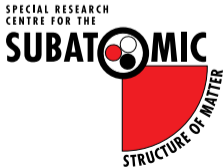


# Testing the Quark Model on the $\Delta$ Baryon Spectrum

Liam Hockley, Curtis Abell, Waseem Kamleh, Derek Leinweber & Anthony Thomas

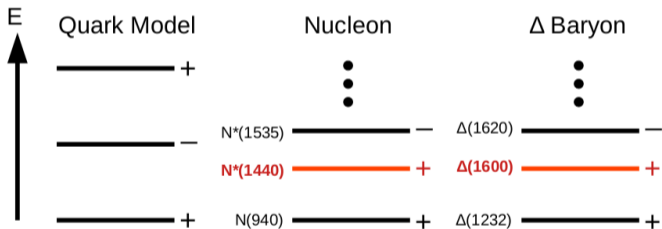
Australian Institute of Physics Congress

13 Dec 2022



# The Quark Model

- Baryons modelled as 3-quark systems; good for understanding the “particle zoo”
- BUT this simple model doesn't produce the entire spectrum of baryons in nature (i.e. the Roper puzzle)



**Q:** Is the  $\Delta(1600)$  resonance *really* a 3-quark state?

# Overview of our methods

---

Our two methods are Lattice QCD and Hamiltonian Effective Field Theory (HEFT):

- Lattice QCD allows us to compute baryon masses from first principles **on a finite volume**
- We can then construct an **infinite volume** Hamiltonian and compute various scattering observables to describe experiment
- The final step is to convert this Hamiltonian to **finite volume** so we can compare experiment with the lattice results

**HEFT allows us to simultaneously compare with results on the finite volume of the lattice and with the infinite volume of the real world.**

# Lattice QCD

---

Lattice QCD is a well established method for calculating hadronic observables in the non-perturbative regime of QCD

We use standard techniques to calculate baryon masses:

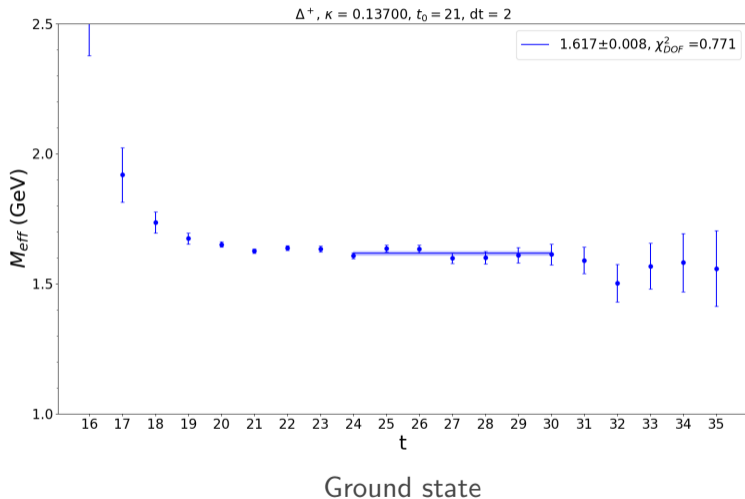
- **3-quark interpolating fields with quantum numbers of the  $\Delta$**
- Positive parity and spin-3/2 projection operators
- 4 by 4 correlation matrix of smearings extends the operator basis
- Compute two-point correlation functions and solve a GEVP to get effective masses

Gauge field ensembles used are from the PACS-CS collaboration

S. Aoki et al. [PACS-CS], Phys. Rev. D 79 (2009) 034503 [arXiv:0807.1661 [hep-lat]]

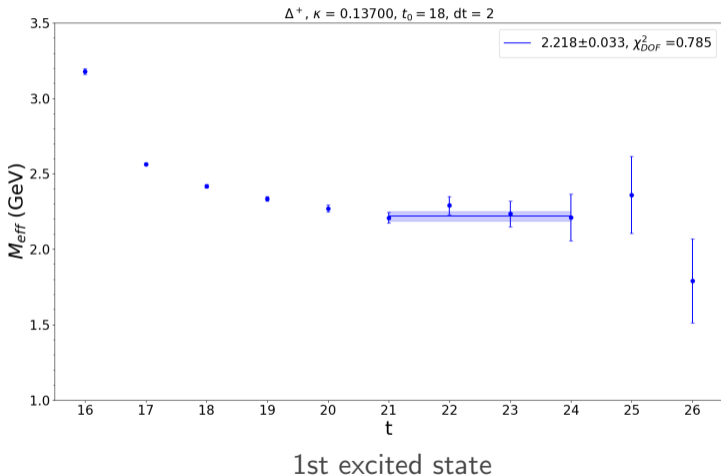
# Effective Mass Fits (preliminary)

- Extract a single projected “effective” mass
- Fit a plateau (using a full covariance matrix) across multiple time slices
- Repeat across all 5 gauge field ensembles for both ground and first excited states

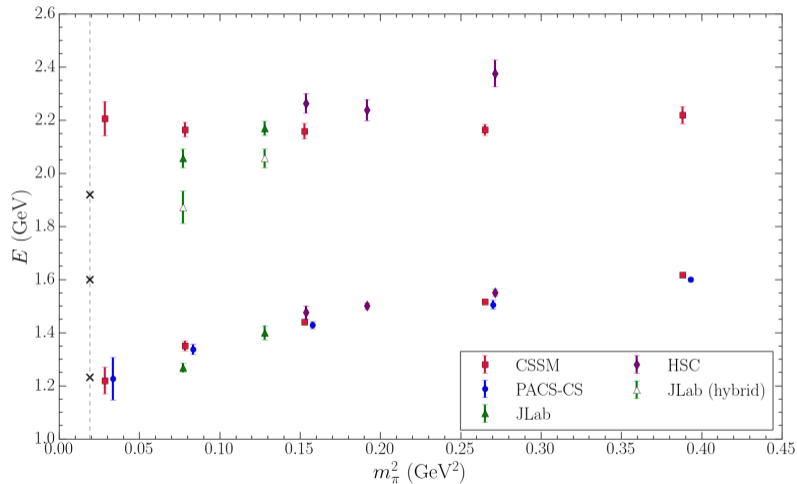


# Effective Mass Fits (preliminary)

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# Lattice Results (preliminary)



## PACS-CS:

S. Aoki et al. [PACS-CS], Phys. Rev. D 79 (2009) 034503 [arXiv:0807.1661 [hep-lat]]

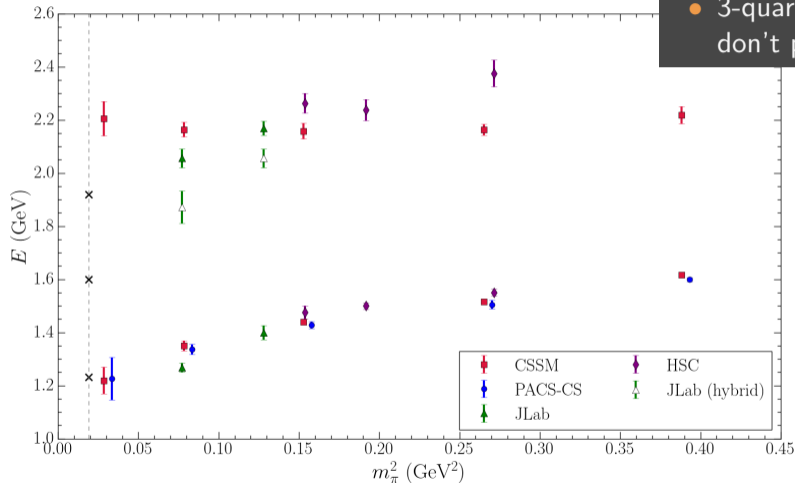
## JLab:

T. Khan, D. Richards and F. Winter, Phys. Rev. D 104 (2021) 034503 [arXiv:2010.03052 [hep-lat]]

## HSC:

J. Bulava, et al., Phys. Rev. D 82 (2010) 014507 [arXiv:1004.5072 [hep-lat]]

# Lattice Results (preliminary)



3-quark interpolating operators don't produce  $\Delta(1600)$

## PACS-CS:

S. Aoki et al. [PACS-CS], Phys. Rev. D 79 (2009) 034503 [arXiv:0807.1661 [hep-lat]]

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# Infinite Volume Hamiltonian

---

Construct an infinite volume Hamiltonian  $H = H_0 + H_I$  with basis states composed of:

- quark model-like, 3-quark states at rest
- non-interacting two-particle meson-baryon states, with some back-to-back momentum

**Non-interacting Hamiltonian:**

$$H_0 = \sum_{B_0} |B_0\rangle m_{B_0} \langle B_0| + \sum_{\alpha} \int d^3k |\alpha(\vec{k})\rangle \omega_{\alpha}(\vec{k}) \langle \alpha(\vec{k})|$$

$|B_0\rangle$  : quark model state       $|\alpha(\vec{k})\rangle$  : 2-particle state

$$\omega_{\alpha}(\vec{k}) \equiv \sqrt{m_{\alpha_B}^2 + \vec{k}^2} + \sqrt{m_{\alpha_M}^2 + \vec{k}^2}$$

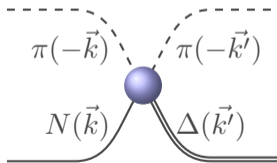
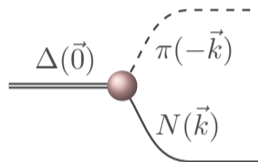
# Infinite Volume Hamiltonian

Interacting Hamiltonian:

$$H_I = g + v$$

$$g = \sum_{\alpha, B_0} \int d^3k \left\{ |\alpha(\vec{k})\rangle G_{\alpha, B_0}^\dagger(k) \langle B_0| + h.c. \right\}$$

$$v = \sum_{\alpha, \beta} \int d^3k d^3k' |\alpha(\vec{k})\rangle V_{\alpha, \beta}(k, k') \langle \beta(\vec{k}')|$$



# Functional Forms

Basis states:

- Quark model-like (motivated by two states on the lattice):  $|\Delta^{(1)}\rangle$ ,  $|\Delta^{(2)}\rangle$
- Meson-baryon (motivated by PDG branching ratios):  
 $|\pi N(\vec{k})\rangle$ ,  $|\pi\Delta(\vec{k})\rangle$  in p-wave  $\equiv |\pi\Delta_p(\vec{k})\rangle$ ,  $|\pi\Delta(\vec{k})\rangle$  in f-wave  $\equiv |\pi\Delta_f(\vec{k})\rangle$

The couplings and potentials are derived from chiral EFT i.e.

$$G_{\pi N}^{\Delta}(k) = \frac{g_{\pi N}^{\Delta}}{m_{\pi}} \frac{k}{\sqrt{\omega_{\pi}(k)}} u(k, \Lambda_{\pi N}) \quad \text{Dipole regulator: } u(k, \Lambda) = (1 + k^2/\Lambda^2)^{-2}$$

$$V_{\pi N, \pi N}(k, k') = \frac{v_{\pi N, \pi N}}{(m_{\pi})^2} \frac{k}{\omega_{\pi}(k)} \frac{k'}{\omega_{\pi}(k')} u(k, \Lambda_{\pi N}^v) u(k', \Lambda_{\pi N}^v)$$

- Need to constrain the free parameters  $m_{\Delta}^{(i)}$ ,  $g_{\alpha}^{(i)}$ ,  $v_{\alpha, \beta}$ ,  $\Lambda_{\alpha}$ ,  $\Lambda_{\alpha}^v$ .

# T-matrix

---

Solve for the T-matrix using the coupled channel integral equations:

$$T_{\alpha\beta}(k, k'; E) = \tilde{V}_{\alpha\beta}(k, k'; E) + \sum_{\gamma} \int dq q^2 \frac{\tilde{V}_{\alpha\gamma}(k, q, E) T_{\gamma\beta}(q, k'; E)}{E - \omega_{\gamma}(q) + i\epsilon}$$

where we've defined the coupled channel potential

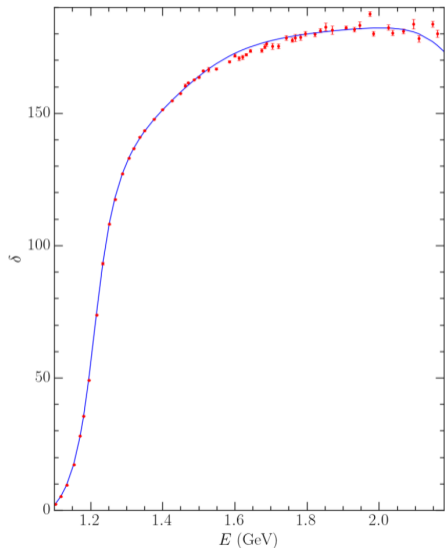
$$\tilde{V}_{\alpha\beta}(k, k', E) = \frac{G_{\alpha}^{B_0\dagger}(k) G_{\beta}^{B_0}(k')}{E - m_{B_0}} + V_{\alpha\beta}(k, k')$$

Extract scattering phase shifts and inelasticities via the S-matrix:

$$S_{\pi N, \pi N}(E) = 1 - 2i\pi \frac{\omega_{\pi}(k_{\text{cm}}) \omega_N(k_{\text{cm}})}{E} k_{\text{cm}} T_{\pi N, \pi N}(k_{\text{cm}}, k_{\text{cm}}; E) = \eta(E) e^{2i\delta(E)}$$

where  $k_{\text{cm}}$  satisfies  $\omega_{\alpha}(k_{\text{cm}}) = \omega_{\alpha_M}(k_{\text{cm}}) + \omega_{\alpha_B}(k_{\text{cm}}) = E$

# Fit to scattering data (preliminary)



- Fit phase shifts and inelasticities to  $\pi N \rightarrow \pi N$  scattering data available from GWU SAID partial wave analysis
- Extract resonance poles to check against the PDG

Resonance	PDG Pole	This Work
$\Delta(1232)$	$1.210(1) - 0.050(1)i$	$1.211 - 0.049i$
$\Delta(1600)$	$1.510(50) - 0.135(35)i$	$1.444 - 0.219i$
$\Delta(1920)$	$1.900(50) - 0.150(100)i$	$2.262 - 0.185i$

# Finite Volume Hamiltonian

- Cube with side length  $L$  with periodic b.c.  $\implies$  discrete momenta

$$k_n = \frac{2\pi}{L} \sqrt{n}, \quad n = n_x^2 + n_y^2 + n_z^2, \quad C_3(n) \text{ counts degeneracy in } n$$

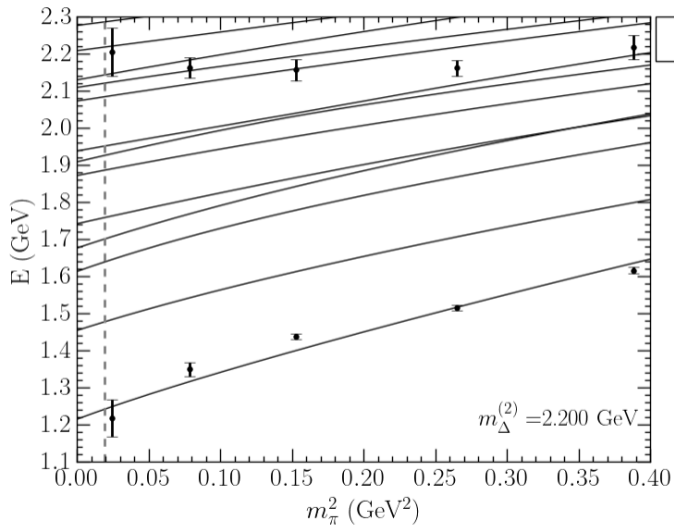
$$\int k^2 dk = \frac{1}{4\pi} \int d^3k \rightarrow \frac{1}{4\pi} \sum_{\vec{n} \in \mathbb{Z}^3} \left(\frac{2\pi}{L}\right)^3 \rightarrow \frac{1}{4\pi} \sum_n C_3(n) \left(\frac{2\pi}{L}\right)^3$$

- Finite volume corrections

$$\bar{G}_\alpha^{B_0}(k_n) = \sqrt{\frac{C_3(n)}{4\pi}} \left(\frac{2\pi}{L}\right)^{3/2} G_\alpha^{B_0}(k_n)$$
$$\bar{V}_{\alpha\beta}(k_n, k_m) = \sqrt{\frac{C_3(n)}{4\pi}} \sqrt{\frac{C_3(m)}{4\pi}} \left(\frac{2\pi}{L}\right)^3 V_{\alpha\beta}(k_n, k_m)$$

- Can solve the now finite volume Hamiltonian matrix for its eigenvalues and eigenvectors

# Finite Volume Energy Spectrum (preliminary)



- Solve the Hamiltonian for eigenvalues and eigenvectors using LAPACK numerical routines
- Eigenvalues will sit at points along the physical pion mass (dashed line)
- Extrapolate masses of hadrons in  $m_\pi^2$ , obtain a finite volume energy spectrum in  $m_\pi^2$

# Eigenvector Analysis

---

Eigenvectors are powerful since they tell us about the composition of energy eigenstates in terms of our basis states

Consider the states which have largest overlaps with:

- 3-quark basis states  $|\Delta^{(1)}\rangle$ ,  $|\Delta^{(2)}\rangle$
- 2-particle basis states  $|\pi N(k_n)\rangle$ ,  $|\pi\Delta_p(k_n)\rangle$ ,  $|\pi\Delta_f(k_n)\rangle$

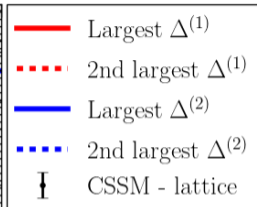
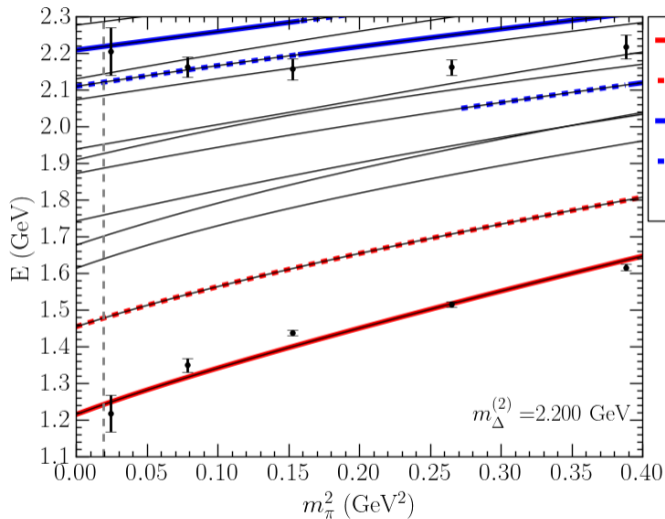
In particular, the overlaps  $\langle\Delta^{(1)}|E_i\rangle$  and  $\langle\Delta^{(2)}|E_i\rangle$  tell us how dominant the quark model-like basis states are in the finite volume energy eigenstate  $|E_i\rangle$

**Hamiltonian eigenstates with the largest overlaps with the 3-quark states should match the lattice QCD results**

- Lattice QCD provides a second type of constraint for the HEFT results

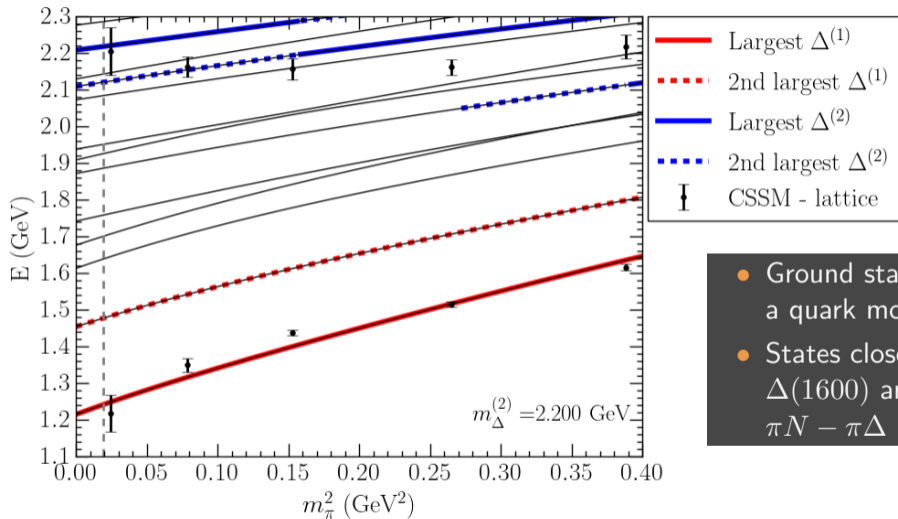


# Comparison with Lattice Data (preliminary)



- Dressed lines show the states with largest overlaps with the 3-quark basis states
- Fit dressed lines to the lattice data
- Search for a scenario which simultaneously describes experimental and lattice data

# Comparison with Lattice Data (preliminary)



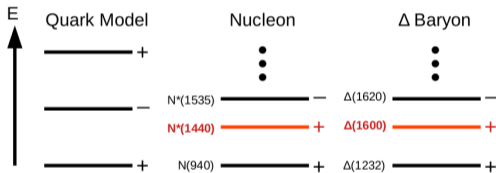
- Ground state dominated by a quark model-like state
- States close to the  $\Delta(1600)$  are predominantly  $\pi N - \pi \Delta$

# Conclusions

HEFT provides a bridge between finite volume lattice results and infinite volume scattering data

LQCD + HEFT results suggest the following:

- $\Delta(1232)$  predominantly composed of a quark model-like state
- $\Delta(1600)$  is predominantly a dynamical resonance generated through  $\pi N$  and  $\pi\Delta$  scattering

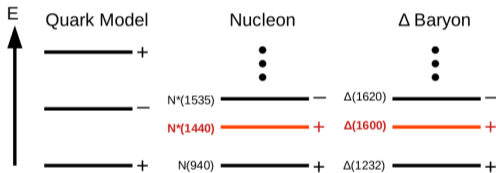


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- $\Delta(1232)$  predominantly composed of a quark model-like state
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**A:** The  $\Delta(1600)$  doesn't appear to be a simple radial excitation of the  $\Delta(1232)$ , it appears to be dynamically generated

# Backup Slides

# Lattice QCD - Gauge field configurations

---

- Well established method for first principles calculations of hadronic observables in the non-perturbative regime of QCD
- Note: we discretise space-time onto a 4D grid  $\implies$  **finite volume**

Gauge field configurations provided by the PACS-CS collaboration

- 2 + 1 flavour, dynamical-fermion configurations
- $\mathcal{O}(\alpha)$  improved Wilson fermion action + Iwasaki-gauge action
- $32^3 \times 64$  lattice, periodic boundary conditions,  $\beta = 1.90$
- 5 ensembles:  $m_\pi = 702, 572, 413, 293, 156$  MeV

S. Aoki et al. [PACS-CS], Phys. Rev. D 79 (2009) 034503 [arXiv:0807.1661 [hep-lat]]

# Lattice Ensembles

---

$\kappa$	$m_\pi$ (MeV)	$a$ (fm)	Number configs
0.13700	702	0.1023(15)	399
0.13727	572	0.1009(15)	397
0.13754	413	0.0961(13)	449
0.13770	293	0.0951(13)	400
0.13781	156	0.0933(13)	198*

\*Lightest ensemble needs additional statistics

# Hadron Spectroscopy on the Lattice

---

$$G_{ij}(t, \vec{p}) = \sum_{\vec{x}} e^{-i\vec{p}\cdot\vec{x}} \langle \Omega | T \{ \chi_i(x) \bar{\chi}_j(0) \} | \Omega \rangle = \sum_{B^\pm} \lambda_i^\pm \bar{\lambda}_j^\pm e^{-m_{B^\pm} t}$$

- $\chi_i$  are 3-quark interpolating fields coupling to  $\Delta$  baryons
- Apply positive parity and spin-3/2 projection operators  
Target  $J^P = \frac{3}{2}^+$   $\Delta$  baryons:  $\Delta(1232)$ ,  $\Delta(1600)$ ,  $\Delta(1920)$
- Use a **variational method** to project excited states

$$[(G(t_0))^{-1} G(t_0 + \Delta t)] u^\alpha = e^{-m_\alpha \Delta t} u^\alpha, \quad v^\alpha [G(t_0 + \Delta t) (G(t_0))^{-1}] = e^{-m_\alpha \Delta t} v^\alpha$$

$$\implies M_{\text{eff}}^\alpha(t) = \ln \left( \frac{v^\alpha G(t, \vec{p} = 0) u^\alpha}{v^\alpha G(t + 1, \vec{p} = 0) u^\alpha} \right), \quad \alpha = \text{energy level}$$



# Spin Projection Operators

---

$$P_{mn}^{3/2}(\vec{p}) = g_{\mu\nu} - \frac{1}{3}\gamma_\mu\gamma_\nu - \frac{1}{3p^2}(\gamma \cdot p\gamma_\mu p_\nu + p_\mu\gamma_\nu\gamma \cdot p)$$

At rest  $\vec{p} = 0$ :

$$P_{mn}^{3/2}(0) = \delta_{mn}\mathbb{I} - \frac{1}{3}\gamma_m^S\gamma_n^S$$

where  $S$  denotes the Sakurai representation:

$$\gamma_S^4 = \gamma_{\text{Dirac}}^0, \quad \gamma_S^i = -i\gamma_{\text{Dirac}}^i$$

Spin projected correlator:

$$G_{\mu\nu}^{3/2}(\vec{p} = 0) = \sum_{\sigma,\rho=1}^4 G_{\mu\sigma}g^{\sigma\rho}P_{\rho\nu}^{3/2}(0)$$

# Chiral Extrapolation

---

Expand the bare  $\Delta$  masses about the physical pion mass:

$$m_{\Delta}^{(i)} = m_{\Delta}^{(i)}|_{\text{phys}} + \alpha_i \left( m_{\pi}^2 - m_{\pi}^2|_{\text{phys}} \right) \quad i = 1, 2$$

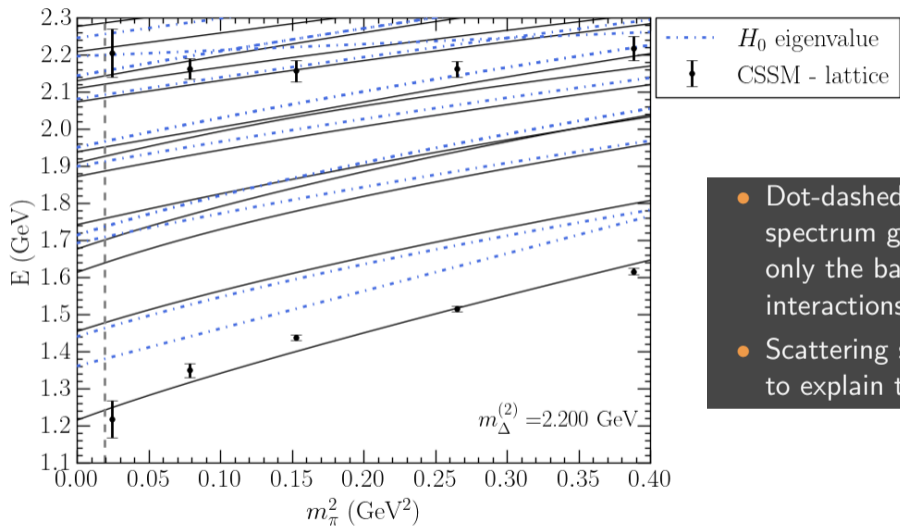
Likewise

$$m_B = m_B|_{\text{phys}} + \alpha_B \left( m_{\pi}^2 - m_{\pi}^2|_{\text{phys}} \right) \quad B = N, \Delta$$

Fix the lattice size at  $L = 3$  fm and solve the matrix eigenvalue problem for unphysical values of  $m_{\pi}^2$

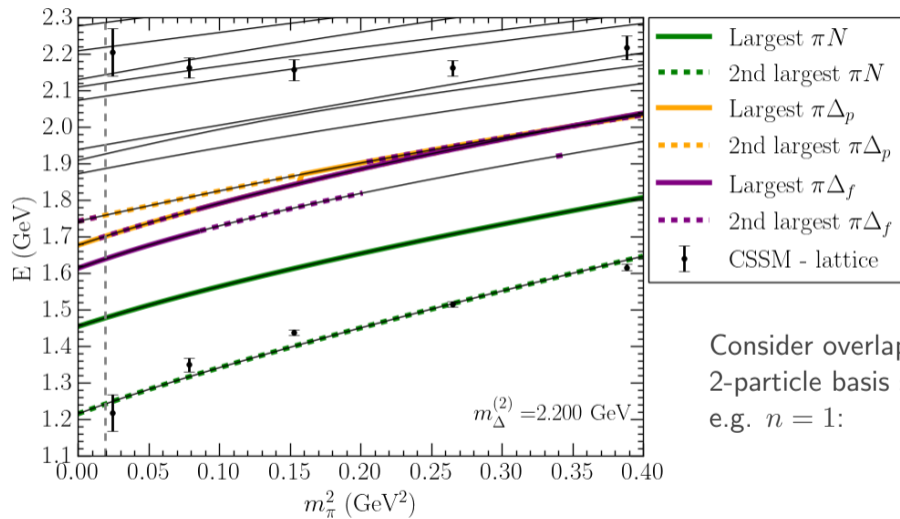
- $\alpha_N$  and  $\alpha_{\Delta}$  have known values from previous lattice studies
- Perform fits to our lattice data to obtain  $\alpha_1$  and  $\alpha_2$

# Non-interacting state spectrum



- Dot-dashed lines show the spectrum generated from only the bare states, i.e. no interactions
- Scattering states necessary to explain the lattice data

# 2-particle dominated states



Consider overlaps with  
2-particle basis states with  
e.g.  $n = 1$ :

# Current Best Fit Parameters

---

$g_{\pi N}^{(1)}$	0.1345		
$g_{\pi\Delta_p}^{(1)}$	0.0965		
$g_{\pi\Delta_f}^{(1)}$	0.0234	$m_{\Delta}^{(1)}$	1.3806
$g_{\pi N}^{(2)}$	0.1209	$m_{\Delta}^{(2)}$	2.2000
$g_{\pi\Delta_p}^{(2)}$	0.1701	$\Lambda_{\pi N}$	0.8350
$g_{\pi\Delta_f}^{(2)}$	0.0755	$\Lambda_{\pi\Delta_p}$	0.8645
$v_{\pi N, \pi N}$	-0.0087	$\Lambda_{\pi\Delta_f}$	0.7932
$v_{\pi N, \pi\Delta_p}$	-0.0739	$\Lambda_{\pi N}^v$	0.6224
$v_{\pi N, \pi\Delta_f}$	-0.0938	$\Lambda_{\pi\Delta_p}^v$	0.7826
$v_{\pi\Delta_p, \pi\Delta_p}$	-0.0199	$\Lambda_{\pi\Delta_f}^v$	0.8302
$v_{\pi\Delta_f, \pi\Delta_p}$	0.0074		
$v_{\pi\Delta_f, \pi\Delta_f}$	-0.0768		