

Long-lived particles in t -channel mediator dark matter models

Jan Heisig

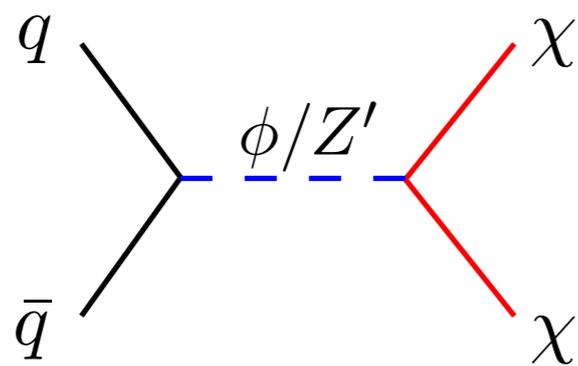
Chargé de
recherches



*Searching for long-lived particles at the LHC and beyond:
Eleventh workshop of the LLP Community
June 2, 2022*

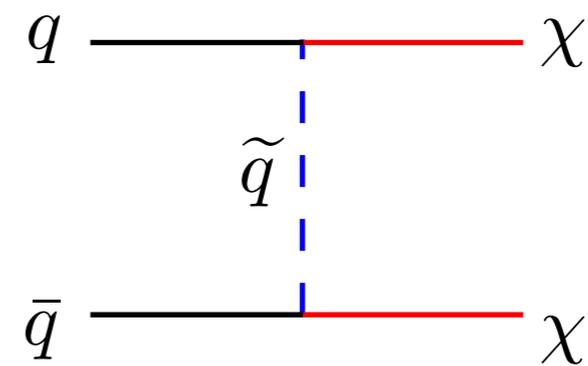
Simplified dark matter models at LHC

s-channel mediator



even

t-channel mediator



odd

under Z_2 -symmetry that stabilizes dark matter

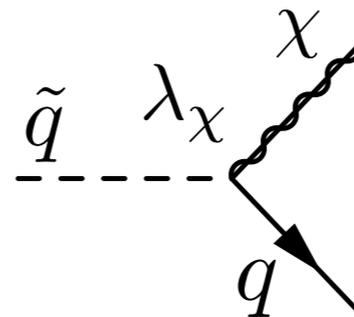
Example: SUSY-like t -channel mediator model

Majorana DM χ $m_\chi < m_{\tilde{q}}$ scalar quark-partner \tilde{q}

λ_χ g_s

$$\mathcal{L}_{\text{int}} = |D_\mu \tilde{q}|^2 - \lambda_\chi \tilde{q} \tilde{q} \frac{1 - \gamma_5}{2} \chi + \text{h.c.}$$

\Rightarrow Yukawa-type interaction:



In general, λ_χ is a free parameter here [see Ibarra et al. 2009 for SUSY realization]

Simplified dark matter models with t -channel mediators

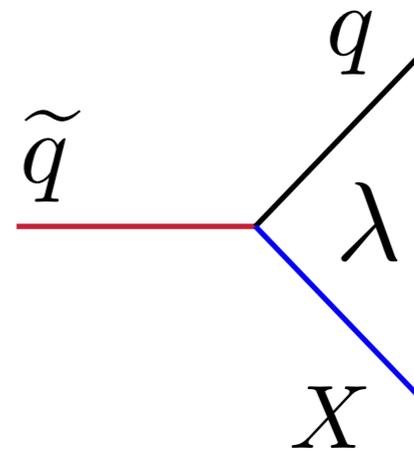
dark matter

Field	Spin	Repr.	Self-conj.
\tilde{S}	0	$(\mathbf{1}, \mathbf{1}, 0)$	yes
S	0	$(\mathbf{1}, \mathbf{1}, 0)$	no
$\tilde{\chi}$	1/2	$(\mathbf{1}, \mathbf{1}, 0)$	yes
χ	1/2	$(\mathbf{1}, \mathbf{1}, 0)$	no
\tilde{V}_μ	1	$(\mathbf{1}, \mathbf{1}, 0)$	yes
V_μ	1	$(\mathbf{1}, \mathbf{1}, 0)$	no

t -channel mediator

$\varphi_Q = \begin{pmatrix} \varphi_Q^{(u)} \\ \varphi_Q^{(d)} \end{pmatrix}$	0	$(\mathbf{3}, \mathbf{2}, \frac{1}{6})$	no
φ_u	0	$(\mathbf{3}, \mathbf{1}, \frac{2}{3})$	no
φ_d	0	$(\mathbf{3}, \mathbf{1}, -\frac{1}{3})$	no
$\psi_Q = \begin{pmatrix} \psi_Q^{(u)} \\ \psi_Q^{(d)} \end{pmatrix}$	1/2	$(\mathbf{3}, \mathbf{2}, \frac{1}{6})$	no
ψ_u	1/2	$(\mathbf{3}, \mathbf{1}, \frac{2}{3})$	no
ψ_d	1/2	$(\mathbf{3}, \mathbf{1}, -\frac{1}{3})$	no

[Arina, Fuks, Mantani 2001.05024]



$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{kin}} + \mathcal{L}_F(\chi) + \mathcal{L}_F(\tilde{\chi})$$

$$+ \mathcal{L}_S(S) + \mathcal{L}_S(\tilde{S}) + \mathcal{L}_V(V) + \mathcal{L}_V(\tilde{V}),$$

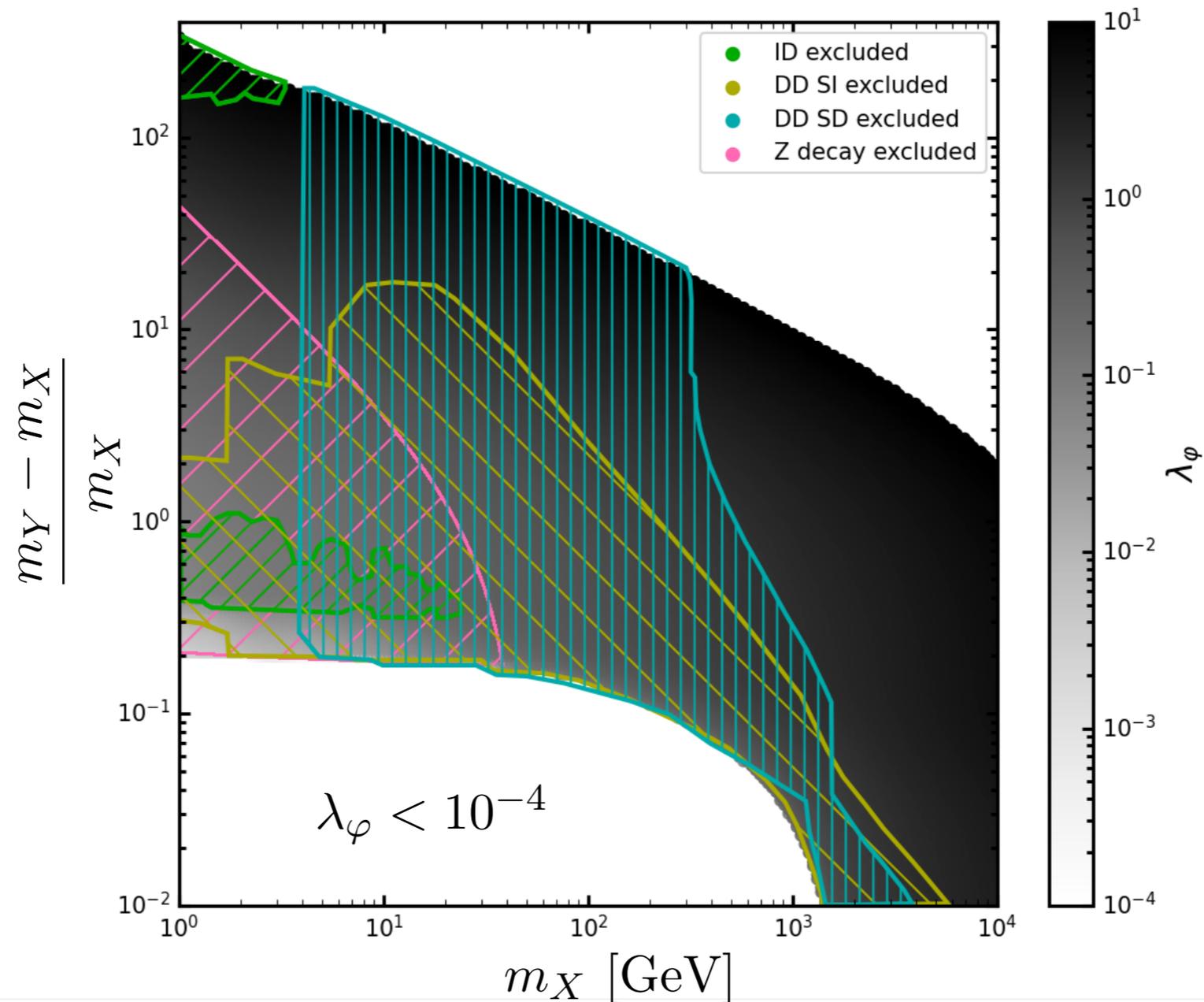
$$\mathcal{L}_F(X) = \left[\lambda_Q \bar{X} Q \varphi_Q^\dagger + \lambda_u \bar{X} u \varphi_u^\dagger + \lambda_d \bar{X} d \varphi_d^\dagger + \text{h.c.} \right],$$

$$\mathcal{L}_S(X) = \left[\hat{\lambda}_Q \bar{\psi}_Q Q X + \hat{\lambda}_u \bar{\psi}_u u X + \hat{\lambda}_d \bar{\psi}_d d X + \text{h.c.} \right],$$

$$\mathcal{L}_V(X) = \left[\hat{\lambda}_Q \bar{\psi}_Q \not{X} Q + \hat{\lambda}_u \bar{\psi}_u \not{X} u + \hat{\lambda}_d \bar{\psi}_d \not{X} d + \text{h.c.} \right].$$

Example: SUSY-like t -channel mediator model

$$\Omega h^2 = 0.12$$



[Arina et al. 2010.07559;
Arina, Fuks, JH, et al., in preparation]

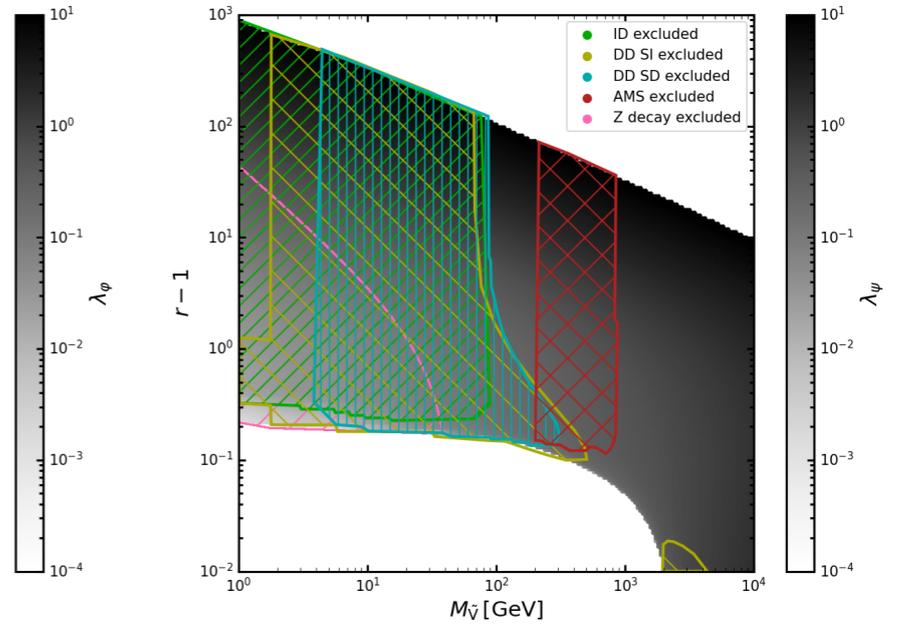
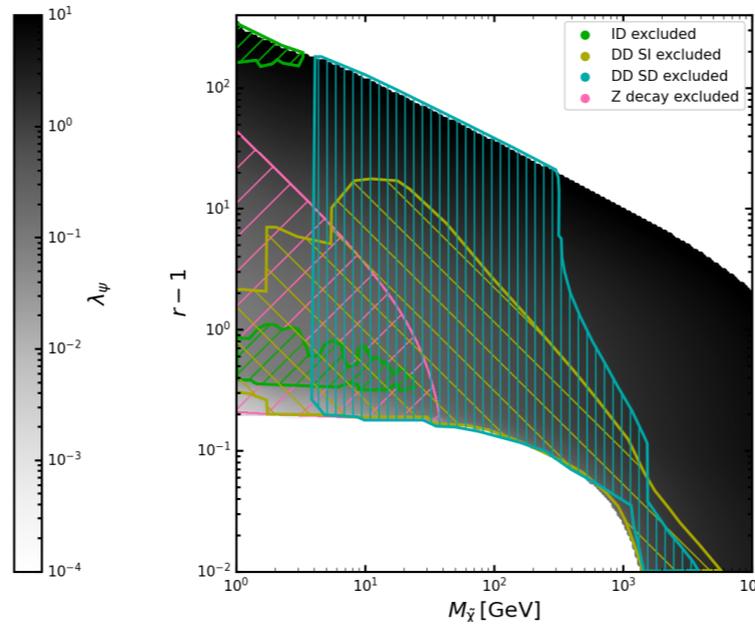
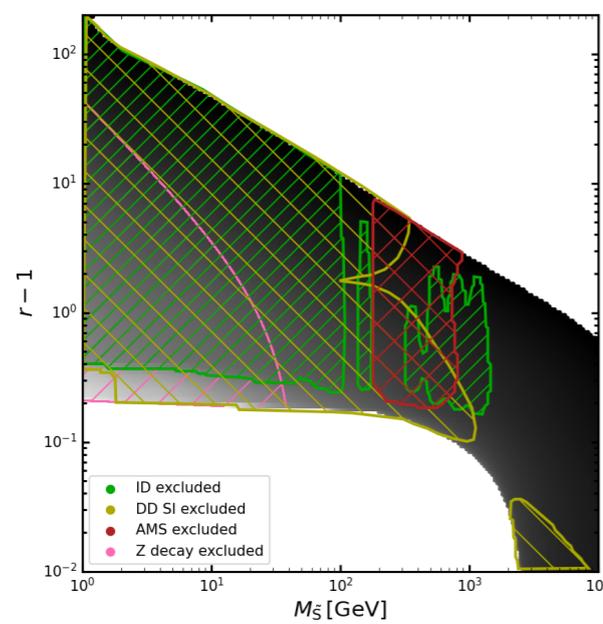
Other choices

scalar DM, fermion med.

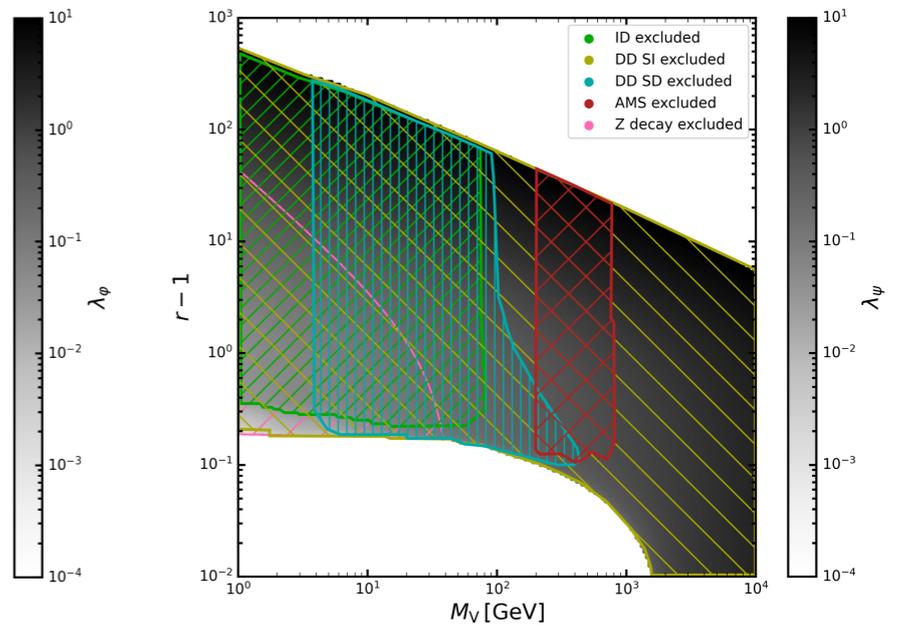
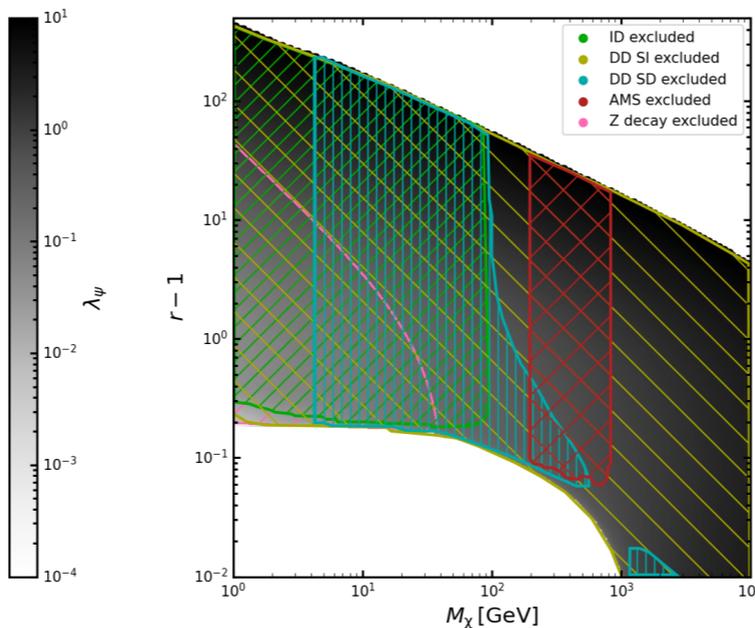
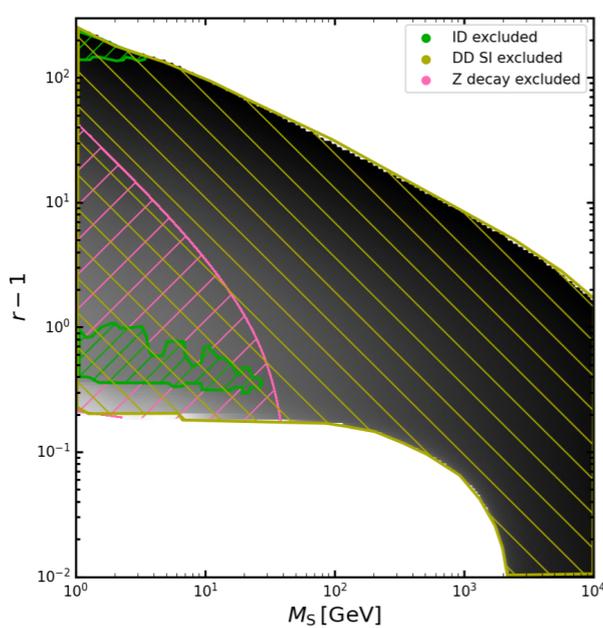
fermion DM, scalar med.

vector DM, fermion med.

self-conjugate



'complex'



[Arina, Fuks, JH, Krämer, Mantani, Panizzi, Salko, *in preparation*]

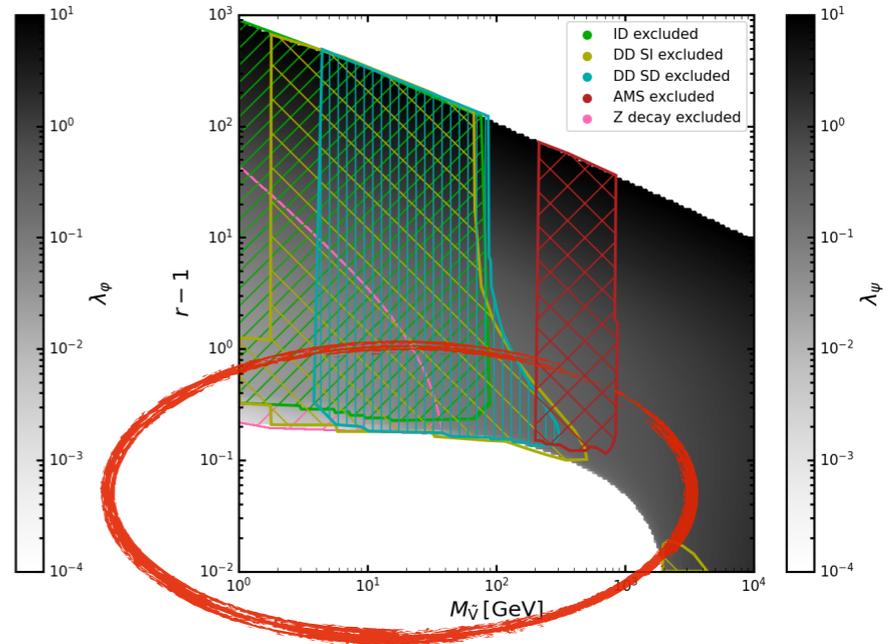
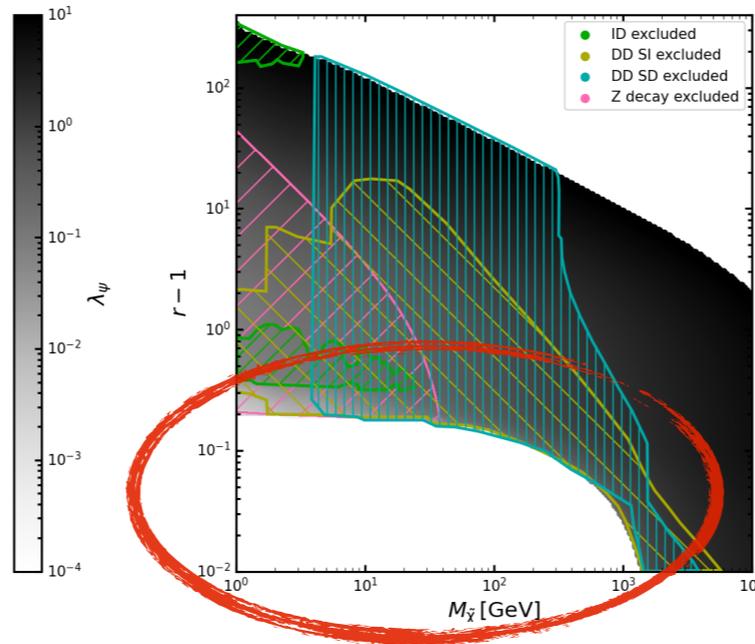
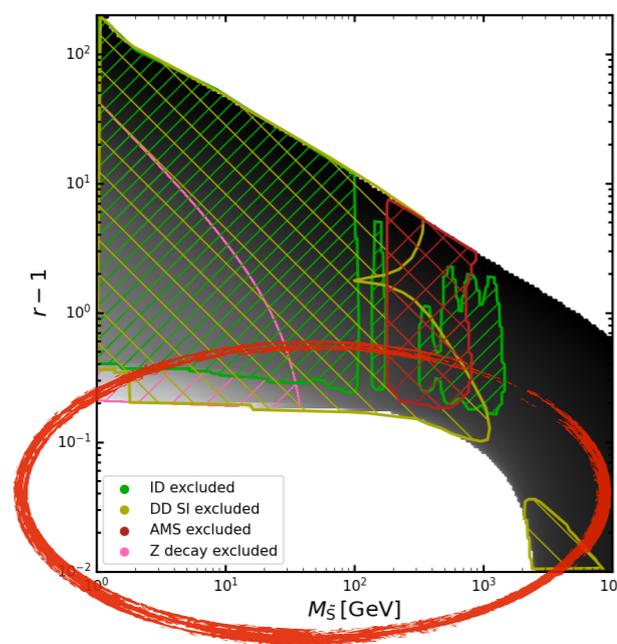
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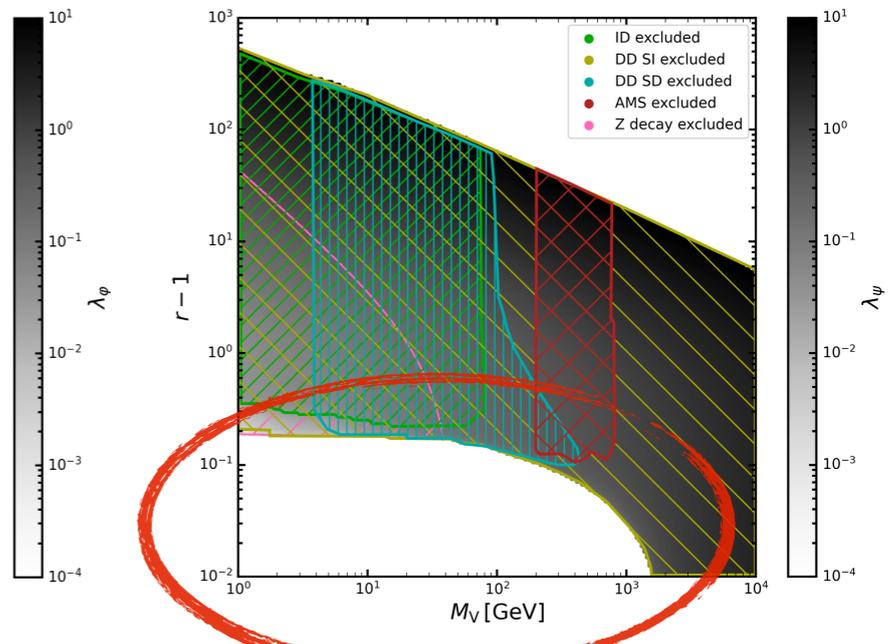
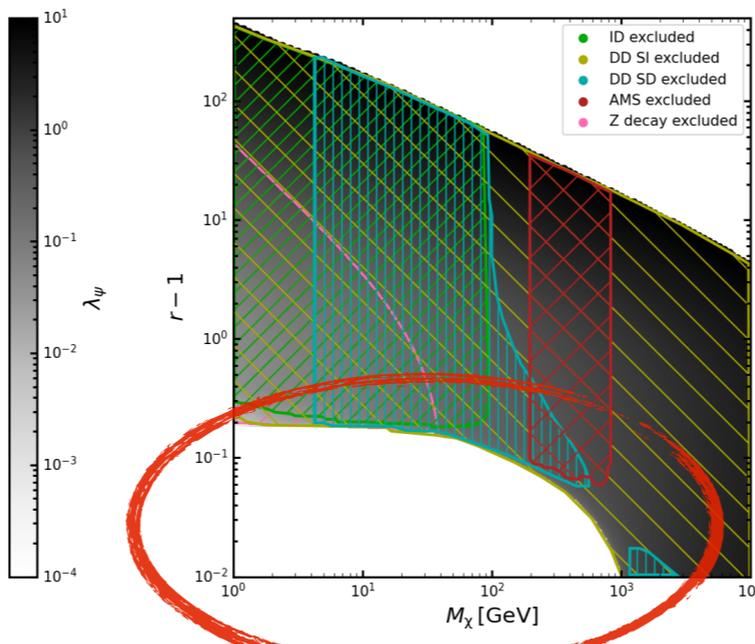
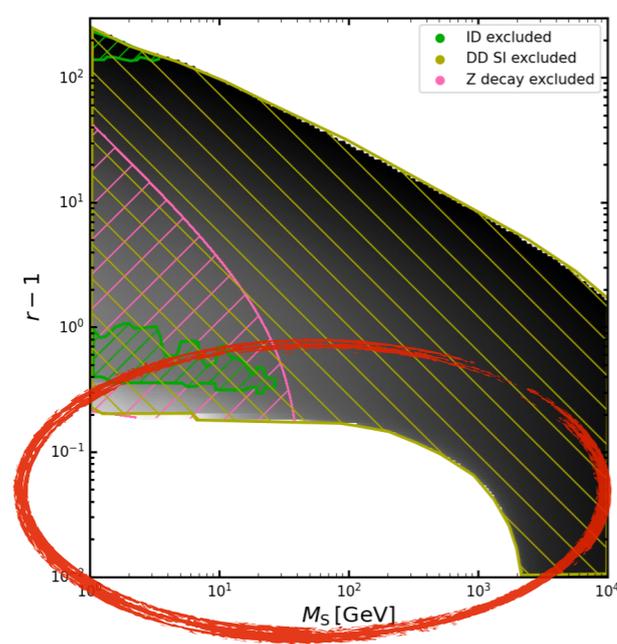
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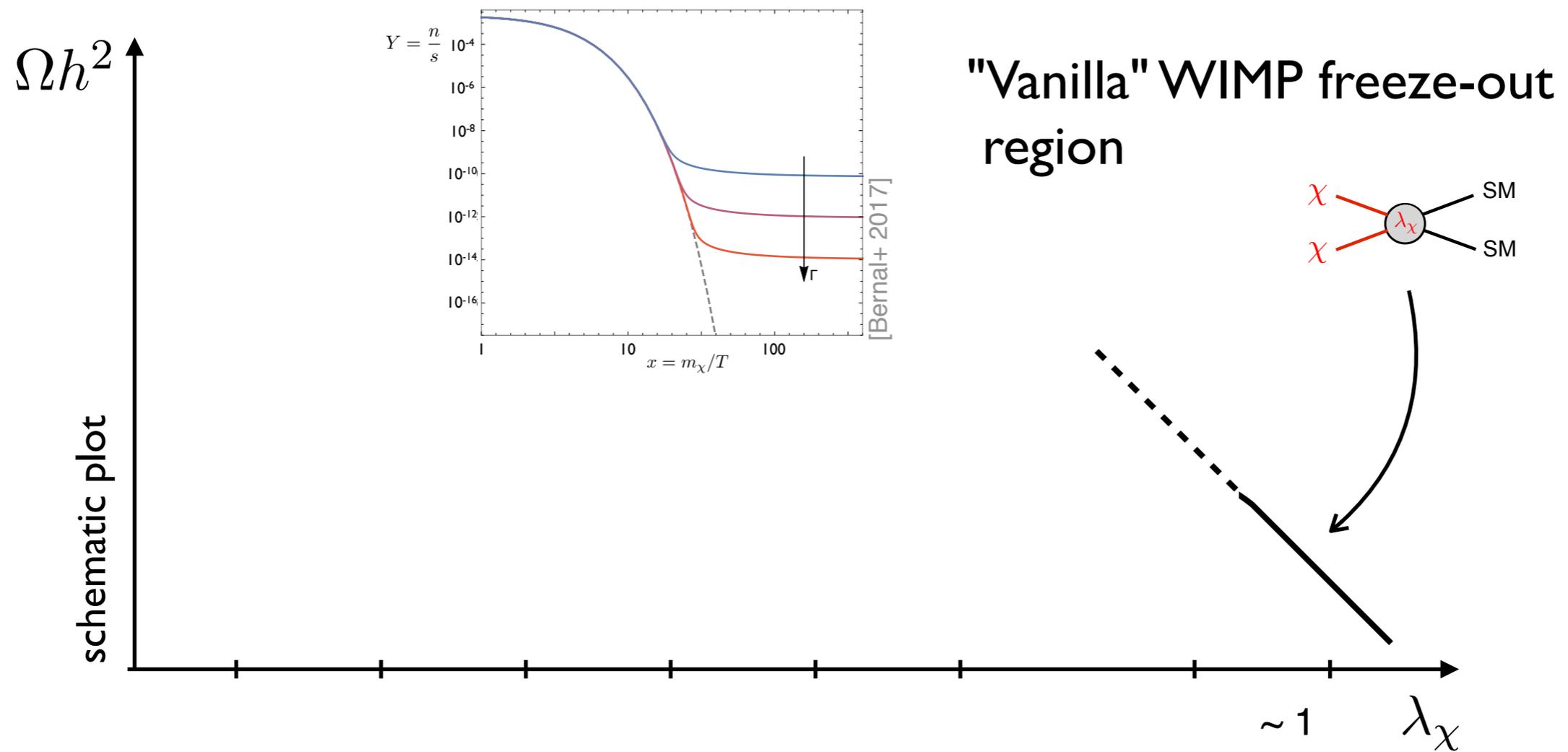


'complex'

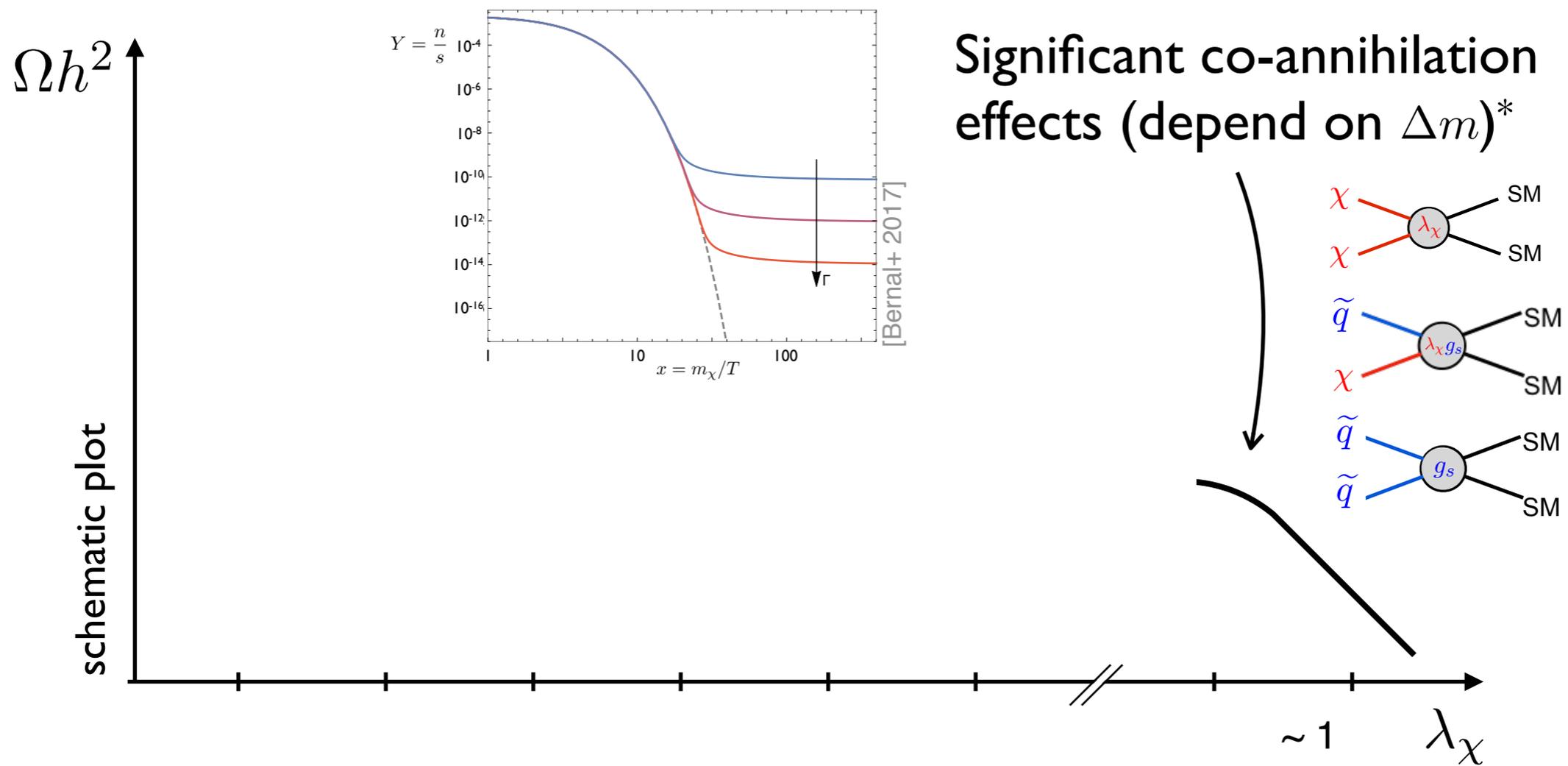


[Arina, Fuks, JH, Krämer, Mantani, Panizzi, Salko, *in preparation*]

How does the relic density work out?

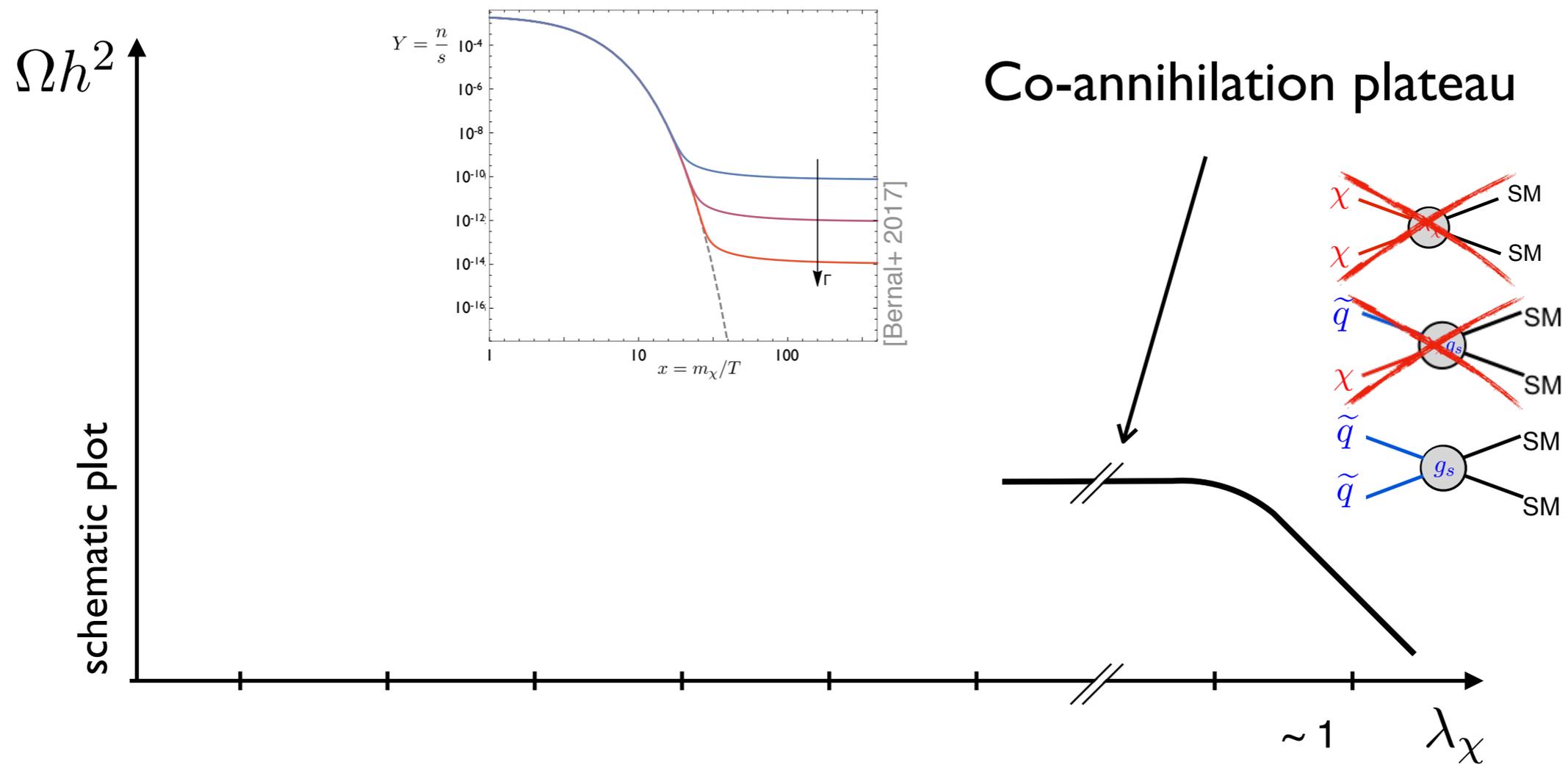


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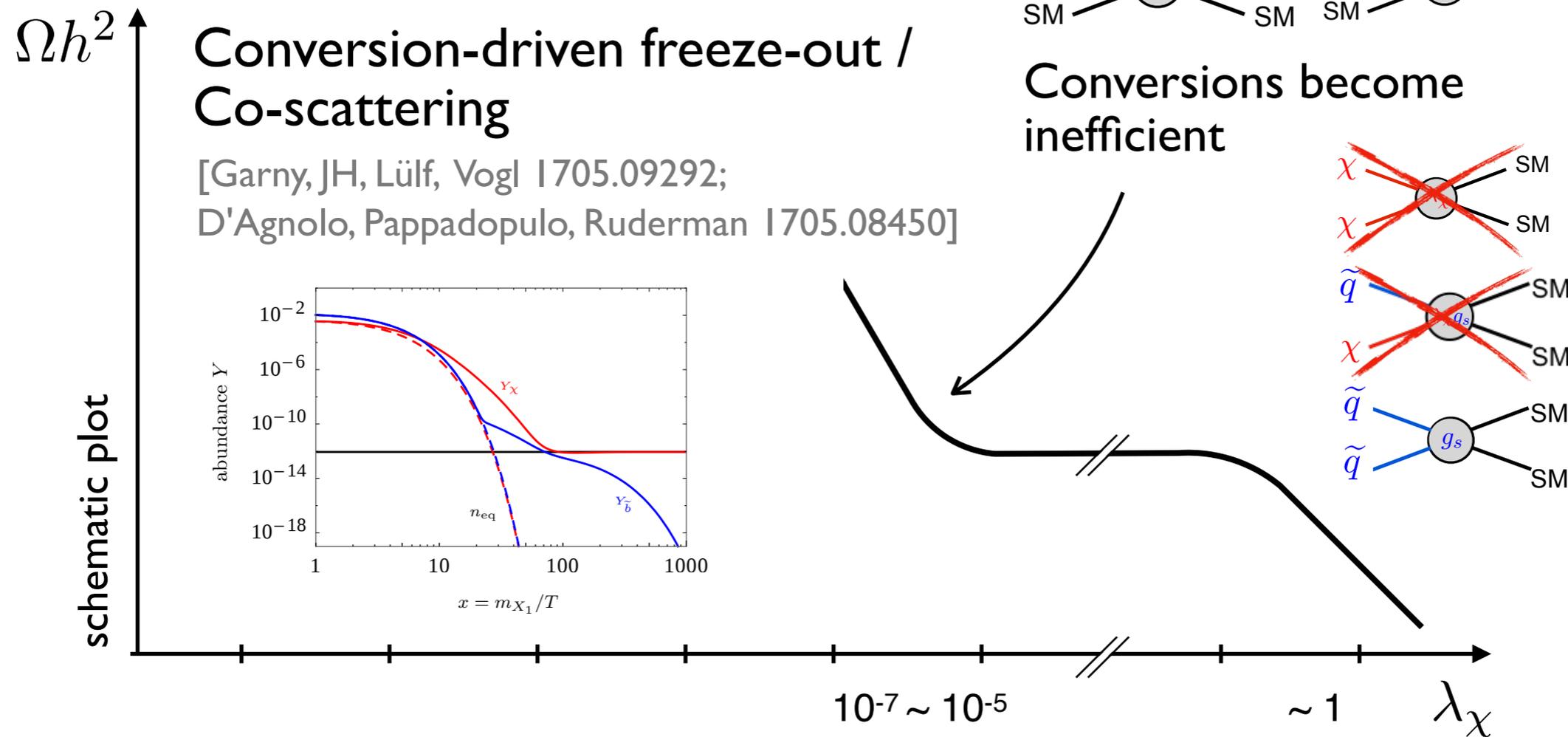


[* see e.g. D'Agnolo *et al.* 1803.02901]

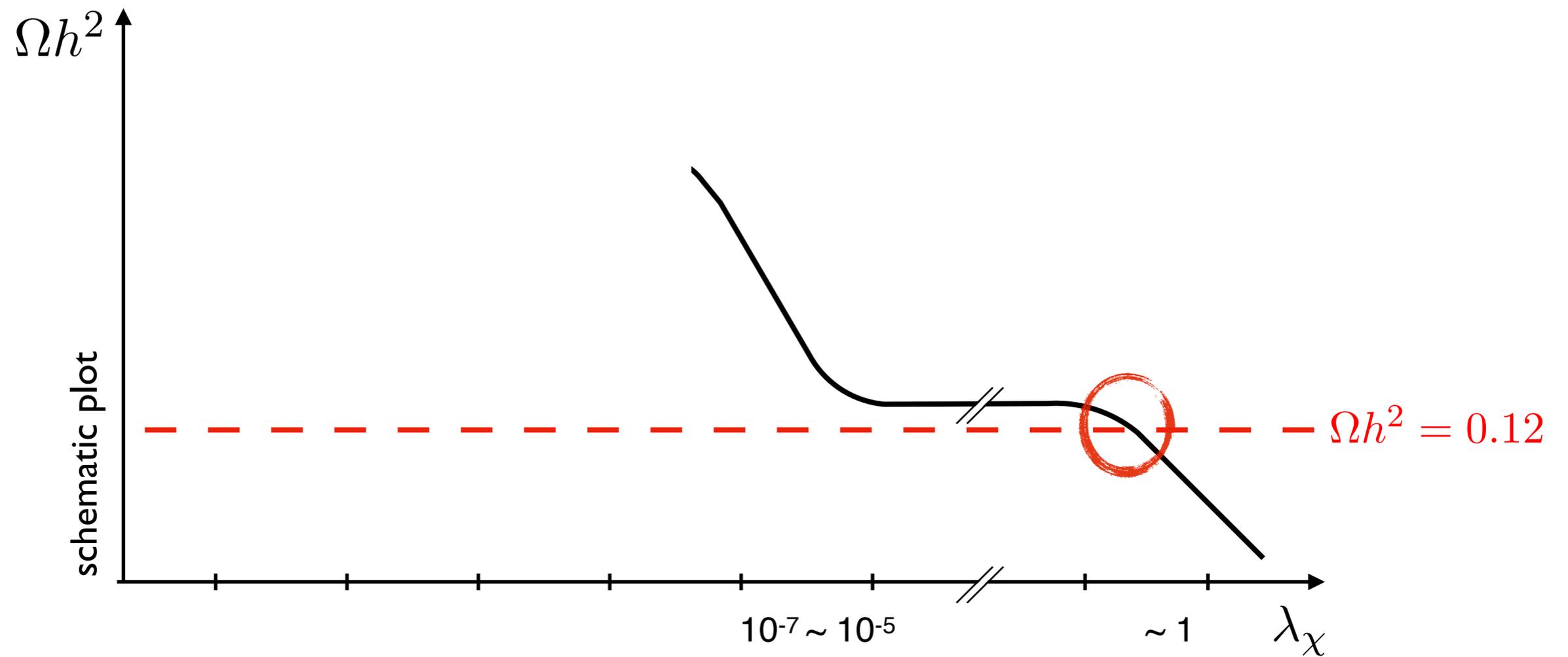
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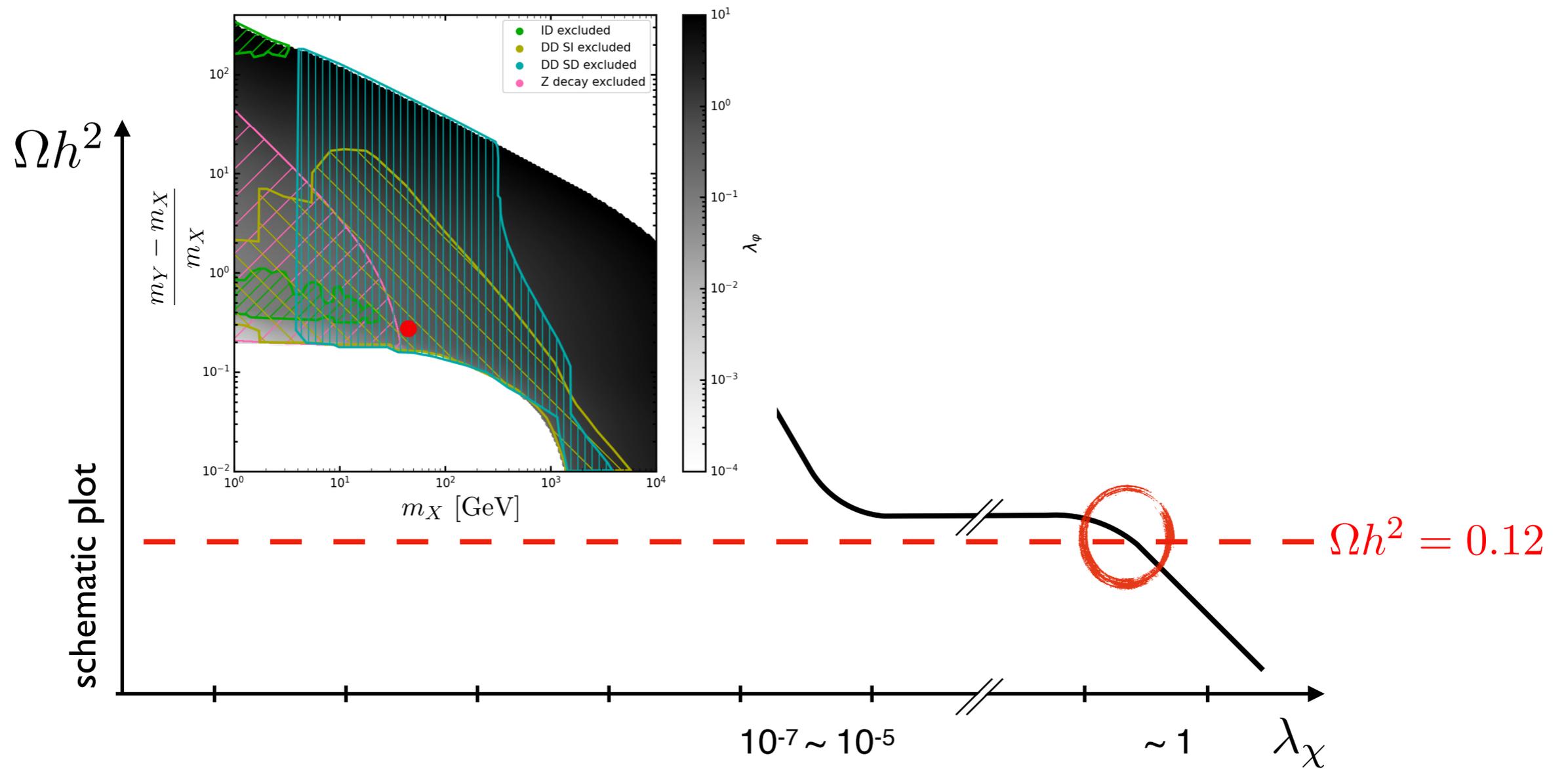
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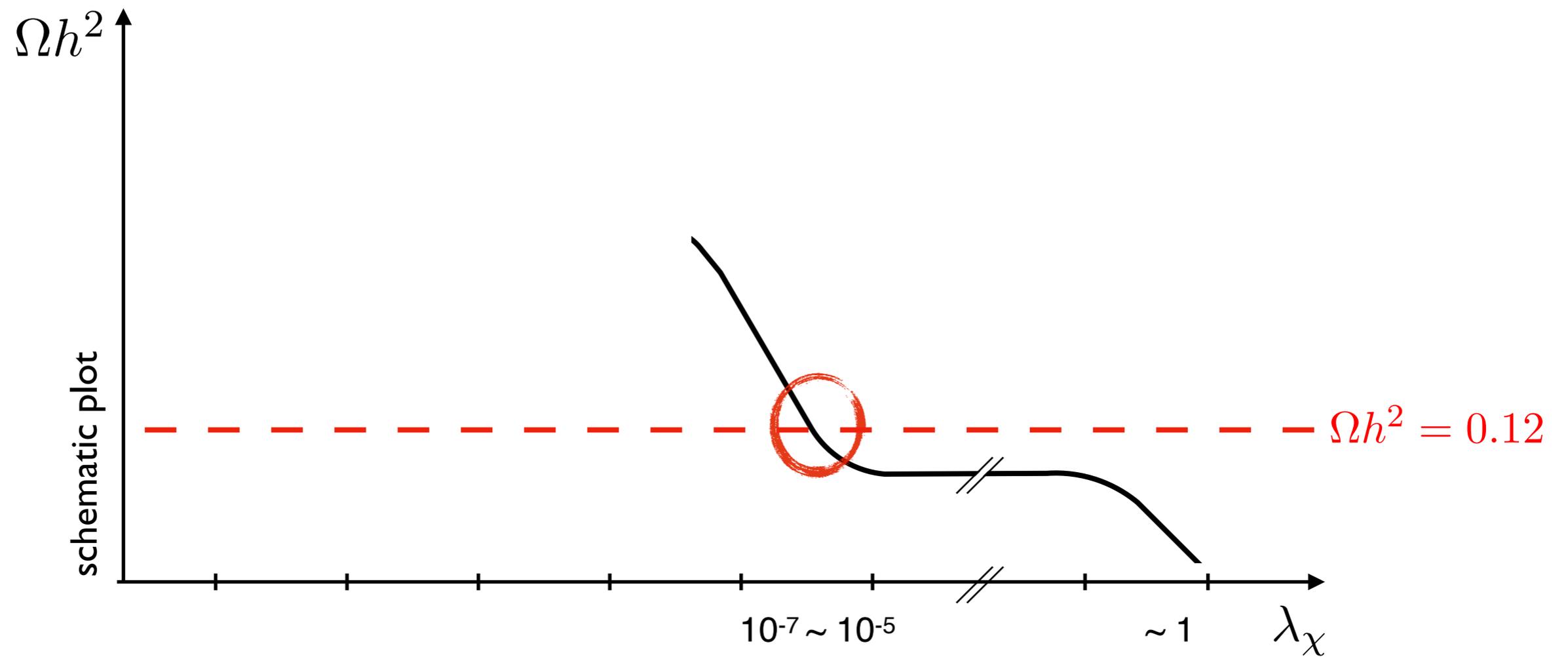
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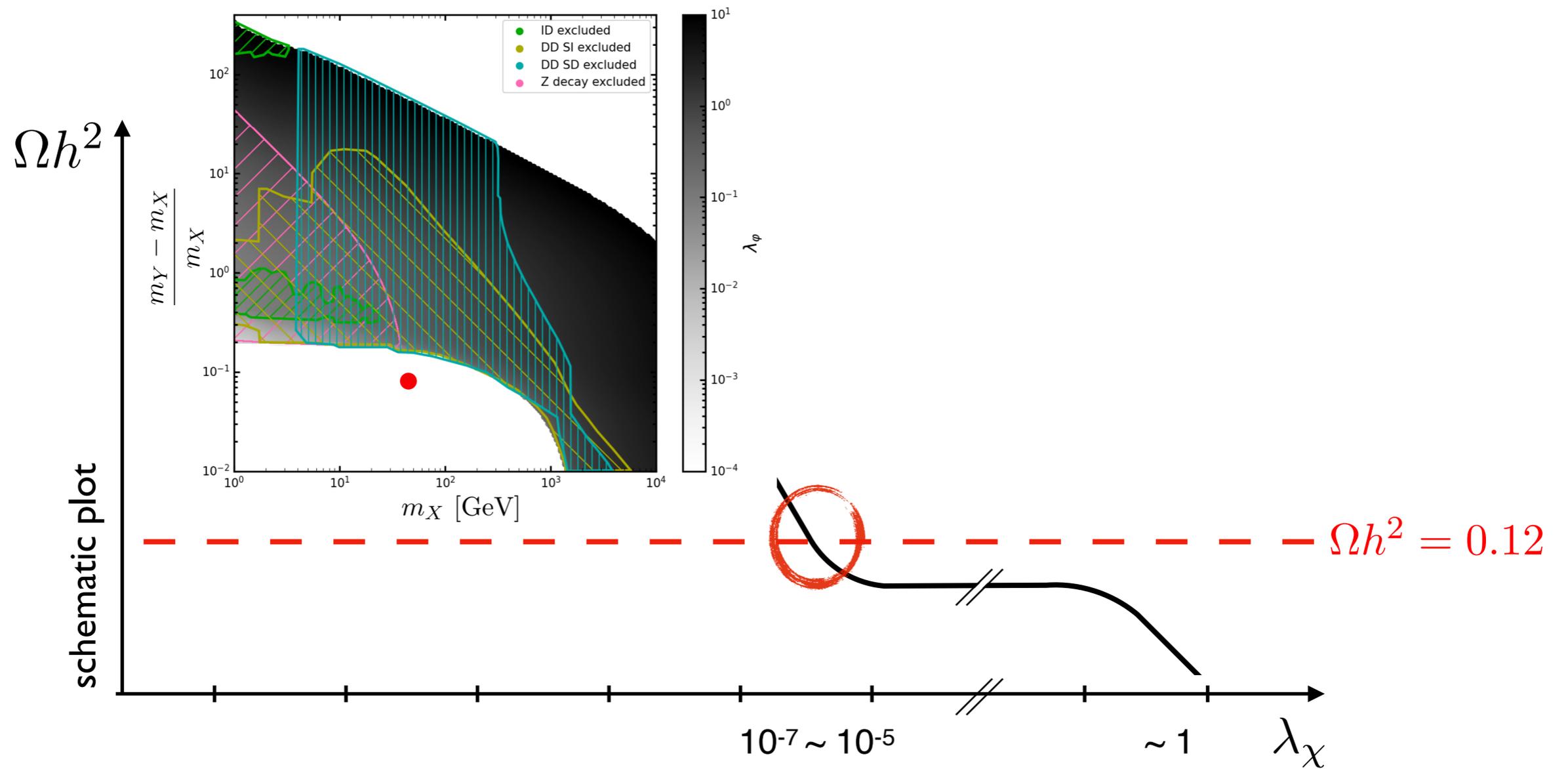
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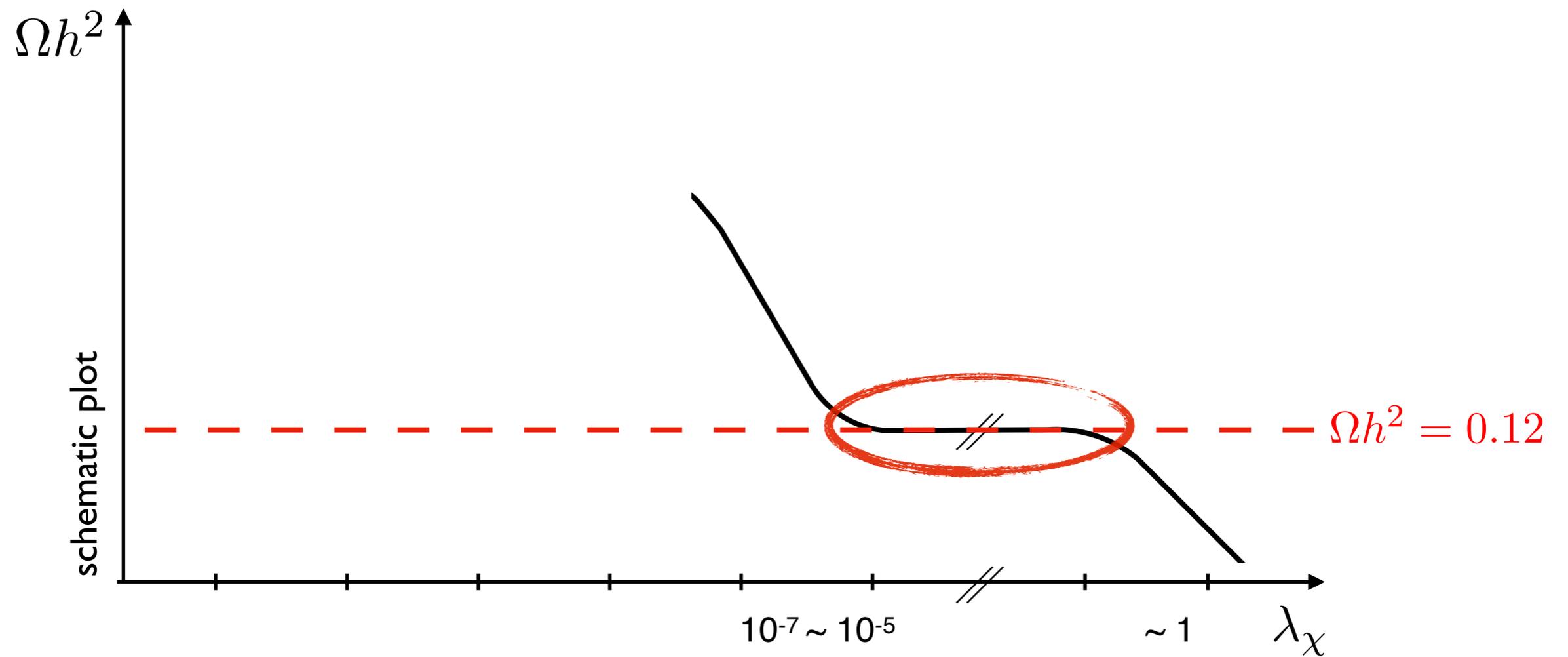
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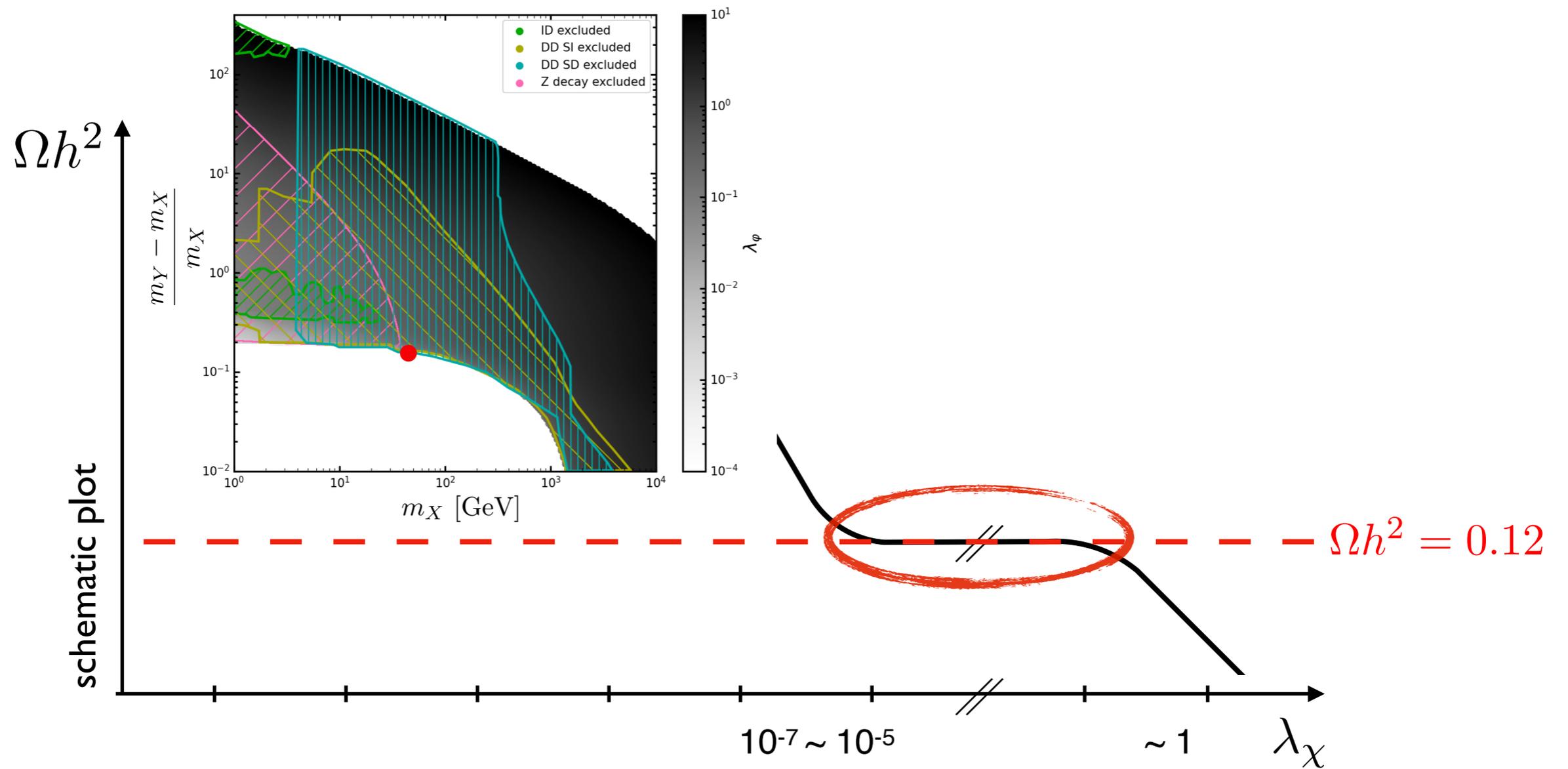
How does the relic density work out?



How does the relic density work out?



How does the relic density work out?



Generic signature?

Conversion rate on the edge of being efficient:

$$\Gamma_{\text{conv}} \sim H$$

$$\Rightarrow \Gamma_{\text{dec}} \lesssim H$$

$$c\tau \gtrsim H^{-1} \simeq 1.5 \text{ cm} \left(\frac{(100 \text{ GeV})^2}{T^2} \right)$$

$$T \lesssim (10-100) \text{ GeV}$$

\Rightarrow Long-lived particles (LLPs) at LHC!

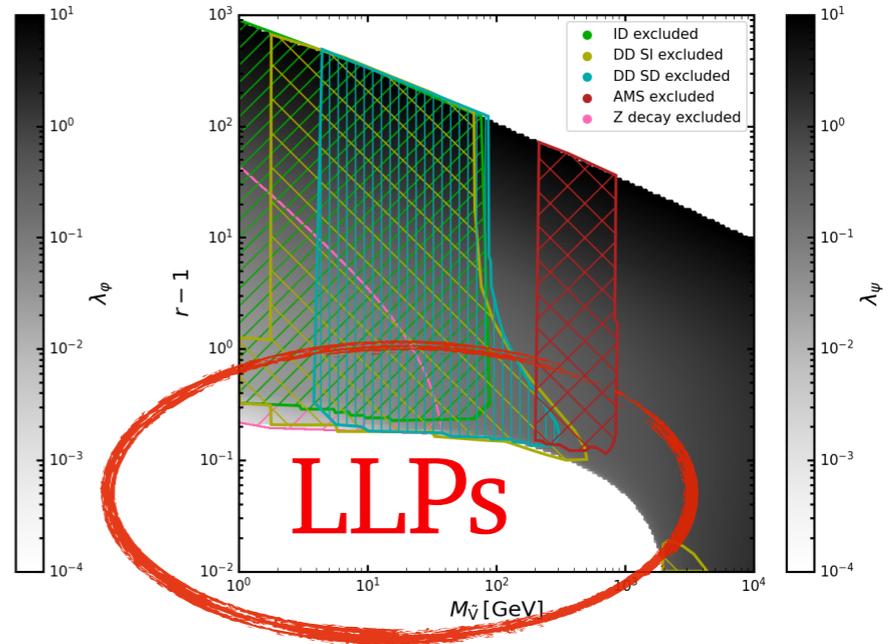
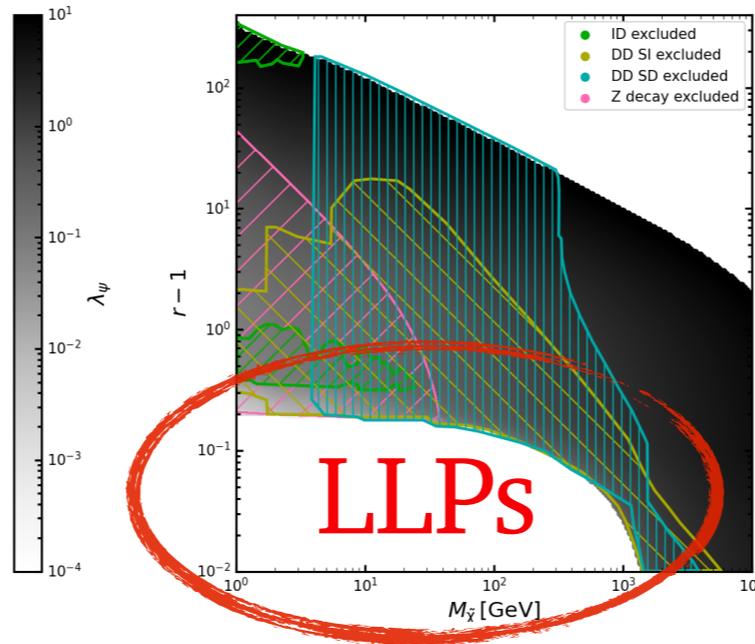
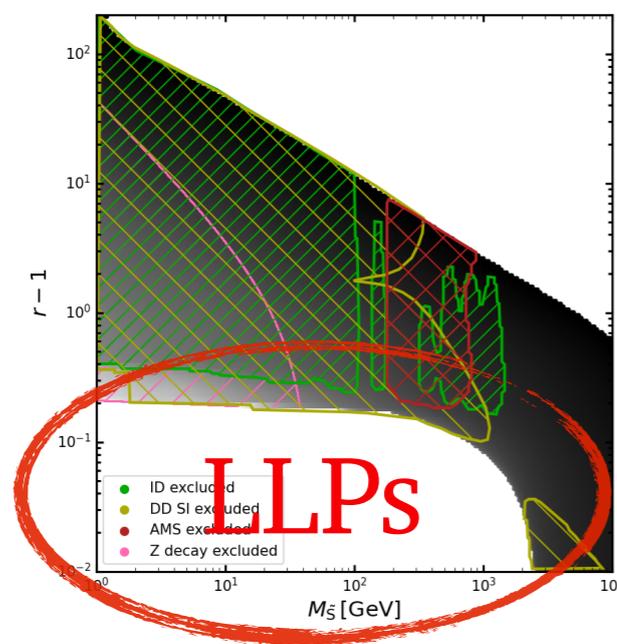
Other choices

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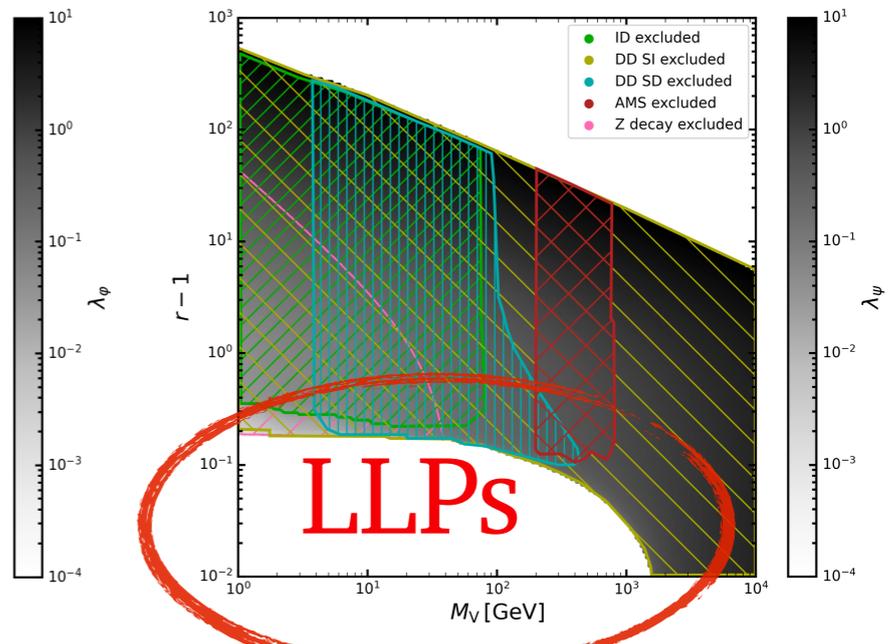
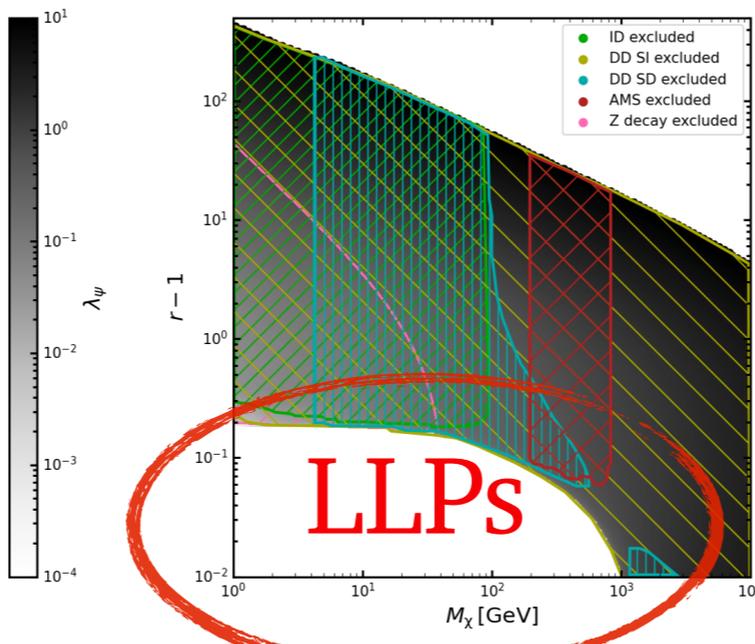
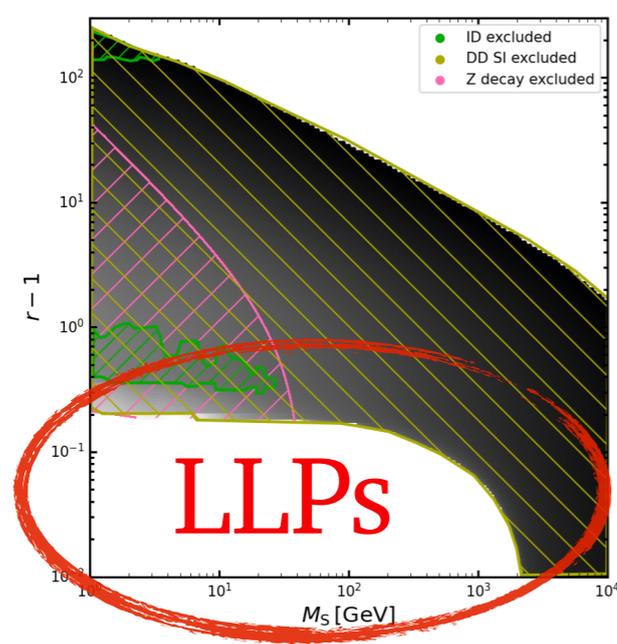
fermion DM, scalar med.

vector DM, fermion med.

self-conjugate



'complex'



Progress in relic density computation

Bound-state effects on dark matter coannihilation: Pushing the boundaries of conversion-driven freeze-out

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James-Franck-Straße 1, D-85748 Garching, Germany*

Jan Heisig

*Institute for Theoretical Particle Physics and Cosmology,
RWTH Aachen University, Sommerfeldstraße 16, D-52056 Aachen, Germany and
Centre for Cosmology, Particle Physics and Phenomenology (CP3),
Université catholique de Louvain, Chemin du Cyclotron 2, B-1348 Louvain-la-Neuve, Belgium*

Bound-state formation can have a large impact on the dynamics of dark matter freeze-out in the early Universe, in particular for colored coannihilators. We present a general formalism to include an arbitrary number of excited bound states in terms of an effective annihilation cross section, taking bound-state formation, decay and transitions into account, and derive analytic approximations in the limiting cases of no or efficient

arXiv:2112.01499

Impact of Sommerfeld Effect and Bound State Formation in Simplified t -Channel Dark Matter Models

Mathias Becker^{*1}, Emanuele Copello^{†1}, Julia Harz^{‡1}, Kirtimaan A. Mohan^{§2}, and Dipan Sengupta^{¶3,4}

¹Physik Department T70, Technische Universität München, 85748 Garching, Germany

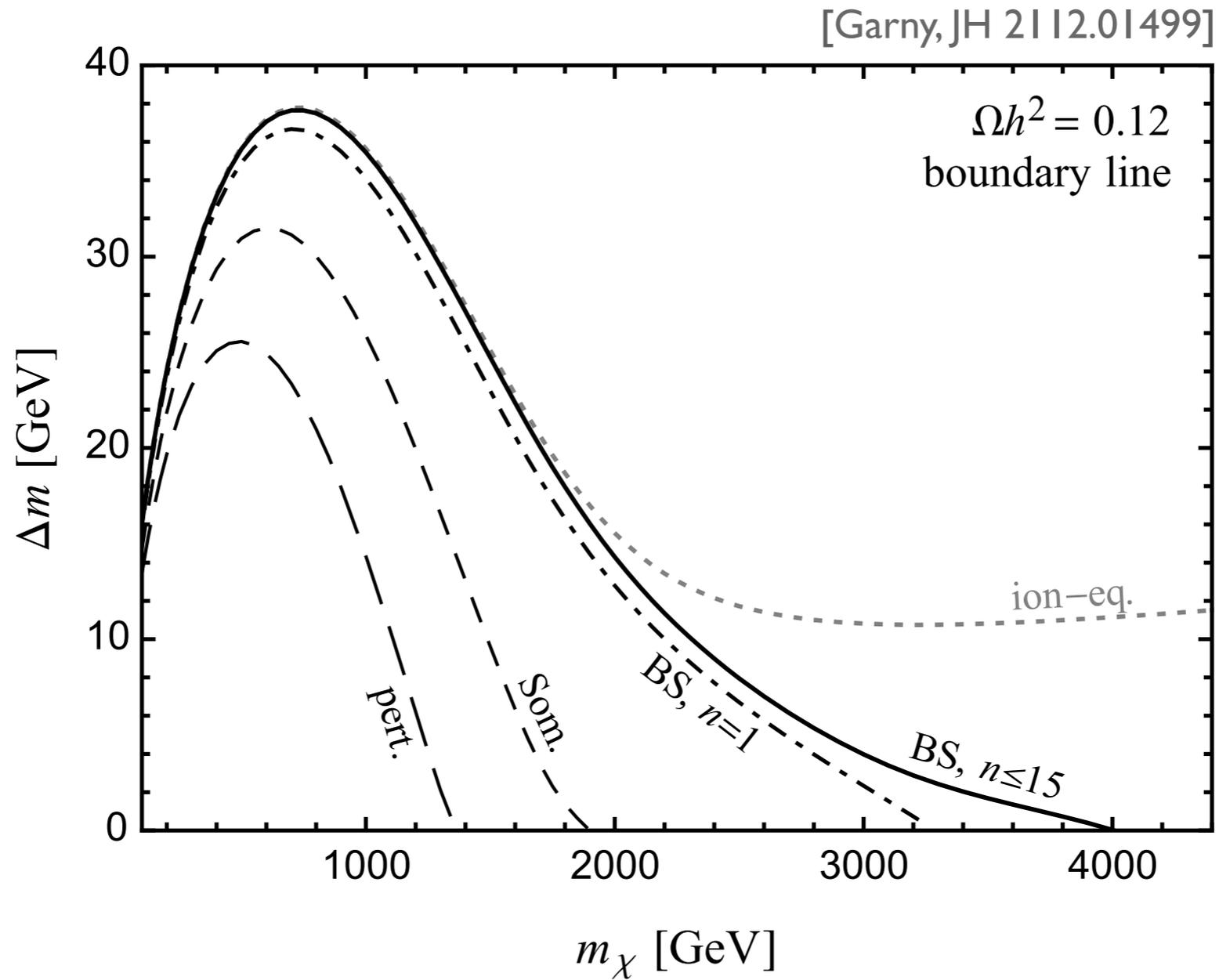
²Department of Physics and Astronomy, 567 Wilson Road, East Lansing, Michigan-48824, USA

³Department of Physics and Astronomy, 9500 Gilman Drive, University of California, San Diego, USA

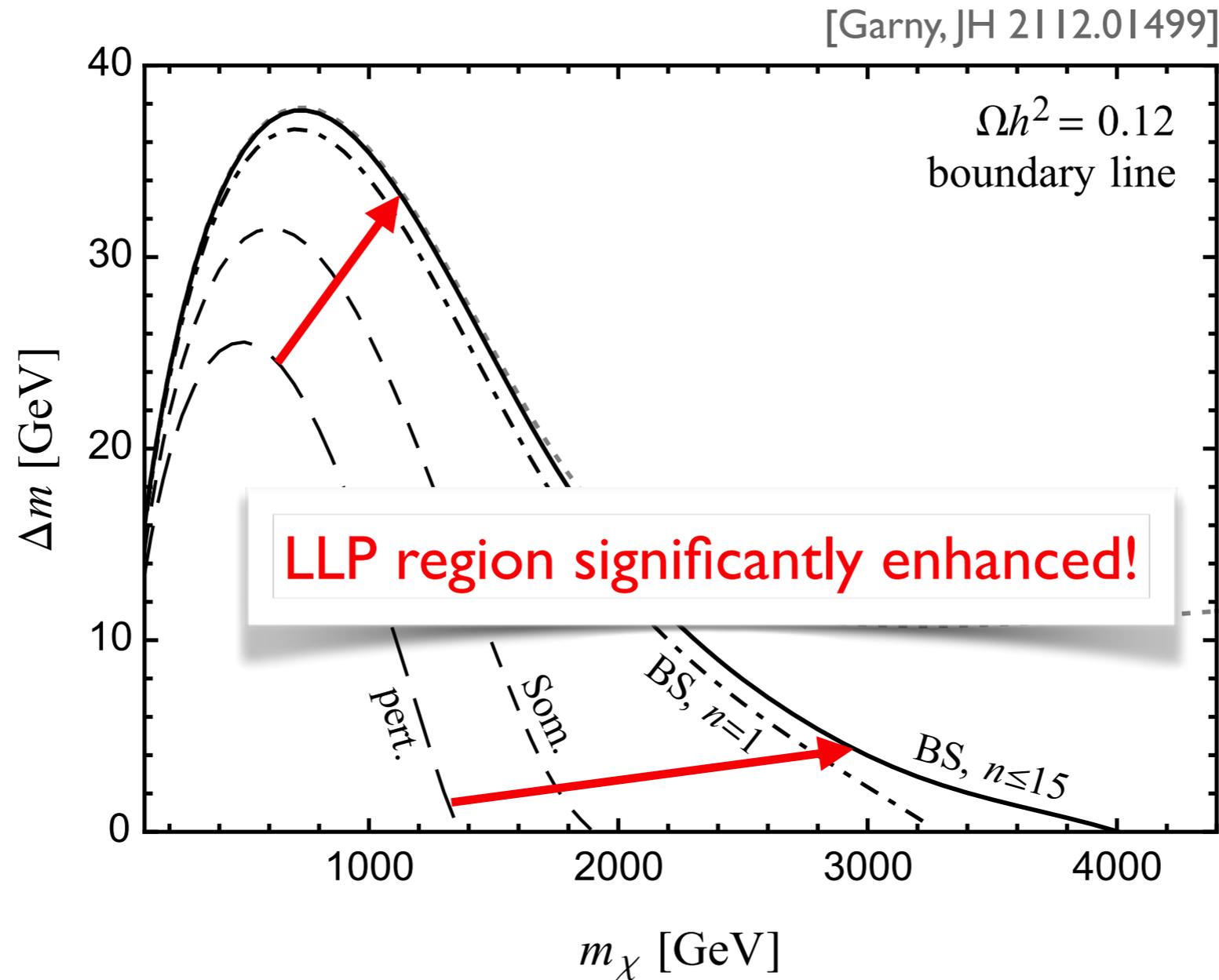
⁴ARC Centre of Excellence for Dark Matter Particle Physics, Department of Physics, University of Adelaide, South Australia 5005, Australia

arXiv:2203.04326

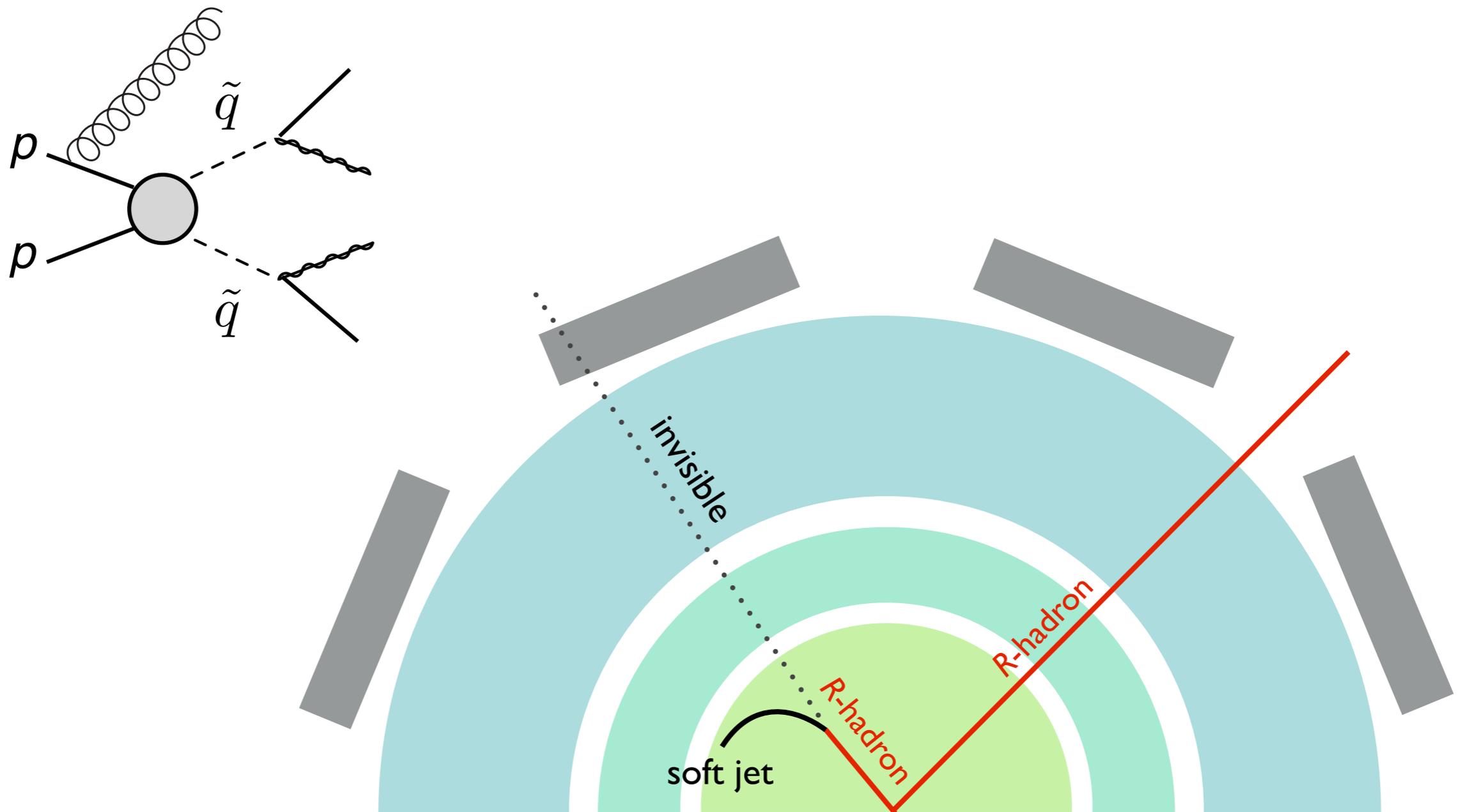
Boundary between WIMP region and conversion-driven region



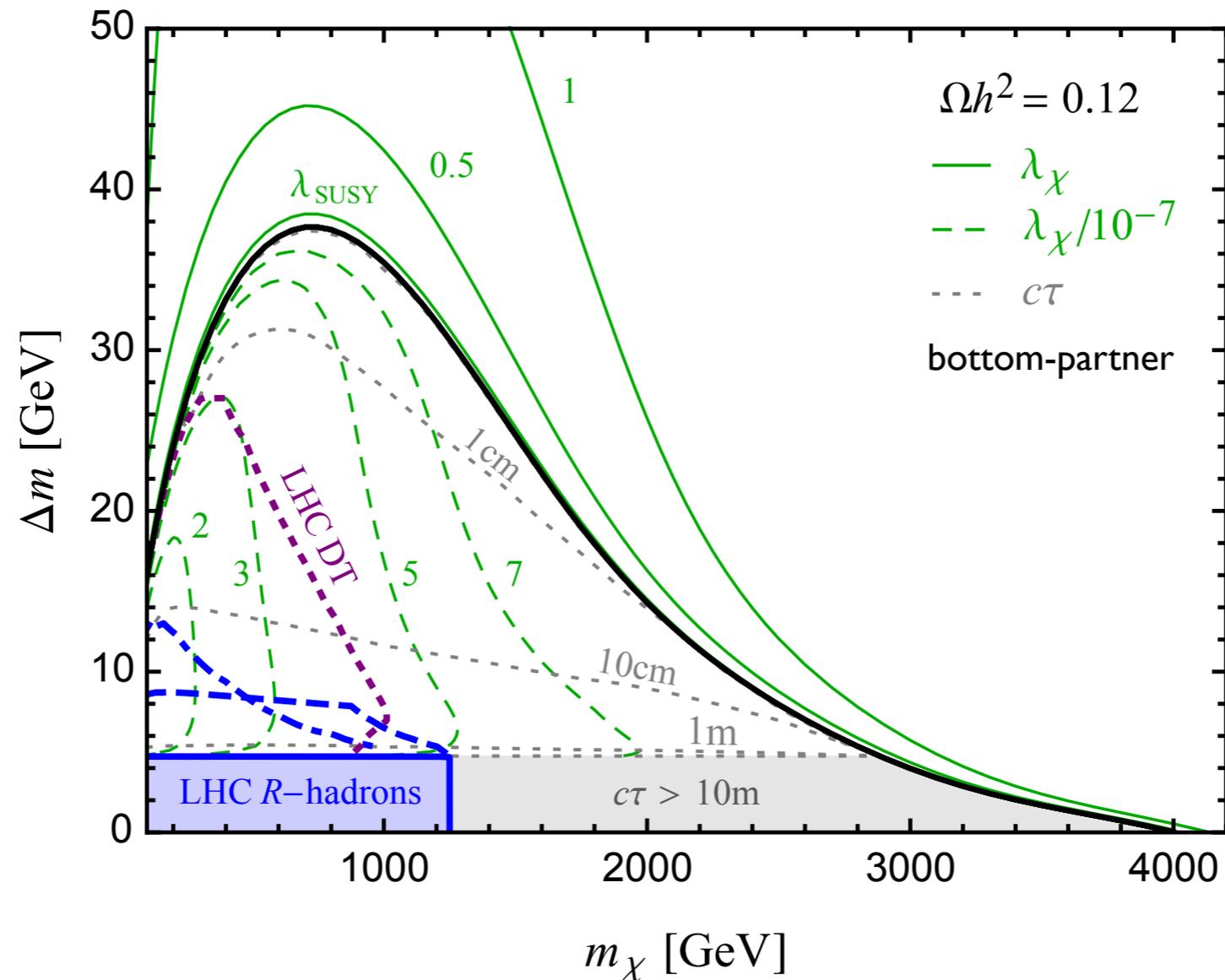
Boundary between WIMP region and conversion-driven region



Long-lived particles at LHC



Current LLP constraints

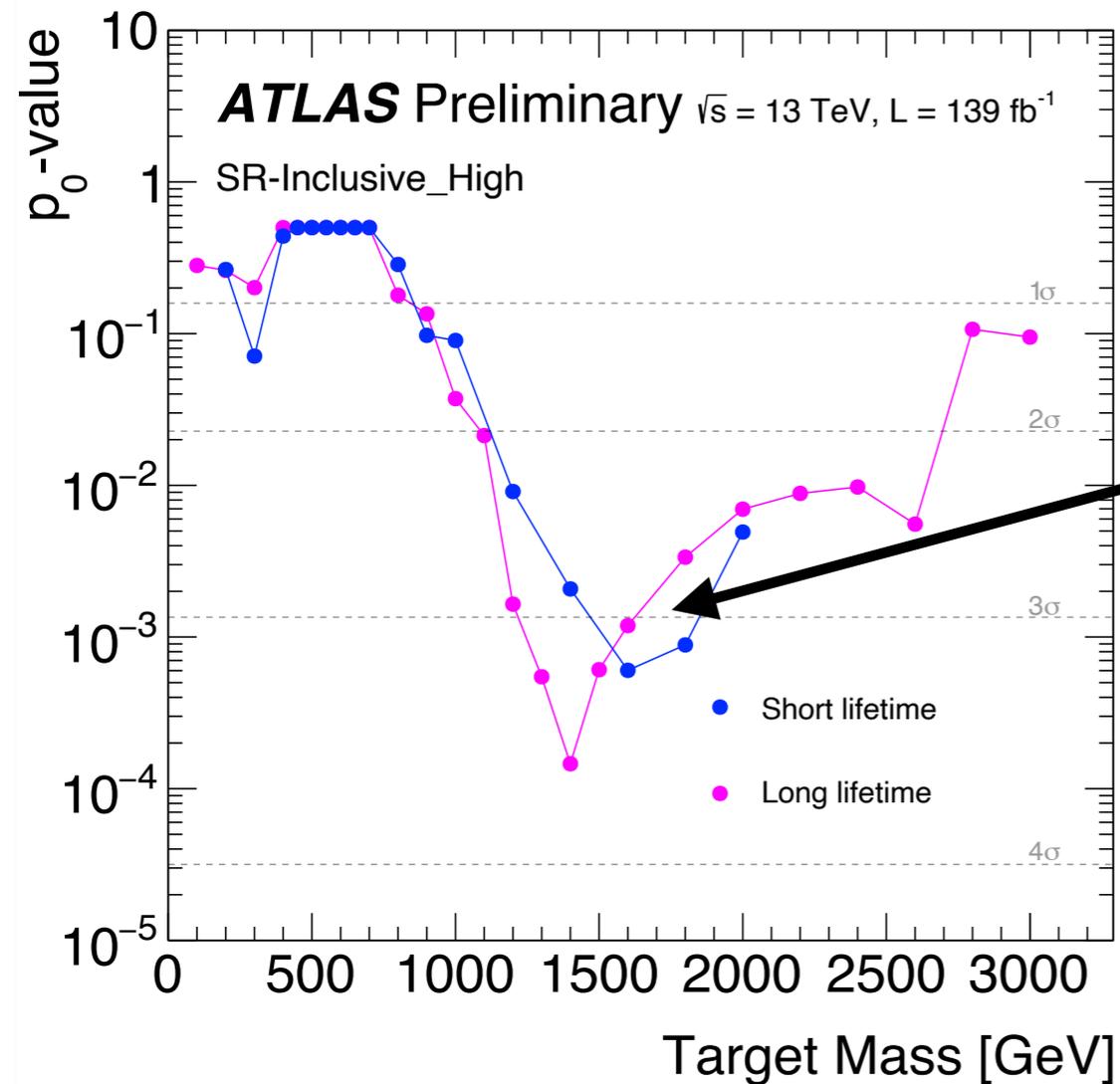


LHC – R-hadrons: ATLAS [1902.01636, 1808.04095 approximate reinterpretation];
 CMS [CMS-PAS-EXO-16-036, recasting from 1705.09292]

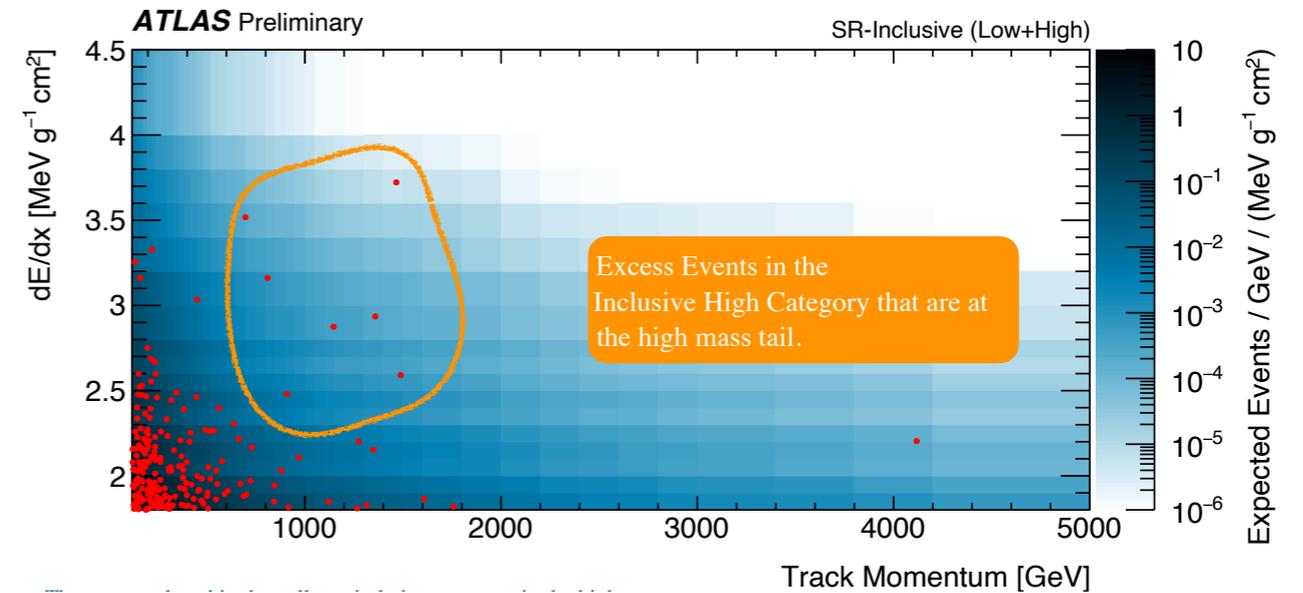
LHC – DT: ATLAS Disappearing-track search [1712.02118, recasting from 2002.12220, 7]

Recent excess in LLP searches

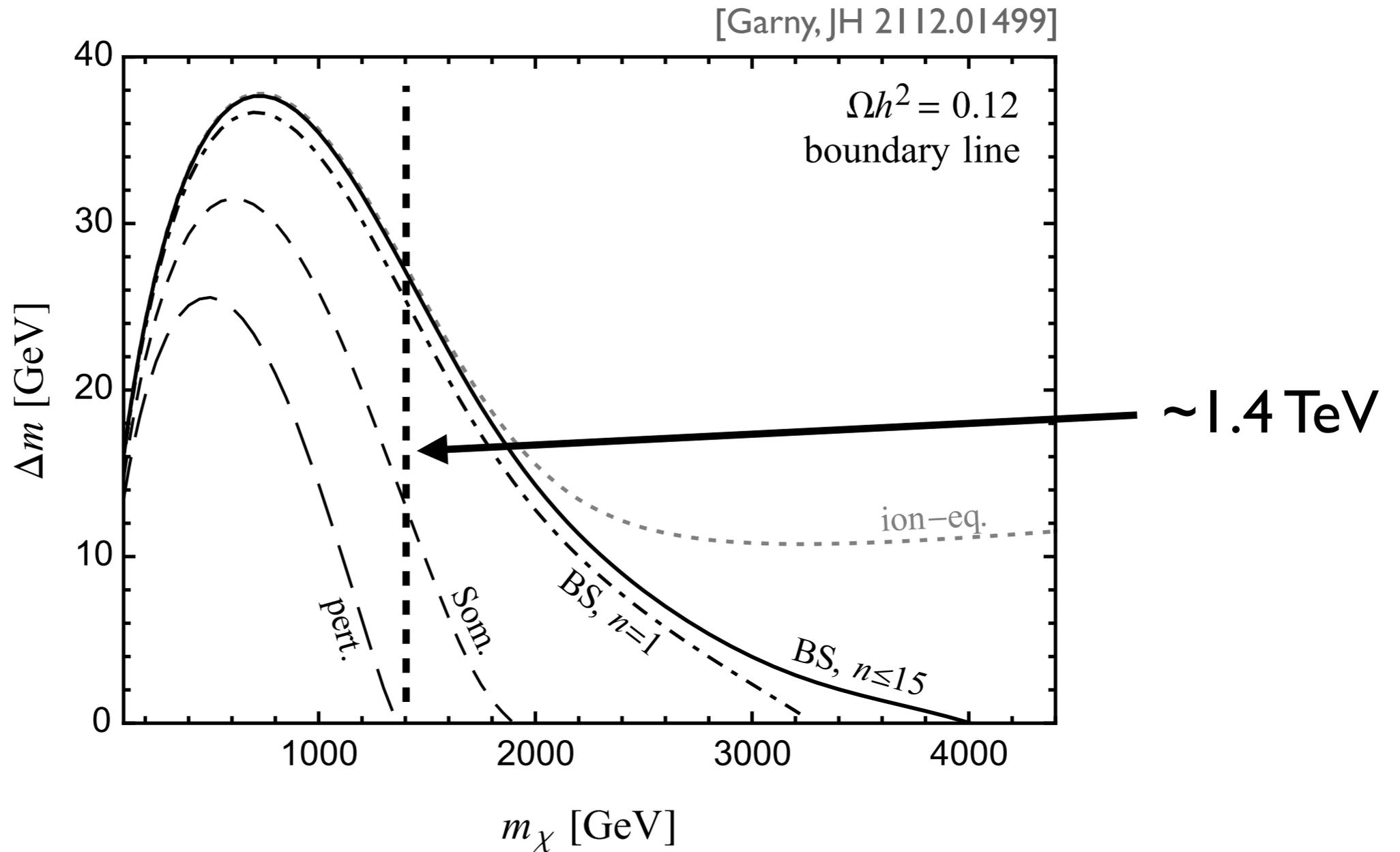
[ATLAS, Talk at Moriond 2022]



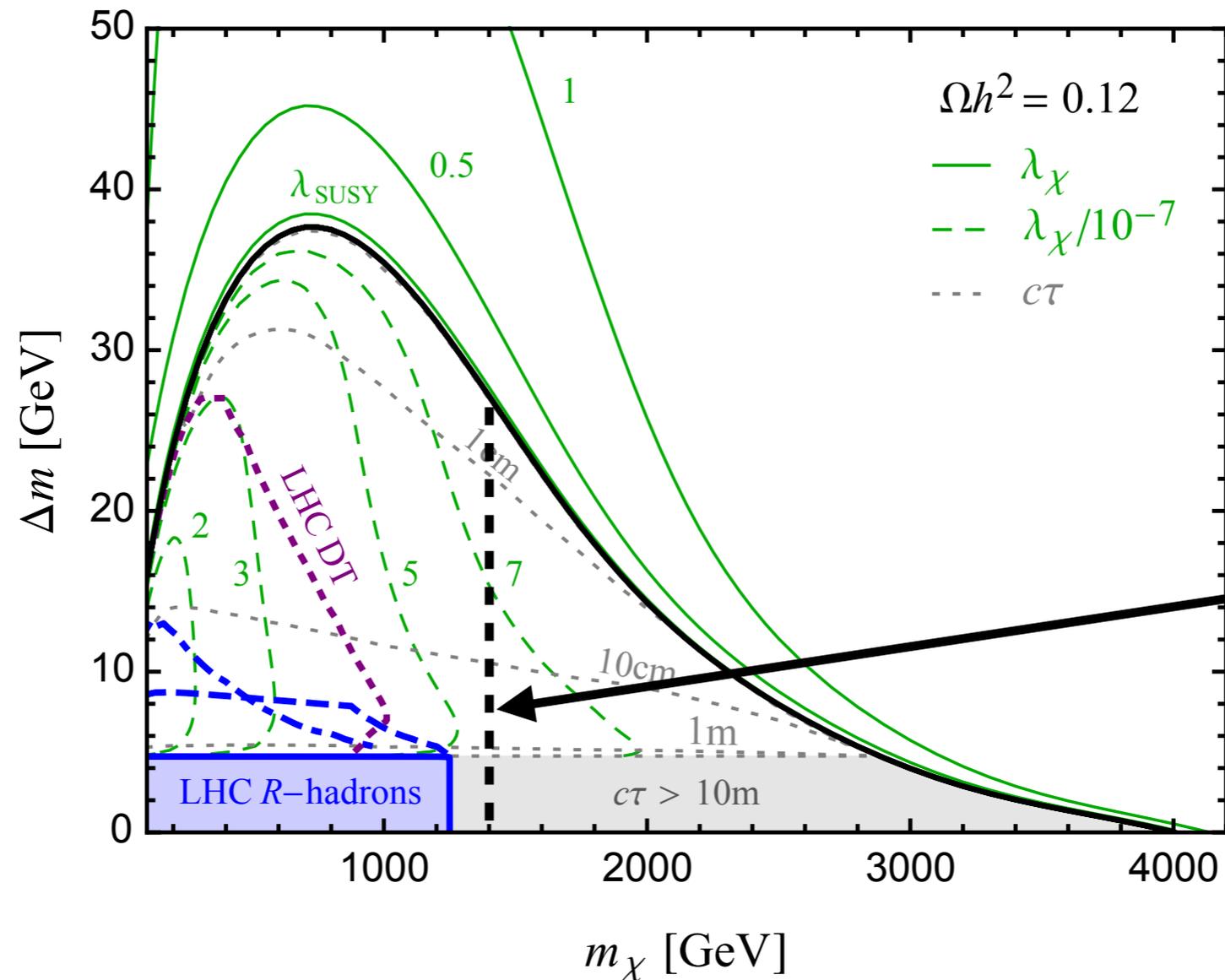
Heavy stable
charge particle
with mass $\sim 1.4 \text{ TeV}$



Boundary between WIMP region and conversion-driven region



Current LLP constraints



$\sim 1.4 \text{ TeV}$
 $\Rightarrow \sigma_s \sim 0.4 \text{ fb}$
 $\sim 50 \text{ events}$
 $\beta > 0.9 ??$

LHC – R-hadrons: ATLAS [1902.01636, 1808.04095 approximate reinterpretation];
 CMS [CMS-PAS-EXO-16-036, recasting from 1705.09292]

LHC – DT: ATLAS Disappearing-track search [1712.02118, recasting from 2002.12220, 7]

Summary

- t-channel mediator models: WIMP parameter space almost excluded
 - Remaining region: Conversion-driven freeze-out less explored terrain
 - Long-lived particles main signature
 $H \sim \Gamma$: Lifetimes naturally $O(1-100\text{cm})$
 - Freeze-out revisited: Bound states very relevant, higher excitation important at low energies
 - LLP parameter space significantly enlarged
 - Kinematics differs from targets of current searches
-