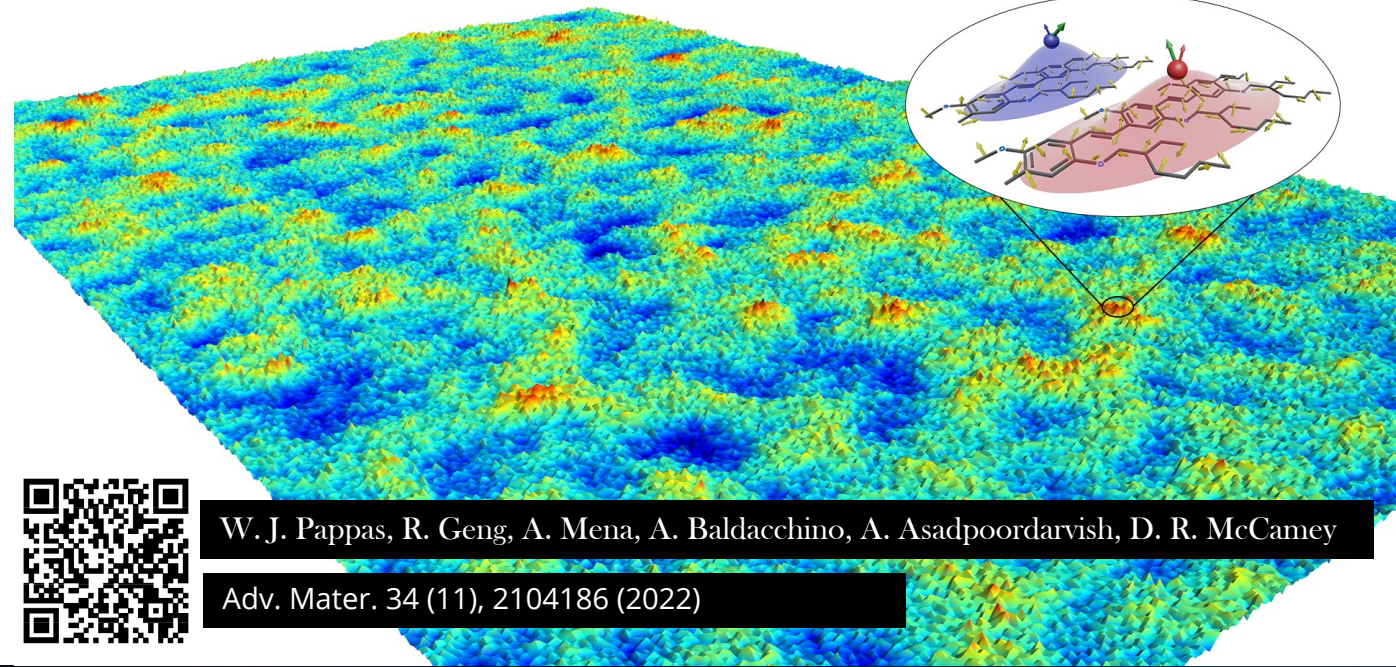


Poster Session II
Thursday 5:30pm



W. J. Pappas, R. Geng, A. Mena, A. Baldacchino, A. Asadpoordarvish, D. R. McCamey

Adv. Mater. 34 (11), 2104186 (2022)

Imaging the Spatial, Temporal
and Energetic Landscape of
Spin Processes in Organic
Optoelectronic Devices

Resolving Spin in Organic Light-Emitting Diodes

William Pappas

McCamey Lab, UNSW Sydney

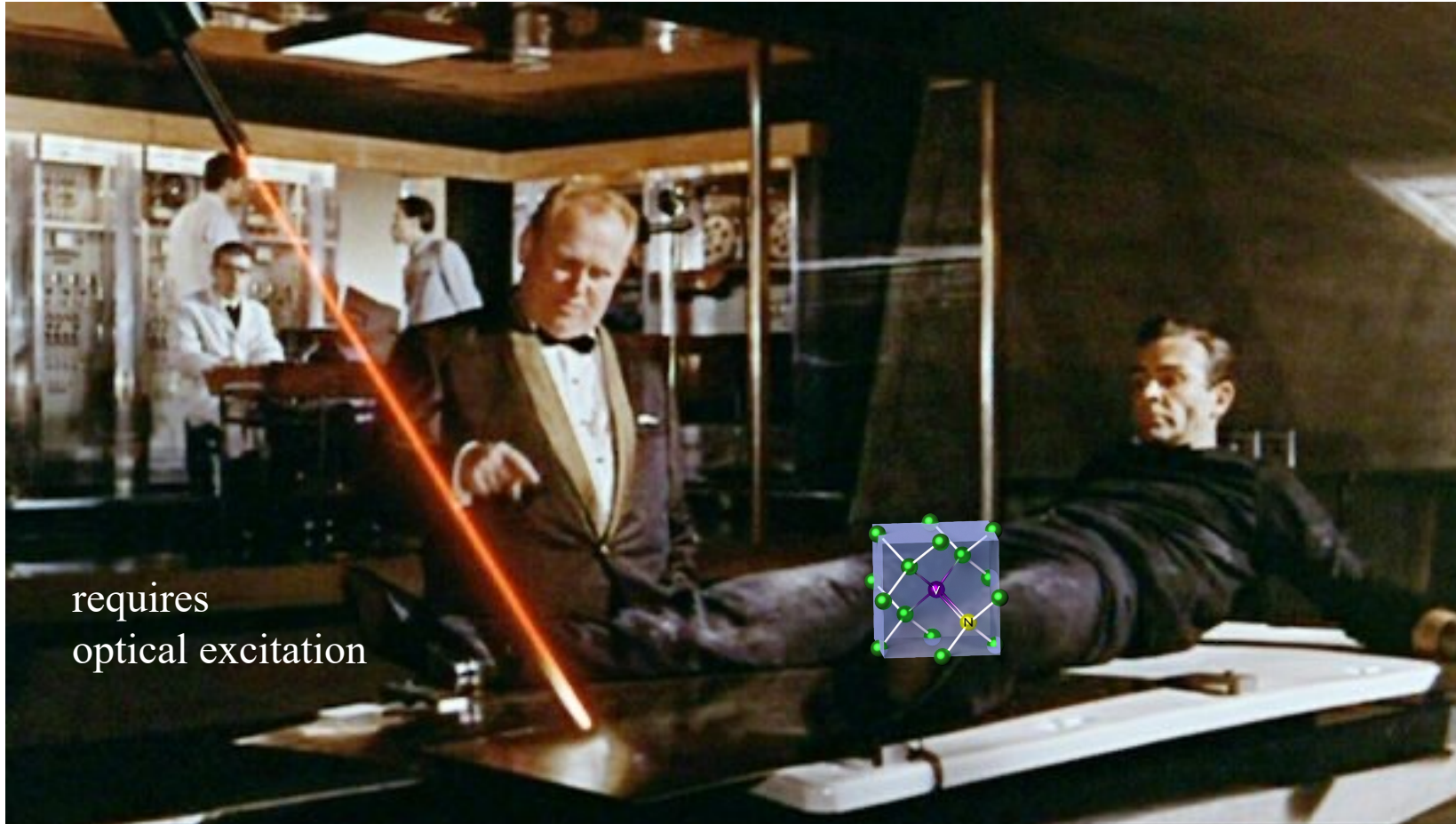
Australian Research Council
Centre of Excellence in



exciton
science

NV Centres are amazing sensors!

high sensitivity, high resolution, versatile, and more...



requires
optical excitation

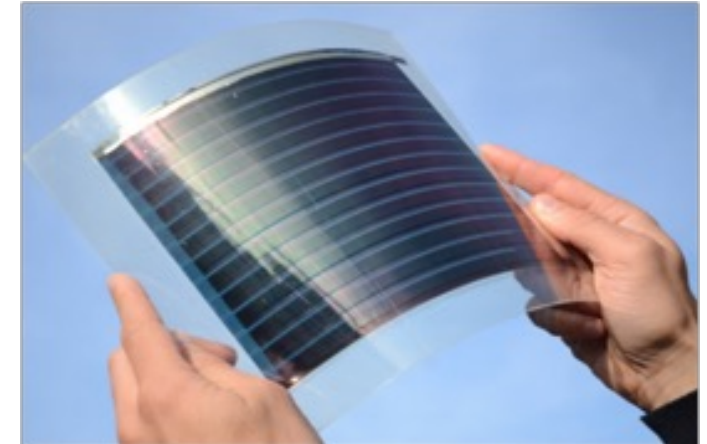
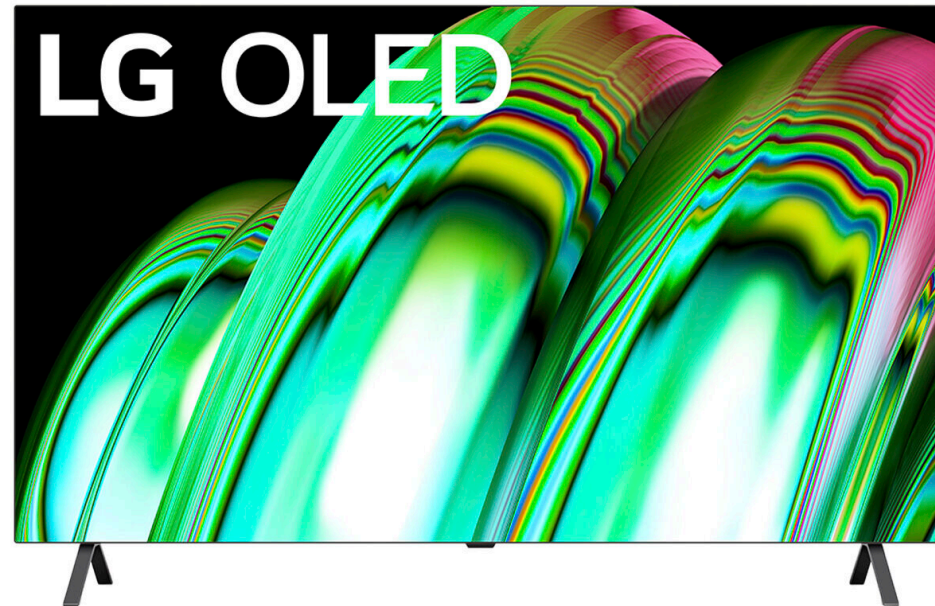
laser limits compactness and integrability, drift, can be expensive, scalability?

Organic Electronics

Commercially available technologies integrating organic electronic materials



© Samsung



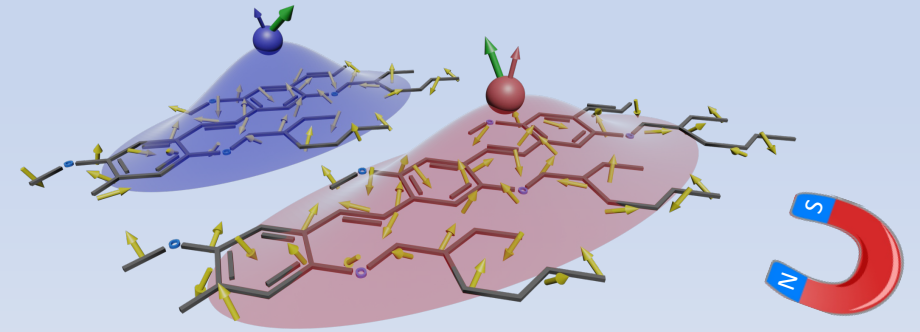
© Eight19

- Electrical operation, laser free
- Cheap, flexible, printable
- Readily modified by synthetic chemistry
- Proven scalability, existing industrial infrastructure

Can spin be used to influence useful properties of these materials?

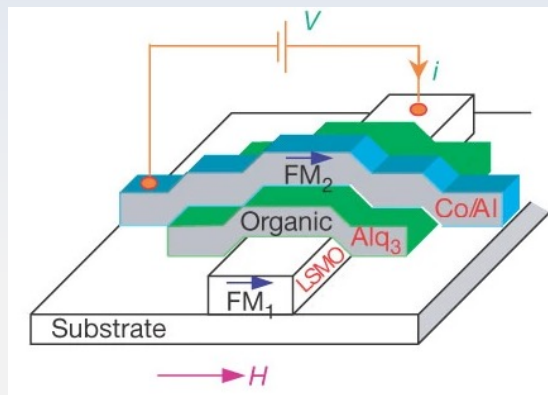
Exploiting the Quantum Properties of Materials in Organic Semiconductors

organic materials exhibit weak spin-orbit coupling – long τ_s and l_s
 spin properties can be tuned using the tools of synthetic chemistry
 understanding **microscopic variation** critical for device performance

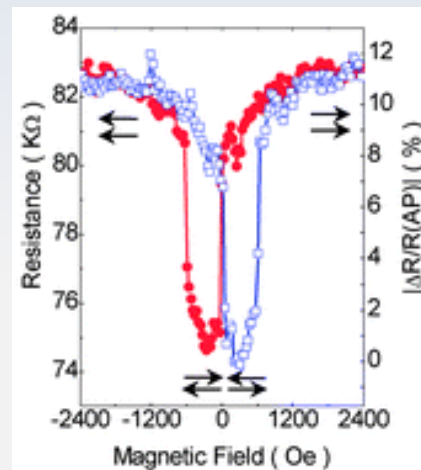


Organic Spintronics Spin Polarisation Effects

Injection, manipulation and transport of spin polarised current

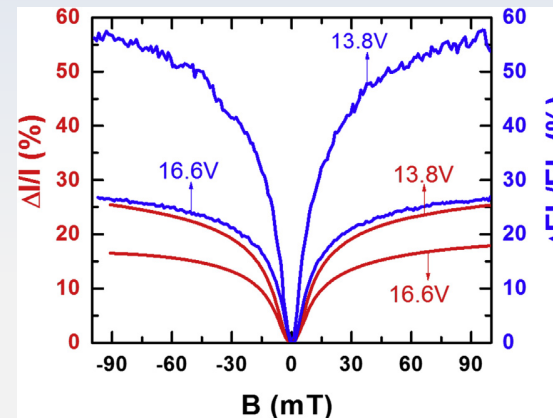


Z. H. Xiong, Di Wu, Z. Vally Vardeny & Jing Shi, *Nature* **427**, 821–824 (2004)

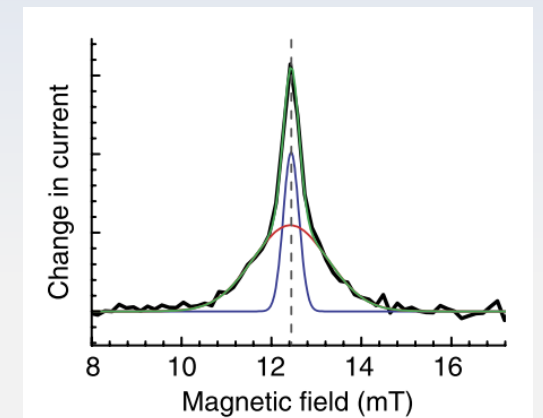


Intrinsic Spin Effects Spin-Mixing Effects

Modulating singlet and triplet excited state populations



T. D. Nguyen *et al.*, *Phys. Rev. B* **77**, 235209 (2008)



W. Baker *et al.*, *Nat Commun* **3**, 898 (2012)

Hybrid Quantum Systems - macroscopic quantum behaviour of molecular systems

Quantum sensing: NV Diamond

- Integrating molecular spins into quantum circuits

D. Kim et al., Nature Electronics, 2 (7), 284–289 (2019)

Collective Spin Modes: Simple Molecular Magnets

- Couple many molecular spins to a single resonant mode

A. W. Eddins et al., Phys. Rev. Lett. 112, 120501 (2014)

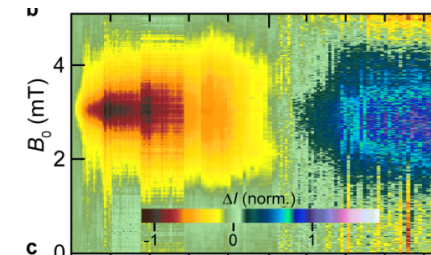
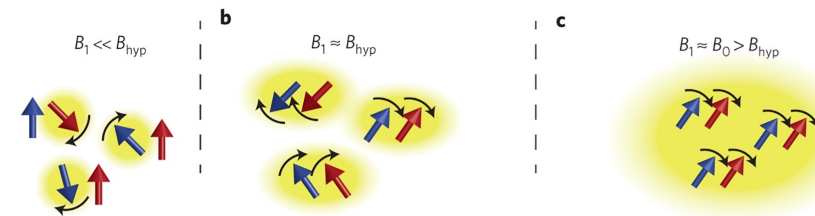
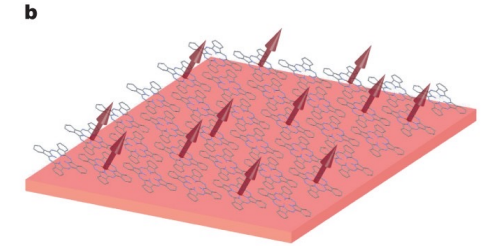
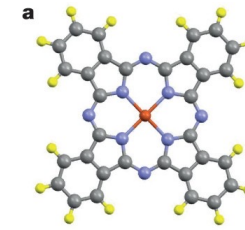
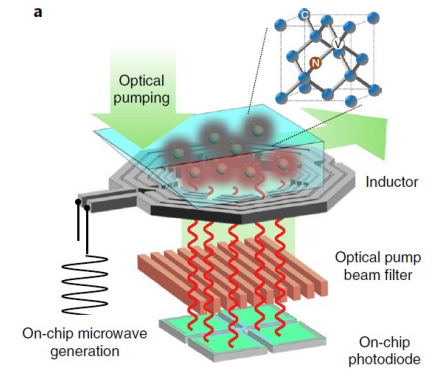
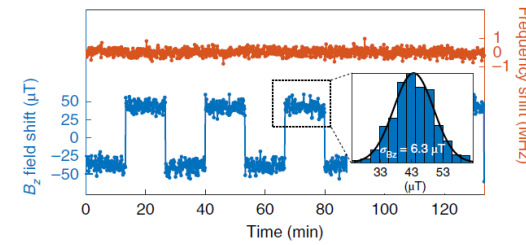
Cooperative Spin Phenomena: Spin-Dicke Effect

- Cooperative Larmor precession of spins under ESR in polymer thin-film OLED

D. P. Waters et al., Nature Physics, 11, 910–914 (2015)

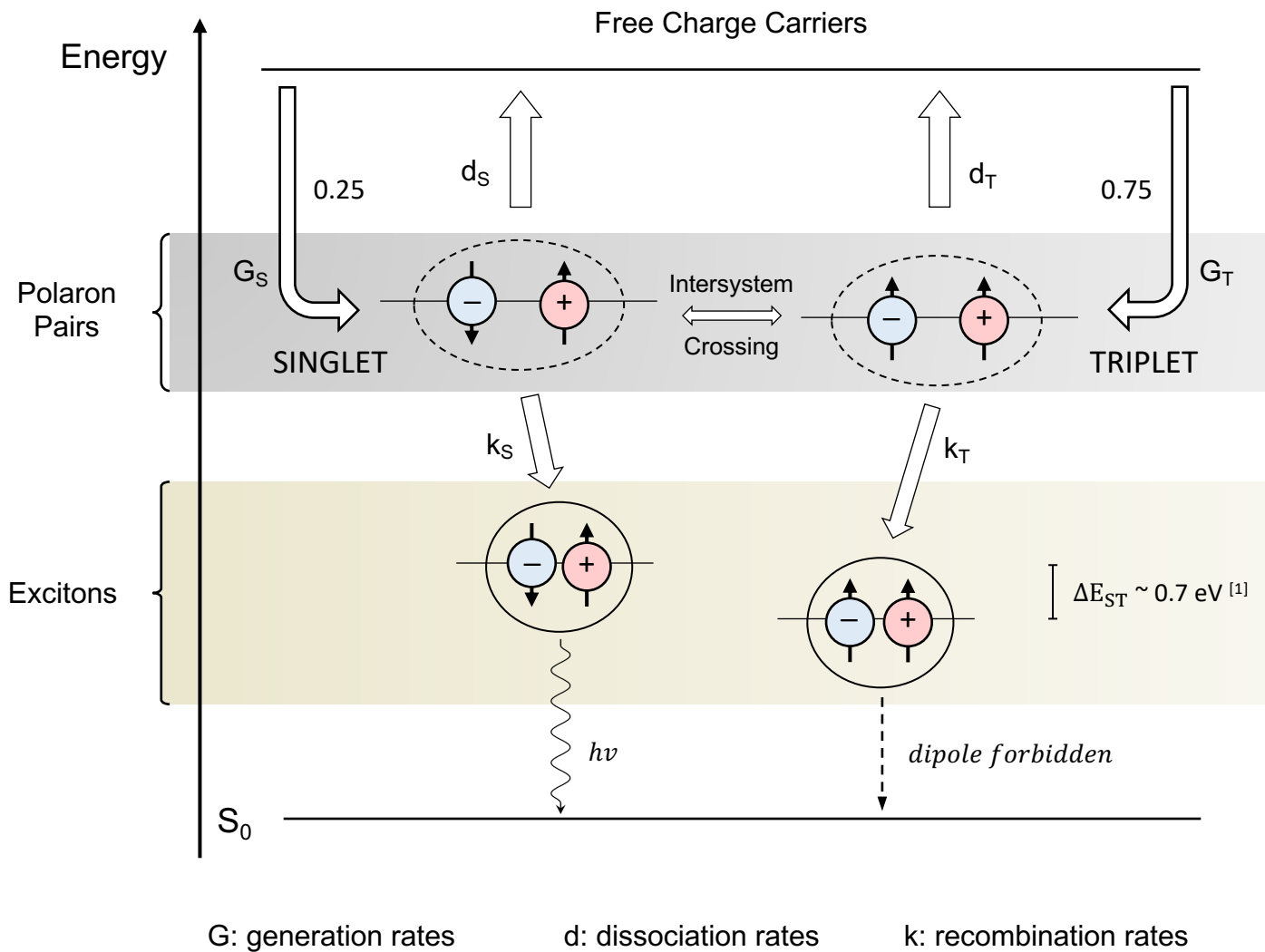
Spin ensembles: deal with **variation** in local spin environment

- distribution in optoelectronic properties
- reproducibility of spin-based organic technologies



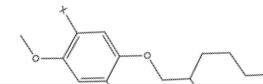
Moving towards the **miniaturisation** and **integration** of organic quantum devices

Spin Permutation Symmetry of Excited States



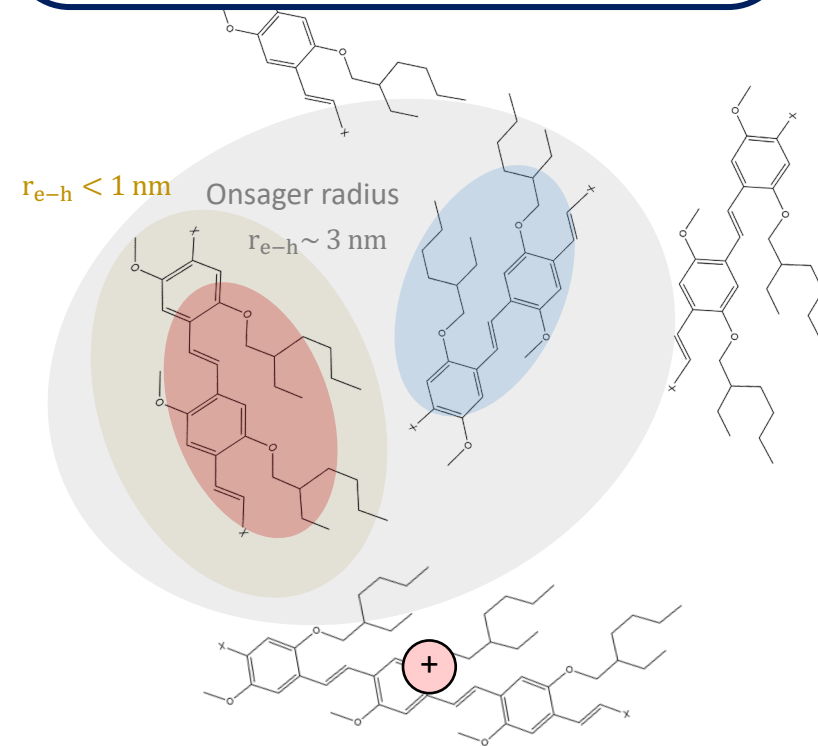
Excitonic Materials

Small-Molecules & Polymer Chromophores
Charge Transport: Thermally Activated Hopping



Intersystem Crossing

- I. Conservation of Energy
- II. Spin Flipping Mechanism



[1] M. Reufer et al., *Nature Mater.* 4, 340 (2005)

Singlet-Triplet Interconversion (Spin Mixing)

Spin Hamiltonian: $H = H_Z + H_{ex} + H_{HF} + H_D + H_{SOC}$

[Distant Pairs]

Interactions

Zeeman: $H_Z = \frac{1}{2} \mu_B B_0 [g_a \hat{\sigma}_z^a + g_b \hat{\sigma}_z^b]$

Exchange: $H_{ex} = -J \hat{S}_a \cdot \hat{S}_b$

Hyperfine: $H_{HF} = \sum_{i=1}^2 \sum_{j=1}^{Ni} [S_i \cdot \tilde{A}_{ij} \cdot I_j]$

Dipolar: $H_D = \hat{S}_a^\dagger \tilde{D}_d \hat{S}_b \quad (\tilde{D}_d)_{ab} \sim r^{-3}$

SOC: $H_{SOC} = \hat{S}_a^\dagger \tilde{D} \hat{S}_b \quad \tilde{D} \sim Z^4$

Singlet-Triplet Interconversion (Spin Mixing)

Traditional Fluorescent Molecules

Spin Hamiltonian: $H = H_Z + H_{ex} + H_{HF} + H_D + H_{SOC}$ [Distant Pairs]

Interactions

Zeeman: $H_Z = \frac{1}{2} \mu_B B_0 [g_a \hat{\sigma}_z^a + g_b \hat{\sigma}_z^b]$ { $T_{+,-}$ splitting from hyperfine field is 0.5 μeV at 5 mT

Exchange: $H_{ex} = -J \hat{S}_a \cdot \hat{S}_b$ { Dictates size where spin-dept. transitions occur
Pair distance $r \sim 3$ nm \rightarrow small ψ overlap

Hyperfine: $H_{HF} = \sum_{i=1}^2 \sum_{j=1}^{N_i} [S_i \cdot \tilde{A}_{ij} \cdot I_j]$ **DOMINANT** – responsible for *intrinsic* MFEs in OSCs

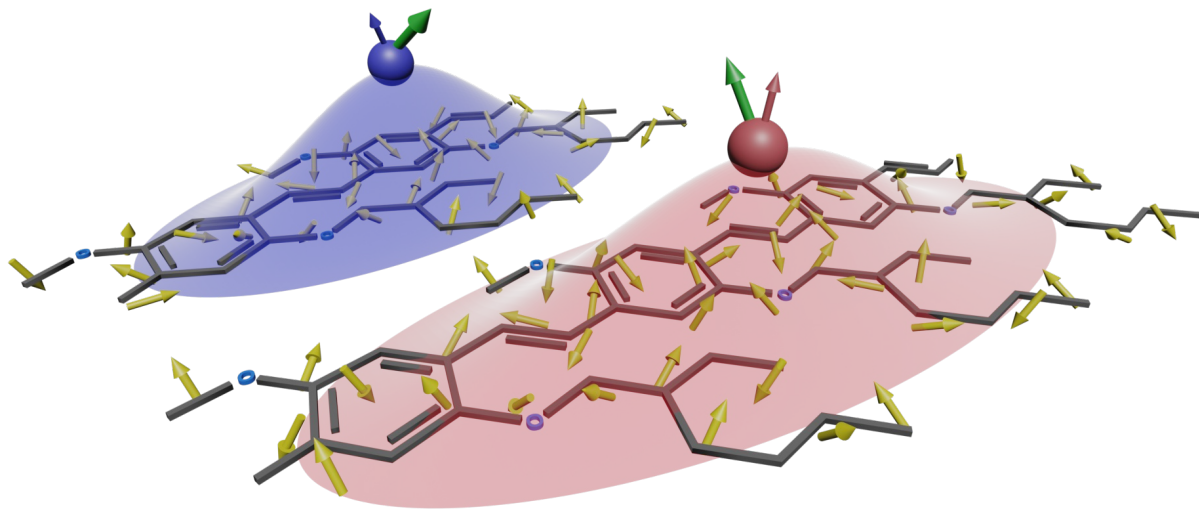
Dipolar: $H_D = \hat{S}_a^\dagger \tilde{D}_d \hat{S}_b$ $(\tilde{D}_d)_{ab} \sim r^{-3}$ { Negligible at large e-h separation distances / USMEL
Outcompeted by exchange at short distances

SOC: $H_{SOC} = \hat{S}_a^\dagger \tilde{D} \hat{S}_b$ $\tilde{D} \sim Z^4$ { Weak due to hydrocarbon composition

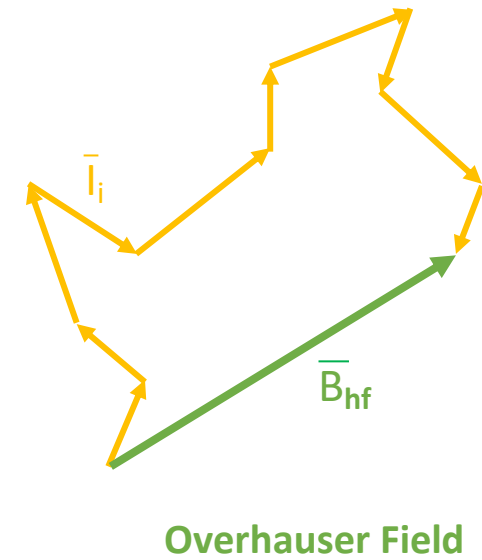
The Hyperfine Interaction

$$H_{\text{hf}} = \sum_{\mathbf{k}} a_{\text{HF},\mathbf{k}} \mathbf{S} \cdot \mathbf{I}_{\mathbf{k}} |\psi(\mathbf{r}_{\mathbf{k}})|^2$$

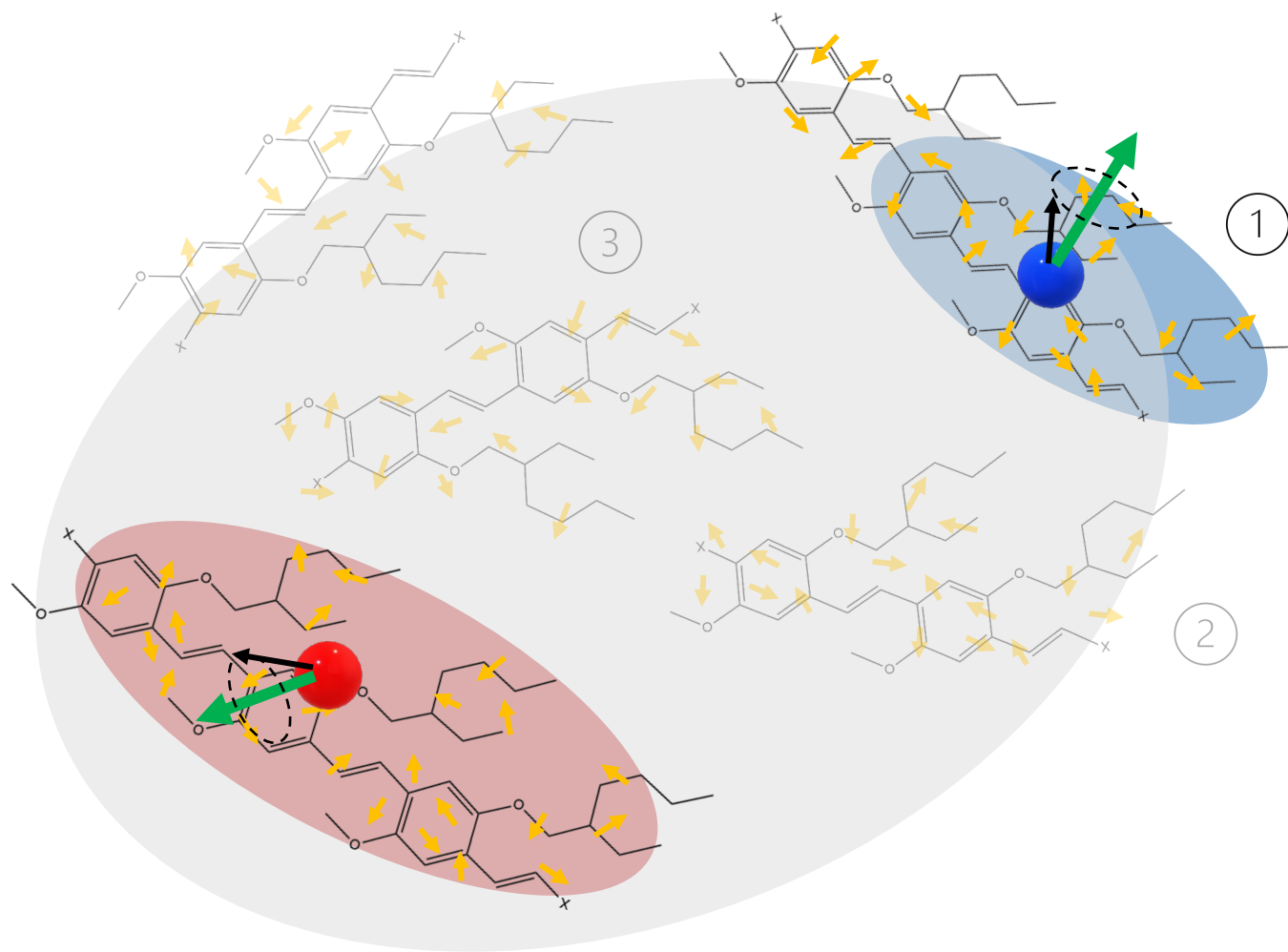
Interaction



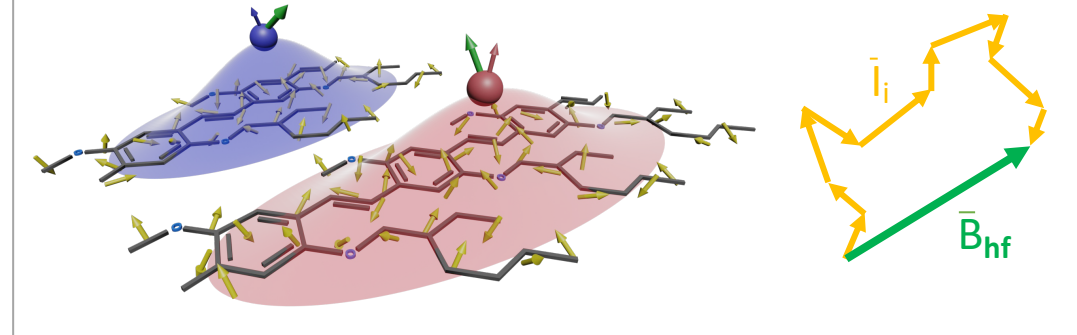
Property



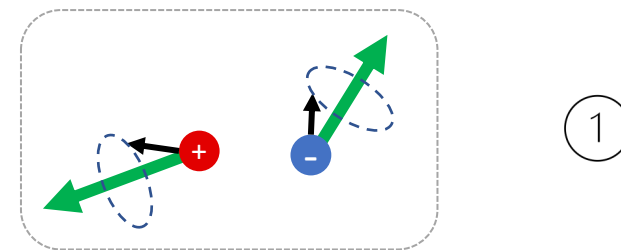
-  = Charge Carrier Spin
-  = ¹H Protium Spin
-  = Onsager Radius
-  = Local Overhauser Field



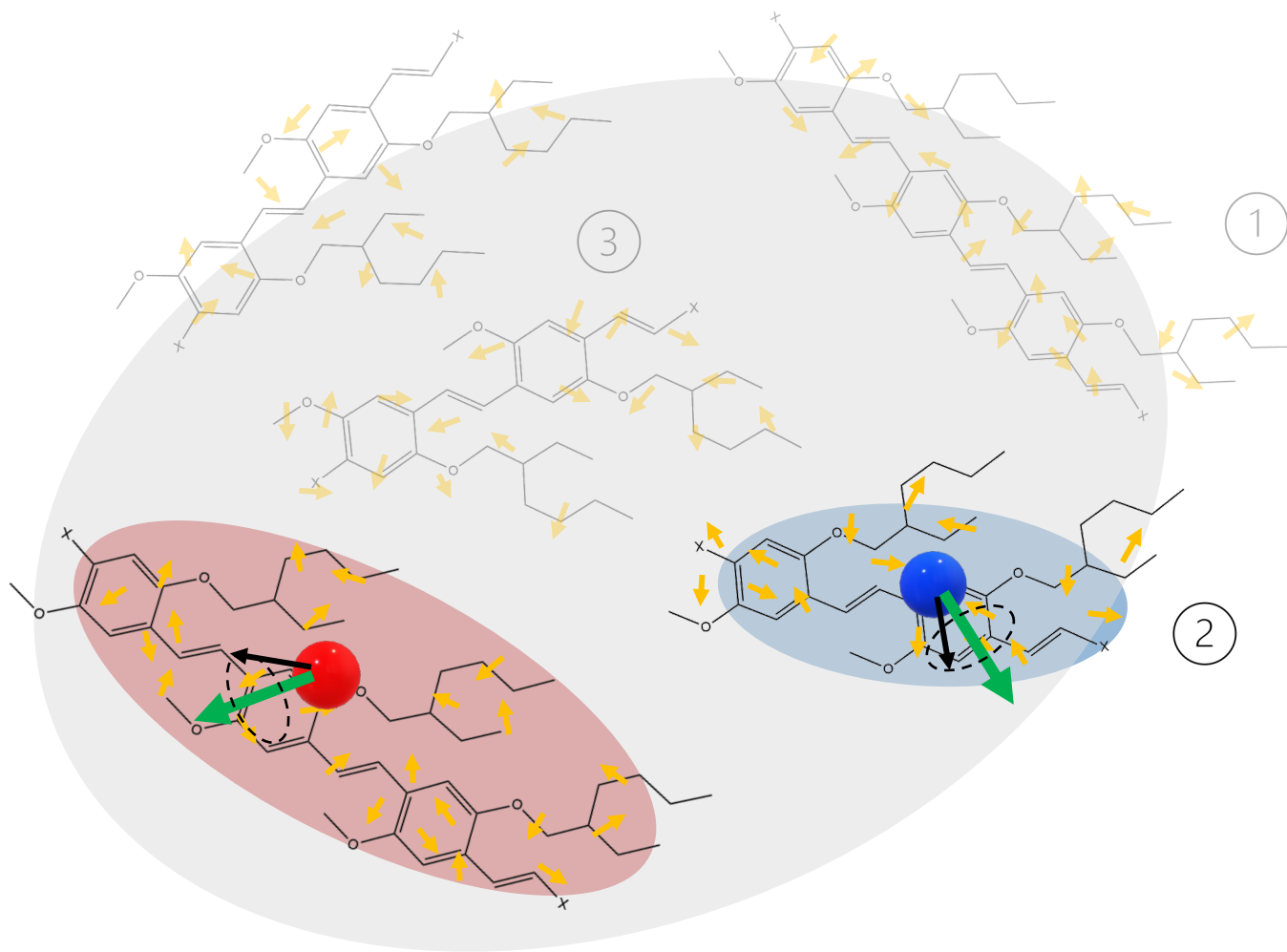
The Hyperfine Interaction

$$H_{hyp} = \sum_k a_{HF,k} \mathbf{S} \cdot \mathbf{I}_k |\psi(\mathbf{r}_k)|^2$$


The diagram shows two polymer chains, one blue and one red, with charge carriers (blue and red spheres) and their spins (black arrows). Yellow arrows represent the spins of the polymer chain segments. A dashed circle is drawn around the charge carrier in the blue chain. To the right, a vector diagram shows a green arrow labeled \bar{B}_{hf} and a yellow arrow labeled \bar{I}_i .

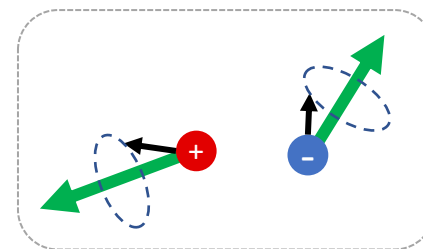
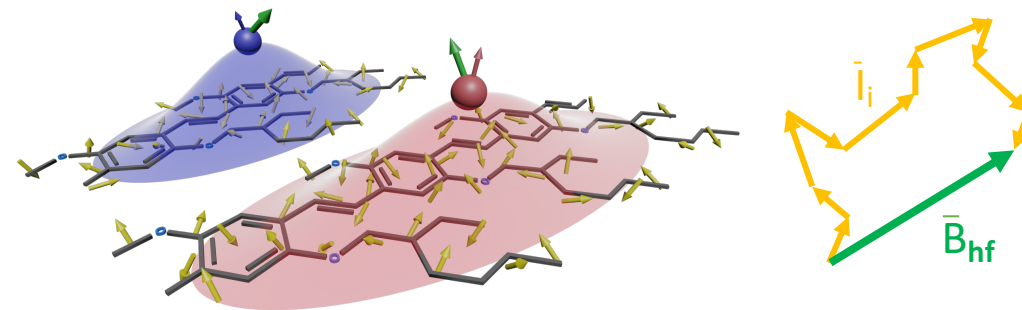


- = Charge Carrier Spin
- = ¹H Protium Spin
- = Onsager Radius
- = Local Overhauser Field

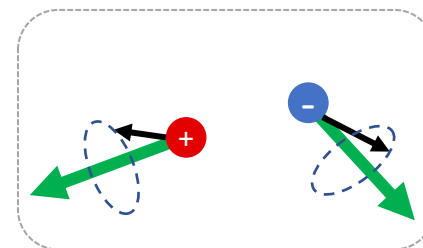


The Hyperfine Interaction

$$H_{hyp} = \sum_k a_{HF,k} \mathbf{S} \cdot \mathbf{I}_k |\psi(\mathbf{r}_k)|^2$$

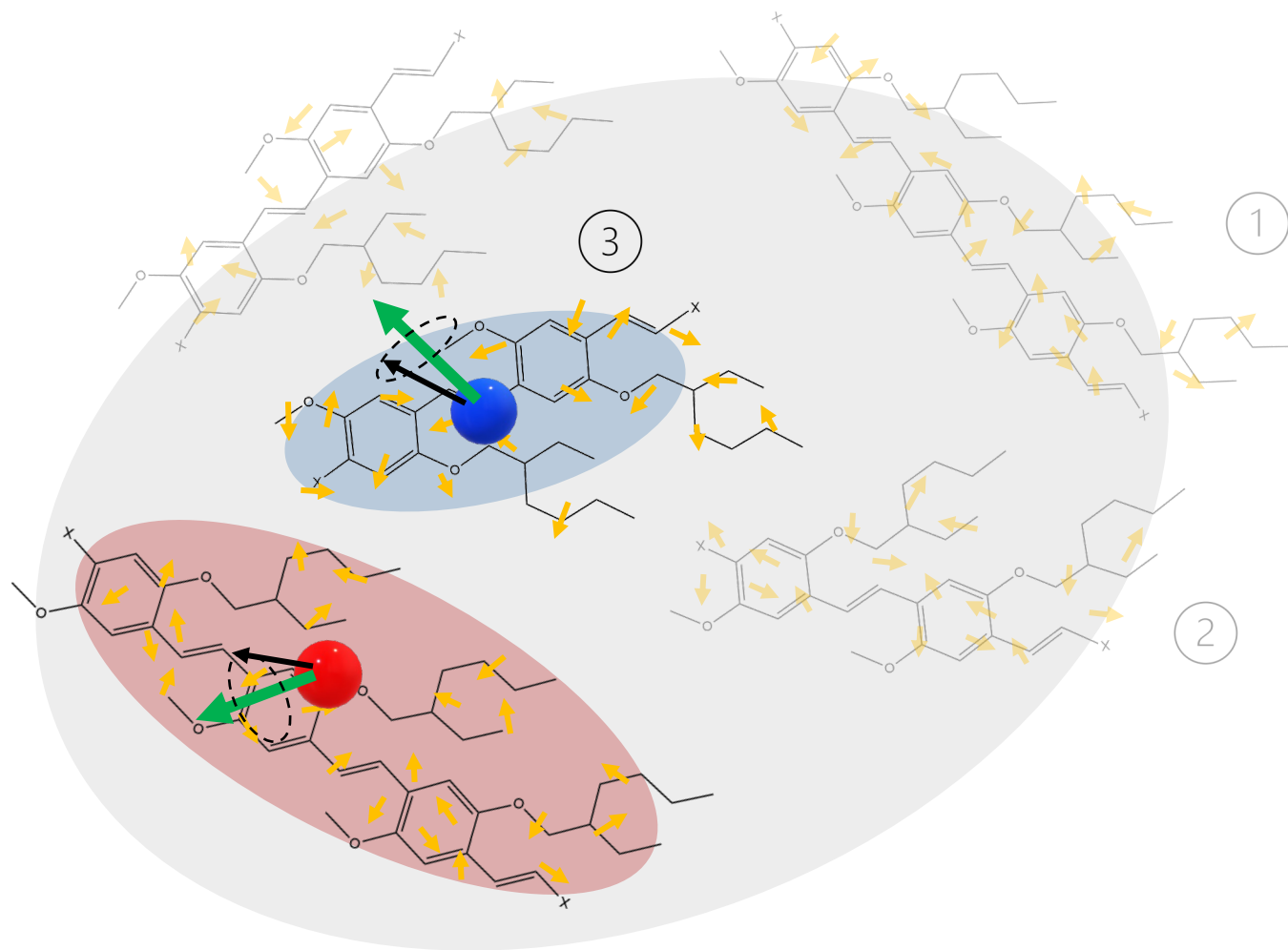


1

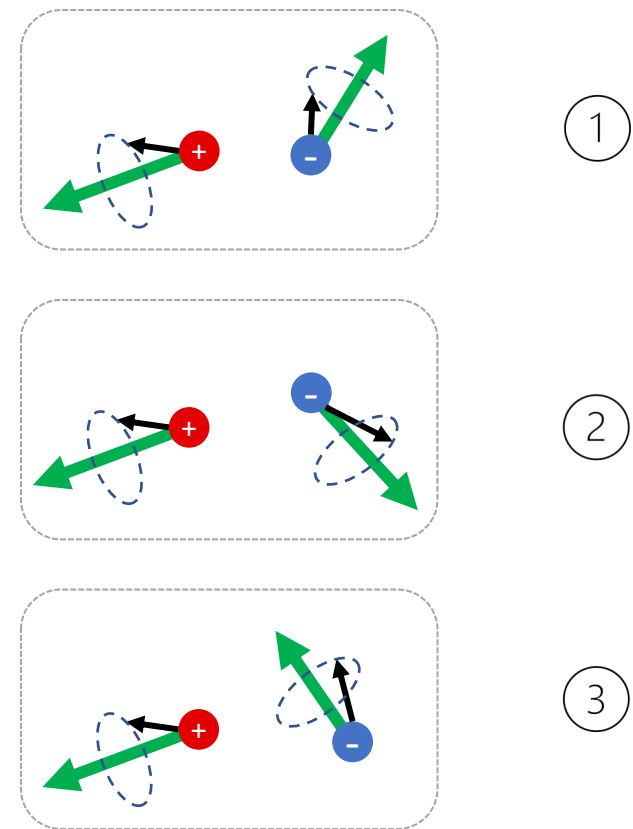


2

- = Charge Carrier Spin
- = ^1H Protium Spin
- = Onsager Radius
- = Local Overhauser Field



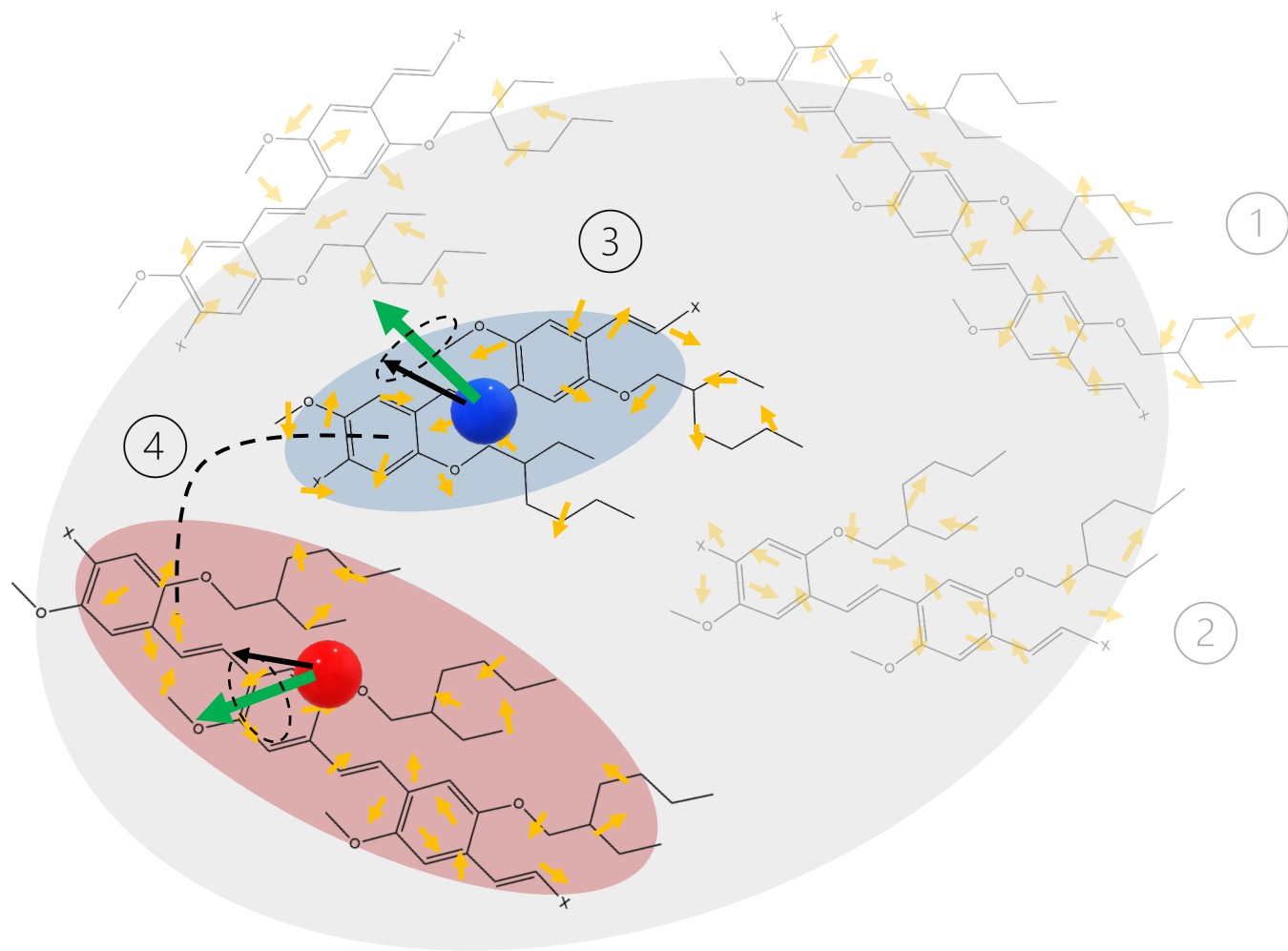
The Hyperfine Interaction

$$H_{hyp} = \sum_k a_{HF,k} \mathbf{S} \cdot \mathbf{I}_k |\psi(\mathbf{r}_k)|^2$$


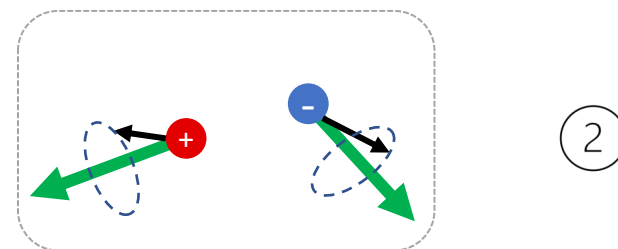
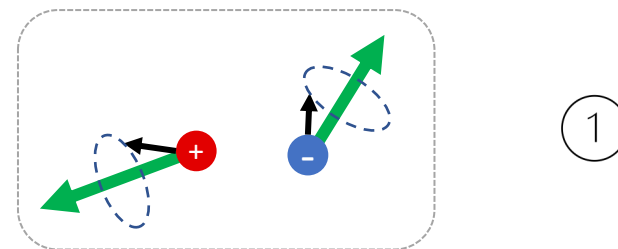
Spin mixing
 $\Delta\omega_{hf} + \Delta\theta$

mixed state: $\alpha|T_0\rangle + \beta|S\rangle$

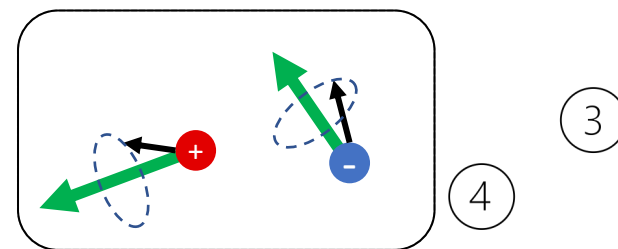
- = Charge Carrier Spin
- = ¹H Protium Spin
- = Onsager Radius
- = Local Overhauser Field



The Hyperfine Interaction $H_{hyp} = \sum_k a_{HF,k} \mathbf{S} \cdot \mathbf{I}_k |\psi(\mathbf{r}_k)|^2$



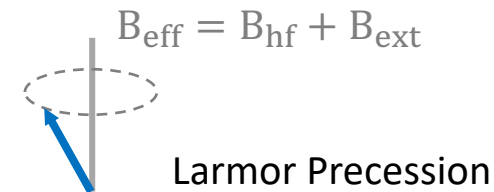
$\Delta E_{ST} \gg 0$



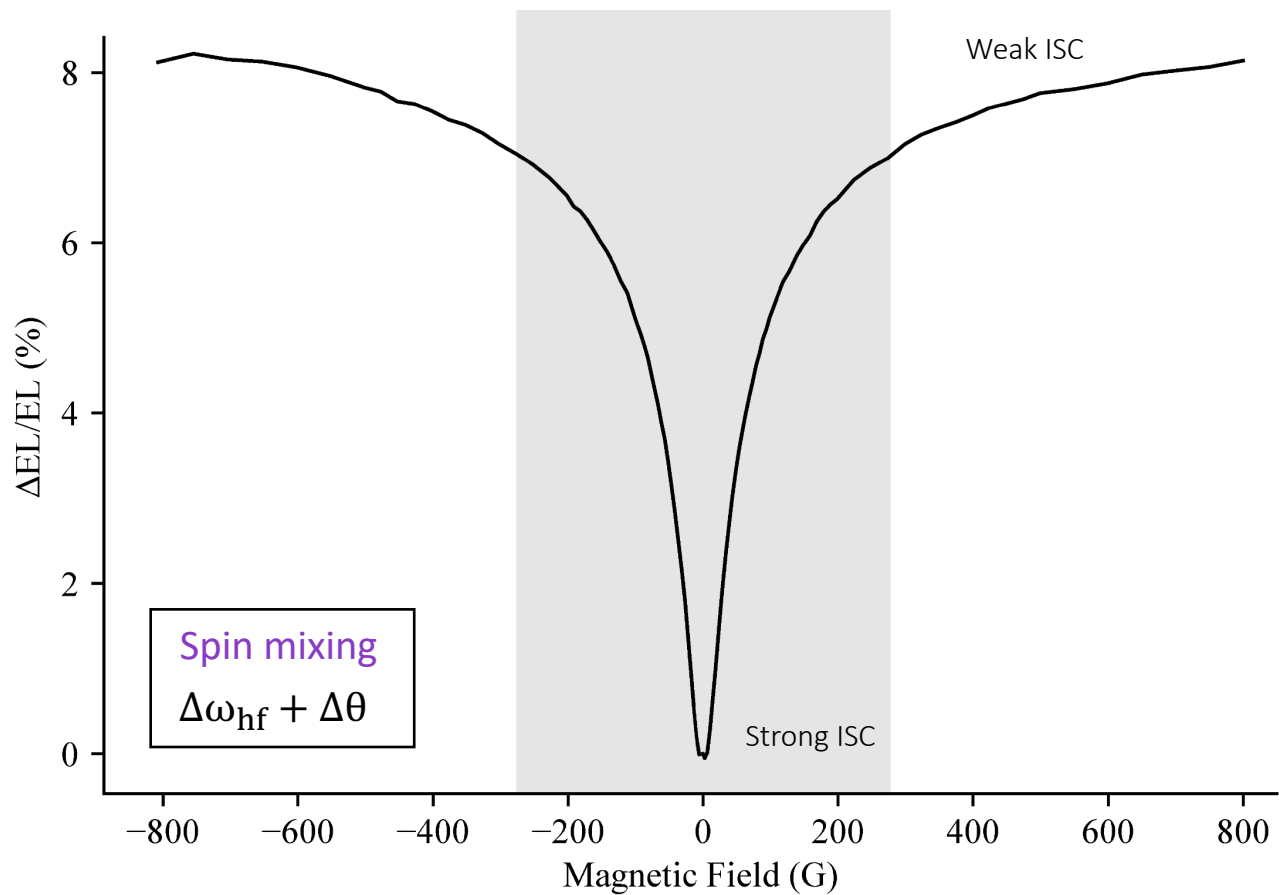
④

Spin symmetry "locked-in"

Probing Hyperfine Fields using Magnetic Fields

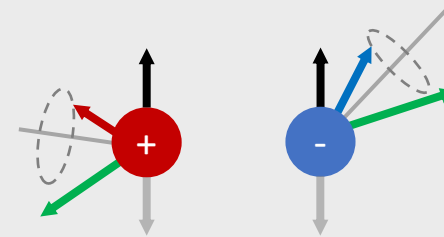


Magneto-electroluminescence (MEL)



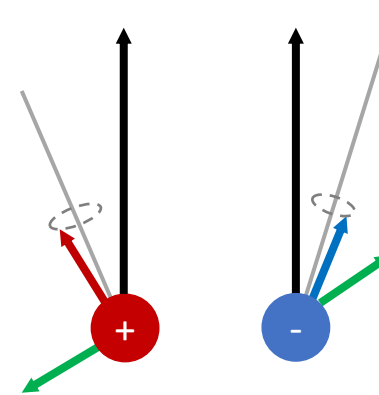
$S_1 \rightarrow T_1$ Intersystem Crossing

$B_{\text{ext}} \sim B_{\text{hf}}$



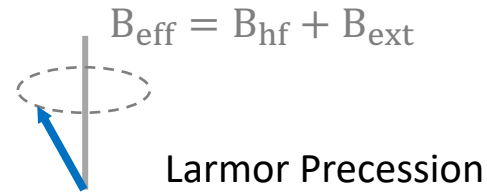
Weak spin mixing

$B_{\text{ext}} \gg B_{\text{hf}}$

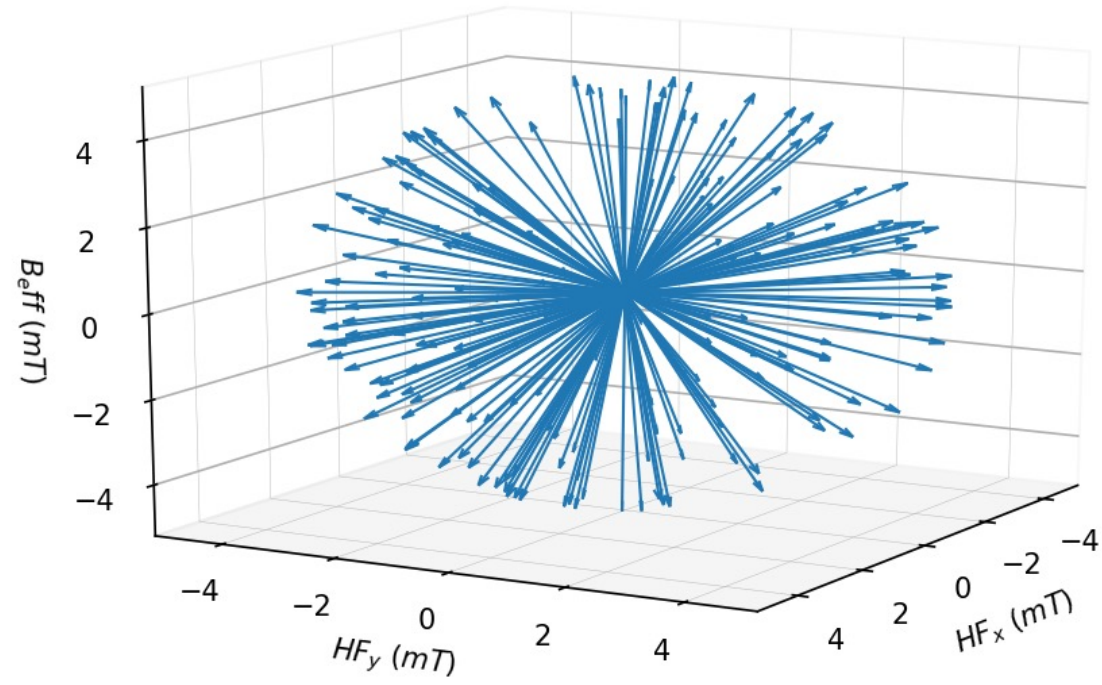


Spin mixing
 $\Delta\omega_{hf} + \Delta\theta$

Overcoming Random Hyperfine Fields



Effective Field Distribution | $B_{ext} = 0.0$ mT



$S_1 \rightarrow T_1$ Intersystem Crossing

$B_{ext} \sim B_{hf}$

Strong spin mixing

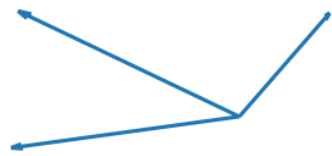
Weak spin mixing

$B_{ext} \gg B_{hf}$

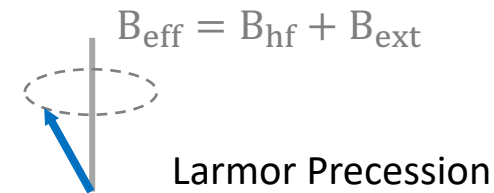
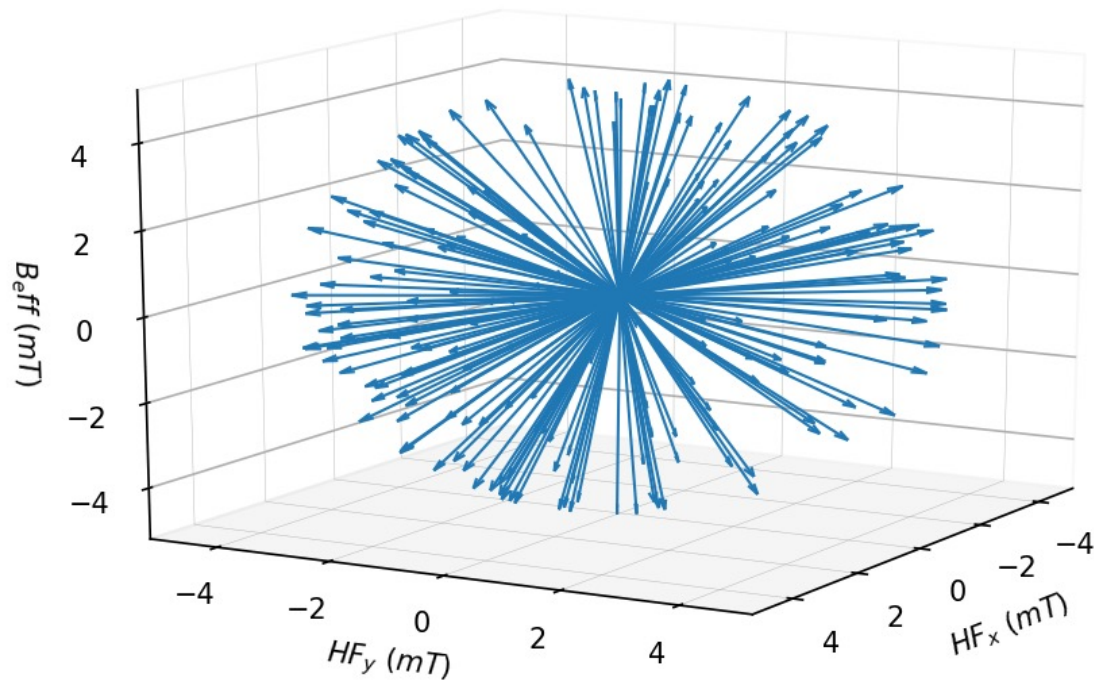
Intersystem Crossing Suppressed

Spin mixing
 $\Delta\omega_{\text{hf}} + \Delta\theta$

Overcoming Random Hyperfine Fields

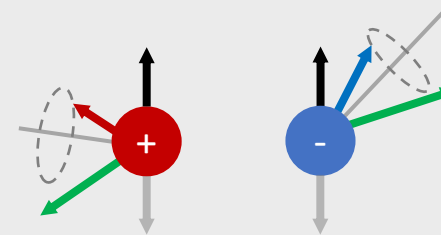


Effective Field Distribution | $B_{\text{ext}} = 0.0$ mT



$S_1 \rightarrow T_1$ Intersystem Crossing

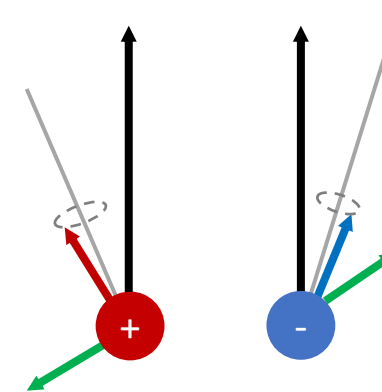
$B_{\text{ext}} \sim B_{\text{hf}}$



Strong spin mixing

Weak spin mixing

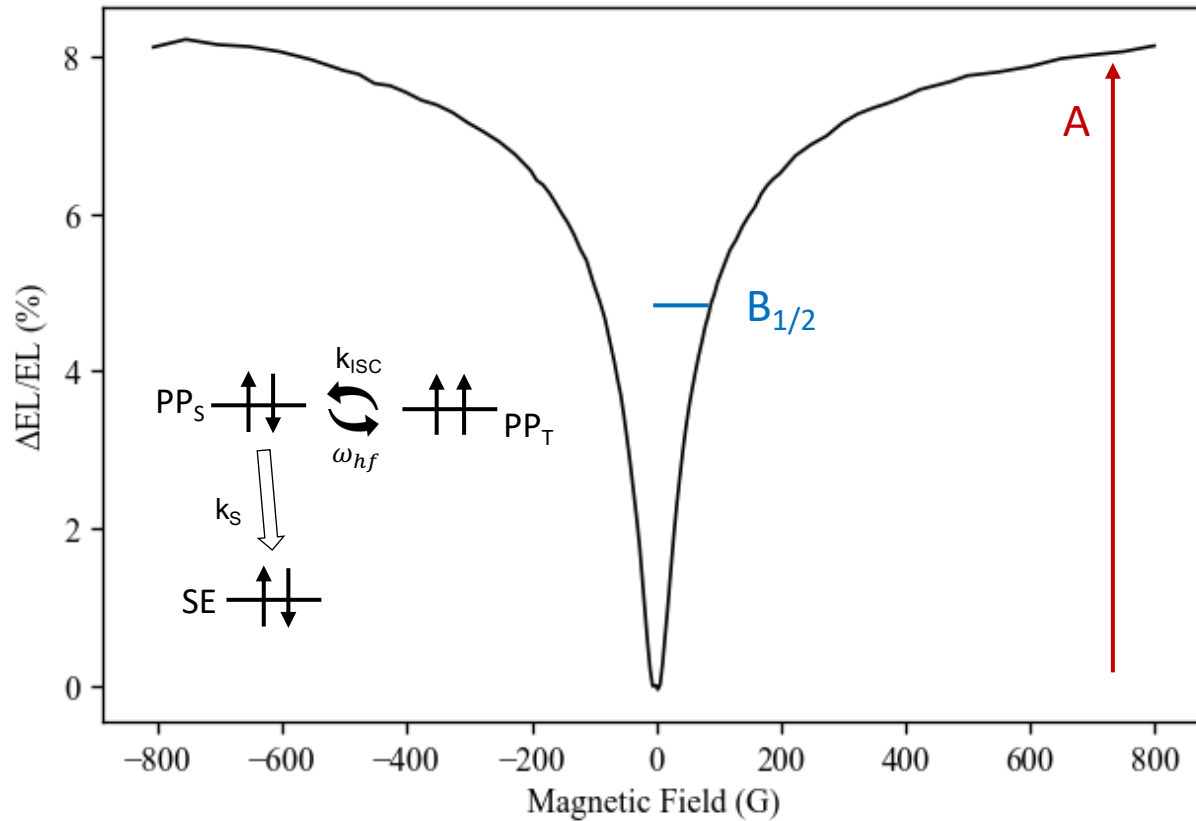
$B_{\text{ext}} \gg B_{\text{hf}}$



Intersystem Crossing Suppressed

Fitting parameters **A** and **B_{1/2}** of the MEL curve give information about spin properties

$$\text{MEL}(B) = A \frac{B^2}{(|B| + B_0)^2} + \delta$$



S.P. Kersten, A.J. Schellekens, B. Koopmans, P. A. Bobbert, *PRL*, 106 (19), 197402 (2011)

T.D. Nguyen *et al.*, *Nature Mater* 9, 345–352 (2010)

Amplitude (A)

A competition between **spin mixing** (ω_{hf}) and **singlet exciton formation rate** (k_S)

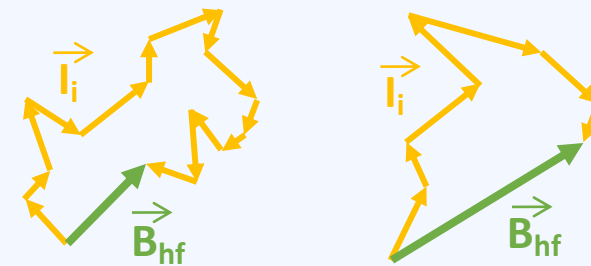
- Large **A** $\Rightarrow k_S / \omega_{hf} \ll 1$
- Small **A** $\Rightarrow k_S / \omega_{hf} \gg 1$

k_S depends on material properties

Half-Width (B_{1/2})

A measure of the **hyperfine interaction**

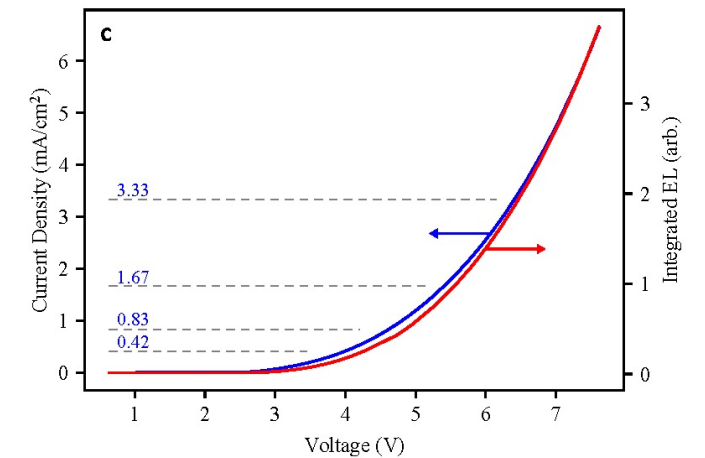
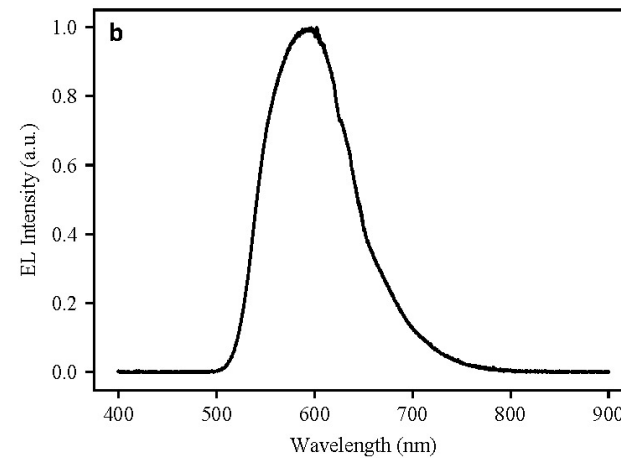
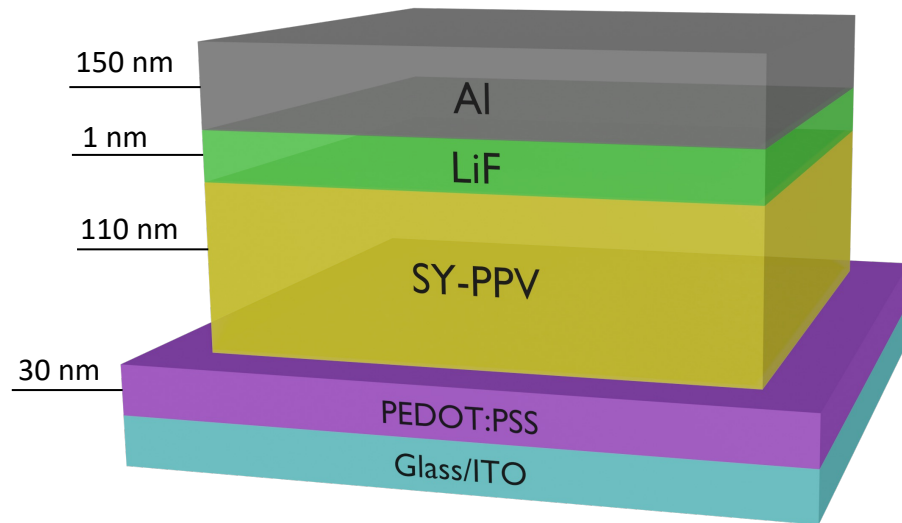
- **characteristic volume** of polarons (CLT)

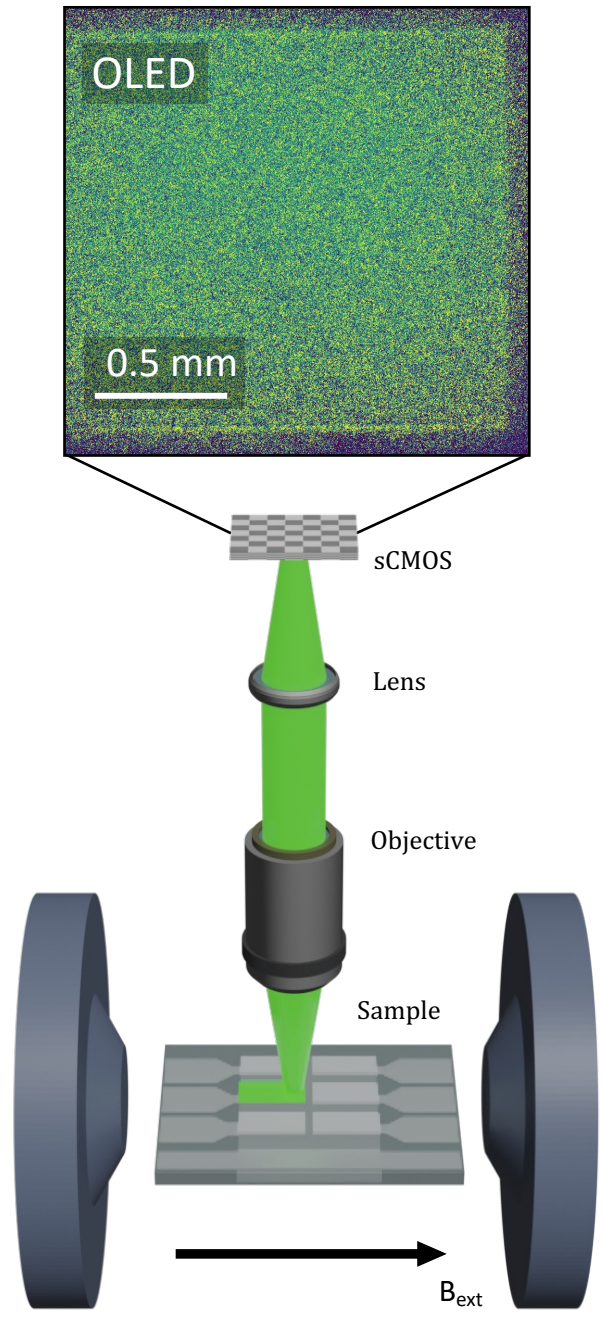


$$B_{hf} \sim 1/\sqrt{N} \sim 1/\sqrt{V}$$

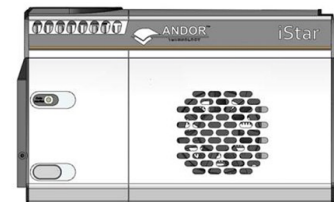
Probing the quantum properties of materials in organic electronics

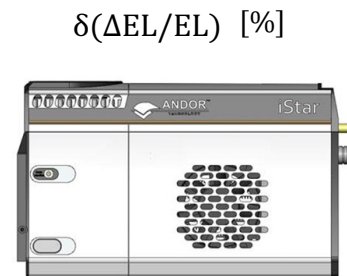
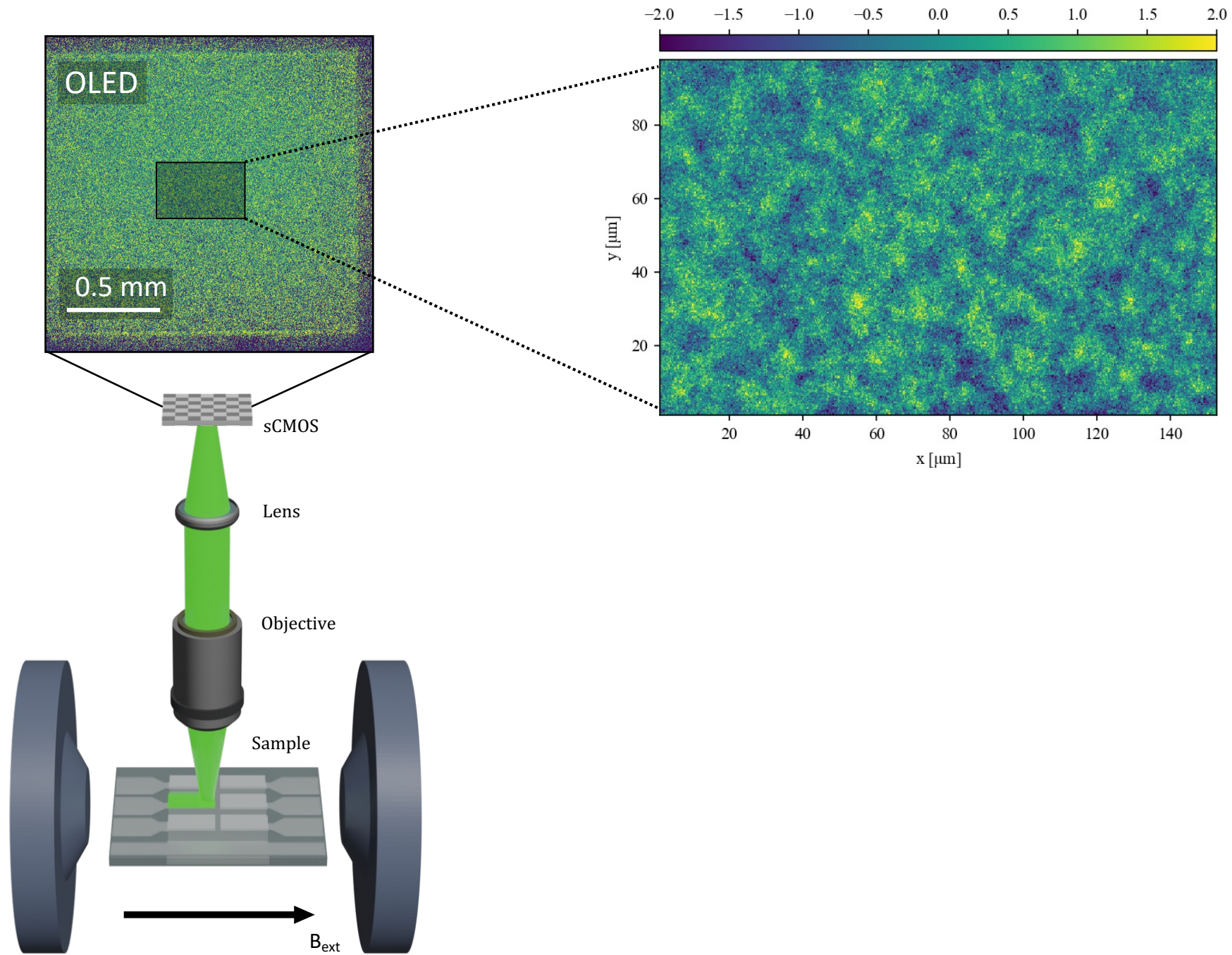
Simple vertical structure polymer OLED

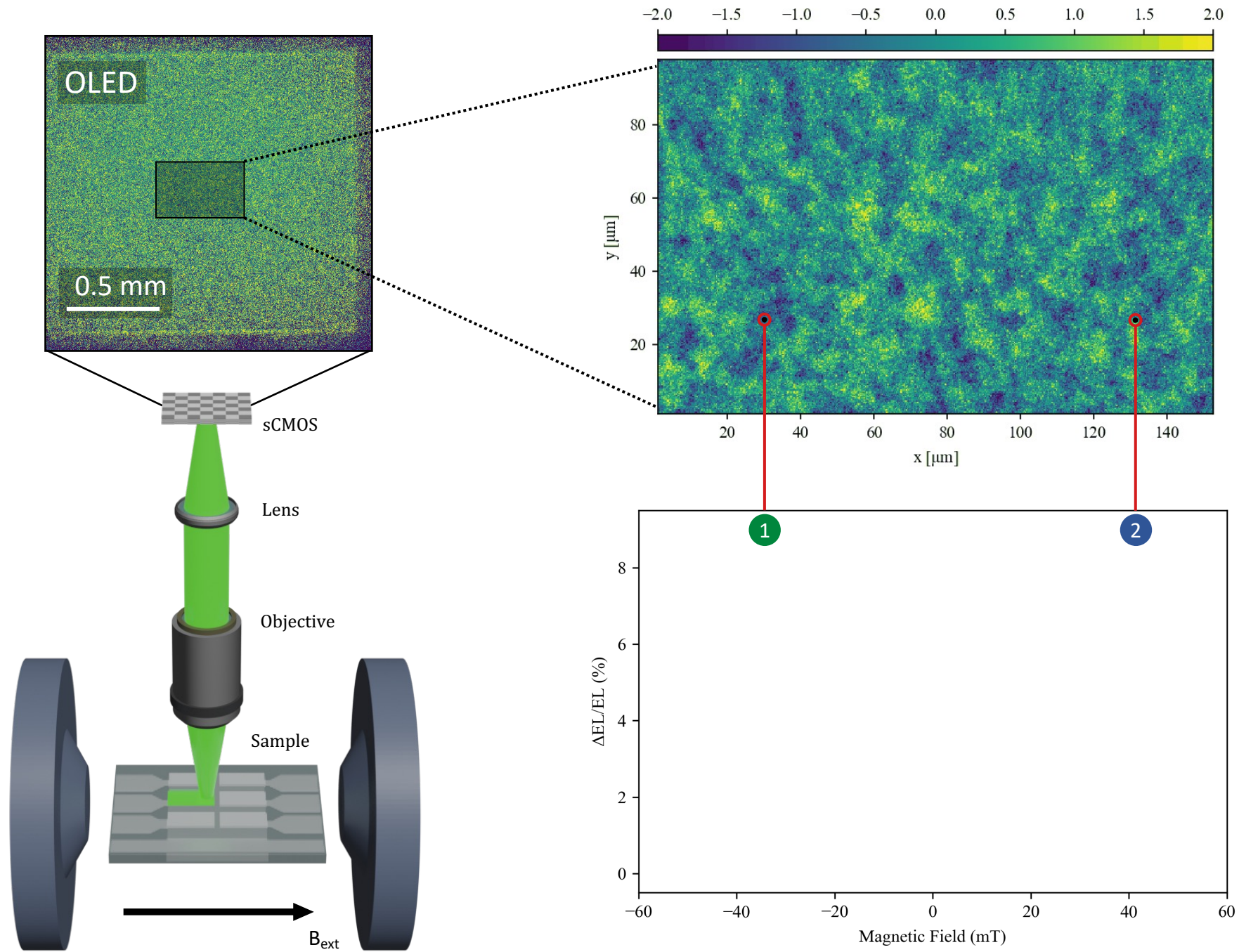




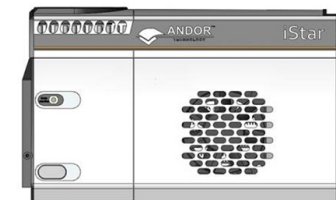
$\delta(\Delta EL/EL)$ [%]





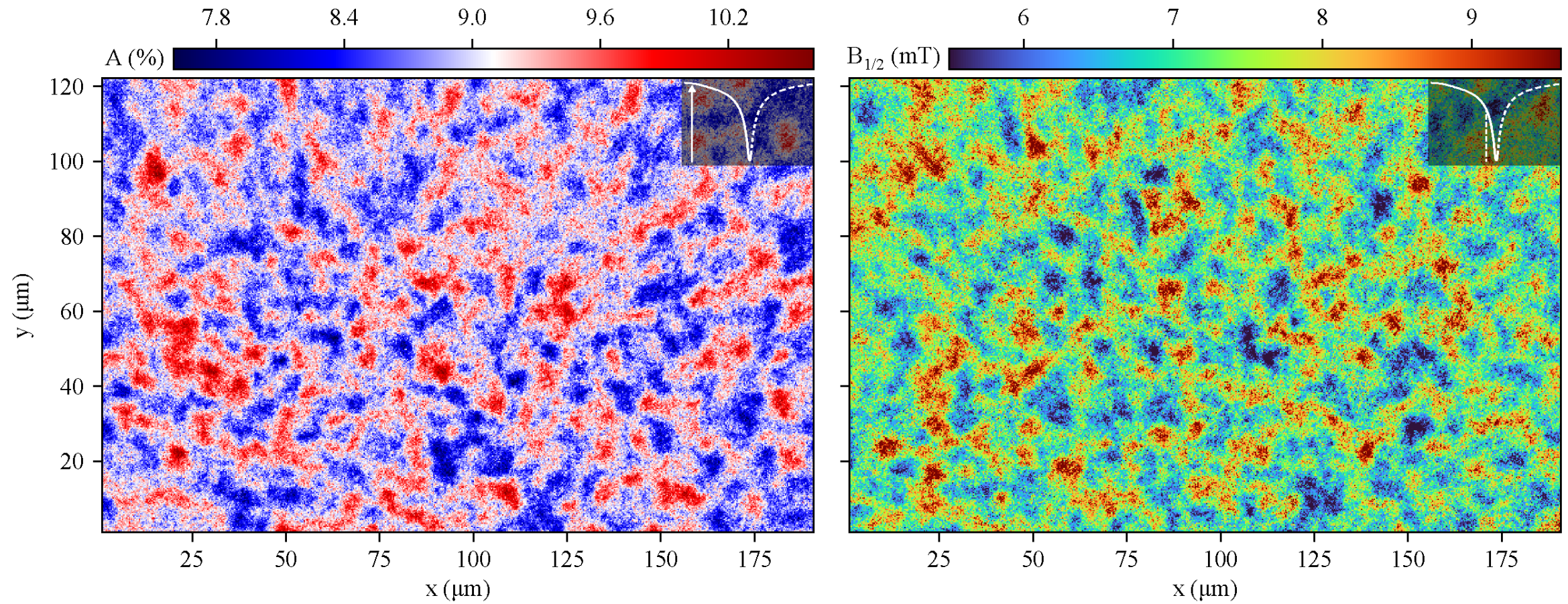


$\delta(\Delta\text{EL}/\text{EL})$ [%]



$$\text{MEL}(B) = A \frac{B^2}{(|B| + B_0)^2} + \delta$$

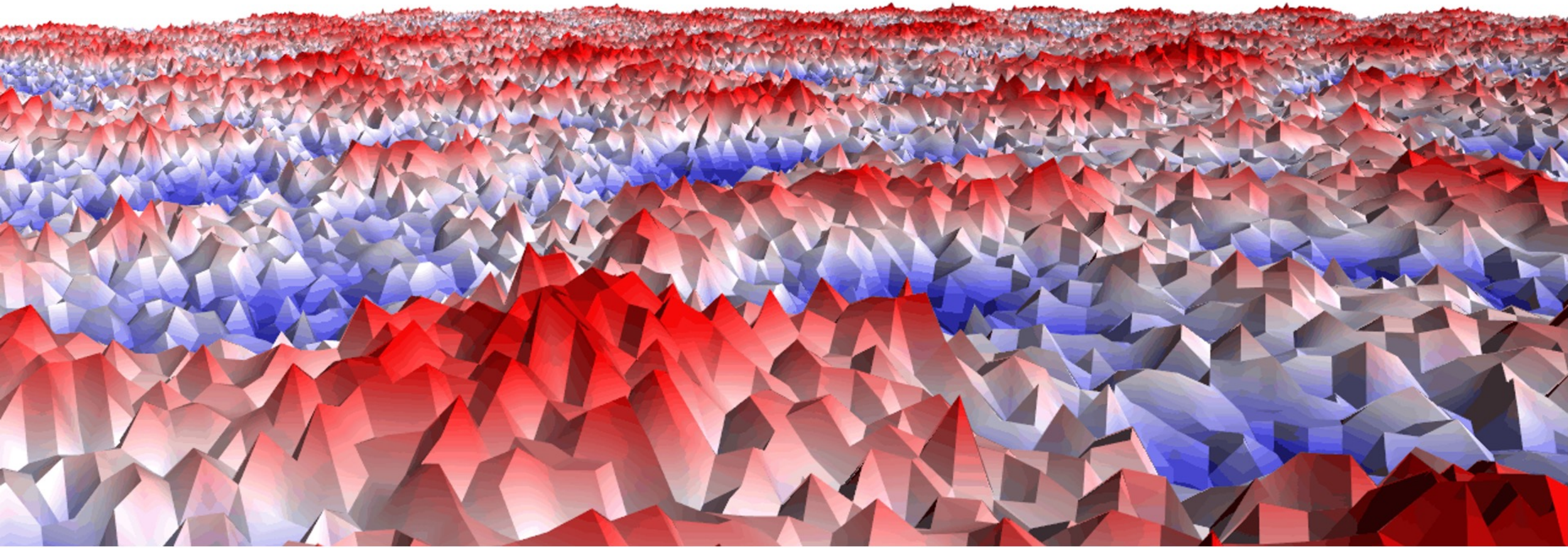
Spatial distributions of MEL parameters (A , $B_{1/2}$)



A competition between **spin mixing** (ω_{hf}) and **singlet exciton formation rate** (k_S)...

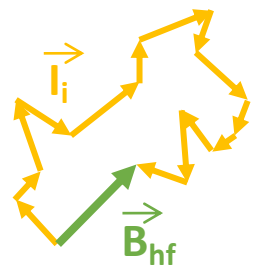
regions of larger **A** $\Rightarrow k_S / \omega_{\text{hf}} \ll 1$

smaller **A** $\Rightarrow k_S / \omega_{\text{hf}} \gg 1$

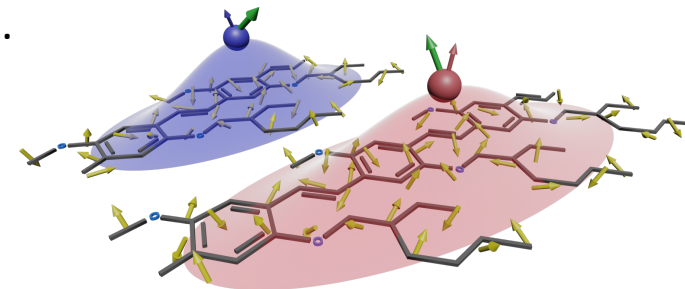
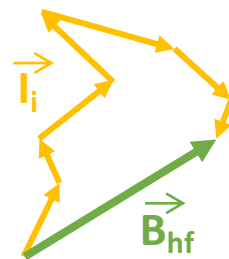


A measure of the **hyperfine interaction** and therefore **characteristic volume** of polarons ...

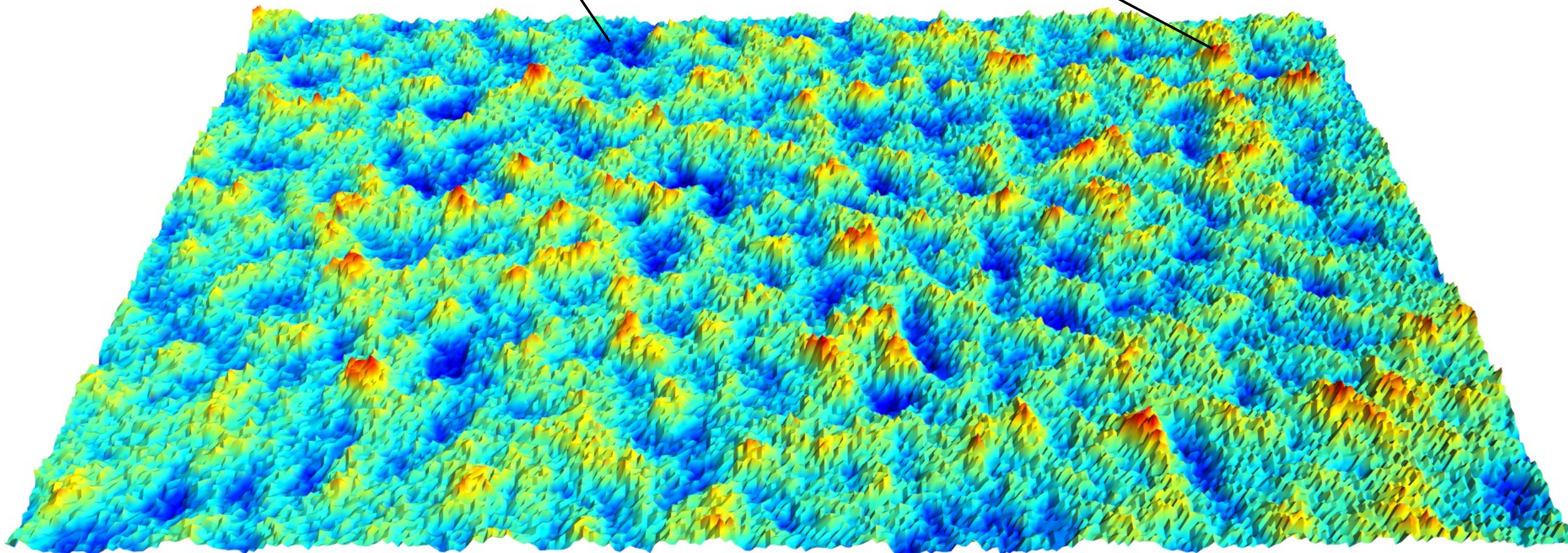
*larger
polarons*



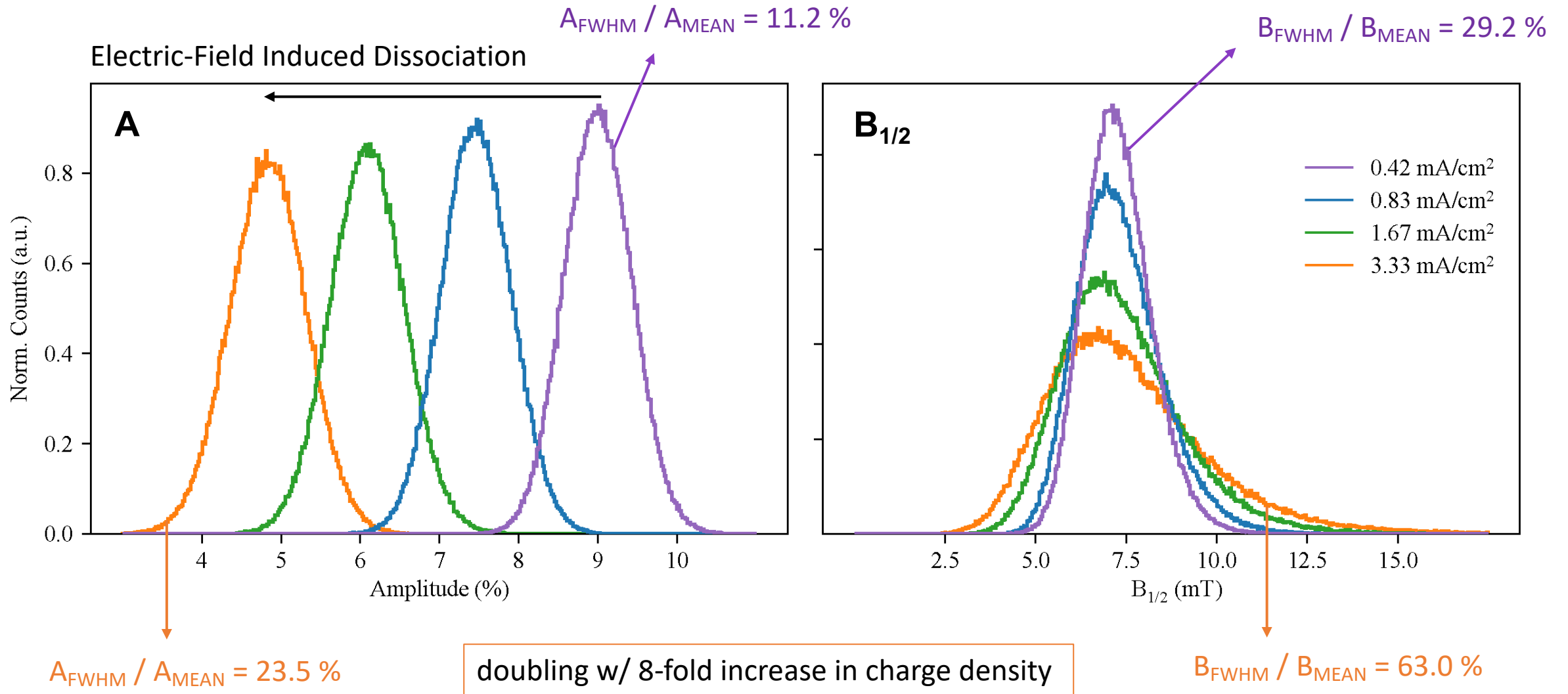
*smaller
polarons*



$$B_{\text{hf}} \sim 1/\sqrt{N} \sim 1/\sqrt{V}$$

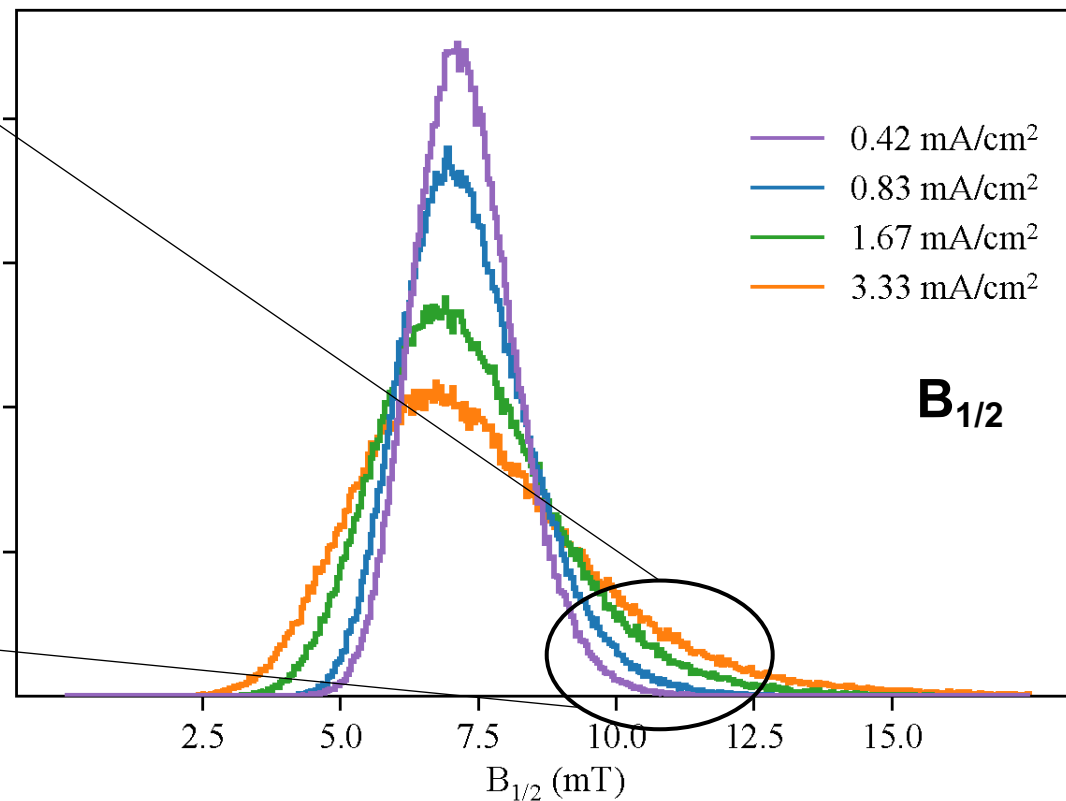
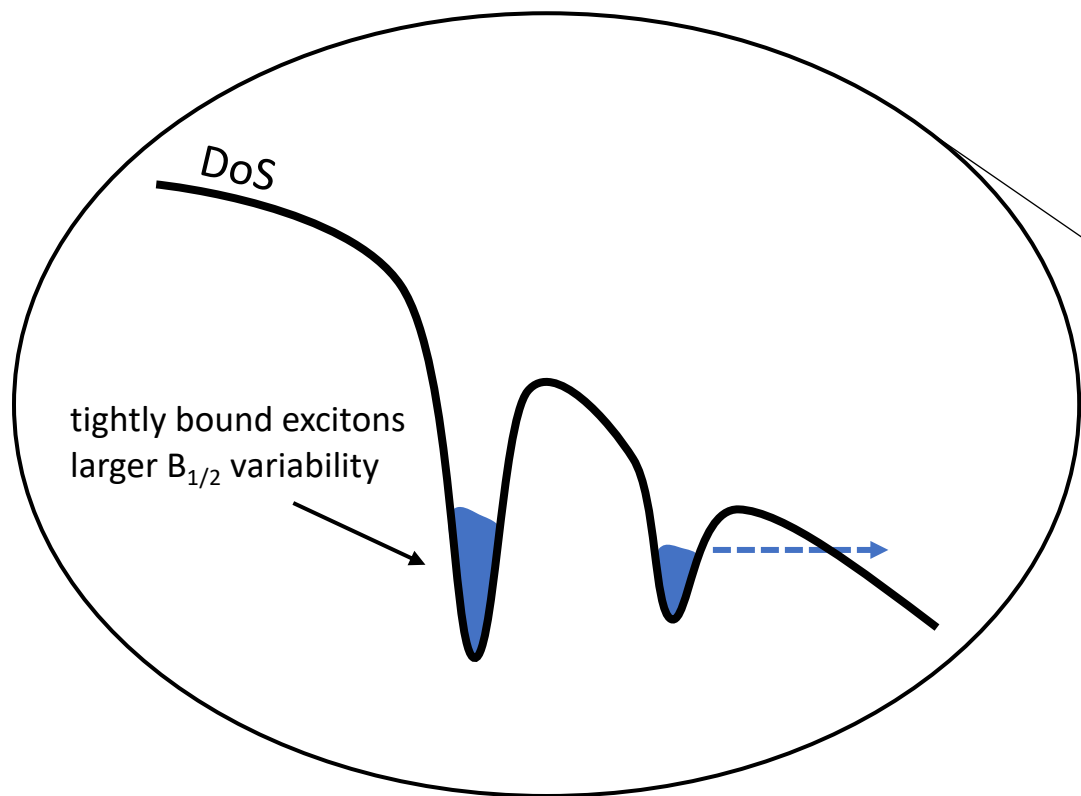
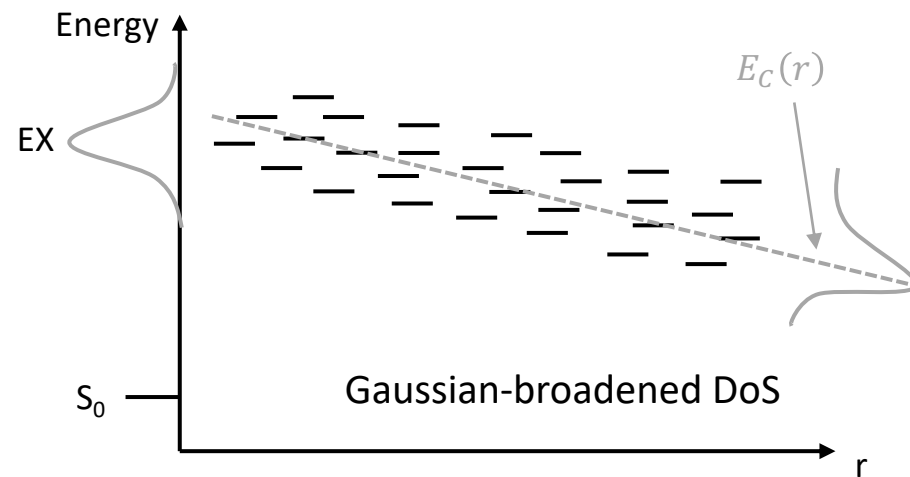


Intra-device variation of MEL parameters (A, B_{1/2})



Electric-Field Induced Dissociation

$$E_C(r) = -\frac{e^2}{4\pi\epsilon_0\epsilon r} - eF.r$$

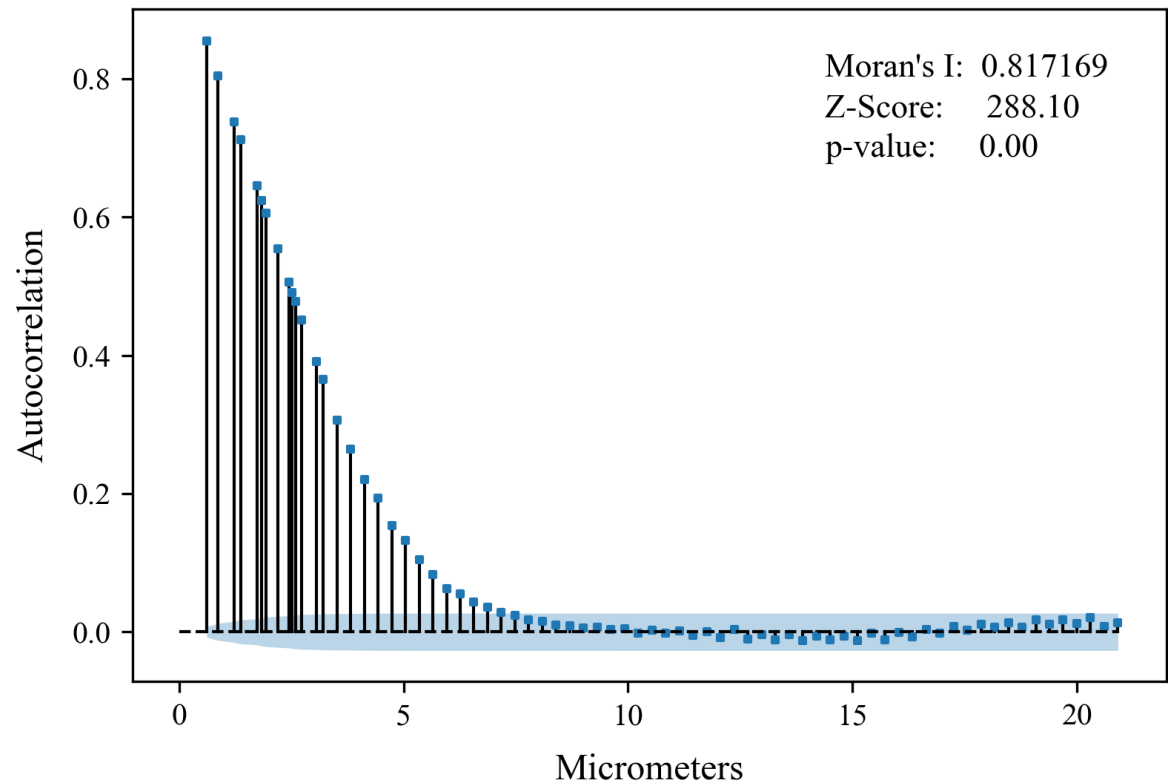
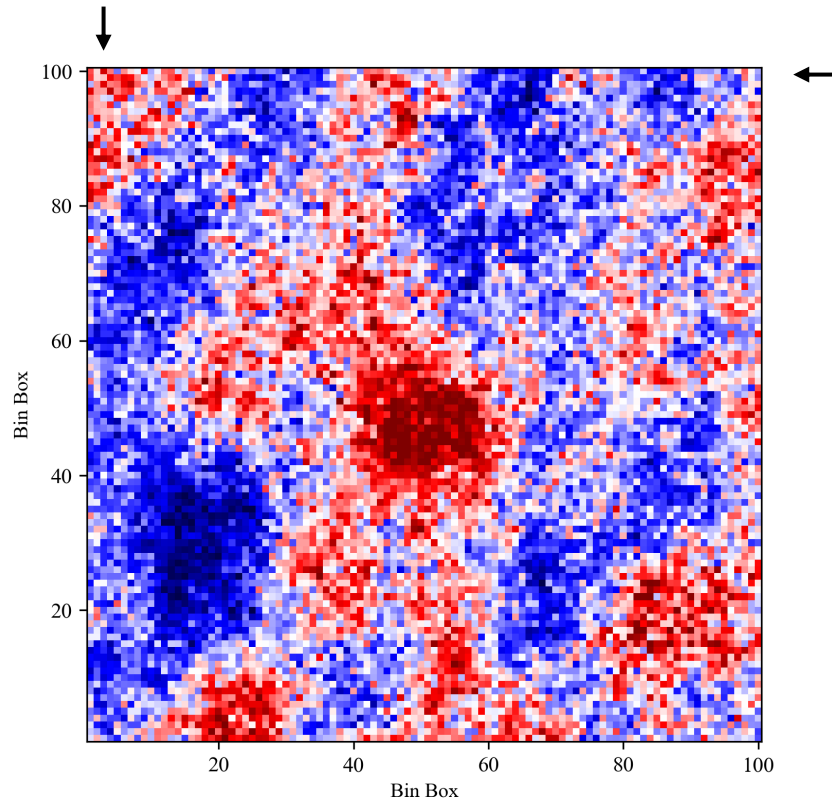


bias-dependent changes in **characteristic** polaron sizes
(smaller binding energy = more rapid dissociation)

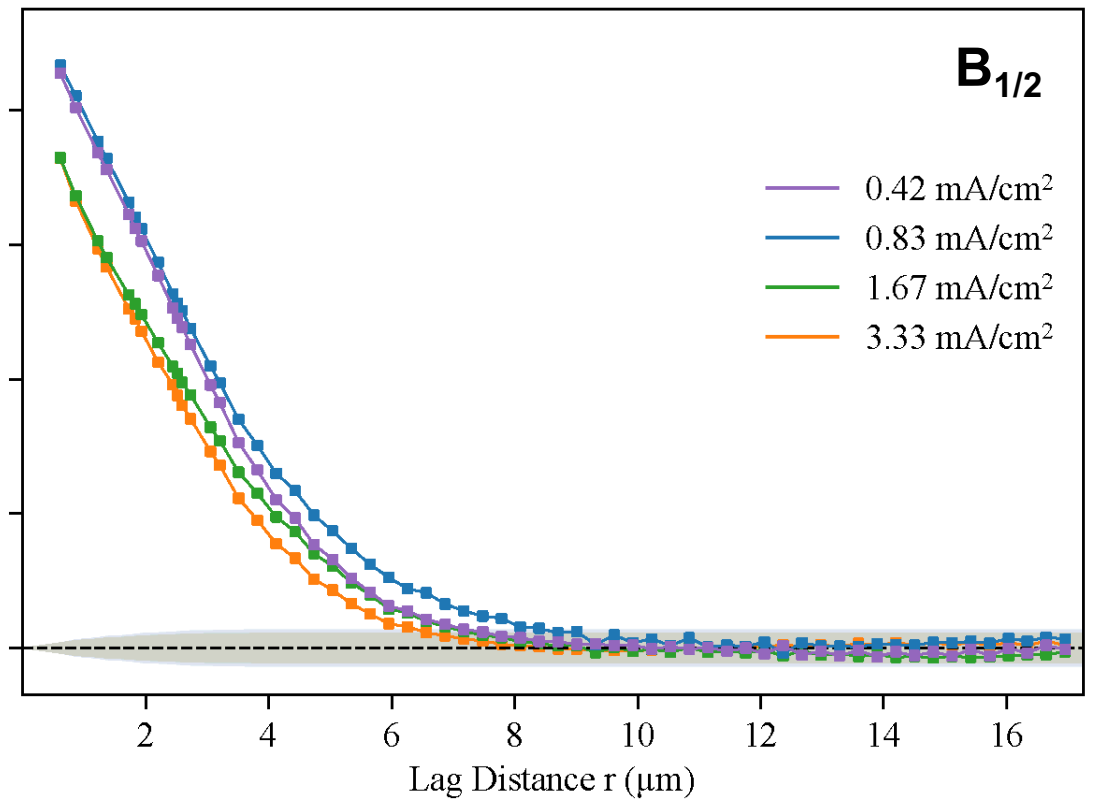
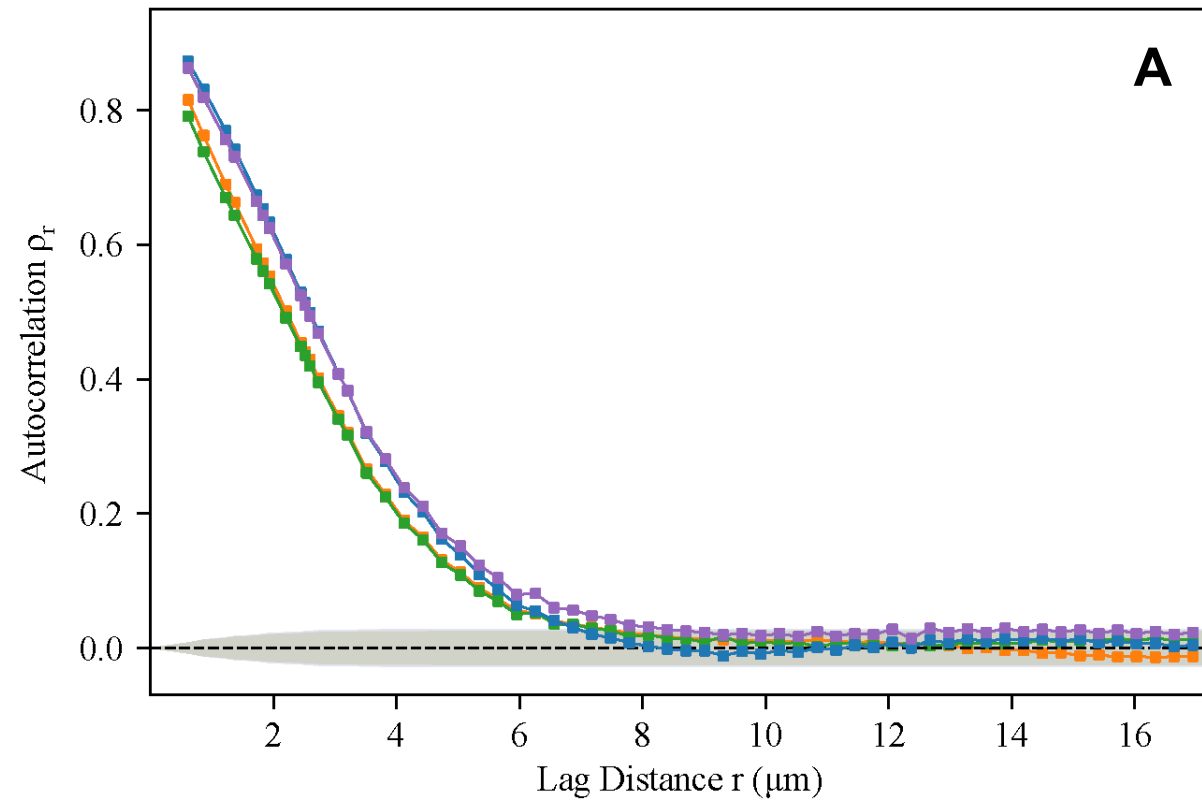
Moran's I – A global measure of spatial autocorrelation

Calculating the similarity of values between neighbours:

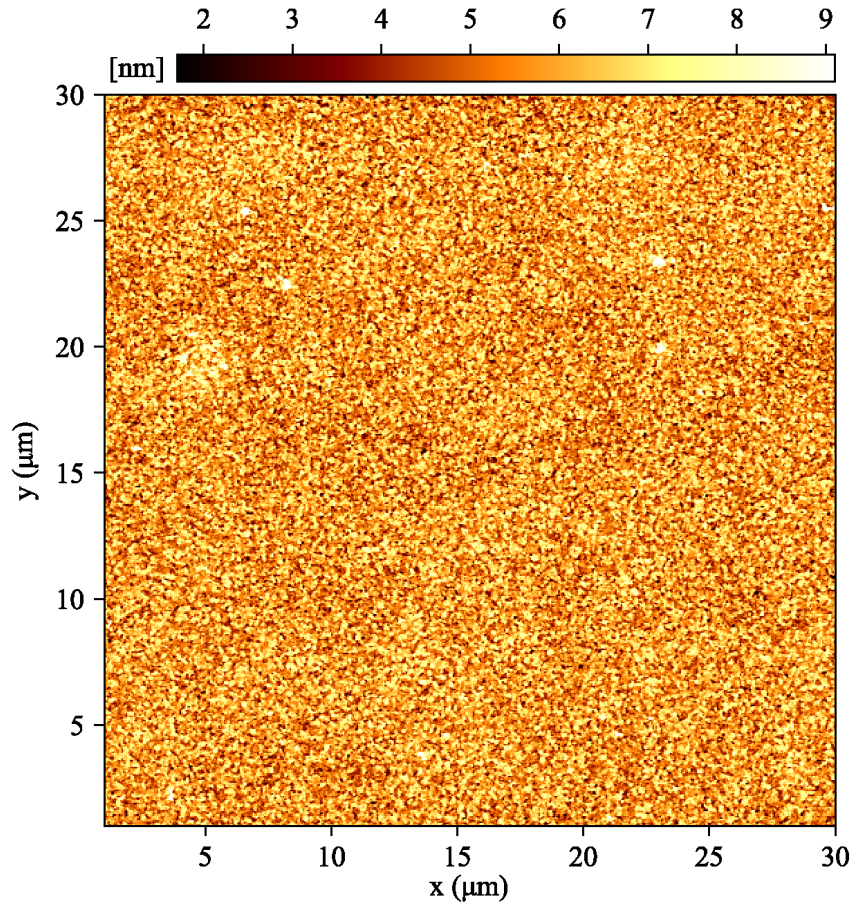
$$\rho_r = \frac{\gamma_r}{\gamma_0} = \frac{\sum_i^M \sum_j^N \sigma_{ij} \sum_k^M \sum_l^N w_{kl} \cdot \sigma_{kl}}{\sum_i^M \sum_j^N \sigma_{ij}^2}$$



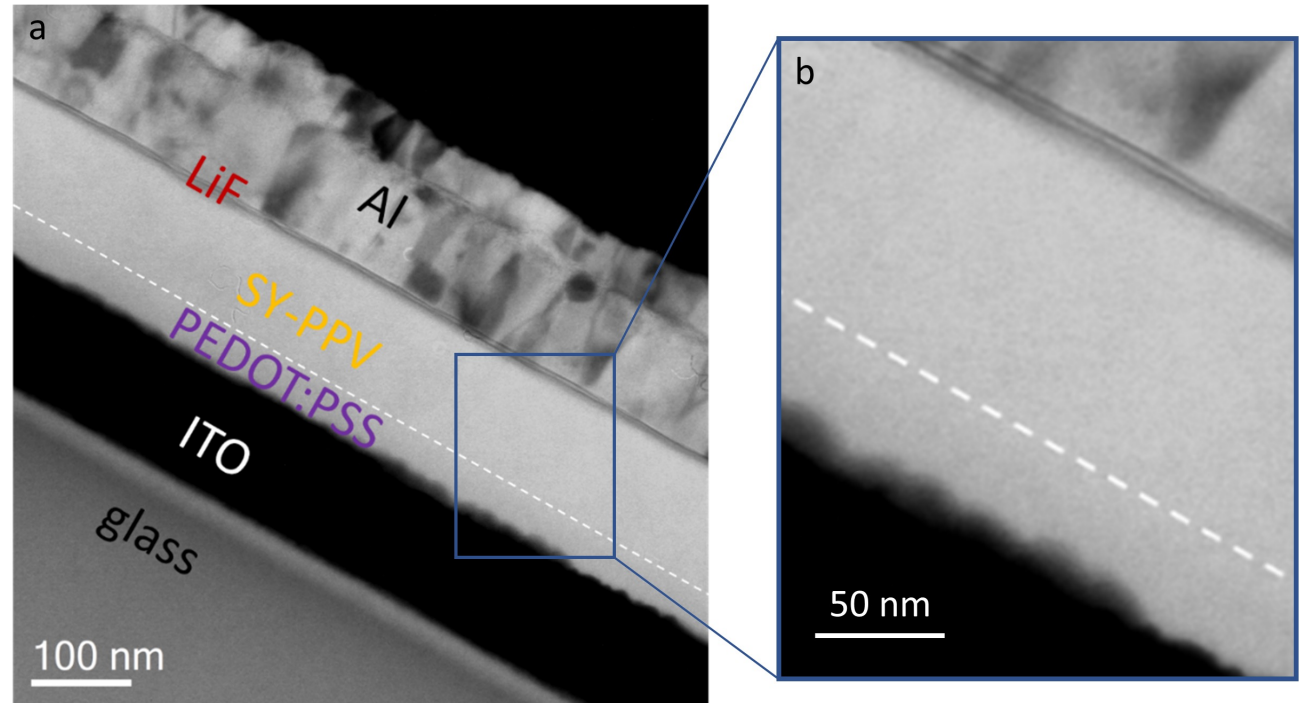
Spatial correlation of MEL parameters (A, $B_{1/2}$)



Film Homogeneity – Uniform and defect free on a microscopic scale

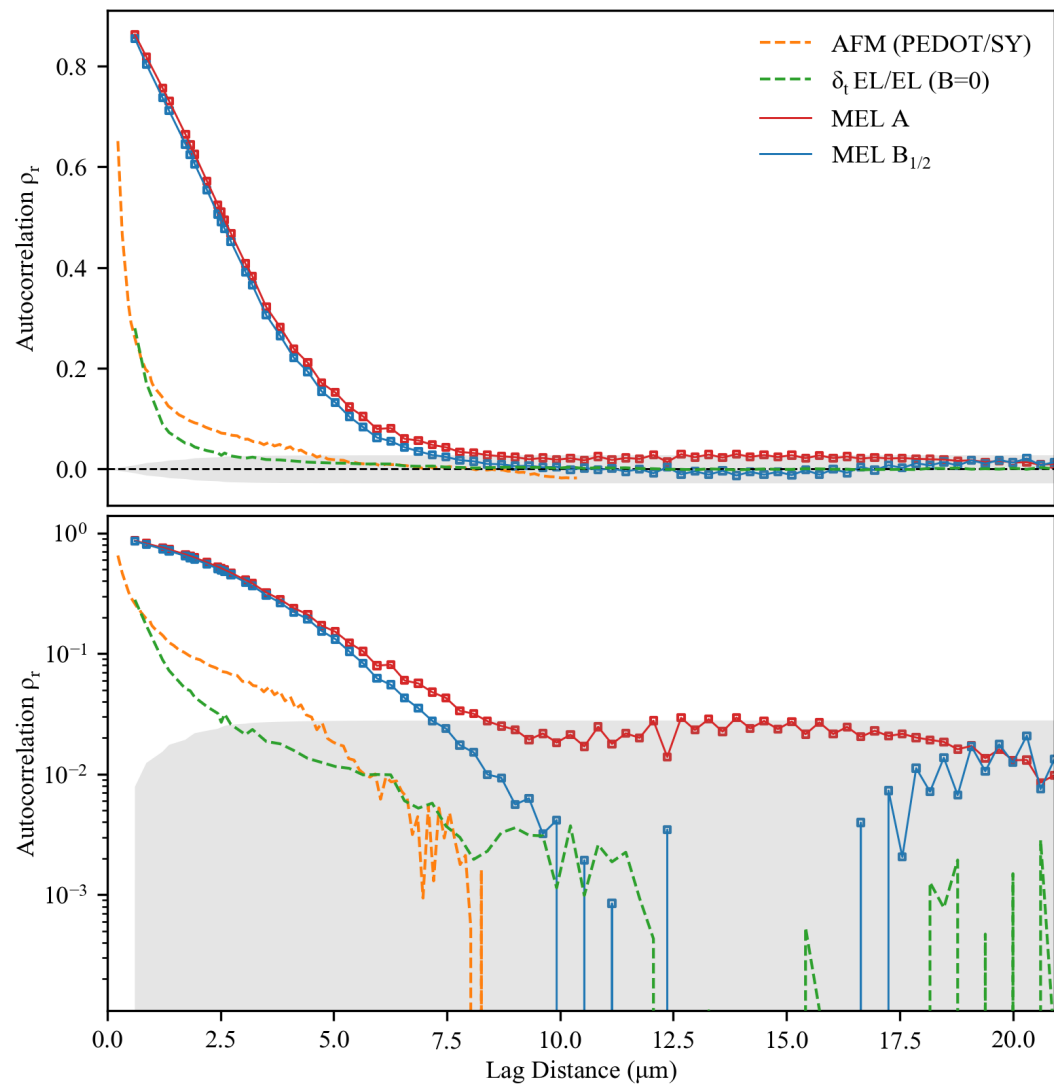


AFM (SY-PPV)

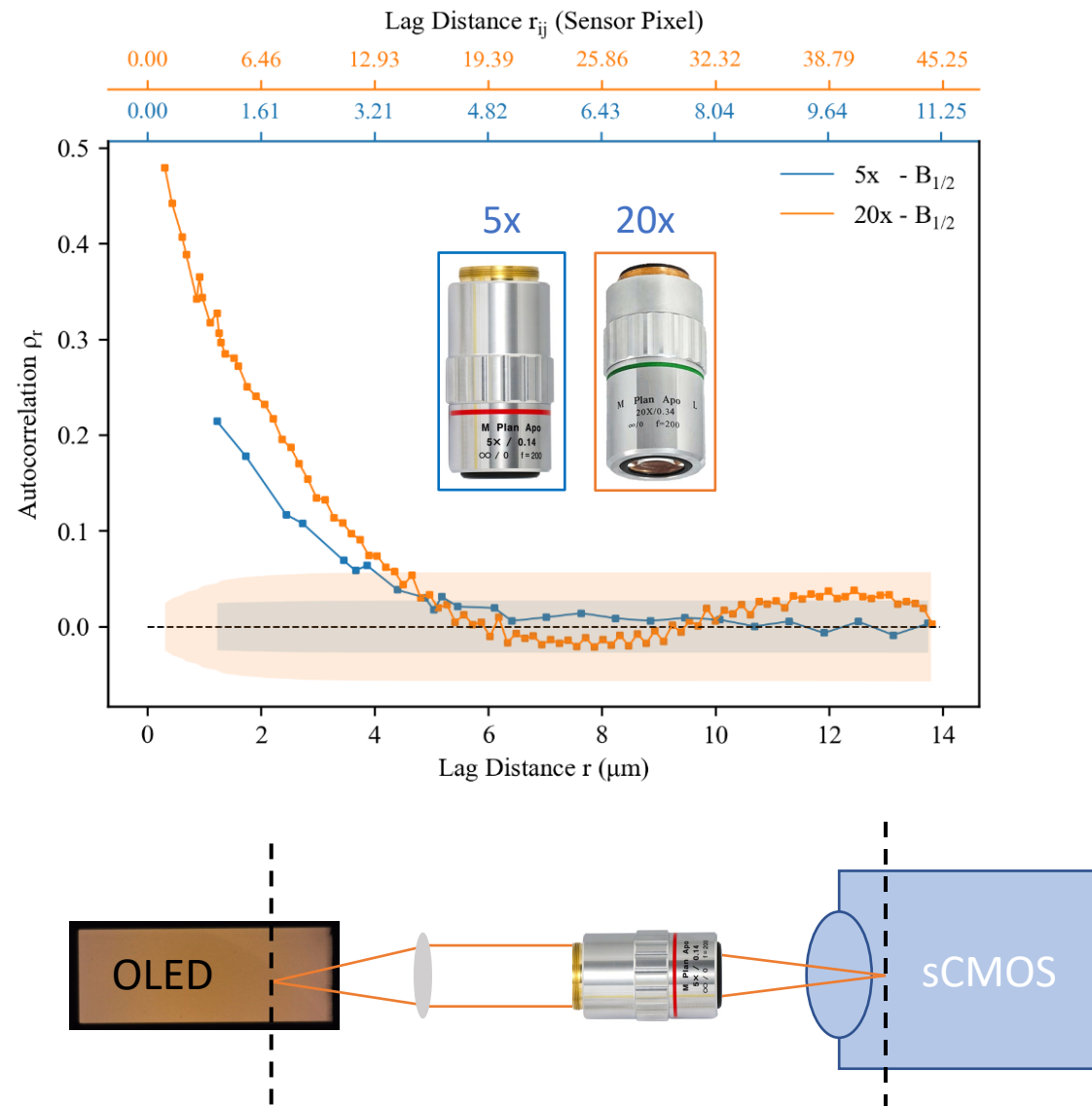


TEM

Separating spatially-correlated features

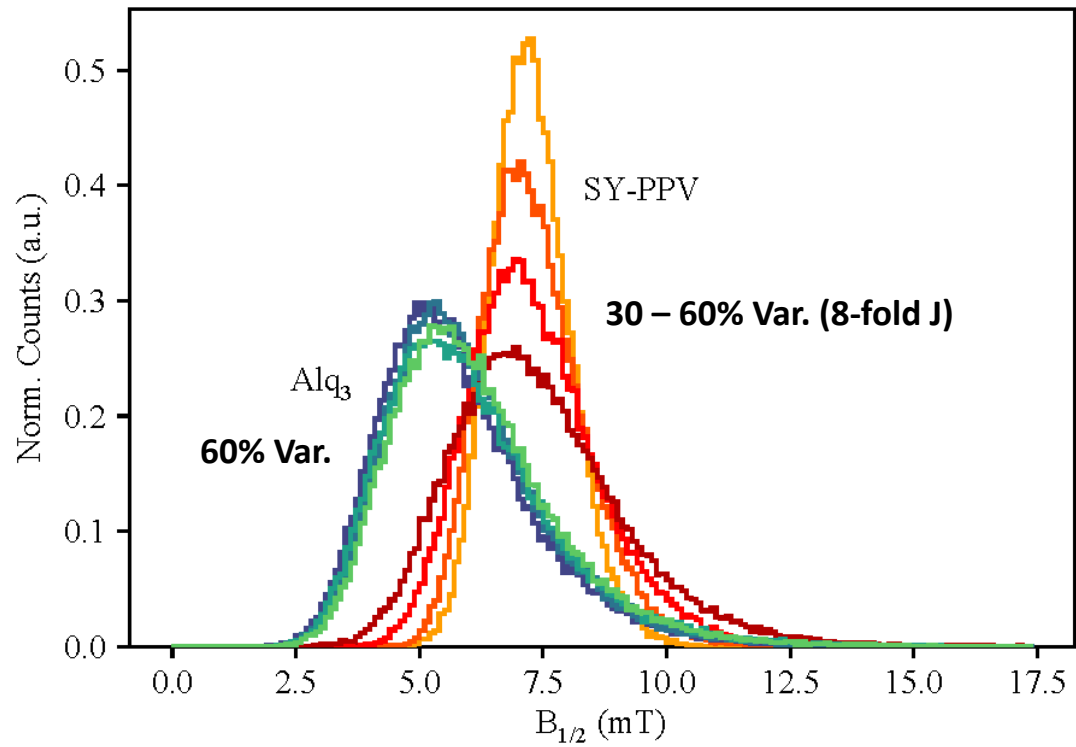


Verifying the length scales of correlation

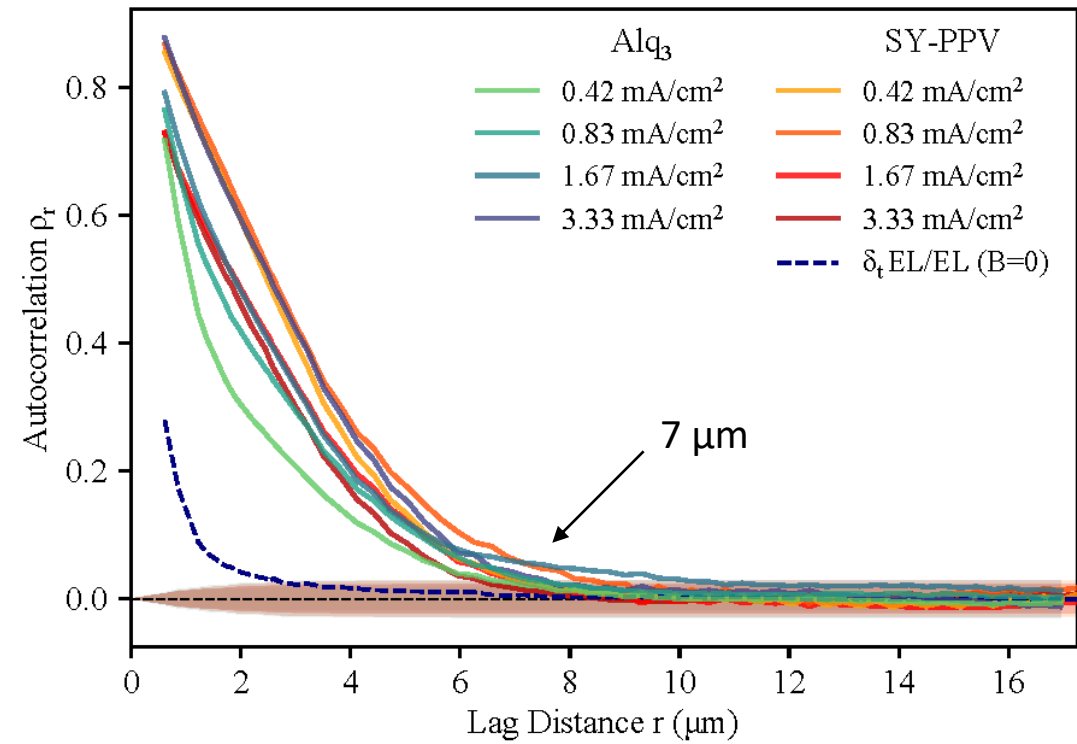


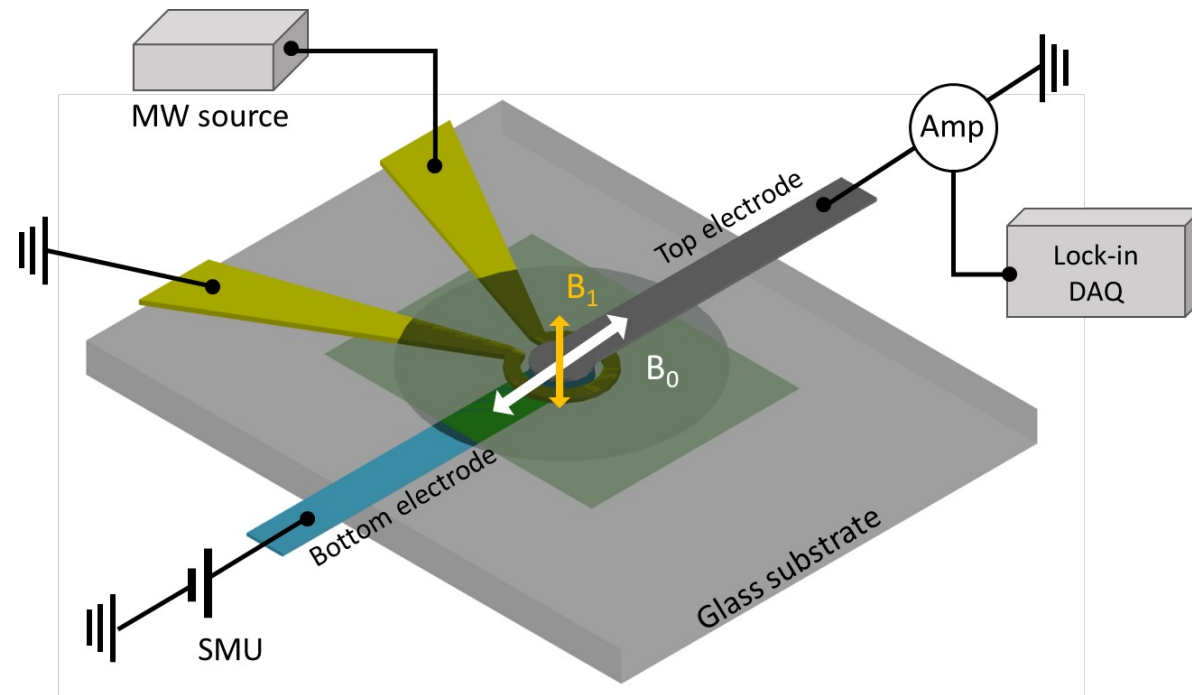
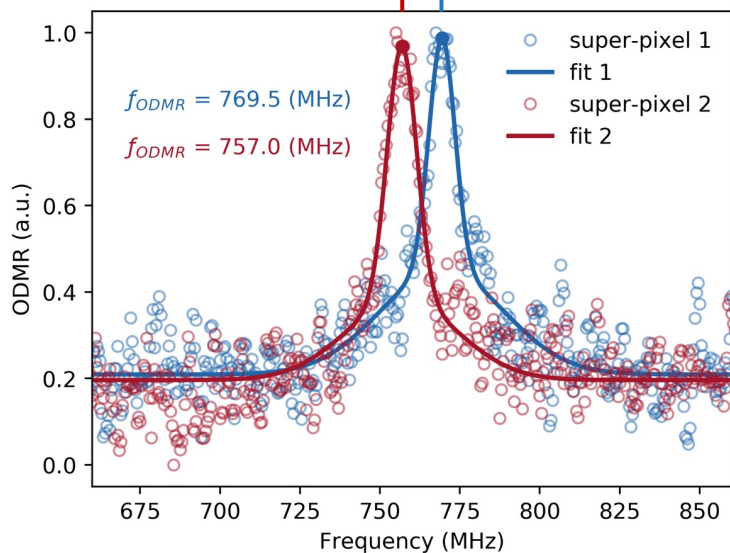
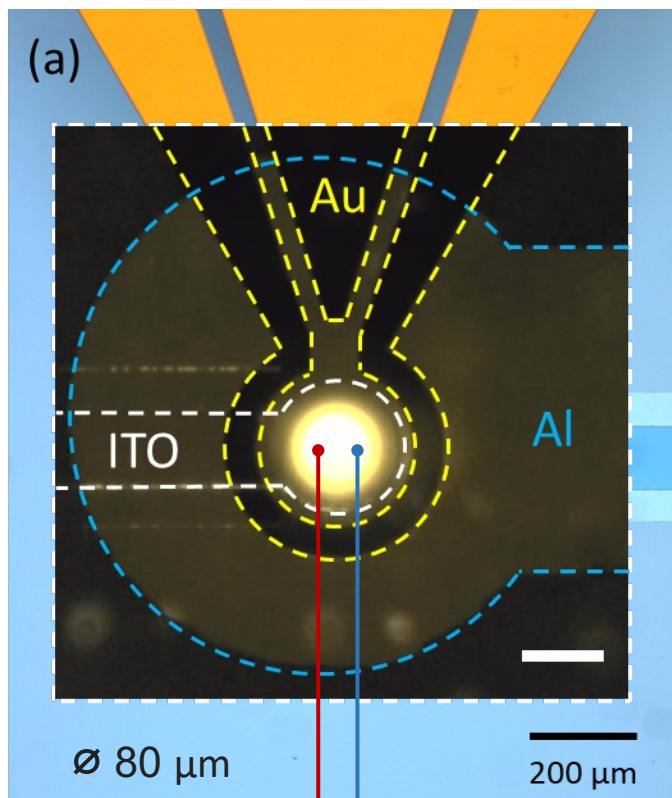
Spin Property Variation: Small Molecule (Alq_3) vs Polymer (SY-PPV)

Intra-device Variation ($B_{1/2}$)



Spatial Correlation ($B_{1/2}$)





Integrated 'on-chip' RF resonator for magnetic sensing

Prototype

Sensitivity: $\sim 160 \mu\text{T Hz}^{-1/2} \mu\text{m}^{-2}$

Improvements Underway

Better material magneto-response

Increase CMOS well depth

Improved coupling to B_1 via resonator design

Manuscript Under Review <https://arxiv.org/abs/2207.02479>



Dr. Rugang Geng

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Adrian Mena

Alex Baldacchino

Spin Research Group UNSW



UNSW
SYDNEY

Australian Research Council Centre of Excellence in

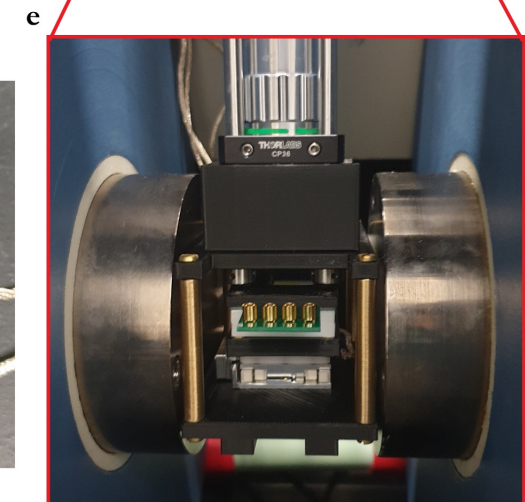
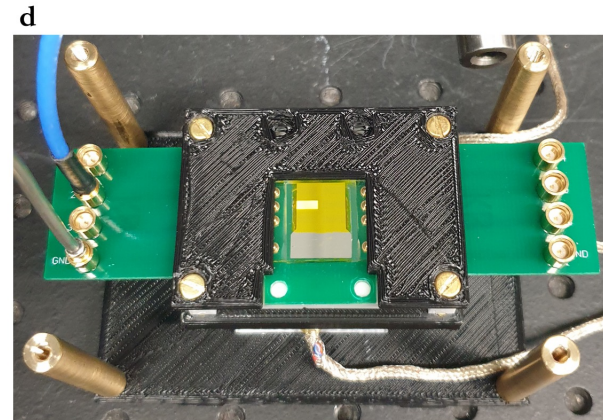
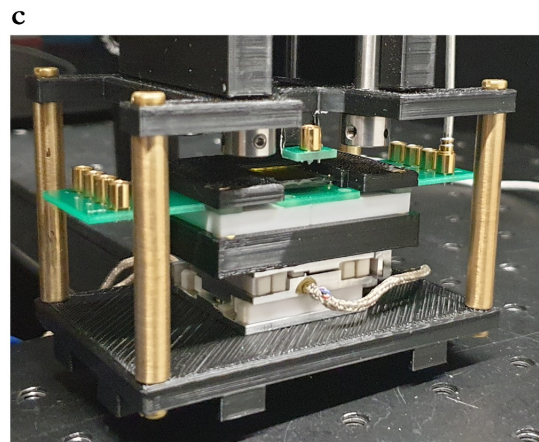
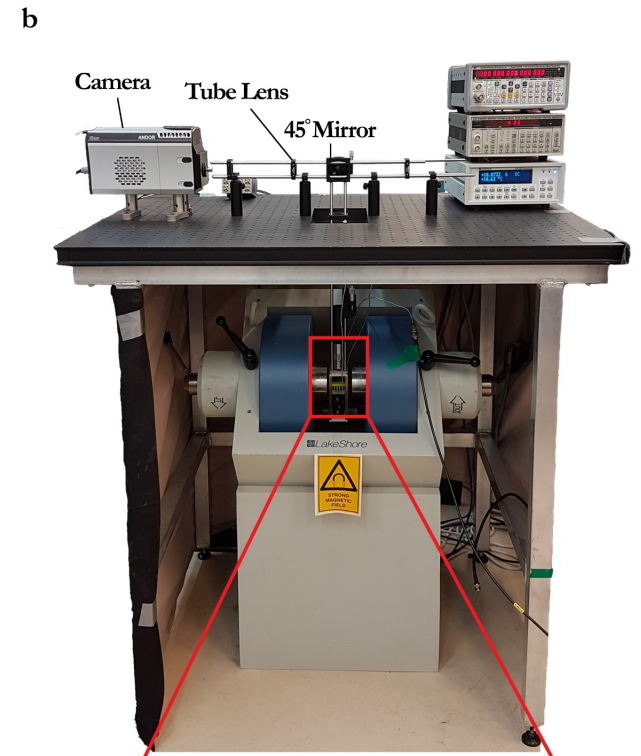
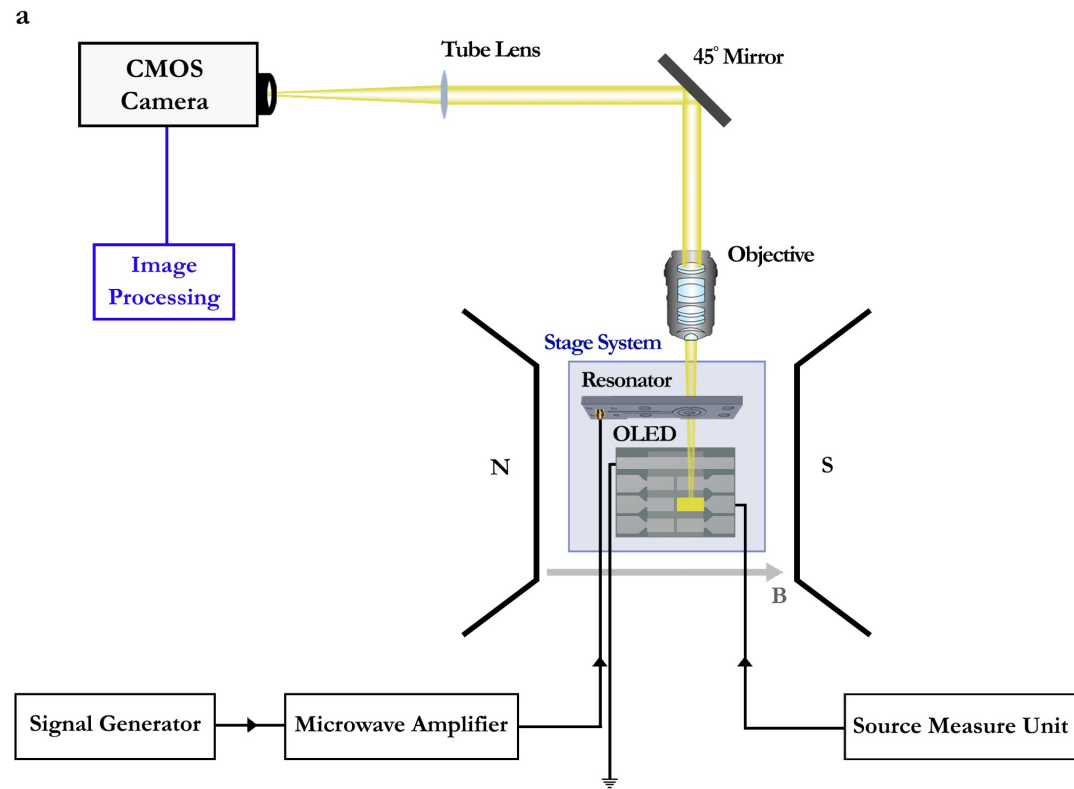
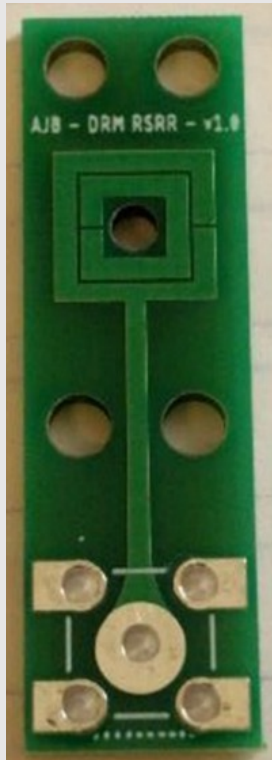


Australian Government
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Supplementary Slides

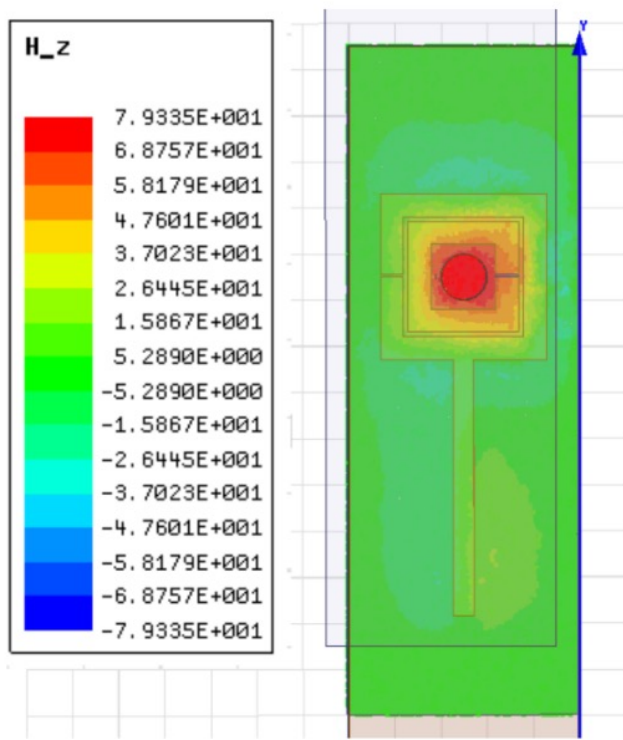
Spatiotemporal Variation

Optically Detected Magnetic Resonance

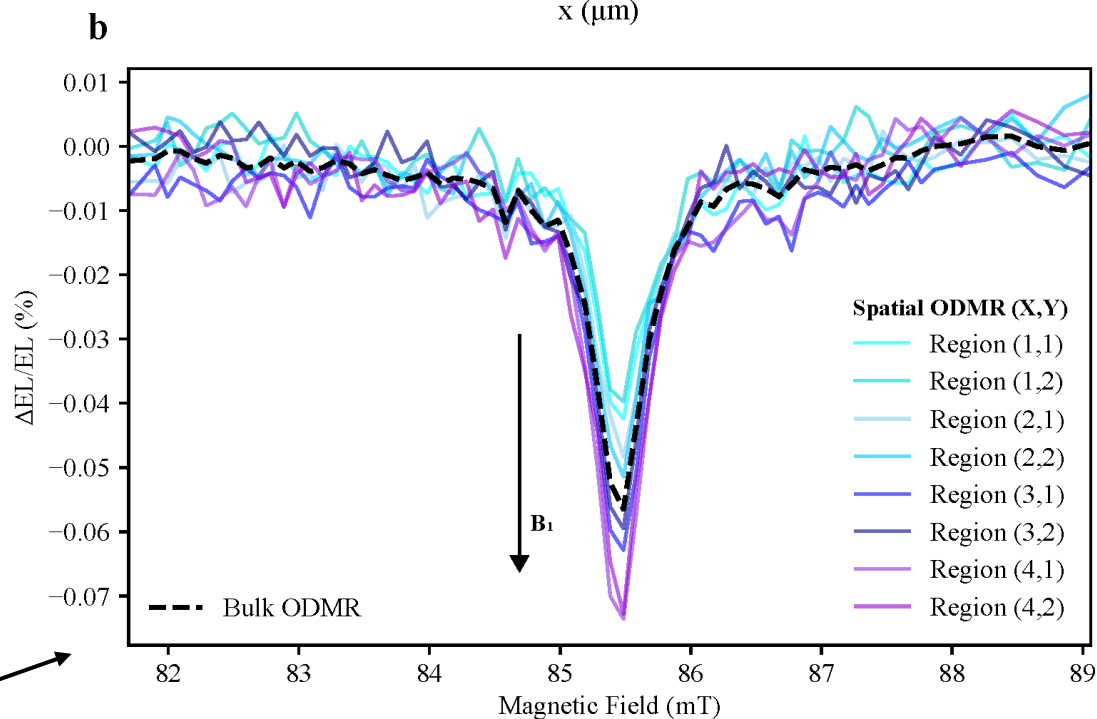
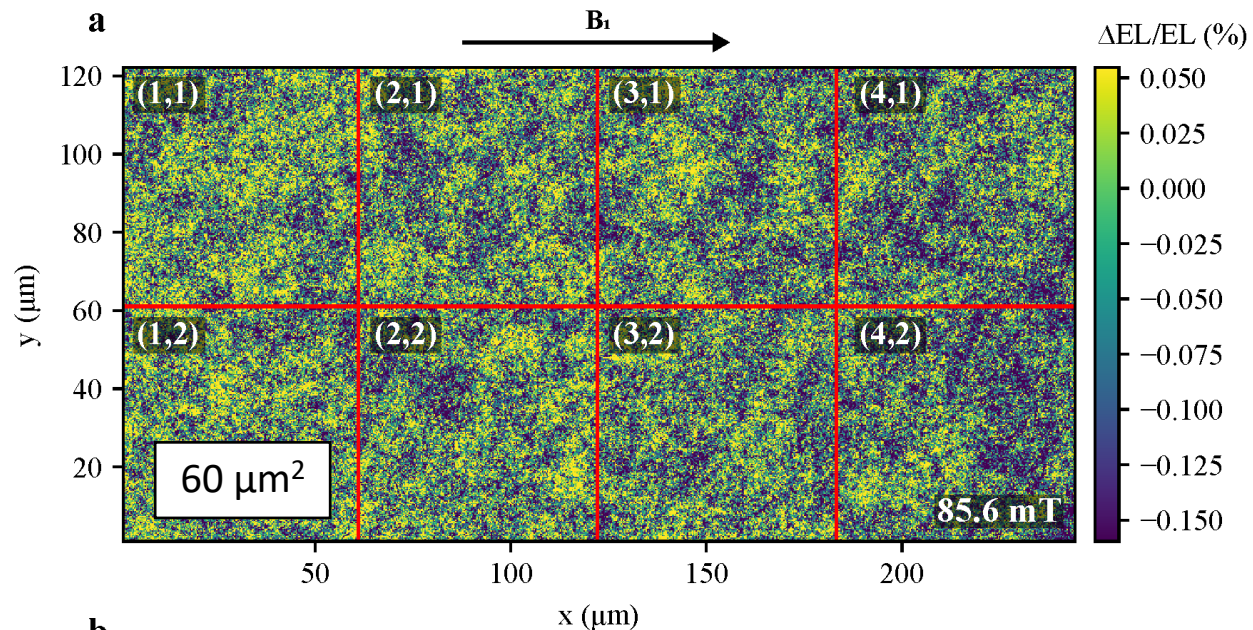


Spatially Resolved Magnetic Resonance

ODMR amplitude $\propto B_1$ magnitude



long time-averaged data

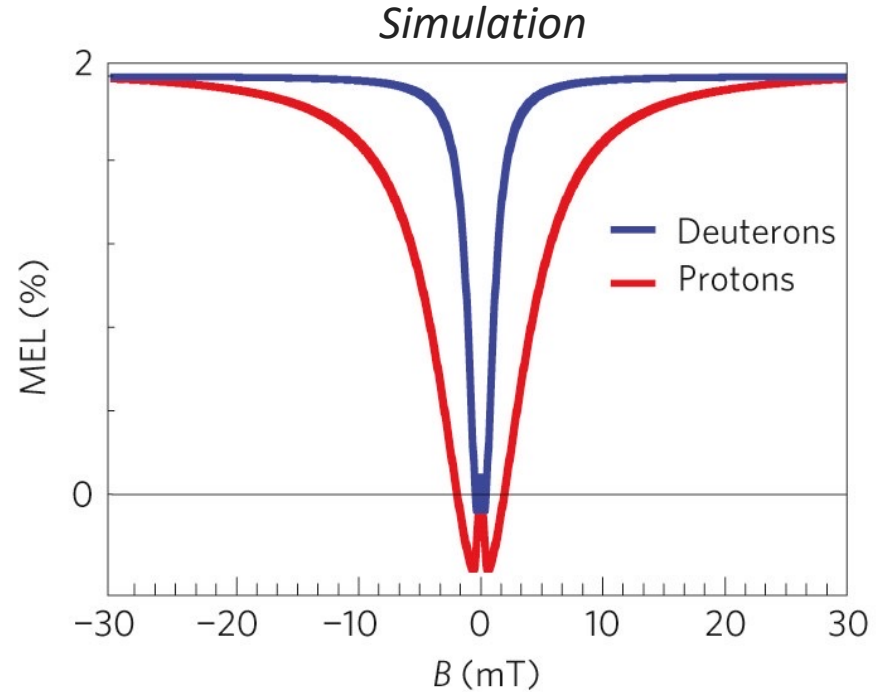
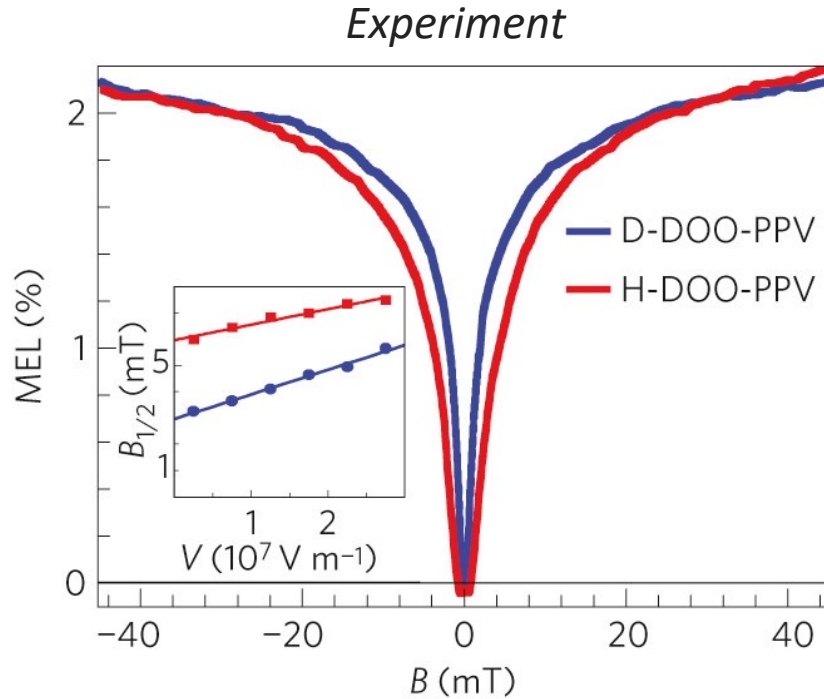


Spin mixing

$$\Delta\omega_{hf} + \Delta\theta$$

Isotope Exchange

replace all strongly coupled hydrogen atoms (^1H , $I=1/2$) with deuterium atoms (^2H , $I=1$)
 have much smaller a_{HF} , $a(\text{D}) = a(\text{H})/6.5$



half-width ($B_{1/2}$) scales with hyperfine interaction strength $\delta B_D \approx 0.5 \delta B_H$

Polymer (DOO-PPV)

T.D. Nguyen *et al.*, *Nature Mater* **9**, 345–352 (2010)

Small-Molecule (Alq_3)

T.D. Nguyen *et al.*, *Phys. Rev. B*, **85** (24), 245437 (2012)

$$\delta B_{hf} = a/\sqrt{N}$$

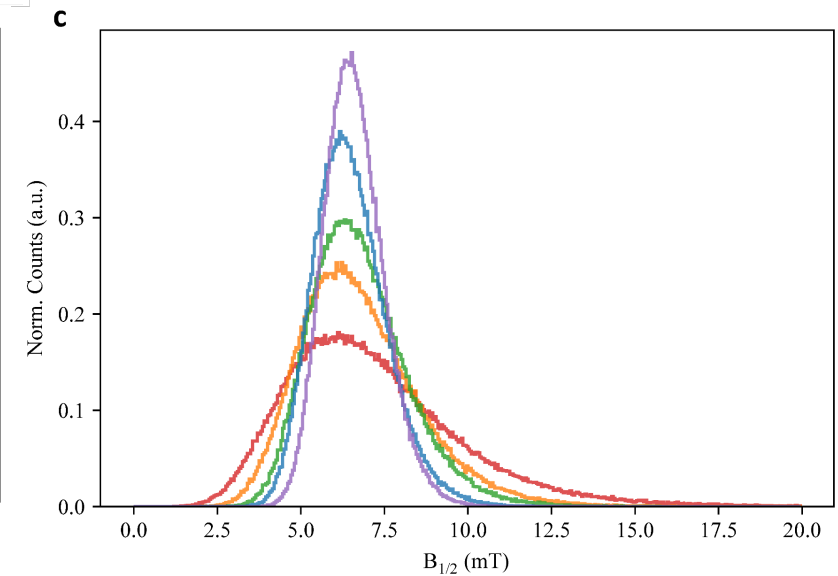
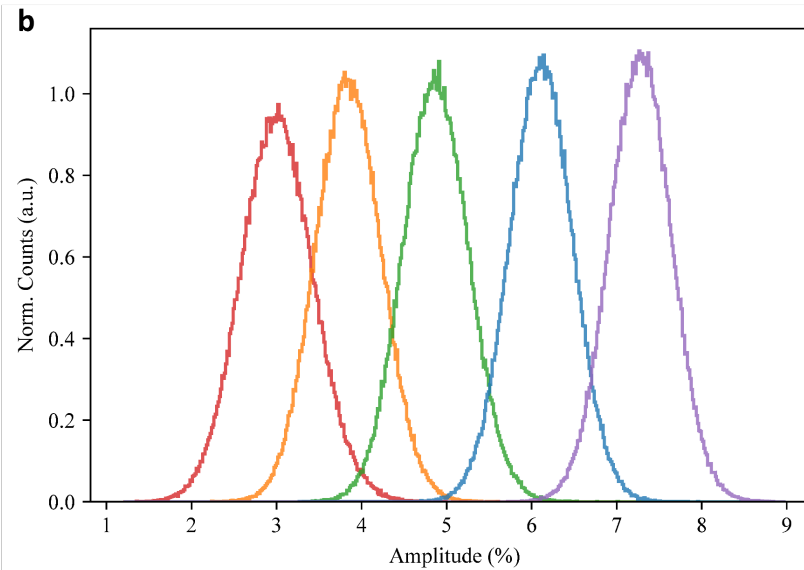
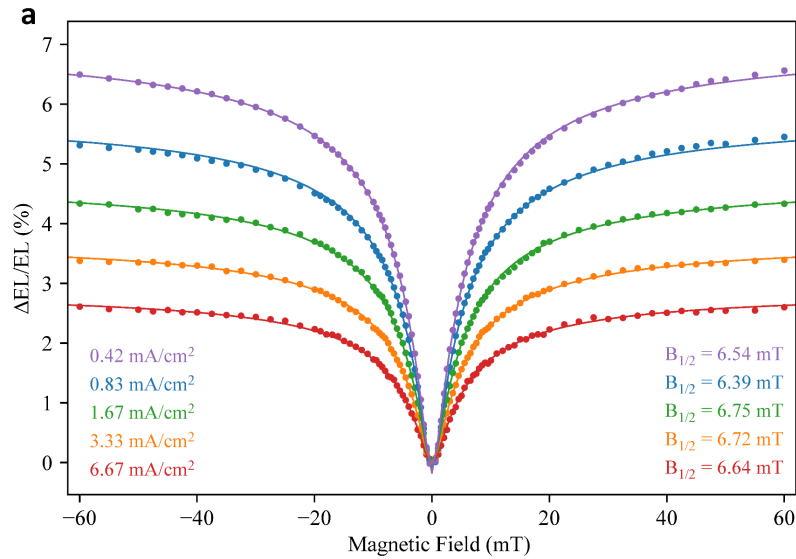
- Spin density
- Spin dephasing

Active Layer Thickness – No Dependence ($B_{1/2}$)

SY-PPV (1500 rpm / **110 nm**)

SY-PPV (3000 rpm / **70 nm**)

0.42 mA cm ⁻²	0.83 mA cm ⁻²	1.67 mA cm ⁻²	3.33 mA cm ⁻²	6.67 mA cm ⁻²
29.2	35.8	49.3*	63.0*	x
31.1	38.5	48.1*	58.8*	79.6*



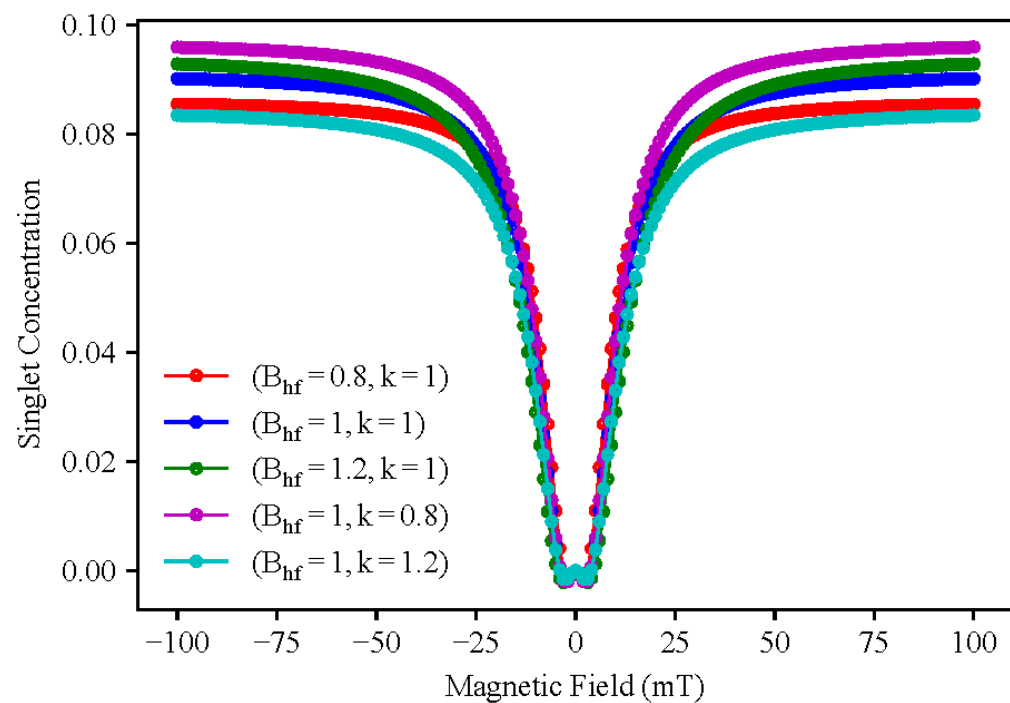
Histogram broadening probably not due to **exciton quenching** processes i.e. **efficiency roll-off**

Rate Variation vs. Hyperfine Variation

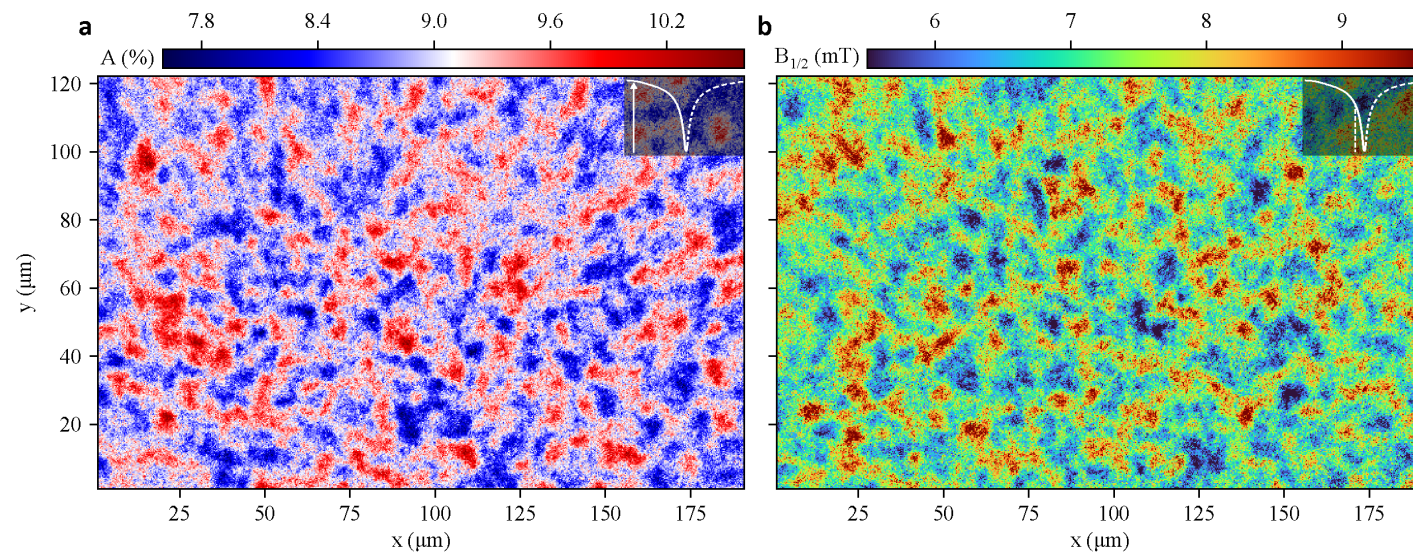
$$\chi_S = \frac{1}{2} \left\{ \frac{k_S \frac{1}{2} \sin^2 \theta/2}{k_S \frac{1}{2} \sin^2 \theta/2 + k_T \left[\frac{1}{2} \sin^2 \theta/2 + \cos^2 \theta/2 \right]} + \frac{k_S \frac{1}{2} \cos^2 \theta/2}{k_S \frac{1}{2} \cos^2 \theta/2 + k_T \left[\frac{1}{2} \cos^2 \theta/2 + \sin^2 \theta/2 \right]} \right\}$$

Kersten *et al.*, *Phys. Rev. Lett.* **106**, 197402 (2011)

a Singlet Yield at Different Mixing and Recombination Rates

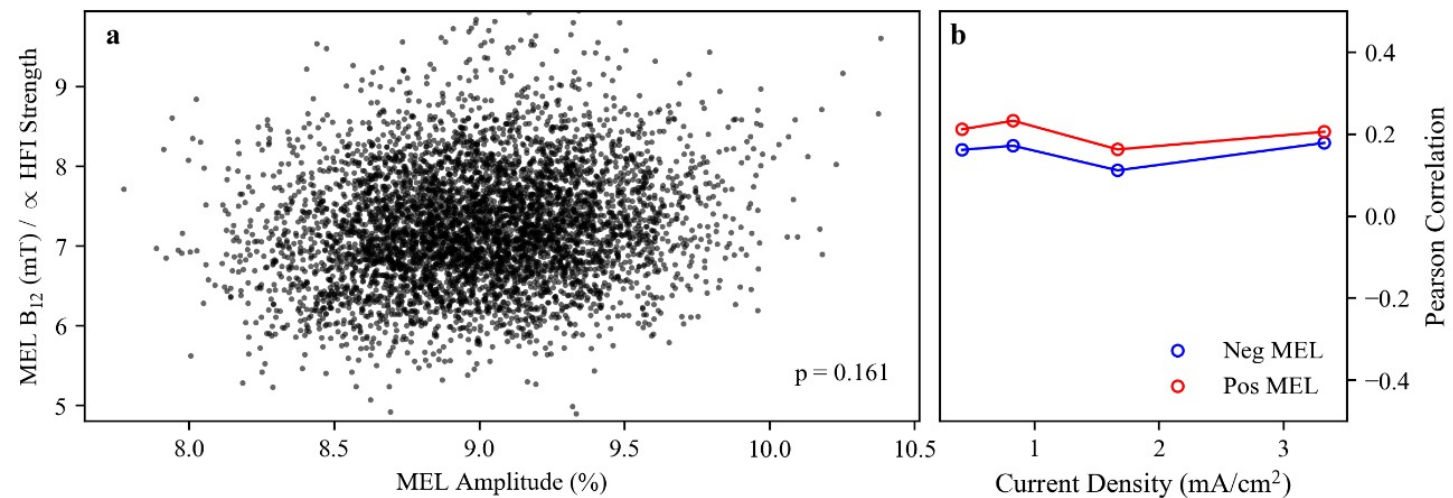


Rate Variation vs. Hyperfine Variation

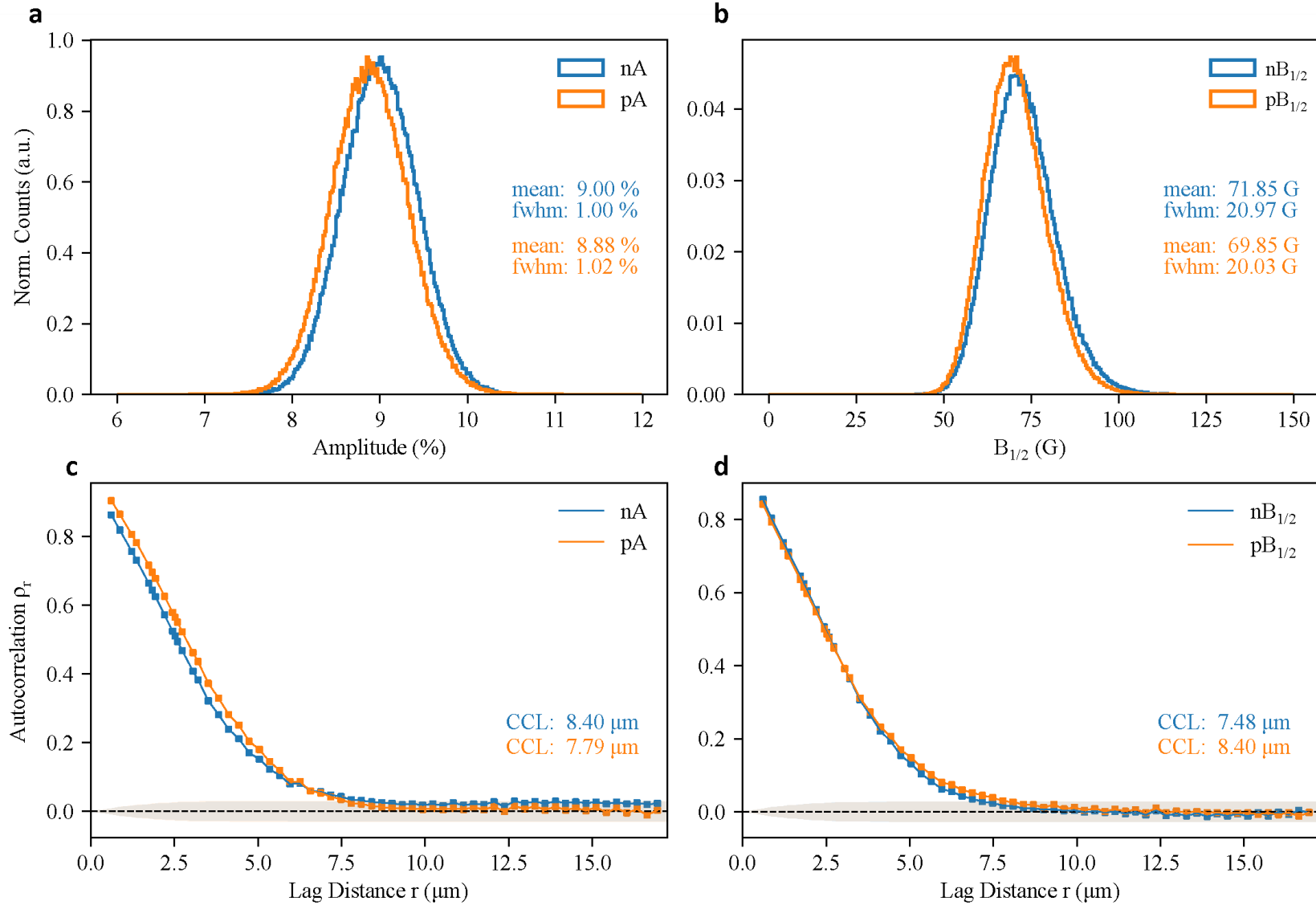


250k points

Correlation coefficients remain **positive** and **significant** (~ 0.2 with a 2-sided p-test $< 10^{-5}$)



Global Spatiotemporal Evolution: Neg / Pos Polarity Field Sweeps



Ruled Out

Spatial Correlation of MEL not adequately explained by:

- Aggregates / Film Homogeneity (AFM, TEM, MEL dynamics)
- Molecular Material (polymer vs small-molecule)
- Exciton Quenching Processes (active layer thickness)
- Spatial Coherence of OLED (time-correlated EL maps)
- Recombination Rate Variation (positive Pearson A vs $B_{1/2}$)

Contributing factors:

- Hyperfine Variation + Diffusion (MEL $B_{1/2}$, A- $B_{1/2}$ correlation, dynamics)
- Exchange Variation (Spin-mixing variation, range of r_{e-h})
- Hopping Rate Variation (MEL lineshape dependence)
- Interface Effect (Bulk vs. Active-LiF/Al effect)

Possible Sources

