

Solving μ and little hierarchy problems of the MSSM with electroweak triplets

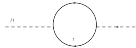
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work in progress with K.Agashe, A.Katz and D.Kim

- SUSY is one of the most motivated scenarios of the Beyond Standard Model Physics.
- SUSY protects Higgs mass from quadratic divergences



A Feynman diagram showing a Higgs boson (represented by a circle) with two external dashed lines. The diagram is labeled with a tilde symbol (~) and the expression $\frac{\Lambda^2 y^2}{16\pi^2}$.

$$\sim \frac{\Lambda^2 y^2}{16\pi^2}$$

In the SM, $\Lambda \sim M_{Pl} \sim 10^{18}\text{GeV}$, so natural scale of the Higgs mass $M_H \sim M_{Pl}$ but in MSSM $\Lambda \sim m_{soft} \sim 100\text{GeV}$

- Other attractive features: dark matter candidate, coupling unification...

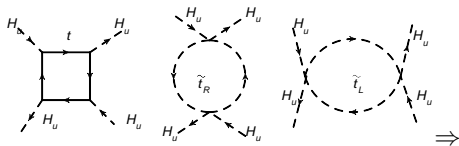
MSSM and little hierarchy problem

- Higgs quartic coupling H^4 is fixed by the gauge interactions, this leads to the

$$m_h < |\cos 2\beta| M_Z \lesssim 91\text{GeV}$$

LEP constraint : $m_h > 115\text{GeV}$

- tree level bound will be modified by the radiative corrections to the quartic interactions:



$$\delta(m_h^2) \sim \frac{3}{4\pi^2} y_t^2 m_t^2 \ln\left(\frac{m_{\tilde{t}_1} m_{\tilde{t}_2}}{m_t^2}\right)$$
$$m_{h^0} \lesssim 135\text{GeV}$$

for all stop masses below 1 TeV

Higgs mass bound in the extensions of the MSSM

- NMSSM $\sim \lambda S H_u H_d \Rightarrow \Delta V \sim \lambda^2 |(H_u H_d)|^2$

$$m_h^2 > M_Z^2 \left(\cos^2 2\beta + \frac{2\lambda^2}{g^2 + g_1^2} \sin^2 2\beta \right)$$

effect of the coupling $\lambda S H_u H_d$ is important only for the small values of $\tan \beta$

- problematic large values of the coupling with perturbativity up to the GUT scale.

Generic Superpotential with trilinear terms

- Let us consider generic lagrangian where N are neutral scalars
(Espinosa, Quiros)

$$W \sim \lambda_{ud}^{(i)} N_{ud} H_u^0 H_d^0 + \lambda_{uu} N_{uu} (H_u^0)^2 + \lambda_{dd} N_{dd} (H_d^0)^2$$
$$m_h^2 \leq M_z^2 \left[\cos^2 2\beta + \frac{2}{g^2 + g'^2} (\lambda_{ud}^2 \sin^2 2\beta + \lambda_{uu}^2 \sin^4 \beta + \lambda_{dd}^2 \cos^4 \beta) \right]$$

- fields that can couple to the H_u^2 in the superpotential will give the correction which will not have $\frac{1}{\tan \beta}$, \Rightarrow
- we need EW triplets

The Model with electroweak triplets (Espinosa, Quiros)



$$W = \mu H_u H_d + \mu_T H_u H_d + \chi_u H_u \bar{T} H_u + \chi_d H_d T H_d$$

- to be consistent with bounds on ρ parameter we need $\langle T \rangle \rightarrow$ small, but F term of the triplets leads to the

$$|F_T|^2 \sim \chi_u \mu_T H_u^2 \bar{T}^*$$

- the tadpole of the triplet, we will have nonzero $\langle T \rangle$ even in the case when $m_T^2 > 0$
- need to suppress vev of the triplet

Supersymmetric vs nonsupersymmetric integration of the triplets

■ Supersymmetric integration of the triplets

(Dine, Seiberg, Thomas; Carena, Kong, Ponton, Zurita)

$$\mu_T \gg m_{T,soft}$$

$$\Delta V \sim \frac{m_{T,soft}^2}{\mu_T^2 + m_{T,soft}^2} \chi_u^2 |H_u^2|^2$$

■ $m_{soft} \gg \mu_T$ limit large correction to the quartic!

$$\Delta V \sim \chi_u^2 |H_u^2|^2$$

EW precision

- bosonic contribution seems under control

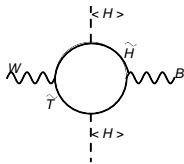
$$\langle T \rangle \sim \frac{\chi_u \mu_T v^2}{m_T^2 + \mu_T^2} \sim 6 \text{ GeV}$$

for $\mu_T = 150$, $m_T = 600$, $\chi_u = 0.5$

- fermionic contribution?

$$\begin{pmatrix} \mu & \sqrt{2}\chi_d v_d \\ -\sqrt{2}\chi_u v_u & \mu_T \end{pmatrix} \begin{pmatrix} H_u^+ \\ T^+ \end{pmatrix},$$

- $S \sim$



Small μ_T term no strong suppression of the S, T parameters

S,T from the tripletino -higgsino mixing

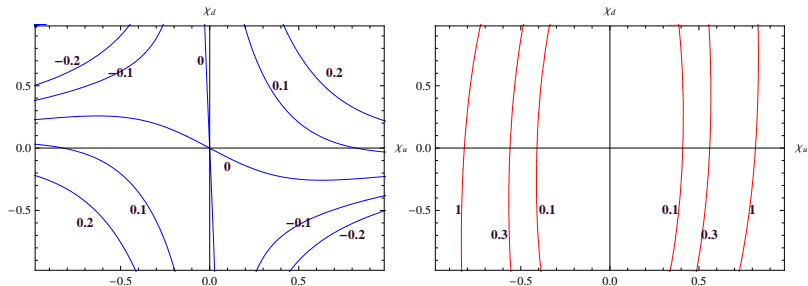


Figure: Contour plots of the regions with different values of the S (blue) and T (red) parameters in the, $\mu = \mu_T = 150\text{GeV}$

minisummary

- The models with triplets can help to solve little hierarchy problem
- Need large soft masses of the triplets and small μ_T term to get viable spectra with suppressed triplet vev.
 - Although vev of the triplet is small it still not enough to get neutrino masses.
- Generically gauge mediation leads to the spectrum where the soft masses of the triplets are larger than the doublets due to the larger Casimir invariant, can be a good candidate for the UV theory
- Need small μ_T for triplets, new μ problem ?

Combining NMSSM with triplets

- one of the easiest solution to the μ problem comes from NMSSM where the μ is vev of the scalar singlet $\langle S \rangle$.
- NMSSM + gauge mediation is problematic, needs some generation of negative $m_{\tilde{G}}^2$, solutions:
 - extra matter that couples to singlets (*Dine, Nelson, Shirman; Agashe, Graesser; de Gouvea, Friedland, Murayama*)
- MSSM+EW triplets we need generation of the μ terms for both doublets and triplets
- can these models solve each others problem?

RGE analysis

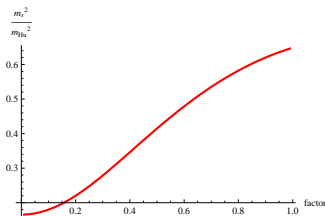
- Using triplets to generate tachyonic m_S^2

- NMSSM

$$\frac{dm_S^2}{dt} \sim 2[\lambda^2(m_{H_u}^2 + m_{H_d}^2 + m_S^2)] + 2[3\kappa^2 m_S^2]$$

- NMSSM+triplets

$$\frac{dm_S^2}{dt} \sim 2[\lambda^2(m_{H_u}^2 + m_{H_d}^2 + m_S^2)] + 2[3\kappa^2 m_S^2] + 3[\lambda_T^2(m_T^2 + m_{\bar{T}}^2 + m_S^2)] \Rightarrow$$



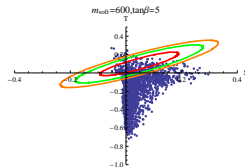
Tension in the parameter space

- EW fit prefers smaller values of the triplet coupling λ_T

$$\langle T \rangle \sim \frac{\chi_u \mu_T v^2}{m_T^2 + \mu_T^2}, \quad \mu_T \sim \lambda_T \langle S \rangle$$

- To get large contribution to the singlet mass squared we need large values of λ_T

$$\frac{dm_S^2}{dt} \sim 3[\lambda_T^2(m_T^2 + m_{\bar{T}}^2 + m_S^2)]$$



To verify whether we can satisfy both requirements we checked whether points passing EW constraints can have zero soft mass for the singlet in the UV, keeping all the couplings perturbative.

R axion

- In the limit when all the A terms vanish the model like usual MSSM is invariant under

$$H_u, H_d, T, \bar{T}, S \rightarrow e^{i\phi_R} H_u, H_d, T, \bar{T}, S.$$

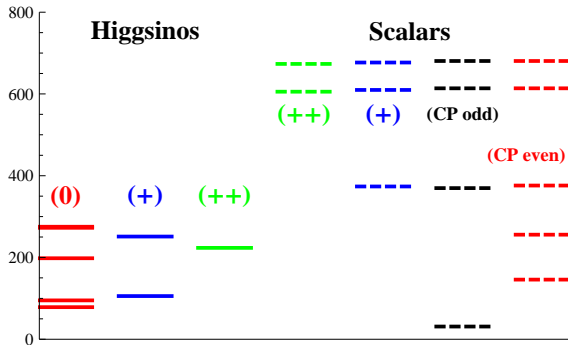
- In the limit when $\langle T(\bar{T}) \rangle \rightarrow 0$ we recover the usual NMSSM expression for the R axion in the terms of the field components

$$R_{axion} \approx \frac{v_s S_I + 2v c_\beta s_\beta (c_\beta H_{uI} + s_\beta H_{dI})}{\sqrt{v_s^2 + v^2 \sin^2 2\beta}}$$

- mass of the axion will be proportional to the A terms

$$m_{ax}^2 \sim -3A_S v_s - \frac{14A v_u v_d}{v_s} + \dots$$

■ Typical spectrum of the Higgs-Higgsino sector



Summary

- NMSSM + triplets is a viable model
 - electroweak triplets help to solve little hierarchy problem and also help to drive the mass of the singlet to be negative
 - vev of the singlet field is used to explain both μ terms of the doublet and triplet Higgs fields
- Model has an interesting spectrum of the new physics excitations, with higgsinos and tripletinos being the lightest states and the scalar triplets including doubly charged particles in the $O(600)$ GeV range.
- Questions:
 - phenomenology of light higgsino-tripletinos?
 - realistic UV completion?