### Searching for dark radiation at LHC

Based on <u>2204.01759</u>

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### A loving relationship...

LHC: power and versatility DM hints to new physics

LHC can look for DM DM can guide LHC



#### ...with some LLP issues

#### Simple observation:

 $H(T_{\text{EW}}) \leftrightarrow \text{LHC length}$ 

Interactions effective at the EW scale lead to macroscopic decay lengths!

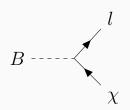
But  $\Omega h_{\rm DM}^2$  is not compatible with that...

But there are other cosmological observations!

 $\Delta N_{\rm eff}$  in the near future:  $\sigma_{\rm CMB-S4}=0.03$ 

### Model

New particle content: B,  $\chi$   $B \text{ decays into almost massless } \chi$   $B^T = (B_e, B_\mu, B_\tau) \text{ charged under SM}$ 



$$\mathcal{L}_{\mathsf{NP}} \supset B^T \cdot y_l \cdot (\bar{l}_R \chi) + h.c.$$

$$y_l = \begin{pmatrix} y & 0 & 0 \\ 0 & y & 0 \\ 0 & 0 & y \end{pmatrix} \quad \text{with } y \lesssim 10^{-6}$$

# Calculating $\Delta N_{\rm eff}$ (standard)

 $\Delta N_{
m eff}$  is the extra radiation added on top of SM

$$\Delta N_{\text{eff}}(x) = \frac{\rho_{\chi}(x)}{\rho_{1\nu}(x)} = \frac{Z_{\chi}(x) \, s_0^{4/3}}{\frac{7}{8} \left(\frac{4}{11}\right)^{4/3} \rho_{\gamma,0}},\tag{1}$$

$$Z_{\chi}(x) \equiv \frac{\rho_{\chi}(x)}{s^{4/3}(x)} \tag{2}$$

Z(x) can be derived by Boltzmann equation! IR freeze-in via parent decay: pretty easy!

# Calculating $\Delta N_{\rm eff}$ (refined calculation)

Usual assumptions:

- B decays while non-relativistic
- Backreaction  $\chi \, \mathsf{SM} \, \to B$  is negligible

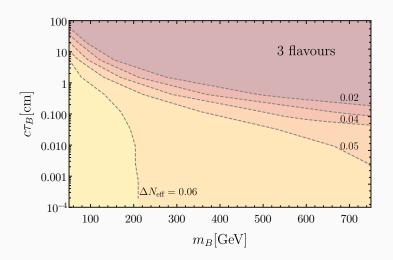
We relax these assumptions to get a better determination of the parameter space!

$$\tilde{H}xs^{4/3}(x)\frac{\mathrm{d}Z_{\chi}}{\mathrm{d}x} = \frac{m_B^4\Gamma_B}{8\pi^2}\mathcal{I}(x, T_{\chi}, \mathsf{spins}) \tag{3}$$

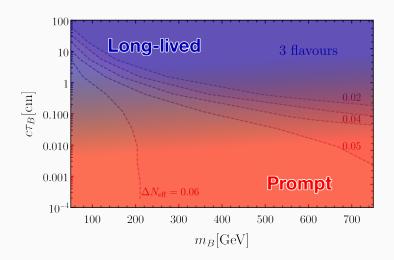
Also here there are some simplifications...

 $\mathcal{I}$  can be tabulated and is provided at the arXiv link.

## $\Delta N_{ m eff}$ result and LHC parameter space



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## Prompt searches

Recast SUSY searches for  $l l + E_T$  1908.08215, 2012.08600

Lifetime effect enters in the impact parameter cuts: different for ATLAS and CMS!

**ATLAS**:  $|d_0| < 3(5) \ \sigma(d_0)$  for electrons (muons)

**CMS**:  $|d_0| < 0.5$ mm for electrons and muons

Trick: the cut on  $d_0$  can be extrapolated for different lifetimes without re-simulating every time!

### LLP searches

"Recast" SUSY searches for displaced leptons 2011.07812, 2110.04809

Limits provided as function of  $m_B$  and  $c au_B$ 

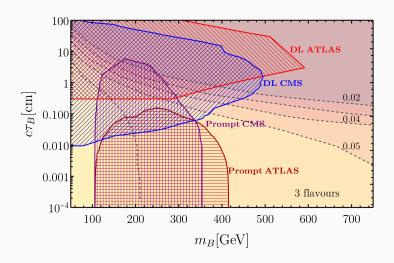
Still different cuts for ATLAS and CMS on the impact parameter

**ATLAS**:  $|d_0| \in [3\text{mm}, 300\text{mm}]$ 

**CMS**:  $|d_0| \in [0.1 \text{mm}, 100 \text{mm}]$ 

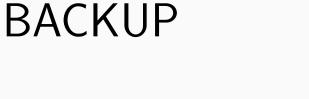
Warning: recasting here is more complicated for the single-flavour scenario

# LHC constraints on $\Delta N_{ m eff}$

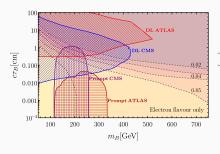


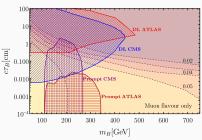
### Conclusions

- ullet Calculations of  $\Delta N_{
  m eff}$  has been improved to better determine the decay lengths
- The interesting parameter space lies at the boundary of prompt and long-lived searches → complementarity!
- ATLAS and CMS have different cuts which result in differences in parameter space probed

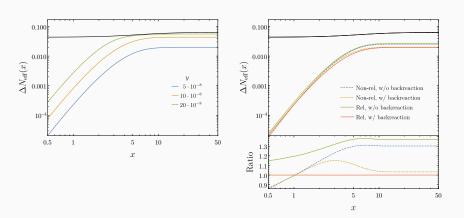


# Single flavour case



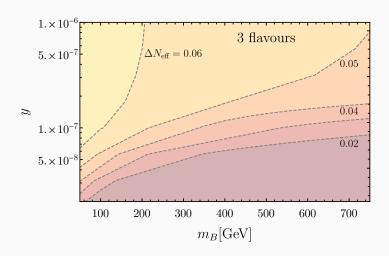


# Effect of approximations



Not portrayed here: magnitude of corrections depends sensitively on the parameters!

# Coupling parameter space



### Other ATLAS-CMS differences

#### **Prompt ATLAS:**

bins in  $m_{T2}$ ,  $e\mu$  as signal region

#### **Prompt CMS**:

bins in  $p_T^{\rm miss}$ ,  $e\mu$  as control region

**LLP CMS** does not provide limits on  $\sigma(m_B,c\tau_B)$  for the single flavour scenario, so an approximation on the mass dependence is used.