

Optimisation of electron spin qubits in electrically driven multi-donor quantum dots

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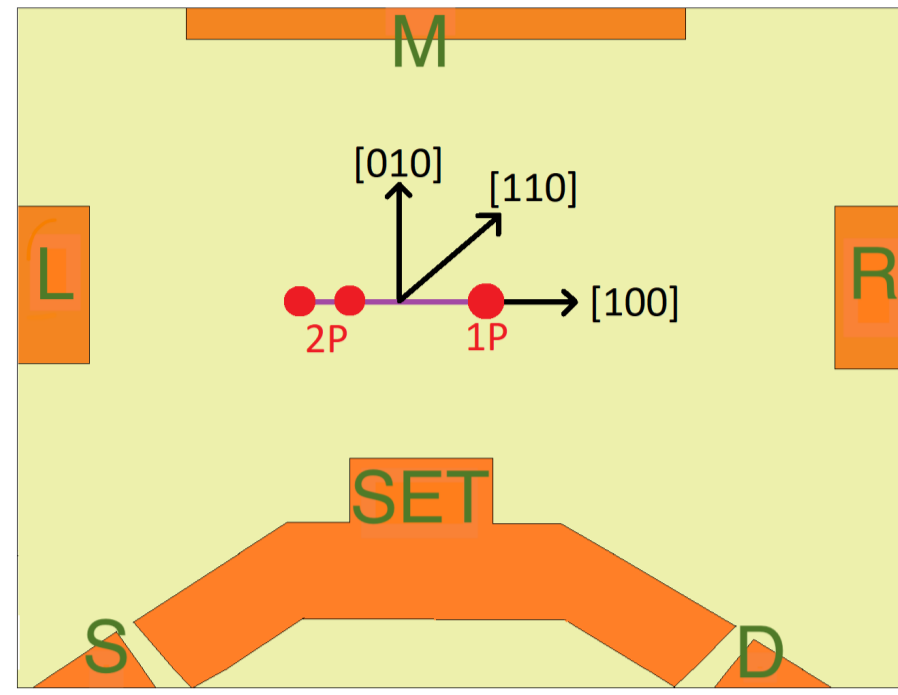
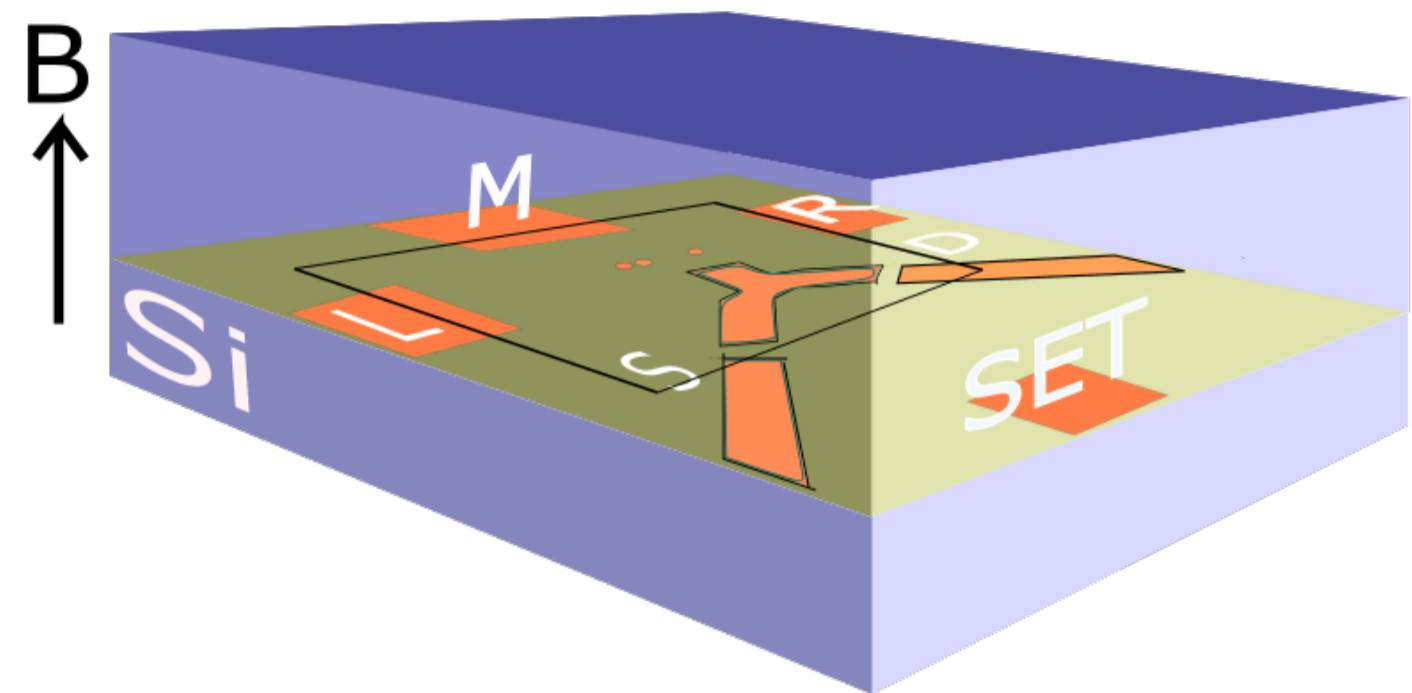


Fig 1: a) STM imaging of the qubit, b) in-plane gate control

- Proposal**
- 2P (phosphorus double donor): 1P (single donor) spin qubit
 - Describe the theory of the multi donor quantum dot: wave function, electrical operation, noise
 - Electron dipole spin resonance (EDSR)
 - Minimal computational expense: effective theory.

Motivation • 2P:1P qubit with Rabi freq. 1.2 MHz, $T_2^* = 295$ ns, 0.8 ns two-qubit gate¹

- Donors offer atomic-scale precision placement, orientation specified properties (e.g. [110] qubit axis sees suppression of exchange oscillation)
- Scalable semiconductor multi-qubit chips possible.²
- Hyperfine interaction between the qubit electron spin and P-donor nucleus provides strong spin-orbit coupling (advantage over micromagnet engineering): $H_{hf} = \mathcal{A}_0 \sum_i \delta(\mathbf{r} - \mathbf{R}_i) \vec{I}_i \cdot \vec{S}$ ³

Fig.2: a) STM imaging of 2P-1P qubit, b) Rabi oscillation, c) hyperfine mediated EDSR

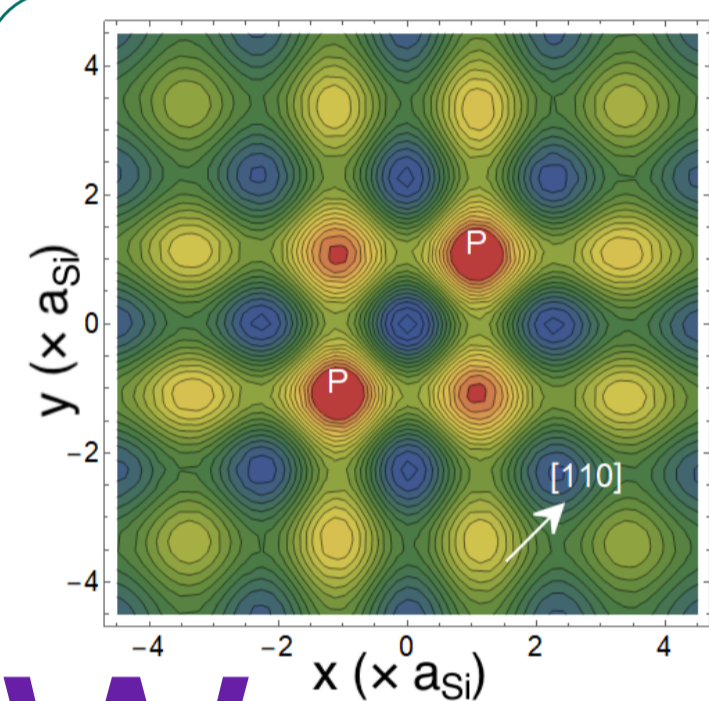
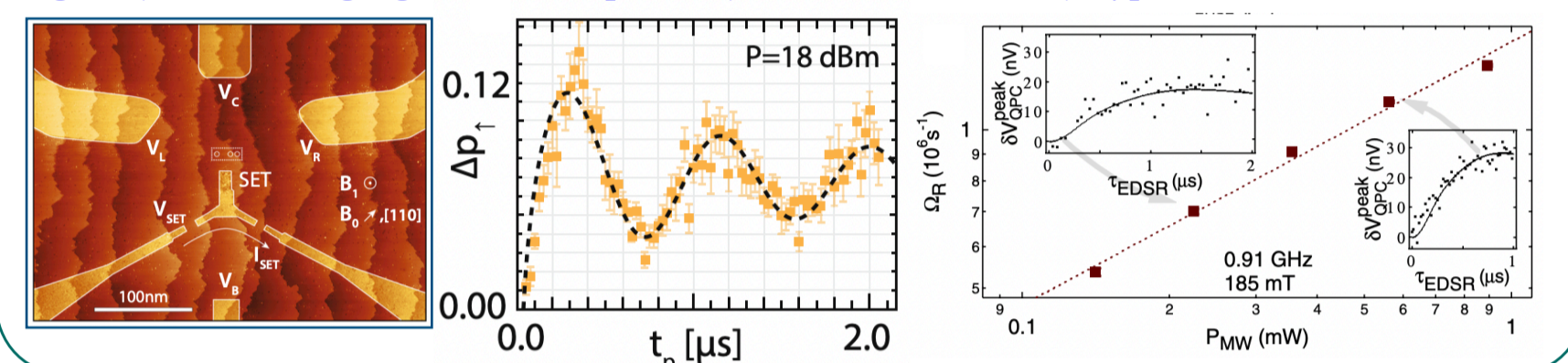


Fig.3: Theoretical result for the 2P wave function from the EMA approach

- Wavefunction**
- 2P: effective mass theory, Slater's formalism⁴
 - Indirect bandgap in Si leads to multi valley physics, LCAO method gives ground-state wave function for double donor QD, consisting symmetric combination of valleys from the two donors. $\Psi_{2D}(\mathbf{r}) = \frac{1}{\sqrt{\sum_{\xi} w_{\xi}^2}} \sum_{\xi} w_{\xi} S_{\xi} u_{\xi}(\mathbf{r})$

EDSR • 2P-1P overlap introduces tunneling of e^- between QDs.

- The 2P:1P GS is obtained using Hund-Mulliken approach.
- A global magnetic field resolves the spin-states into \uparrow and \downarrow states.
- Hyperfine interaction : $(1,0) \uparrow \leftrightarrow (0,1) \downarrow$
- Rabi ratio: $T_1/T_{\pi} \sim 10^6$, million operations in one relaxation time.
- When an ac $\tilde{E}(t)$ is applied, spin-flip takes place owing to time-dependent modulation of difference in hyperfine between left and right dots. (EDSR)

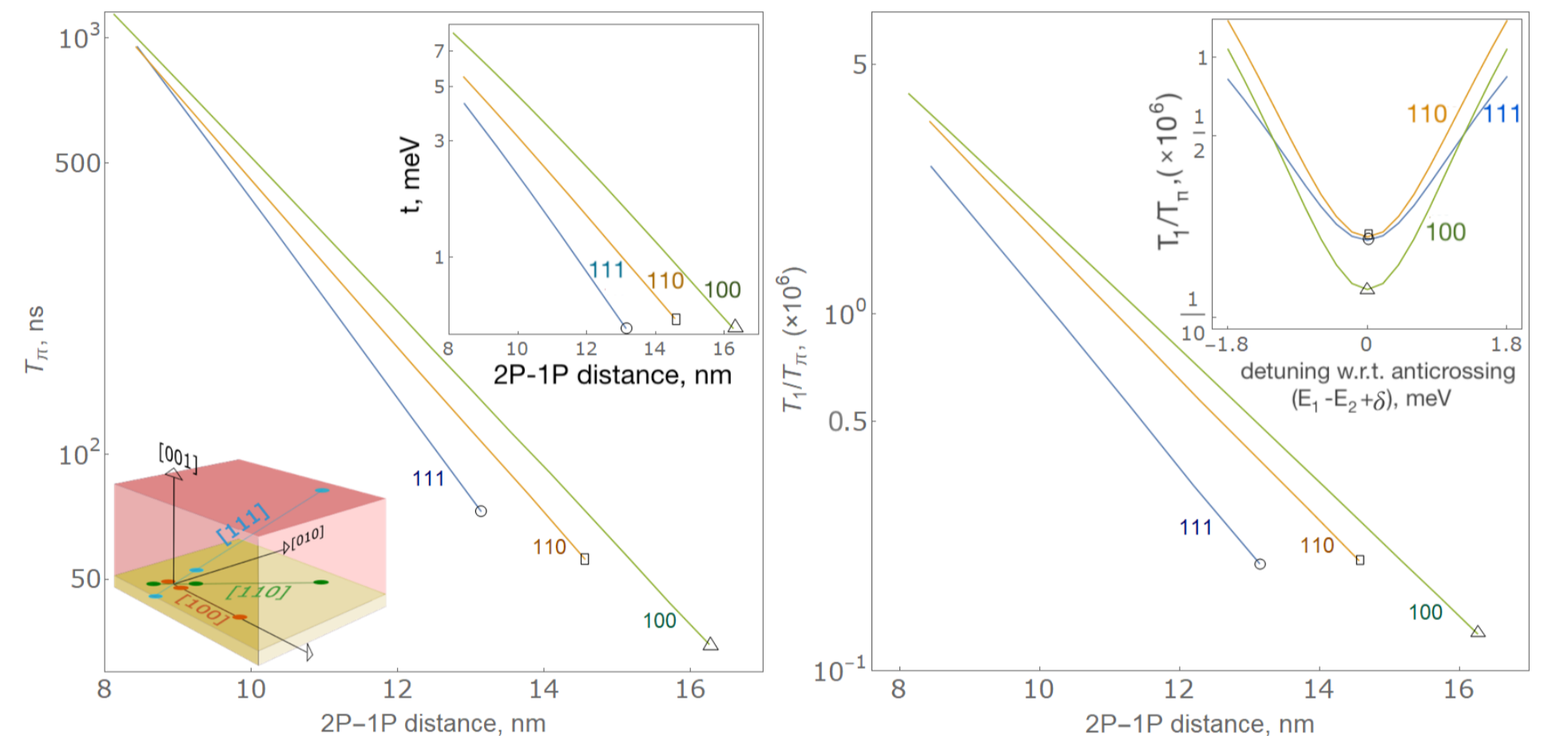


Fig.4: a) EDSR Rabi time vs 2P-1P distance for different qubit geometry, b) Rabi ratio vs 2P-1P distance

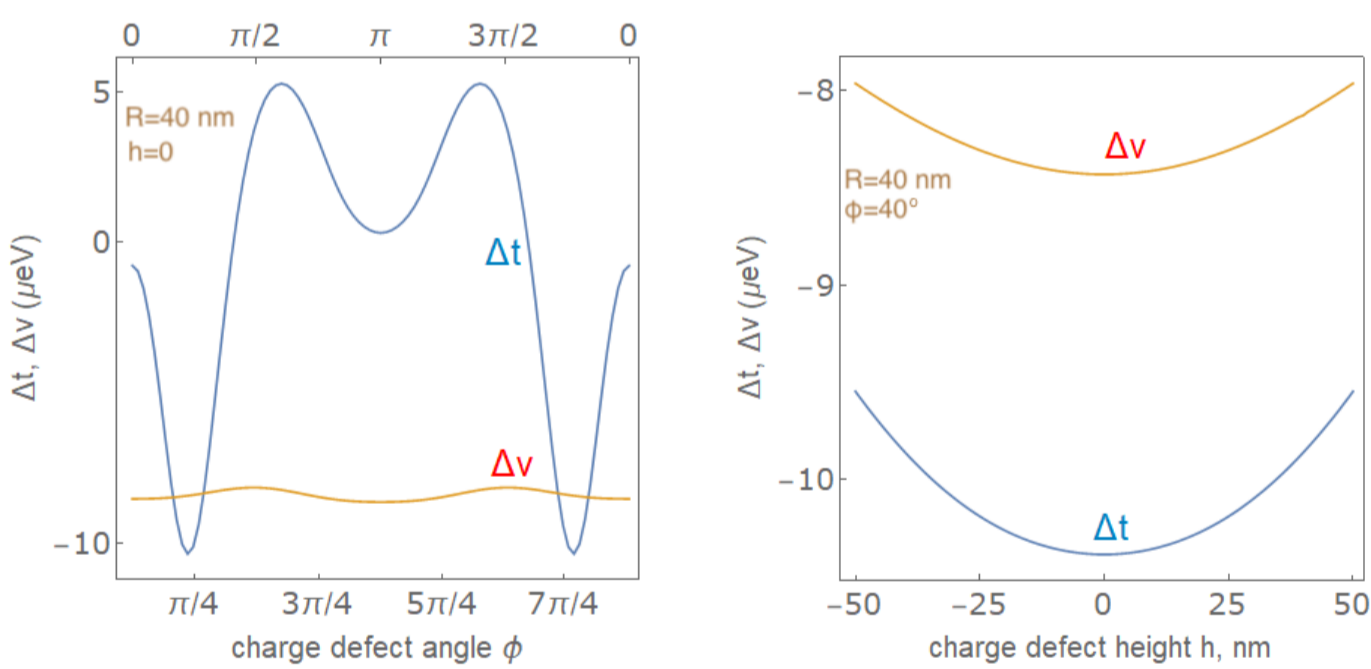
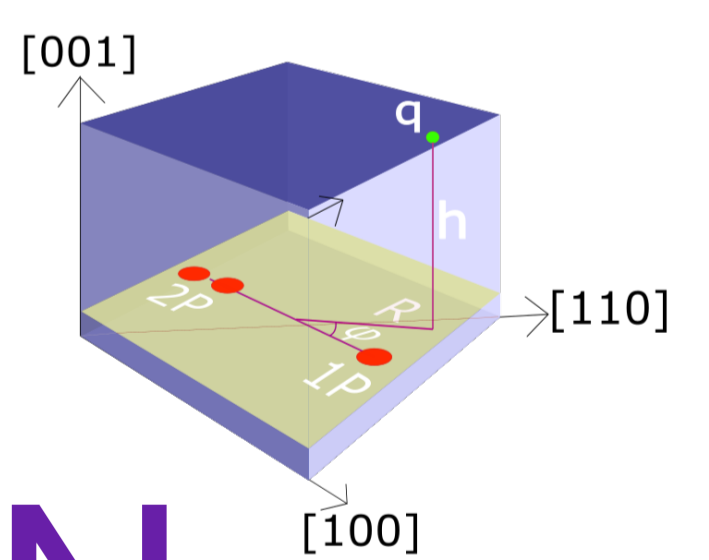


Fig.5: a) charge defect schematic b) angular noise effect of charge noise c) noise effect vs. height of charge defect.

- Noise**
- Charge defect in the vicinity of the physical qubit can modify the tunneling between the 2P and 1P dots as well as the self-energies of the QDs.
 - Such charge noise lead to decoherence of the qubit state $\uparrow \rightarrow \downarrow$.
 - Time of decoherence (T_2^*) from charge noise can be modelled as Random Telegraph noise: $\frac{1}{T_2^*} \Big|_{RTN} = \frac{32\zeta^2 \tau k_{ge}^4}{\hbar^2 \delta \epsilon^4} \left(\frac{4t\Delta t}{\delta \epsilon^2} + \frac{\epsilon \Delta v}{\delta \epsilon^2} \right)^2$
 - Noise is the most detrimental at particular angular orientation, and is minimal at the qubit plane, supporting the in-plane gated control.

- Perorations**
- 2P ground state energy of -121 meV, well-benchmarked against tight binding (TB), NEMO numerical models.⁵⁻⁷
 - 2P ground state valley weights explained by competing effects: effective mass anisotropy (EMA) and valley-orbit coupling (VOC).

- Complete in-plane gate control is possible, under a global out-of-plane magnetic field \vec{B}_z an in-plane ac electric field $\vec{E}(t)$ induces qubit spin-flip.
- The spin-flip EDSR mechanism relies on contact hyperfine coming from the donor nuclear spins.
- $T_{\pi} \sim 50$ ns, fast qubit rotation; $T_1/T_{\pi} \sim 10^6$, excellent Rabi ratio.
- Qubit geometry determines the best EDSR for when qubit axis \parallel [111] and best Rabi ratio for when qubit axis \parallel [100].
- Noise simulation produces the detrimental defect position to be $\theta = 40^\circ$, $h=0$, $R=30$ nm.
- Charge noise leads to EDSR single qubit gate error of 2%
- Efficient two qubit exchange gate of 3 GHz.

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