



Chirally-Enhanced Muon $g - 2$ and Its Implications to Higgs-Related Observables

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Based on

Phys. Rev. Lett. 26, 191801 (2021), Phys. Rev. D 104, 055033 (2021), Phys. Rev. D 104, L091301 (2021)
w/ Radovan Dermisek and Navin McGinnis

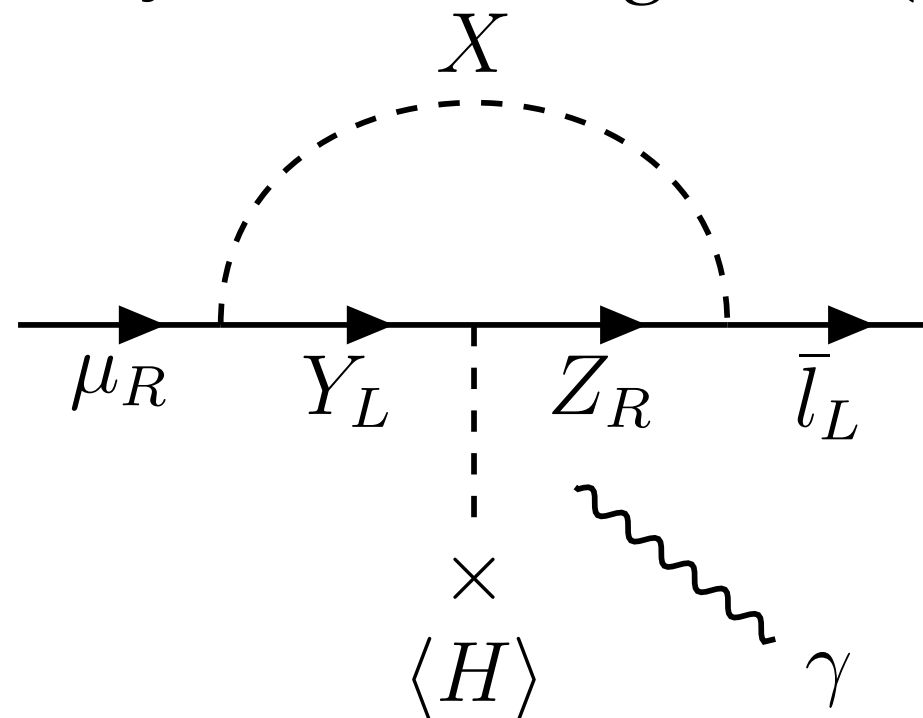


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Chiral Enhancement in $g - 2$

- Chirally-enhanced $g - 2$ (Δa_μ) can be generated by new particles coupling to the muon and Higgs inside the loop:



$$\Delta a_\mu \simeq \frac{1}{16\pi^2} \left(\frac{m_\mu v \lambda_{NP}^3}{m_{NP}^2} \right)$$

Muon g-2 Collab., Phys. Rev. Lett. 126, 141801 (2021)

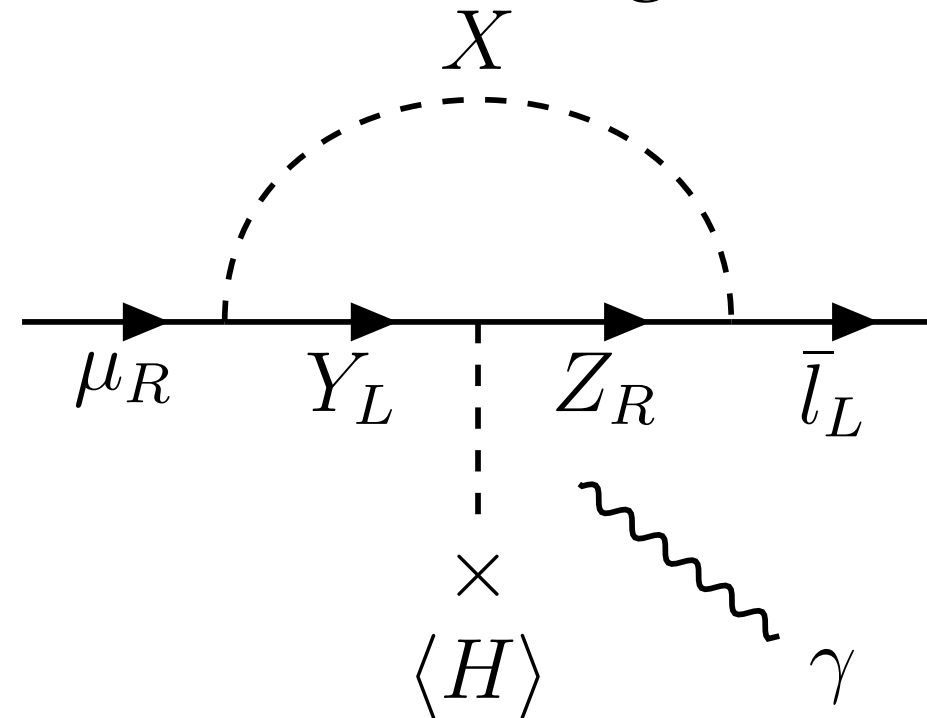
$$\Delta a_\mu^{exp} = (2.51 \pm 0.59) \times 10^{-9} \text{ requires } \sim 10 \text{ (50) TeV scale for couplings } 1 \text{ } (\sqrt{4\pi})$$

- We can do this by working in a type-II two-Higgs-Doublet Model (H_u and H_d couple exclusively to the up- and down-type sectors (2HDM-II)) extended with *vectorlike leptons (VLs)*
- X can be SM bosons or new *Higgses* and Y, Z are the new *VLs*



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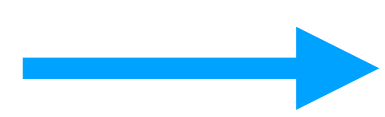
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	l_L	μ_R	H_u	H_d	$L_{L,R}$	$E_{L,R}$
$SU(2)_L$	2	1	2	2	2	1
$U(1)_Y$	$-\frac{1}{2}$	-1	$-\frac{1}{2}$	$\frac{1}{2}$	$-\frac{1}{2}$	-1
Z_2	$+$	$-$	$+$	$-$	$+$	$-$



$$\mathcal{L} = -y_\mu \bar{l}_L \mu_R H_d - \lambda_E \bar{l}_L E_R H_d - \lambda_L \bar{L}_L \mu_R H_d - \lambda \bar{L}_L E_R H_d - \bar{\lambda} H_d^\dagger \bar{E}_L L_R - M_L \bar{L}_L L_R - M_E \bar{E}_L E_R + h.c.$$

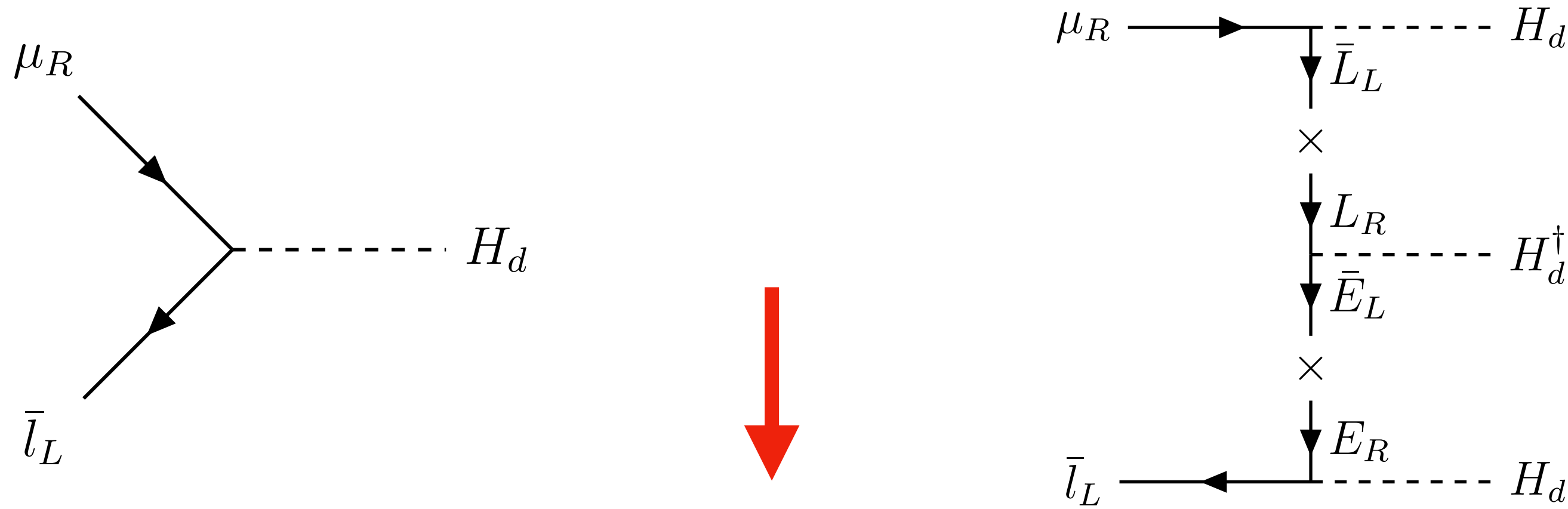
$$\langle H_d^0 \rangle = v_d = v \cos \beta \quad \langle H_u^0 \rangle = v_u = v \sin \beta$$

$$v = 174 \text{ GeV} \quad v_u/v_d = \tan \beta$$



Tree-level Mixing

- For $M_{L,E} \gg v$, the heavy VLs *modify* m_μ and y_μ since the full Lagrangian reduces to



Necessarily modifies
 $h \rightarrow \mu^+ \mu^-$

$$\mathcal{L} = -y_\mu \bar{l}_L \mu_R H_d - \underbrace{\left(\frac{\lambda_L \lambda_E \bar{\lambda}}{M_L M_E} \right)}_{m_\mu^{LE}/v_d^3} \bar{l}_L \mu_R H_d (H_d^\dagger H_d) + h.c.$$

$$m_\mu \simeq y_\mu v_d + m_\mu^{LE} \longrightarrow \lambda_{\mu\mu}^h \simeq (m_\mu + 2m_\mu^{LE})/v \longrightarrow R_\mu \equiv \frac{BR(h \rightarrow \mu^+ \mu^-)}{BR(h \rightarrow \mu^+ \mu^-)_{SM}} = \left(1 + 2 \frac{m_\mu^{LE}}{m_\mu} \right)^2$$

m_μ^{LE} is interpreted to be the muon's mass *if* $y_\mu = 0$

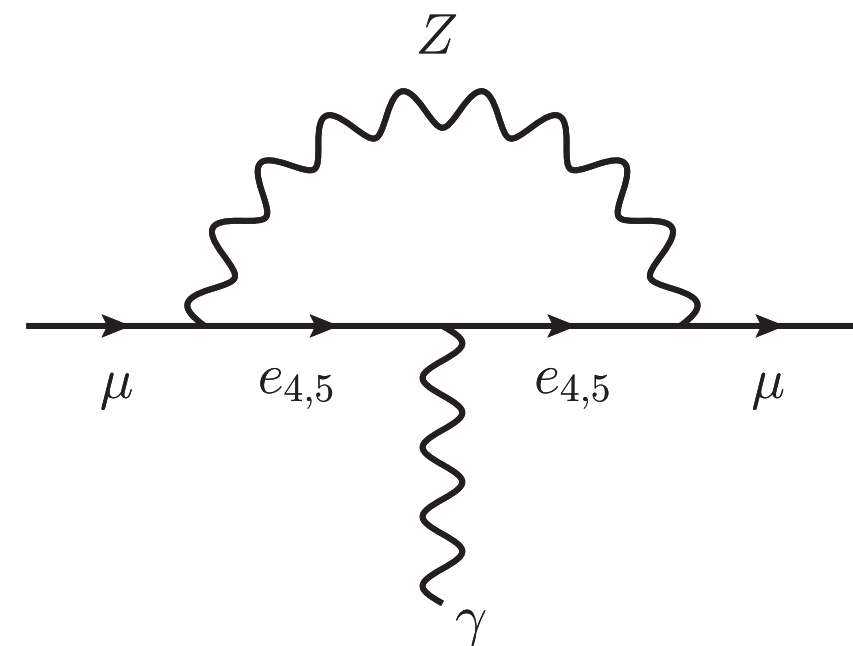
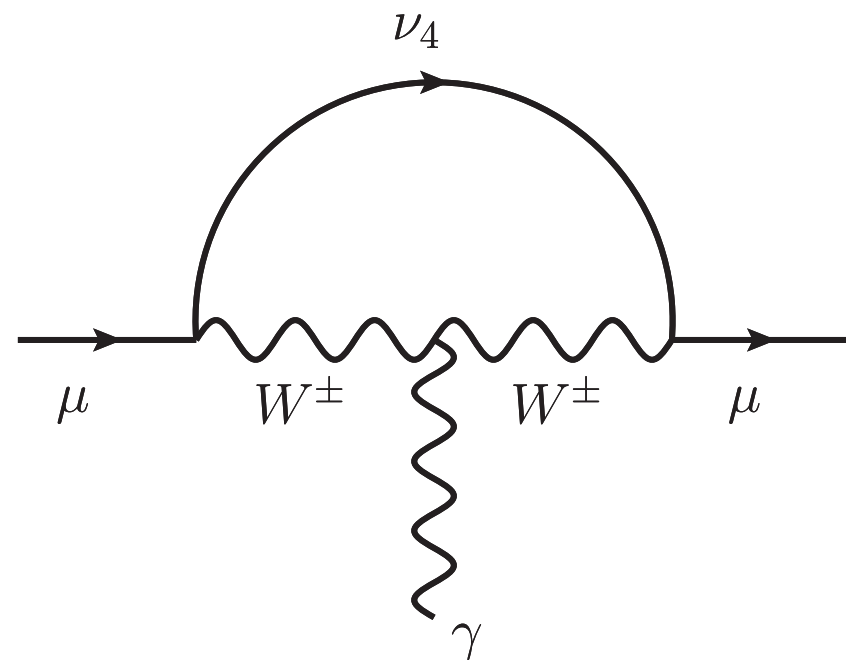


The Anomalous Magnetic Moment

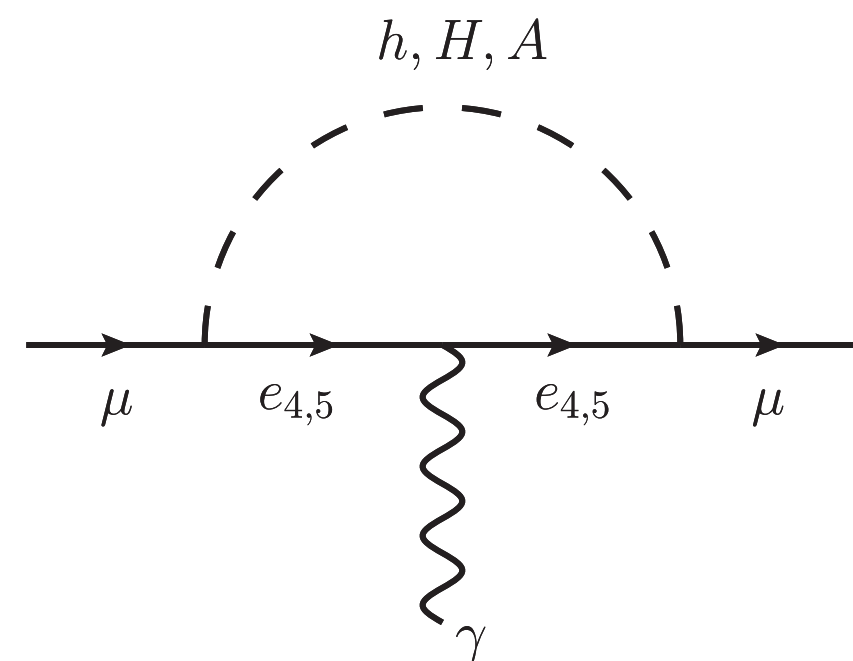
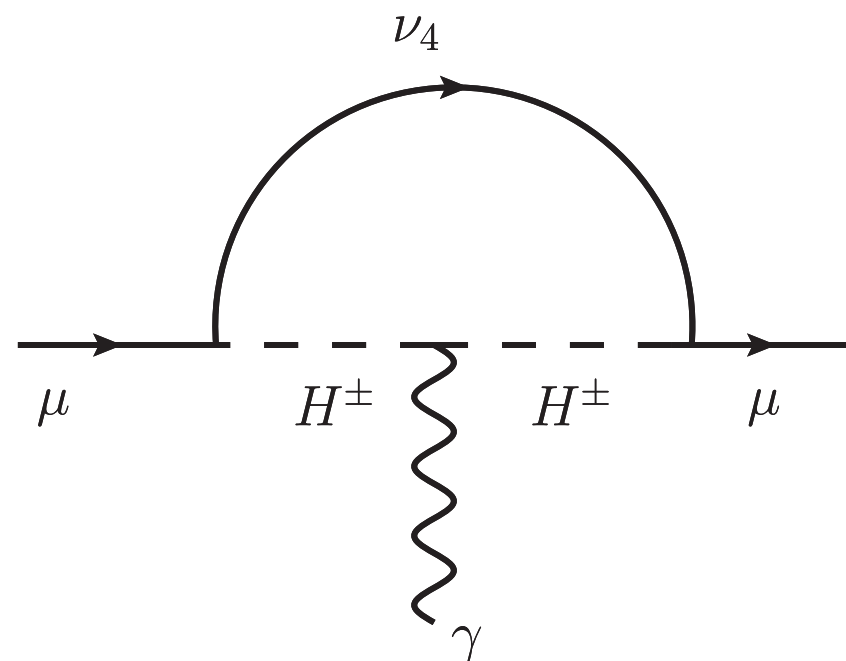
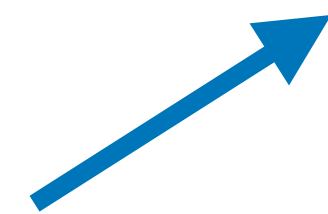
- Rotating the Lagrangian to the physical basis via bi-unitary transformations

$$(\bar{\mu}_L, \bar{L}_L^-, \bar{E}_L) \begin{pmatrix} y_\mu v_d & 0 & \lambda_E v_d \\ \lambda_L v_d & M_L & \lambda v_d \\ 0 & \bar{\lambda} v_d & M_E \end{pmatrix} \begin{pmatrix} \mu_R \\ L_R^- \\ E_R \end{pmatrix} \Rightarrow U_L^{e\dagger} M_e U_R^e = \text{diag} (m_{e_2}, m_{e_4}, m_{e_5})$$

- We can calculate Δa_μ in the mass eigenstate basis and take the heavy mass limit $M_{L,E} \sim m_{H,A,H^\pm} \gg m_Z$



$$\Delta a_\mu = -\frac{1}{16\pi^2} \left(\frac{m_\mu m_\mu^{LE}}{v^2} \right) (1 + \tan^2 \beta)$$



- Heavy Higgses (H, A, H^\pm) give additional $\tan^2 \beta$ enhancement compared to the (Z, W, h) mediators
- Notice that m_μ^{LE} is present. We'll see why later!

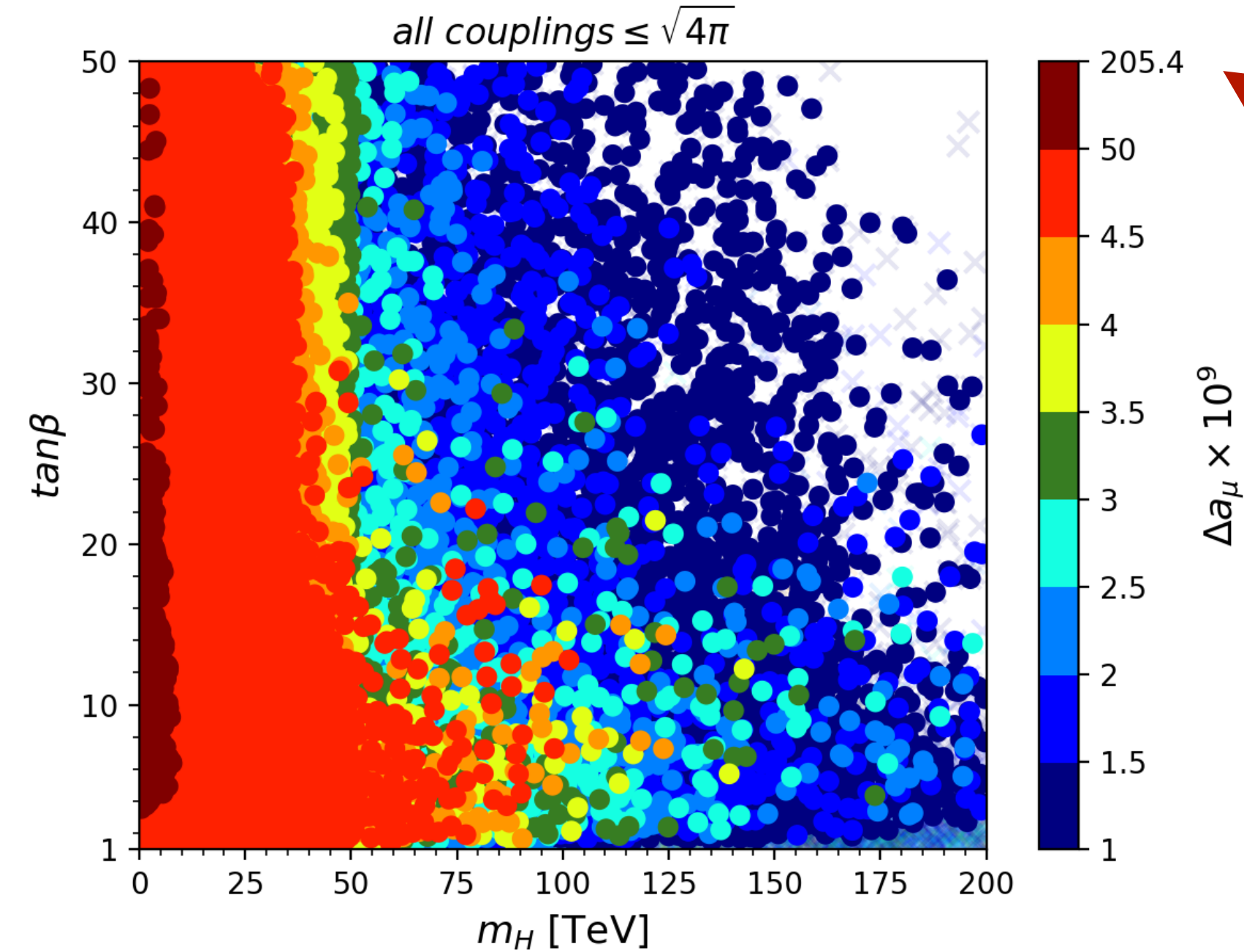
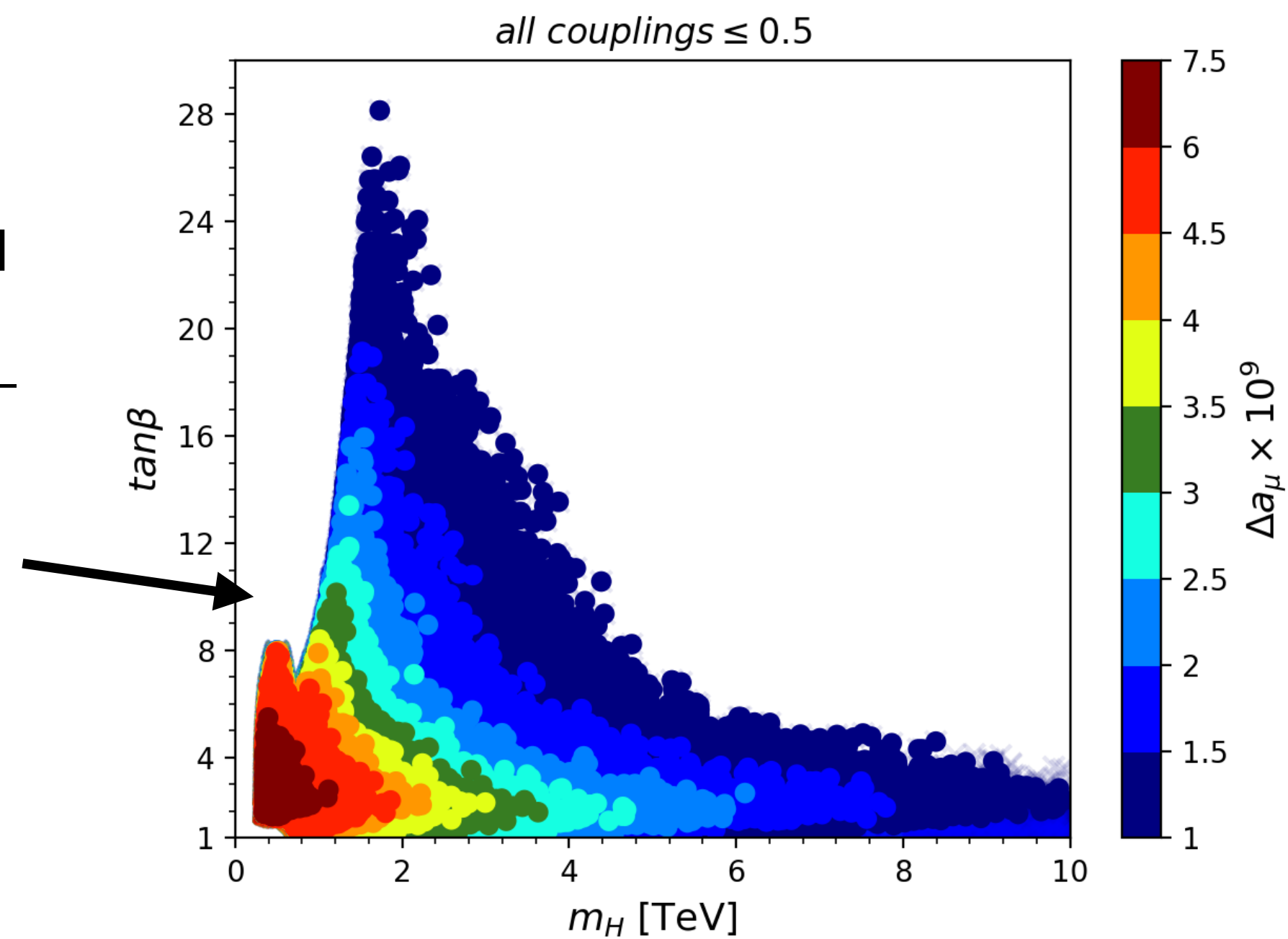


Δa_μ in 2HDM-II + VLS

Region ruled out by

$$H(A) \rightarrow \tau^+ \tau^-$$

ATLAS Collab.,
Phys. Rev. Lett. 125,
051801 (2020)



Couplings up to $\sqrt{4\pi}$ give almost 100 times Δa_μ^{exp} !

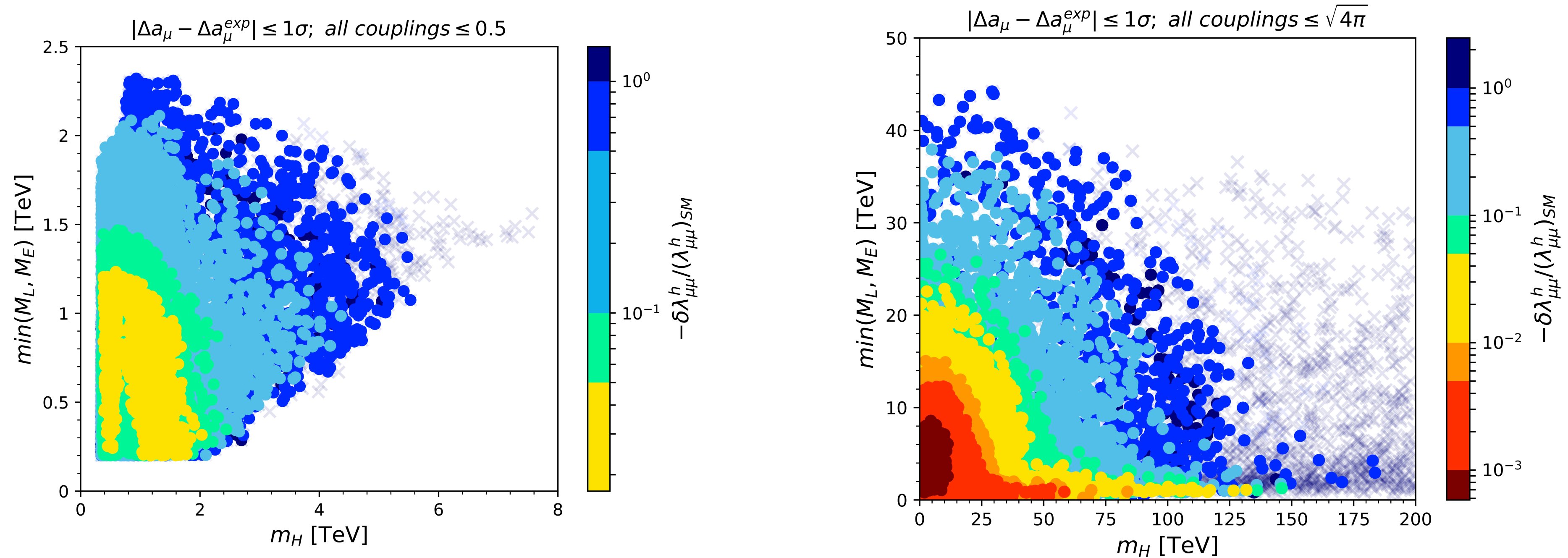
$$\Delta a_\mu = - \left(\frac{1}{16\pi^2} \right) \left(\frac{m_\mu m_\mu^{LE}}{v^2} \right) (1 + \tan^2 \beta) \quad m_\mu^{LE} = v_d^3 \left(\frac{\lambda_L \lambda_E \bar{\lambda}}{M_L M_E} \right)$$

- Largest contributions occur for intermediate $\tan \beta$ values and light Higgs masses ($\tan^2 \beta$ enhancement from heavy Higgs loops)
- All relevant constraints like EW precision constraints $|\lambda_L| \lesssim 0.04 \times M_L/v_d$ and $|\lambda_E| \lesssim 0.03 \times M_E/v_d$ are satisfied in these scenarios



Modifying Yukawa Couplings in 2HDM-II + VLs

Bright points w/
X correspond to
 $\Delta a_\mu^{H,A,H^\pm} / \Delta a_\mu > 50\%$



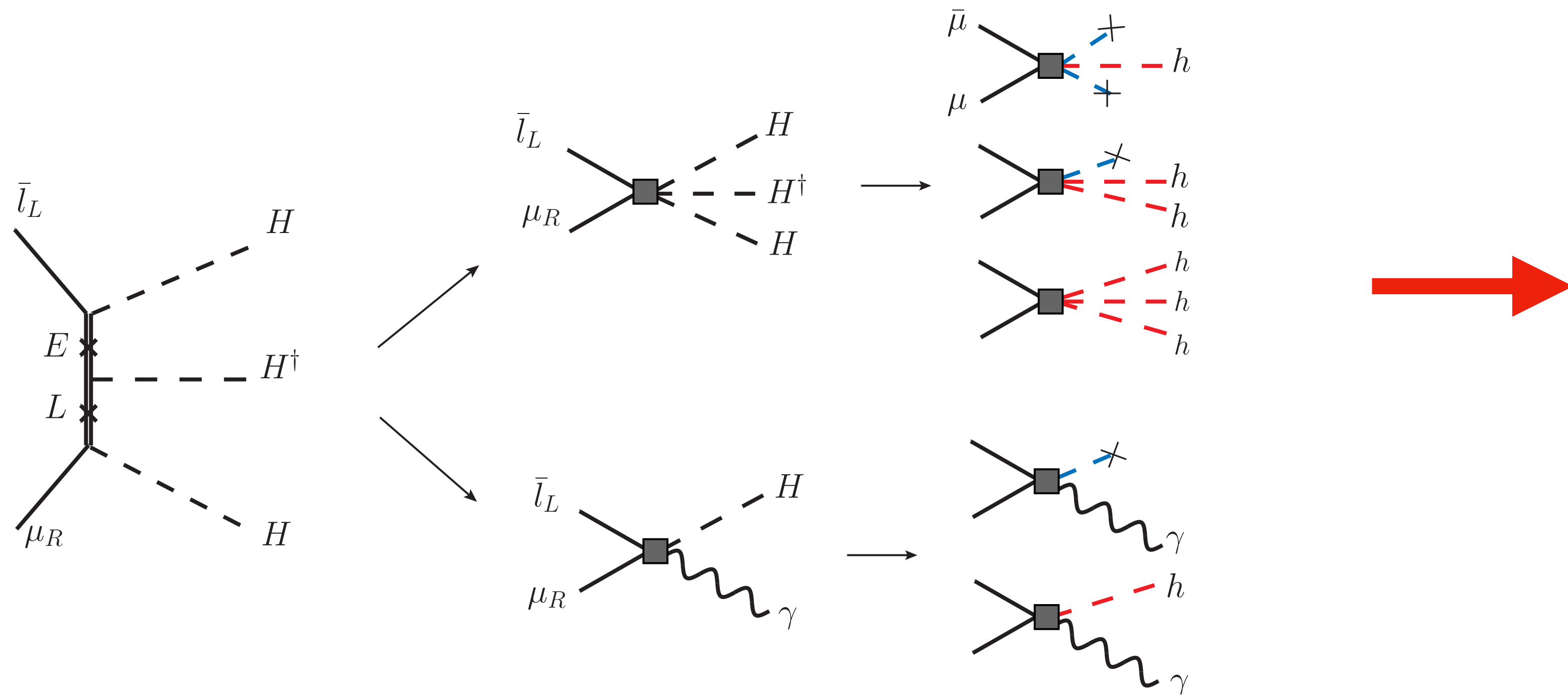
$$R_{h \rightarrow \mu^+ \mu^-} = \frac{|\lambda_{\mu\mu}^h|^2}{|(\lambda_{\mu\mu}^h)_{SM}|^2} = \left(1 + 2 \frac{m_\mu^{LE}}{m_\mu}\right)^2$$

- The mass spectrum of leptons goes from $M_{L,E} > 2.5$ (45) TeV for 0.5 ($\sqrt{4\pi}$) and couplings
- $-\delta\lambda_{\mu\mu}^h / (\lambda_{\mu\mu}^h)_{SM} = 2$ corresponds to $m_\mu^{LE} = -m_\mu$ and predicts the same $h \rightarrow \mu^+ \mu^-$ decay rate as the SM
- At LHC (FCC-hh), $\lambda_{\mu\mu}^h$ is expected to be measured at 5×10^{-2} (5×10^{-3}) precision



SM Effective Field Theory

- Let's look at the dimension-6 operator $\bar{l}_L H \mu_R (H^\dagger H)$



$$\lambda_{\mu\mu}^h = (m_\mu + 2m_\mu^{LE})/v$$

$$\lambda_{\mu\mu}^{hh} = 3 \frac{m_\mu^{LE}}{v^2}$$

$$\lambda_{\mu\mu}^{hhh} = \frac{3}{\sqrt{2}} \frac{m_\mu^{LE}}{v^3}$$

$$\Delta a_\mu = - \left(\frac{1}{16\pi^2} \right) \left(\frac{m_\mu m_\mu^{LE}}{v^2} \right)$$

In the 2HDM-II, this Δa_μ is multiplied by $1 + \tan^2 \beta$; Same result as earlier!

- Because of this connection, the experimental value $\Delta a_\mu^{exp} = (2.51 \pm 0.59) \times 10^{-9}$ fixes

$$\frac{m_\mu^{LE}}{m_\mu} = -1.07 \pm 0.25$$

$$R_{h \rightarrow \mu^+ \mu^-} \equiv \frac{BR(h \rightarrow \mu^+ \mu^-)}{BR(h \rightarrow \mu^+ \mu^-)_{SM}} = \left(1 + 2 \frac{m_\mu^{LE}}{m_\mu} \right)^2 = 1.32_{-0.90}^{+1.40}$$

We can make a prediction for di- and ti-Higgs signals!
How can we test this model?

Current limit is 2.2



Testing at a Muon Collider

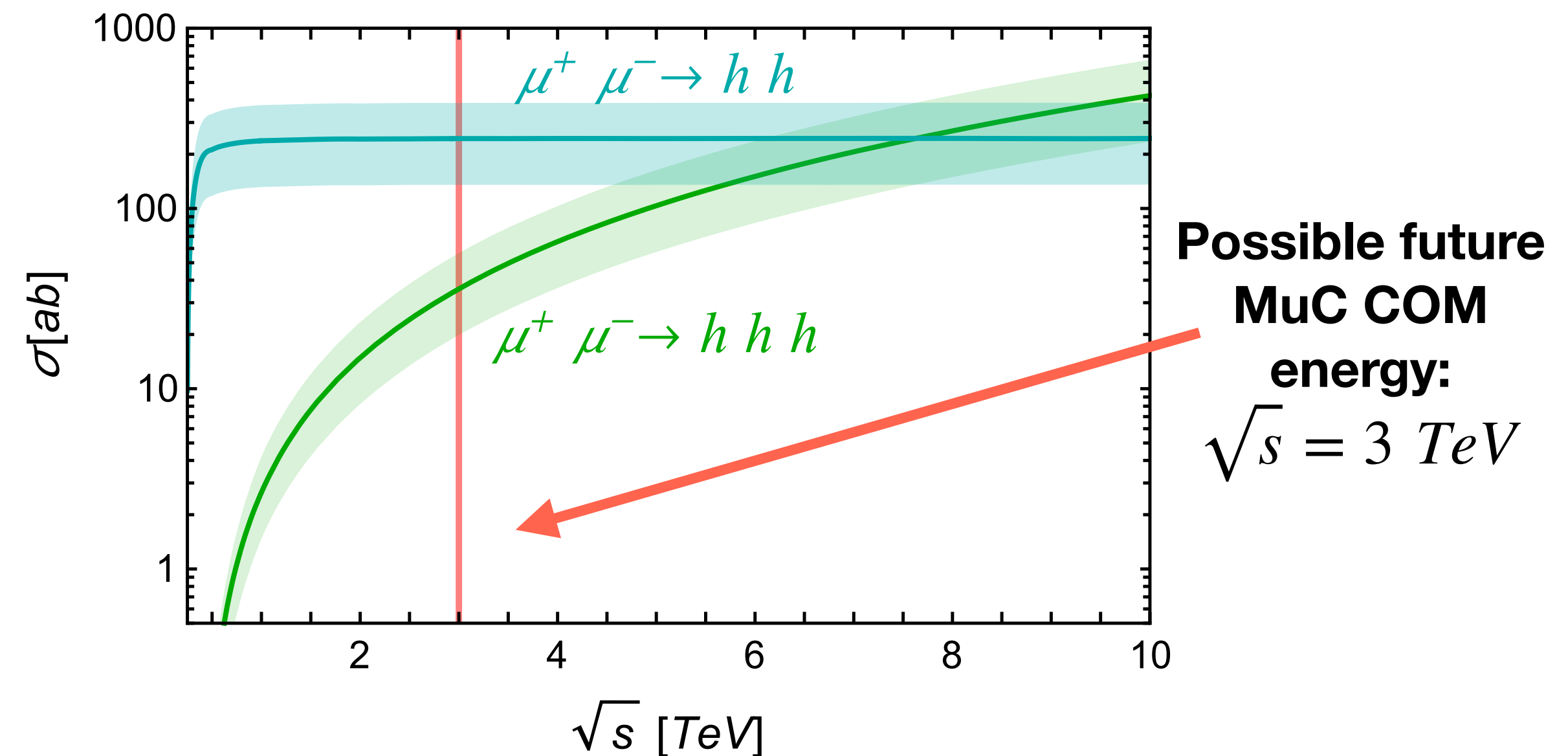
- The effective cross-sections for $\mu^+\mu^- \rightarrow hh$ and $\mu^+\mu^- \rightarrow hhh$ can be written within $\Delta a_\mu^{exp} \pm 1\sigma$

$$\sigma_{\mu^+\mu^- \rightarrow hh} = \frac{|\lambda_{\mu\mu}^{hh}|^2}{64\pi} = \frac{9}{64\pi} \left(\frac{m_\mu^{LE}}{v^2} \right)^2$$

$$\sigma_{\mu^+\mu^- \rightarrow hhh} = \frac{|\lambda_{\mu\mu}^{hhh}|^2}{6144\pi^3} s = \frac{3}{4096\pi^3} \left(\frac{m_\mu^{LE}}{v^3} \right)^2 s$$

- A $\sqrt{s} = 3$ TeV MuC with 1 ab^{-1} luminosity predicts 244 and 36 events for di- and tri-Higgs processes at the central value of Δa_μ^{exp}
- At $\sqrt{s} > 1$ TeV, SM predictions for di- and tri-Higgs signals are small by *at least* 3-4 orders

T. Han *et al.* JHEP12(2021)162



Different representations of VLs coupling to the muon can decrease $\mu^+\mu^- \rightarrow hh$, $\mu^+\mu^- \rightarrow hhh$ dramatically; sharp distinction between models!



Takeaways

- 2HDM-II + VLs generates chiral and $\tan^2 \beta$ enhancements that can generate even up to two orders larger contributions to Δa_μ while satisfying constraints
- Precision measurements of $\lambda_{\mu\mu}^h$ can indirectly probe the parameter space of heavy VL masses and Higgses
- In the SM with VLs, the dimension-6 operator $\bar{l}_L H \mu_R (H^\dagger H)$ connects Δa_μ with $h \rightarrow \mu^+ \mu^-$, $\mu^+ \mu^- \rightarrow hh$, $\mu^+ \mu^- \rightarrow hhh$
- 1 ab^{-1} luminosity at $\sqrt{s} = 3 \text{ TeV}$ predicts 244 events for $\mu^+ \mu^- \rightarrow hh$ and 36 events for $\mu^+ \mu^- \rightarrow hhh$ at the central value of Δa_μ^{exp} . Other representations may lower the rate up to a factor of 25

Thank you for listening!



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