

Stability of the dark matter from non-abelian discrete flavor symmetry

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Hirsch, Morisi, Peinado, Valle, ***PRD 82 (2010)***

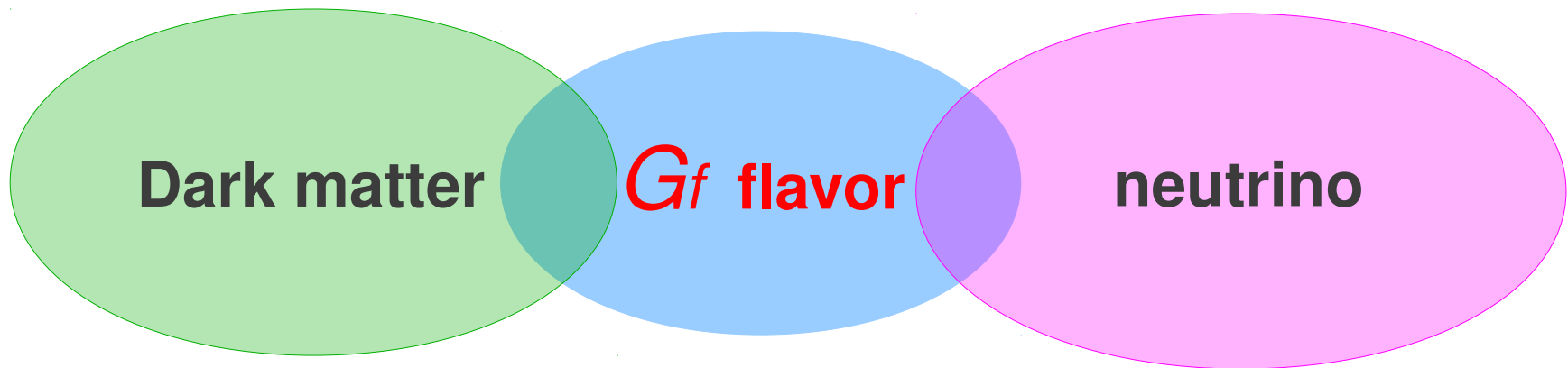
Meloni, Morisi, Peinado, ***PLB697 (2011)***

Boucenna, Hirsch, Morisi, Peinado, Taoso, Valle, ***JHEP1105 (2011)***

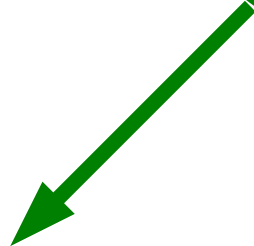
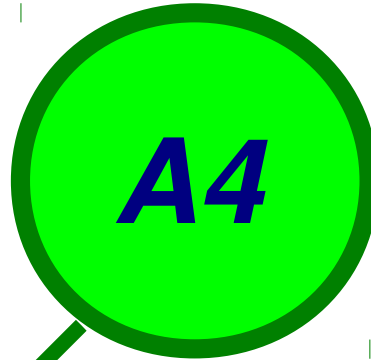
Meloni, Morisi, Peinado *1104.0178*

Adelhart, Bazzocchi, Morisi *1104.5676*

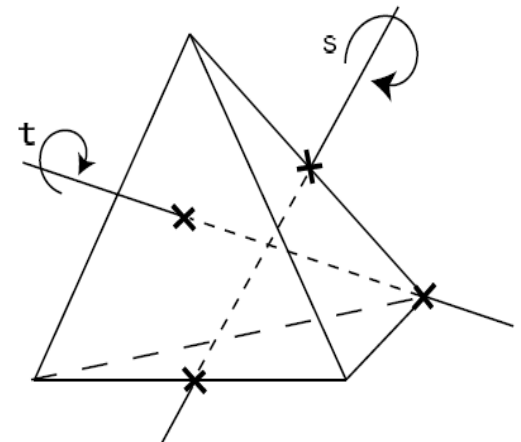
Introduction



The flavor symmetry



Discrete group
even permutations
four objects



tetrahedron

Z2 , Z3 subgroups

Stability of DM from A4

$$A_4 \xrightarrow{\text{spontaneously}} Z_2$$

triplet of A4

$$\langle \eta \rangle \sim (1, 0, 0)$$

$$S\Psi = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix} = \begin{pmatrix} a_1 \\ -a_2 \\ -a_3 \end{pmatrix}$$



$S^2 = 1$ generator of Z_2 subgroup of A_4

The Model

Hirsch, Morisi, Peinado, Valle, *PRD 82 (2010)*

	L_e	L_μ	L_τ	l_e^c	l_μ^c	l_τ^c	N_T	N_4	H	η
$SU(2)$	2	2	2	1	1	1	1	1	2	2
A_4	1	$1'$	$1''$	1	$1''$	$1'$	3	1	1	3

**Quarks are
singlets of A_4**

4 R-nu **4 Higgs**

$$\eta = (\eta_1, \eta_2, \eta_3)$$

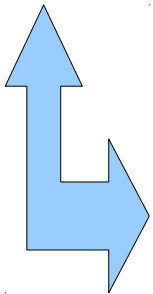
The lagrangian

$$\mathcal{L} = y_e L_e l_e^c \hat{H} + y_\mu L_\mu l_\mu^c \hat{H} + y_\tau L_\tau l_\tau^c \hat{H} +$$

3 **3**

↓ ↓

$$+ y_1^\nu L_e (N_T \eta)_1 + y_2^\nu L_\mu (N_T \eta)_{1''} + y_3^\nu L_\tau (N_T \eta)_{1'} +$$
$$+ y_4^\nu L_e N_4 \hat{H} + M_1 N_T N_T + M_2 N_4 N_4 + \text{h.c.}$$



η has Dirac couplings only through heavy right-handed neutrino

The Model: the DM candidate

Z_2 :

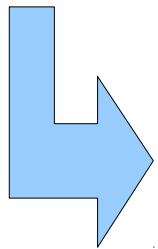
$$\begin{aligned} N_1 &\rightarrow +N_1, & \eta_1 &\rightarrow +\eta_1, \\ N_2 &\rightarrow -N_2, & \eta_2 &\rightarrow -\eta_2, \\ N_3 &\rightarrow -N_3, & \eta_3 &\rightarrow -\eta_3, \end{aligned}$$

$$\eta_2 = \begin{pmatrix} H_2'^+ \\ (H_2' + iA_2')/\sqrt{2} \end{pmatrix}, \quad \eta_3 = \begin{pmatrix} H_3'^+ \\ (H_3' + iA_3')/\sqrt{2} \end{pmatrix}$$

the lightest is the DM candidate

Neutrino phenomenology

$$m_D = \begin{pmatrix} x_1 & 0 & 0 & x_4 \\ x_2 & 0 & 0 & 0 \\ x_3 & 0 & 0 & 0 \end{pmatrix}, \quad M_R = \begin{pmatrix} M_1 & 0 & 0 & 0 \\ 0 & M_1 & 0 & 0 \\ 0 & 0 & M_1 & 0 \\ 0 & 0 & 0 & M_2 \end{pmatrix}$$



$$m_\nu = -m_{D_{3 \times 4}} M_{R_{4 \times 4}}^{-1} m_{D_{3 \times 4}}^T \equiv \begin{pmatrix} y^2 & ab & ac \\ ab & b^2 & bc \\ ac & bc & c^2 \end{pmatrix}$$

IH: $m_3 = 0$
 $0.03 \text{ eV} < 0\nu_{bb} < 0.05 \text{ eV}$

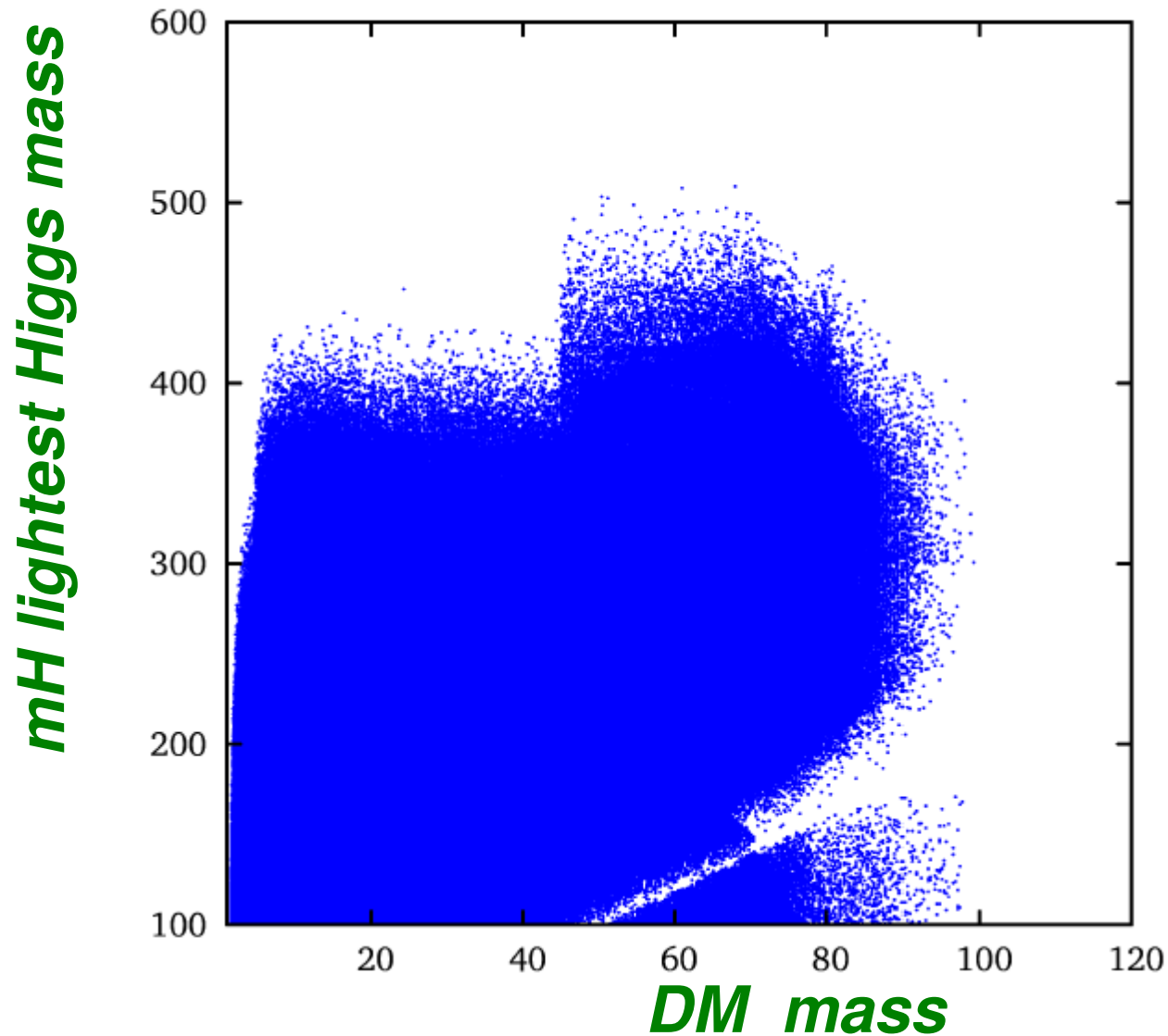
$$V_3 \sim \begin{pmatrix} 0 \\ -b/c \\ 1 \end{pmatrix}$$

$$\theta_{13} = 0$$

Relic density

$$0.09 \leq \Omega h^2 \leq 0.13$$

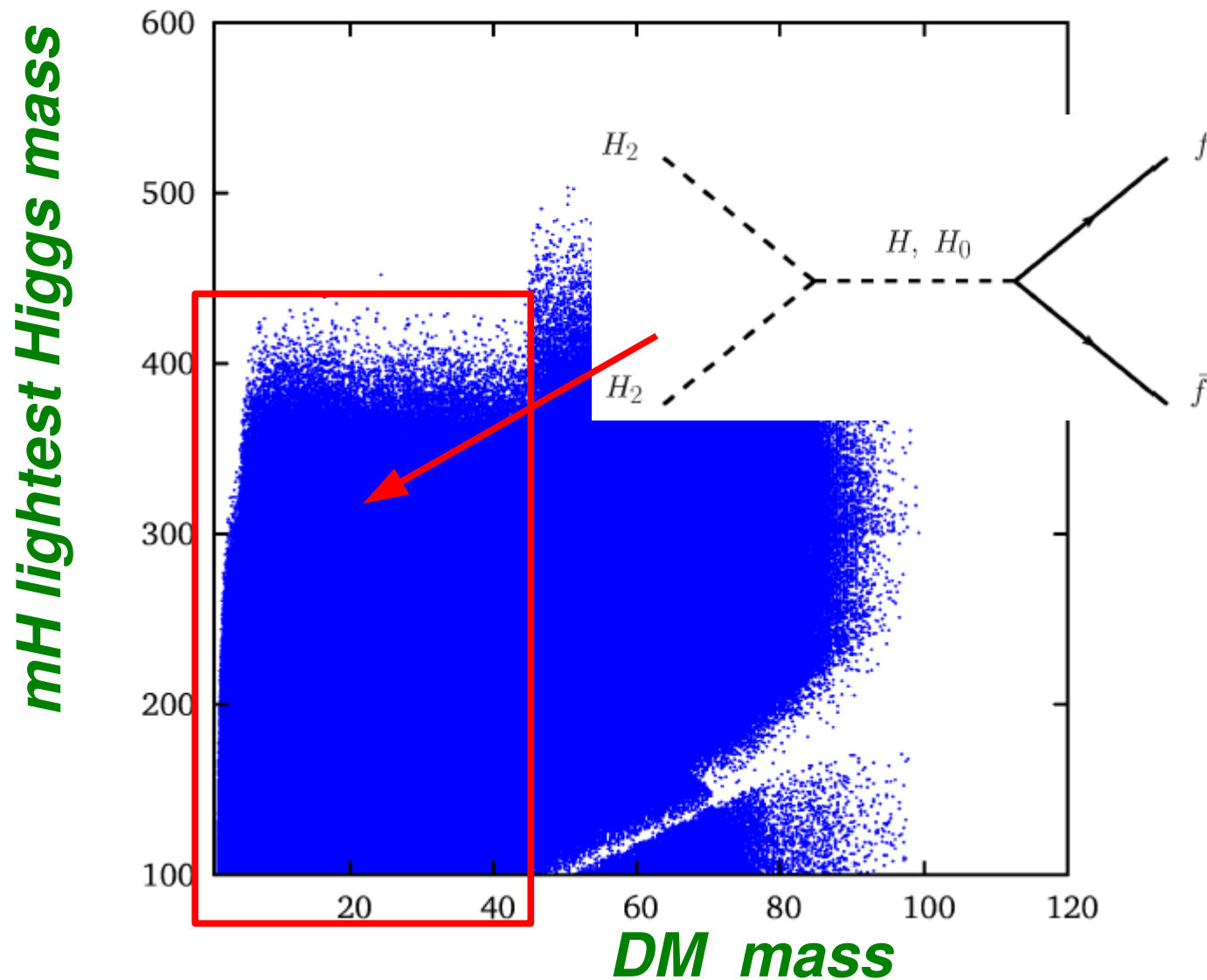
Boucenna *et al.*, *JHEP1105* (2011)



Relic density

$$0.09 \leq \Omega h^2 \leq 0.13$$

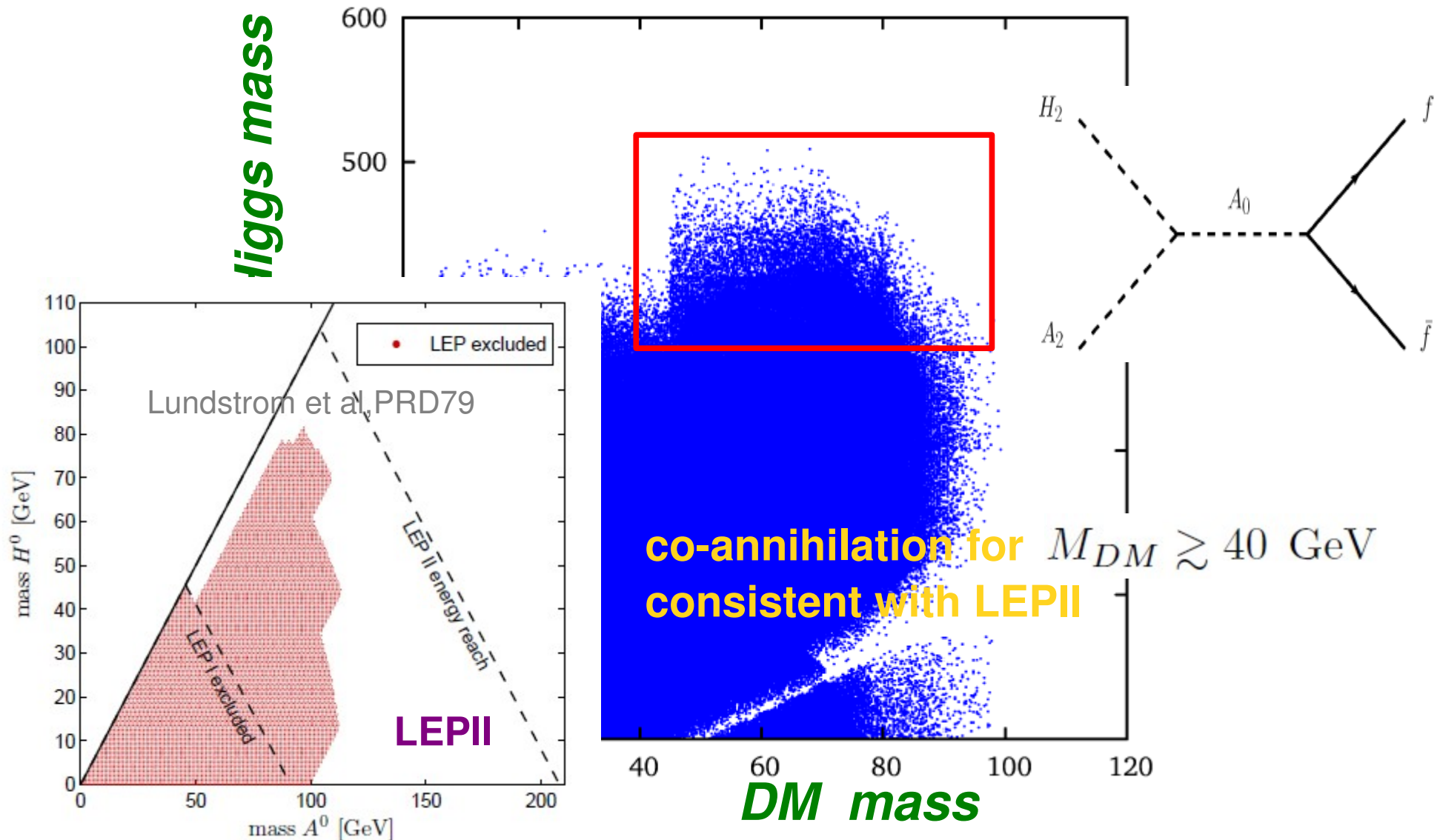
Boucenna *et al.*, *JHEP1105* (2011)



Relic density

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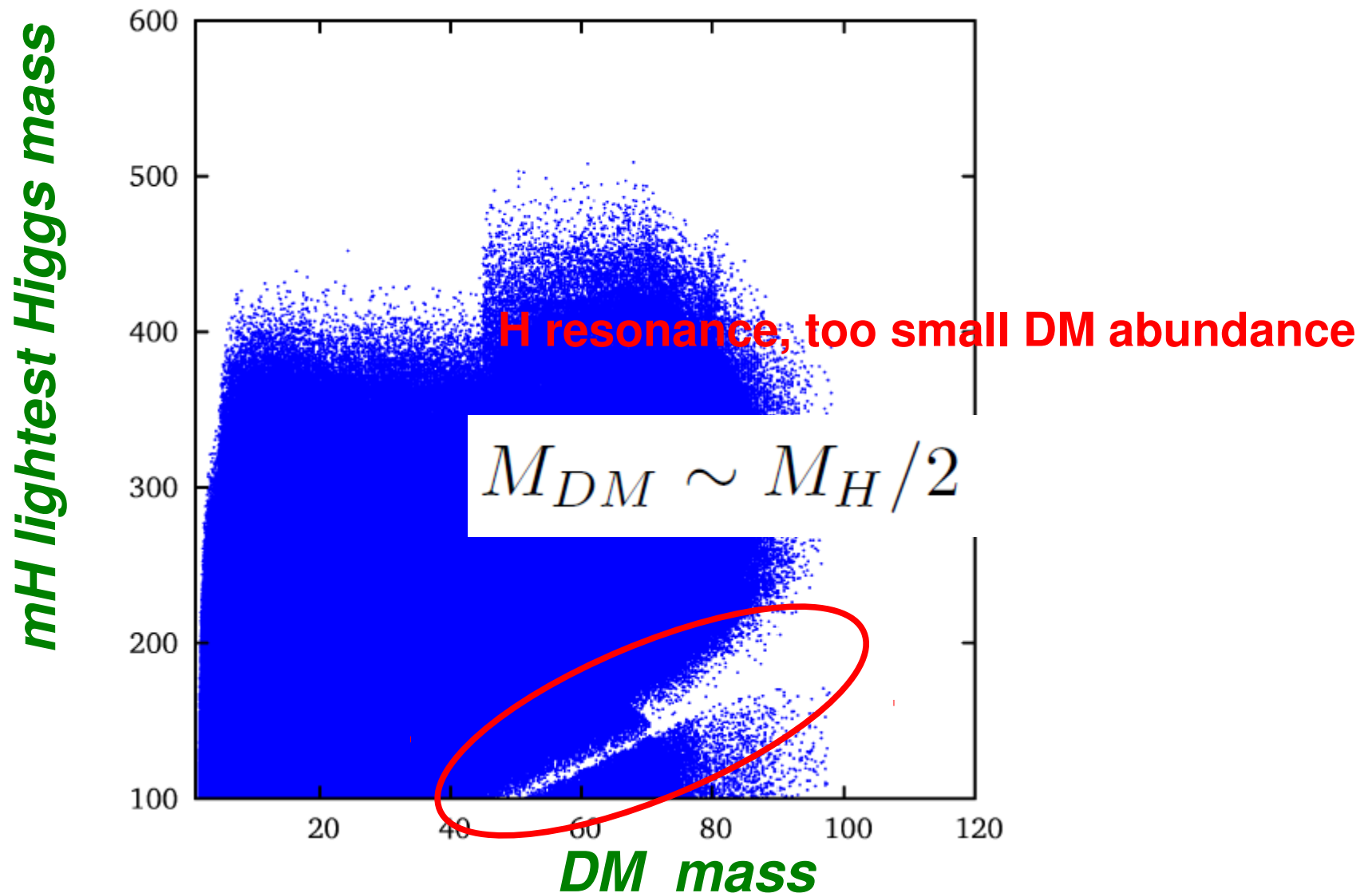
Boucenna *et al.*, *JHEP1105* (2011)



Relic density

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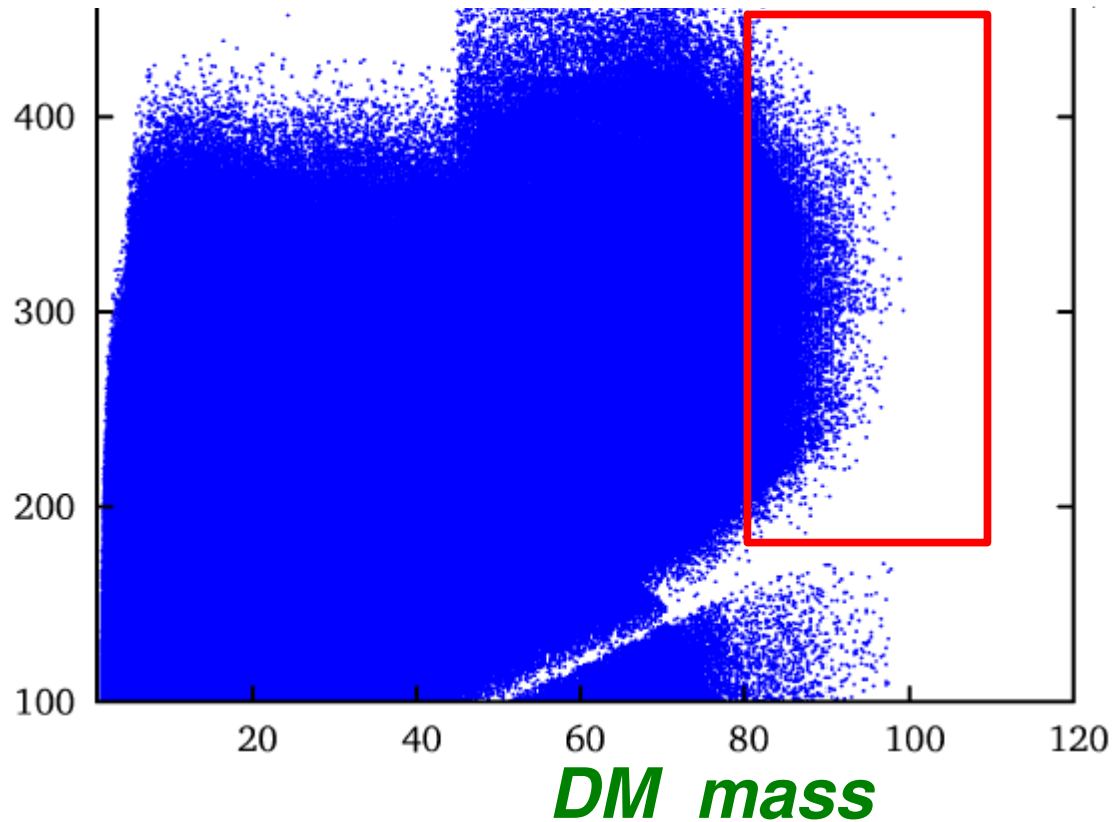
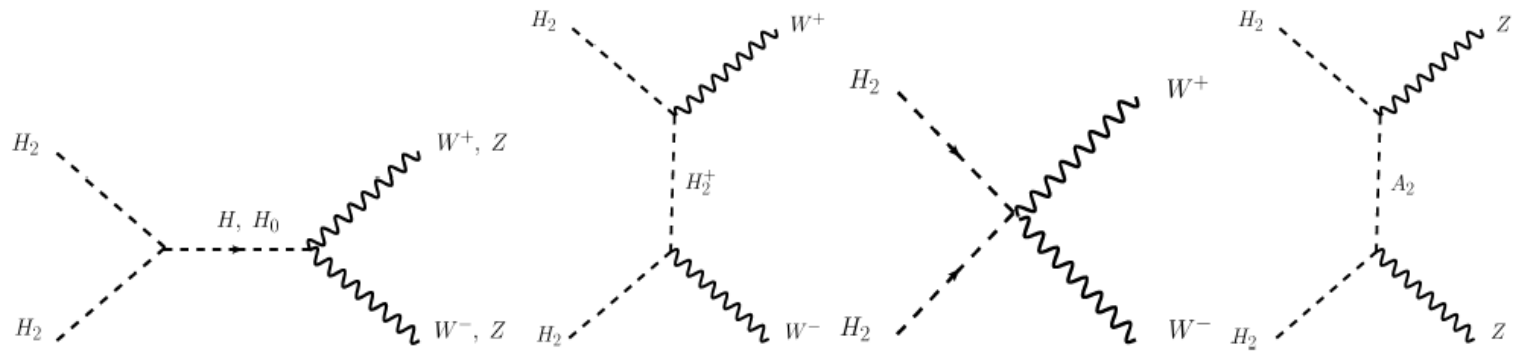
Boucenna *et al.*, *JHEP1105* (2011)



Relic density

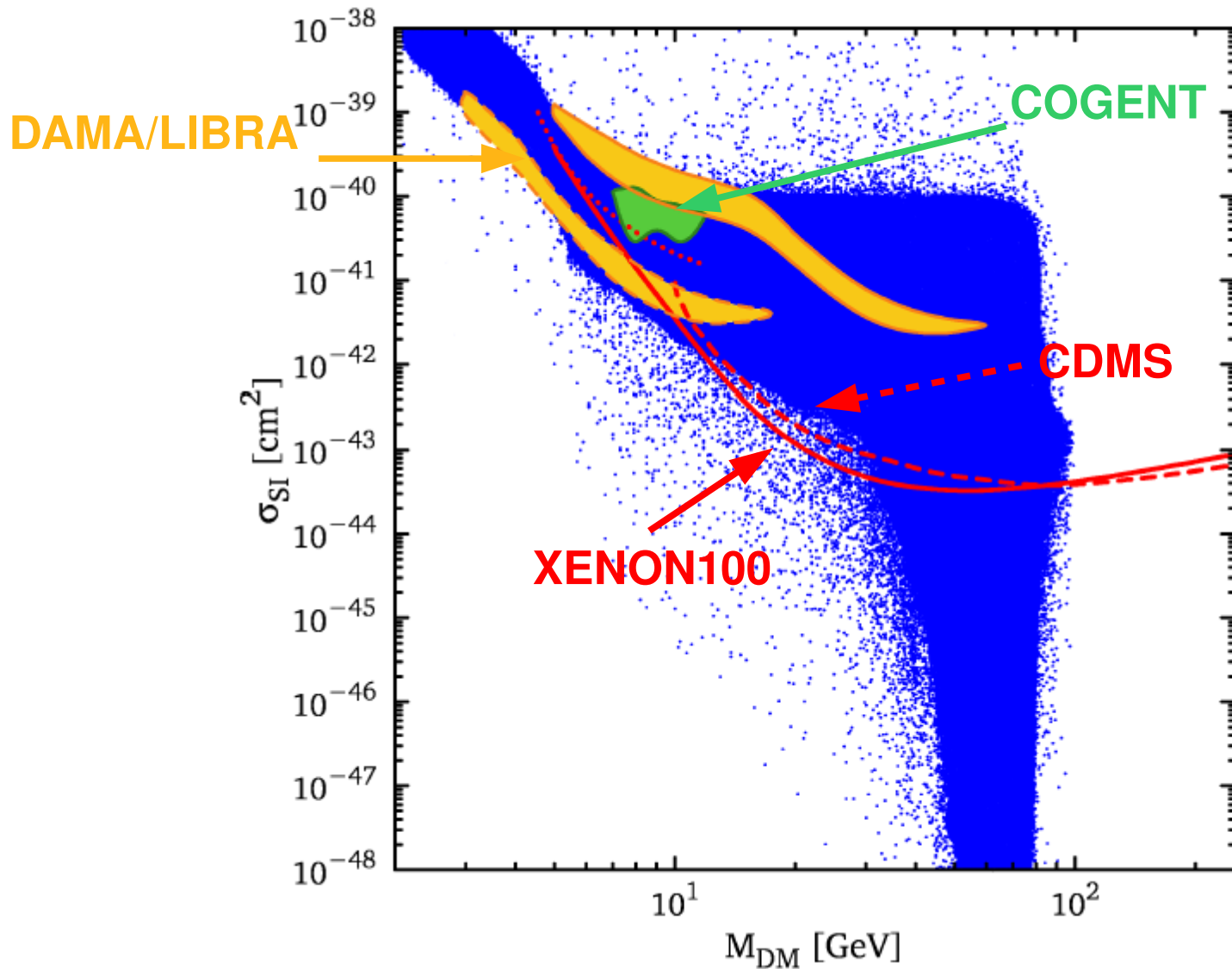
$$0.09 \leq \Omega h^2 \leq 0.13$$

mH lightest Higgs mass



Direct detection

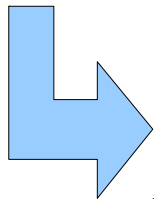
Boucenna *et al.*, *JHEP1105* (2011)



Quarks

Adelhart, Bazzocchi, Morisi 1104.5676

$$\sum \frac{f_{ij}}{\Lambda^2} (\bar{Q}_i \hat{H}) d_j (\eta^\dagger \eta) + \frac{f'_{ij}}{\Lambda^2} (\bar{Q}_i \eta) d_j (\eta^\dagger \hat{H}) + \frac{f''_{ij}}{\Lambda^2} (\bar{Q}_i \eta) d_j (\hat{H}^\dagger \eta)$$

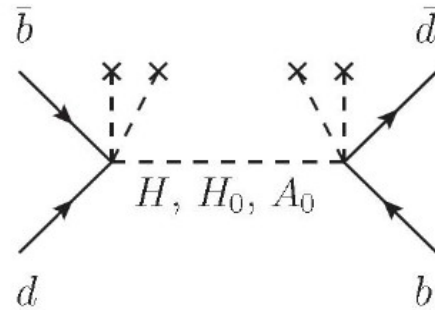
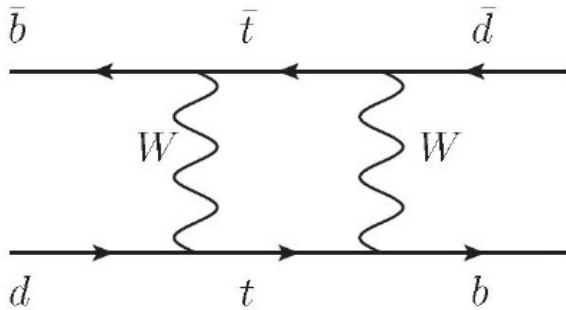


$$M_d = \begin{pmatrix} m_d & 0 & 0 \\ 0 & m_s & 0 \\ 0 & 0 & m_b \end{pmatrix} + \frac{v_H v_\eta^2}{\Lambda^2} \begin{pmatrix} h_{dd} & h_{ds} & h_{db} \\ h_{sd} & h_{ss} & h_{sb} \\ h_{bd} & h_{bs} & h_{bb} \end{pmatrix} + \mathcal{O}(1/\Lambda^4),$$

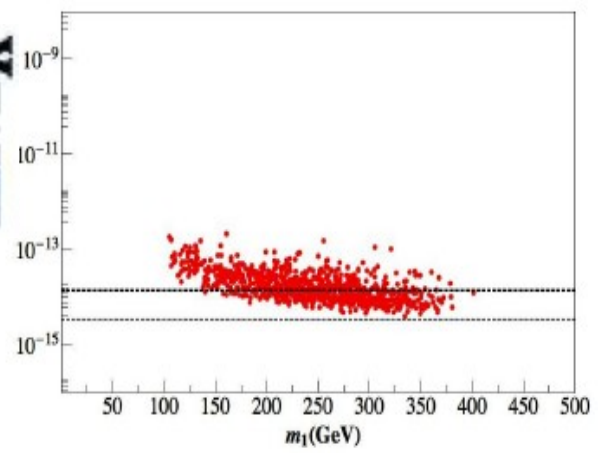
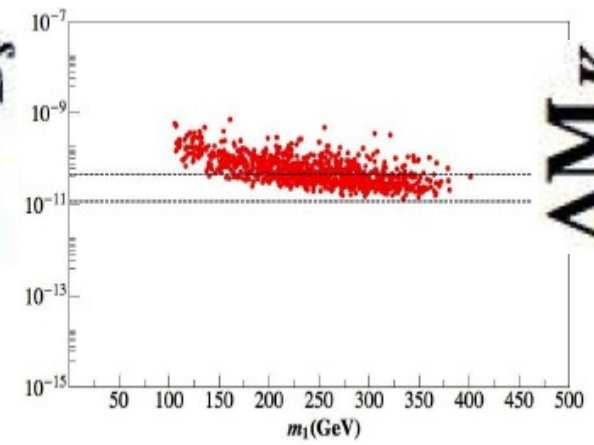
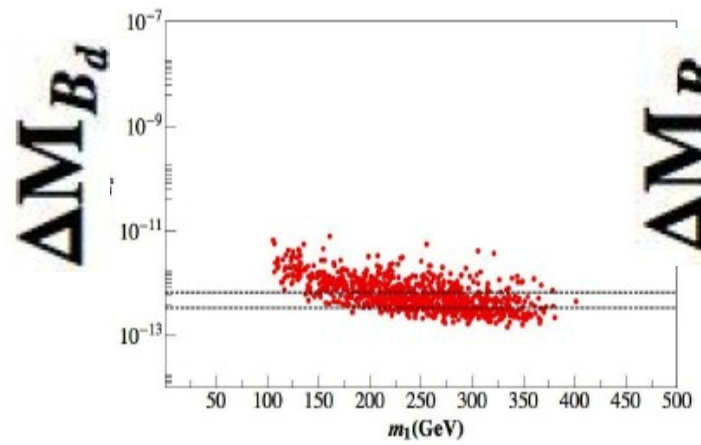
$$h_{ds} \frac{v_H v_\eta^2}{\Lambda^2} = \lambda_C m_s$$

1-10 TeV

Quarks: FCNC



CKM mostly from down sector



lightest Higgs mass



Outlook: embedding into GUT

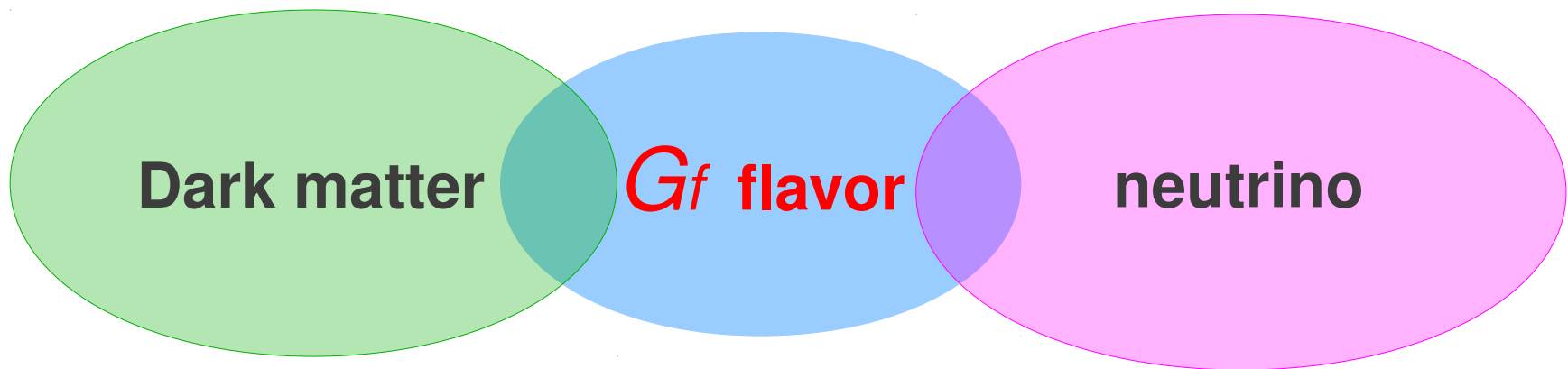
	T_1	T_2	T_3	F_1	F_2	F_3	N_T	N_4
SU(5)	10	10	10	$\bar{5}$	$\bar{5}$	$\bar{5}$	1	1
A_4	1	1'	1''	1	1''	1'	3	1

matter

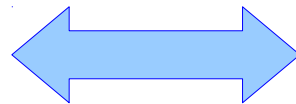
	5_H	$\bar{5}_H$	5_η	45_H
SU(5)	5_H	$\bar{5}_H$	5_η	45_H
A_4	1	1	3	1

scalar

conclusion



stability of DM



In a specific model

reactor angle = 0
inverse hierarchy
 $m_3=0$

Tri-bimaximal neutrino mixing (TBM)

$$\sin^2 \theta_{23} = 0.5$$

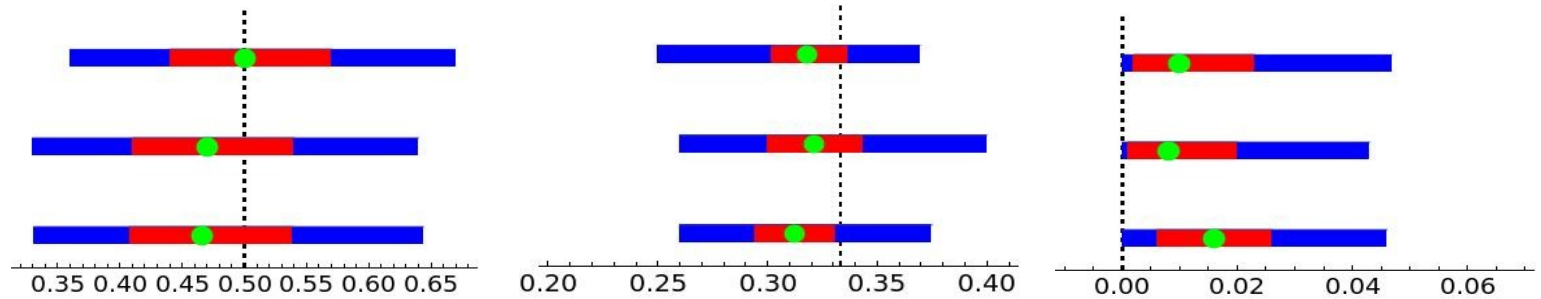
$$\sin^2 \theta_{12} = 1/3$$

$$\sin^2 \theta_{13} = 0$$

Schwetz et al

Gonzalez et al

Fogli et al



$$U_{\text{HPS}} = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -1/\sqrt{6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

Harrison, Perkins & Scott



A4

Babu, Ma, Valle PLB552
Altarelli, Feruglio NPB720

See also Hagedorn, Varzielas, Merlo, Luhn,....

Outlook: embedding into GUT

	T_1	T_2	T_3	F_1	F_2	F_3	N_T	N_4
SU(5)	10	10	10	$\bar{5}$	$\bar{5}$	$\bar{5}$	1	1
A_4	1	1'	1''	1	1''	1'	3	1

matter

	5_H	$\bar{5}_H$	5_η	45_H
SU(5)	5_H	$\bar{5}_H$	5_η	45_H
A_4	1	1	3	1

scalar

$$\begin{aligned} \mathcal{L}_{down} &= y_1^{l,d} T_1 F_1 \bar{5}_H + y_2^{l,d} T_2 F_2 \bar{5}_H + y_3^{l,d} T_3 F_3 \bar{5}_H + y_1^{n,d} T_1 F_1 45_H + y_2^{n,d} T_2 F_2 45_H + y_3^{n,d} T_3 F_3 45_H; \\ \mathcal{L}_{up} &= y_1^u T_1 T_1 5_H + y_2^u T_2 T_3 5_H + y_1'^u T_1 T_1 45_H + y_2'^u T_2 T_3 45_H; \\ \mathcal{L}_\nu &= y_1^\nu T_1 N_4 5_H + y_2^\nu T_2 N_4 5_H + y_3^\nu T_3 N_4 5_H + y_1^\nu T_1 N_T 5_\eta + M_1 N_T N_T + M_2 N_4 N_4. \end{aligned}$$

it is possible to fit the masses

$$m_e = y_1^{l,d} \langle 5_H \rangle - 3y_1^{n,d} \langle 45_H \rangle; \quad m_\mu = y_2^{l,d} \langle 5_H \rangle - 3y_2^{n,d} \langle 45_H \rangle; \quad m_\tau = y_3^{l,d} \langle 5_H \rangle - 3y_3^{n,d} \langle 45_H \rangle;$$

$$m_d = y_1^{l,d} \langle 5_H \rangle + y_1^{n,d} \langle 45_H \rangle; \quad m_s = y_2^{l,d} \langle 5_H \rangle + y_2^{n,d} \langle 45_H \rangle; \quad m_b = y_3^{l,d} \langle 5_H \rangle + y_3^{n,d} \langle 45_H \rangle;$$

$$m_u = y_1^u \langle 5_H \rangle; \quad m_c = y_2^u \langle 5_H \rangle - y_2'^u \langle 45_H \rangle; \quad m_t = y_2^u \langle 5_H \rangle + y_2'^u \langle 45_H \rangle.$$

CKM?