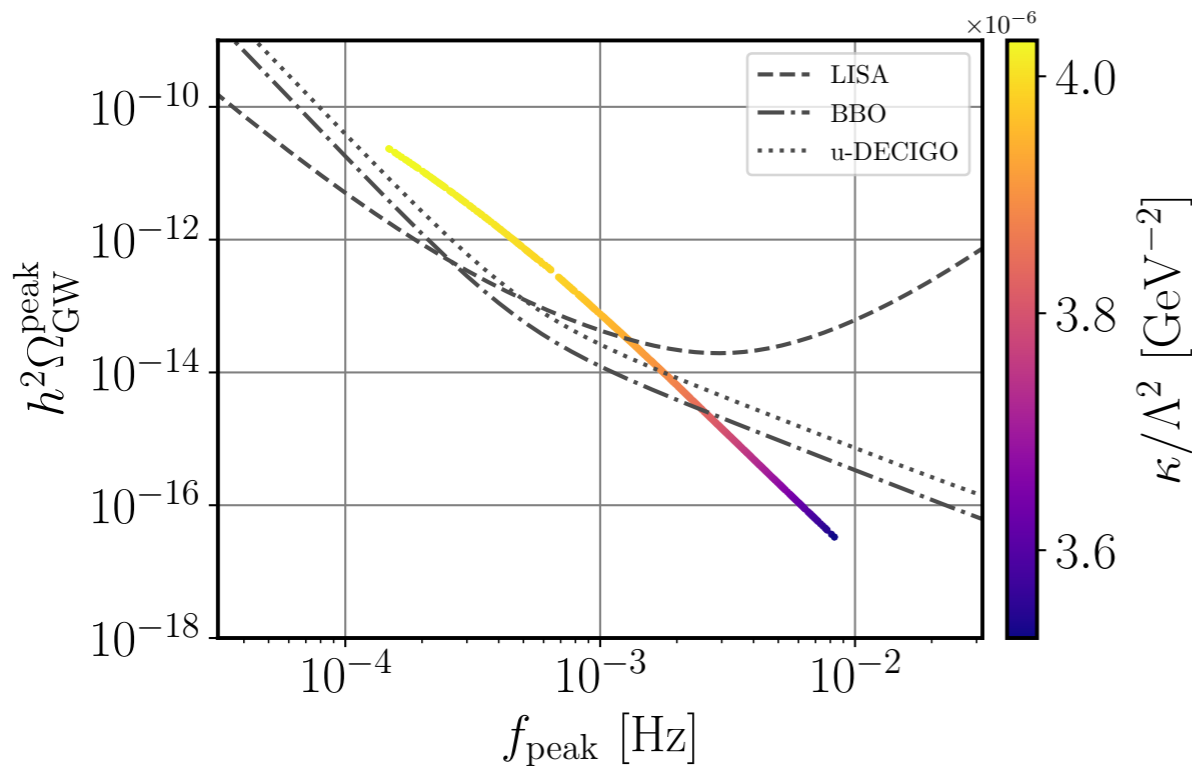
Primordial GWs power spectrum:

$$h^2 \Omega_{\text{GW}} = h^2 \Omega_{\text{GW}}^{\text{peak}} \left(\frac{4}{7} \right)^{-\frac{7}{2}} \left(\frac{f}{f_{\text{peak}}} \right)^3 \left[1 + \frac{3}{4} \left(\frac{f}{f_{\text{peak}}} \right) \right]^{-\frac{7}{2}}$$

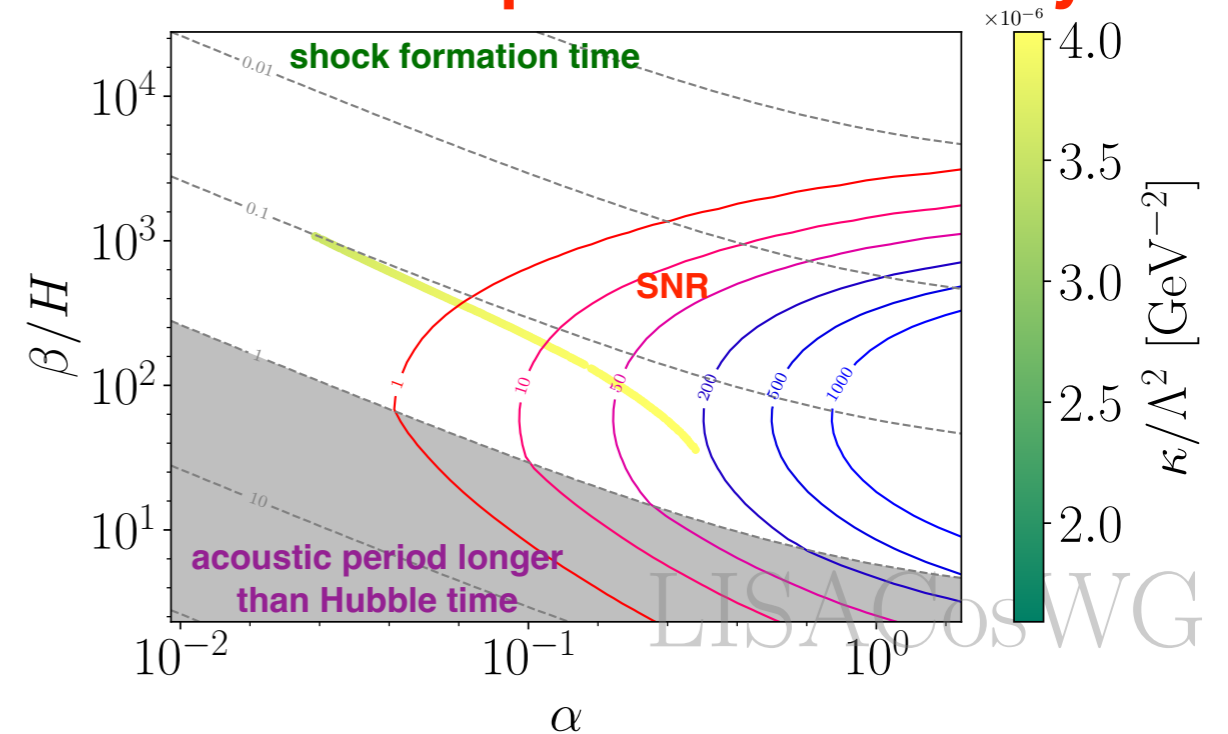
C. Caprini *et al.*, *JCAP* **2003**, 024 (2020), 1910.13125

Primordial GWs in a minimal triplet model

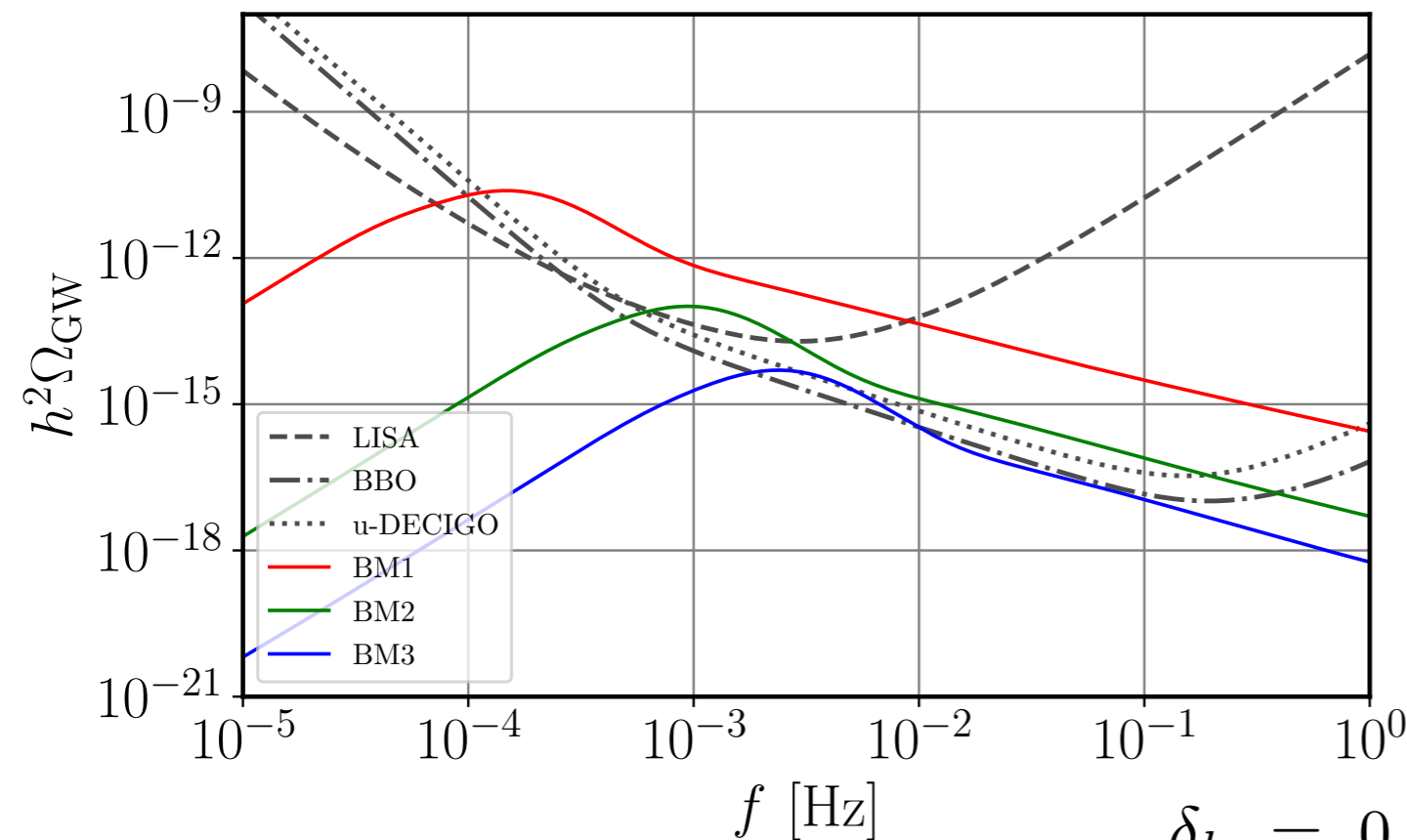
peak-amplitude vs energy scale



duration & SNR for points detectable by LISA



GWs power spectrum



Benchmarks:

T_* (GeV)	α	β/H_*	$\kappa^{-1/2} \Lambda$ (GeV)	$\delta_{\sigma_{hz}}$ (%)
43.8	0.30	36.37	498.12	1.8
55.6	0.12	180.94	502.40	2.1
64.2	0.07	394.14	508.38	2.2



- Consistent with LHC bounds
- Can be probed in future measurements of trilinear Higgs coupling:

$$\lambda_{3h} = -(1 + \delta_h) \frac{Ah^3}{6} \quad A = 3m_h^2/v$$

$$\delta_h = 0.66 \div 2 \quad \longleftrightarrow \quad \delta_{\sigma_{hz}} = \delta\sigma_{hz}/\sigma_{hz}$$

Summary

- **Recent W mass measurement offers new opportunities for New Physics searches**
- **A simple model with heavy scalar triplet provides potentially observable new signatures (W mass correction, triple-Higgs coupling, FOPTs & GWs) and addresses some of the fundamental questions (e.g. neutrino mass)**
- **Primordial gravitational waves represent a complimentary source of information to the collider measurements (such as HE-LHC and Circular e^+e^- Collider)**
- **Combining W mass, future measurements of triple Higgs coupling and primordial GWs would provide strong case of probing such a class of models BSM**