Beam Loading



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CERN



Advanced Accelerator Physics

9 November 2022

Outline

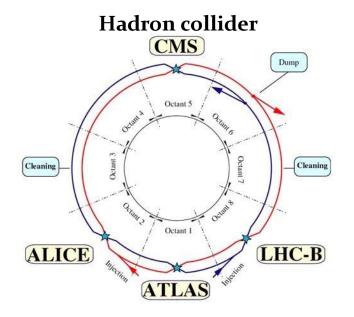
- Introduction
- RF cavity parameters
 - Shunt impedance, beam loading, power coupling
- Fundamental theorem of beam loading
- Passage of a bunches through a cavity
 - Single passage or bunches with large spacing
 - Multiple bunch passages

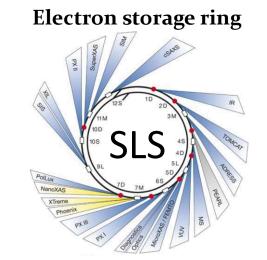
• Steady state beam loading and partial filling

- Few bunches with large spacing
- Summary

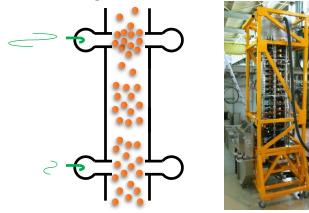
Introduction

What do these devices in common?





Klystron amplifier



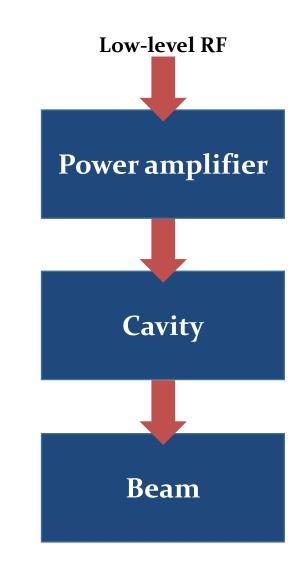
Microwave oven



→ They all **suffer** from or **make use** of beam loading

Introduction

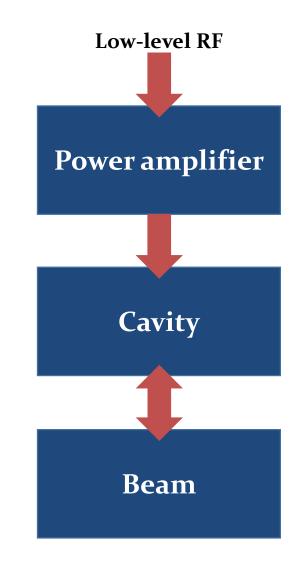
- The radiofrequency (RF) system should provide
 → Energy to the beam
 →Longitudinal focusing
- Intended energy flow usually from cavity to beam
- But beam also likes to influence the field in the cavity



Introduction

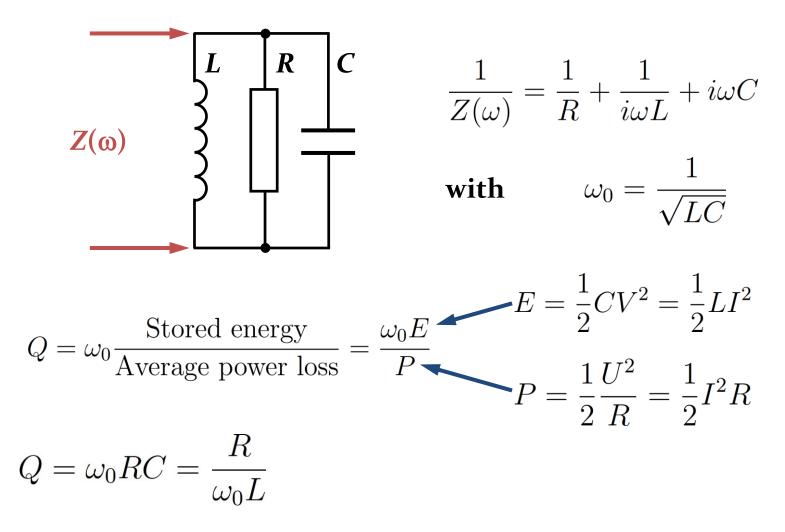
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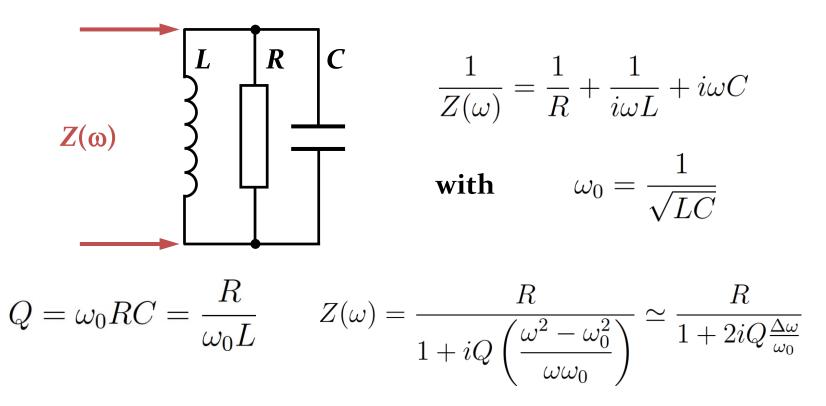


RF cavity

• The resonance of a cavity can be understood as simple parallel resonant circuit described by *R*, *L*, *C*

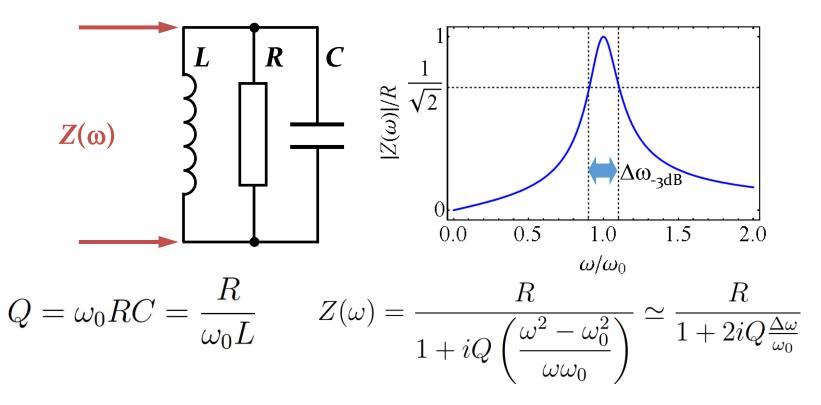


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→ Resonant circuit can also be described by R, R/Q, ω_o or any other set of three parameters

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→ Resonant circuit can also be described by R, R/Q, ω_o or any other set of three parameters

- Most common choice by cavity designers ω_0 , R, R/Q why?
- **Resonance frequency**, ω_o

 \rightarrow Exactly defined for given application, e.g. $h\omega_{rev}$

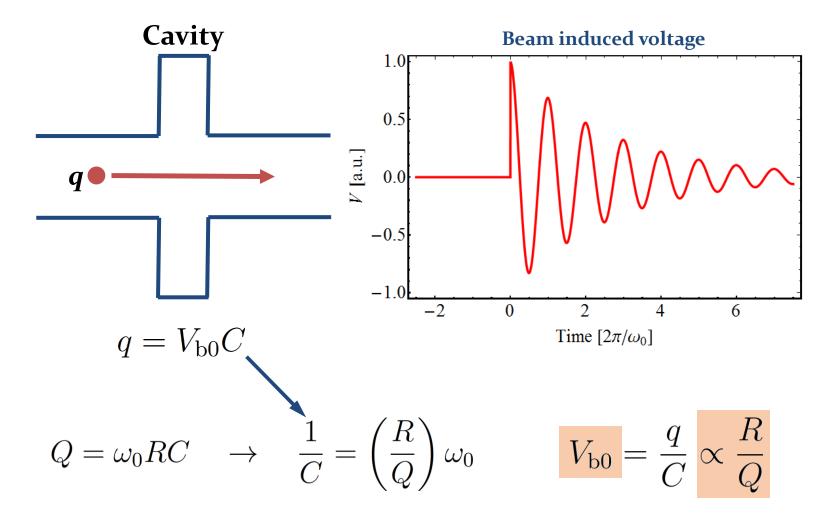
• Shunt impedance, *R*

→ Power required to produce a given voltage without beam

- **"R-upon-Q"**, *R/Q*
 - \rightarrow Defined only by the cavity geometry
 - \rightarrow Criterion to optimize a geometry
 - \rightarrow Detuning with beam proportional to R/Q

Why R/Q?

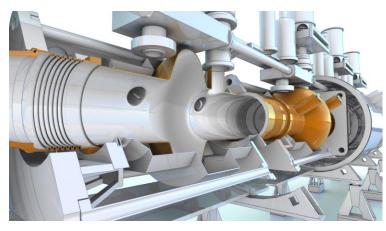
→ Charged particle experiences cavity gap as capacitor



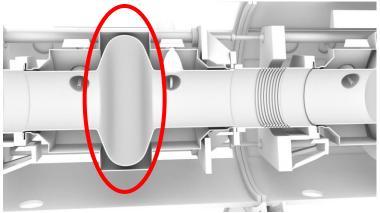
 \rightarrow Cavity geometry with small R/Q to reduce beam loading

Example: 400 MHz cavities in LHC

- → **Reduce beam loading in RF cavities**
- → Shunt impedance, R, low for small R/Q with normal conducting cavities → superconducting cavities in LHC

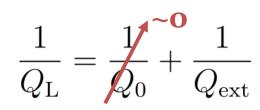


Bell shape: $R/Q \sim 44 \Omega$, 400 MHz



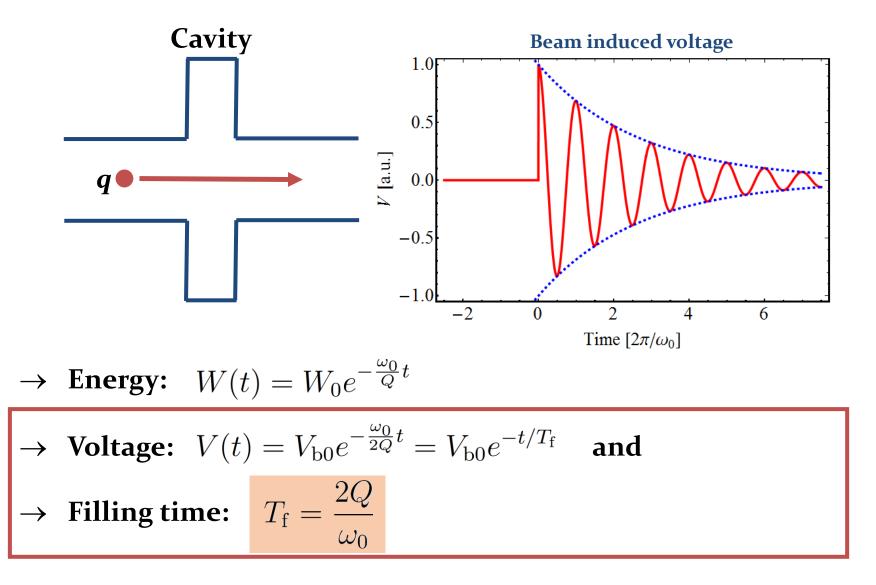


 \rightarrow 2×8 cavities, 5.3 MV/m

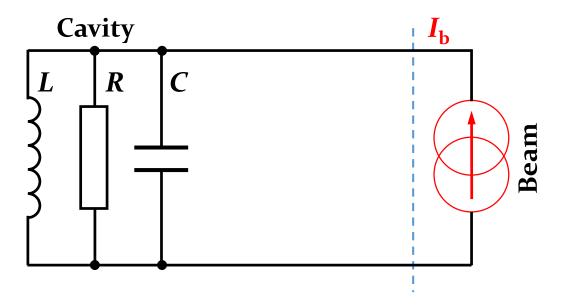


Field decay in cavity

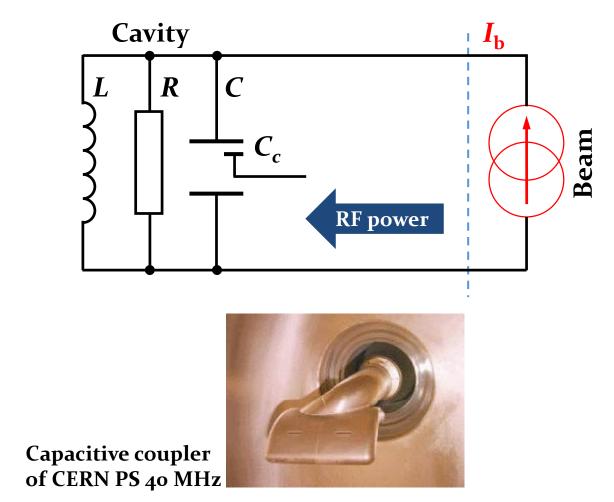
→ After passage of charge: energy and fields decay exponentially



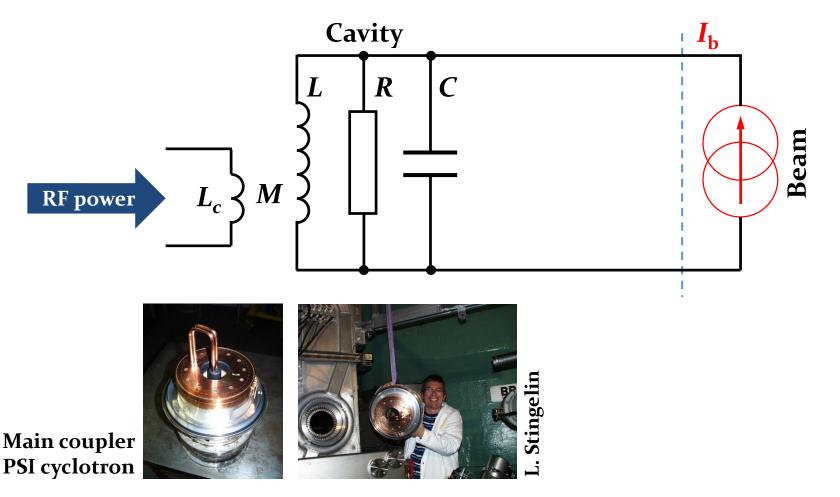
- Connection of cavity to power amplifier
 - \rightarrow **Capacitive:** Capacitor coupling electrically to the gap
 - → **Inductive:** Coupling loop in region of large magnetic field



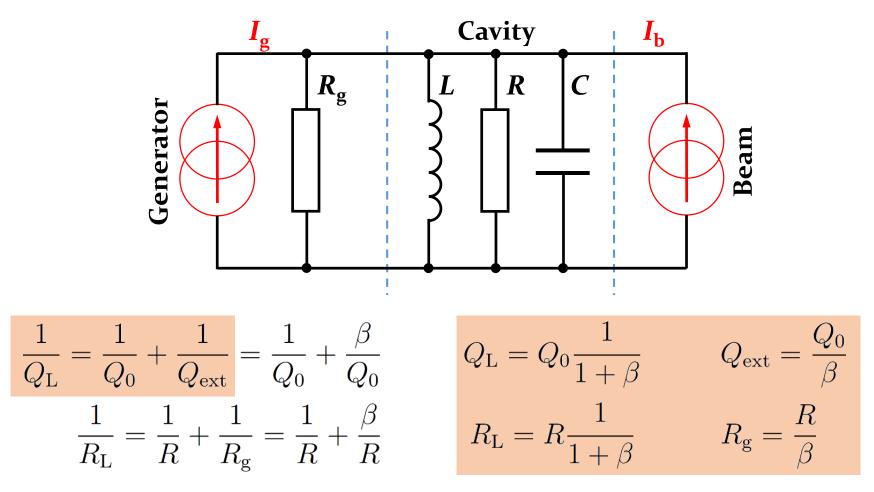
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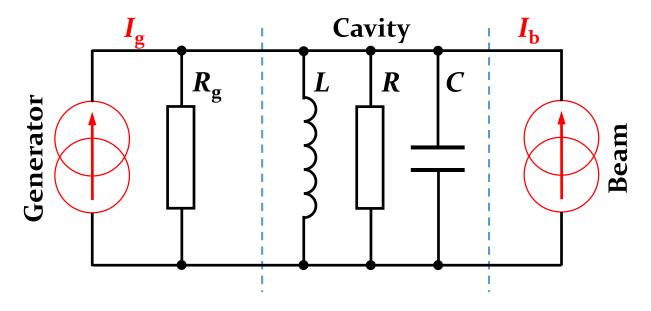


- \rightarrow Output impedance loads the resonant circuit: $R_{g} || R$
- → **Reduction of quality factor:**
- \rightarrow Coupling coefficient, β , defines coupling ratio



 $Q_0 \rightarrow Q_L$

- \rightarrow Output impedance loads the resonant circuit: $R_{g} \parallel R$
- → **Reduction of quality factor:**
- \rightarrow Coupling coefficient, β , defines coupling ratio



- 1. Generator output impedance is not a physical resistor
 - \rightarrow Generator does not experience own output impedance
- 2. Beam experiences output impedance of generator as resistor
 - $\rightarrow R_{g} \mid \mid R$ relevant for beam loading

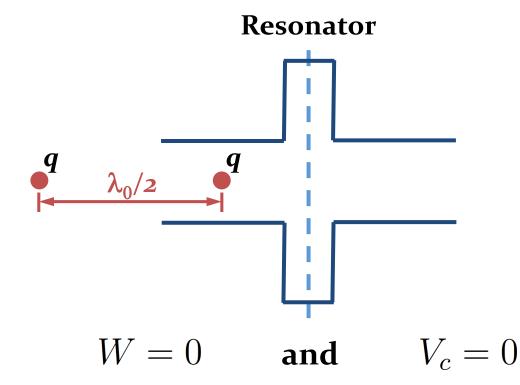
 $Q_0 \rightarrow Q_1$

Fundamental theorem of beam loading

Initially empty cavity

Which fraction does a charge experience of its induced voltage?

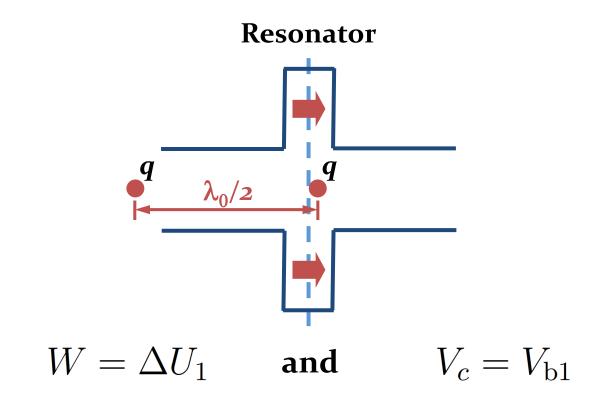
- Equal charges passing through cavity at distance $\lambda_o/2 = \pi c_o/\omega_o$
- → Principles: energy conservation and superposition



After passage of first charge

- 1st charge passes through the cavity and induces voltage
- Fraction, *r* describes part of induced voltage affecting itself:

$$\Delta U_1 = r \cdot q V_{\rm b1}$$

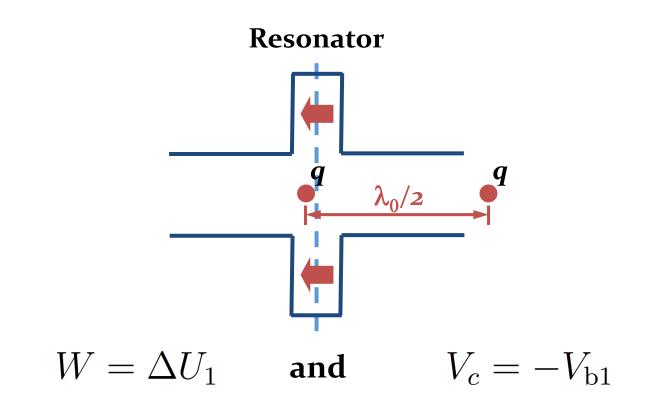


Before passage of 2nd charge

• 2nd charge passes through the cavity

 \rightarrow

• Affected by induced field of 1st charge



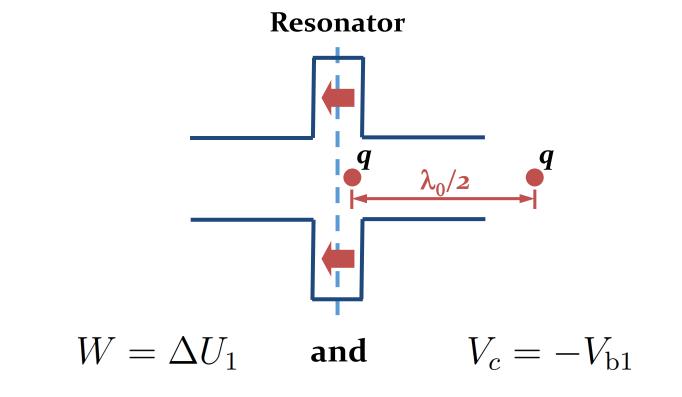
Passage of 2nd charge

• 2nd charge passes through the cavity

 \rightarrow

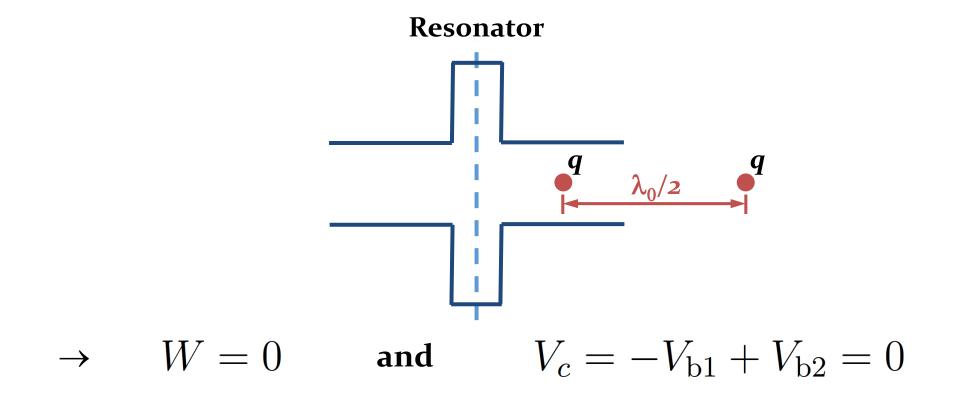
• Affected by induced field of 1st charge and its own induced

$$\Delta U_2 = -qV_{\rm b1} + r \cdot qV_{\rm b2}$$



After passage of first bunch

- After passage of 2nd charge through the cavity
- \rightarrow Takes the same energy as brought into cavity by 1st charge



Ratio of induced field

→ Total energy brought in and taken out of cavity must be zero

$$\Delta U_1 + \Delta U_2 = 0$$

$$r \cdot qV_{b1} - qV_{b1} + r \cdot qV_{b2} = 0$$

$$r (V_{b1} + V_{b2}) = V_{b1}$$

$$2r \cdot V_{b0} = V_{b0}$$

$$\rightarrow r = \frac{1}{2}$$

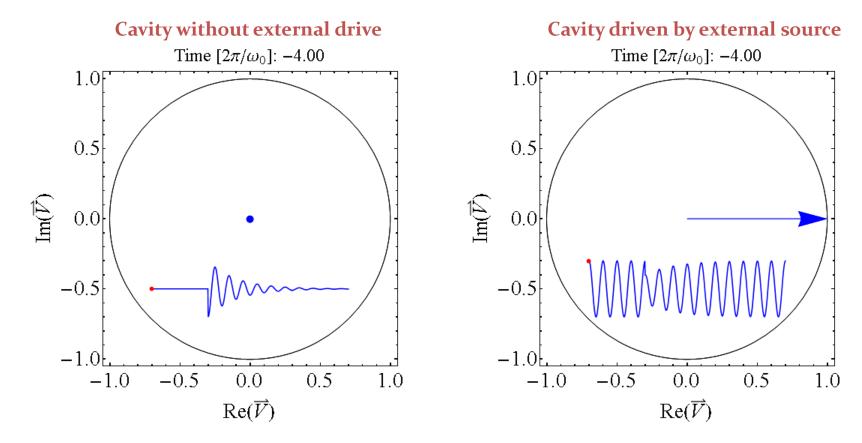
→ Fundamental theorem of beam loading:

Charge passing through a resonator sees $\frac{1}{2}$ of its induced voltage: $V_{\rm b} = \frac{1}{2} V_{\rm bo}$

Single passage through a cavity

Vector representation

- Passing charge induces voltage
- Voltage vector rotates with resonance frequency of cavity

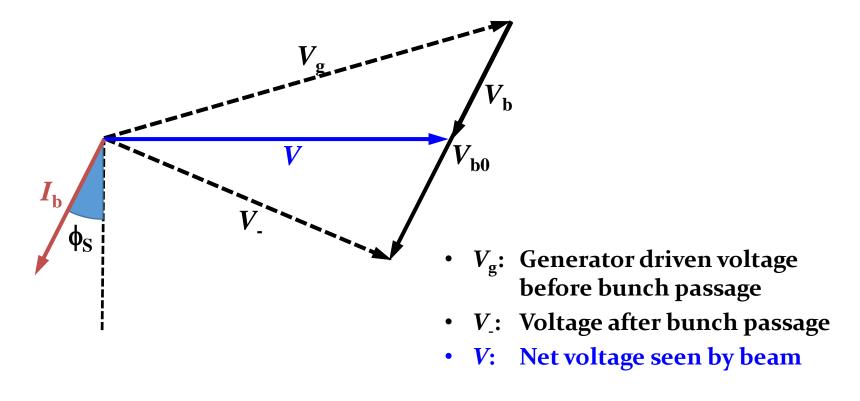


 \rightarrow Vector rotation with ω_o not relevant

→ Need cavity voltage at arrival of next charge

Single passage

• Vector diagram at the instant of the bunch passage:



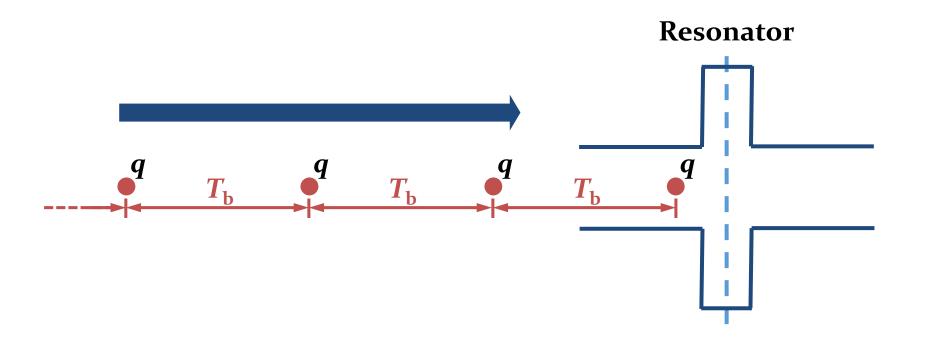
$$\rightarrow$$
 Vector sum: $\vec{V} = \vec{V}_g + \vec{V}_b = \vec{V}_g + \frac{1}{2}\vec{V}_{b0}$

- → Induced voltage changes cavity phase: detuning
- → **De-phase** generator to obtain expected net voltage

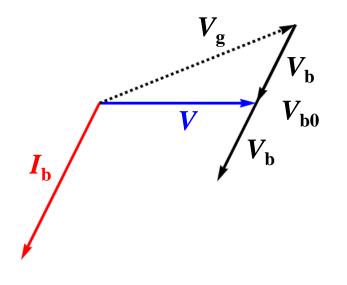
Multiple passages through a cavity

Multiple passage of bunches

- Resonator excited by chain of charges or particle bunches
- 1. Fields in resonator decay from one charge to the next
 - → Single passage case
- 2. Field from previous still present
 - → Accumulation of induced voltages



- Arrange generator phase and voltage for real net voltage
- After 1st bunch passage



$$\rightarrow \quad \vec{V} = \vec{V}_{\rm g} + \frac{1}{2}\vec{V}_{\rm b0}$$

- Arrange generator phase and voltage for real net voltage
 - After 2nd bunch passage V_{g} V_{b} V_{b0} V_{b0} V_{b0}

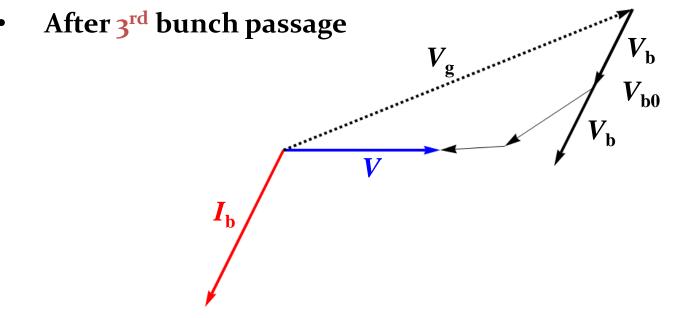
$$\rightarrow \quad \vec{V} = \vec{V}_{g} + \frac{1}{2}\vec{V}_{b0} + \vec{V}_{b0}e^{-\delta}e^{i\Psi}$$

- \rightarrow Induced voltage of 1st passage decayed: $e^{-\delta}$ with $\delta = \frac{T_b}{T}$
- \rightarrow Phase advance between two bunches:

$$^{-i\Psi}$$
 with
 $\Psi = \omega_0 T_{\rm b} - 2\pi h_{\rm b}$

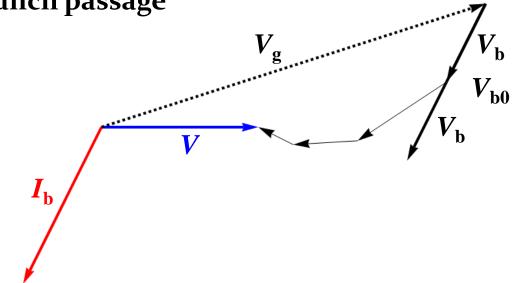
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• Arrange generator phase and voltage for real net voltage



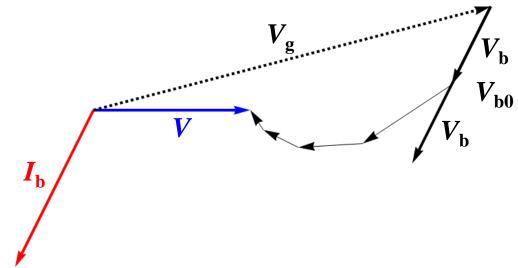
$$\rightarrow \quad \vec{V} = \vec{V}_{g} + \frac{1}{2}\vec{V}_{b0} + \vec{V}_{b0}e^{-\delta}e^{i\Psi} + \vec{V}_{b0}e^{-2\delta}e^{2i\Psi}$$

- Arrange generator phase and voltage for real net voltage
- After 4th bunch passage



$$\rightarrow \quad \vec{V} = \vec{V}_{g} + \frac{1}{2}\vec{V}_{b0} + \vec{V}_{b0} \left(e^{-\delta}e^{i\Psi} + e^{-2\delta}e^{2i\Psi} + \ldots\right)$$

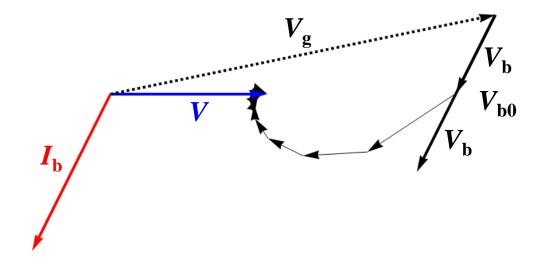
- Arrange generator phase and voltage for real net voltage
- After 5th bunch passage



$$\rightarrow \quad \vec{V} = \vec{V}_{g} + \frac{1}{2}\vec{V}_{b0} + \vec{V}_{b0} \left(e^{-\delta}e^{i\Psi} + e^{-2\delta}e^{2i\Psi} + \ldots\right)$$

Multiple passages

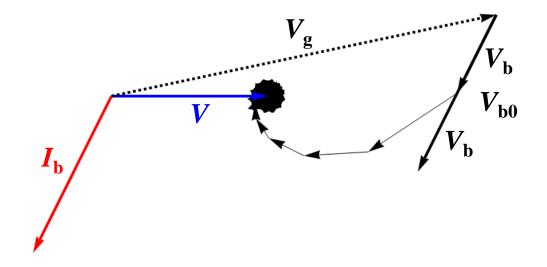
- Arrange generator phase and voltage for real net voltage
- After 10th bunch passage



$$\rightarrow \quad \vec{V} = \vec{V}_{g} + \frac{1}{2}\vec{V}_{b0} + \vec{V}_{b0} \left(e^{-\delta}e^{i\Psi} + e^{-2\delta}e^{2i\Psi} + \ldots\right)$$

Multiple passages

- Arrange generator phase and voltage for real net voltage
- After 100th bunch passage



- $\rightarrow \quad \vec{V} = \vec{V}_{g} + \frac{1}{2}\vec{V}_{b0} + \vec{V}_{b0} \left(e^{-\delta}e^{i\Psi} + e^{-2\delta}e^{2i\Psi} + \ldots\right)$
- → Infinite passages: $1 + e^{-\delta}e^{i\Psi} + e^{-2\delta}e^{2i\Psi} + \ldots = \frac{1}{1 e^{-\delta}e^{i\Psi}}$

General beam induced voltage

$$\vec{V}_{b} = \vec{V}_{b0} \left(\frac{1}{1 - e^{-\delta} e^{i\Psi}} - \frac{1}{2} \right)$$
Separate real and
imaginary part:

$$V_{b} = V_{b0} \left[F_{1}(\delta, \Psi) + iF_{2}(\delta, \Psi) \right]$$

$$F_{1}(\delta, \Psi) = \frac{1 - e^{-2\delta}}{2(1 - 2e^{-\delta} \cos \Psi + e^{-2\delta})}$$

$$F_{2}(\delta, \Psi) = \frac{e^{-\delta} \sin \Psi}{1 - 2e^{-\delta} \cos \Psi + e^{-2\delta}}$$

Change of variables

- Variables for damping, δ, and bunch-by-bunch phase advance, Ψ, not very practical
- → New variables with RF system parameters:

1. Coupling coefficient, β

$$1 + \beta = \frac{Q_0}{Q_{\rm L}}$$

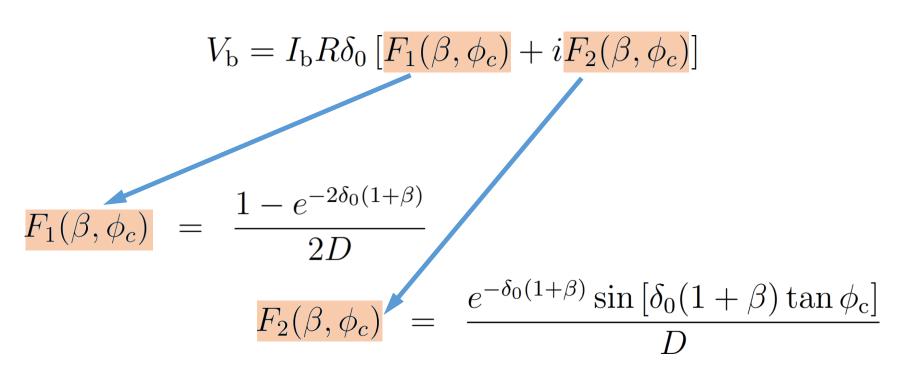
2. Cavity tuning angle, ϕ_c

$$\tan \phi_c = 2Q_{\rm L} \frac{\omega_0 - \omega}{\omega_0}$$
$$Z_{\rm L}(\omega) = \frac{R}{1 + 2iQ_{\rm L} \frac{\Delta \omega}{\omega_0}}$$
$$Z_{\rm L}(\phi_{\rm c}) = \frac{R}{1 - i \tan \phi_{\rm c}}$$

$$\delta = \delta_0 (1 + \beta)$$
$$\Psi = \delta_0 (1 + \beta) \tan \phi_c$$

Beam induced voltage in new variables

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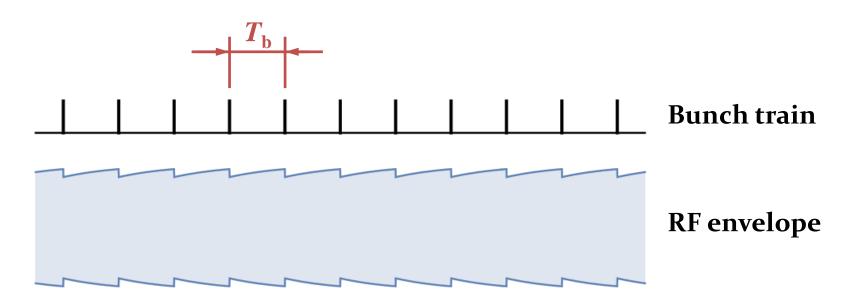
with denominator

 $D = 1 - 2e^{-\delta_0(1+\beta)} \cos \left[\delta_0(1+\beta) \tan \phi_c\right] + e^{-2\delta_0(1+\beta)}$

 \rightarrow Numerical computations required for analysis \rightarrow Let us look at a particularly relevant approximation: $\delta_0 \simeq 0$

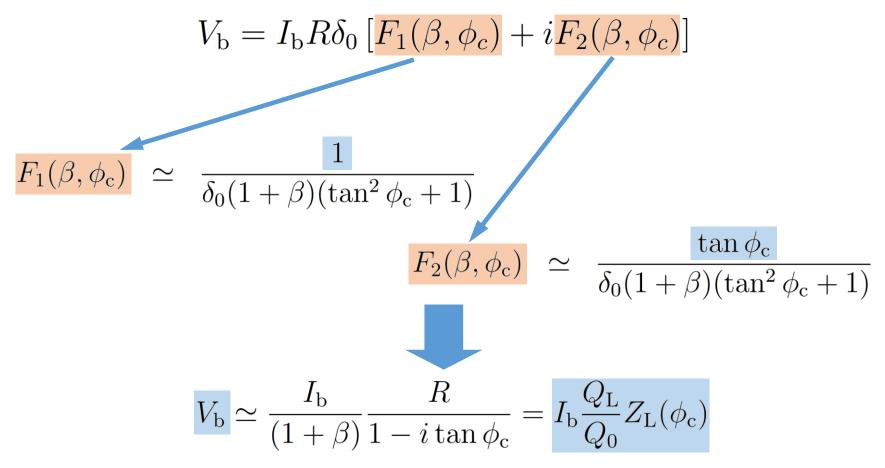
Approximation

- \rightarrow Bunch distance short compared filling time: $\delta_0 \simeq 0$
- \rightarrow Approximate terms including $\mathcal{O}(\delta_0^2)$



Approximation

- ightarrow Bunch distance short compared filling time: $\delta_0\simeq 0$
- ightarrow Approximate terms including $\mathcal{O}(\delta_0^2)$

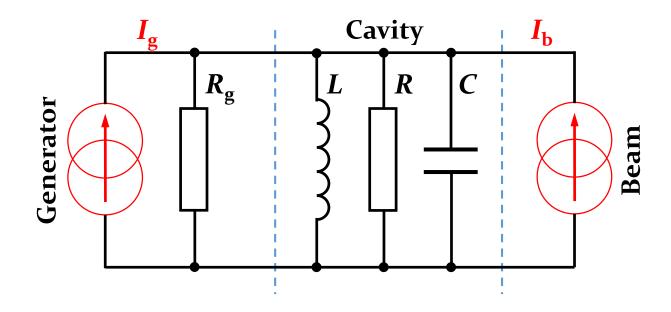


→ Ohm's law for the loaded cavity impedance: steady state case

Steady state beam loading

Equivalent circuit model

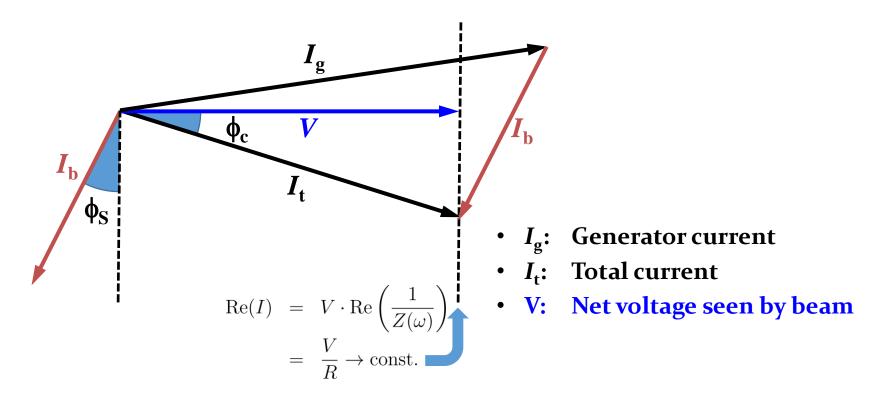
• Lumped element circuit model for steady state case



- \rightarrow Total current: $\vec{I}_{t} = \vec{I}_{g} + \vec{I}_{b}$
- \rightarrow Power required from generator: $P_{\rm g} = \frac{1}{2} R_{\rm g} I_{\rm g}^2$

Steady state

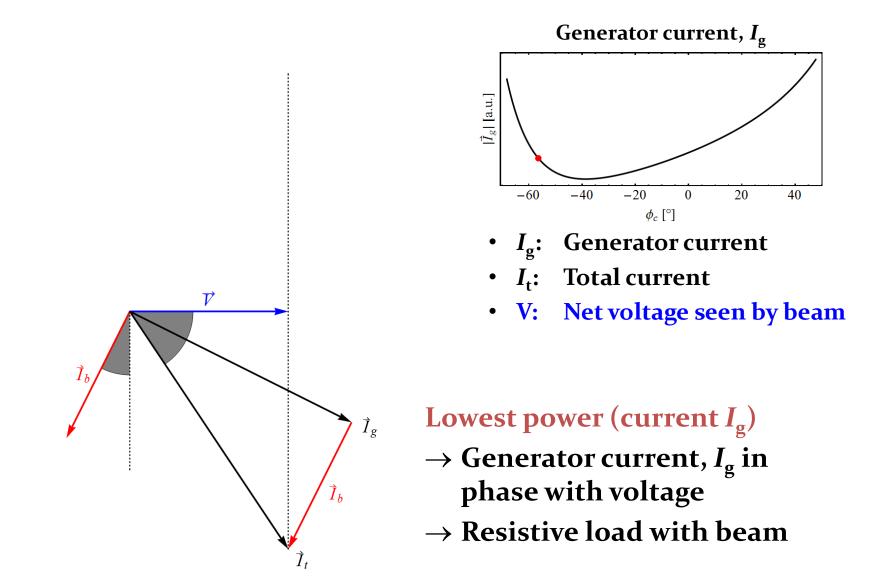
• Vector diagram for passage of continuous bunch train



→ Parameters to achieve minimum generator current?

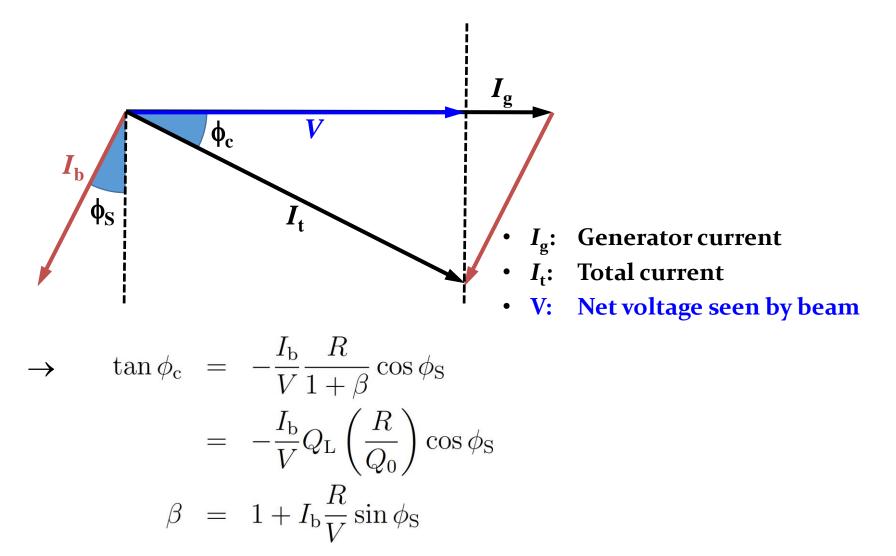
Steady state: minimum generator current

• Vector diagram for passage of continuous bunch train

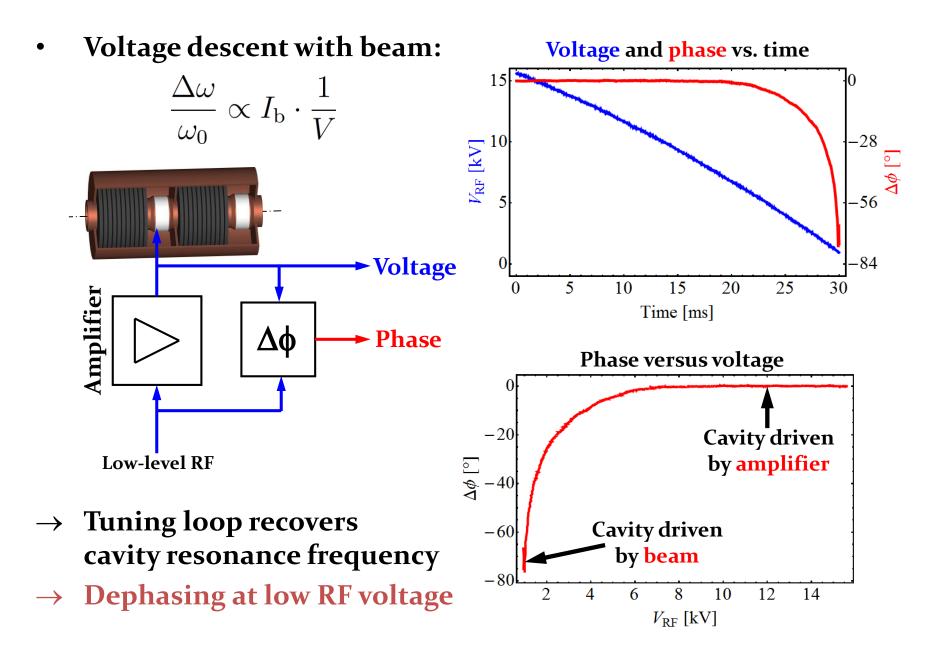


Steady state: minimum generator current

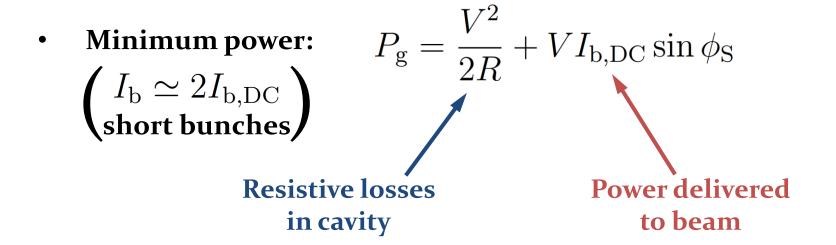
• Vector diagram for passage of continuous bunch train



Example: Cavity dephasing in PS



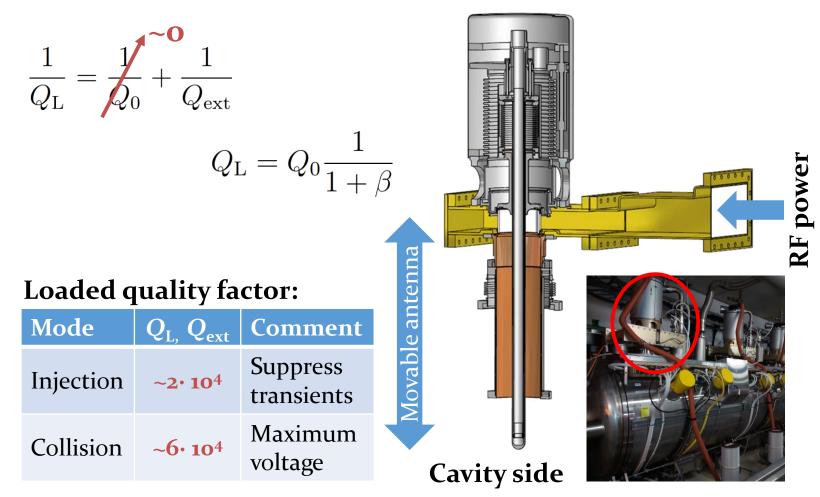
Steady state: minimum generator current



- 1. **Optimum detuning:** $\frac{\omega \omega_0}{\omega_0} = \frac{\Delta \omega}{\omega_0} = \frac{1}{2} \frac{I_{\rm b}}{V} \left(\frac{R}{Q_0}\right) \cos \phi_{\rm S}$
 - \rightarrow Cavity and beam appear as resistive load to generator
 - → Automatically adjusted by cavity tuning loop
- 2. Optimum coupling: $\beta = 1 + I_{\rm b} \frac{R}{V} \sin \phi_{\rm S}$
 - → Usually mechanically fixed by construction

Example: LHC power coupler

- Control of both cavity resonance frequency and coupling
- Optimize quality through Q_{ext} for injection and storage



Filling pattern with gaps

Why leaving a gap and not filling full ring?

- → **Electron** storage rings:
- → Hadron accelerator:

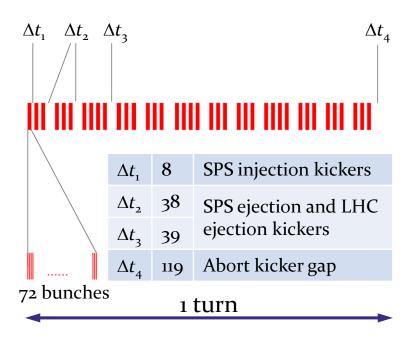
Clear ions attracted by electron beam Leave gap for kicker magnets at injection/ejection

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ESRF: 7/8 + 1 filling mode

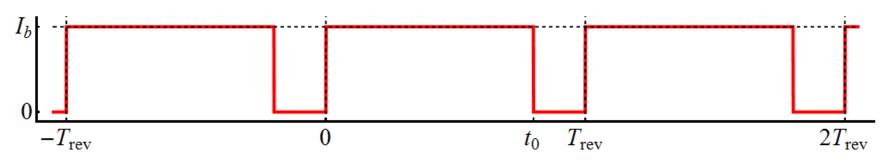
868	23 uA/b	200 mA in 7/8 train
2	1 mA/b	Marker bunches
1	2 mA	Single bunch
2×62	<2 pA/b	Gap
Revolution of	lock	
aldra manara analatina Nyaji kalenge program	esta fara esta a serie se presenta de series en el fara esta esta esta esta esta esta esta est	200 mA in 868 Bunches
		Tir

LHC: original nominal

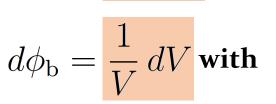


Beam loading with gaps

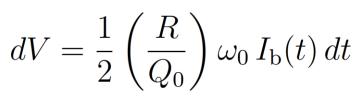
• Limitations: $\delta_0 \simeq 0$, no acceleration, lossless cavity

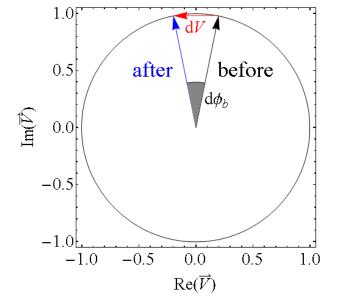


- Phase change due to cavity detuning:
- Phase change due to induced voltage:



 $d\phi_{\rm a} = \Delta \omega \, dt$





$$d\phi = \left[\Delta\omega - \frac{1}{2}\left(\frac{R}{Q_0}\right)\frac{\omega_0}{V}I_{\rm b}(t)\right]\,dt$$

Beam loading with gaps

→ **Periodicity condition** $\int_{1 \text{turn}} d\phi = 0$ **to get average detuning**

$$\Delta\omega_0 = \frac{1}{2} \left(\frac{R}{Q_0}\right) \frac{\omega_0}{V} \frac{1}{T_{\rm rev}} \int_0^{T_{\rm rev}} I_{\rm b}(t) \, dt = \frac{1}{2} \left(\frac{R}{Q_0}\right) \frac{\omega_0}{V} \bar{I}_{\rm b}$$

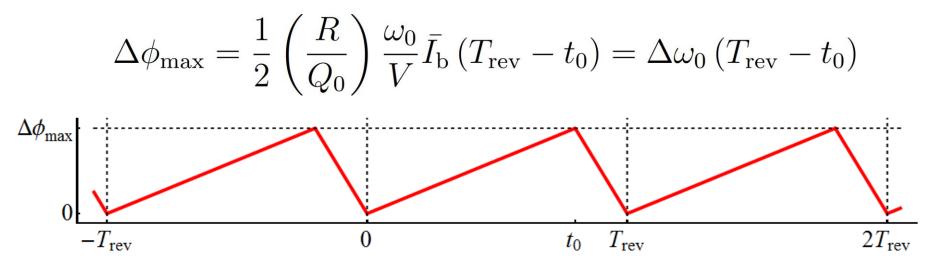
 \rightarrow and phase along the circumference

$$\phi(t) = \int_0^t d\phi = \frac{1}{2} \left(\frac{R}{Q_0}\right) \frac{\omega_0}{V} \int_0^t \left[\bar{I}_{\rm b} - I_{\rm b}(t)\right] dt$$

 \rightarrow Phase changes linearly for $I_{\rm b}(t)$ = const. during beam region

Maximum phase excursion

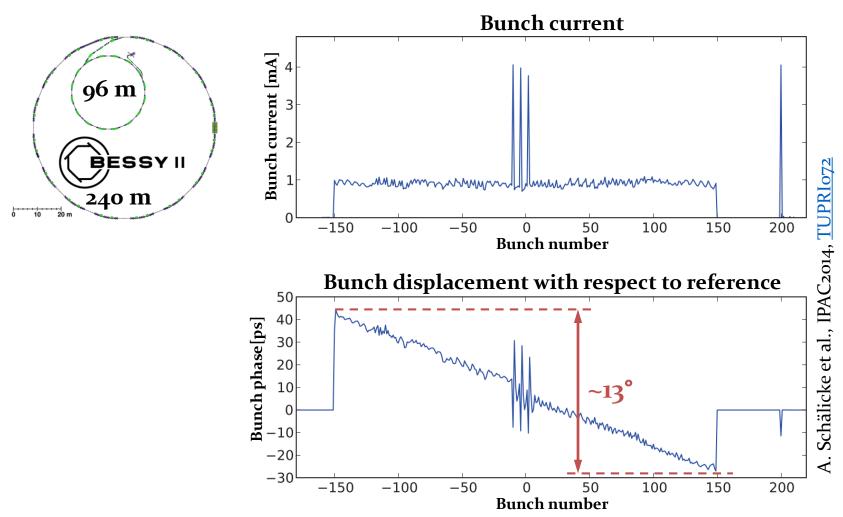
• Maximum phase excursion



- → Displaces timing of synchrotron radiation pulses
- → Longitudinally moves collision point in collider
- → **Compromise** between RF power and collision point

Example: Electron storage ring

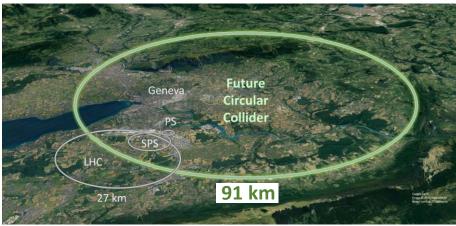
• Transient beam loading in electron storage ring BESSY II



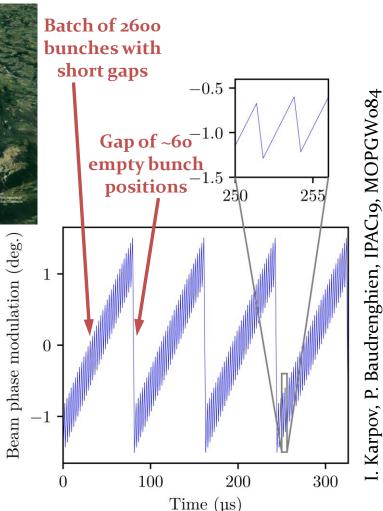
→ Synchrotron radiation light pulses slightly shifted in time

Example: FCC-hh (hadron-hadron)

• Proposed future circular collider

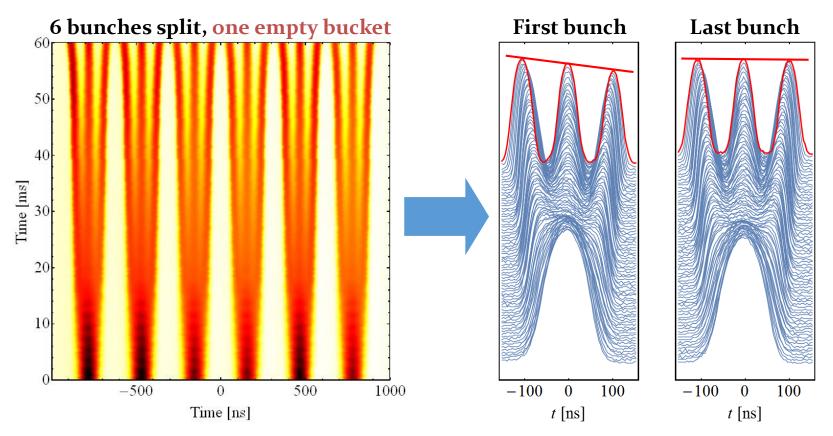


- Machine protection requires
 - → Four batches per turn
 - \rightarrow Gaps of ~1.5 µs
- → Full-detuning causes a bunch phase modulation of ~2°
- → Position of collision point modulated



Transient beam loading between RF systems

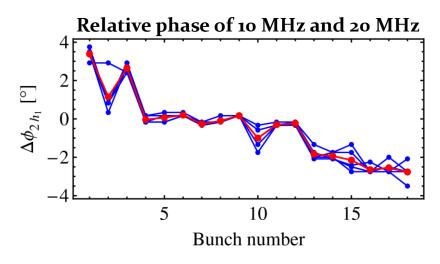
 Triple splitting of LHC-type beams in CERN PS requires three RF systems (h = 7, 14 and 21) in phase at degree level

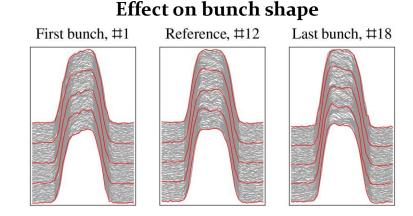


- \rightarrow Transient beam loading: relative phases different for 1st bunch
- → Bunch-by-bunch intensity variations in LHC

Transient beam loading between RF systems

 \rightarrow Fast phase measurement to directly observe relative changes



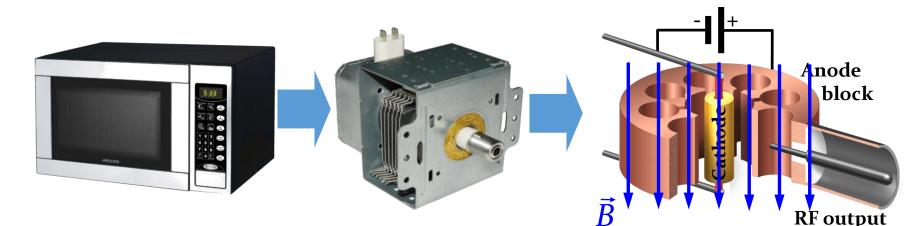


- Cavity detuning not an option
 - → Would even enhance phase modulation along batch
- Feedback systems
 - → Counteract beam loading with additional RF power
 - → Stabilize phase

Beam loading in microwave oven?

Beam loading in microwave oven?

• Microwave ovens use magnetrons as RF power source



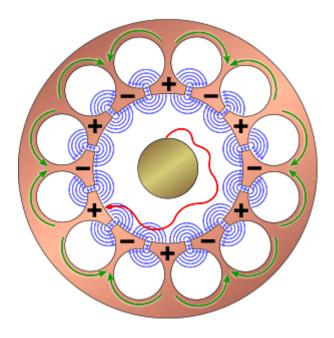
- Anode block consists of ring of cavity resonators
- Electrons from the cathode accelerated toward anode (cavities)
- Perpendicular magnetic field causes cyclotron motion

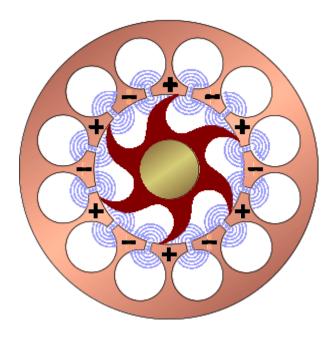


Ch. Wolff, http://www.radartutorial.eu/o8.transmitters/Magnetron.en.html

Beam loading in microwave oven?

• Magnetron as RF power source





- → Electron flow from cathode to anode self-bunched under influence of oscillating fields in anode resonators
- → Bunched electrons excite RF fields → beam loading!
- \rightarrow Food gets heated

Summary

- RF cavity parameters
 → System of cavity, coupling and amplifier
- Single and multi-passage of bunches through a cavity
 → Fundamental theorem of beam loading
 → Multiple passages limiting case of steady state
- Steady state beam loading
 → Minimize RF power by detuning and coupling
- Partial filling

 \rightarrow Modulation of bunch phase and RF voltage

• Magnetron principle

 \rightarrow Heating food with beam loading

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Thank you very much for your attention!

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Approximations

$\rightarrow~\mathbf{2^{nd}}~\mathbf{order}~\mathbf{Taylor}~\mathbf{expansion}~\mathbf{for}~\delta_0\simeq 0$

$$e^{-\delta_0(1+\beta)} \simeq 1 - \delta_0(1+\beta) + \frac{1}{2}\delta_0^2(1+\beta)^2$$
$$e^{\delta_0(1+\beta)} \simeq 1 + \delta_0(1+\beta) + \frac{1}{2}\delta_0^2(1+\beta)^2$$
$$\cos\left[\delta_0(1+\beta)\tan\phi_{\rm c}\right] \simeq 1 - \frac{1}{2}\delta_0^2(1+\beta)^2\tan^2\phi_{\rm c}$$
$$\sin\left[\delta_0(1+\beta)\tan\phi_{\rm c}\right] \simeq \delta_0(1+\beta)\tan\phi_{\rm c}$$

Approximations: *F*₁

\rightarrow Simplification of real part $F_1(\beta, \phi_c)$ for $\delta_0 \simeq 0$

$$\begin{split} F_1 &= \frac{1 - e^{-2\delta_0(1+\beta)}}{2\{1 - 2e^{-\delta_0(1+\beta)}\cos\left[\delta_0(1+\beta)\tan\phi_c\right] + e^{-2\delta_0(1+\beta)}\}} \\ &= \frac{e^{\delta_0(1+\beta)} - e^{-\delta_0(1+\beta)}}{2\{e^{\delta_0(1+\beta)} - 2\cos\left[\delta_0(1+\beta)\tan\phi_c\right] + e^{-\delta_0(1+\beta)}\}} \\ &\simeq \frac{1 + \delta_0(1+\beta) + \frac{1}{2}\delta_0^2(1+\beta)^2 - 1 + \delta_0(1+\beta) - \frac{1}{2}\delta_0^2(1+\beta)^2}{2\{1 + \delta_0(1+\beta) + \frac{1}{2}\delta_0^2(1+\beta)^2 - 2\left[1 - \frac{1}{2}\delta_0^2(1+\beta)^2 \tan^2\phi_c\right] + 1 - \delta_0(1+\beta) + \frac{1}{2}\delta_0^2(1+\beta)^2\}} \\ &= \frac{2\delta_0(1+\beta)}{2\{2 + \delta_0^2(1+\beta)^2 - 2 + \delta_0^2(1+\beta)^2 \tan^2\phi_c\}} \\ &= \frac{\delta_0(1+\beta)}{\delta_0^2(1+\beta)^2 + \delta_0^2(1+\beta)^2 \tan^2\phi_c} \\ &= \frac{1}{\delta_0(1+\beta)} \end{split}$$

Approximations: F_{2}

→ Simplification of real part $F_2(\beta, \phi_c)$ for $\delta_0 \simeq 0$

$$F_{2} = \frac{e^{-\delta_{0}(1+\beta)} \sin [\delta_{0}(1+\beta) \tan \phi_{c}]}{1-2e^{-\delta_{0}(1+\beta)} \cos [\delta_{0}(1+\beta) \tan \phi_{c}] + e^{-2\delta_{0}(1+\beta)}}$$

$$= \frac{\sin [\delta_{0}(1+\beta) \tan \phi_{c}]}{e^{\delta_{0}(1+\beta)} - 2\cos [\delta_{0}(1+\beta) \tan \phi_{c}] + e^{-\delta_{0}(1+\beta)}}$$

$$\approx \frac{\delta_{0}(1+\beta) \tan \phi_{c}}{1+\delta_{0}(1+\beta) + \frac{1}{2}\delta_{0}^{2}(1+\beta)^{2} - 2[1-\frac{1}{2}\delta_{0}^{2}(1+\beta)^{2} \tan^{2}\phi_{c}] + 1-\delta_{0}(1+\beta) + \frac{1}{2}\delta_{0}^{2}(1+\beta)^{2}}$$

$$= \frac{\delta_{0}(1+\beta) \tan \phi_{c}}{2+\delta_{0}^{2}(1+\beta)^{2} - 2+\delta_{0}^{2}(1+\beta)^{2} \tan^{2}\phi_{c}}$$

$$= \frac{\delta_{0}(1+\beta) \tan \phi_{c}}{\delta_{0}^{2}(1+\beta)^{2} + \delta_{0}^{2}(1+\beta)^{2} \tan^{2}\phi_{c}}$$

$$= \frac{\tan \phi_{c}}{\delta_{0}(1+\beta)(\tan^{2}\phi_{c}+1)}$$

Frequency and wavelength ranges

