

Dynamical branes on expanding orbifold and complex projective space

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arXiv:2211.13501 [hep-th]

[1] Introduction and our results

- We construct a dynamical p -brane solution on orbifold (the complex line bundle over CP^n space) as a solution.
- There are interesting properties when branes are located at an orbifold.
(M.R. Douglas and G.W. Moore, hep-th/9603167)
(M.R. Douglas, B.R. Greene and D.R. Morrison, Nucl. Phys.B 506 (1997) 84 [hep-th/9704151])
- ~ Orbifold singularities are resolved in the D-brane world volume theory.
- ~ One can expect that the spacetime itself becomes regular without any naked singularity once p -branes are placed at the orbifold singularities.



- This is important when one finds hints that the gravity theory giving a p-brane on the orbifold is known and well-defined supergravity, thus giving a handle on the strong coupling dynamics of string theory.
- **Black hole on the Eguchi-Hanson space**
(H. Ishihara, M. Kimura, K. Matsuno and S. Tomizawa, Phys. Rev. D 74 (2006) 047501 [hep-th/0607035])
(Hideki Ishihara, Masashi Kimura, Shinya Tomizawa, Class.Quant.Grav.23 (2006) L89 [hep-th/0609165])
(Chul-Moon Yoo, Hideki Ishihara, Masashi Kimura, Ken Matsuno, Shinya Tomizawa, Class.Quant.Grav. 25 (2008) 095017 [0708.0708 [gr-qc]])
- **Black holes on the complex line bundle over CP^n space**
(M. Nitta, K. Uzawa, Eur.Phys.J.C 81 (2021) 6, 513, [arXiv:2011.13316 [hep-th]])
- **(Static) Black p-brane on the orbifolds**
(M. Nitta, K. Uzawa, JHEP 03 (2021) 018, [arXiv:2012.13285 [hep-th]])

[2] Dynamical p-brane on orbifolds

- Metric of dynamical p-brane in D dimensions

(Pierre Binetruy, Misao Sasaki, Kunihito Uzawa, Phys.Rev.D 80 (2009) 026001, 0712.3615 [hep-th])

(Kei-ichi Maeda, Nobuyoshi Ohta, Kunihito Uzawa, JHEP 06 (2009) 051, 0903.5483 [hep-th])

(Masato Minamitsuji, Nobuyoshi Ohta, Kunihito Uzawa, Phys. Rev. D 82 (2010) 086002)

$$ds^2 = h^a(x, y) q_{\mu\nu}(X) dx^\mu dx^\nu + h^b(x, y) u_{ij}(Y) dy^i dy^j,$$

$$a = -\frac{D - p - 3}{D - 2}, \quad b = \frac{p + 1}{D - 2},$$

 orbifold

- (p+2)-form field strength

$$F_{(p+2)} = d(h^{-1}) \wedge dt \wedge dx^1 \wedge \dots \wedge dx^p$$

- It is straightforward to check that with such an ansatz the Bianchi identity is trivially satisfied.
- The field equation for the antisymmetric tensor becomes

(Kei-ichi Maeda, Nobuyoshi Ohta, Kunihiro Uzawa, JHEP 06 (2009) 051)

(Masato Minamitsuji, Nobuyoshi Ohta, Kunihiro Uzawa, Phys. Rev. D 82 (2010) 086002)

$$\partial_\mu \partial_i h = 0, \quad \Delta_Y h = 0,$$

$$\rightarrow h(x, y) = h_0(x) + h_1(y), \quad \Delta_Y h_1 = 0,$$

- Einstein equations

$$R_{\mu\nu}(X) - h^{-1} D_{\mu} D_{\nu} h_0 - \frac{a}{2} h^{-1} q_{\mu\nu} (\Delta_X h_0 + h^{-1} \Delta_Y h_1) = 0,$$

$$R_{ij}(Y) - \frac{b}{2} q_{\mu\nu} (\Delta_X h_0 + h^{-1} \Delta_Y h_0) = 0,$$

$$\Rightarrow R_{\mu\nu}(X) = 0, \quad R_{ij}(Y) = \frac{1}{2} b (p + 1) \lambda u_{ij}(Y), \quad D_{\mu} D_{\nu} h_0 = \lambda q_{\mu\nu}(X),$$

- The space Y is not Ricci flat, but the Einstein space such as CP^n if $\lambda \neq 0$, and the function h can be more non-trivial.

(Pierre Binétruy, Misao Sasaki, Kunihiro Uzawa, Phys.Rev.D 80 (2009) 026001 0712.3615 [hep-th])

(Kei-ichi Maeda, Nobuyoshi Ohta, Kunihiro Uzawa, JHEP 06 (2009) 051, arXiv: 0903.5483 [hep-th])

If we set

$$q_{\mu\nu}(X) = \eta_{\mu\nu}(X), \quad D_\mu h_0 \neq 0, \quad (D_\mu h_0)(D^\mu h_0) \neq 0,$$

the solution for h_0 is given by

$$h_0(x) = \frac{\lambda}{2} x_\mu x^\mu + \bar{a}_\mu x^\mu + \bar{a}$$

- When the space Y is Ricci flat like orbifold, the function h_0 is linear in the coordinates x^μ because of $\lambda=0$.

◆ Dynamical p-brane on the orbifold

cf) Black holes on Eguchi-Hanson space :

(H. Ishihara, M. Kimura and S. Tomizawa,
Class.Quant.Grav. 23 (2006) L89 [hep-th/0609165])

Eguchi-Hanson space – complex line bundle over CP^1 (2-sphere)

- p-brane on orbifold

$$u_{ij}(Y)dy^i dy^j = dr^2 + r^2 \left[\left\{ d\rho + \sin^2 \xi_{n-1} \left(d\psi_{n-1} + \frac{1}{2(n-1)} \omega_n \right) \right\}^2 + ds_{CP^{n-1}}^2 \right]$$

If we impose $h_1 = h_1(r)$, the field equations become for

$$\Delta_Y h_1(r) = 0, \quad \Rightarrow \quad h_1(r) = b_1 + \frac{b_2}{r^{D-p-3}}$$

the solution for $h(x, r)$ is given by

$$h(x, r) = \bar{a}_\mu x^\mu + b_1 + \frac{b_2}{r^{D-p-3}}$$

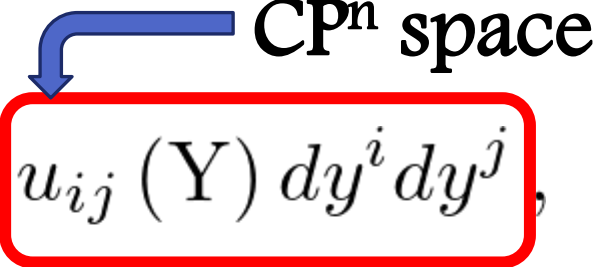
- There is a naked singularity at $h=0$.

(1) $r \rightarrow 0$: static p-brane solution.

(2) $r \rightarrow \infty$: asymptotically Kasner geometry

[3] Dynamical solution on the $\mathbb{C}\mathbb{P}^n$ space

- Metric

$$ds^2 = h^a(x, y) q_{\mu\nu}(X) dx^\mu dx^\nu + h^b(x, y) u_{ij}(Y) dy^i dy^j,$$


$$a = -\frac{D - p - 3}{D - 2}, \quad b = \frac{p + 1}{D - 2},$$

- $\mathbb{C}\mathbb{P}^n$ is the Einstein space (NOT Ricci flat).
- Our universe depends on the quadratic function of time.

For the \mathbb{CP}^1 space

$$ds_{\mathbb{CP}^1}^2 = (1 + \tilde{r}^2)^{-2} (d\tilde{r}^2 + \tilde{r}^2 d\tilde{\theta}^2)$$

• Solution

$$h(x, \tilde{r}, \tilde{\theta}) = \frac{\lambda}{2} x_\mu x^\mu + \bar{a}_\mu x^\mu + \tilde{c}_1 \ln \tilde{r} + \tilde{c}_2 \tilde{\theta} + \tilde{c}_3$$

For the \mathbb{CP}^2 space

$$ds_{\mathbb{CP}^2}^2 = (1 + \bar{\rho}^2)^{-2} d\bar{\rho}^2 + \frac{\bar{\rho}^2}{4} (1 + \bar{\rho}^2)^{-2} (d\psi + \cos \theta d\phi)^2 \\ + \frac{\bar{\rho}^2}{4} (1 + \bar{\rho}^2)^{-1} (d\theta^2 + \sin^2 \theta d\phi^2)$$

• **Solution**

$$h(x, \rho, \theta) = \frac{\lambda}{2} x_\mu x^\mu + \bar{a}_\mu x^\mu + \bar{c}_1 \left(-\frac{1}{2\bar{\rho}^2} + \ln \bar{\rho} \right) + \bar{c}_2 \ln \tan \frac{\theta}{2} + \bar{c}_3$$

[4] Discussion and remarks

(1) Extension to dynamical p-brane

- Dynamical p-brane on the complex line bundle over CP^n space.
- We will be able to describe p-brane collision.

(2) Black hole

- Black hole on orbifolds (complex line bundle over CP^n space).
- Time dependent black hole.