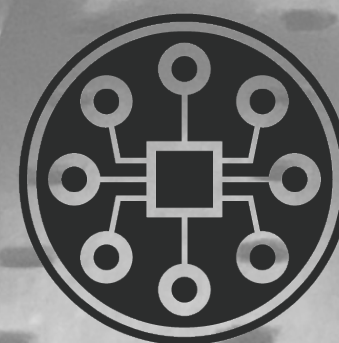


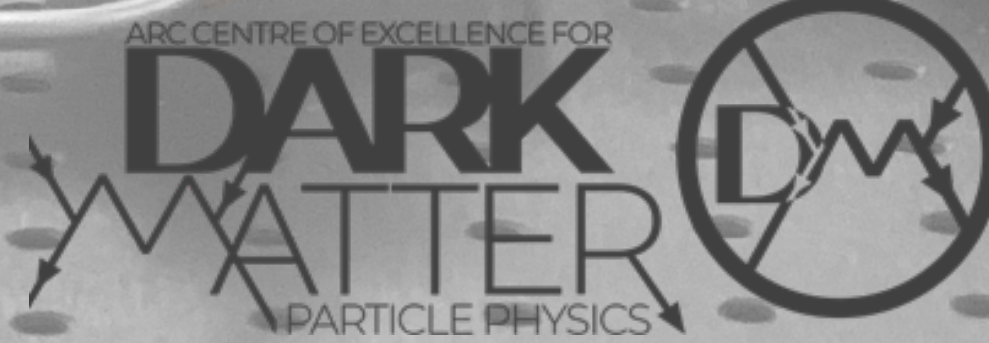
Precision Metrology with Photons, Phonons and Spins: Answering Major Unsolved Problems in Physics and Advancing Translational Science



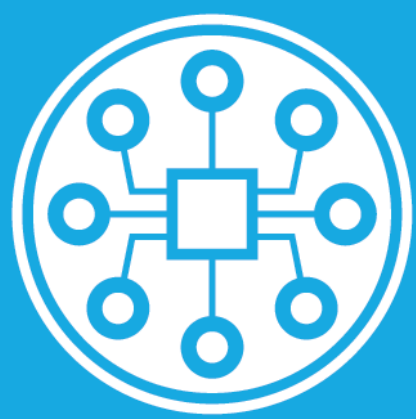
IEEE UFFC Distinguished Lecturer Program



EQUS
Australian Research Council
Centre of Excellence for
Engineered Quantum Systems



THE UNIVERSITY OF
**WESTERN
AUSTRALIA**



EQUS

Australian Research Council
Centre of Excellence for
Engineered Quantum Systems

The QDM Lab: <https://www.qdmlab.com/>
QUANTUM TECHNOLOGIES AND DARK
MATTER RESEARCH LAB



THE UNIVERSITY OF
**WESTERN
AUSTRALIA**



Our Team

HDR/PHD STUDENTS

Catriona Thomson
William Campbell
Aaron Quiskamp
Elrina Hartman

UNDERGRAD STUDENTS

Steven Samuels (Hons)
Emma Paterson (Hons)
Campbell Millar (MPE)
Ishaan Goel (MPE)
Deepali Rajawat (MPE)
Michael Hatzon (BPhil)
Emily Waterman (BPhil)
Ashley Johnson (BPhil)

ACADEMIC

Michael Tobar
Eugene Ivanov
Maxim Goryachev

POSTDOCS

Ben McAllister
Cindy Zhao
Jeremy Bourhill
Graeme Flower

ADJUNCT

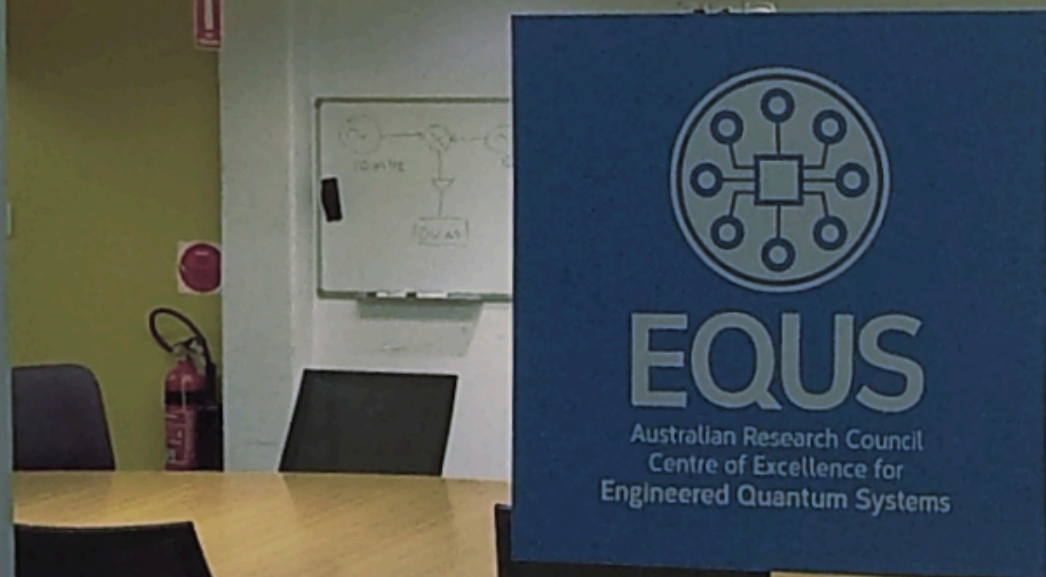
Alexey Veryaskin (Trinity Labs)

<https://www.qdmlab.com/>

QDM Laboratory

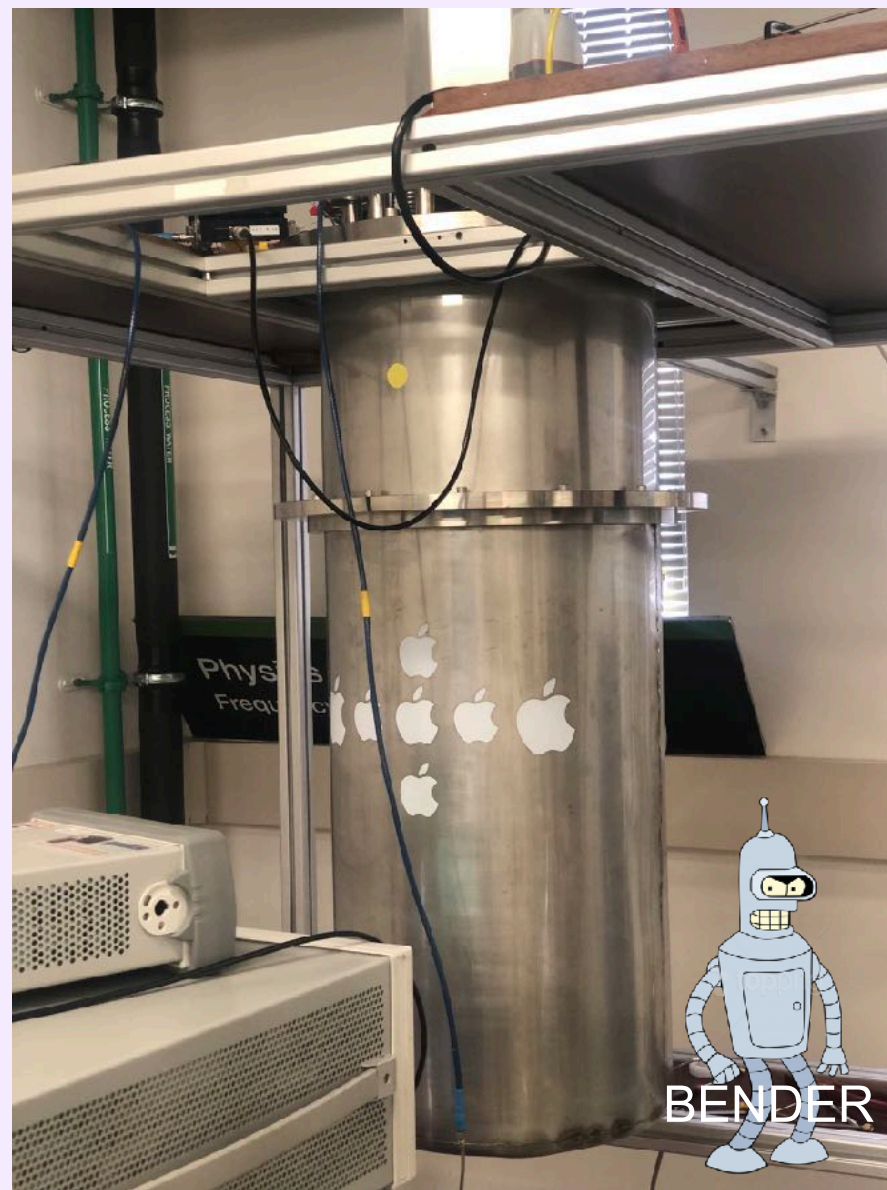
Welcome to the Quantum Technology and Dark Matter Laboratory at UWA!

Come inside and check out our world class facilities, home to nodes of the ARC Centre of Excellence for Engineered Quantum Systems, and the ARC Centre of Excellence for Dark Matter Particle Physics.

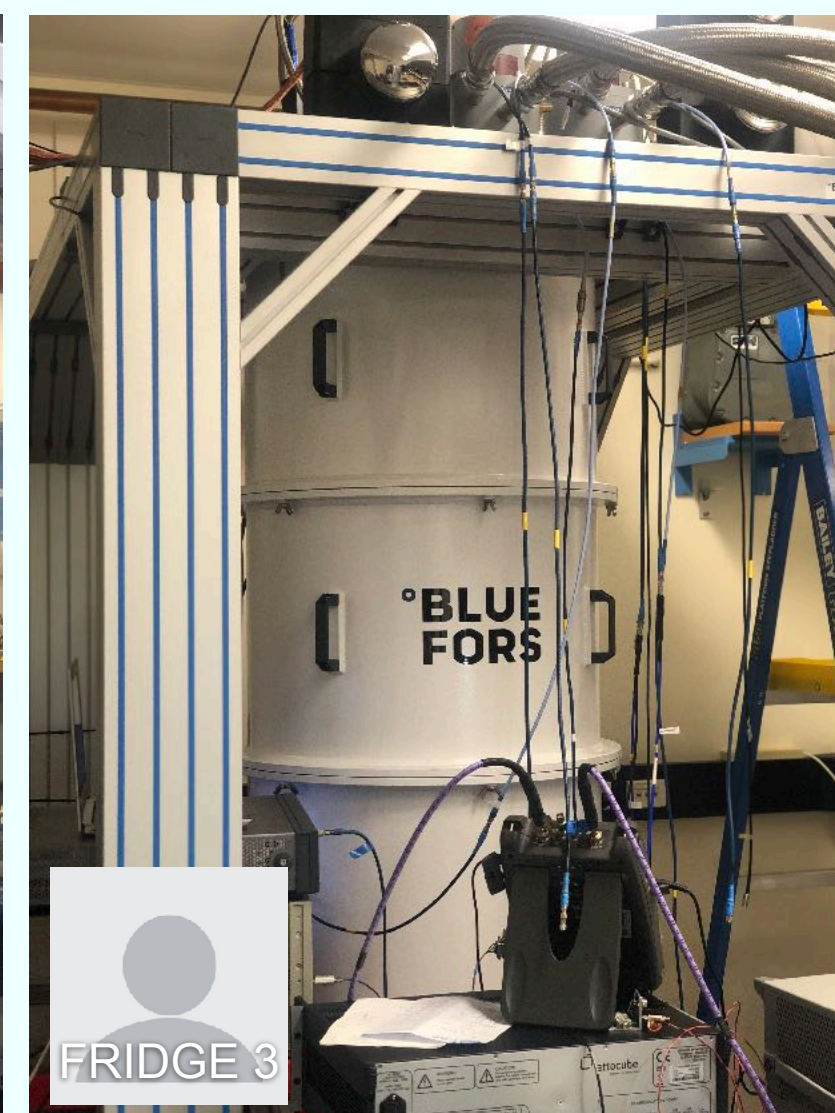
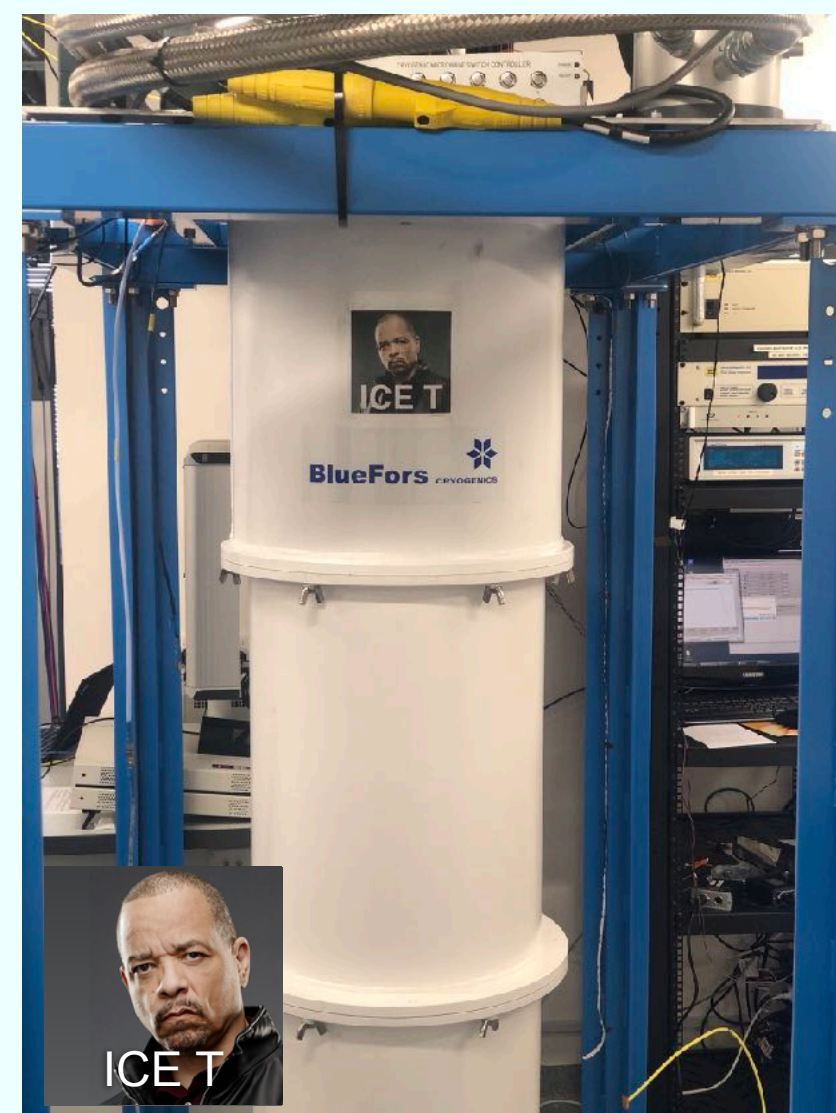


QDM Laboratory Meeting Room and Foyer Quantum Systems Laboratory

4 Kelvin Systems



- Extensive experience with cryogenic systems
- 3, 7 and 12 T superconducting magnets
- Large collection of microwave (and a some optical) diagnostic equipment and hardware
- Expertise with precision frequency metrology



Dilution Systems

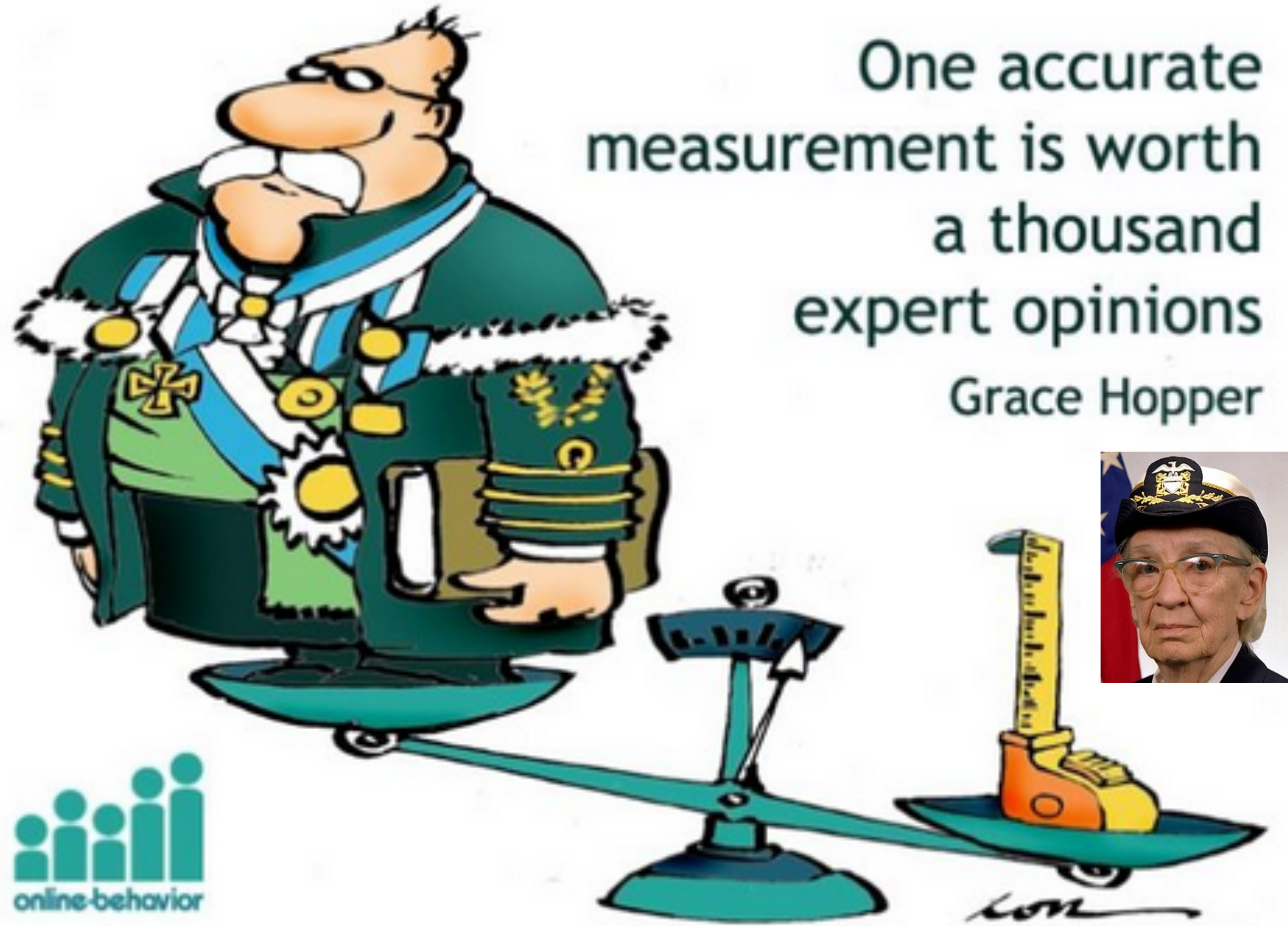
Physics: An experimental science

White, Harvey Elliott

Note: This is not the actual book cover

Physics: An experimental science

White, Harvey Elliott

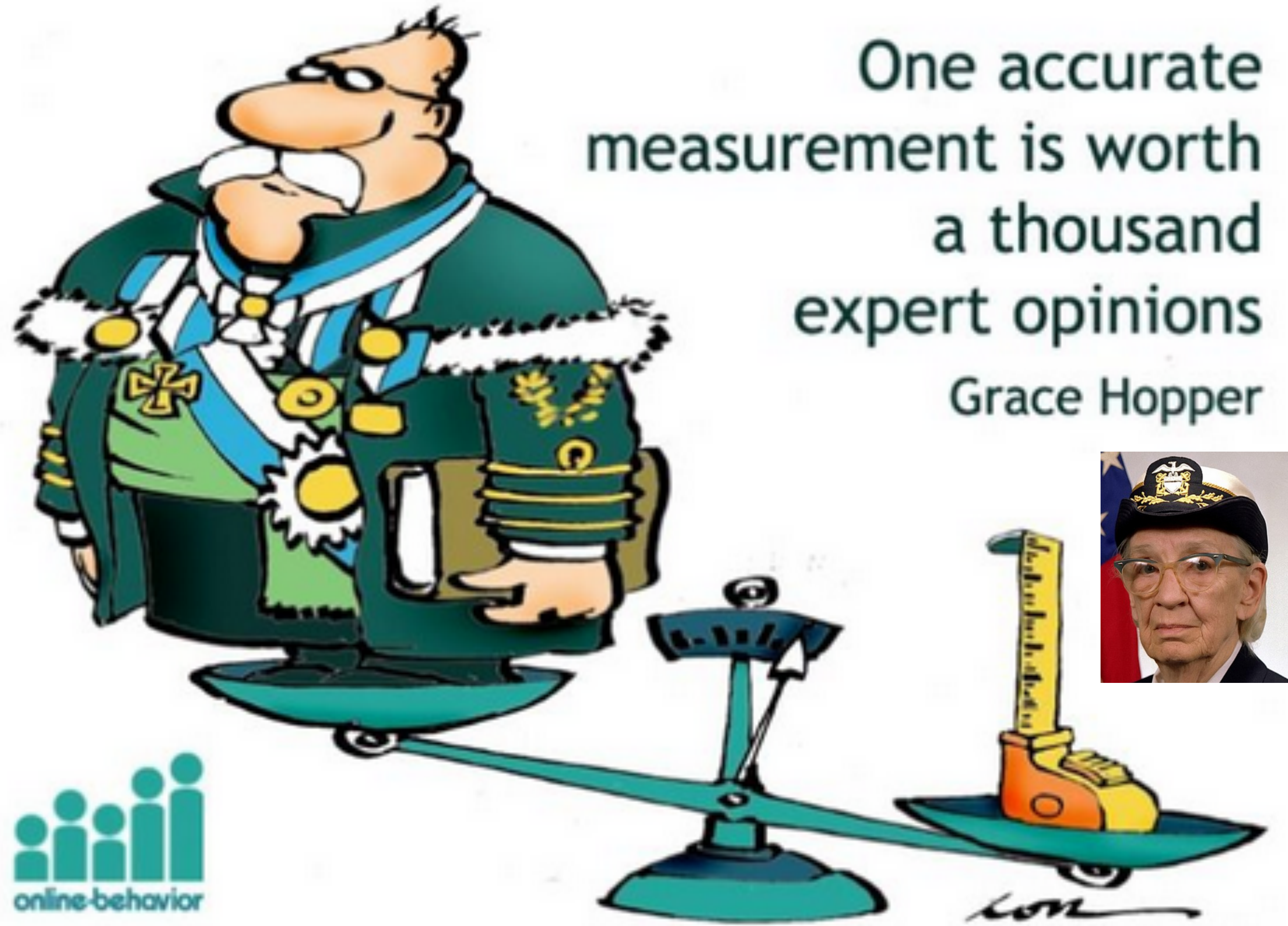


Note: This is not the actual book cover

Physics: An experimental science

White, Harvey Elliott

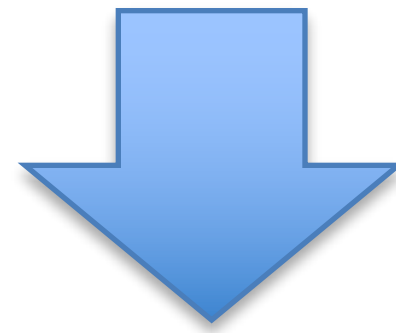
Note: This is not the actual book cover



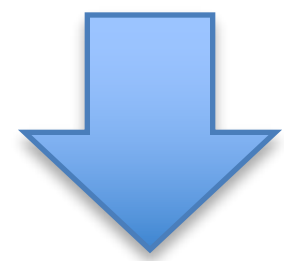
Computer Scientist
Invented COBOL

Precision Measurements @ QDM Lab Uni of Western Australia

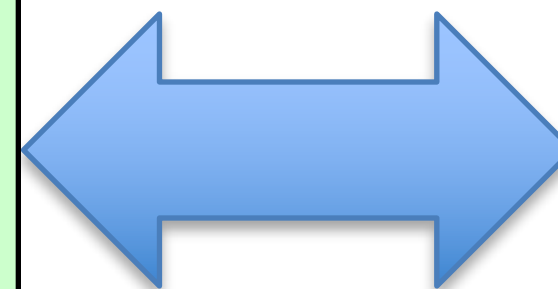
Precision measurement =>
Phase, Frequency, Energy, Time



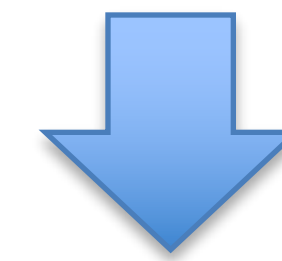
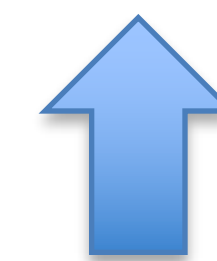
Technology: High-Q -> Narrow Line Width Systems:
Low Noise Techniques Classical and Quantum (SQL)



New Tests of
Fundamental Physics



Applications: Sensors,
Clocks, Radar etc.



**QDM has realised many best tests on
Dark Matter, Quantum Gravity and
Relativity**

**12 Patents: Radar, Gradiometer, Sensing,
Oscillators**

Organisation

Governing Board

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- Andrea de Marchi
- James C. Bergquist
- Patrick Gill
- Lute Maleki
- Fritz Riehle

Symposium Chair

- Michael Tobar

Symposium Secretariat

- Angela Bird

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- Tingrei Tan (USyd)
- Katrina Tune (UQ)
- Robert Williams (NMI)
- Magdalena Zych (UQ)
- Virginia Escudero: Non-Academic Events

International Steering Committee

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- Sebastian Bize (France)
- Davide Calonico (Italy)
- Scott Diddams (USA)
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- Pierre Dubé (Canada)
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- Kurt Gibble (USA)
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- Hidetoshi Katori (Japan)
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- Ekkehard Peik (Germany)
- Fritz Riehle (Germany)
- Patrizia Tavella (France)
- Thomas Udem (Germany)
- Peter Wolf (France)
- Michael Wouters (Australia)
- Jun Ye (USA)
- Nan Yu (USA)

The 9th Symposium
on Frequency Standards
and Metrology

Summer School on Frequency Standards, Precision and Quantum Measurement

Rydges Gold Coast Airport Hotel,
Gold Coast, QLD, Australia

9–13 October 2023

*The week before 9FSM
we are presenting a
Summer School on
Frequency Standards,
Precision and Quantum
Measurement for students,
early career researchers,
and those professionals
who want to learn more
on the subject.*

<https://www.qdmlab.com/9fsm/summer-school>

(+61) 0404872944

9fsm@uwa.edu.au



Confirmed Speakers

John Close, Professor of Physics, Leader of the Atom Laser Group, Australian National University, Canberra, Australia
Topic: Using Atoms for Precision and Quantum Measurements

Tara Fortier, Project Leader, Time and Frequency Division, National Institute of Standards and Technology, Colorado, USA
Topic: Optical and Microwave Metrology

Anna Grassellino, Director, Superconducting Quantum Materials and Systems Center, Fermilab, Illinois, USA
Topic: High-Q Super Conducting Cavities, and Application to Quantum Measurements and Tests of Fundamental Physics

Eugene Ivanov, Research Professor, University of Western Australia, Dept. Physics, Perth, Australia.
Topic: Low Noise Frequency Stable Microwave Oscillators and Metrology

John Kitching, Group Leader and Fellow, Time and Frequency Division, National Institute of Standards and Technology, Colorado, USA.
Topic: Vapor cell clocks, CSACs and quantum sensing

Helen Margolis, Head of Science for Time & Frequency and NPL Fellow in Optical Frequency Standards and Metrology, National Physical Laboratory, London, UK.
Topic: Optical clocks for international timekeeping

Ekkehard Peik, Head of Time and Frequency Department, Physikalisch-Technische Bundesanstalt, Braunschweig, Germany.
Topic: Nuclear Clocks

Piet Schmidt, Physikalisch-Technische Bundesanstalt, Braunschweig, Germany.
Topic: Highly Charged Ion Clocks and Testing Fundamental Physics.

Patrizia Tavella, Director of Department of Time, Bureau International des Poids et Mesures, Paris, France.
Topic: UTC time scales or mathematical and statistical methods for time and frequency standards

Nora Tischler, ARC DECRA Fellow, Centre for Quantum Dynamics, Griffith University, Brisbane, Australia.
Topic: The physics of light detection and correlation measurements.

Michael Wouters, Head of Time and Frequency National Measurement Institute Australia, Sydney, NSW, Australia.
Topic: Advanced time and frequency, time-transfer methods, analysis of Time and Frequency data

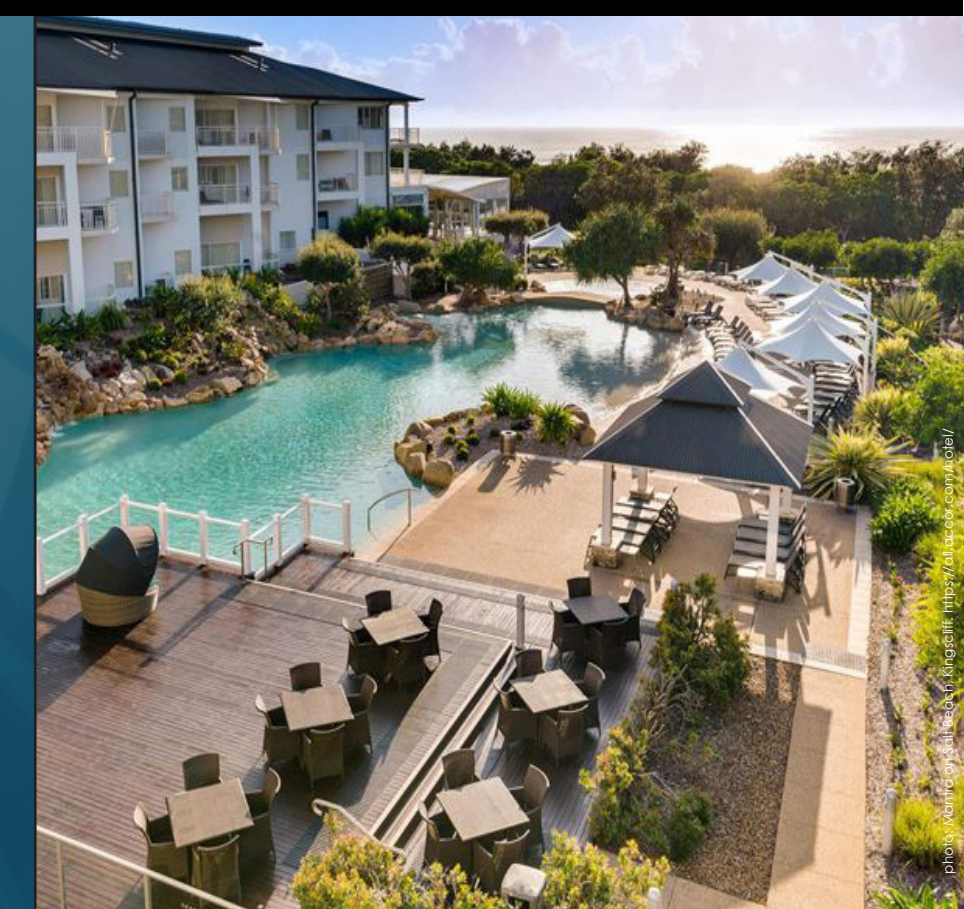
Sébastien Bize, Head of the optical frequency group, Systèmes de Référence Temps Espace, Observatoire de Paris.
Topic: Optical lattice clocks.

Tanya Zeevinsky, Professor of Physics, Columbia University, New York City, USA.
Topic: Precision molecular clocks and systems

The 9th Symposium on Frequency Standards and Metrology

Mantra on Salt Beach Resort,
Kingscliff, NSW, Australia

16–20 October 2023



Confirmed Invited Speakers

Alexander Romaneko, Fermilab, Illinois, USA.

Andrei Derevianko, University of Nevada, Reno, USA.

Andrew Ludlow, National Institute of Standards and Technology, Boulder, USA.

Andrey Matsko, Jet Propulsion Laboratory, Pasadena, USA.

Anna Grassellino, Fermilab, Illinois, USA.

Anne Amy-Klein, Université Sorbonne, Paris Nord, France.

Davide Calonico, Istituto Nazionale di Ricerca Metrologica, Torino, ITALY.

Dmitry Budker, Helmholtz Institute, Mainz, Germany.

Ekkehard Peik, Physikalisch-Technische Bundesanstalt, Braunschweig, Germany.

Eric Burt, Jet Propulsion Laboratory, Pasadena, USA.

Eugene Ivanov, University of Western Australia, Perth, Australia.

Helen Margolis, National Physical Laboratory, Teddington, London, UK.

Hidetoshi Katori, RIKEN, Saitama, Japan.

Jacques Vanier, University of Montreal, Canada.

James Chin-wen Chau, National Institute of Standards and Technology, Boulder, USA.

Jian-Wei Pan, University of Science and Technology of China, Hefei, China.

John Kitching, National Institute of Standards and Technology, Boulder, USA.

Jun Ye, JILA, Boulder Colorado, USA.

Krzysztof Szymaniec, National Physical Laboratory, Teddington, London, UK.

Mark Kasevich, Stanford, USA.

Maxim Goryachev, University of Western Australia, Perth, Australia.

Murray Barrett, National University of Singapore, Singapore.

Nan Yu, Jet Propulsion Laboratory, Pasadena, USA.

Nathan Newbury, National Institute of Standards and Technology, Boulder, USA.

Patrick Gill, National Physical Laboratory, Teddington, London, UK.

Piet Schmidt, Physikalisch-Technische Bundesanstalt, Braunschweig, Germany.

Scott Papp, National Institute of Standards and Technology, Boulder, USA.

Sébastien Bize, Systèmes de Référence Temps Espace, Paris Observatory, France

Tanja Mehlstäubler, Physikalisch-Technische Bundesanstalt, Braunschweig, Germany.

Tanya Zeevinsky, Columbia University, New York City, USA

Tara Fortier, National Institute of Standards and Technology, Boulder, USA.

Tetsuya Ido, National Institute of Information and Communications Technology, Japan.

Thomas Udem, Max Planck Institute of Quantum Optics, Garching, Germany.

Uwe Sterr, Low noise optical cavities and transportable clock laser systems

Victor Flambaum, University of New South Wales, Sydney, Australia.

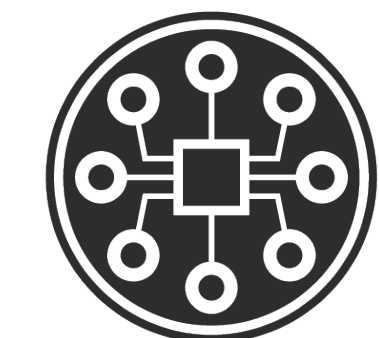
Yannick Bidet, ONERA, The French Aerospace Lab, France.

An international
discussion forum on
precision frequency
standards throughout
the electromagnetic
spectrum, and
associated precision
and quantum
metrology.

<https://www.qdmlab.com/9fsm>

(+61) 0404872944

9fsm@uwa.edu.au



EQUS

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Centre of Excellence for
Engineered Quantum Systems

QDM Lab Precision Metrology: See www.qdmlab.com

Science of precise measurement

Physics at low energies

Metrological Systems:

Photonic

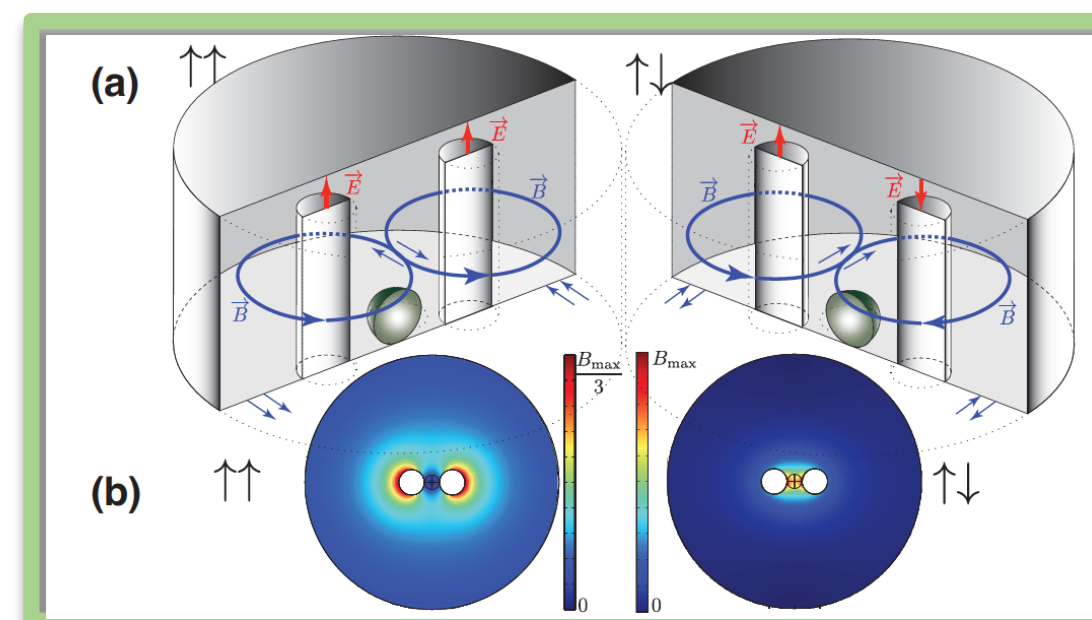
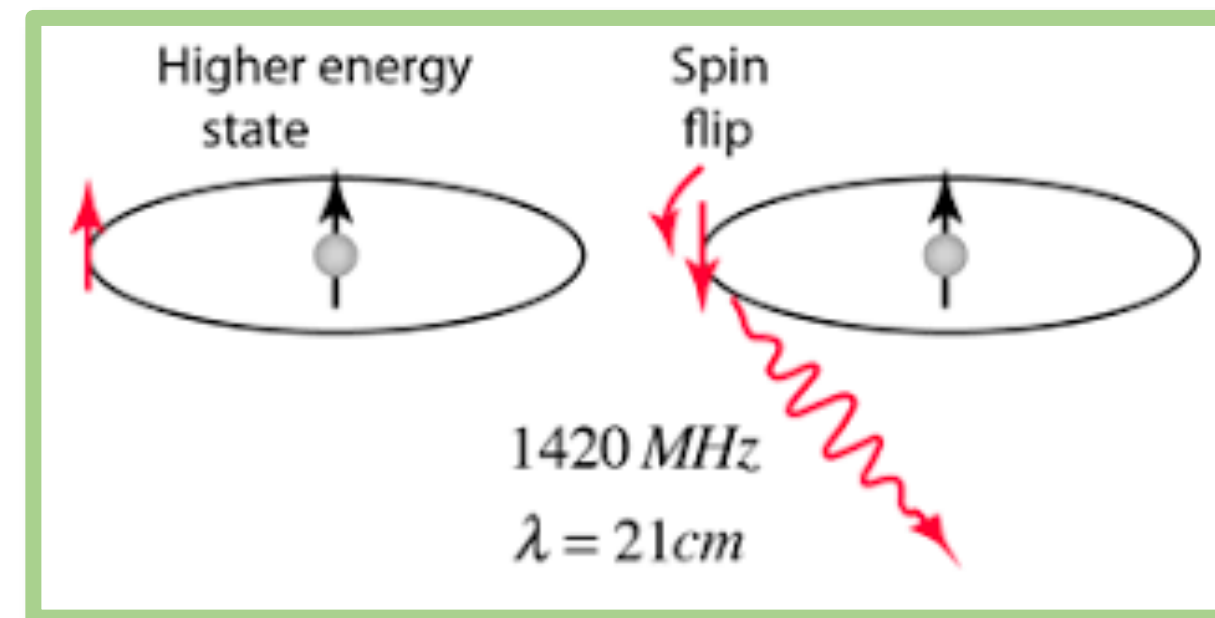
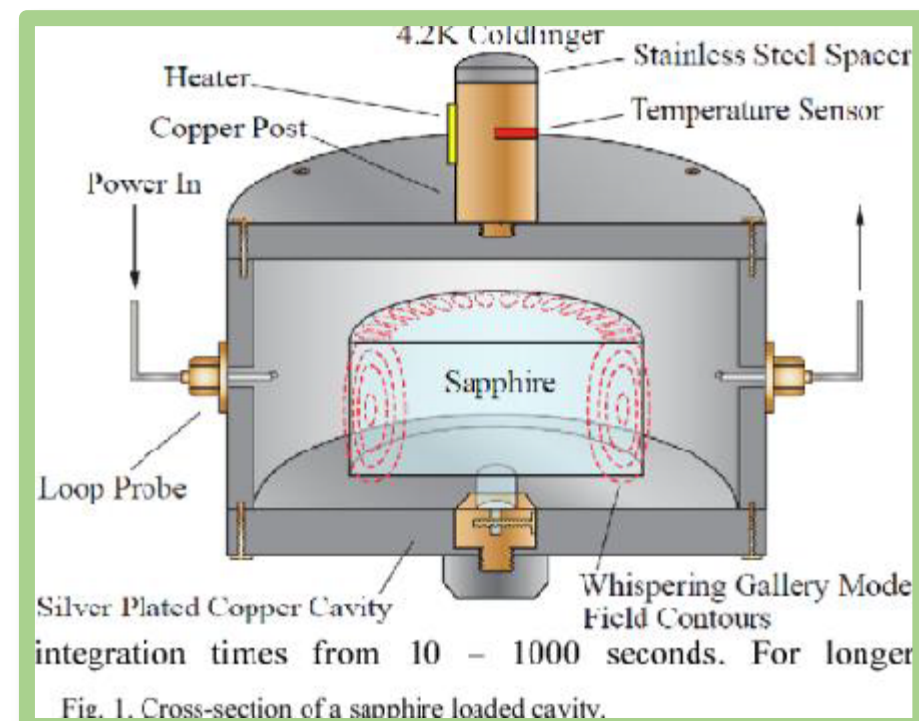
- WGM Resonators
- Novel Microwave Cavities

Atomic/Spins

- H - Maser
- Atomic Clocks
- Spin Waves

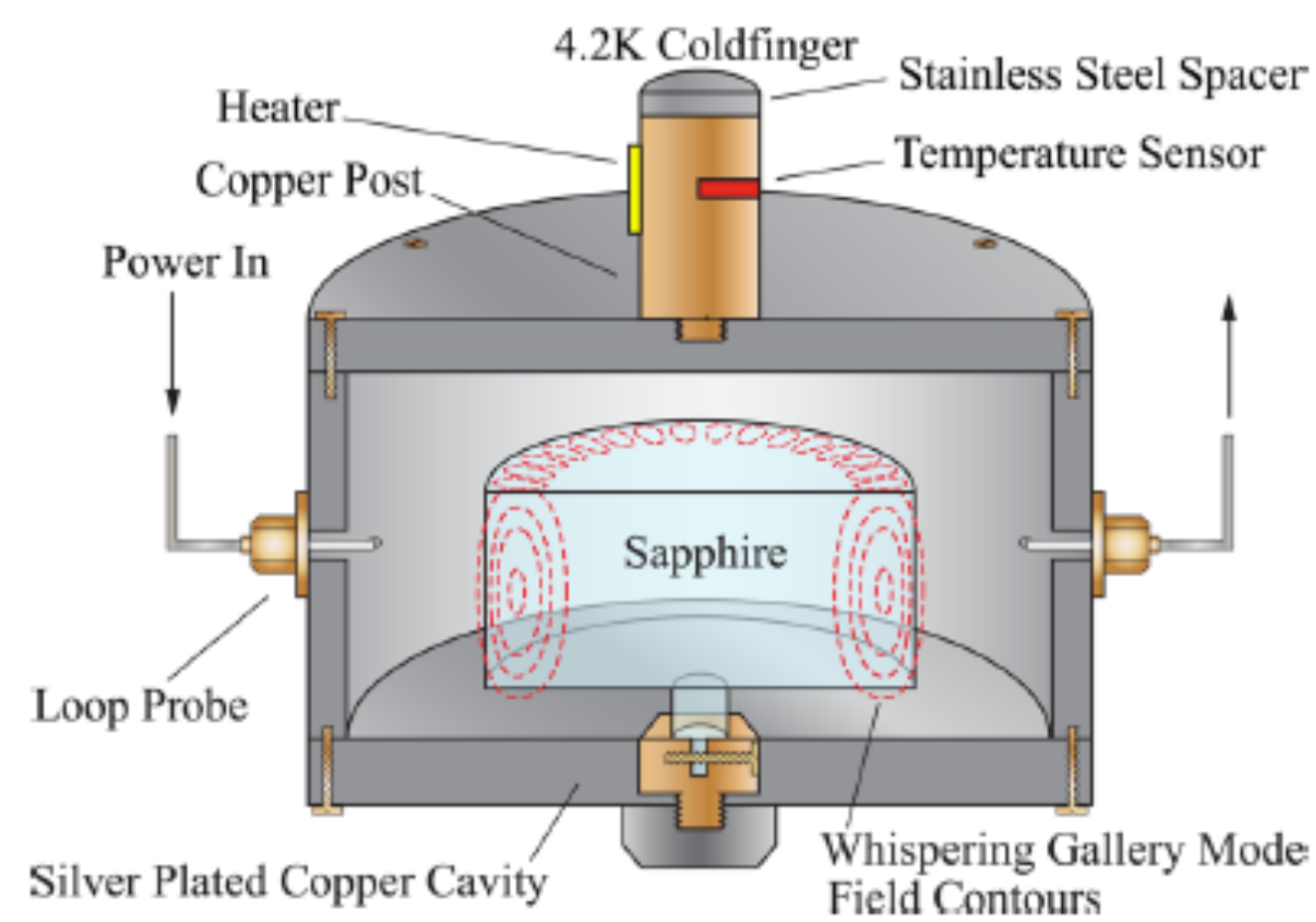
Acoustic

- Superfluid
- BAW Resonator

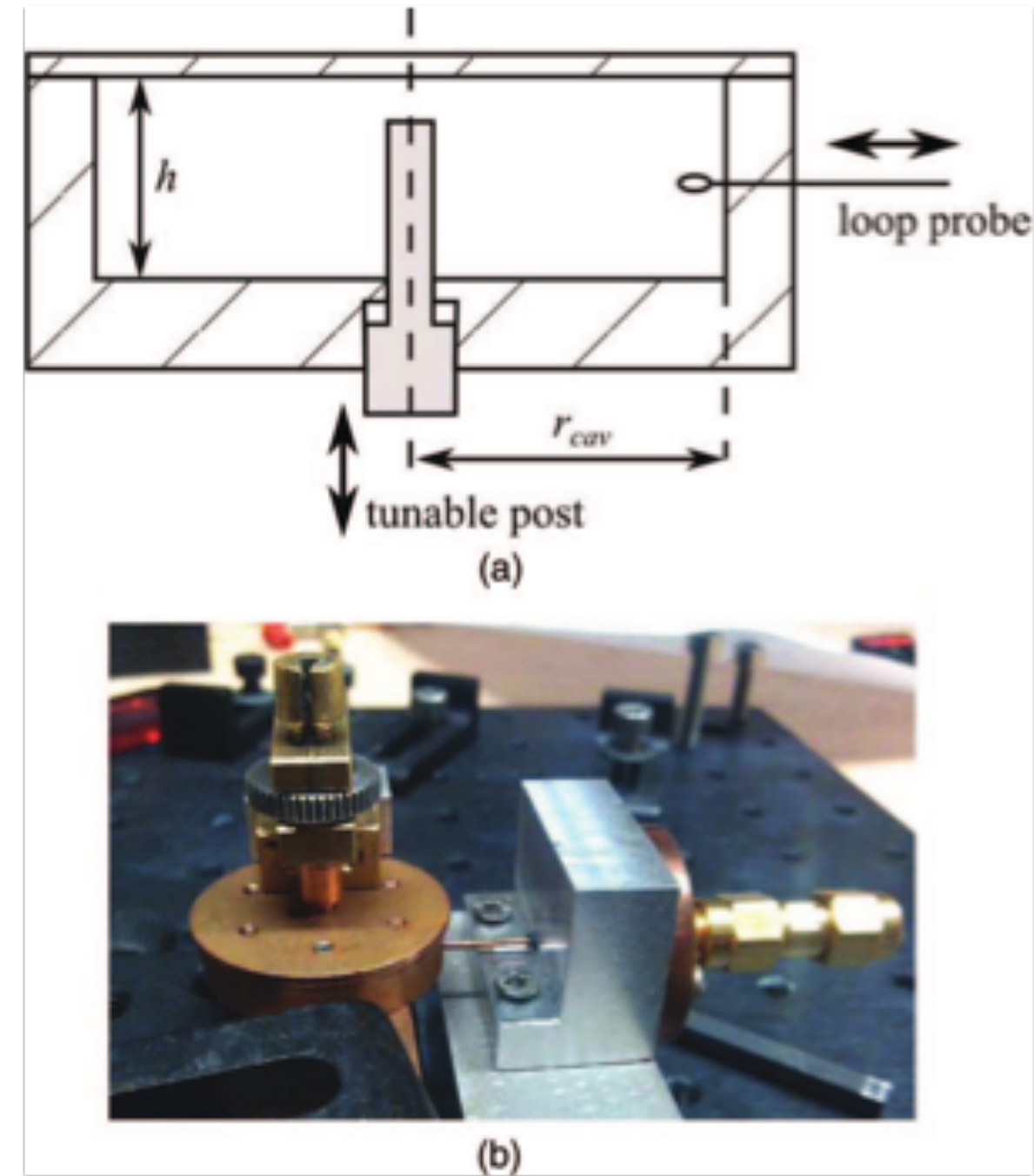


Types of Photonic Cavities to Couple to Spins, Phonons etc.

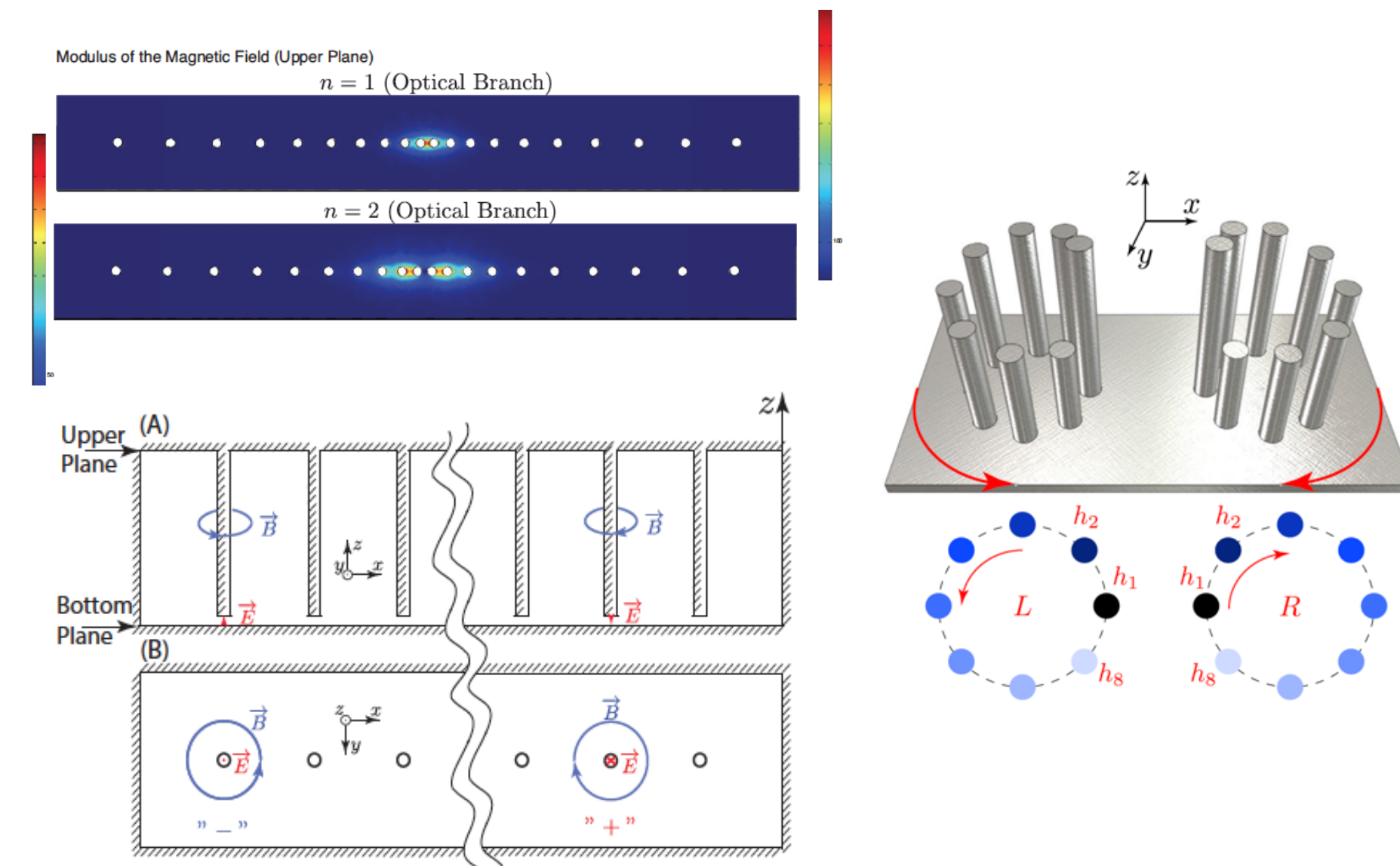
WG Modes



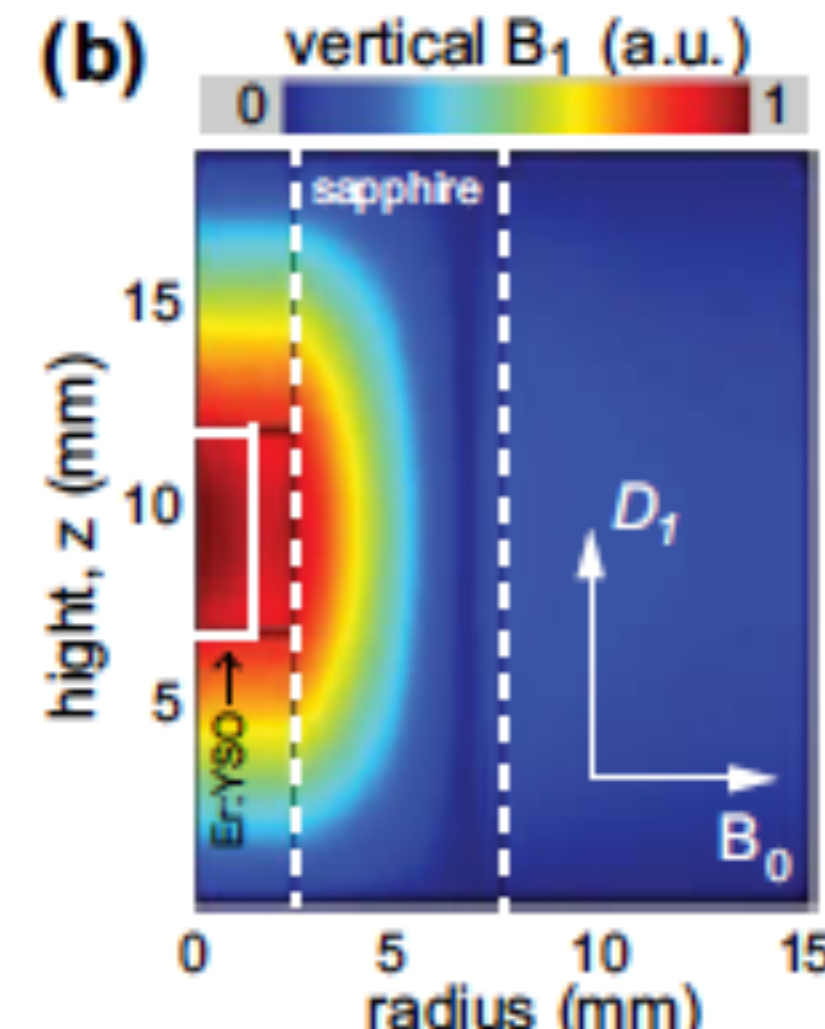
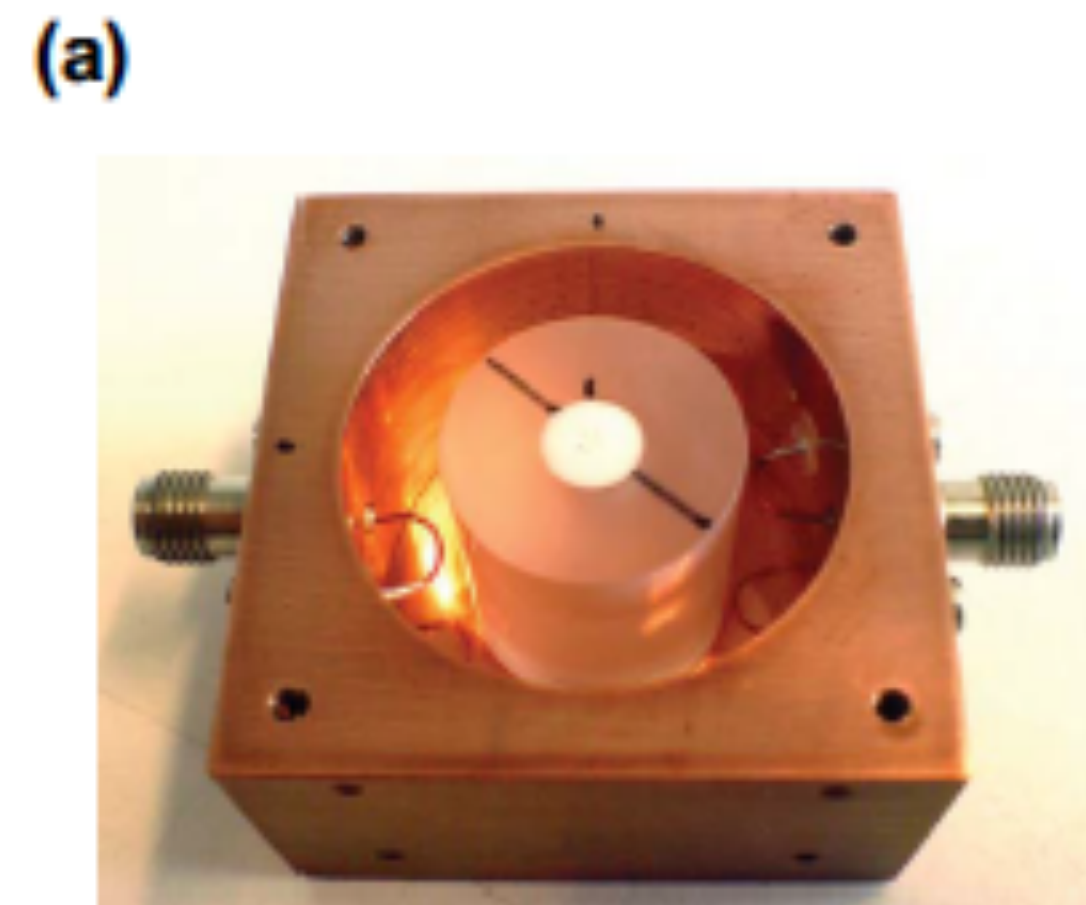
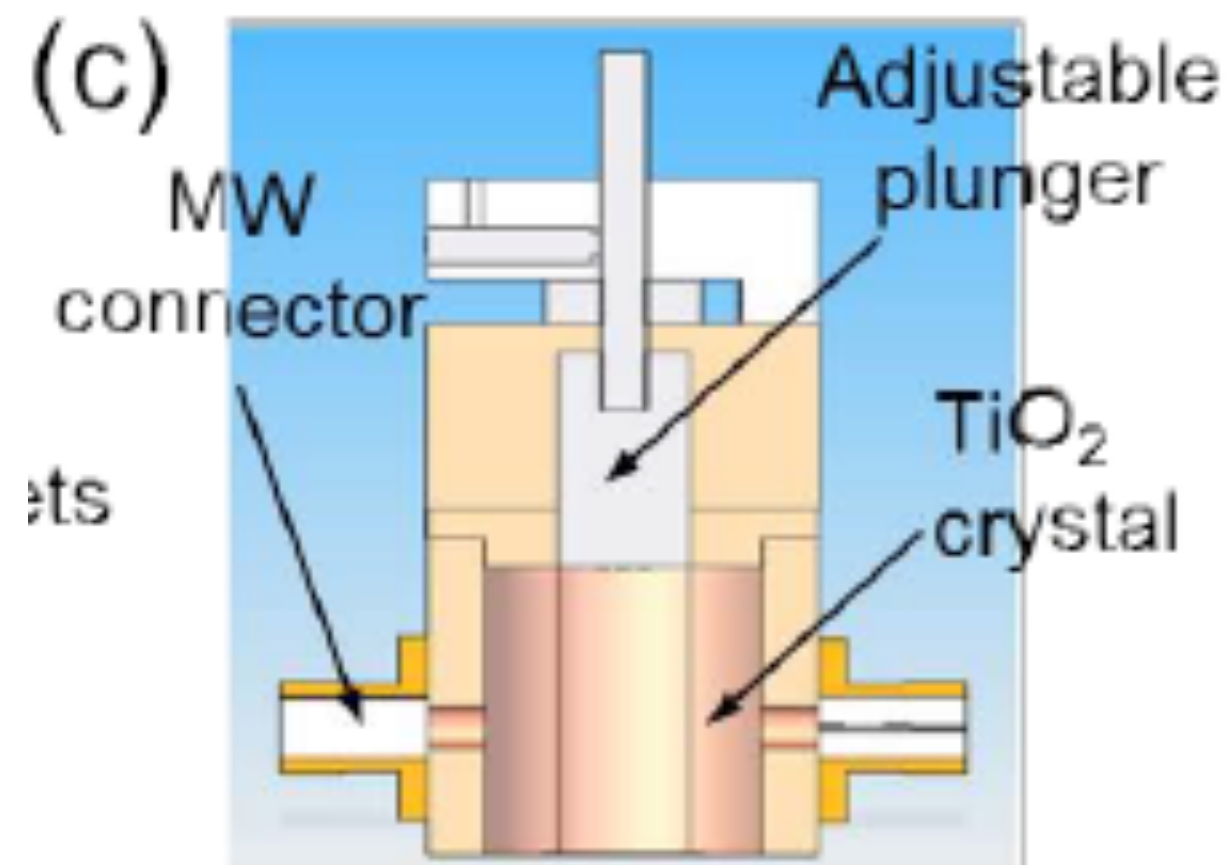
Reentrant: 3D LC



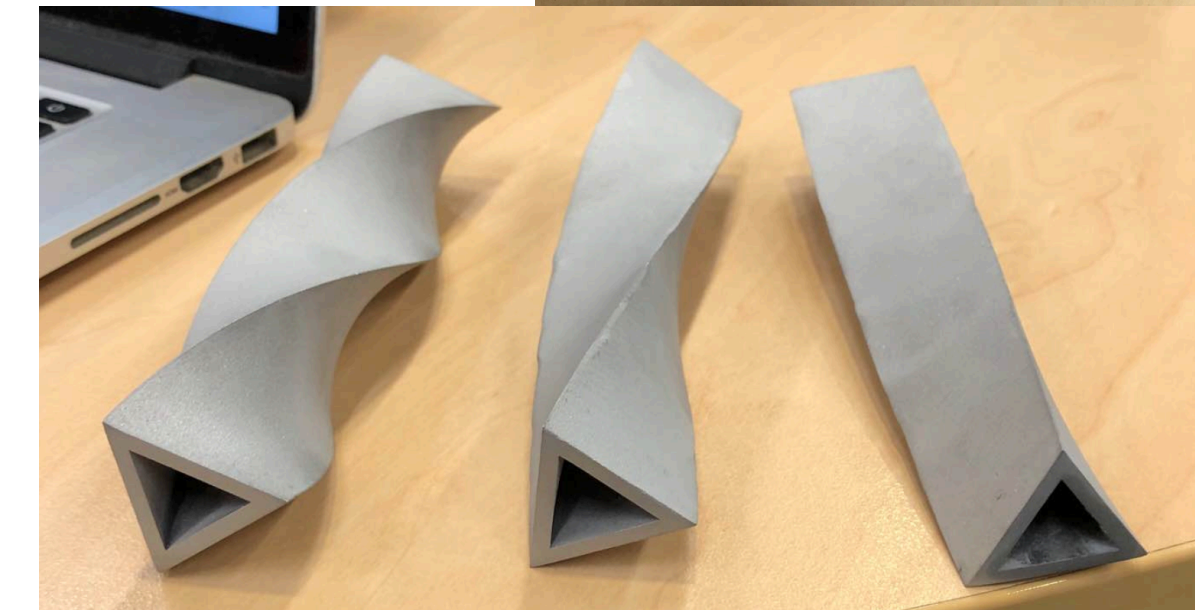
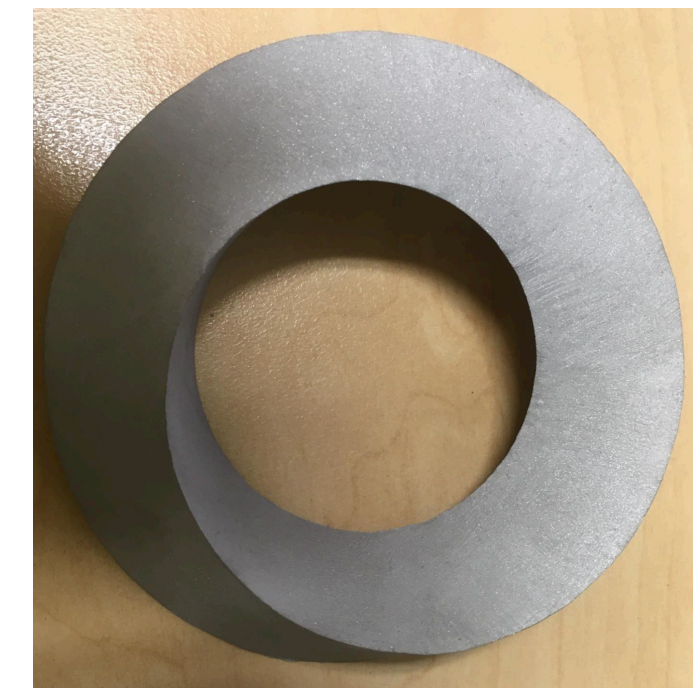
Reentrant Lattice



TE + TM Cylindrical modes

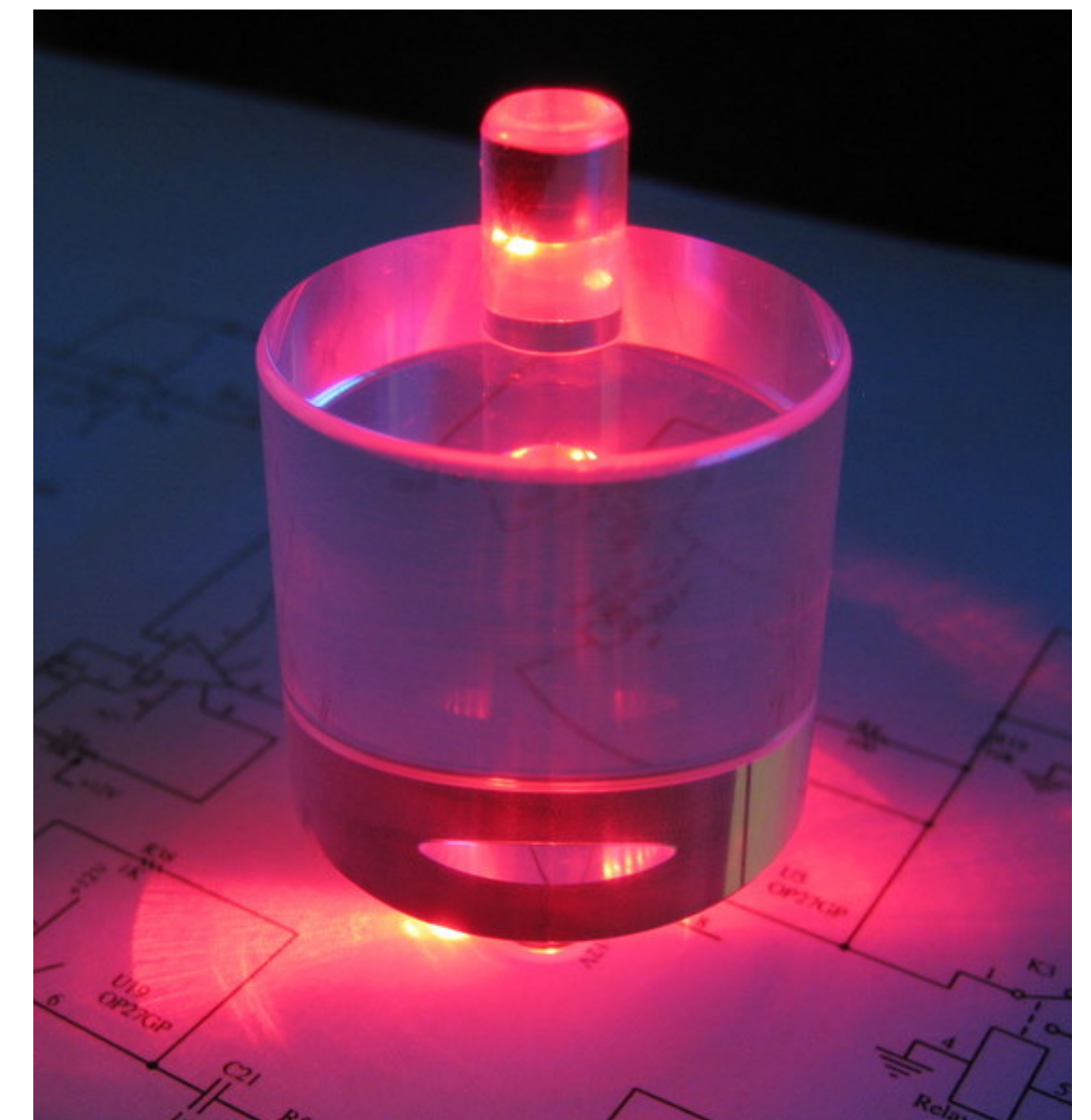
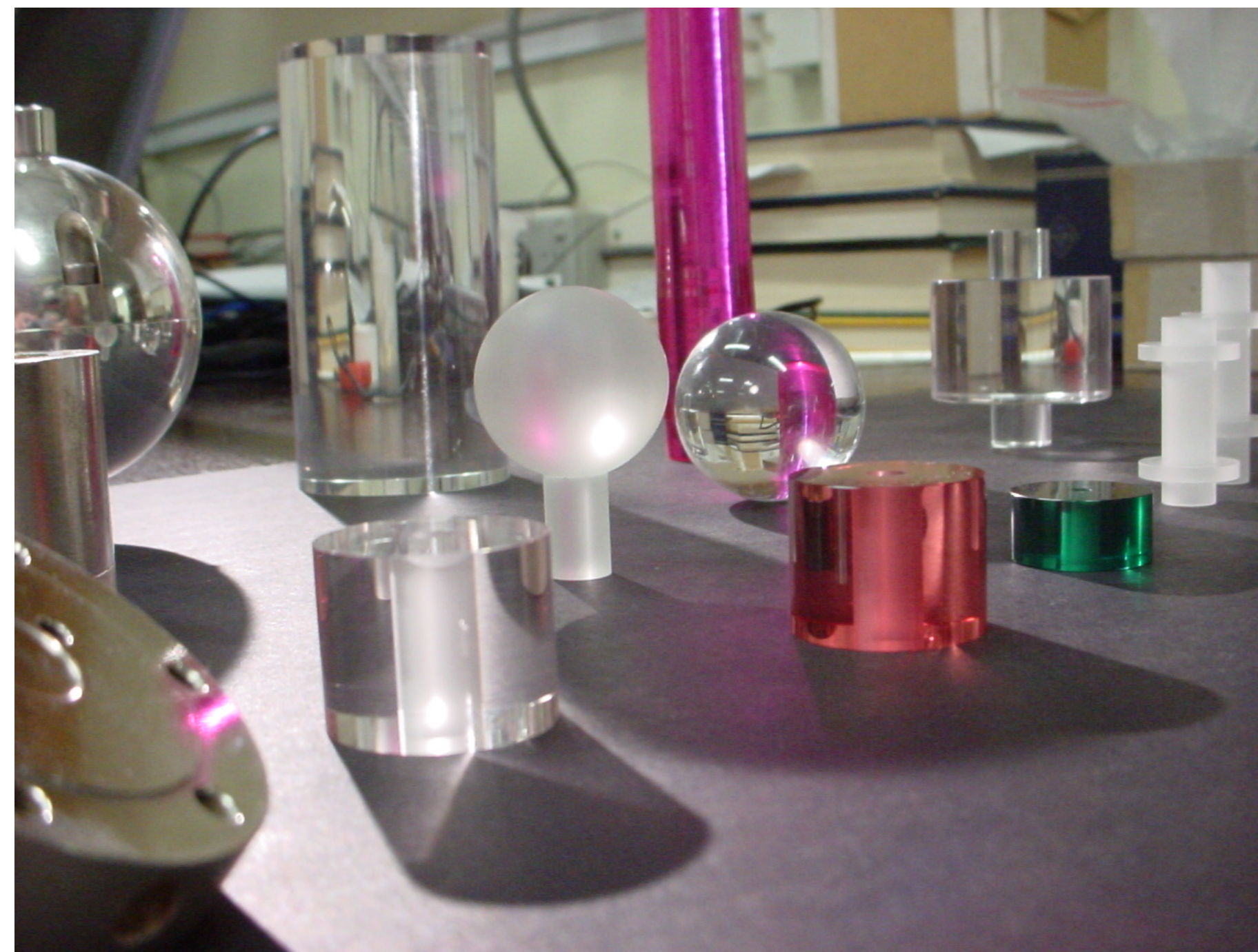
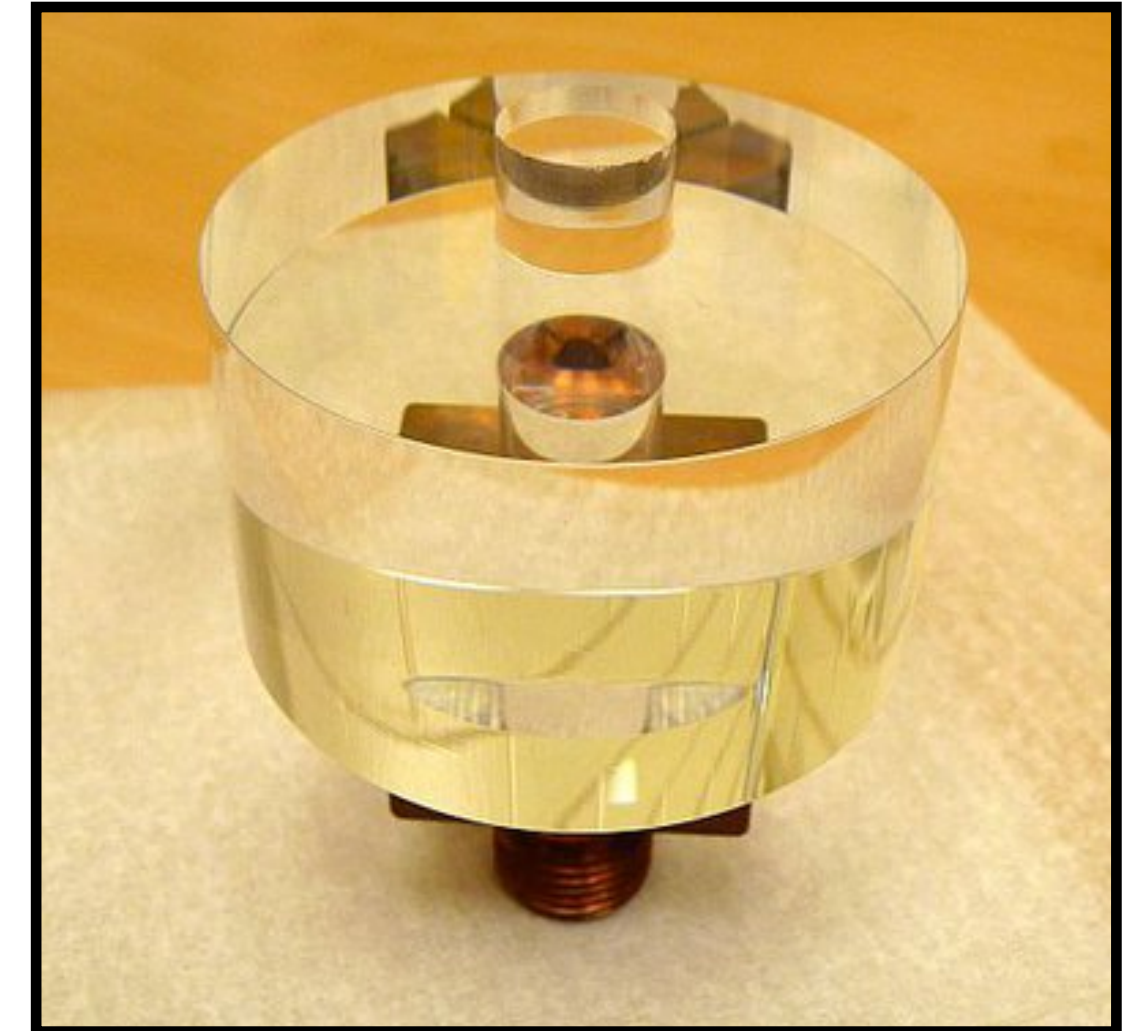


The Anyon Cavity Resonator

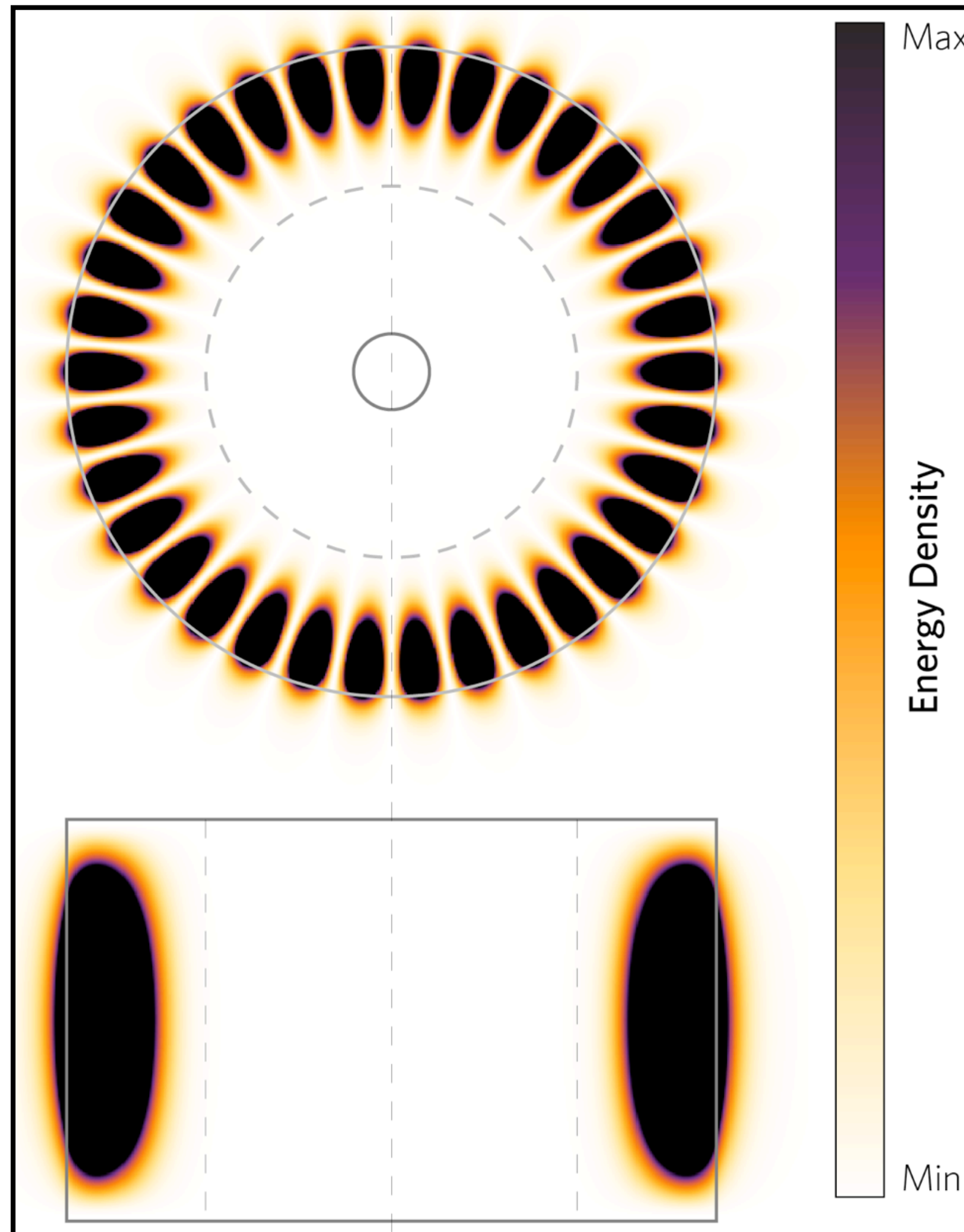


@ UWA: Photonic High-Q Sapphire Resonators

- Crystal Dielectric Resonators
 - In general low loss tangent and anisotropic
 - High-Q
 - Permittivity ~ 10 \rightarrow Whispering Gallery Mode (WGM)



Whispering Gallery Modes: $m > 0$



$WGH_{m,1,1}$ – transverse magnetic, E_z and H_r

$WGE_{m,1,1}$ – transverse electric, H_z and E_r

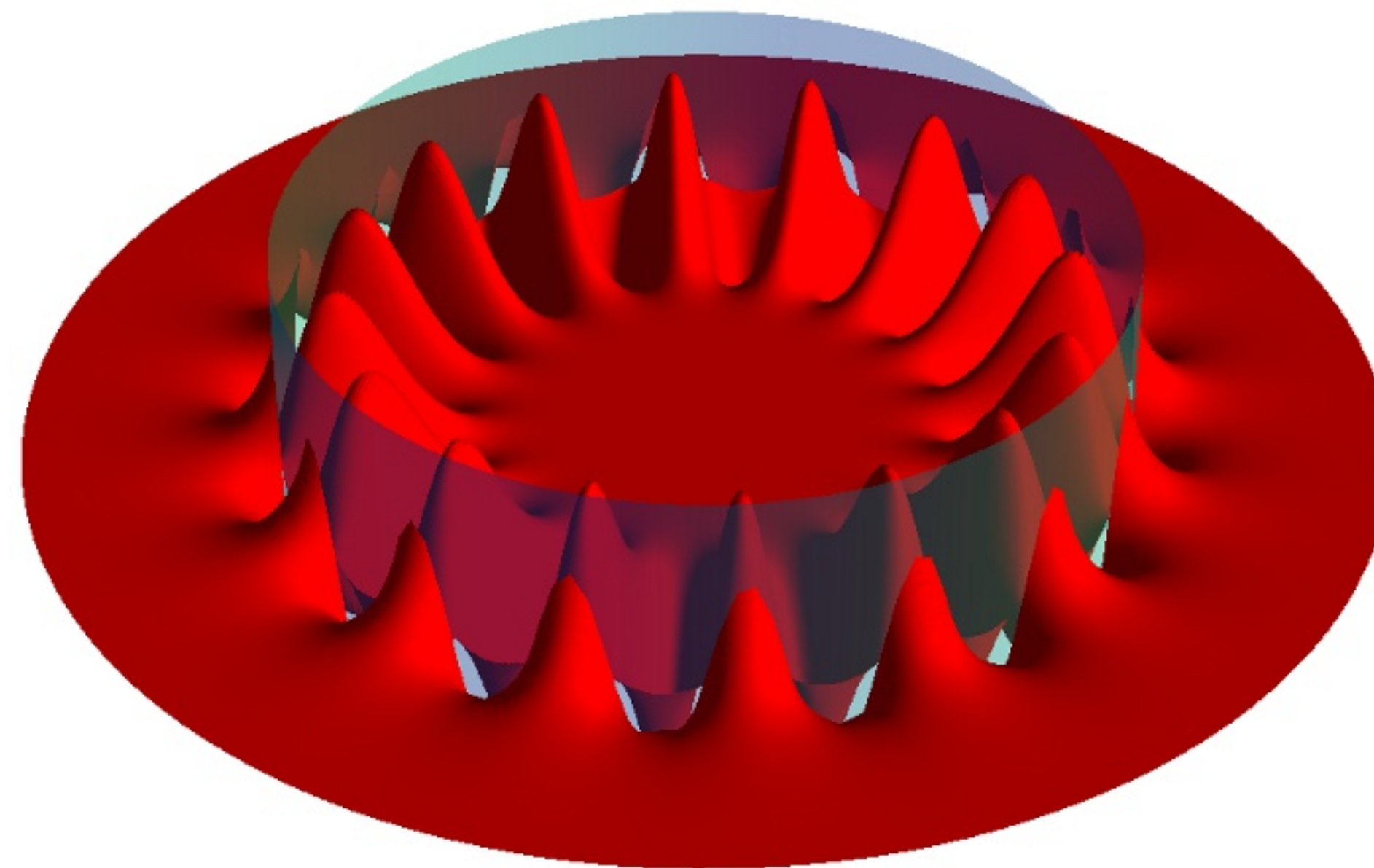
High-Q low temperature sensitivity

-> High stability oscillators

-> Low phase noise oscillators

-> Test fundamental physics

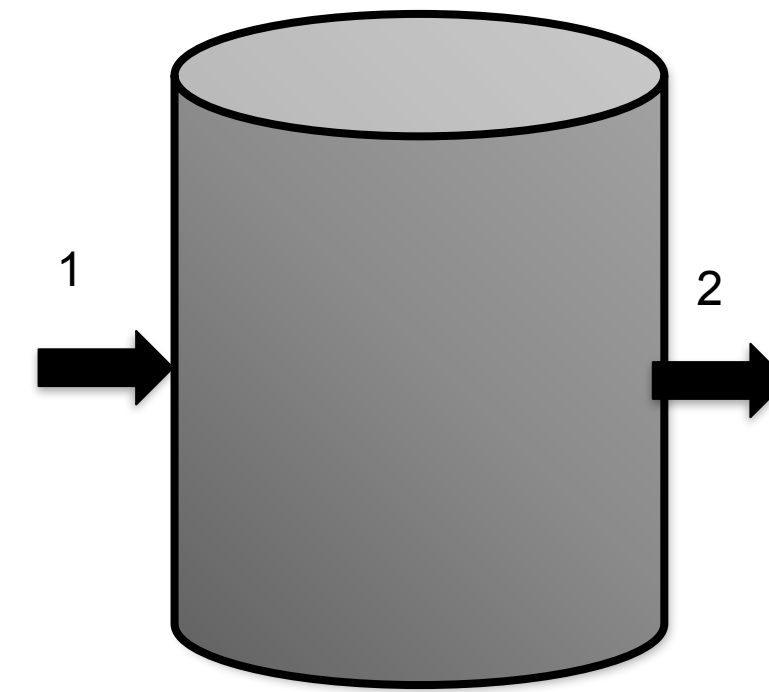
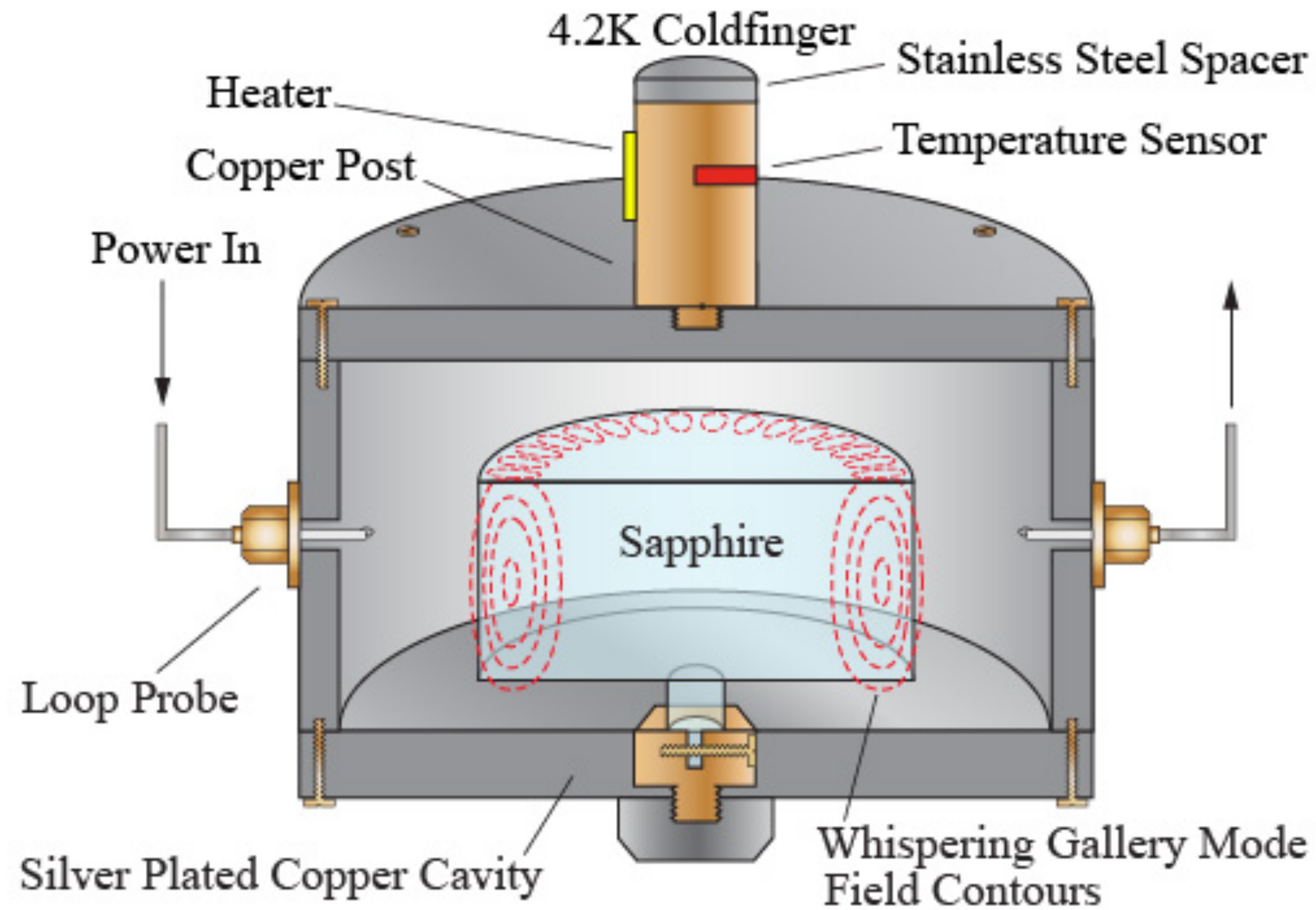
-> Engineered Quantum Systems



ME Tobar and AG Mann, "Resonant frequencies of higher order modes in cylindrical anisotropic dielectric resonators" IEEE Trans. on MTT, vol 39, no. 12, pp. 2077-2083, Dec. 1991.

Example: Cryogenic Sapphire Oscillators (CSO) based on WGM

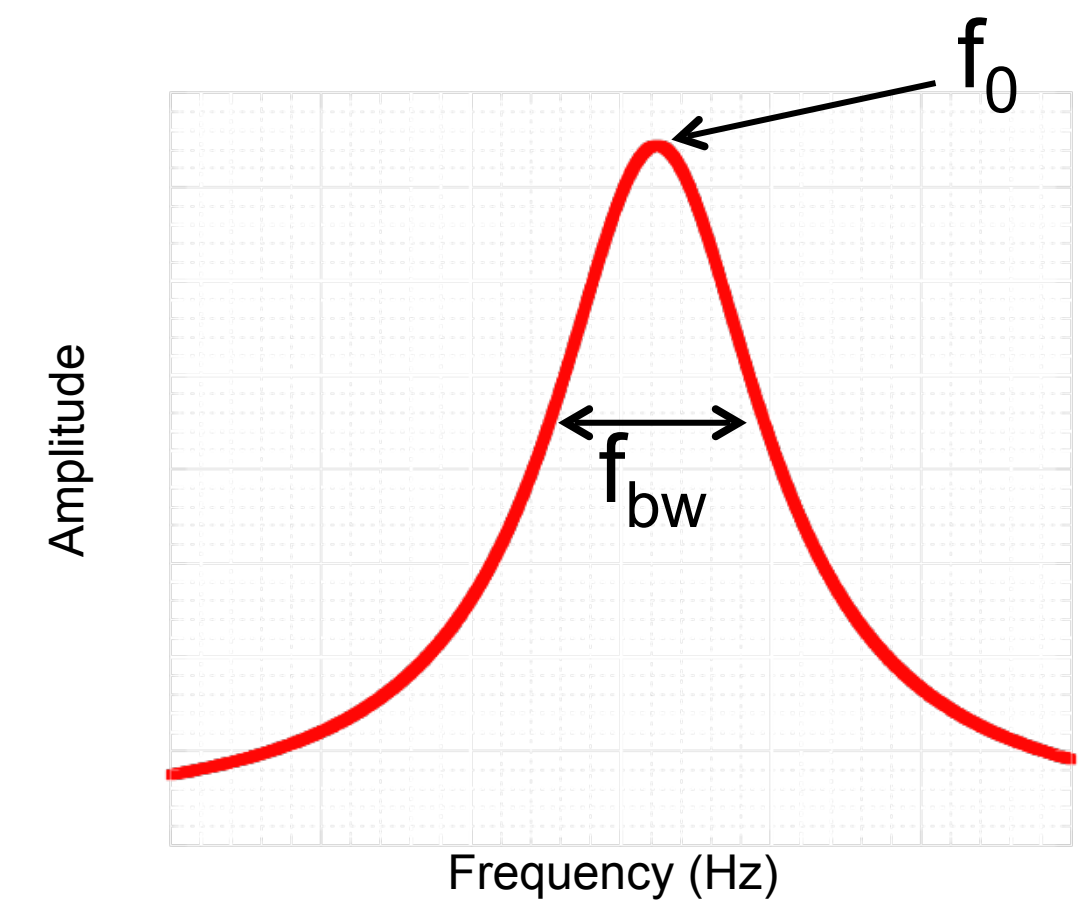
Precision Frequency, Phase and Time



$$Q_L = f_0 / f_{bw}$$

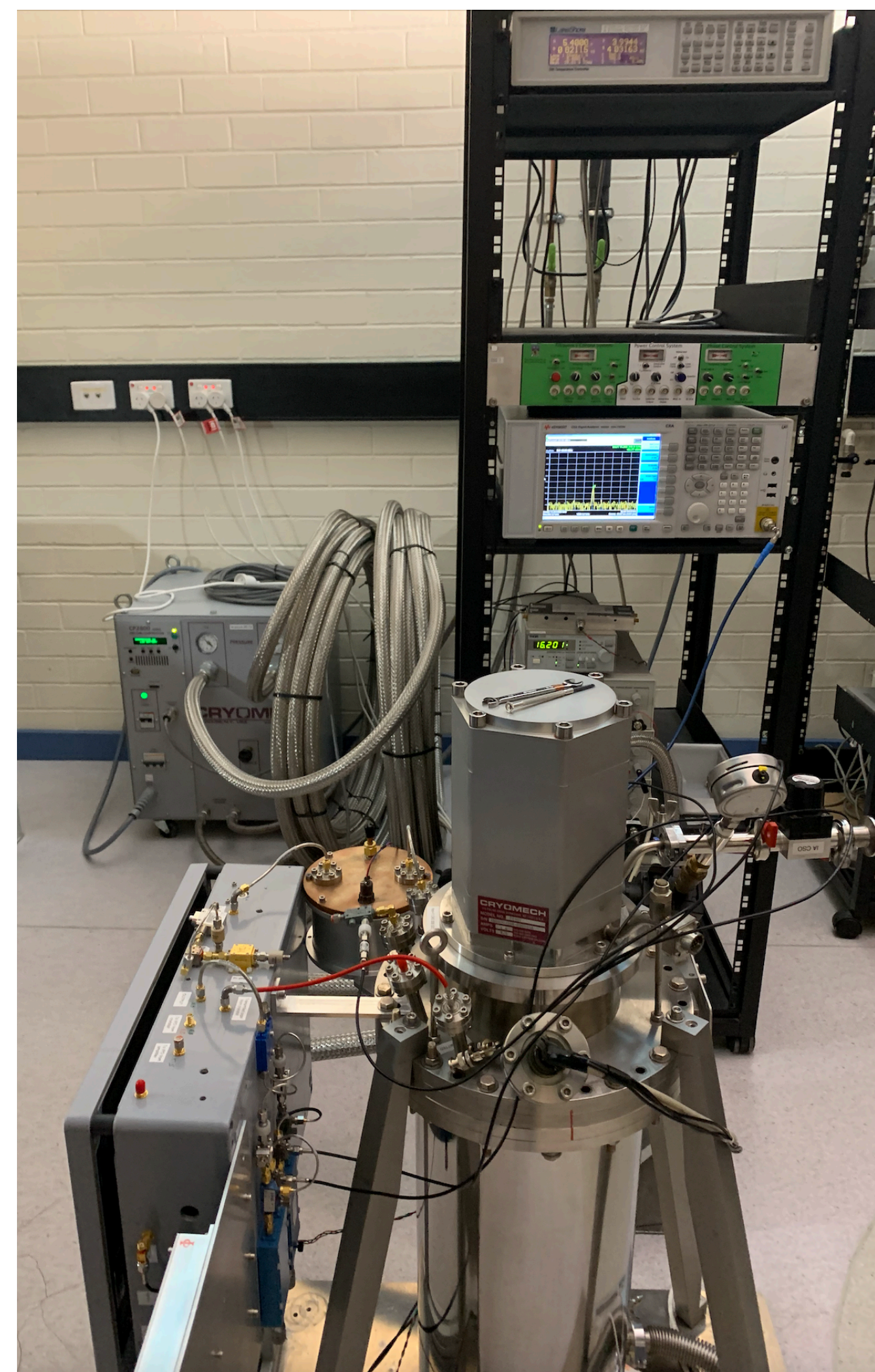
$$Q_L = 10^9 \quad f_0 = 10 \text{ GHz} \quad f_{bw} = 10 \text{ Hz}$$

S_{21} : Transmission Spectrum

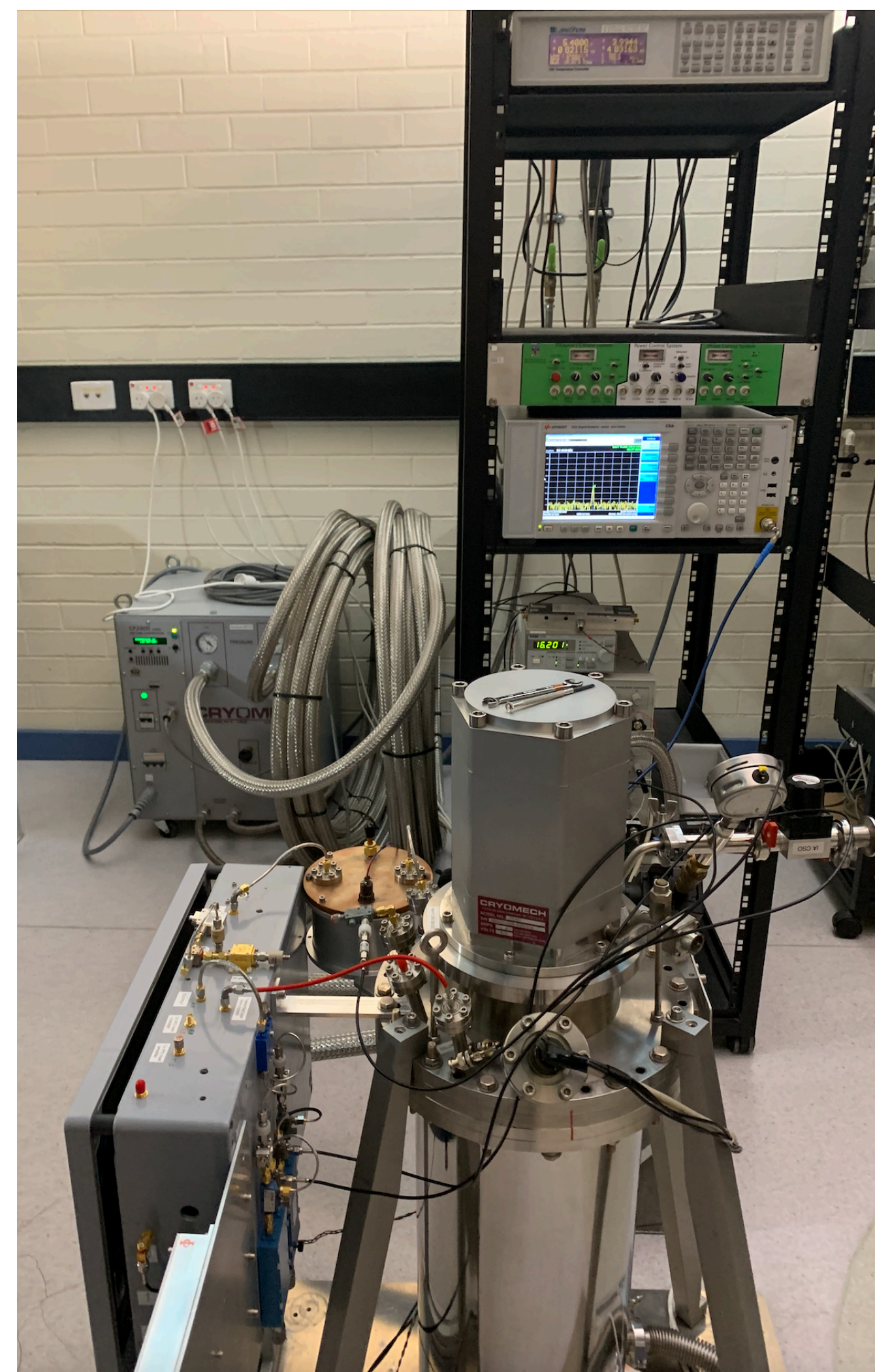
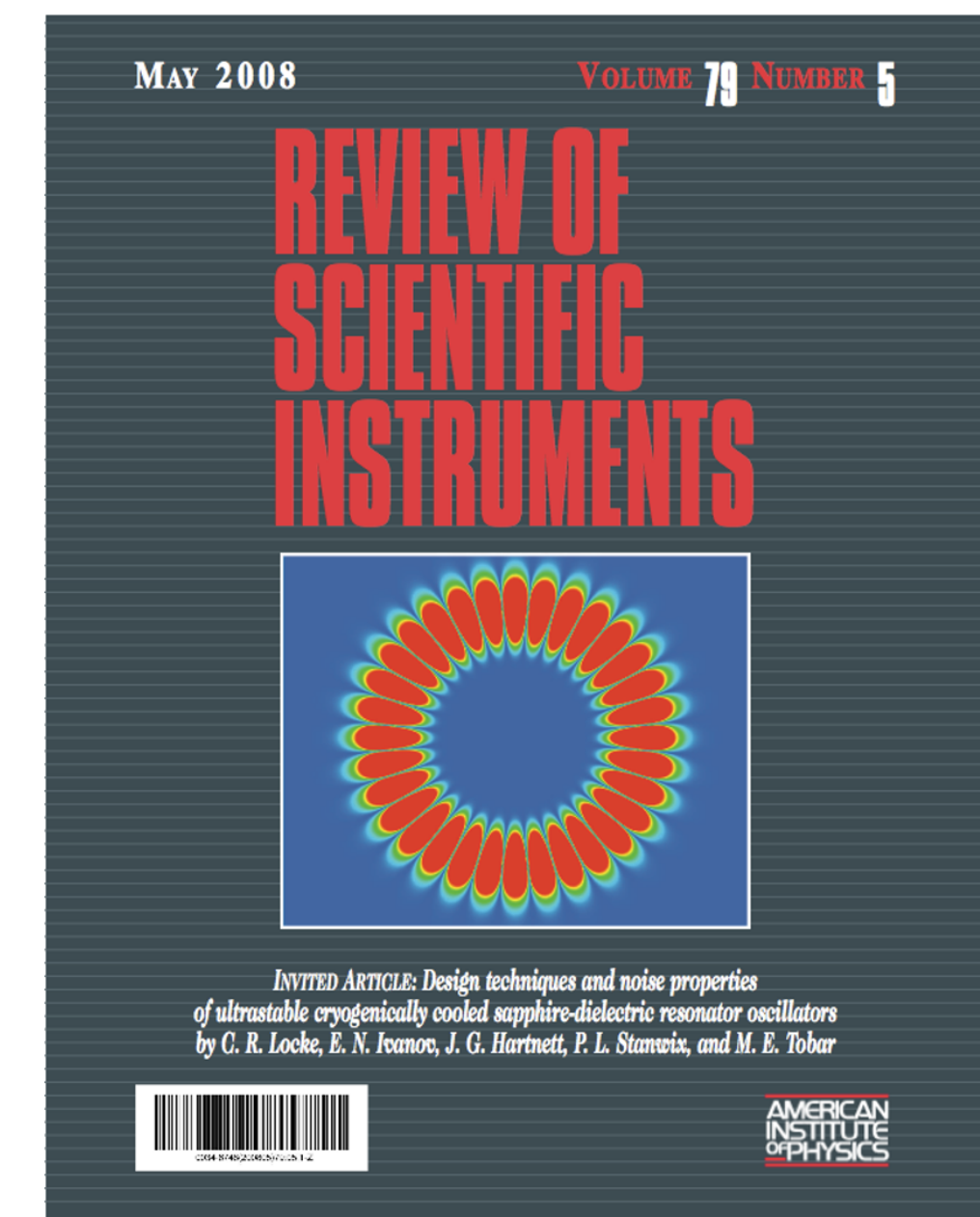


Q-factor Greater than one billion Q-factor at microwave frequencies

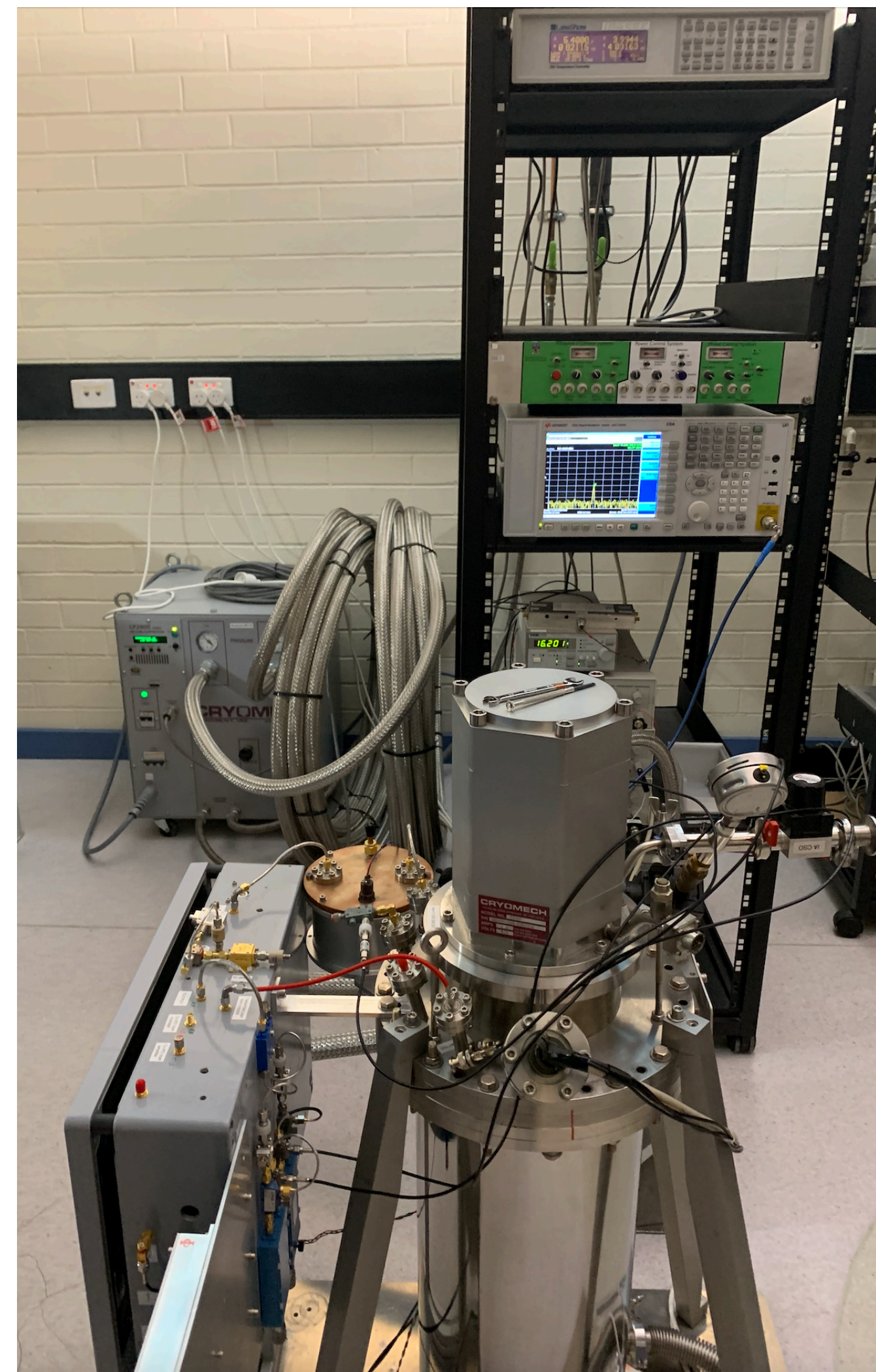
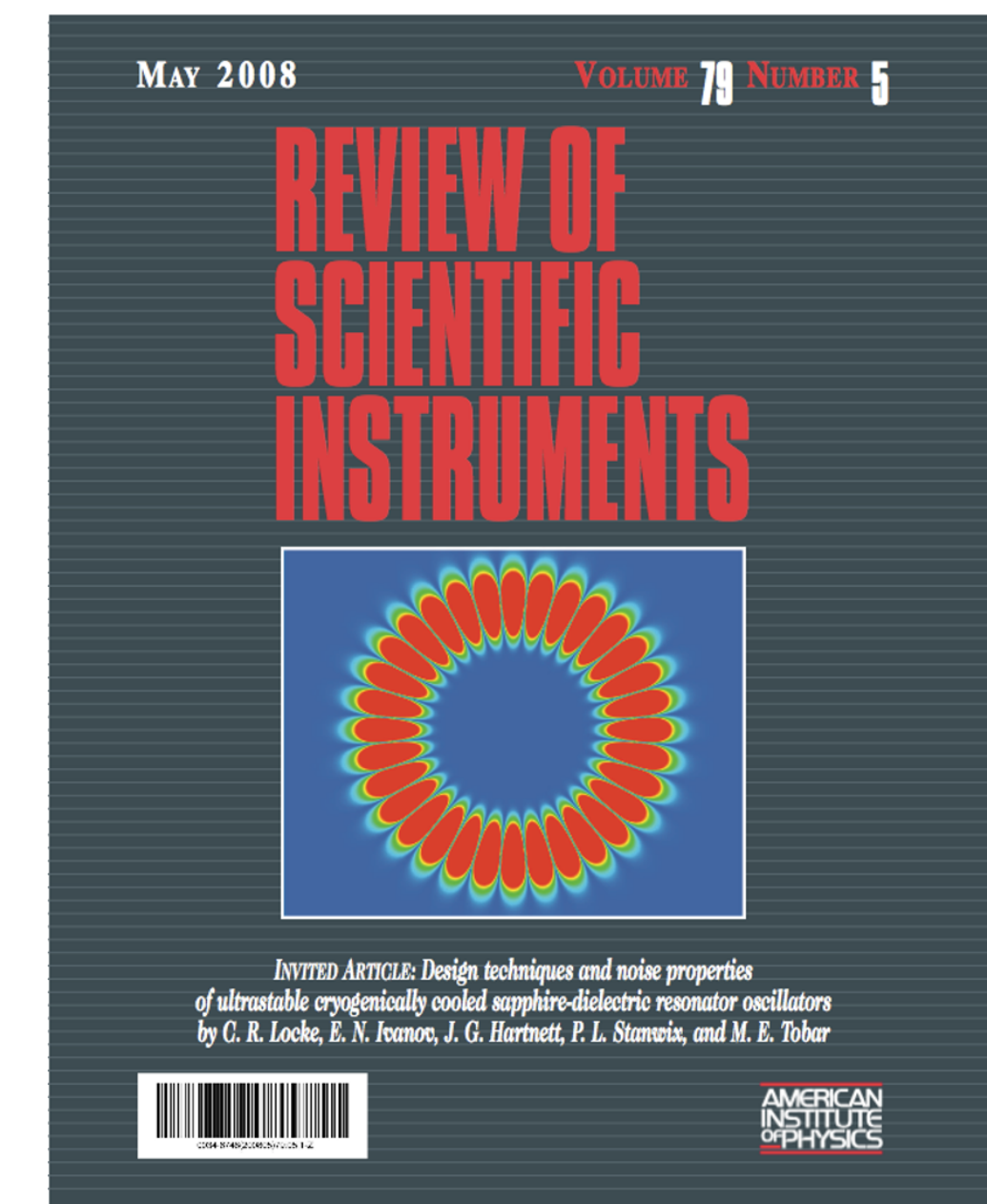
Cryogenic Sapphire Oscillators (CSO) 1989-Now



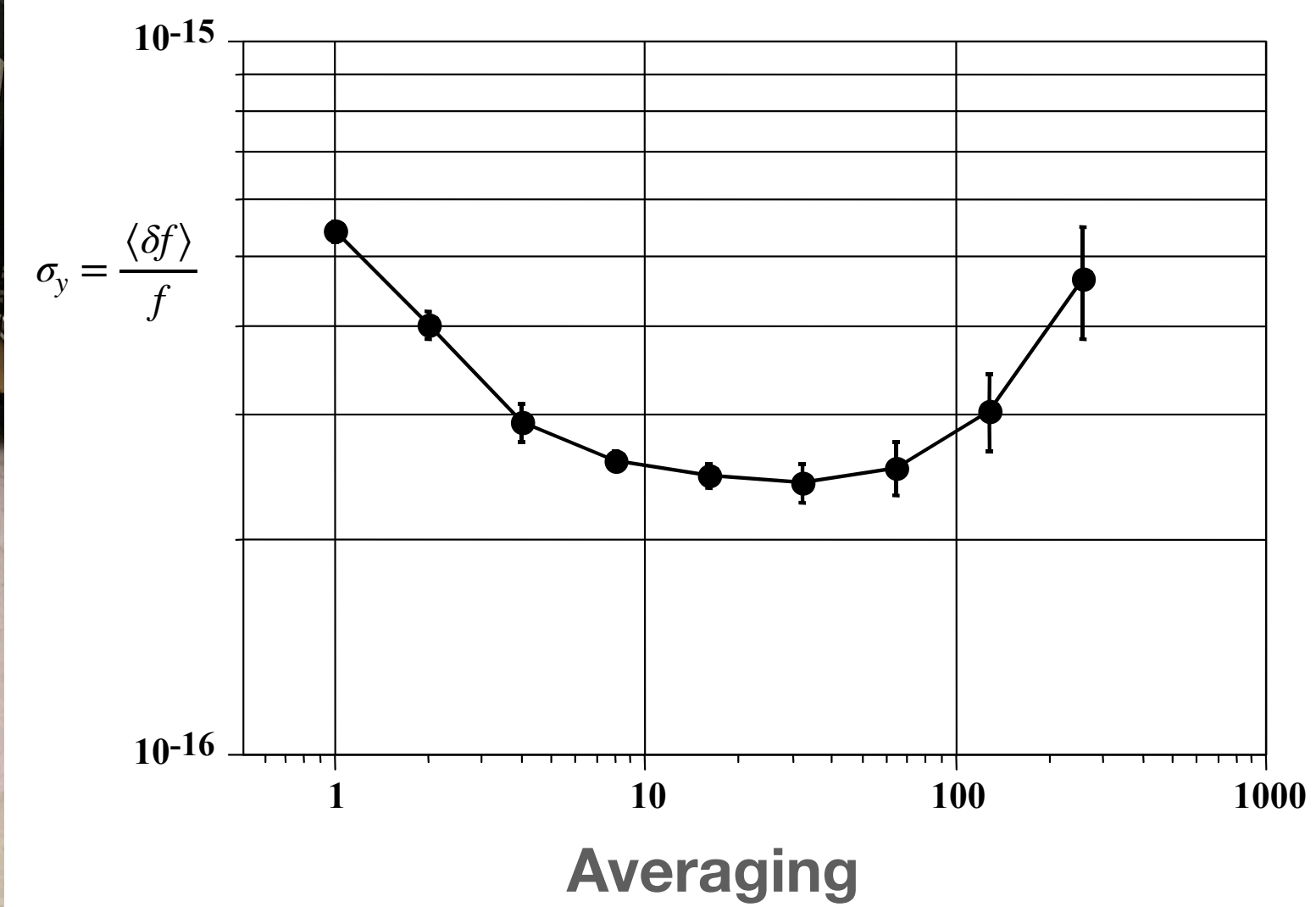
Cryogenic Sapphire Oscillators (CSO) 1989-Now



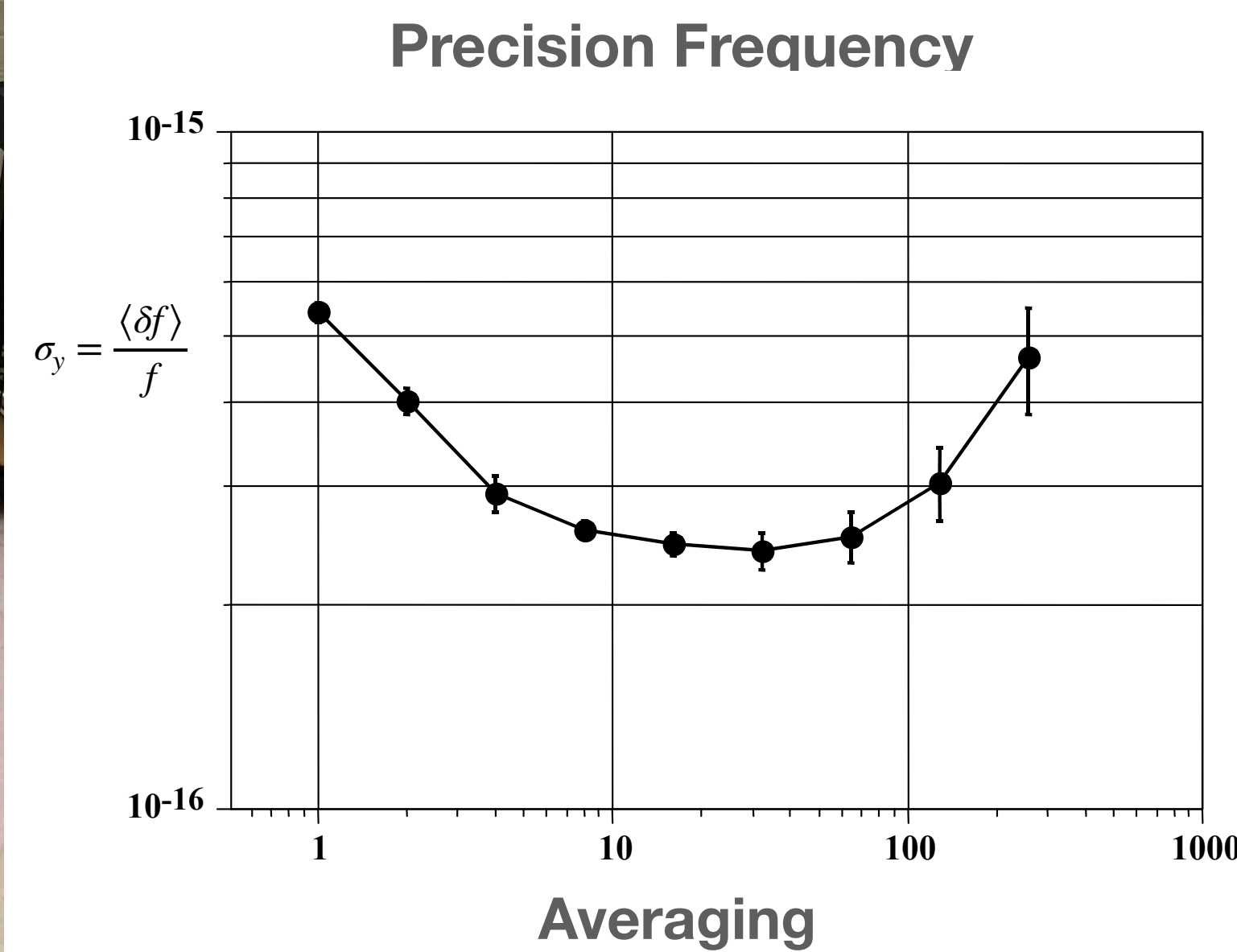
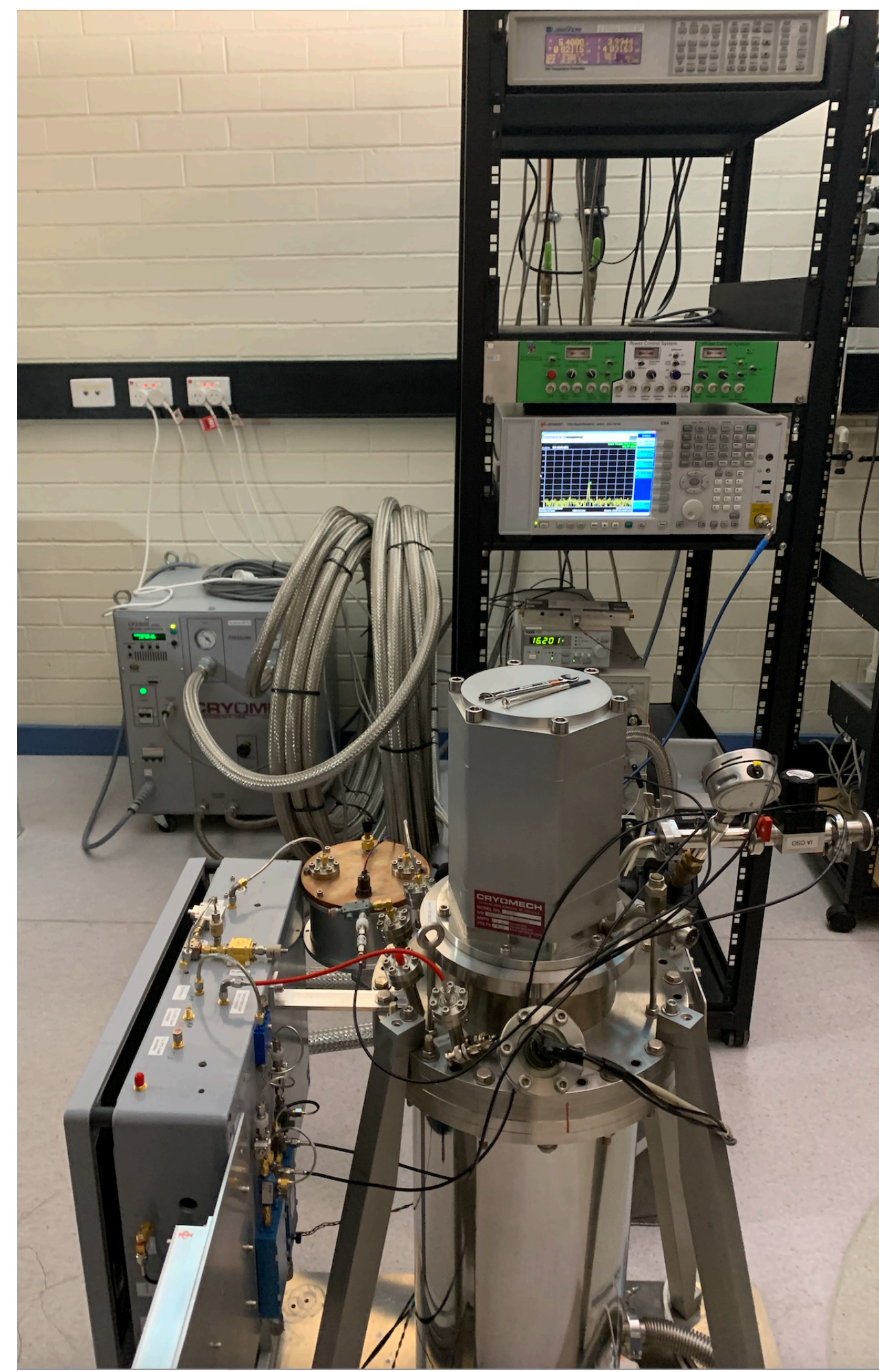
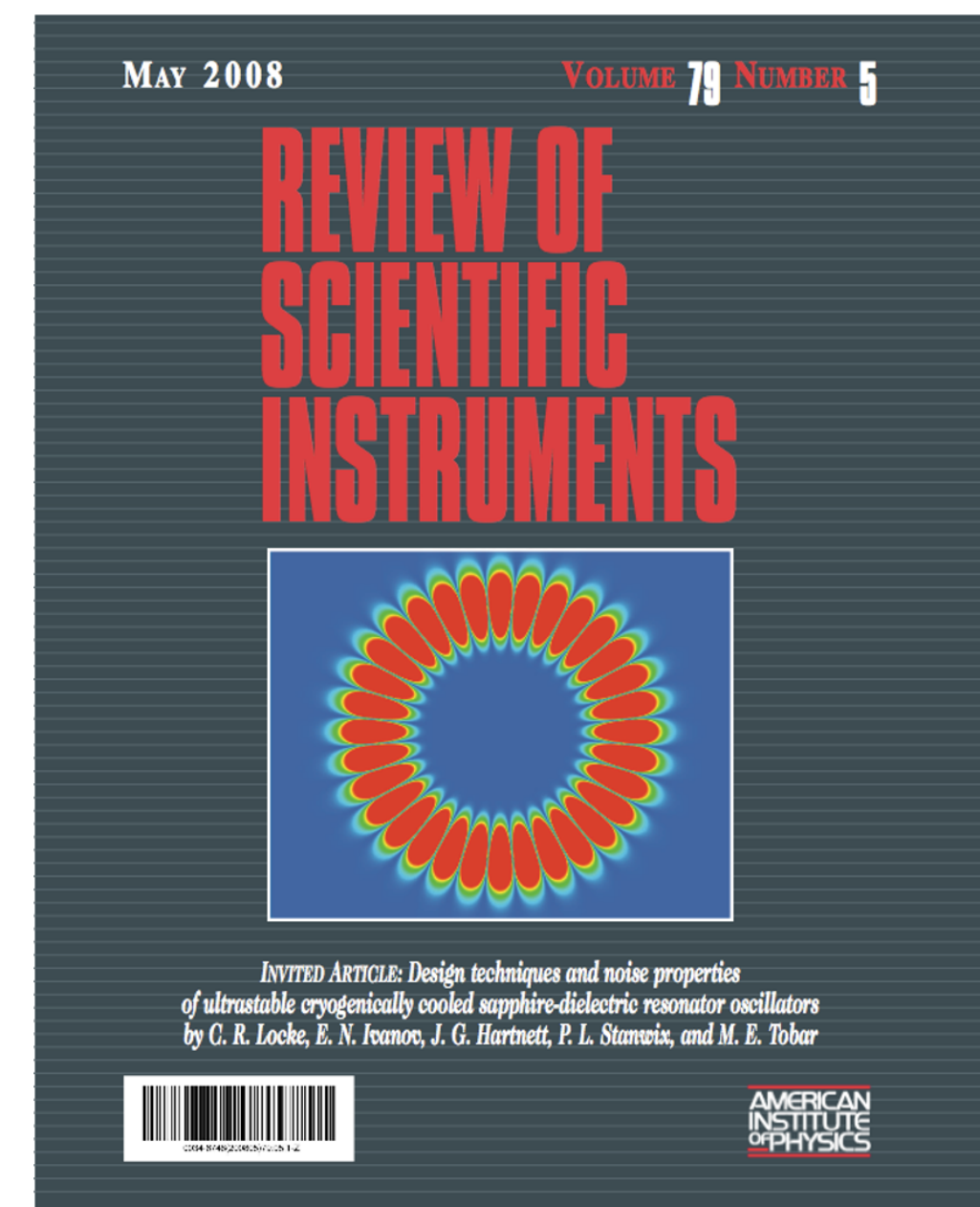
Cryogenic Sapphire Oscillators (CSO) 1989-Now



Precision Frequency



Cryogenic Sapphire Oscillators (CSO) 1989-Now



QUANTX

HOME INDUSTRIES PRODUCTS COMPANY RESOURCES

CRYOCLOCK

The purest frequency signal source

CRYOCLOCK

Ultra-stable time and frequency signals

Cryoclock is the culmination of 20 years of research and now delivers the world's purest microwave and RF signals in a reliable and autonomous package that is user friendly. Its output signals possess both ultra-low noise phase noise as well as ultra-high short and medium term frequency stability. Our 'turn-key' products are continuously-operating, fully autonomous, self-contained, and do not require any cryogenics knowledge. It has a low

femto ENGINEERING
from science to society

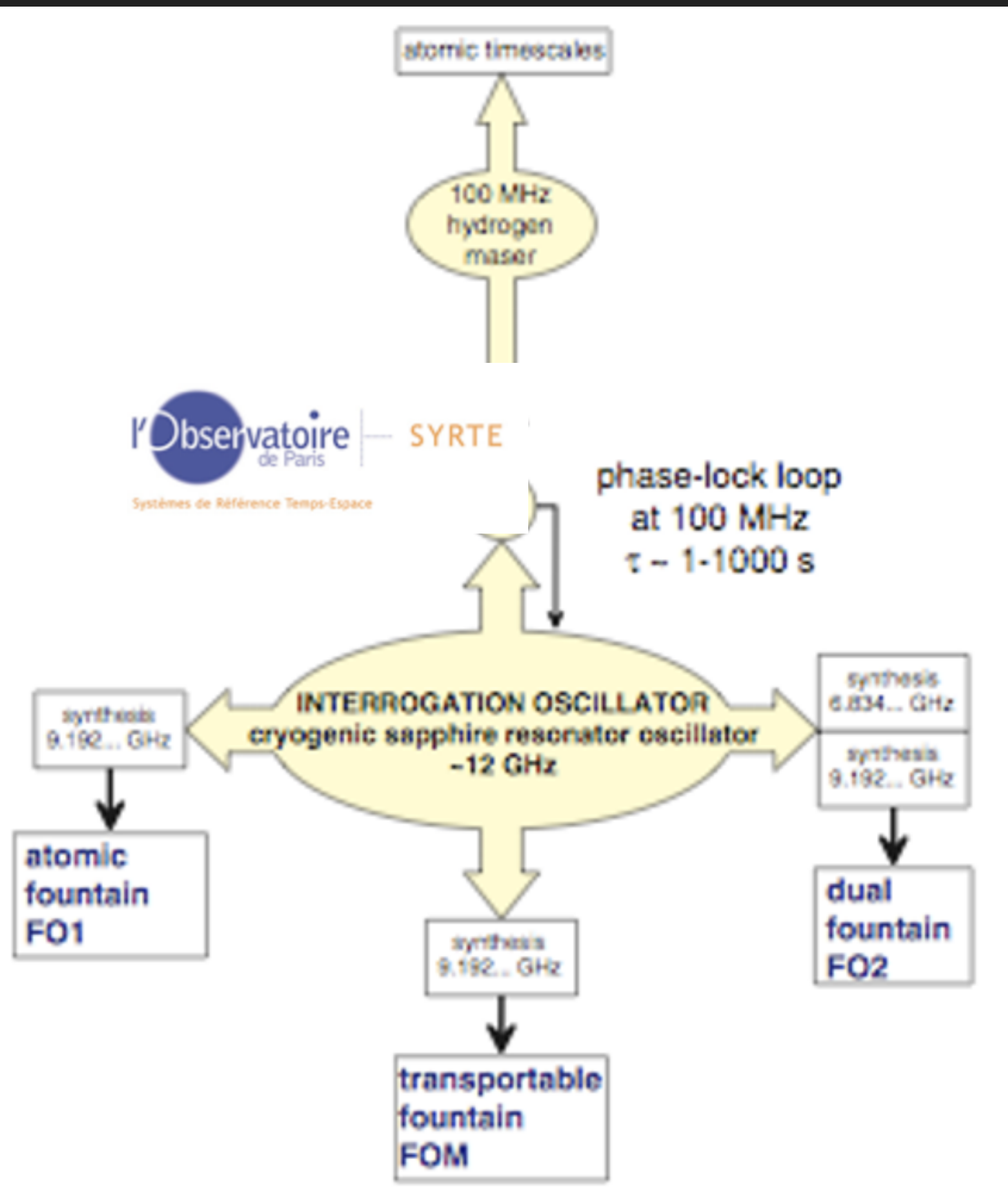
HOME PAGE ABOUT US OUR ACTIVITIES NEWS CONTACT

Nos activités : Radio frequency

ULISS Cryogenic Sapphire Oscillator

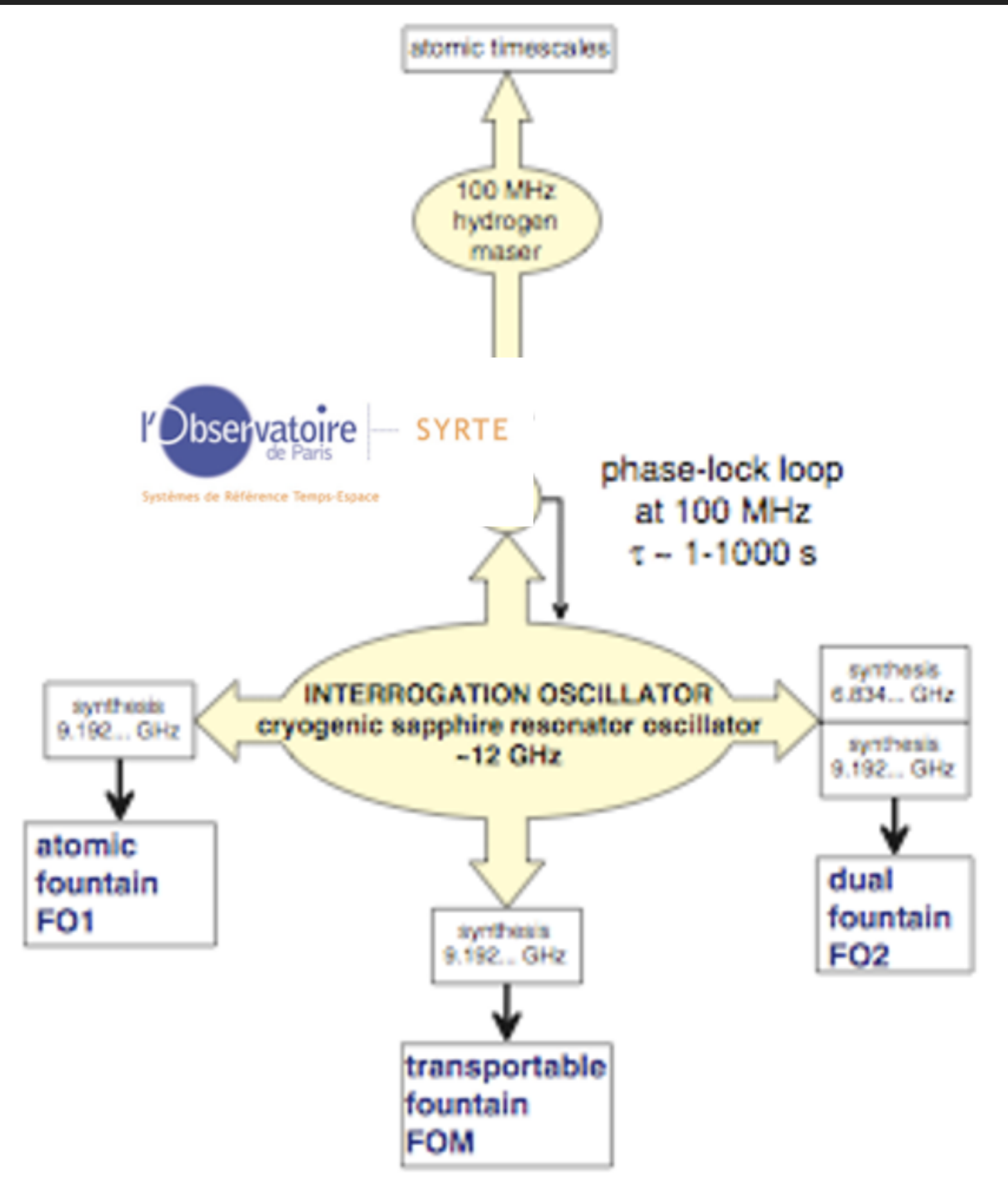
ULISS Cryogenic Sapphire Oscillator (CSO) offers unprecedented frequency stability performance thanks to the exceptional regularity of its heartbeat: a high purity sapphire crystal placed at low temperature in a controlled environment.

CSO Operates Atomic Clock @ the Projection Noise Limit: Paris Observatory SYRTE since 1999

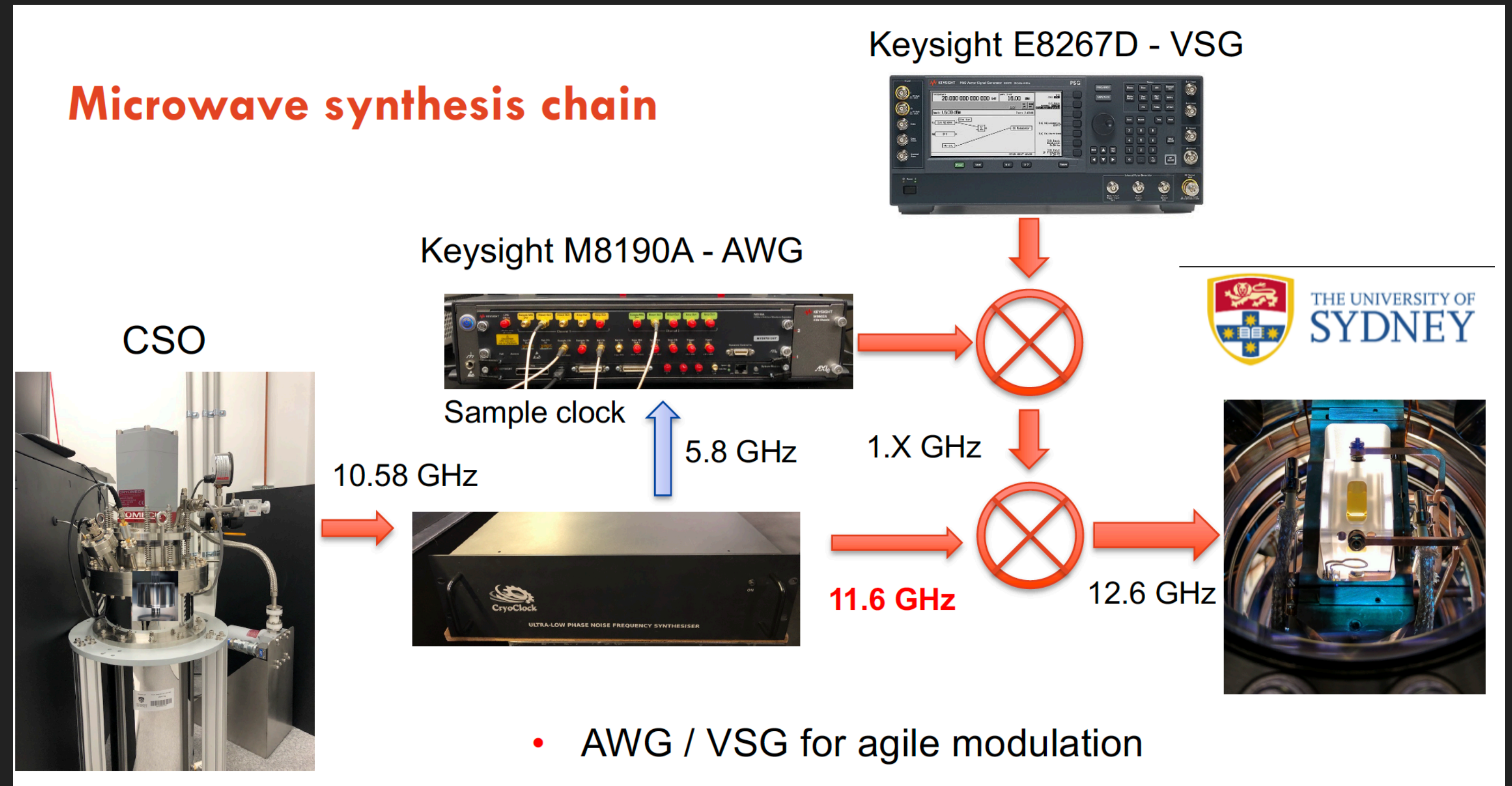


CSO Operates Atomic Clock @ the Projection Noise Limit: Paris Observatory SYRTE since 1999

Now: Atomic or Spin Qubits at the Quantum Limit

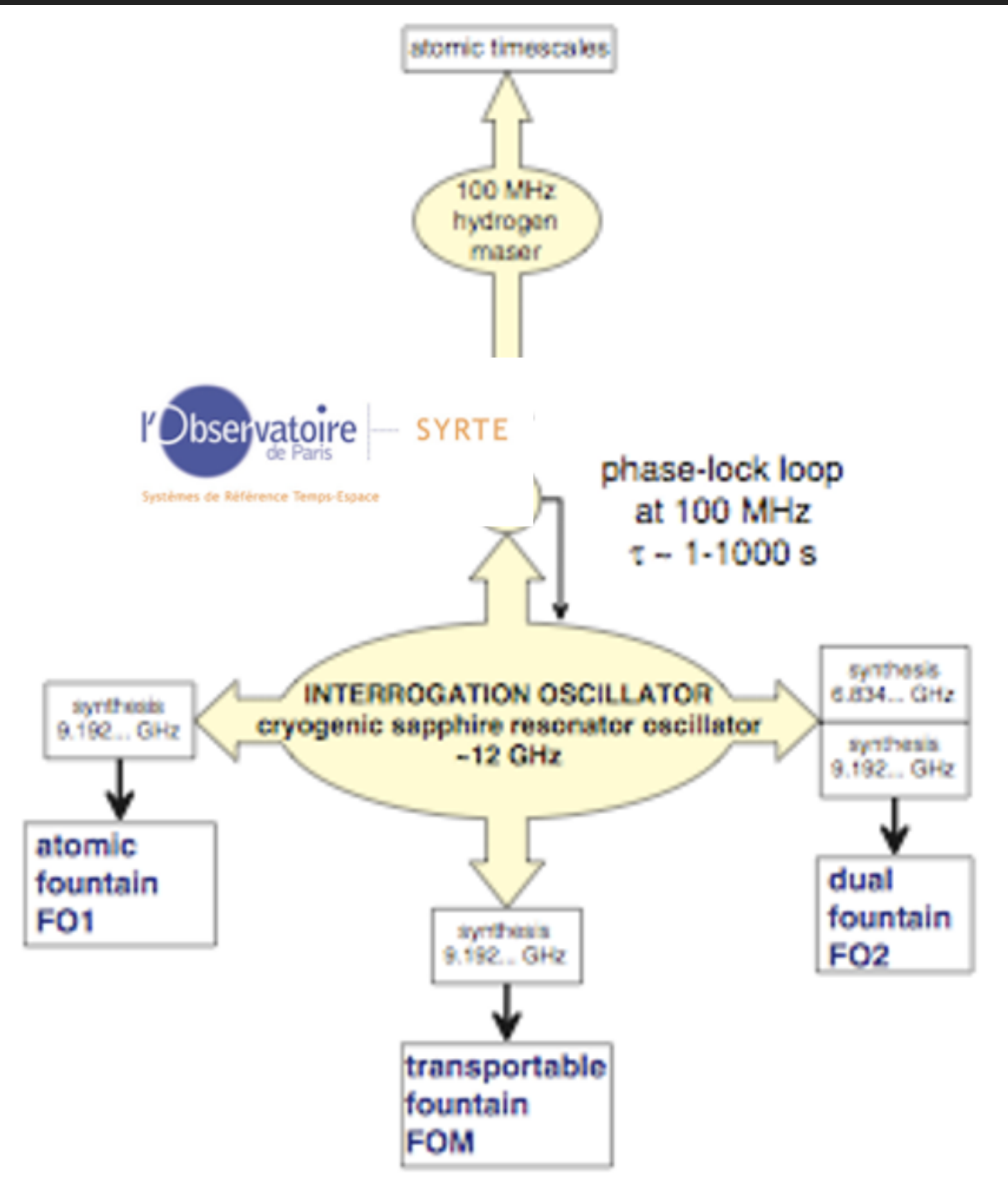


Microwave synthesis chain

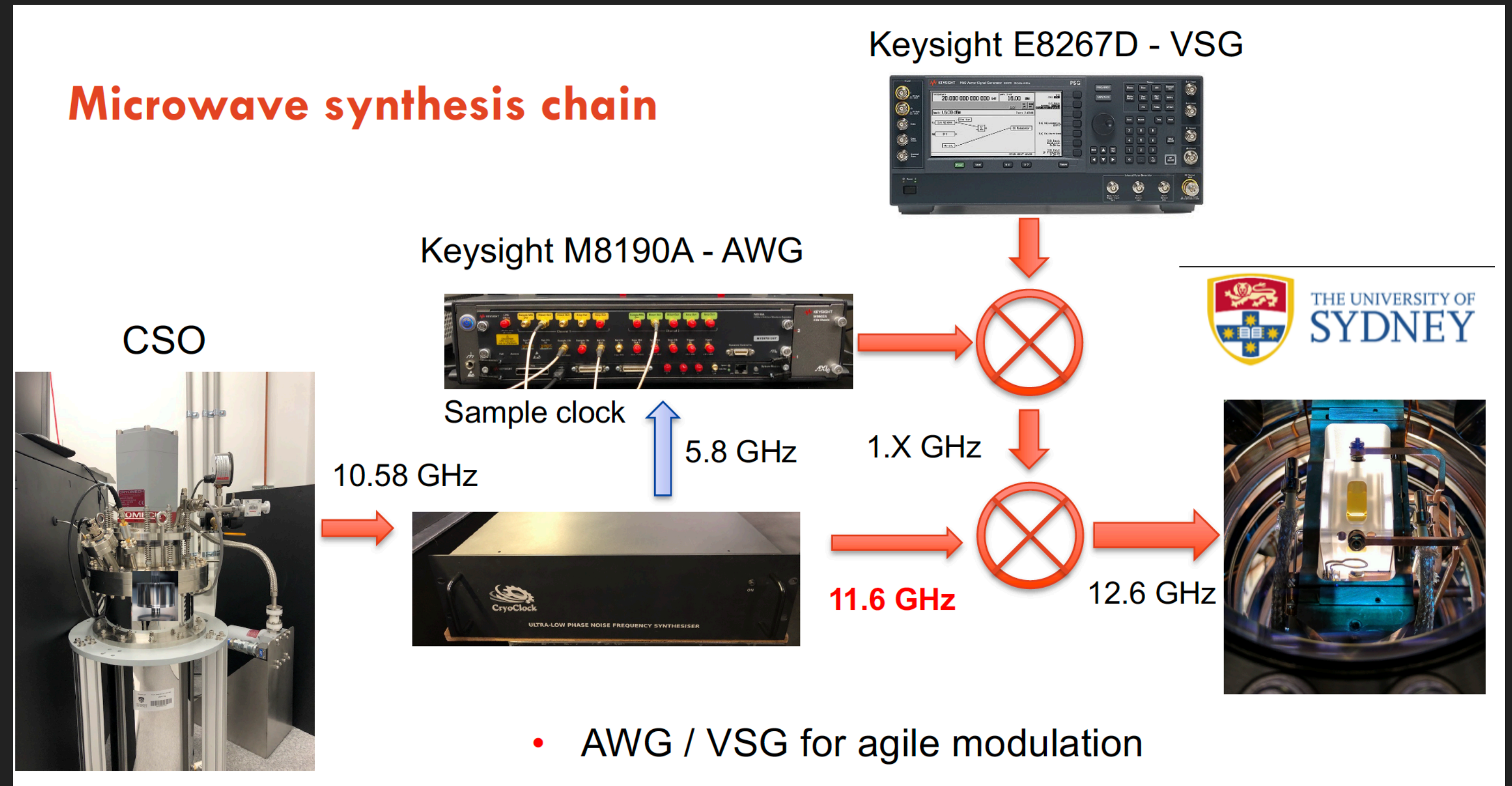


CSO Operates Atomic Clock @ the Projection Noise Limit: Paris Observatory SYRTE since 1999

Now: Atomic or Spin Qubits at the Quantum Limit



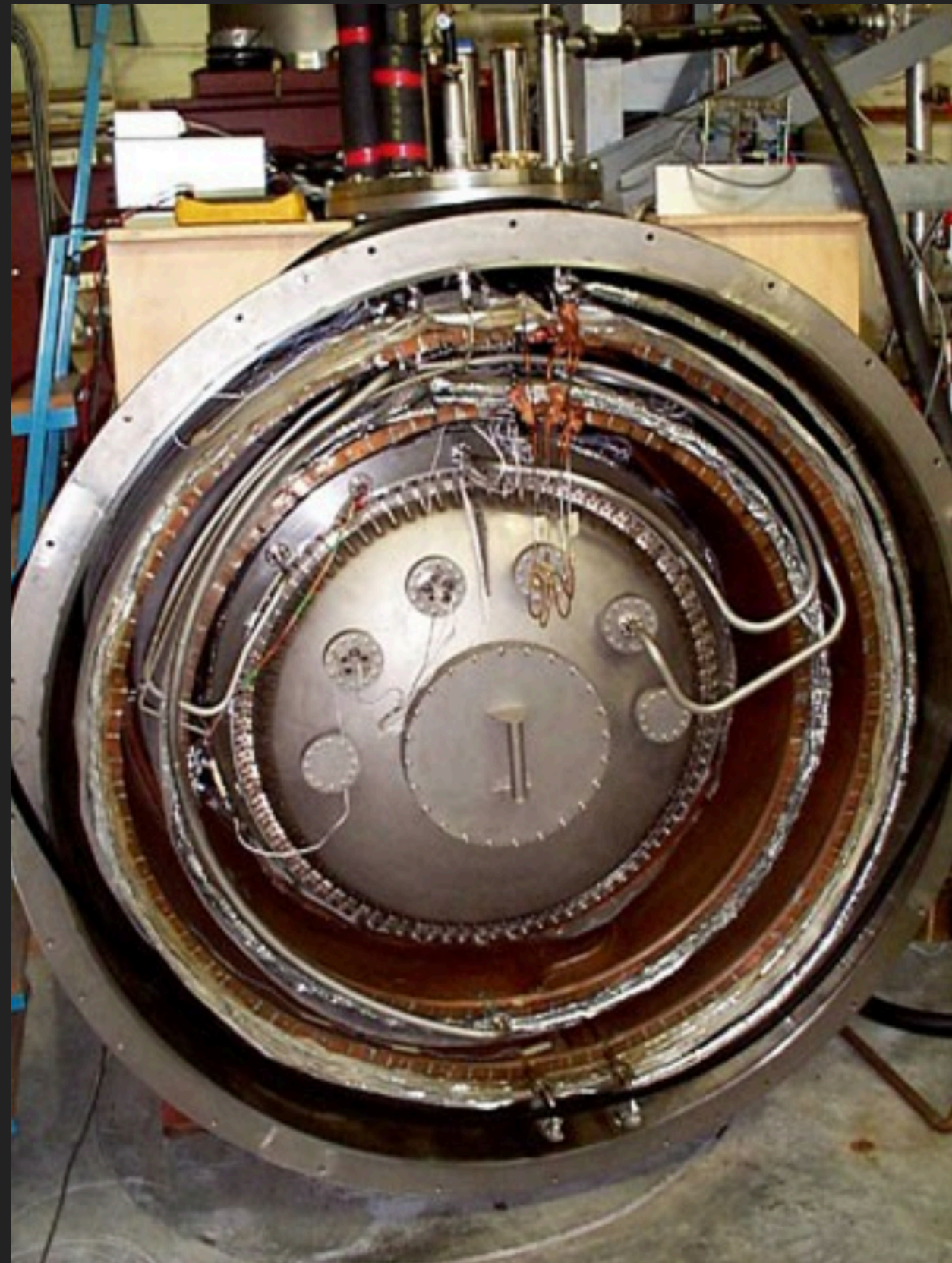
Microwave synthesis chain



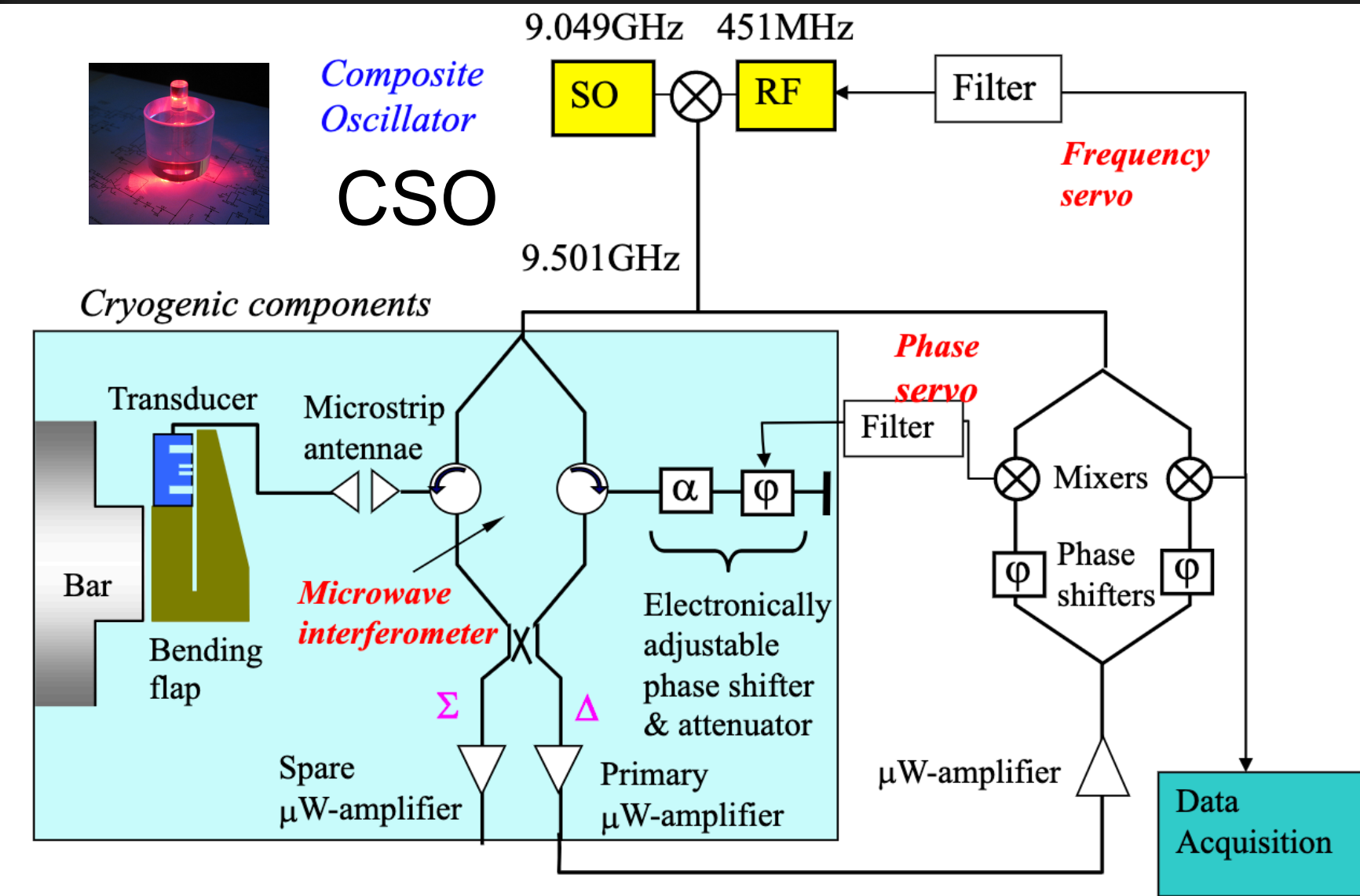
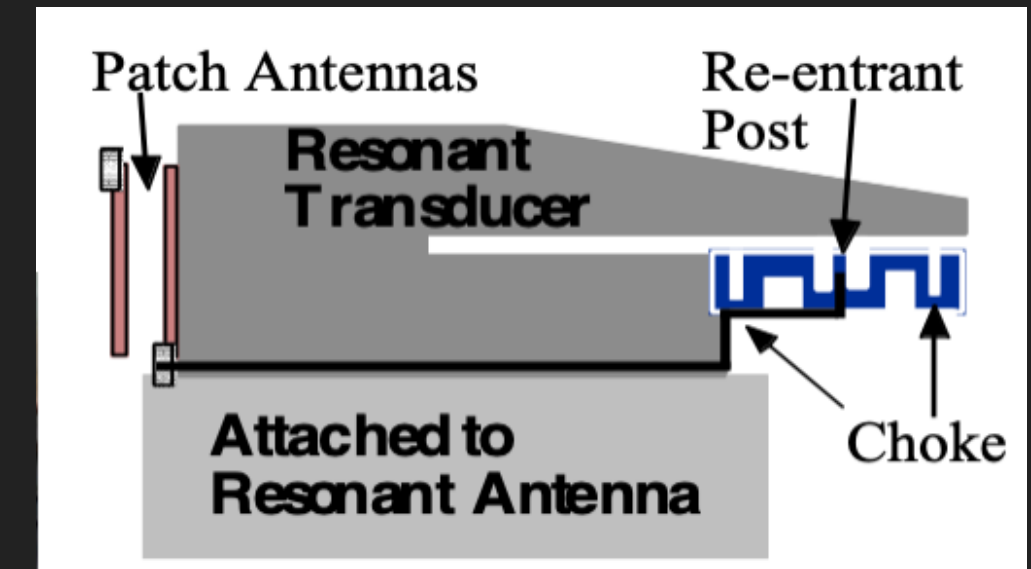
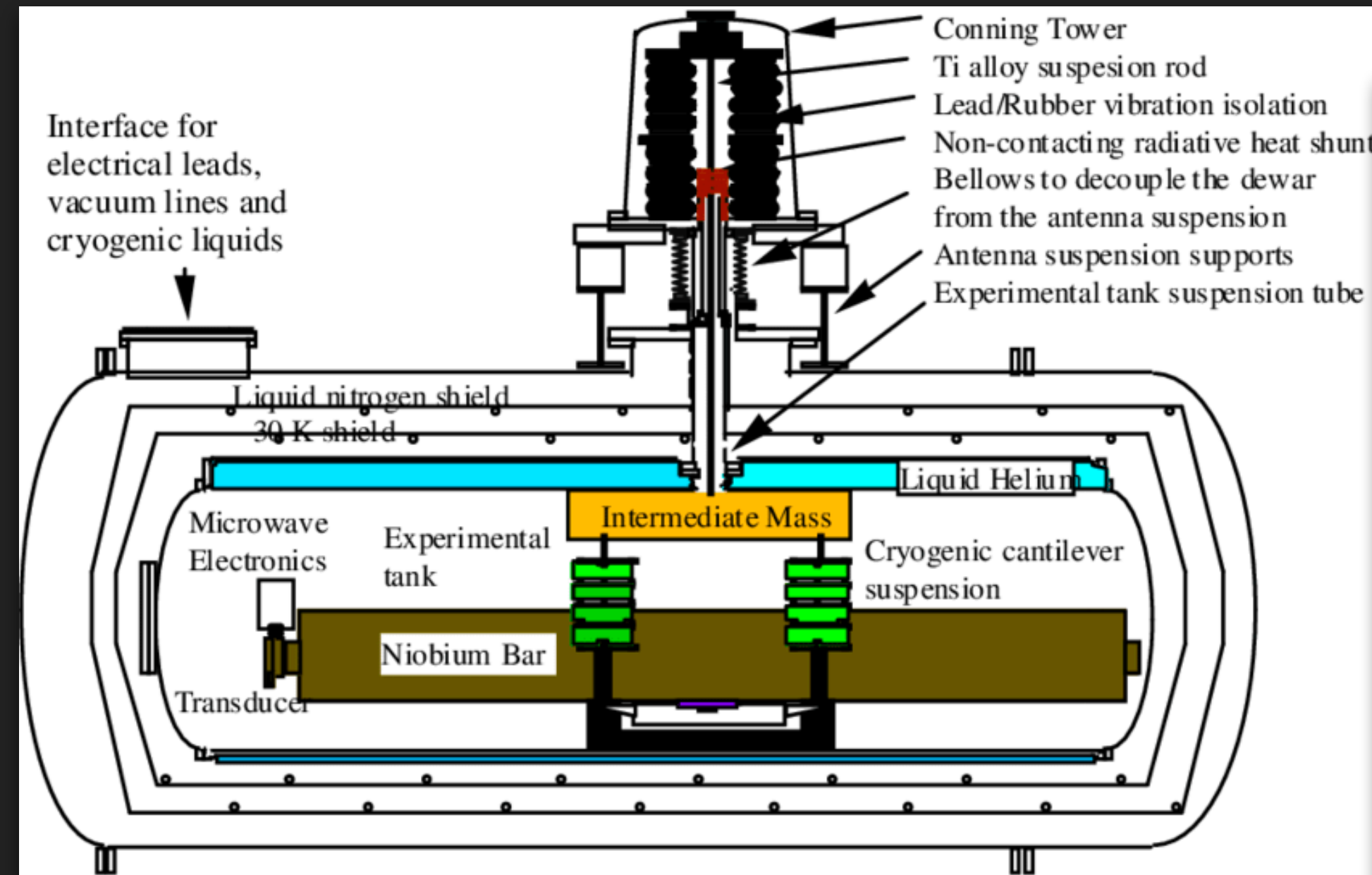
Yb ion qubits.

Also, other systems will need low phase noise pump oscillators

Cryogenic Resonant Bar Gravity Wave Detectors: ultra precise optomechanical displacement measurement: needs low phase noise oscillator



Phys. Rev. Lett 74, 1908 (1995)



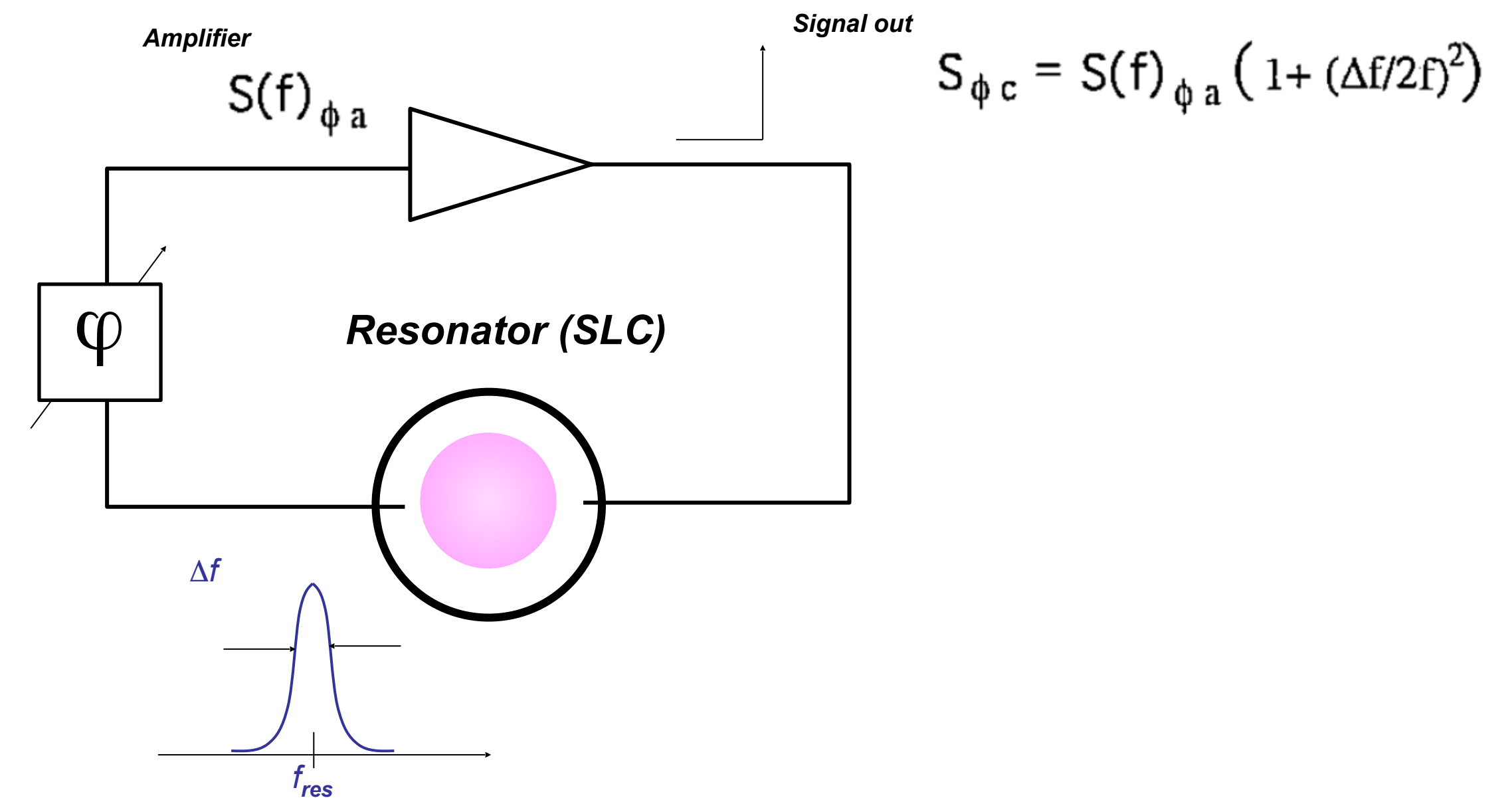
Low Phase Noise Oscillators

Lesson's Model

Noise in oscillator depends on the amplifier phase noise

and Cavity Q-factor = $f_0/\Delta f$

(f is Fourier frequency)



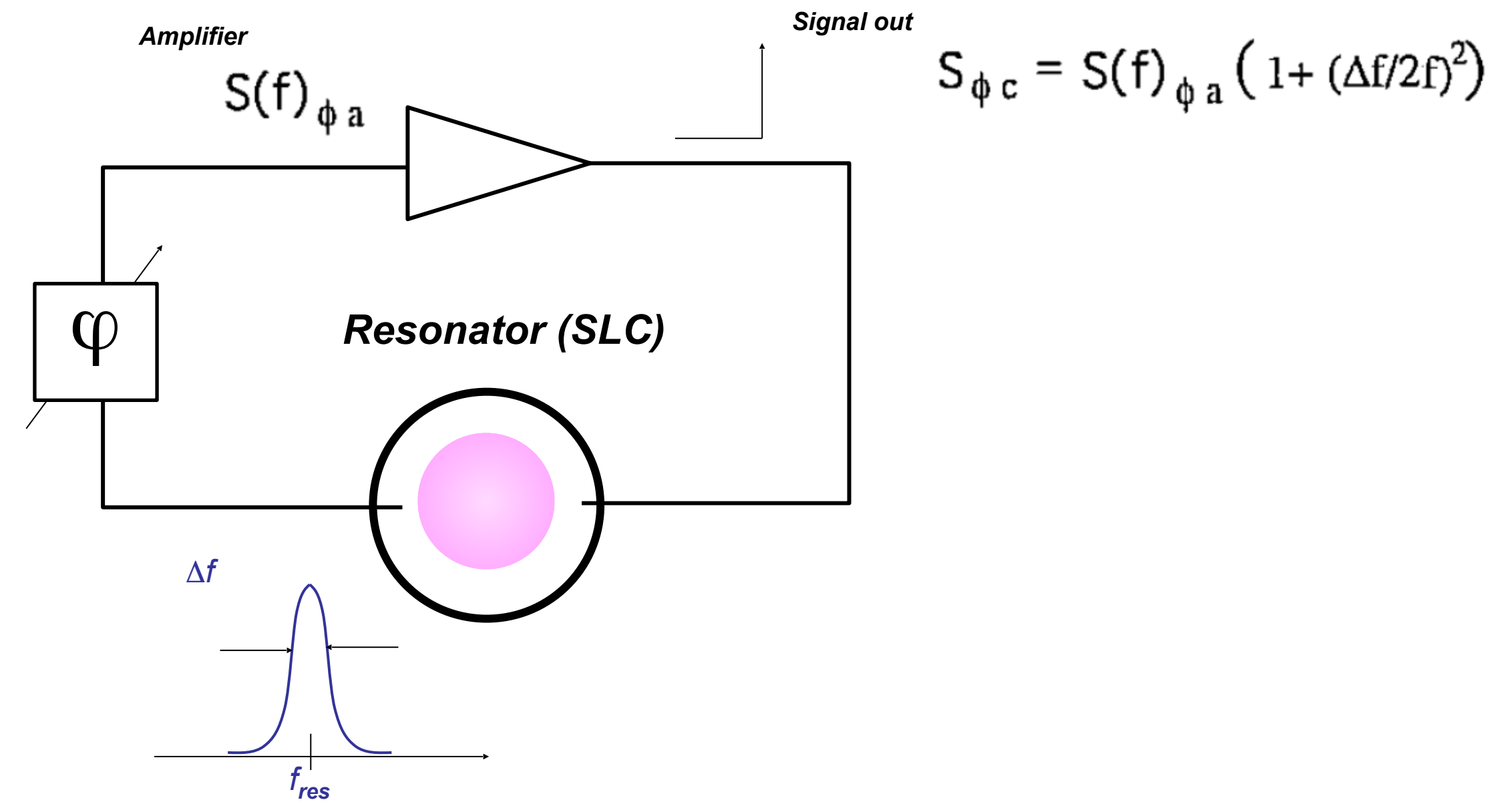
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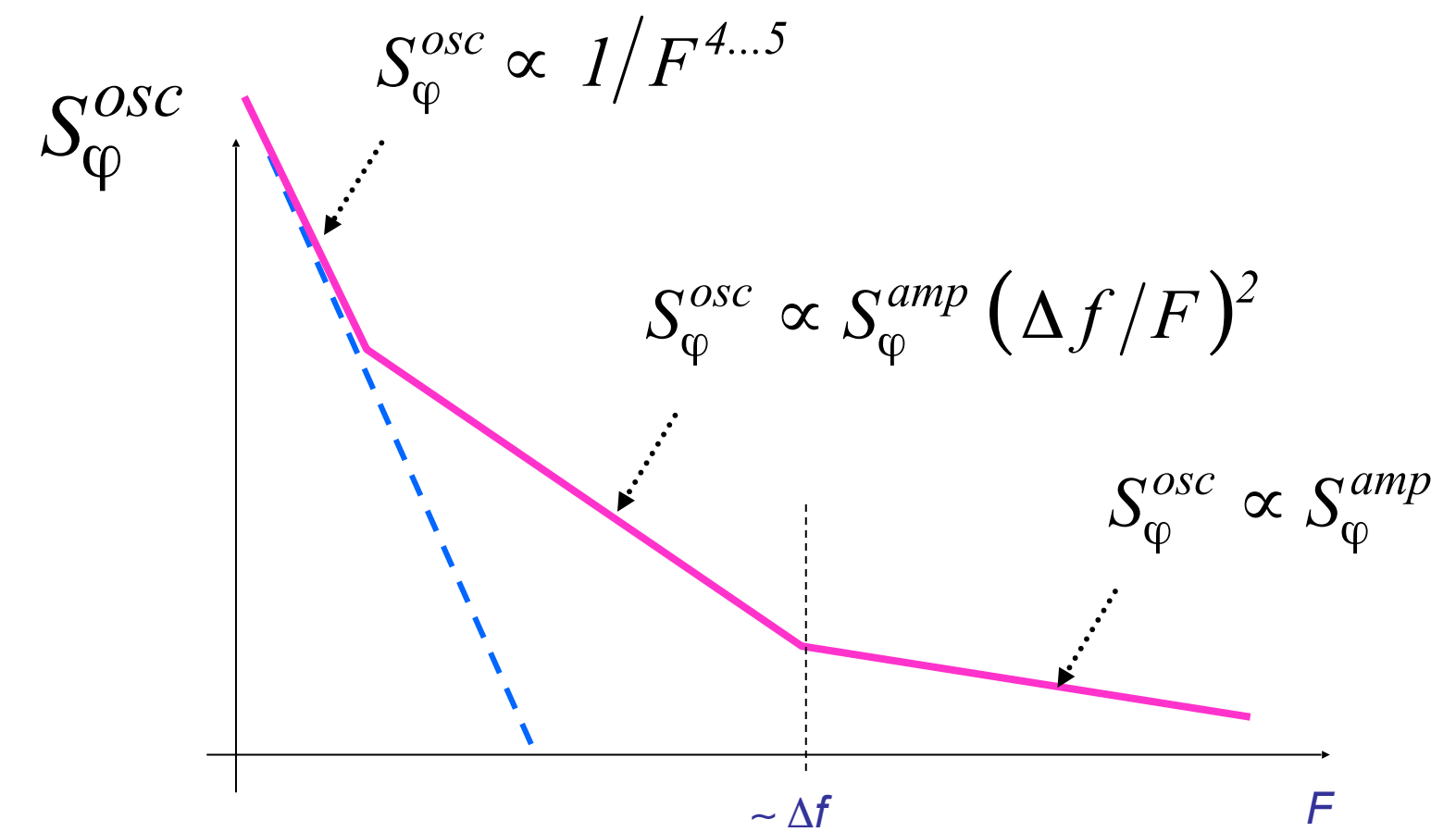
(f is Fourier frequency)

Lesson's Model



$$S_{\phi c} = S(f)_{\phi a} (1 + (\Delta f/2f)^2)$$

Oscillator phase noise spectrum



Low Phase Noise Oscillators

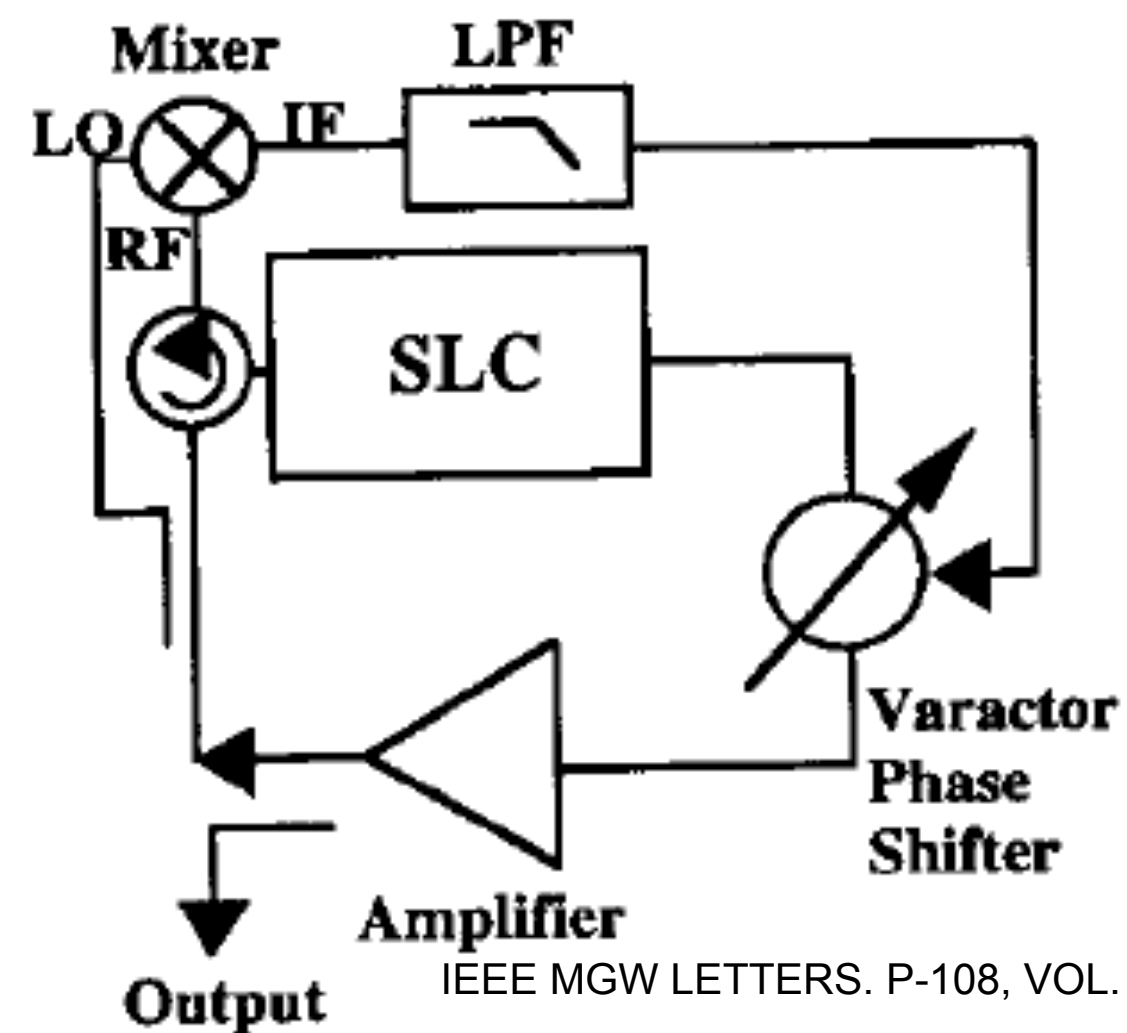
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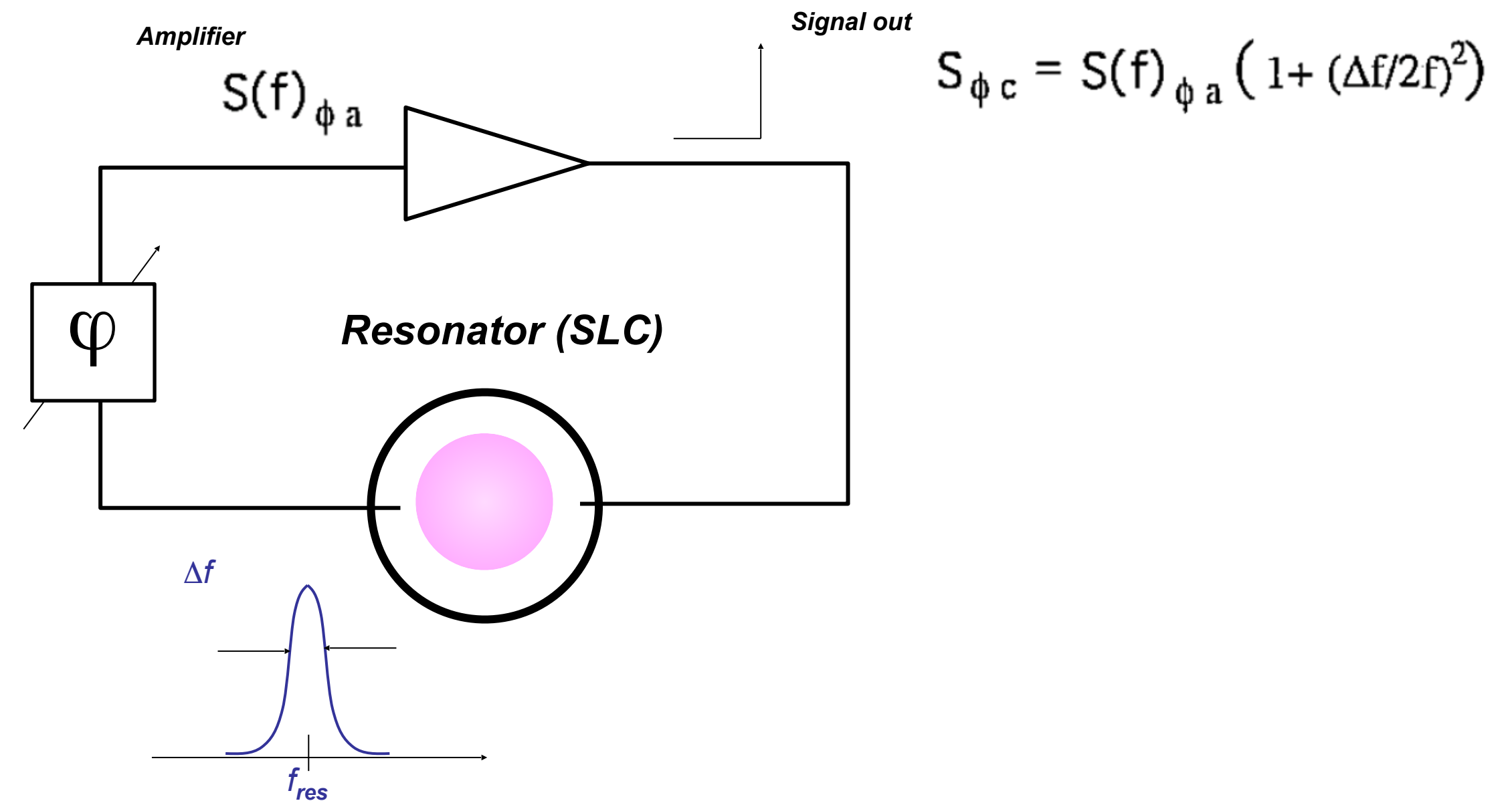
Cancel Noise with Phase Detector

Mixer phase detector

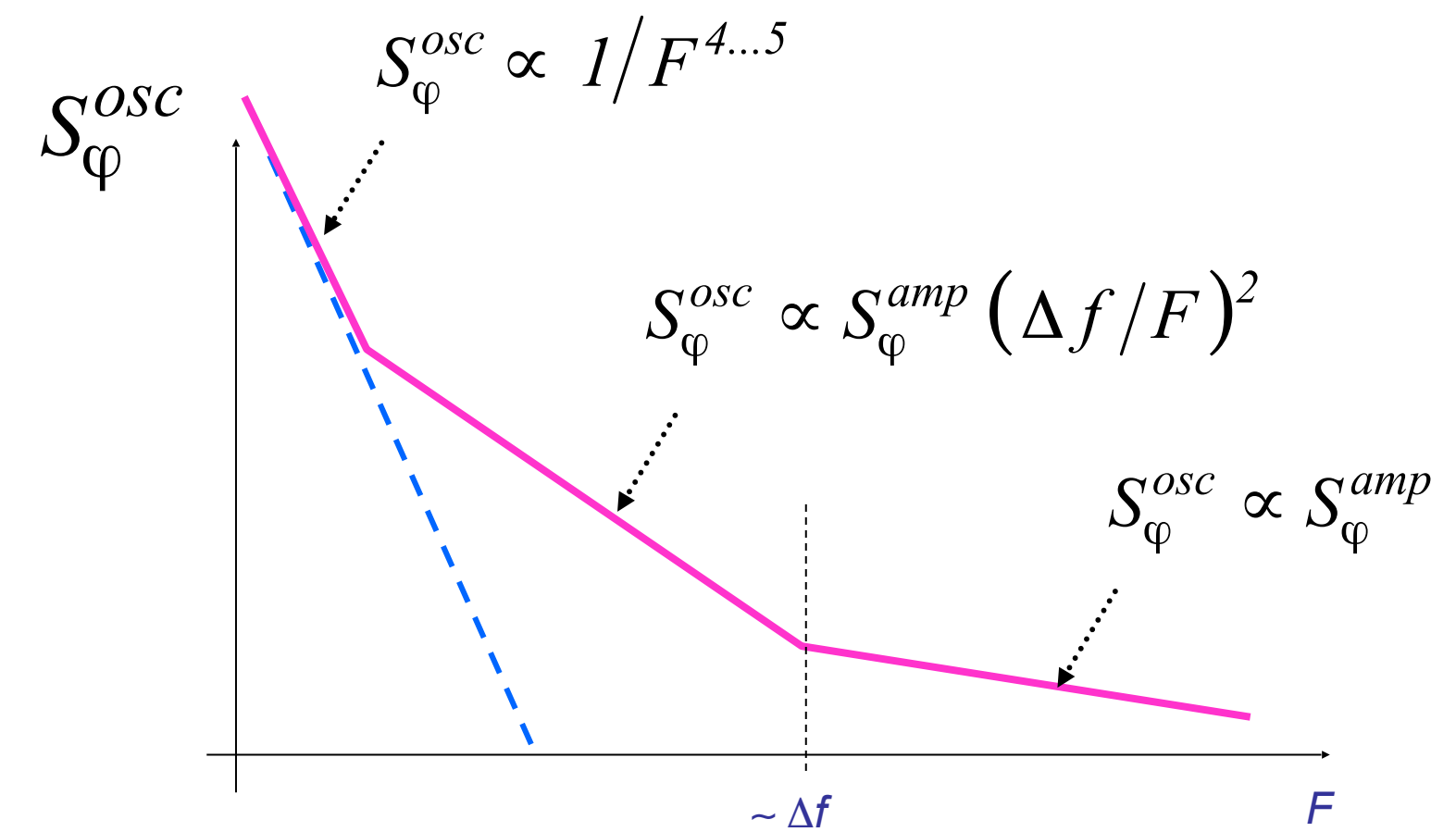


IEEE MGW LETTERS. P-108, VOL. 5. NO. 4, APRIL 1995

Lesson's Model



Oscillator phase noise spectrum



Low Phase Noise Oscillators

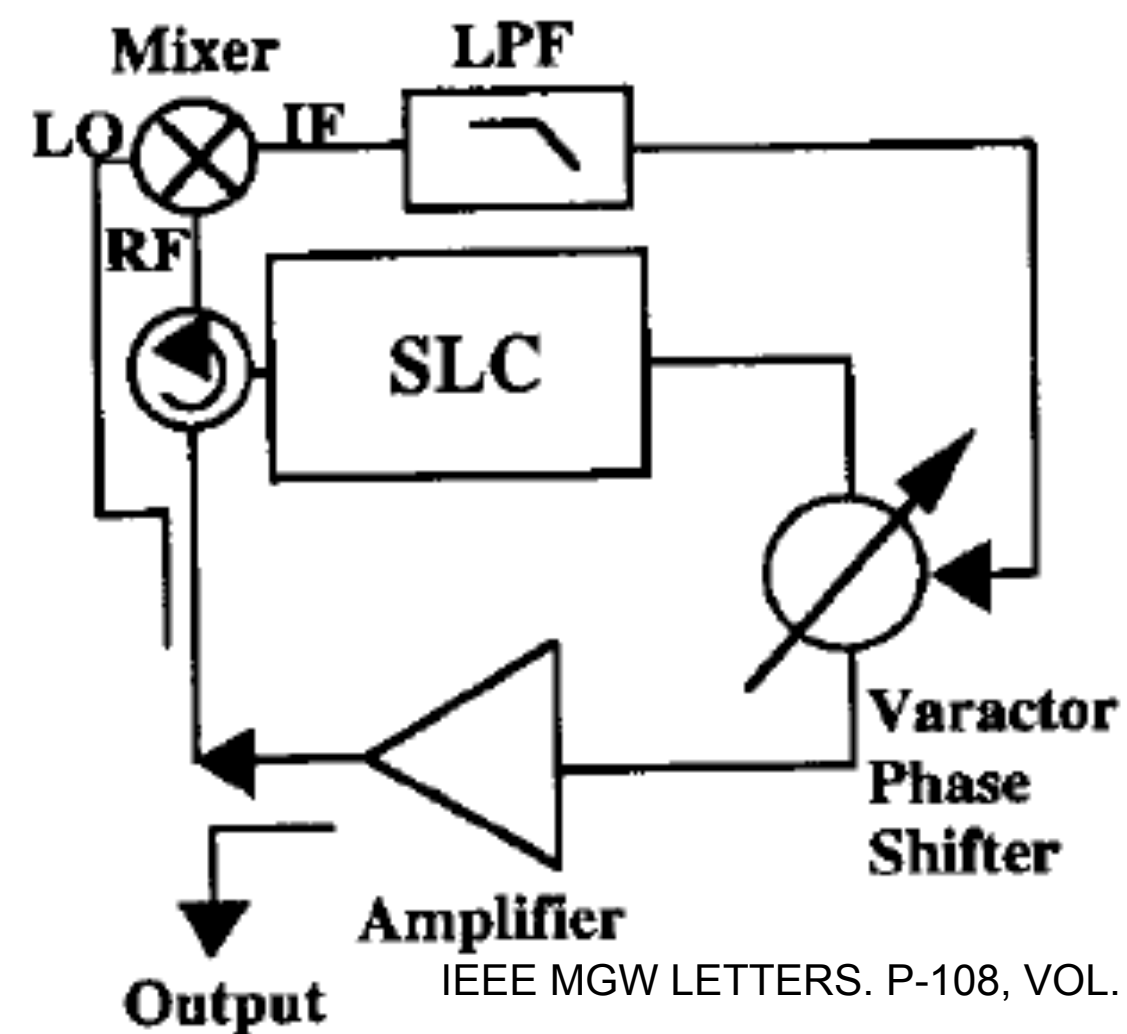
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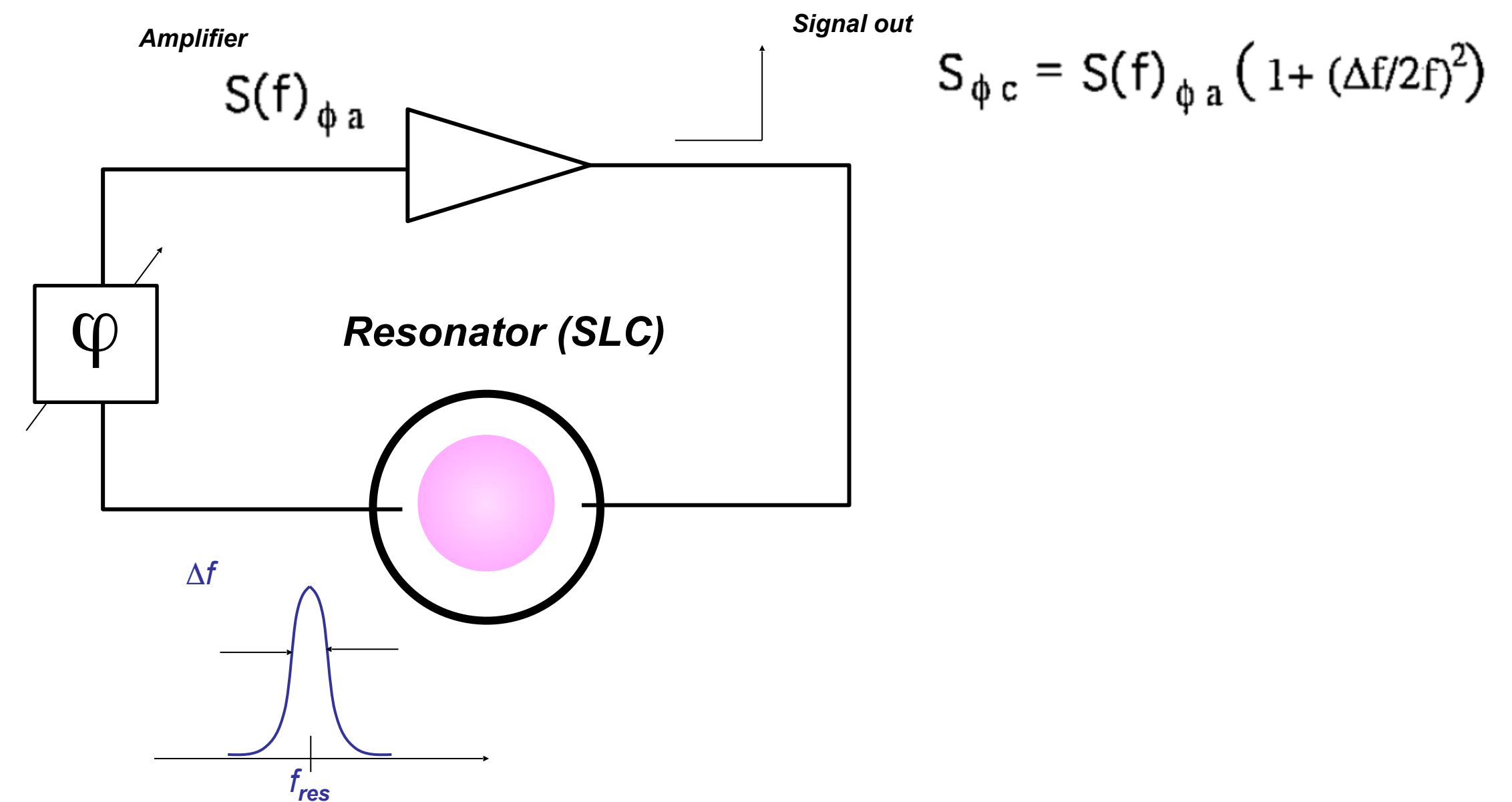
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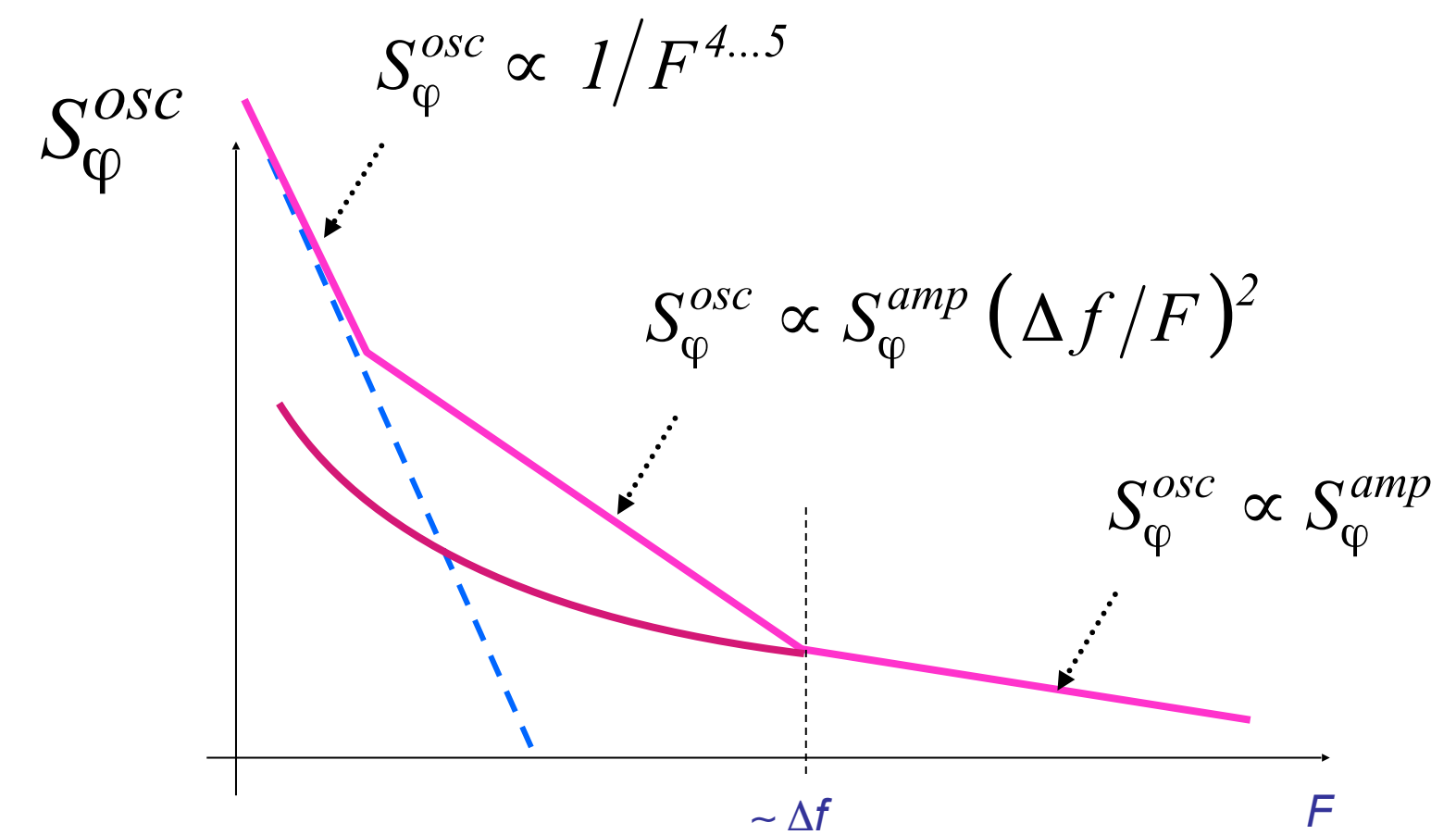


IEEE MGW LETTERS. P-108, VOL. 5. NO. 4, APRIL 1995

Lesson's Model



Oscillator phase noise spectrum



A MAJOR BREAK THROUGH IN MICROWAVE SENSING AND LOW-NOISE OSCILLATORS

CARRIER SUPPRESSION INTERFEROMETER

In 1993 we developed the interferometer as a phase detector, only limited by basic fundamental thermal noise if designed properly

The novel technique is covered by several international patent applications and granted patents -> UWA licensed technology to PSI Pty. Ltd. delayed publication.

1526

IEEE TRANSACTIONS ON ULTRASONICS, FERROELECTRICS, AND FREQUENCY CONTROL, VOL. 45, NO. 6, NOVEMBER 1998

Microwave Interferometry: Application to Precision Measurements and Noise Reduction Techniques

Eugene N. Ivanov, M. E. Tobar, *Member, IEEE*, and R. A. Woode

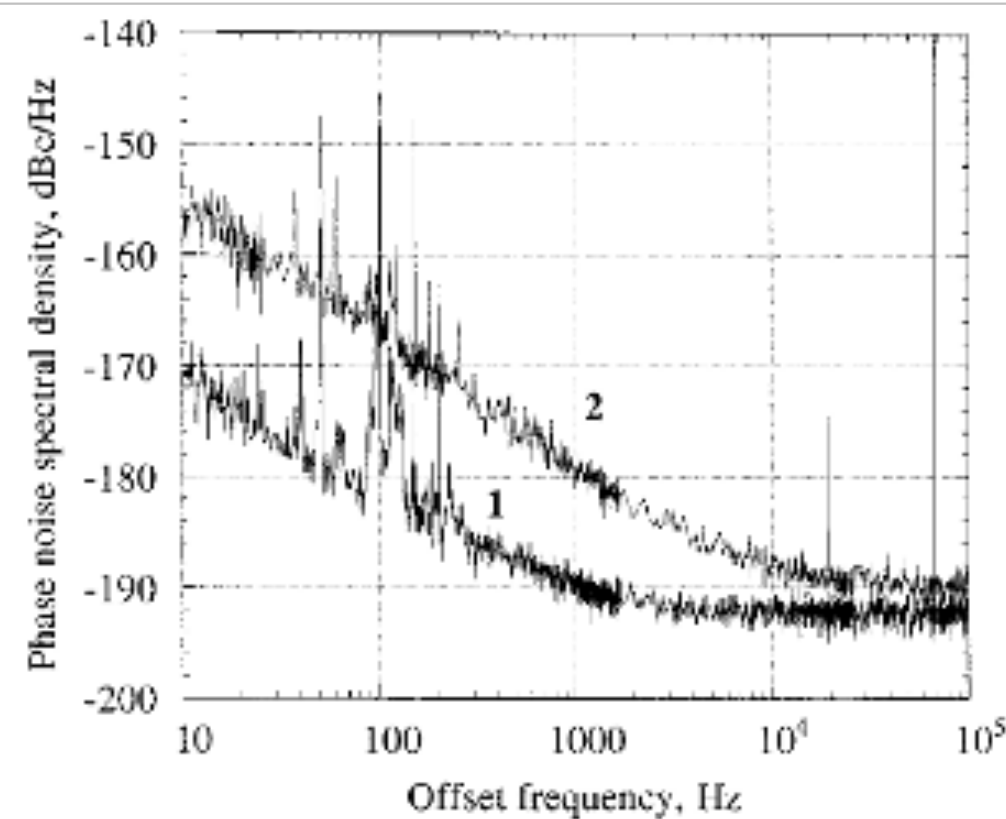


Fig. 2. The phase noise floor of interferometric noise measurement system (curve 1), phase noise of 6 microwave isolators connected in series (curve 2). Input power is 20 dBm, carrier frequency 9 GHz.

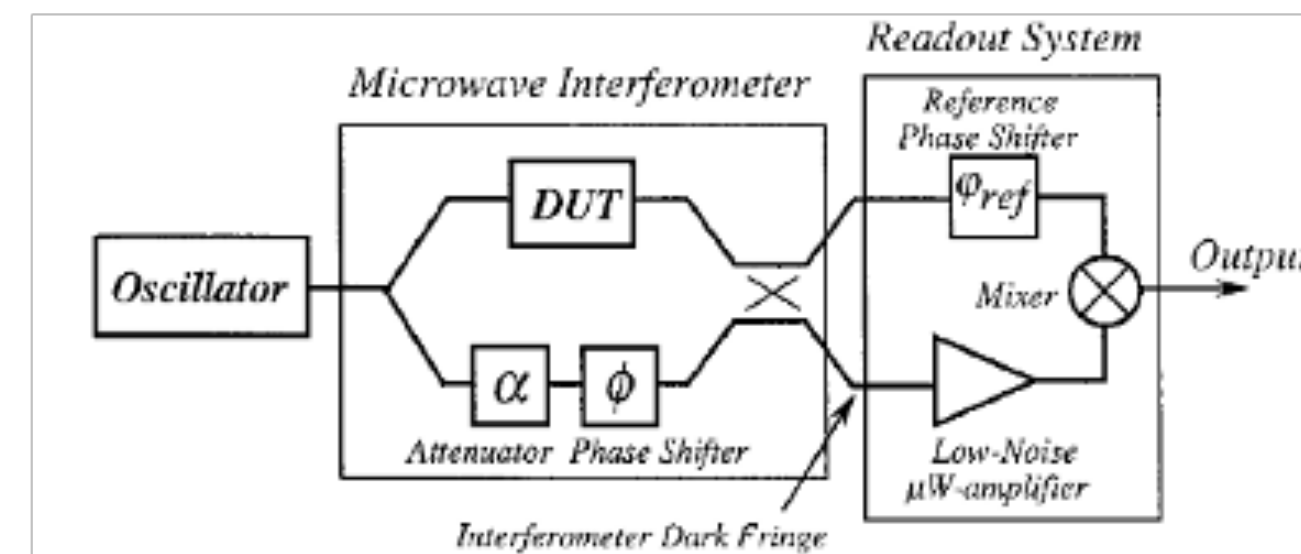


Fig. 1. Interferometric noise measurement system.

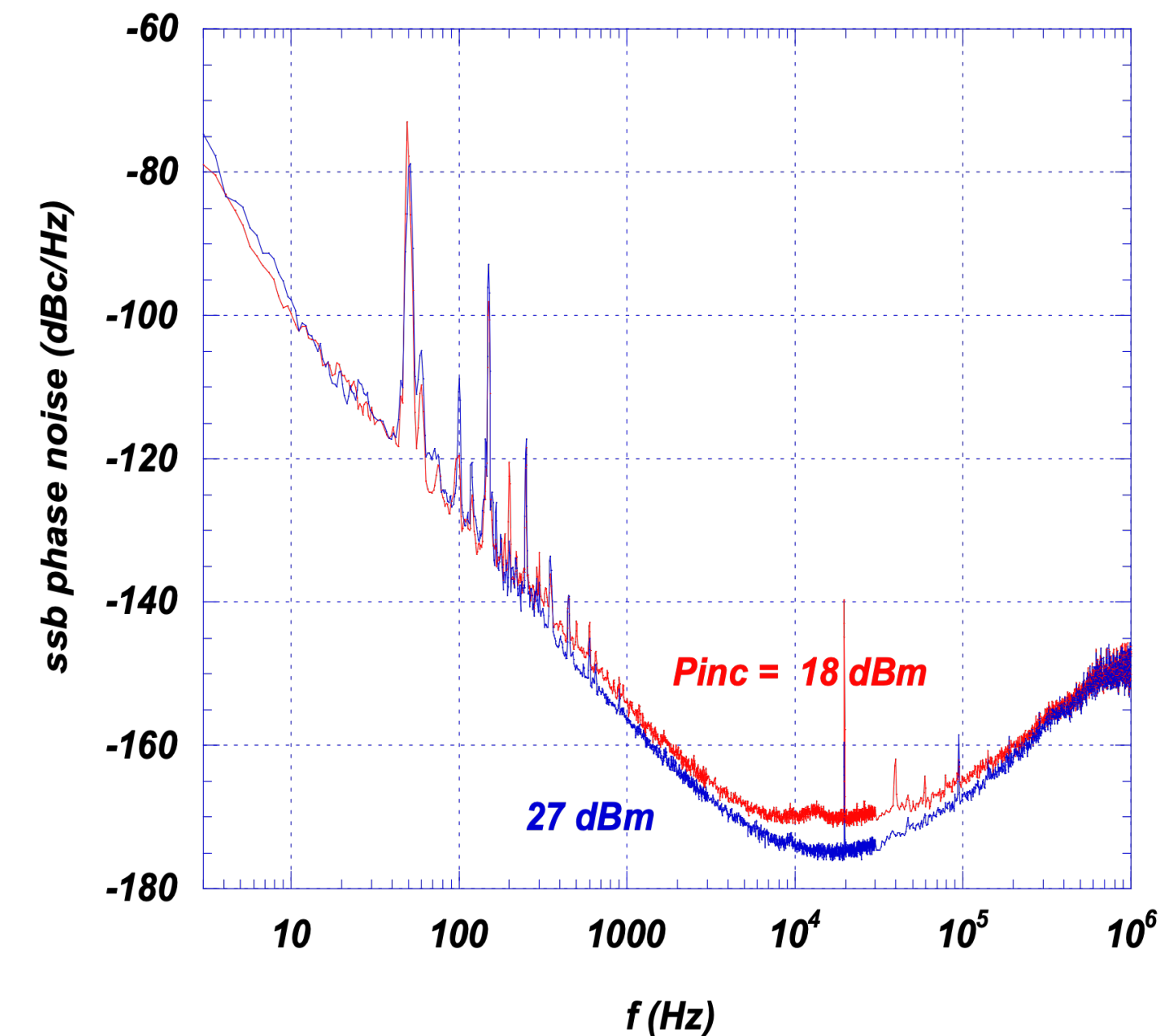
$$\mathcal{L}_{\varphi}^{n/f(1)}(f) = \mathcal{L}_{AM}^{n/f(1)}(f) = \frac{k_B T_{RS}}{P_{in} L_{DUT}}, \quad (1)$$

IEEE TRANSACTIONS ON ULTRASONICS, FERROELECTRICS, AND FREQUENCY CONTROL, VOL. 56, NO. 2, FEBRUARY 2009

263

Low Phase-Noise Sapphire Crystal Microwave Oscillators: Current Status

Eugene N. Ivanov and Michael E. Tobar, *Senior Member, IEEE*

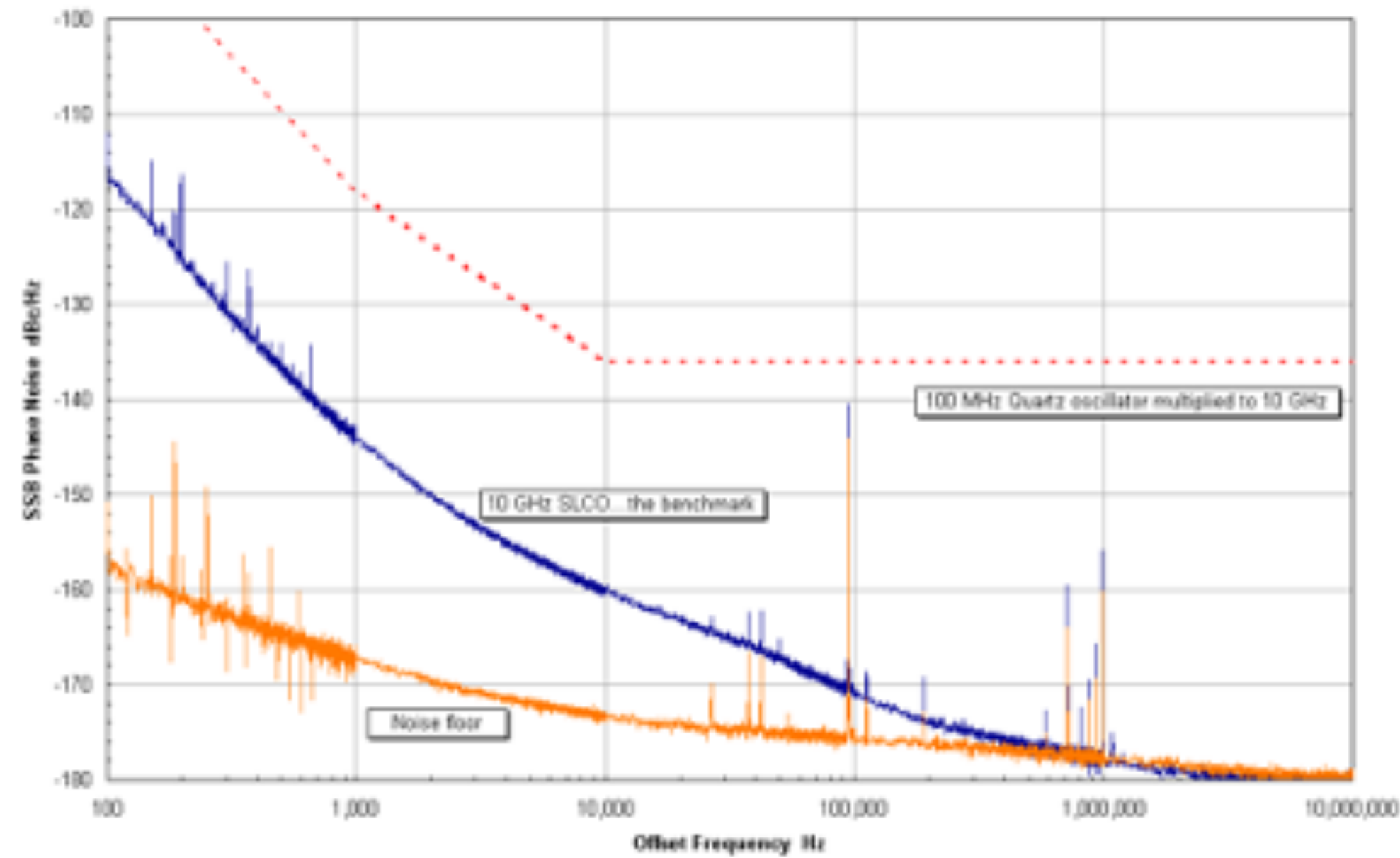


ROOM TEMPERATURE LOW NOISE SAPPHIRE OSCILLATORS: CAN BE MADE COMPACT OR RACK MOUNT



PoseidonScientificInstruments

- Lowest phase and amplitude noise
- Exceptional spectral purity
- Low spurious content
- Low vibrational sensitivity
- Unequaled short-term stability.



SAPPHIRE LOADED CAVITY OSCILLATORS

Compact low noise oscillator: Defence Radar Applications

Microwave Stripline Components: Compact Design

History of Commercialization

- 10 UWA Patents to do with low noise oscillators

Raytheon Australia: PSI Program

http://www.raytheon.com.au/businesses/integrated_solutions/cap...

Raytheon Australia

Raytheon Australia Acquires Poseidon Scientific Instruments

Raytheon Australia has acquired Poseidon Scientific Instruments in order to enhance Raytheon's suite of world leading technology and defence capabilities. Poseidon Scientific Instruments brings to Raytheon a depth of technical expertise in ultra low phase noise signal generation and measurement. Terms of the transaction are not being disclosed.

Products include:

- The SBO Class of Ultra Low Phase Noise Oscillators
 - SBO-XPL Compact Sapphire Oscillator
 - SBO-HS Compact and High-Speed Sapphire Oscillator
 - SBO Accessories
- The SKO Class of Sapphire Loaded Oscillator
- The SLCO Class of Sapphire Loaded Cavity Oscillator
 - Sapphire Loaded Cavity Oscillator (SLCO-BCS)
 - Sapphire Loaded Cavity Oscillator (SLCO NCS)
- Low Noise Dielectric Resonator Oscillators (DRO)
- Low Noise Divider Ensemble
- Low Noise Regenerative Divider
- DENA-5A Rack Mount Divider
- Oscillator Accessories
- Phase Noise Analysers and Receiver Modules
 - ODIN-320AS Phase Noise Analyser
 - OR-101A – 6GHz to 12GHz Receiver Module
 - OR-102A – 5MHz to 1GHz Receiver Module
 - OR-105A – 1GHz to 18GHz Receiver Module
 - OD-103B – Delay Line
 - OC-104A – Calibration Modules

For more information, please contact: PSIProgram@raytheon.com.au

Corporate Communications
Raytheon Australia
4 Brindabella Circuit
Brindabella Business Park
Canberra Airport ACT 2609
[Contact Us](#)

History of Commercialization • 10 UWA Patents to do with low noise oscillators

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http://www.raytheon.com.au/businesses/integrated_solutions/cap...

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- Resurgence -> Now low noise oscillators necessary to drive quantum qubits to work at Quantum Limit

History of Commercialization • 10 UWA Patents to do with low noise oscillators

Raytheon Australia: PSI Program

http://www.raytheon.com.au/businesses/integrated_solutions/cap...

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- Resurgence -> Now low noise oscillators necessary to drive quantum qubits to work at Quantum Limit



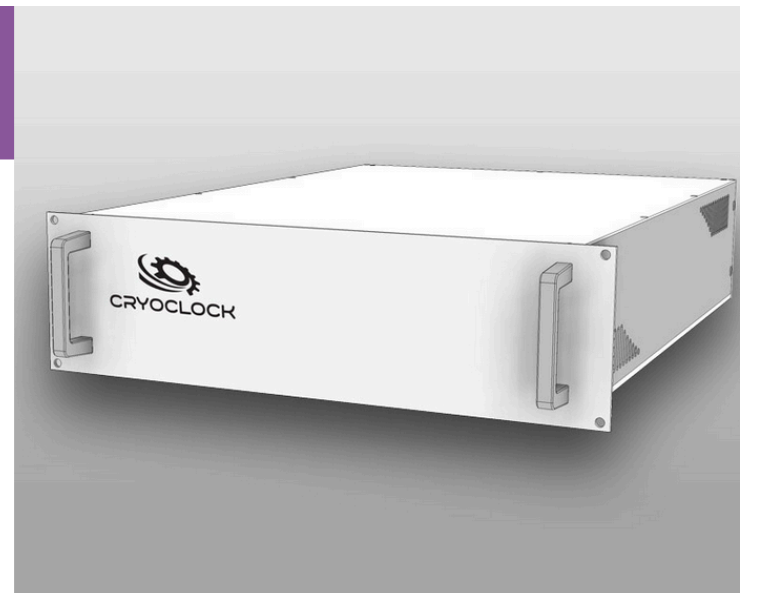
X-LNO

Ultra-Low-Noise Microwave Oscillator

The X-LNO is a microwave reference oscillator that produces world-leading ultra-low phase noise reference signal in the X-band region. By exploiting the remarkably high Q of sapphire, the oscillator delivers a +10 dBm signal with phase noise below -165 dBc/Hz at 10 kHz offset.

The X-LNO has a standard 3U package that is suited to rack-mounting although other OEM configurations are available on request. A key application for the X-LNO is the master oscillator in microwave communications and radar systems, including Precision Approach Radars and surface detection radars. The ultra-low phase noise of the X-LNO will enable significantly greater sensitivity in these radar systems when compared to quartz-based systems

This product can be configured to provide any frequency outputs between 8 to 12 GHz and wider ranges are possible on request.



KEY APPLICATIONS:

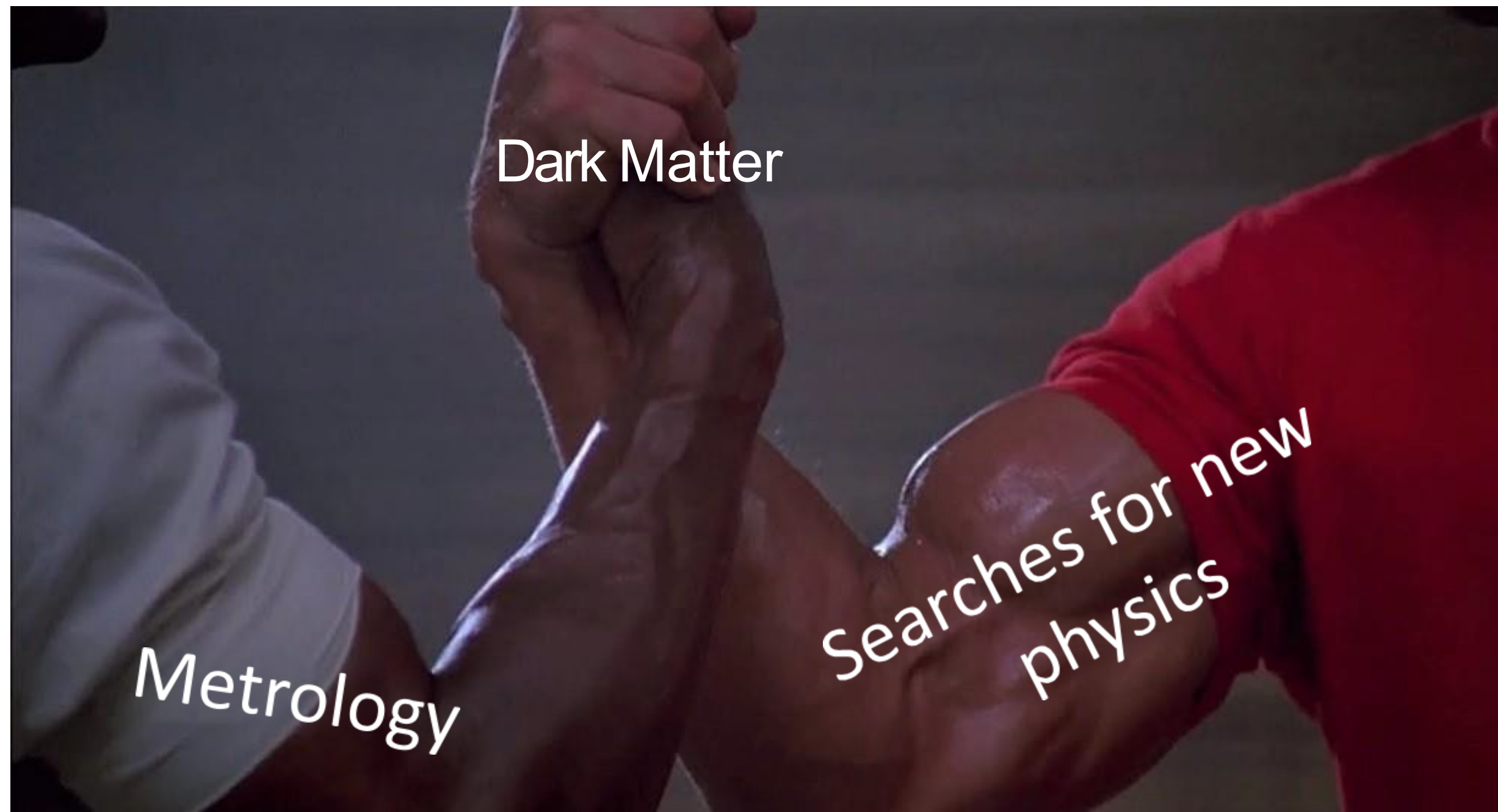
- Civilian and military radar
- Wind/gust monitoring near civilian airports
- Detection and tracking of fast-moving objects at a distance

Precision Metrology

Science of precise measurement



Physics at low energies

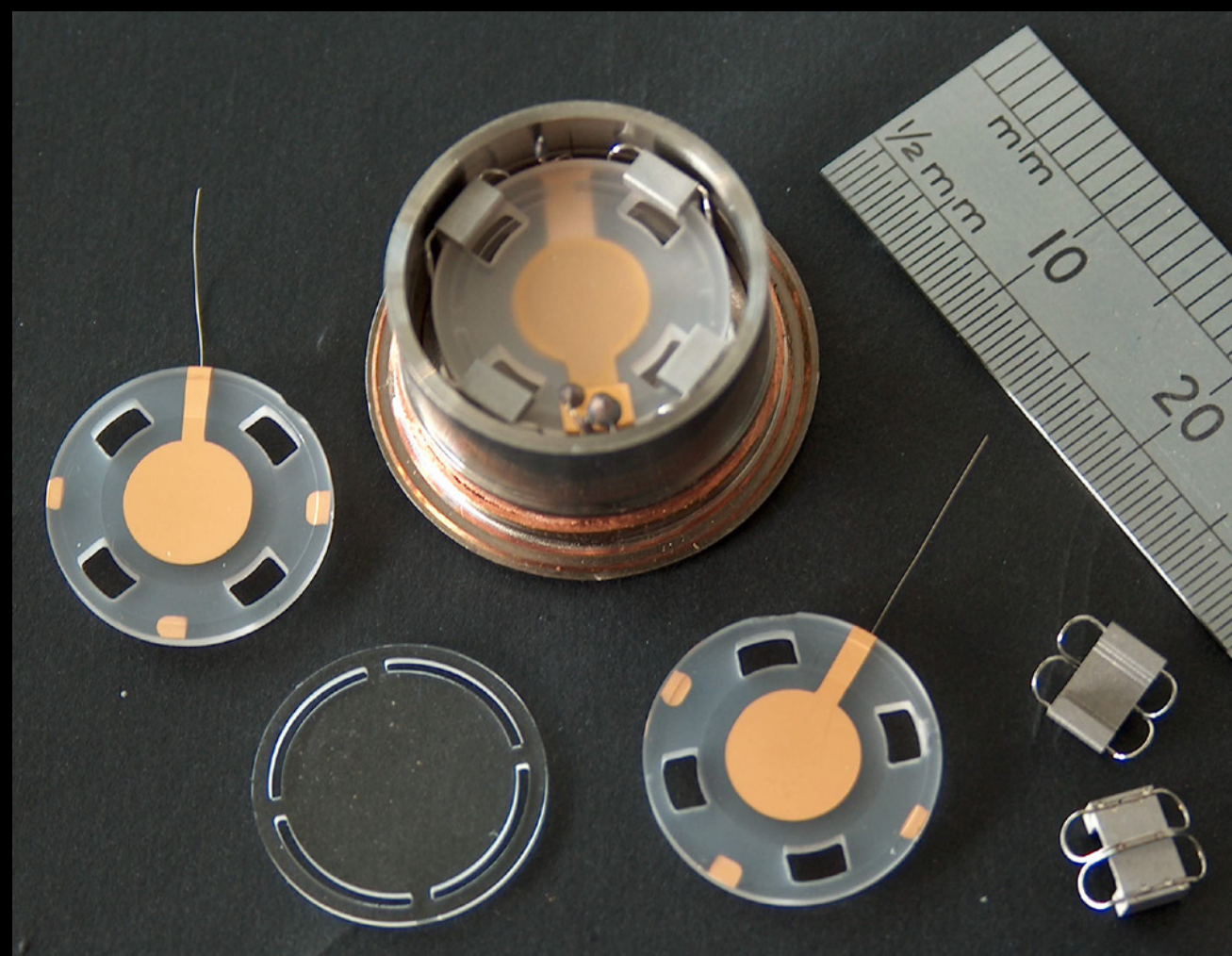


Metrology helps us search for physics beyond the standard model

Lorentz violation, fundamental constant variation, tests of general relativity & gravitation, violations of quantum statistics
+ more

Bulk Acoustic Wave High-Frequency GW Detectors (A Resonant-Mass Detector)

Prof. Michael Tobar

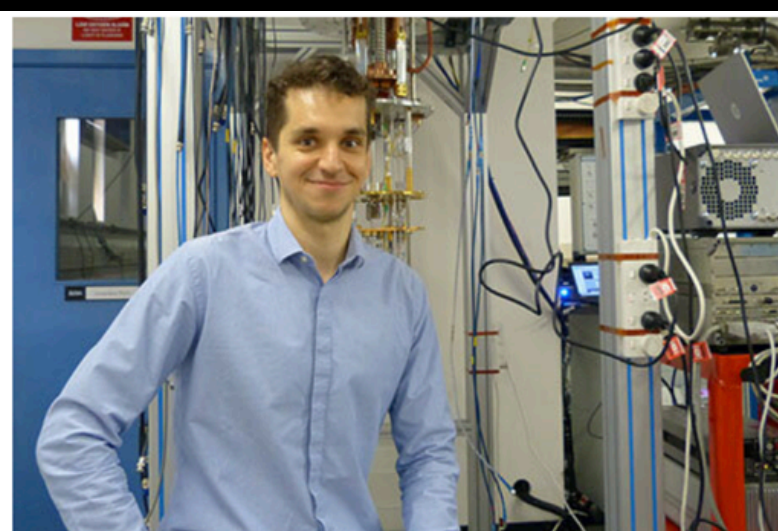


Will Campbell

PhD Student



Professor Mike Tobar



Dr Maxim Goryachev



Professor Eugene Ivanov

UWA Staff



Ik Siong Heng
Prof. Glasgow



Serge Galliou
Prof. Franche-Comté



Australian Government

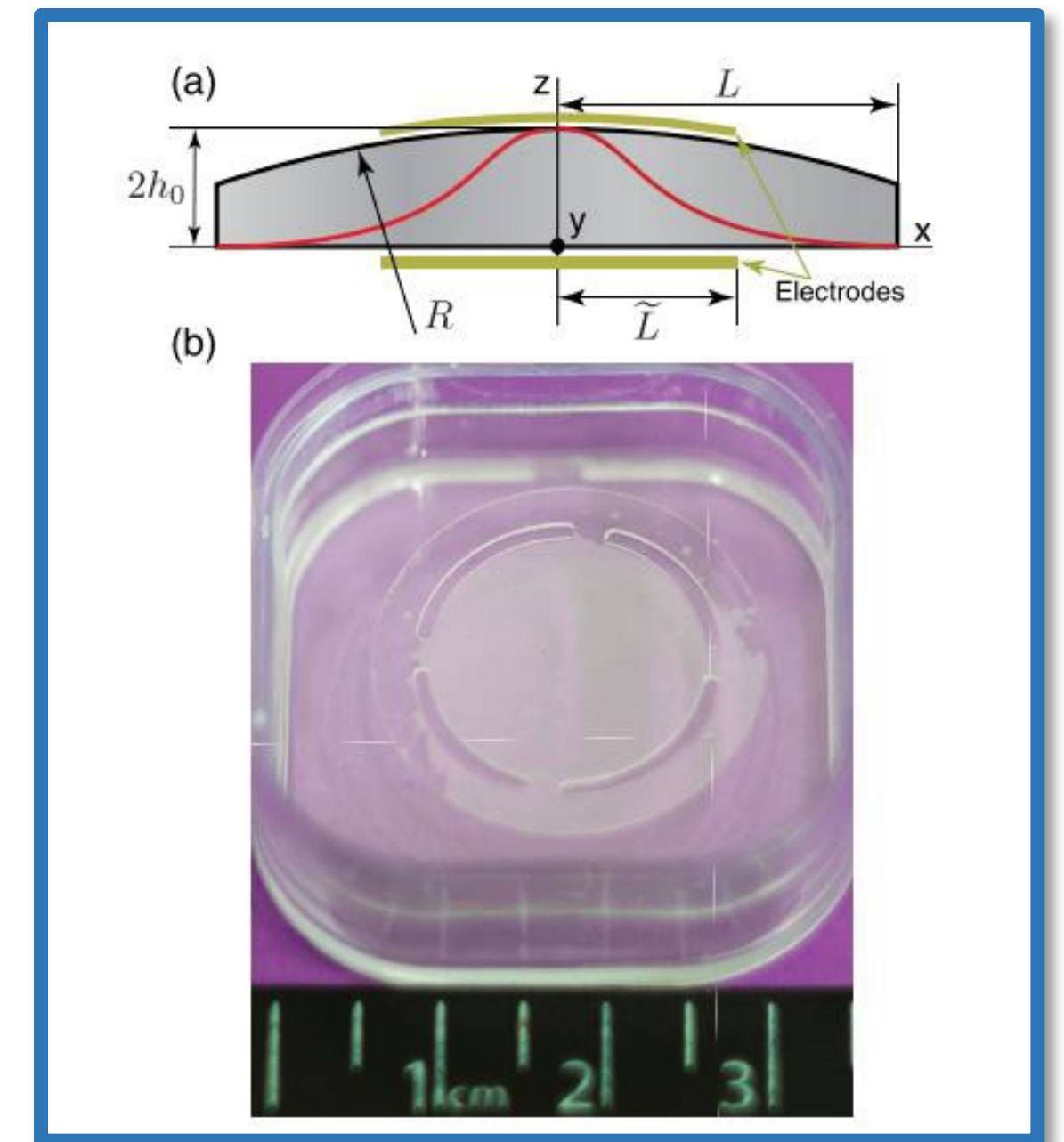
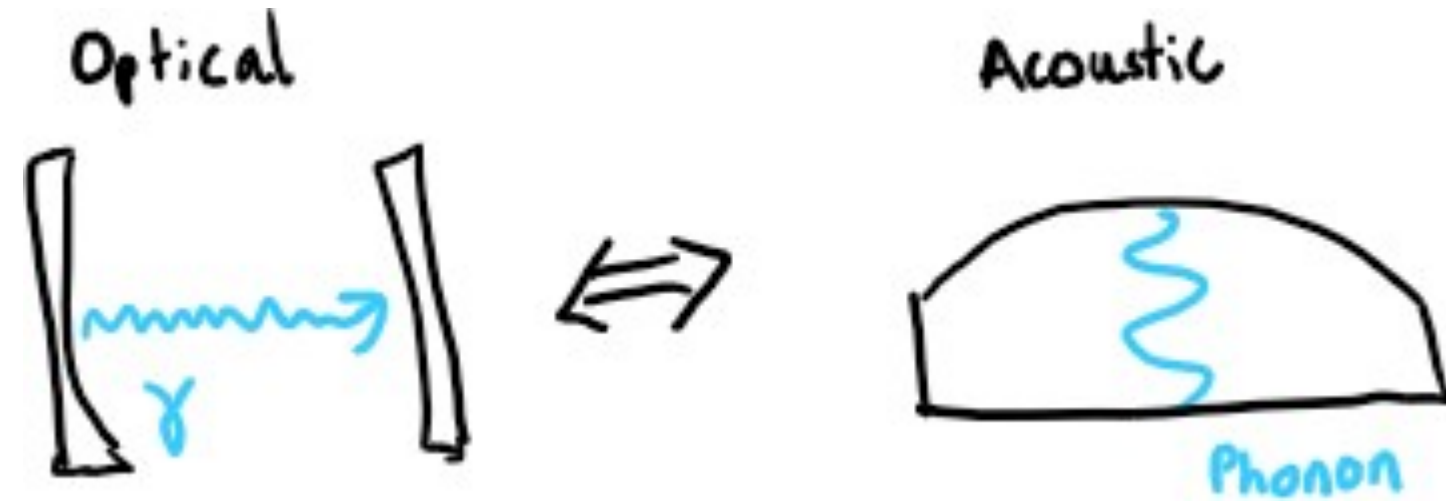
Australian Research Council



THE UNIVERSITY OF
**WESTERN
AUSTRALIA**

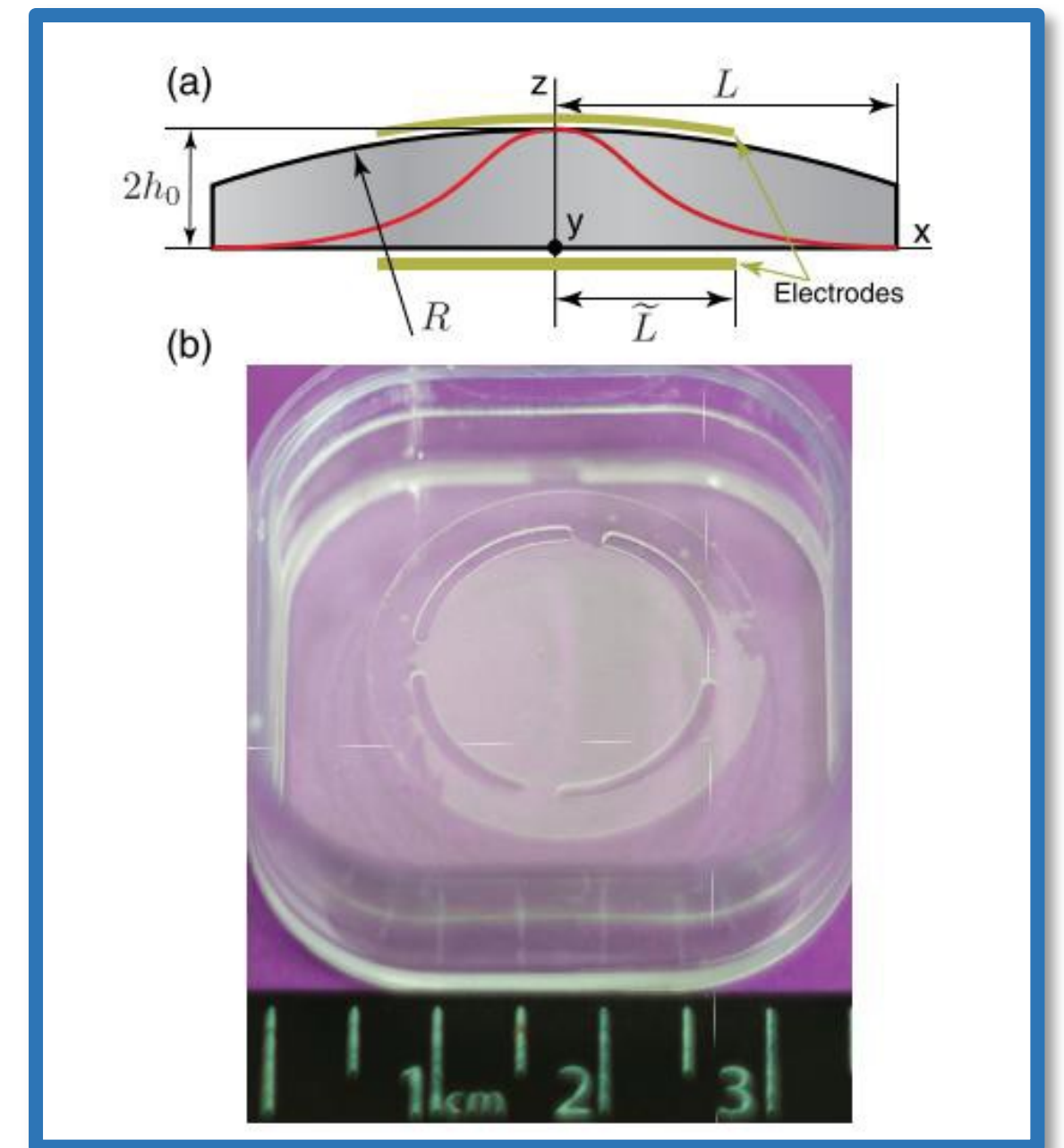
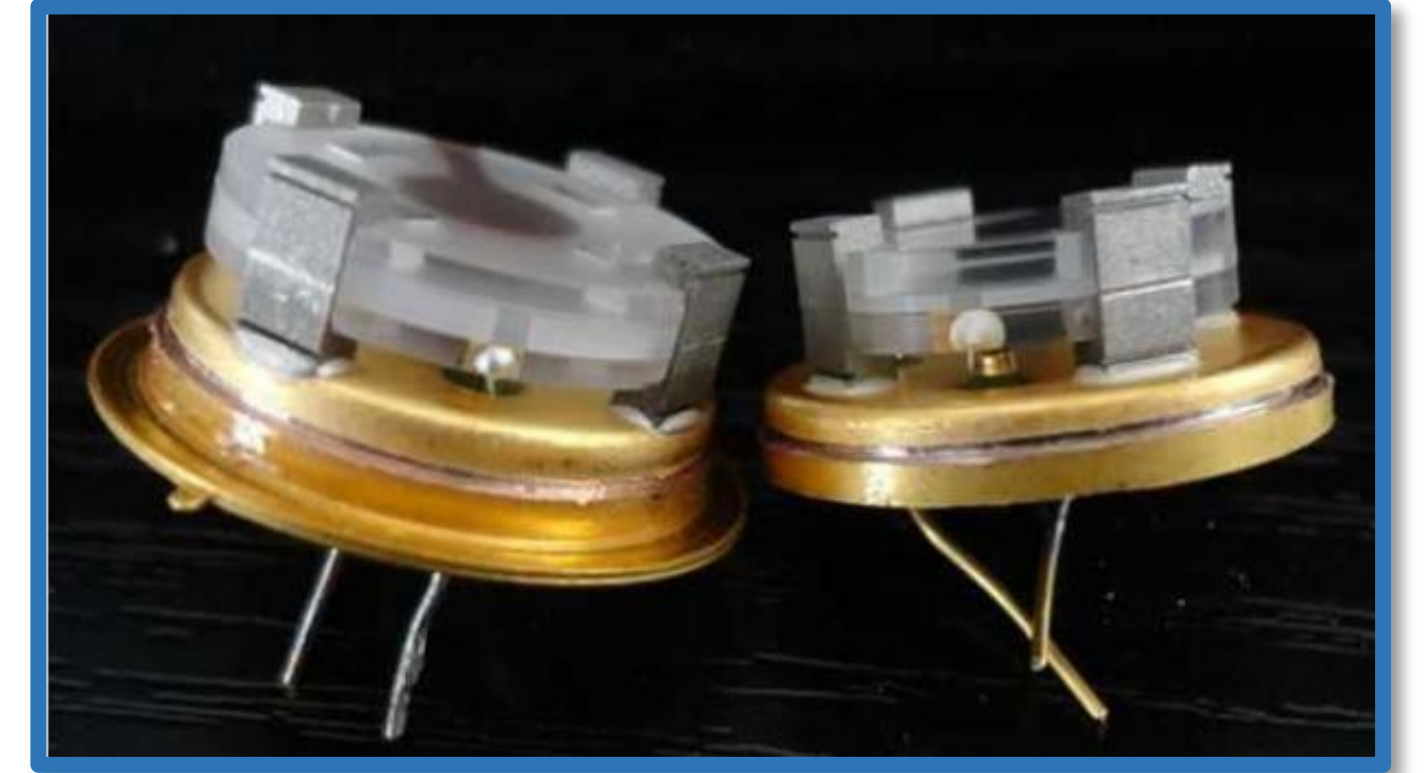
Quartz Bulk Acoustic Wave Resonators

- Acoustic analogue to a Optical Fabry-Perot cavity.



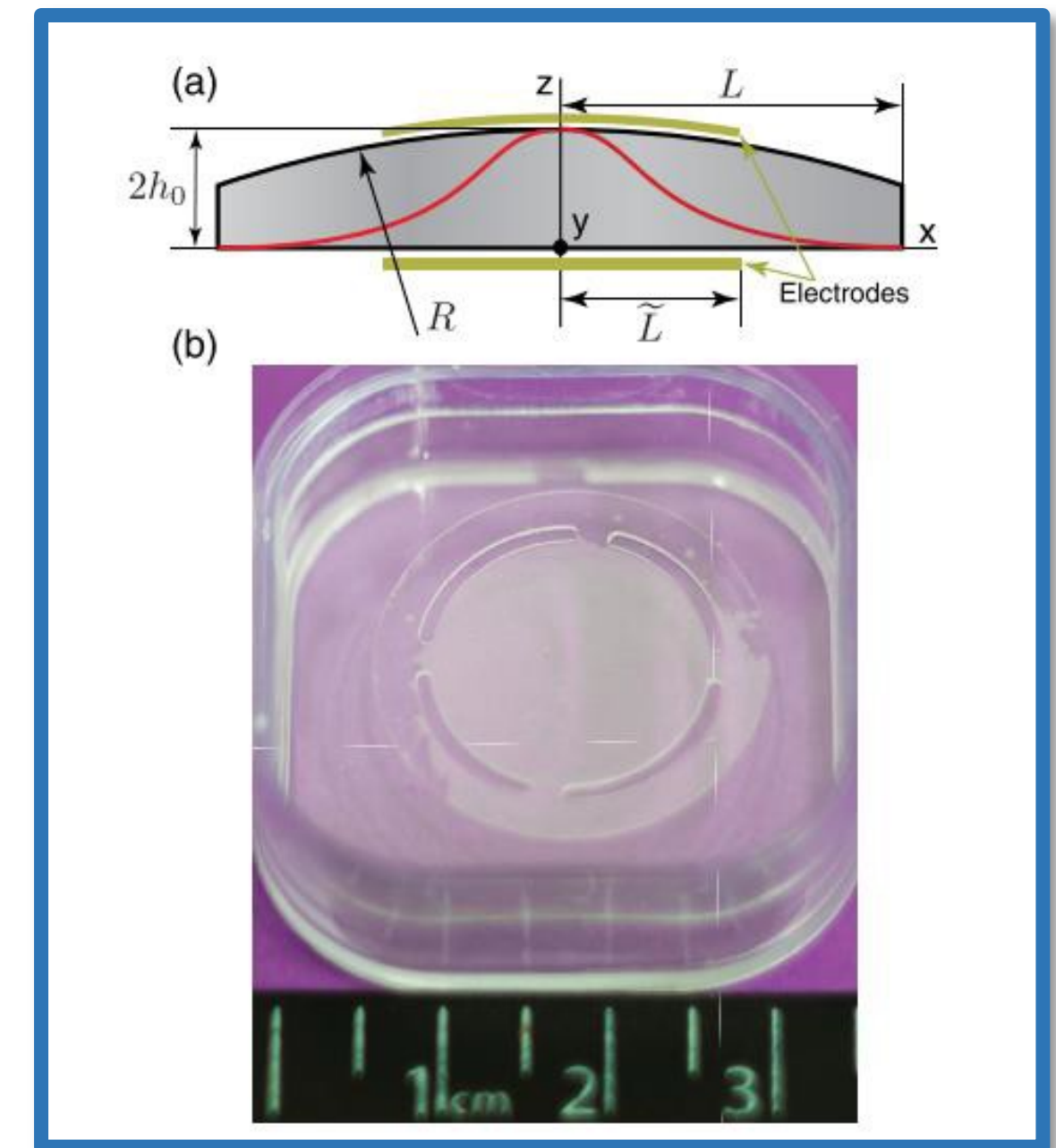
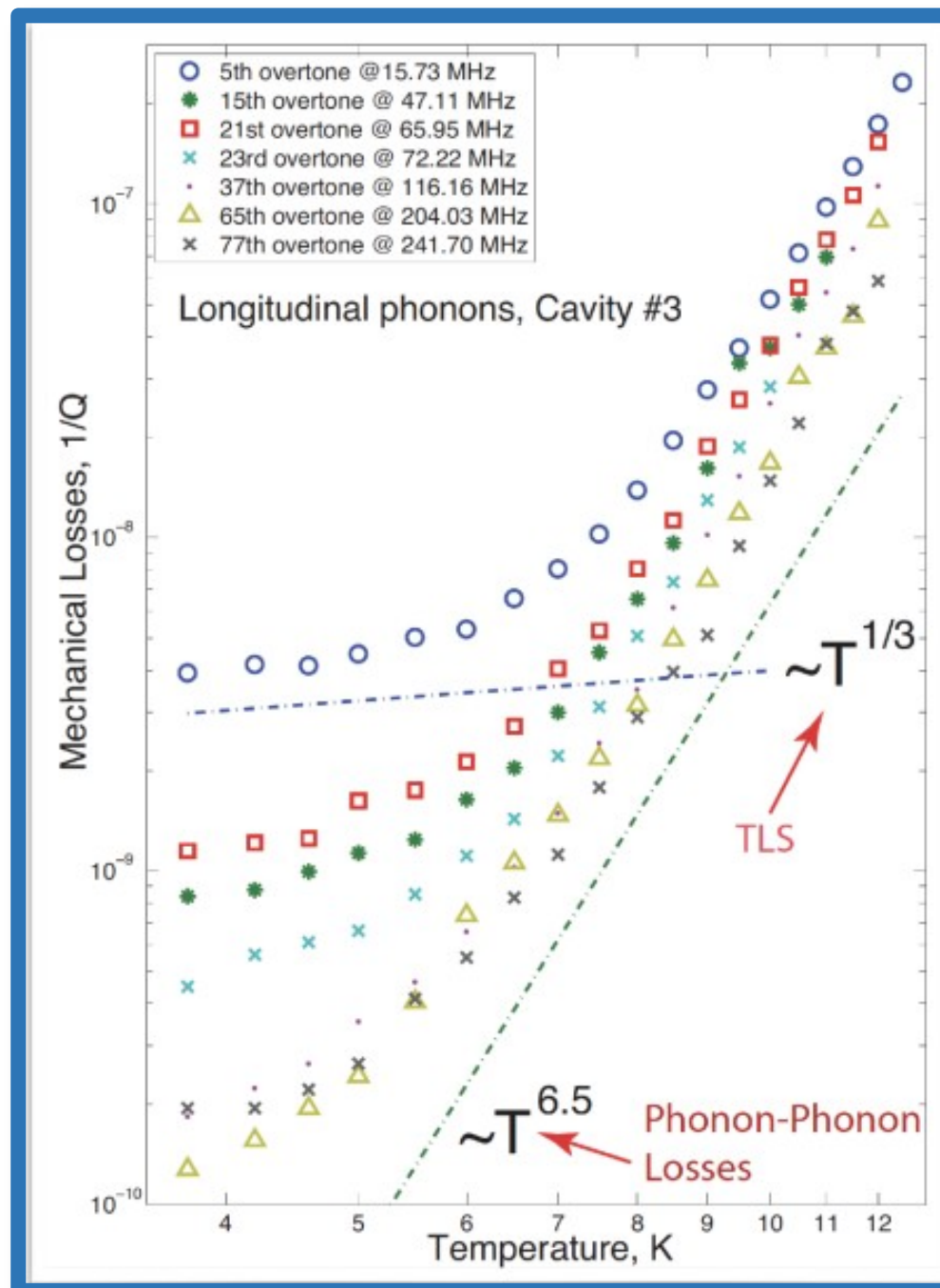
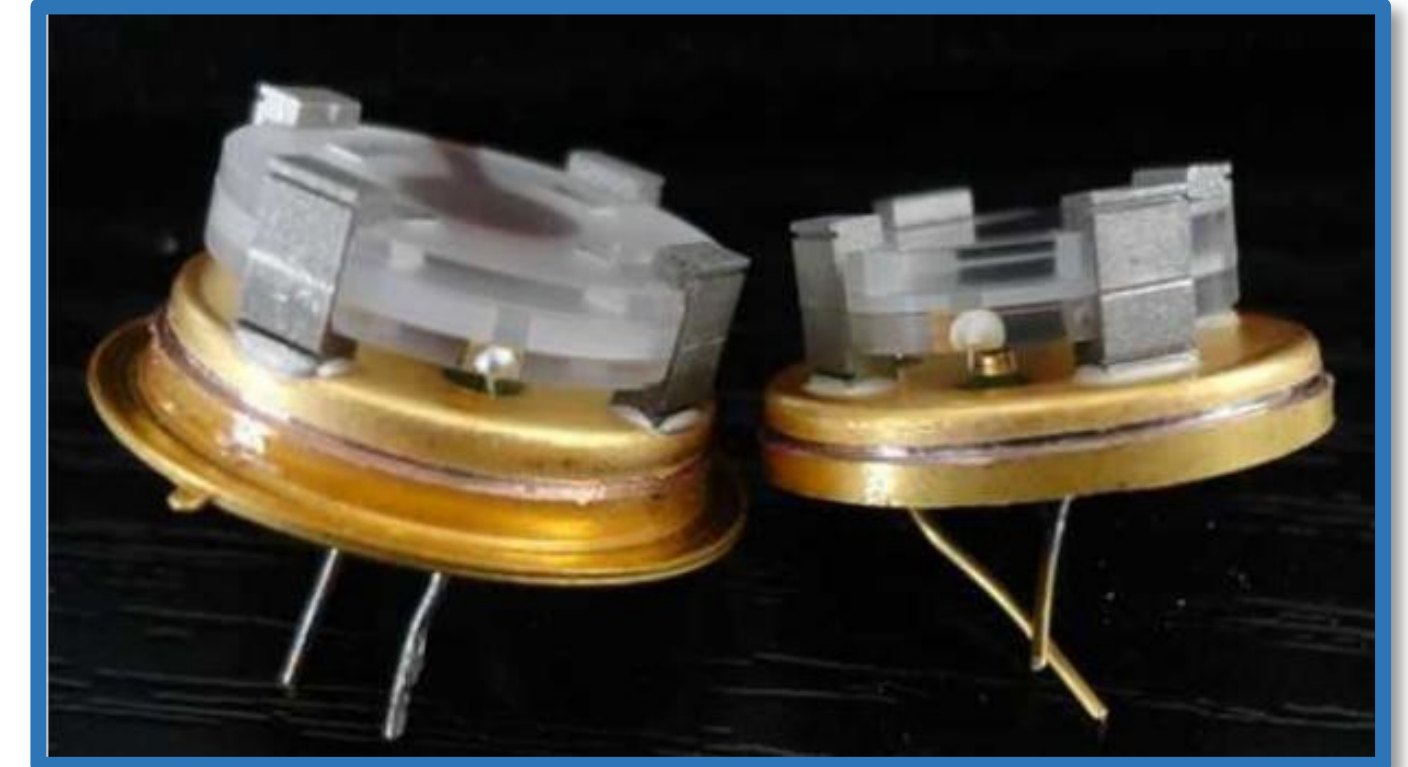
Quartz Bulk Acoustic Wave Resonators

- Acoustic analogue to a Optical Fabry-Perot cavity.
- Already a well established technology
- Gram scale mode mass, macroscopic resonator



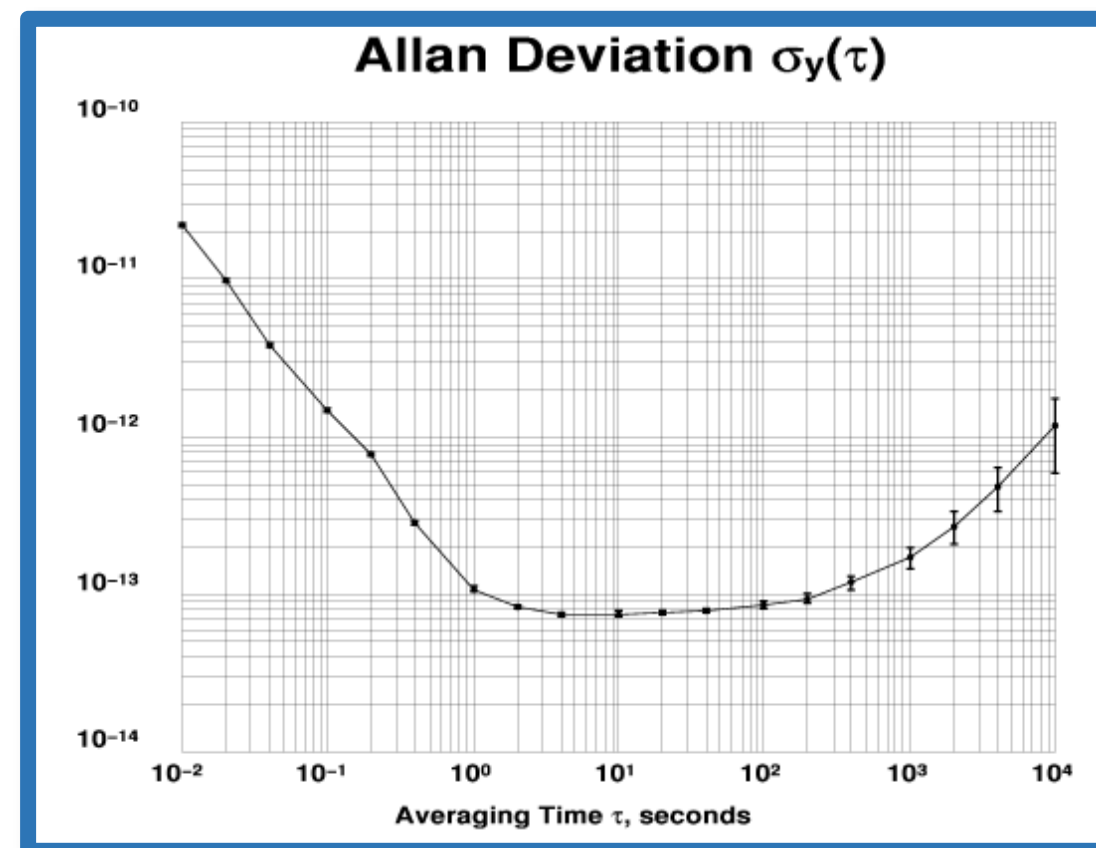
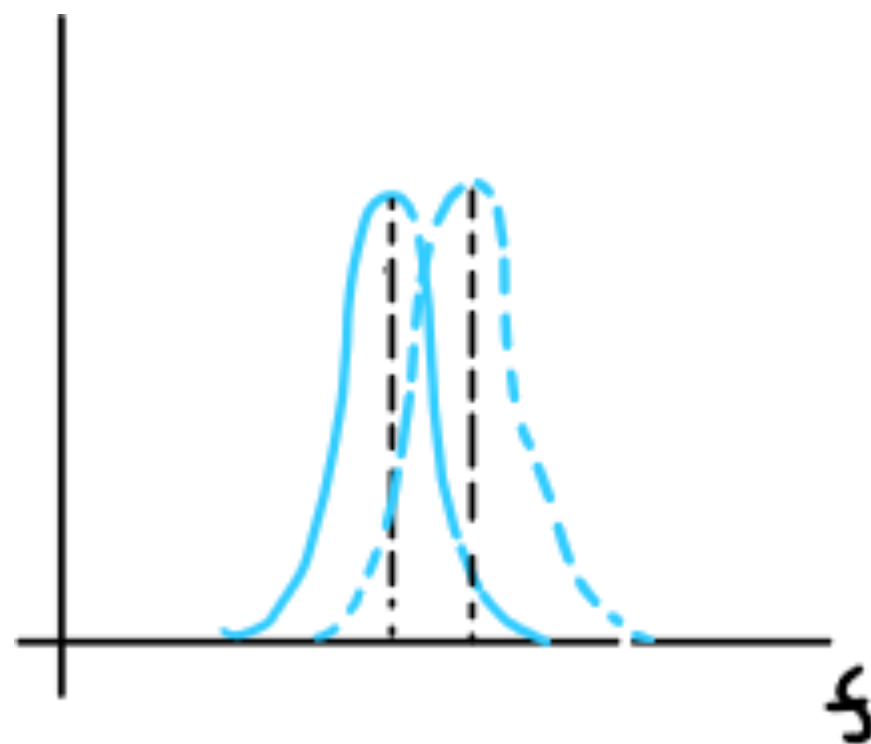
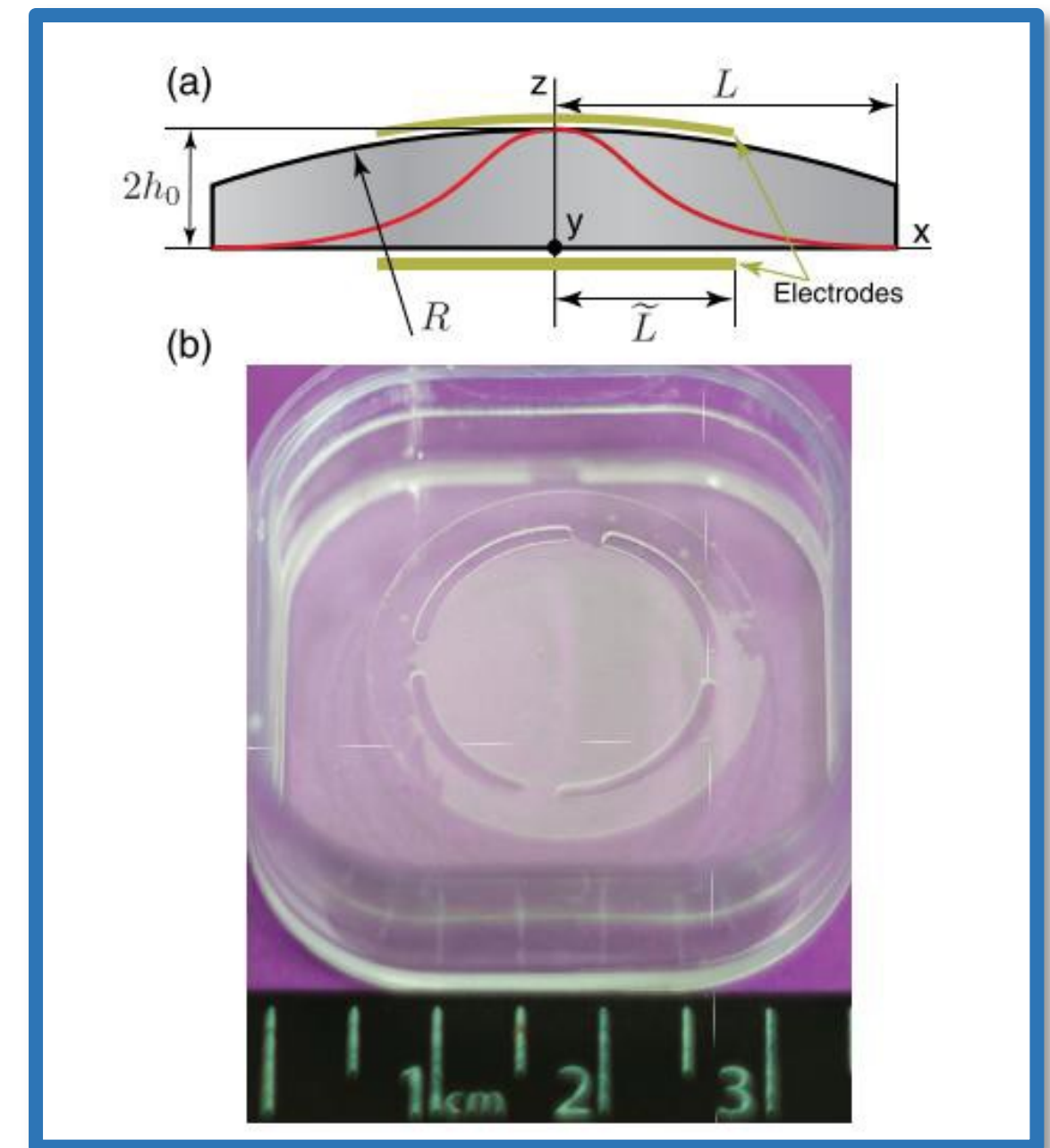
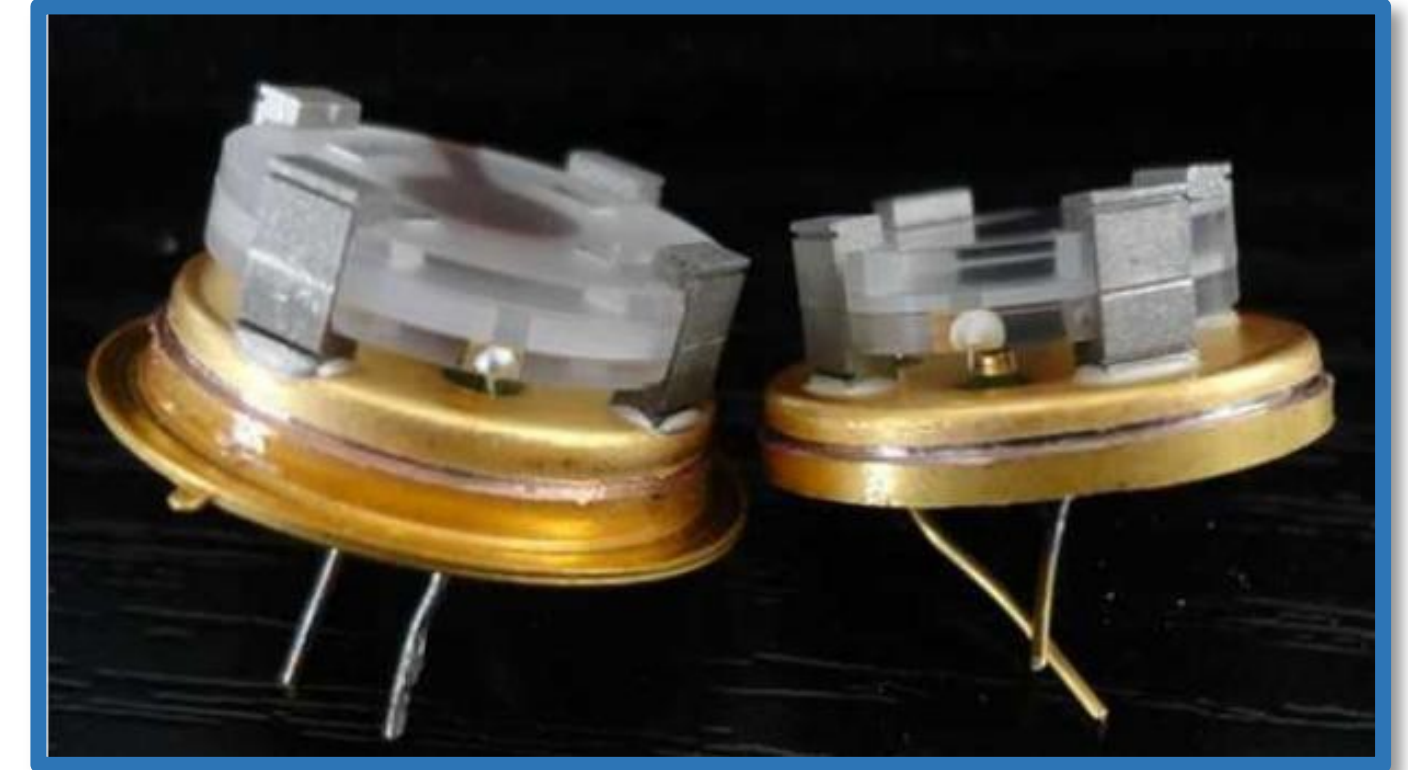
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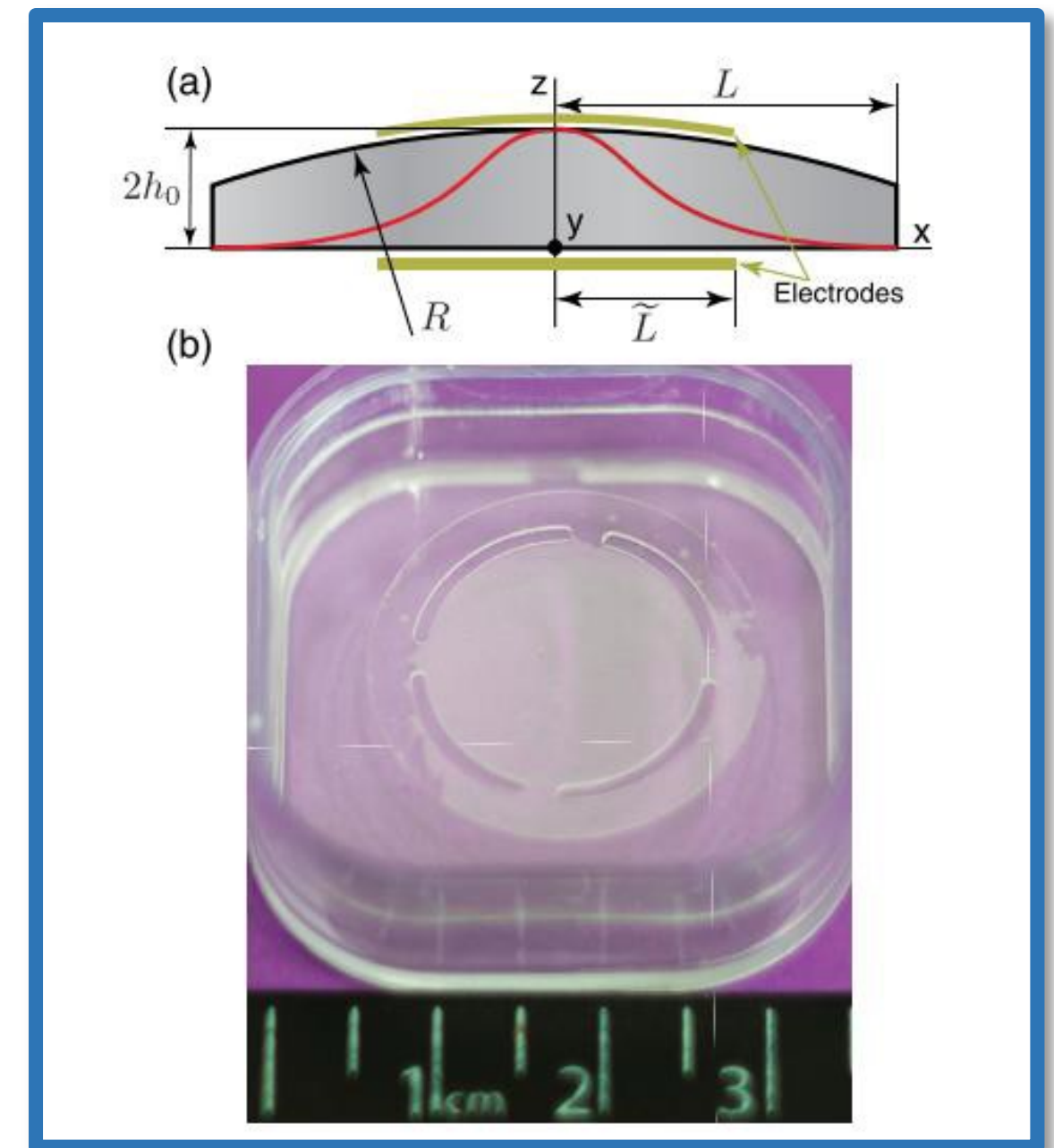
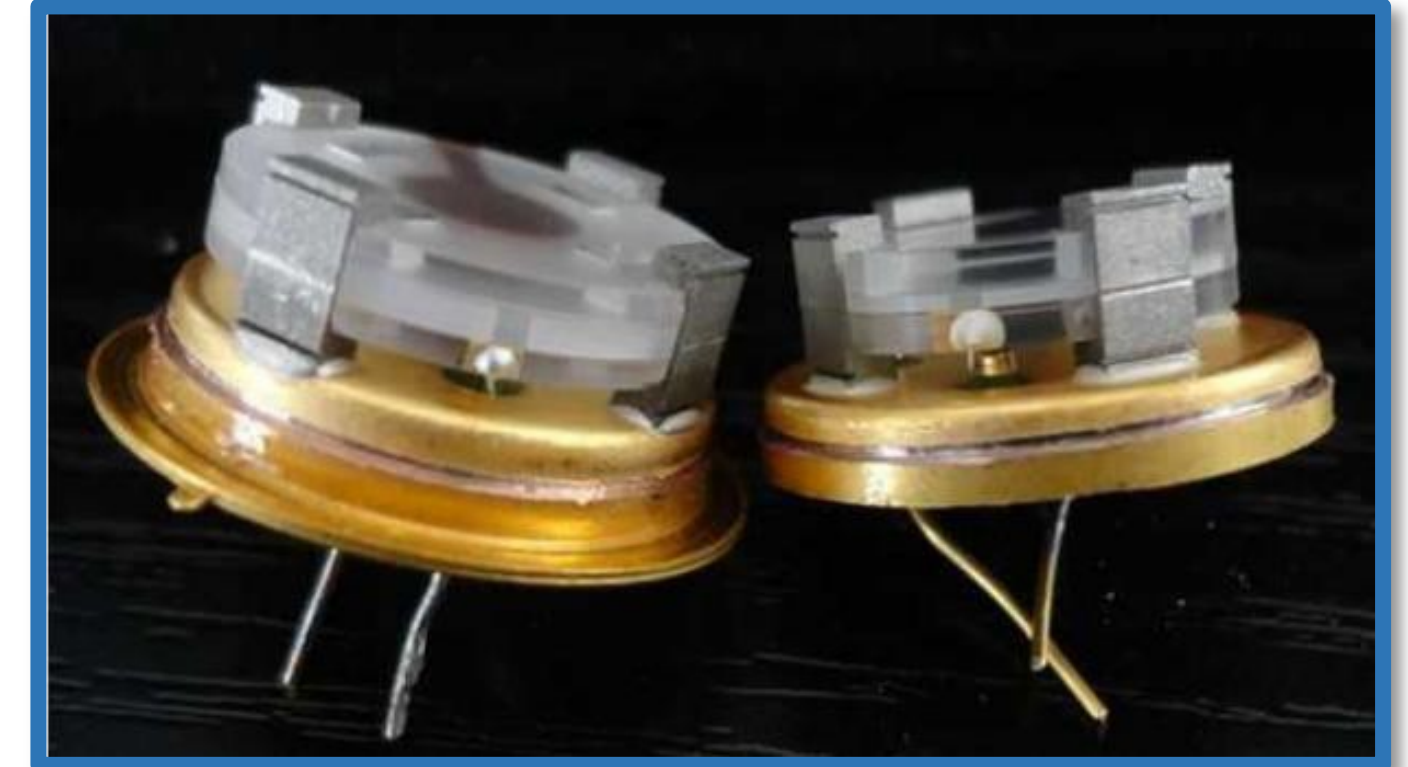
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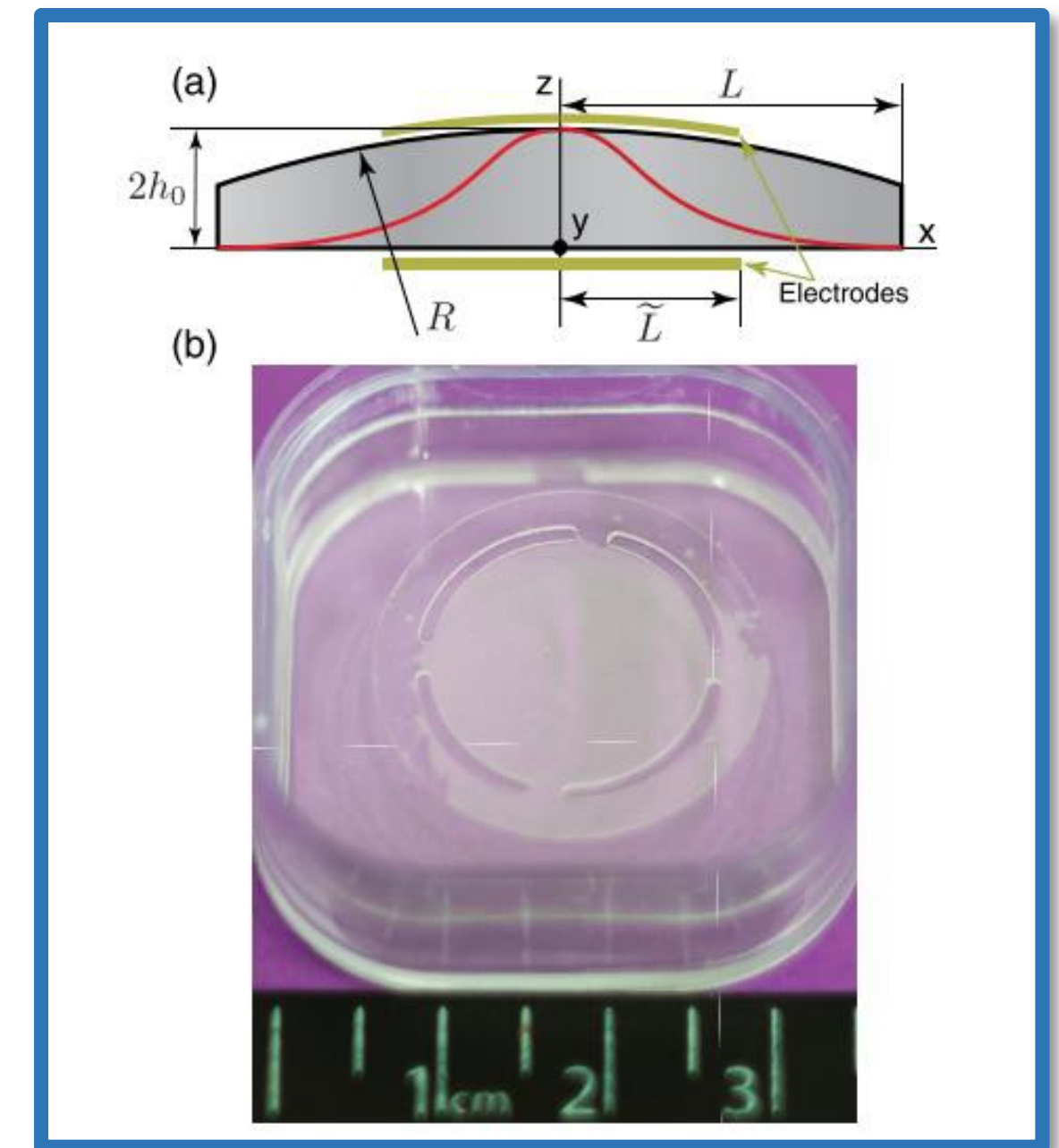
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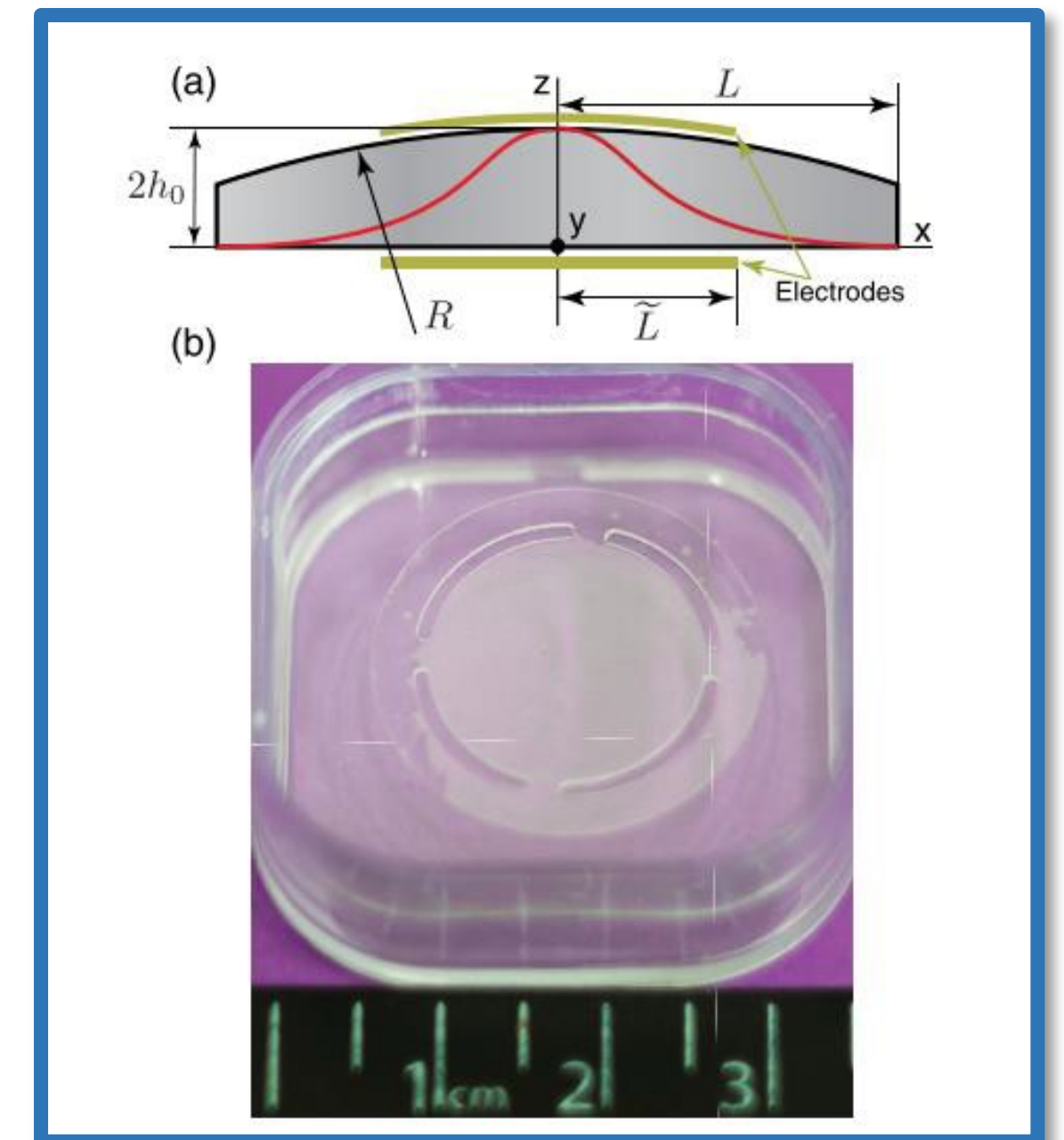
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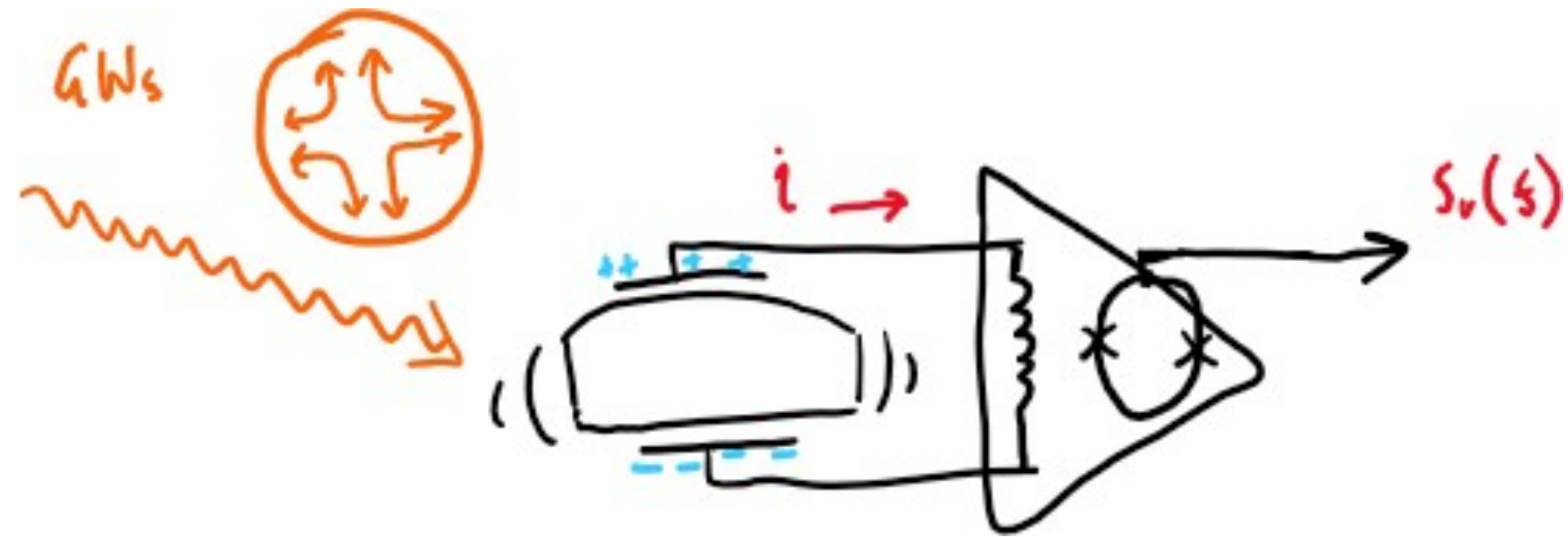
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- High density of modes from 1-1000 MHz
- Ongoing studies of behaviour at cryogenic temperatures



MAGE – Searching for new physics

Quartz BAW coupled to a DC SQUID amplifier \longrightarrow Highly sensitive resonant mass antenna

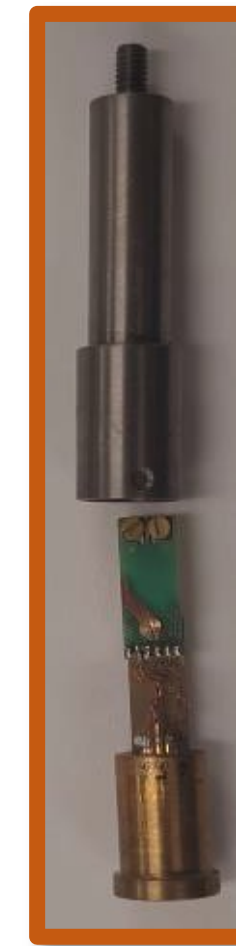


$$\ddot{B}_\lambda + \tau_\lambda^{-1} \dot{B}_\lambda + \omega_\lambda^2 B = -c^2 R_{i0j0} \int_V dv \frac{\rho}{m_\lambda} U_\lambda^i(\mathbf{x}) x^j$$

Primary target:

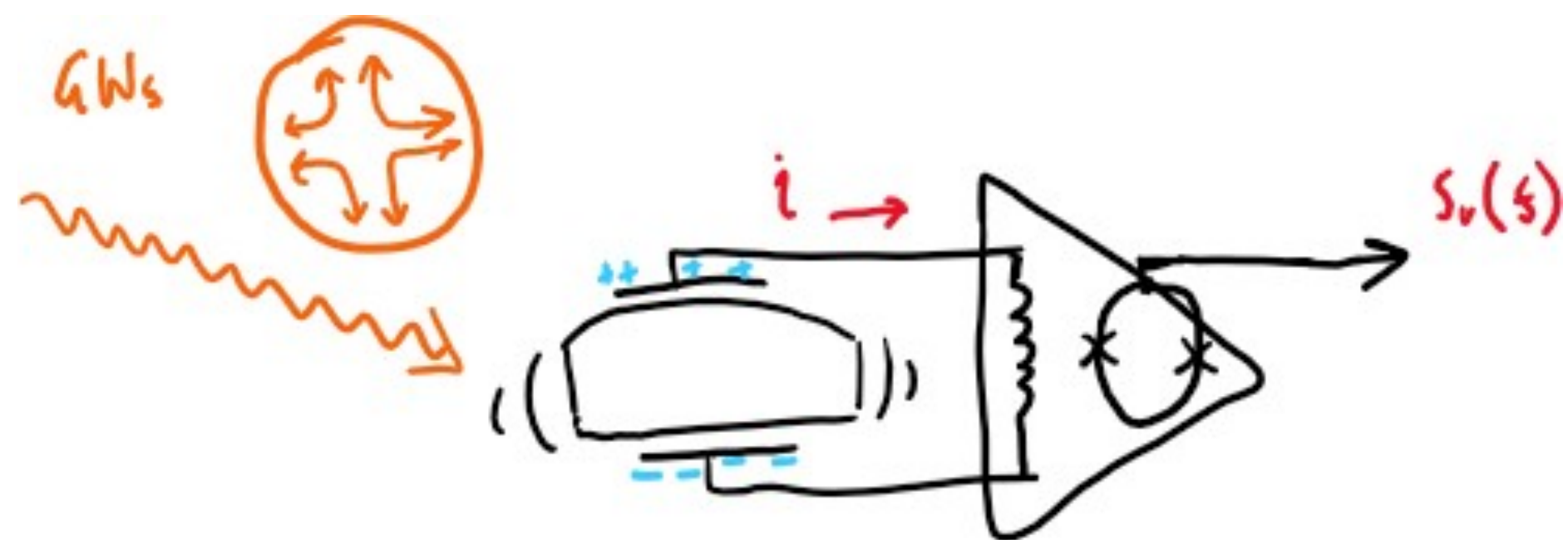
High frequency gravitational waves (MHz)

PRD 90, 102005 (2014)



MAGE – Searching for new physics

Quartz BAW coupled to a DC SQUID amplifier \longrightarrow Highly sensitive resonant mass antenna



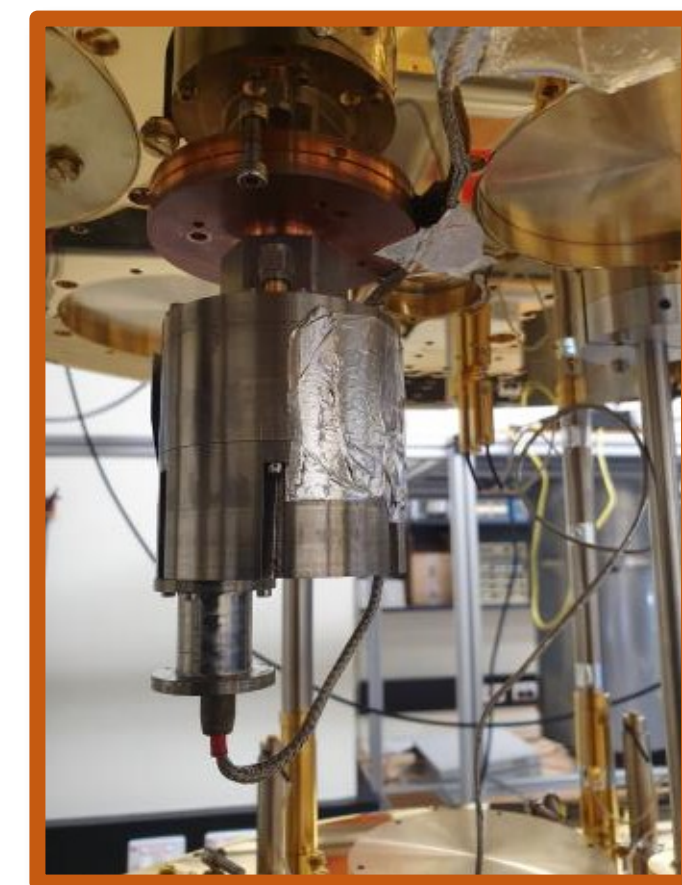
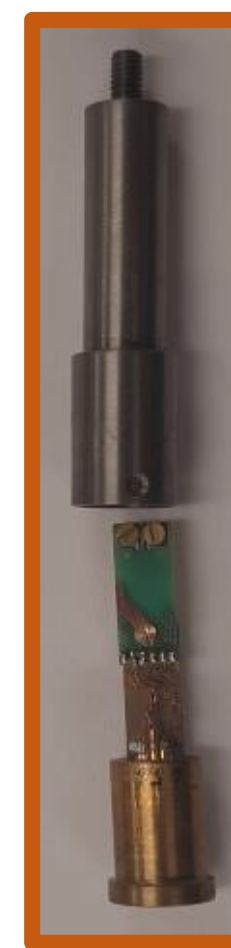
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Primary target:

High frequency gravitational waves (MHz) \longrightarrow

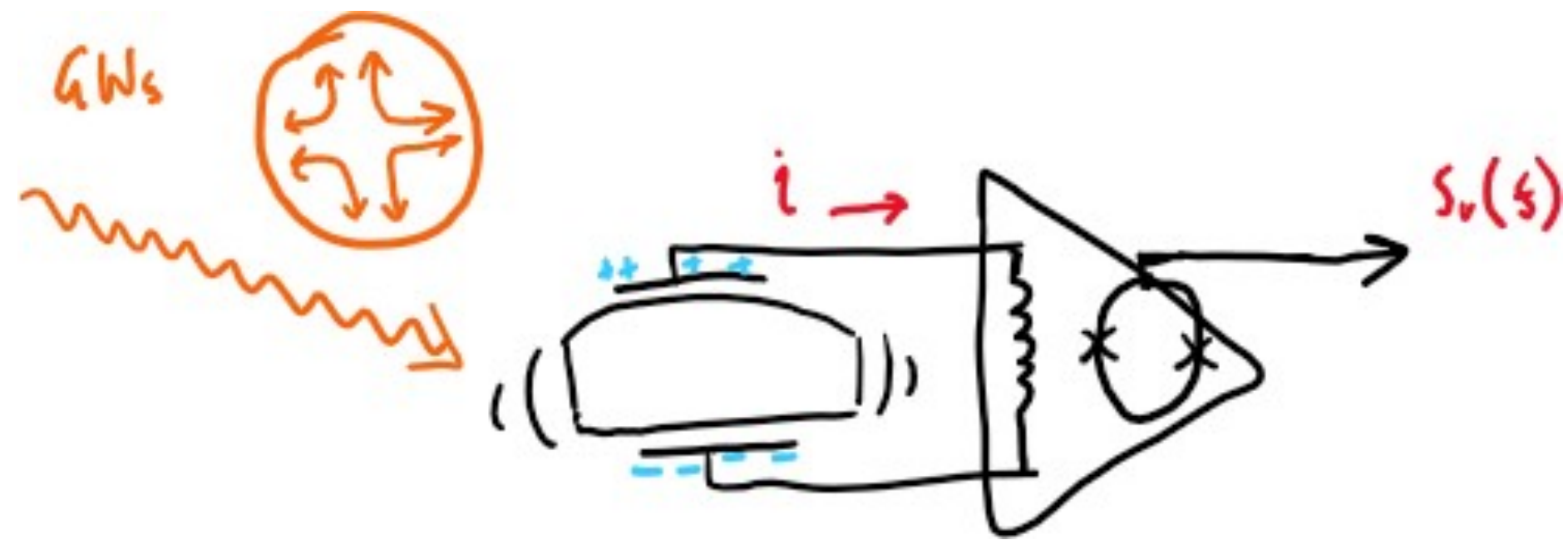
No known astrophysical sources exist at these frequencies

PRD 90, 102005 (2014)



MAGE – Searching for new physics

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$$\ddot{B}_\lambda + \tau_\lambda^{-1} \dot{B}_\lambda + \omega_\lambda^2 B = -c^2 R_{i0j0} \int_V dv \frac{\rho}{m_\lambda} U_\lambda^i(\mathbf{x}) x^j$$

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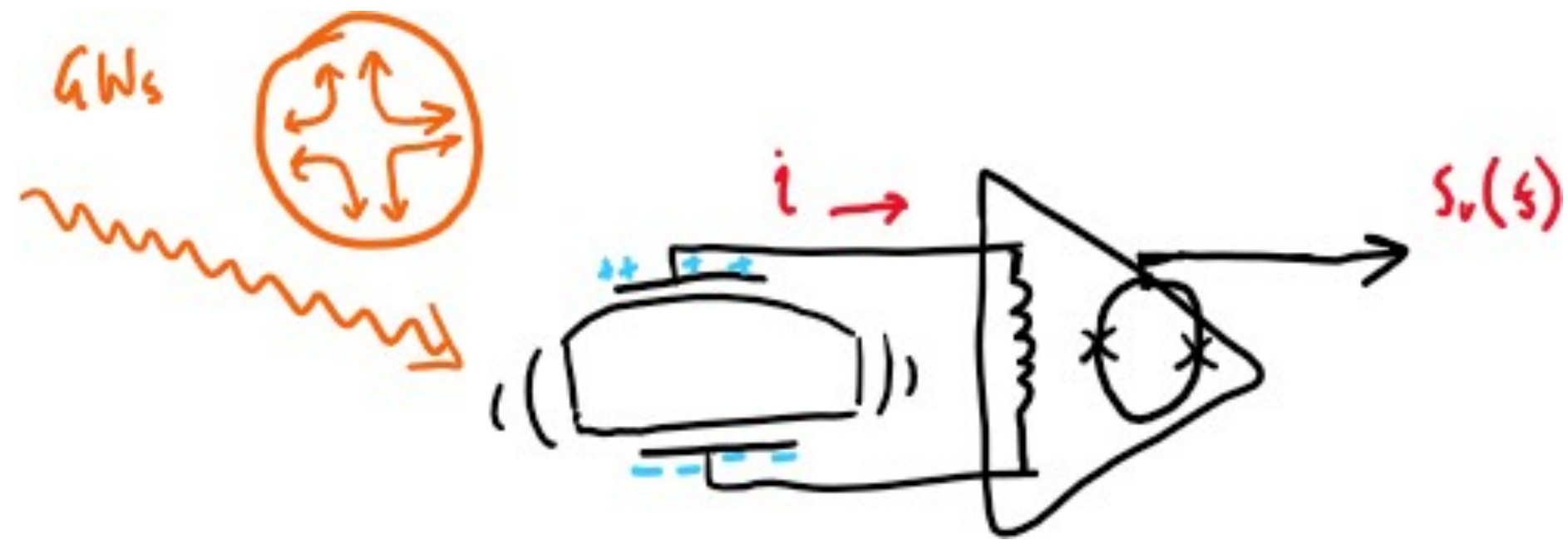
No known astrophysical sources exist at these frequencies

Any potential detection points to new physics outside the standard model!

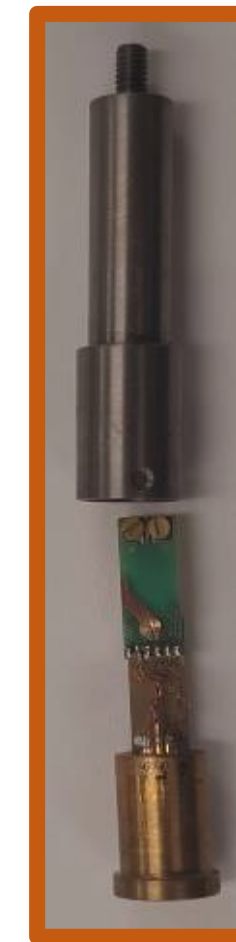


MAGE – Searching for new physics

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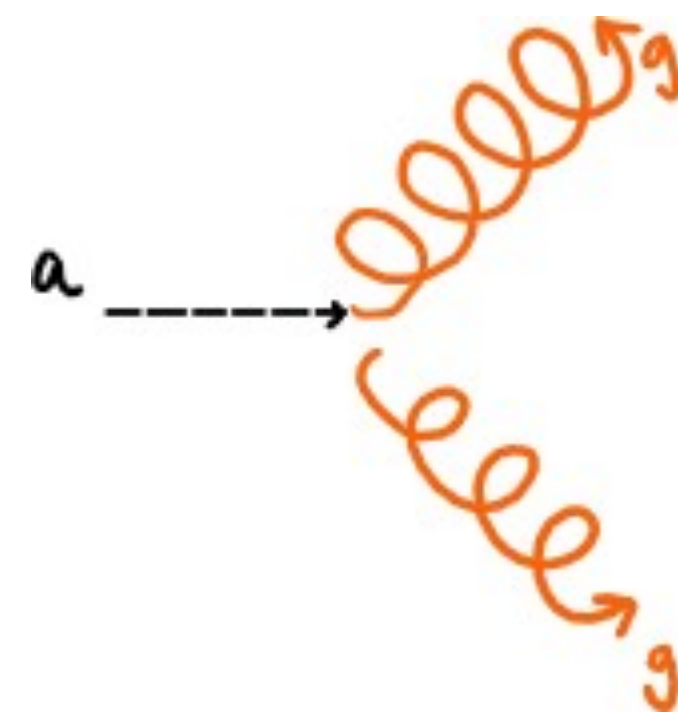
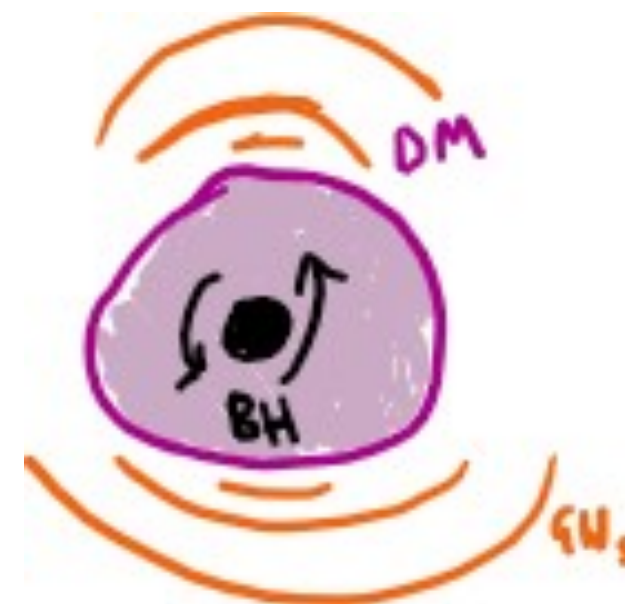


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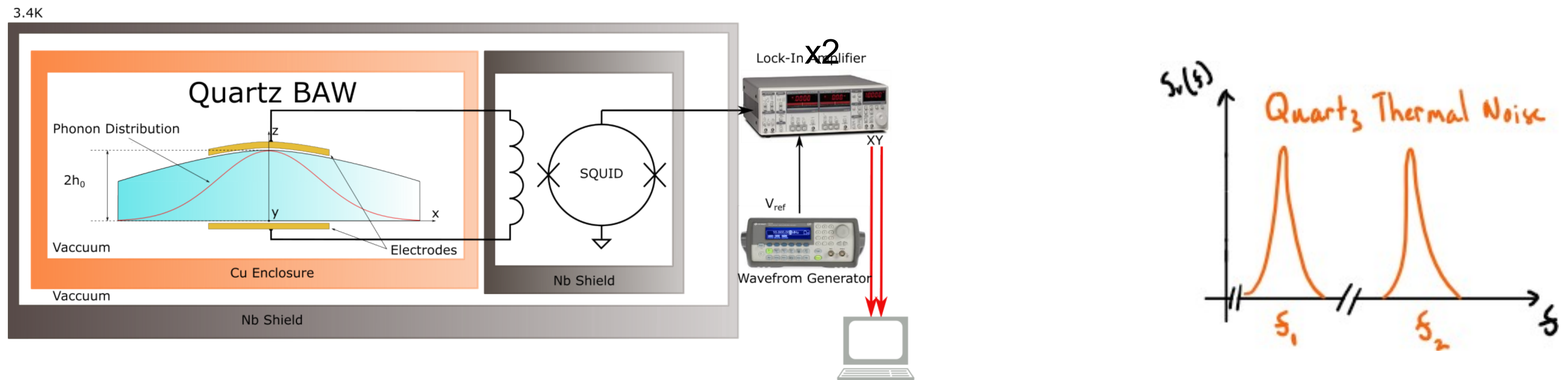
HFGWs due to DM:

Sub – solar black hole mergers, black hole super radiance, axion decay into gravitons



MAGE – Searching for new physics

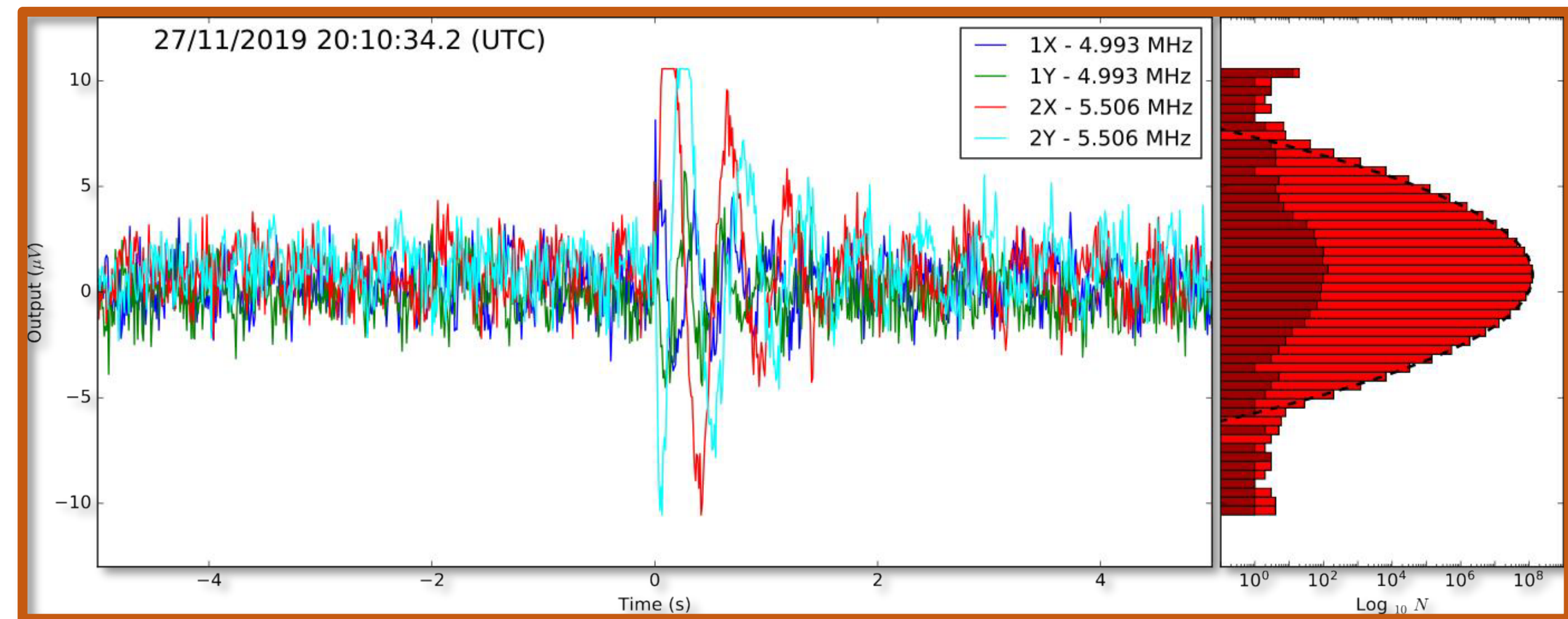
First Observational Period → GEN 1 & GEN 2, 153 days of data, two modes



Data Analysis:

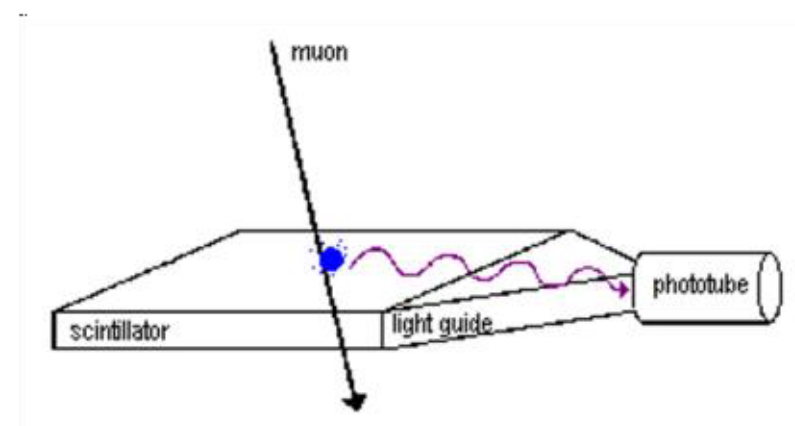
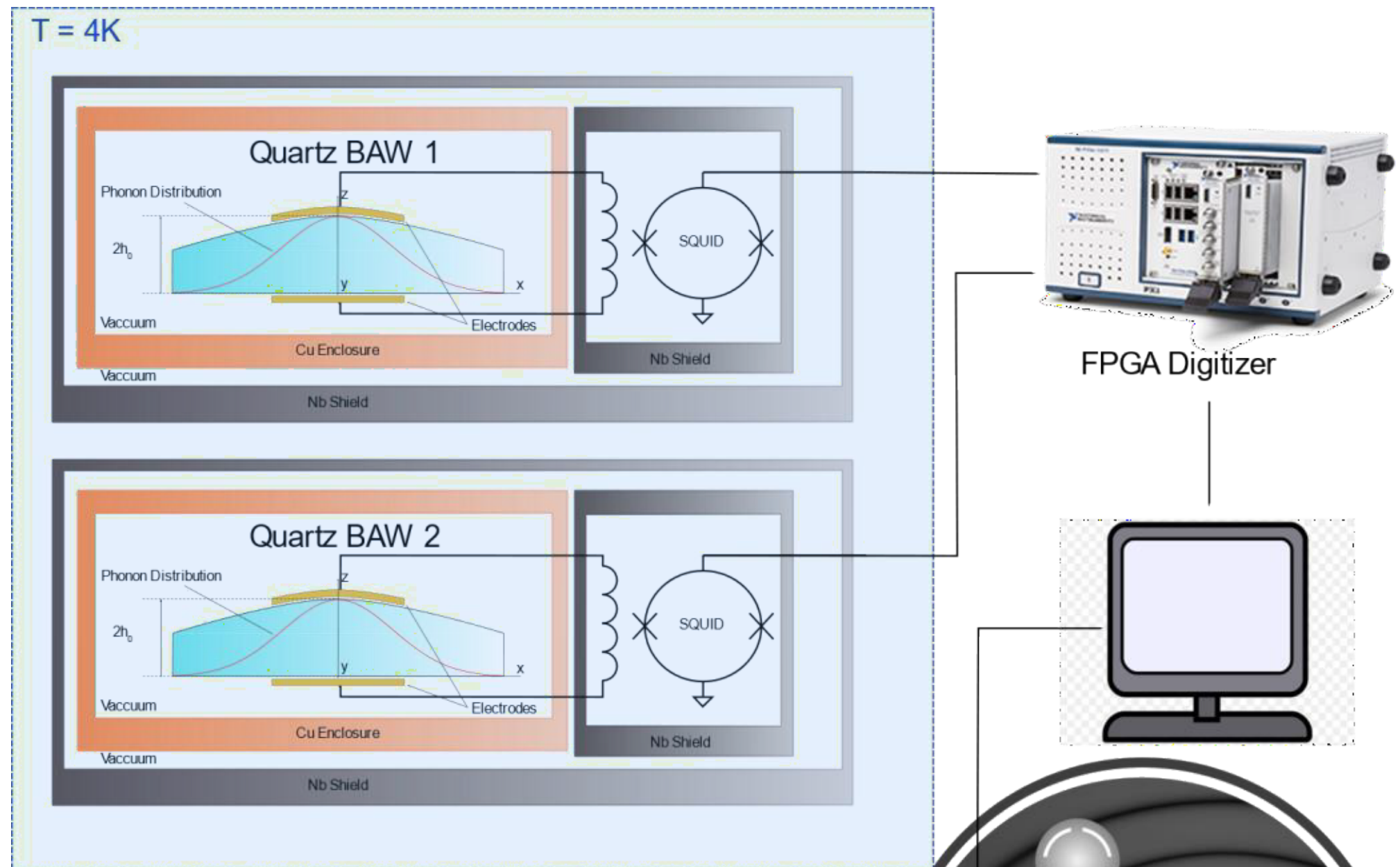
Two significantly strong, rare events

Phys. Rev. Lett. **127**, 071102



MAGE – Searching for new physics

What's next ? → Multimode Acoustic Gravitational Wave Experiment



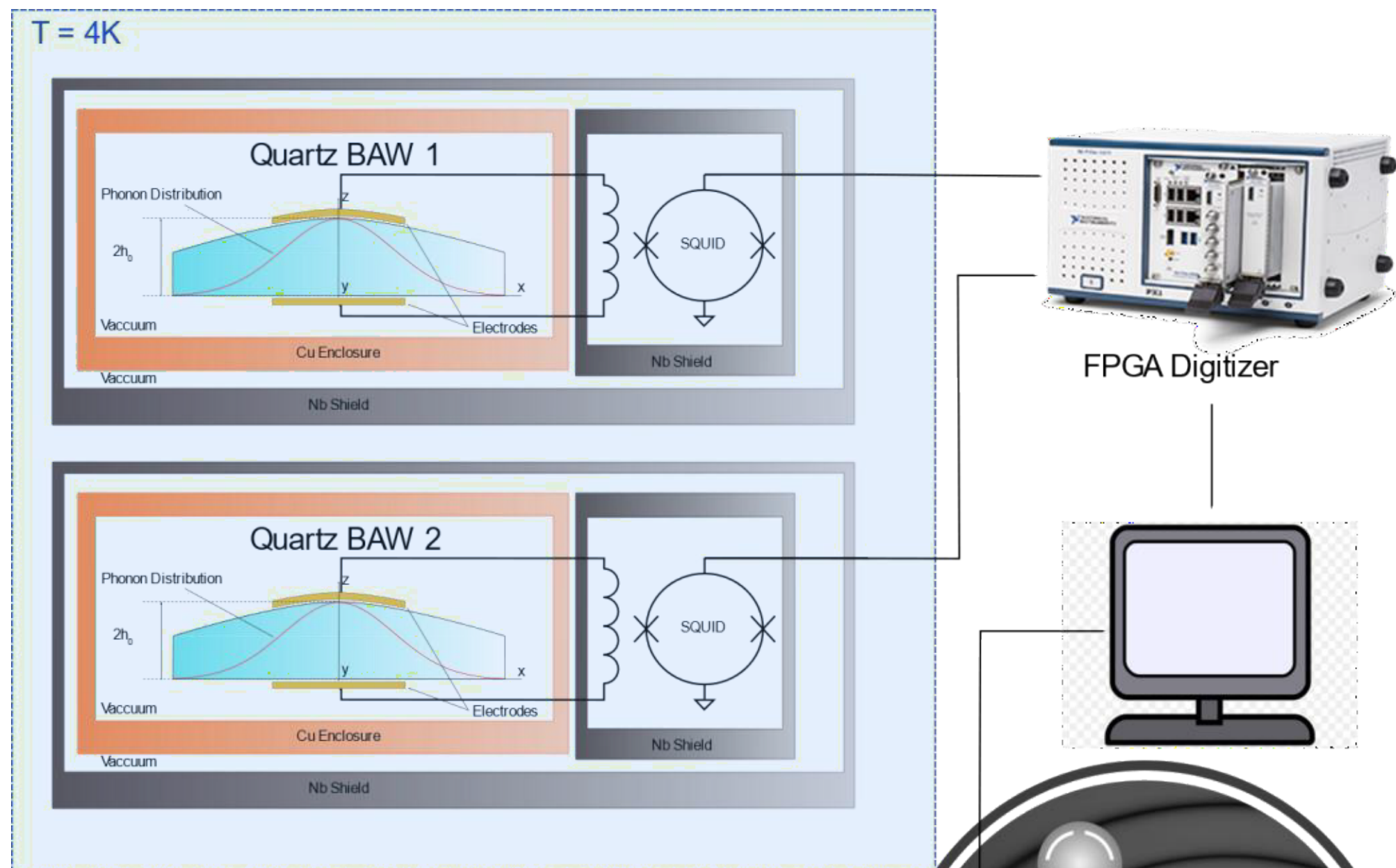
Muon / Cosmic Particle Veto Detector



MAGE – Searching for new physics

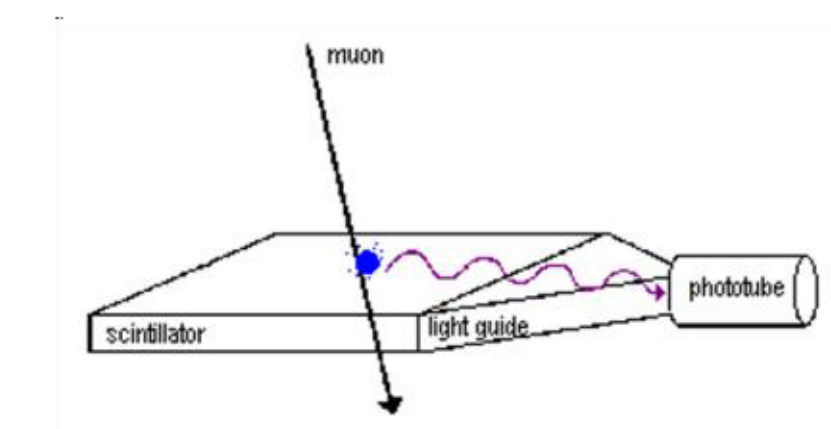
What's next ?

→ **M**ultimode **A**coustic **G**ravitational Wave **E**xperiment

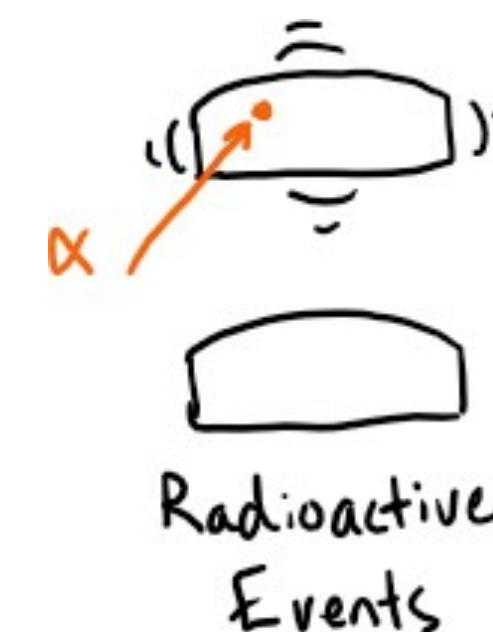
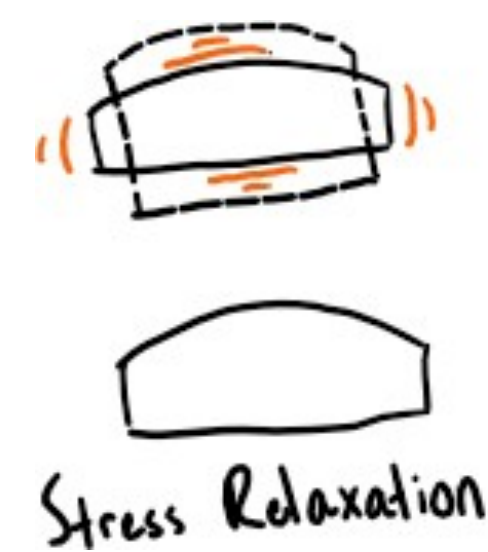


- 2 x Quartz BAW crystals
- 2 x DC SQUID amplifiers
- FPGA DAQ
- Cosmic particle veto (coming soon)

Exclude potential sources of events:



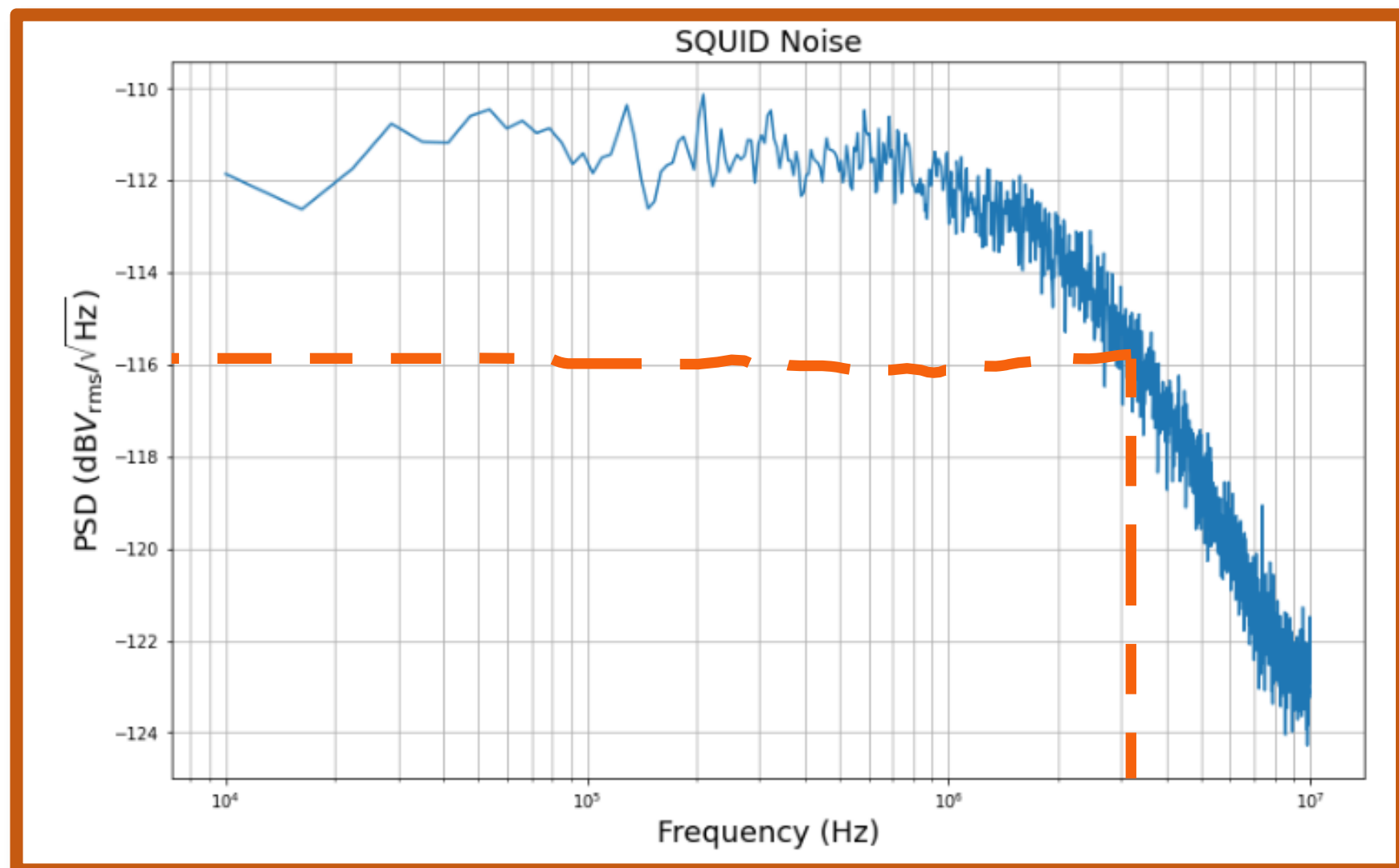
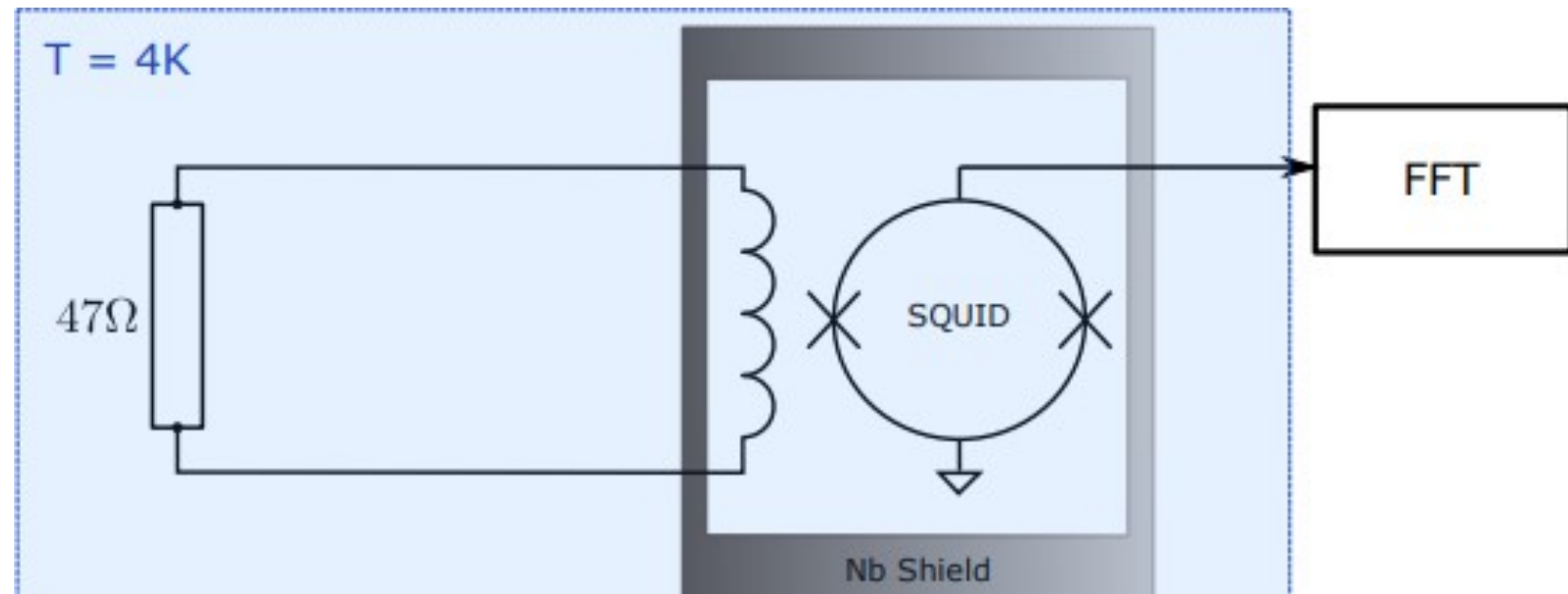
Muon / Cosmic Particle Veto Detector





MAGE – Searching for new physics

Calibration of 2nd detector



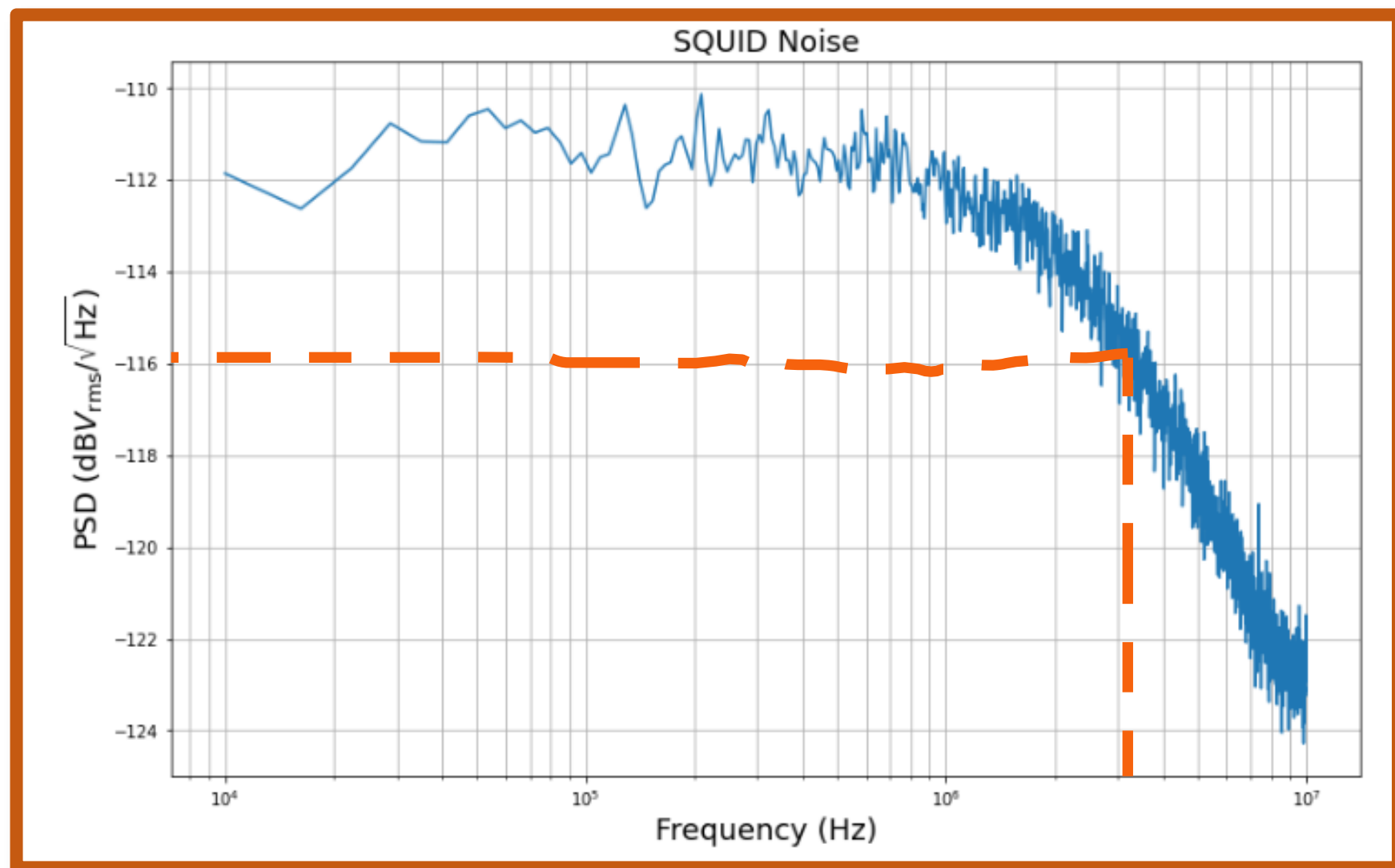
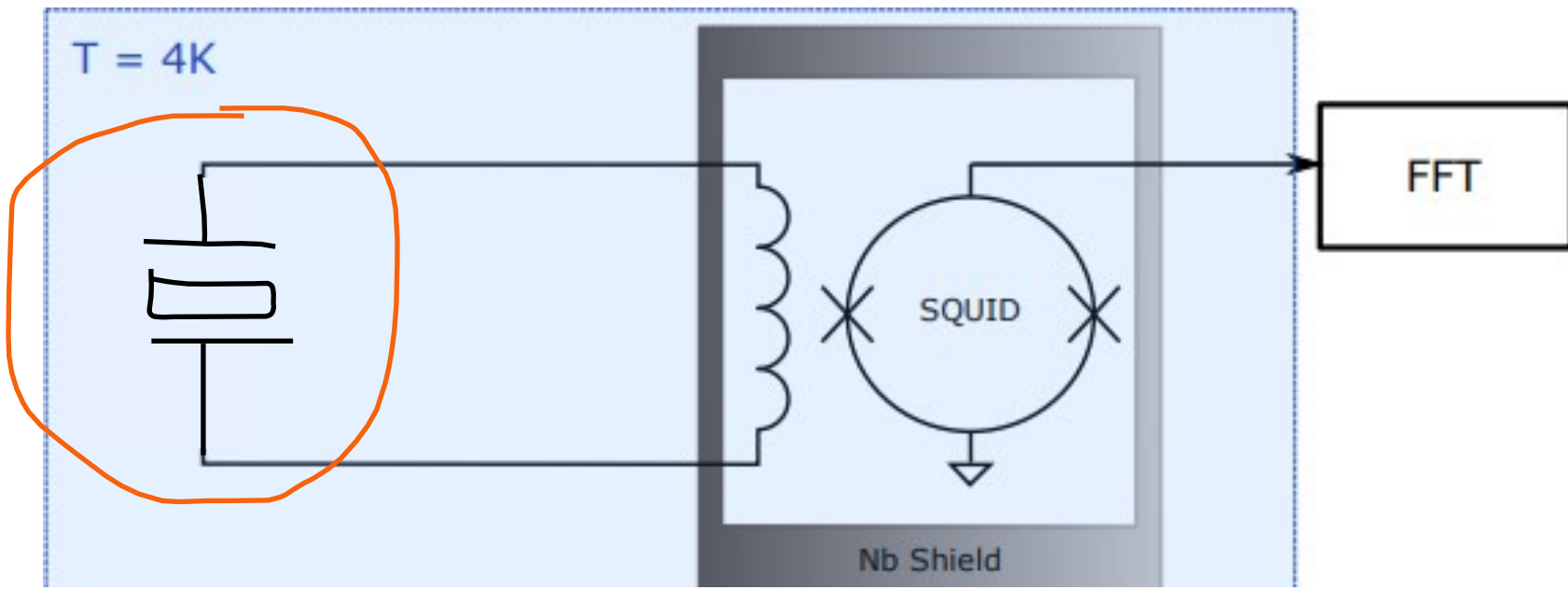
Limited by SQUID electronics

$f_{3dB} \sim 3 \text{ MHz}$



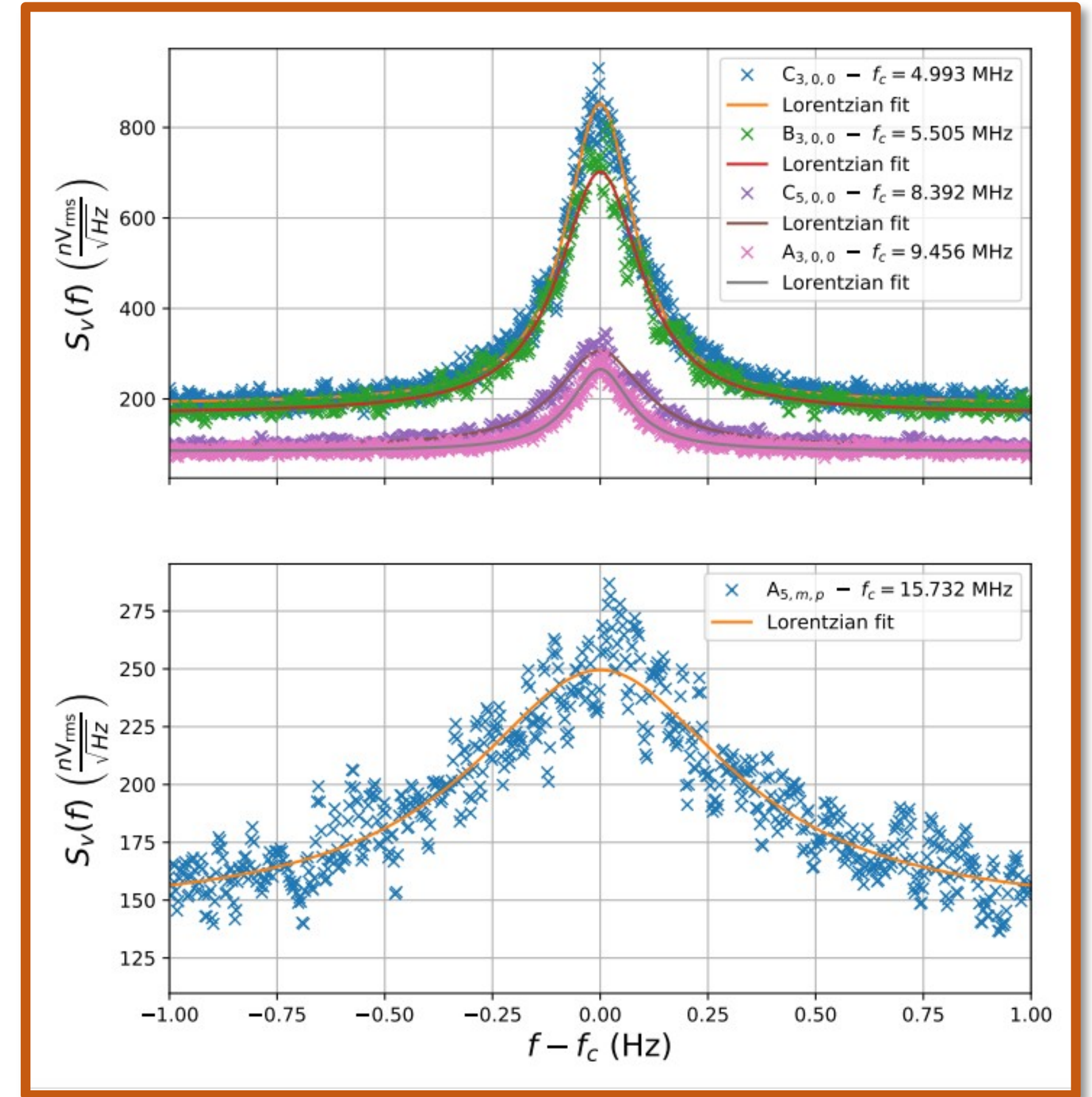
MAGE – Searching for new physics

Calibration of 2nd detector



$f_{3dB} \sim 3$ MHz

Modes up to 20 MHz are still observable





MAGE – Searching for new physics

Development of FPGA data acquisition

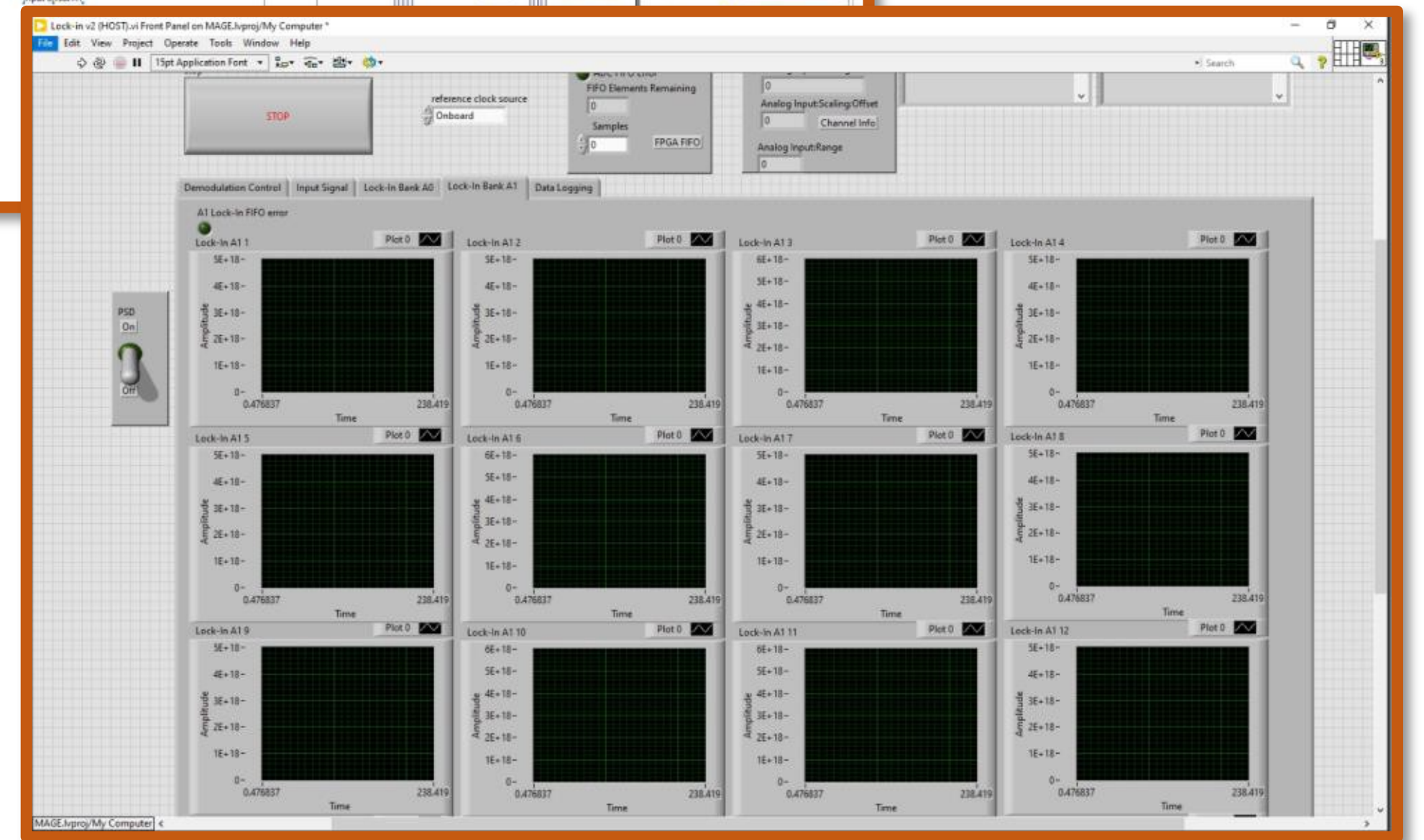
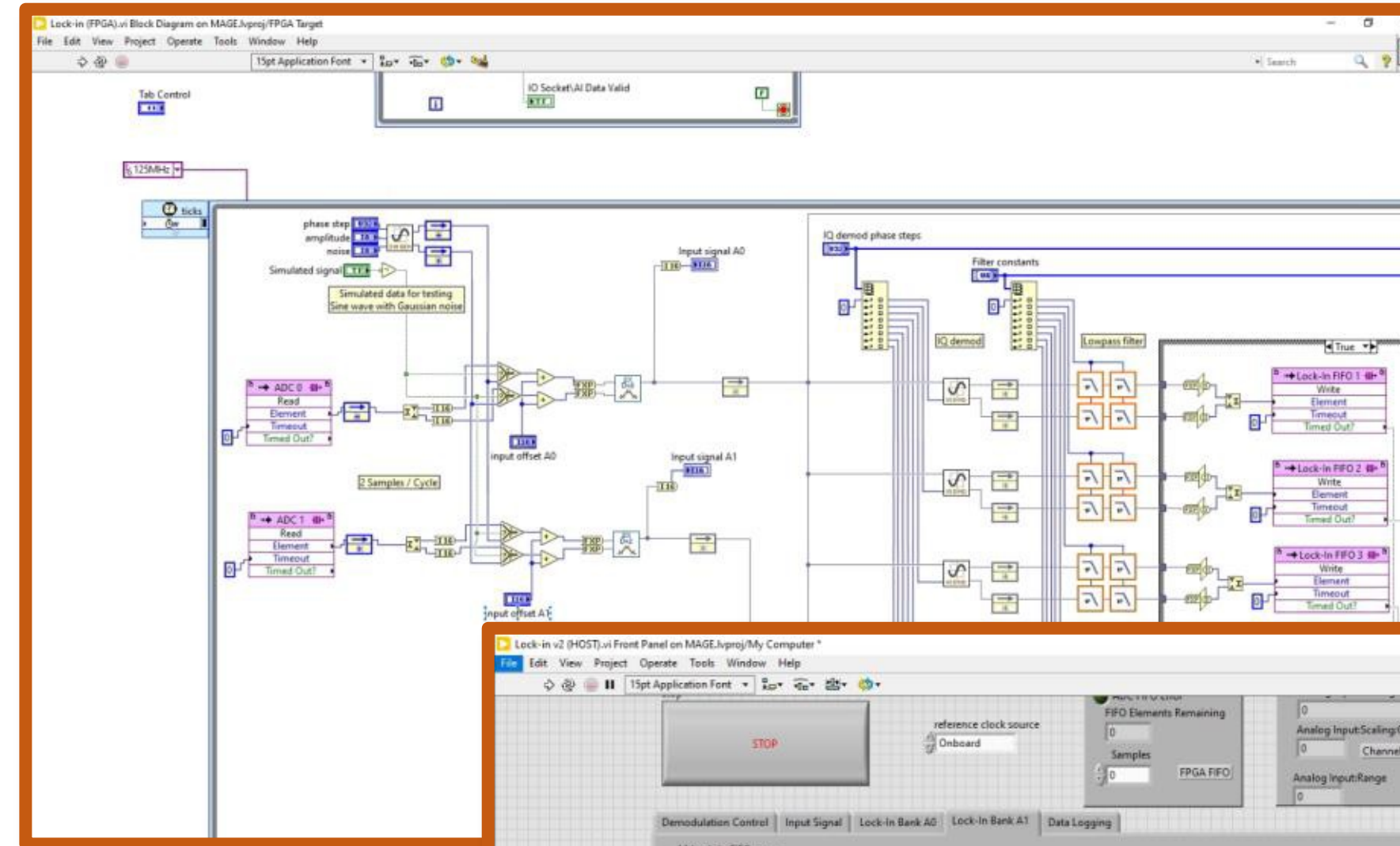
National Instruments – 5763 Digitizer
LabVIEW

32 Lock-in amplifiers across two inputs

Continuous data streaming & acquisition

In real time w/ strict timing & zero data loss

Yet to reach hardware limitation of device



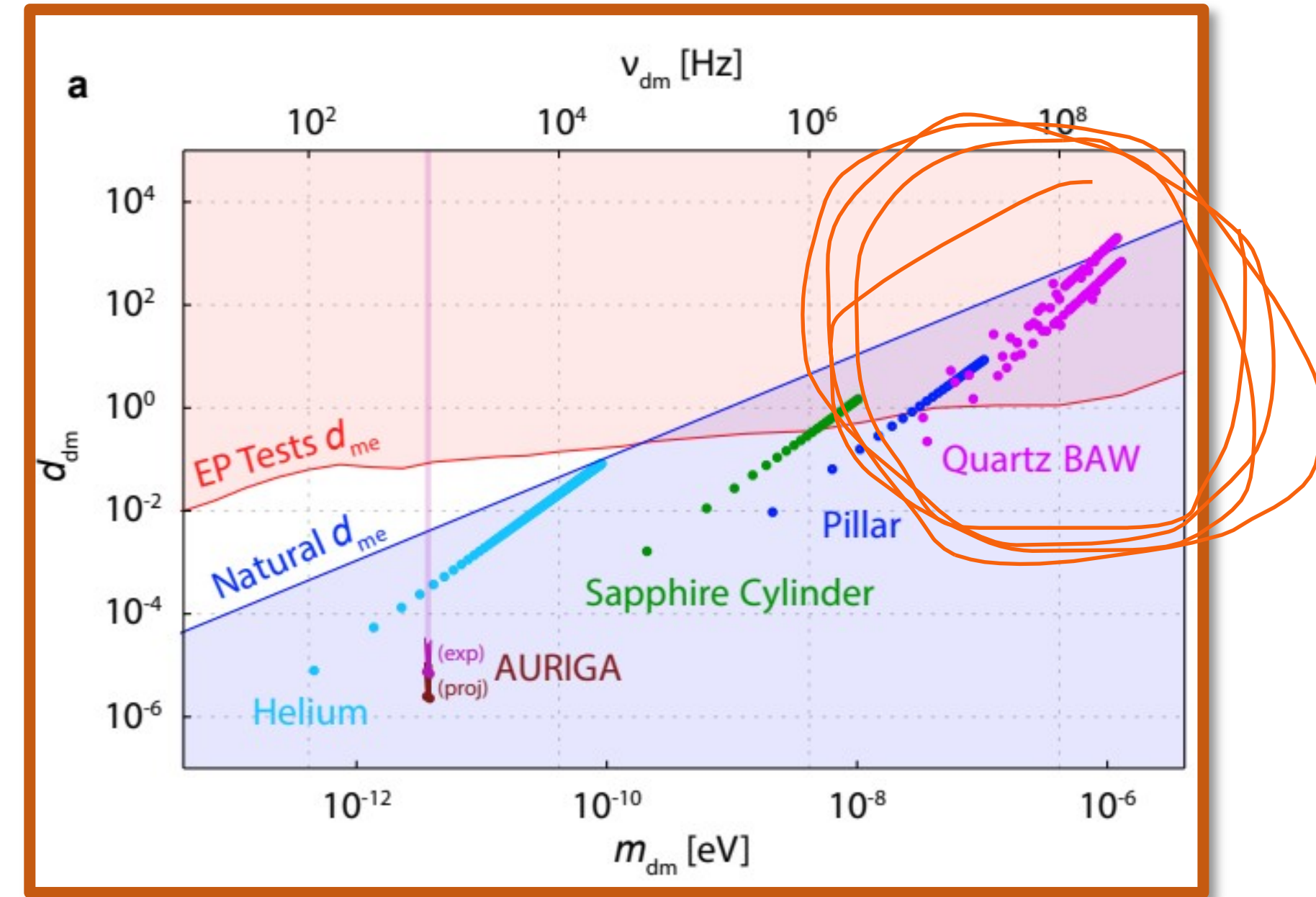
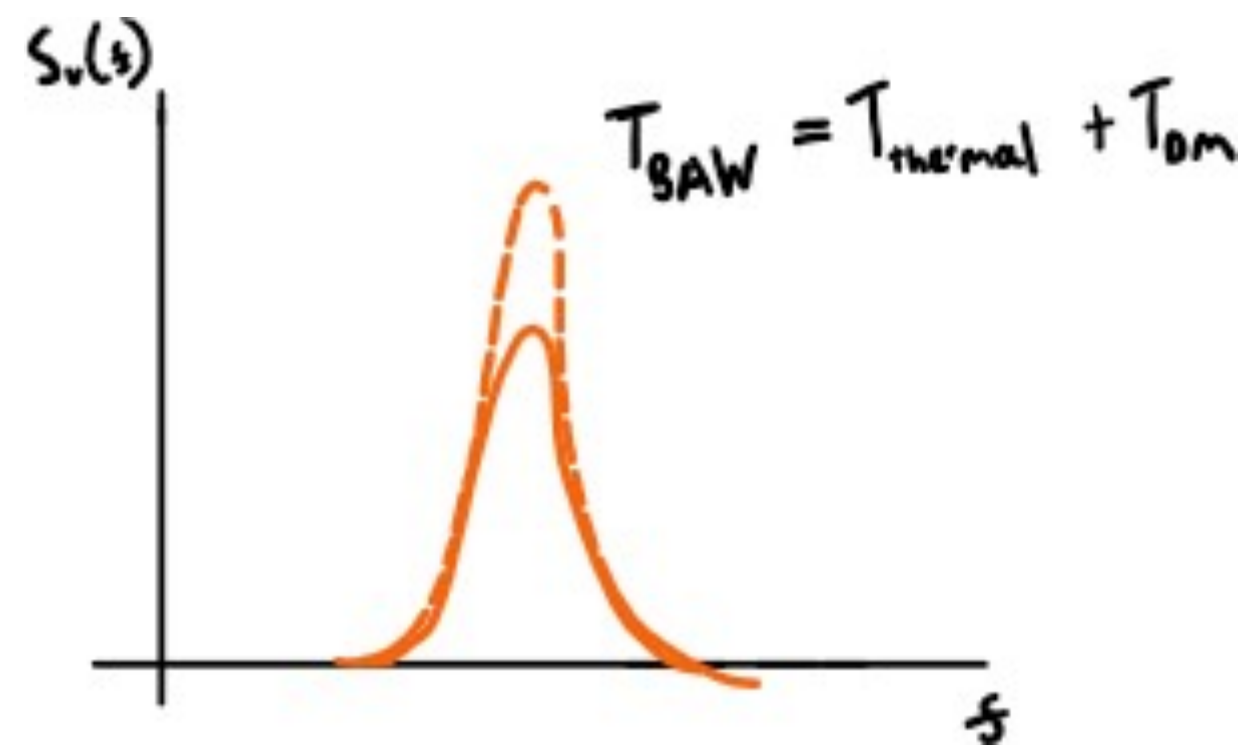
16 modes in each crystal
Currently taking data



MAGE – Searching for new physics

Other possibilities for MAGE:

Scalar DM \rightarrow Isotropic strain signal



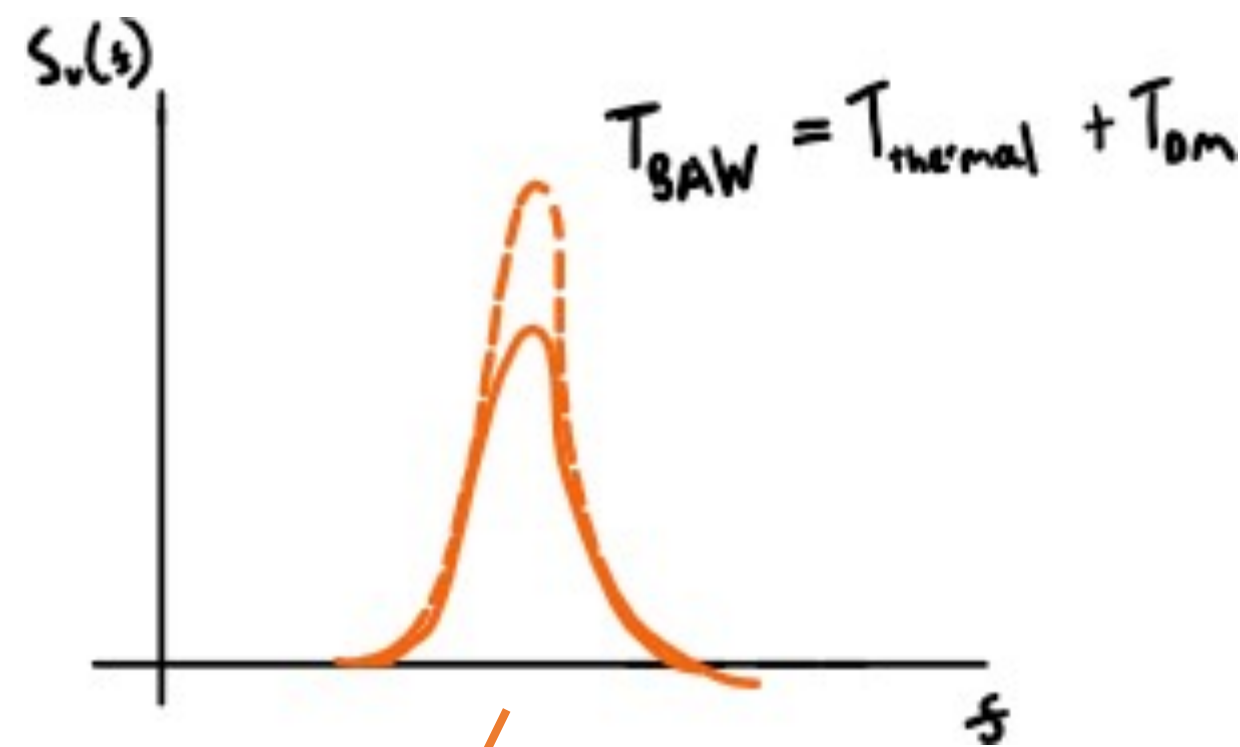
Phys. Rev. Lett. **124**, 151301



MAGE – Searching for new physics

Other possibilities for MAGE:

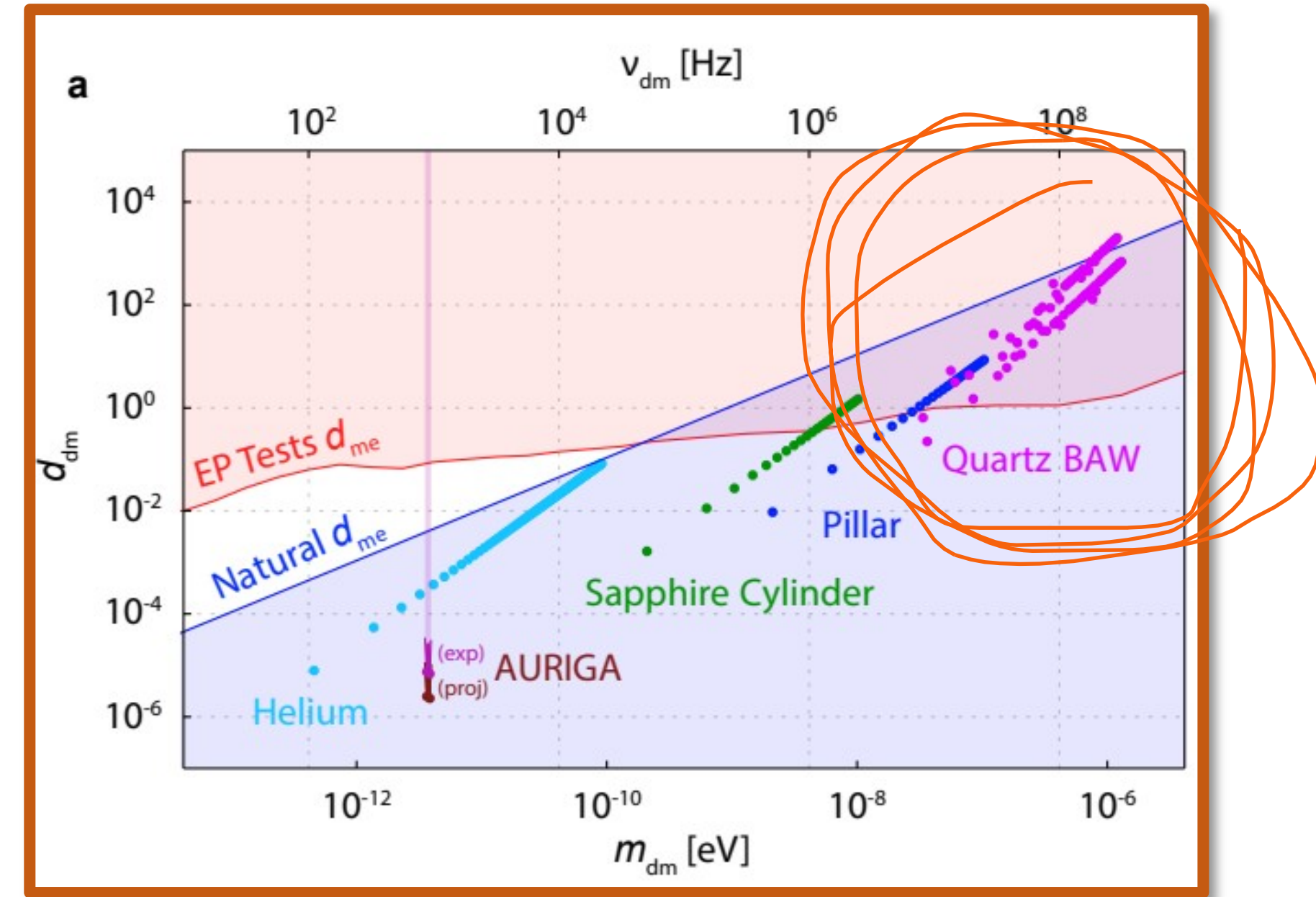
Scalar DM \rightarrow Isotropic strain signal



Nontrivial DM signal

Can exclude:
 Transient flows
 Daily modulation

Ongoing work:
 Resonance tuning



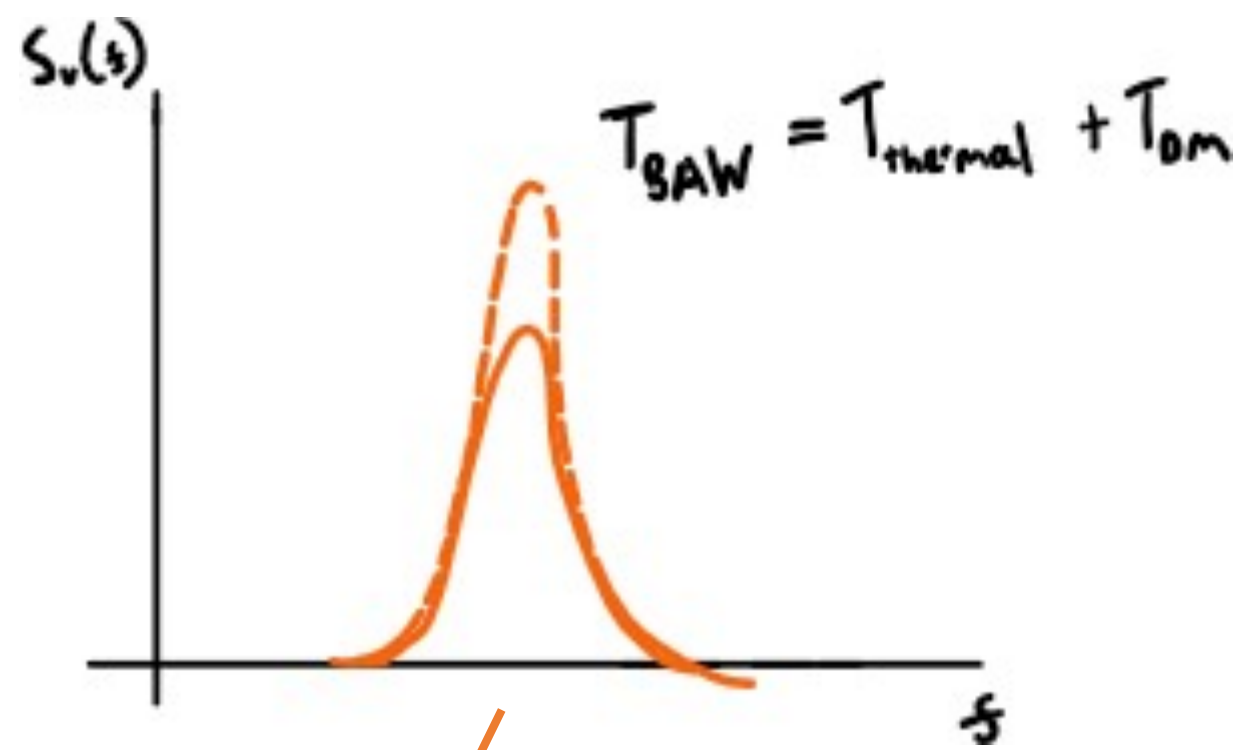
Phys. Rev. Lett. **124**, 151301



MAGE – Searching for new physics

Other possibilities for MAGE:

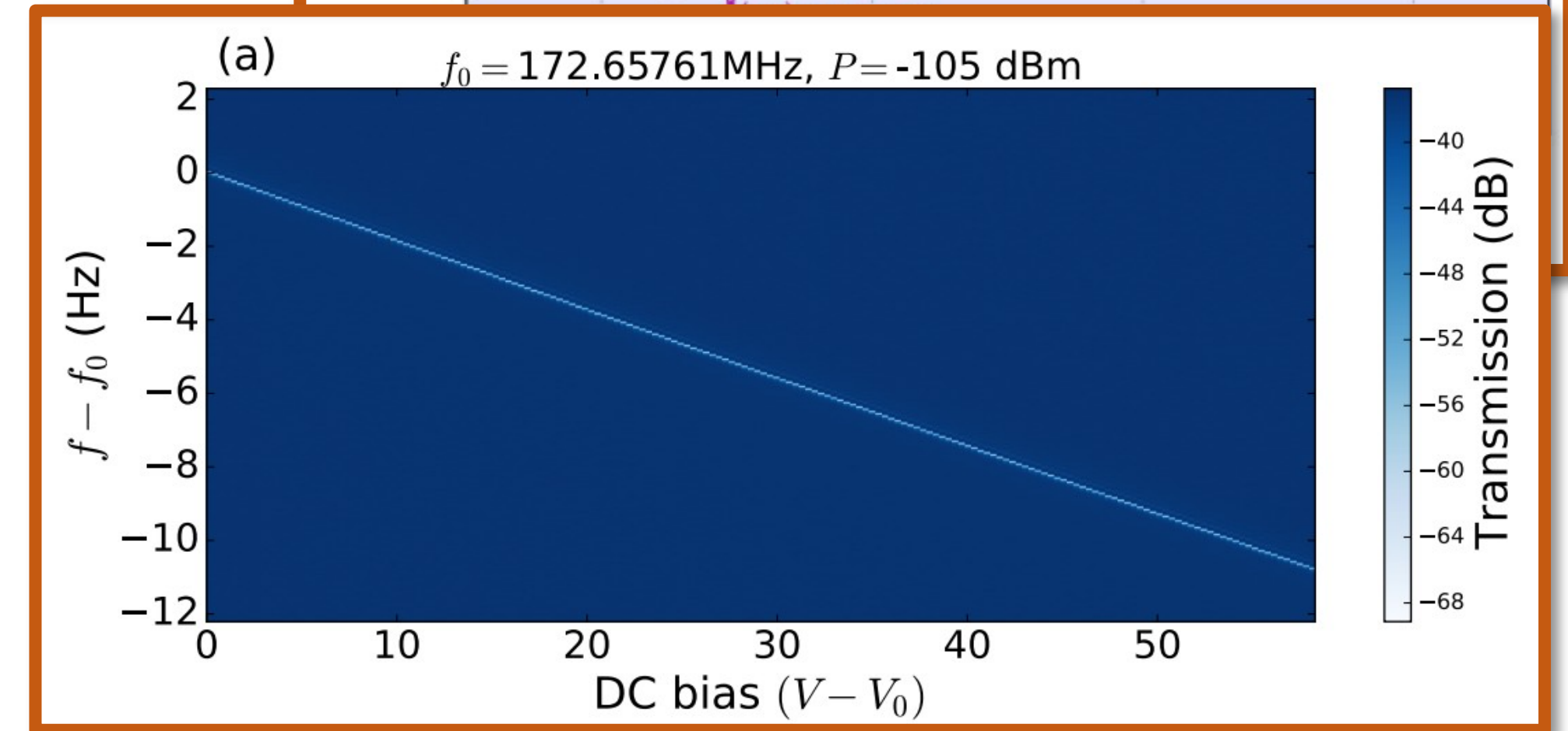
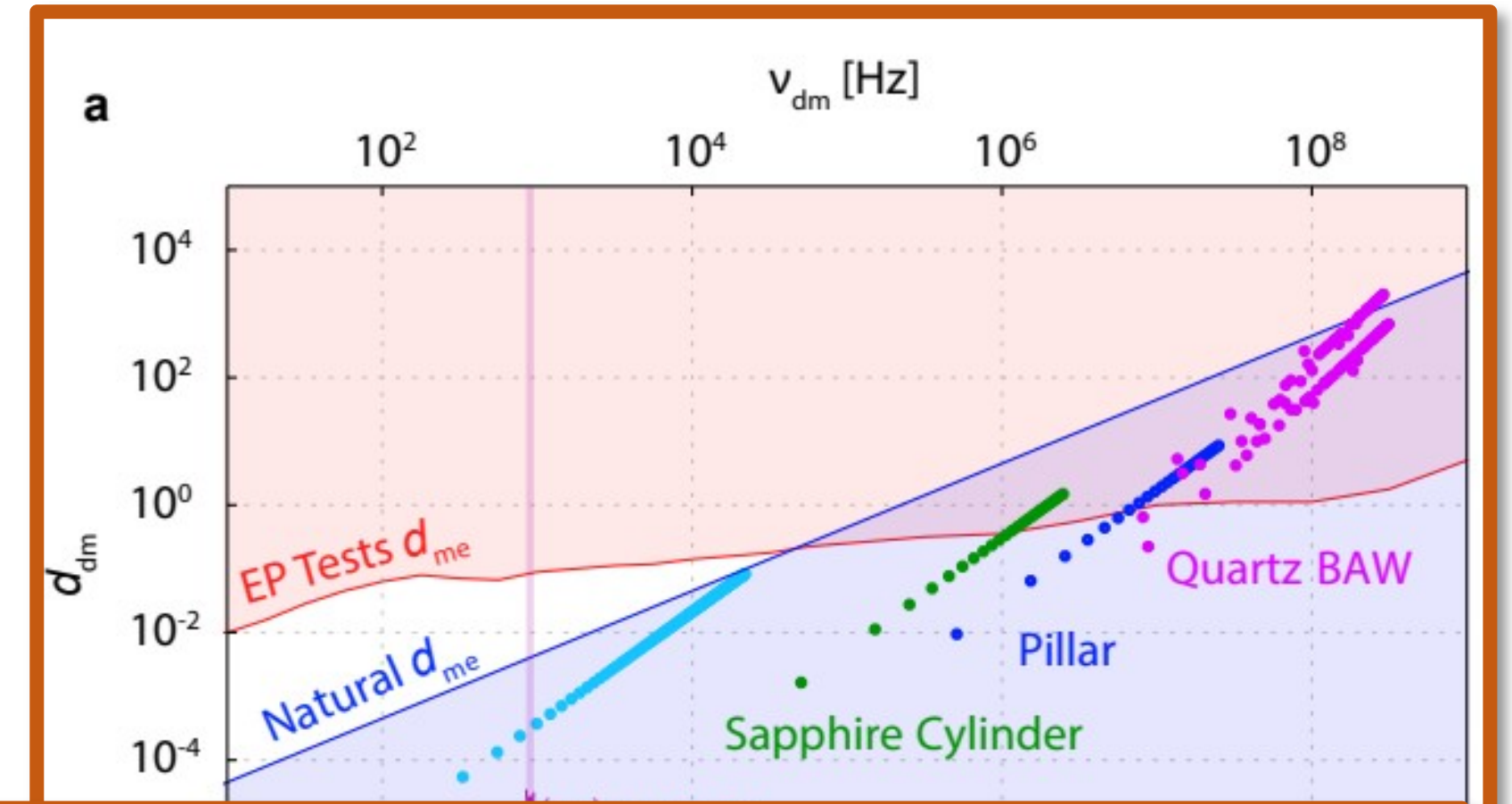
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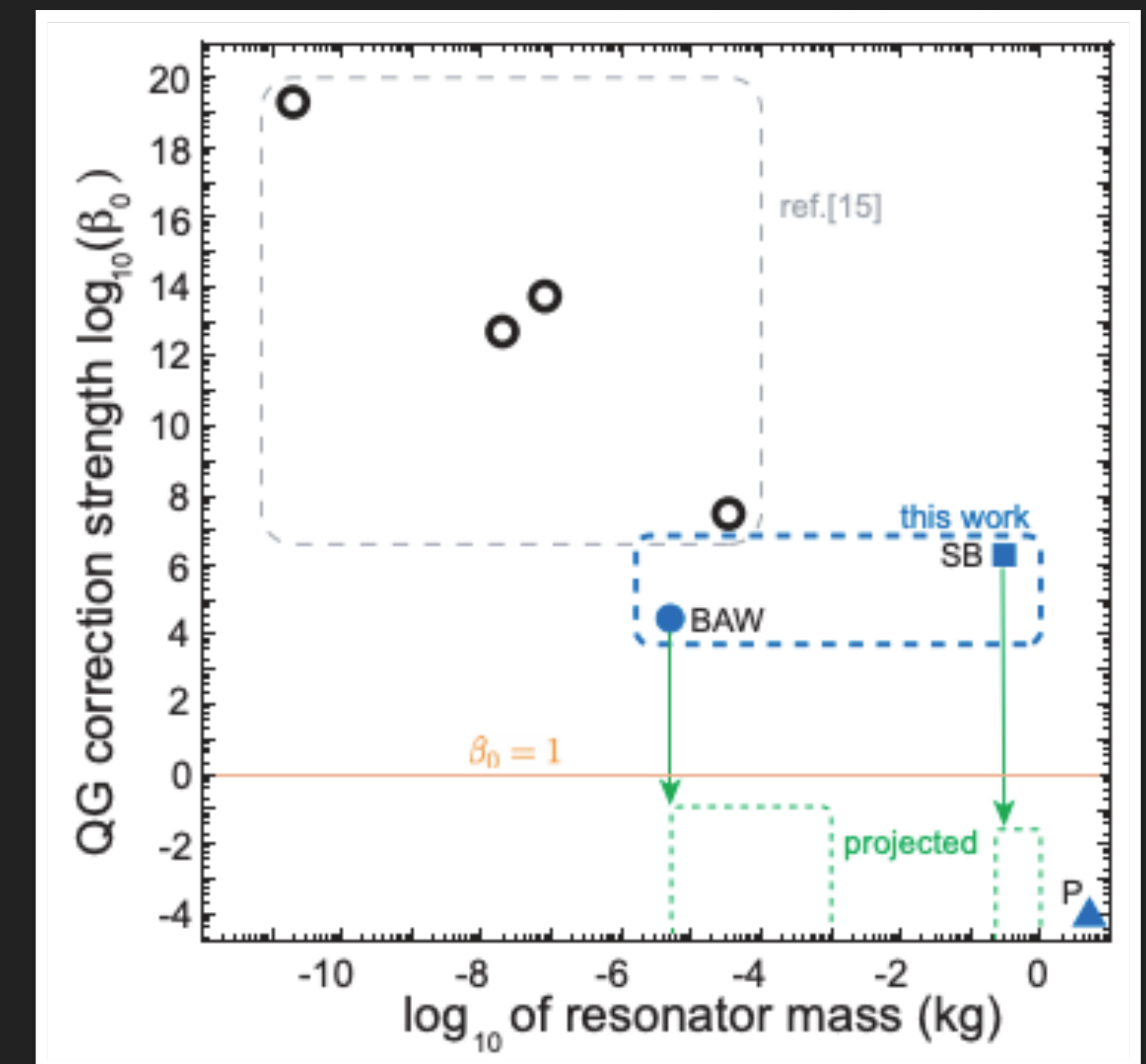
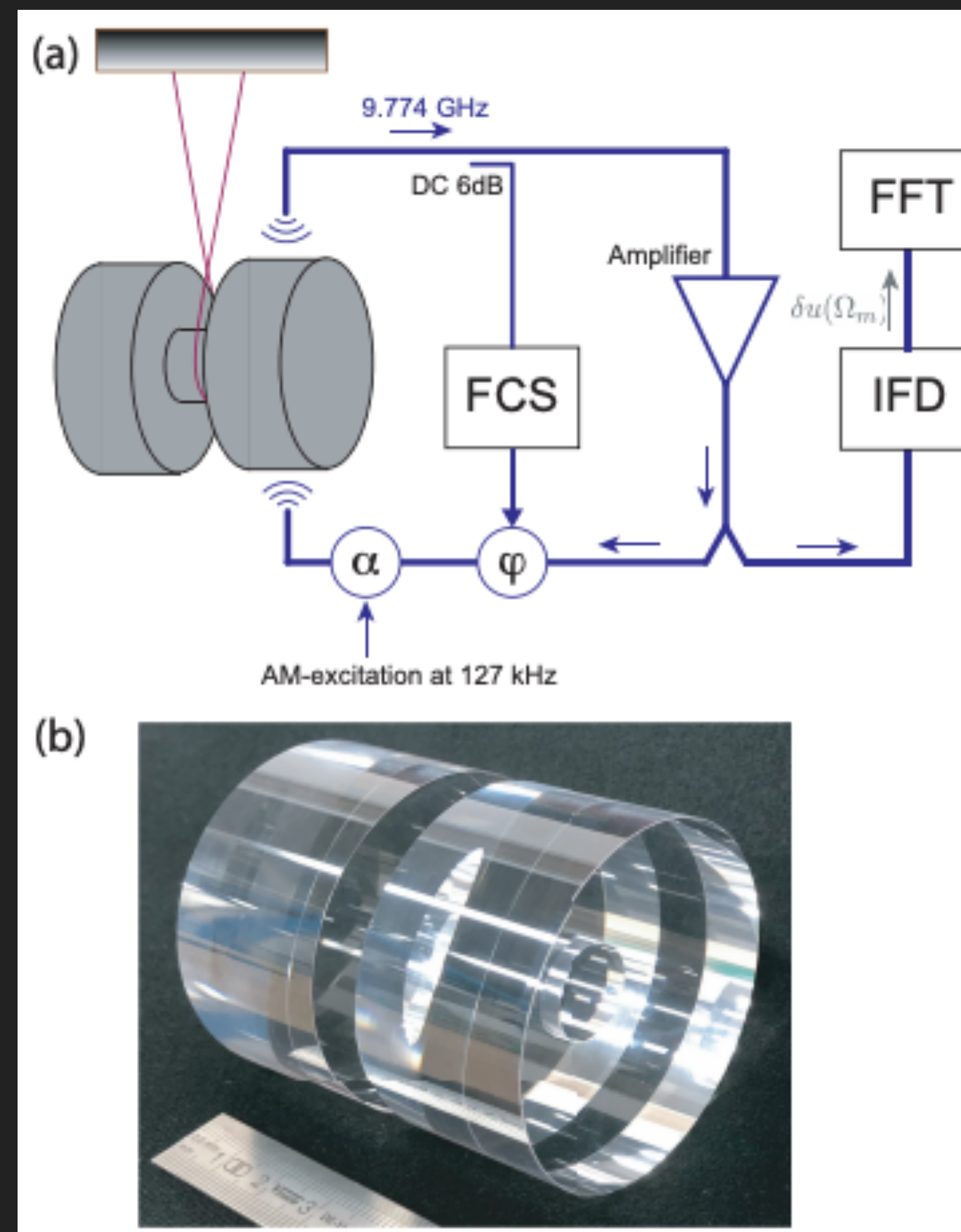
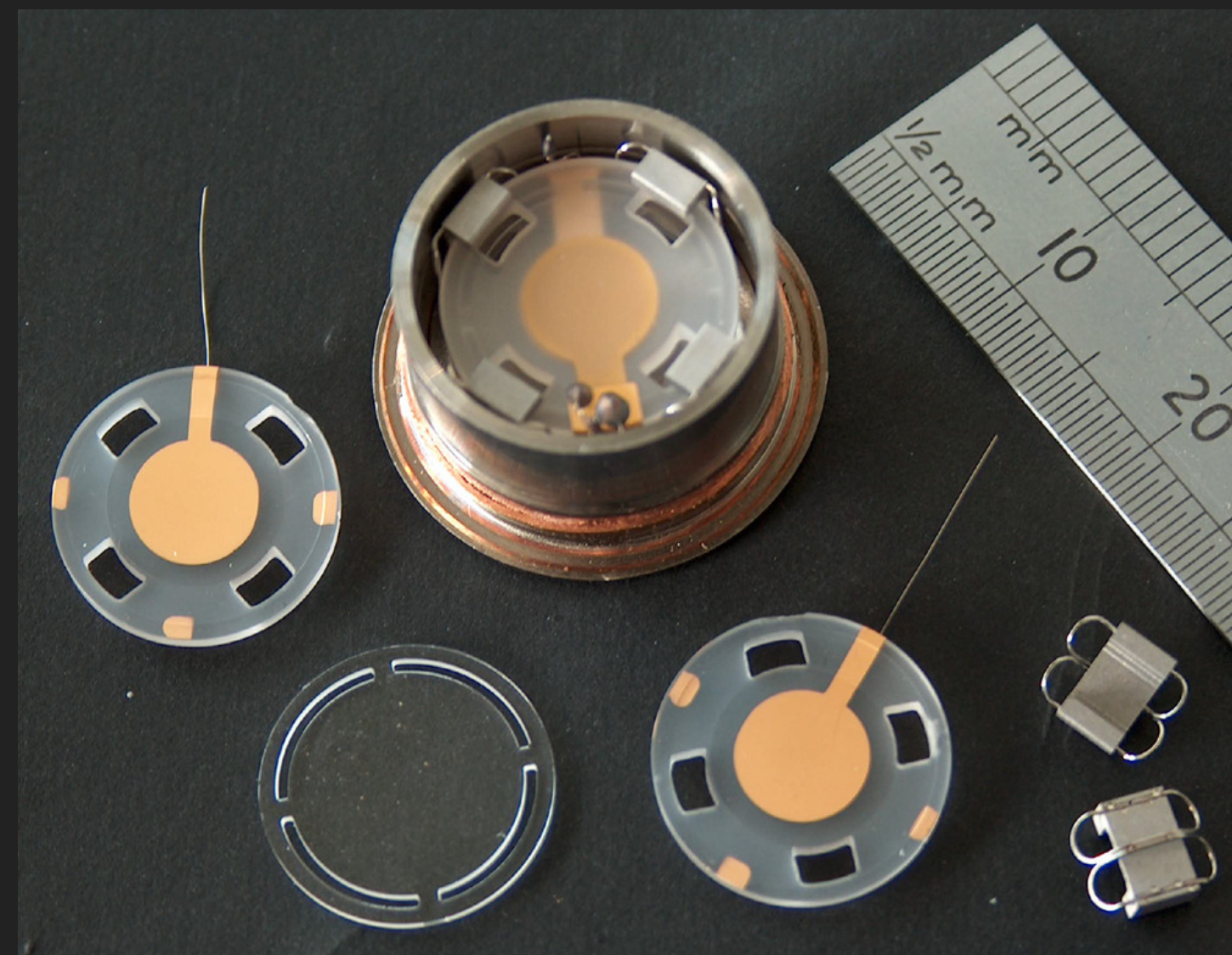
Ongoing work:
 Resonance tuning \rightarrow 2207.01176



Acoustic Tests of Quantum Gravity

Testing the generalized uncertainty principle with macroscopic mechanical oscillators and pendulums

P. A. Bushev, J. Bourhill, M. Goryachev, N. Kukharchyk, E. Ivanov, S. Galliou, M. E. Tobar, and S. Danilishin
Phys. Rev. D **100**, 066020 – Published 20 September 2019



Minimum Length

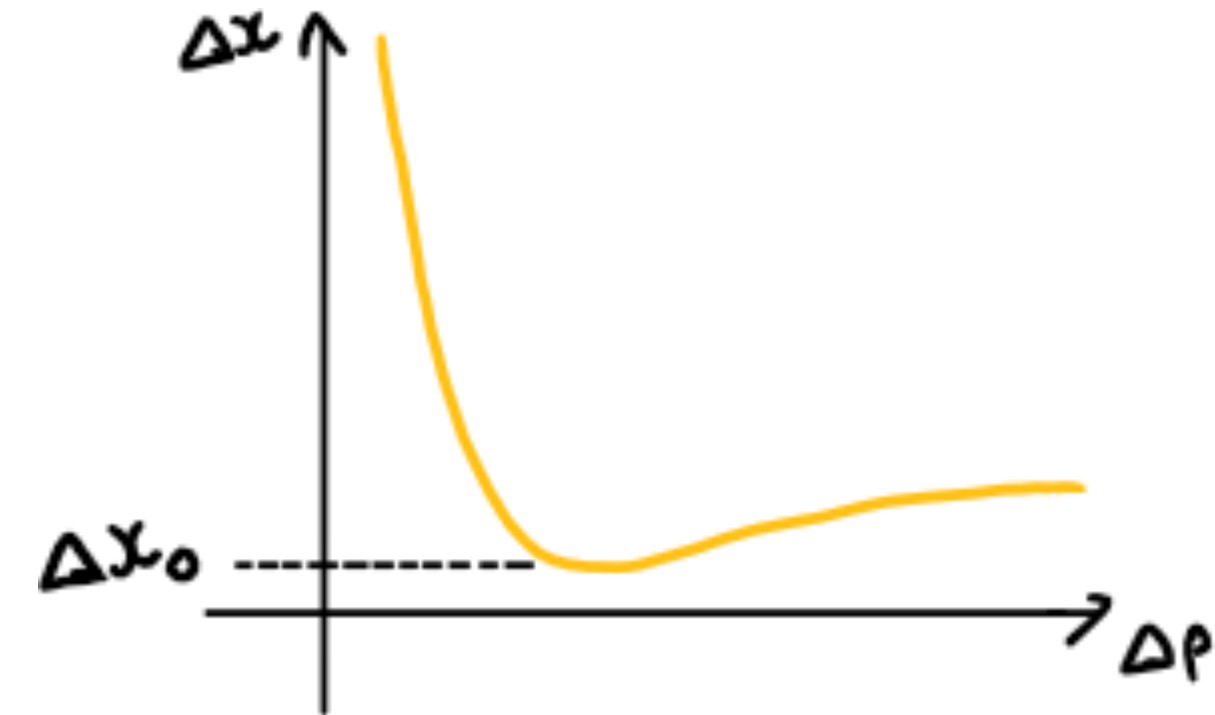
Quantum gravity models



Fundamental minimum length scale $\sim l_{\text{planck}}$



Generalized uncertainty principle $\longrightarrow \Delta x \Delta p \geq \frac{\hbar}{2} \left[1 + \beta_0 \left(\frac{l_p \Delta p}{\hbar} \right)^2 \right],$

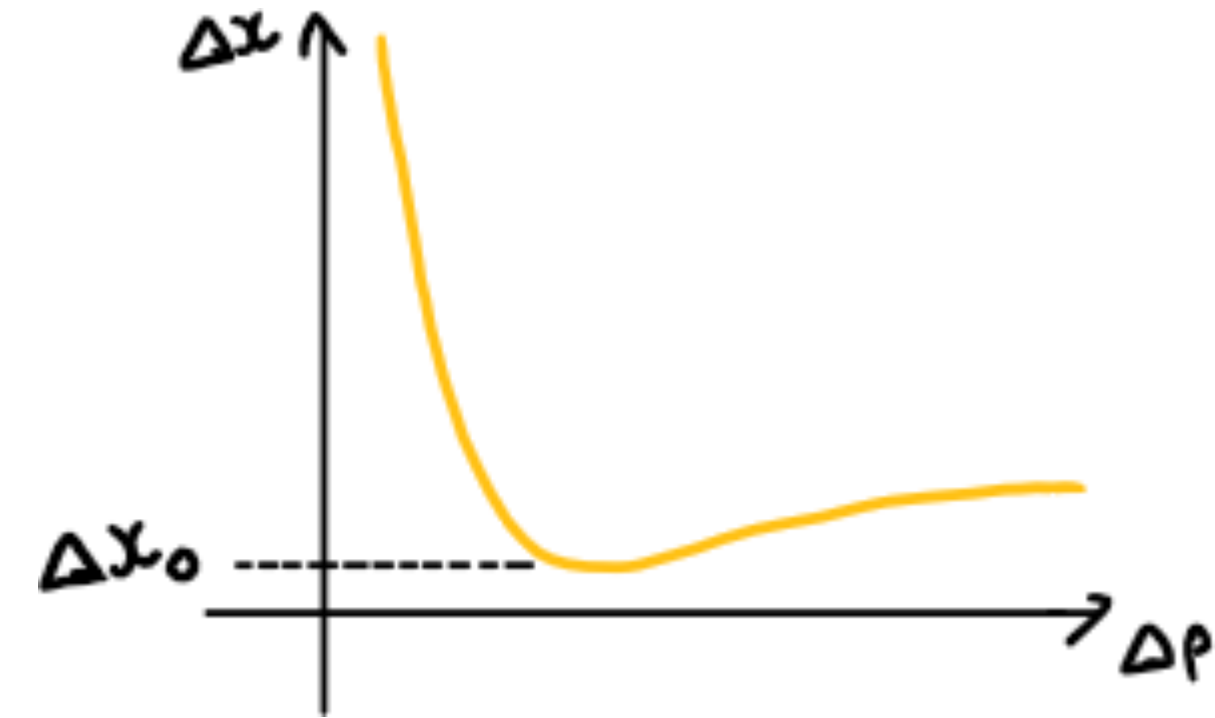


Minimum Length

How can we test it ?

$$l_{planck} = 1.62 \times 10^{-35} \text{ m}$$

$$E_{planck} = 1.2 \times 10^9 \text{ GeV}$$



Minimum Length

How can we test it ?

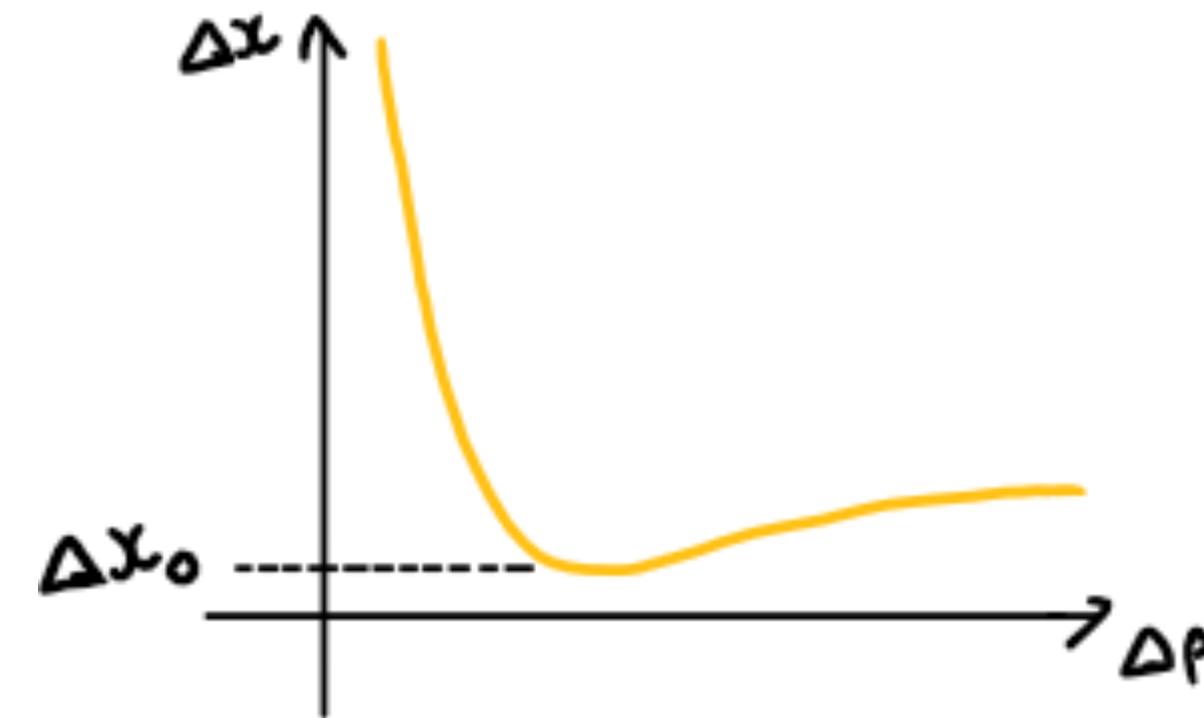
Consider *deformed* Heisenberg algebra of harmonic oscillator

$$[\hat{x}, \hat{p}] = i\hbar \left(1 + \beta_0 \left(\frac{l_p}{\hbar} \hat{p} \right)^2 \right)$$

Equations of motion of harmonic oscillator

$$Q = Q_0 \left[\sin(\tilde{\omega}t) + \frac{\beta}{8} Q_0^2 \sin(3\tilde{\omega}t) \right],$$

$$\tilde{\omega} = \left(1 + \frac{\beta}{2} Q_0^2 \right) \omega_0.$$



Minimum Length

How can we test it?

Consider *deformed* Heisenberg algebra of harmonic oscillator

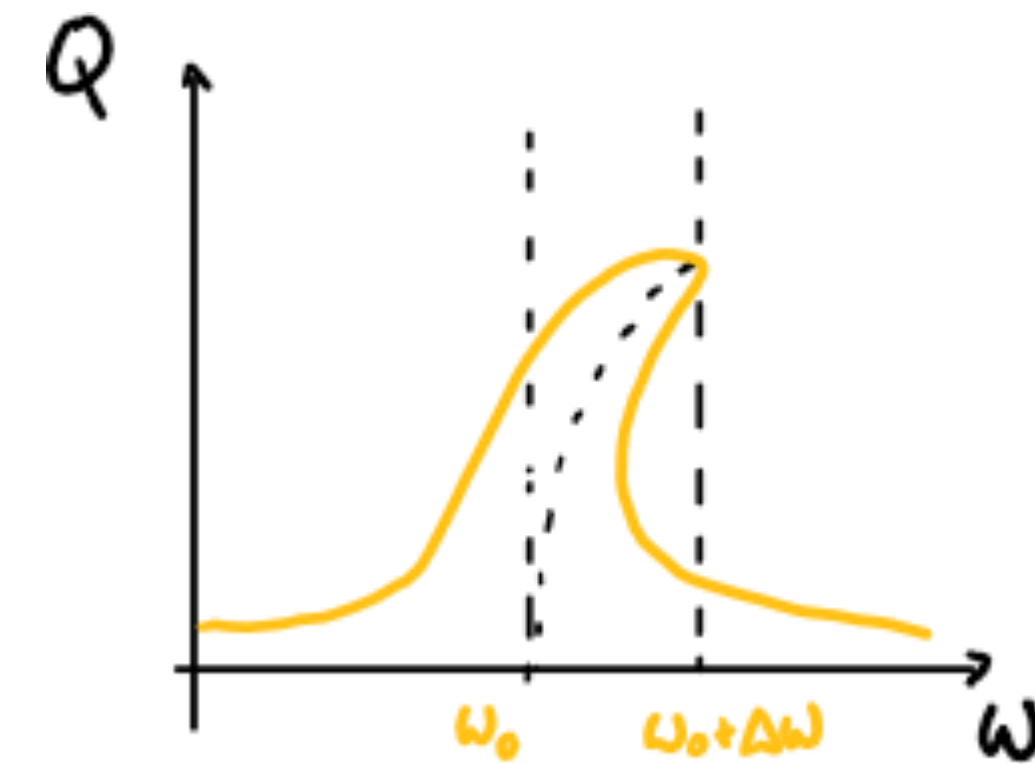
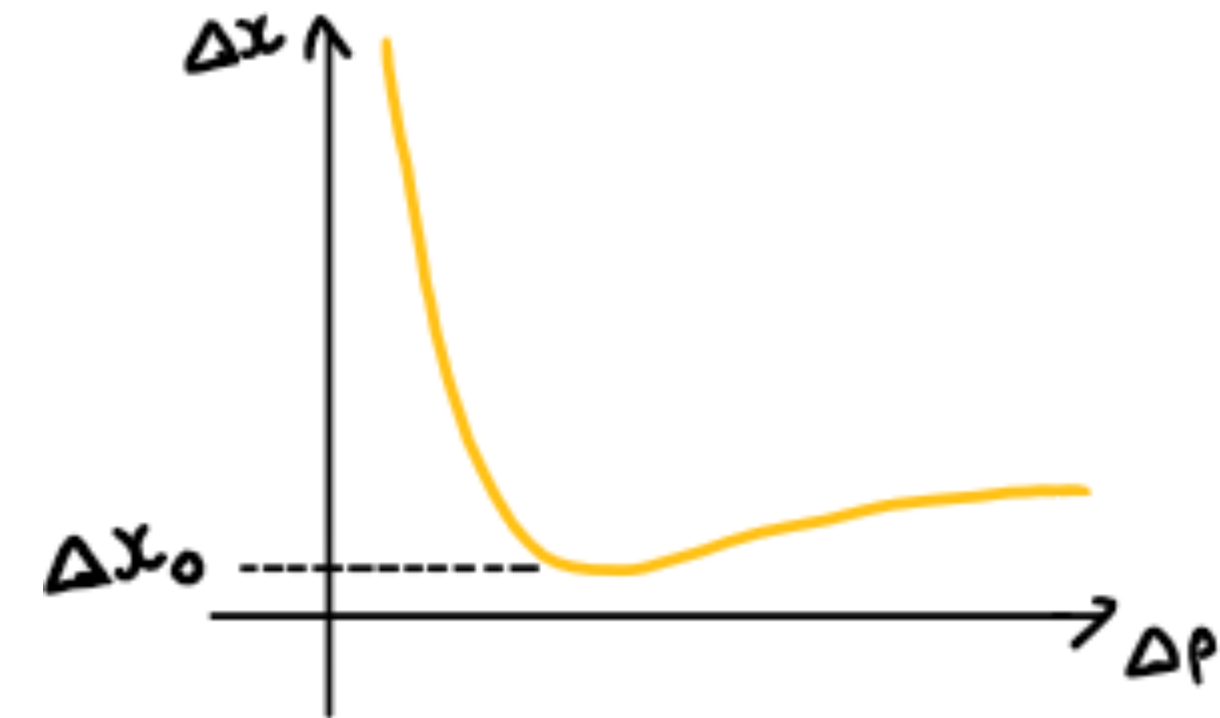
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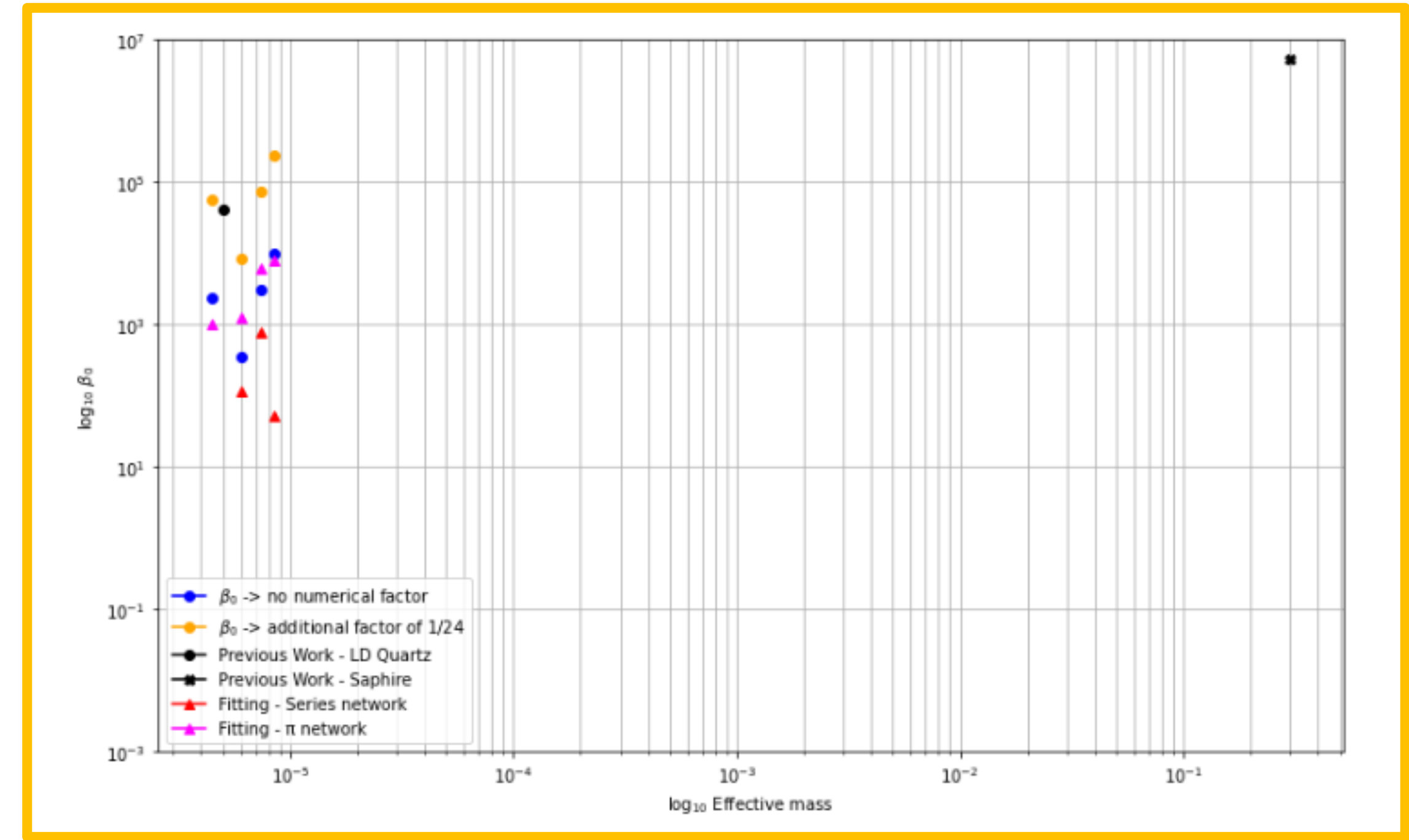
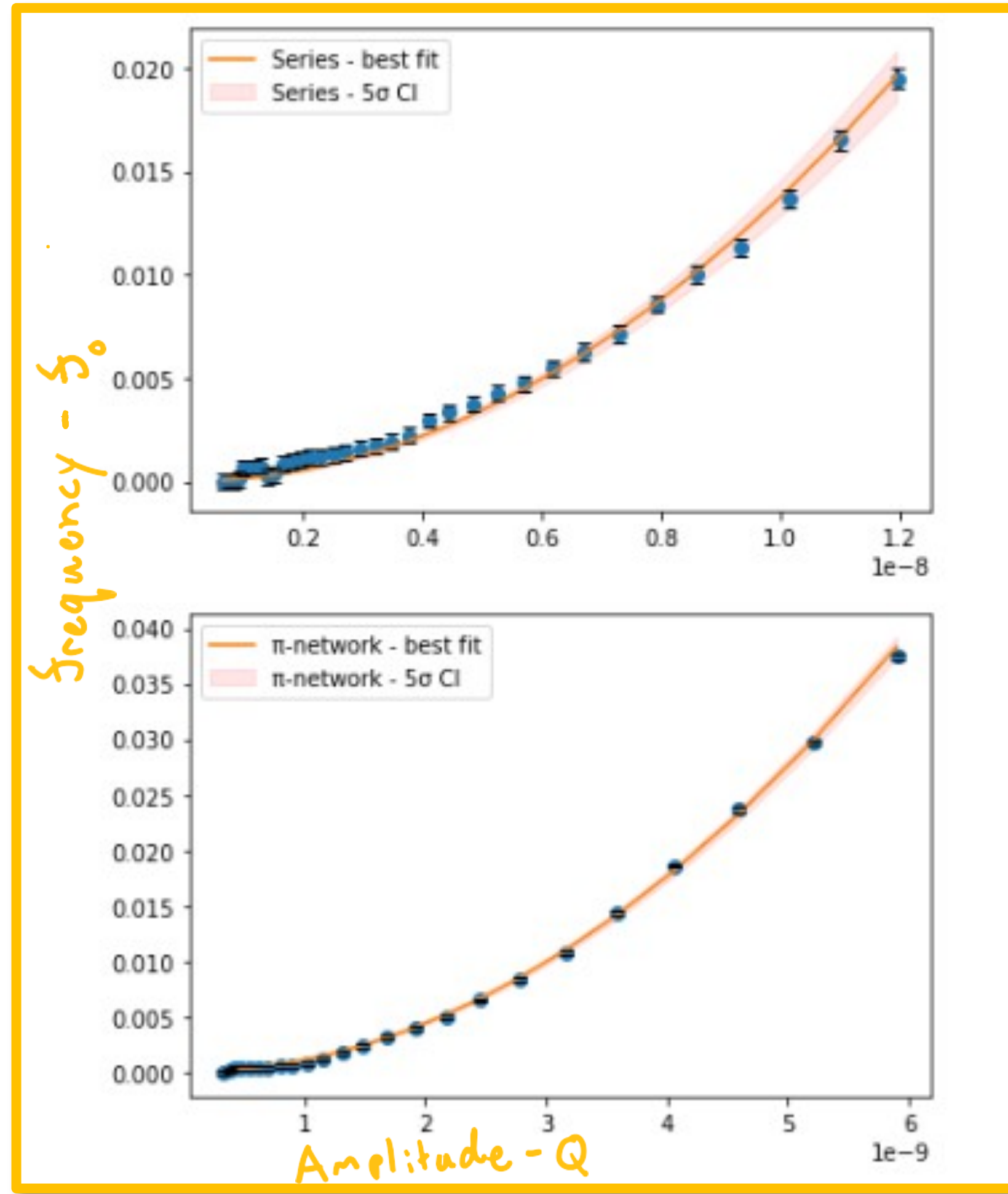
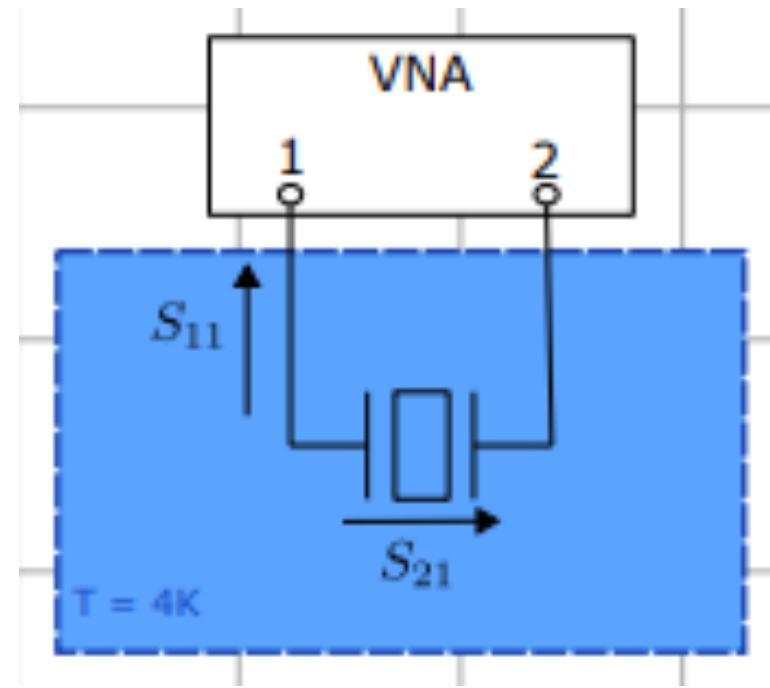
$$\tilde{\omega} = \left(1 + \frac{\beta}{2} Q_0^2 \right) \omega_0.$$

$\Delta\omega_0$



Minimum Length

Using quartz resonator



Twisted “anyon” microwave cavities

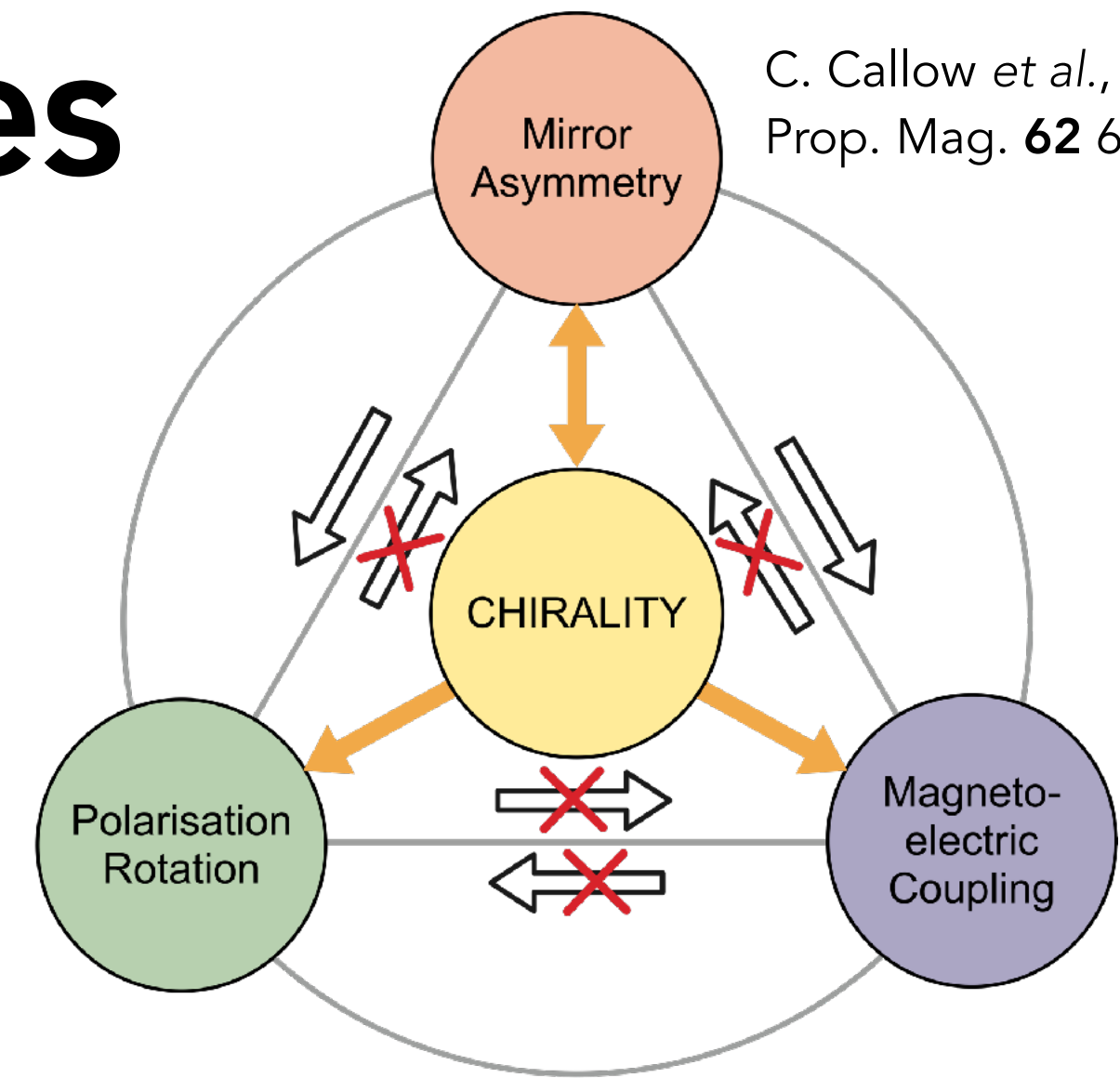
Engineered electromagnetic helicity *in vacuo*

Twisted “anyon” microwave cavities

Engineered electromagnetic helicity *in vacuo*

C. Callow *et al.*, IEEE Ant. And Prop. Mag. **62** 60 (2020)

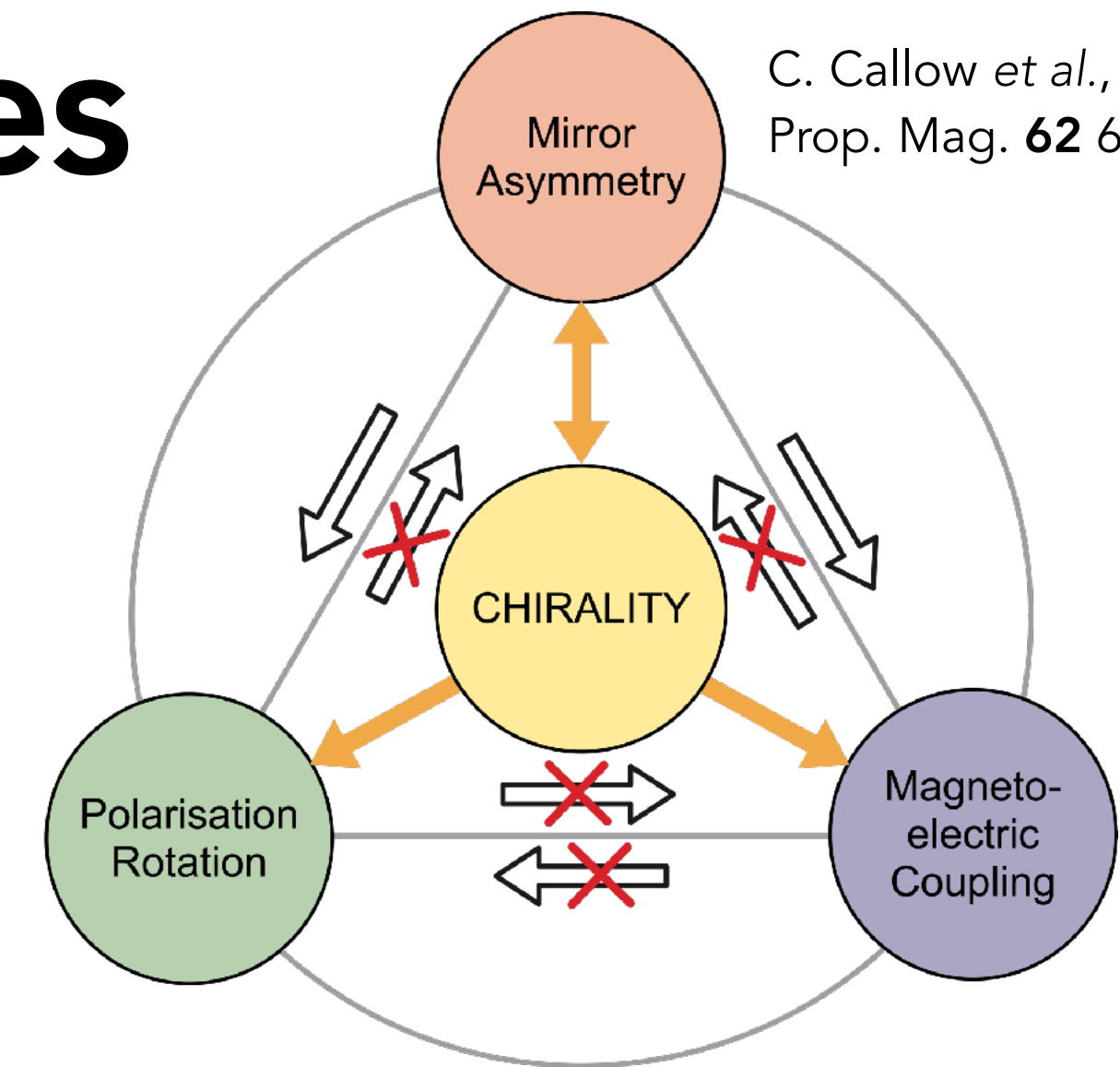
- **Chirality** is an inherent property of a particle, just like your left and right hands



Twisted “anyon” microwave cavities

Engineered electromagnetic helicity *in vacuo*

C. Callow *et al.*, IEEE Ant. And Prop. Mag. **62** 60 (2020)



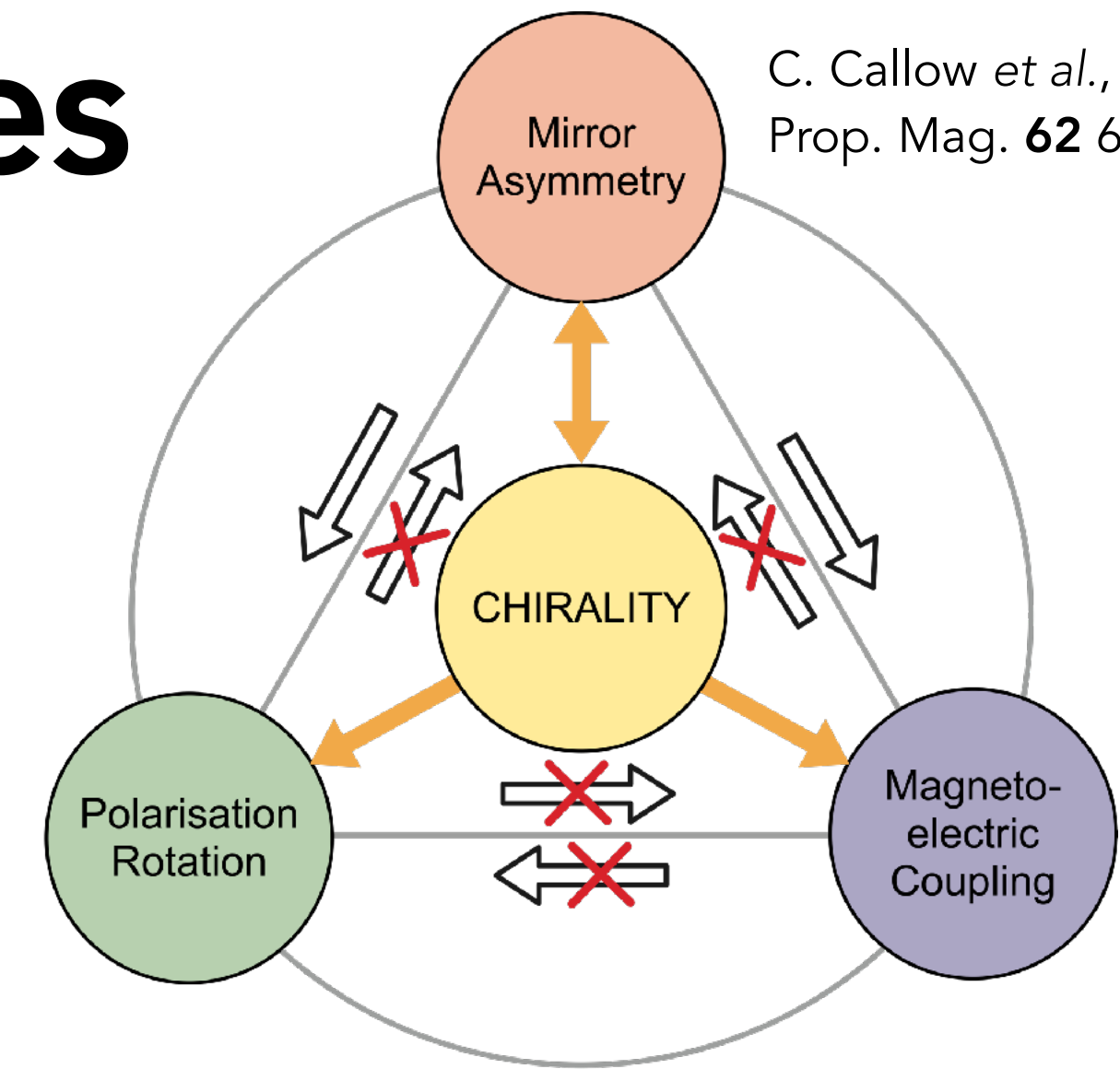
- **Chirality** is an inherent property of a particle, just like your left and right hands
- **Helicity** depends on if a particle’s spin is aligned or anti-aligned with its momentum

$$\mathcal{H}_p = \frac{2\text{Im}[\int \mathbf{B}_p(\vec{r}) \cdot \mathbf{E}_p^*(\vec{r}) d\tau]}{\sqrt{\int \mathbf{E}_p(\vec{r}) \cdot \mathbf{E}_p^* d\tau \int \mathbf{B}_p(\vec{r}) \cdot \mathbf{B}_p^*(\vec{r}) d\tau}}$$

Twisted “anyon” microwave cavities

Engineered electromagnetic helicity *in vacuo*

C. Callow *et al.*, IEEE Ant. And Prop. Mag. **62** 60 (2020)



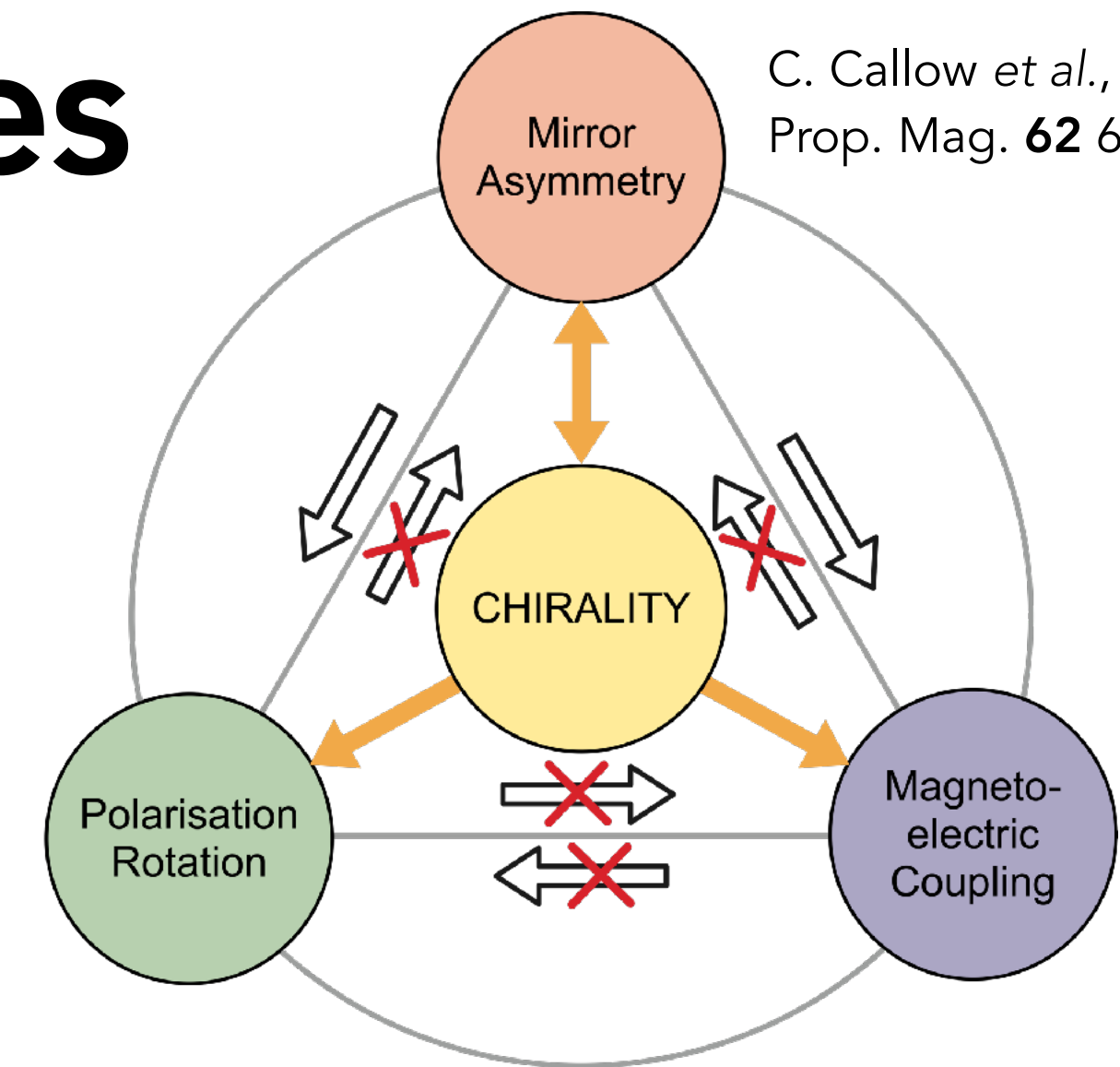
- **Chirality** is an inherent property of a particle, just like your left and right hands
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- Massless particle -> unable to alter helicity by changing frames -> If particle is chiral right-handed it will have right handed helicity

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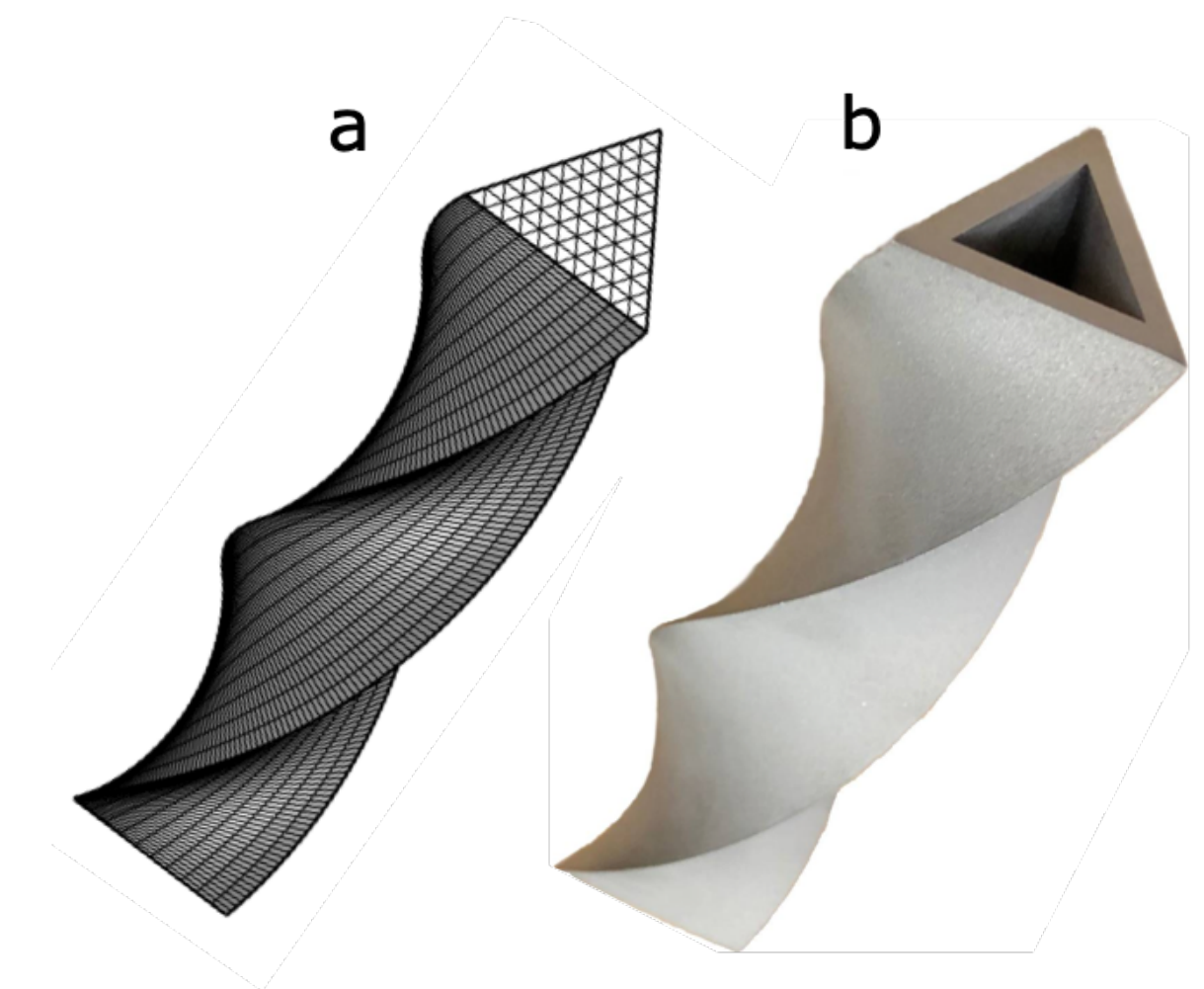
Twisted “anyon” microwave cavities

Engineered electromagnetic helicity *in vacuo*

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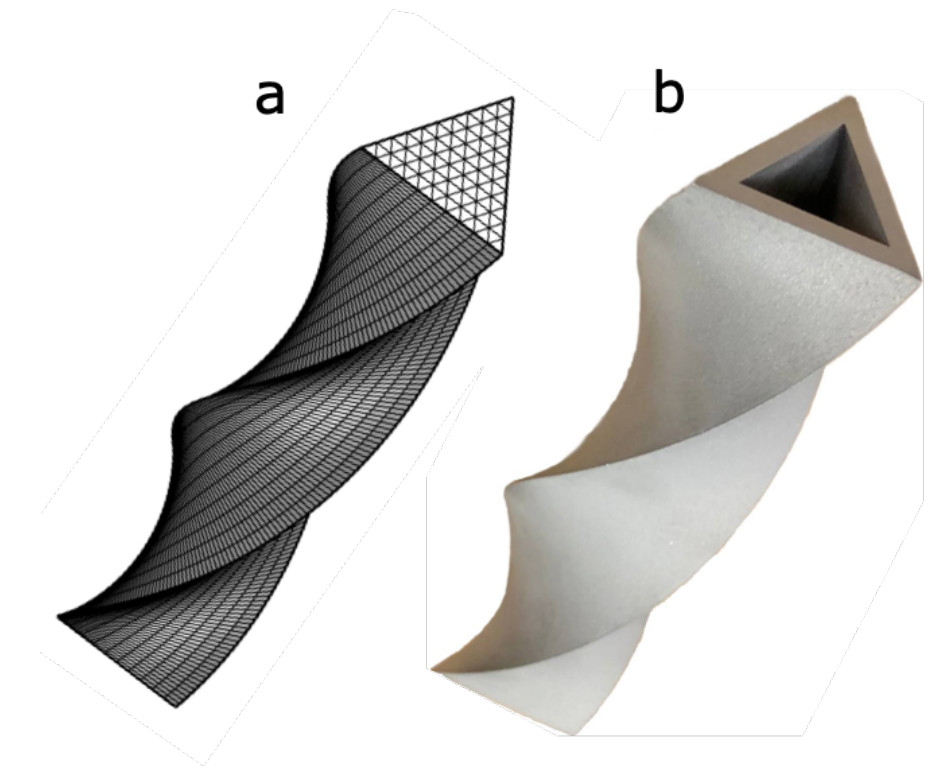
Twisted “anyon” microwave cavities

So why a triangle cross-section?

Twisted “anyon” microwave cavities

So why a triangle cross-section?

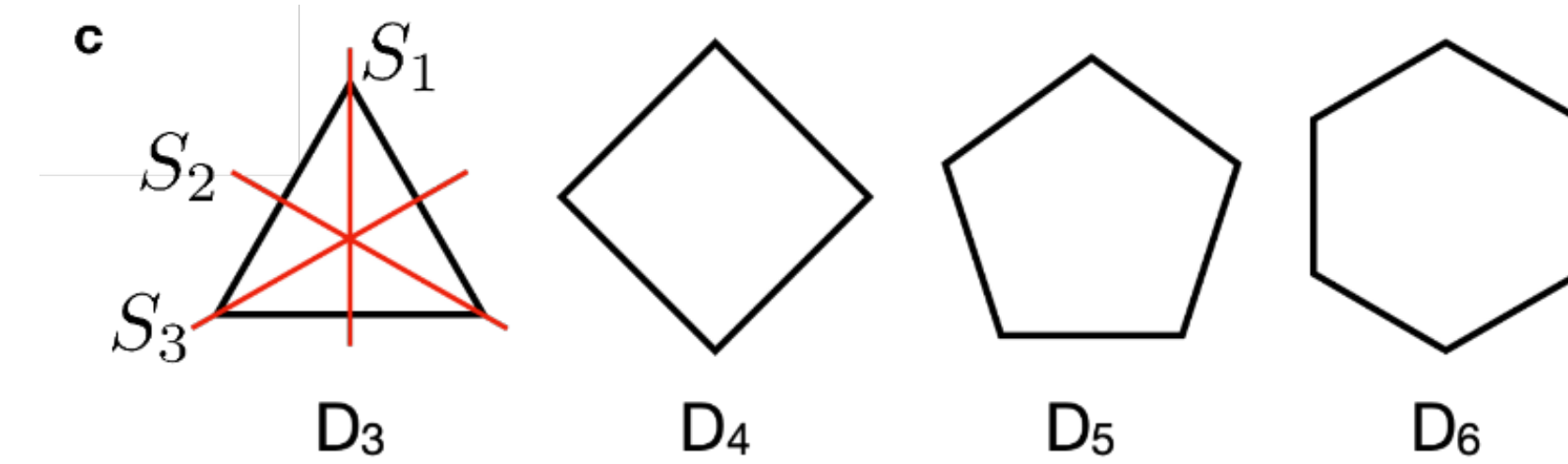
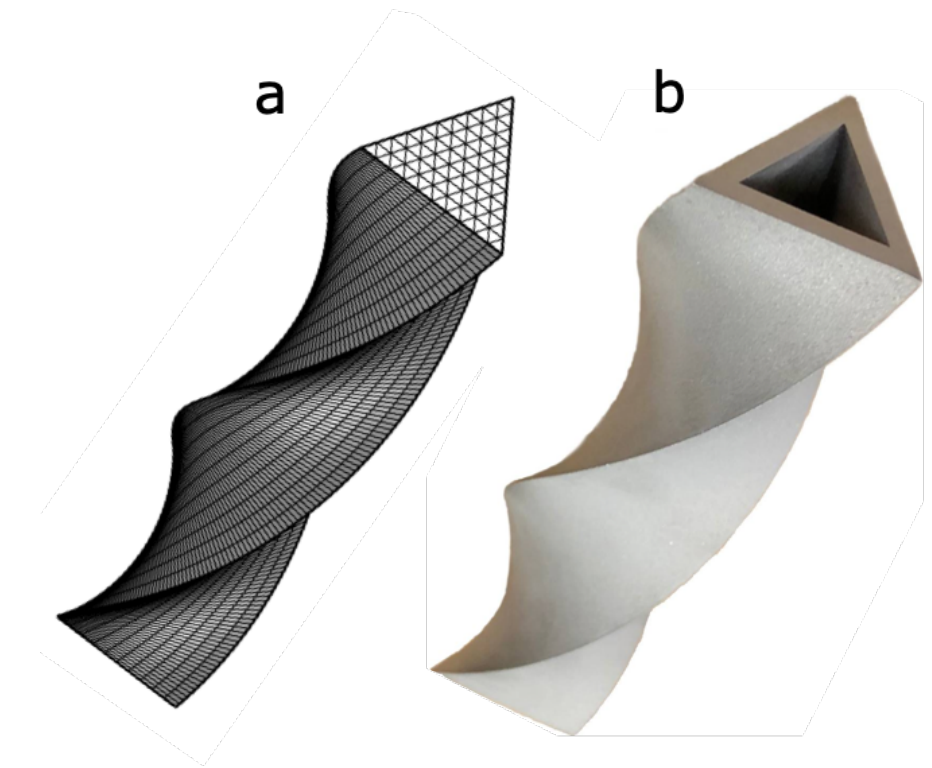
- Cavity **mirror asymmetry** from the clockwise or counter-clockwise twist



Twisted "anyon" microwave cavities

So why a triangle cross-section?

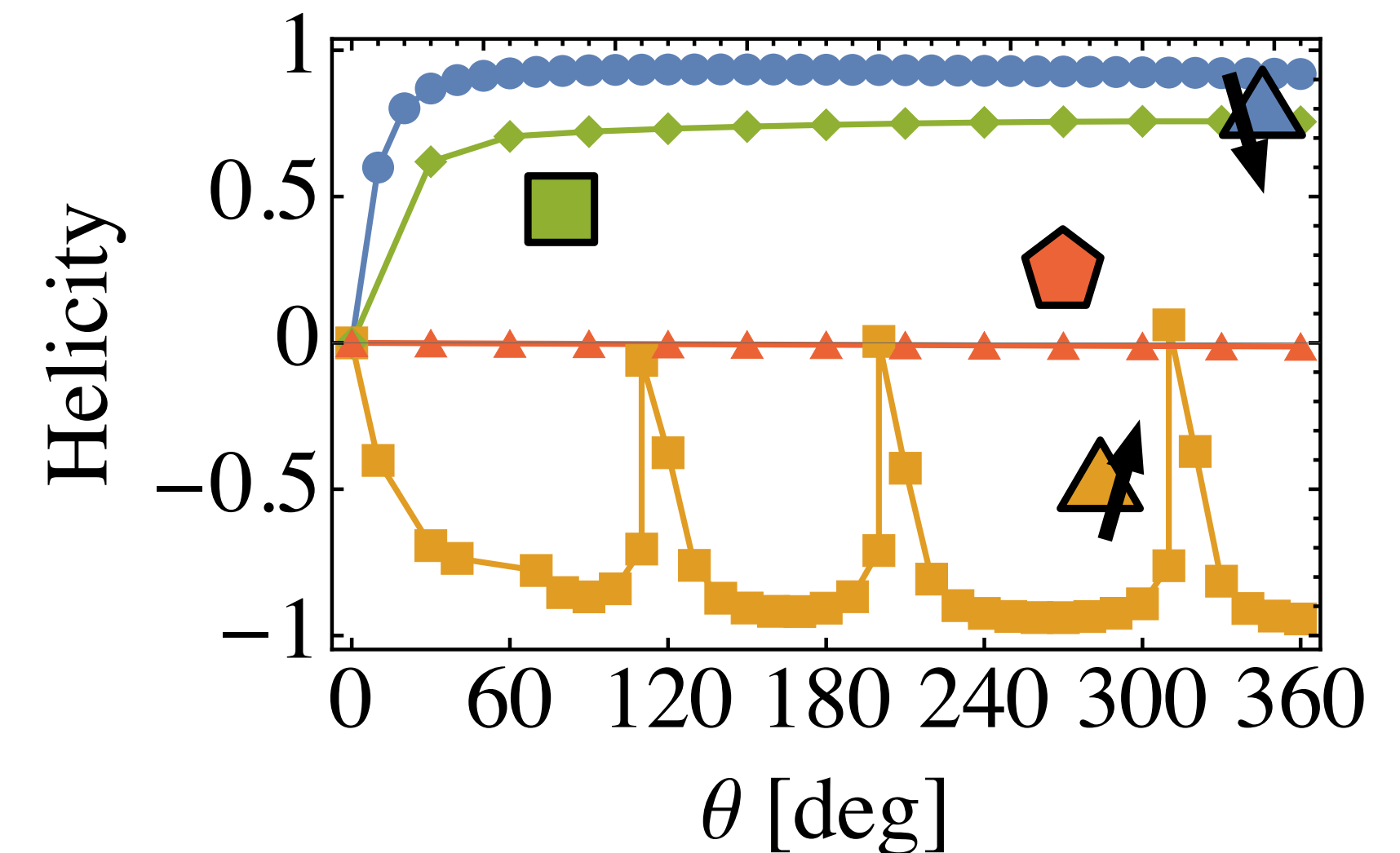
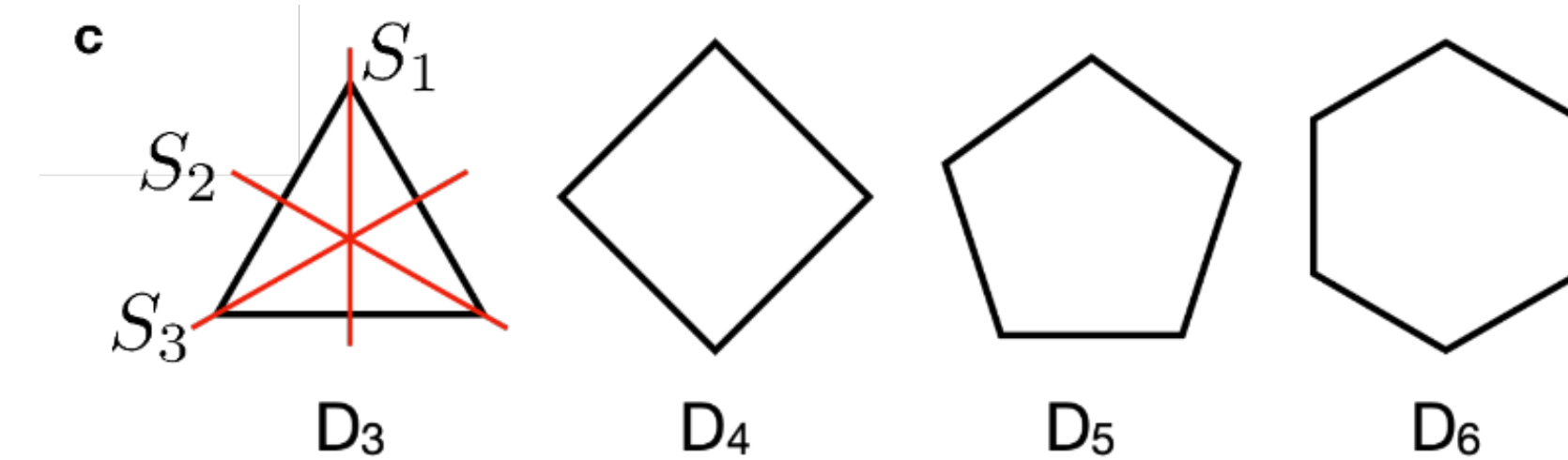
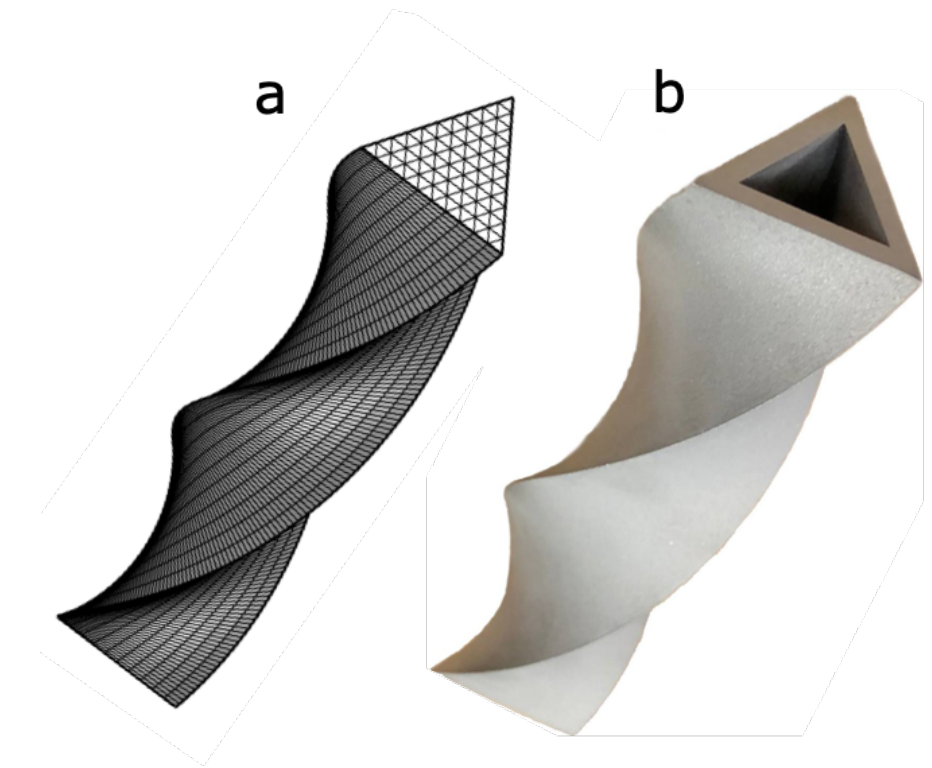
- Cavity **mirror asymmetry** from the clockwise or counter-clockwise twist
- If cross-section a regular polygon, as we increase the number of planes of symmetry, we further approximate a circle -> which shows no mirror asymmetry



Twisted "anyon" microwave cavities

So why a triangle cross-section?

- Cavity **mirror asymmetry** from the clockwise or counter-clockwise twist
- If cross-section a regular polygon, as we increase the number of planes of symmetry, we further approximate a circle -> which shows no mirror asymmetry
- The least amount of symmetry lines -> greatest helicity



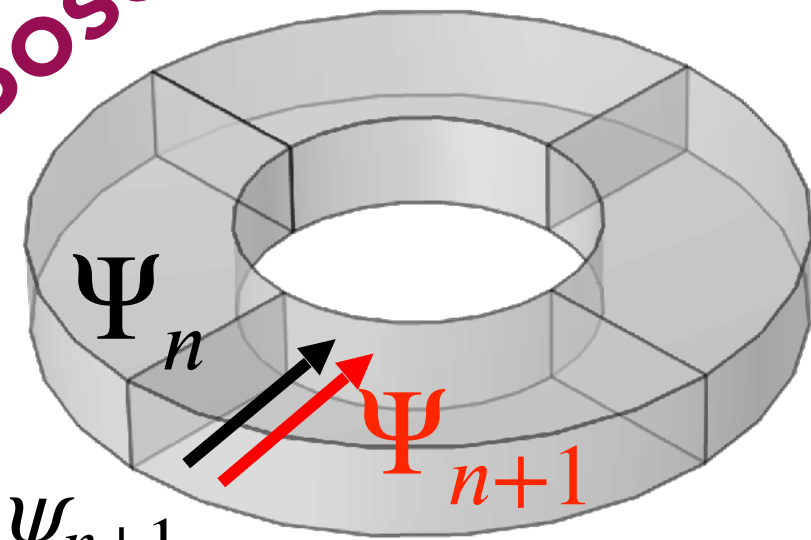
Twisted “anyon” microwave cavities

Why is it called an “anyon” cavity?

Twisted "anyon" microwave cavities

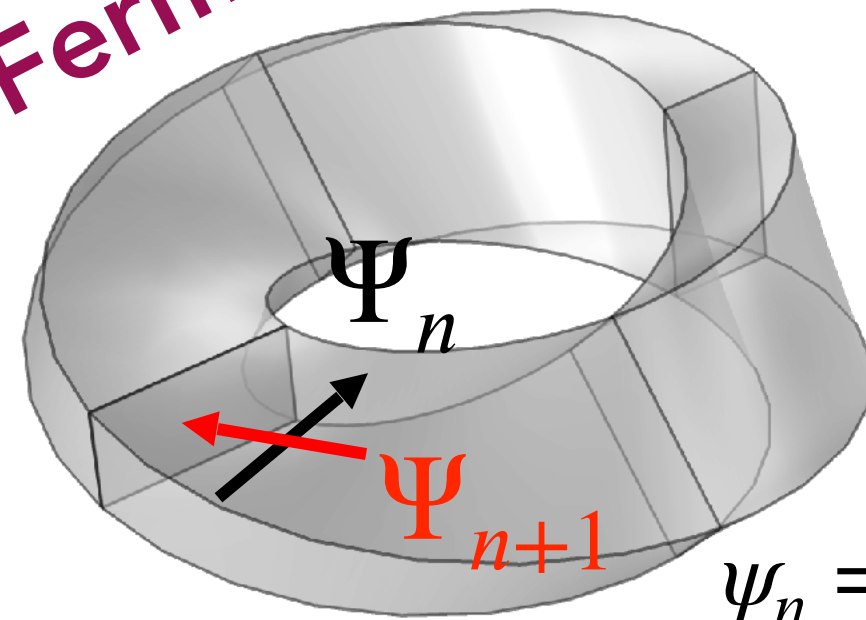
Why is it called an "anyon" cavity?

"Boson"



$$\begin{aligned}\psi_n &= \psi_{n+1} \\ \psi_n &= \psi_{n+N} \\ \theta &= 0\end{aligned}$$

"Fermion"

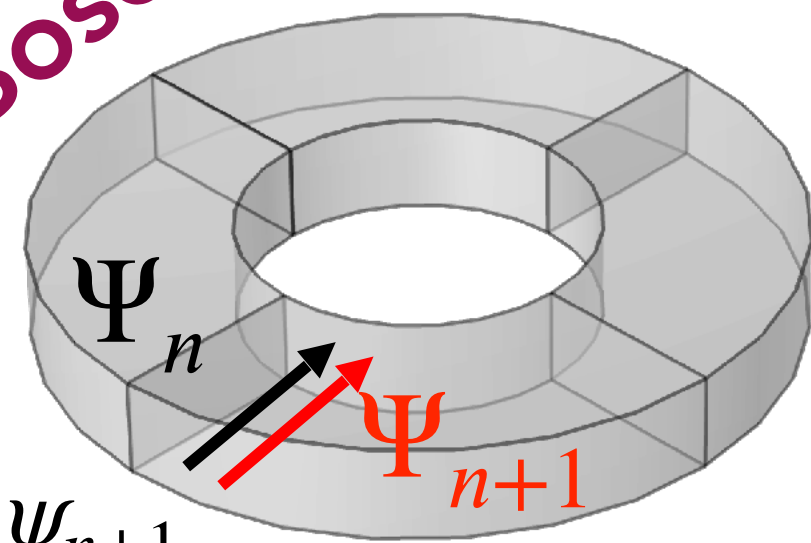


$$\begin{aligned}\psi_n &= -\psi_{n+1} \\ \psi_n &= \psi_{n+2N} \\ \theta &= \pm \pi\end{aligned}$$

Twisted "anyon" microwave cavities

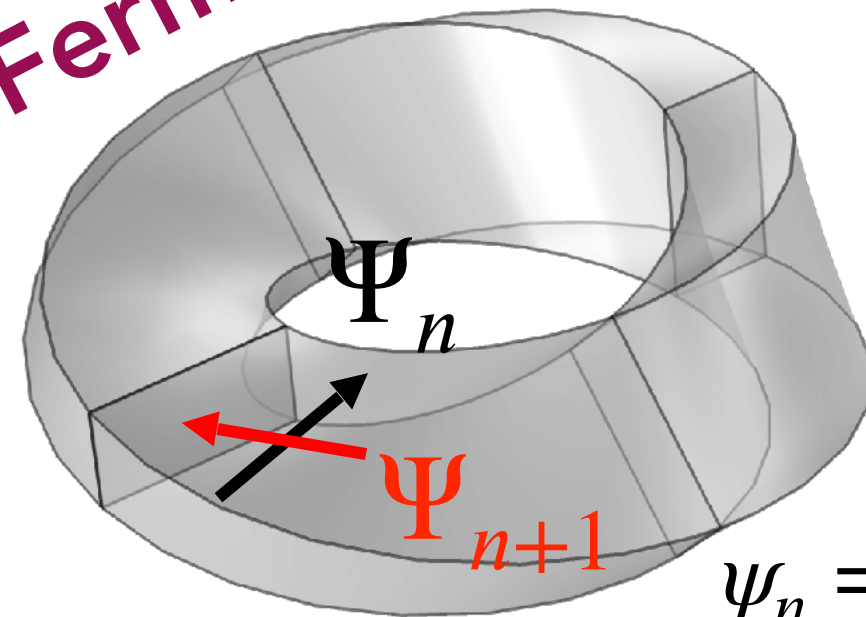
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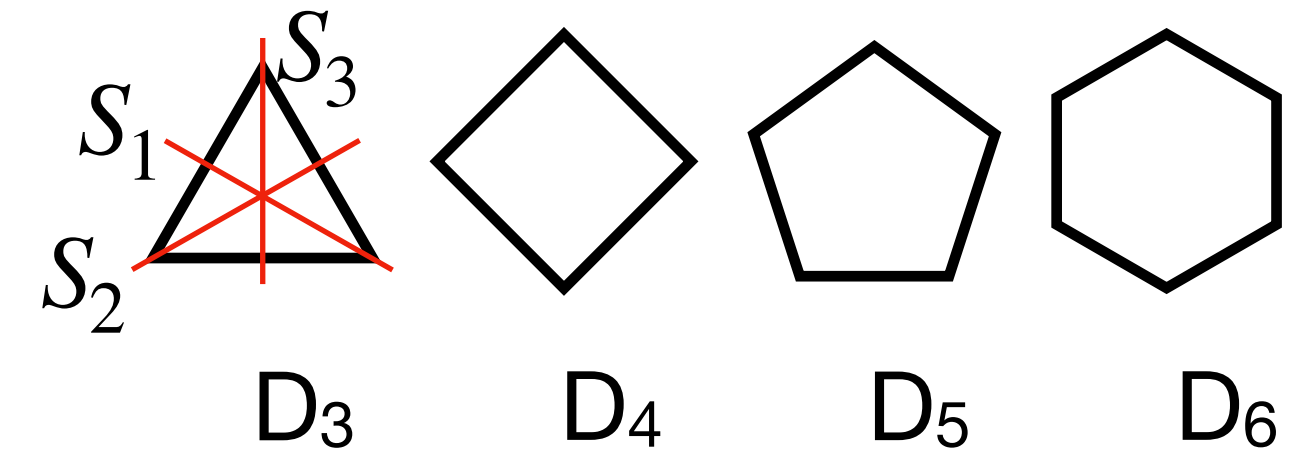
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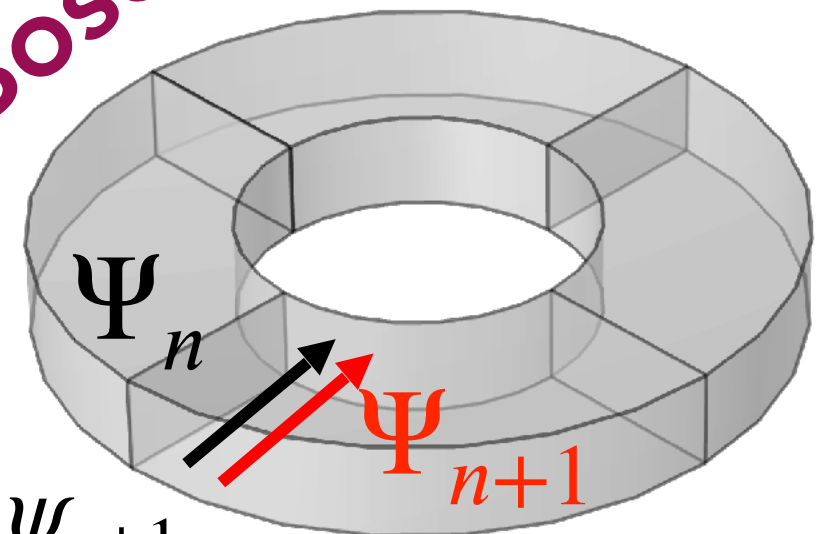
Dihedral group of regular convex polygons: D_p



Twisted "anyon" microwave cavities

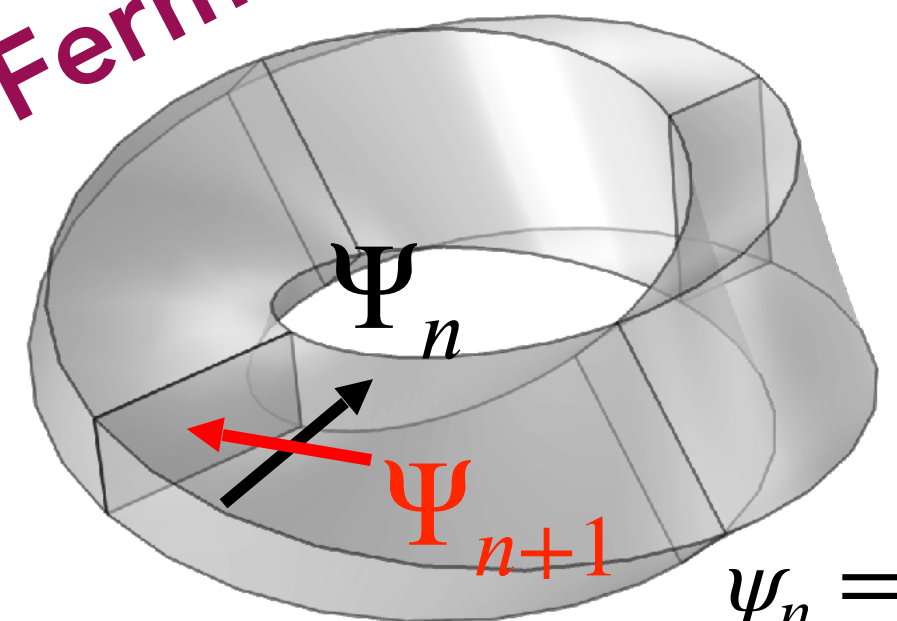
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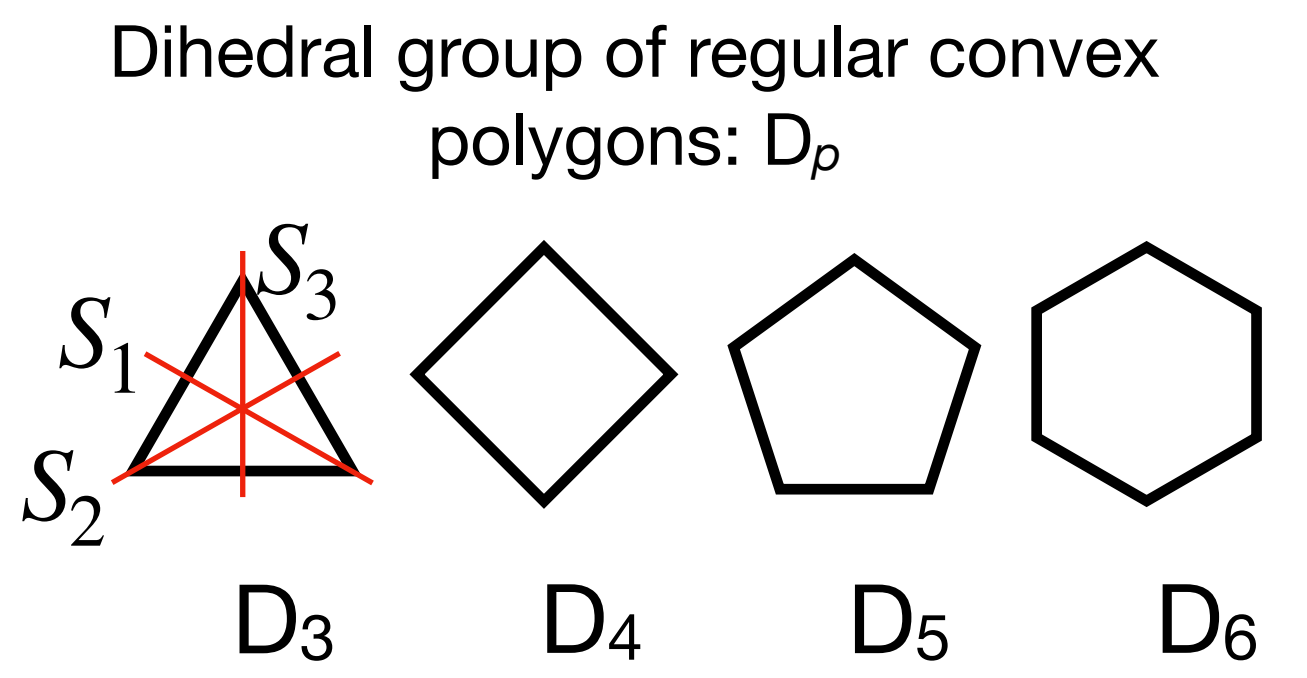
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2p symmetries:
p rotational + p reflection



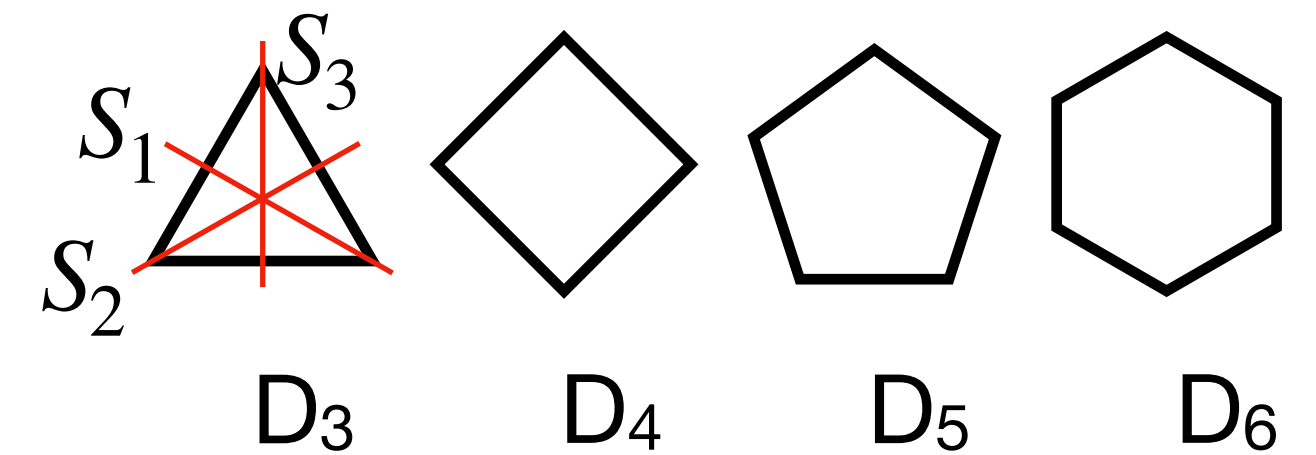
Rotation by $2\pi/p$ preserves the object

Twisted "anyon" microwave cavities

Why is it called an "anyon" cavity?

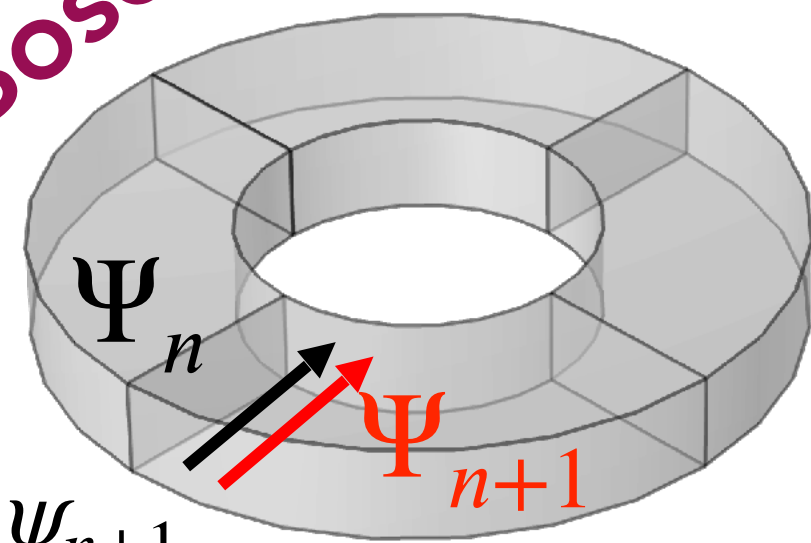
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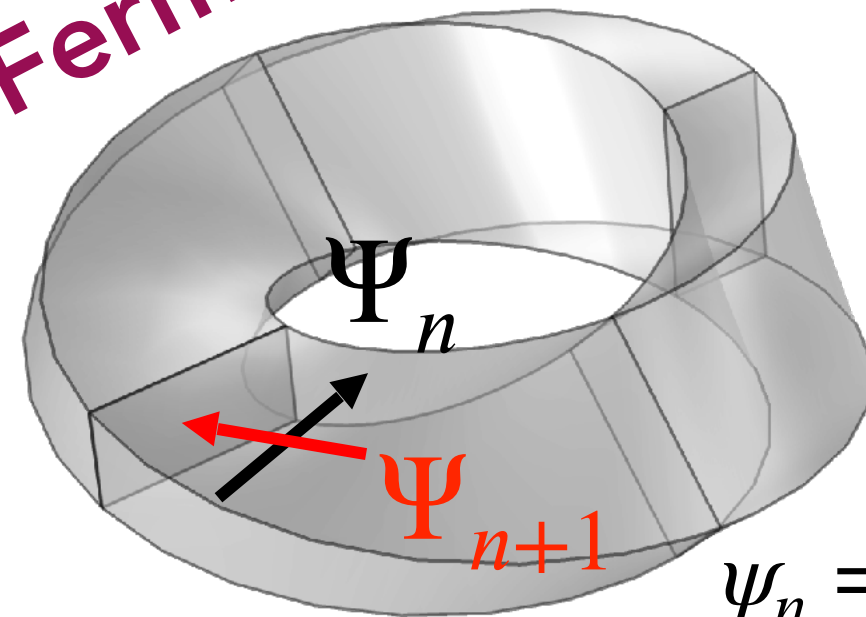
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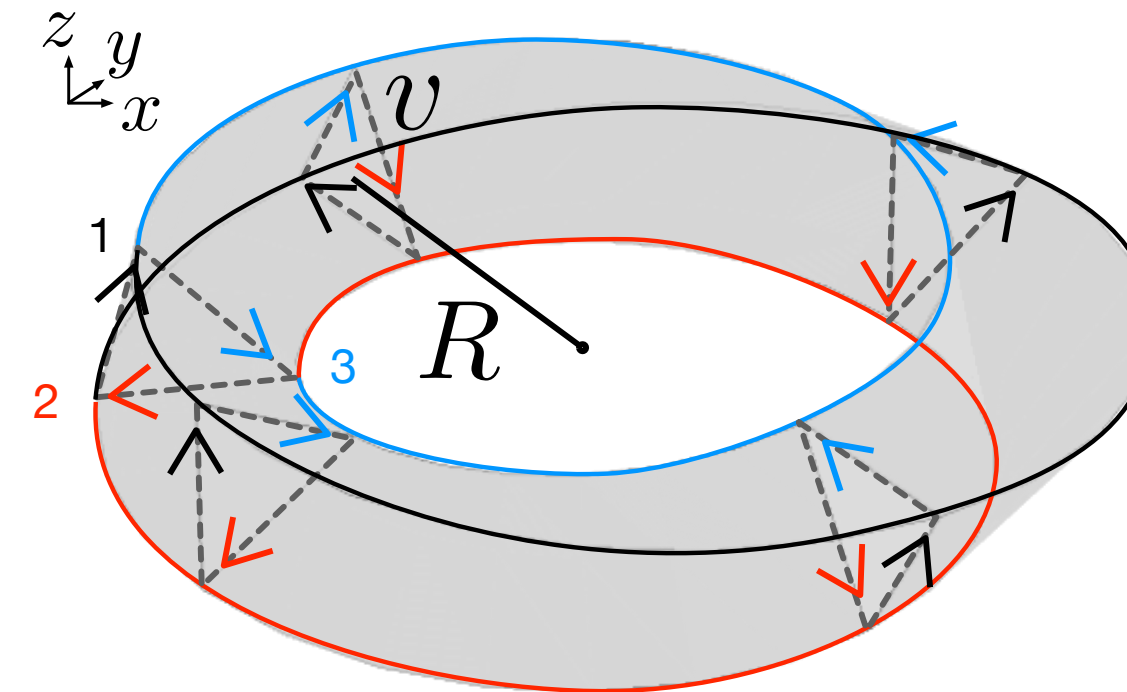


$$\begin{aligned} \psi_n &= \psi_{n+1} \\ \psi_n &= \psi_{n+N} \\ \theta &= 0 \end{aligned}$$

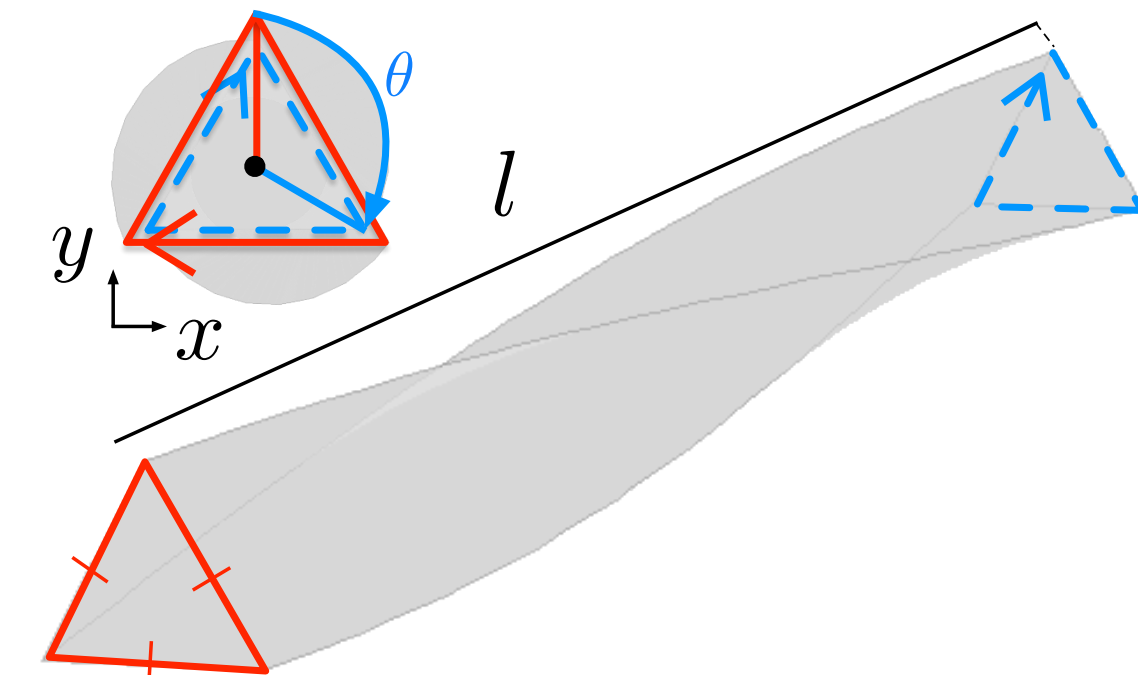
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$$\begin{aligned} \psi_n &= -\psi_{n+1} \\ \psi_n &= \psi_{n+2N} \\ \theta &= \pm \pi \end{aligned}$$



$$\begin{aligned} \psi_n &= e^{i\theta} \psi_{n+1} \\ \theta &= (2\pi/p)Z \\ Z &\in \pm \mathbb{Z} \end{aligned}$$

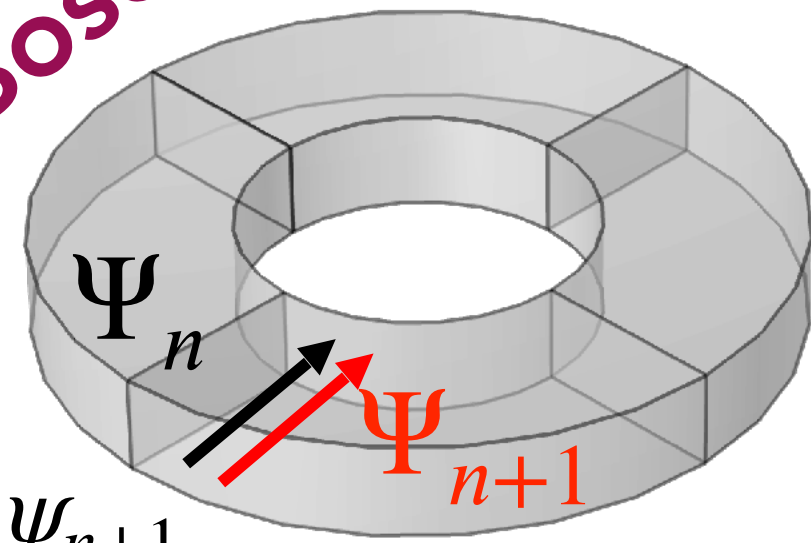


$$\begin{aligned} \psi_n &= e^{i\theta} \psi_{n+1} \\ \theta &\in \mathbb{R} \end{aligned}$$

Twisted "anyon" microwave cavities

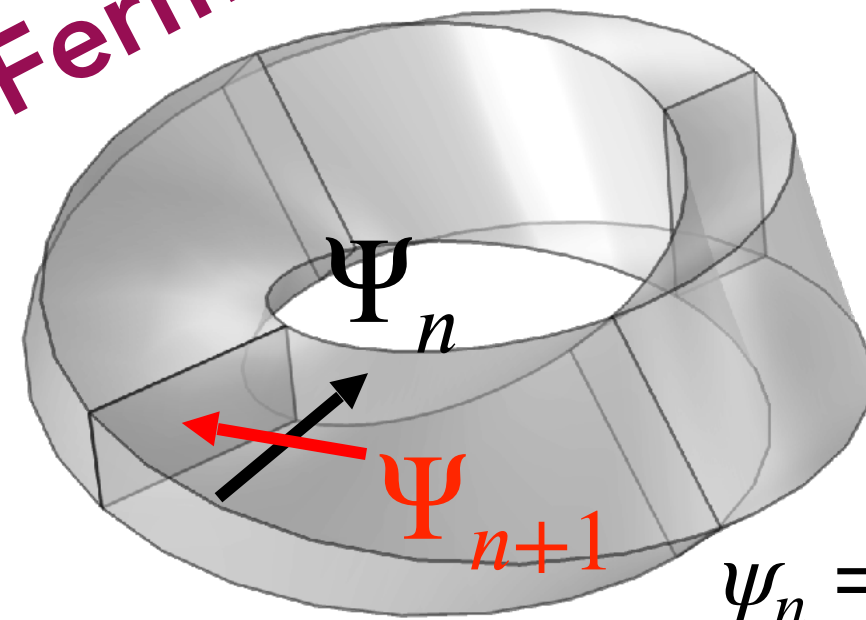
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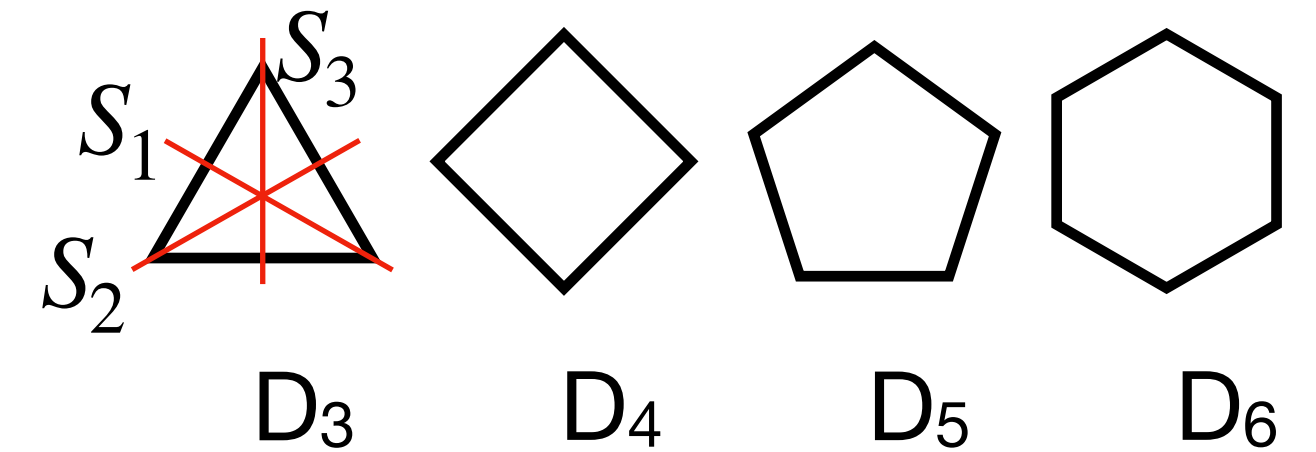
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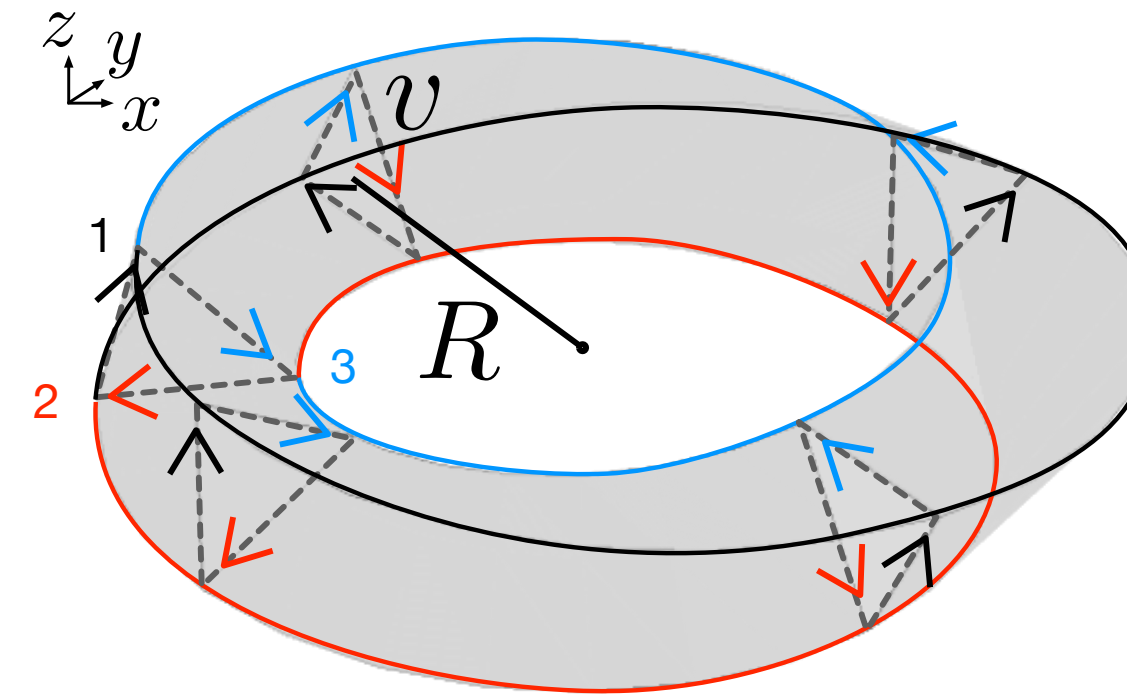
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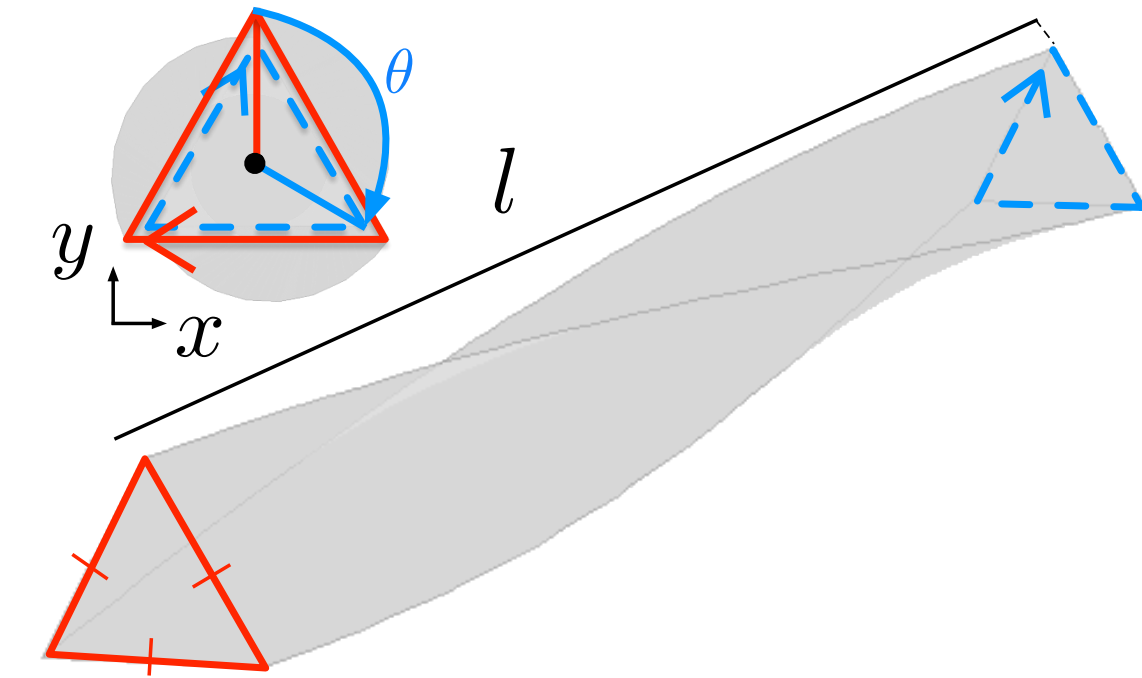
Dihedral group of regular convex polygons: D_p



Rotation by $2\pi/p$ preserves the object



$$\begin{aligned} \psi_n &= e^{i\theta} \psi_{n+1} \\ \theta &= (2\pi/p)Z \\ Z &\in \pm \mathbb{Z} \end{aligned}$$



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PHYSICAL REVIEW LETTERS

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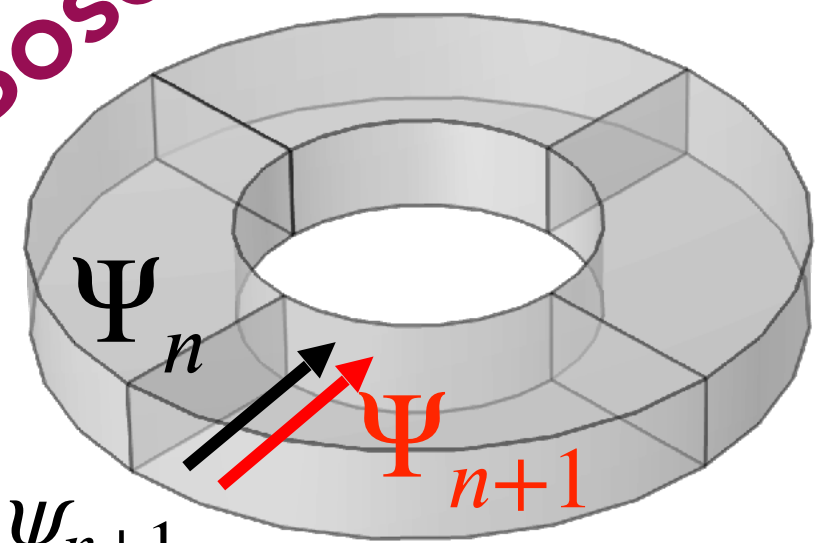
Classical Möbius-Ring Resonators Exhibit Fermion-Boson Rotational Symmetry

Douglas J. Ballon and Henning U. Voss
Phys. Rev. Lett. **101**, 247701 – Published 9 December 2008

Twisted "anyon" microwave cavities

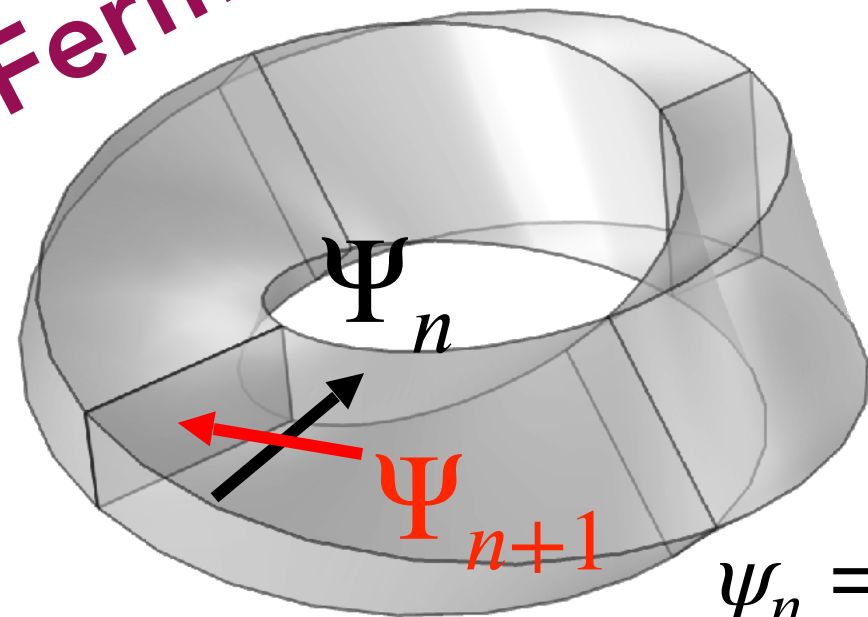
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$$\begin{aligned} \psi_n &= \psi_{n+1} \\ \psi_n &= \psi_{n+N} \\ \theta &= 0 \end{aligned}$$

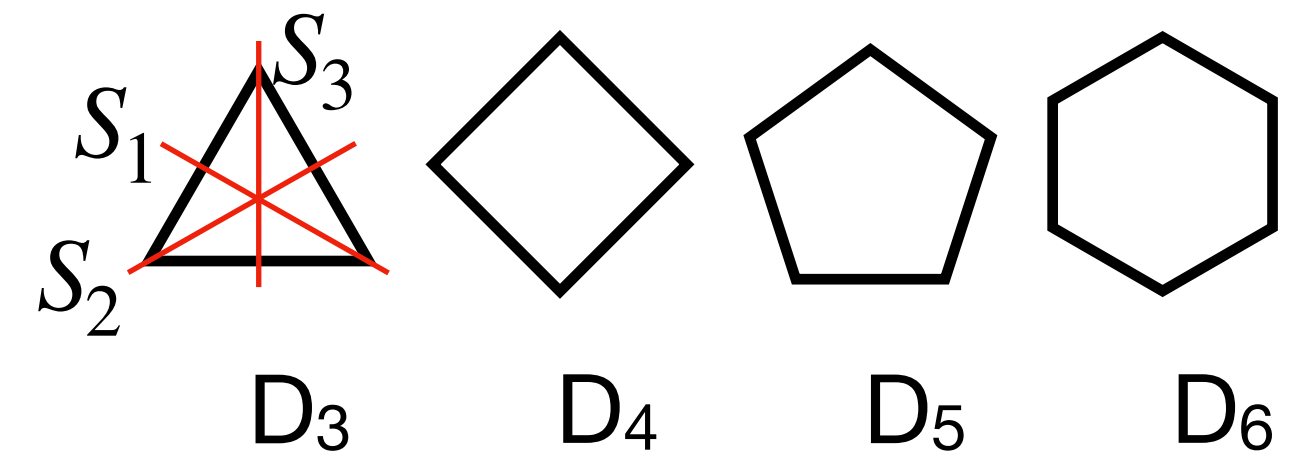
"Fermion"



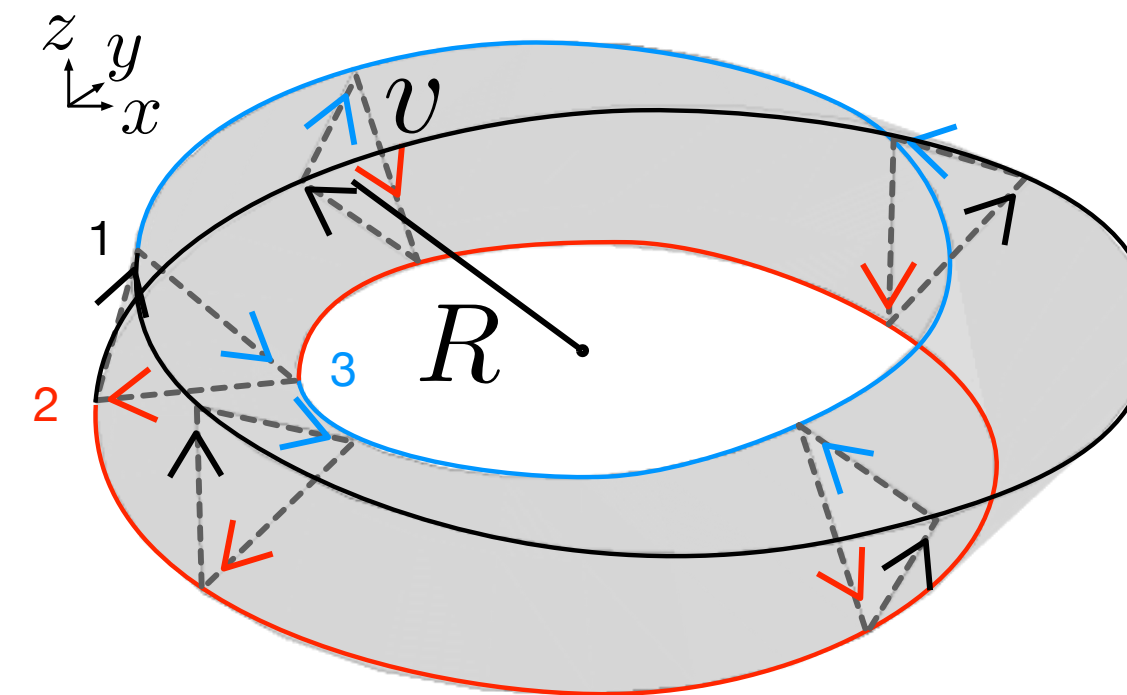
$$\begin{aligned} \psi_n &= -\psi_{n+1} \\ \psi_n &= \psi_{n+2N} \\ \theta &= \pm \pi \end{aligned}$$

$2p$ symmetries:
 p rotational + p reflection

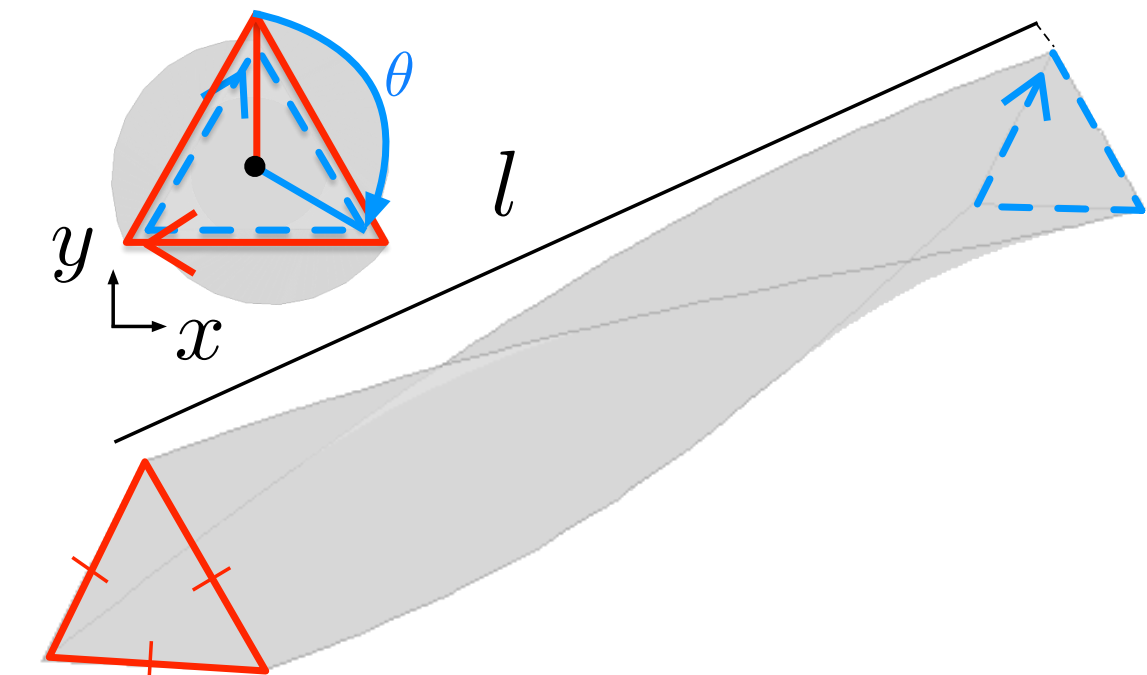
Dihedral group of regular convex polygons: D_p



Rotation by $2\pi/p$ preserves the object



$$\begin{aligned} \psi_n &= e^{i\theta} \psi_{n+1} \\ \theta &= (2\pi/p)Z \\ Z &\in \pm \mathbb{Z} \end{aligned}$$



$$\begin{aligned} \psi_n &= e^{i\theta} \psi_{n+1} \\ \theta &\in \mathbb{R} \end{aligned}$$

"Anyon"

PHYSICAL REVIEW LETTERS

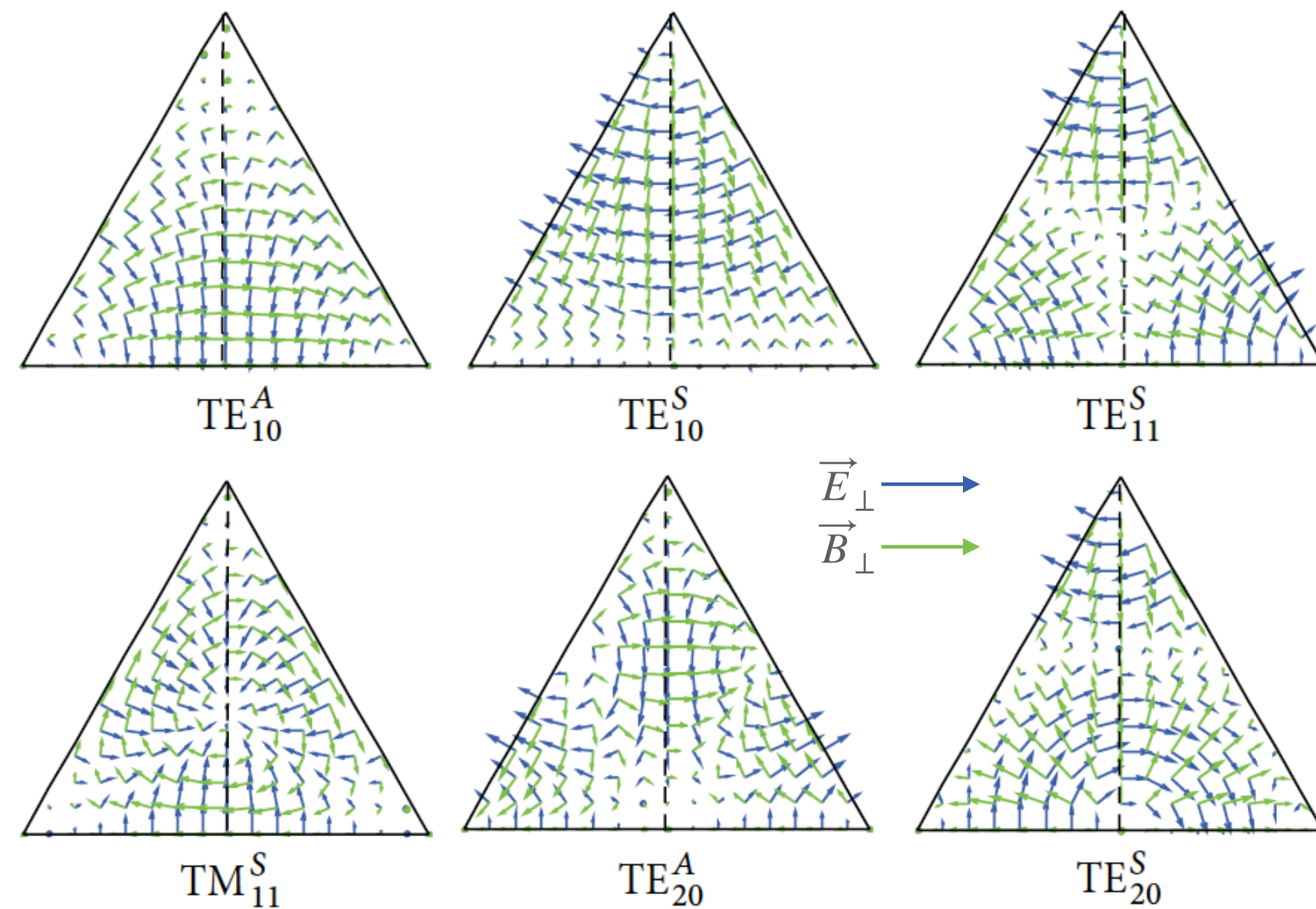
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Classical Möbius-Ring Resonators Exhibit Fermion-Boson Rotational Symmetry

Douglas J. Ballon and Henning U. Voss
Phys. Rev. Lett. **101**, 247701 – Published 9 December 2008

Twisted "anyon" microwave cavities

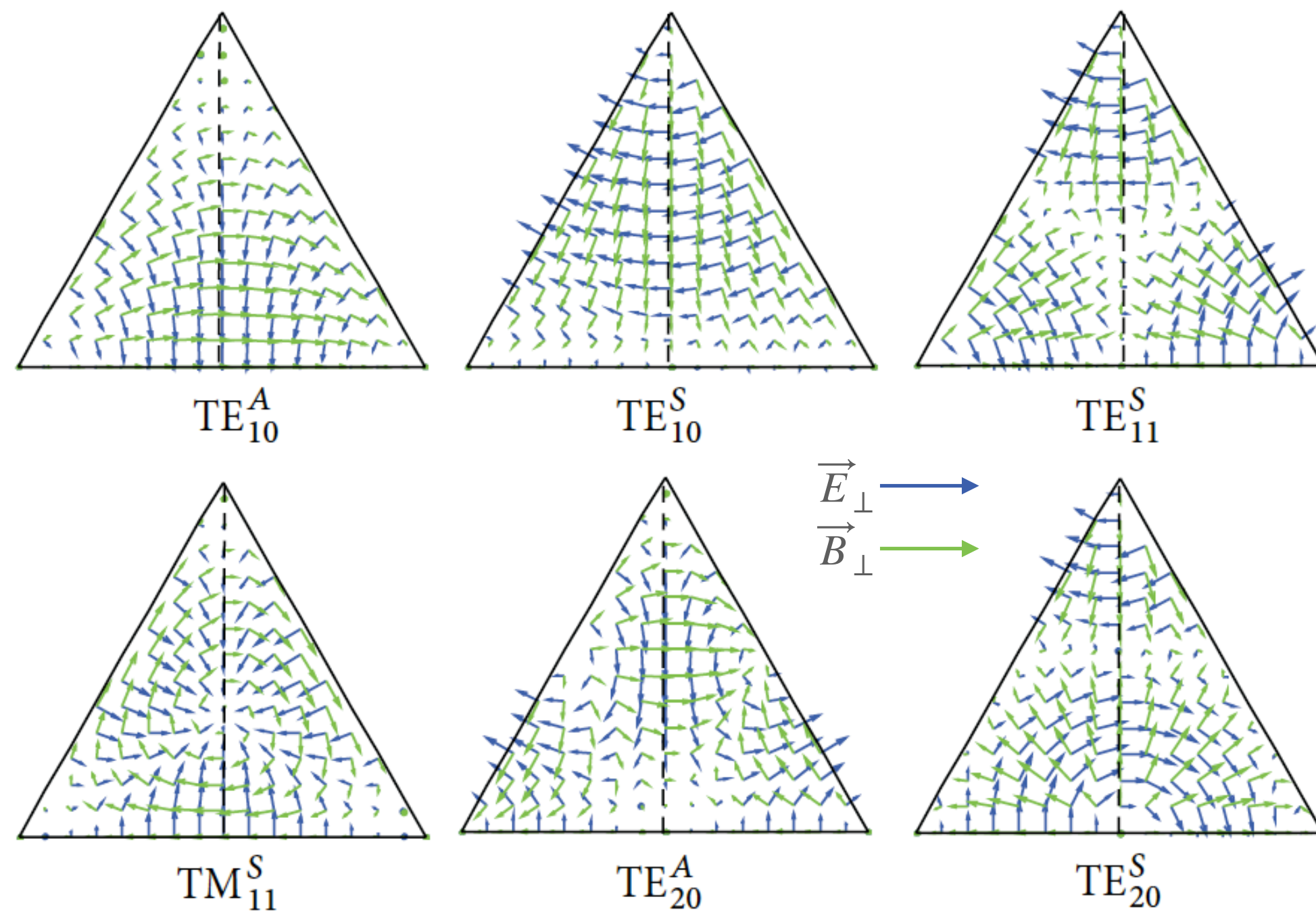
How do these geometries generate helicity?



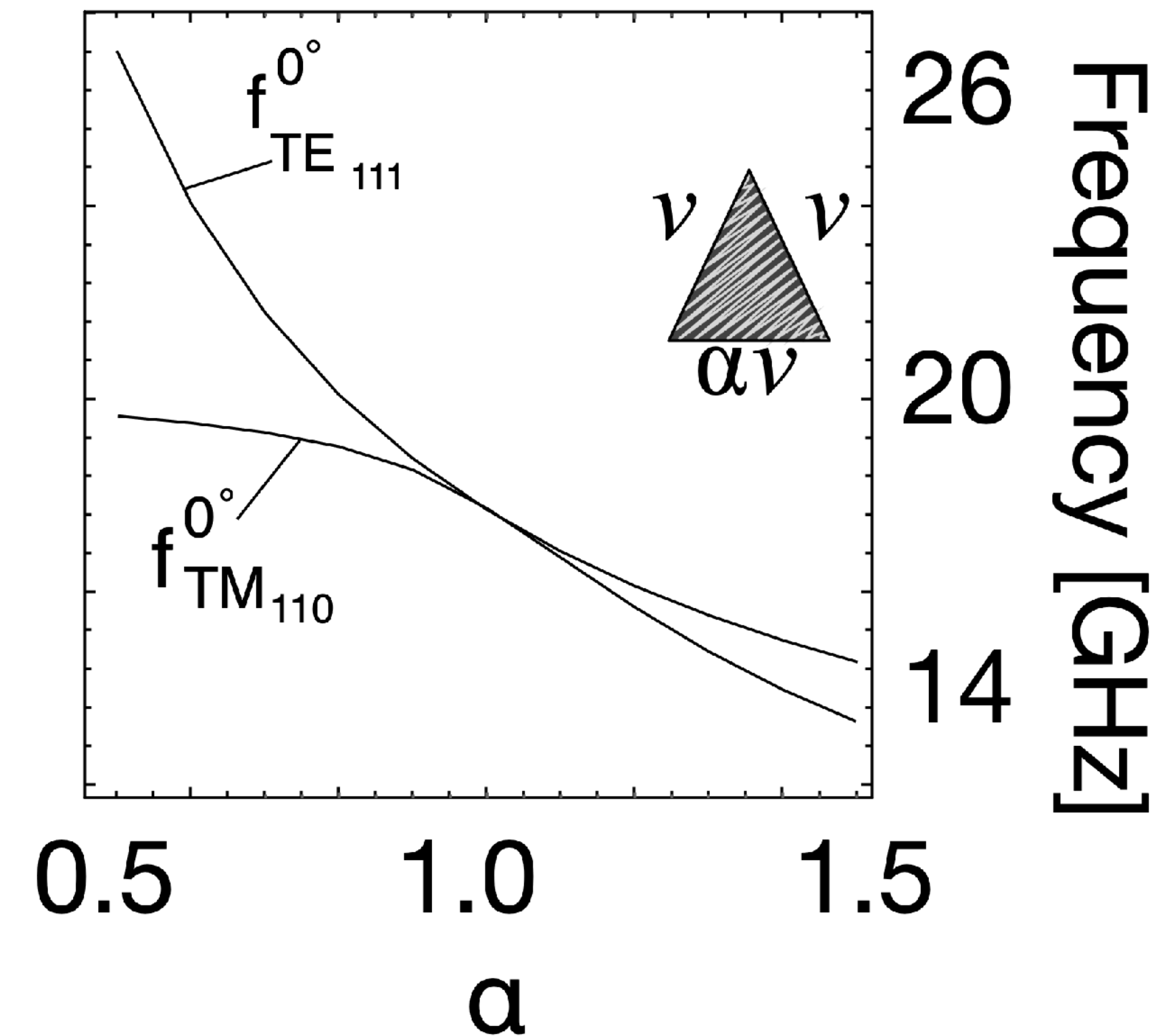
A. Moran-Lopez et al. Adv. Math. Phys., 2016, 2974675 (2016)

Twisted "anyon" microwave cavities

How do these geometries generate helicity?

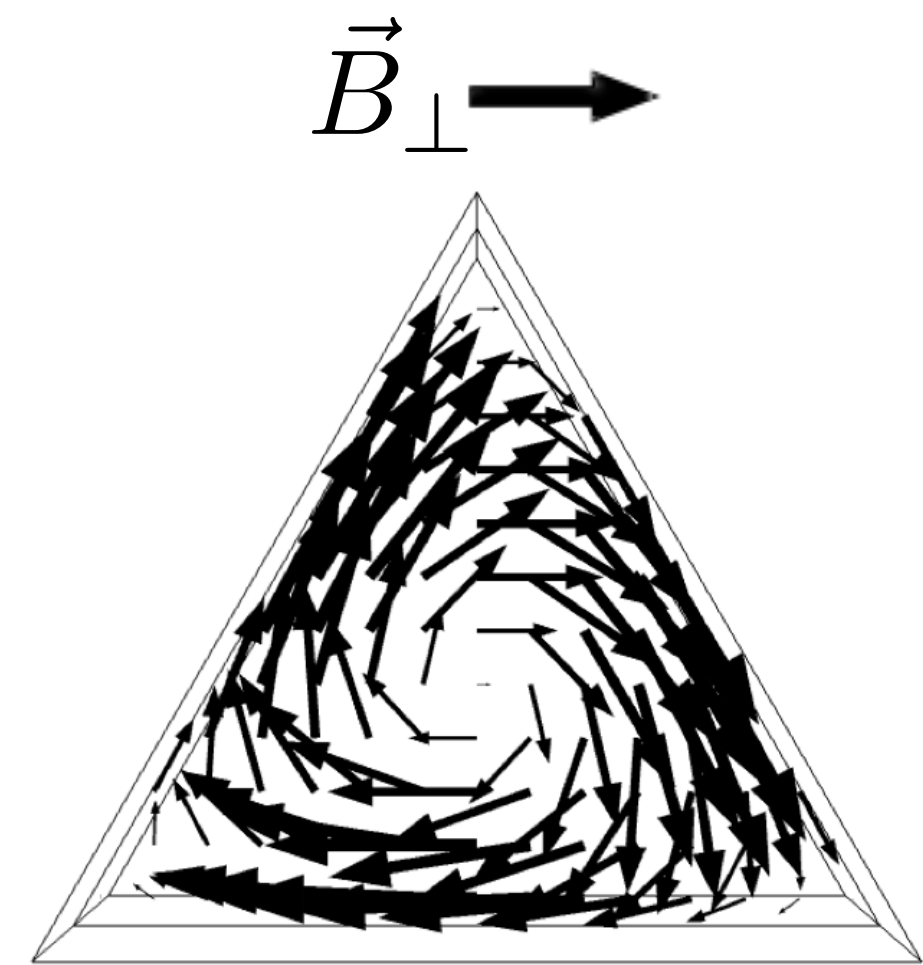


A. Moran-Lopez et al. Adv. Math. Phys., 2016, 2974675 (2016)

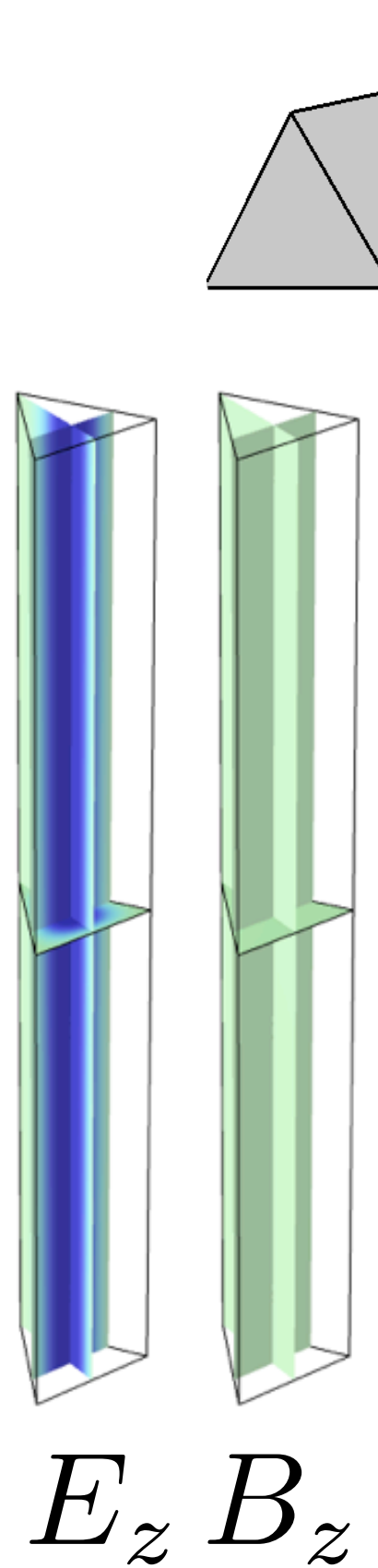


Twisted "anyon" microwave cavities

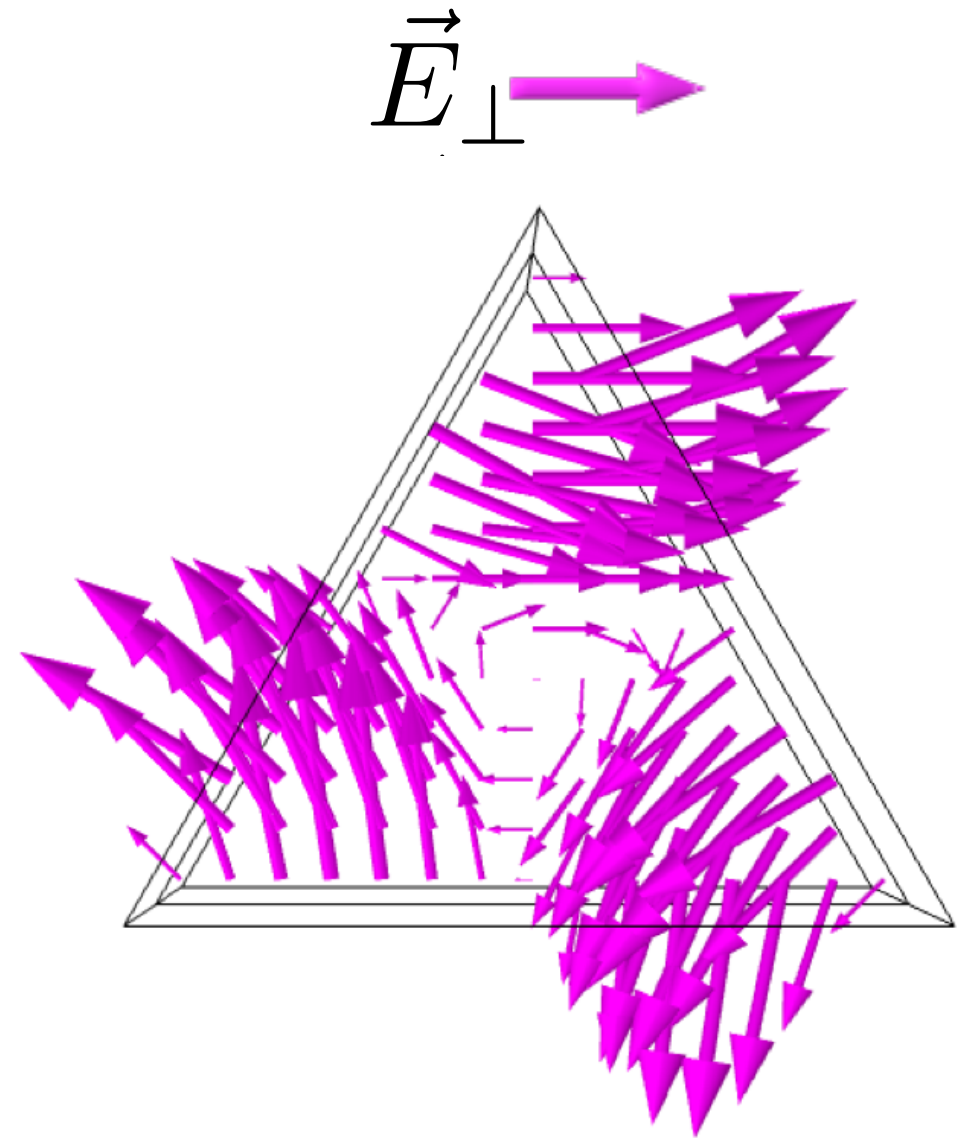
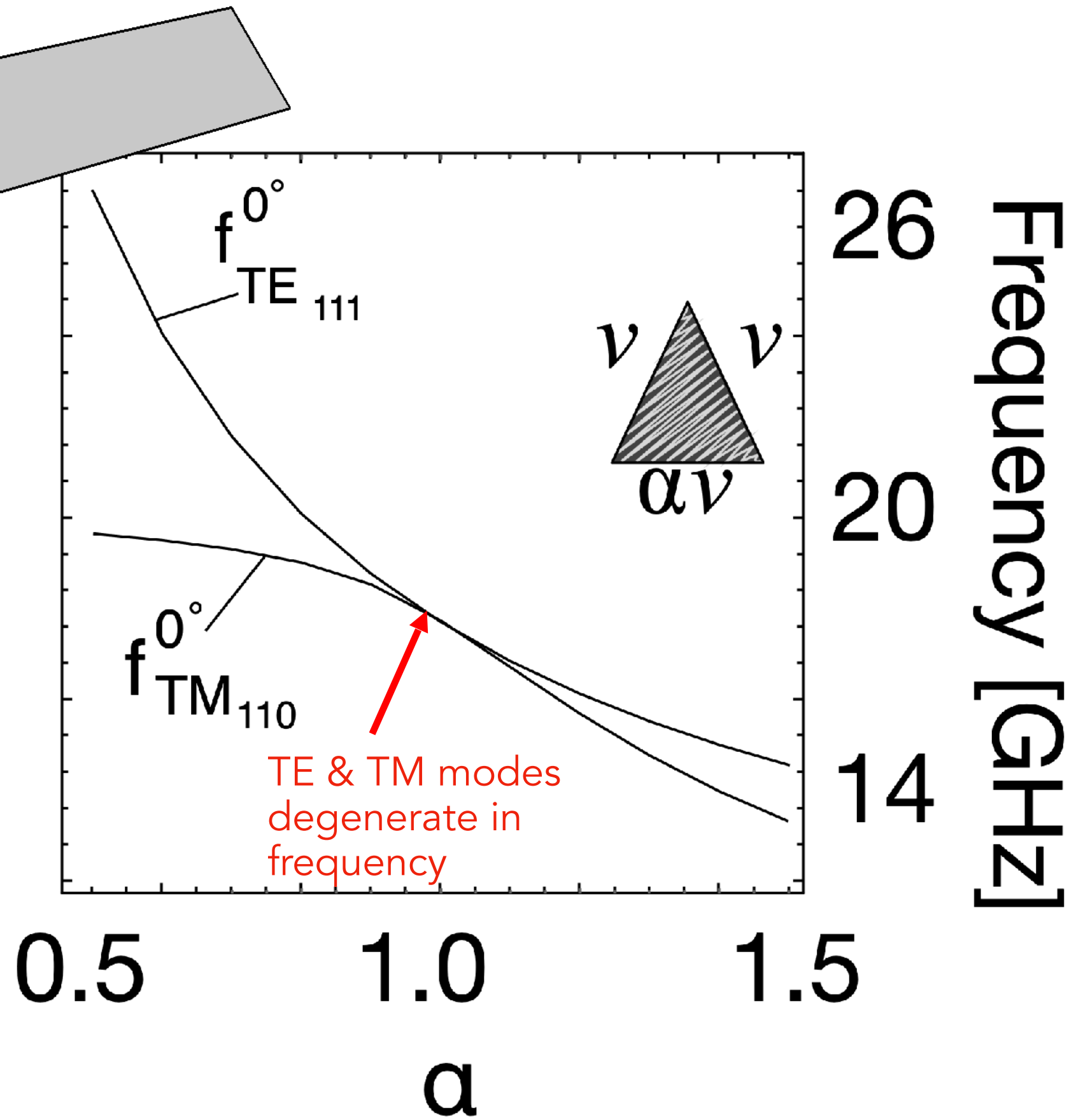
How do these geometries generate helicity?



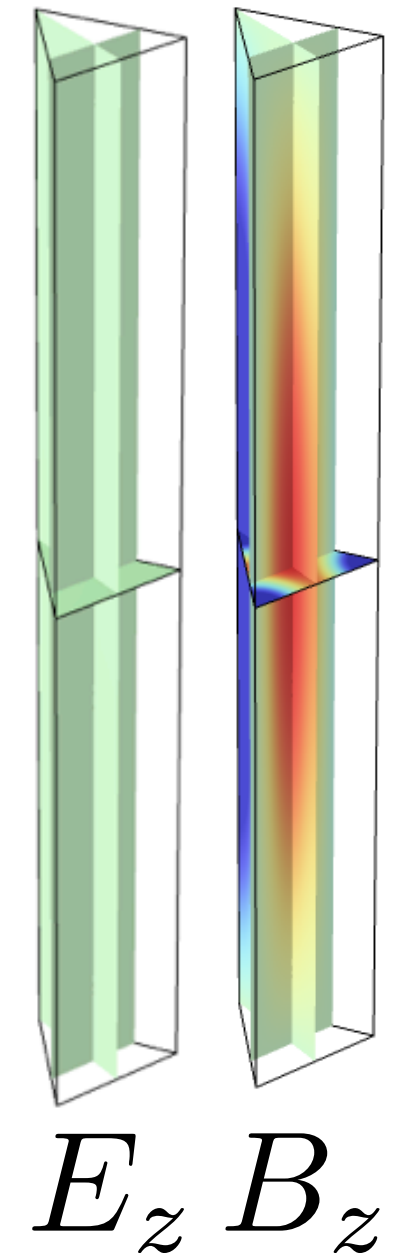
Transverse Magnetic (TM)



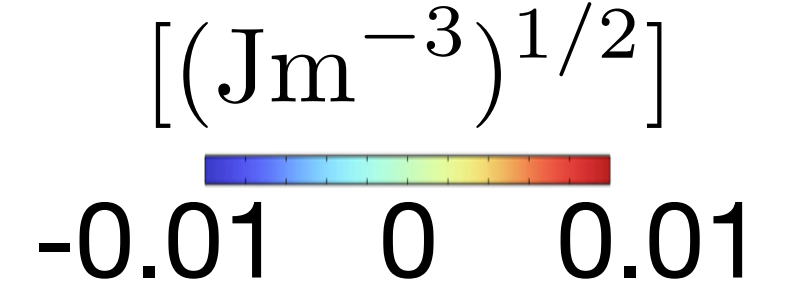
E_z B_z



Transverse Electric (TE)

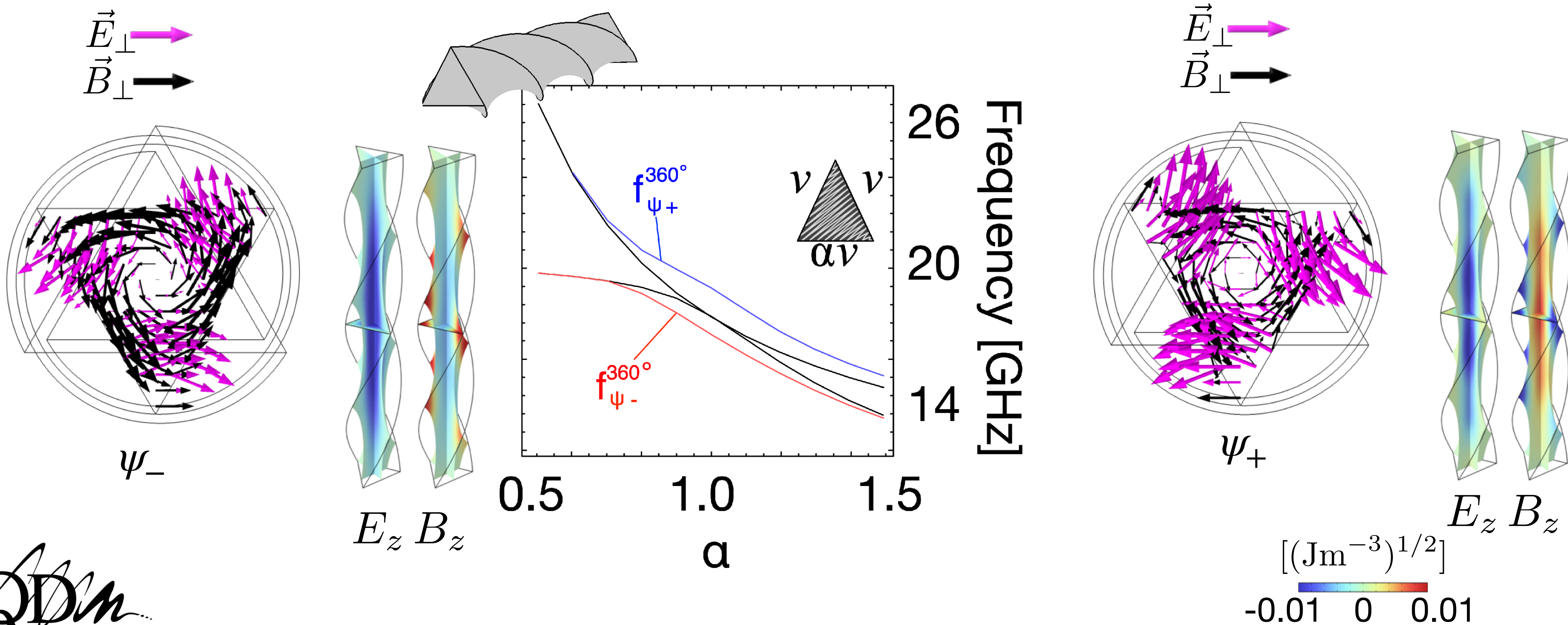


E_z B_z



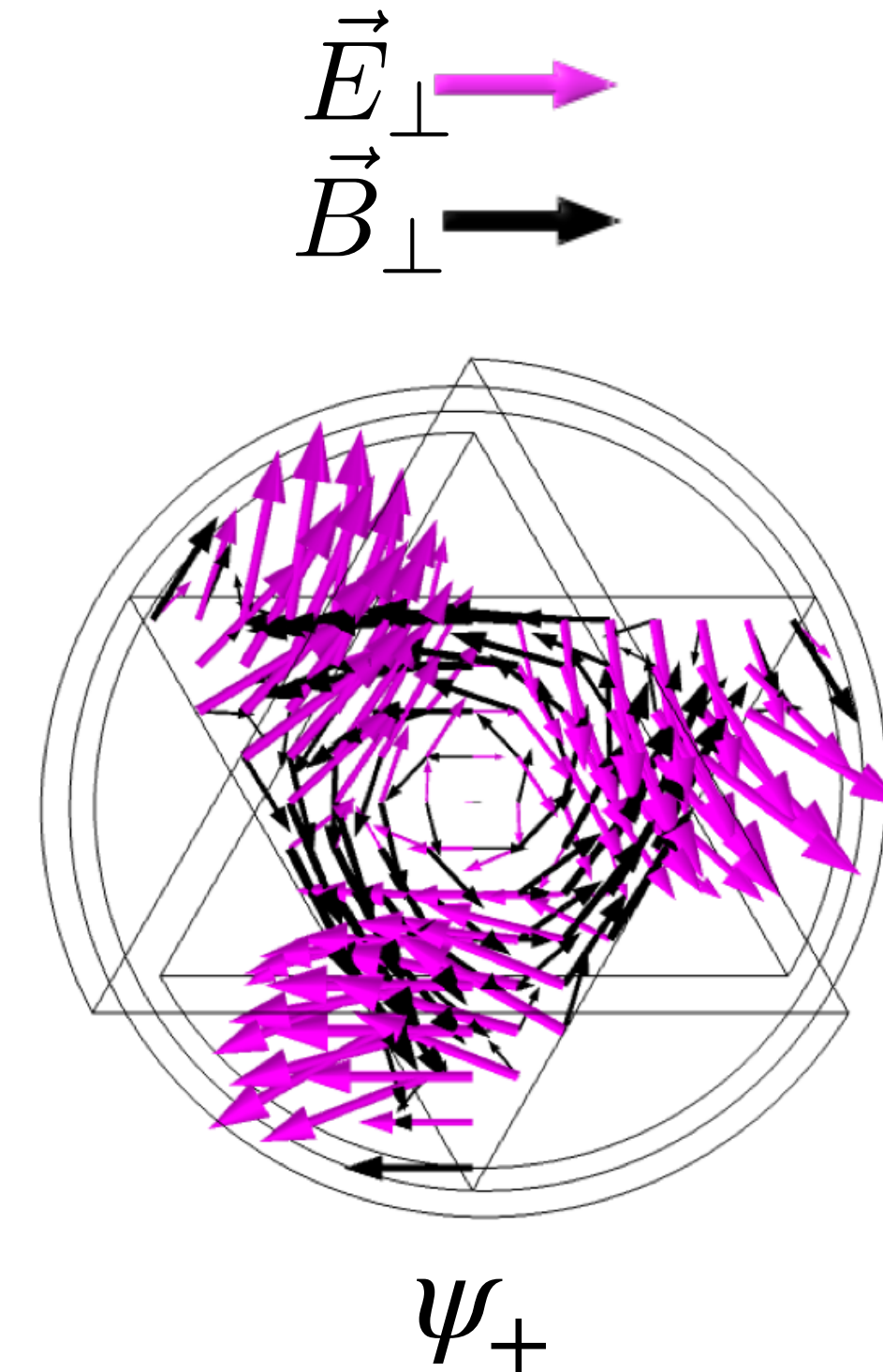
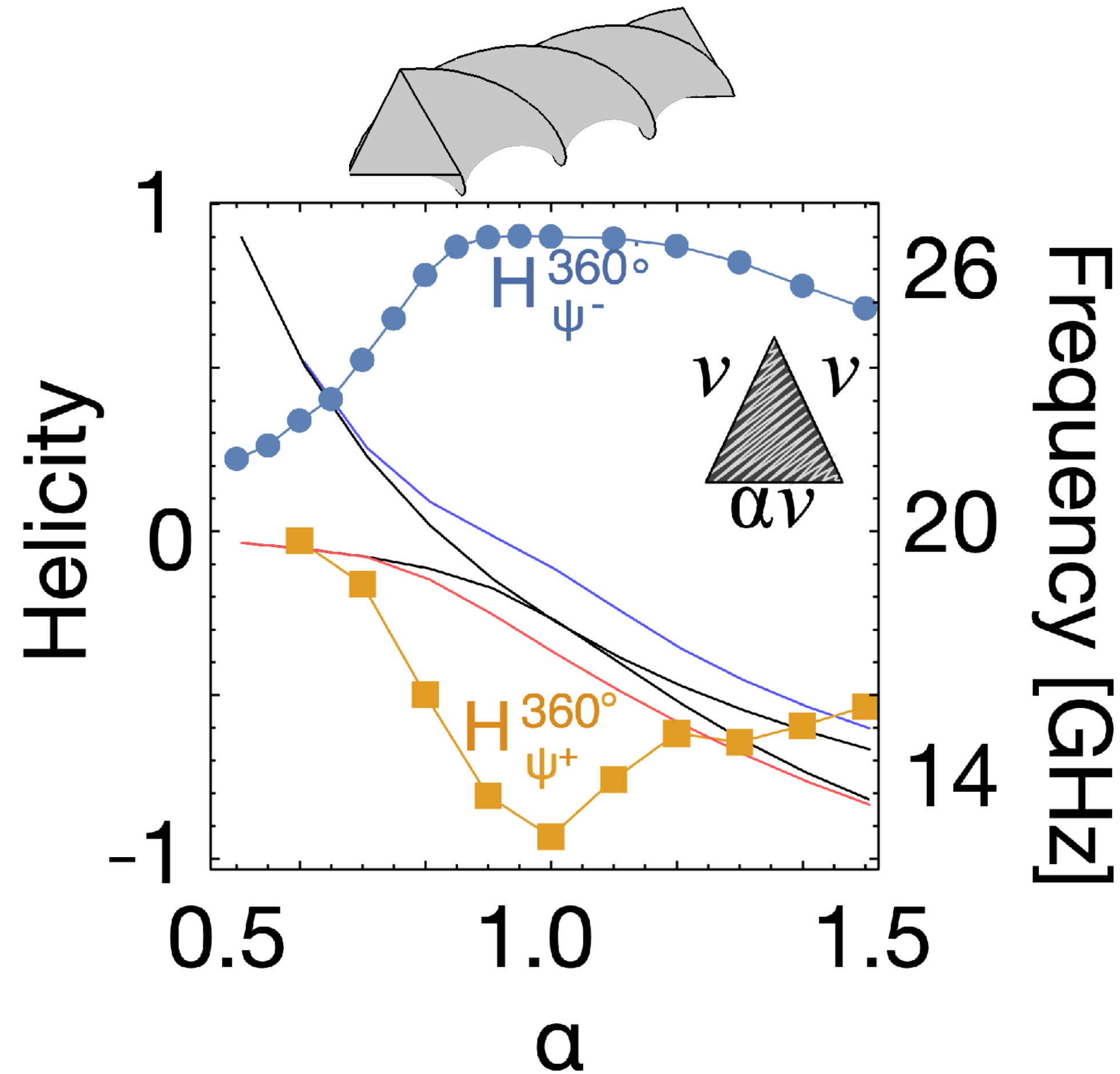
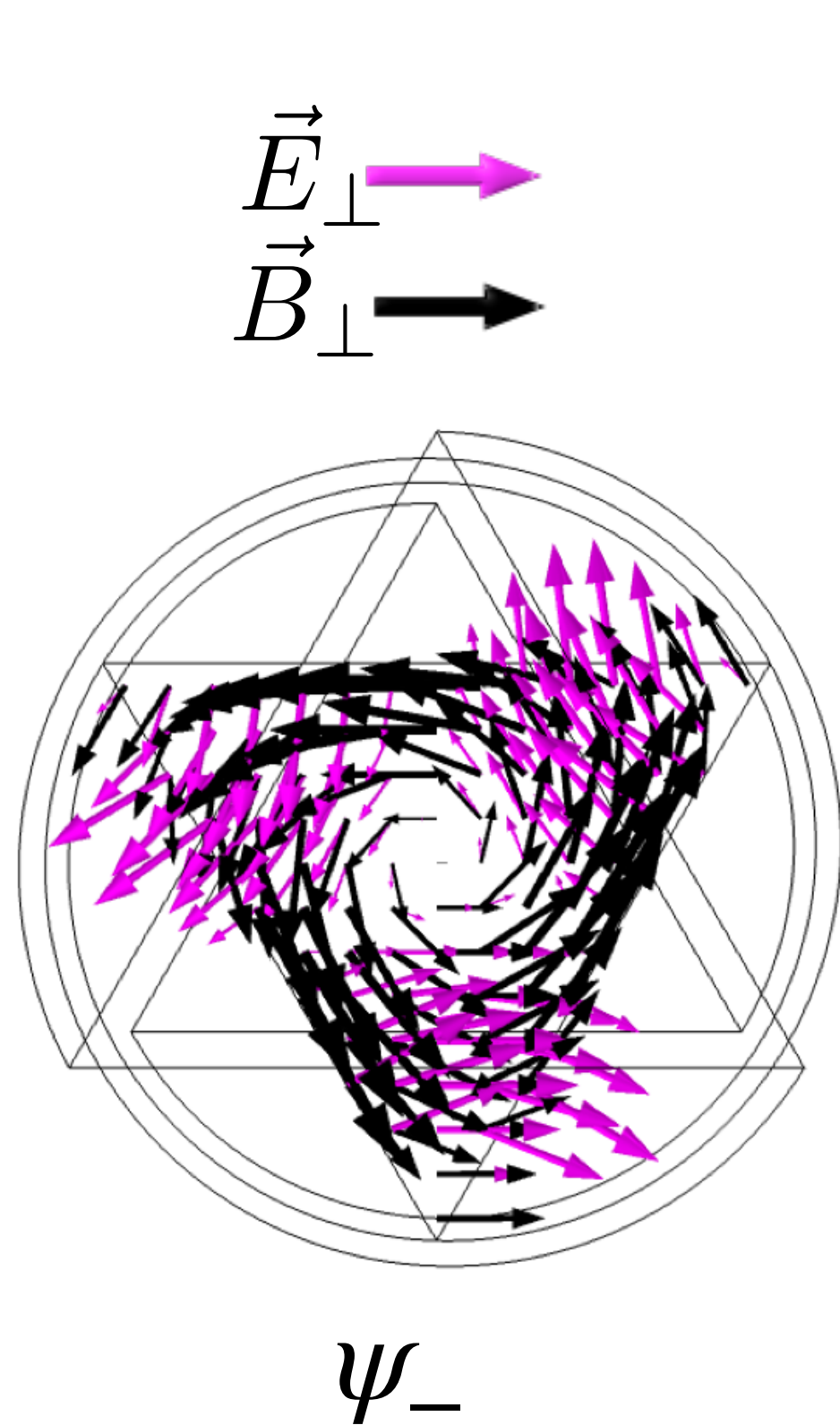
Twisted "anyon" microwave cavities

How do these geometries generate helicity?



Twisted "anyon" microwave cavities

How do these geometries generate helicity?



Twisted “anyon” microwave cavities

Can we build these geometries?

Twisted "anyon" microwave cavities

Can we build these geometries?

- 3D printed using Selective Laser Melting printers and Al-Si powder



Twisted "anyon" microwave cavities

Can we build these geometries?

- 3D printed using Selective Laser Melting printers and Al-Si powder



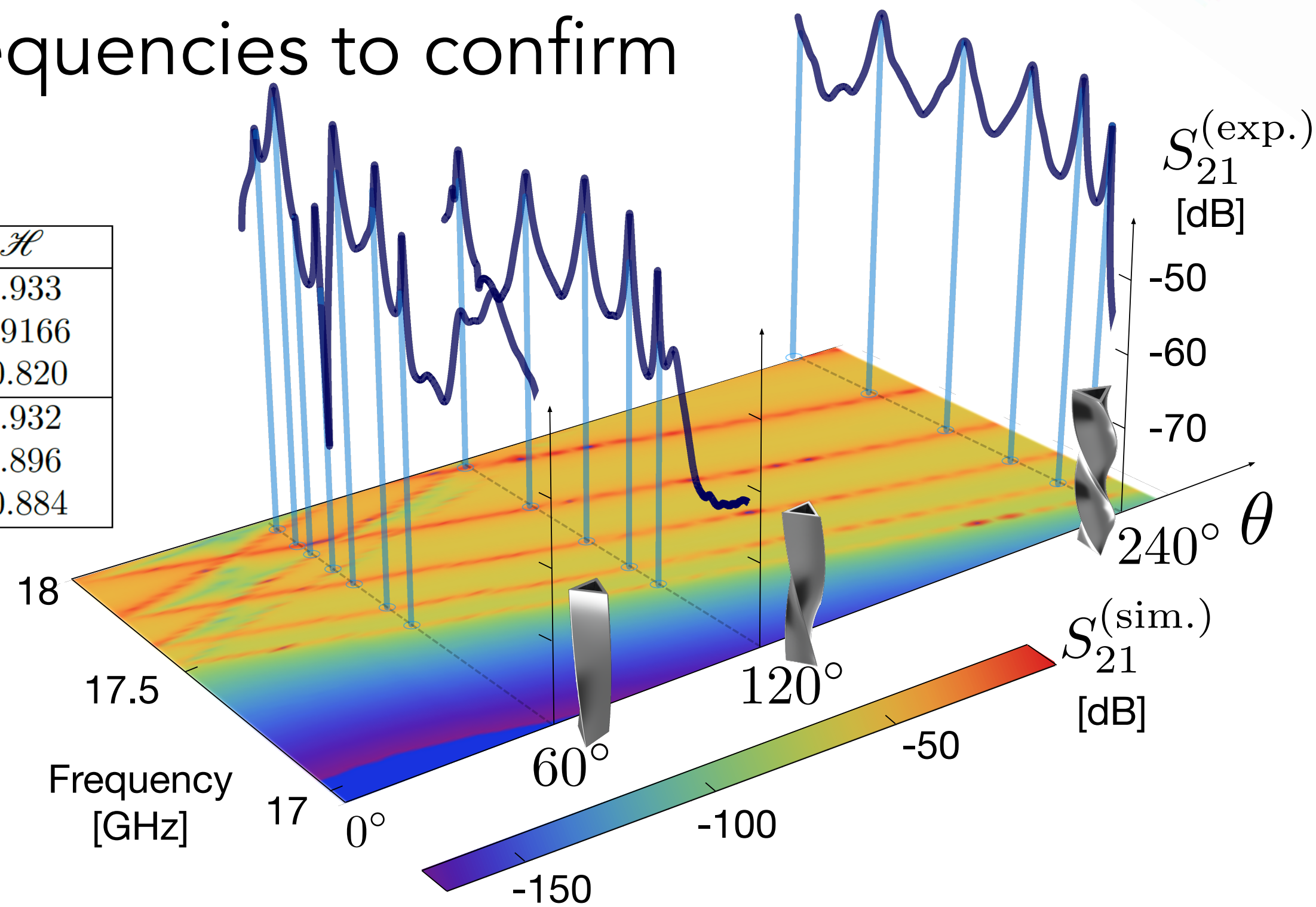
Resonator	Mode	f (GHz)	G (Ω)	\mathcal{H}
Ring	ψ_0^-	17.221	6200	0.933
Ring	ψ_1^-	17.297	6570	0.9166
Ring	ψ_0^+	17.895	7290	-0.820
Linear	ψ_0^-	17.214	1950	0.932
Linear	ψ_1^-	17.278	2030	0.896
Linear	ψ_0^+	17.859	1920	-0.884

Twisted "anyon" microwave cavities

Can we build these geometries?

- 3D printed using Selective Laser Melting printers and Al-Si powder
- Measured mode frequencies to confirm simulation results

Resonator	Mode	f (GHz)	G (Ω)	\mathcal{H}
Ring	ψ_0^-	17.221	6200	0.933
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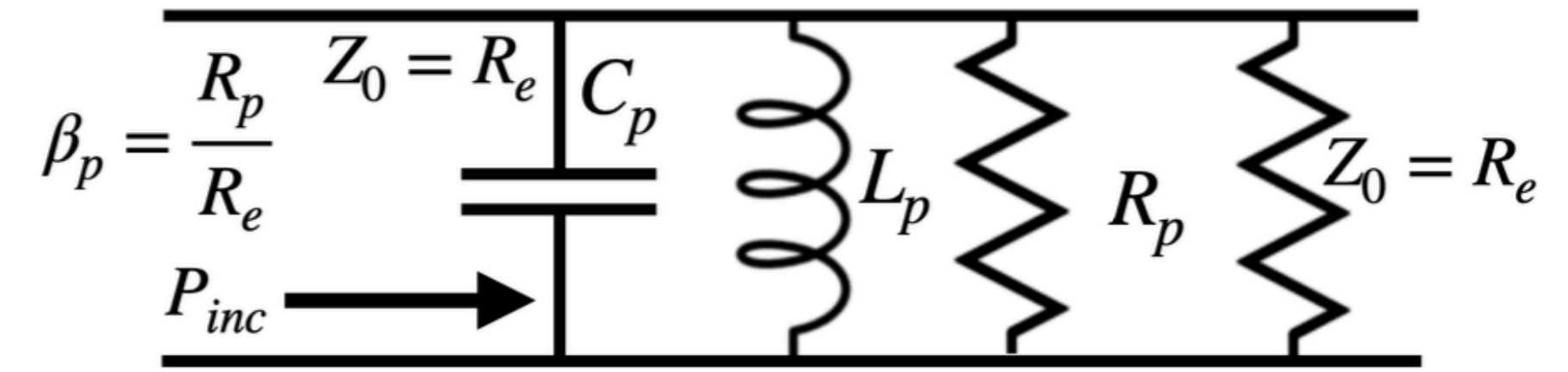
Twisted “anyon” microwave cavities

Dark matter detection thanks to helicity

Twisted "anyon" microwave cavities

Dark matter detection thanks to helicity

- Due to the helicity, **ultra-light dark matter axions**, whose mass range falls within the cavity bandwidth will amplitude modulate the cavity mode



$$P_p = \frac{\beta_p P_d}{\beta_p + 1} = \frac{4\beta_p^2}{(1 + \beta_p)^2} P_{inc}.$$

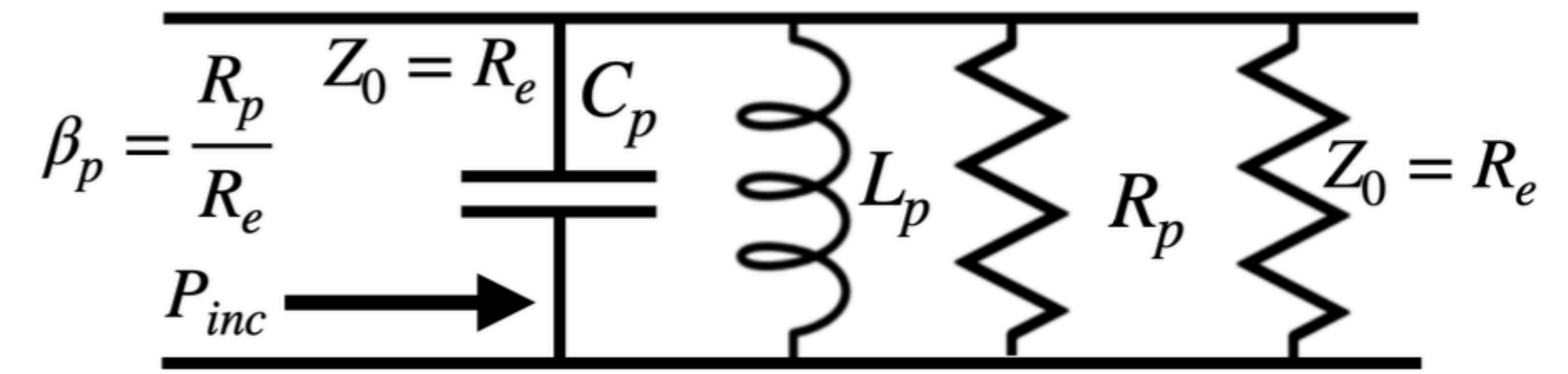
$$\frac{P_{am}}{P_{inc}} = \frac{m_{am}^2 P_p}{P_{inc}} = Q_p^2 \frac{4\beta_p^2}{(1 + \beta_p)^2} \left(\frac{\omega_a}{\omega_p}\right)^2 \frac{\langle \theta_0 \rangle^2}{8} \mathcal{H}_p^2.$$

$$SNR = \frac{g_{a\gamma\gamma} \beta_p |\mathcal{H}_p|}{\sqrt{2}(1 + \beta_p)} \frac{Q_p}{\sqrt{1 + 4Q_p^2 \left(\frac{\omega_a}{\omega_p}\right)^2}} \frac{\left(\frac{10^6 t}{\omega_a}\right)^{\frac{1}{4}} \sqrt{\rho_a c^3}}{\omega_p \sqrt{S_{am}}}$$

Twisted "anyon" microwave cavities

Dark matter detection thanks to helicity

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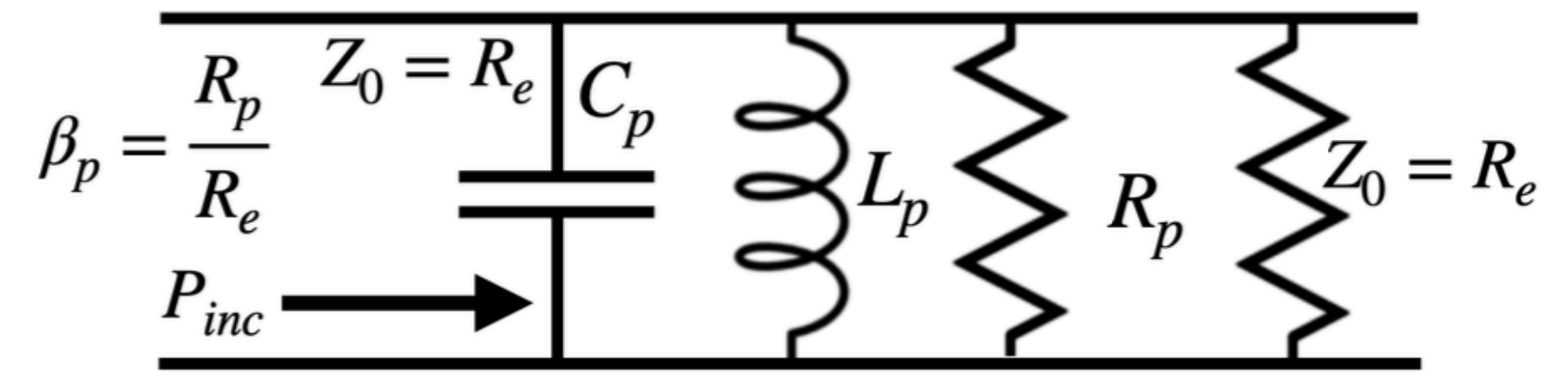
Labels for the equation:

- $g_{a\gamma\gamma}$: Axion Photon Coupling
- β_p : Microwave Probe Coupling
- \mathcal{H}_p : Helicity
- Q_p : Q factor
- ω_a : Axion Frequency
- $10^6 t$: Measurement time (1 week)
- ρ_a : Cold dark matter density ($8 \times 10^{-22} \text{kgm}^{-3}$)
- c : Speed of light ($3 \times 10^8 \text{ms}^{-1}$)
- ω_p : Cavity frequency (1 GHz)
- S_{am} : Amplitude noise (-160dBcHz^{-1})

Twisted "anyon" microwave cavities

Dark matter detection thanks to helicity

- Due to the helicity, **ultra-light dark matter axions**, whose mass range falls within the cavity bandwidth will amplitude modulate the cavity mode
- The frequency of the AM will be proportional to the axion mass



$$P_p = \frac{\beta_p P_d}{\beta_p + 1} = \frac{4\beta_p^2}{(1 + \beta_p)^2} P_{inc}$$

$$\frac{P_{am}}{P_{inc}} = \frac{m_{am}^2 P_p}{P_{inc}} = Q_p^2 \frac{4\beta_p^2}{(1 + \beta_p)^2} \left(\frac{\omega_a}{\omega_p}\right)^2 \frac{\langle \theta_0 \rangle^2}{8} \mathcal{H}_p^2$$

$$SNR = \frac{\overset{\text{Axion Photon Coupling}}{g_{a\gamma\gamma}} \overset{\text{Microwave Probe Coupling}}{\beta_p} \overset{\text{Helicity}}{\mathcal{H}_p}}{\sqrt{2}(1 + \beta_p)} \frac{\overset{\text{Q factor}}{Q_p}}{\sqrt{1 + 4Q_p^2 \left(\frac{\omega_a}{\omega_p}\right)^2}} \frac{\overset{\text{Measurement time (1 week)}}{10^6 t}^{\frac{1}{4}} \overset{\text{Cold dark matter density (8 \times 10^{-22} \text{ kg m}^{-3})}}{\sqrt{\rho_a c^3}}}{\underset{\text{Cavity frequency (1 GHz)}}{\omega_p} \underset{\text{Amplitude noise (-160 dBcHz}^{-1})}}{\sqrt{S_{am}}}}$$

Speed of light (3x10⁸ ms⁻¹)

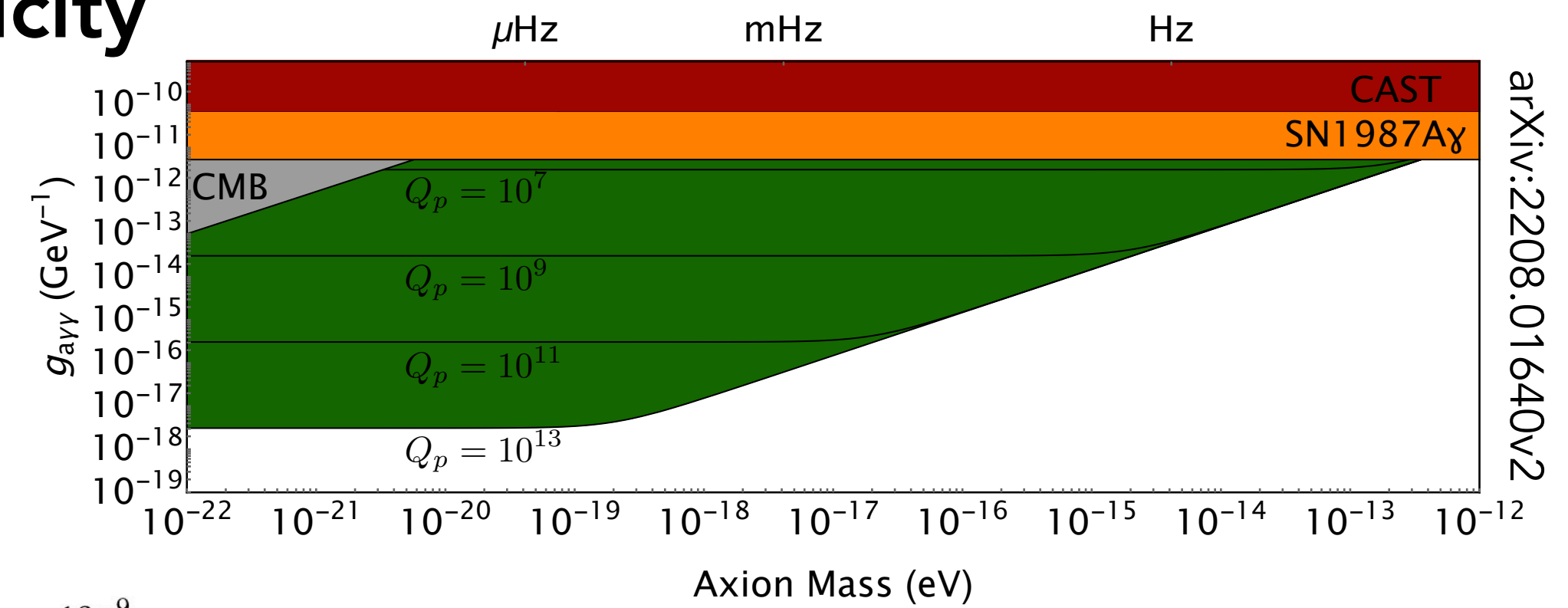
Twisted “anyon” microwave cavities

Dark matter detection in a single mode thanks to helicity

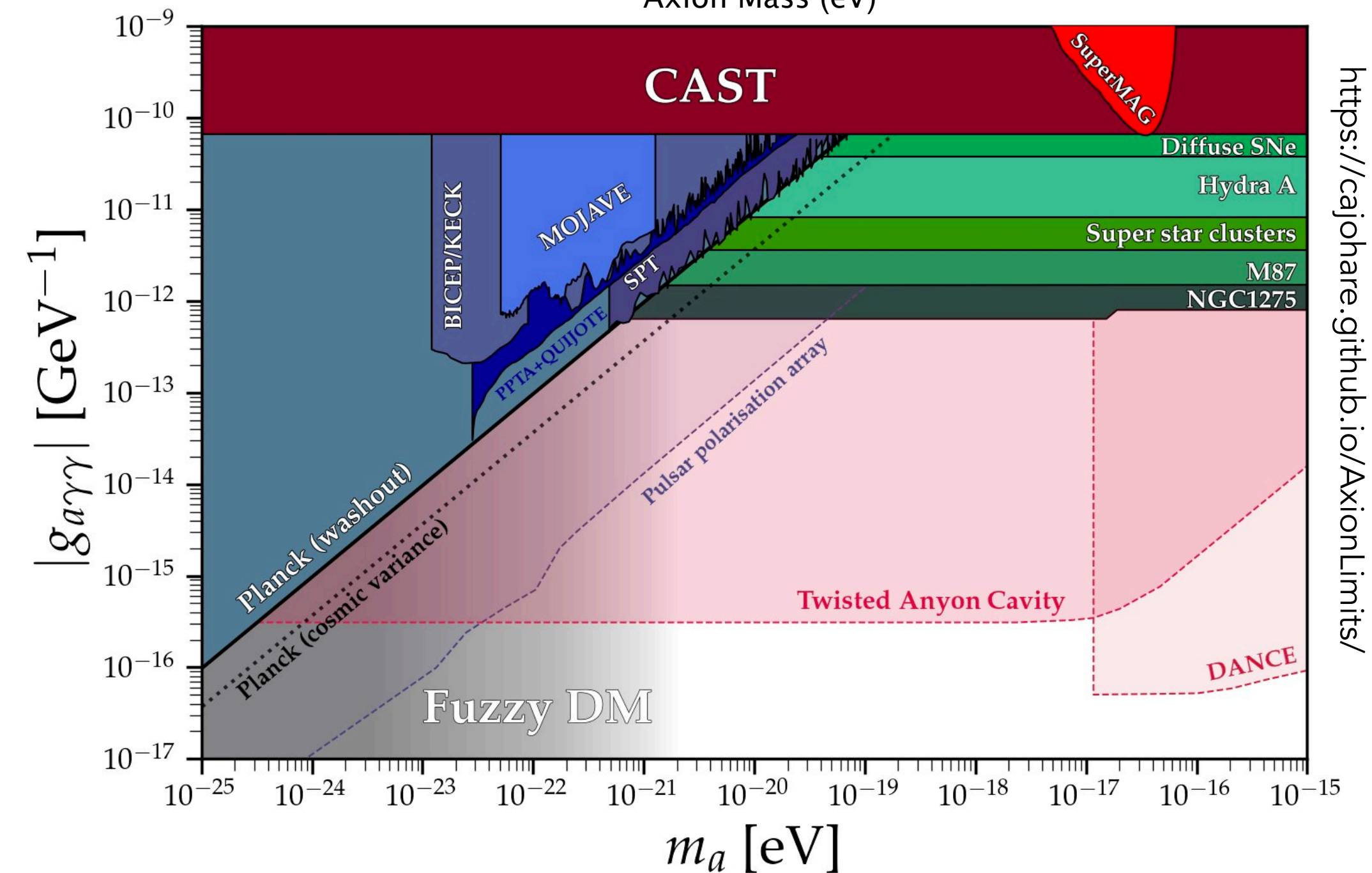
Twisted "anyon" microwave cavities

Dark matter detection in a single mode thanks to helicity

- Accesses an axion mass range very difficult to search



arXiv:2208.01640v2

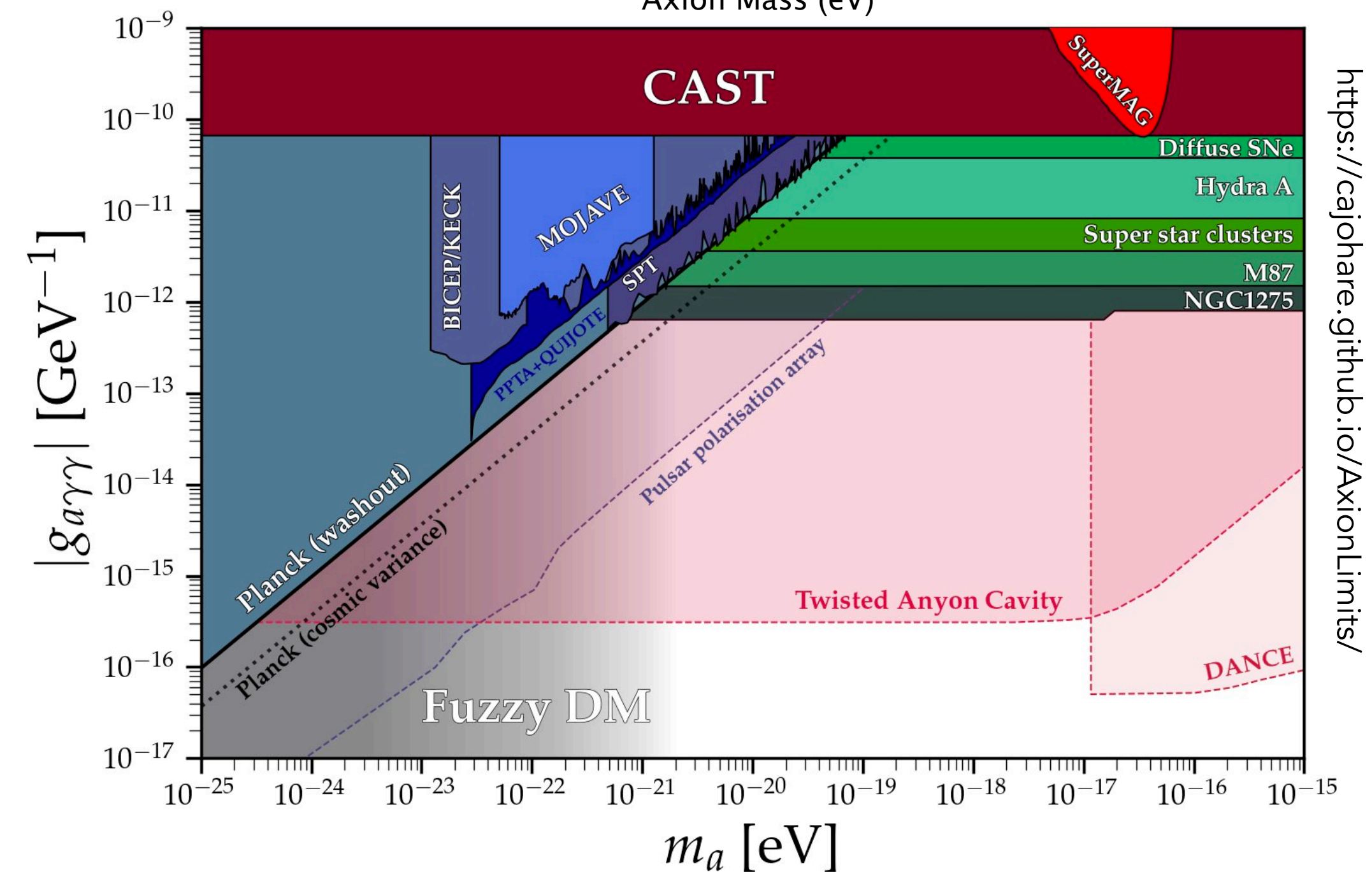
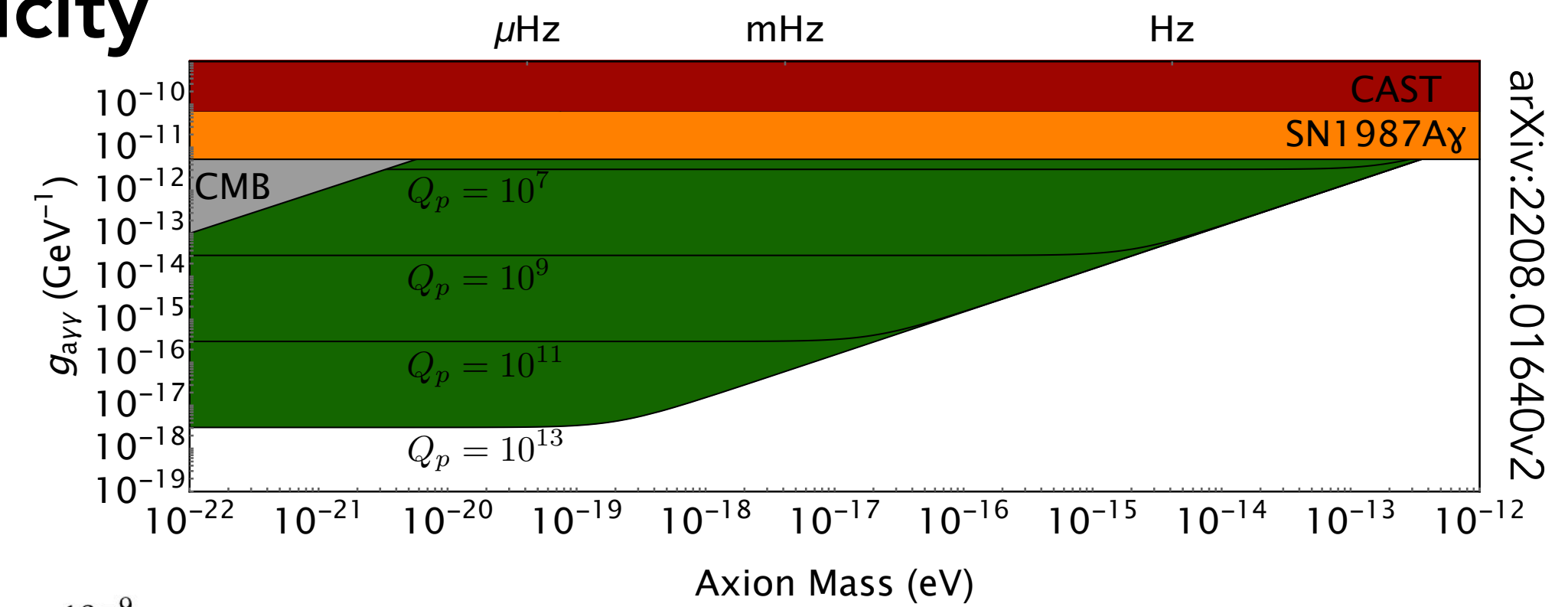


<https://cajohare.github.io/AxionLimits/>

Twisted "anyon" microwave cavities

Dark matter detection in a single mode thanks to helicity

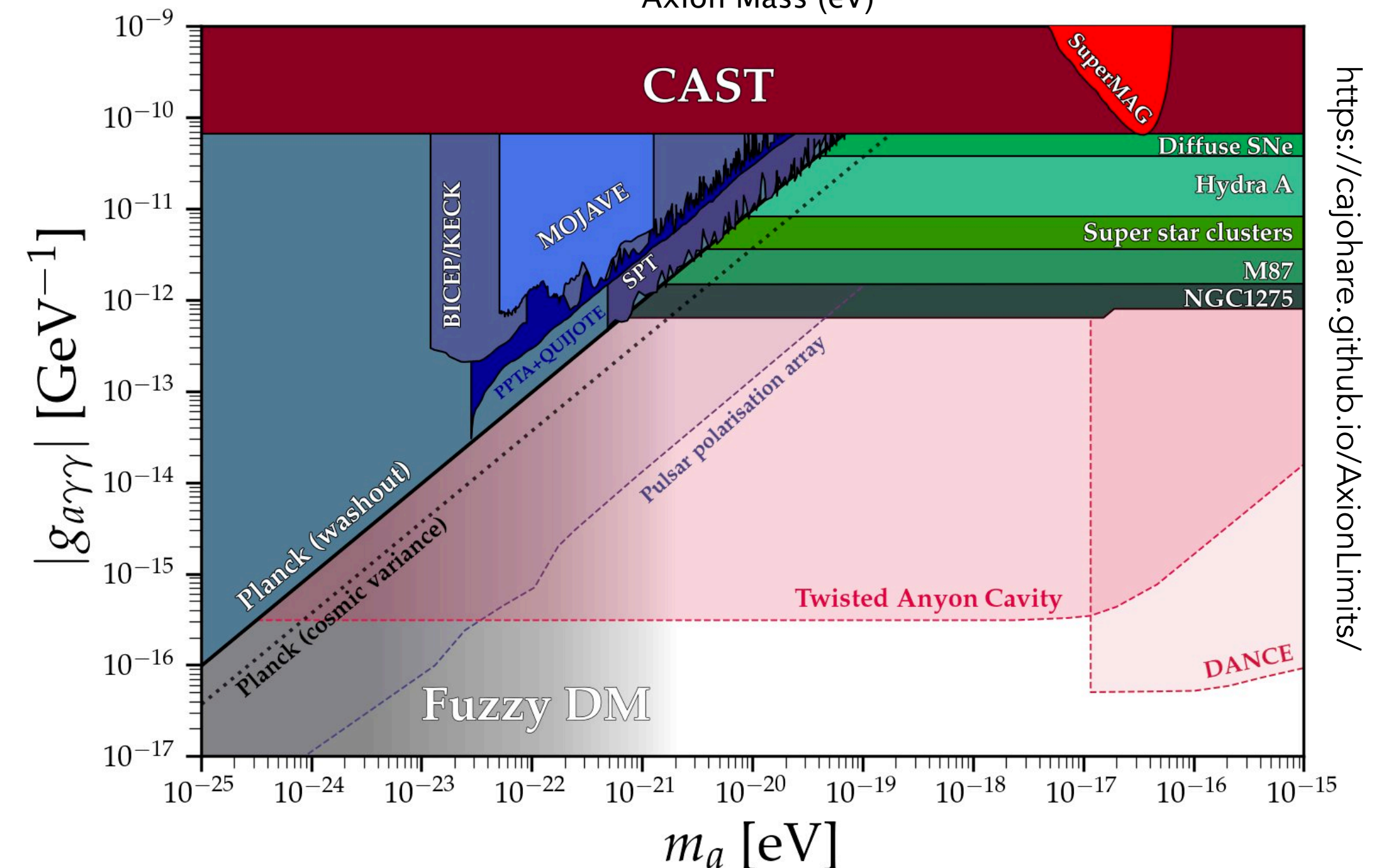
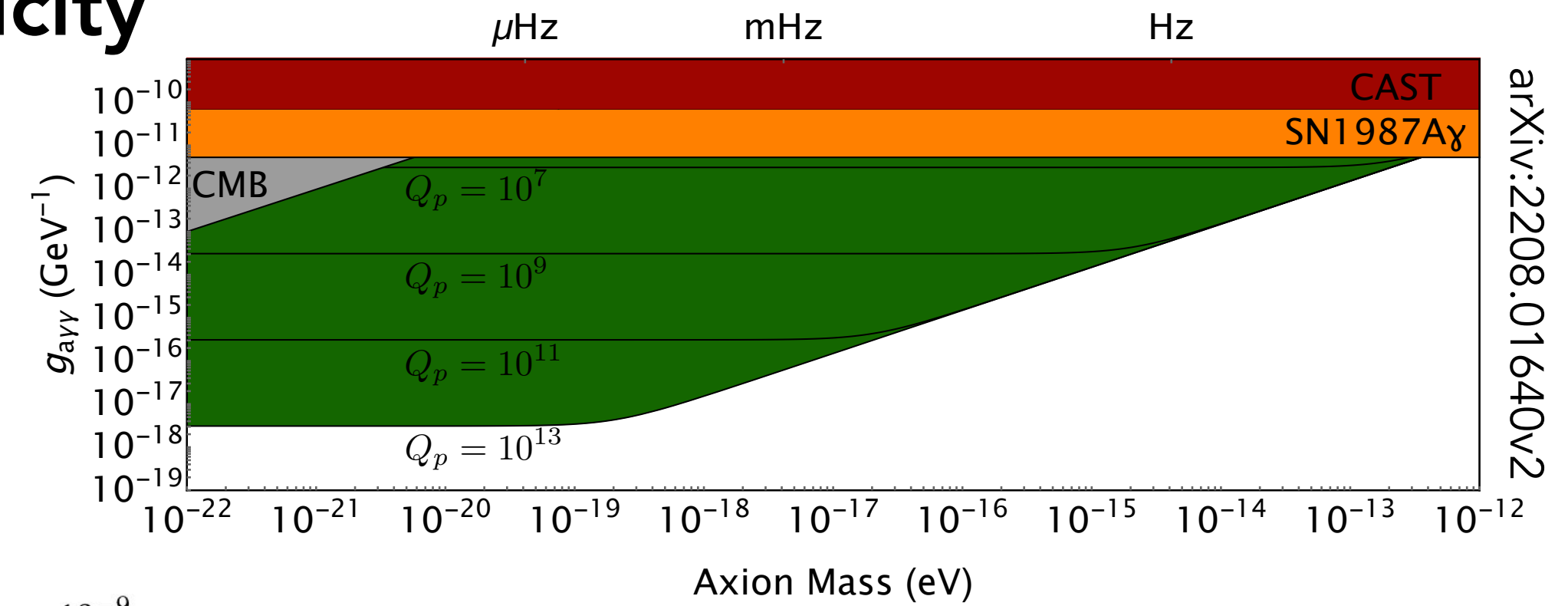
- Accesses an axion mass range very difficult to search
- **No external magnetic field needed**



Twisted "anyon" microwave cavities

Dark matter detection in a single mode thanks to helicity

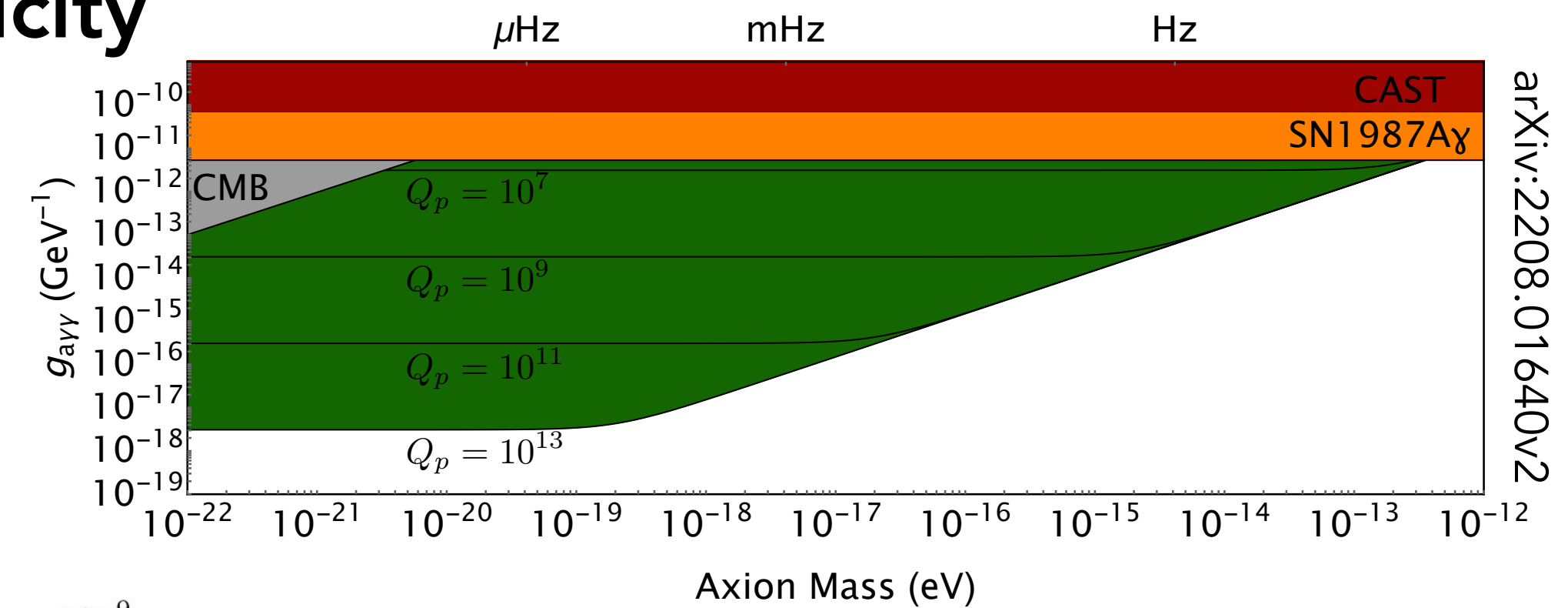
- Accesses an axion mass range very difficult to search
- **No external magnetic field needed**
- Ability to use **superconducting** materials



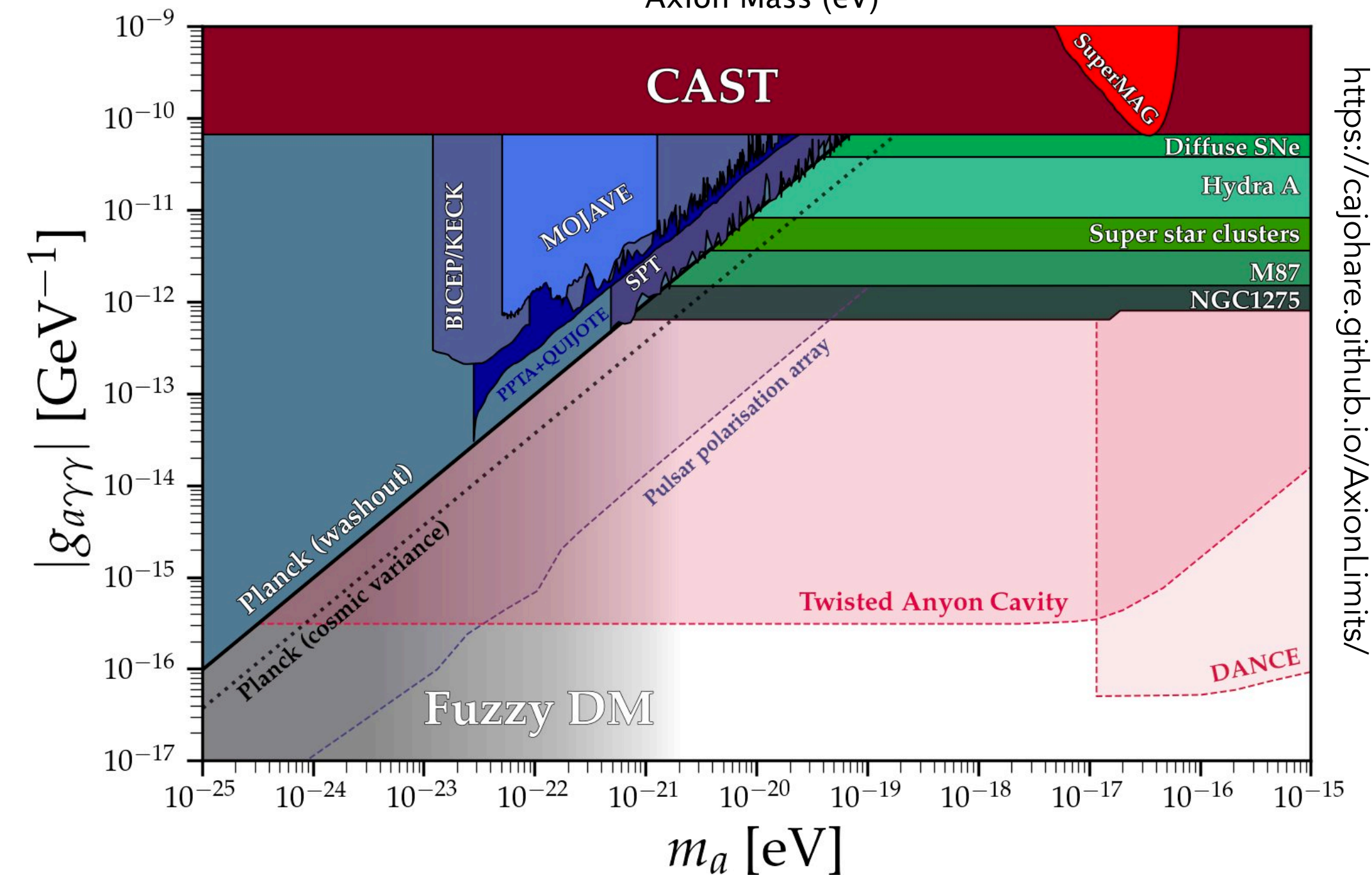
Twisted "anyon" microwave cavities

Dark matter detection in a single mode thanks to helicity

- Accesses an axion mass range very difficult to search
- **No external magnetic field needed**
- Ability to use **superconducting** materials
- Allows high Q-factors and improved sensitivity



arXiv:2208.01640v2

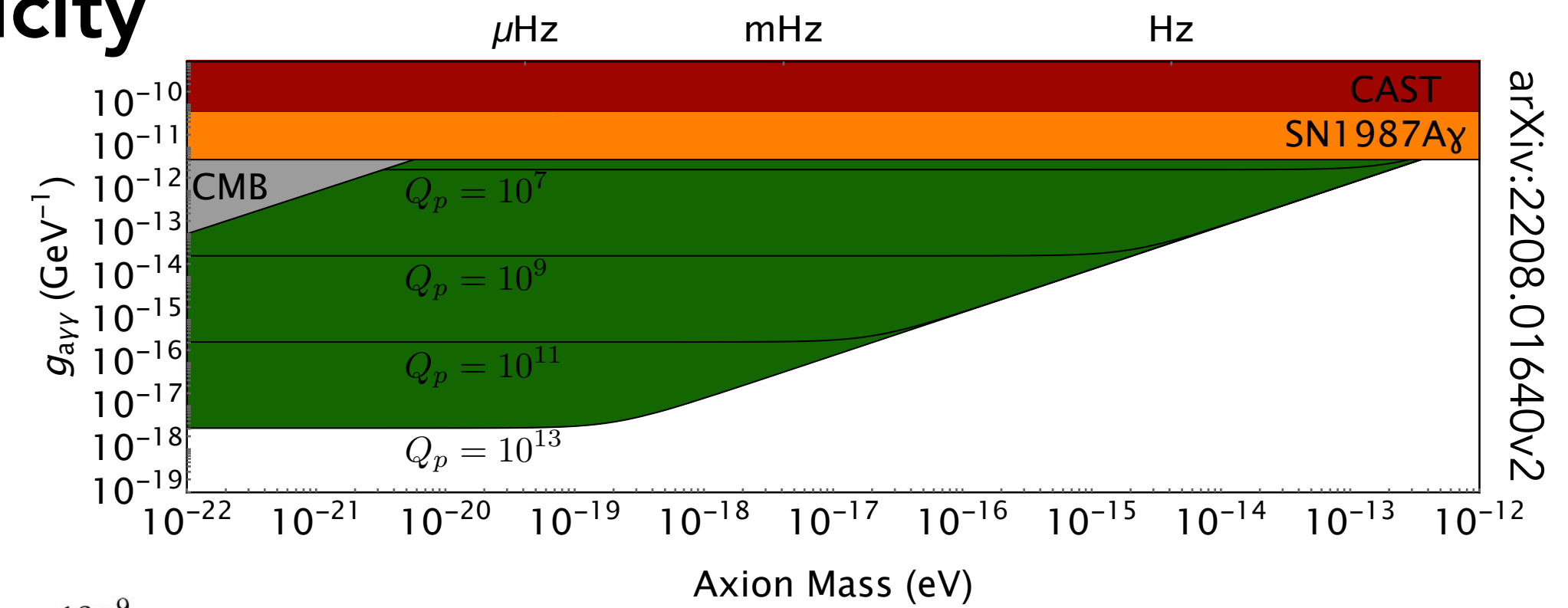


<https://cajohare.github.io/AxionLimits/>

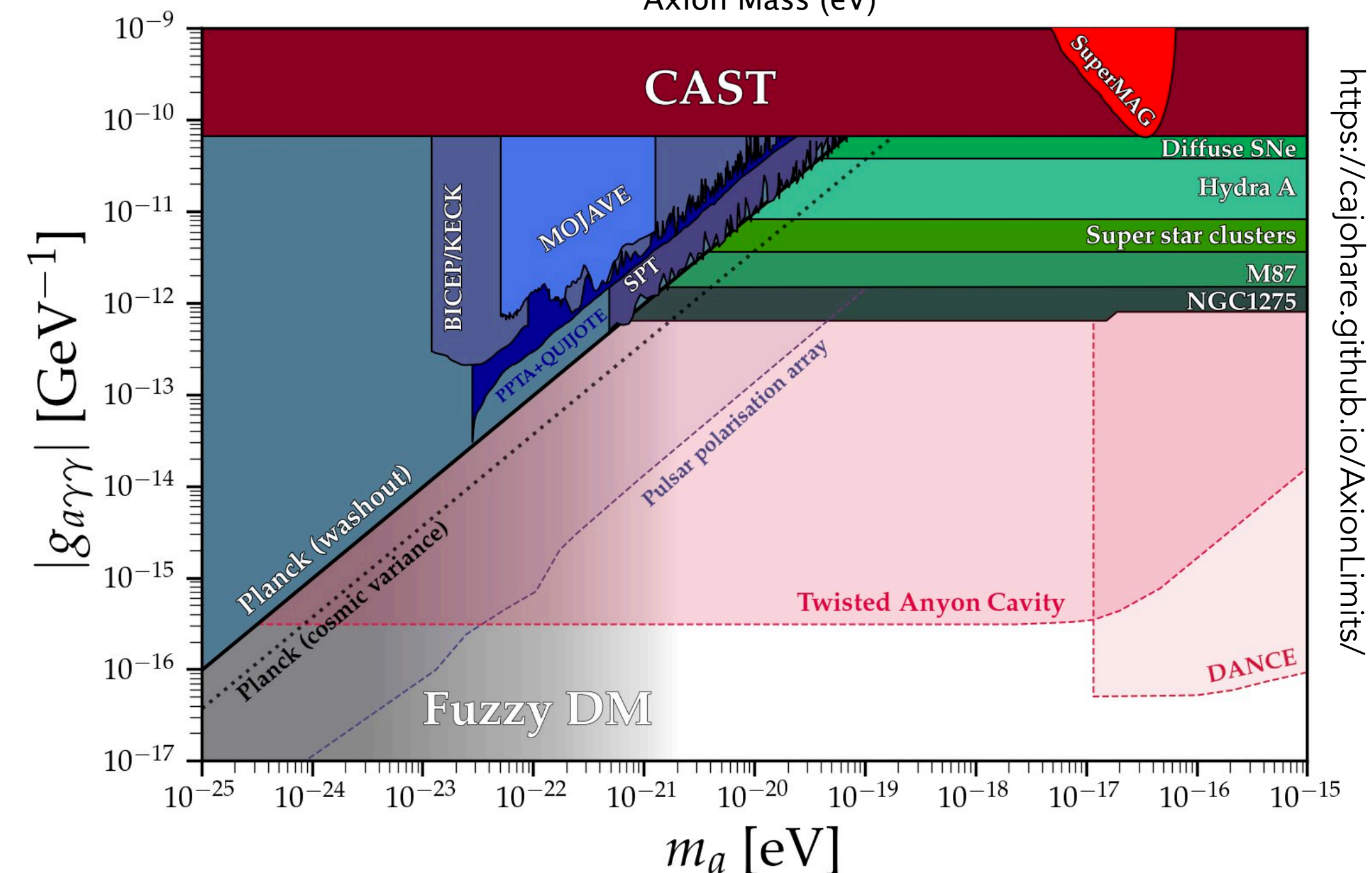
Twisted "anyon" microwave cavities

Dark matter detection in a single mode thanks to helicity

- Accesses an axion mass range very difficult to search
- **No external magnetic field needed**
- Ability to use **superconducting** materials
- Allows high Q-factors and improved sensitivity
- Next: Optimising Q-factors and minimising read-out amplitude modulation noise for a detection run



arXiv:2208.01640v2



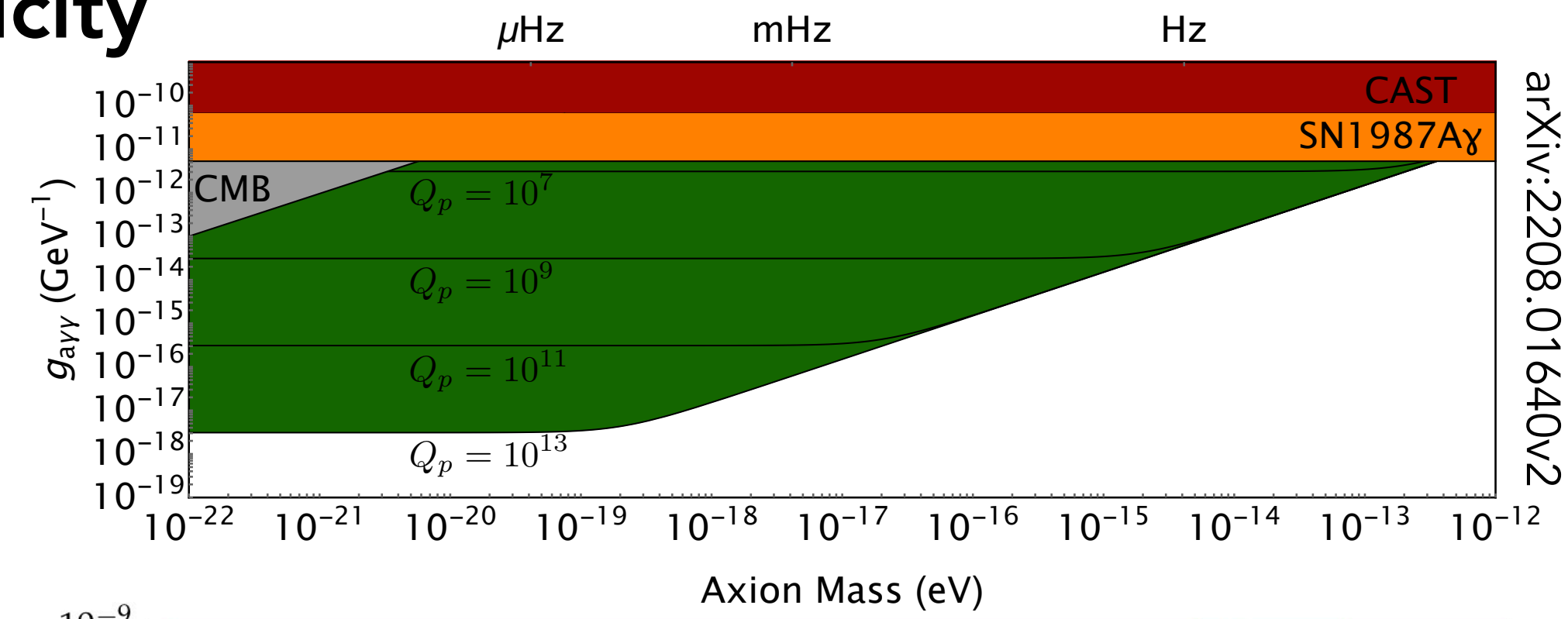
<https://cajohare.github.io/AxionLimits/>

Twisted "anyon" microwave cavities

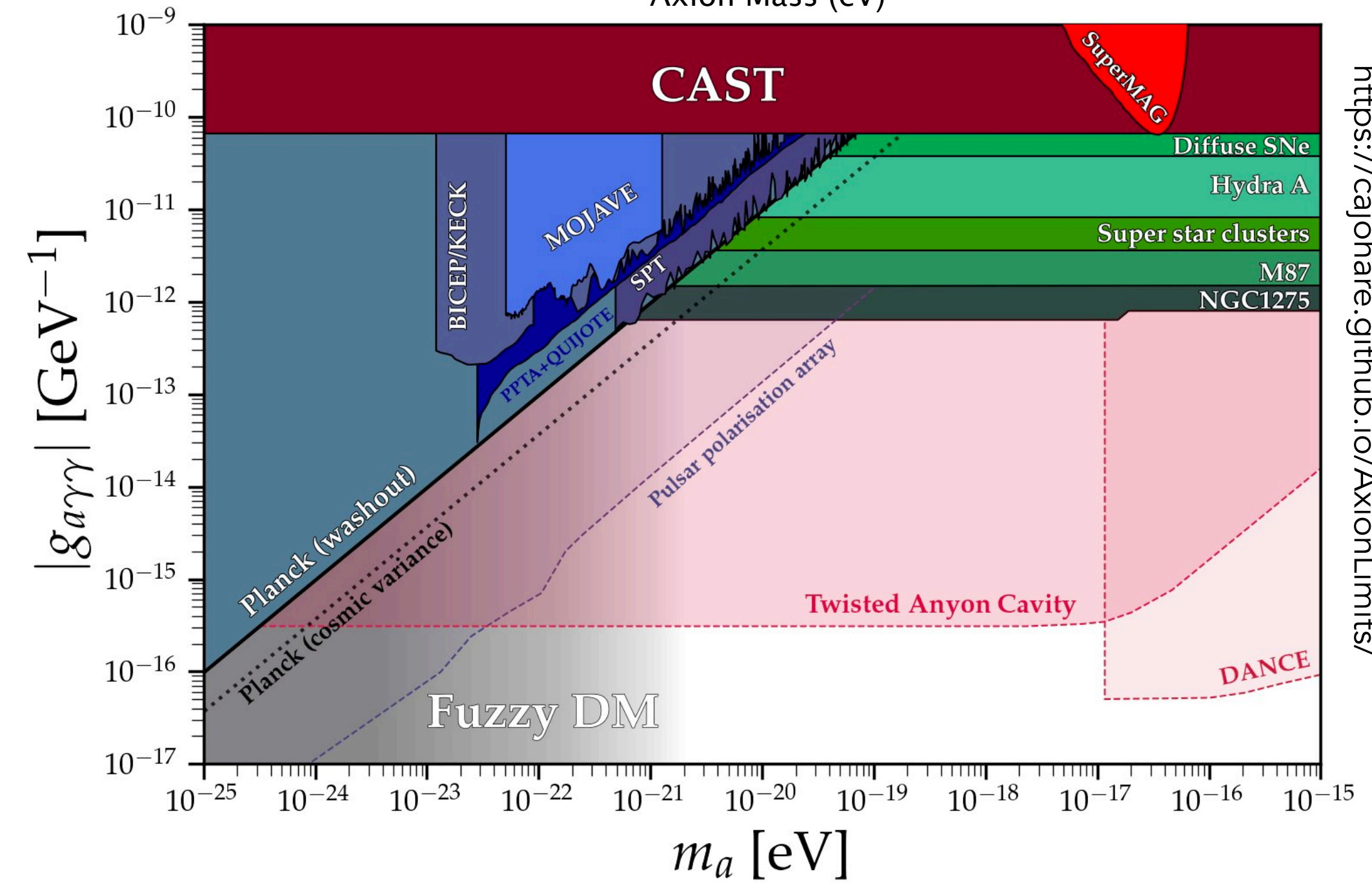
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See poster #82 by
Jeremy Bourhill



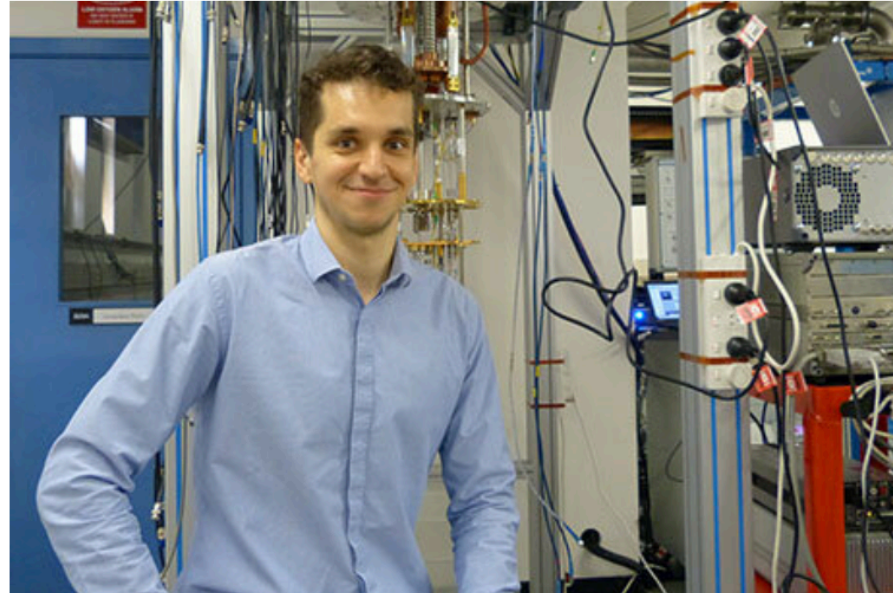
arXiv:2208.01640v2



https://cajohare.github.io/AxionLimits/



Professor Mike Tobar
Director



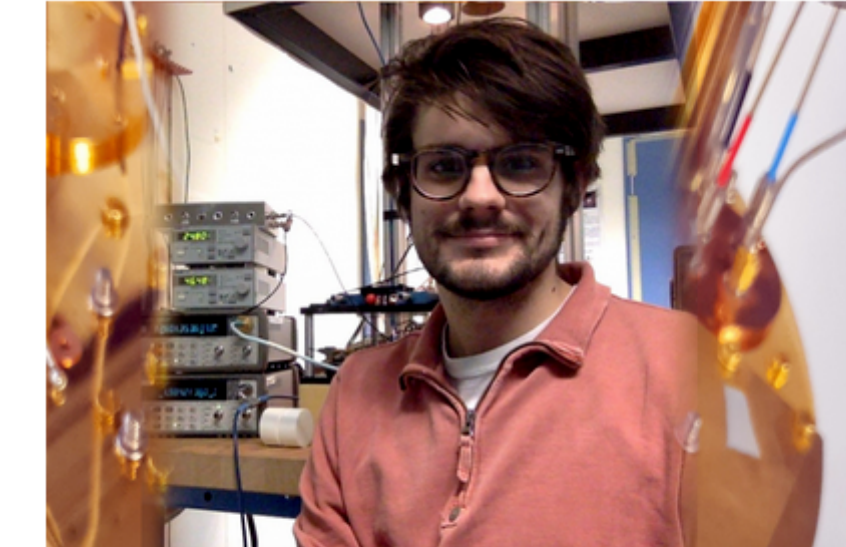
Dr Maxim Goryachev
Research Associate



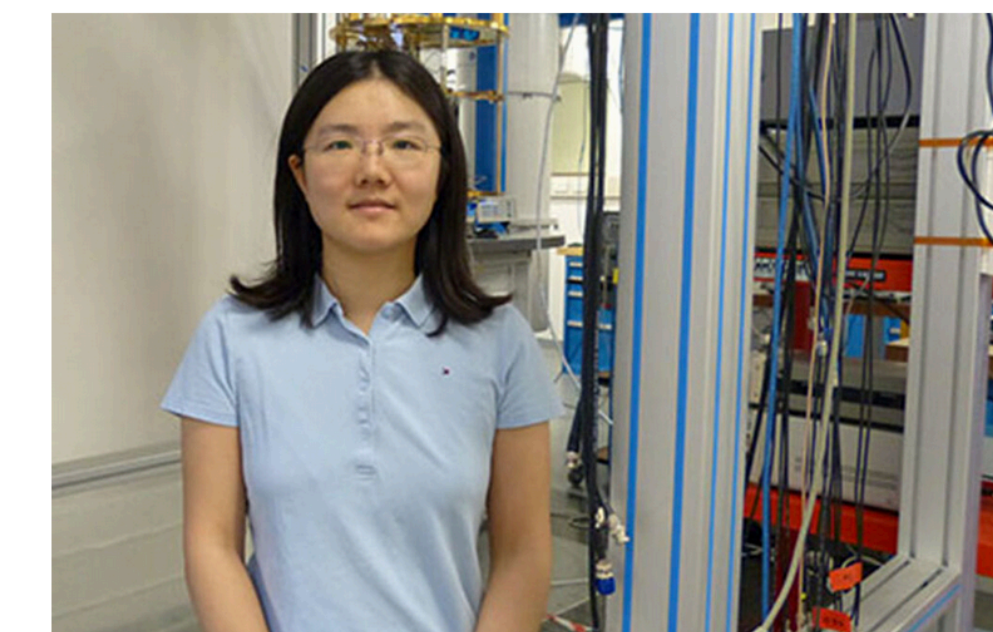
Dr Ben McAllister
Research Associate



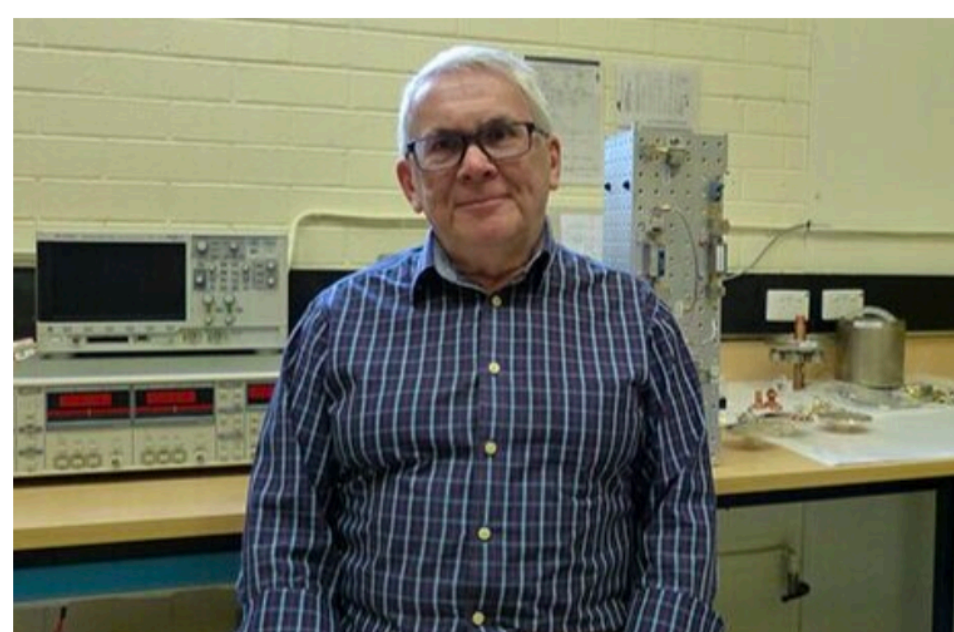
Professor Eugene Ivanov
Winthrop Research Professor—Dept of Physics



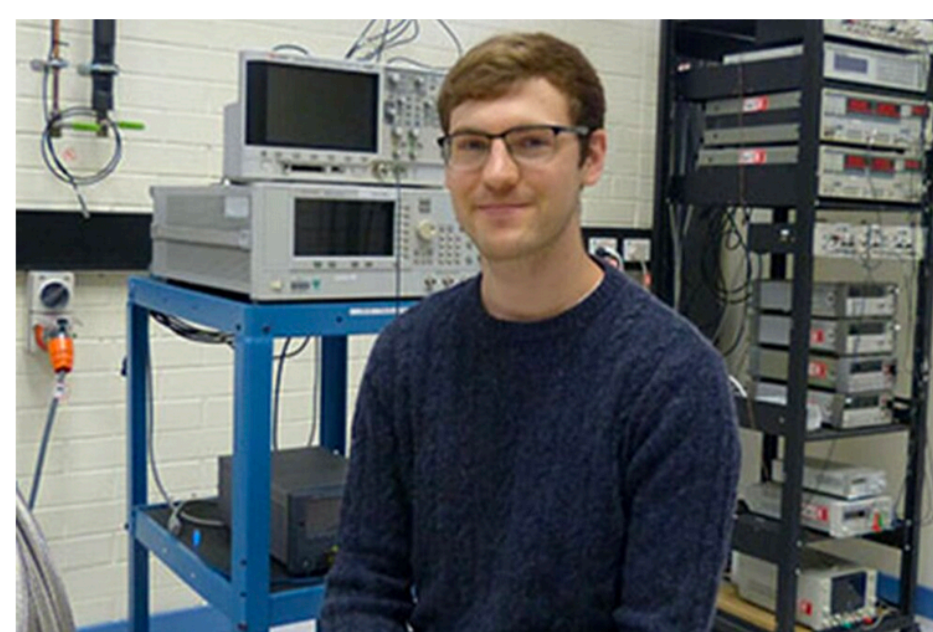
Dr Jeremy Bourhill
Postdoctoral Research Associate



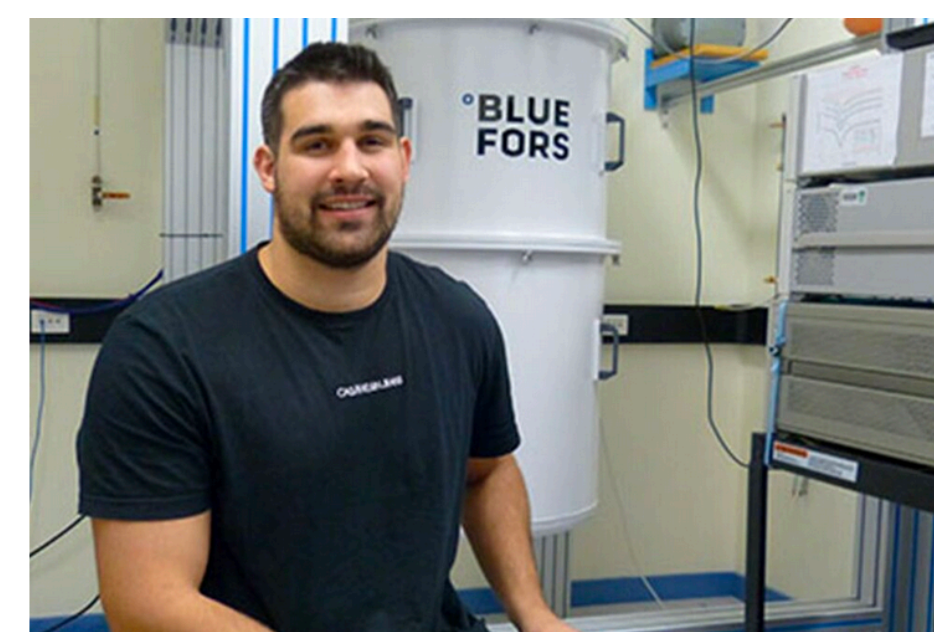
Dr Cindy Zhao
Deborah Jin Fellow—EQUS



Professor Alexey Veryaskin
Adjunct Professor



Graeme Flower
PhD



Aaron Quiskamp
PhD



Will Campbell
PhD



Catriona Thomson
PhD



Elrina Hartman
PhD



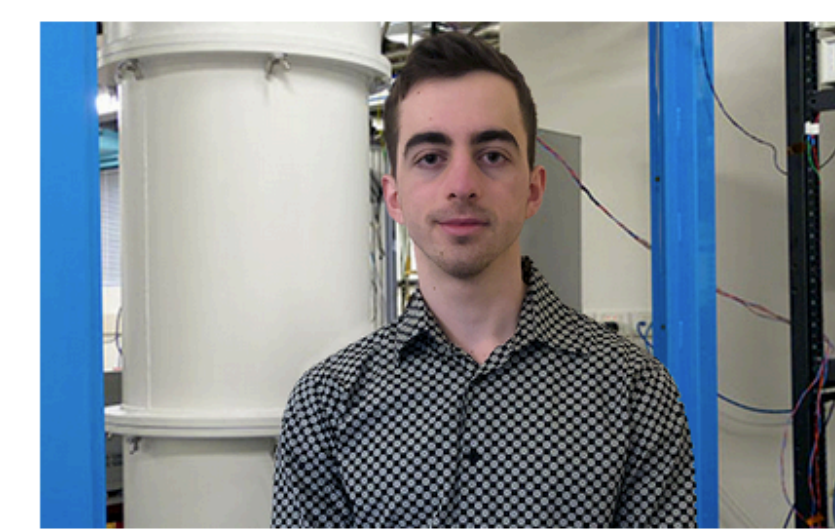
Jay Mummery
Masters



Robert Crew
BPhil (Hons) Placement



Daniel Tobar
BPhil (Hons) Placement



Michael Hatzon
BPhil (Hons) Placement



Steve Osborne
Technician

**THE END
THE
TEAM**

