On-Shell Diagrams, Recursion Relations, Combinatorics

Jacob L. Bourjaily

Amplitudes 2022 Summer School Charles University, Prague, Czech Republic





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Prague, Czech Republic

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The Combinatorics and Geometry of On-Shell Physics



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Organization and Outline

- On-Shell, Tree-Level Recursion Relations for Scattering Amplitudes
- From On-Shell Physics to the (Positive) Grassmannian From the Bottom-Up:
 - (Combinatorially) Constructing and Computing On-Shell Functions
 - Asymptotic Symmetries of the S-Matrix: the Yangian

From the Top-Down:

- Grassmannian Geometry of (Generalized) Parke-Taylor 'Amplitudes'
- 3 The *Positroid* Stratification of the Grassmannian
 - A Combinatorial Stratification of the Grassmannian

Recall from Last Lecture: the BCFW Recursion Relations

Diagrams are characterized by 'm'—the number of "minus-helicity" gluons:

$$m \equiv 2n_B + n_W - n_I.$$

For the bridge terms, we have:

 $m_L + m_R = m + 1.$



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Recall the Example Amplitude Discussed: $\mathcal{A}_6^{(3)}$

The BCFW recursion relations realize an incredible fantasy: they **directly** produces the **Parke-Taylor** formula for all amplitudes with m=2, $\mathcal{A}_n^{(2)}$! And it generates **very concise** formulae for all other amplitudes—*e.g.* $\mathcal{A}_6^{(3)}$:

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Observations regarding recursed representations of scattering amplitudes:

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Is there any way to invariantly characterize the on-shell functions associated with on-shell diagrams?

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Combinatorial Characterization of On-Shell Diagrams

On-shell diagrams can be altered without changing their associated functions

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On-shell diagrams can be altered without changing their associated functions

- chains of equivalent three-particle vertices can be arbitrarily connected
- any four-particle 'square' can be drawn in its two equivalent ways



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These moves leave invariant a permutation defined by 'left-right paths': Starting from any leg a, turn:



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• *left* at each white vertex;



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- *left* at each white vertex;
- *right* at each blue vertex.



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left-right permutation σ

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left-right permutation σ

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 $\sigma: \begin{pmatrix} 1 & 2 & 3 & 4 & 5 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \\ 3 & 5 & 6 & 1 & 2 \end{pmatrix}$

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Notice that the merge and square moves leave the number of 'faces' of an on-shell diagram invariant.

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Combinatorial Characterization of On-Shell Diagrams

Notice that the merge and square moves leave the number of 'faces' of an on-shell diagram invariant. Diagrams with different numbers of faces can be related by 'reduction'—also known as 'bubble deletion':

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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Such factors of $d\alpha/\alpha$ arising from bubble deletion encode loop integrands!



Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Canonical Coordinates for Computing On-Shell Functions

Recall that attaching 'BCFW bridges' can lead to very rich on-shell diagrams.

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Recall that attaching 'BCFW bridges' can lead to very rich on-shell diagrams. Read the other way, we can 'peel-off' bridges and thereby decompose a permutation into transpositions according to $\sigma = (ab) \circ \sigma'$



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On-Shell, Tree-Level Recursion Relations for Scattering Amplitudes From On-Shell Physics to the (Positive) Grassmannian The Positroid Stratification of the Grassmannian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

> 'Bridge' Decomposition $1 \ 2 \ 3 \ 4 \ 5 \ 6$

 f_8 {7 8 3 10 5 6 }

Canonical Coordinates for Computing On-Shell Functions

There are many ways to decompose a permutation into transpositions—*e.g.*, always choose the first transposition $\tau \equiv (a b)$ such that $\sigma(a) < \sigma(b)$:

 $f_0 = \frac{d\alpha_1}{\alpha_1} \frac{d\alpha_2}{\alpha_2} \frac{d\alpha_3}{\alpha_3} \frac{d\alpha_4}{\alpha_4} \frac{d\alpha_5}{\alpha_5} \frac{d\alpha_6}{\alpha_6} \frac{d\alpha_7}{\alpha_7} \frac{d\alpha_8}{\alpha_8} f_8$

$$f_8 = \prod_{a=\sigma(a)+n} \left(\delta^4(\widetilde{\eta}_a) \delta^2(\widetilde{\lambda}_a) \right) \prod_{b=\sigma(b)} \left(\delta^2(\lambda_b) \right)$$

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$$C = \left(\begin{array}{cccc} \frac{1}{2} & \frac{2}{3} & \frac{4}{5} & \frac{5}{6} \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \end{array} \right)$$

$$f_{8} \left\{ 7 \ 8 \ 3 \ 10 \ 5 \ 6 \right\}$$

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$$= \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

$$C = \begin{pmatrix} \frac{1}{2} & \frac{3}{2} & \frac{4}{2} & \frac{5}{6} & \frac{6}{2} \\ \frac{1}{2} & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \end{pmatrix}$$

$$f_{8} \{7, 8, 3, 10, 5, 6\}$$

 f_8

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$$f_{0} = \frac{d\alpha_{8}}{\alpha_{8}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

$$C = \begin{pmatrix} \frac{1}{2} & \frac{3}{2} & \frac{4}{2} & \frac{5}{6} & \frac{6}{2} \\ \frac{1}{2} & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & \alpha_{8} \end{pmatrix}$$

$$(4 \, 6): \ c_{6} \mapsto c_{6} + \alpha_{8} \, c_{4}$$

$$F_{1} \{7 \ 8 \ 3 \ 6 \ 5 \ 10\} (4 \, 6)$$

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$$f_{0} = \frac{d\alpha_{1}}{\alpha_{1}} \frac{d\alpha_{2}}{\alpha_{2}} \frac{d\alpha_{3}}{\alpha_{3}} \frac{d\alpha_{4}}{\alpha_{4}} \frac{d\alpha_{5}}{\alpha_{5}} \frac{d\alpha_{6}}{\alpha_{6}} \frac{d\alpha_{7}}{\alpha_{7}} \frac{d\alpha_{8}}{\alpha_{8}} f_{8}$$

$$f_{0} = \frac{d\alpha_{7}}{\alpha_{7}} \frac{d\alpha_{8}}{\alpha_{8}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

$$C = \begin{pmatrix} \frac{1}{2} & \frac{2}{3} & \frac{3}{4} & \frac{5}{5} & \frac{6}{6} \\ 0 & 1 & 0 & \alpha_{7} & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & \alpha_{8} \end{pmatrix}$$

$$(24): c_{4} \mapsto c_{4} + \alpha_{7} c_{2}$$

$$f_{0} = \begin{pmatrix} \frac{1}{2} & \frac{3}{2} & \frac{4}{2} & \frac{5}{2} & \frac{6}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{3}{2} & \frac{4}{2} & \frac{5}{2} & \frac{6}{2} \\ \frac{1}{2} & \frac{1}{2} &$$

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$$f_{5} = \frac{d\alpha_{6}}{\alpha_{6}} \frac{d\alpha_{7}}{\alpha_{7}} \frac{d\alpha_{8}}{\alpha_{8}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

$$C = \begin{pmatrix} \frac{1}{2} & \frac{2}{3} & \frac{3}{4} & \frac{5}{6} & \frac{6}{2} \\ \frac{1}{2} & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & \alpha_{7} & \alpha_{6} \alpha_{7} & 0 \\ 0 & 0 & 0 & 1 & \alpha_{6} & \alpha_{8} \end{pmatrix}$$

$$f_{5} \{7 \ 6 \ 3 \ 5 \ 8 \ 10\} (45)$$

$$f_{6} \{7 \ 6 \ 3 \ 8 \ 5 \ 10\} (24)$$

$$f_{7} \{7 \ 8 \ 3 \ 6 \ 5 \ 10\} (24)$$

$$f_{8} \{7 \ 8 \ 3 \ 10 \ 5 \ 6\}$$

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Canonical Coordinates for Computing On-Shell Functions

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$$C = \begin{pmatrix} \frac{1}{2} & \frac{2}{\alpha_{5}} & \frac{3}{\alpha_{5}} & \frac{4}{\alpha_{5}} & \frac{5}{\alpha_{6}} & \frac{6}{\alpha_{7}} & \frac{6}{\alpha_{8}} \\ 0 & 1 & 0 & \alpha_{7} & \alpha_{6} \alpha_{7} & 0 \\ 0 & 0 & 0 & 1 & \alpha_{6} & \alpha_{8} \end{pmatrix}$$

$$f_{4} \{6 \ 7 \ 3 \ 5 \ 8 \ 10\} (12)$$

$$f_{5} \{7 \ 6 \ 3 \ 5 \ 8 \ 10\} (45)$$

$$f_{6} \{7 \ 6 \ 3 \ 8 \ 5 \ 10\} (45)$$

$$f_{7} \{7 \ 8 \ 3 \ 6 \ 5 \ 10\} (24)$$

$$f_{7} \{7 \ 8 \ 3 \ 6 \ 5 \ 10\} (46)$$

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$$f_{3} = \frac{d\alpha_{4}}{\alpha_{4}} \cdots \frac{d\alpha_{8}}{\alpha_{8}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

$$f_{3} = \frac{d\alpha_{4}}{\alpha_{4}} \cdots \frac{d\alpha_{8}}{\alpha_{8}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

$$f_{3} = \frac{d\alpha_{4}}{\alpha_{4}} \cdots \frac{d\alpha_{8}}{\alpha_{8}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

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$$f_{4} = \frac{d\alpha_{4}}{\alpha_{4}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times$$

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$$f_{2} = \frac{d\alpha_{3}}{\alpha_{3}} \cdots \frac{d\alpha_{8}}{\alpha_{8}} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

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$$f_{2} \{ 5 \ 6 \ 3 \ 7 \ 8 \ 10 \} (12)$$

$$f_{3} \{ 6 \ 5 \ 3 \ 7 \ 8 \ 10 \} (12)$$

$$f_{3} \{ 6 \ 5 \ 3 \ 7 \ 8 \ 10 \} (12)$$

$$f_{4} \{ 6 \ 7 \ 3 \ 5 \ 8 \ 10 \} (12)$$

$$f_{5} \{ 7 \ 6 \ 3 \ 5 \ 8 \ 10 \} (12)$$

$$f_{6} \{ 7 \ 6 \ 3 \ 8 \ 5 \ 10 \} (45)$$

$$f_{7} \{ 7 \ 8 \ 3 \ 6 \ 5 \ 10 \} (24)$$

$$f_{7} \{ 7 \ 8 \ 3 \ 6 \ 5 \ 10 \} (24)$$

$$f_{8} \{ 7 \ 8 \ 3 \ 10 \ 5 \ 6 \}$$

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$$f_{1} = \frac{d\alpha_{1}}{\alpha_{1} + \alpha_{3} + \alpha_{5}} \frac{d\alpha_{6}}{\alpha_{2} + \alpha_{5}} \frac{d\alpha_{6}}{\alpha_{4} + \alpha_{7}} \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

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$$f_{2} = \frac{d\alpha_{1}}{\alpha_{1} + \alpha_{3} + \alpha_{5}} \frac{d\alpha_{6}}{\alpha_{2} + \alpha_{7}} \frac{d\alpha_{6}}{\alpha_{6} + \alpha_{7}} \frac{d\alpha_{6}}{\alpha_{6} + \alpha_{7}} \frac{d\alpha_{6}}{\alpha_{6} + \alpha_{7}} \delta^{2 \times 3} \frac{d\alpha_{6}}{\alpha_{7} + \alpha_{7}} \frac{d\alpha_{7}}{\alpha_{6} + \alpha_{7}} \frac{d\alpha_{7}}{\alpha_{6} + \alpha_{7}} \frac{d\alpha_{7}}{\alpha_{6} + \alpha_{7}} \frac{d\alpha_{7}}{\alpha_{7} + \alpha_{7}} \frac{d\alpha_{7}}{\alpha_{6} + \alpha_{7}} \frac{d\alpha_{7}}{\alpha_{7} + \alpha_{7}} \frac{d\alpha_{7}}$$

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$$f_{1} = \frac{d\alpha_{1}}{\delta} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})$$

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Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

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$$f_{2} = \frac{d\alpha_{1}}{\delta} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{4$$

Amplitudes 2022 PhD Summer School

Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Canonical Coordinates for Computing On-Shell Functions



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'Bridge' Decomposition

 $\begin{array}{c} 2 & 5 & 4 & 5 & 0 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \\ 5 & 6 & 7 & 8 & 10 \end{array}$

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$$f = \int \frac{d\alpha_1}{\alpha_1} \wedge \dots \wedge \frac{d\alpha_d}{\alpha_d} \, \delta^{k \times 4} \big(C(\vec{\alpha}) \cdot \widetilde{\eta} \big) \delta^{k \times 2} \big(C(\vec{\alpha}) \cdot \widetilde{\lambda} \big) \delta^{2 \times (n-k)} \big(\lambda \cdot C(\vec{\alpha})^{\perp} \big)$$

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Measure-preserving diffeomorphisms leave the function invariant

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The Grassmannian Geometry of Scattering Amplitudes

Consider all the *iterated factorizations* of the Parke-Taylor Amplitude:

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$$\mathcal{A}_{4}^{(2)}(1,2,3,4):$$

$$\frac{\delta^{2\times4}(\lambda\cdot\tilde{\eta})\delta^{2\times2}(\lambda\cdot\tilde{\lambda})}{\langle 1\,2\rangle\langle 2\,3\rangle\langle 3\,4\rangle\langle 4\,1\rangle}$$

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\frac{\lambda\in G(2,4)\quad\lambda_{a}\in\mathbb{P}^{1}}{\left(\begin{array}{ccc}\lambda_{1}^{1}&\lambda_{2}^{1}&\lambda_{3}^{1}&\lambda_{4}^{1}\\\lambda_{1}^{2}&\lambda_{2}^{2}&\lambda_{3}^{2}&\lambda_{4}^{2}\end{array}\right)}
\end{array}$$

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\hline
\lambda\in G(2,4)\quad\lambda_{a}\in\mathbb{P}^{1}\\
\begin{pmatrix}\lambda_{1}^{1}\quad\lambda_{2}^{1}\quad\lambda_{3}^{1}\quad\lambda_{4}^{1}\\
\lambda_{1}^{2}\quad\lambda_{2}^{2}\quad\lambda_{3}^{2}\quad\lambda_{4}^{2}\end{pmatrix}
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\hline \lambda \in G(2,4) \quad \lambda_{a} \in \mathbb{P}^{1}\\
\left(\begin{array}{ccc}\lambda_{1}^{1} & \lambda_{2}^{1} & \lambda_{3}^{1} & \lambda_{4}^{1}\\ \lambda_{1}^{2} & \lambda_{2}^{2} & 0 & 0\end{array}\right)\\
\hline \lambda_{3} \longrightarrow \alpha_{3\,4}\lambda_{4}
\end{array}$$



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$$\lambda\in G(2,4)\quad\lambda_{a}\in\mathbb{P}^{1}$$

$$\begin{pmatrix}\lambda_{1}^{1}\quad\lambda_{2}^{1}\quad\lambda_{3}^{1}\quad\lambda_{4}^{1}\\\lambda_{1}^{2}\quad0\quad0\quad\lambda_{4}^{2}\end{pmatrix}$$

$$\lambda_{2}\longrightarrow\alpha_{2\,3}\lambda_{3}$$



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\left(\begin{array}{ccc}\lambda_{1}^{1}&\lambda_{2}^{1}&\lambda_{3}^{1}&\lambda_{4}^{1}\\0&\lambda_{2}^{2}&\lambda_{3}^{2}&0\end{array}\right)\\
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$$\frac{\delta^{2\times4}(\lambda\cdot\tilde{\eta})\delta^{2\times2}(\lambda\cdot\tilde{\lambda})}{\alpha_{12}\langle 23\rangle\alpha_{34}}$$

$$\lambda \in G(2,4) \quad \lambda_a \in \mathbb{P}^1$$

$$\begin{pmatrix} 0 & 0 & \lambda_3^1 & \lambda_4^1 \\ \lambda_1^2 & \lambda_2^2 & 0 & 0 \end{pmatrix}$$

$$\lambda_1 \longrightarrow \alpha_{12}\lambda_2$$

$$\lambda_3 \longrightarrow \alpha_{34}\lambda_4$$



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\frac{\lambda_{2}\longrightarrow\alpha_{2,3}\lambda_{3}}{\lambda_{4}\longrightarrow\alpha_{4,1}\lambda_{1}}
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$$\mathcal{\lambda}\in G(2,4) \quad \lambda_{a}\in\mathbb{P}^{1}$$

$$\begin{pmatrix}\lambda_{1}^{1}\ \lambda_{2}^{1}\ \lambda_{3}^{1}\ \lambda_{4}^{1}\\\lambda_{1}^{2}\ 0\ 0\ 0\end{pmatrix}$$

$$\lambda_{2}\longrightarrow\alpha_{23}\lambda_{3}$$

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\underline{\delta^{2\times4}(\lambda\cdot\tilde{\eta})\delta^{2\times2}(\lambda\cdot\tilde{\lambda})}{\langle 12\rangle\,\alpha_{2\,3}\,\,\alpha_{3\,4}}\\
\begin{array}{l}
\overline{\lambda\in G(2,4)\quad\lambda_{a}\in\mathbb{P}^{1}}\\
\left(\begin{array}{ccc}\lambda_{1}^{1}\quad\lambda_{2}^{1}\quad\lambda_{3}^{1}\quad\lambda_{4}^{1}\\\lambda_{1}^{2}\quad0\quad0&0\end{array}\right)\\
\end{array}$$

$$\begin{array}{l}
\lambda_{2}\longrightarrow\alpha_{2\,3}\lambda_{3}\\
\lambda_{3}\longrightarrow\alpha_{3\,4}\lambda_{4}
\end{array}$$



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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

The Grassmannian Geometry of Scattering Amplitudes

Consider all the *iterated factorizations* of the Parke-Taylor Amplitude:

$$\mathcal{A}_{4}^{(2)}(1,2,3,4):$$

$$\frac{\delta^{2\times4}(\lambda\cdot\tilde{\eta})\delta^{2\times2}(\lambda\cdot\tilde{\lambda})}{\langle 1\,2\rangle\langle 2\,3\rangle\langle 3\,1\rangle}$$

$$\lambda \in G(2,4) \quad \lambda_{a} \in \mathbb{P}^{1}$$

$$\begin{pmatrix} \lambda_{1}^{1} \quad \lambda_{2}^{1} \quad \lambda_{3}^{1} \quad 0 \\ \lambda_{1}^{2} \quad \lambda_{2}^{2} \quad \lambda_{3}^{2} \quad 0 \end{pmatrix}$$

$$\lambda_{4} \longrightarrow \alpha_{4\,1}\lambda_{1}$$

$$\lambda_{4} \longrightarrow \alpha_{4\,3}\lambda_{3}$$

$$(\therefore \lambda_{4} = 0)$$



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$$\lambda \in G(2,4) \quad \lambda_{a} \in \mathbb{P}^{1}$$

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The Grassmannian Geometry of Scattering Amplitudes

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$$\mathcal{A}_{4}^{(2)}(1,2,3,4):$$

$$\frac{\delta^{2\times4}(\lambda\cdot\tilde{\eta})\delta^{2\times2}(\lambda\cdot\tilde{\lambda})}{\langle 12\rangle\alpha_{23}}$$

$$\lambda \in G(2,4) \quad \lambda_{a} \in \mathbb{P}^{1}$$

$$\begin{pmatrix} \lambda_{1}^{1} \quad \lambda_{2}^{1} \quad \lambda_{3}^{1} \quad 0 \\ \lambda_{1}^{2} \quad 0 \quad 0 \quad 0 \end{pmatrix}$$

$$\lambda_{4} \longrightarrow \alpha_{41}\lambda_{1}$$

$$\lambda_{4} \longrightarrow \alpha_{43}\lambda_{3}$$

$$\lambda_{2} \longrightarrow \alpha_{23}\lambda_{3}$$



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Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

The Grassmannian Geometry of Scattering Amplitudes

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The Grassmannian Geometry of Scattering Amplitudes

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$$\begin{array}{l}
\mathcal{A}_{4}^{(2)}(1,2,3,4):\\
\frac{\delta^{2\times4}(\lambda\cdot\tilde{\eta})\delta^{2\times2}(\lambda\cdot\tilde{\lambda})}{\alpha_{12} \quad \langle 31\rangle}\\
\begin{array}{l}
\overline{\lambda \in G(2,4) \quad \lambda_{a} \in \mathbb{P}^{1}}\\
\left(\begin{array}{ccc}\lambda_{1}^{1} \quad \lambda_{2}^{1} \quad \lambda_{3}^{1} \quad 0\\
0 \quad 0 \quad \lambda_{3}^{2} \quad 0\end{array}\right)\\
\begin{array}{c}
\lambda_{4} \longrightarrow \alpha_{41}\lambda_{1}\\
\lambda_{4} \longrightarrow \alpha_{43}\lambda_{3}\\
\lambda_{1} \longrightarrow \alpha_{12}\lambda_{2}
\end{array}$$



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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

The Grassmannian Geometry of Scattering Amplitudes

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$$\mathcal{A}_{4}^{(2)}(1,2,3,4):$$

$$\frac{\delta^{2\times4}(\lambda\cdot\tilde{\eta})\delta^{2\times2}(\lambda\cdot\tilde{\lambda})}{\langle 12\rangle\,\alpha_{32}}$$

$$\lambda \in G(2,4) \quad \lambda_{a} \in \mathbb{P}^{1}$$

$$\begin{pmatrix} \lambda_{1}^{1} \quad \lambda_{2}^{1} \quad \lambda_{3}^{1} \quad 0\\ \lambda_{1}^{2} \quad 0 \quad 0 \quad 0 \end{pmatrix}$$

$$\lambda_{4} \longrightarrow \alpha_{41}\lambda_{1}$$

$$\lambda_{4} \longrightarrow \alpha_{43}\lambda_{3}$$

$$\lambda_{3} \longrightarrow \alpha_{32}\lambda_{2}$$



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$$\begin{pmatrix} \lambda_{1}^{1} \quad \lambda_{2}^{1} \quad \lambda_{3}^{1} \quad 0\\ \lambda_{1}^{2} \quad 0 \quad 0 \quad 0 \end{pmatrix}$$

$$\lambda_{4} \longrightarrow \alpha_{41}\lambda_{1}$$

$$\lambda_{4} \longrightarrow \alpha_{43}\lambda_{3}$$

$$\lambda_{3} \longrightarrow \alpha_{32}\lambda_{2}$$



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$$\lambda \in G(2,4) \quad \lambda_{a} \in \mathbb{P}^{1}$$

$$\begin{pmatrix} \lambda_{1}^{1} \quad \lambda_{2}^{1} \quad \lambda_{3}^{1} \quad 0\\ \lambda_{1}^{2} \quad 0 \quad 0 \quad 0 \end{pmatrix}$$

$$\lambda_{4} \longrightarrow \alpha_{41}\lambda_{1}$$

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\left(\begin{array}{ccc}\lambda_{1}^{1} & \lambda_{2}^{1} & \lambda_{3}^{1} & 0\\
\lambda_{1}^{2} & 0 & 0 & 0\end{array}\right)\\
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Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

Amplitudes with 2 negative-helicity gluons

$$\mathcal{A}_{n}^{(2)} = \frac{\delta^{2 \times 4} (\lambda \cdot \widetilde{\eta}) \delta^{2 \times 2} (\lambda \cdot \widetilde{\lambda})}{\langle 1 \, 2 \rangle \langle 2 \, 3 \rangle \langle 3 \, 4 \rangle \cdots \langle n \, 1 \rangle}$$

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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$$\boldsymbol{\lambda} \equiv \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \cdots & \lambda_n^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \cdots & \lambda_n^2 \end{pmatrix}$$

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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Parke-Taylor 'Amplitudes' and Grassmannian Residues

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

Amplitudes with *m* negative-helicity gluons:

$$\mathcal{A}_{n}^{(m)} \stackrel{?}{=} \frac{\delta^{m \times 4} (\mathbf{C} \cdot \widetilde{\eta}) \delta^{m \times 2} (\mathbf{C} \cdot \widetilde{\lambda})}{\langle 1 \cdots m \rangle \langle 2 \cdots m + 1 \rangle \cdots \langle n \cdots m - 1 \rangle}$$



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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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In order for momentum conservation, $\delta^{2\times 2}(\lambda \cdot \tilde{\lambda})$, to be part of the constraints, we must have that $C \supset \lambda$

Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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In order for momentum conservation, $\delta^{2\times 2}(\lambda \cdot \widetilde{\lambda})$, to be part of the constraints, we must have that $C \supset \lambda$, imposed via $\delta^{2\times (n-m)}(\lambda \cdot C^{\perp})$

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

Amplitudes with 3 negative-helicity gluons-e.g.,

$$\mathcal{A}_{6}^{(3)} \stackrel{?}{=} \frac{\delta^{3 \times 4} (C \cdot \widetilde{\eta}) \delta^{3 \times 2} (C \cdot \widetilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^{\perp})}{\langle 1 \, 2 \, 3 \rangle \langle 2 \, 3 \, 4 \rangle \langle 3 \, 4 \, 5 \rangle \langle 4 \, 5 \, 6 \rangle \langle 5 \, 6 \, 1 \rangle \langle 6 \, 1 \, 2 \rangle}$$

$$C \equiv \begin{pmatrix} c_1^1 & c_2^1 & c_3^1 & \cdots & c_n^1 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ c_1^m & c_2^m & c_3^m & \cdots & c_n^m \end{pmatrix}$$



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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Parke-Taylor 'Amplitudes' and Grassmannian Residues

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Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

Amplitudes with 3 negative-helicity gluons-e.g.,

$$\oint \frac{d\tau}{\langle 123 \rangle (\tau) \cdot \langle 234 \rangle (\tau) \cdot \langle 345 \rangle (\tau) \cdot \langle 456 \rangle (\tau) \cdot \langle 561 \rangle (\tau) \cdot \langle 612 \rangle (\tau) \rangle} \delta^{2 \times 2} (\lambda \cdot \widetilde{\lambda})$$

$$C \equiv \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \lambda_4^1 & \lambda_5^1 & \lambda_6^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ c_1^3 & c_2^3 & c_3^3 & c_4^3 & c_5^3 & c_6^3 \end{pmatrix}$$

$$\dim(C) = 3 \times 6 - 3 \times 3 = 9 = 8 + 1$$

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

Amplitudes with 3 negative-helicity gluons-e.g.,

$$\oint_{\substack{\langle 123\rangle=0}} \frac{d\tau}{\langle 123\rangle(\tau)\cdot\langle 234\rangle(\tau)\cdot\langle 345\rangle(\tau)\cdot\langle 456\rangle(\tau)\cdot\langle 561\rangle(\tau)\cdot\langle 612\rangle(\tau)}$$

$$C = \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \lambda_4^1 & \lambda_5^1 & \lambda_6^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ 0 & 0 & 0 & c_4^3 & c_5^3 & c_6^3 \end{pmatrix}$$

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$$\oint \frac{d\tau}{\langle 123\rangle(\tau) \cdot \langle 234\rangle(\tau) \cdot \langle 345\rangle(\tau) \cdot \langle 456\rangle(\tau) \cdot \langle 561\rangle(\tau) \cdot \langle 612\rangle(\tau)} \begin{pmatrix} \varphi \\ \langle 123\rangle(\tau) \cdot \langle 234\rangle(\tau) \cdot \langle 345\rangle(\tau) \cdot \langle 456\rangle(\tau) \cdot \langle 561\rangle(\tau) \cdot \langle 612\rangle(\tau) \\ \langle 123\rangle(\tau) \cdot \langle 234\rangle(\tau) \cdot \langle 345\rangle(\tau) \cdot \langle 456\rangle(\tau) \cdot \langle 561\rangle(\tau) \cdot \langle 612\rangle(\tau) \\ \end{pmatrix}$$

$$C \equiv \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \lambda_4^1 & \lambda_5^1 & \lambda_6^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 & \lambda_6^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 & \lambda_6^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 & \lambda_6^2 & \lambda_6^2 \\ \lambda_1^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 & \lambda_$$



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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

Amplitudes with 3 negative-helicity gluons-e.g.,

$$\oint \frac{d\tau}{\langle 123\rangle(\tau) \cdot \langle 234\rangle(\tau) \cdot \langle 345\rangle(\tau) \cdot \langle 456\rangle(\tau) \cdot \langle 561\rangle(\tau) \cdot \langle 612\rangle(\tau)} \begin{pmatrix} 1 & | \tau \\ \bullet & | \\ \bullet & | \\ \bullet & | \\ 0 \end{pmatrix} C = \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \lambda_4^1 & \lambda_5^1 & \lambda_6^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ \lambda_1^3 & c_2^3 & 0 & 0 & 0 & c_6^3 \end{pmatrix} \qquad (1) \qquad (1) \qquad (2) \qquad (3)$$



Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Parke-Taylor 'Amplitudes' and Grassmannian Residues

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$$\oint_{\substack{\langle 561\rangle = 0}} \frac{d\tau}{\langle 123\rangle(\tau) \cdot \langle 234\rangle(\tau) \cdot \langle 345\rangle(\tau) \cdot \langle 456\rangle(\tau) \cdot \langle 561\rangle(\tau) \cdot \langle 612\rangle(\tau)}$$



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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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Parke-Taylor 'Amplitudes' and Grassmannian Residues

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$$(1) \Leftrightarrow \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \lambda_4^1 & \lambda_5^1 & \lambda_6^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ 0 & 0 & 0 & [56] & [64] & [45] \end{pmatrix}$$

$$(1) \qquad (1) \qquad (2) \qquad (6) \qquad (6$$

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Parke-Taylor 'Amplitudes' and Grassmannian Residues

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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$$\frac{1}{\langle 23 \rangle [56]}$$

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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$$(1) \Leftrightarrow \begin{pmatrix} \lambda_{1}^{1} & \lambda_{2}^{1} & \lambda_{3}^{1} & \lambda_{4}^{1} & \lambda_{5}^{1} & \lambda_{6}^{1} \\ \lambda_{1}^{2} & \lambda_{2}^{2} & \lambda_{3}^{2} & \lambda_{4}^{2} & \lambda_{5}^{2} & \lambda_{6}^{2} \\ 0 & 0 & 0 & [56] & [64] & [45] \end{pmatrix}$$

$$\frac{1}{\langle 23 \rangle [56] [6|(5+4)|3 \rangle}$$

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Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics
Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

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$$(1) \Leftrightarrow \begin{pmatrix} \lambda_{1}^{1} & \lambda_{2}^{1} & \lambda_{3}^{1} & \lambda_{4}^{1} & \lambda_{5}^{1} & \lambda_{6}^{1} \\ \lambda_{1}^{2} & \lambda_{2}^{2} & \lambda_{3}^{2} & \lambda_{4}^{2} & \lambda_{5}^{2} & \lambda_{6}^{2} \\ 0 & 0 & 0 & | [56] & [64] & [45] \end{pmatrix}$$

$$\frac{1}{\langle 23 \rangle [56] [6|(5+4)|3 \rangle s_{456}}$$

Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

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$$(1) \Leftrightarrow \begin{pmatrix} \lambda_{1}^{1} \\ \lambda_{2}^{1} \\ \lambda_{2}^{2} \\ \lambda_{3}^{2} \\ \lambda_{3}^{2} \\ \lambda_{4}^{2} \\ \lambda_{2}^{2} \\ \lambda_{3}^{2} \\ \lambda_{4}^{2} \\ \lambda_{5}^{2} \\ \lambda_{5}^{2} \\ \lambda_{6}^{2} \\ \lambda_{7}^{2} \\ \lambda_{7}^{2$$

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Parke-Taylor 'Amplitudes' and Grassmannian Residues

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$$(1) \Leftrightarrow \begin{pmatrix} \lambda_{1}^{1} & \lambda_{2}^{1} & \lambda_{3}^{1} & \lambda_{4}^{1} & \lambda_{5}^{1} & \lambda_{6}^{1} \\ \lambda_{1}^{2} & \lambda_{2}^{2} & \lambda_{3}^{2} & \lambda_{4}^{2} & \lambda_{5}^{2} & \lambda_{6}^{2} \\ 0 & 0 & [56] & [64] & [45] \end{pmatrix}$$

$$\frac{1}{\langle 23 \rangle [56] [6|(5+4)|3 \rangle_{s_{4}56} \langle 1|(6+5)|4] [45] \langle 12 \rangle}$$

Parke-Taylor 'Amplitudes' and Grassmannian Residues

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$$(1) \Leftrightarrow \begin{pmatrix} \lambda_{1}^{1} & \lambda_{2}^{1} & \lambda_{3}^{1} & \lambda_{4}^{1} & \lambda_{5}^{1} & \lambda_{6}^{1} \\ \lambda_{1}^{2} & \lambda_{2}^{2} & \lambda_{3}^{2} & \lambda_{4}^{2} & \lambda_{5}^{2} & \lambda_{6}^{2} \\ 0 & 0 & 0 & [56] & [64] & [45] \end{pmatrix}$$

$$\frac{1}{\langle 23 \rangle [56] [6|(5+4)|3 \rangle_{s_{4}56} \langle 1|(6+5)|4] [45] \langle 12 \rangle}$$

Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

Parke-Taylor 'Amplitudes' and Grassmannian Residues

Recall the natural desire to generalize the Parke-Taylor formula according to:

Amplitudes with 3 negative-helicity gluons-e.g.,

$$\oint_{\langle 123\rangle=0} \frac{d\tau \quad (\langle 246\rangle^4 \, \tilde{\eta}_2^4 \tilde{\eta}_4^2 \tilde{\eta}_6^4 + \dots \,) \delta^{2\times 2} (\lambda \cdot \tilde{\lambda})}{\langle 123\rangle(\tau) \cdot \langle 234\rangle(\tau) \cdot \langle 345\rangle(\tau) \cdot \langle 456\rangle(\tau) \cdot \langle 561\rangle(\tau) \cdot \langle 612\rangle(\tau)}$$

$$(1) \Leftrightarrow \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \lambda_4^1 & \lambda_5^1 & \lambda_6^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ 0 & 0 & 0 & [56] & [64] & [45] \end{pmatrix}$$

$$\frac{\langle 2|(4+6)|5|^4}{\langle 23\rangle [56][6|(5+4)|3\rangle s_{456}\langle 1|(6+5)|4][45]\langle 12\rangle}$$



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Prague, Czech Republic Part II: On-

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Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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$$(1) \Leftrightarrow \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \lambda_3^1 & \lambda_4^1 & \lambda_5^1 & \lambda_6^1 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \lambda_4^2 & \lambda_5^2 & \lambda_6^2 \\ 0 & 0 & 0 & [5\,6] & [6\,4] & [4\,5] \end{pmatrix} \Leftrightarrow f_{\{3,5,6,7,8,10\}} (!)$$

$$\mathcal{A}_{6}^{(3)}(+,-,+,-,+,-) = (1+r^{2}+r^{4})\frac{\langle 2|(4+0)|3|}{\langle 2|3\rangle[5|6|(5+4)|3\rangle_{54|5|6|}(6+5)|4|[4|5|\langle 1|2\rangle|}$$

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Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

Canonical Coordinates and Combinatorial Computation Asymptotic Symmetries of the S-Matrix: the Yangian Grassmannian Geometry of Generalized Parke-Taylor 'Amplitudes'

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$$\mathcal{A}_{6}^{(3)}(+, -, +, -, +, -) = (1) + (3) + (5)$$

$$\overset{6}{4}$$

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$$\overset{6}{4}$$

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$$\overset{6}{4}$$

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Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

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$$\mathcal{A}_{6}^{(3)}(+, -, +, -, +, -) = (1) + (3) + (5)$$

$$\overset{6}{4}$$

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$$\overset{6}{4}$$

$$\overset{1}{4}$$

$$\overset{6}{4}$$

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Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

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Amplitudes with *m* negative-helicity gluons:

$$\mathcal{A}_{n}^{(m)} = \oint \frac{d^{m \times n} C}{\operatorname{vol}(GL(m))} \frac{\delta^{m \times 4} (C \cdot \widetilde{\eta}) \delta^{m \times 2} (C \cdot \widetilde{\lambda}) \delta^{2 \times (n-m)} (\lambda \cdot C^{\perp})}{\langle 1 \ 2 \ \cdots \ m \rangle \langle 2 \ 3 \ \cdots \ m+1 \rangle \cdots \langle n \ 1 \ \cdots \ m-1 \rangle}$$

$$C \equiv \begin{pmatrix} c_1^1 & c_2^1 & c_3^1 & \cdots & c_n^1 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ c_1^m & c_2^m & c_3^m & \cdots & c_n^m \end{pmatrix}$$

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Classifying the Iterated Residues of the Volume-Form $\mathcal{L}_{n,k}$

A complete, GL(k)-invariant description of any contour of $\mathcal{L}_{n,k}$ would be a list of all the ranks of spaces spanned by all consecutive chains of columns.

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Geometry and Combinatorics of the Positroid Stratification

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Amplitudes 2022 PhD Summer School

Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

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Ranks of Consecutive Chains of Columns					
1 … 1					

A (10) < A (10) < A (10) </p>

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Ranks of Consecutive Chains of Columns					
1 2					
1 ··· 1					

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Ranks of Consecutive Chains of Columns					
1 <i>n</i> -2 ;					
1 2					
11 1					

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Ranks of Consecutive Chains of Columns				
$ 1\cdots n-1 $				
$ 1\cdots n-2 $				
:				
1 2				
1 ··· 1				

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Ranks of Consecutive Chains of Columns
$ 1 \cdots n $
1···n–1
1···n-2
:
1 ··· 2
1 1

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	$ 2 \cdots n+1 $				
$ 1 \cdots n $	$ 2 \cdots n $				
$ 1 \cdots n-1 $	$ 2 \cdots n-1 $				
$ 1 \cdots n-2 $	2···· <i>n</i> -2				
:	:				
1 ··· 2	2 2				
1 ··· 1					

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Ranks of Consecutive Chains of Columns					
				$ n \cdots 2n $	
	$ 2\cdots n+1 $:	
$ 1 \cdots n $	$ 2 \cdots n $		$ n-1 \cdots n $	$ n \cdots n $	
$ 1 \cdots n-1 $	$ 2 \cdots n-1 $. · *	$ n-1\cdots n-1 $		
$ 1 \cdots n-2 $	$ 2 \cdots n-2 $	· · '			
:	÷	· · [·]			
$ 1 \cdots 2 $	2 2				
$ 1 \cdots 1 $					

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Ranks of	f Consecu	utive	Chains of	Columns
$ 1 \cdots 2n $	$ 2 \cdots 2n $	····	$ n-1\cdots 2n $	$ n \cdots 2n $
:	•	• •	:	:
$ 1 \cdots n $	$ 2 \cdots n $	[.]	$ n-1 \cdots n $	$ n \cdots n $
$ 1 \cdots n-1 $	$ 2 \cdots n-1 $	· · '	$ n-1\cdots n-1 $	
$ 1 \cdots n-2 $	$ 2 \cdots n-2 $			
÷	÷			
$ 1 \cdots 2 $	2 2			
$ 1 \cdots 1 $				

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Ranks o	f Consec	utive	Chains of	Columns
$ 1 \cdots 2n $	$ 2 \cdots 2n $	 	$ n-1\cdots 2n $	$ n \cdots 2n $
$ 1 \cdots n $	$ 2 \cdots n $		$ n-1 \cdots n $	$ n \cdots n $
$ 1 \cdots n-1 $	$ 2 \cdots n-1 $.·*	$ n-1\cdots n-1 $	
$ 1 \cdots n-2 $	$ 2 \cdots n-2 $	·.'		
÷	:	[.]		
1 ··· 2	2 2			
1 1				

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Ranks of Consecutive Chains of Columns						
k	k		k	k		
÷		[.]	:	÷		
k	$ 2 \cdots n $	·	$ n-1 \cdots n $	$ n \cdots n $		
$ 1 \cdots n-1 $	$ 2 \cdots n-1 $	[*]	$ n-1\cdots n-1 $			
$ 1 \cdots n-2 $	$ 2 \cdots n-2 $					
÷	÷					
$ 1 \cdots 2 $	2 2					
$ 1 \cdots 1 $						

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Classifying the Iterated Residues of the Volume-Form $\mathcal{L}_{n,k}$

A complete, GL(k)-invariant description of any contour of $\mathcal{L}_{n,k}$ would be a list of all the ranks of spaces spanned by all consecutive chains of columns. Pierre Deligne suggested organizing this data in the following way:

- Let us use $|a \cdots b|$ to denote the rank of the space spanned by columns $\{c_a, \ldots, c_b\}$
- Write the rank $|a \cdots b|$ in the a^{th} column and b^{th} row (from the bottom)
- for any generic configuration, $|a \cdots a + n \cdots| = k$ for all a
- let us conventionally declare $|a \cdots a-1| \rightarrow 0$

Ranks of Consecutive Chains of Columns					
k	k		k	k	
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k	$ 2 \cdots n $		$ n-1 \cdots n $	$ n \cdots n $	
$ 1 \cdots n-1 $	$ 2 \cdots n-1 $	· · [·]	$ n-1\cdots n-1 $		
$ 1 \cdots n-2 $	$ 2 \cdots n-2 $	·			
÷	÷				
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1 1					

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÷			:	÷
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$ 1 \cdots n-1 $	$ 2 \cdots n-1 $	·.'	$ n-1\cdots n-1 $	$ n \cdots n-1 $
$ 1 \cdots n-2 $	$ 2 \cdots n-2 $		$ n-1\cdots n-2 $	
÷	÷			
$ 1 \cdots 2 $	2 2			
$\begin{array}{ccc} \ 1 \ \cdots \ 1 \\ \ 1 \ \cdots \ 0 \end{array}$	2 1			

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÷	÷			
$ 1 \cdots 2 $	2 2	·		
$egin{array}{ccc} 1 & \cdots & 1 \\ 0 \end{array}$	0			

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Geometry and Combinatorics of the Positroid Stratification

Exempli Gratia: Iterated Residues of $\mathcal{L}_{8,4}$

Take a generic $C \in G(4, 8), c_a \in \mathbb{P}^3$

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Prague, Czech Republic Part II: On-Shell Diagrams, Recursion Relations, and Combinatorics

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Geometry and Combinatorics of the Positroid Stratification

The Positroid Stratification of the Grassmannian

The dimensionality of the configuration is encoded by its permutation label as follows:



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The dimensionality of the configuration is encoded by its permutation label as follows:



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The dimensionality of the configuration is encoded by its permutation label as follows:

Dimensionality of a configuration labelled by σ :

$$\dim(\sigma) = \left(\sum_{a} |a \cdots \sigma(a)|\right) - k^2$$



For each column *a*, there is a unique, **nearest** column $b \ge a$ such that $a \in \text{span}\{a+1, \dots, b-1, b\}$



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$$\partial \sigma \equiv \left\{ \sigma' = (a b) \circ \sigma | \dim(\sigma') = \dim(\sigma) - 1 \right\}$$

Geometric meaning of the permutation $\sigma(a) = b$

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 $\{5,4,8,7,11,10,9,14\}$

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GRASSMANNIAN **GEOMETRY OF** SCATTERING AMPLITUDES NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV

ALEXANDER **POSTNIKOV** JAROSLAV **TRNKA**





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