# Landau damping



Dr. Xavier Buffat

Beams Department – Accelerator and Beam Physics Collective Effects and Impedances

CERN, Switzerland, Geneva

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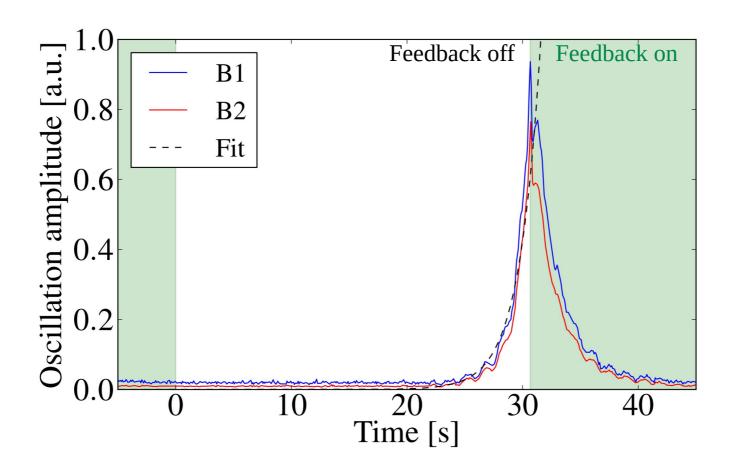
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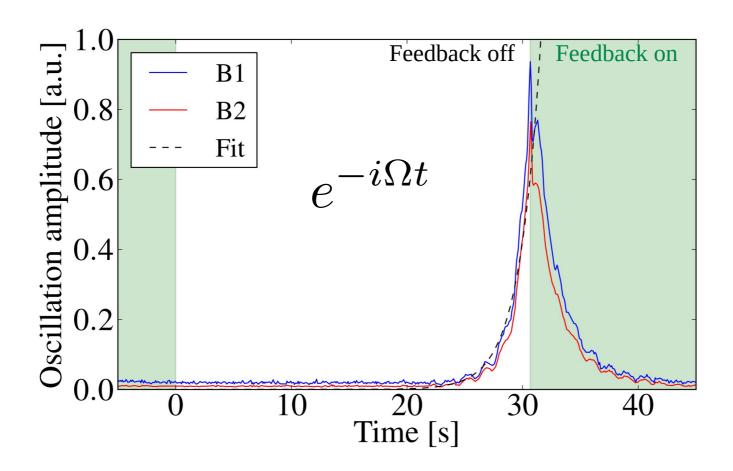
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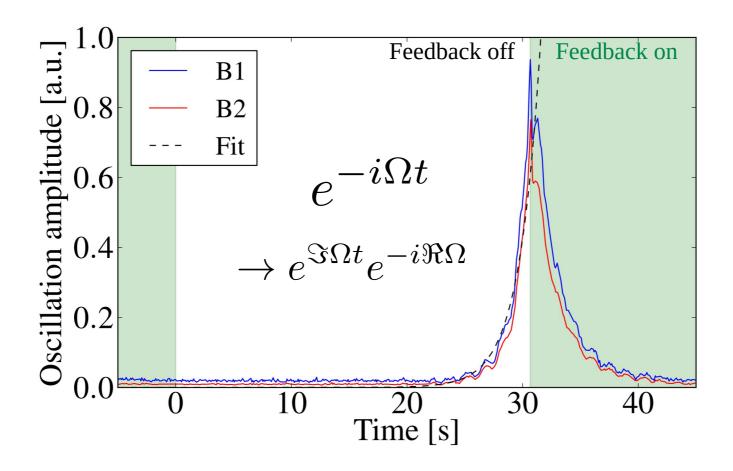
Beams tend to self-destruct via self-amplified oscillations



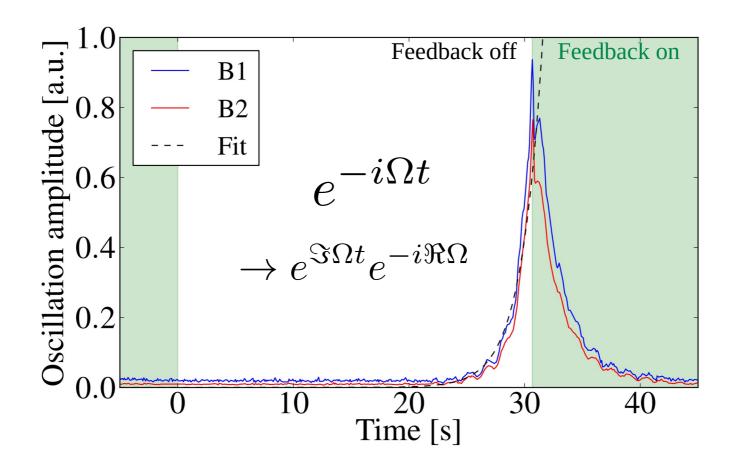
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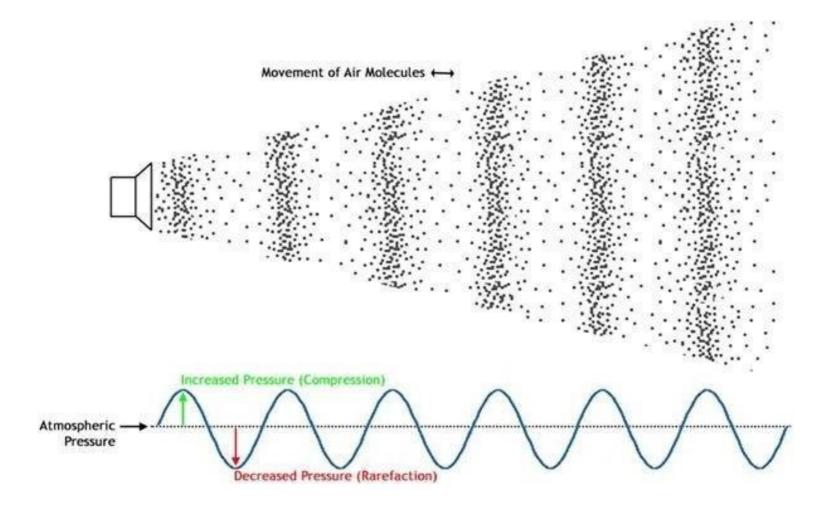


→ Landau damping is (almost) **always needed** to obtain good quality beams

#### Content

- Part I (concept)
  - Wave particle interaction
  - Decoherence
  - Landau damping using Van Kampen approach
  - Stability diagram and beam transfer function
- Part II (applications)
  - Longitudinal and transverse Landau damping in unbunched and bunched beams
  - Non-linear collective forces
  - Advanced Landau damping techniques

#### Sound Propagation



# Interaction of particle with the collective force



# Interaction of particle with the collective force



• Surfers catch the wave when they have a similar velocity

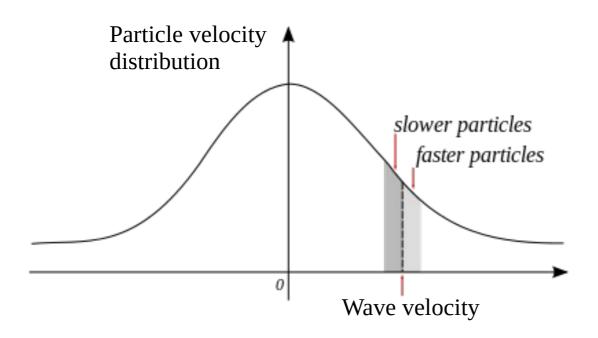
### Interaction of particle with the collective force

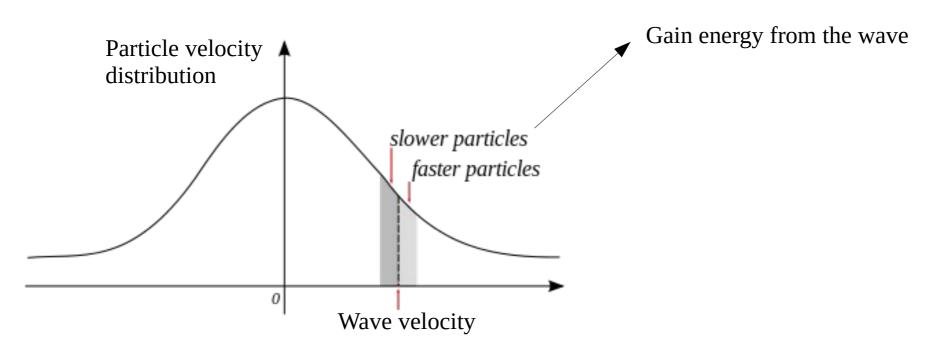


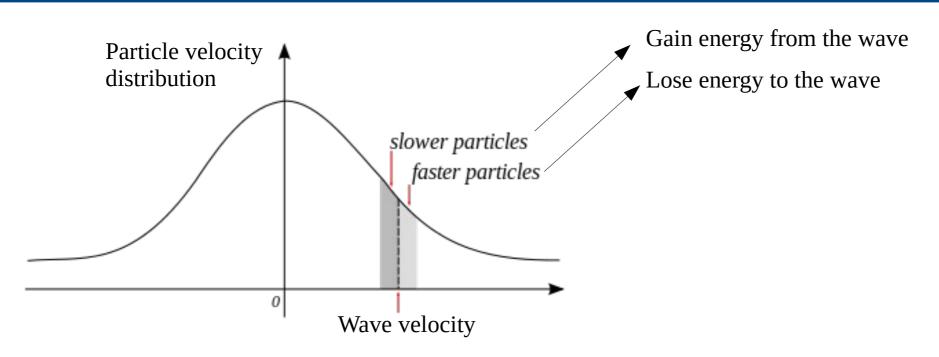
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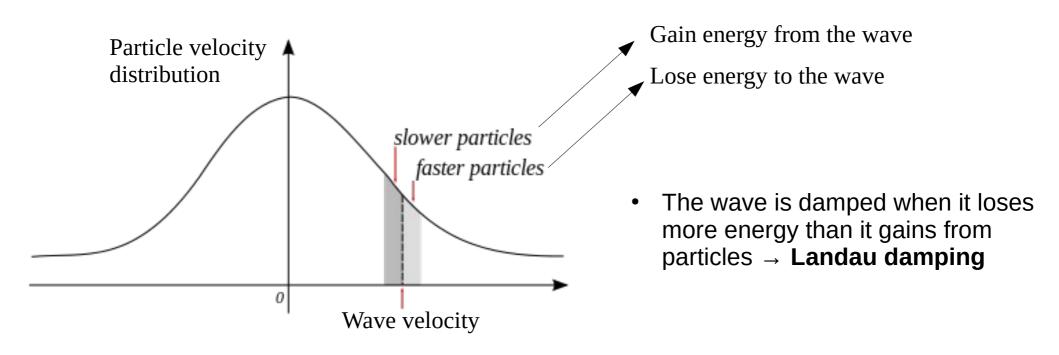


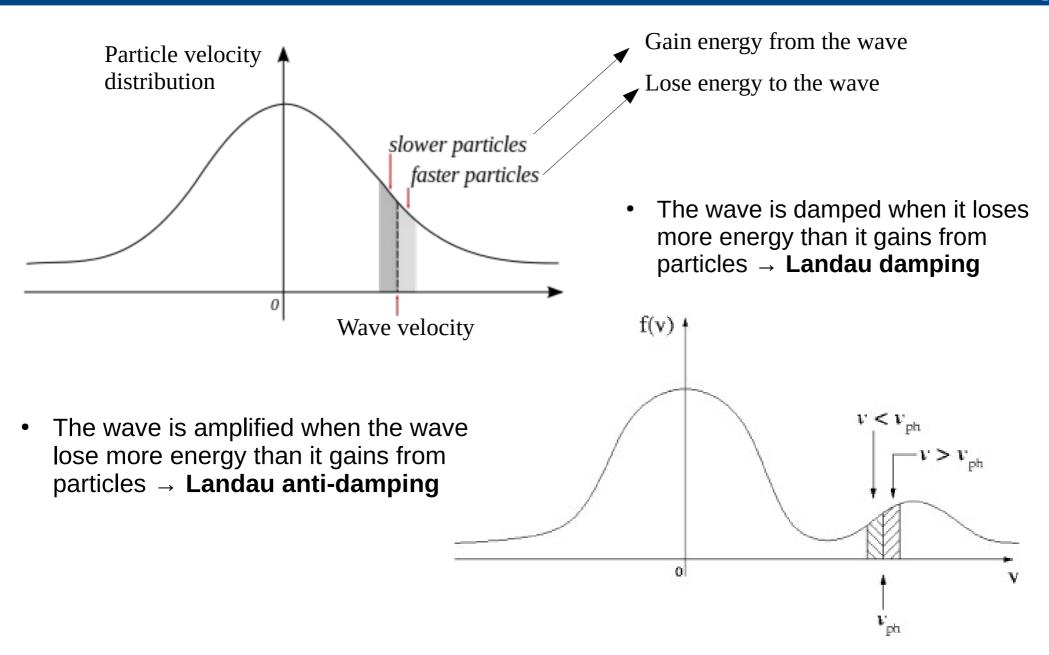
Particles can exchange energy with a wave when they have a similar velocity

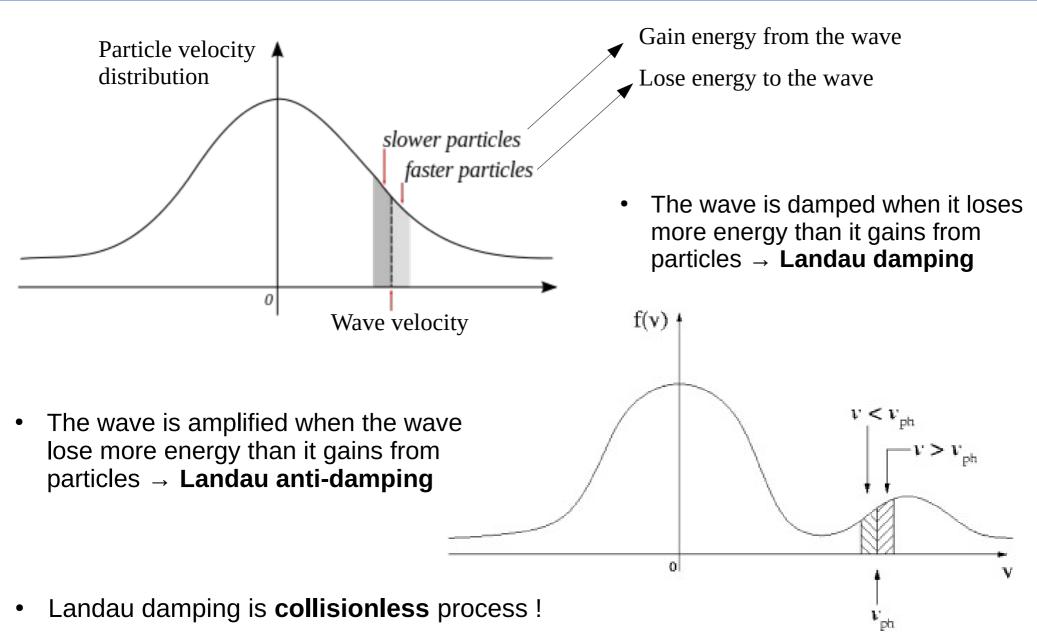












The interaction between the particles and the wave occures only via the collective force (e.g. electromagnetic fields)

A little subtlety for accelerators

[Distribution]

Landau damping prevents instabilities to happens

#### [Distribution]

# Damping of collective motion A little subtlety for accelerators

### Landau damping prevents instabilities to happens

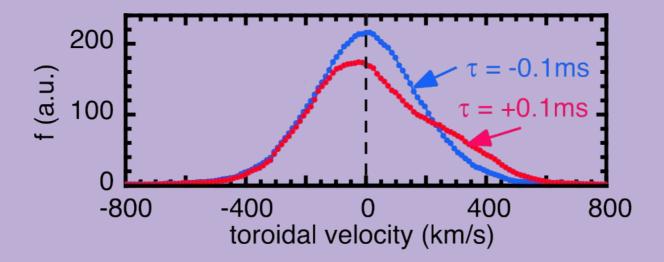
'If a small perturbation occurs, it is immediately damped preventing its self-amplification'  $\rightarrow$  No energy exchange

# Damping of collective motion A little subtlety for accelerators

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'If a small perturbation occurs, it is immediately damped preventing its self-amplification'  $\rightarrow$  No energy exchange

 When an external force drives the collective motion, the energy input is absorbed by the particles via Landau damping

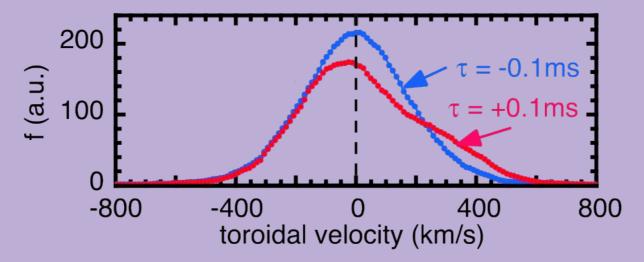


# Damping of collective motion A little subtlety for accelerators

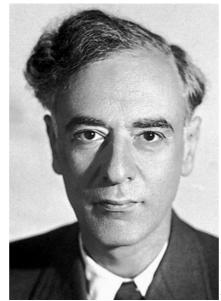
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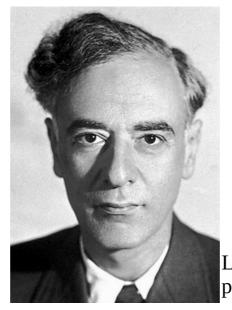
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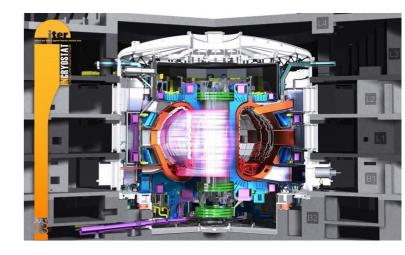


- In accelerators we refer to this effect as decoherence or filamentation
  - → The main difference with Landau damping is the corresponding emittance growth

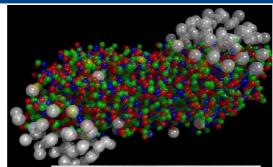


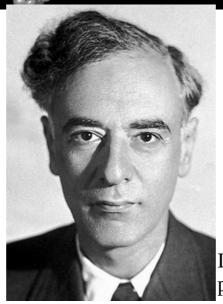
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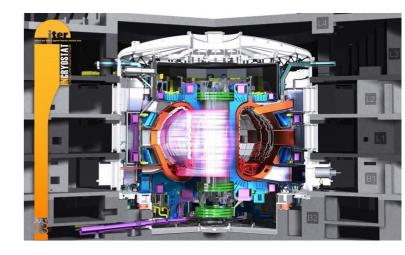




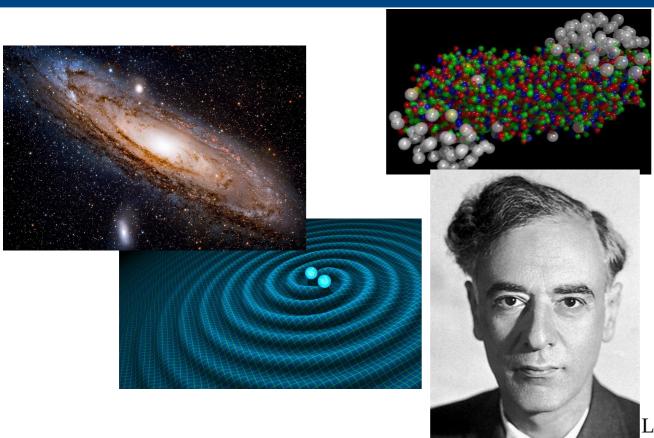
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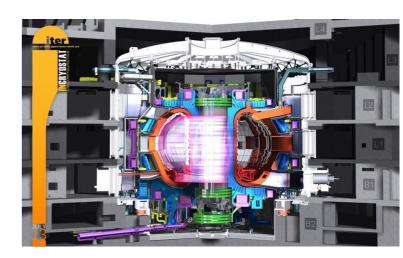




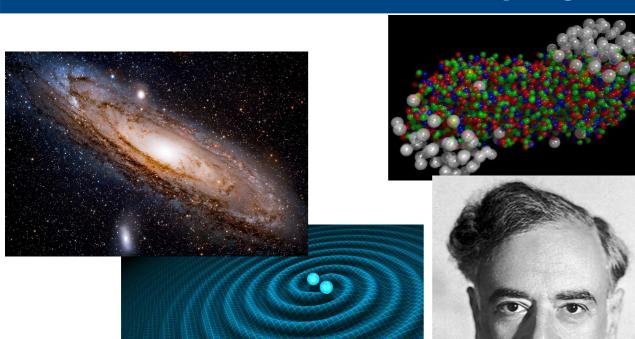


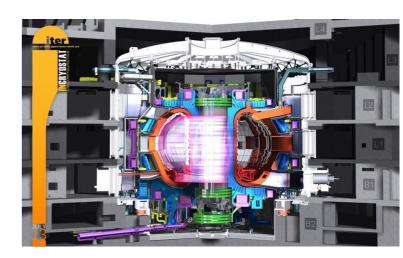
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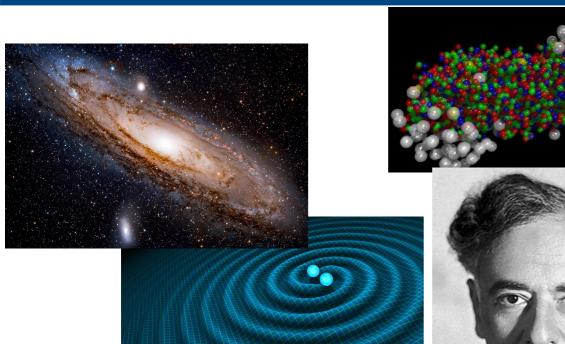
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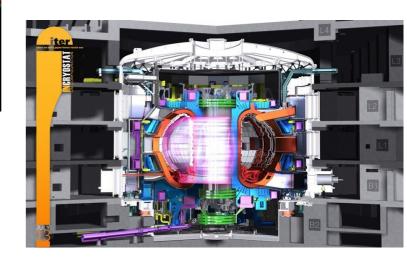




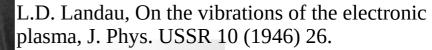
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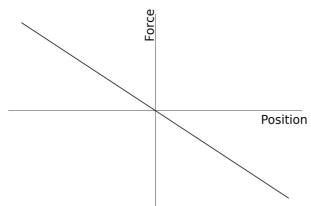






• The velocity spread is usually small in particle beams → an analogous effect occurs thanks to the **tune spread** 





#### Linear force

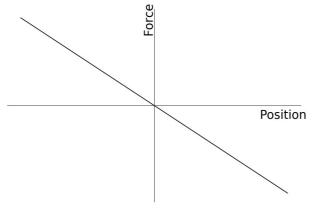
→ Fixed oscillation frequency

$$\omega = \omega_0 = 2\pi Q_0$$

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Linear force

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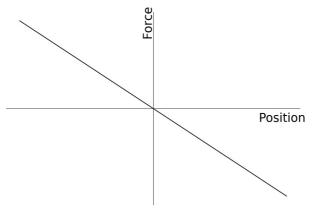
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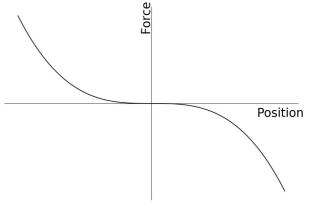




Linear force

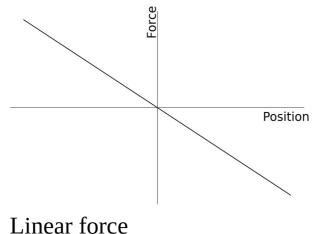
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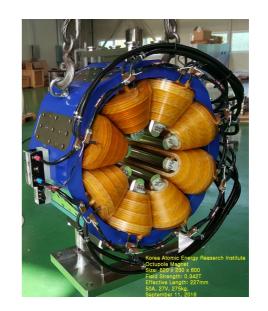
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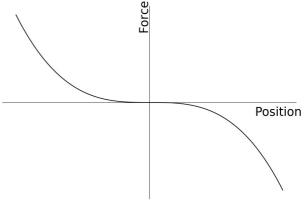




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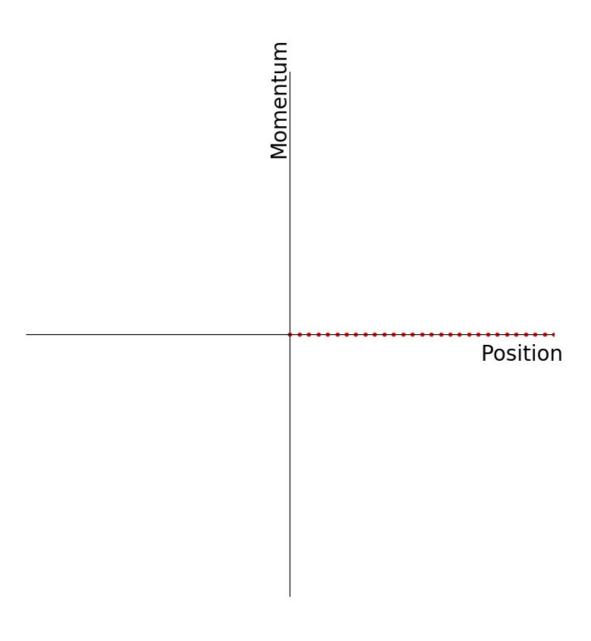


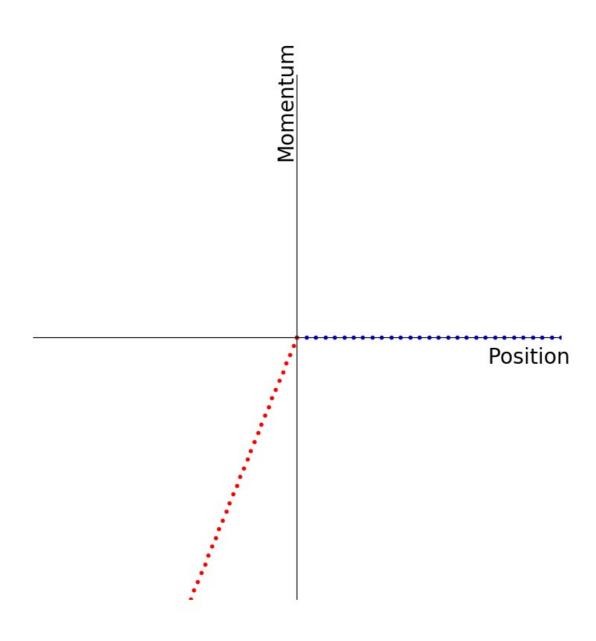


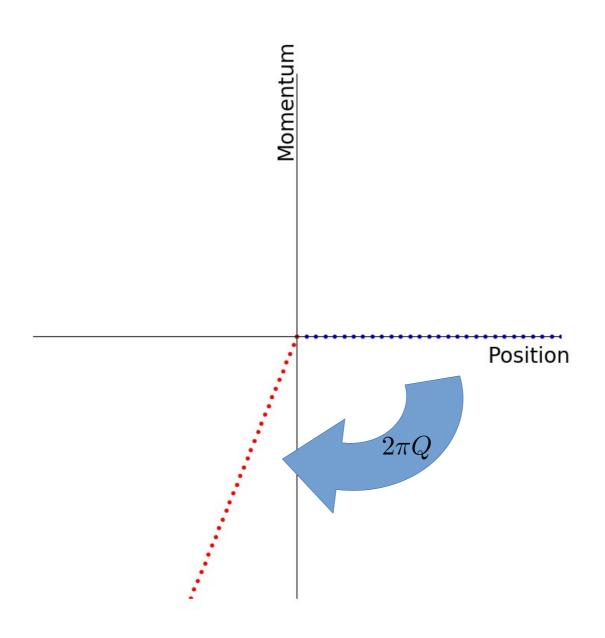
Non linear force

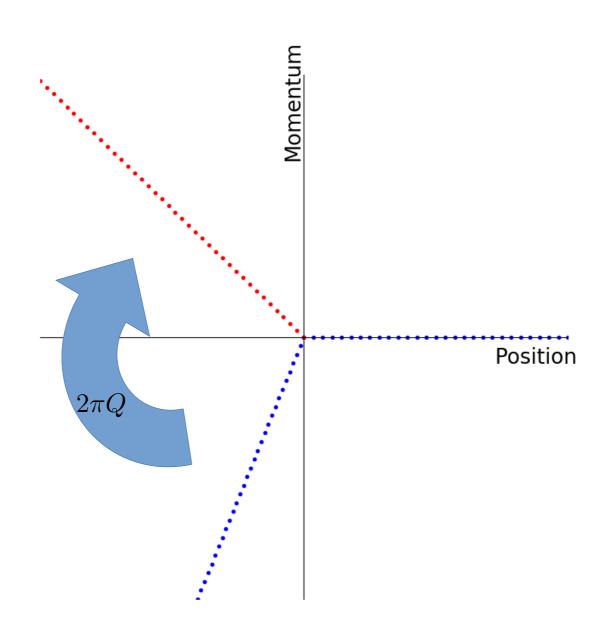
 $\rightarrow$  Amplitude dependent frequency / **detuning** 

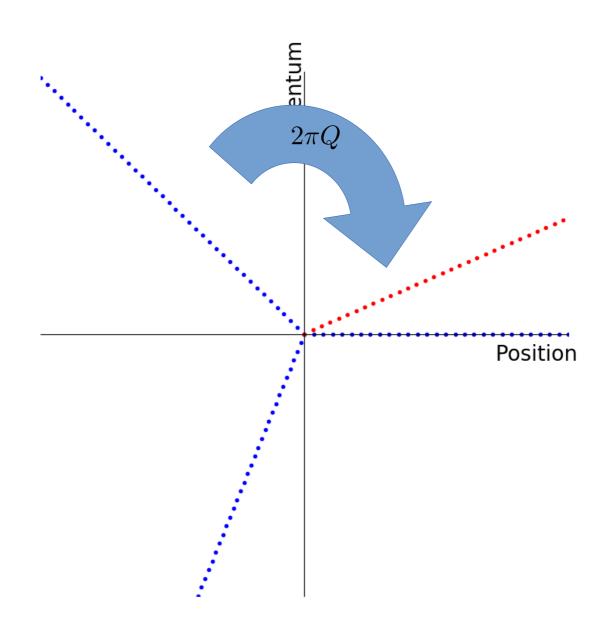
$$\omega(J) = 2\pi(Q_0 + aJ)$$

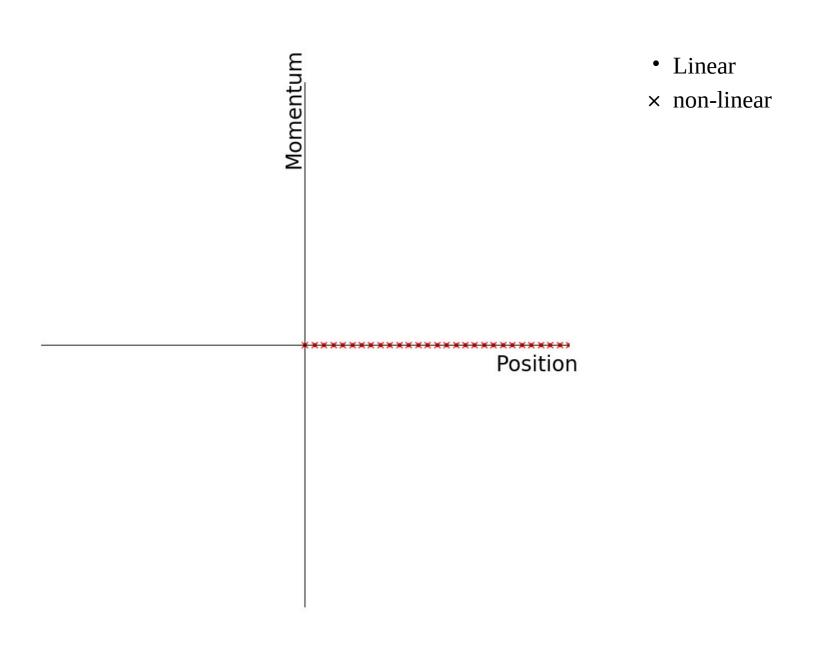


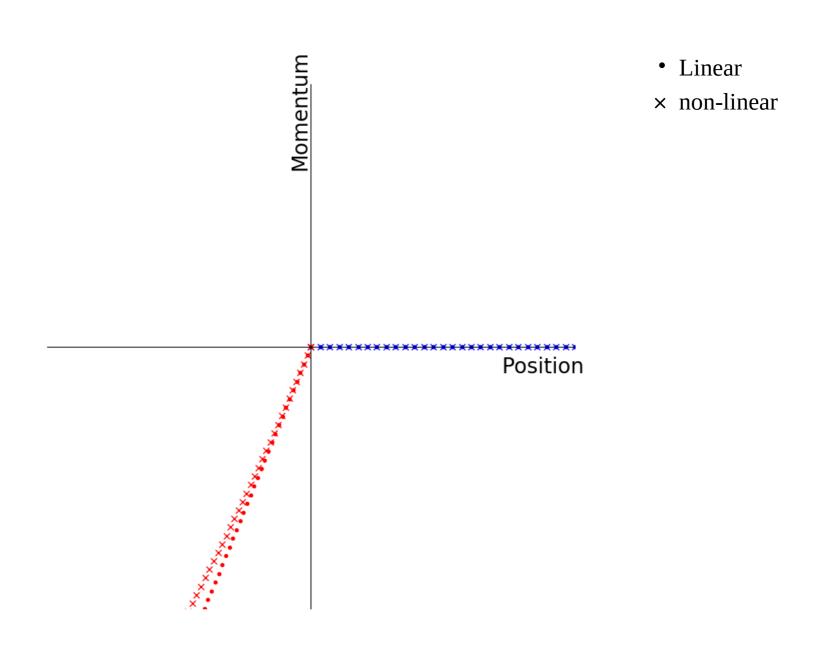


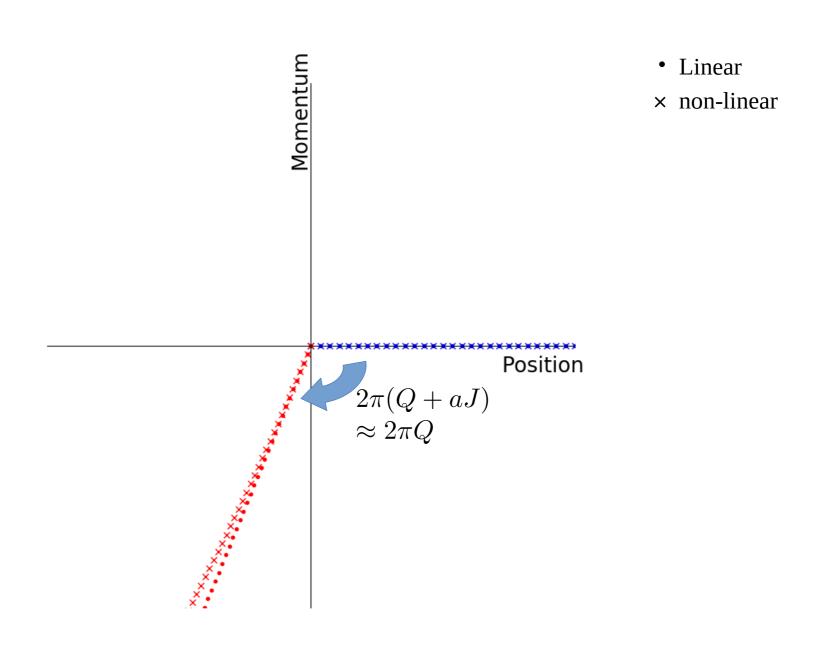


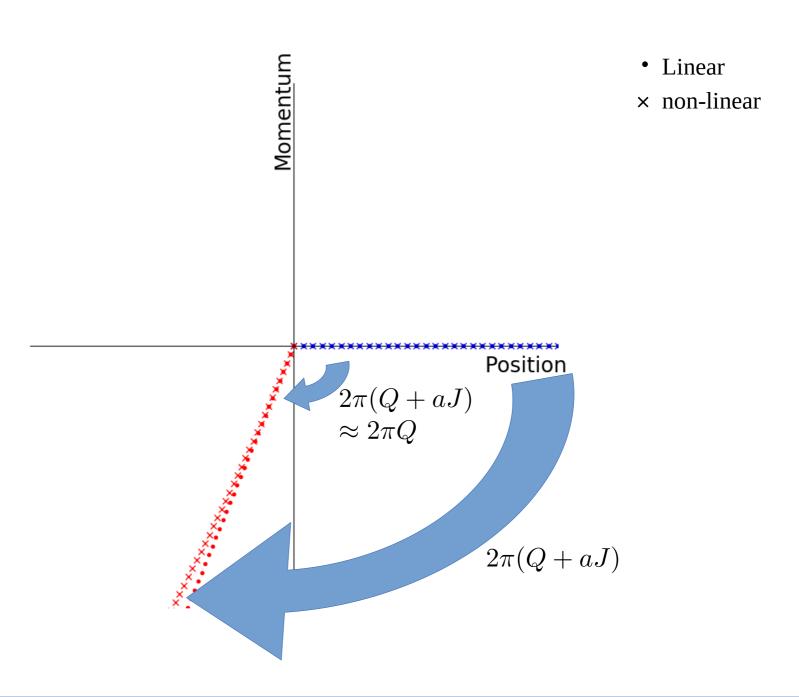


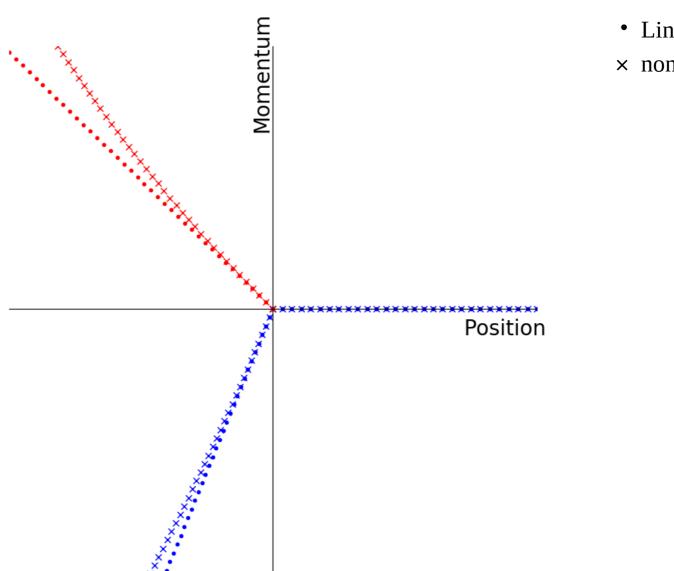




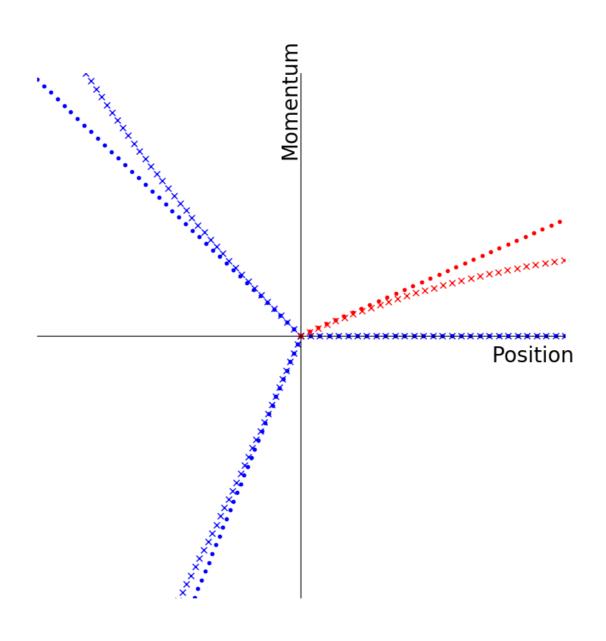




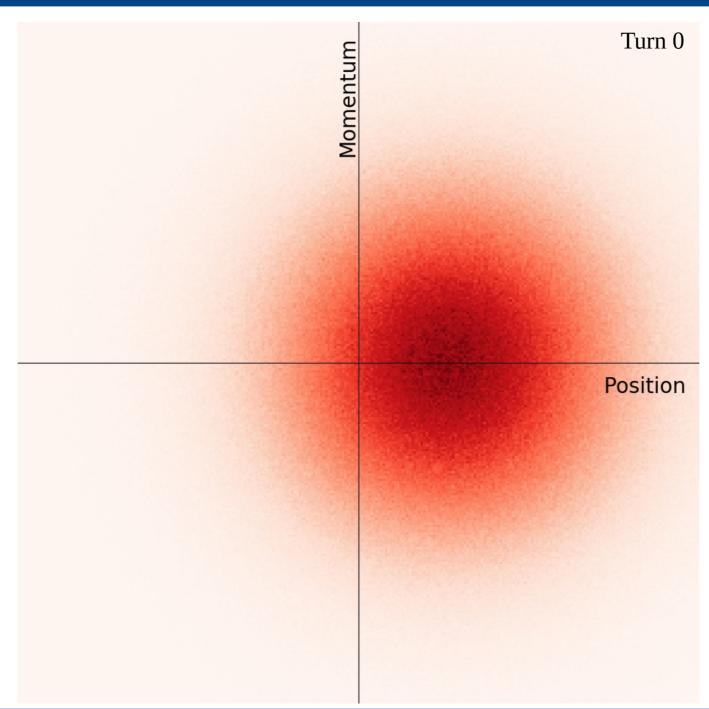


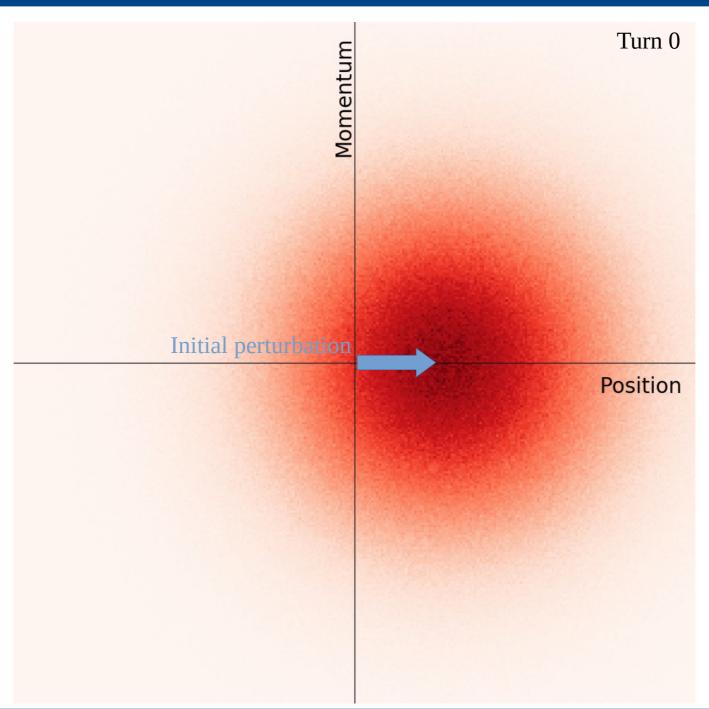


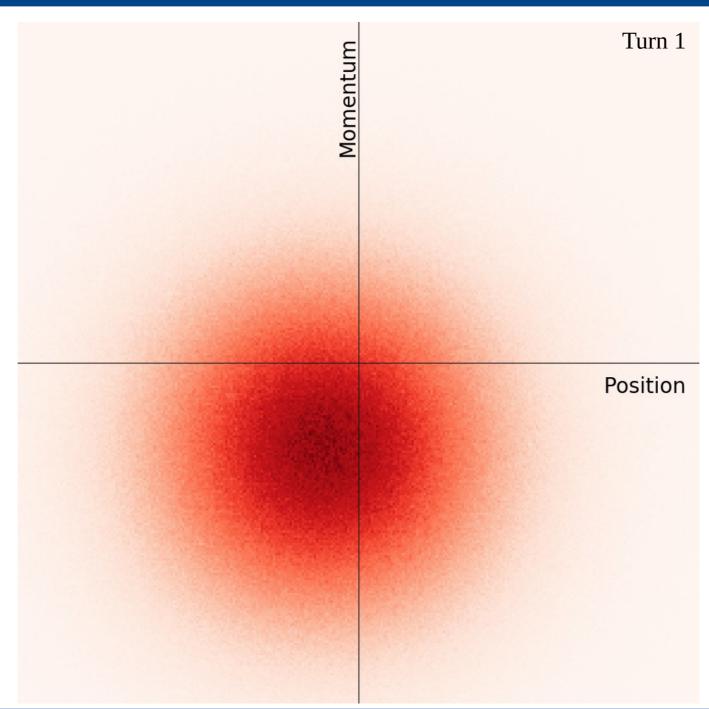
- Linear
- $\times$  non-linear

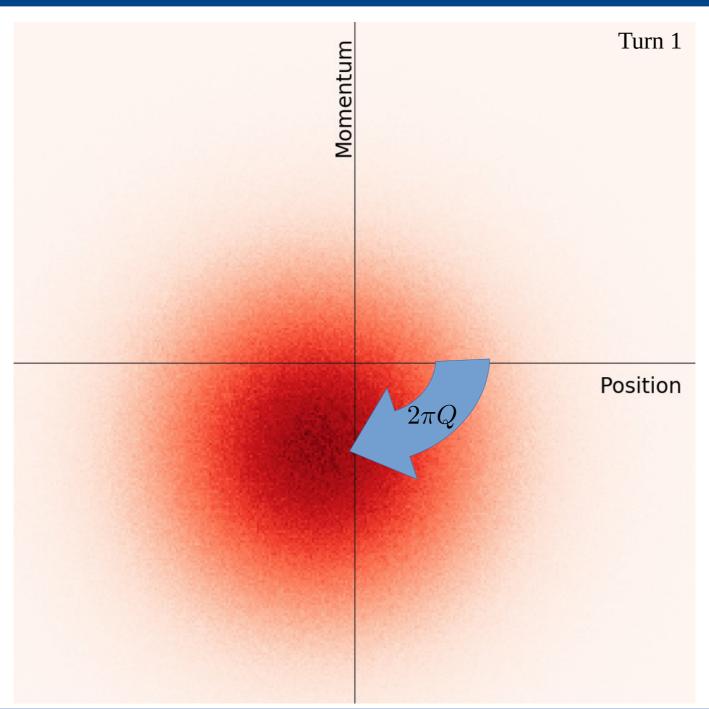


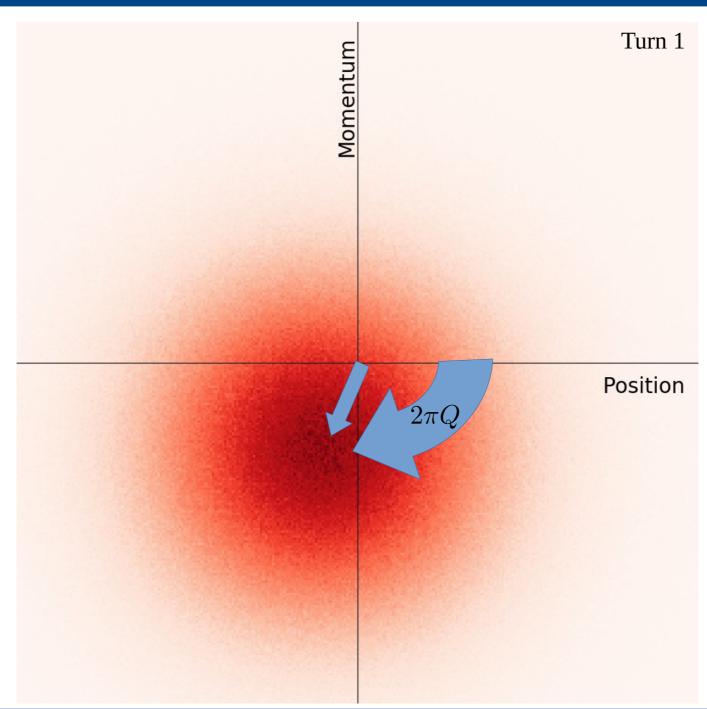
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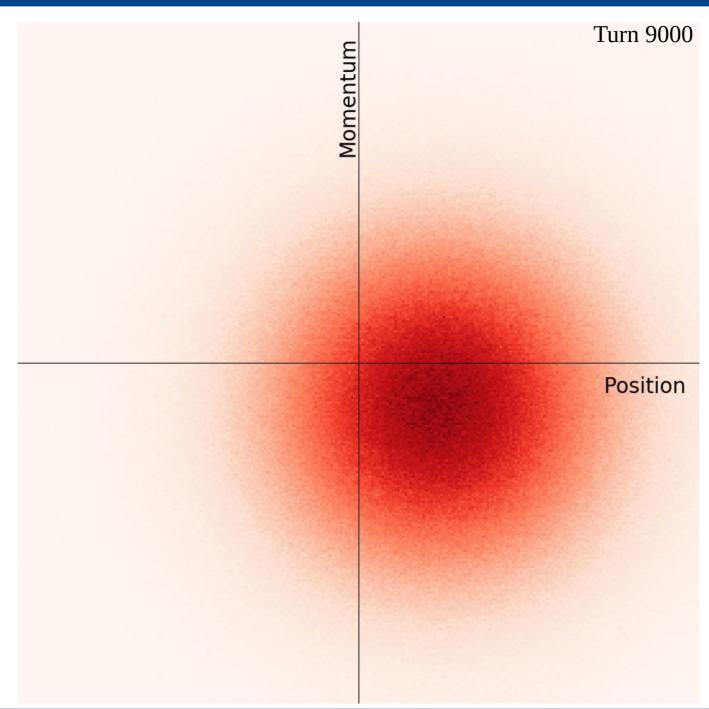


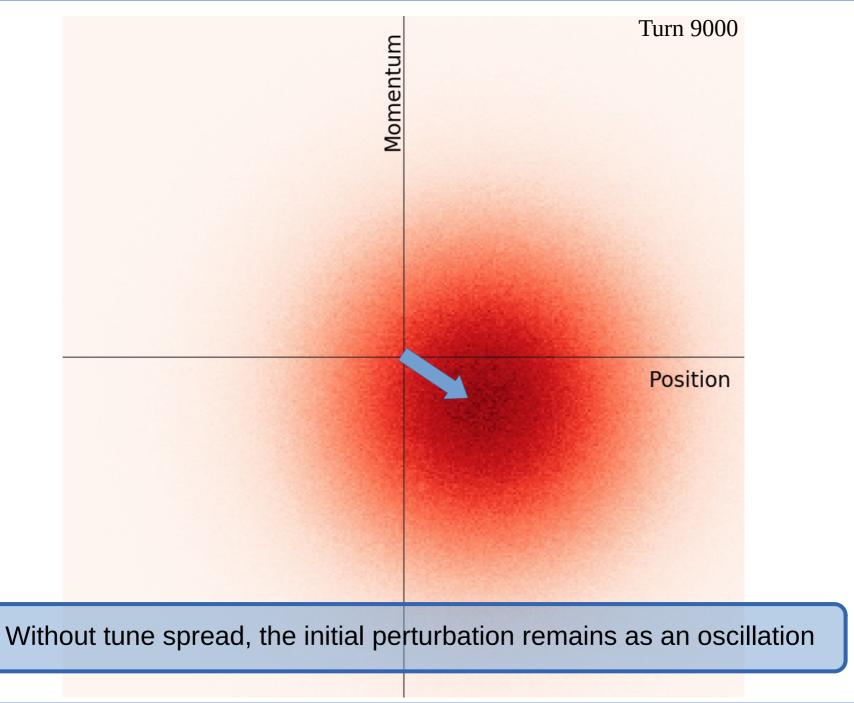


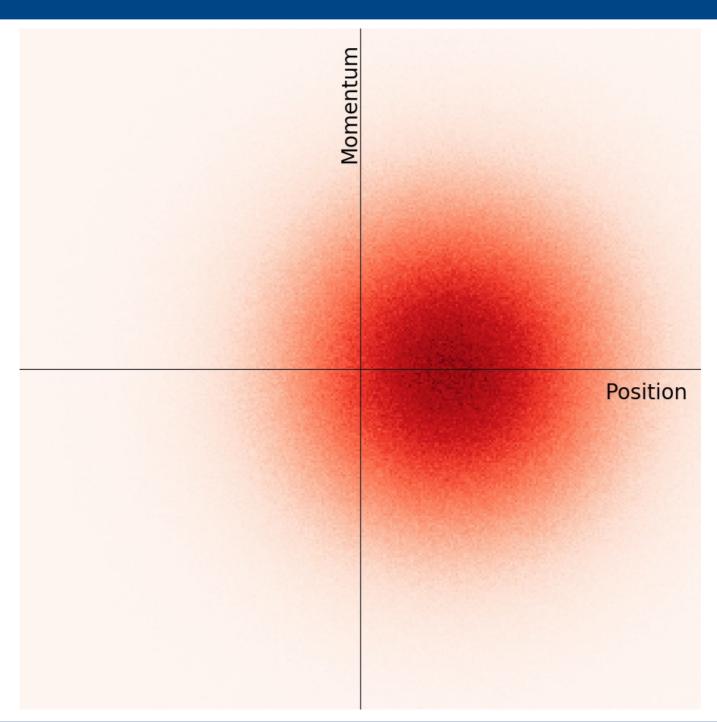


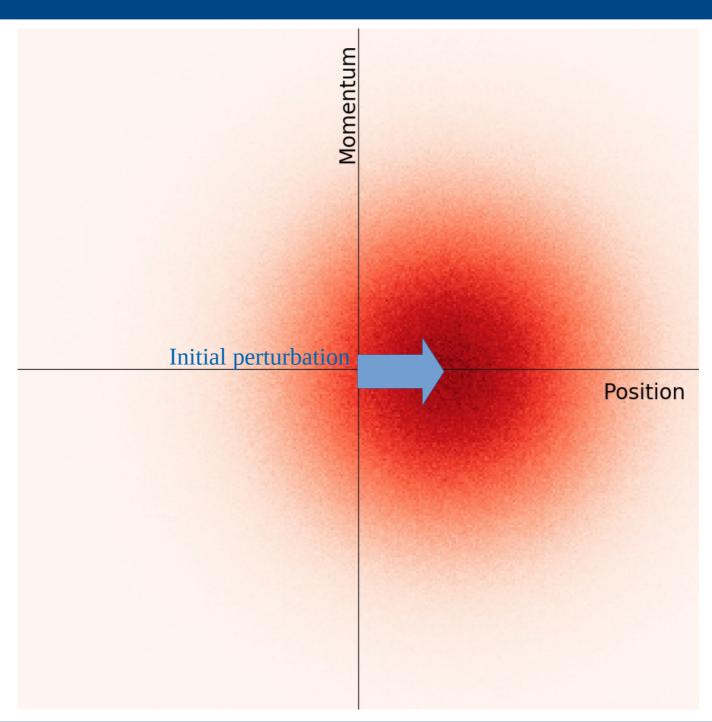


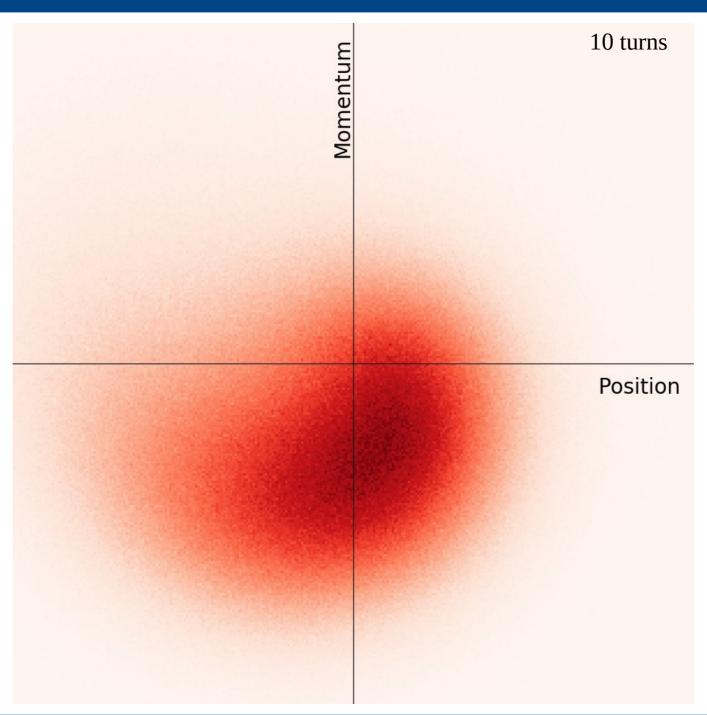


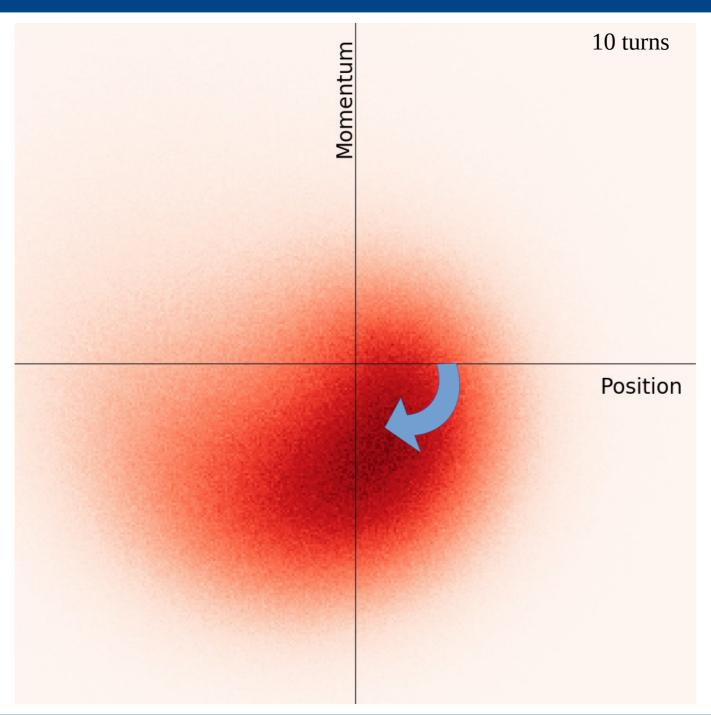


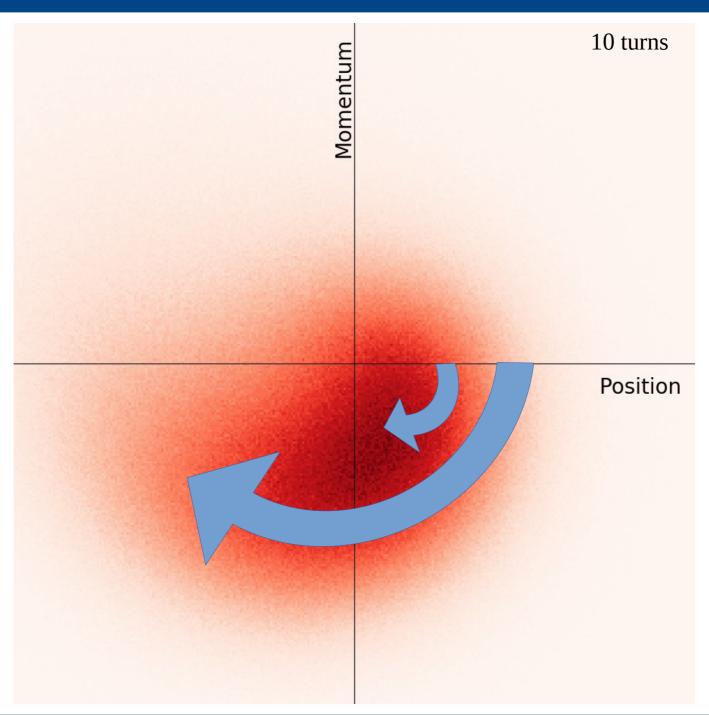


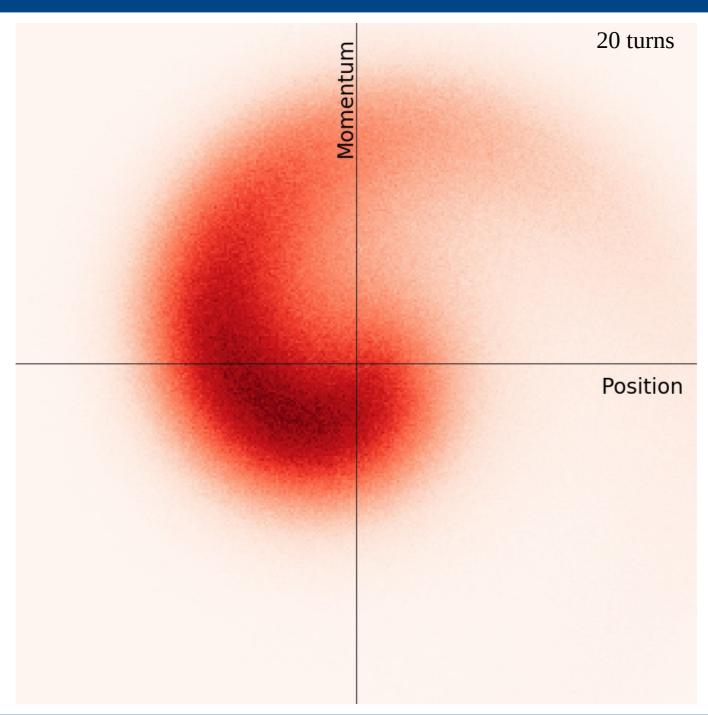


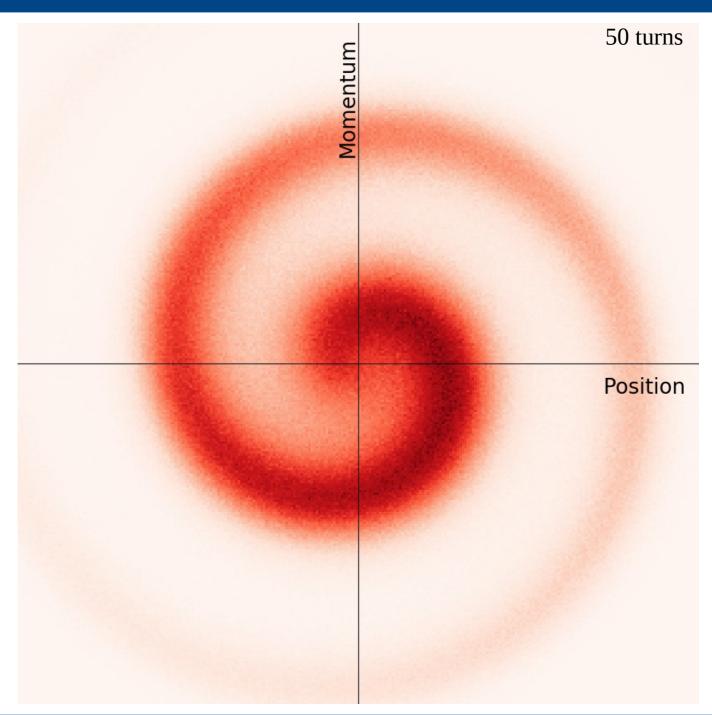


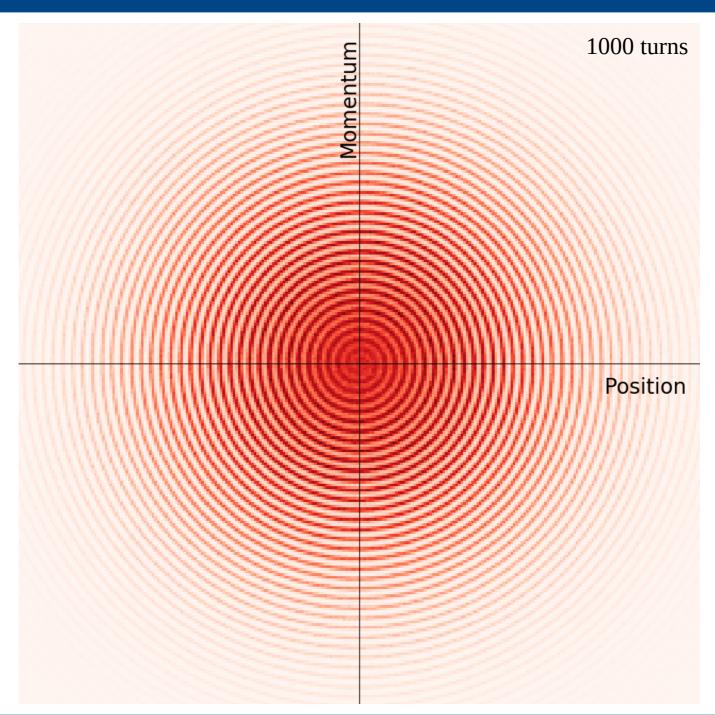


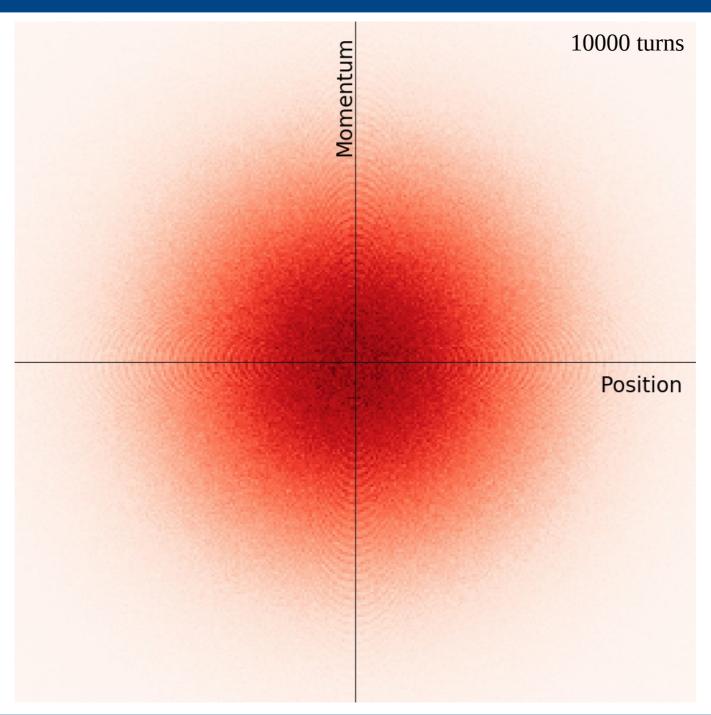


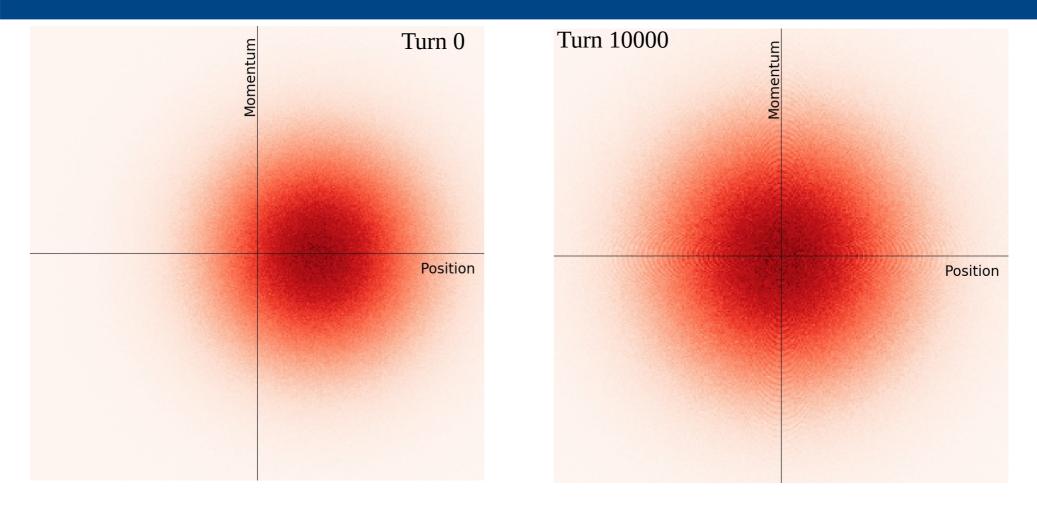




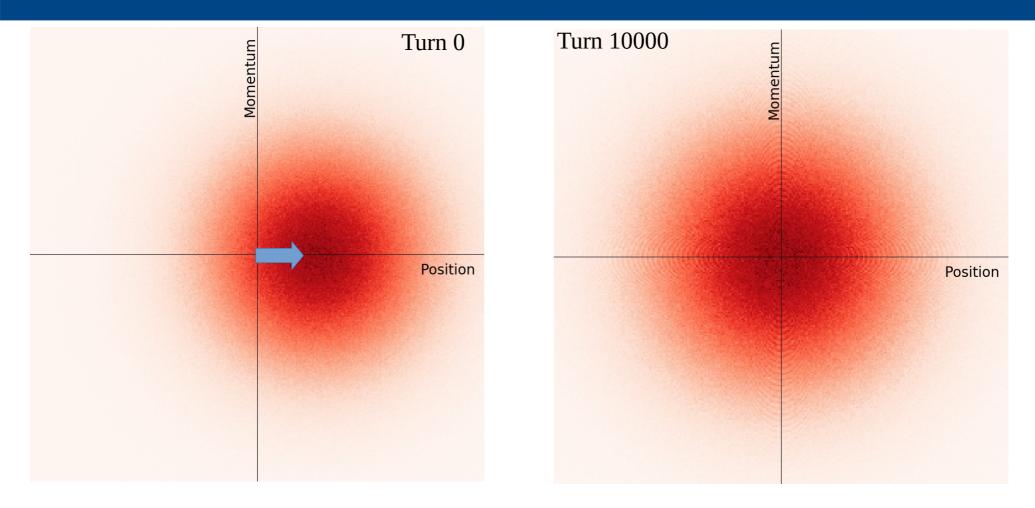




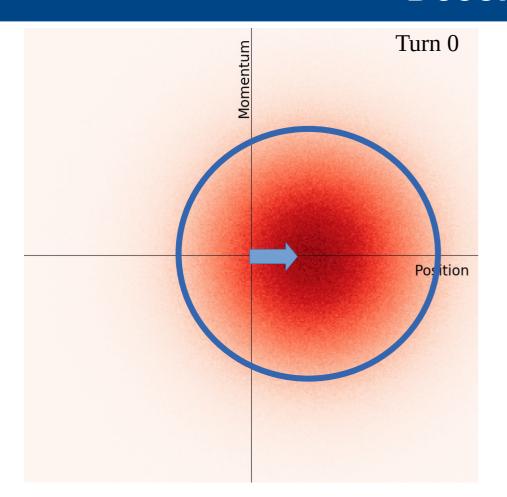


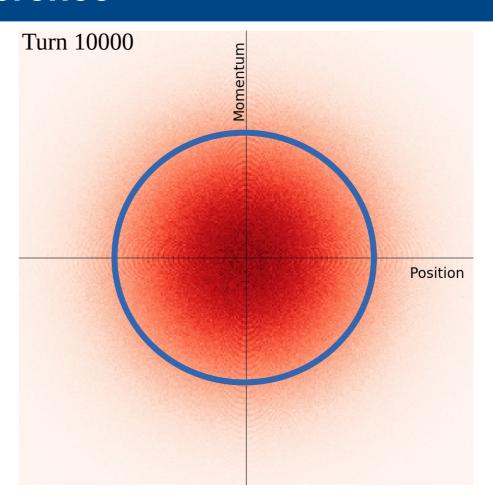


• Due to the tune spread, the initial perturbation is damped at the expense of a change of distribution → **emittance growth** 

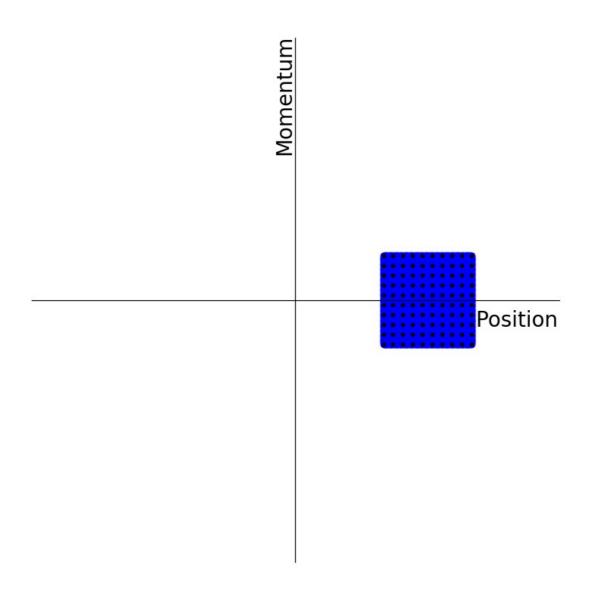


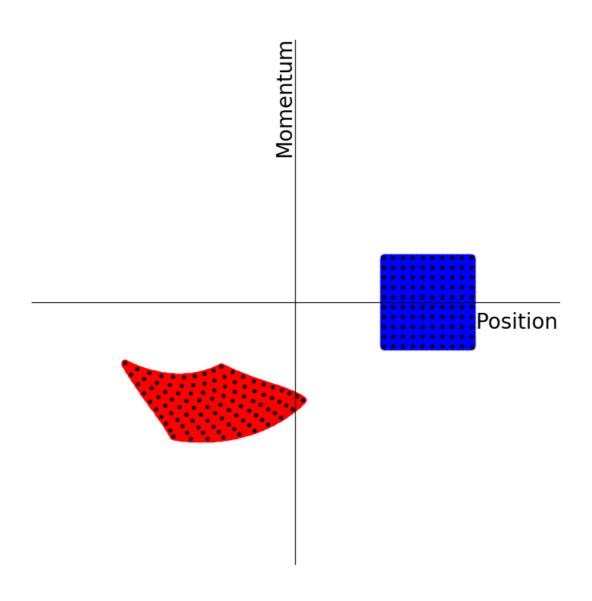
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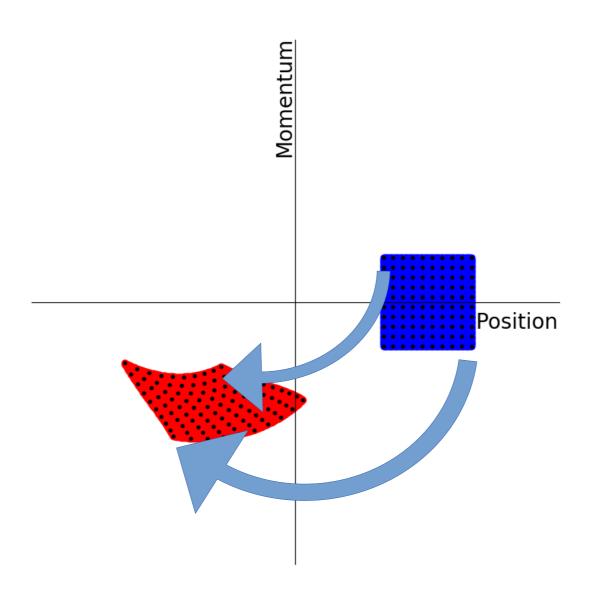


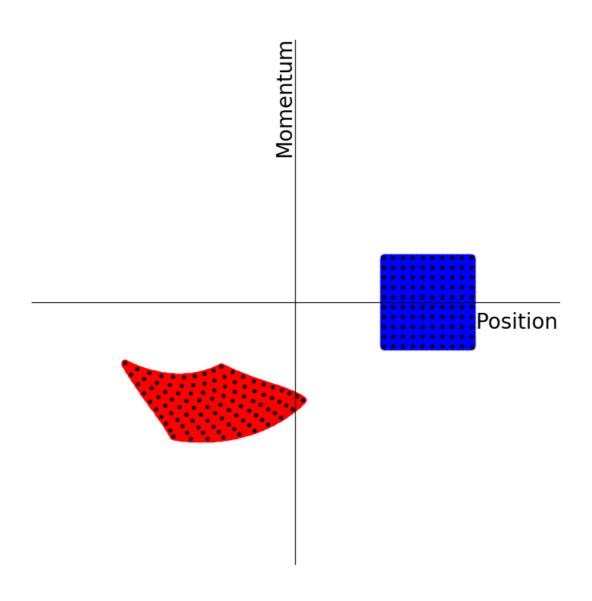


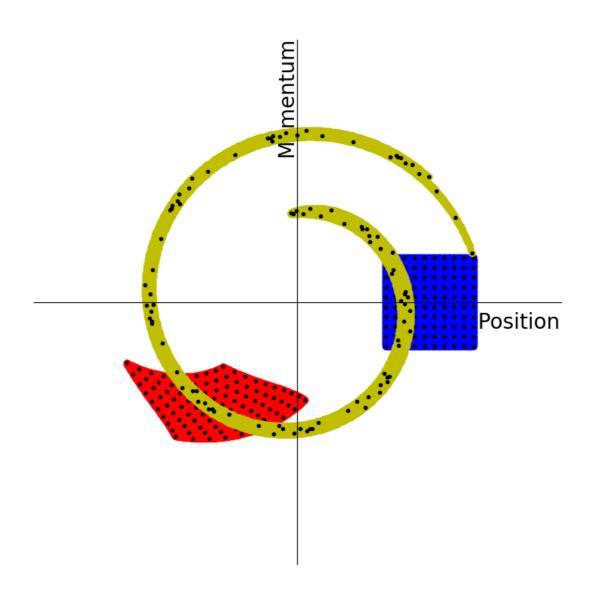
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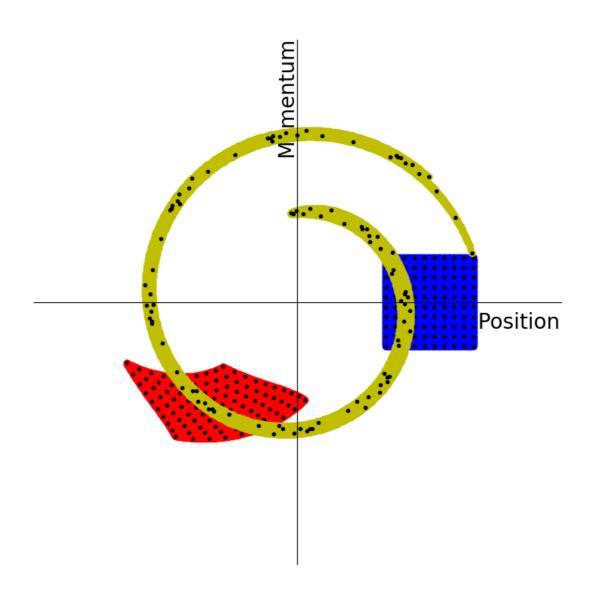


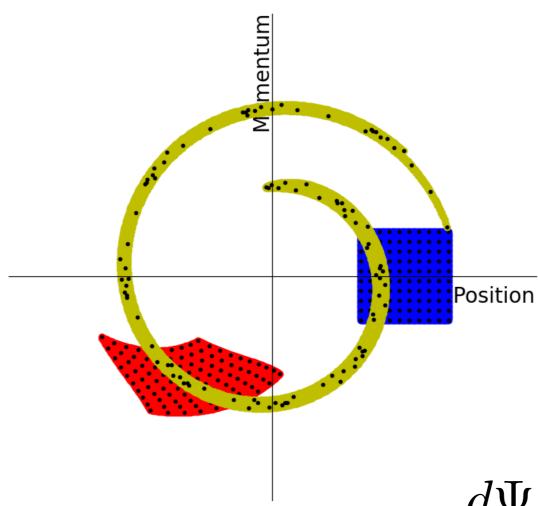






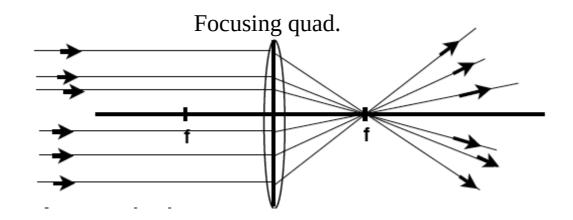


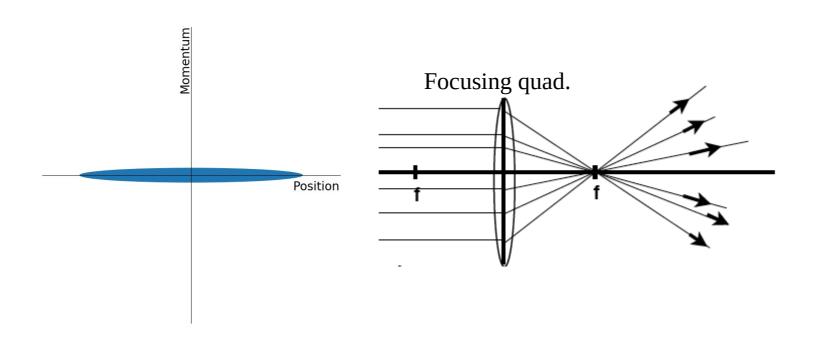


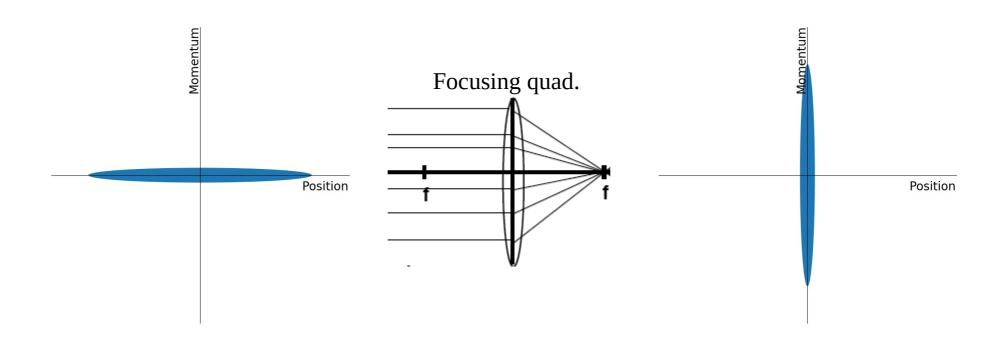


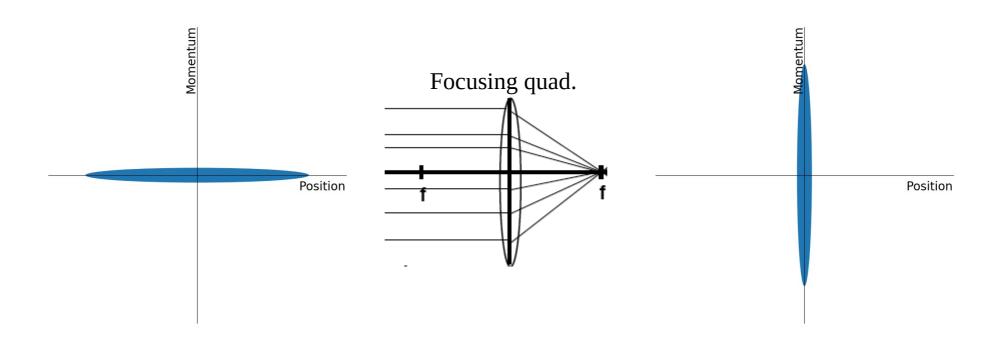
• Even with distorted trajectories,the phase-space density is preserved:

$$\frac{d\Psi}{dt} = 0$$



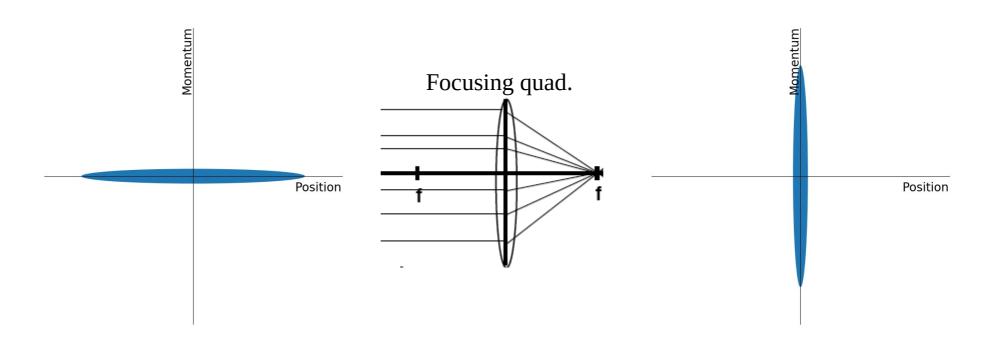






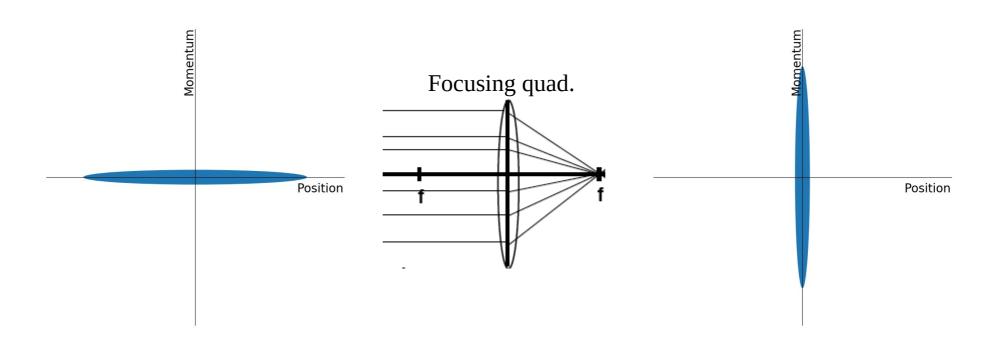
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# Liouville theorem: A simple illustration



- The conservation of the emittance is a consequence of Liouville theorem
  - Liouville is more general: The phase-space density is conserved even in the presence of non-linear forces, provided that the system can be described with Hamilton's equation

### Liouville theorem: A simple illustration



- The conservation of the emittance is a consequence of Liouville theorem
  - Liouville is more general: The phase-space density is conserved even in the presence of non-linear forces, provided that the system can be described with Hamilton's equation
    - → Non-conservative forces such as intrabeam scattering or the emission of synchrotron radiation cannot be described with Hamilton's equation: Liouville theorem does not apply

 Vlasov equation can be derived from Liouville theorem. It is a special case for plasmas (i.e. charged particles interacting 'long-range' via electromagnetic fields)

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  - → Very similar to particle beams! We can write:

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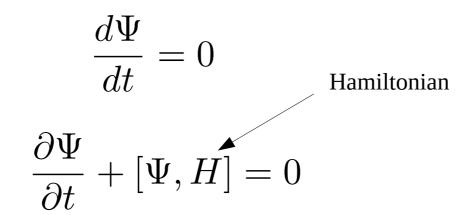
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$$\frac{\partial \Psi}{\partial t} + [\Psi, H] = 0$$

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 Hamiltonian 
$$\frac{\partial\Psi}{\partial t} + [\Psi, H] = 0$$

$$\frac{\partial \Psi}{\partial t} + \sum_{i} \frac{\partial H}{\partial p_{i}} \frac{\partial \Psi}{\partial q_{i}} - \frac{\partial H}{\partial q_{i}} \frac{\partial \Psi}{\partial p_{i}} = 0$$

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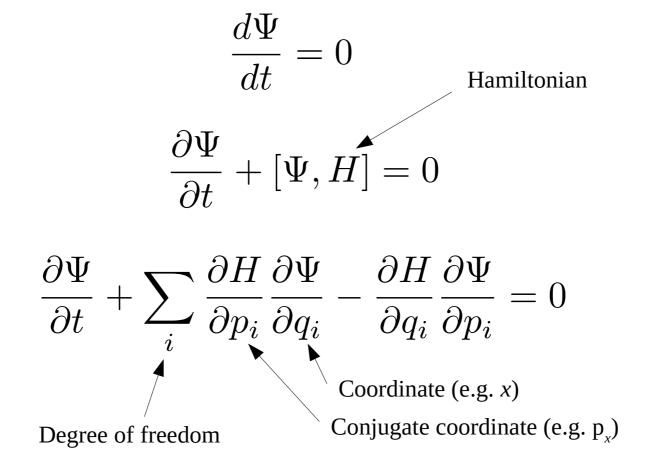
Degree of freedom

$$\frac{d\Psi}{dt} = 0$$
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 Degree of freedom Conjugate coordinate (e.g.  $\mathbf{p}_{\mathbf{x}}$ )

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$$\begin{array}{rcl}
x & = & \sqrt{2J}\cos(\theta) \\
p_x & = & \sqrt{2J}\sin(\theta)
\end{array}$$

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x & = & \sqrt{2J}\cos(\theta) \\
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\end{array} \qquad \frac{\partial\Psi}{\partial t} + \frac{\partial H_0}{\partial J}\frac{\partial\Psi}{\partial \theta} - \frac{\partial H_0}{\partial \theta}\frac{\partial\Psi}{\partial J} = 0$$

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$$H_0 = \omega_0 \left(Q_0 J + \frac{a}{2} J^2\right) \qquad \text{Hamiltonian of a harmonic oscillator with a 'simple' non-linear force}$$
 Revolution frequency

$$\begin{array}{rcl}
x & = & \sqrt{2J}\cos(\theta) \\
p_x & = & \sqrt{2J}\sin(\theta)
\end{array}$$

Linear detuning coefficient

$$H_0 = \omega_0 \left( Q_0 J + \frac{a}{2} J^2 \right)$$
Revolution Tune frequency

$$\frac{\partial \Psi}{\partial t} + \frac{\partial H_0}{\partial J} \frac{\partial \Psi}{\partial \theta} - \frac{\partial H_0}{\partial \theta} \frac{\partial \Psi}{\partial J} = 0$$

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→ Hamiltonian of a harmonic oscillator with a 'simple' non-linear force

$$\frac{\partial H_0}{\partial J} = \omega_0 \left( Q_0 + aJ \right) \equiv \omega(J)$$

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$$\rightarrow$$
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$$\frac{\partial H_0}{\partial J} = \omega_0 \, (Q_0 + aJ) \equiv \omega(J)$$

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$$\frac{\partial H_0}{\partial J} = \omega_0 \left( Q_0 + aJ \right) \equiv \omega(J)$$

Example of solution: Exponential distribution in action (Gaussian in x,  $p_x$ ):

$$\Psi_0 = \frac{1}{2\pi\epsilon} e^{-\frac{J}{\epsilon}}$$

# Perturbation theory with an external force

[Ruggiero]

Let's consider a first order perturbation of the distribution:

$$\Psi = \Psi_0 + \Psi_1$$

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 And a first order perturbation of the Hamiltonian by an external force:

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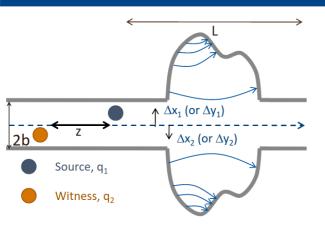
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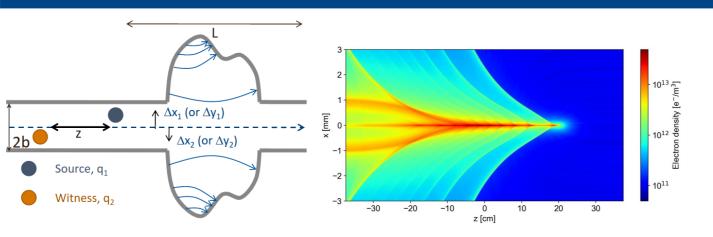
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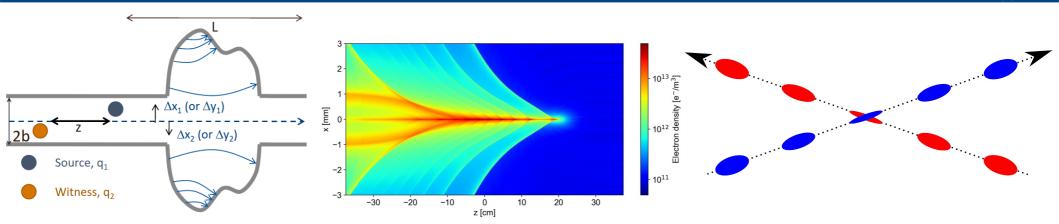
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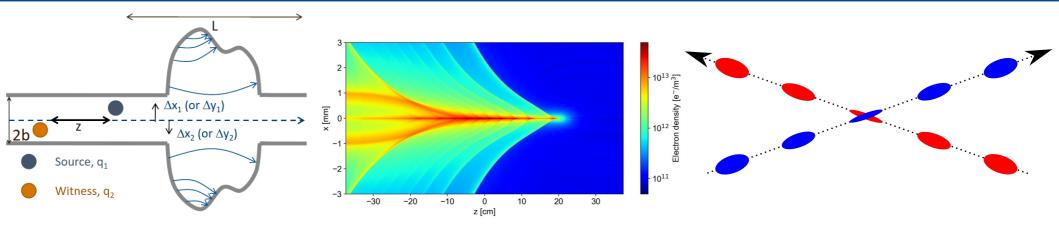
First order perturbation of Vlasov equation:

$$\frac{\partial \Psi_1}{\partial t} + \omega(J) \frac{\partial \Psi_1}{\partial \theta} - \sqrt{2J} \sin(\theta) F_{ext} \frac{\partial \Psi_0}{\partial J} = 0$$



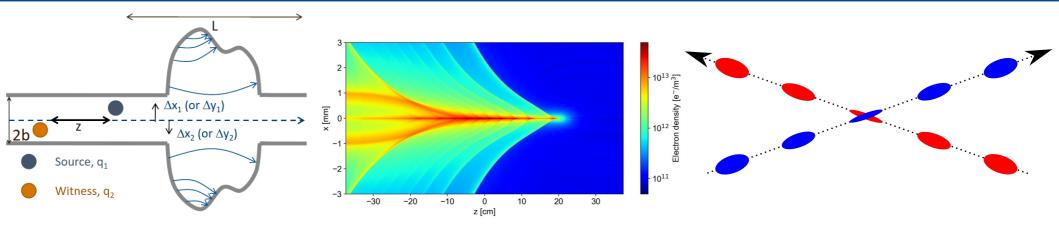






Simple model for the collective force:

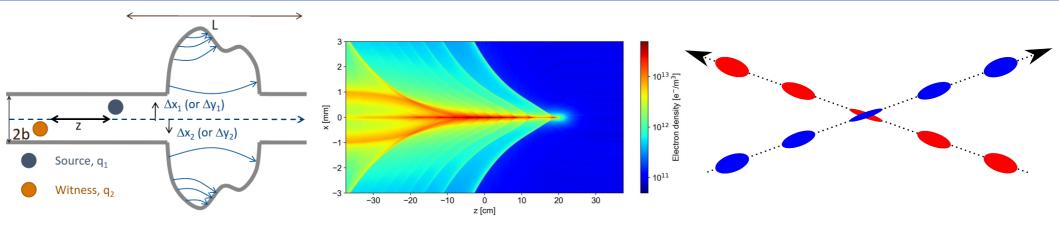
$$F_{ext} = -2\Delta\Omega_{ext}\langle x\rangle$$



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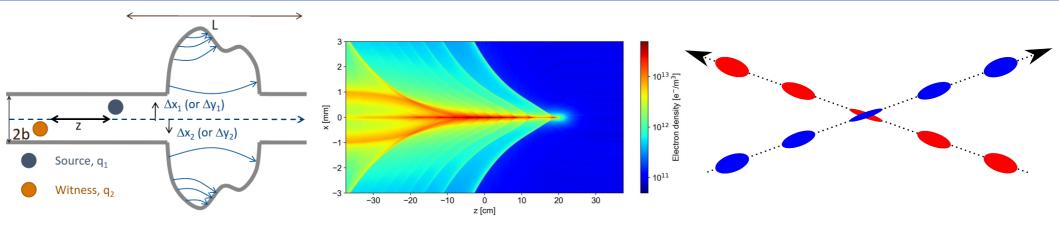
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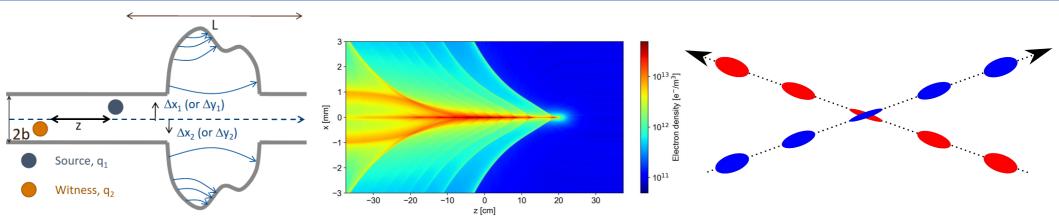
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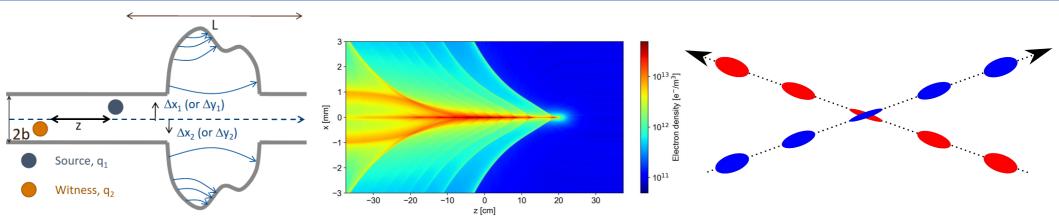
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$$g_c = \frac{-1}{2} \Delta \Omega_{ext} \frac{\sqrt{2J} \frac{df_0}{dJ}}{\Omega_c - \omega}$$

Coherent mode 
$$(\Omega-\omega)g=\frac{-1}{2}\Delta\Omega_{ext}\frac{df_0}{dJ}\sqrt{2J}\int\int dJ\sqrt{2J}g \int_{=1}^{\infty}dJ$$

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 The dispersion relation links the coherent mode frequency with the frequency shift due to the collective force via the tune spread

$$\frac{-1}{\Delta\Omega_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega_c - \omega(J)}$$

$$f_0 = \frac{1}{\epsilon} e^{-\frac{J}{\epsilon}} \qquad \omega(J) = \omega_0 \left( Q_0 + aJ \right)$$

## The stability diagram

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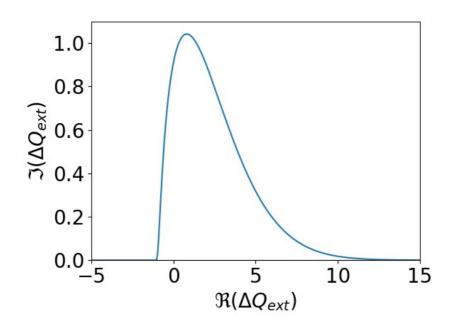
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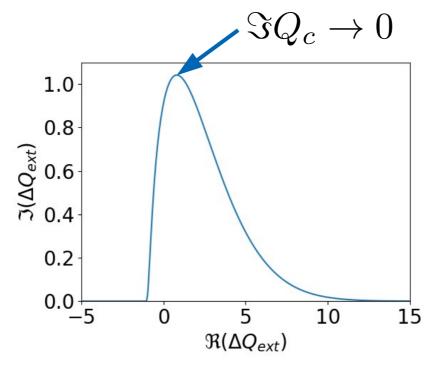
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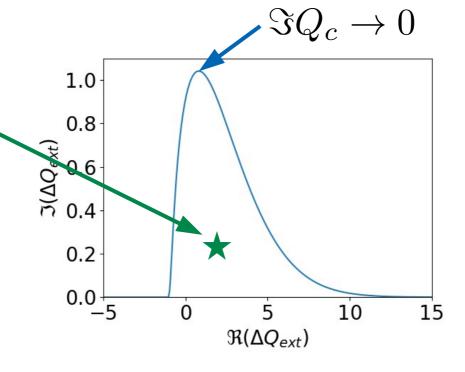
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If the tune shift due to the collective force is **inside** of the stability diagram

$$\Im Q_c < 0$$

The beam is stable



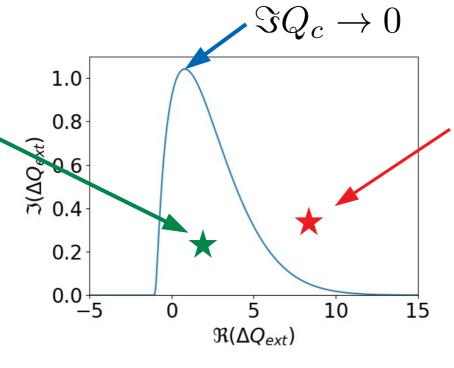
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The beam is unstable

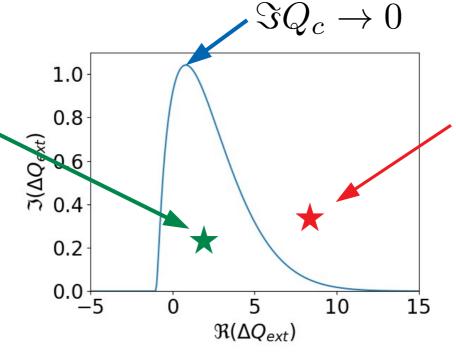
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 The stability diagram is a very common way of representing Landau damping when the impact of the collective force can be represented by a complex tune shift

$$F_{ext} = A_{ext}e^{-i\Omega t}$$

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 We look for harmonic solutions resonant with the excitation:

$$\Psi_1 = g(J)e^{i(\theta - \Omega t)}$$

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$$\frac{\langle x \rangle}{A_{ext}} = \int dJ d\theta x g = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)}$$

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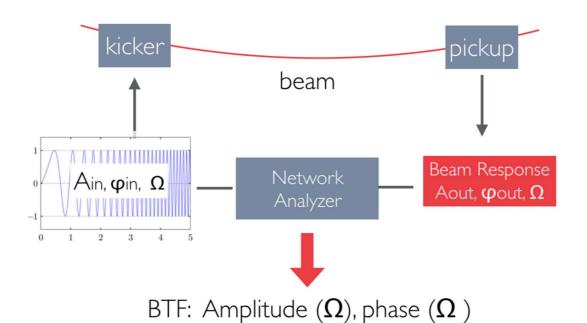


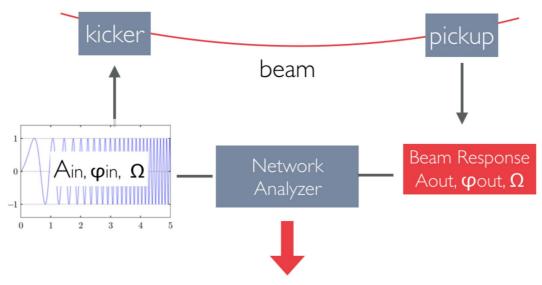
$$\frac{g}{A_{ext}} = \frac{1}{2} \frac{\sqrt{2J} \frac{df_0}{dJ}}{\Omega - \omega(J)}$$



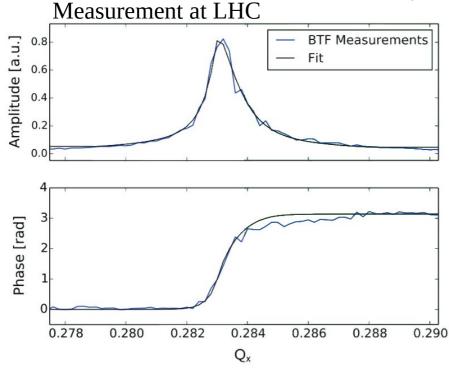
$$\frac{\langle x \rangle}{A_{ext}} = \int dJ d\theta x g = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)}$$

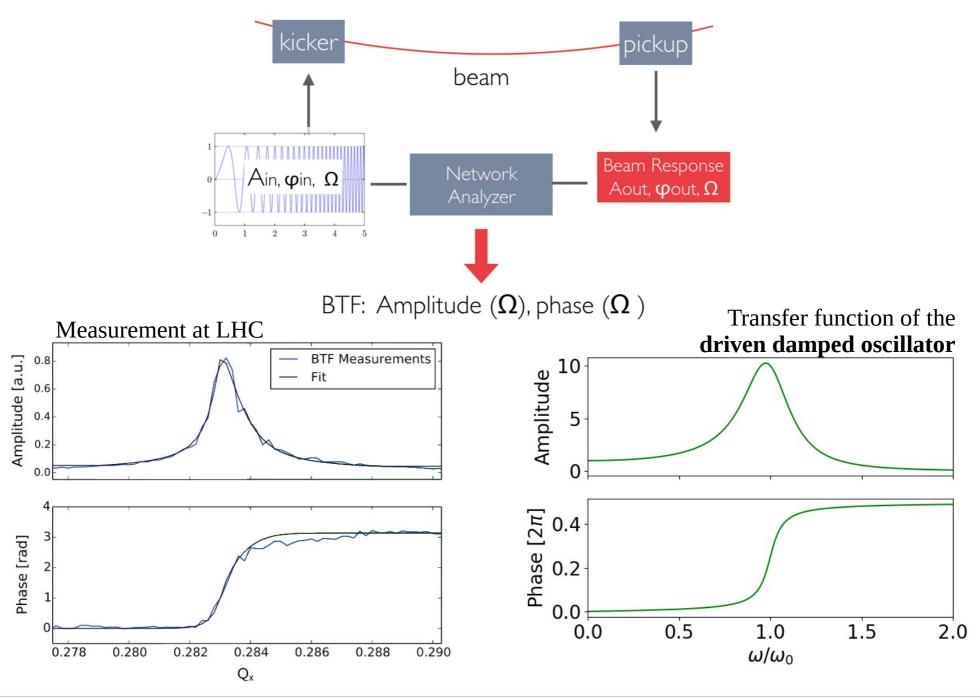
 The beam oscillation amplitude normalised to the excitation amplitude is called the beam transfer function → A measurable quantity that directly relates to the stability diagram





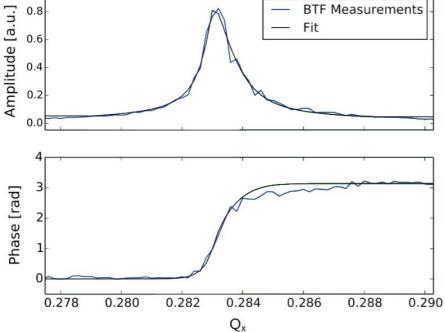
BTF: Amplitude  $(\Omega)$ , phase  $(\Omega)$ 





$$\frac{\langle x \rangle}{A_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)}$$

# Measurement at LHC



## Beam transfer function and stability diagram

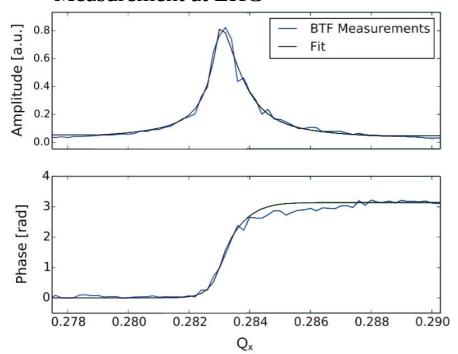
[Tambasco]

$$\frac{\langle x \rangle}{A_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)}$$



$$\frac{\langle x \rangle}{A_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)} \qquad \qquad \frac{-1}{\Delta \Omega_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega_c - \omega(J)}$$

#### Measurement at LHC

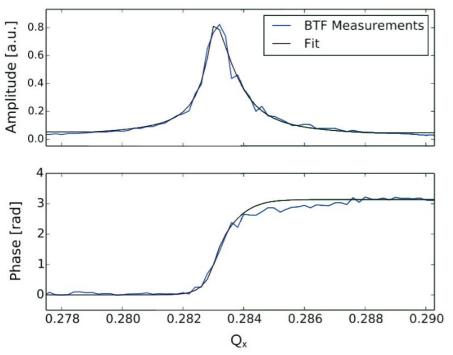


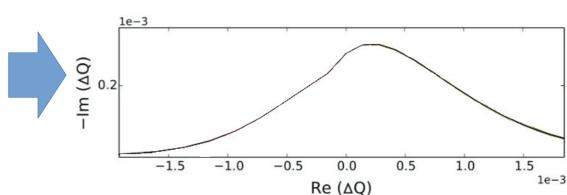
$$\frac{\langle x \rangle}{A_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)}$$



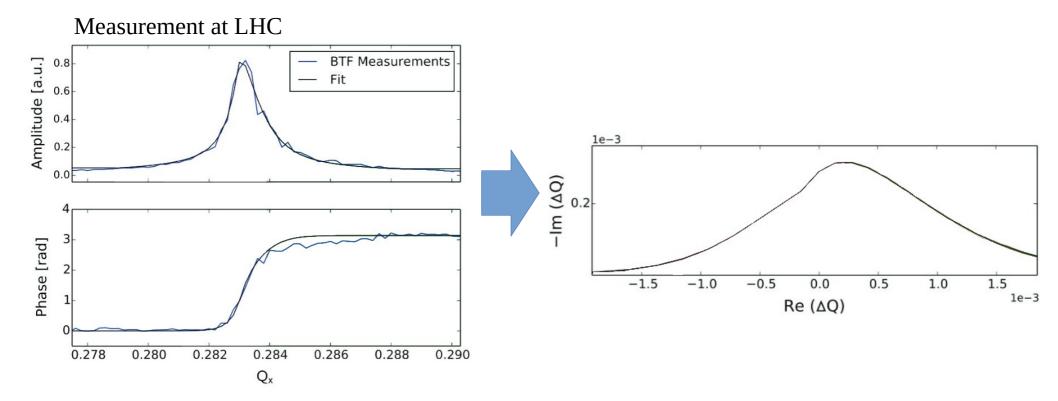
$$\frac{\langle x \rangle}{A_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)} \qquad \qquad \frac{-1}{\Delta \Omega_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega_c - \omega(J)}$$

#### Measurement at LHC





$$\frac{\langle x \rangle}{A_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)} \qquad \qquad \frac{-1}{\Delta \Omega_{ext}} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega_c - \omega(J)}$$



The BTF is an interesting way to quantify experimentally Landau damping

- Landau damping stems from the interaction of single particles with waves
  - A necessary condition for Landau damping is the a comparable velocity / frequency of the wave and the particles motion

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- While collective forces such as wake fields or electron clouds tend to generate unstable modes of oscillation, Landau damping stabilises them without emittance growth
  - An external perturbation may also decay through a similar phenomenon, we rather talk about decoherence or filamentation. This mechanism leads to emittance growth
- Landau damping originates in the spread of oscillation frequencies of the particles in the beam
  - It is a linear mechanism, as in plasmas. However in accelerators the frequency spread often originates from non-linear forces

### "Now what?"

– Fuego,a down-to-earth rabbit



• Ok, in the second part we'll address practical applications...

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