Tops Beyond the Asymmetry

Andreas Weiler CERN&DESY

> Planck 2011 Lisbon

> > 3/6/2011

Naturalness vs. flavor triviality

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EWSB & Flavor

 $\frac{\Lambda^2}{\Lambda^2_{\text{Flavor}}} (\bar{\psi}\psi)(\bar{\psi}\psi)$

 $y_{ij}\, \bar{\psi}_i \mathcal{H} \psi_j$



E

naturalness







 composite Higgs vs. MFV (EWPT, top mass, EDMs) w/ Michele Redi (CERN/INFN) and ideas by R. Rattazzi
 chromo-electric dipole moment of the top
 w/ Jernej Kamenik (IJS, Ljubljana) Michele Papucci (CERN/LBL)

3) Flavor & Naturalness & susy breaking w/ Michele Papucci (CERN/LBL)

if time allows

Old Flavor problem of composite Higgs Higgs as bound state, naively $D_{\mathcal{H}=\langle \bar{\psi}\psi \rangle} \approx 3$ $\frac{1}{\Lambda^{D_{\mathcal{H}}-1}} y_{ij} \bar{\psi}_i \mathcal{H}\psi_j + \frac{1}{\Lambda^2} c_{ijkl} \bar{\psi}_i \psi_j \bar{\psi}_k \psi_l$

A can not be too large, because want top mass

 $\Lambda = \mathcal{O}(\text{TeV})$

=> talk by Rychkov

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A can not be too large, because want top mass

A must be very large because this leads to FCNCs

K

 $\Lambda = \mathcal{O}(\text{TeV}) \quad \Lambda > 10^5 \text{ TeV}$

=> talk by Rychkov

Two ways of giving mass to fermions...

Bi-linear (like SM):

 $\mathcal{L} = y f_L \mathcal{O}_H f_R, \quad \mathcal{O}_H \sim (1,2)_{\frac{1}{2}}$



Linear:

D.B. Kaplan '91 $\mathcal{L} = yf_L\mathcal{O}_R + y_Rf_R\mathcal{O}_L + m\mathcal{O}_L\mathcal{O}_R, \quad \stackrel{q_i}{\mathcal{O}_R} \sim (3,2)_{\frac{1}{6}H}$

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 q_i

 $\mathcal{L} = y f_L \mathcal{O}_R + y_R f_R \mathcal{O}_L + m \mathcal{O}_L \mathcal{O}_R, \quad \overset{q_i}{\mathcal{O}_R} \sim (3, 2)_{\frac{1}{6}}$

Quarks & Leptons mix with strong sector

mass \propto compositeness



Partial compositeness $|SM\rangle = \cos\phi|elem.\rangle + \sin\phi|comp.\rangle$ $|heavy\rangle = -\sin\phi|elem.\rangle + \cos\phi|comp.\rangle$ Composites are heavy ($m_{\rho} \approx \text{TeV}$). Light quarks have very little composite admixture.

mixing
$$\sim$$
 mass
strong sector
Higgs&EVVSB
top
resonances
 ρ_{μ}
elementary fields
 u, d, c, s, b, A_{μ}

 $g_*, m_{
ho}$

 $1 \lesssim g_* \lesssim 4\pi$

Kaplan; Contino, Kramer, Son, Sundrum

high p_T

Resonance production (option 1)



 $\sim g_*^2 \sin^2 \theta_{u_R}$

strongly suppressed for light quarks!

high p_T Resonance production (option 2) U ρ gluon $\sim rac{g_s}{g*}$ U analog to $\gamma - \rho$ mixing

NB, gluon-rho-rho = 0

high pt

t,b

t,b

Resonance decay

ρ

decays dominantly into 3rd generation! (tt, bt, bb)

CP problems

Csaki, Falkowski, AW; Buras et al; Casagrande et al $g_* \approx Y_* \approx 3 \dots 6$

$\Delta F = 2$ (strongest from ϵ_K)



$$M_* \gtrsim 10 \left(\frac{g_*}{Y_*}\right) \text{TeV}$$

 $\Delta F = 1$ (strongest constraint from ϵ'/ϵ)

Gedalia et. al



 $\Delta F = 0$ neutron EDM => talk by R. Ziegler

$$M_* \gtrsim 1.3 Y_* \,\mathrm{TeV}$$

$$M_* \ge 2.5 Y_* \,\mathrm{TeV}$$

Agashe et. al, Delaunay et. al, Redi, AW

Flavor breaking external Michele Redi, AW Postulate flavor agnostic strong sector

Usually tension between large top mass & universal mixings and EWPT Cacciapaglia, Csaki, Galloway, Marandella, Terning, A.W., '08,

either some flavor breaking or special
 flavor dynamics ("shining") => talk by C. Delaunay

Found a simpler, easily discoverable model that avoids these problems and doesn't require a flavorful strong sector!

MFV with split LH doublets Michele Redi, AW Main Idea: mixing w/ split LH doublets*

 $Q_{Lu} (2,2)_{2/3} Q_{Ld} (2,2)_{-1/3} u_R (1,1)_{2/3} d_R (1,1)_{-1/3}$

Strong sector $SU(3)_F$

 $egin{aligned} \lambda_{Lu} \propto y_u, & \lambda_{Ld} \propto y_d \ \lambda_{Ru} \propto Id & \lambda_{Rd} \propto Id \end{aligned}$

In the limit that $y_u, y_d \rightarrow 0 => U(3)^3$ aka **GIM**

* inspired by composite Higgs. NB, in anarchic scenario this is completely unsafe and leads generically to large FCNC's.

EWPT&compositeness Michele Redi, AW

$$\delta g \approx \frac{Y^2 v^2}{2 m_{\rho}^2} \sin^2 \varphi_{q_L} \left(T_{3L}'(Q) - T_{3L}(q_L) \right) + g_2^2 \frac{v^2}{4 m_{\rho}^2} \sin^2 \varphi_{q_L} (T_{3R} - T_{3L})$$

Hadronic width



Dijet search (by ATLAS/CMS)

$$\sin^2 arphi_{q_{L,R}} \lesssim rac{1}{g_
ho} \left(rac{m_
ho}{3\,{
m TeV}}
ight)$$

RH quarks can be easily fully composite!



high pt in MFV

Resonance production



=> O(I) coupling to light quarks!

"Ultra-visible"

Michele Redi, AW

production cross-section of color octet spin-I:

(similar conclusions for spin 1/2)

$g_{ ho}$	$\sin \varphi_{u_R}$	$\sin \varphi_{d_R}$	$\sigma(\mathrm{pb})$	$\Gamma(\text{GeV})$	${ m Br}(uar{u})$	$\operatorname{Br}(t_L \bar{t}_L)$	${ m Br}(t_R ar{t}_R)$
3	RS	RS	3.3	425	0.05	0.045	0.76
5	RS	RS	2.7	1075	0.006	0.0005	0.97
3	0.7	0.7	41	825	0.17	0.0004	.16
5	0.5	0.5	24	645	0.17	0.01	0.15
2	1	.25	74	990	0.30	0.006	0.27
3	1	.25	375	2200	0.33	0.0008	0.032

> 10 x anarchic (RS) model !



$$\delta \sigma_{\rm 700-800\,GeV}^{\rm NP} / \sigma_{\rm 700-800\,GeV}^{\rm SM} \lesssim 47\%$$

 $\delta \sigma_{t\bar{t}}^{\rm NP} / \sigma_{t\bar{t}}^{\rm SM} \lesssim 10\%$

MFGB <900 GeV, geff ~O(1) $A_{FB}^{t\bar{t}}(M_{inv} > 450 GeV) \stackrel{<}{\sim} 10\%$

@ σ_{NP}/σ_{SM}(pT>400GeV): 2-3



2) chromo-magnetic EDM of the top

I I 06.xxxx with Jernej Kamenik (IJS, Ljubljana) & Michele Papucci (CERN/LBL)

Measure CPV of the top?

Chromo-electric and chromo-magnetic dipole

 $\mathcal{L}_{t\bar{t}G} = -g_s \,\bar{t}\gamma^{\mu}G_{\mu}t - i\frac{d'_t}{2} \,\bar{t}\sigma^{\mu\nu}\gamma_5 G_{\mu\nu}t - \frac{\mu'_t}{2} \,\bar{t}\sigma^{\mu\nu}G_{\mu\nu}t$

CEDM

CMDM

Can be sizable in partial compositeness, RS, susy (light top partners required by naturalness)

ttbar cross-section

P. Haberl, O. Nachtmann, A. Wilch

$qq \rightarrow tt$

$$\frac{d\hat{\sigma}_{q\bar{q}}}{d\hat{t}} = \frac{\pi\alpha_s^2}{\hat{s}^2} \frac{8}{9} \left(\frac{1}{2} - v + z - 2\hat{\mu}_t' + (\hat{\mu}_t'^2 - \hat{d}_t'^2) + (\hat{\mu}_t'^2 + \hat{d}_t'^2)\frac{v}{z}\right)$$

and similar for gg initial state

see e.g. Hioki et al, Peskin et al, many more

Experimental constraints

 $\sigma_{\rm obs}^{1.96 \text{ TeV}} = 7.5 \pm 0.31 (\text{stat}) \pm 0.34 (\text{syst}) \pm 0.15 (\text{lumi}) \text{ pb}$

Tevatron

ATLAS

CMS

$$\sigma_{700} \equiv \sigma^{t\bar{t}}(700 \text{ GeV} < M_{t\bar{t}} < 800 \text{ GeV}) = 80 \pm 37 \text{ fb},$$

five-channel combined 180 ± 9

 $180 \pm 9 \text{ (stat.)} \pm 15 \text{ (syst.)} \pm 6 \text{ (lumi.)}$

$$158 \pm 10 \pm \frac{15}{15} \pm 6$$

(36 pb⁻¹)

Ahrens et al. ,Kidonakis

	Teva	itron	LHC7		
	MSTW	CTEQ	MSTW	CTEQ	
LO	$6.66^{+2.95+(0.34)}_{-1.87-(0.27)}$	$5.45^{+2.16+0.33(0.29)}_{-1.42-0.27(0.24)}$	$122^{+49+(6)}_{-32-(7)}$	$100^{+35+9(7)}_{-24-8(7)}$	
NLO	$6.72^{+0.41+0.47(0.37)}_{-0.76-0.45(0.24)}$	$6.77_{-0.74-0.40(0.34)}^{+0.40+0.50(0.43)}$	$159^{+20+14(8)}_{-21-13(9)}$	$148^{+18+13(11)}_{-19-12(10)}$	
NNLO approx.	$6.63^{+0.07+0.63(0.33)}_{-0.41-0.48(0.25)}$	$6.91^{+0.09+0.53(0.46)}_{-0.44-0.43(0.36)}$	$155^{+8+14(8)}_{-9-14(9)}$	$153^{+8+13(11)}_{-8-12(10)}$	

vs. SM theory:

Top Chromo-electric dipole



mathematica + mstw08 pdf's, SM normalized to ''N''NLO

Top Chromo-electric dipole



Weinberg op. mixes into chromo-magnetic top

But chromo-magnetic top does not mix into Weinberg operator (which would lead to neutron EDM)

QFT: no mixing of higher dim op.'s in lower dim.

$$O_{e}^{q} = -\frac{i}{2}e Q_{q}m_{q} \bar{q}\sigma^{\mu\nu}\gamma_{5}q F_{\mu\nu},$$

$$O_{c}^{q} = -\frac{i}{2}m_{q}\bar{q}\sigma^{\mu\nu}t^{a}\gamma_{5}q G_{\mu\nu}^{a},$$

$$O_{G} = -\frac{1}{6}f^{abc}G_{\mu\rho}^{a}G_{\nu}^{b\rho}G_{\lambda\sigma}^{c}\epsilon^{\mu\nu\lambda\sigma}$$

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$$O_{G} = -\frac{1}{6}f^{abc}G_{\mu\rho}^{a}G_{\nu}^{b\rho}G_{\lambda\sigma}^{c}\epsilon^{\mu\nu\lambda\sigma}$$

But there's a finite part at loop!



$$C_G(m_Q^{-}) = C_G(m_Q^{+})$$
$$+ C_Q(M) \left(\frac{g_s(m_Q)}{g_s(M)}\right)^{\gamma_{qq}/\beta} \frac{1}{8\pi} \frac{\alpha_s(m_Q)}{m_Q}$$

Chang et. al, Braaten et. al 90's

Finite threshold correction

O m_{top} $C_{\text{Weinberg}} = \frac{g_s^2}{32\pi^2}$

$$\frac{1}{2}$$
 CEDM_{top}

Weinberg 89, Pospelov et. al



$$C_G(m_Q^{-}) = C_G(m_Q^{+})$$

+ $C_Q(M) \left(\frac{g_s(m_Q)}{g_s(M)}\right)^{\gamma_{qq}/\beta} \frac{1}{8\pi} \frac{\alpha_s(m_Q)}{m_Q}$

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Finite threshold correction

0 m_{top} $C_{\text{Weinberg}} = rac{g_s^2}{32\pi^2} \operatorname{CEDM_{top}}$

Weinberg 89, Pospelov et. al

→ Weinberg op contributes to neutron EDM!

 $d_n(w) = e(10 - 30) \text{ MeV} w(1 \text{ GeV})$

Top Chromo-electric dipole



Top Chromo-electric dipole Kamenik, Papucci, W Our new constraint is ~ 100x stronger than direct collider constraints and somehow has been missed.

Cross-section limits are now not sensitive any more, but CPV observables $(a.(b \times c))$ might be stronger.

Interesting to study in any non CPV-conserving scenarios with light top partners (any natural model...)!

3) Naturalness vs. flavor blind susy breaking

Fast forward to Fall 2011 (or even 2012?) and nothing new yet

in progress w/ Michele Papucci (CERN/LBL)
susy is the prime example for flavor triviality, use flavor blind susy breaking (GMSB) and decouple flavor genesis.

Still natural after 5 1/fb LHC7?



Friday, May 27, 2011

2 ESSISVER ANDESS

In the Higgs decoupling limit of the MSSM, the lower bound on the mass of O Gluinos enter the higgs potential @ 2100p, with the 114.4 GeV Griass is bounded too reaction at ATTIMOLOOPS to be cannot be too HEAVY EITHER:

o The MSSM is tuned, tension from LEP Higgs bound O(1 in 100) = 1741 GeV = 28 talk by Djouadi

$$m_h^2 \simeq m_Z^2 \cos^2 2\beta + \frac{3}{4\pi^2} \frac{m_t^4}{v^2} \left[\log \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{X_t^2}{12m_{\tilde{t}}^2} \right) \right] \quad X_t = A_t - \mu \cot \beta$$

loop quaequaticorr(et) ionshouting first toptaildustop hoops(28;28; 80a.ss1, 32), 6 first term 33]. Herijggis tuning may be reduced by extending ousquaged the Higgs stop masses and out ingrow the stop Geass where and values bes(2) the light of exact coupling spars with the masgrow the stop Geass where and values bes(2) the light of exact files way 27, 2011 assumes the stop of the second files of the stop of the second files of the second files way 27, 2011 assumes the stop of the second for a stop of the second for a stop of the second files way 27, 2011 assumes the stop of the second for a stop of the second for a stop of the second files way 27, 2011 assumes the stop of the second for a stop of the second for

A natural (light) 3rd gen? A LIGHT 3RD GENERATIKS by Hoecker, Boccali o Present searches do not exclude light stops/



o Limits well below 250-350 GeV if gluinos, squarks_{1,2}, ... are decoupled

sbottoms, susy can still be natural

Friday, May 27, 2011

Anatural (light) 3rd gen? talks by Hoecker, Boccali

o Present searches do not exclude light stops/ sbottoms, susy can still be natural



o Limits well below 250-350 GeV if gluinos, squarks_{1,2}, ... are decoupled

LIMITEIMIts On Squarks & Glubosnos talks by Hoecker, Boccali Early (Jets+MET+(0,1) lepton) searches put strong limits on gluinos first two generations of squarks



Already have to be > 500-600 GeV!

Friday, May 27, 2011

Splitting the generations stop mass

 $\begin{array}{ll} & \underset{M_{\tilde{t}}}{\underset{\tilde{t}}{\sim}} & \underset{M_{t}}{\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}(A_{t} - \mu/\tan\beta)}{\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}(A_{t} - \mu/\tan\beta)}{\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}}}{\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}(A_{t} - \mu/\tan\beta)}{\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}(A_{t} - \mu/\tan\beta)}{\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}}}{\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}}} & \underset{M_{t}^{2} + D_{u_{R}}} {\overset{(m_{t}^{2})}{\sim}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}} & \underset{M_{t}^{2} + \Delta/3 + m_{t}^{2} + D_{u_{R}}} & \underset{M_{t}^$

$$\Delta \simeq \frac{3y_t^2}{4\pi^2} \left(f_{33}(t) M_3^2 + f_{3t}(t) M_3 A_t \qquad t = \ln \left(M_{mess} / m_{\tilde{t}} \right) + f_{tt}(t) A_t^2 + f(t) (m_{H_u}^2 + m_{Q_3}^2 + m_{U_3}^2) + \dots \right)$$

Friday, May 27, 2011 Flavor universal boundary conditions: $M_3, A_t, m_{H_u}, M_{Q_3}, M_{u_3}$ bounded from above \rightarrow splitting between 1,2 vs. 3 is bounded from above

Fall 2011 (2012?)

Difficult to extrapolate experimental constraints, (analysis in final stages), **BUT**

Stop/Bottom: most holes closed this summer, increased limit $(350 \sim 400 \text{ GeV})$

Gluinos & light squarks limits well above I TeV, likely around 1.5 TeV.

Already some tension but still not high tuning!

Split spectrum necessary for naturalness

Fall 2011 (2012?)

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Stop/Bottom: most hole **versal** increased limit (350 **universal** Gluinos & **both addition** likely a **contreation** well above I TeV, likely a **contreation** but still not high tuning!

Split spectrum necessary for naturalness

General message

Natural Susy surviving this fall (w/ no observed signal) likely involves flavor non-universal susy breaking (''flavorful susy''). => talk by Isidori, Lalak

Caveats:

o squeezing the spectrum (tuning?) o engineer small missing E_T (not looking for it) o R-parity violation ...

Conclusions

MFV with split doublets solves flavor and EWPT problems of partial compositeness, early discovery at LHC.

A new & very strong bound on the CEDM tough to see at LHC. BUT: Very important to test models with large CPV in the top sector!

Supersymmetry can still be natural but if nothing is seen until 2011 ('12?), flavor blind susy breaking is tuned.

The SM flavor puzzle

 $Y_D \approx \operatorname{diag} \left(\begin{array}{ccc} 2 \cdot 10^{-5} & 0.0005 & 0.02 \end{array} \right)$ $Y_U \approx \left(\begin{array}{ccc} 6 \cdot 10^{-6} & -0.001 & 0.008 + 0.004i \\ 1 \cdot 10^{-6} & 0.004 & -0.04 + 0.001 \\ 8 \cdot 10^{-9} + 2 \cdot 10^{-8}i & 0.0002 & 0.98 \end{array} \right)$

Origin of this structure?

Other dimensionless parameters of the SM: $g_s \approx I$, $g \approx 0.6$, $g' \approx 0.3$, $\lambda_{Higgs} \approx I$, $|\theta| < 10^{-9}$

Operator	Bounds on Λ in TeV $(c_{ij} = 1)$		Bounds on c_{ij} ($\Lambda = 1$ TeV)		Observables	
	Re	Im	${ m Re}$	Im		
$\overline{(ar{s}_L\gamma^\mu d_L)^2}$	$9.8 imes 10^2$	$1.6 imes 10^4$	9.0×10^{-7}	3.4×10^{-9}	$\Delta m_K; \epsilon_K$	
$(\bar{s}_R d_L)(\bar{s}_L d_R)$	$1.8 imes 10^4$	$3.2 imes 10^5$	6.9×10^{-9}	2.6×10^{-11}	$\Delta m_K; \epsilon_K$	
$(ar{c}_L \gamma^\mu u_L)^2$	1.2×10^{3}	$2.9 imes 10^3$	5.6×10^{-7}	1.0×10^{-7}	$\Delta m_D; q/p , \phi_D$	
$(\bar{c}_R u_L)(\bar{c}_L u_R)$	6.2×10^3	$1.5 imes 10^4$	5.7×10^{-8}	1.1×10^{-8}	$\Delta m_D; q/p , \phi_D$	
$(\overline{b}_L \gamma^\mu d_L)^2$	$5.1 imes 10^2$	$9.3 imes 10^2$	3.3×10^{-6}	1.0×10^{-6}	$\Delta m_{B_d}; S_{\psi K_S}$	
$(\overline{b}_R d_L) (\overline{b}_L d_R)$	1.9×10^3	$3.6 imes 10^3$	5.6×10^{-7}	1.7×10^{-7}	$\Delta m_{B_d}; S_{\psi K_S}$	
$\overline{(ar{b}_L \gamma^\mu s_L)^2}$	1.1×10^{2}		7.6×10^{-5}		Δm_{B_s}	
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Very strong suppression! New flavor violation must either approximately (exactly?) follow SM pattern...

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Very strong suppression! New flavor violation must either approximately (exactly?) follow SM pattern...

... or exist only at very high scales ($10^2 - 10^5 \text{ TeV}$)

$$Y_D \approx \operatorname{diag} \left(\begin{array}{ccc} 2 \cdot 10^{-5} & 0.0005 & 0.02 \end{array} \right)$$
$$Y_U \approx \left(\begin{array}{ccc} 6 \cdot 10^{-6} & -0.001 & 0.008 + 0.004i \\ 1 \cdot 10^{-6} & 0.004 & -0.04 + 0.001 \\ 8 \cdot 10^{-9} + 2 \cdot 10^{-8}i & 0.0002 & 0.98 \end{array} \right)$$

Log(SM flavor puzzle)

$$-\log|Y_D| \approx \operatorname{diag}(11 \ 8 \ 4)$$
$$-\log|Y_U| \approx \begin{pmatrix} 12 \ 7 \ 5 \\ 14 \ 6 \ 3 \\ 18 \ 9 \ 0 \end{pmatrix}$$

If $Y = e^{-\Delta}$, then the Δ don't look crazy.

anarchic ("structure-less") $Mass_{ij} \propto Y_{ij}e^{-MR(c_i+c_j)}$ split fermions/RS $\propto Y_{ij} \left(rac{\mu_{
m low}}{\mu_{
m high}}
ight)^{\gamma^i + \gamma^j}$ strong dynamics $\propto Y_{ij} \left(\frac{\langle \Phi \rangle}{M_{\rm mess}}\right)^{Q^i - Q^j}$ Froggatt-Nielsen Hierarchy { => hierarchical masses & mixing angles



split fermions/RS

strong dynamics

 $\propto Y_{ij} \left(\frac{\langle \Phi \rangle}{M_{\rm mess}}\right)^{Q^i - Q^j}$ Froggatt-Nielsen

Hierarchy { => hierarchical masses & mixing angles

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Flavorgenesis scale?

 $\frac{1}{\Lambda_{\rm Flavor}^2} (\bar{\psi}\psi)(\bar{\psi}\psi)$

 $y_{ij}\, \bar{\psi}_i \mathcal{H} \psi_j$



E

naturalness









Example: MSSM is MFV before susy breaking. If flavor is generated well above messenger scale, TeV theory flavor trivial (= MFV).

 $S = \int \mathrm{d}^4 x \left(\mathrm{d}^2 \theta \mathrm{d}^2 \bar{\theta} \, \Phi_i^* \exp\left(2g_A T_A^a V_A^a\right) \Phi_i + \left\{ \mathrm{d}^2 \theta \left[\mathcal{W}(\{\Phi_i\}) + \frac{1}{4} W_A^a W_A^a \right] + \mathrm{h.c.} \right\} \right)$



Model independent constraints

Minimal flavor violation

UTfit, Buras et. al, Hurth et al

	Tree	1
Operator	Bound on Λ	Observables
$H^{\dagger}\left(\overline{D}_{R}Y^{d\dagger}Y^{u}Y^{u\dagger}\sigma_{\mu\nu}Q_{L}\right)\left(eF_{\mu\nu}\right)$	$6.1 { m TeV}$	$B \to X_s \gamma, B \to X_s \ell^+ \ell^-$
$\frac{1}{2} (\overline{Q}_L Y^u Y^{u\dagger} \gamma_\mu Q_L)^2$	$5.9~{\rm TeV}$	$\epsilon_K, \Delta m_{B_d}, \Delta m_{B_s}$
$H_D^{\dagger} \left(\overline{D}_R Y^{d\dagger} Y^u Y^{u\dagger} \sigma_{\mu\nu} T^a Q_L \right) \left(g_s G^a_{\mu\nu} \right)$	$3.4 { m TeV}$	$B \to X_s \gamma, B \to X_s \ell^+ \ell^-$
$\left(\overline{Q}_L Y^u Y^{u\dagger} \gamma_\mu Q_L\right) \left(\overline{E}_R \gamma_\mu E_R\right)$	$2.7 { m ~TeV}$	$B \to X_s \ell^+ \ell^-, B_s \to \mu^+ \mu^-$
$i\left(\overline{Q}_{L}Y^{u}Y^{u\dagger}\gamma_{\mu}Q_{L}\right)H_{U}^{\dagger}D_{\mu}H_{U}$	$2.3 { m TeV}$	$B \to X_s \ell^+ \ell^-, B_s \to \mu^+ \mu^-$
$\left(\overline{Q}_L Y^u Y^{u\dagger} \gamma_\mu Q_L\right) \left(\overline{L}_L \gamma_\mu L_L\right)$	$1.7 { m TeV}$	$B \to X_s \ell^+ \ell^-, B_s \to \mu^+ \mu^-$
$\left(\overline{Q}_L Y^u Y^{u\dagger} \gamma_\mu Q_L\right) \left(e D_\mu F_{\mu\nu}\right)$	$1.5 { m TeV}$	$B \to X_s \ell^+ \ell^-$

If I-loop suppressed like in MSSM < TeV !

$$\Lambda_{\rm loop} \approx (\frac{\alpha}{4\pi})^{\frac{1}{2}} \Lambda_{\rm tree} \approx \frac{1}{10} \Lambda_{\rm tree}$$

Alignment vs. MFV



Flavour violating				Λ/Λ_{MFV}			
	dimension six operator	Ex. 1	Ex. 2	Ex. 3	$U(1)^2$	N-A	F
$\mathcal{O}_0 =$	$\frac{1}{2}(\bar{Q}_L X_{LL}^Q Q_L)^2$	ϵ^{-4}	ϵ^{-4}	1	1	ϵ^{-2}	1
$\mathcal{O}_{F1} =$	$H^{\dagger}\left(\bar{D}_{R}X_{LR}^{D\dagger}\sigma_{\mu\nu}Q_{L}\right)F_{\mu\nu}$	$x\epsilon^{-2}$	$x\epsilon^{-3/2}$	$x\epsilon^{-2}$	$x\epsilon$	$x\epsilon^{-2}$	$x\epsilon^{-2}$
$\mathcal{O}_{G1} =$	$H^{\dagger}\left(\bar{D}_{R}X_{LR}^{D\dagger}\sigma_{\mu\nu}T^{a}Q_{L}\right)G_{\mu\nu}^{a}$	$x\epsilon^{-2}$	$x\epsilon^{-3/2}$	$x\epsilon^{-2}$	$x\epsilon$	$x\epsilon^{-2}$	$x\epsilon^{-2}$
$\mathcal{O}_{\ell 1} =$	$(\bar{Q}_L X_{LL}^Q \gamma_\mu Q_L) (\bar{L}_L \gamma_\mu L_L)$	ϵ^{-2}	ϵ^{-2}	1	1	ϵ^{-1}	1
$\mathcal{O}_{\ell 2} =$	$(\bar{Q}_L X^Q_{LL} \gamma_\mu \tau^a Q_L) (\bar{L}_L \gamma_\mu \tau^a L_L)$	ϵ^{-2}	ϵ^{-2}	1	1	ϵ^{-1}	1
$\mathcal{O}_{H1} =$	$(\bar{Q}_L X^Q_{LL} \gamma_\mu Q_L) (H^\dagger i D_\mu H)$	ϵ^{-2}	ϵ^{-2}	1	1	ϵ^{-1}	1
$\mathcal{O}_{q5} =$	$(\bar{Q}_L X^Q_{LL} \gamma_\mu Q_L) (\bar{D}_R \gamma_\mu D_R)$	ϵ^{-2}	ϵ^{-2}	1	1	ϵ^{-1}	1

 $\epsilon = \frac{\text{flavon vev}}{\text{messenger mass}} \ll 1 \qquad \qquad x = (m_t/m_b)^{\frac{1}{2}}$

Here only MFV operators, flavorgenesis scale from first two generations

$\begin{array}{c} Combination \ of \ K-K \ and \ D-D \\ & \ Nir \ 07; \ Blum \ et. \ al \ '09 \\ Can \ not \ simultaneously \ evade \ constraints \ from \ DD \ \& \ KK \end{array}$



no effect in K-K mixing

 $\frac{1}{\Lambda_{\rm MD}^2} (\overline{Q}_{Li}(X_Q)_{ij} \gamma_{\mu} Q_{Lj}) (\overline{Q}_{Li}(X_Q)_{ij} \gamma^{\mu} Q_{Lj})$

Combination of K-K and D-D Can not simultaneously evade constraints from DD & KK K-K mixing ! VCKM $Y_D Y_D^{\dagger}$

no effect in D-D mixing

 $\frac{1}{\Lambda_{\rm ND}^2} (\overline{Q}_{Li}(X_Q)_{ij} \gamma_{\mu} Q_{Lj}) (\overline{Q}_{Li}(X_Q)_{ij} \gamma^{\mu} Q_{Lj})$



 $\frac{1}{\Lambda^2_{VII}} (\overline{Q}_{Li}(X_Q)_{ij} \gamma_\mu Q_{Lj}) (\overline{Q}_{Li}(X_Q)_{ij} \gamma^\mu Q_{Lj})$

A particular class of models: partial compositeness (geometric alignment vs. MFV)

Weak scale is unstable

elementary scalar Higgs



Inspiration by QCD



 ρ, \ldots

mass protected by global symmetry

 $\pi \to \pi + \alpha$

Inspired by QCD



 ho,\ldots

Potential tilted: due to quark masses and gauging of EM $GB \rightarrow pGB$



Fermions get masses by coupling to this new sector

MFV or not MFV?
Old Flavor problem of composite Higgs Higgs as bound state, naively $D_{\mathcal{H}=\langle \bar{\psi}\psi \rangle} \approx 3$ $\frac{1}{\Lambda^{D_{\mathcal{H}}-1}} y_{ij} \, \bar{\psi}_i \mathcal{H}\psi_j + \frac{1}{\Lambda^2} c_{ijkl} \, \bar{\psi}_i \psi_j \bar{\psi}_k \psi_l$

A can not be too large, because want top mass

 $\Lambda = \mathcal{O}(\text{TeV})$

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A can not be too large, because want top mass

A must be very large because this leads to FCNCs

K

 $\Lambda = \mathcal{O}(\text{TeV}) \quad \Lambda > 10^5 \text{ TeV}$

Two ways of giving mass to fermions...

Bi-linear (like SM):

 $\mathcal{L} = y f_L \mathcal{O}_H f_R, \quad \mathcal{O}_H \sim (1,2)_{\frac{1}{2}}$



Linear:

D.B. Kaplan '91 $\mathcal{L} = yf_L\mathcal{O}_R + y_Rf_R\mathcal{O}_L + m\mathcal{O}_L\mathcal{O}_R, \quad \stackrel{q_i}{\mathcal{O}_R} \sim (3,2)_{\frac{1}{6}H}$

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Quarks & Leptons mix with strong sector

mass \propto compositeness



Partial compositeness $|SM\rangle = \cos\phi|elem.\rangle + \sin\phi|comp.\rangle$ $|heavy\rangle = -\sin\phi|elem.\rangle + \cos\phi|comp.\rangle$ Composites are heavy ($m_{\rho} \approx \text{TeV}$). Light quarks have very little composite admixture.

mixing
$$\sim$$
 mass
strong sector
Higgs&EVVSB
top
resonances
 ρ_{μ}
elementary fields
 u, d, c, s, b, A_{μ}

 $g_*, m_{
ho}$

 $1 \lesssim g_* \lesssim 4\pi$

Kaplan; Contino, Kramer, Son, Sundrum



RGE of the mixing UV III IR



IR

RGE of the mixing UV



IR

RGE of the mixing UV



IR

RGE of the mixing UV

high p_T

Resonance production (option 1)



 $\sim g_*^2 \sin^2 \theta_{u_R}$

strongly suppressed for light quarks!

high pt Resonance production (option 2) U ρ gluon $\sim rac{g_s}{g*}$ U similar to $\gamma - \rho$ mixing

NB, gluon-rho-rho = 0

high pt

t,b

t,b

Resonance decay

ρ

decays dominantly into 3rd generation! (tt, bt, bb)

Top FCNCs

SM $Br(t \rightarrow q(Z, \gamma, G)) \sim 10^{-12}$ partial compositeness/ warped flavor $Br(t \rightarrow c_R Z) \propto |U_R|_{23} \times \delta g_Z \sim 10^{-5}$

LHC (100 1/fb)

 $Br(t \to (Z, \gamma)) \ge 10^{-5}$

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Resonances decay to Tops top jets, road to KK's discovery Agashe et a

Agashe et al, Lillie et al





Collimation poses challenge (m_{KK} ~ 3 TeV vs. m_{top})

high p_T - flavor interplay!

Resonances decay to Tops top jets, road to KK's discovery Agashe et al, Lillie et al

KK gluon t 000000 t w



Collimation poses challenge $(m_{KK} \sim 3 \text{ TeV vs. } m_{top})$

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FCNC protection

Gherghetta, Pomarol; Huber; Agashe, Perez, Soni;



masses from mixing in composites $m_d \sim v \sin \theta_{d_L} Y^* \sin \theta_{d_R}$

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FCNC protection

 d_{r} SdL $f \ll H >$ d_{R} SdR



masses from mixing in composites

 $m_d \sim v \sin \theta_{d_L} Y^* \sin \theta_{d_R}$



 $\bar{K}^{0} - \bar{K}^{0}$

FCNCs suppressed by the same mixings $\sim \frac{g_*^2}{M_\rho^2} s_{dL} s_{dR} s_{sL} s_{sR}$ $\sim \frac{g_*^2}{M_\rho^2} \frac{m_d m_s}{vY^2} \text{ RS-G}$

Little CP problem

Csaki, Falkowski, AW; Buras et al; Casagrande et al $\Delta F = 2 \; (\text{strongest from } \epsilon_K) \; g_* \approx Y_* \approx 3 \dots 6$



$$M_* \gtrsim 10 \left(\frac{g_*}{Y_*}\right) \text{TeV}$$

 $\Delta F = 1$ (strongest constraint from ϵ'/ϵ)

Gedalia et. al



 $\Delta F=0$ neutron EDM

$$M_* \gtrsim 1.3 \, Y_* \, {\rm TeV}$$

$$M_* \ge 2.5 Y_* \,\mathrm{TeV}$$

Agashe et. al, Delaunay et. al, Redi, AW



generate $Y_{U,D}$ at high scale

new physics dynamics can depend non-trivially on $Y_{U,D}$

Flavor triviality: dynamical MFV

Cacciapaglia, Csaki, Galloway, Marandella, Terning, A.W.

strong sector $SU(3)_Q \times SU(3)_u \times SU(3)_d$



mixing can be large & universal

MFV-RS allows for sizable A_{FB} (Small asymmetry in anarchic warped flavo^Bauer et al

plot from Blum et al



MFV-RS allows for sizable A_{FB} (Small asymmetry in anarchic warped flavoBauer et al

plot from Blum et al



measurement at LHCb?

Flavor gauge bosons at LHC

Csaki, Kagan, Lee, Perez, AW

 $g_{\rm eff} G^{(1)\rm KK}_{\mu} \bar{\psi} \psi$

Flavor gauge bosons do not have massless modes (flavor is broken)



no $\gamma - \rho$ mixing ! But quark composite mixing can be flavor universal & large $\sim g_*^2 \sin^2 \theta_{u_R}$

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FGBs at the LHC (preliminary)

The flavor gauge bosons & scalars might be observable.



FGBs at the LHC (preliminary)

The flavor gauge bosons & scalars might be observable.



Flavor Gauge Boson @ Tevatron? $\mathcal{L} = g_{eff} \bar{u}_R V_{\mu}^A \frac{T^A}{2} \gamma_{\mu} u_R + h.c.$

Can partially explain
 A_{FB} with the usual
 constraints:

$$\delta \sigma_{700-800\,{\rm GeV}}^{\rm NP} / \sigma_{700-800\,{\rm GeV}}^{\rm SM} \lesssim 47\%$$

 $\delta \sigma_{t\bar{t}}^{\rm NP} / \sigma_{t\bar{t}}^{\rm SM} \lesssim 10\%$

MFGB <900 GeV, geff ~O(1) $A_{FB}^{t\bar{t}}(M_{inv} > 450 GeV) \stackrel{\checkmark}{\sim} 10\%$

@ σ_{NP}/σ_{SM}(pT>400GeV): 2-3



Conclusions

Most well-motivated models of NP at the TeV predict experimentally resolvable deviations from the SM

Discovery of non-MFV new physics might give insight in origin of Yukawas

high p_T can also offer window into flavor (see explanations of the top FB anomaly)