



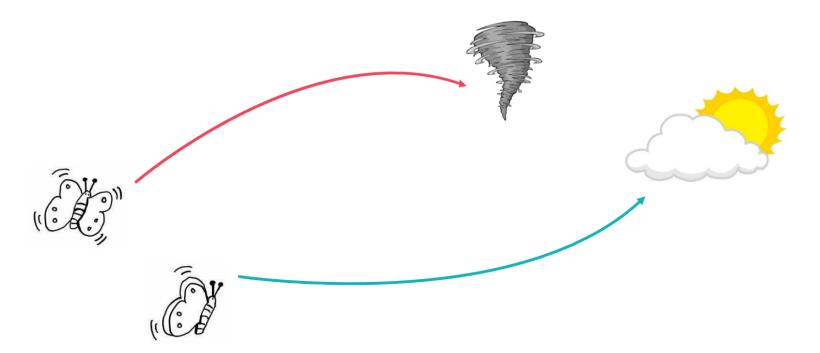
Quantum Chaos = Volume-Law Spatiotemporal Entanglement

Neil Dowling

Neil Dowling and Kavan Modi (2022), arXiv:2210.14926

Classical Chaos: The Butterfly Effect

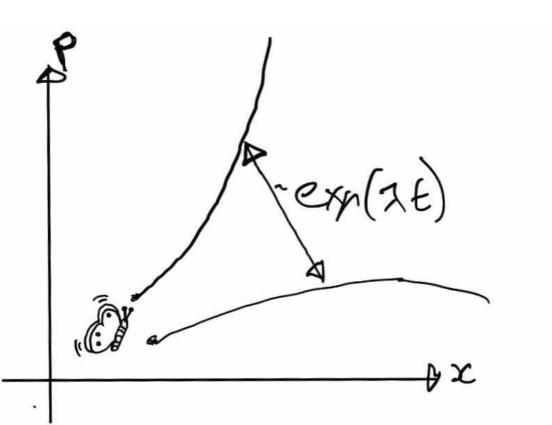
"The phenomenon that a small alteration in the state of a dynamical system will cause subsequent states to differ greatly from the states that would have followed without the alteration"



Lorenz, Edward (1993). The Essence of Chaos.

Classical Chaos: Lyapunov Exponents

Classical chaos determined by a exponential sensitivity to perturbation



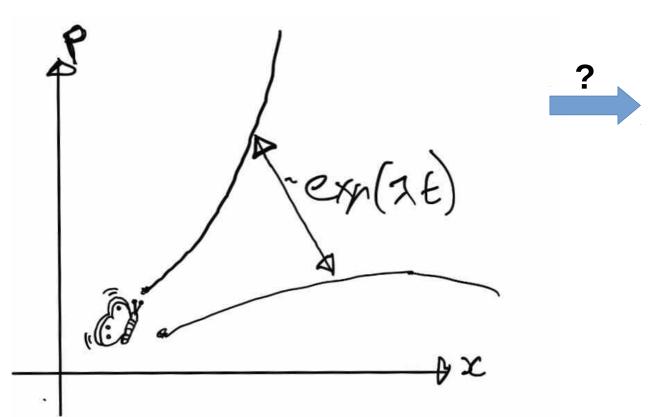
Classical Chaos: Lyapunov Exponents

Classical chaos determined by a exponential sensitivity to perturbation

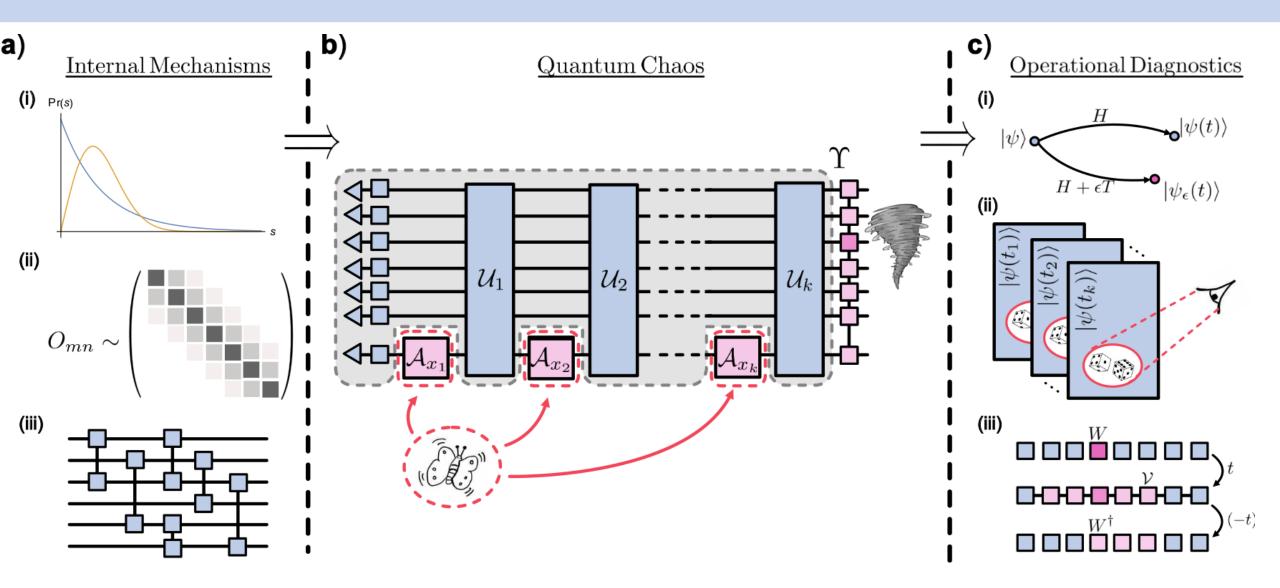
Linearity of isolated quantum mechanics – how can there be chaos?!

 $\langle \psi_t | \phi_t \rangle = \langle \psi | U_t^{\dagger} U_t | \phi \rangle$

 $=\langle \psi | \phi \rangle$

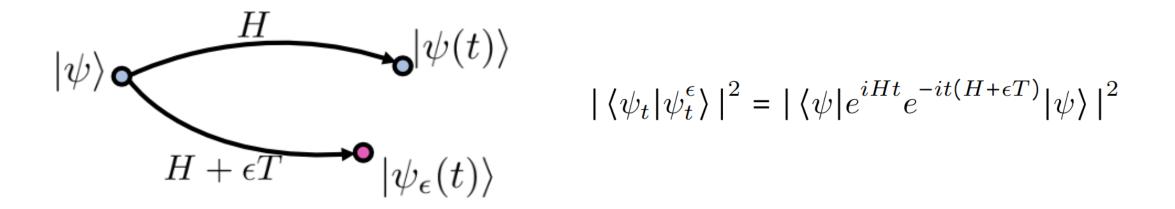


Quantum Chaos and its Causes and Effects



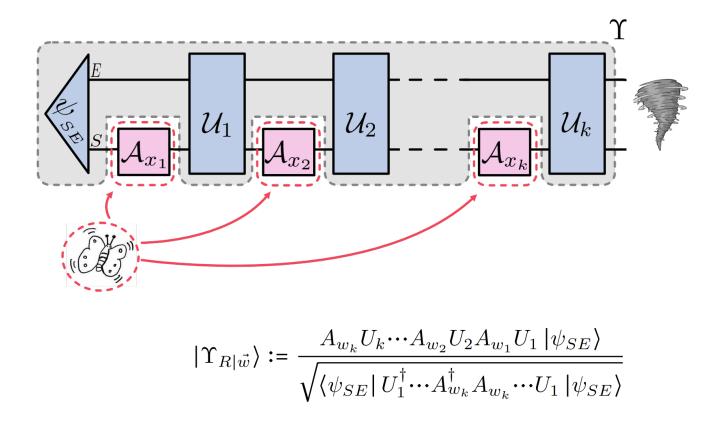
Example: Peres-Loschmidt Echo

 How does a perturbation to the Hamiltonian affect pure state dynamics?

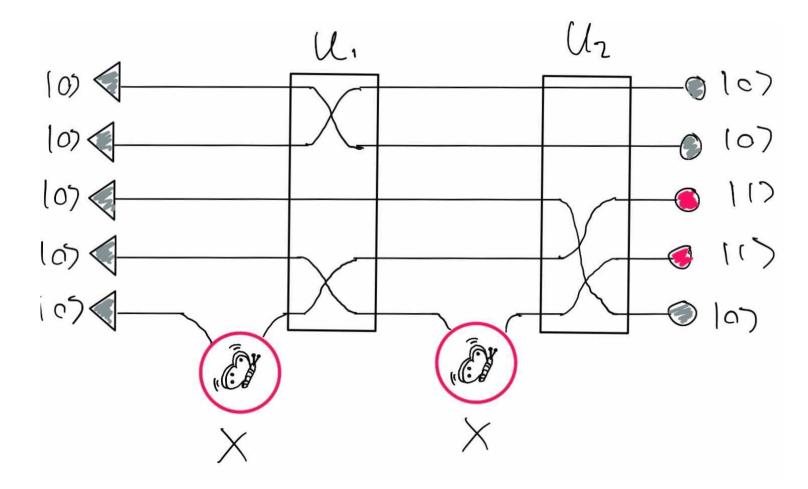


- Conjectured that it decays exponentially for chaotic systems.
- But can show this is not sufficient for chaos.

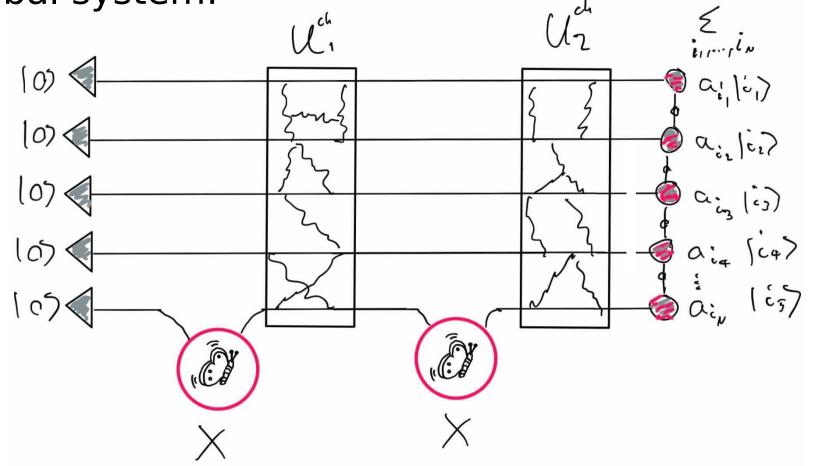
How well does a sequence of local operations orthogonalise the global pure state of a system?



But, a circuit of swaps should not be chaotic:

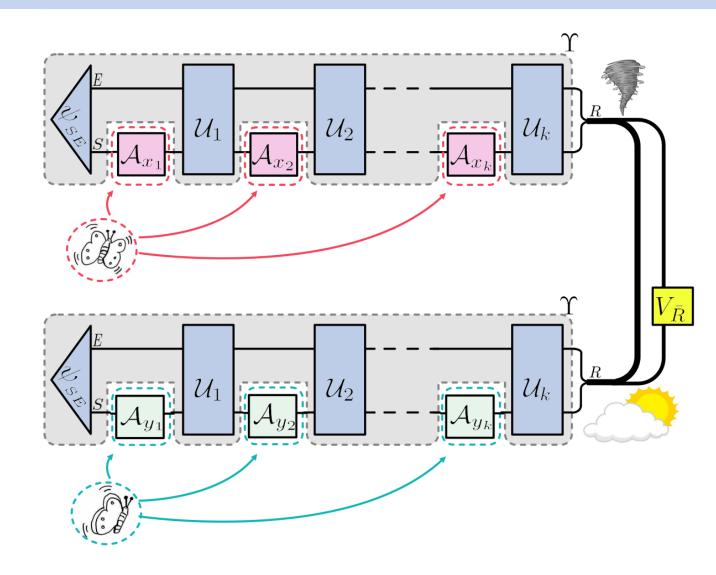


The orthogonality of the perturbation should spread through the global system:



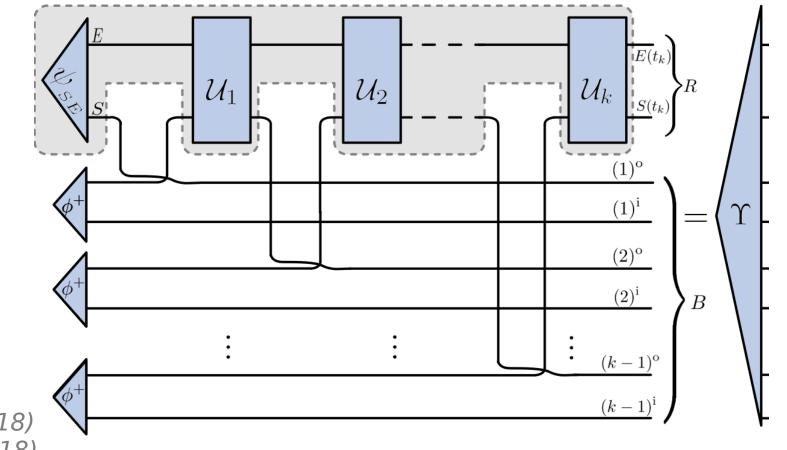
So enforce that this effect is highly non-local – that the process is *scrambling*.

$$\zeta(\Upsilon) := \sup_{\bar{R}, V, \langle x|y \rangle = 0} \left(|\langle \Upsilon_{R|x} | V_{\bar{R}} | \Upsilon_{R|y} \rangle|^2 \right)$$



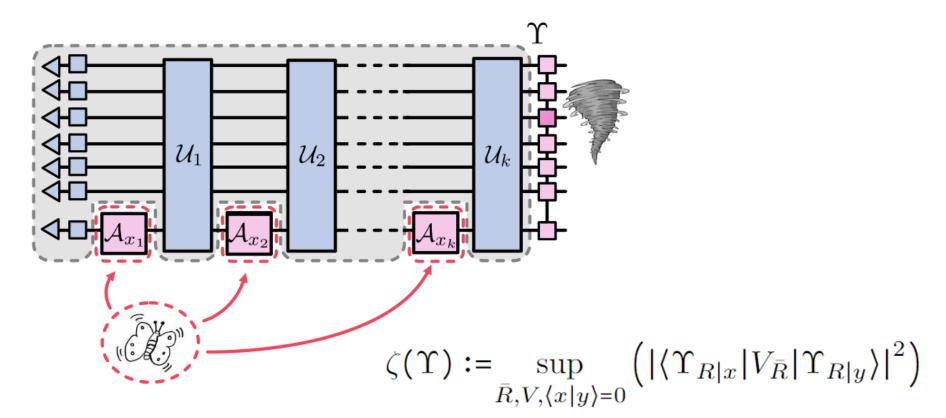
Tool: Process-State Duality

Choi–Jamiołkowski Isomorphism allows one to map any process oneto-one to a quantum state.

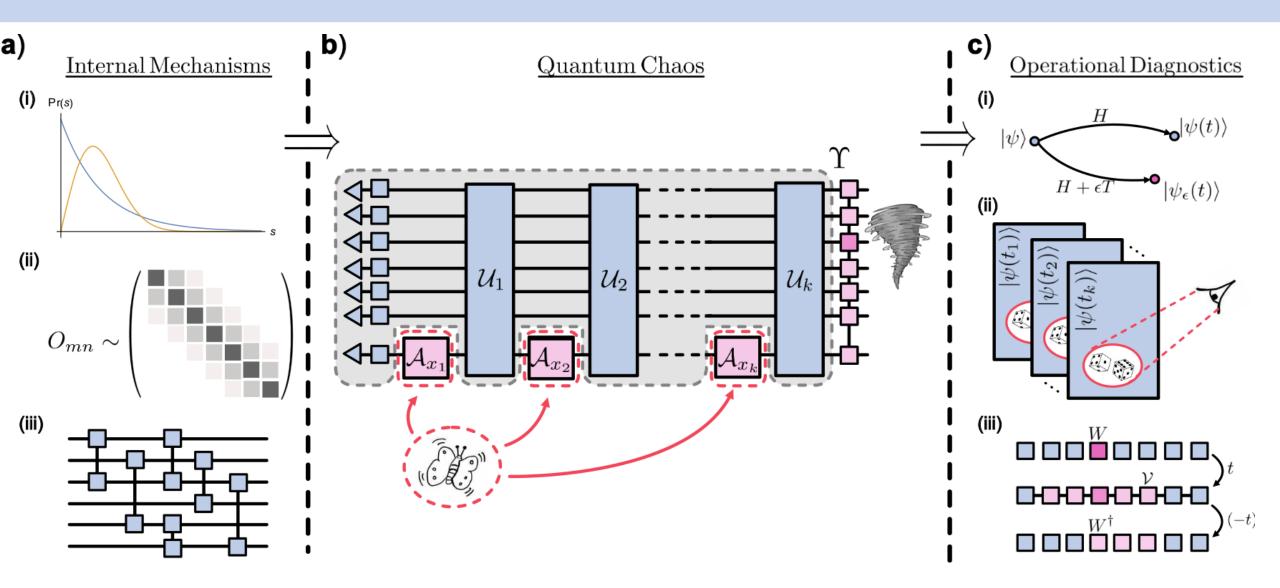


Pollock et al., PRL (2018) Pollock et al., PRA (2018).

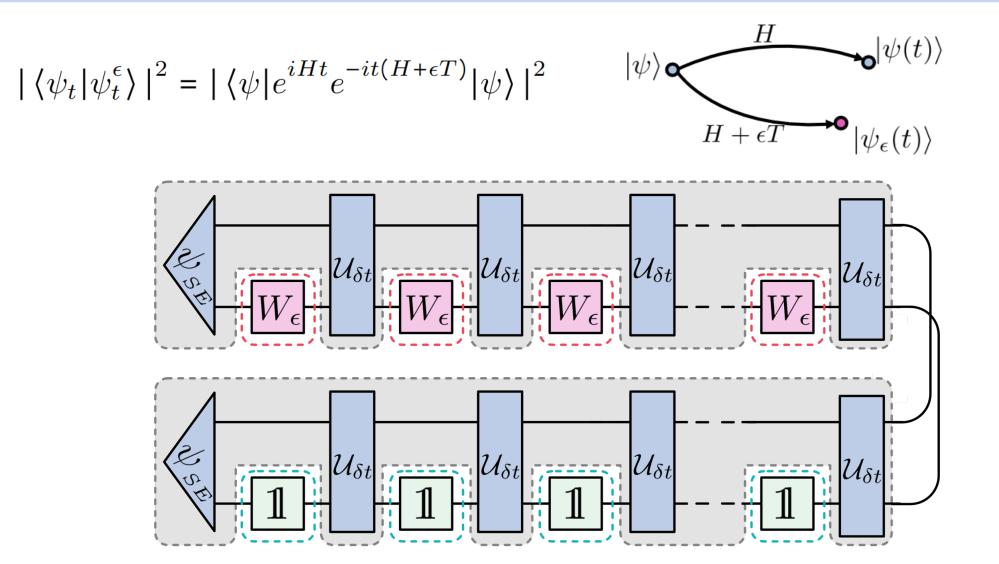
Quantum Chaos = Volume-Law spatiotemporal entanglement



Quantum Chaos and its Causes and Effects

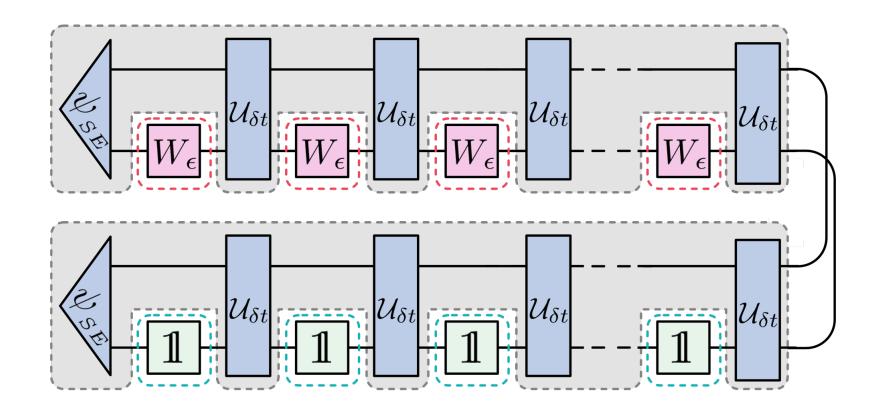


The Peres-Loschmidt Echo

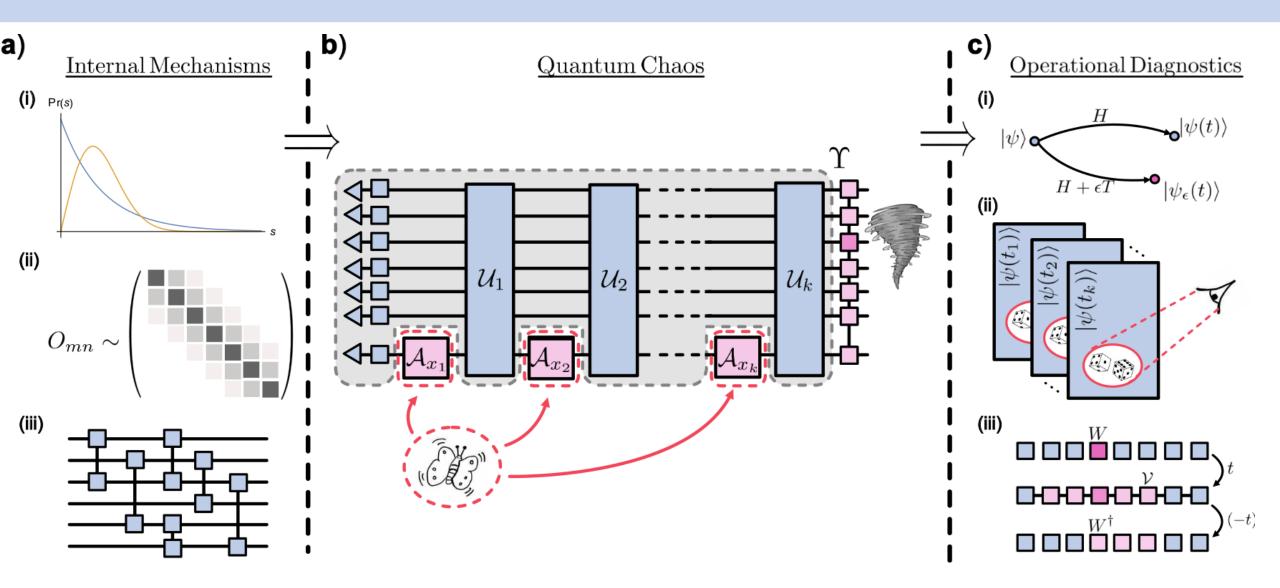


The Peres-Loschmidt Echo

Can show that for very chaotic systems (volume-entangled), this decays exponentially.



Quantum Chaos and its Causes and Effects



Mechanisms of Chaos

Random Dynamics are highly likely to generate a chaotic process.

- Exponentially so for Haar random unitary evolution.
- Polynomially so for t-designs.

$$\mathrm{tr}_{R}[\Upsilon_{\mathrm{H}}] :\approx \frac{\mathbbm{1}}{d_{S}^{2k}}, \text{ for } d_{B} \ll d_{R}$$

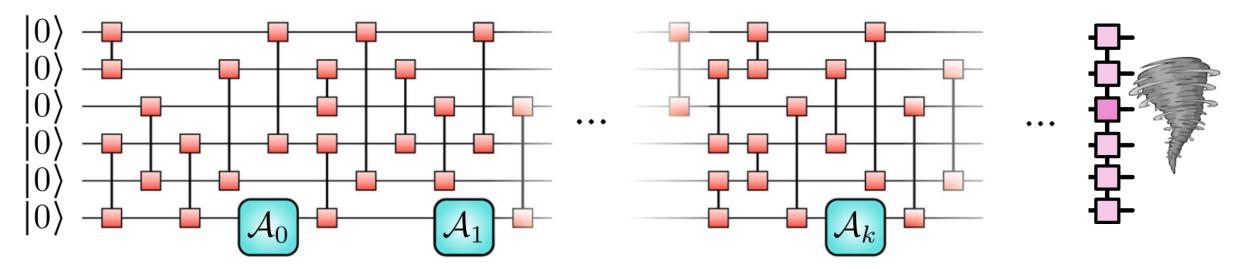


Fig. From: Figueroa-Romero et al., Nature Comm. Phys. (2021).

Summary and Applications

- We have constructed a quantum butterfly effect from a theoryindependent principle of chaos.
- This allows us to identify chaos as the volume-law spatiotemporal entanglement property of a quantum process.
- Other operational diagnostics derive from this, and we show many mechanisms lead to it.
- Can systematically approach problems in many body physics.
- Other mechanisms of chaos: Random Matrix Theory?
- Emergence of a quantum Lyapunov exponent?

Thank you!

The Peres-Loschmidt Echo: decay

$$\zeta'(\Upsilon^{(\mathrm{H})}, \mathcal{X}_{LE}) \approx \frac{|\langle \vec{x} | 1/d_B | \vec{y} \rangle|^2}{\langle \vec{x} | 1/d_B | \vec{x} \rangle \langle \vec{y} | 1/d_B | \vec{y} \rangle}$$

$$= (1 - \epsilon)^{2k} d_B^2 (1/d_B^2)$$

$$= (1 - \epsilon)^{2k}$$

$$\approx e^{-2k\epsilon}, \text{ for small } \epsilon,$$

$$\approx 0, \text{ for large } k.$$

Mechanisms of Chaos

Random Dynamics are highly likely to generate a chaotic process.

$$\mathbb{P}_{U_i \sim \mu} \left\{ \tilde{\zeta}(\Upsilon, \mathcal{X}_{\mathrm{b}}) \geq \frac{\mathcal{J}_{\mu}(\delta)}{(d_{B\bar{R}} - 1)(1 - \sqrt{\mathcal{J}_{\mu}(\delta)})^2} \right\} \leq \mathcal{G}_{\mu}(\delta).$$

Haar evolution:

ε-approximate t-deisgn:

$$\mathcal{J}_{\mathrm{H}}(\delta) = d_{B\bar{R}}(\mathcal{B} + \delta) \approx d_{B\bar{R}}(\frac{1}{d_{R}} + \delta), \text{ and}$$
$$\mathcal{G}_{\mathrm{H}}(\delta) = \exp[-\mathcal{C}\delta^{2}] \approx \exp[-\frac{(k+1)d_{R}}{8d_{B}}\delta^{2}],$$

$$\mathcal{J}_{\mu_{\epsilon-t}}(\delta) = d_{B\bar{R}}\delta, \text{ and}$$
$$\mathcal{G}_{\mu_{\epsilon-t}}(\delta) = \frac{\mathcal{F}(d_B, d_R, m, t, \epsilon)}{\delta^m}$$