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Internal magnetic fields of white dwarfs: a solution to the cooling anomaly?

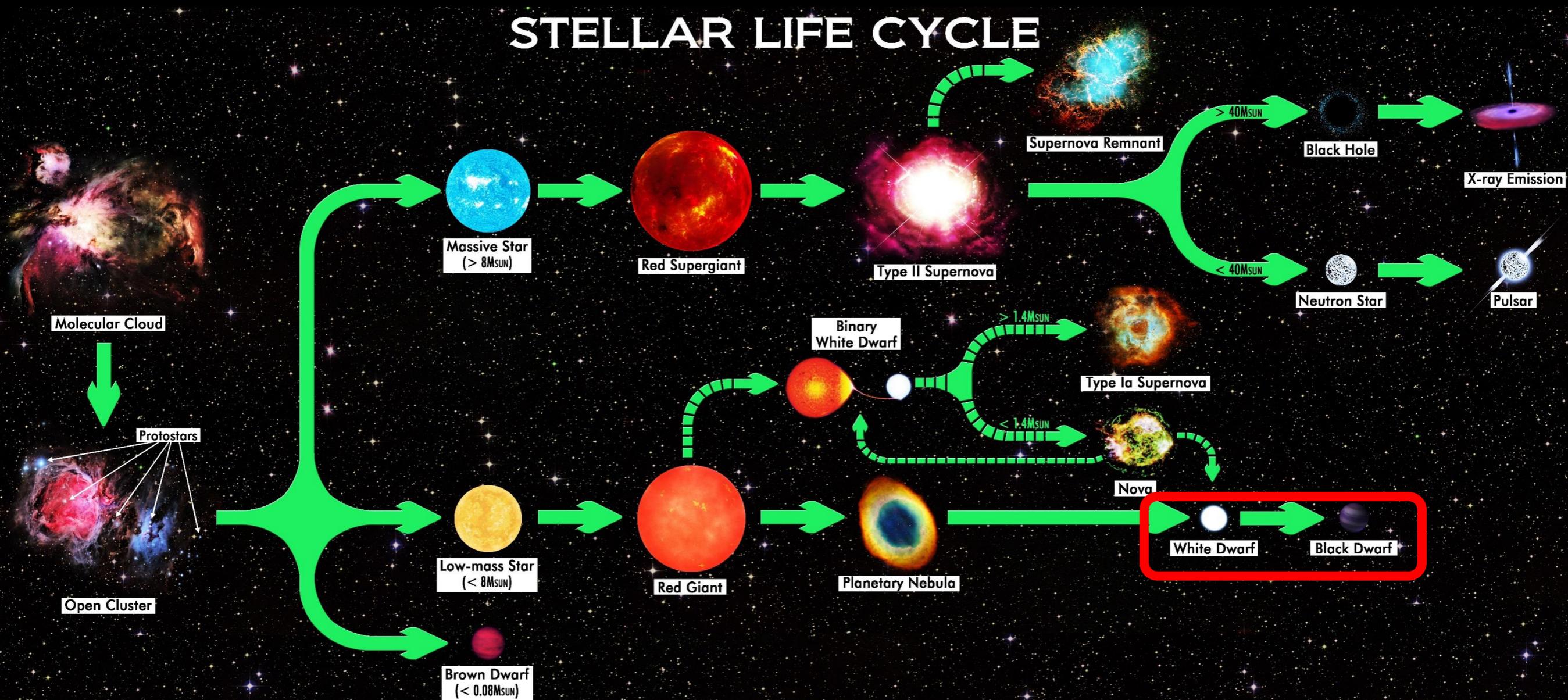
01. 06. 2022

LLP11 Workshop

Based on [arXiv:2109.06158](https://arxiv.org/abs/2109.06158) [astro-ph.SR]

A Star's Life

STELLAR LIFE CYCLE



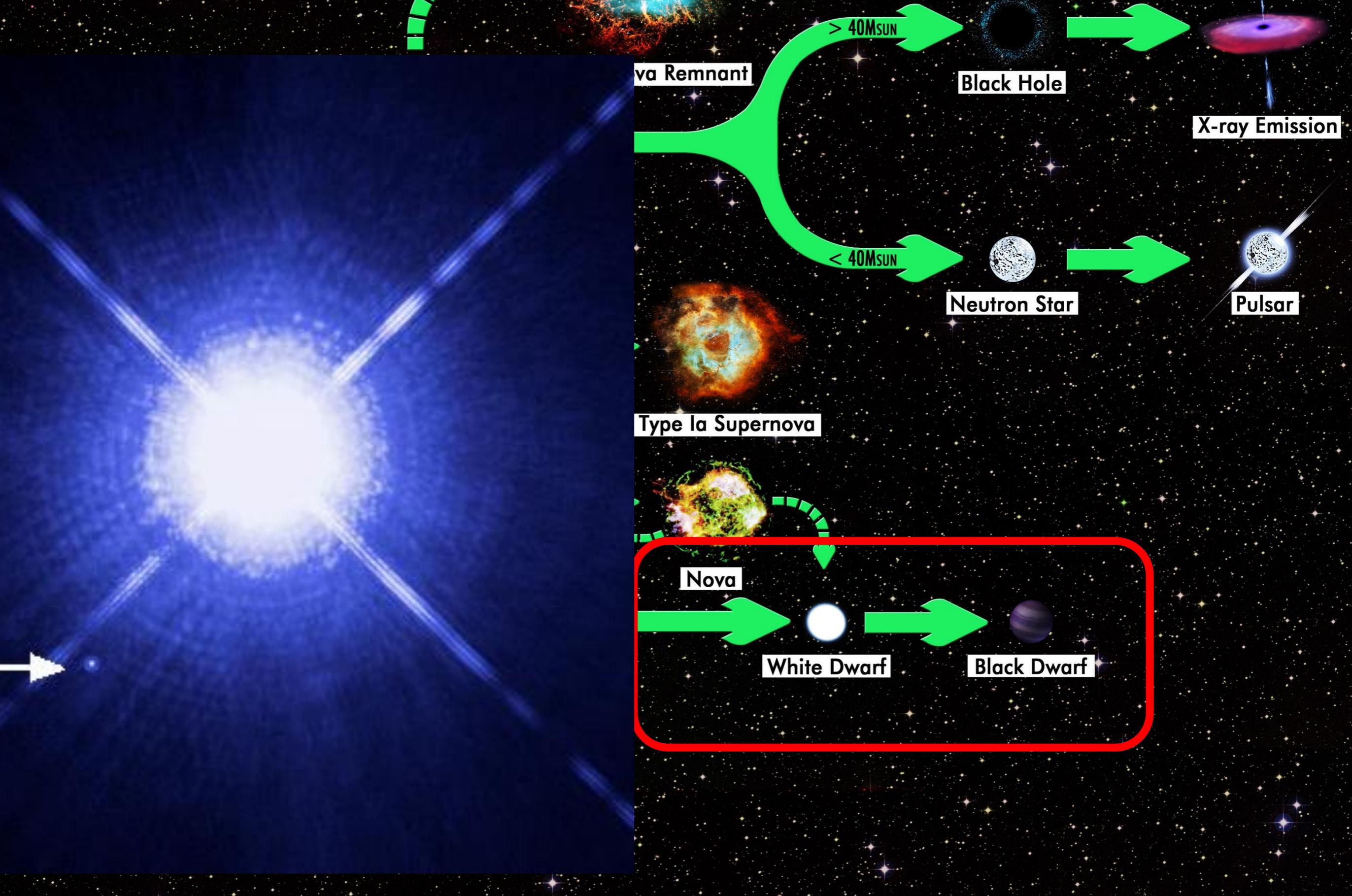
Birth

Main Sequence

Old Age

Death

Remnant



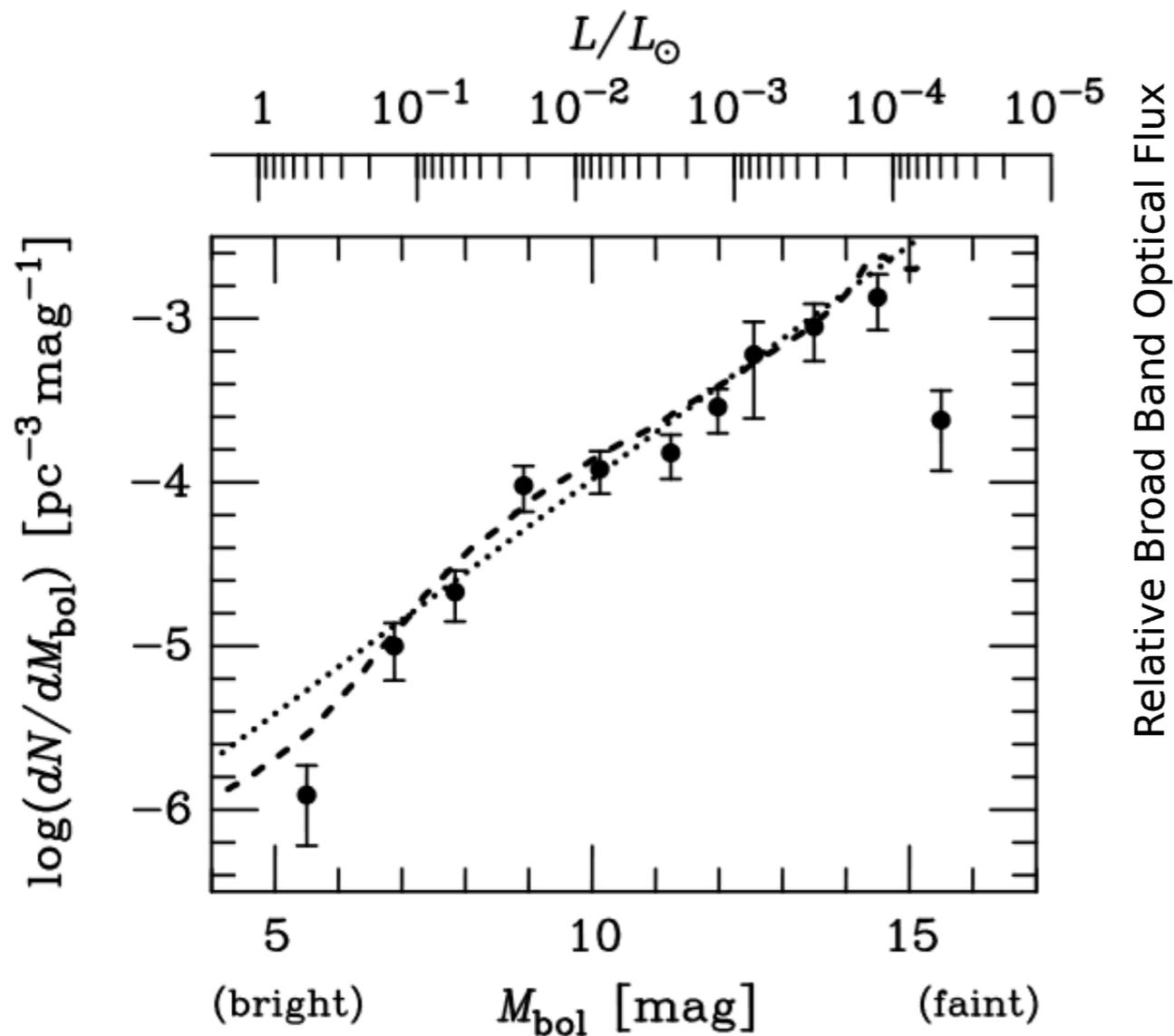
Age

Death

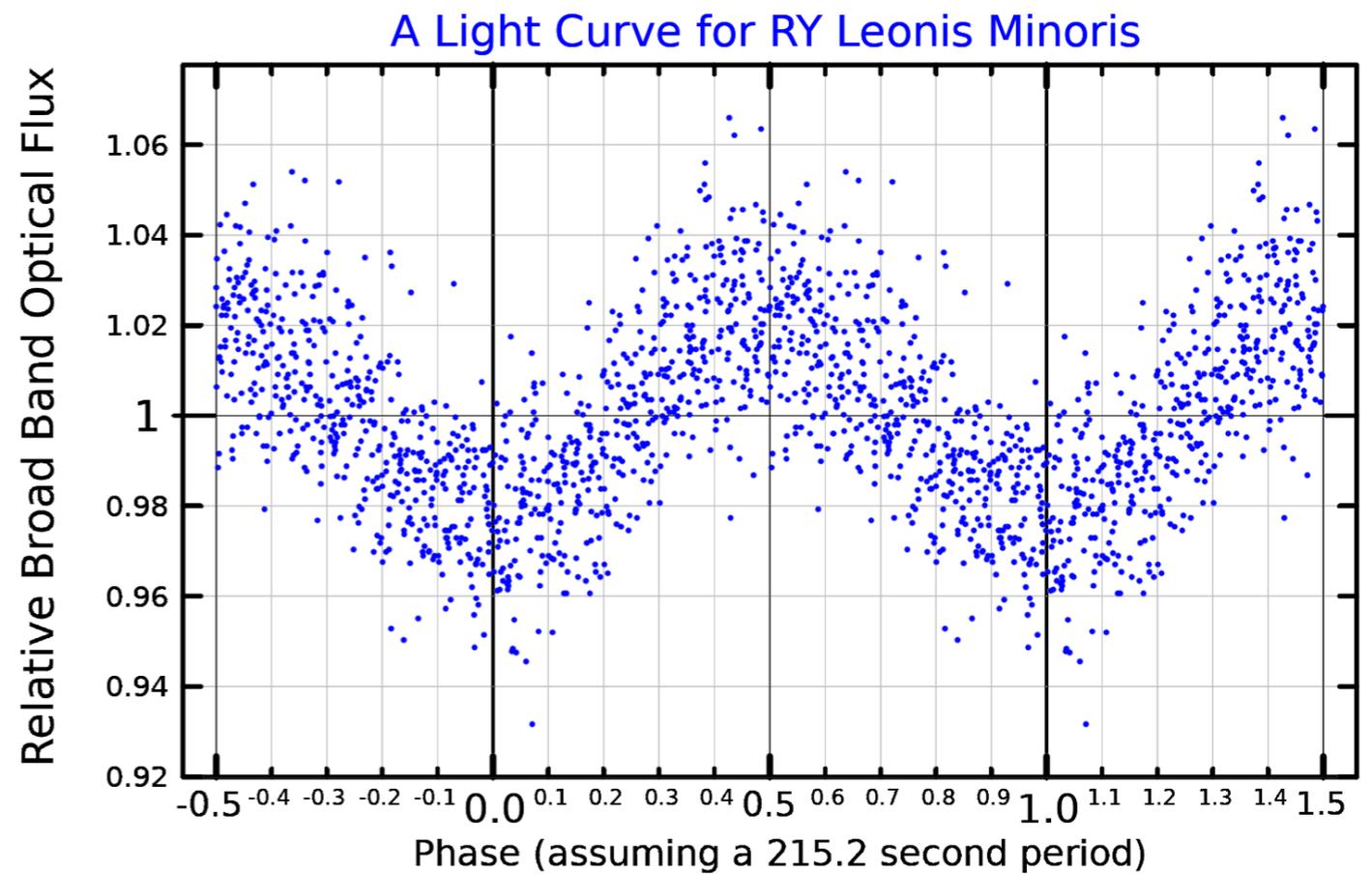
Remnant

Observing WD Cooling

White Dwarf Luminosity Function
(WDLF):
Population



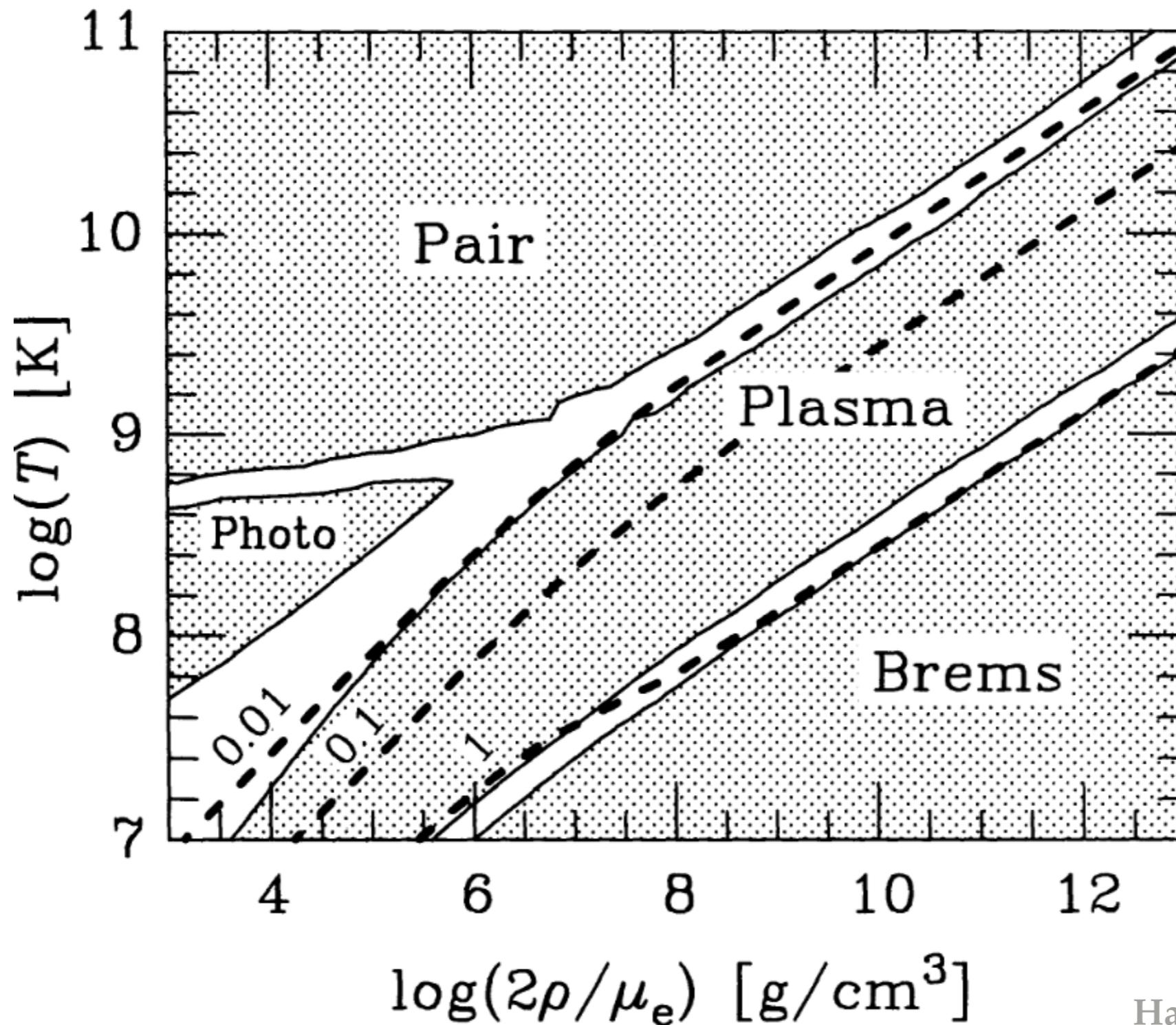
Pulsating WDs:
Individual stars



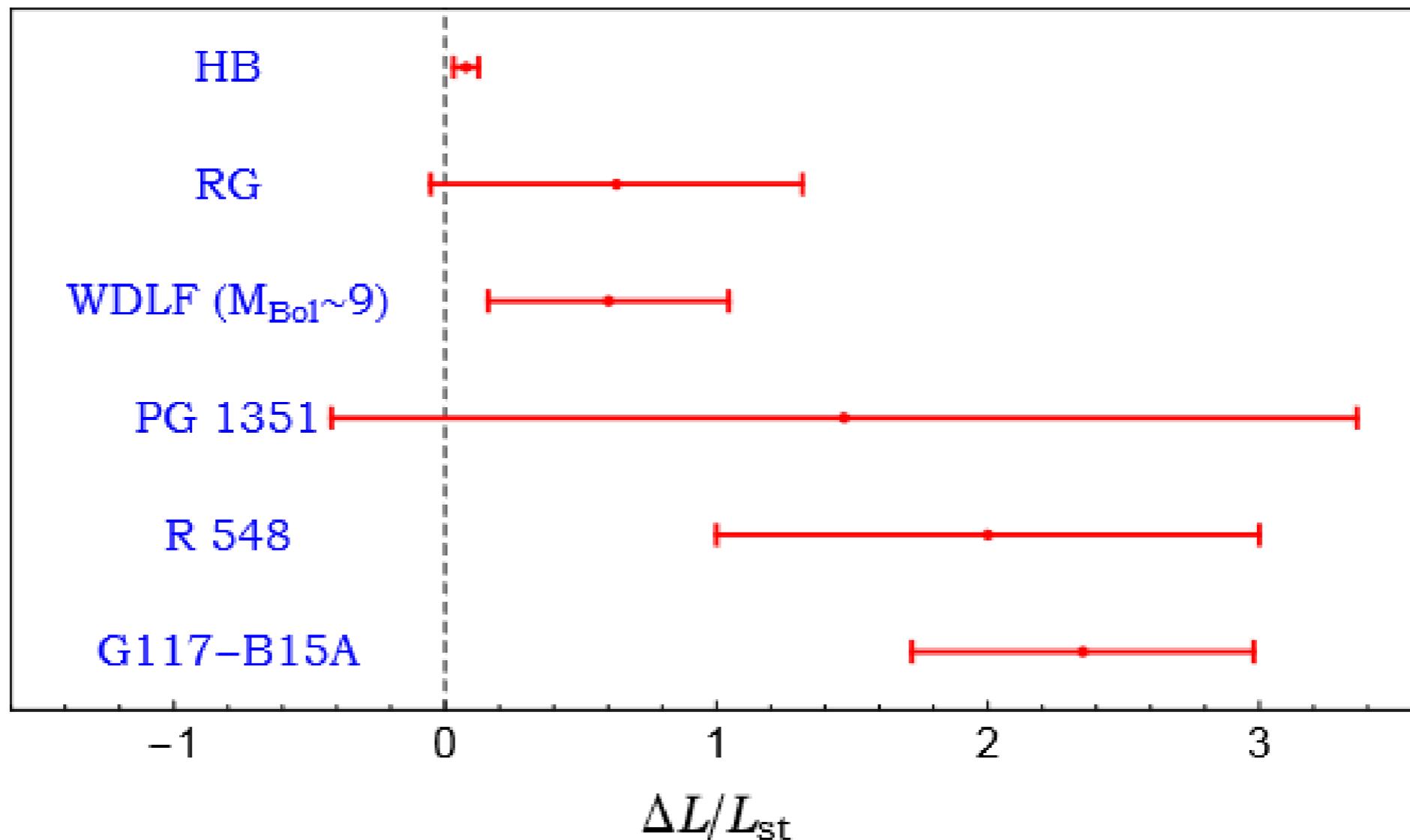
G117-B15A

$$T = 1.2 \times 10^7 \text{ K} \quad Y_e \rho \simeq 10^6 \text{ g/cm}^3$$

Theory of WD Cooling



Cooling Anomaly

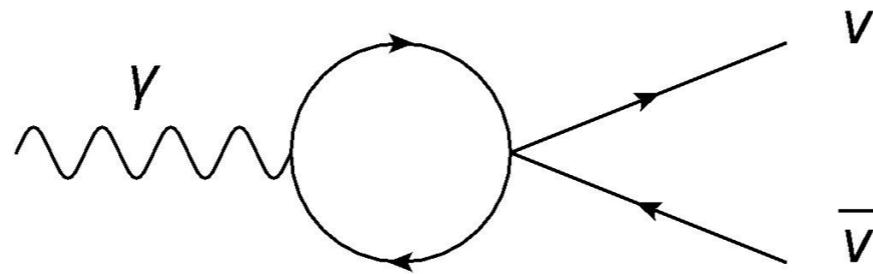


- Some WDs appear to be cooling too fast....
- Do they emit LLPs (axions, ALPs, ...)?

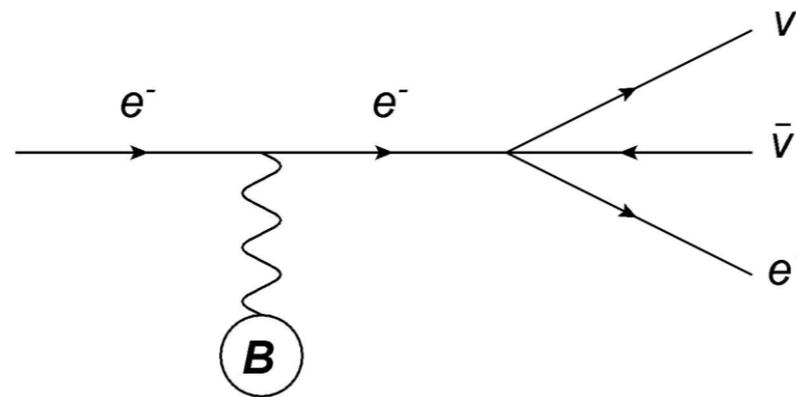
Impact of B-fields

Can internal magnetic fields explain this within the SM?

- Modify plasma processes $\gamma \rightarrow \nu\nu$



- Enable synchrotron radiation $e \rightarrow e\nu\nu$



- Heating through Ohmic decay

Plasma Processes

- $\gamma \rightarrow \nu\nu$ possible in dense medium due to modified photon and plasmon dispersion relations, roughly characterised by the plasma frequency

$$\omega_p^2 = \frac{4\pi\alpha n_e}{m_e} \left[1 + \frac{1}{m_e^2} (3\pi^2 n_e)^{2/3} \right]^{-1/2} \simeq \left(20 \text{ keV } \rho_6^{1/2} \right)^2 \quad \text{cf. e.g. Braaten/Segel [9302213](#)}$$

- Refractive index (“thermal mass”) is determined by electron density, relevant scale is the frequency

cf. e.g. Kennet/Melrose [astro-ph/9901156](#)

$$\omega_B = \frac{eB}{m_e} = m_e \frac{B}{B_c} \simeq 11.5 B_{12} \text{ keV} \quad B_c = m_e^2/e = 4.41 \times 10^{13} \text{ G}$$

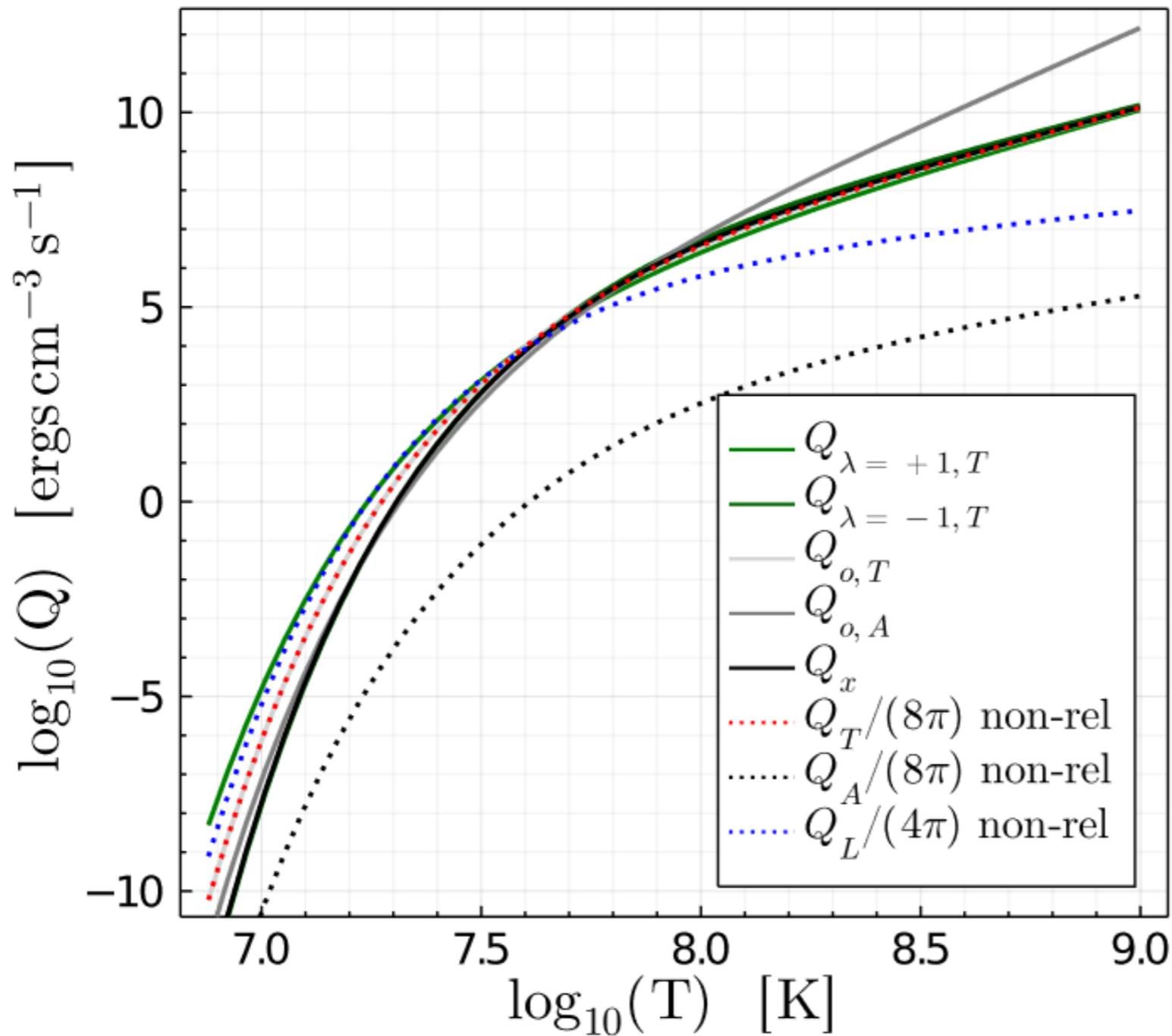
- Magnetic fields force electrons on Landau levels, modify refractive index

$$E_\nu = \sqrt{p_{\parallel}^2 + m_e^2 \left(1 + 2\nu \frac{B}{B_c} \right)}$$

- Other effects (Schwinger-like pair creation, modification of wave function...) are negligible or sub-dominant

Plasma Processes

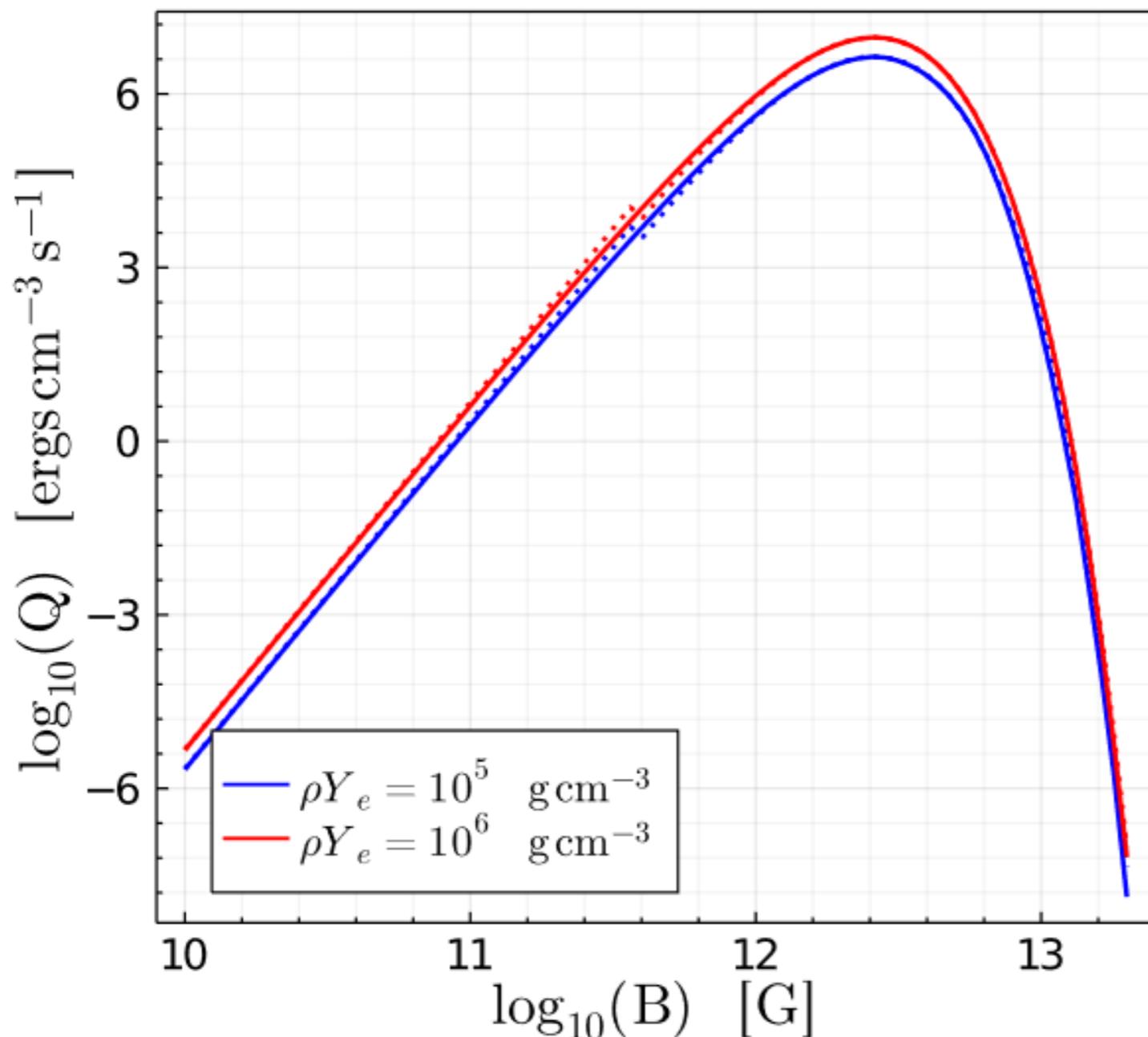
$$\rho Y_e = 10^6 \text{ g cm}^{-3}$$



- For typical WD parameters, impact of B fields significant...
- ...but only at temperature where other processes are more important

Synchrotron Radiation

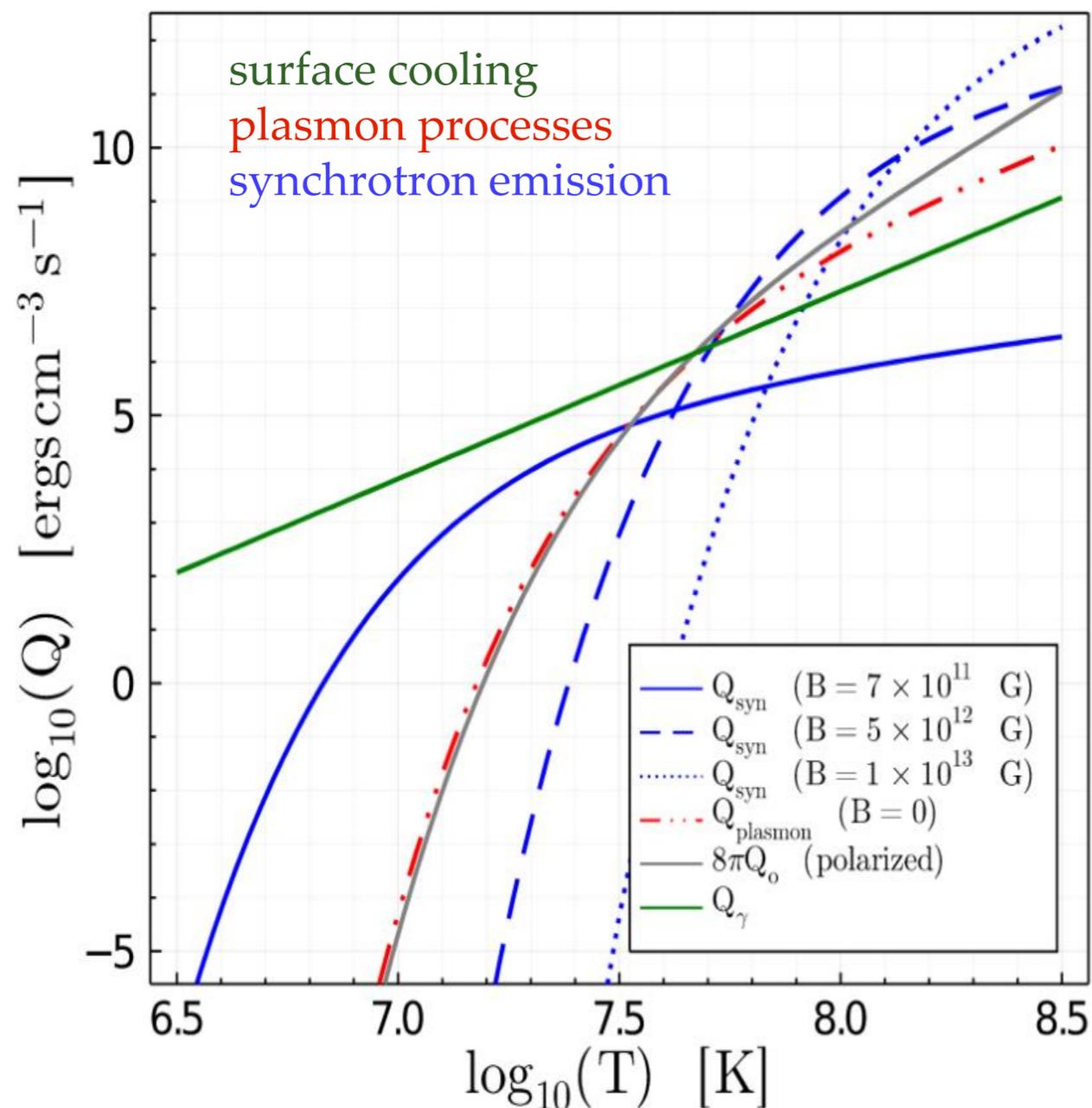
$$T = 5.0 \times 10^7 \text{ K}$$



- B-fields open up new cooling channel $e \rightarrow e\nu\nu$
- In relevant regime the effect grows with B
- For very large B: suppression because next Landau level becomes inaccessible

Comparing Mechanisms

$$\rho Y_e = 10^6 \text{ g cm}^{-3}$$



- Synchrotron emission can dominate for large temperatures
- Requires comparably large B fields
- Can potentially solve the anomaly for

$$B \sim 3 \times 10^{11} \text{ G}$$

- But how to generate these fields?
- Non-observation of stronger anomaly imposes upper bound

$$B < 6 \times 10^{11} \text{ G}$$

Final Remark: Ohmic Heating

- We so far assumed the B fields to be stable, but in reality they decay
- This decay releases energy into the plasma. Let's estimate

B-field energy density $U_{\text{em}} = B^2/2$

Exponential decay $B = B_0 \exp(-t/t_{\text{Ohm}})$

Heating rate Q $dU_{\text{em}}/dt \simeq B_0^2/t_{\text{Ohm}} \simeq 750 \text{ erg/cm}^3/\text{s}$

- For cooling time $t_{\text{Ohm}} \simeq 3 \times 10^{11}$ this yields

$$|Q_{\text{Ohm}}/Q_{\text{syn}}| \simeq 7.6 \quad \text{implying} \quad B < 2.73 \times 10^{11} \text{ G}/\sqrt{Y_e}$$

- Ohmic heating dominates over cooling! But depends on unknown cooling time...

Conclusions

- B-fields can in principle explain the WD cooling anomaly without the need for new LLPs...
- ...but requires rather large magnetic fields

$$B \sim 3 \times 10^{11} \text{ G}$$

- Non-observation of larger anomalies allows to impose a new upper bound in internal magnetic fields of WDs,

$$B < 6 \times 10^{11} \text{ G}$$

Even if this is not the solution to the cooling anomaly, it demonstrates the use of neutrino emission as a diagnostic tool for the internal structure of WDs!